



Light Dirac neutrino portal dark matter with gauged $U(1)_{B-L}$ symmetry

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Overview



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Introduction and Motivation

Motivation

- Search for beyond Standard model (BSM) particles is one of the main objective in particle physics.
- BSM particles can have weak coupling, difficult to detect them in terrestrial experiments.
- Precision cosmology, with various current and upcoming cosmological experiments, provide excellent opportunity to study BSM physics.
- BSM particles in early universe can leave their signatures.
- These signatures can be detected via cosmological experiments like **Cosmic Microwave Background** experiments.
- Dark matter (DM) can be weakly (WIMP) interacting or Feebly interacting (FIMP).
- Nature of neutrino : Dirac or Majorana.

Motivation

- If Dirac, CMB experiment can help in detection prospect due to light nature.
- Define -

$$\rho_r = \rho_\gamma + \rho_{\nu_L} + \rho_{\text{BSM}} = \left(1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right) \rho_\gamma. \quad (1)$$

- N_{eff} \implies effective number of relativistic species.
- Incorporate all possible light BSM particles.
- $N_{\text{eff}} = N_{\text{eff}}^{\text{SM}} + \Delta N_{\text{eff}}$.
- For the Standard Model, $N_{\text{eff}} = N_{\text{eff}}^{\text{SM}} = 3.045$ [1606.06986].
- Current CMB bound (2σ) from Planck 2018 data,
 $N_{\text{eff}} = 2.99^{+0.34}_{-0.33} \implies \Delta N_{\text{eff}} = 0.284$ [1807.06209].
- Future expected CMB bound (2σ) from CMB-S4, $\Delta N_{\text{eff}} = 0.06$ [1907.04473].



The Model

The Model



- Here origin of DM is related to Dirac nature of neutrino.
- Dirac neutrino takes the role of mediating interactions with DM and SM bath.
- To realise, we consider a UV complete gauged B-L model.
- SM is extended by - ϕ_1, ϕ_2, ψ and ν_R with B-L charges 1, 3, 0, -1.



The Model

Scalar part

$$\mathcal{L}_s = (D^\mu \Phi)^\dagger D_\mu \Phi + (D^\mu \phi_1)^\dagger D_\mu \phi_1 + (D^\mu \phi_2)^\dagger D_\mu \phi_2 - V(\Phi, \phi_1, \phi_2), \quad (2)$$

$$D_\mu \Phi = \left(\partial_\mu - i\frac{g}{2} \tau_\sigma W_\mu^\sigma - i\frac{g'}{2} B_\mu \right) \Phi; \quad D_\mu \phi_1 = \left(\partial_\mu - ig_{BL} B'_\mu \right) \phi_1; \quad D_\mu \phi_2 = \left(\partial_\mu - i3g_{BL} B'_\mu \right) \phi_2,$$

$$\begin{aligned} V(\Phi, \phi_1, \phi_2) &= -\mu^2 (\Phi^\dagger \Phi) + \mu_1^2 (\phi_1^\dagger \phi_1) - \mu_2^2 (\phi_2^\dagger \phi_2) \\ &+ \lambda (\Phi^\dagger \Phi)^2 + \lambda_1 (\phi_1^\dagger \phi_1)^2 + \lambda_2 (\phi_2^\dagger \phi_2)^2 \\ &+ \lambda_{H\phi_1} (\Phi^\dagger \Phi) (\phi_1^\dagger \phi_1) + \lambda_{H\phi_2} (\Phi^\dagger \Phi) (\phi_2^\dagger \phi_2) + \lambda_{\phi_1\phi_2} (\phi_1^\dagger \phi_1) (\phi_2^\dagger \phi_2). \end{aligned} \quad (3)$$

Yukawa part

$$-\mathcal{L}_Y \supset Y_\nu \bar{L} \tilde{\Phi} \nu_R + \textcolor{red}{y_{\phi_1}} \bar{\psi} \phi_1 \nu_R + m_\psi \bar{\psi} \psi \quad (4)$$

