Multi-Component Dark Matter: Identifying at Collider

Purusottam Ghosh

IACS Kolkata

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What we know about DM ?

- $\checkmark\,$ Non-luminous and non-baryonic.
- $\checkmark \sim 24\%$ of our Universe made of DM.
- $\checkmark\,$ Massive and interact gravitationally.
- $\checkmark\,$ Stable on cosmological time scale. SM fails to accomodate DM.



Thermal DM : WIMP



- Kinetic Eqlbm. $_{(T_{\rm DM}\,=\,T_{\rm SM})}$ DM SM \leftrightarrow DM SM
- Chemical Eqlbm. $(n_{\text{DM}}^{\text{eq.}} = n_{\text{SM}}^{\text{eq.}})$ DM DM \leftrightarrow SM SM

when $\Gamma \ll H$: DM DM $\not\rightarrow$ SM SM DM becomes relic.

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- larger allowed parameter space .. ref. SB, PG.., JHEP 03 (2020) 090

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Subhaditya Bhattacharya, **P Ghosh**, Jayita Lahiri and Biswarup Mukhopadhyaya

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The peak of the ME distribution depends on both $m_{\rm DM}$ and Δm :



• Two component DM with $m_{\text{DM1}} \neq m_{\text{DM2}}$ and $\Delta m_1 \neq \Delta m_2$

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• Fermion DM (ψ^0): Lepton doublet $\Psi = \begin{pmatrix} \psi & \psi^- \end{pmatrix}^T + \text{Lepton}$ Singlet χ_R ; $(\Psi, \chi) \xrightarrow{\mathcal{Z}'_2} (-\Psi, -\chi)$;

with $m_{\psi^0} < m_{\psi^{\pm}} < m_{\psi_2} < m_{\psi_3}$.



Signal and Background: $\ell^+\ell^- + 0j + ME$



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Subhaditya Bhattacharya, PG, Jayita Lahiri, Biswarup Mukhopadhyaya JHEP12(2022)049

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- The signal+background distribution with S/B = 3 $S = 11\sigma$ for BP1.
- The distribution retain the double hump behaviour → signature of two component DM.
- The separation of the peaks depends on Δm ; while height depnd on production crosssection.

 $\frac{\text{Two peak Gaussian Fitting}: y_H = G(t) = A_1 e^{-\frac{(t-\mu_1)^2}{2\sigma_1^2}} + A_2 e^{-\frac{(t-\mu_2)^2}{2\sigma_2^2}} + \mathcal{B}}{\text{with } \chi^2(\mu_1, \sigma_1; \mu_2, \sigma_2) = \sum_{i=1}^n \frac{\left(G(\mu_1, \sigma_1; \mu_2, \sigma_2)[t_H^i] - y_H^i\right)^2}{y_H^i}.$

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C1:
$$\Delta N_1 = \int_t^{t_1} y dt, \quad \Delta N_2 = \int_{t_1}^{t'} y dt \quad R_{C1} = \frac{|\Delta N_2 - \Delta N_1|}{\sqrt{\Delta N_1}} > 2.$$

C2:
$$R_{C2} = \frac{y(t_{-}) - y'(t_{-})}{\sqrt{y'(t'')}} > 2; \quad C4: R_{C4} = \frac{y(t_{2}) - y(t_{\min})}{\sqrt{y(t_{\min})}} > 2$$

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• $R_{C_{1-4}}(\mathcal{L}) > 2\sigma \rightarrow$ There is definitely the presence of second peak.

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mono-X signal $(X = h, Z, \gamma)$



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DM1: $O_3^s = \frac{c}{\Lambda^2} (B_{\mu\nu} B^{\mu\nu} + W_{\mu\nu} W^{\mu\nu})(\chi^2)$ DM2: $O_1^f = \frac{c}{\Lambda^2} (\bar{L} \gamma^{\mu} L + \bar{\ell}_R \gamma^{\mu} \ell_R) (\bar{\chi} \gamma_{\mu} \chi)$

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The large SM background cannot be reduced beyond a certain limit.Signal+Background distribution can barely show two peak behaviour.

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• We propose to look for bin wise signal significance!



The bin-wise significance comes to the rescue, where one ends up with two-peak only at the presence of the two DM components with $R_{C_{3-4}} > 2\sigma$.

Can we observe two peak distribution at LHC ? Signal: $\ell^-\ell^+ + 0j + X$



- Two peaks can be observed in the signal.
- The signal encounter a huge QCD background.

Works in Progress

The paper **JHEP 04 (2010) 086** by **Partha Konar et al.** studied the signature of multicomponent dark matter at LHC using the M_{T2} topology.

