



DEPARTMENT OF  
PHYSICS

# A New Purpose for the W-mass Measurement: Searching for New Physics via $l+mET$

Particle Physics on the Plains at University of Kansas

[2310.xxxx] by Kaustubh Agashe, Sagar Airen, Roberto Franceschini, Doojin Kim, Ashutosh Kotwal, Lorenzo Ricci, DS

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Deepak Sathyan  
October 14, 2023

# W-mass Measurement Results

## High-precision measurement of the $W$ boson mass with the CDF II detector

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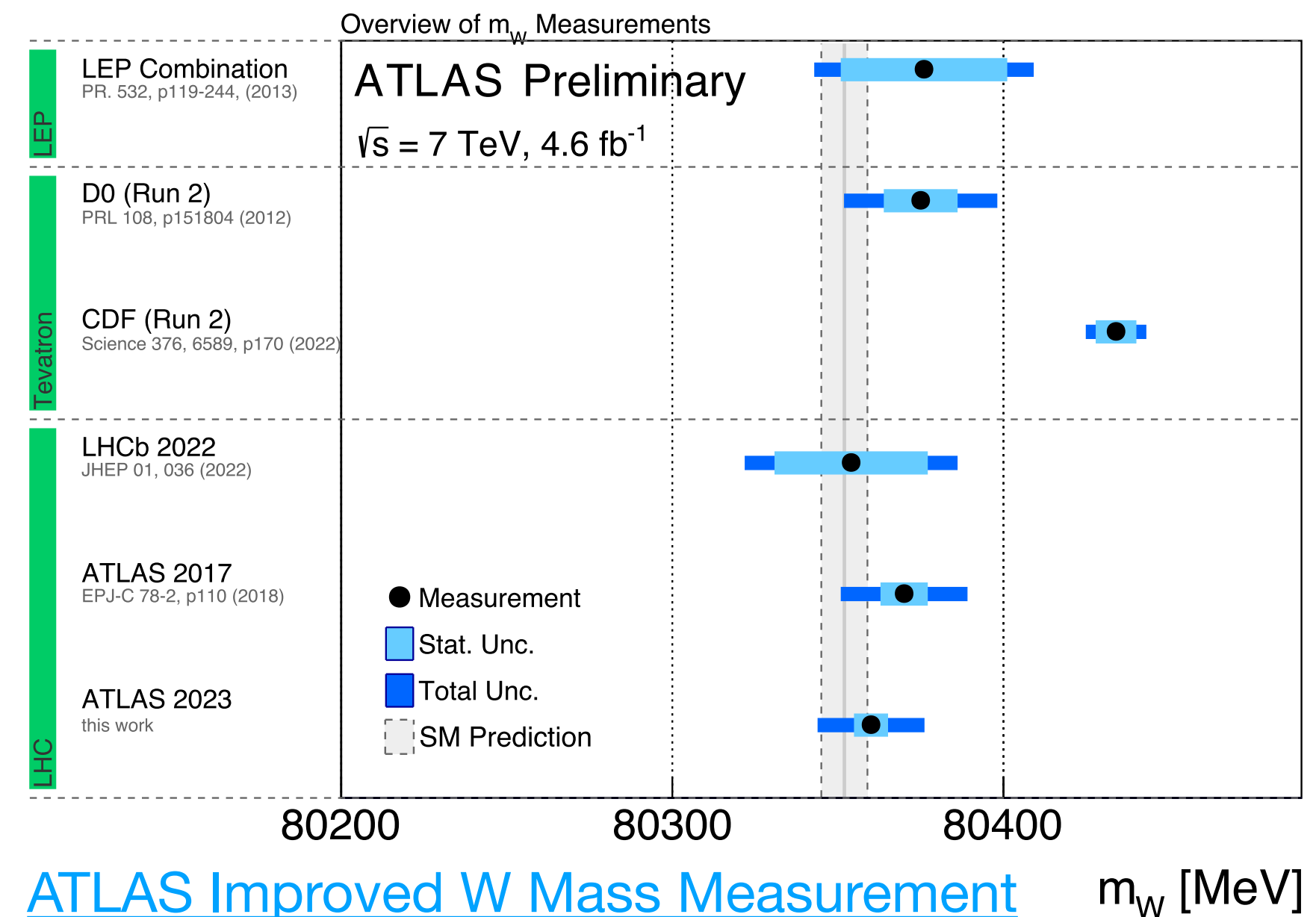
[Authors Info & Affiliations](#)

April 2022

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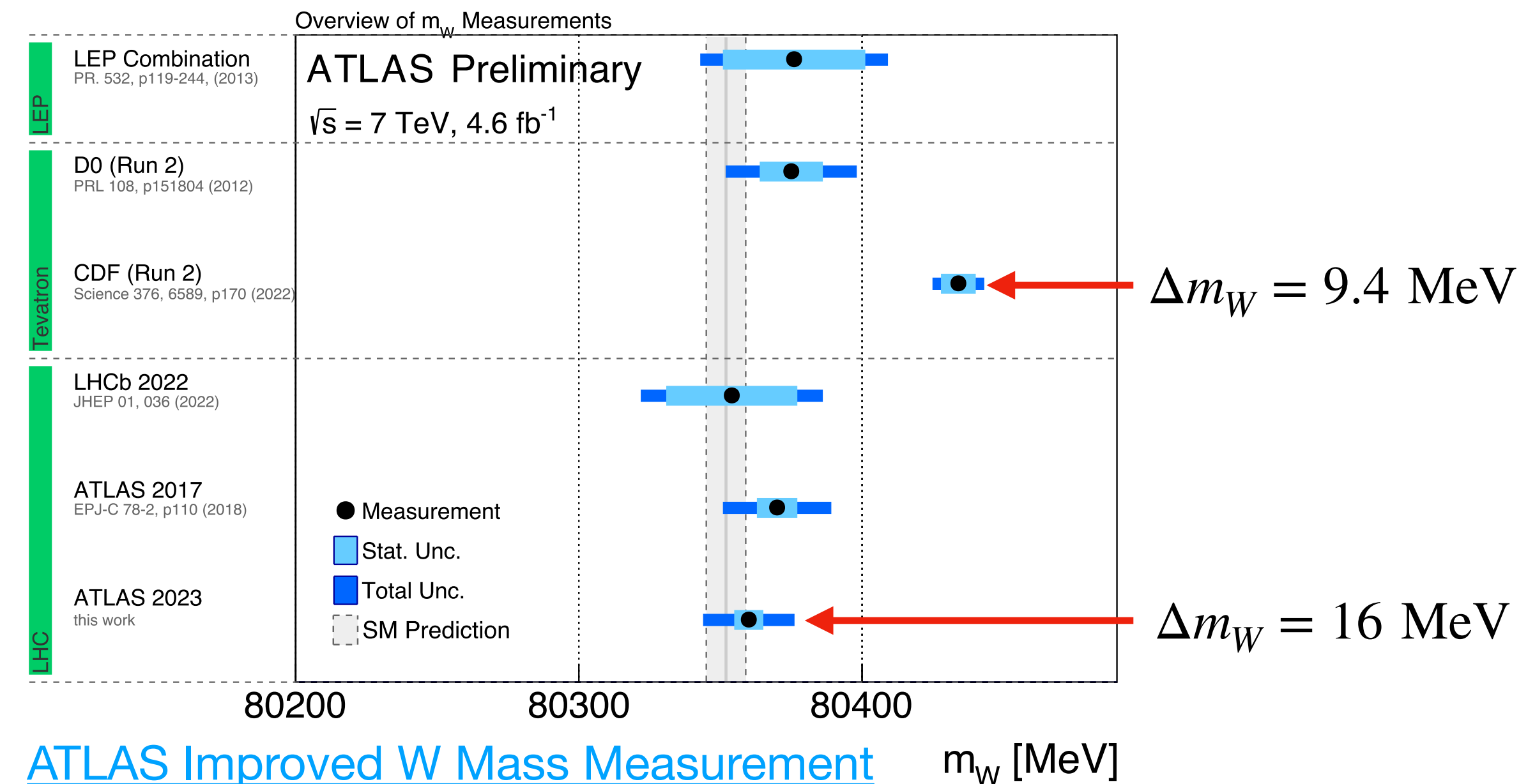
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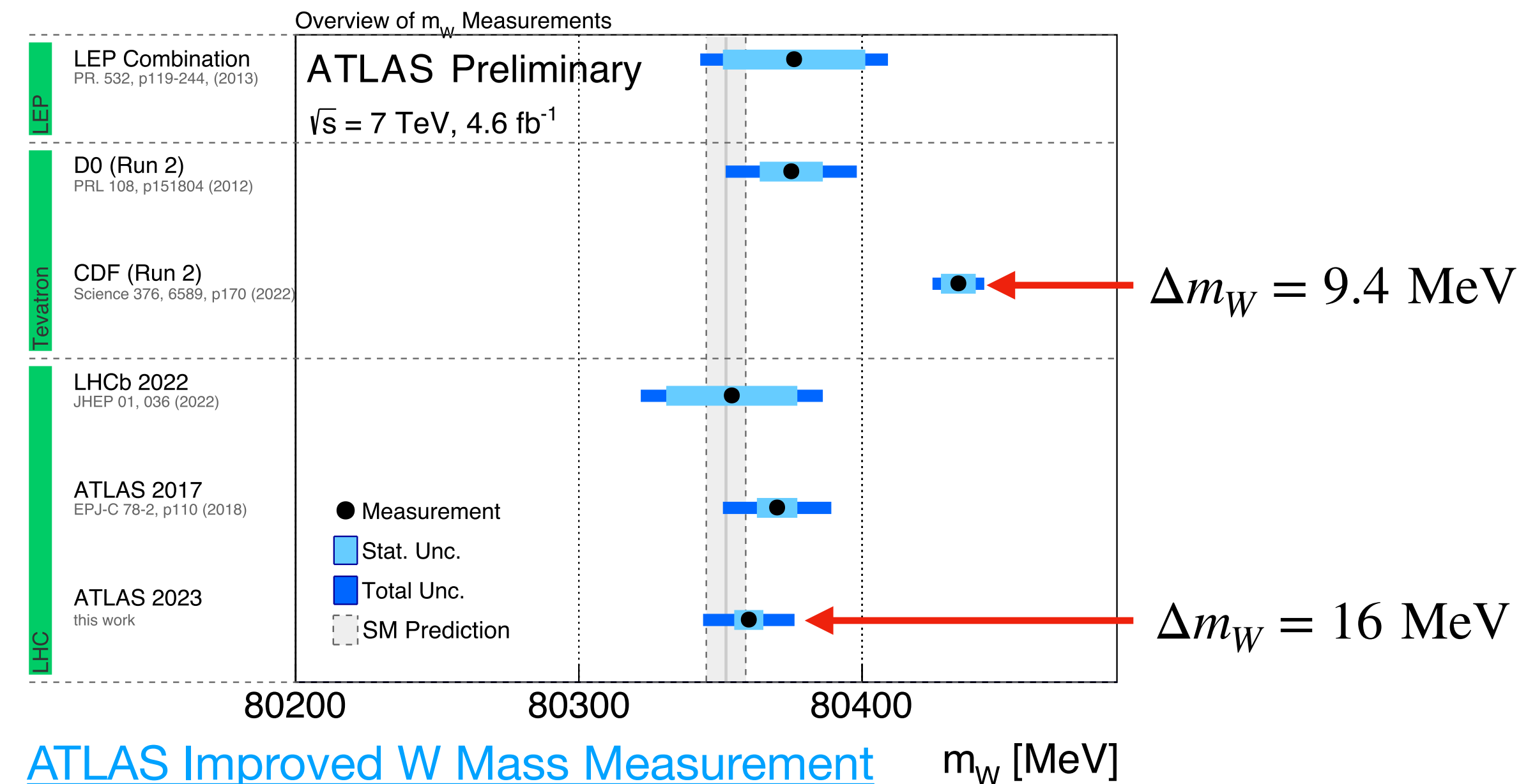
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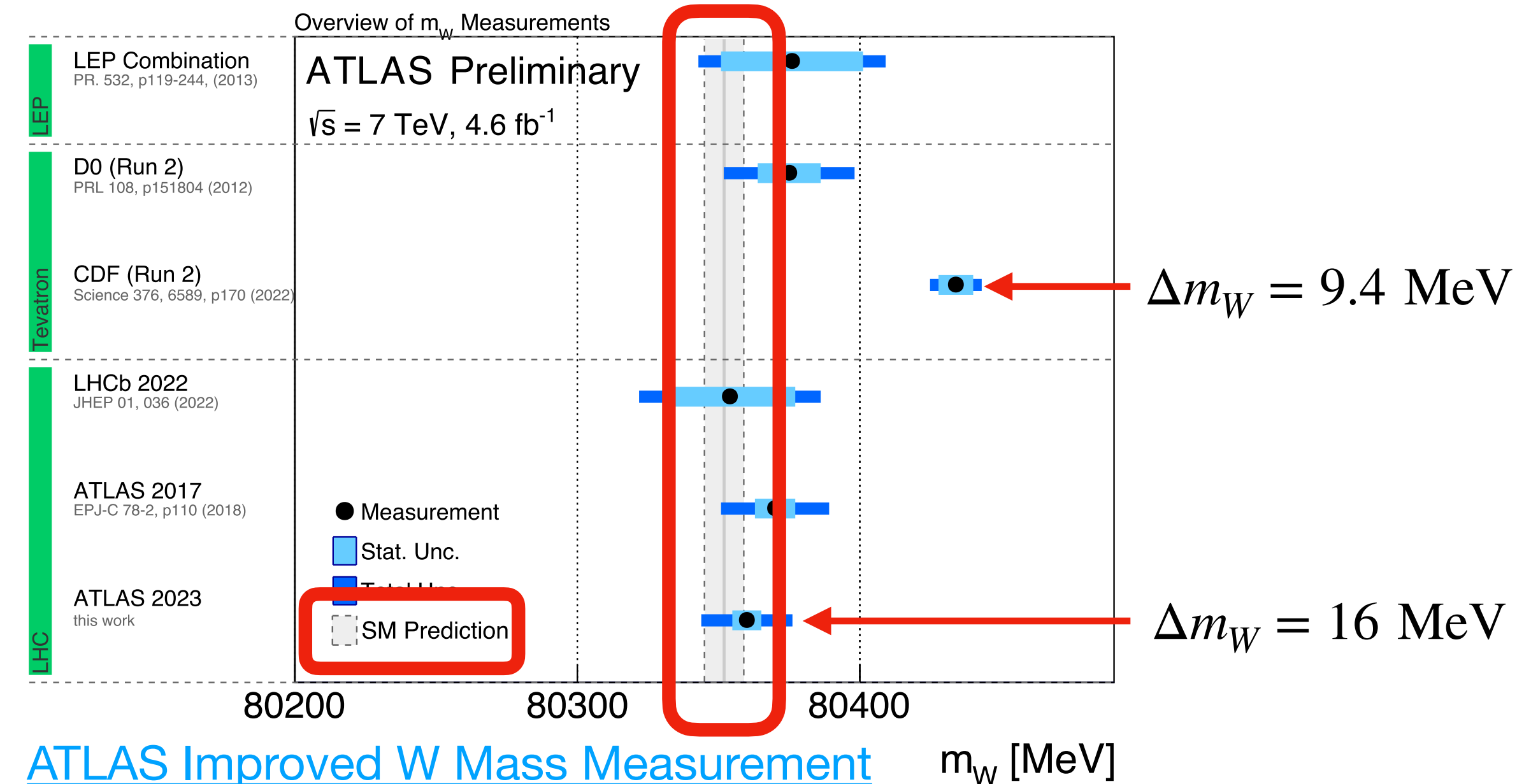
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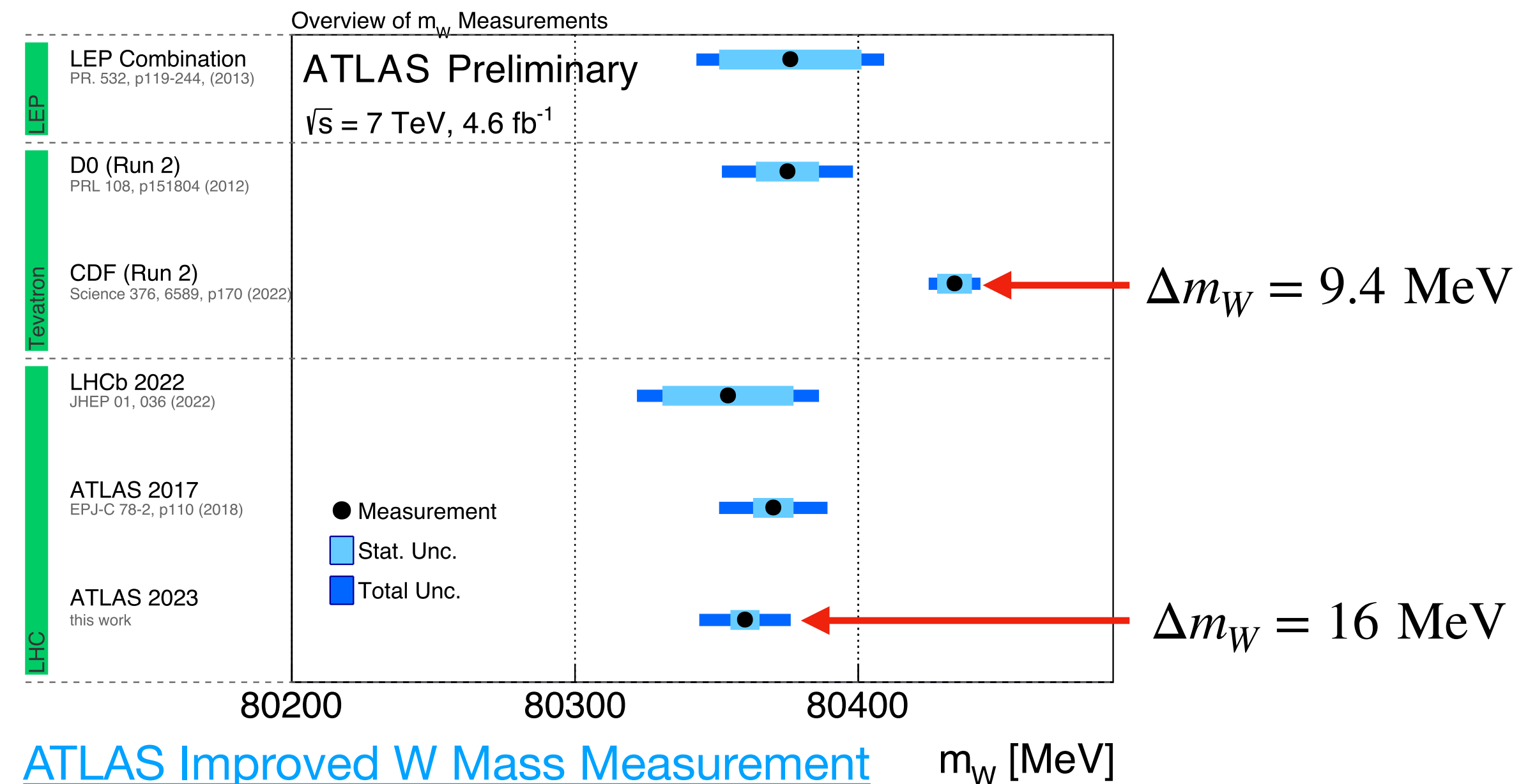
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- Consistency check of Standard Model
- Probe for new physics (NP):
  - Comparing measured  $m_W$  to prediction from EW fit tests indirect effects of heavy new physics (NP)
  - Can  $m_W^{\text{meas.}}$  alone without EW fit probe direct effects of NP? (removing EW fit from the picture)