The Model



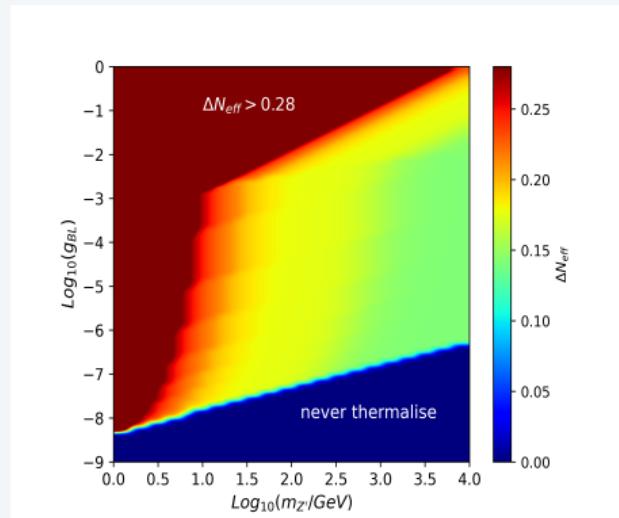
- ϕ_2 accquires vev and give mass to B-L gauge boson.
- ϕ_1 does not get vev, remain heavier than ψ making ψ stable.
- ψ is a fermionic dark matter.
- ϕ_1 , ψ and ν_R interact through yukawa coupling y_{ϕ_1} .
- Relevant parameters are Relevant parameters - m_{ϕ_1} , $m_{Z'}$, m_ψ , $\lambda_{H\phi_1}$, g_{BL} , y_{ϕ_1} .
- Depending on y_{ϕ_1} , we can have both FIMP or WIMP type DM.



FIMP type DM

FIMP type DM

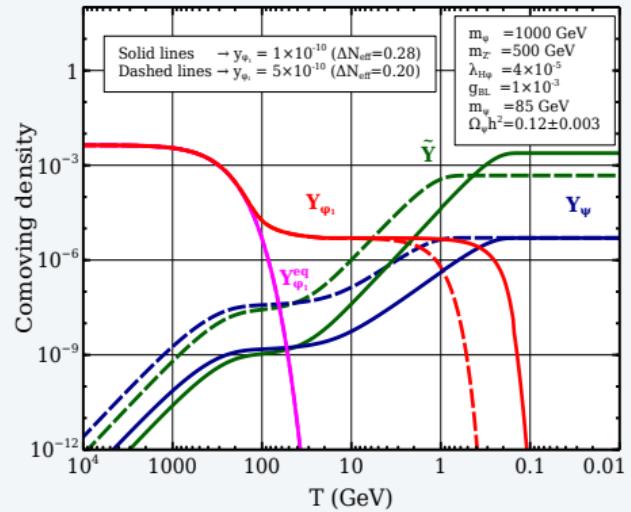
- Due to small yukawa, DM and ν_R are produced non-thermally.
- Thermal ν_R is also possible via $f\bar{f} \leftrightarrow \nu_R \bar{\nu}_R$ (resonant) via Z' gauge boson.
- $\Delta N_{\text{eff}} = \Delta N_{\text{eff}}^{\text{th}} + \Delta N_{\text{eff}}^{\text{non-th}}$.
- $\Delta N_{\text{eff}}^{\text{th}} = \frac{7}{8} \times 2 \times 0.027 \left(\frac{106.75}{g_*(T_{\text{dec}})} \right)^{4/3}$.



