Scotogenic $U(1)_{L_{\mu}-L_{\tau}}$ origin of $(g-2)_{\mu}$, W-mass anomaly, and 95 GeV excess

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<u>CMS result on low mass Higgs search in $pp \rightarrow \phi \rightarrow \gamma \gamma$ </u>



CMS-PAS-HIG-20-002

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<u>CMS result on low mass Higgs search in $pp \rightarrow \phi \rightarrow \gamma \gamma$ </u>



$$\mu_{\gamma\gamma} = \frac{\sigma(gg \to H_{95})}{\sigma_{\rm SM}(gg \to H_{95})} \times \frac{{\rm BR}(H_{95} \to \gamma\gamma)}{{\rm BR}_{\rm SM}(H_{95} \to \gamma\gamma)}$$

CMS-PAS-HIG-20-002

<u>CMS result on low mass Higgs search in $pp \rightarrow \phi \rightarrow \gamma \gamma$ </u>



CMS-PAS-HIG-20-002

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ATLAS result on low mass Higgs search in $pp \rightarrow \phi \rightarrow \gamma \gamma$



https://indico.cern.ch/event/1281604/

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ATLAS result on low mass Higgs search in $pp \rightarrow \phi \rightarrow \gamma \gamma$



https://indico.cern.ch/event/1281604/

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Anomalaous magnetic moment of muon

$$\overrightarrow{\mu_{\mu}} = g_{\mu} \left(\frac{q}{2m}\right) \overrightarrow{S} , \quad a_{\mu} = \frac{(g_{\mu} - 2)}{2}$$



Figure: Experimental value of a_{μ}



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Anomalaous magnetic moment of muon

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Figure: Experimental value of a_{μ}

$$a_{\mu}^{\rm SM} = 116591810(43) \times 10^{-11}$$

 $a_{\mu}^{\exp} = 116592059(22) \times 10^{-11}$

Anomalaous magnetic moment of muon

$$\overrightarrow{\mu_{\mu}} = g_{\mu} \left(\frac{q}{2m}\right) \overrightarrow{S}, \quad a_{\mu} = \frac{(g_{\mu} - 2)}{2}$$



Figure: Experimental value of a_{μ}

$$a_{\mu}^{\text{SM}} = 116591810(43) \times 10^{-11} \qquad a_{\mu}^{\text{exp}} = 116592059(22) \times 10^{-11}$$
$$\Delta a_{\mu} = a_{\mu}^{\text{exp}} - a_{\mu}^{\text{SM}} = 249(48) \times 10^{-11} (5.1 \ \sigma)$$

2308.06230

CDF-II W-mass anomaly



 $M_W(SM) = 80357 \pm 6 \text{ MeV}$

$M_W(\text{CDF}) = 80433.5 \pm 9.4 \text{ MeV}$

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CDF-II W-mass anomaly



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 7σ away from the SM prediction

Science 376 (2022) 6589, 170-176

Gauge	Fermion Fields			Sc	Scalar Field		
Group	N _e	N_{μ}	N_{τ}	Φ_1	Φ_2	η	
$S U(2)_L$	1	1	1	1	1	2	
$U(1)_Y$	0	0	0	0	0	$\frac{1}{2}$	
$U(1)_{L_{\mu}-L_{\tau}}$	0	1	-1	1	2	Ō	
Z ₂	-1	-1	-1	+1	+1	-1	



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 $\mathcal{L}_{Fermion} \supseteq \overline{N_e} i \gamma^{\mu} \partial_{\mu} N_e + \overline{N_{\mu}} i \gamma^{\mu} \mathfrak{D}_{\mu} N_{\mu} + \overline{N_{\tau}} i \gamma^{\mu} \mathfrak{D}_{\mu} N_{\tau} - \frac{M_{ee}}{2} \overline{N_e^C} N_e - M_{\mu\tau} \overline{N_{\mu}^C} N_{\tau} - Y_{e\mu} \Phi_1^{\dagger} \overline{N_e^C} N_{\mu} - Y_{e\tau} \Phi_1 \overline{N_e^C} N_{\tau} - Y_{\mu \mu} \Phi_2^{\dagger} \overline{N_{\mu}^C} N_{\mu} - Y_{\tau} \Phi_2 \overline{N_{\tau}^C} N_{\tau} - Y_{De} \overline{L_e} \overline{\eta} N_e - Y_{D\mu} \overline{L_{\mu}} \overline{\eta} N_{\mu} - Y_{D\tau} \overline{L_{\tau}} \overline{\eta} N_{\tau} - Y_{le} \overline{L_e} H e_R - Y_{l\mu} \overline{L_{\mu}} H \mu_R - Y_{l\tau} \overline{L_{\tau}} \overline{L_{\tau}} H \tau_R + \text{h.c.}$ (3)



