

Should we consider scalar extensions seriously?

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Plan of the talk:

To give an affirmative answer to the question.

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- A possible cold dark matter candidate
 - A possible source of CP violation
 - A possible way to make the EW phase transition first-order
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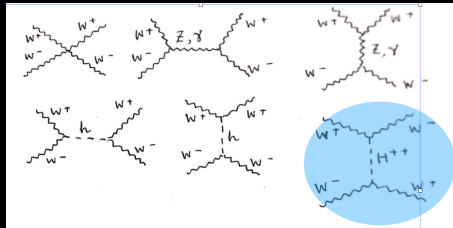
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SM: A renormalisable gauge theory

- The Higgs mechanism makes the EW gauge bosons massive but keeps the theory renormalisable

(Veltman and 't Hooft)

- This means that the scattering of gauge bosons does not violate unitarity, the amplitude does not grow monotonically with \sqrt{s} . Consider $W_L W_L \rightarrow W_L W_L$:





$$\mathcal{M} = 16\pi \sum_{ell} (2\ell + 1) a_\ell(s) P_\ell(\cos \theta),$$
$$a_\ell = A(q/M_W)^4 + B(q/M_W)^2 + C,$$

where q is the W momentum in the CM frame. Only $\ell = 0, 1, 2$ are relevant.

- The A term cancels with 3- and 4-point gauge vertices but cancellation of B needs the Higgs boson

(Lee, Quigg, and Thacker)

Notation

- SM Higgs: h

Any other scalar of mass M : h_M or H_M

A SM-like scalar (identical couplings with fermions and gauge bosons) but at a mass M : ϕ_M

Coupling strength of particles i, j, k : g_{ijk}

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$$\kappa_X = \frac{g_{h_M X \bar{X}}}{g_{h X \bar{X}}}$$

$$\mu_X = \frac{\sigma(pp \rightarrow h_M \rightarrow X \bar{X})}{\sigma(pp \rightarrow \phi_M \rightarrow X \bar{X})}$$

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Theoretical constraints on extra scalars

- $V(\Phi_1, \Phi_2, \dots, \Phi_n)$ must be a gauge and Lorentz scalar
- All dimensionless couplings in V should better be perturbative
- Must be bounded from below along every direction in the field space

$$\begin{aligned} V(\Phi_1, \Phi_2) &= \lambda_1 \left(\Phi_1^\dagger \Phi_1 \right)^2 + \lambda_2 \left(\Phi_2^\dagger \Phi_2 \right)^2 + \lambda_3 \left(\Phi_1^\dagger \Phi_1 \right) \left(\Phi_2^\dagger \Phi_2 \right) \\ &= \left(\sqrt{\lambda_1} \Phi_1^\dagger \Phi_1 - \sqrt{\lambda_2} \Phi_2^\dagger \Phi_2 \right)^2 \\ &\quad + \underbrace{\left(\lambda_3 + 2\sqrt{\lambda_1 \lambda_2} \right)}_{>0} \Phi_1^\dagger \Phi_1 \Phi_2^\dagger \Phi_2 \end{aligned}$$

- Partial wave unitarity must be respected in all scattering channels:
 $\Re(a_\ell) \leq 1/2$

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Theoretical constraints on extra scalars

Custodial symmetry

If an $SU(2)$ gauge symmetry undergoes SSB in a way that the neutral and the charged gauge bosons remain mass-degenerate at the tree-level, there is an $SU(2)$ symmetry in the gauge boson mass terms and is present in the pure gauge terms even in the broken phase.

$$\begin{aligned}\rho_{\text{tree}} &= \frac{M_W^2}{M_Z^2 \cos^2 \theta_W} = 1, && \text{(Broken by fermions)} \\ \rho_0 &= 1.00038 \pm 0.00020, && \text{(Includes top contrib., SM=1)}\end{aligned}$$

Condition

$$\begin{aligned}\sum c [T(T+1) - n^2] v^2 &= 2 \sum n^2 v^2, \\ c &= \begin{cases} 1 & \text{if the multiplet is complex,} \\ \frac{1}{2} & \text{if the multiplet is real.} \end{cases}\end{aligned}$$

T = isospin of the multiplet, n = T_3 -value of the neutral. Identically satisfied for any singlet ($T = 0$) and doublet ($T = \frac{1}{2}$)

Types of Custodial Symmetry

- CS is a symmetry of the entire \mathcal{L} . No correction to ρ is generated even in loops¹, disregarding the fermions. Scalar custodial multiplets are mass-degenerate.

Example: SM, gauge singlets, constrained 2HDM, canonical Georgi-Machacek.

- The kinetic terms after SSB respect CS but the scalar potential does not. The scalar custodial multiplets are not degenerate.

Example: Most general 2HDM, Generalised Georgi-Machacek.

(AK, Mondal, Pal, PRD 2022)

¹In SU(2) theories, no $\times U(1)$.