FIMP type DM

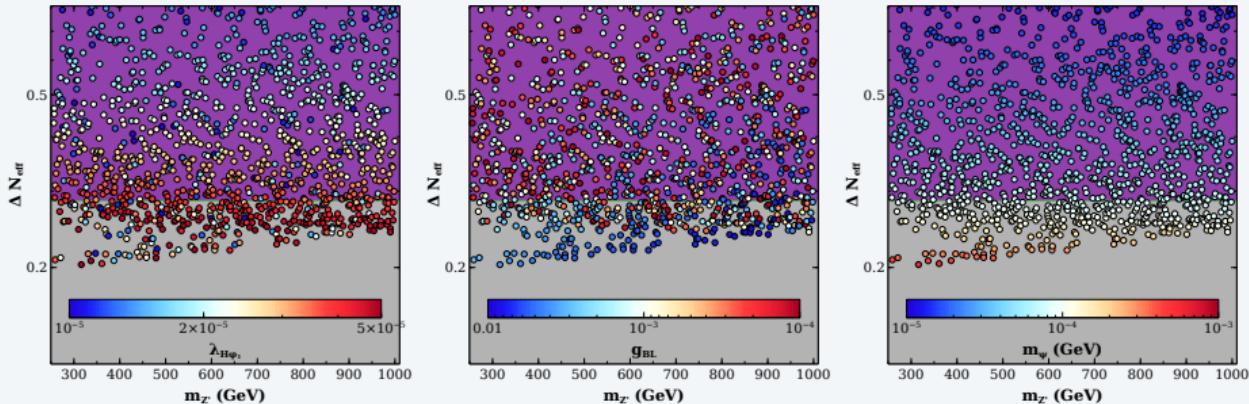


- Non - thermal production : ϕ_1 decays to DM, ψ and ν_R .
- Solve Boltzmann equation for $Y_{\phi_1} = \frac{n_{\phi_1}}{s}$, $Y_\psi = \frac{n_\psi}{s}$ and $Y_{\nu_R} = \frac{\rho_{\nu_R}}{s^{4/3}}$.
- ϕ_1 freeze-out is determined by : $\lambda_{H\phi_1}$ and g_{BL} .
- Decay width : y_{ϕ_1} .



FIMP type DM

- $250 \text{ GeV} < m_{Z'} < 1000 \text{ GeV},$ $10^{-5} < \lambda_{H\phi_1} < 5 \times 10^{-5},$
 $10^{-4} < g_{BL} < 10^{-2},$ $10 \text{ keV} < m_\psi < 1000 \text{ keV}.$



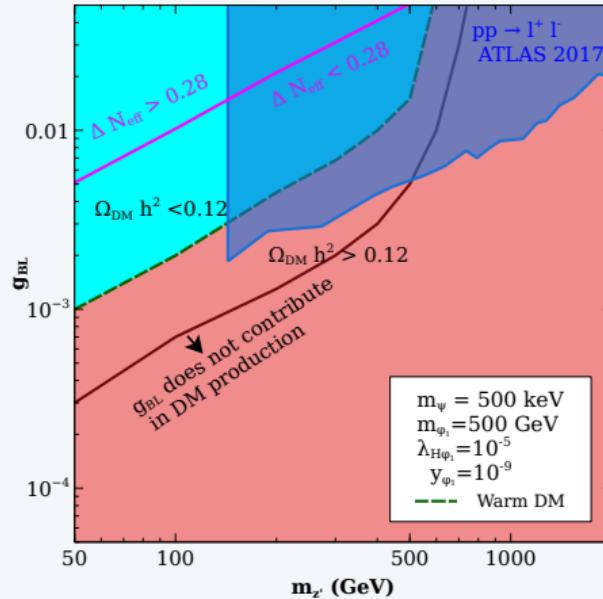
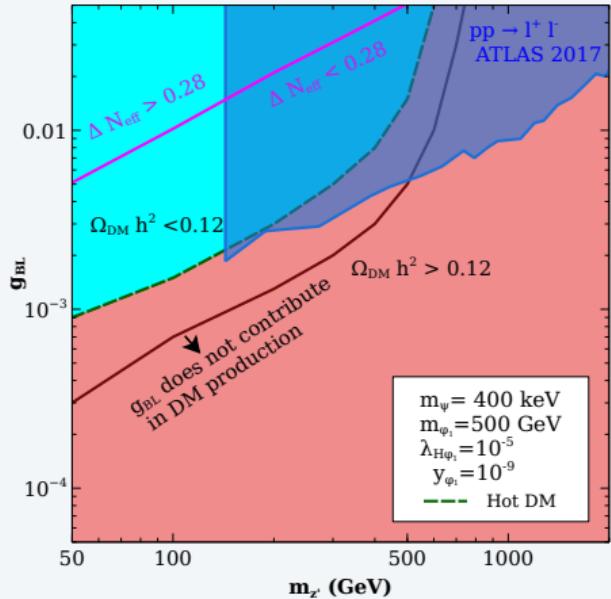
FIMP type DM : Structure formation