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Gauge	Fermion Fields			Sc	Scalar Field		
Group	Ne	N_{μ}	N_{τ}	Φ_1	Φ_2	η	
$SU(2)_L$	1	1	1	1	1	2	
$U(1)_Y$	0	0	0	0	0	$\frac{1}{2}$	
$U(1)_{L_{\mu}-L_{\tau}}$	0	1	-1	1	2	0	
Z ₂	-1	-1	-1	+1	+1	-1	

 $\mathcal{L}_{Fermion} \supseteq \overline{N_e} i \gamma^{\mu} \partial_{\mu} N_e + \overline{N_{\mu}} i \gamma^{\mu} \mathfrak{D}_{\mu} N_{\mu} + \overline{N_{\tau}} i \gamma^{\mu} \mathfrak{D}_{\mu} N_{\tau} - \frac{M_{ee}}{2} \overline{N_e^C} N_e - M_{\mu\tau} \overline{N_{\mu}^C} N_{\tau} - Y_{e\mu} \Phi_1^{\dagger} \overline{N_e^C} N_{\mu} - Y_{e\tau} \Phi_1 \overline{N_e^C} N_{\tau} \\ - Y_{\mu} \Phi_2^{\dagger} \overline{N_{\mu}^C} N_{\mu} - Y_{\tau} \Phi_2 \overline{N_{\tau}^C} N_{\tau} - Y_{De} \overline{L_e} \tilde{\eta} N_e - Y_{D\mu} \overline{L_{\mu}} \tilde{\eta} N_{\mu} - Y_{D\tau} \overline{L_{\tau}} \tilde{\eta} N_{\tau} - Y_{le} \overline{L_e} H e_R - Y_{l\mu} \overline{L_{\mu}} H \mu_R \\ - Y_{l\tau} \overline{L_{\tau}} H \tau_R + \text{h.c.}$

$$\mathcal{L}_{\text{Gauge}} = -\frac{1}{4} (Z_{\mu\tau})_{\mu\nu} Z_{\mu\tau}^{\mu\nu} - \frac{\epsilon}{2} (Z_{\mu\tau})_{\mu\nu} B^{\mu\nu}$$
(4)



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Group	N _e	N_{μ}	N_{τ}	Φ_1	Φ_2	η	
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Z_2	-1	-1	-1	+1	+1	-1	

 $\mathcal{L}_{Fermion} \supseteq \overline{N_e} i \gamma^{\mu} \partial_{\mu} N_e + \overline{N_{\mu}} i \gamma^{\mu} \mathfrak{D}_{\mu} N_{\mu} + \overline{N_{\tau}} i \gamma^{\mu} \mathfrak{D}_{\mu} N_{\tau} - \frac{M_{ee}}{2} \overline{N_e^C} N_e - M_{\mu\tau} \overline{N_{\mu}^C} N_{\tau} - Y_{e\mu} \Phi_1^{\dagger} \overline{N_e^C} N_{\mu} - Y_{e\tau} \Phi_1 \overline{N_e^C} N_{\tau} \\ - Y_{\mu} \Phi_2^{\dagger} \overline{N_{\mu}^C} N_{\mu} - Y_{\tau} \Phi_2 \overline{N_{\tau}^C} N_{\tau} - Y_{De} \overline{L_e} \tilde{\eta} N_e - Y_{D\mu} \overline{L_{\mu}} \tilde{\eta} N_{\mu} - Y_{D\tau} \overline{L_{\tau}} \tilde{\eta} N_{\tau} - Y_{le} \overline{L_e} H e_R - Y_{l\mu} \overline{L_{\mu}} H \mu_R \\ - Y_{l\tau} \overline{L_{\tau}} H \tau_R + \text{h.c.}$

$$\mathcal{L}_{\text{Gauge}} = -\frac{1}{4} (Z_{\mu\tau})_{\mu\nu} Z_{\mu\tau}^{\mu\nu} - \frac{\epsilon}{2} (Z_{\mu\tau})_{\mu\nu} B^{\mu\nu}$$
$$\mathcal{L}_{scalar} = |\mathcal{D}_{\mu}H|^2 + |\mathcal{D}_{\mu}\eta|^2 + |\mathfrak{D}_{\mu}\Phi_i|^2 - V(H, \Phi_i, \eta)$$

