

Hawking Radiation, Superradiance, and Dark Sector

in collaboration with

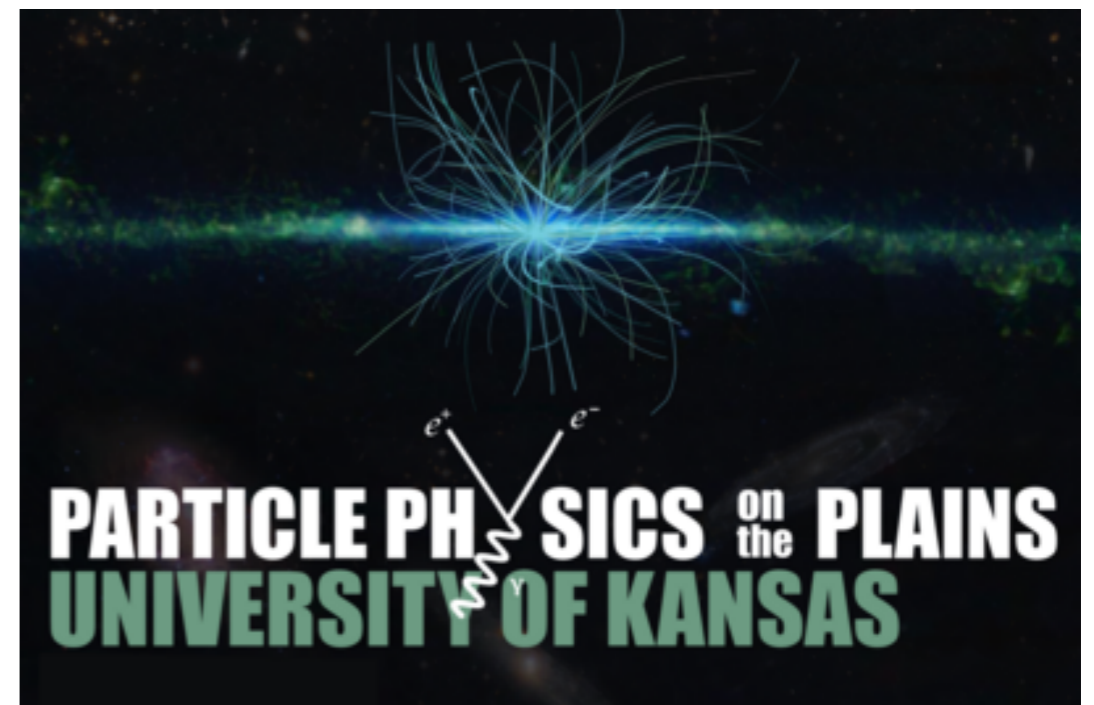
**Kaustubh Agashe, Jae Hyeok Chang,
Steven J. Clarks, Bhaskar Dutta, Yuhsin Tsai**
arXiv: 2212.11980

James B. Dent, Bhaskar Dutta
in preparation

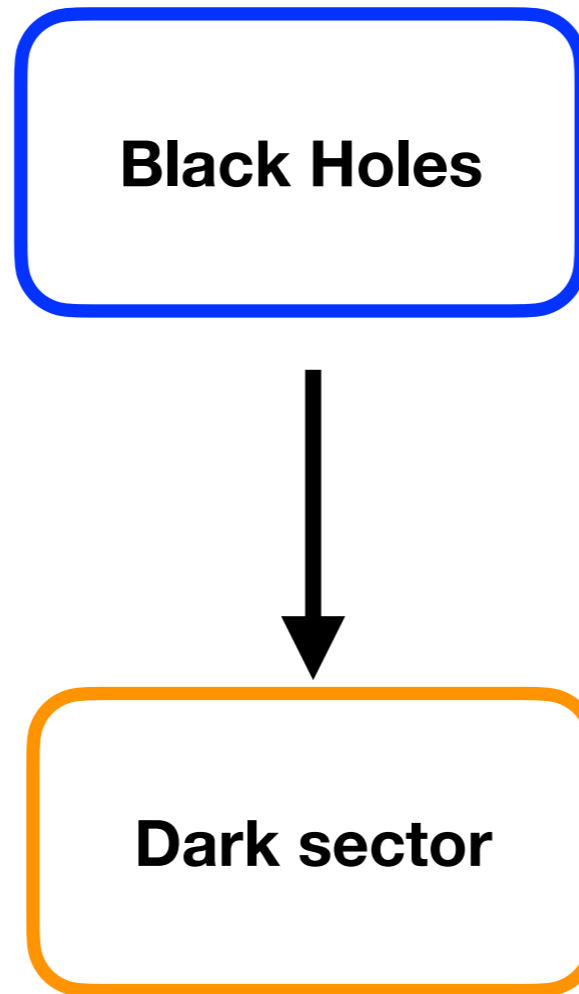


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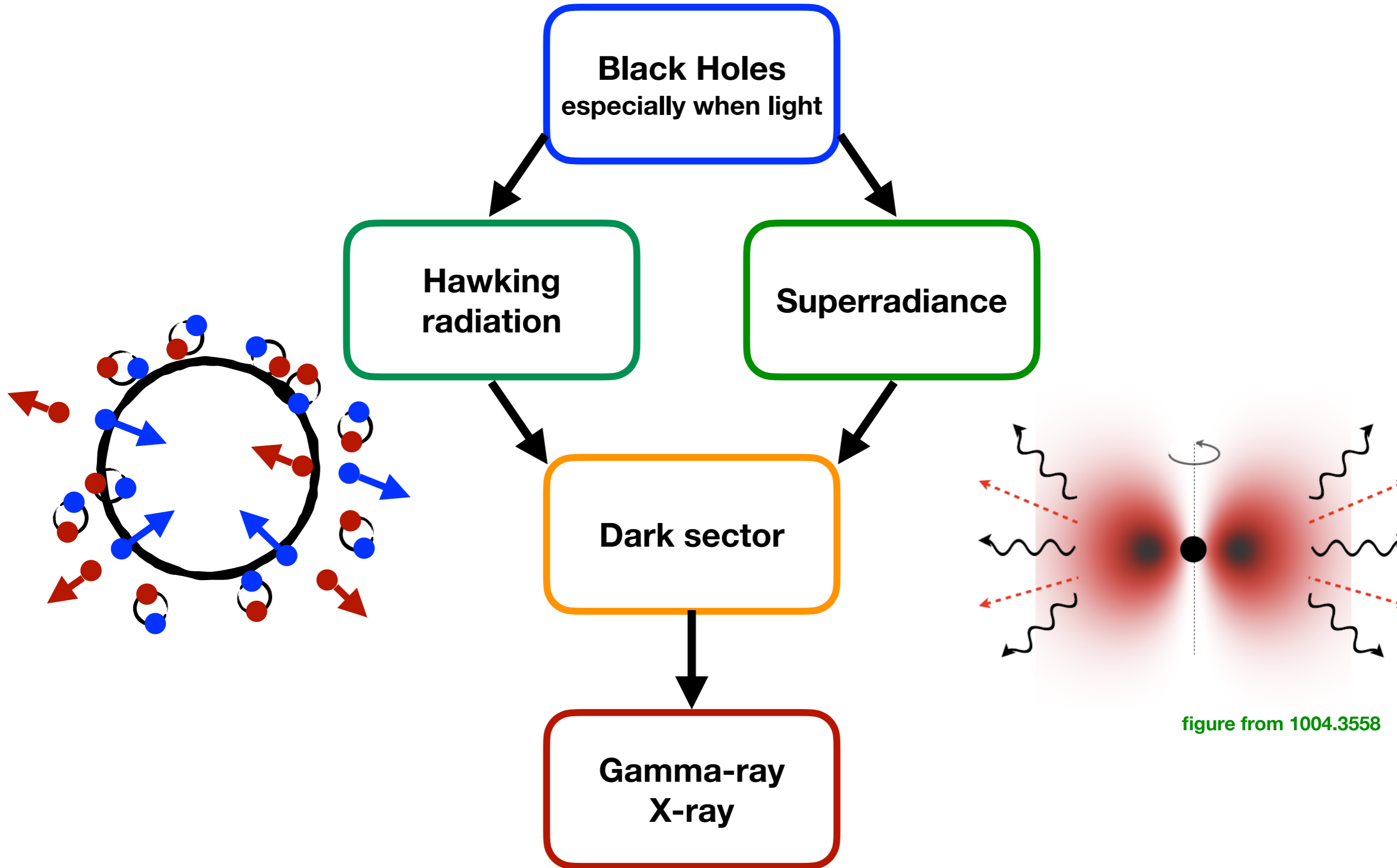


Motivations

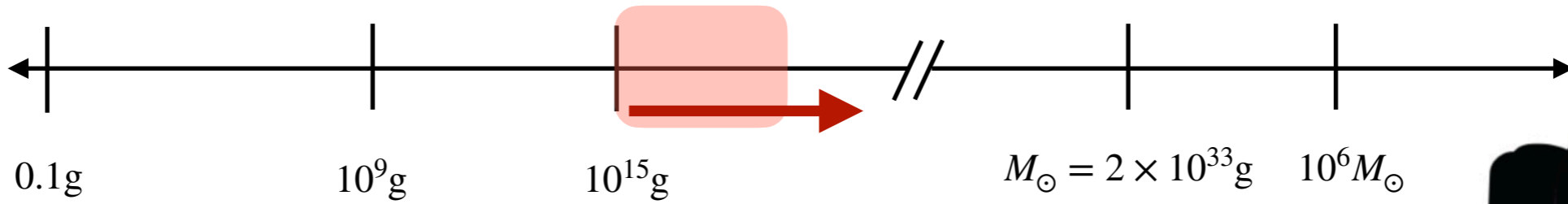


Can we use black holes to probe dark sector?