- Ability to form structure \implies Depend on FSL of DM.
- FSL \implies Distribution function \implies Production mechanism.
- Cold Dark Matter (CDM) $\iff \lambda_{\text{FSL}} < 0.01 \text{Mpc}$.
- Warm Dark Matter (WDM) $\iff 0.01 \text{Mpc} < \lambda_{\text{FSL}} < 0.1 \text{Mpc}$.
- Hot Dark Matter (HDM) $\iff \lambda_{\text{FSL}} > 0.1 \text{Mpc}$. [[1112.0330](#)]

$$\bullet \quad \lambda_{\text{fs}} = \int_{T_{\text{prod}}}^{T_{\text{eq}}} \frac{\langle v(T) \rangle}{a(T)} \frac{dt}{dT} dT, \quad \langle v(T) \rangle = \frac{\int \frac{p_1}{E_1} \frac{d^3 p_1}{(2\pi)^3} f_\psi(p_1, T)}{\int \frac{d^3 p_1}{(2\pi)^3} f_\psi(p_1, T)}.$$

	Parameters						$\Omega_{\text{DM}} h^2$	ΔN_{eff}	FSL(Mpc)
	m_{ϕ_1} (GeV)	$\lambda_{H\phi_1}$	y_{ϕ_1}	$m_{Z'}$ (GeV)	g_{BL}	m_ψ (keV)			
BP I	1000	5×10^{-5}	10^{-10}	500	0.001	126	0.12	0.251	3.09
BP II	500	5×10^{-5}	10^{-10}	500	0.001	233	0.12	0.210	1.32
BP III	1000	2×10^{-4}	10^{-9}	500	0.001	970	0.12	0.184	0.06
BP IV	500	2×10^{-4}	10^{-9}	500	0.001	938	0.12	0.184	0.05

FIMP type DM : Summary plots





WIMP type DM

WIMP type DM

- Here interaction between ϕ_1 , ψ and ν_R is large (large y_{ϕ_1}).
- Form a dark sector (DS) among themselves.
- DS interact with SM via $\phi_1 X \rightarrow \phi_1 X$, $\phi_1 Z' \rightarrow \phi_1 Z'$ and $\nu_R X \rightarrow \nu_R X$.
- T_{dec} can be found -

$$\frac{\Gamma_{\text{total}}}{\mathcal{H}} = \frac{1}{\mathcal{H}} [n_X^{\text{eq}} (\langle \sigma v \rangle_{\phi_1 X \rightarrow \phi_1 X} + \langle \sigma v \rangle_{\phi_1 Z' \rightarrow \phi_1 Z'}) + n_{\nu_R}^{\text{eq}} \langle \sigma v \rangle_{\nu_R \bar{\nu}_R \rightarrow X \bar{X}}] . \quad (5)$$

WIMP type DM

- After decoupling, DS evolves and have separate temperature.
- Define $\xi = \frac{T_{\text{DS}}}{T}$.

$$\frac{dY}{dx} = \frac{1}{2} \frac{\beta s}{Hx} \langle \sigma v \rangle_{\text{eff}} ((Y^{\text{eq}})^2 - Y^2) \quad (6)$$