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thank you

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The minimal renormalizable Lagrangian for this model then reads,

$$\mathcal{L} \supset \mathcal{L}^{\text{SDM}} + \mathcal{L}^{\text{FDM}}.$$
 (1)

The Lagrangian for the SDM sector, having inert scalar doublet Φ can be written as :

$$\mathcal{L}^{\text{SDM}} = \left| \left(\partial^{\mu} - ig_2 \frac{\sigma^a}{2} W^{a\mu} - ig_1 \frac{Y}{2} B^{\mu} \right) \Phi \right|^2 - V(\Phi, H);$$

$$V(\Phi, H) = \mu_{\Phi}^2 (\Phi^{\dagger} \Phi) + \lambda_{\Phi} (\Phi^{\dagger} \Phi)^2 + \lambda_1 (H^{\dagger} H) (\Phi^{\dagger} \Phi) + \lambda_2 (H^{\dagger} \Phi) (\Phi^{\dagger} H)$$

$$+ \frac{\lambda_3}{2} [(H^{\dagger} \Phi)^2 + h.c.] .$$

The minimal renormalizable Lagrangian for FDM having one vector-like doublet (Ψ) and one right-handed singlet (χ_R) reads:

$$\mathcal{L}^{\text{FDM}} = \overline{\Psi}_{L(R)} \left[i \gamma^{\mu} (\partial_{\mu} - i g_2 \frac{\sigma^a}{2} W^a_{\mu} - i g_1 \frac{Y'}{2} B_{\mu}) \right] \Psi_{L(R)} + \overline{\chi_R} \left(i \gamma^{\mu} \partial_{\mu} \right) \chi$$
$$- m_{\psi} \overline{\Psi} \Psi - \left(\frac{1}{2} m_{\chi} \overline{\chi_R} (\chi_R)^c + h.c \right) - \frac{Y}{\sqrt{2}} \left(\overline{\Psi_L} \widetilde{H} \chi_R + \overline{\Psi_R} \widetilde{H} \chi_R^c \right)$$

where $\Psi_{L(R)} = P_{L(R)}\Psi$; $P_{L/R} = \frac{1}{2}(1 \mp \gamma_5)$.

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BPs	SDM sector $\{m_{\phi^0}, \Delta m_1, \lambda_L\}$	FDM sector $\{m_{\psi_1}, \Delta m_2, \sin \theta\}$	$\Omega_{\phi^0} h^2$	$\Omega_{\psi_1}h^2$	$\sigma^{\rm eff}_{\phi^0}~({\rm cm}^2)$	$\sigma_{\psi_1}^{\rm eff}~({\rm cm}^2)$	$\mathrm{BR}(H_{\mathrm{inv}})\%$
BP1	100, 10, 0.01	60.5, 370, 0.022	0.00221	0.1195	3.45×10^{-46}	2.03×10^{-47}	0.25
BP2	100, 10, 0.01	58.91, 285, 0.032	0.00221	0.10962	3.45×10^{-46}	5.38×10^{-47}	1.60
BP3	100, 10, 0.01	58.87, 176, 0.04	0.00221	0.11941	3.45×10^{-46}	5.00×10^{-47}	1.50
BP4	100, 10, 0.01	58.48, 190, 0.042	0.00221	0.1114	3.45×10^{-46}	7.01×10^{-47}	2.4

Table 2. Benchmark points of the model; contribution to relic density, spin-independent direct detection cross-section as well as that of invisible Higgs decay branching ratios of the DM components ϕ^0 and ψ_1 are mentioned.

Benc	hmarks	Collider cross-section (fb)								
		$\sigma_{\text{total}}(\text{OSD})$			$\sigma_{\phi^+\phi^-}(OSD)$			$\sigma_{\psi^+\psi^-}(OSD)$		
\sqrt{s}	Points	P1	P2	P3	P1	P2	P3	P1	P2	P3
1000	BP1	232(10.8)	115(5.5)	58.5(2.75)	57.4(2.9)	28.9(1.5)	14.5(0.75)	173(8.4)	83.0(4.0)	44.0(2.0)
1000	BP2	276(13.4)	141(6.6)	70.0(3.3)	57.4(2.9)	28.9(1.5)	14.5(0.75)	218(10.4)	111(5.3)	55.5(2.7)
500	BP3	686(33.0)	339(15.9)	168.1(7.8)	180(8.9)	90.3(4.5)	44.3(2.3)	494(22.2)	253(11.3)	123.8(5.5)
500	BP4	345(16.7)	170(8.4)	83.5(3.9)	180(8.9)	90.3(4.5)	44.3(2.3)	171.4(7.4)	82.4(3.9)	39.2(1.9)

Table 3. Signal cross-sections for HDSP pair production (OSD final state) at ILC. Total crosssection (σ_{total}), as well as individual contributions from SDM (σ_{q+q-}) and FDM (σ_{q+q-}) are mentioned. Three choices of beam polarisation are used: $P1 \equiv \{P_{e^-}: -0.8, P_{e^+}: +0.3\}, P2 \equiv \{P_{e^-}: 0, P_{e^+}: 0\}$ and $P3 \equiv \{P_{e^-}: +0.8, P_{e^+}: -0.3\}$. (Menergy ($\langle \phi \rangle$) is in the units of GeV.

Backg	rounds	Cross-section(fb)			
\sqrt{s}	Processes	P1	P2	P3	
1 TeV	WW	296	128	18.3	
	ZZ	7.5	4.4	3.5	
	WWZ	1.2	0.5	0.08	
500 C-V	WW	802	342	51	
300 Gev	ZZ	21	12	9.6	
	WWZ	0.8	0.37	0.06	

Table 4. Production cross-sections for $W^+(\ell^+\nu)W^-(\ell^-\overline{\nu})$, $Z(\ell^+\ell^-)Z(\nu\overline{\nu})$ and $W^+(\ell^+\nu)W^-(\ell^-\overline{\nu})Z(\nu\overline{\nu})$ background at $\sqrt{s} = 1$ TeV and 500 GeV for various polarization combinations P1, P2 and P3 (see caption of Table 3).

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