# How is $W$ -mass measured?

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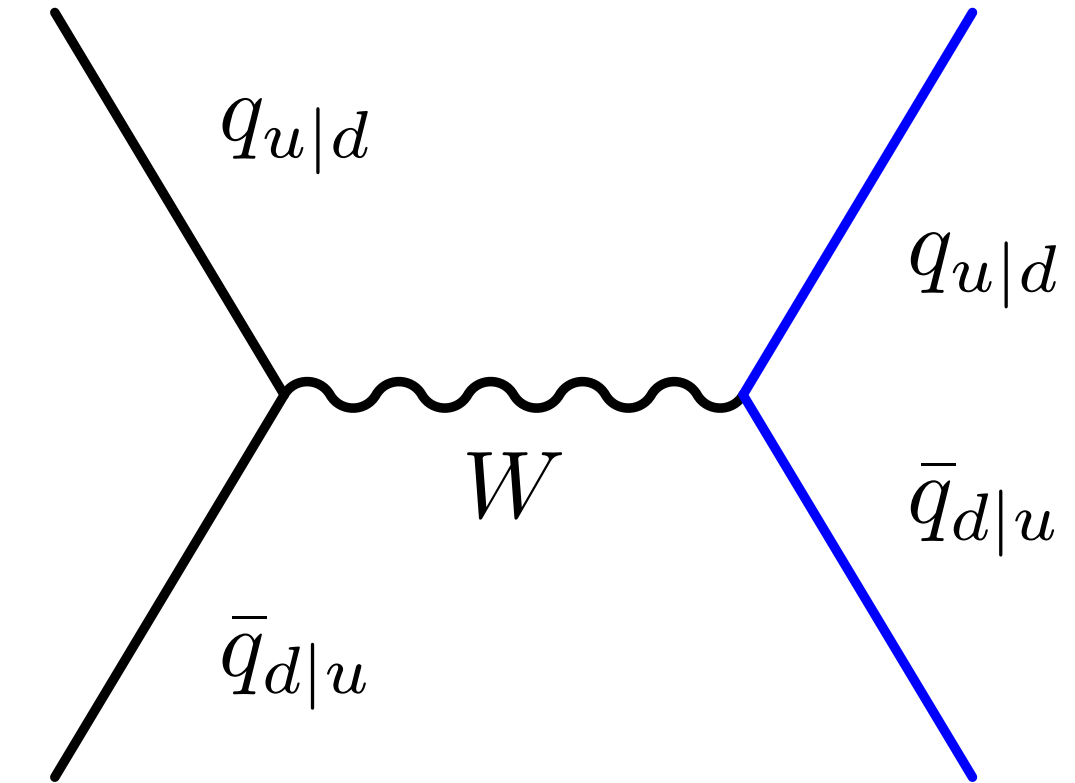
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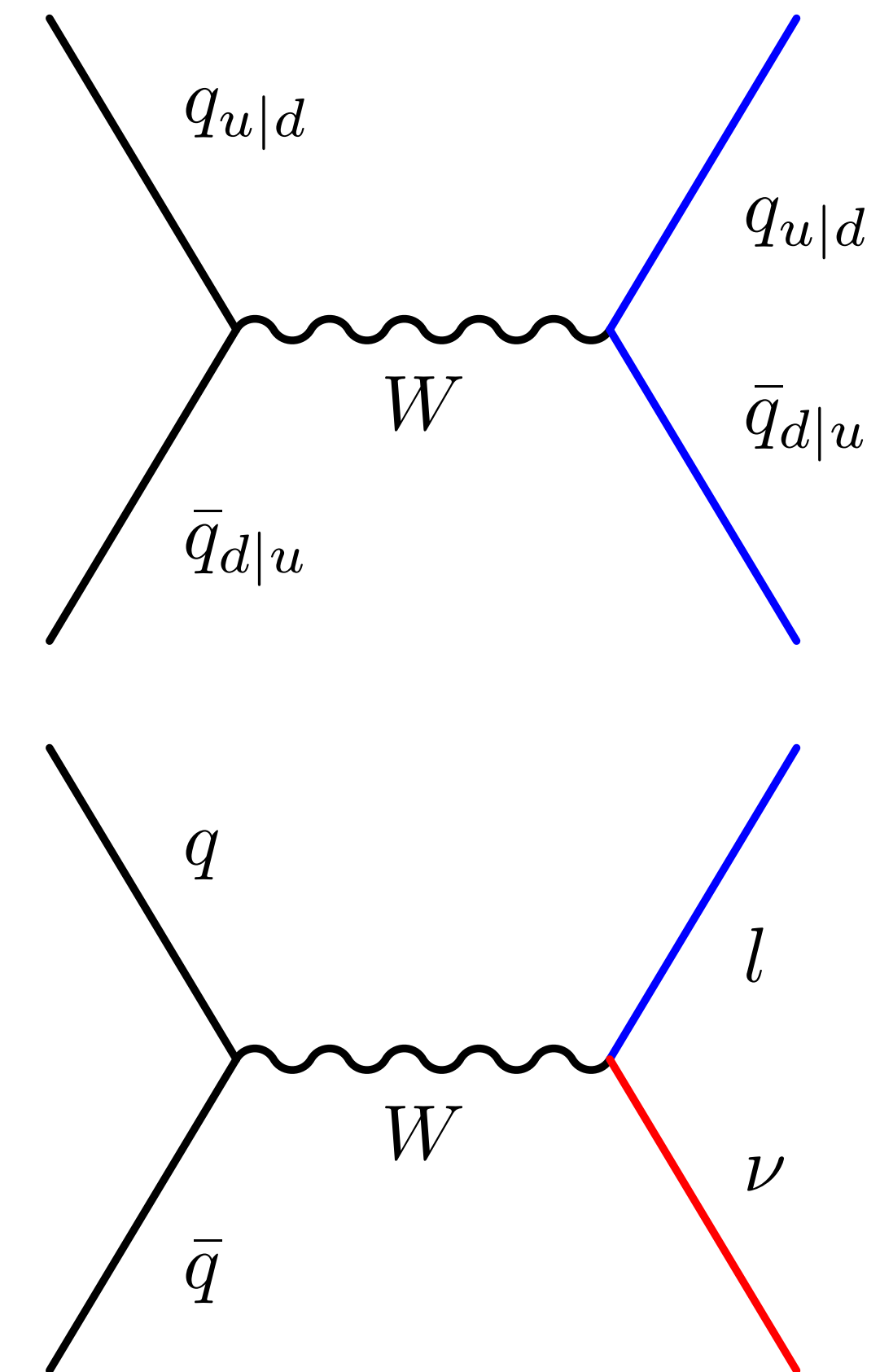
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  - fully reconstructible, but...
  - jet energy scale uncertainty plagues reconstruction of invariant mass



# How is W-mass measured?

Two types of W decay:

- hadronic: 2 jet final state
  - fully reconstructible, but...
  - jet energy scale uncertainty plagues reconstruction of invariant mass
- leptonic:  $\ell + mET$  final state
  - cleaner channel than hadronic decay, but...
  - $\nu$  is invisible  $\rightarrow$  not fully reconstructible
  - good hideout for new physics



Blue: SM visible  
Red: SM invisible

# How is W-mass measured?

- ATLAS and CDF both use two kinematic observables to fit  $m_W$ :

- lepton transverse momentum  $p_T^\ell$ 
  - clean measurement, but affected by nonzero  $p_T^W$

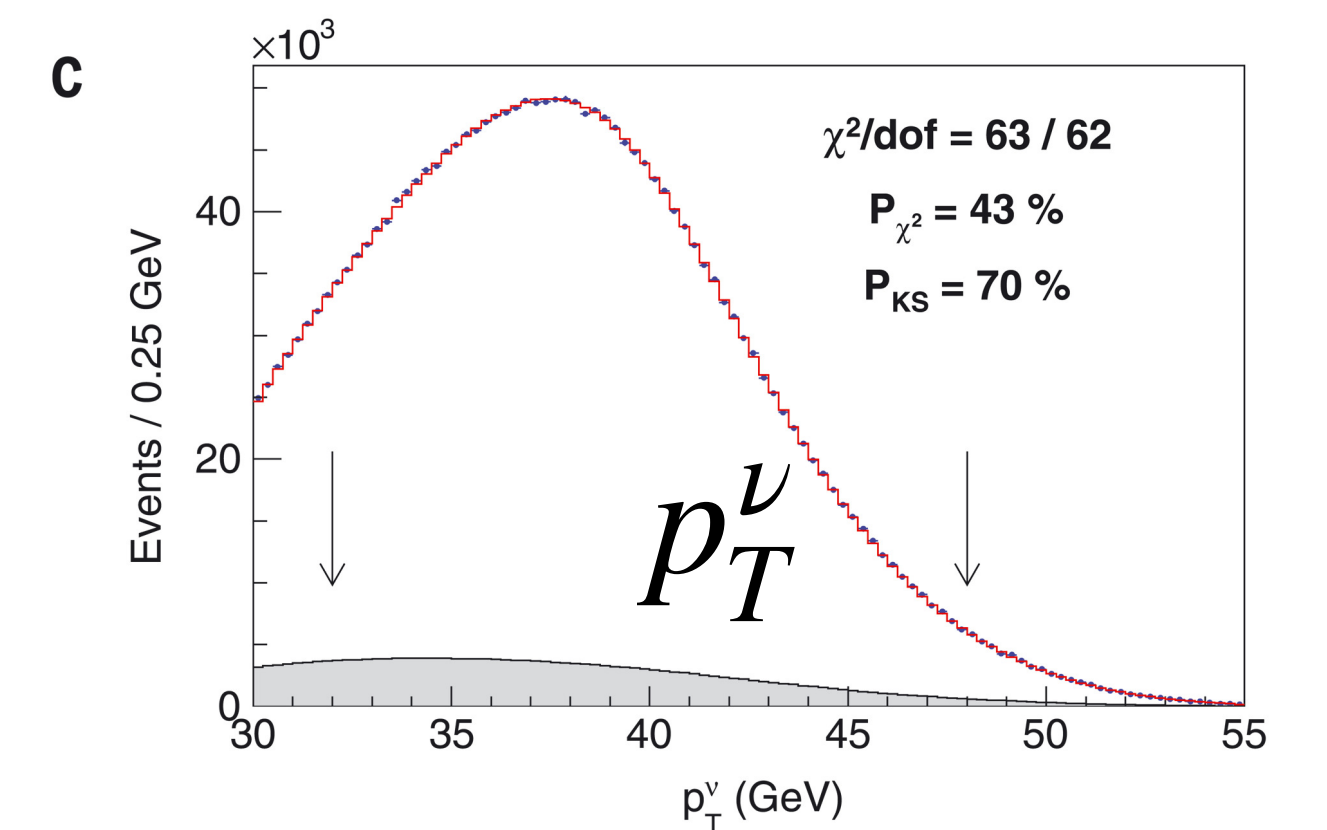
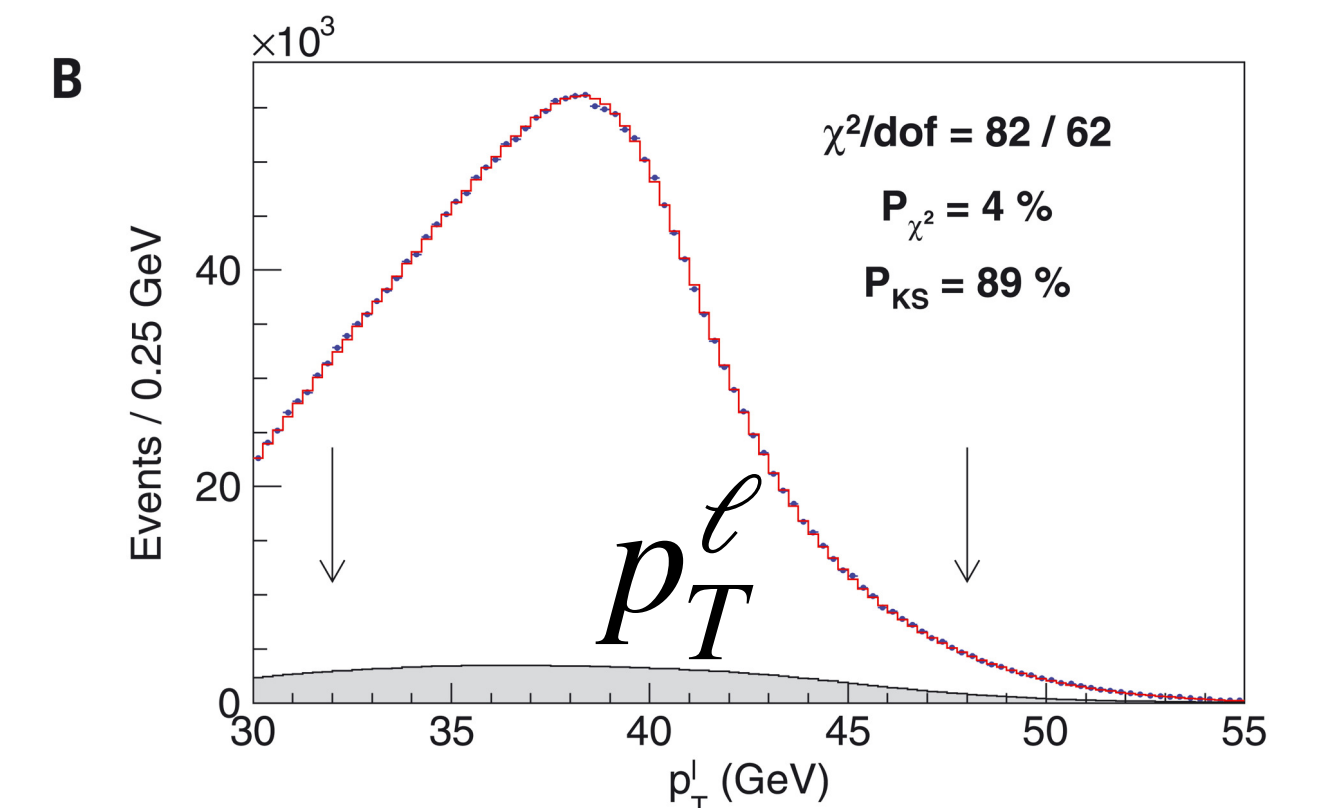
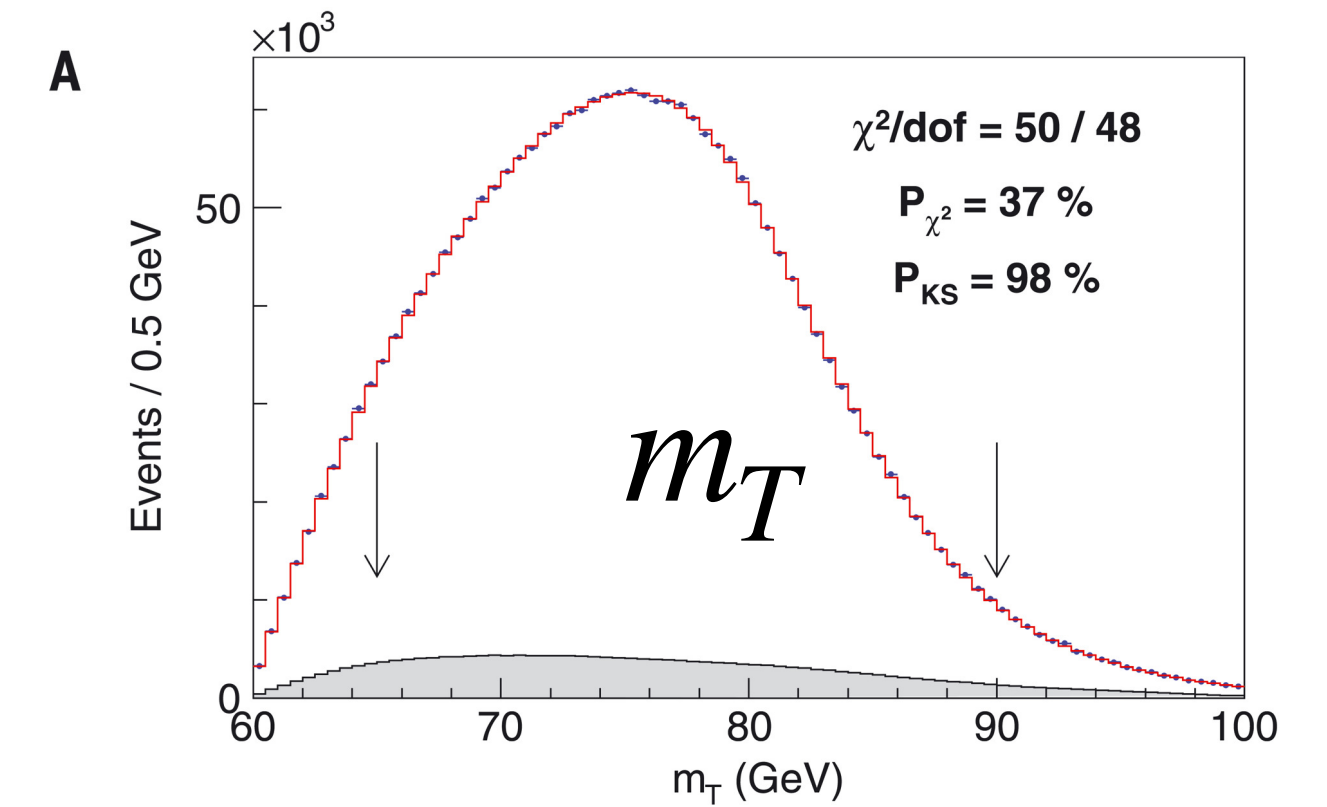
- transverse mass

$$m_T, m_T^2 = (E_T^\ell + E_T^{\text{miss}})^2 - (\vec{p}_T^\ell + \vec{p}_T^{\text{miss}})^2$$

- Provides complementary systematics to  $p_T^\ell$

- CDF also includes  $p_T^\nu$

[High-precision measurement of the W boson mass with the CDF II detector | Science](#)



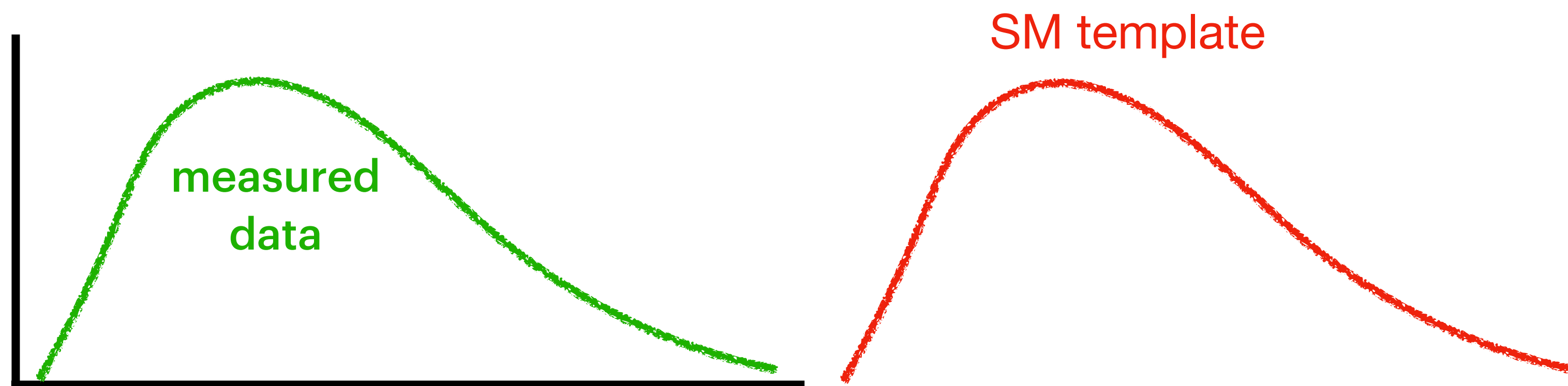
# Probing New Physics

- How can a mass measurement alone probe NP?
  - new physics signals are an irreducible background to the measurement
  - measured kinematic distributions are in perfect agreement with SM predictions with great precision
  - exclude NP signals producing different kinematic distributions