$$\begin{aligned} V(H,\Phi_i,\eta) &= -\mu_H^2 \left(H^{\dagger}H\right) + \lambda_H \left(H^{\dagger}H\right)^2 - \mu_{\Phi_i}^2 (\Phi_i^{\dagger}\Phi_i) + \lambda_{\Phi_i} (\Phi_i^{\dagger}\Phi_i)^2 + \lambda_{H\Phi_i} (H^{\dagger}H) (\Phi_i^{\dagger}\Phi_i) + m_{\eta}^2 (\eta^{\dagger}\eta) + \\ \lambda_2 (\eta^{\dagger}\eta)^2 + \lambda_3 (\eta^{\dagger}\eta) (H^{\dagger}H) + \lambda_4 (\eta^{\dagger}H) (H^{\dagger}\eta) + \frac{\lambda_5}{2} [(H^{\dagger}\eta)^2 + (\eta^{\dagger}H)^2] + \lambda_{\eta\Phi_i} (\eta^{\dagger}\eta) (\Phi_i^{\dagger}\Phi_i) + \\ \lambda_{\Phi_1\Phi_2} (\Phi_1^{\dagger}\Phi_1) (\Phi_2^{\dagger}\Phi_2) + [\mu \Phi_1^2 \Phi_2^{\dagger} + \text{h.c.}] \end{aligned}$$

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Scalar Mixing

The VEV alignments of the scalars are given as,

$$H = \begin{pmatrix} 0\\ \frac{\nu+h}{\sqrt{2}} \end{pmatrix}, \Phi_1 = \frac{1}{\sqrt{2}}(\phi_1 + \nu_1), \Phi_2 = \frac{1}{\sqrt{2}}(\phi_2 + \nu_2).$$



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$$\begin{pmatrix} h \\ \phi_1 \\ \phi_2 \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & c_{13}s_{12} & s_{13} \\ -c_{12}s_{13}s_{23} - c_{23}s_{12} & c_{12}c_{23} - s_{12}s_{13}s_{23} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}c_{23}s_{13} & -c_{12}s_{23} - c_{23}s_{12}s_{13} & c_{13}c_{23} \end{pmatrix} \begin{pmatrix} H_1 \\ H_2 \\ H_3 \end{pmatrix}$$

where $\cos \theta_{ij} = c_{ij}$ and $\sin \theta_{ij} = s_{ij}$.

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where $\cos \theta_{ij} = c_{ij}$ and $\sin \theta_{ij} = s_{ij}$.

$$M_{Z_{\mu\tau}} = g_{\mu\tau} \sqrt{v_1^2 + 4v_2^2}.$$

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Fermion Mixing

$$\begin{pmatrix} N_1 \\ N_2 \\ N_3 \end{pmatrix} = \begin{pmatrix} \cos\beta_{12} & \frac{\sin\beta_{12}}{\sqrt{2}} & \frac{\sin\beta_{12}}{\sqrt{2}} \\ -\sin\beta_{12} & \frac{\cos\beta_{12}}{\sqrt{2}} & \frac{\cos\beta_{12}}{\sqrt{2}} \\ 0 & -\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} . \begin{pmatrix} N_e \\ N_\mu \\ N_\tau \end{pmatrix}$$



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Fermion Mixing

$$\begin{pmatrix} N_1 \\ N_2 \\ N_3 \end{pmatrix} = \begin{pmatrix} \cos\beta_{12} & \frac{\sin\beta_{12}}{\sqrt{2}} & \frac{\sin\beta_{12}}{\sqrt{2}} \\ -\sin\beta_{12} & \frac{\cos\beta_{12}}{\sqrt{2}} & \frac{\cos\beta_{12}}{\sqrt{2}} \\ 0 & -\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} \cdot \begin{pmatrix} N_e \\ N_\mu \\ N_\tau \end{pmatrix}$$

The mixing angle is given by,

$$\beta_{13} = 0; \ \beta_{23} = \frac{\pi}{4}; \ \tan(2\beta_{12}) = \frac{2\sqrt{2}Y_{e\mu}v_1}{M_{ee} - Y_{\mu}v_2 - M_{\mu\tau}}$$

