





# BUBBLE-ASSISTED LEPTOGENESIS

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#### **BARYON ASYMMETRY OF THE UNIVERSE**

$$Y_B \equiv \frac{n_B - n_{\bar{B}}}{s} \approx 0.9 \times 10^{-10}$$

Planck TT, TE, EE+lowE+lensing +BAO		
$\Omega_{ m b}h^2$	$0.02237 \pm 0.00015$	$0.02242 \pm 0.00014$
$\Omega_{ m c} h^2$	$0.1200 \pm 0.0012$	$0.1193 \pm 0.0009$
$100 heta_{ m MC}$	$1.0409 \pm 0.0003$	$1.0410 \pm 0.0003$
$n_{ m s}$	$0.965 \pm 0.004$	$0.966 \pm 0.004$
au	$0.054 \pm 0.007$	$0.056 \pm 0.007$
$\ln(10^{10}\Delta_{\mathcal{R}}^2)$	$3.044 \pm 0.014$	$3.047 \pm 0.014$
1	0.674 + 0.005	0.677   0.004
n	$0.074 \pm 0.005$	$0.677 \pm 0.004$
$\sigma_8$	$0.811 \pm 0.006$	$0.810 \pm 0.006$
$\Omega_{ m m}$	$0.315 \pm 0.007$	$0.311 \pm 0.006$
$\Omega_{\Lambda}$	$0.685 \pm 0.007$	$0.689 \pm 0.006$



#### LEPTOGENESIS

Fukugita, Yanagida, 1986

- Light LH neutrinos with heavy RH neutrinos:
- RHN decays to produce asymmetry in lepton number which converts to baryon asymmetry by EW sphaleron process.
- Dynamical generation of baryon asymmetry:
  - ✓ L and B + L violation:  $M_N$  and EW sphaleron
  - ✓ **C** & **CP** violation:  $\Im(\overline{Y}^2 Y^2) \neq 0$
  - ✓ Out of Equilibrium: decay process

$$\mathcal{L} = Y_D \, \bar{l} \, \overline{H}N + \frac{1}{2} M_N \overline{N^c}N + h. \, c$$

 $\Rightarrow M_{\nu} = Y_D M_N^{-1} Y_D^T v_{EW}^2$ 

$$\epsilon_N \equiv \frac{\Gamma(N \to lH) - \Gamma(N \to \overline{lH})}{\Gamma(N \to lH) + \Gamma(N \to \overline{lH})} \neq 0$$

 $Y_B = Y_N^{\text{eq}} \epsilon_N \kappa_{\text{sph}} \kappa_{\text{eff}}$ 

#### **EFFICIENCY FACTOR**

- RHN decay becomes active at  $T \sim M_N$  to produce the asymmetry.
- Inverse decay could remain in equilibrium till  $T \ll M_N$  to wash out the produced asymmetry (Strong washout regime:  $K \gg 1$ ).

$$\Gamma_{N} = \frac{Y_{D}^{2} M_{N}}{8\pi} \qquad \widetilde{m}_{\nu} \equiv \frac{Y_{D}^{2} v_{EW}^{2}}{M_{N}}$$
$$K \equiv \frac{\Gamma_{N}}{H(M_{N})} = \frac{\widetilde{m}_{\nu}}{\text{meV}}$$

 $K \approx (9, 50)$  for  $\tilde{m}_{\nu} = (8.7, 50)$  meV

$$\kappa_{\rm eff} \approx \frac{1}{2K \, {\rm ln}(K)} \sim (0.027, 0.0026)$$

#### **DAVIDSON-IBARRA BOUND**

#### • For non-degenerate RHNs,

$$\epsilon_N \lesssim \frac{3}{8\pi} \frac{M_N m_{\nu_3}}{v_{EW}^2}$$

$$Y_B \approx 10^{-10} \Rightarrow M_N \gtrsim 3 \cdot 10^{11} \text{ GeV}\left(\frac{0.0026}{\kappa_{\text{eff}}}\right) \left(\frac{0.05 \text{eV}}{m_{\nu_3}}\right)$$

• It can be relaxed significantly in Bubble-assisted Leptogenesis.

#### SEESAW & GW

• Gravitational Wave from cosmic strings as a probe of seesaw associated with  $U(1)_{B-L}$  breaking at  $v_{B-L}$ :



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-4

#### GW FROM BUBBLE-ASSISTED LEPTOGENESIS

•  $U(1)_{B-L}$  breaking could be of strong first-order, which enhances the efficiency of leptogenesis and generates observable gravitational wave background:



#### **LEPTOGENESIS IN FOPT**

How is the vanilla leptogenesis modified by the bubble dynamics?

2305.10759

Huang, Xie, 2206.04691 Sh Dasgupta, et.al., 2206.07032 Ba

Shuve, Tamarit, 1704.01979 Baldes, et.al., 2106.15602

Study strong FOPT in classically scale-invariant models.