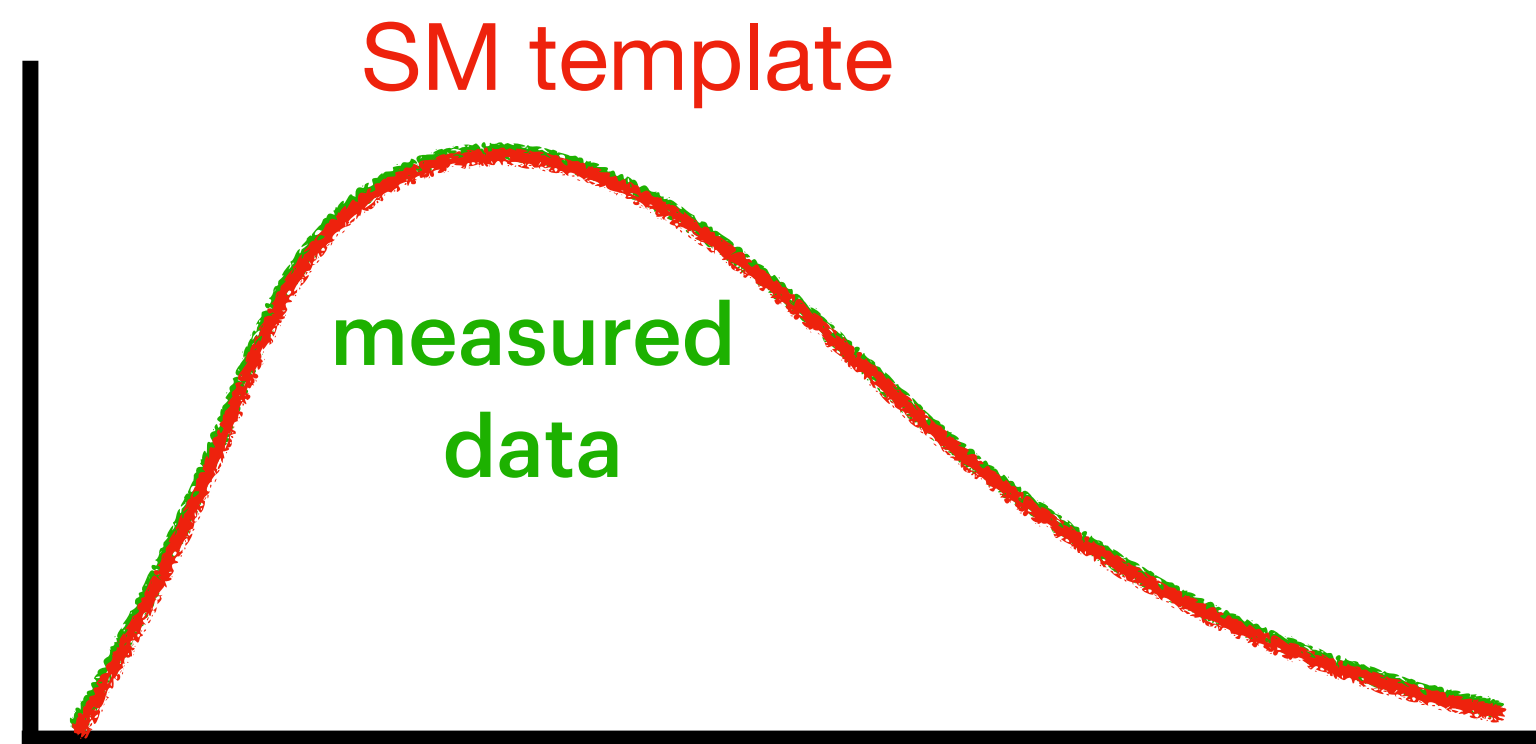


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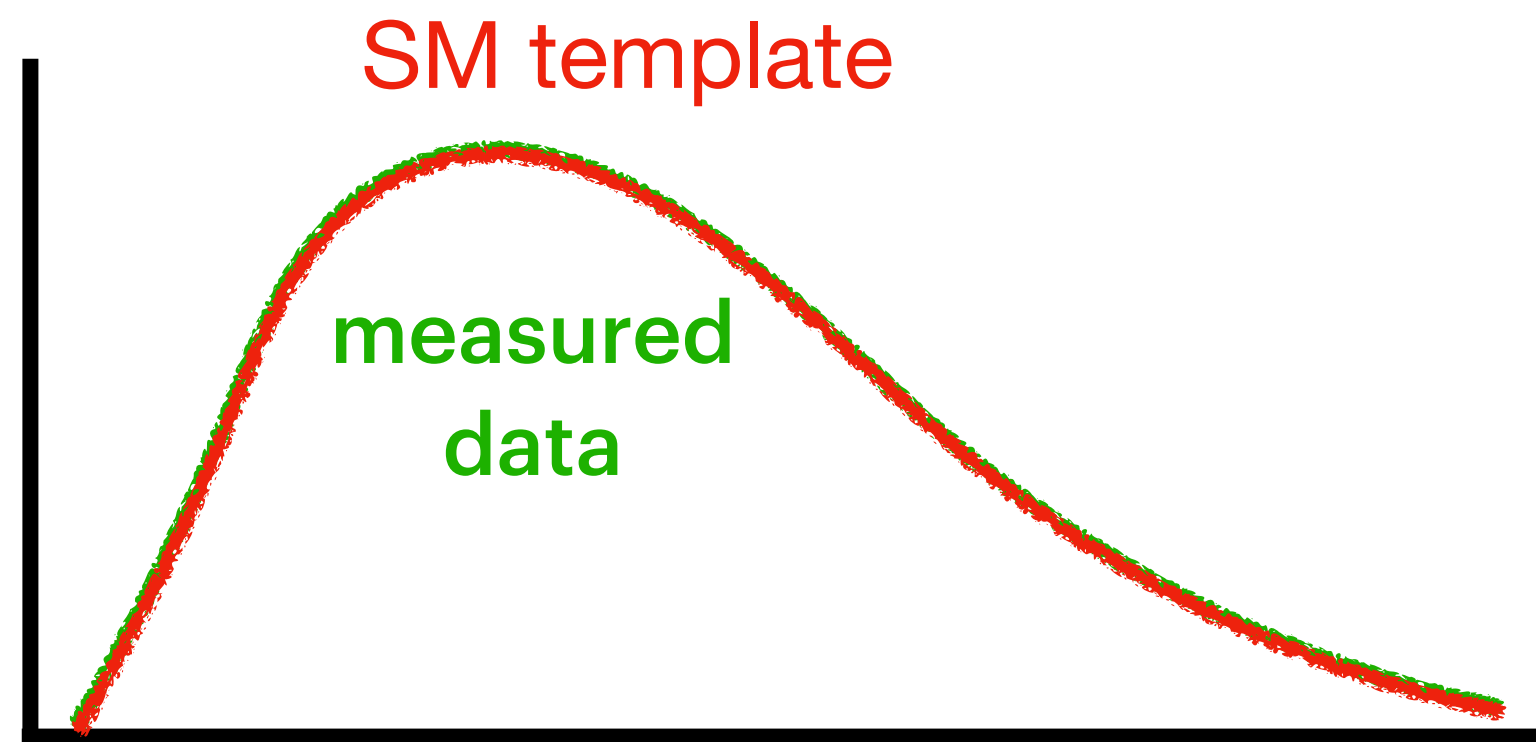


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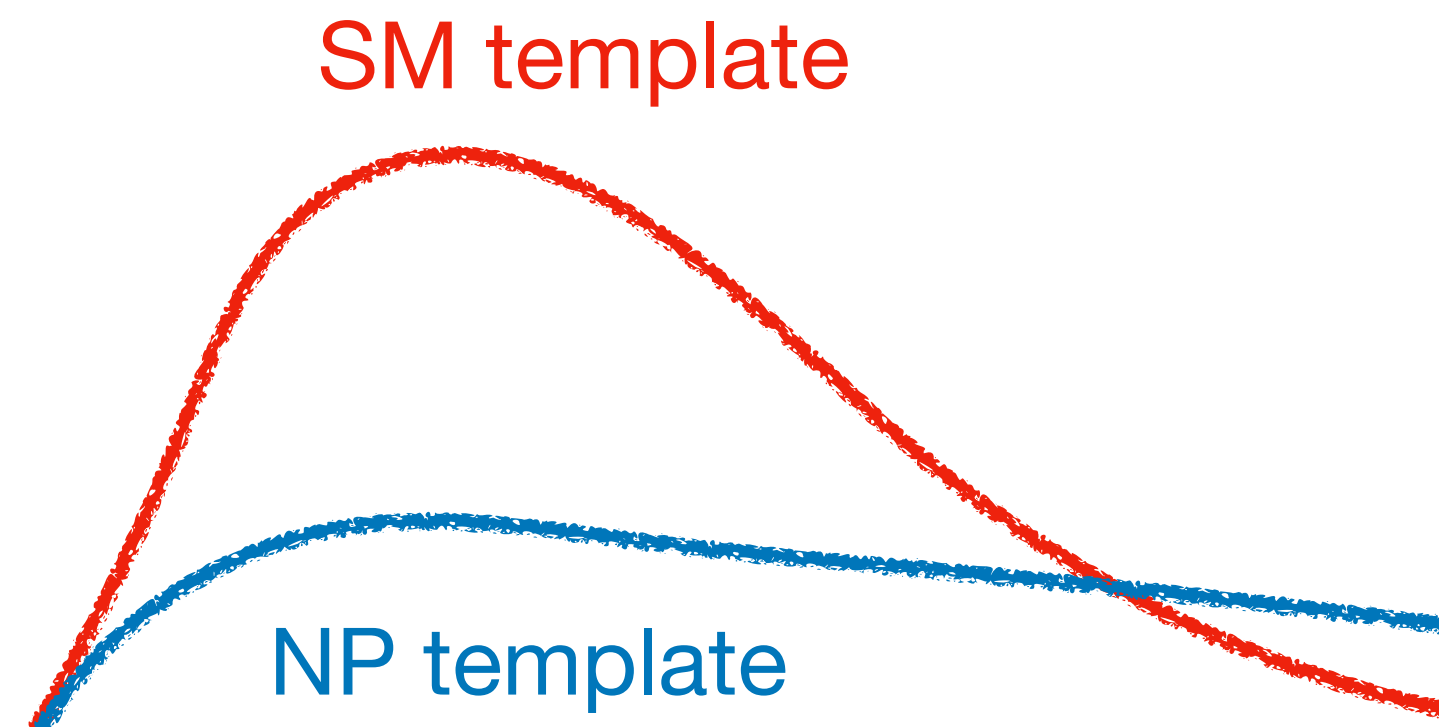
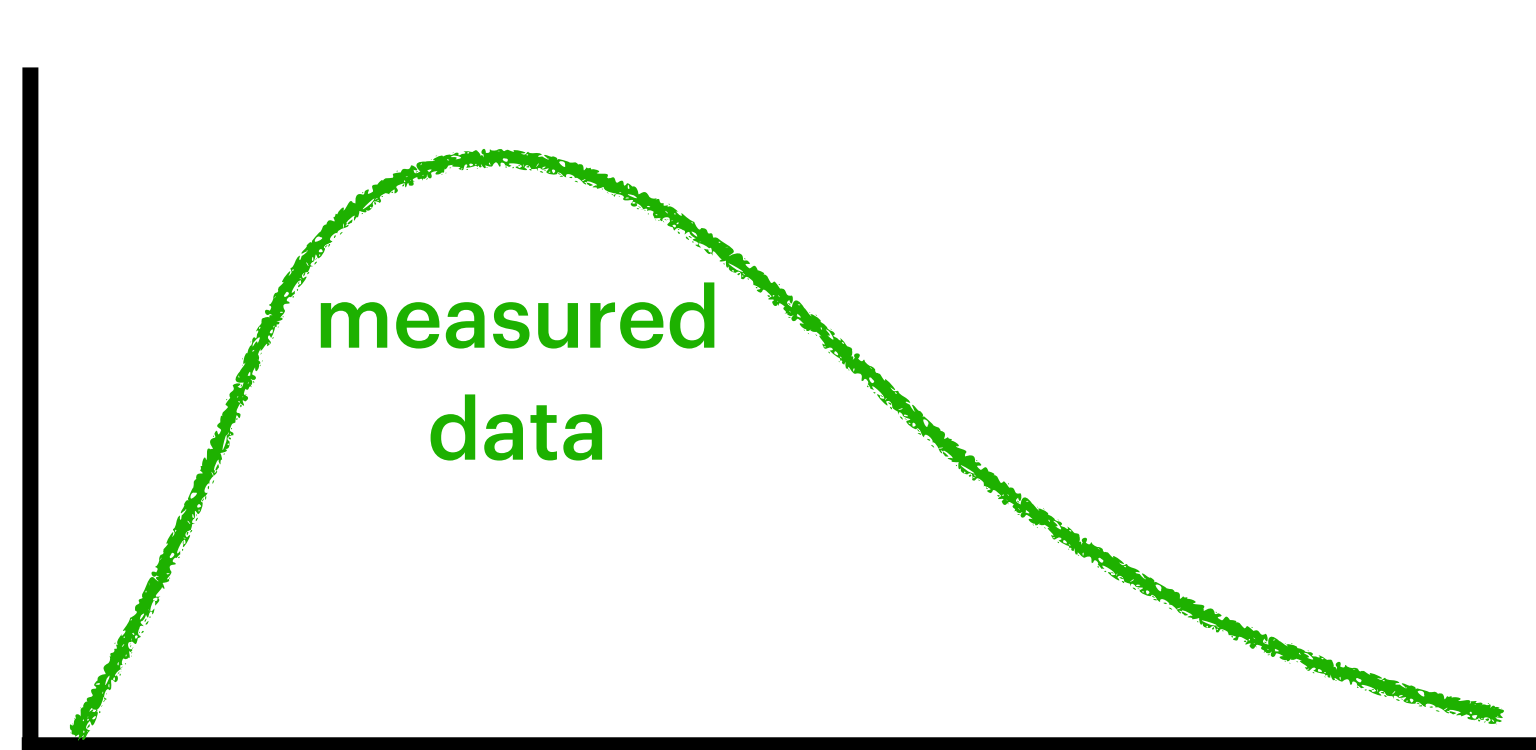


perfect agreement! obtain SM parameter from fit

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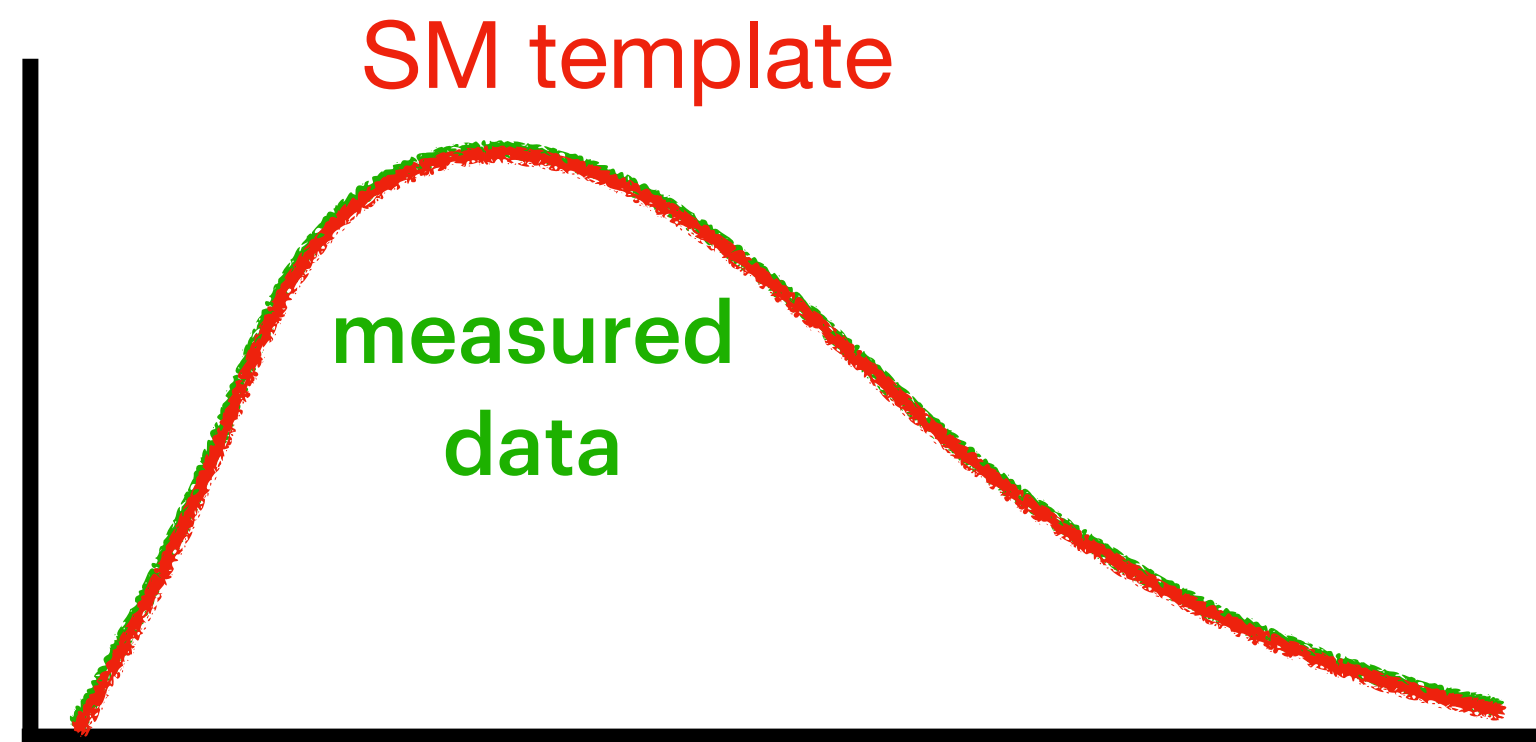


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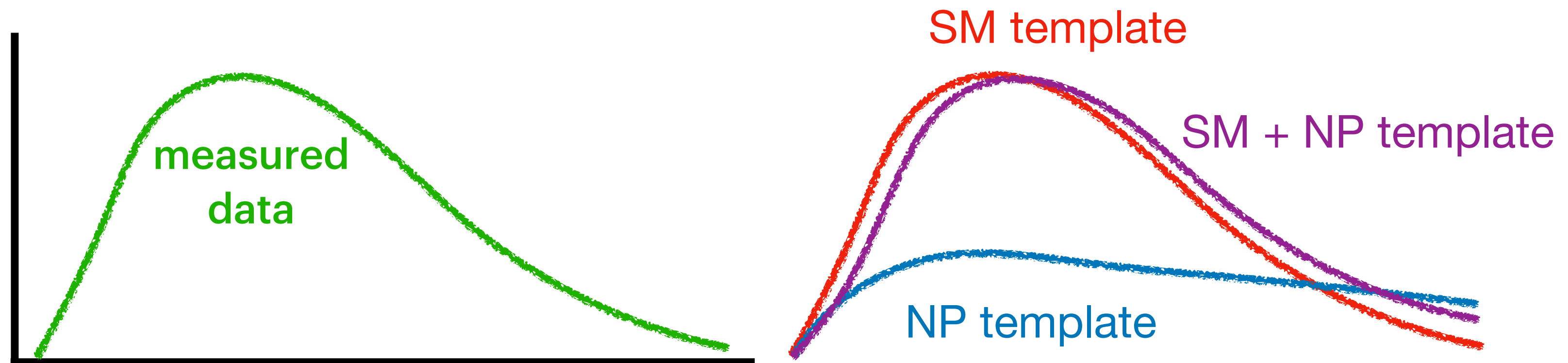