Outline



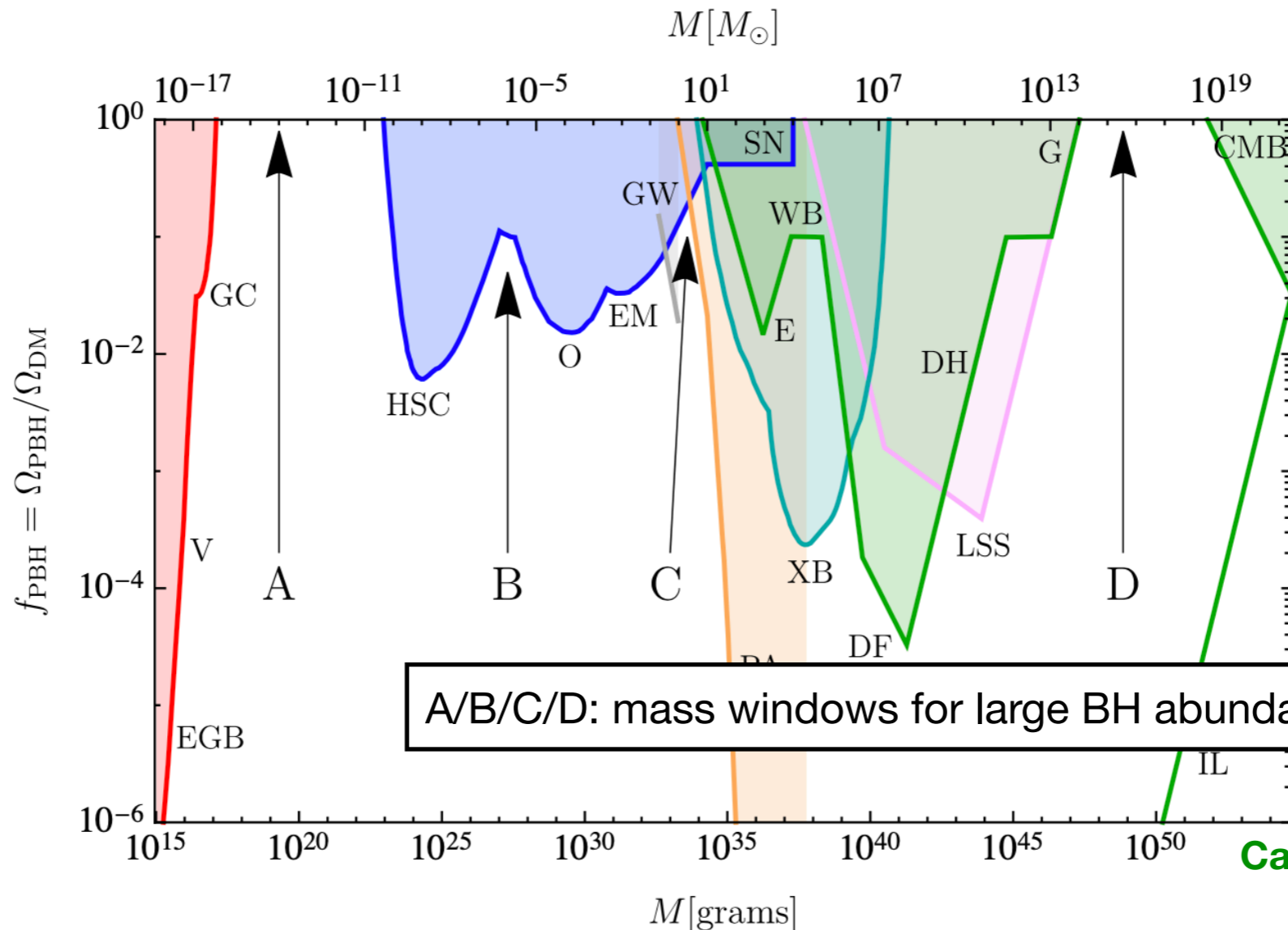
Black Holes



Asteroid mass

fraction of DM made of PBHs

$$f_{\text{PBH}} = \frac{\Omega_{\text{PBH}}}{\Omega_{\text{DM}}}$$



Carr, Kuhnel, 2006.02838

constraints from

evaporation, lensing, gravitational waves, dynamical effects, accretion, CMB distortion, large scale structure

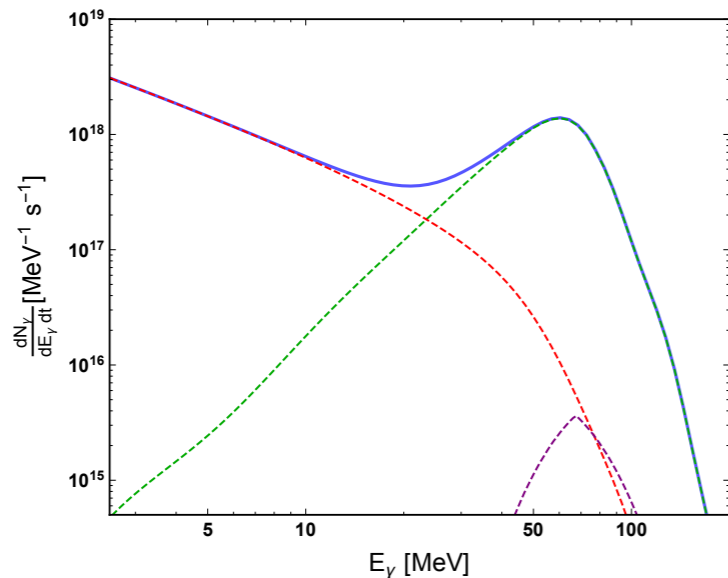
Detect BHs via Hawking Radiation

- Hawking temperature:

$$T_{\text{PBH}} = \frac{1}{8\pi GM_{\text{PBH}}} \simeq 10.5 \left(\frac{10^{15} \text{ g}}{M_{\text{PBH}}} \right) \text{ MeV}$$

- Particle emission nearly **blackbody**

$$\frac{dN_i}{dE_i dt} = \frac{g_i}{2\pi} \frac{\Gamma_i(E_i, m, m_i)}{e^{E_i/T_{\text{PBH}}} \pm 1}$$