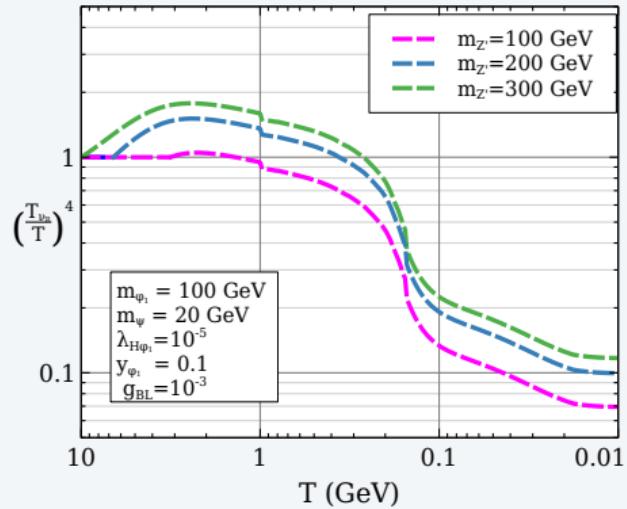
$$\frac{d\xi}{dx} = \frac{1}{x} \left(-\frac{1}{2} \frac{\beta x^4 s^2}{4\alpha \xi^3 H m_{\phi_1}^4} \langle E \sigma v \rangle_{\text{eff}} ((Y^{\text{eq}})^2 - Y^2) - (\beta - 1)\xi \right).$$

$$Y = Y_{\phi_1} + Y_{\psi}.$$

WIMP type DM



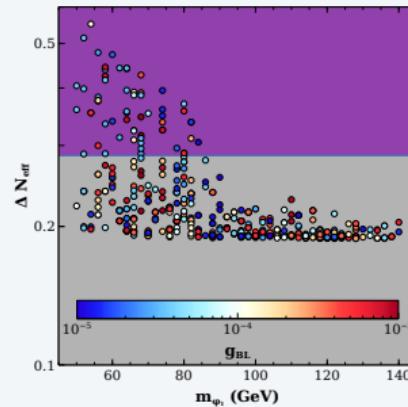
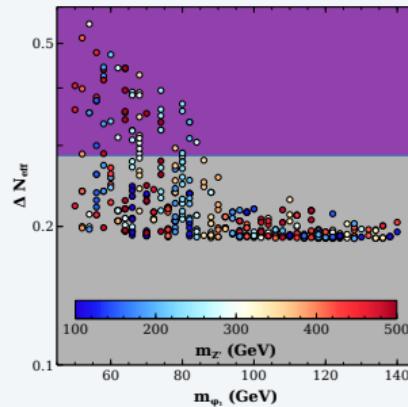
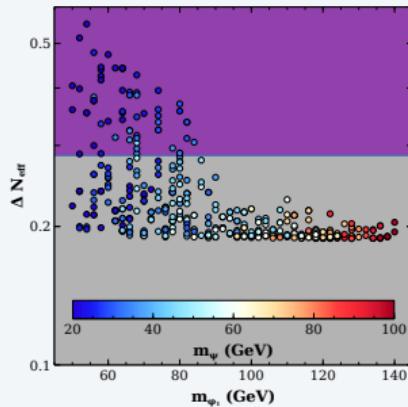
- We find increase in T_{DS} compared to T .
- Reason : annihilation of heavier DS particles to lighter DS particles.
- Gives strong constraint from ΔN_{eff} .