Sum rules

[Gunion, Haber, Wudka, PRD 1991]

- ⇒ The underlying theory is renormalisable, so all scattering amplitudes must respect the unitarity bound
- ⇒ Consider $WW \rightarrow WW$, which can be mediated through both neutral and doubly-charged scalars
- ⇒ Similarly, $WZ \rightarrow WZ$ goes through neutral and singly-charged scalars

$$\underbrace{g^2 (4M_W^2 - 3 \cos^2 \theta_W M_Z^2)}_{\approx g^2 M_W^2} = \underbrace{\sum_k \eta_k g_{WW}^2 \phi_k^0 - \sum_k \eta_k g_{WW}^2 \phi_k^{++}}_{g^2 M_W^2 \text{ in SM; } \eta = +1(-1) \text{ for CP-even(odd)}},$$

$$\frac{g^2 \cos^2 \theta_W M_Z^4}{M_W^2} = \sum_k \eta_k g_{WW} \phi_k^0 g_{ZZ} \phi_k^0 - \sum_k \eta_k g_{WZ}^2 \phi_k^+.$$

- LHC tells us that $g_{WW h_{125}} \approx g_{WW h}$, saturating the first sum rule.
- So any other neutral scalar with a significant coupling to WW definitely indicates a doubly charged scalar in the spectrum.
- Similarly, any neutral scalar with a significant coupling to both WW and ZZ indicates at least one singly charged scalar.

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Any evidence, or empty talks?

- 1 Only h_{125} so far, with properties exactly that of h , within experimental accuracy : $h_{125} \approx h$ (CMS, Nature 2022)

$$\kappa_W = 1.02 \pm 0.08, \quad \kappa_Z = 1.04 \pm 0.07, \quad \kappa_t = 1.01_{-0.10}^{+0.11}, \quad \kappa_b = 0.99_{-0.16}^{+0.17}$$

- 2 Beware: this can be a conspiracy, like 2HDM in alignment limit.

- 3 Are there any new scalars?

— None announced, but *there are indications*, which I'll talk about
— I'll use only published numbers and plots from either ATLAS or CMS, no unofficial combinations (*i.e.*, done by the theoreticians)

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- When should one feel excited, and what is a discovery?

3σ is 99% CL, encouraging, 5σ is a discovery

But be careful, only global significances, not local (LEE)

- For every signal, the p -value gives the probability with which the null hypothesis is rejected

The higher the σ , the smaller the p -value ($5\sigma \Rightarrow p \approx 3 \times 10^{-7}$)

Combination is tricky, it is not just $p = p_1 p_2$ and so on

For example, to combine two results with p_1 and p_2 :

$$p = p_1 p_2 [1 - \ln(p_1 p_2)]$$

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h_{95}

CMS [1811.08459, PLB; 1803.06553, JHEP; CMS-PAS-HIG-21-001]

$$\mu_{\gamma\gamma} = \frac{\sigma(pp \rightarrow h_{95} \rightarrow \gamma\gamma)}{\sigma(pp \rightarrow \phi_{95} \rightarrow \gamma\gamma)}$$

$$= 0.6 \pm 0.2,$$

$$\mu_{\tau^+\tau^-} = \frac{\sigma(pp \rightarrow h_{95} \rightarrow \tau^+\tau^-)}{\sigma(pp \rightarrow \phi_{95} \rightarrow \tau^+\tau^-)}$$

$$= 1.2 \pm 0.5,$$

- ggF works, strong $t\bar{t}h_{95}$
- Stronger destructive interference between top and W loop

LEP [hep-ex/0306033, PLB]

$$\mu_{b\bar{b}} = \frac{\sigma(e^+e^- \rightarrow Zh_{95} \rightarrow Zb\bar{b})}{\sigma(e^+e^- \rightarrow Z\phi_{95} \rightarrow Zb\bar{b})}$$

$$= 0.117 \pm 0.057.$$

- Either ZZh_{95} or $b\bar{b}h_{95}$ is suppressed

Combined significance = 3.9σ

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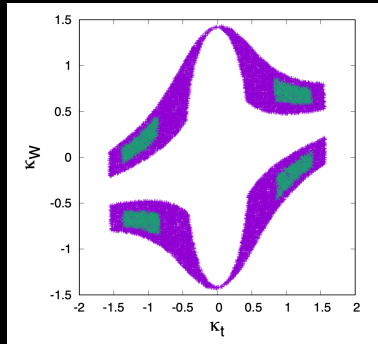
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Let's start with a 2HDM Type-I, so that $\kappa_t = \kappa_\tau$

Type-I: one doublet couples to all fermions, the other to none. Talk by J. Song yesterday

κ_t - κ_W plane with $\mu_{\tau\tau}$ and $\mu_{\gamma\gamma}$ (ggF:VBF = 90:10)



$\kappa_W^{\max.}(h_{95}) = 0.34 (1\sigma), 0.51 (2\sigma)$, without H^{++} , still consistent

H_{650}

CMS-PAS-HIG-20-016

$H_{650} \rightarrow W^+ W^-$, both ggF and VBF productions looked at

CMS-PAS-HIG-21-011

$H_{650} \rightarrow h_{95} h_{125} \rightarrow b\bar{b}\gamma\gamma$

ATLAS, 2103.01918 (JHEP); 2009.14791 (EPJC)