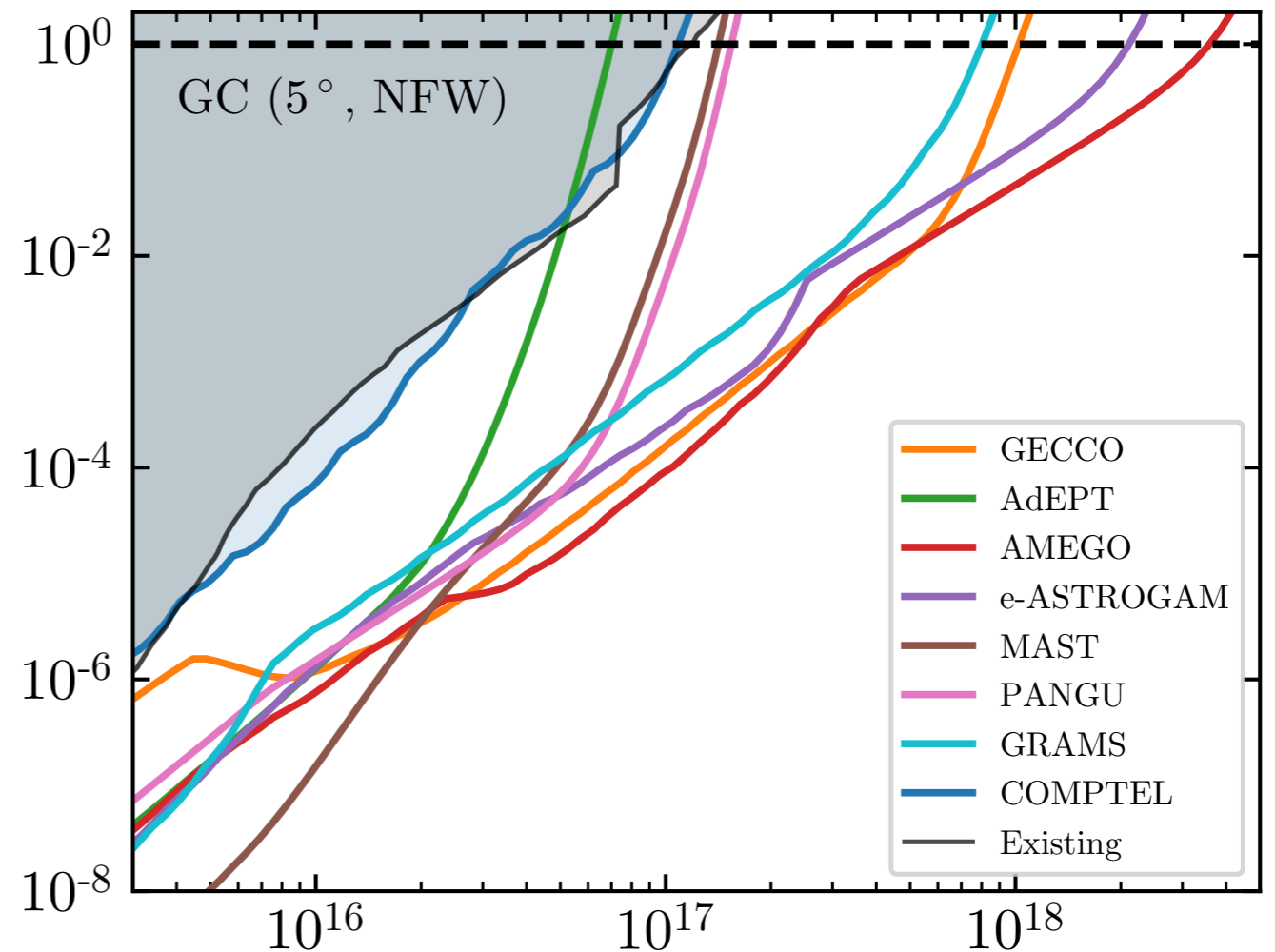


- We test light BHs with particles produced at **MeV-GeV** energy for example, gamma-rays

Future constraint

or

Opportunity of observation?



A. Coogan, L. Morrison, S. Profumo, 2010.04797

Hawking radiation rate of particle i from a non-rotating BH:

$$\frac{\partial N_i}{\partial E_i \partial t} = \frac{g_i}{2\pi} \frac{\Gamma_i}{e^{E_i/T_{\text{PBH}}} \pm 1}$$

- particle mass **kinematically allowed** $m_i \lesssim E_i \lesssim T_{\text{PBH}}$

Asteroid-mass BHs can produce MeV or lighter particles

- production via gravity, only depends on **degree of freedom** g_i , **not coupling**

Hawking radiation is good at producing weakly coupled particles in the spectrum

- how to detect the effect?
 - energy scale determined by Hawking temperature
 - large BSM particle production rate, modify the radiation spectrum
 - clear SM “background” spectrum from Hawking radiation calculation

ALP from BHs

- If exists an **Axion-Like-Particle** in the particle spectrum
- Gamma-ray spectrum is modified by ALPs: double peak

$$\frac{\partial N_{\gamma,\text{tot}}}{\partial E_{\gamma} \partial t} = \frac{\partial N_{\gamma,\text{primary}}}{\partial E_{\gamma} \partial t}$$

primary photon

$$+ \sum_{i=e^{\pm}, \mu^{\pm}, \pi^{\pm}} \int dE_i \frac{\partial N_{i,\text{primary}}}{\partial E_i \partial t} \frac{dN_{i,\text{FSR}}}{dE_{\gamma}}$$

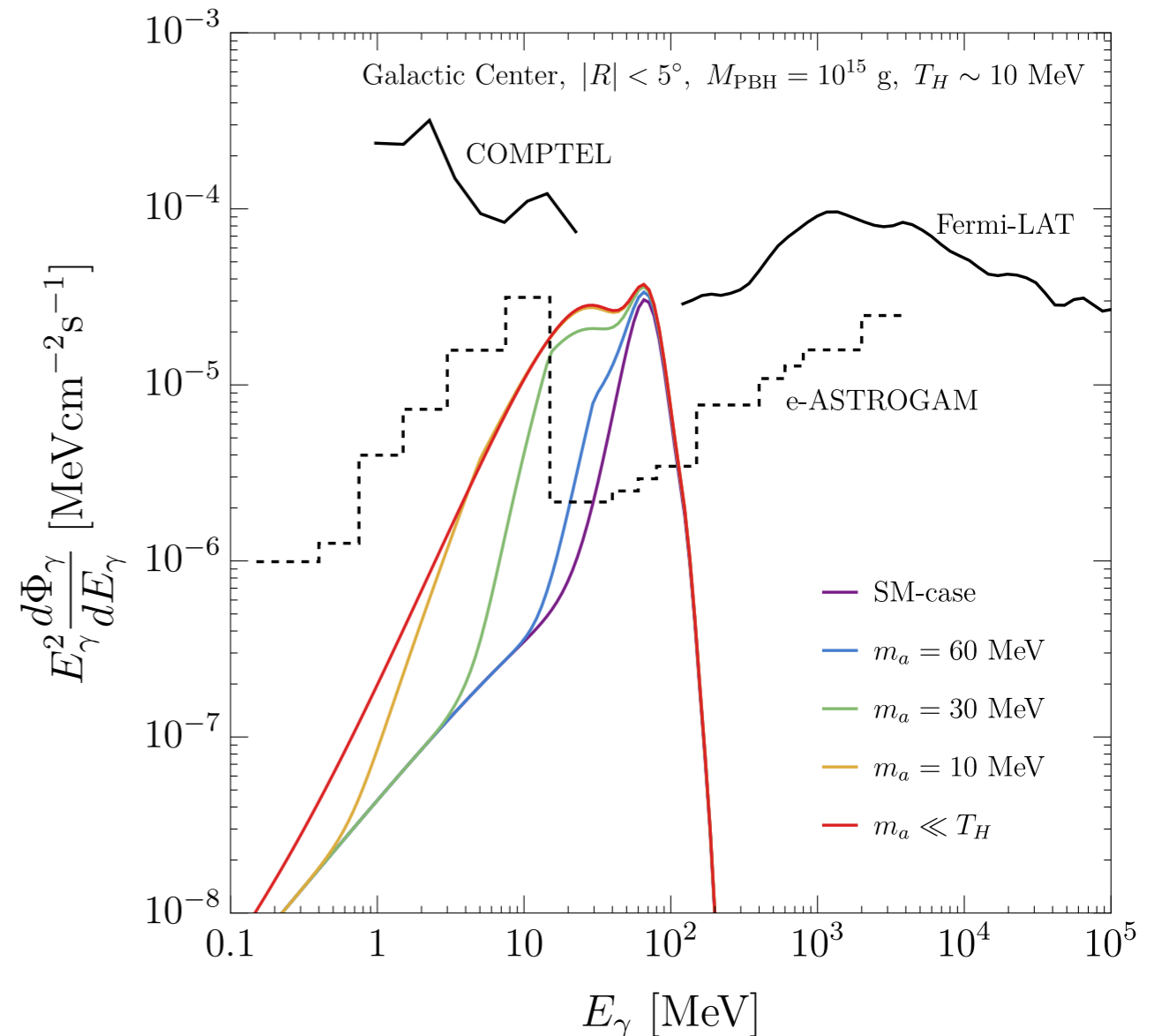
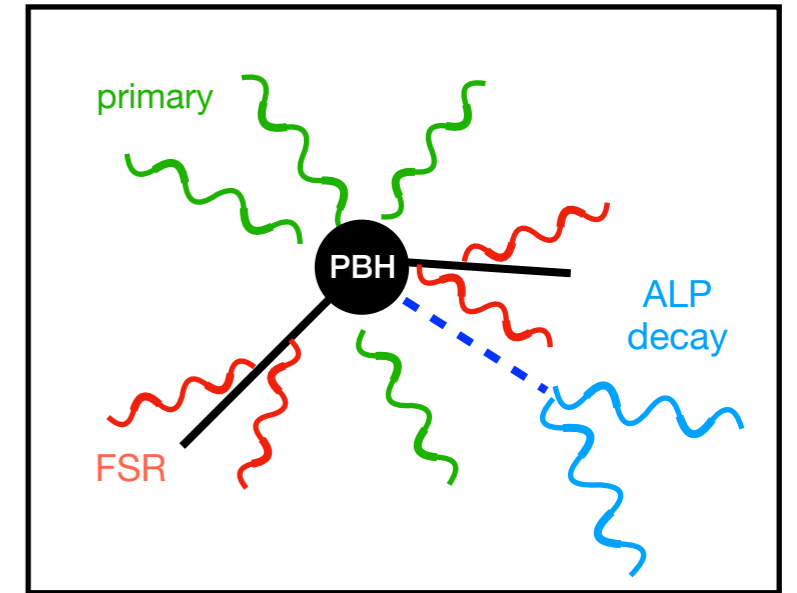
final-state radiation

$$+ \sum_{i=\pi^0} \int dE_i 2 \frac{\partial N_{i,\text{primary}}}{\partial E_i \partial t} \frac{dN_{i,\text{decay}}}{dE_{\gamma}}$$

pion decay

$$+ \int dE_a 2 \frac{\partial N_{a,\text{primary}}}{\partial E_a \partial t} \frac{dN_{a,\text{decay}}}{dE_{\gamma}}$$