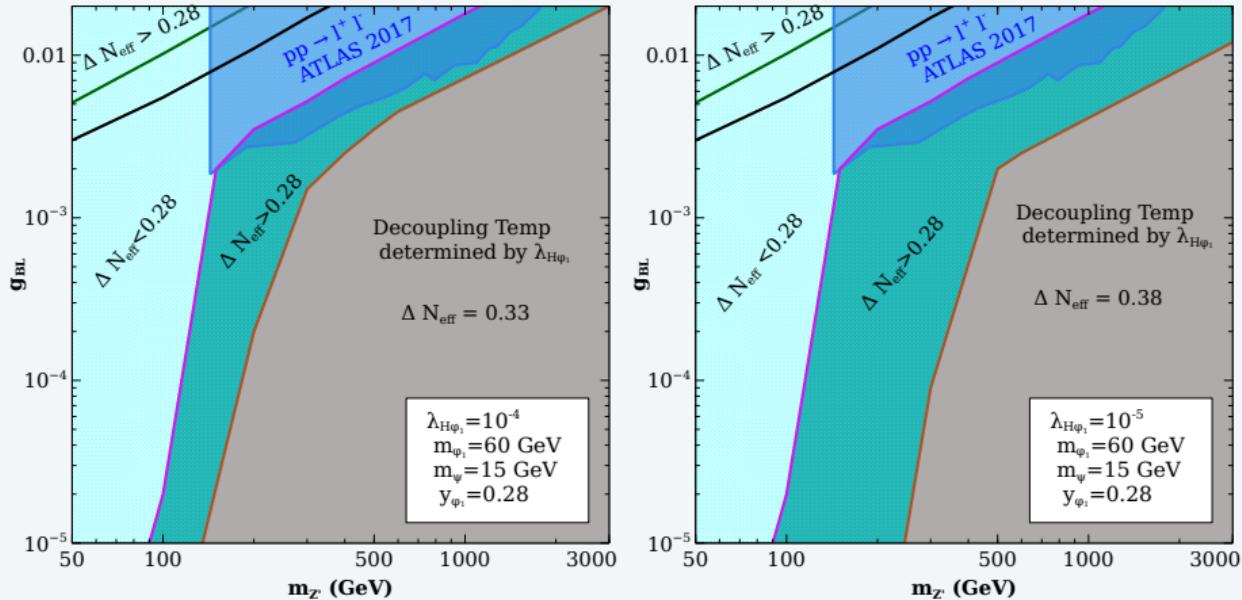
WIMP type DM



- $50\text{GeV} < m_{\phi_1} < 150\text{GeV}$,
 $10^{-5} < \lambda_{H\phi_1} < 10^{-2}$,
- $100\text{GeV} < m_{Z'} < 500\text{GeV}$,
 $10^{-5} < g_{BL} < 10^{-3}$,
- $10\text{GeV} < m_\psi < 100\text{GeV}$
 $0.2 < y_{\phi_1} < 0.3$.



WIMP type DM : Summary plots



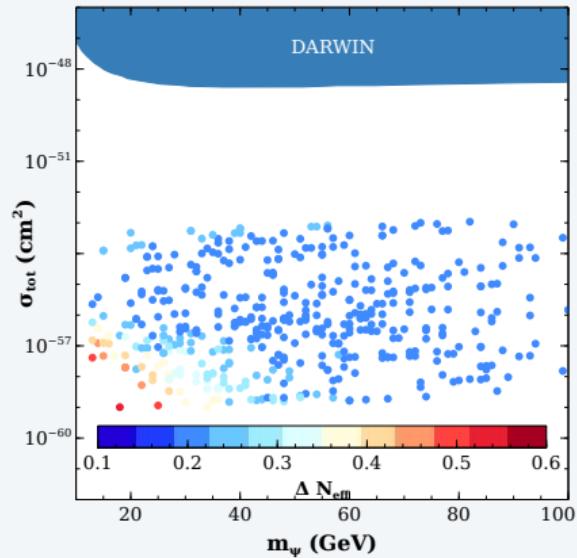
WIMP type DM : Direct Detection



- No tree level diagram.
- One loop - Higgs mediated and Z' mediated.
- Total DM-nucleon cross-section

$$\sigma_{\text{tot}} = \frac{1}{\pi} \frac{m_N^2 m_\psi^2}{(m_N + m_\psi)^2} \left[\frac{m_N}{v} \frac{1}{m_h^2} g_{\psi \bar{\psi} h} f_N + \frac{g_{BL}}{3} \frac{1}{m_{Z'}^2} g_{\psi \bar{\psi} Z'} f_{Z'} \right]^2.$$

- Out of reach from future DD experiment. [1606.07001]





Conclusion



Conclusion

- We study a possible UV completion of light DNPDM.
- Due to gauge $B - L$ interaction, detectable contribution to ΔN_{eff} .
- We study both thermal and non-thermal production of DM.
- We find that cosmological constraint (from structure formation and CMB) put stringent bound to the parameter space of the model: stronger than direct detection bound.



Thank you for your attention

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