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$$Y_D = \begin{pmatrix} Y_{De} & 0 & 0 \\ 0 & Y_{D\mu} & 0 \\ 0 & 0 & Y_{D\tau} \end{pmatrix}, \ M_R = \begin{pmatrix} M_{ee} & Y_{e\mu}v_1 & Y_{e\tau}v_1 \\ Y_{e\mu}v_1 & Y_{\mu}v_2 & M_{\mu\tau} \\ Y_{e\tau}v_1 & M_{\mu\tau} & Y_{\tau}v_2 \end{pmatrix}$$



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2

Neutrino mass



$$Y_D = \begin{pmatrix} Y_{De} & 0 & 0 \\ 0 & Y_{D\mu} & 0 \\ 0 & 0 & Y_{D\tau} \end{pmatrix}, \ M_R = \begin{pmatrix} M_{ee} & Y_{e\mu}v_1 & Y_{e\tau}v_1 \\ Y_{e\mu}v_1 & Y_{\mu}v_2 & M_{\mu\tau} \\ Y_{e\tau}v_1 & M_{\mu\tau} & Y_{\tau}v_2 \end{pmatrix}$$

$$(M_{\nu})_{\alpha\beta} = \sum_{k} \frac{Y_{\alpha k} Y_{k\beta} M_{k}}{32\pi^{2}} \left[L_{k}(M_{\eta_{R}}^{2}) - L_{k}(M_{\eta_{I}}^{2}) \right]$$



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Neutrino mass



$$Y_D = \begin{pmatrix} Y_{De} & 0 & 0 \\ 0 & Y_{D\mu} & 0 \\ 0 & 0 & Y_{D\tau} \end{pmatrix}, \ M_R = \begin{pmatrix} M_{ee} & Y_{e\mu}v_1 & Y_{e\tau}v_1 \\ Y_{e\mu}v_1 & Y_{\mu}v_2 & M_{\mu\tau} \\ Y_{e\tau}v_1 & M_{\mu\tau} & Y_{\tau}v_2 \end{pmatrix}$$

$$(M_{\nu})_{\alpha\beta} = \sum_{k} \frac{Y_{\alpha k} Y_{k\beta} M_{k}}{32\pi^{2}} \left[L_{k}(M_{\eta_{R}}^{2}) - L_{k}(M_{\eta_{I}}^{2}) \right]$$

$$Y_{\alpha k} = \left(U D_{\nu}^{1/2} R^{\dagger} \Lambda^{1/2} \right)_{\alpha k}$$

$$\Lambda_{k} = \frac{2\pi^{2}}{\lambda_{5}} \zeta_{k} \frac{2M_{k}}{v^{2}}, \quad \zeta_{k} = \left(\frac{M_{k}^{2}}{8(M_{\eta_{R}}^{2} - M_{\eta_{I}}^{2})} \left[L_{k}(M_{\eta_{R}}^{2}) - L_{k}(M_{\eta_{I}}^{2})\right]\right)^{-1}, \quad L_{k}(m^{2}) = \frac{m^{2}}{m^{2} - M_{k}^{2}} \ln \frac{m^{2}}{M_{k}^{2}}$$

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$$\Delta a_{\mu} = \frac{\alpha_{\mu\tau}}{2\pi} \int_{0}^{1} dx \frac{2m_{\mu}^{2}x^{2}(1-x)}{x^{2}m_{\mu}^{2} + (1-x)M_{Z_{\mu\tau}}^{2}} ; \quad \alpha_{\mu\tau} = g_{\mu\tau}^{2}/(4\pi)$$



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W-mass anomaly



$$T = \frac{\Theta(m_{\eta^+}^2, m_{\eta_R}^2) + \Theta(m_{\eta^+}^2, m_{\eta_I}^2) - \Theta(m_{\eta_R}^2, m_{\eta_I}^2)}{16\pi^2 \alpha_{\rm em}(M_Z) v^2} \ ; \ \Theta(x, y) \ \equiv \ \frac{1}{2}(x+y) - \frac{xy}{x-y} \ln\left(\frac{x}{y}\right)$$

$$S = \frac{1}{12\pi} \log \left[\frac{M_{\eta_R}^2 + M_{\eta_I}^2}{2M_{\eta^+}^2} \right].$$

$$M_W \simeq M_W^{SM} \left[1 - \frac{\alpha_{\rm em}(M_Z)(S - 2\,\cos^2\theta_W\,T)}{4(\cos^2\theta_W - \sin^2\theta_W)} \right]$$