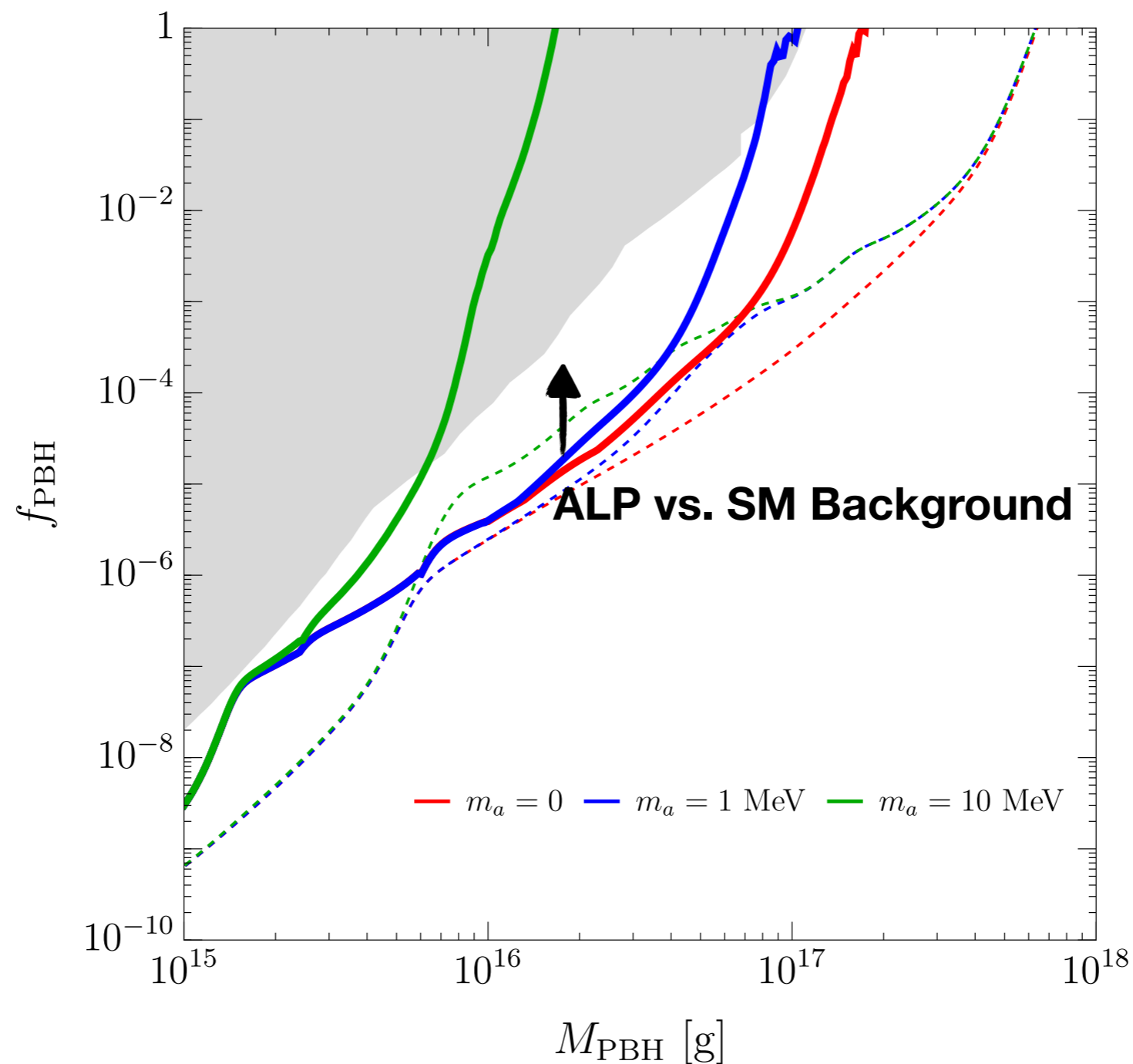
add ALP decay



Identification of ALPs

If f_{PBH} is larger than the detection limit, enough statistics to **distinguish** the ALP.

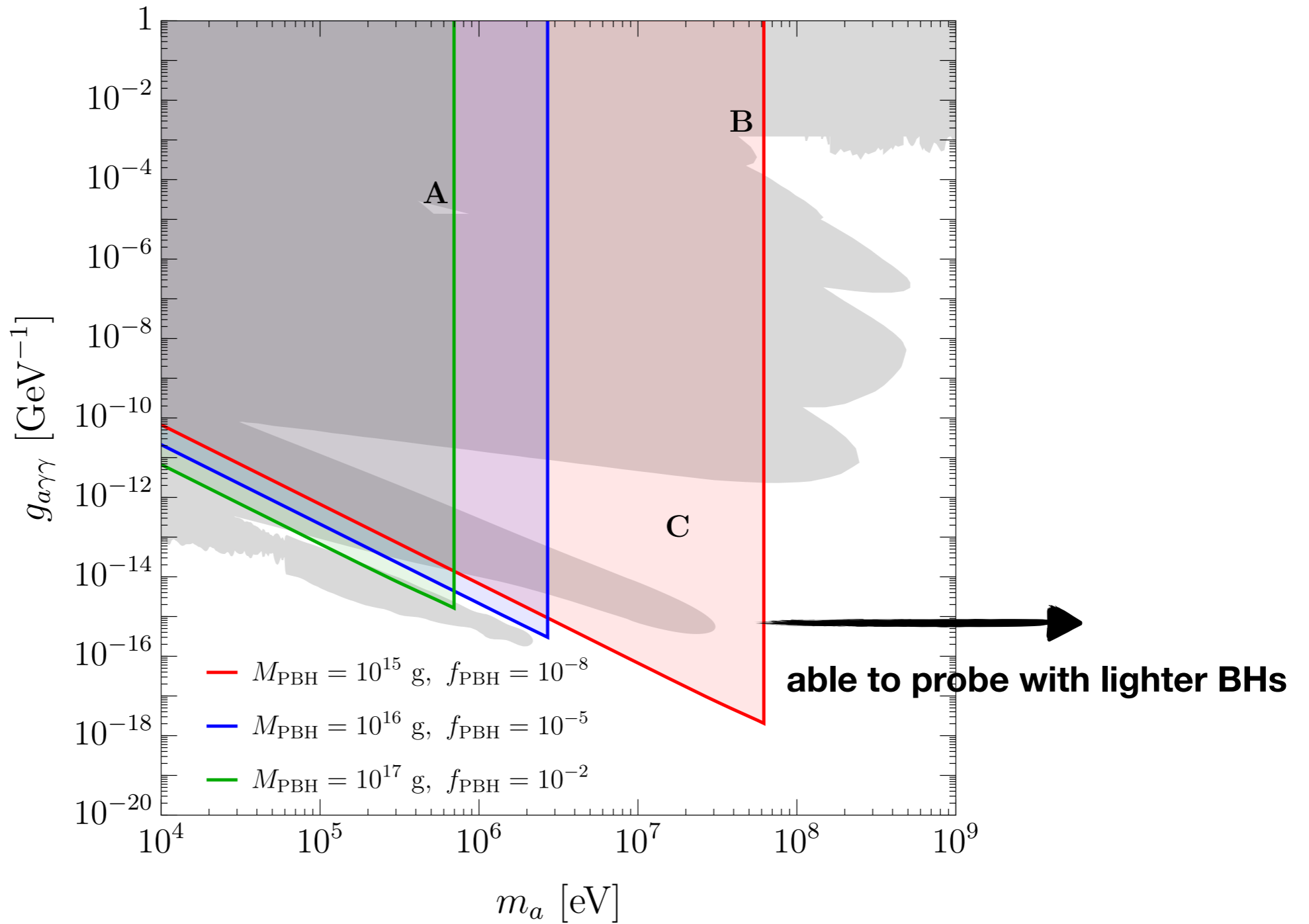
We will be able to know if ALP exists from the shape of gamma-ray spectrum.



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Identification of ALPs

ALP parameter space that can be probed with BHs.

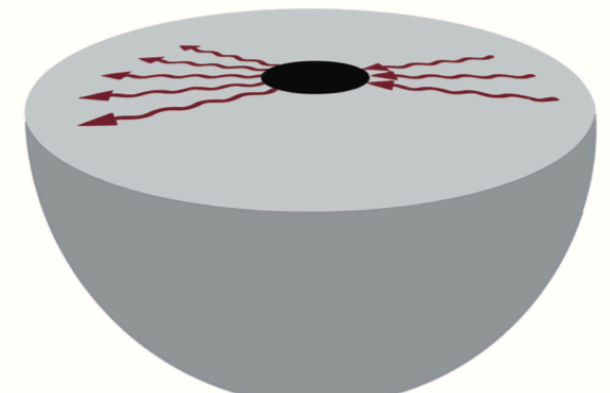
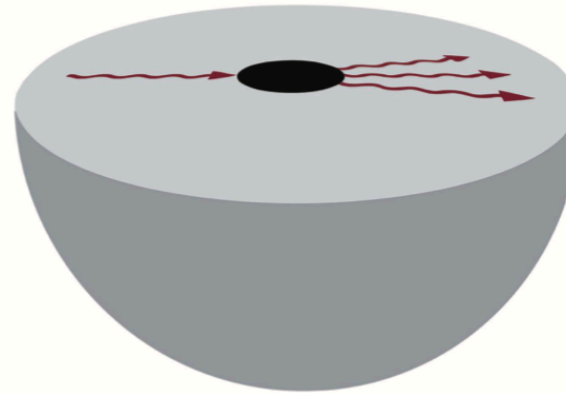
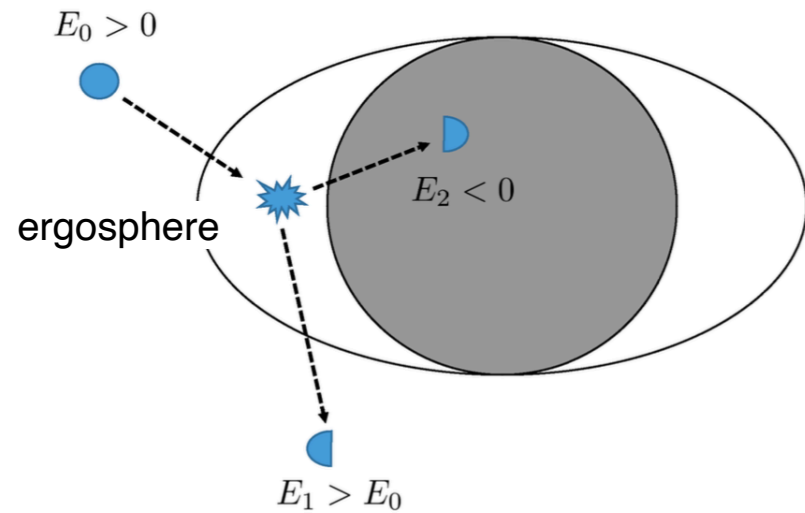