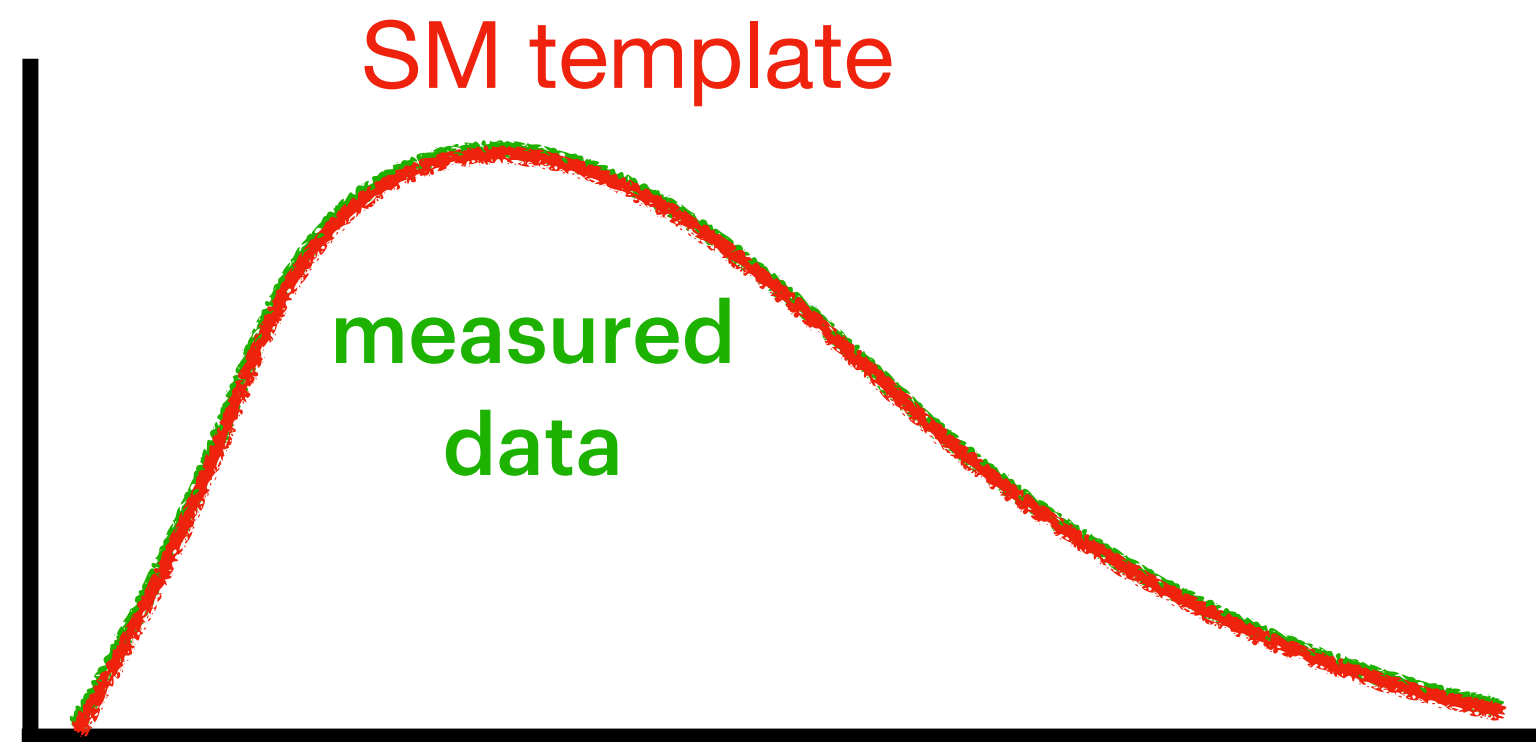
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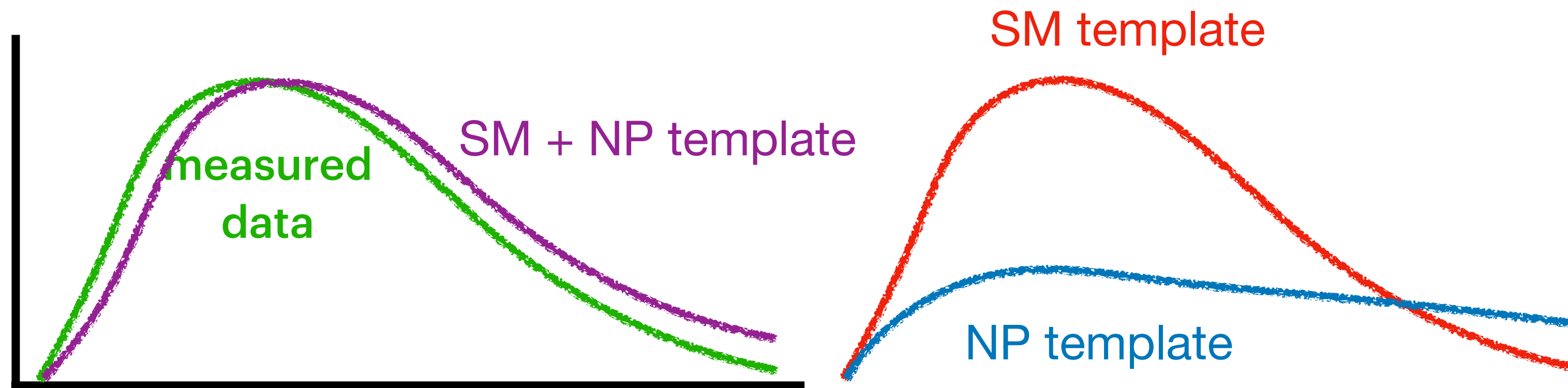
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# Cartoon picture



perfect agreement! obtain SM parameter from fit



bad fit! NP can be excluded

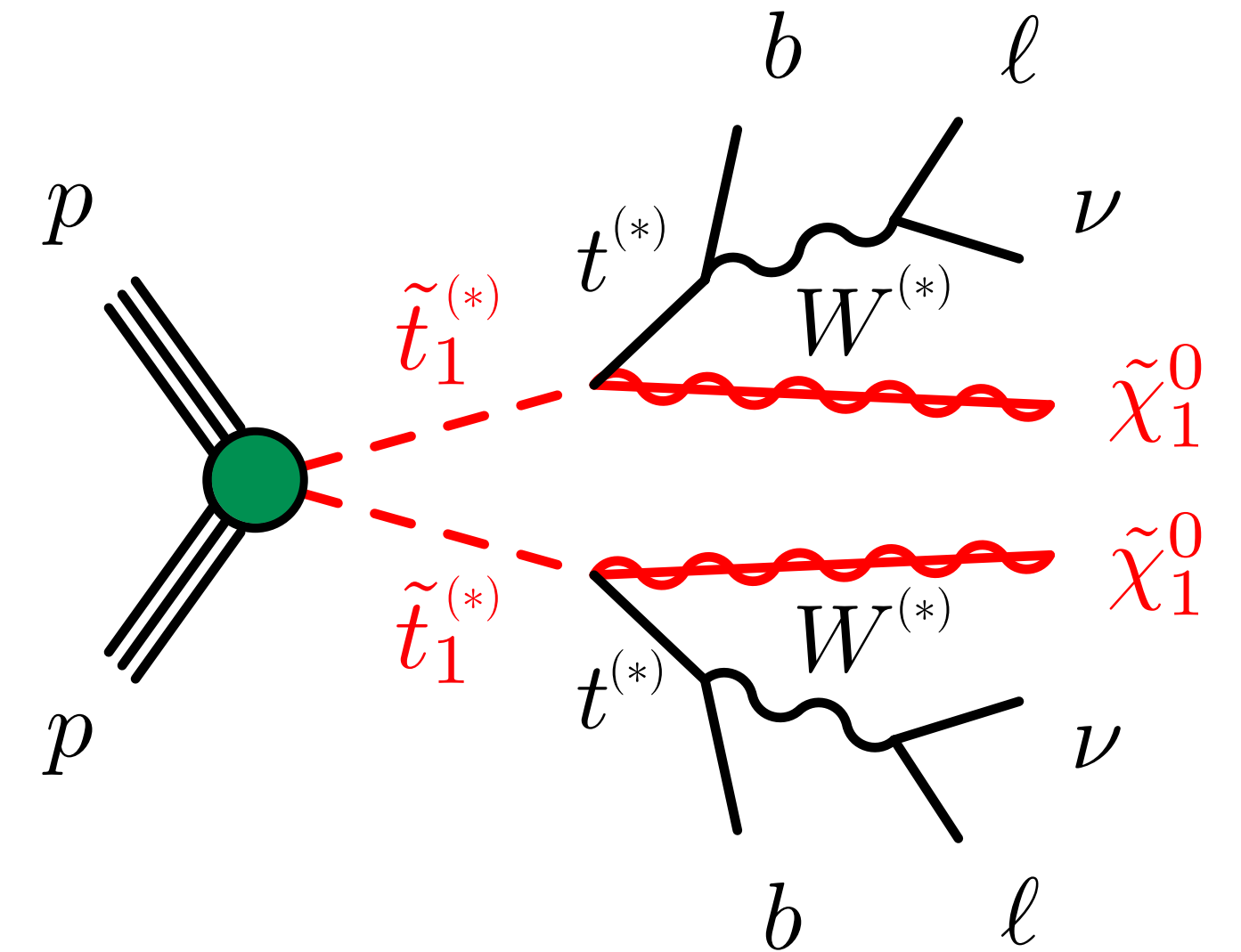
# Examples of Probing NP

- Idea is general to any precision measurement
  - searching for SUSY stop production using top quark mass

$$\Delta m_t / m_t \sim 10^{-2} - 10^{-3}$$

$$\Delta m_W / m_W \sim \mathcal{O}(10^{-4})$$

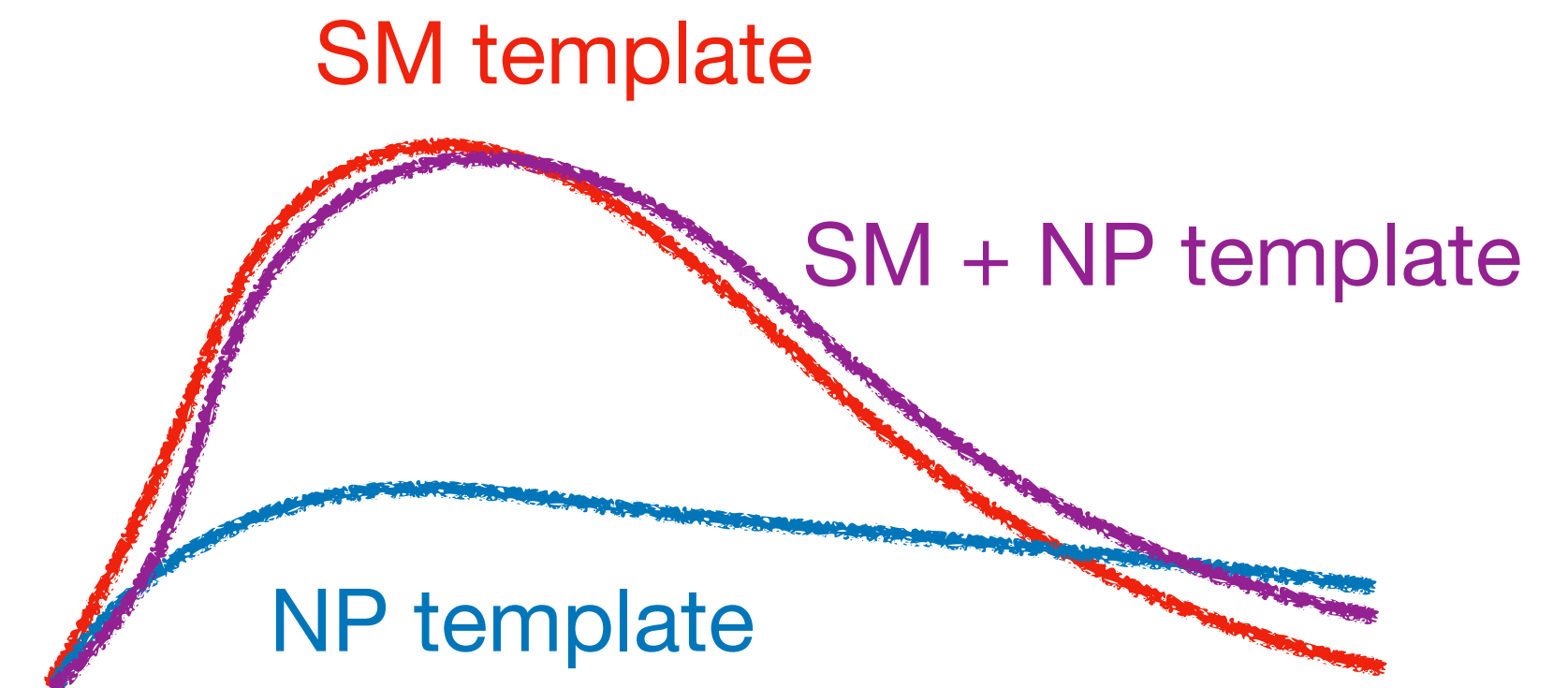
- not yet implemented for  $\ell + mET$  final state
  - higher precision allows probing of fainter NP signals
  - sensitive to NP charged only under EW interactions



[Cohen, Majewski, Ostdiek, Zheng \[1909.09670\]](#)  
[Eifert, Nachman \[1410.7025\]](#)  
[Franceschini \[1601.02684\]](#)  
[Agashe, Airen, Franceschini, Incandela, Kim, DS \[2212.03929\]](#)

# New Purpose for Precision W-mass Data

- Use these precise measurements to probe NP signals that contaminate W-mass data
- Fit templates of kinematic observables to data
  - templates now should include NP + SM W compared to what ATLAS and CDF do
  - Sensitive to regions of parameter space where NP distributions differ noticeably from SM templates



# New Purpose for Precision W-mass Data

- Use these precise measurements to probe NP signals that contaminate W-mass data
- Since the measurement is probing NP parameters *and*  $m_W$ , must perform a global fit to both
  - NP contamination could add shift in  $m_W$  and add additional uncertainty in the measurement
- independent of electroweak fit
- What kinds of new physics can be probed?

# Types of NP Contamination

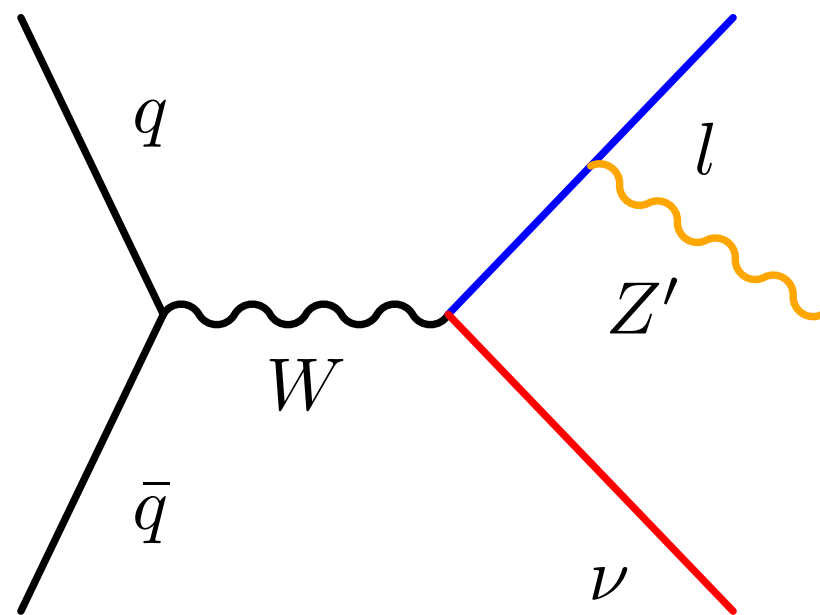
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Three categories to populate SM W region, requiring final state of  $\ell + mET$ :

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- Modify decay of SM W



Blue: SM visible

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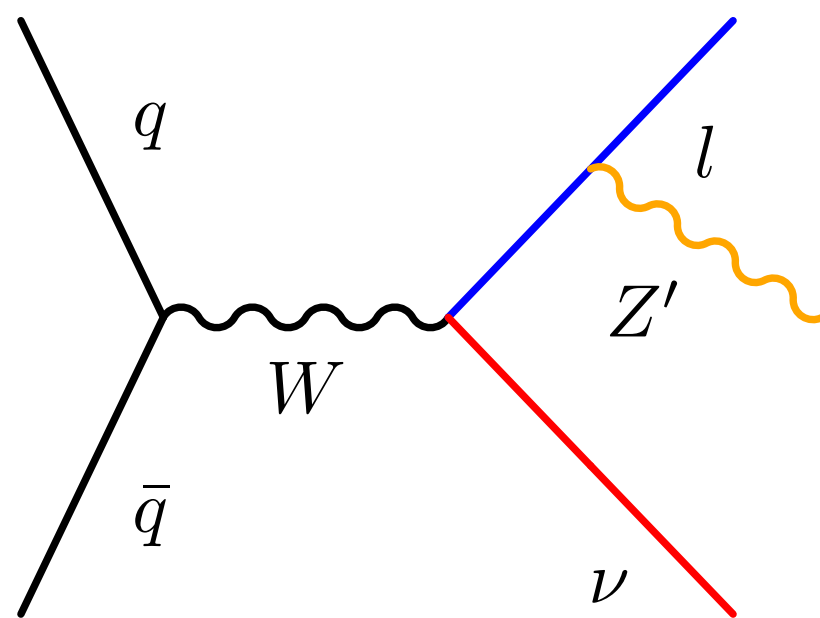
Orange: BSM invisible



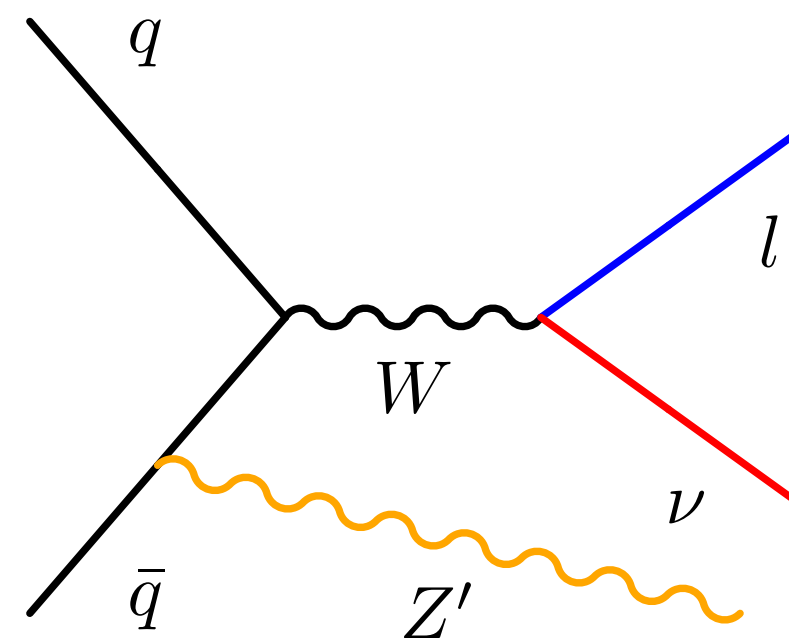
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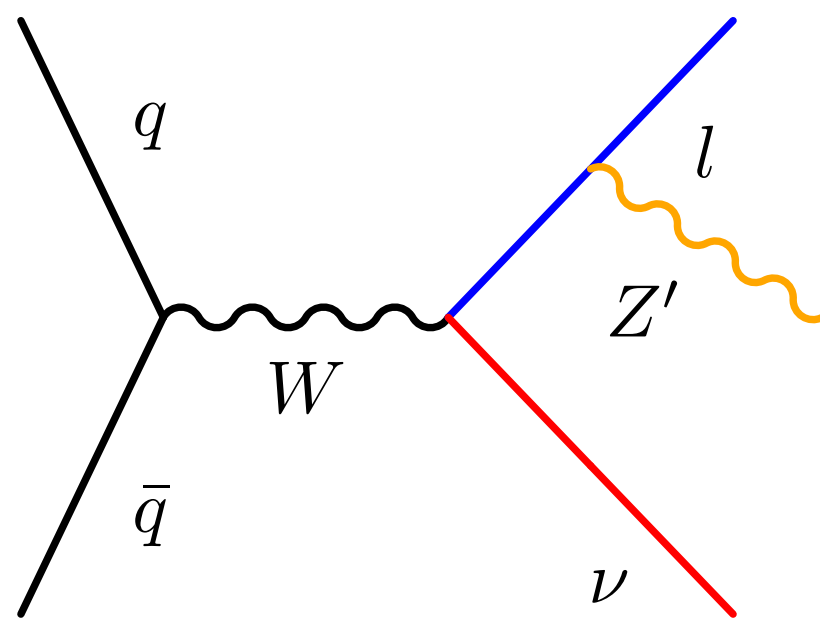
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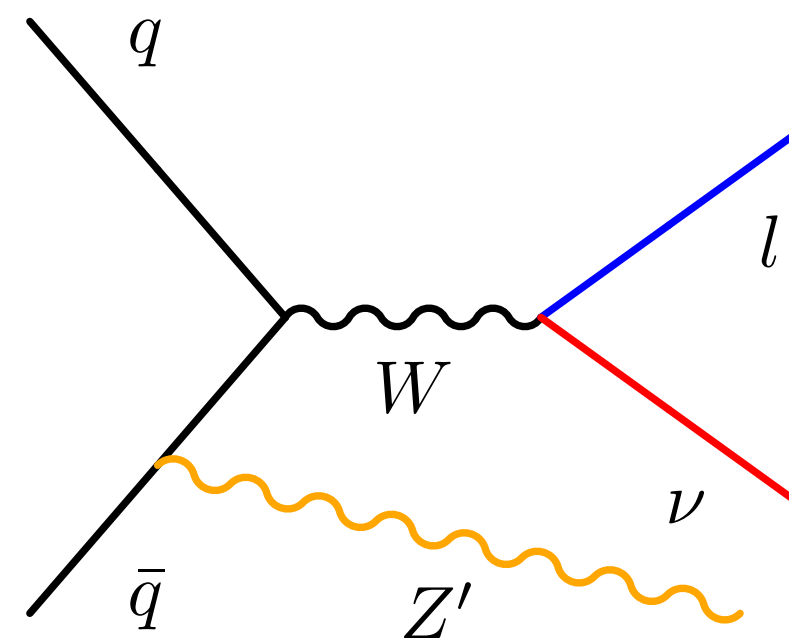
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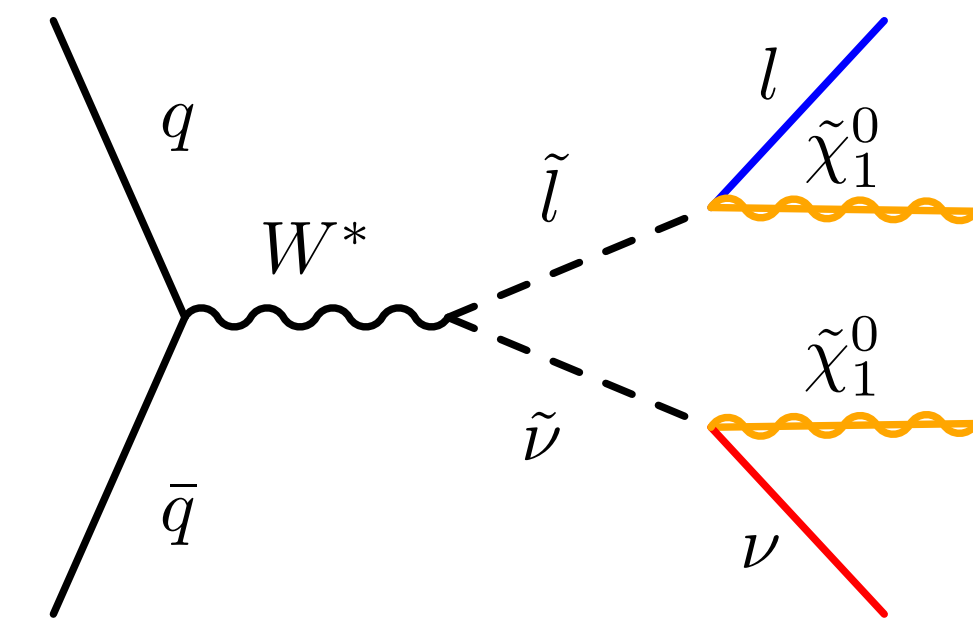
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- Without on-shell W