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$$H_1 \equiv H_{\text{SM}}, H_2 \equiv H_{95}, H_3 \equiv H_{\text{heavy}}.$$



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$$H_1 \equiv H_{\text{SM}}, H_2 \equiv H_{95}, H_3 \equiv H_{\text{heavy}}.$$

$$\mu_{\gamma\gamma} = \sin^2 \theta_{12} \cos^2 \theta_{13} \quad \frac{\mathrm{BR}(H_{95} \to \gamma\gamma)}{\mathrm{BR}_{\mathrm{SM}}(H_{95} \to \gamma\gamma)} \,.$$



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The diphoton decay width of Z_2 -even scalars is given by,

$$\Gamma(H_i \to \gamma \gamma) = \frac{G_F \alpha^2 m_{H_i}^3}{128 \sqrt{2} \pi^3} \left| \sum_f N_c Q_f^2 g_{H_i f f} A_{1/2}(\tau_f) + g_{H_i W W} A_1(\tau_W) + \frac{\nu}{2 M_{\eta^+}^2} C_{H_i \eta \eta} A_0(\tau_\eta) \right|^2$$

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$$H_1 \equiv H_{\text{SM}}, H_2 \equiv H_{95}, H_3 \equiv H_{\text{heavy}}.$$

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Coupling between the neutral scalars and the inert doublet is given by,

$$\begin{array}{lll} C_{H_1\eta\eta} & = & c_{12}c_{13}\lambda_3v + (-c_{23}s_{12} - c_{12}s_{13}s_{23})\lambda_{\eta\Phi_1}v_1 + (-c_{12}c_{23}s_{13} + s_{12}s_{23})\lambda_{\eta\Phi_2}v_2 \,, \\ C_{H_2\eta\eta} & = & s_{12}c_{13}\lambda_3v + (c_{23}c_{12} - s_{12}s_{13}s_{23})\lambda_{\eta\Phi_1}v_1 + (-s_{12}c_{23}s_{13} + c_{12}s_{23})\lambda_{\eta\Phi_2}v_2 \,, \\ C_{H_3\eta\eta} & = & s_{13}\lambda_3v + c_{13}s_{23}\lambda_{\eta\Phi_1}v_1 + c_{13}c_{23}\lambda_{\eta\Phi_2}v_2 \,. \end{array}$$

Range of parameter for scan

The range in which the free parameters are randomly varied for the numerical analysis.

Parameter	Scanned range		
M_{η_R} (GeV)	45,150		
$M_{\eta^+} - M_{\eta_R} (\text{GeV})$	1,100		
λ_3	$1, 4\pi$		
λ_5	$10^{-8}, 10^{-2}$		
$\lambda_{\eta\Phi_1}$	$0.01, 4\pi$		
$\lambda_{\eta\Phi_2}$	$0.01, 4\pi$		
$\sin \theta_{12}$	0.1, 0.7		
$\sin \theta_{23}$	0.1, 0.7		
$\sin \theta_{13}$	0.001, 0.1		
$\sin\beta_{12}$	$10^{-4}, 0.7$		
v_1 (GeV)	20,60		
m_{H_3} (GeV)	200, 500		

We also impose the LEP limits on the doublet scalar as $M_{\eta_R} + M_{\eta_I} > M_Z$ and a conservative limit on the charged scalar $M_{\eta^+} > 100 \text{ GeV}$

One generation of η

Parameter space satisfying CDF-II W-mass anomaly, 95 GeV excess in the plane of $M_{\eta_R} - M_{\eta^+}$ with one generation of η .





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One generation of η

Parameter space satisfying CDF-II W-mass anomaly, 95 GeV excess in the plane of $M_{\eta_R} - M_{\eta^+}$ with one generation of η .



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One generation of η

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Two generation of η

Parameter space satisfying CDF-II W-mass anomaly, 95 GeV excess in the plane of $M_{\eta_R} - M_{\eta^+}$ with two generations of η





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Two generation of η

Parameter space satisfying CDF-II W-mass anomaly, 95 GeV excess in the plane of $M_{\eta_R} - M_{\eta^+}$ with two generations of η



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Three generation of η

Parameter space satisfying CDF-II W-mass anomaly, 95 GeV excess and SM Higgs signal strength in the plane of $M_{\eta_R} - M_{\eta^+}$ with three generations of η .