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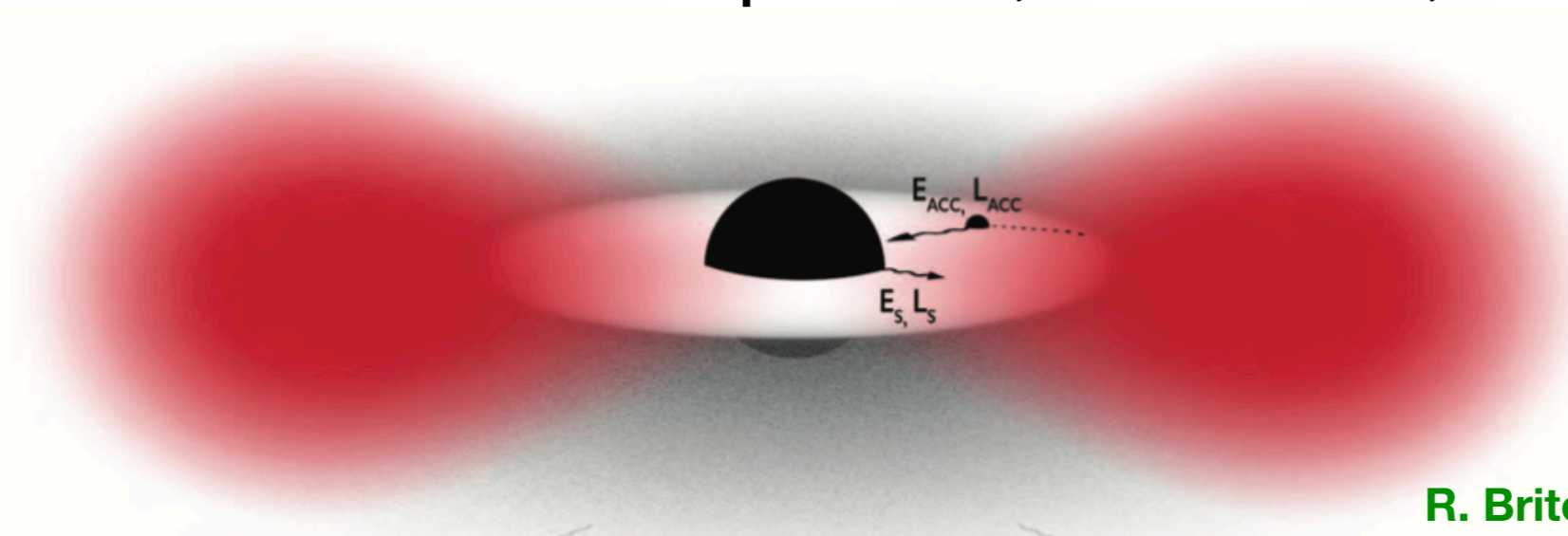
2212.11980

Superradiance

So far we only talked about BH mass, how about BH spin?



amplified \longrightarrow confined \longrightarrow amplified again



R. Brito, V. Cardoso, P. Pani
1501.06570

BH can produce **massive bosons** with BH angular momentum when $\omega_a < m \Omega$

Ω : BH angular velocity
 m : azimuthal quantum number

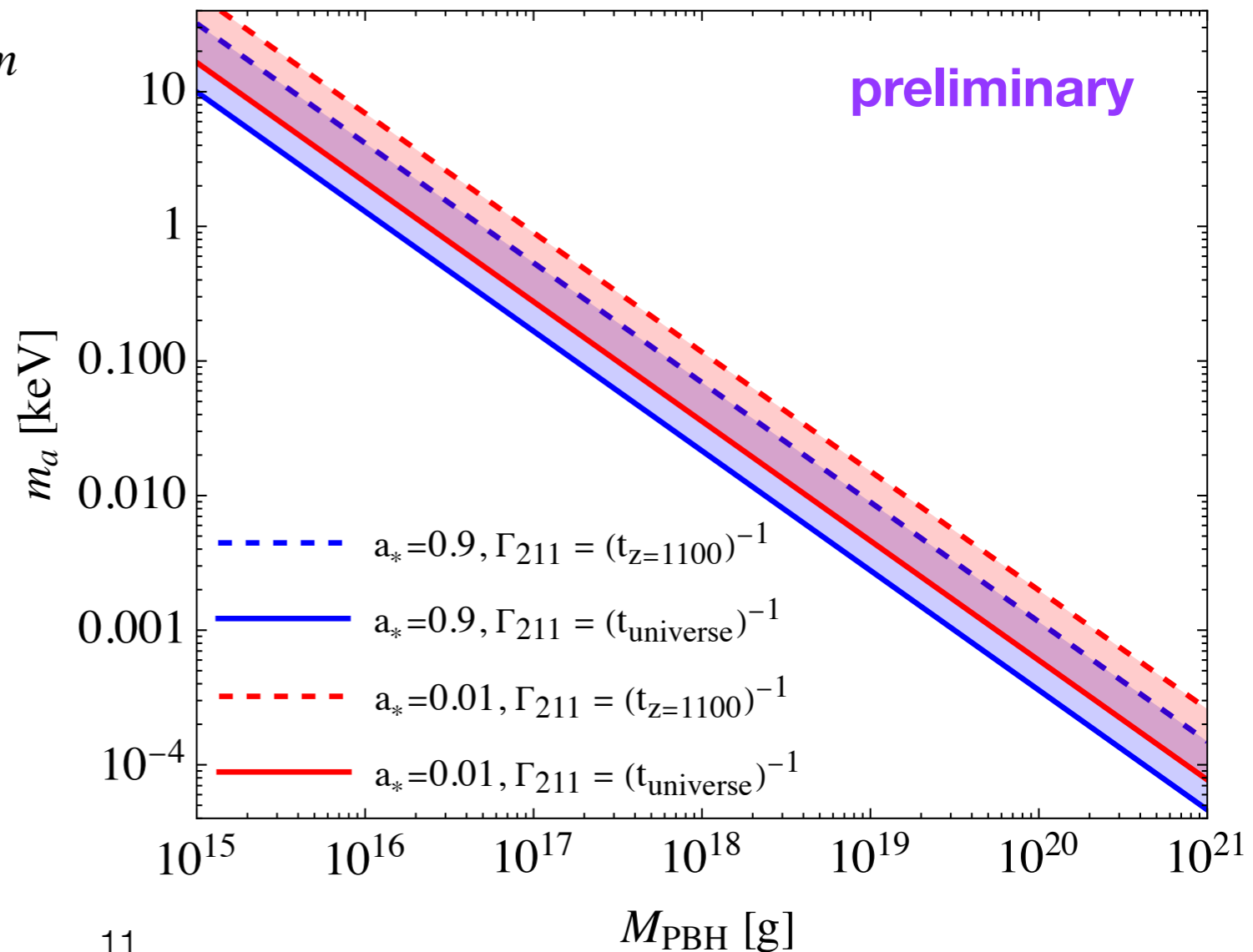
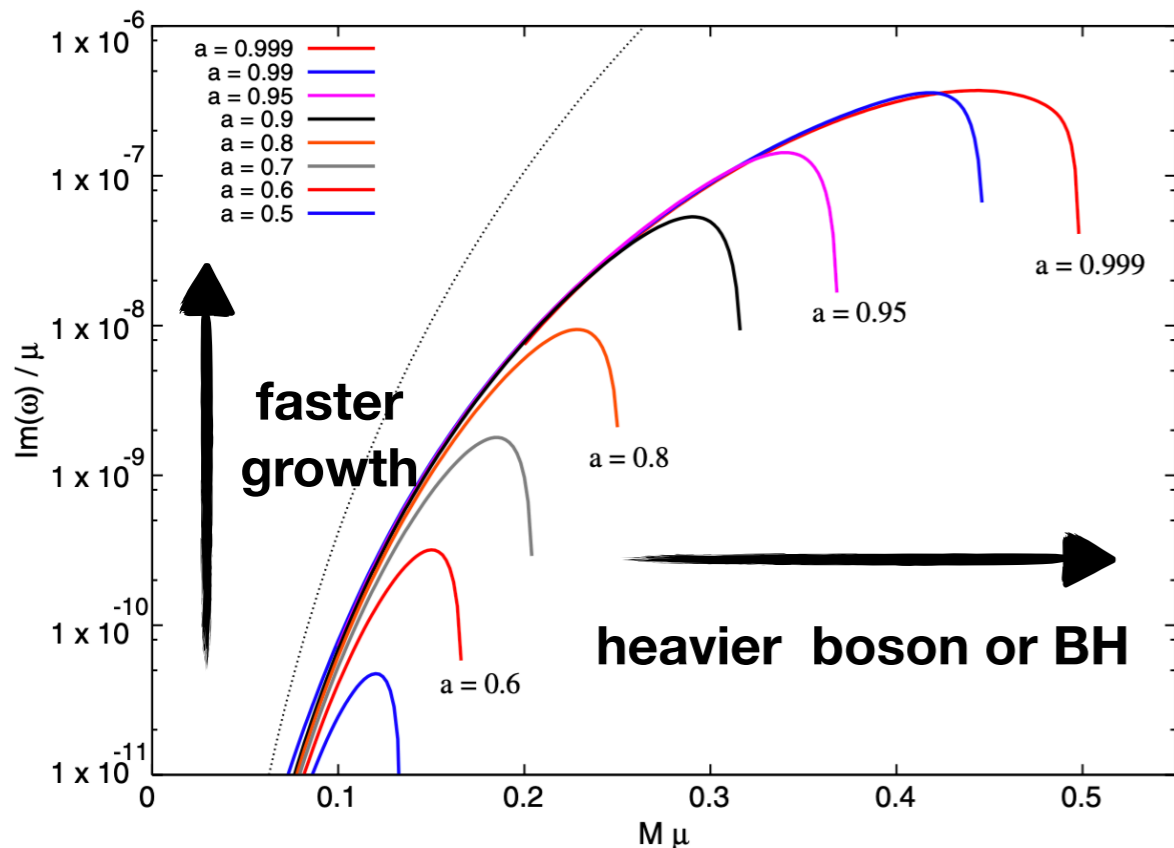
Superradiance Condition and Rate