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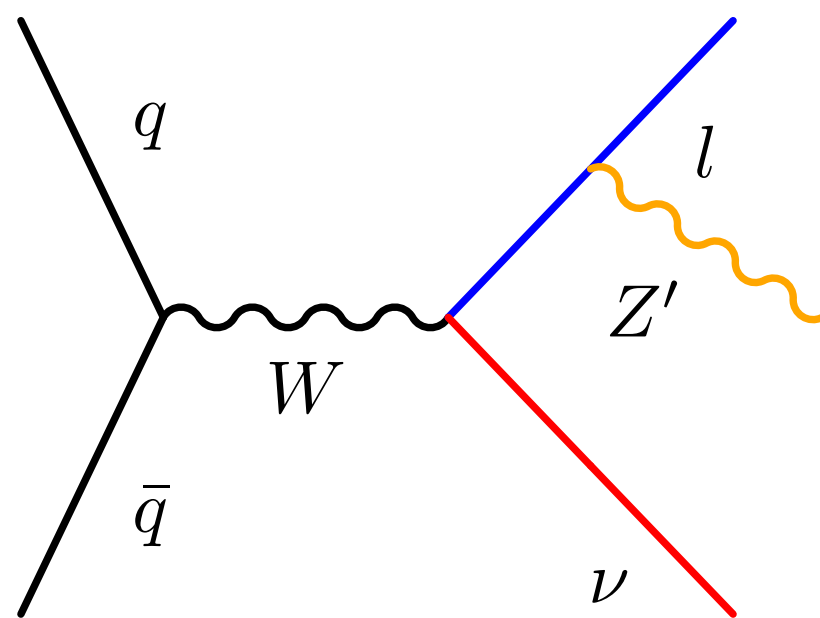
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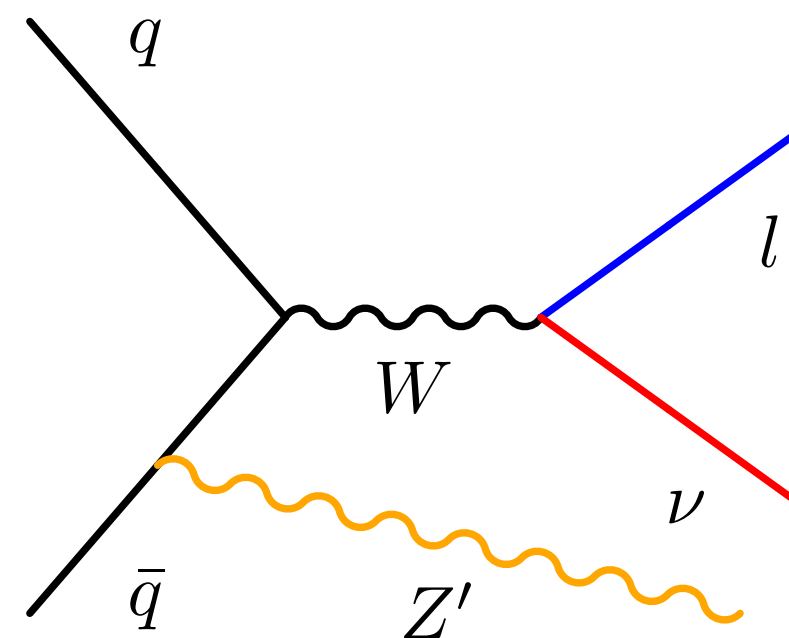
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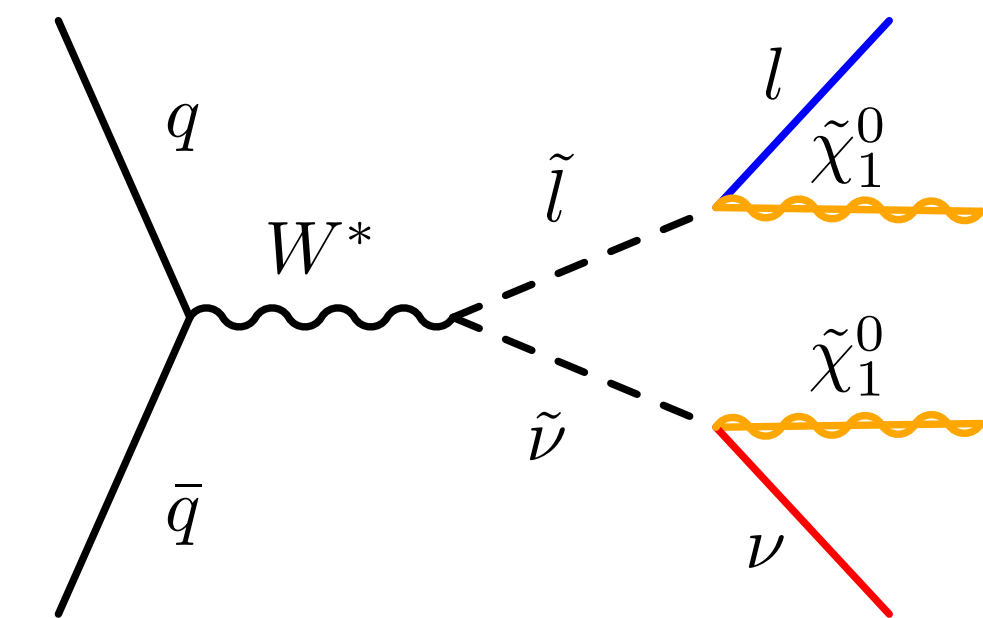


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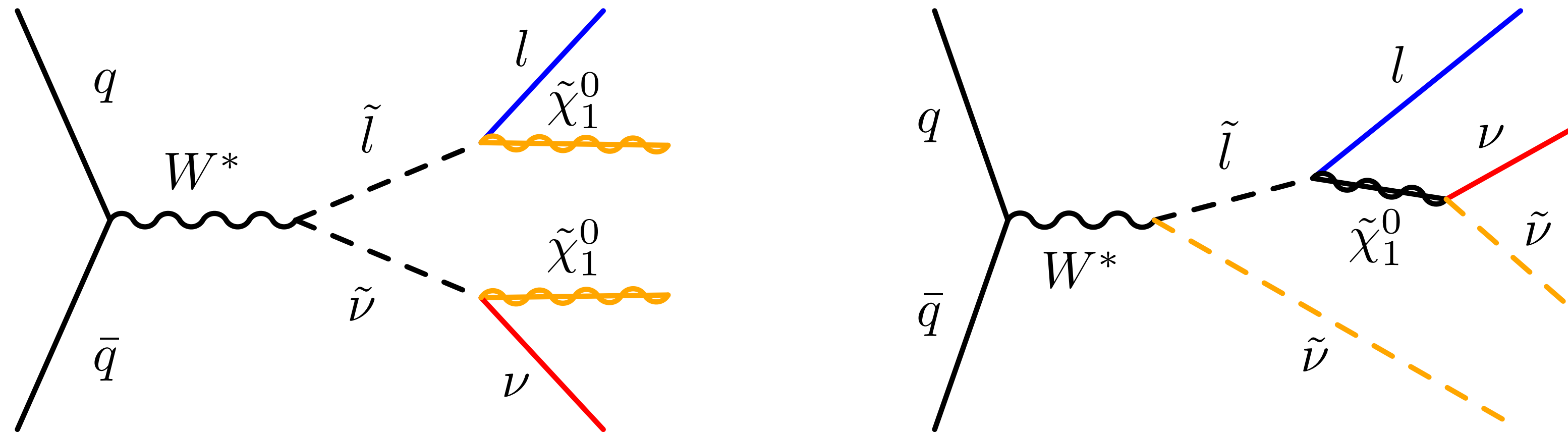
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Focus of today's talk

# NP without on-shell W

## MSSM Slepton-Sneutrino Production



# $\ell + mET$ in MSSM

NP without on-shell W

- Supersymmetry is well-known and well-studied
- Simple production mechanism for  $\ell + mET$ :

$$pp \rightarrow \tilde{\ell} (\rightarrow l \tilde{\chi}_1^0) \tilde{\nu}$$

- Assume all other superpartners are heavy
  - light sleptons are motivated in Sleptonic SUSY

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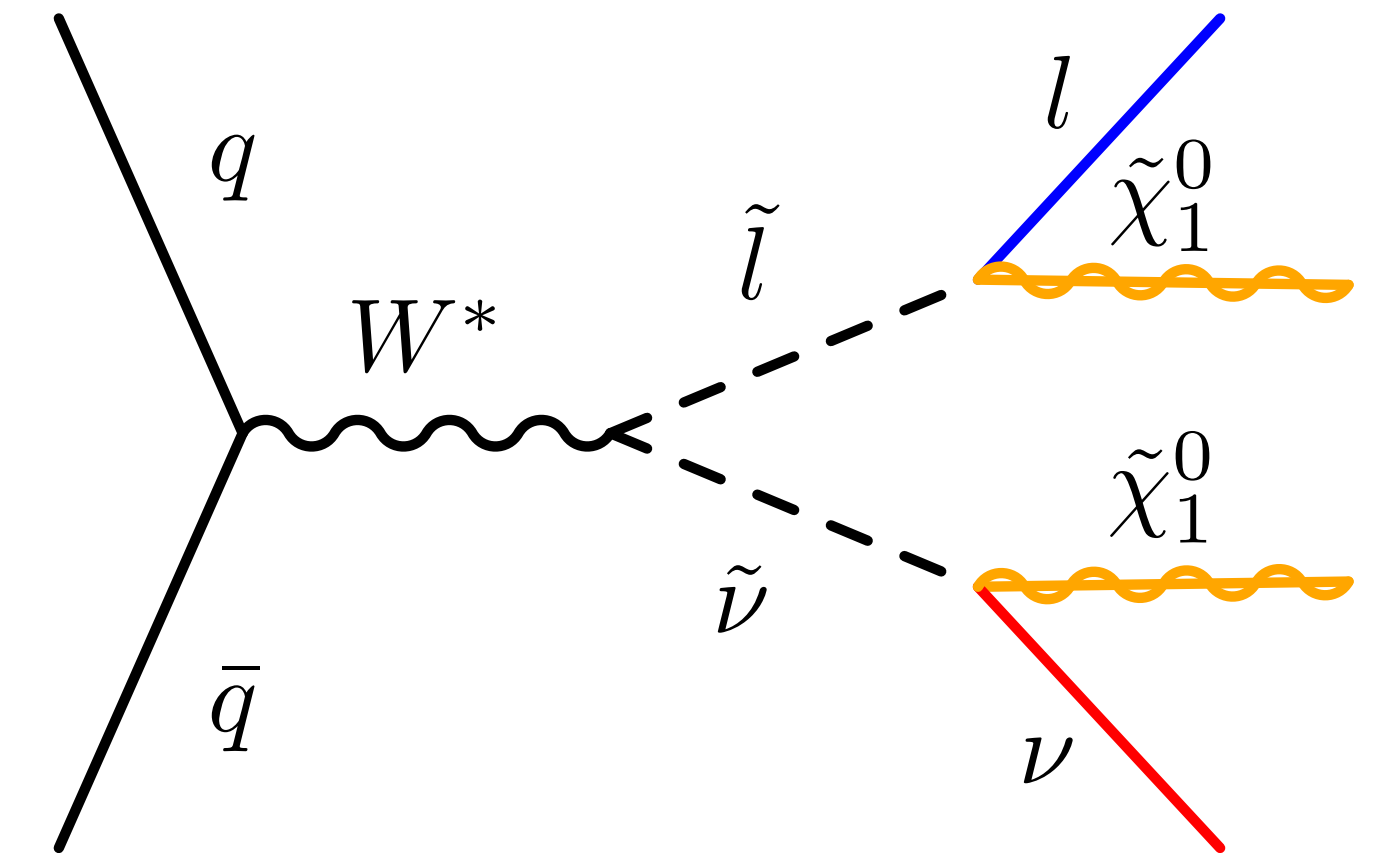
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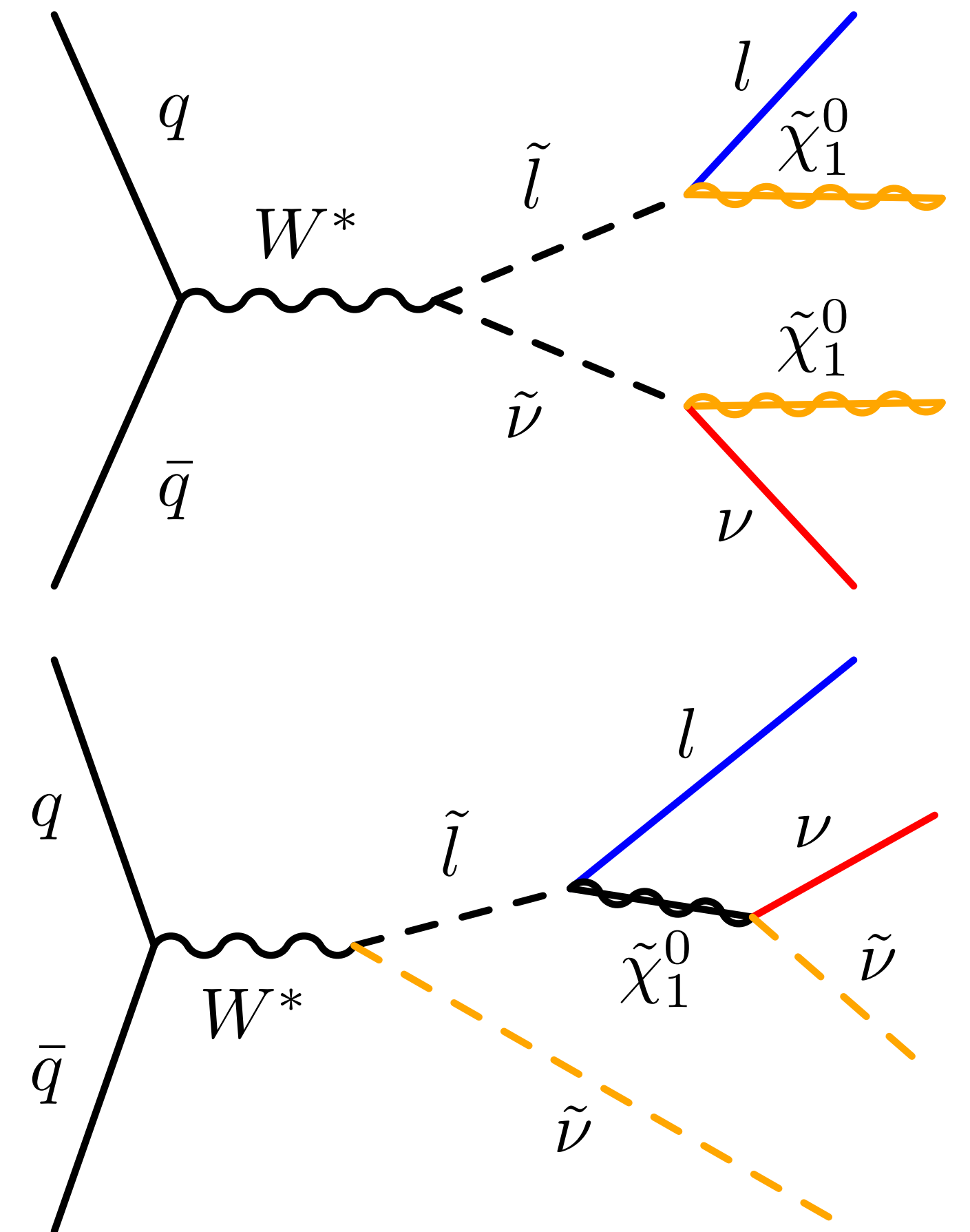
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- We consider two regimes:
  - lightest neutralino is LSP:  $\tilde{\nu} \rightarrow \nu \tilde{\chi}_1^0$
  - sneutrino is LSP:  $\tilde{\chi}_1^0 \rightarrow \tilde{\nu} \nu$



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# MSSM Masses

NP without on-shell W

- slepton-sneutrino process contains 3 NP masses:  $m_{\tilde{l}}, m_{\tilde{\nu}}, m_{\tilde{\chi}_1^0}$

- slepton-sneutrino mass splitting relation fixed by D-term:

$$m_{\tilde{\nu}_\ell}^2 = m_{\tilde{\ell}}^2 + \cos(2\beta) m_W^2$$

- setting  $\tan \beta \rightarrow \infty$  ensures lightest sneutrinos

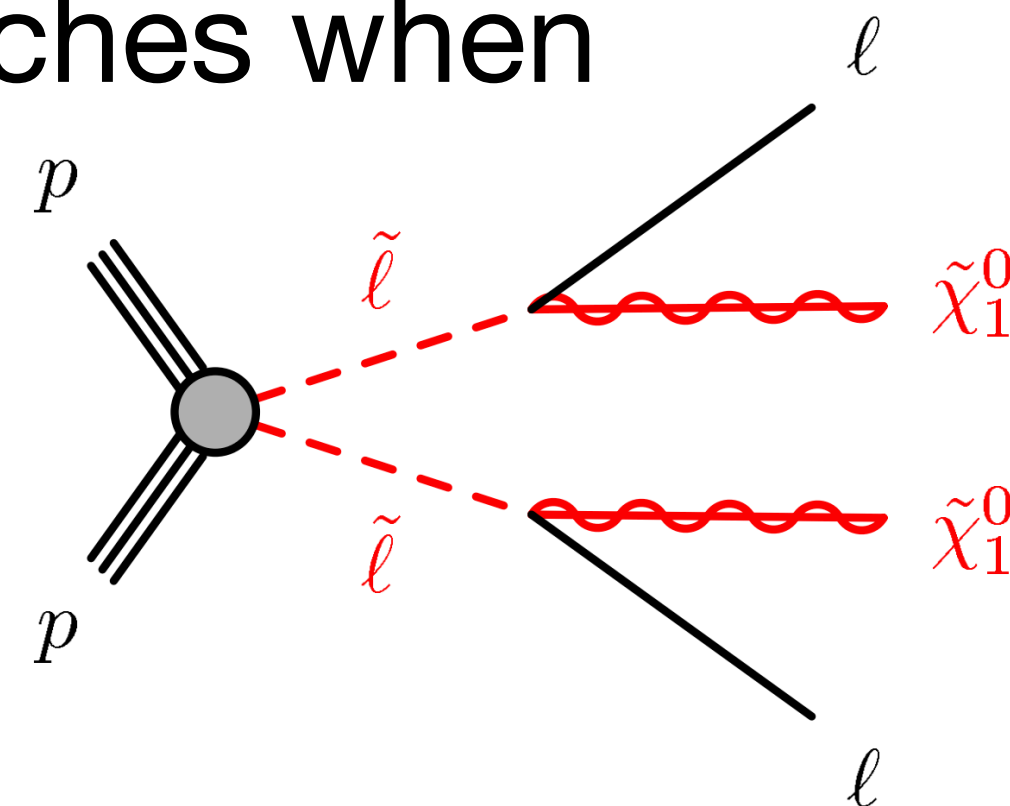
$$m_{\tilde{\nu}_\ell}^2 = m_{\tilde{\ell}}^2 - m_W^2$$

# MSSM Masses

NP without on-shell W

- LEP rules out  $m_{\tilde{l}} < 100$  GeV
- Heavier sleptons have negligible cross sections at Tevatron
- LHC searches for di-slepton production using ML techniques
- Gap in the searches when