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Three generation of η

Parameter space satisfying CDF-II W-mass anomaly, 95 GeV excess and SM Higgs signal strength in the plane of $M_{\eta_R} - M_{\eta^+}$ with three generations of η .



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DM phenomenology

DM relic density as a function of DM mass



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DM phenomenology

DM relic density as a function of DM mass



 $M_{DM} \in [40, 85]$ GeV.

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Direct Detection of DM





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Direct Detection of DM



Summary

- We have explored the potential to identify a shared connection between the $(g - 2)_{\mu}$ anomaly, the CDF-II *W*-mass anomaly, and the CMS 95 GeV excess, within the context of the scotogenic $L_{\mu} - L_{\tau}$ model.
- The minimal model can effectively accounts for the (g 2)_μ and CDF-II W-mass anomalies, it falls short in generating the necessary diphoton signal strength for the 95 GeV scalar while being consistent with SM Higgs diphoton signal strength.
- We introduce two extra scalar doublets solely to contribute radiatively to diphoton decay of neutral scalars.
- These new scalar doublets can be motivated from neutrino mass point of view if we have only one right-handed neutrino.
- The DM is constrained in range \in [40, 85] GeV.
- This scotogenic $U(1)_{L_{\mu}-L_{\tau}}$ model can explain all these simultaneously.

Thank You



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3

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Backup

$$\begin{split} M_{\eta_R}^2 &= m_{\eta}^2 + \frac{v^2}{2}(\lambda_3 + \lambda_4 + \lambda_5) + \frac{v_1^2}{2}\lambda_{\eta\Phi_1} + \frac{v_2^2}{2}\lambda_{\eta\Phi_2} \\ M_{\eta_I}^2 &= m_{\eta}^2 + \frac{v^2}{2}(\lambda_3 + \lambda_4 - \lambda_5) + \frac{v_1^2}{2}\lambda_{\eta\Phi_1} + \frac{v_2^2}{2}\lambda_{\eta\Phi_2} \\ M_{\eta^+}^2 &= m_{\eta}^2 + \frac{v^2}{2}\lambda_3 + \frac{v_1^2}{2}\lambda_{\eta\Phi_1} + \frac{v_2^2}{2}\lambda_{\eta\Phi_2} \end{split}$$

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Backup

The loop functions involved in the calculation of $\Gamma(h_i \rightarrow \gamma \gamma)$ are given by

$$A_0 = -[\tau - f(\tau)]/\tau^2,$$

$$A_{1/2} = 2[\tau + (\tau - 1)f(\tau)]/\tau^2,$$

$$A_1 = -[2\tau^2 + 3\tau + 3(2\tau - 1)f(\tau)]/\tau^2,$$

where the function $f(\tau)$ is defined as

$$f(\tau) = \begin{cases} \arcsin^2 \sqrt{\tau} ; & \tau \le 1 \\ \\ -\frac{1}{4} \left[\log \frac{1 + \sqrt{1 - \tau^{-1}}}{1 - \sqrt{1 - \tau^{-1}}} - i\pi \right]^2 ; & \tau > 1 \end{cases}$$

Backup LFV

$$\operatorname{Br}(\mu \to e\gamma) = \frac{3(4\pi)^3 \alpha}{4G_F^2} (|A_{e\mu}^M|^2 + |A_{e\mu}^E|^2) \operatorname{Br}(\mu \to e\nu_{\mu}\overline{\nu_e}).$$

$$A^{M}_{e\mu} = \frac{-1}{(4\pi)^2} \sum_{k} (Y^*_{ek} Y_{\mu k} I^{++}_{k} + Y^*_{ek} Y_{\mu k} I^{+-}_{k})$$

$$A_{e\mu}^{E} = \frac{-i}{(4\pi)^{2}} \sum_{k} (-Y_{ek}^{*} Y_{\mu k} I_{k}^{-+} - Y_{ek}^{*} Y_{\mu k} I_{k}^{--})$$

$$I_{k}^{(\pm)_{1}(\pm)_{2}} = \int d^{3}X \frac{x(y + (\pm)_{1}z\frac{m_{e}}{m_{\mu}} + (\pm)_{2}\frac{M_{k}}{m_{\mu}})}{-xym_{\mu}^{2} - xzm_{e}^{2} + (1 - x)M_{\eta^{+}}^{2} + xM_{k}^{2}}$$



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