#### SCALE-INVARIANT $U(1)_{B-L}$ MODEL

• Setup for strong FOPT of  $U(1)_{B-L}$  breaking:

 $-\mathcal{L} \ni \frac{1}{2} y_N \Phi NN + Y_D H lN + h. c. + V(\Phi, T)$   $V(\Phi, T) = \frac{1}{4} \lambda_{\phi} \phi^4 + V_{CW}(\phi) + V_T(\phi) \qquad \Phi = \frac{\phi}{\sqrt{2}}$   $V_{CW}(\phi) = \sum_i (-1)^{2s_i} g_i \frac{m_i^2(\phi)}{64\pi^2} \Big[ \ln\left(\frac{m_i^2(\phi)}{\mu^2}\right) - c_i \Big] \qquad m_N^2(\phi) = \frac{1}{2} y_N^2 \phi^2$   $m_{Z_{B-L}}^2(\phi) = 4g_{B-L}^2 \phi^2$   $V_T(\phi) = \pm \frac{g_i}{2\pi} T^4 J_{B,F}\left(\frac{m_i^2(\phi)}{T^2}\right) \qquad m_{\phi}^2(\phi) = 6\lambda_{\phi} \phi^2$ 

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#### **BUBBLE DYNAMICS**



• Nucleation rate per unit volume per unit time:  $\Gamma_{\text{nuc}}(T) \approx T^4 \left(\frac{S_3}{2\pi T}\right)^{\frac{3}{2}} e^{-\frac{S_3}{T}} = H(T)^4 \text{ at } T = T_{\text{nuc}} \qquad H^2 = \frac{\rho(T) + \Delta V}{3M_P^2}$ 

• Time scale between  $T_{nuc}$  and  $T_{PT}$ :

$$\Delta t_{\rm PT} = -\left(\frac{d(S_3/T)}{dt}\right)_{T_{\rm nuc}} \qquad \beta_{\rm PT}^{-1} \equiv \frac{\Delta t_{\rm PT}}{t_{\rm reh}} \ll 1$$

• Strength of FOPT:  $\alpha_n \equiv \frac{\Delta V}{\rho(T_{\text{nuc}})} = \frac{\rho(T_{\text{reh}}) - \rho(T_{\text{nuc}})}{\rho(T_{\text{nuc}})} \gg 1$ 

#### **SUDDEN MASS GAIN & PENETRATION RATE**

#### 2010.02590



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## $U(1)_{B-L}$ MODLE AND BUBBLE PARAMETERS



#### LEPTOGENESIS INSIDE BUBBLE



Require  $\gamma_w T_{nuc} \gg M_N \gg T_{reh} > T_{nuc}$ 

#### LEPTOGENESIS INSIDE BUBBLE

- RHNs inside bubble decay immediately:
- Unavoidable depletion  $\kappa_{dep}$  by annihilation  $\Gamma(NN \rightarrow \phi \phi) > H$ .
- Partial washout  $\kappa_{wash}$  by the inverse decay.
- Dilution  $\kappa_{dil}$  by the reheat:

$$\frac{t_{\rm decay}}{t_{\rm PT}} \approx 0.1 \left(\frac{5}{M_N/T_{\rm reh}}\right)^2 \left(\frac{0.05 \text{ eV}}{\widetilde{m}_v}\right) \left(\frac{\beta_{\rm PT}}{100}\right)$$

$$n_{N}^{0} = \kappa_{\text{pen}} n_{N}^{eq} (M_{N} = 0)$$
  

$$\kappa_{\text{pen}} \rightarrow 1 \text{ vs. } \kappa_{\text{dil}} \downarrow$$
  

$$\kappa_{\text{dil}} = \left(\frac{T_{\text{nuc}}}{T_{\text{reh}}}\right)^{3} = (1 + \alpha_{n})^{-\frac{3}{4}}$$
  

$$\kappa_{\text{pen}} \rightarrow 1 \text{ vs. } \kappa_{\text{dil}} \downarrow$$
  
for  $\alpha_{n} \uparrow$ 

#### **DECAY, INVERSE-DECAY & ANNIHILATION**



#### **SCHEMATIC BEHAVIOR**



 $Y_B^{\text{bubble}} = \kappa_{\text{pen}} Y_N^{eq}(0) \epsilon_N \kappa_{\text{sph}} \kappa_{\text{dep}} \kappa_{\text{wash}} \kappa_{\text{dil}}$ 

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#### **BUBBLE VS. VANILLA**



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#### **GRAVITATIONAL WAVE SIGNAL**



### CONCLUSION

• Leptogenesis in FOPT can allow for a strong departure from thermal equilibrium:

- Conventional washout can be circumvented.
- $\checkmark$  New annihilation channel  $NN \rightarrow \phi \phi$  opens up to deplete the asymmetry.
- $\checkmark$  Dilution from reheating can be sizable.
- In the strong washout regime, the bubbles help Leptogensis to enhance the efficiency upto  $3 \sim 30$  for  $M_N = 10^8 \sim 10^{10}$  GeV.

• Observable GW signals are predicted for  $M_N \lesssim 5 \cdot 10^9$  GeV.