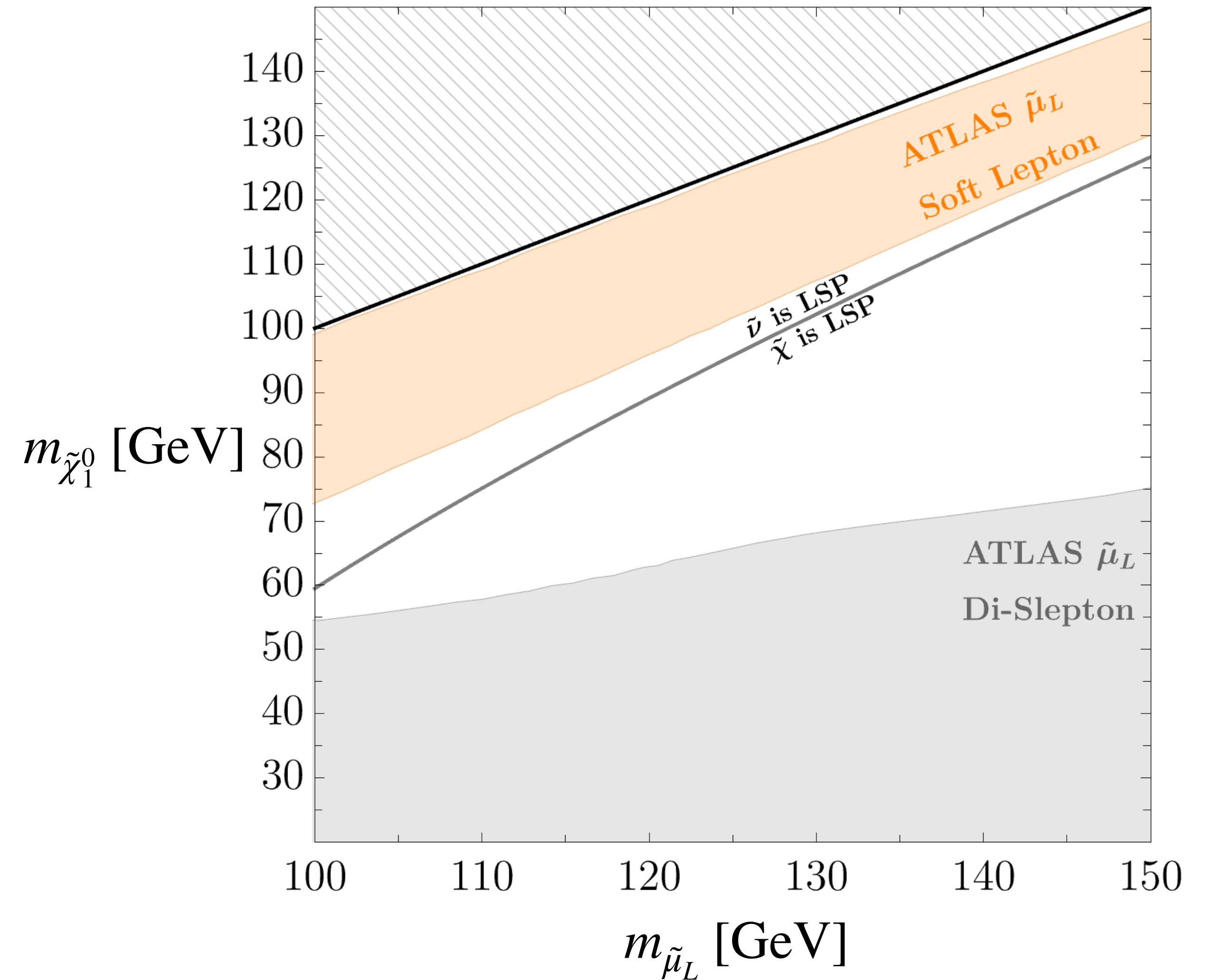
$$m_{\tilde{l}} \sim m_{\tilde{\chi}_1^0}$$



[LEP2 SUSY WORKING GROUP Sleptons \(cern.ch\)](https://cds.cern.ch/record/2209139/files/ATLAS-CONF-2019-027)

[\[2209.13935\] ATLAS dislepton search \(arxiv.org\)](https://arxiv.org/abs/2209.13935)

[CMS-SUS-21-008](#)



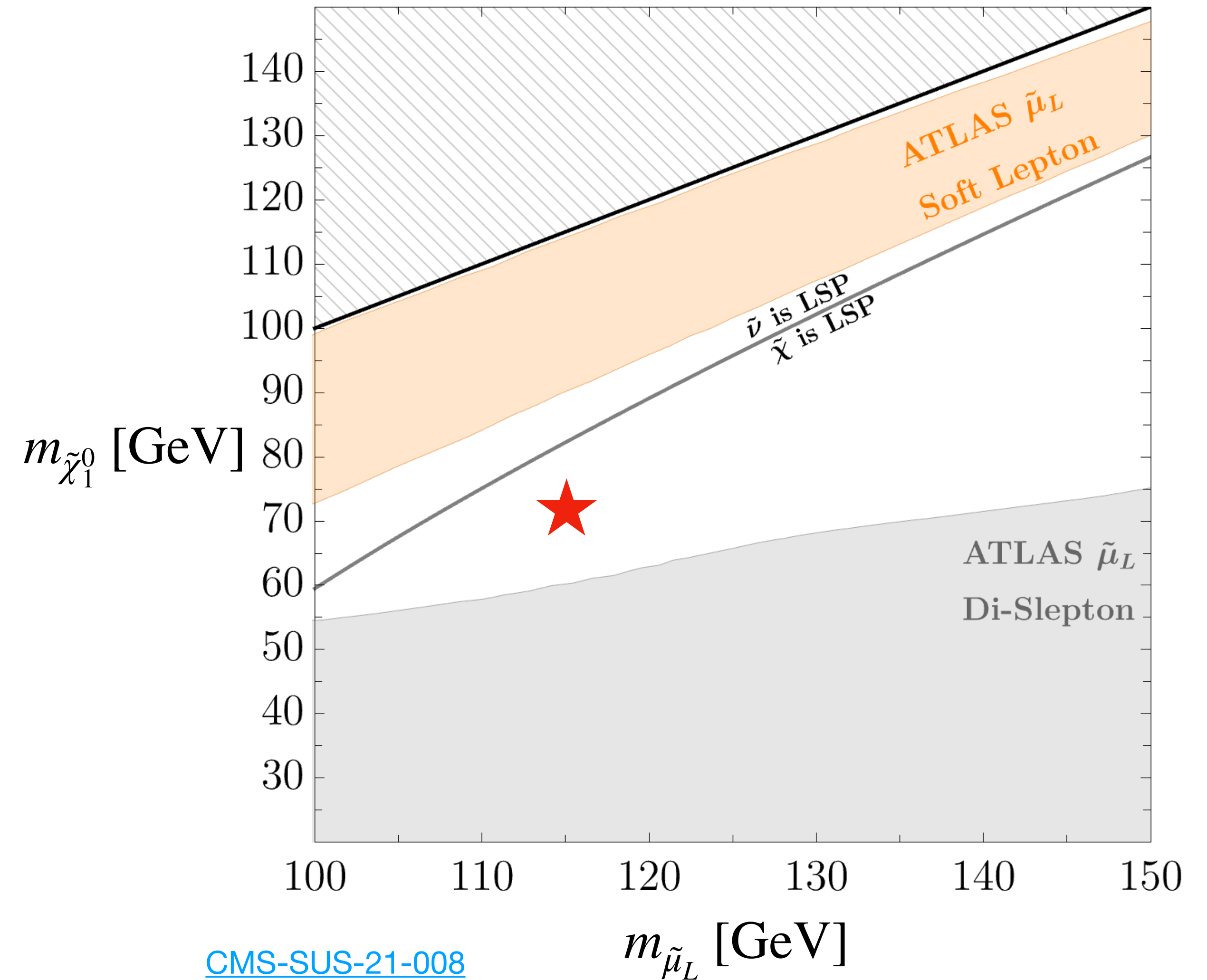
# MSSM Masses

NP without on-shell W

- Proposals to use precision  $WW$  data to disentangle dilepton events
- gap still exists
- 1D fit: assumes  $m_W$  from independent measurements, whereas we float  $m_W$

[Curtin, Jaiswal, Meade \[1206.6888\]](#)

[Curtin, Jaiswal, Meade, Tien \[1304.7011\]](#)



[CMS-SUS-21-008](#)

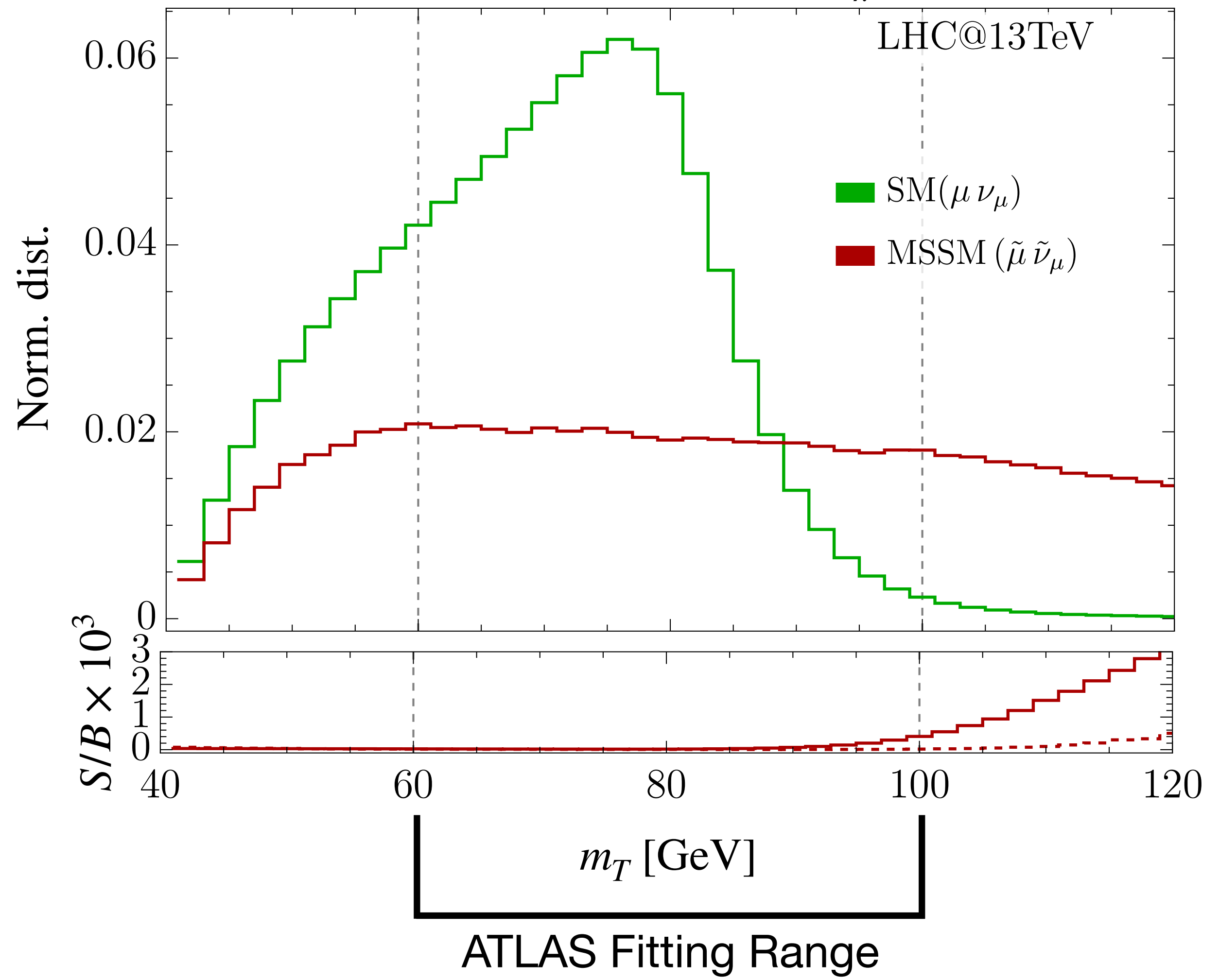
[\[2209.13935\] ATLAS dilepton search \(arxiv.org\)](#)

# Slepton-Sneutrino Kinematic Distributions

NP without on-shell W

$m_{\tilde{l}} = 115 \text{ GeV}, m_{\tilde{\nu}} = 83 \text{ GeV}, m_{\tilde{\chi}} = 70 \text{ GeV}$  ★

- SUSY slepton  $m_T$  (red) vs SM W (green)
- SUSY distributions are flat
- These distributions don't shift  $m_W$ :
  - Marginalize over  $m_W$  obtained from min.  $\chi^2$ , but we find negligible shift
  - for simplicity, fix  $\Delta_{m_W} = 0$ , focus on sensitivity to NP

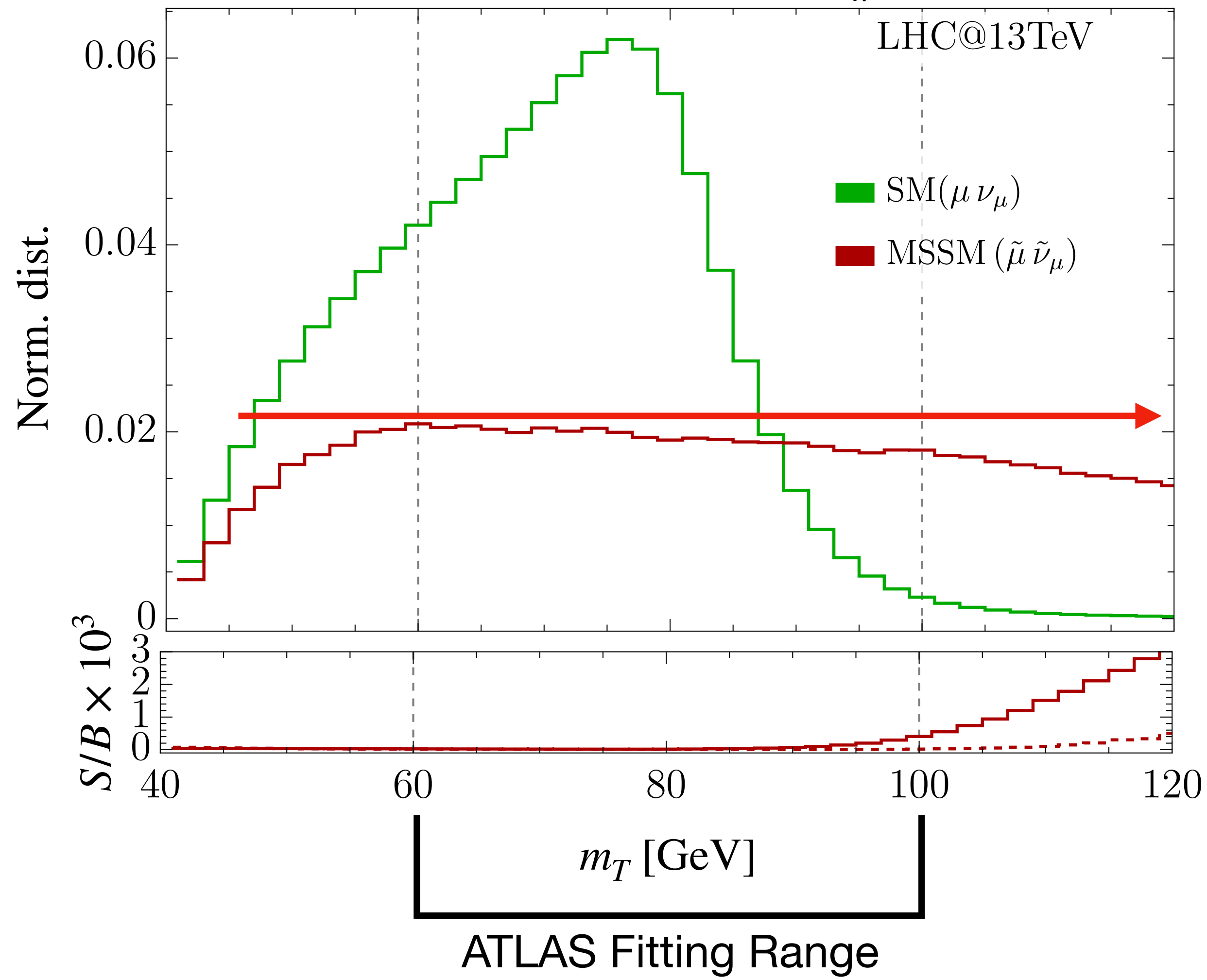


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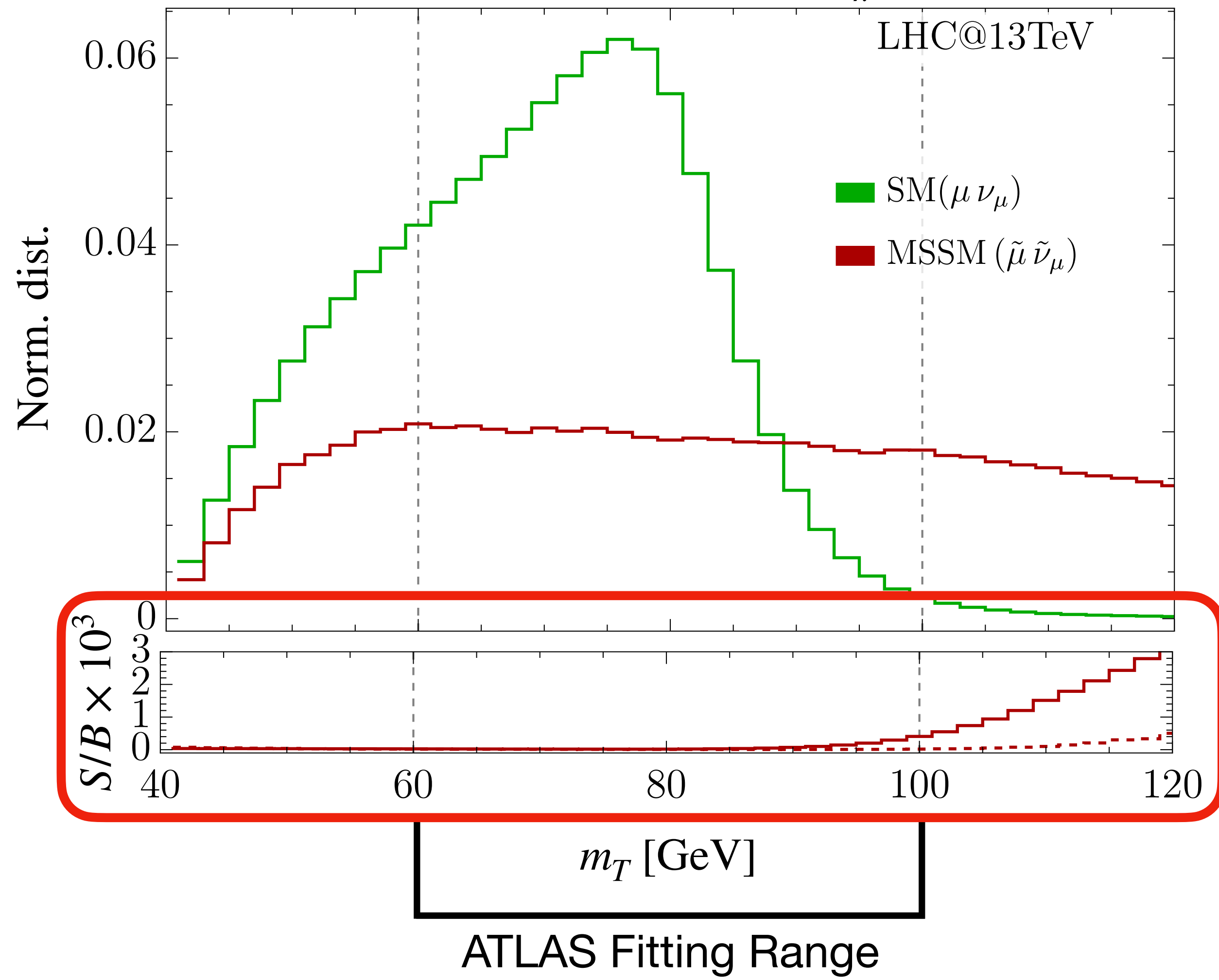