- Gravitational coupling between BH and axion: $\alpha = G_N M m_a$
- Frequency of axion mode bounded by BH: $\omega = \omega_R + i\omega_I$

$$\omega_R = m_a \left(1 - \frac{\alpha^2}{2n^2}\right) \quad N_a(t) \simeq N_0 e^{2\Gamma_{nlm} t} \quad \Gamma_{nlm} = 2\omega_I$$

$$\frac{\omega_I}{m_a} \simeq \alpha^{4l+4} (m\Omega - \omega_R) 2r_+ \mathcal{C}_{nlm}$$

superradiance condition

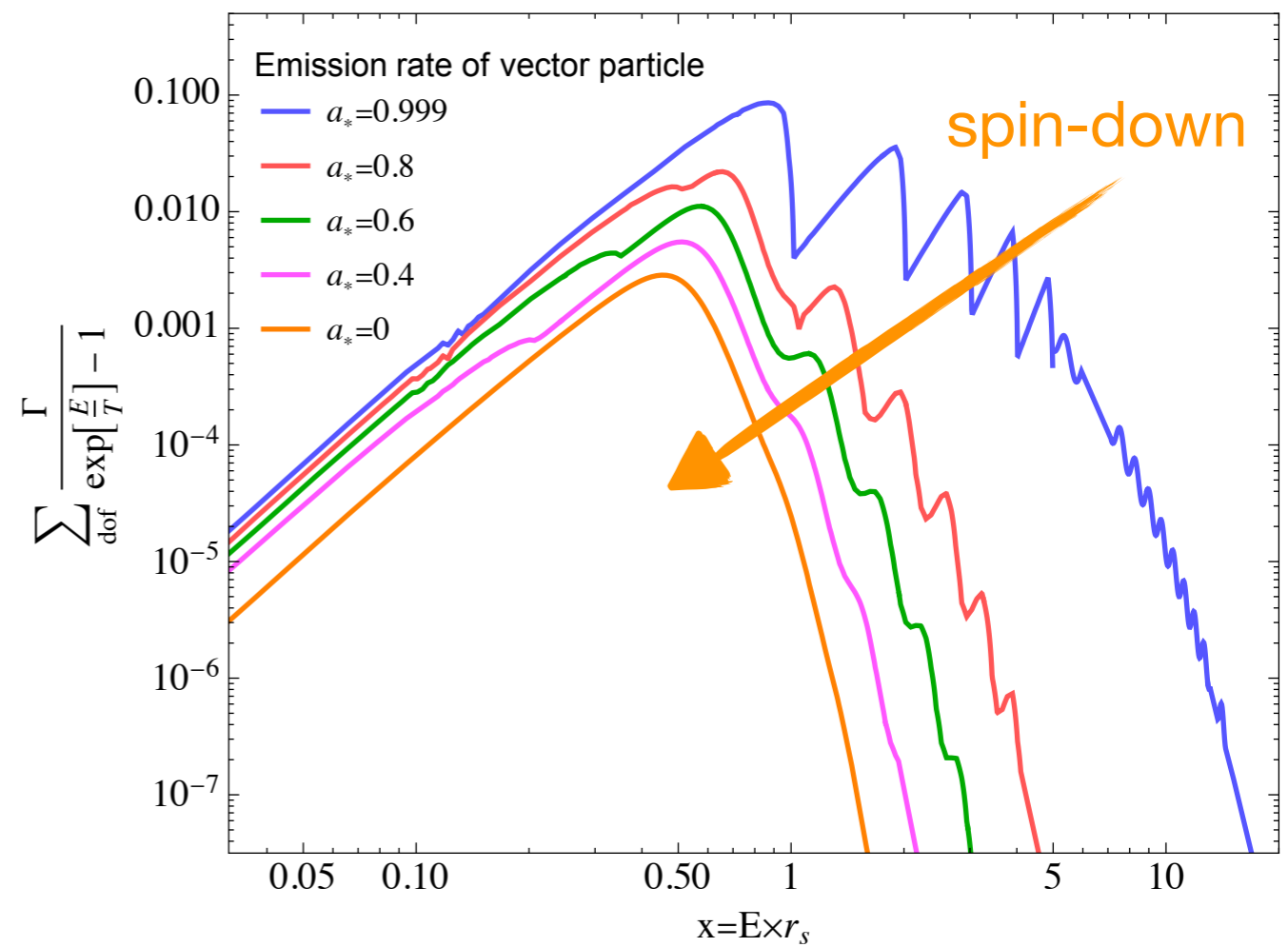
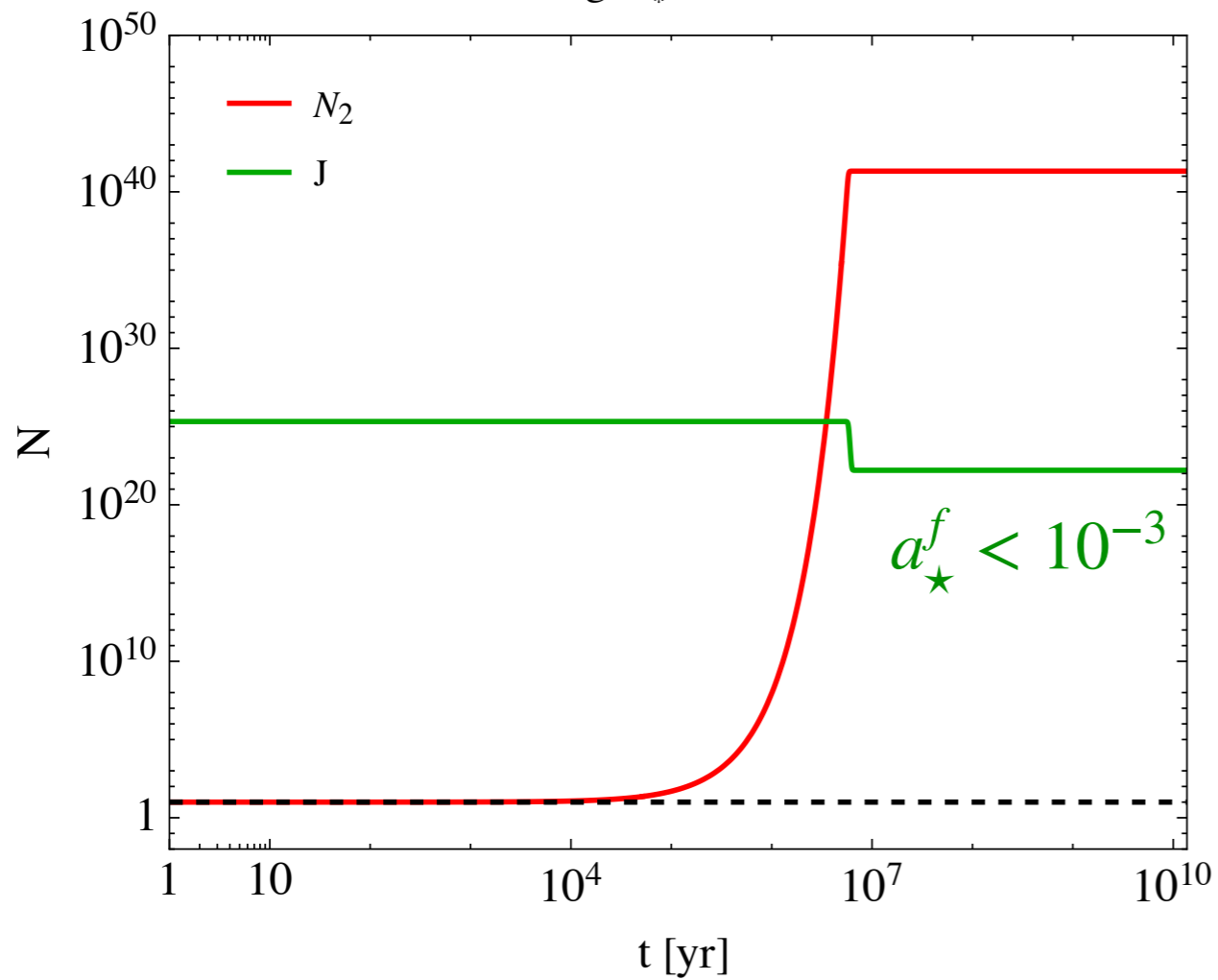