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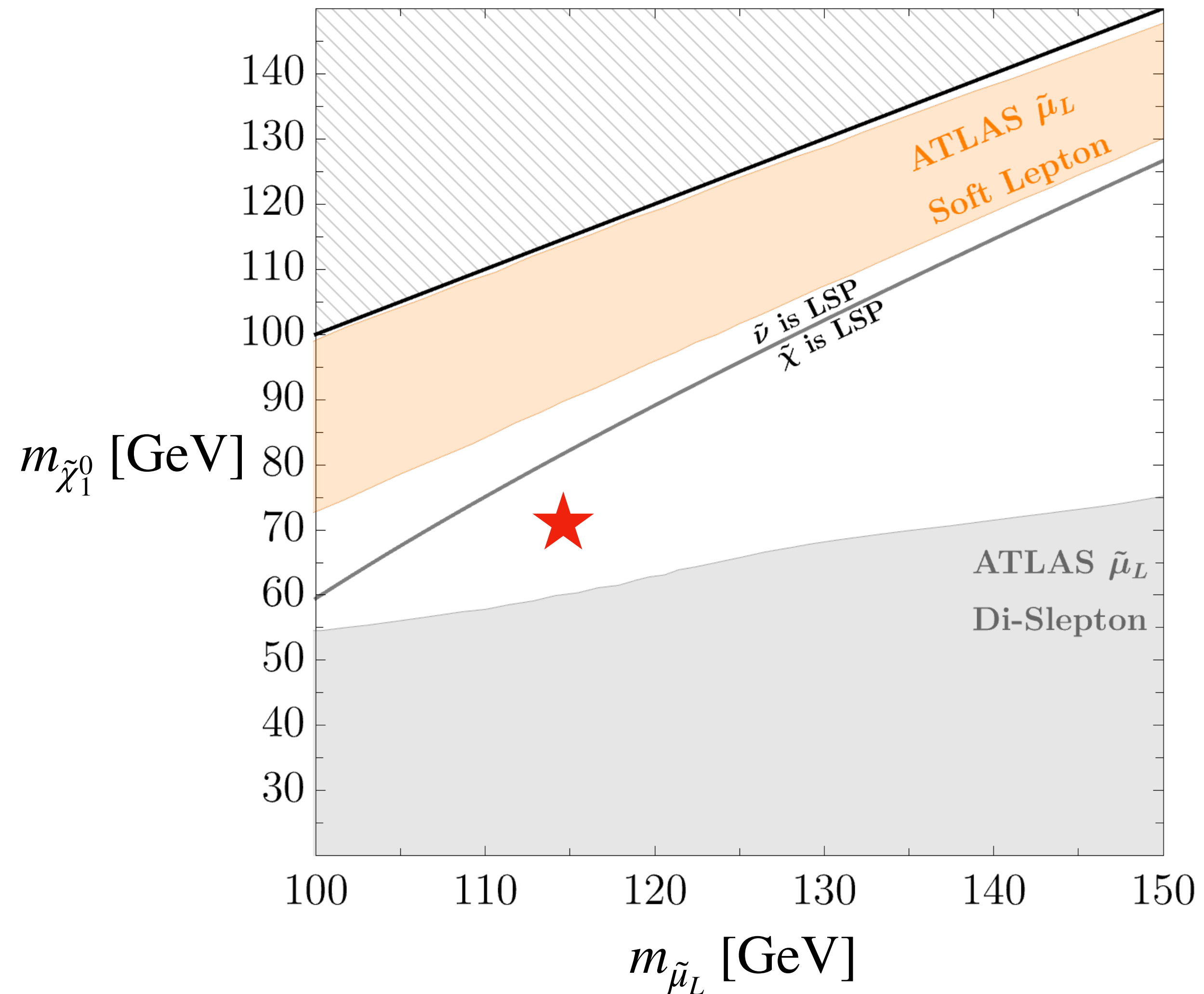
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- Bottom figure:
  - S/B (solid) increases considerably beyond Jacobian peak
    - motivates extending fitting range to increase sensitivity to NP
  - Z contamination (dashed, bottom figure) negligible, even beyond fitting range
  - Important to consider, as Z sample is used in data-driven approach



# SUSY Slepton-Sneutrino Search

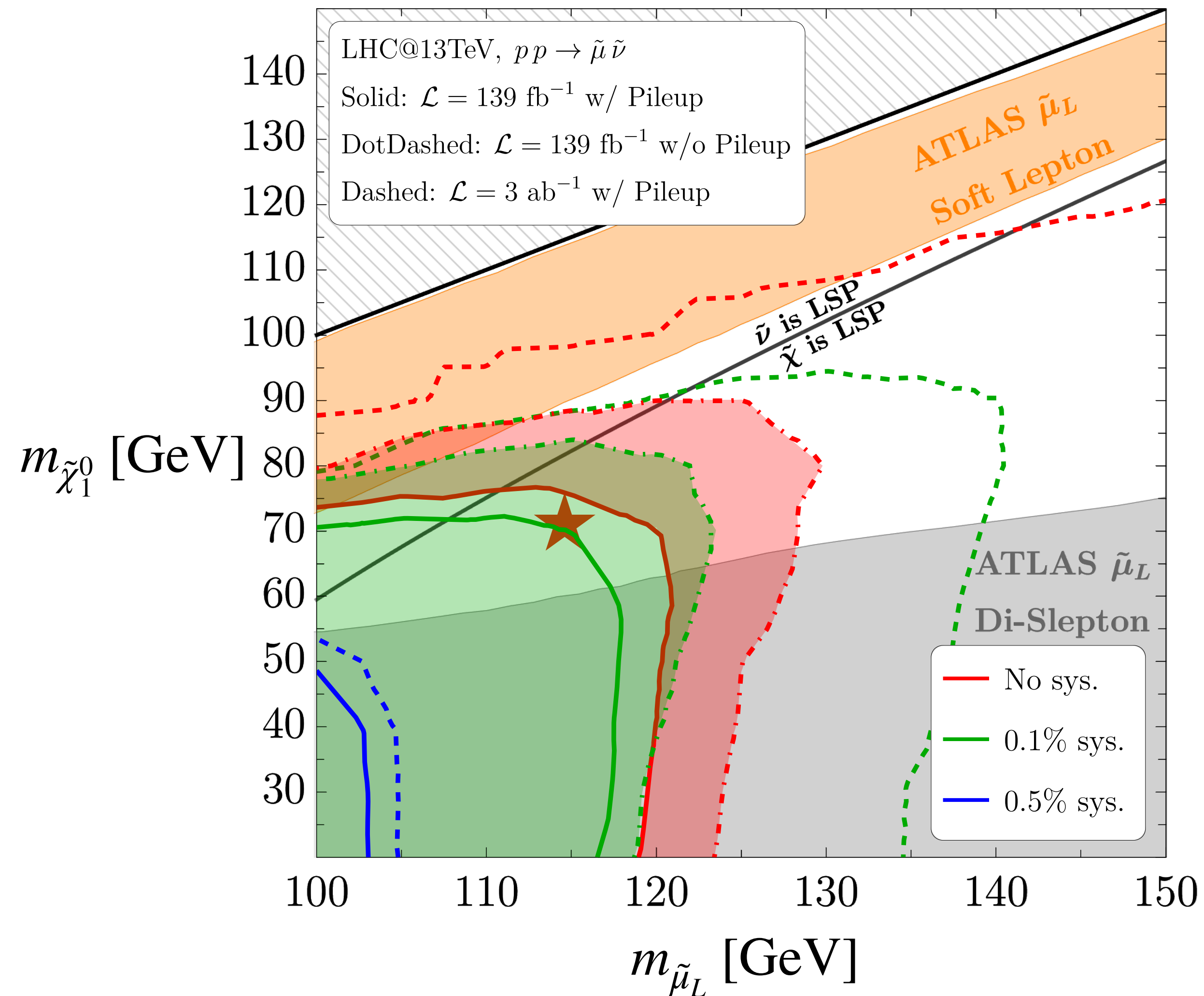
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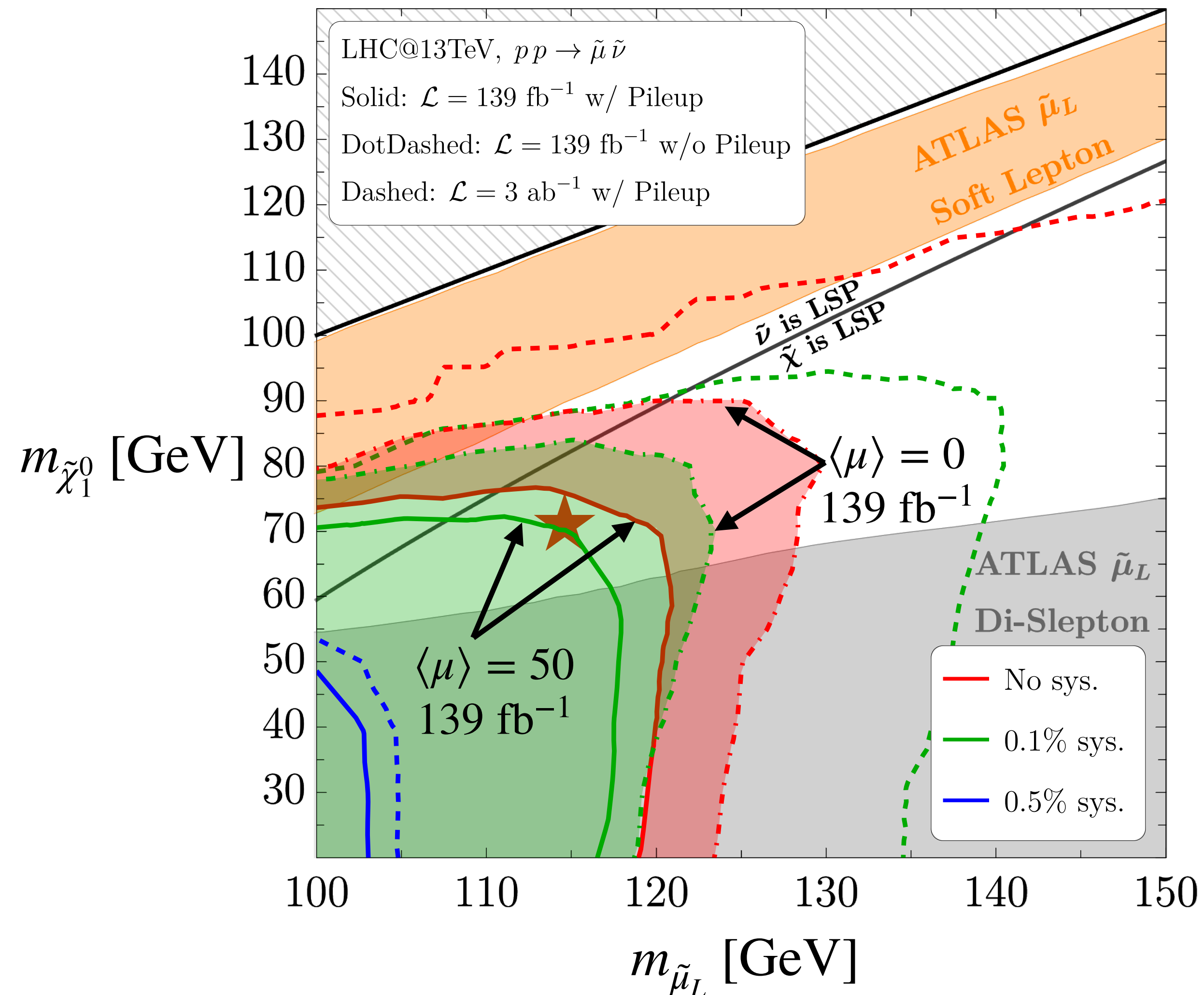
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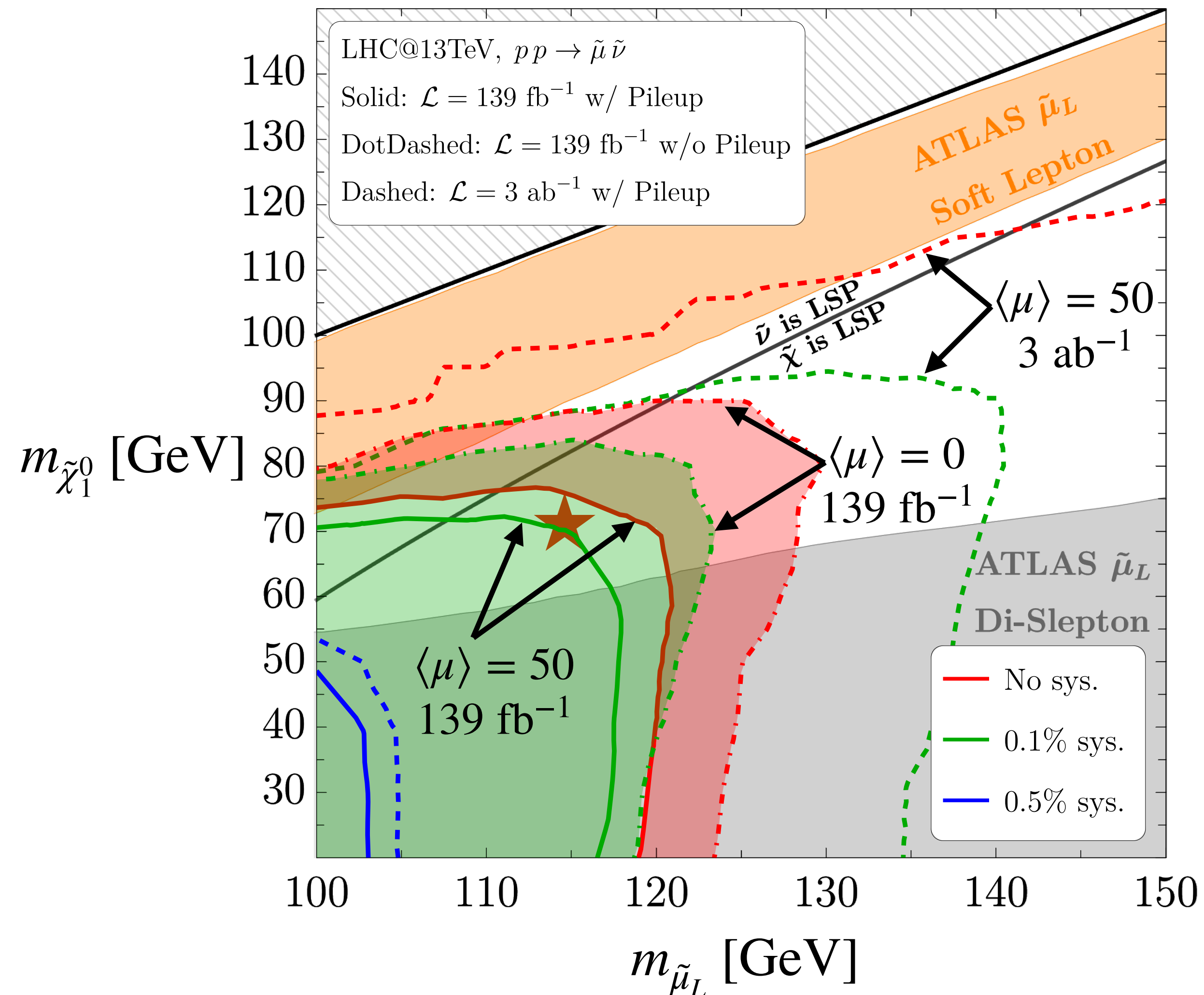
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- Run-2 projections shown
  - with pileup (solid)
  - without pileup (dot-dashed)
  - No-pileup exclusions reached with smaller extension beyond ATLAS fitting range

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NP without on-shell W



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- Run-2 projections shown
  - with pileup (solid)
  - without pileup (dot-dashed)
    - No-pileup exclusions reached with smaller extension beyond ATLAS fitting range
- HL-LHC projections shown (dashed) with pileup
  - saturated at larger systematics

# Extending the fitting range

NP without on-shell W

	ATLAS ( $\mu$ )	$\tilde{\ell}_\mu \tilde{\nu}_\mu$
$p_T^\ell$ (GeV)	> 30 (analysis) > 18 (trigger)	> 30
$m_T$ (GeV)	> 60	> 60
$ \vec{u}_T $ (GeV)	< 30	< 30
$m_T$ range (GeV)	[60, 100]	[60, 120]* [60, 140]
$p_T^\ell$ range (GeV)	[30, 50]	[30, 60]* [30, 70]

\*without pileup

[ATLAS Improved W Mass Measurement](#)

# Extending the fitting range

NP without on-shell W

- Excluding light sleptons with slight extensions in  $p_T^\ell$ ,  $m_T$  fitting ranges

	ATLAS ( $\mu$ )	$\tilde{\ell}_\mu \tilde{\nu}_\mu$
$p_T^\ell$ (GeV)	> 30 (analysis) > 18 (trigger)	> 30
$m_T$ (GeV)	> 60	> 60
$ \vec{u}_T $ (GeV)	< 30	< 30
$m_T$ range (GeV)	[60, 100]	[60, 120]* [60, 140]
$p_T^\ell$ range (GeV)	[30, 50]	[30, 60]* [30, 70]

\*without pileup

[ATLAS Improved W Mass Measurement](#)

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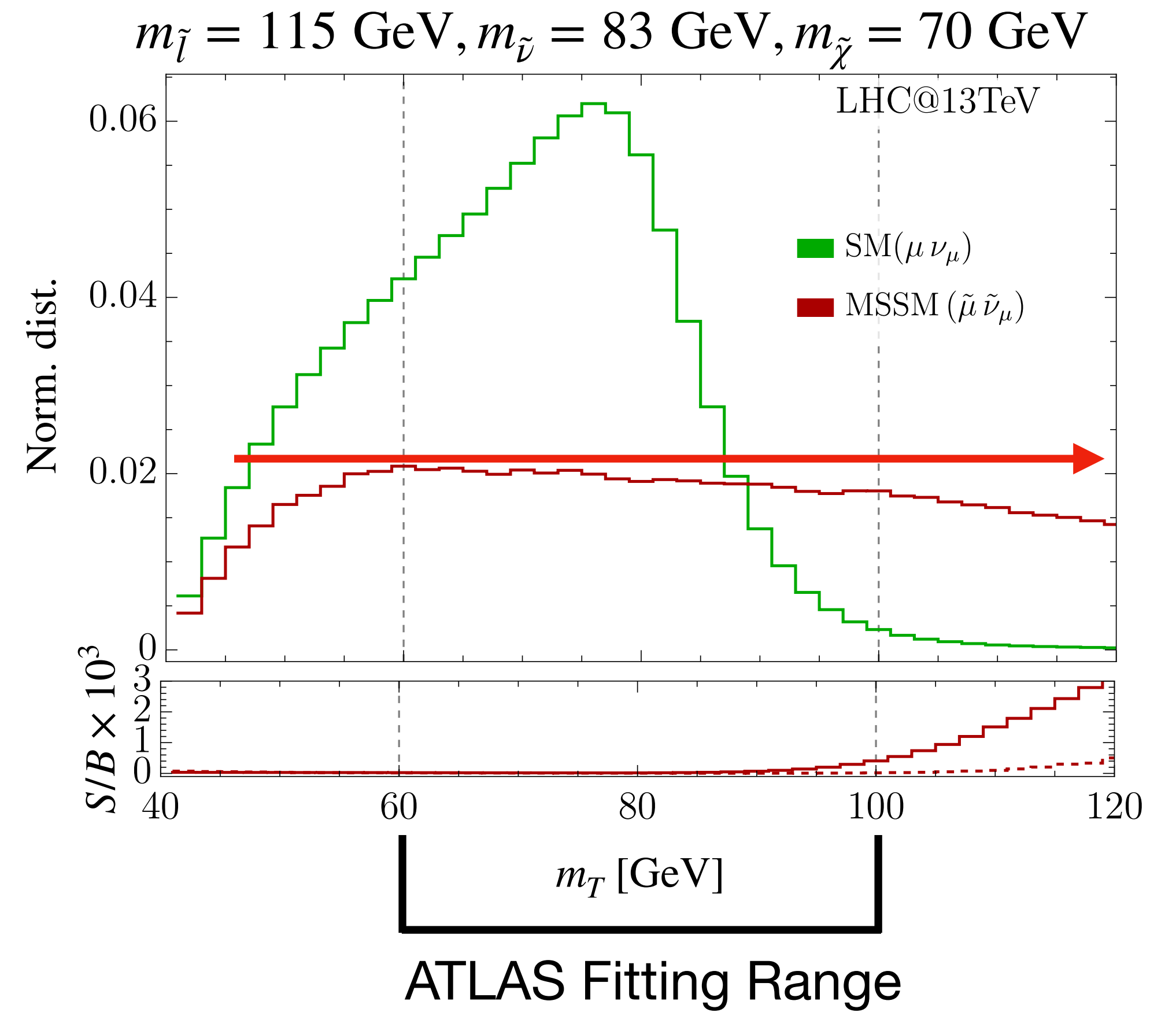
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- Added sensitivity extending to  
 $\sim 250 \text{ GeV } p_T, 500 \text{ GeV } m_T?$