Spin-down of BH

The rotational energy of a BH is **depleted** into the axion cloud,

rotating PBH $\xrightarrow{\text{superradiance}}$ non-rotating PBH

$M_{\text{PBH}}=10^{16}\text{g}, a_*^i=0.99, m_a=5\text{keV}$

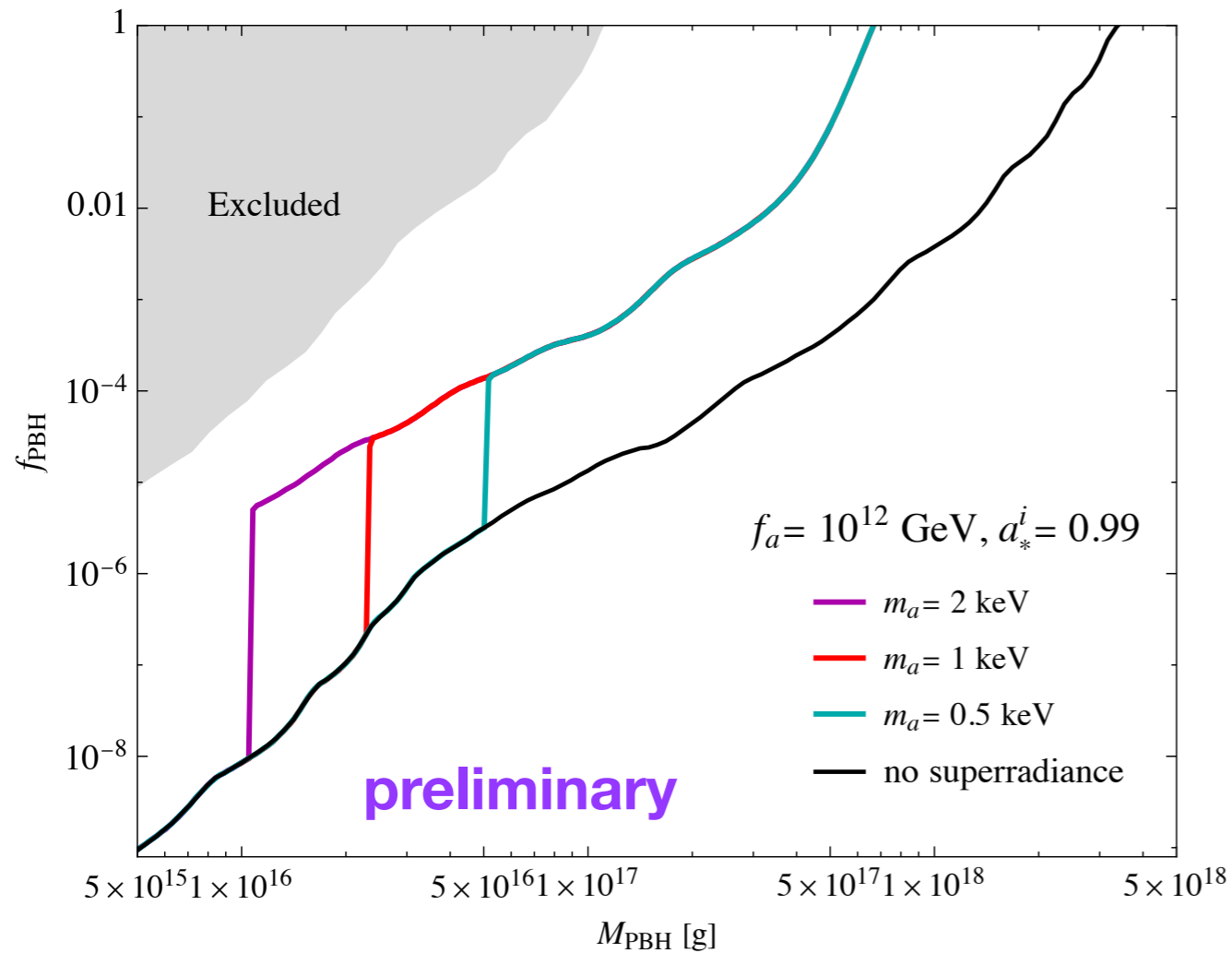


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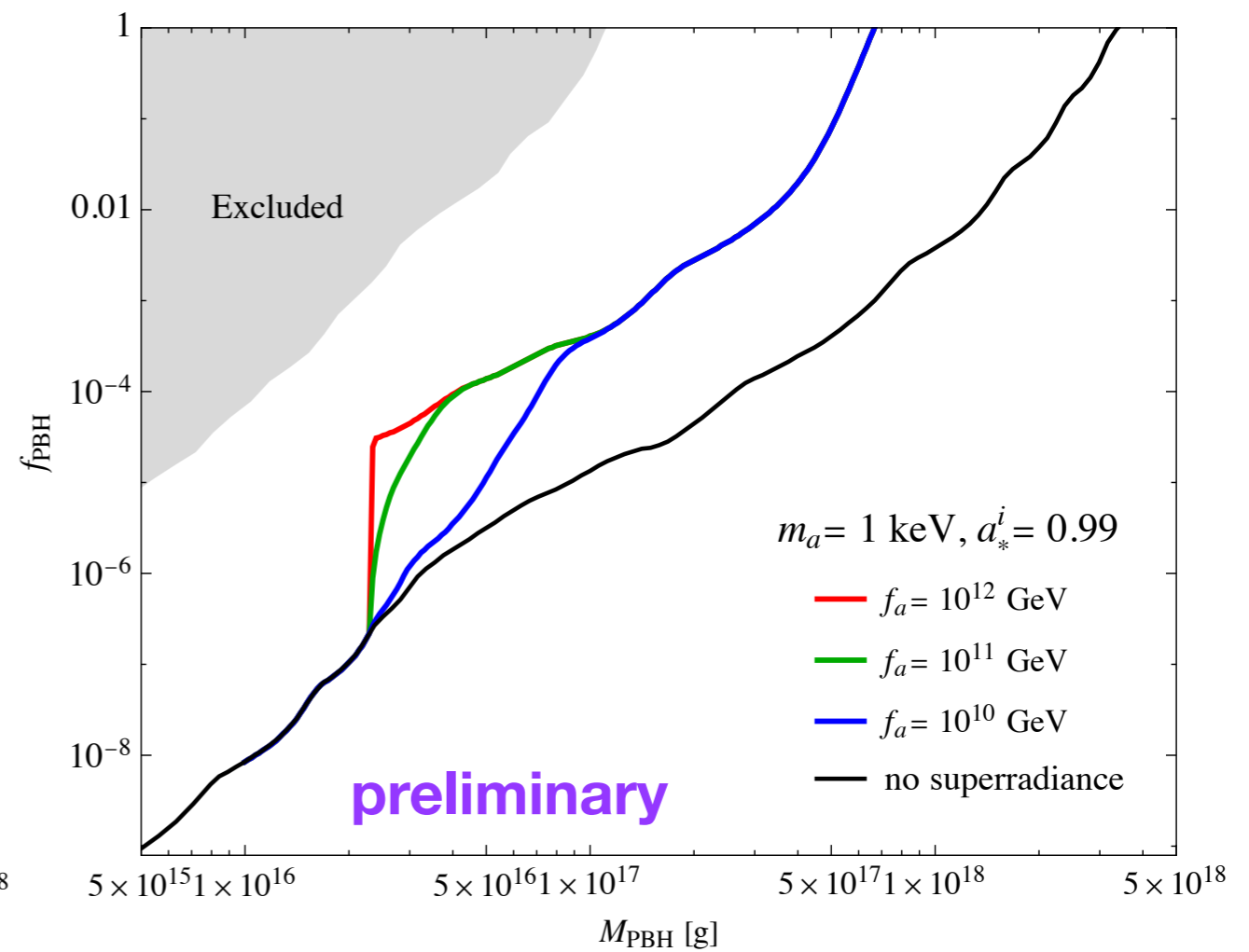
Spin-down of BH

Galactic Center searches for BHs are only sensitive to the **final spin** of the Kerr BH.

Constraints becomes weaker with superradiance.



different axion mass

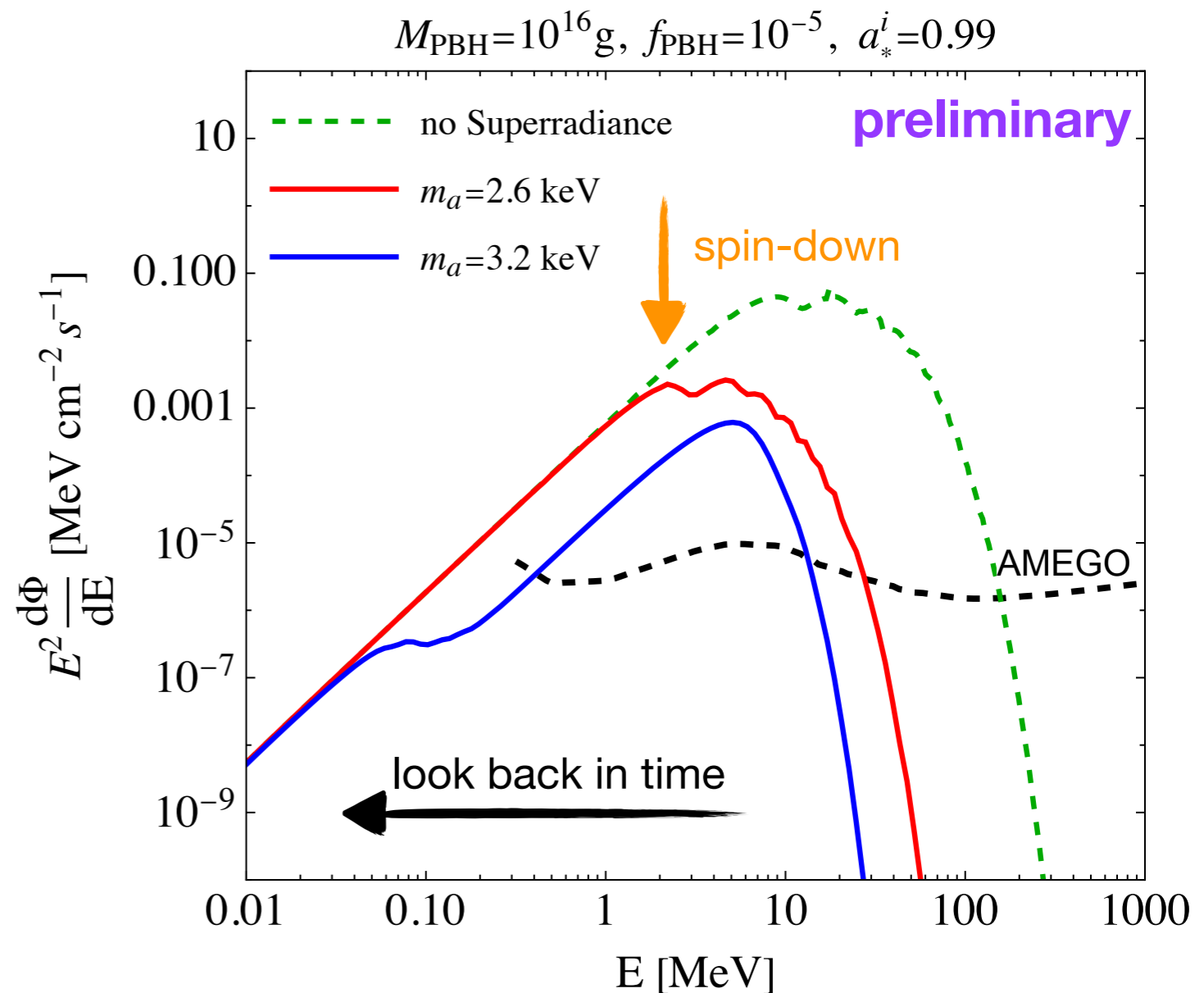


different axion self-interaction

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Extra-galactic flux can be correlated to test BH spin evolution at higher redshift.

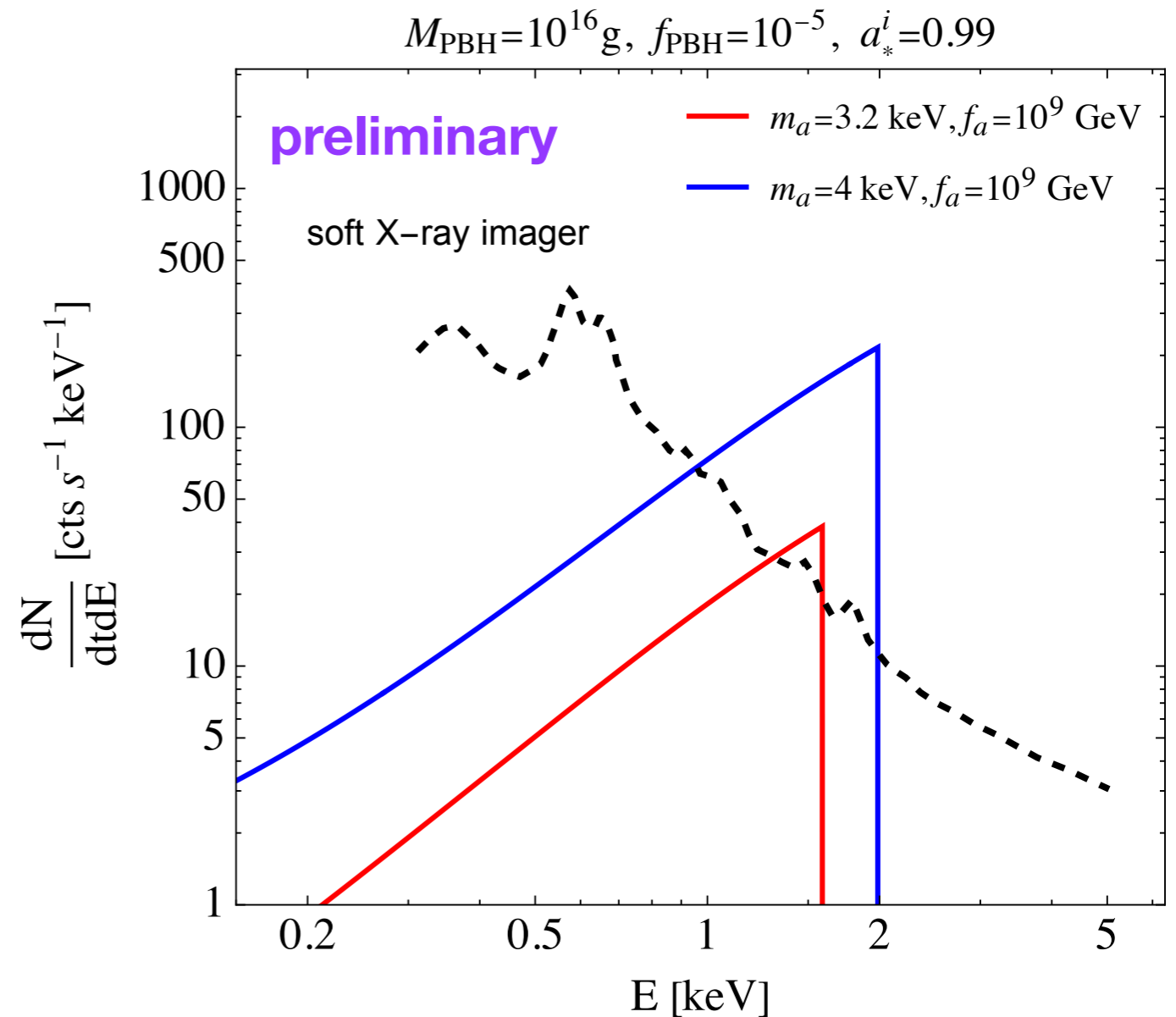
- Reduction in gamma-ray flux happens as BHs lose spin due to superradiance process
- Heavier axion spin down sooner, $\Gamma \propto \alpha^{4l+4}$, flux drops earlier at low energies



Axions produced with BH superradiance can contribute to keV lines with decay.

$$g_{a\gamma\gamma} = \frac{\alpha_{\text{EM}}}{2\pi f_a}$$

- Direct axion production signal within reach of future X-ray observations
- Signal strength depends on f_a for both axion decay width and superradiance dynamics



Summary

- Black holes can be particle factories with gravitational production. Hawking radiation and superradiance combined create interesting searches for new physics.
- To correlate a particle physics energy scale to BH physics
 - Hawking radiation: Particle mass comparable to Hawking temperature or below.
 - Superradiance: Particle wavelength matches BH radius.
Timescale of superradiance in the observation window.
- To test the idea with asteroid-mass BHs, we use ALP as an example
 - Modification of Hawking radiation spectrum from galactic center
 - Spin evolution of BHs from galactic and extra-galactic signals
 - Decay signal from the superradiance cloud
 - Outlook: a variety of possible correlative channels to probe dark sector.

Thank you!