# SUSY Tail Analysis

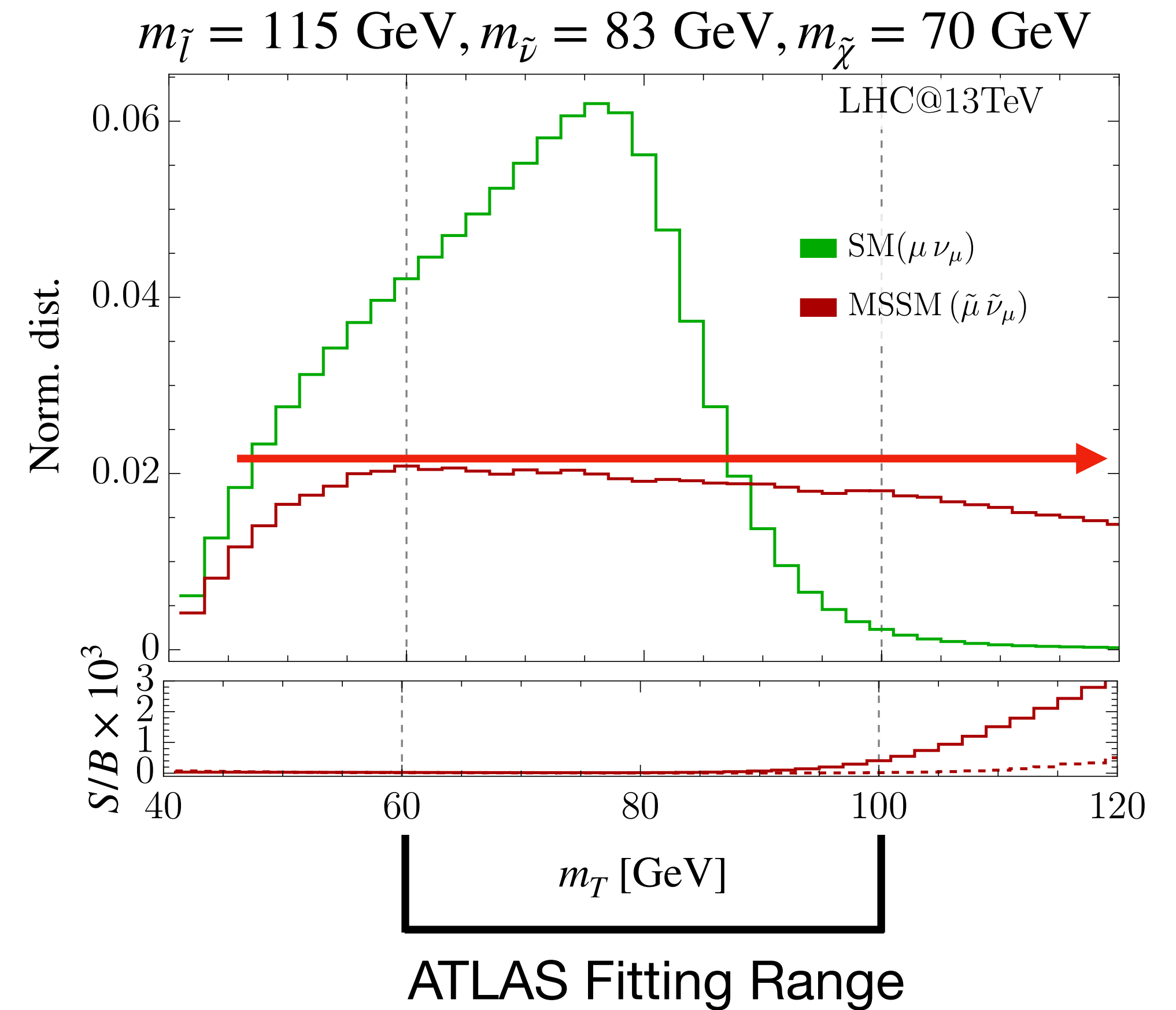
NP without on-shell W



# SUSY Tail Analysis

## NP without on-shell W

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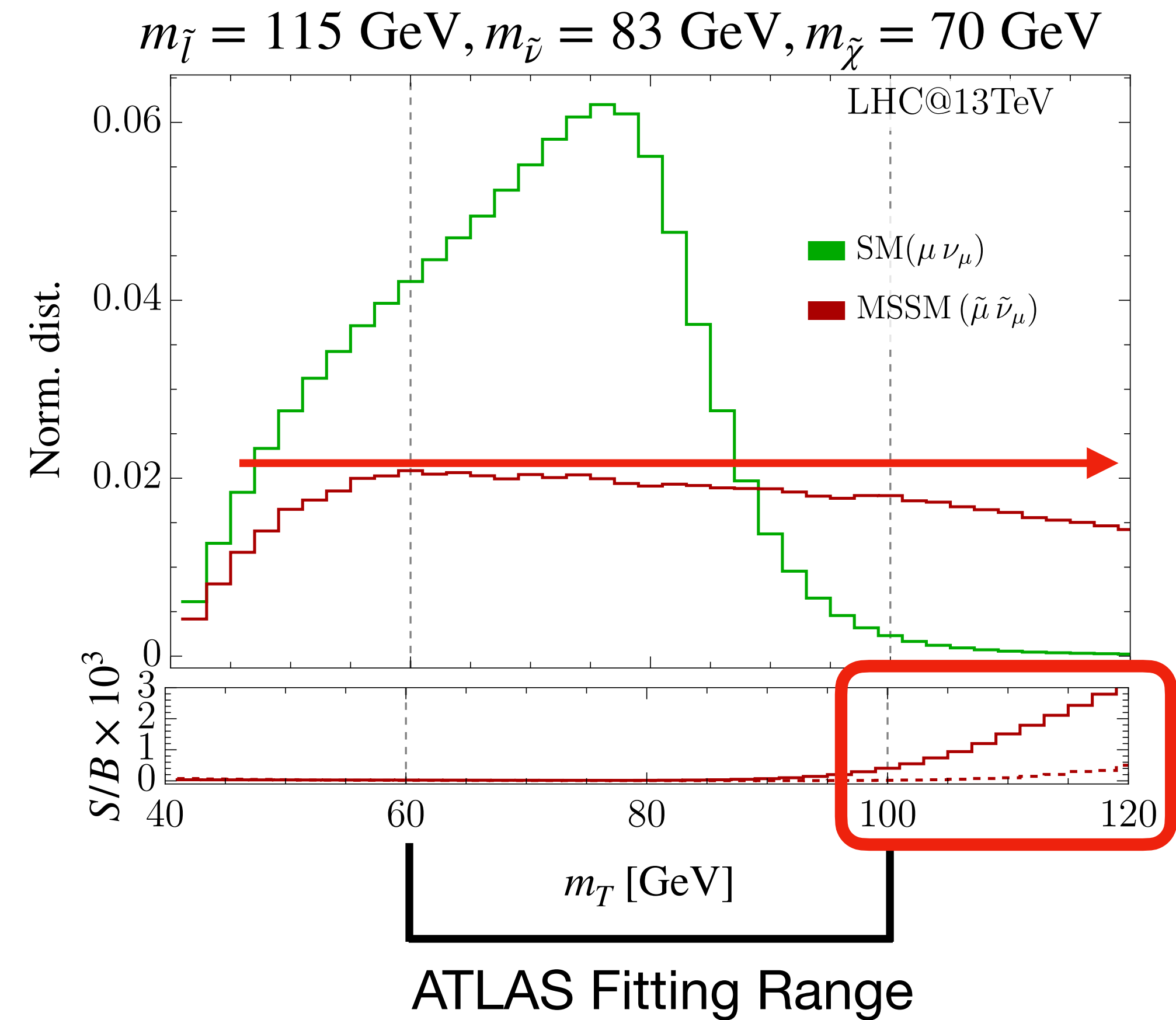




# SUSY Tail Analysis

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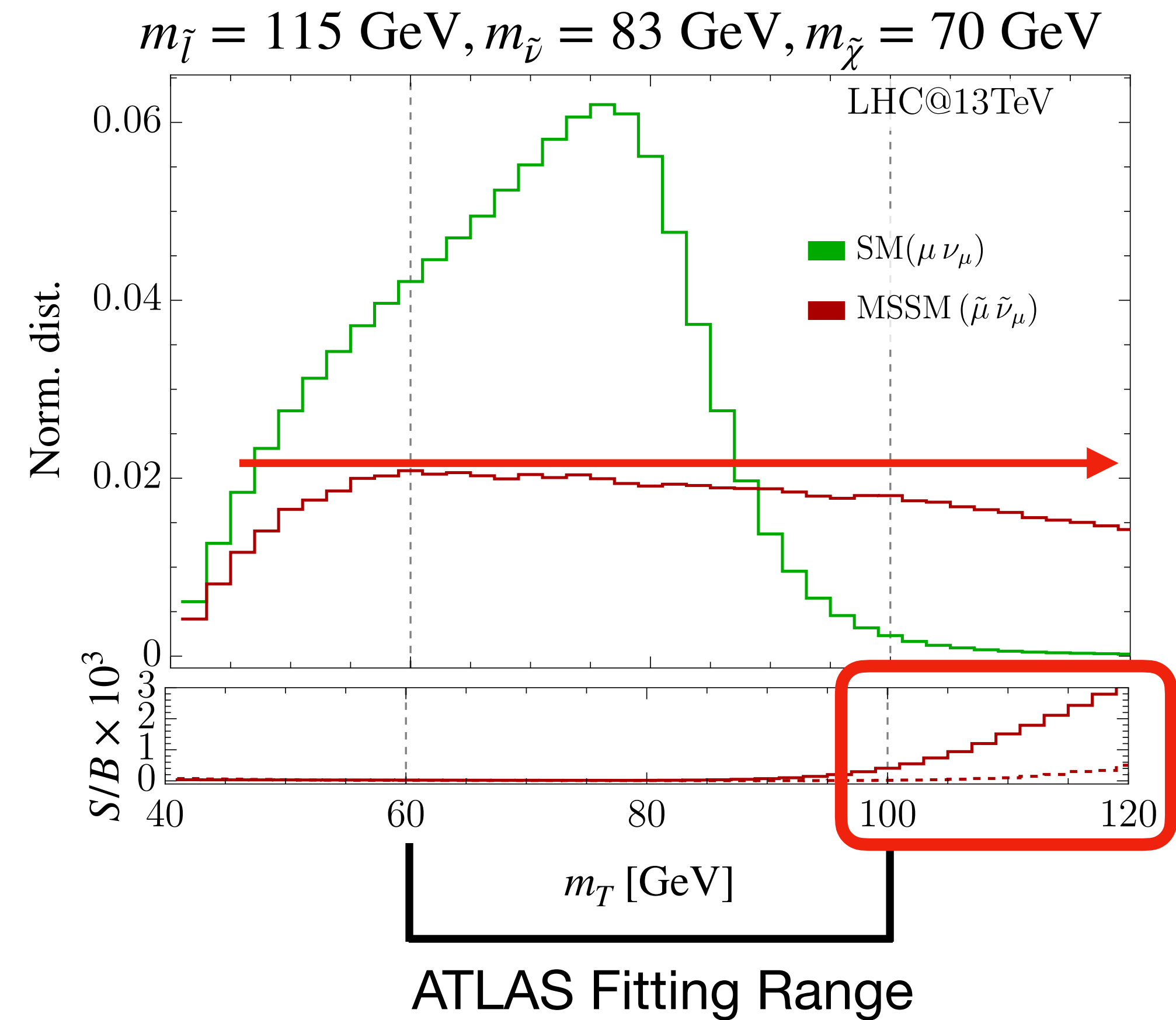
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  - backgrounds are not as under control as in SM W analysis region
  - larger per-bin systematics
  - error on luminosity (correlated systematics)
  - other backgrounds



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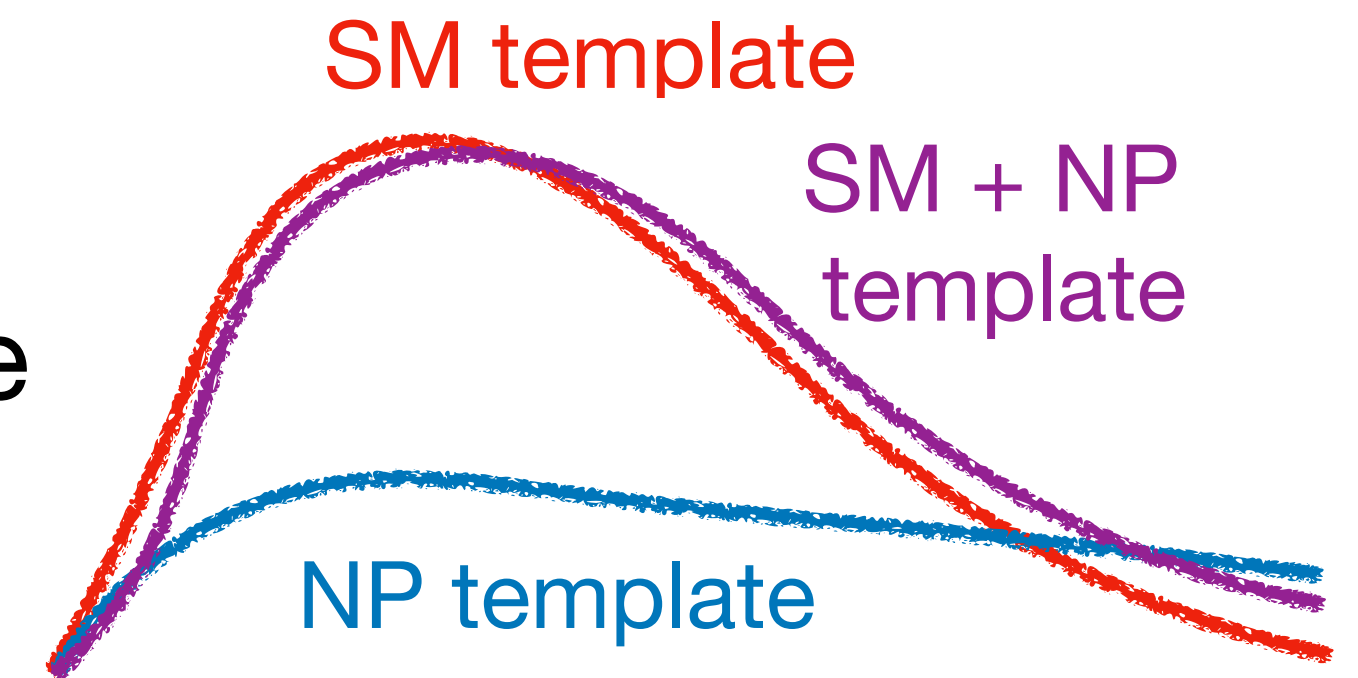
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- Sensitivity greatly improves since background is smaller, but...
  - backgrounds are not as under control as in SM W analysis region
  - larger per-bin systematics
  - error on luminosity (correlated systematics)
  - other backgrounds
- Could a recasted  $W'$  search work for SUSY sleptons?
  - Still investigating



# Brief summary of our idea

- Experiments performed precise measurements of  $m_W$  using kinematic observables  $p_T^\ell, m_T$ 
  - Measured distributions agree with SM templates
- Any NP producing  $\ell + mET$  final state could contaminate the measured sample, and should be added to the template
- Sensitive to NP that produces different kinematics in  $p_T^\ell, m_T$  compared to measured sample
- Adding NP requires a simultaneous fit to NP parameters and  $m_W$
- 3 categories of NP contamination: anomalous W decay, anomalous W production, and  $\ell + mET$  without on-shell W



# Conclusions

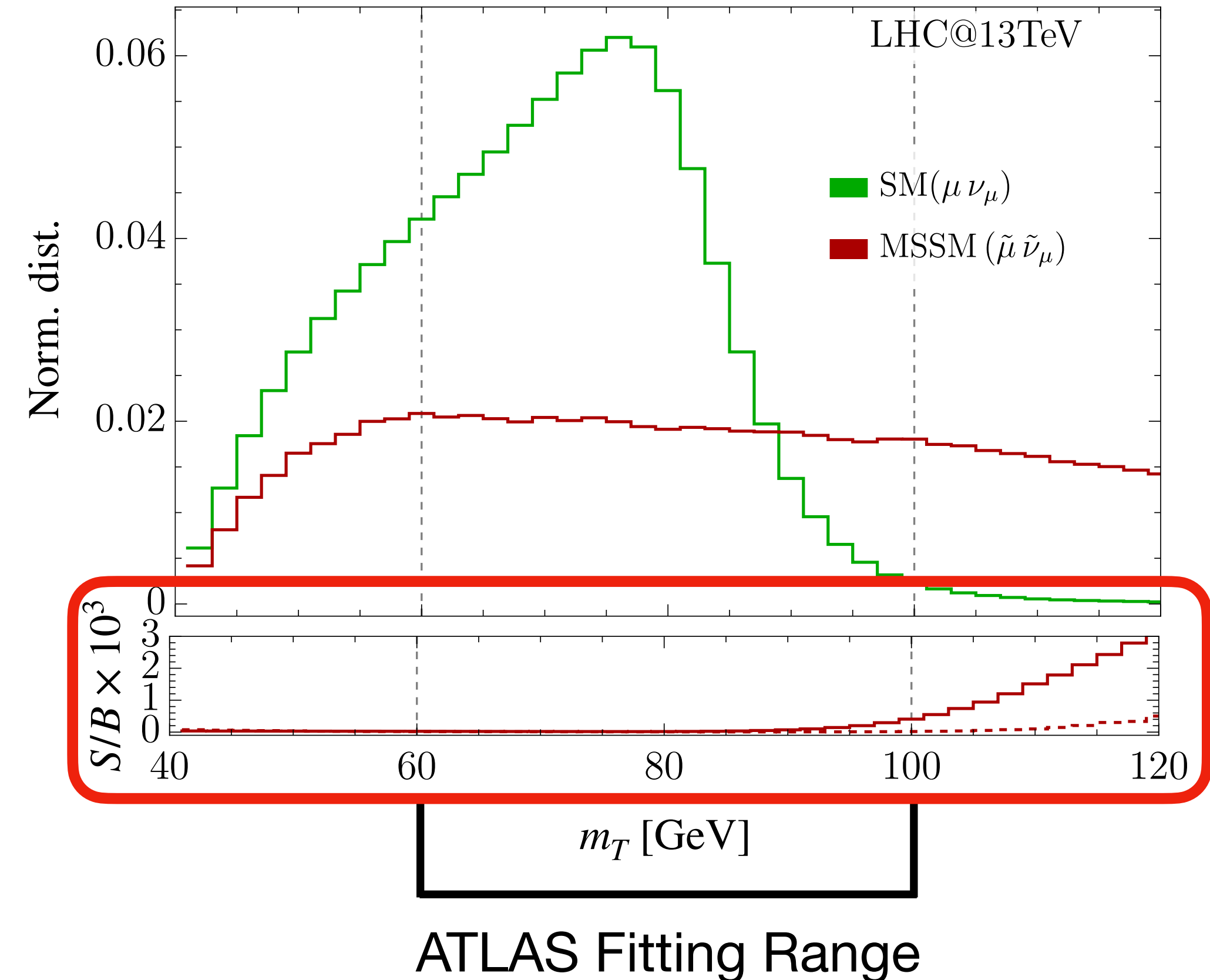
- A new purpose for precision  $W$ -mass measurements: constraining NP that produces the same  $\ell + mET$  final state yet modifying the measured kinematic observables
- Sensitivity to unexplored parameter space for MSSM sleptons!!
- Ongoing work:
  - Further analysis on other models
  - Extending the range of  $p_T^\ell, m_T$  beyond SM  $W$  possibly constrains SUSY sleptons further

**Thank you!**

# Backup Slides

# More details on $Z$ contamination

- Estimating the  $Z$  contamination
  - toy model approach, not what experimentalists would actually do
  - Simulate NP in  $Z \rightarrow l\bar{l}$  sample, then impose invariant mass cut  $80 < m_{l\bar{l}} [\text{GeV}] < 100$ 
    - cut kills most, but not always all of the signal
  - Treat one of the leptons as invisible, and compute  $m_T$  of that event
  - Show  $S/B$  of the  $m_T$  from contamination of  $Z \rightarrow l\bar{l}$  to  $m_T$  of  $W$



# W Decay Width

- CDF fixed  $W$  decay width when measuring  $m_W$
- ATLAS treats it as a nuisance parameter
  - reports uncertainty in  $m_W$  from  $\Gamma_W$  to be 2-7 MeV

Obs.	Mean [MeV]	Elec. Unc.	PDF Unc.	Muon Unc.	EW Unc.	PS & $A_i$ Unc.	Bkg. Unc.	$\Gamma_W$ Unc.	MC stat. Unc.	Lumi Unc.	Recoil Unc.	Total sys.	Data stat.	Total Unc.
$p_T^\ell$	80360.1	8.0	7.7	7.0	6.0	4.7	2.4	2.0	1.9	1.2	0.6	15.5	4.9	16.3
$m_T$	80382.2	9.2	14.6	9.8	5.9	10.3	6.0	7.0	2.4	1.8	11.7	24.4	6.7	25.3

- for simplicity, we fix the decay width like CDF, expecting any shift in gamma to be negligible