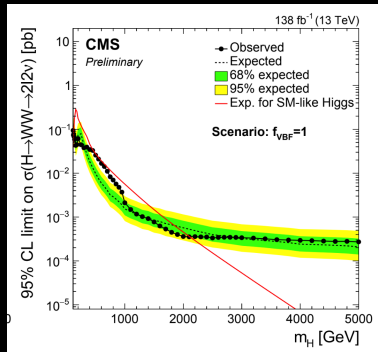
$H_{650} \rightarrow ZZ \rightarrow 4\ell$

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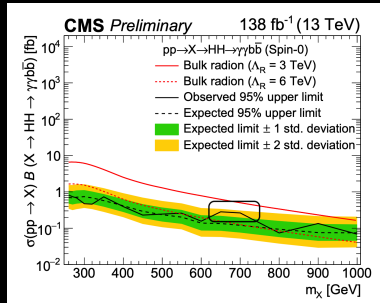
Scenario	Mass [GeV]	ggF cross sec. [pb]	VBF cross sec. [pb]	Local signi. [σ]	Global signi. [σ]
SM f_{VBF}	800	0.16	0.057	3.2	1.7 ± 0.2
$f_{VBF} = 1$	650	0.0	0.16	3.8	2.6 ± 0.2
$f_{VBF} = 0$	950	0.19	0.0	2.6	0.4 ± 0.6
floating f_{VBF}	650	2.9×10^{-6}	0.16	3.8	2.4 ± 0.2

[CMS-PAS-HIG-20-016]

Local: 3.8σ , Global: 2.6σ LE factor = p -value (global)/ p -value (local) ≈ 64



$$H_{650} \rightarrow h_{95} h_{125}$$

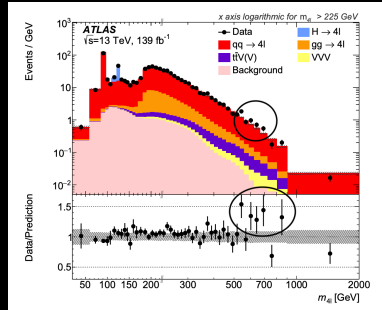
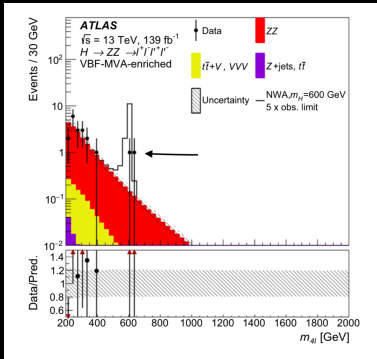


[CMS-PAS-HIG-21-011]

0.90 – 0.04 fb whereas the expected limits are 0.79 – 0.05 fb, with the considered mass ranges in m_χ and m_Y . The largest deviation from background-only hypothesis with local (global) significance of 3.8 (2.8) standard deviations is observed for $m_\chi = 650$ GeV and $m_Y = 90$ GeV. The HH limits are compared with predictions in the warped extra

Local: 3.8 σ , Global: 2.8 σ LE factor \approx 30

$H_{650} \rightarrow ZZ$



[ATLAS, EPJC 2021, 2009.14791; JHEP 2021, 2103.01918]

production separately. For the ggF production, the maximum deviation is for a signal mass hypothesis around 240 GeV, with a local significance of 2.0 standard deviations and a global significance of 0.5 standard deviation. For the VBF production, the maximum deviation is for a signal mass hypothesis around 620 GeV, with a local significance of 2.4 standard deviations and a global significance of 0.9 standard deviation.

H_{650} in summary

- Indicated by both CMS and ATLAS, a broad resonance
 - Definitely through VBF (WW and ZZ fusion), possibly also through ggF
 - Decays, possibly, to WW , $h_{95}h_{125}$, and to ZZ but suppressed
 - $t\bar{t}$ mode not yet observed, but background is huge
- $H_{650} W^+ W^-$ coupling is large:

$$\kappa_W = WW/SM = 0.85 \pm 0.15, \quad \kappa_Z = ZZ/SM = 0.40 \pm 0.15$$

(This is only a lower limit)

Combined significance = 4.2σ

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This is the most conservative number

3.6 σ

- These two plus ATLAS without paying the LE price

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■

$$\frac{g_{W^+W^+H^{--}}^2}{g_{hWW, \text{SM}}^2} = 1.30 \pm 0.57, \quad \frac{g_{W^+ZH^-}^2}{g_{hWW} g_{hZZ}} = 0.60 \pm 0.42,$$

■ This leads to

$$\Gamma(H^{++} \rightarrow W^+W^+) = 16 \pm 5 \text{ GeV}, \quad \Gamma(H^+ \rightarrow W^+Z) = 13 \pm 4 \text{ GeV}.$$

■ This is only with h_{125} and H_{650} . Stronger $W^+W^+H^{--}$ if we consider h_{95} too

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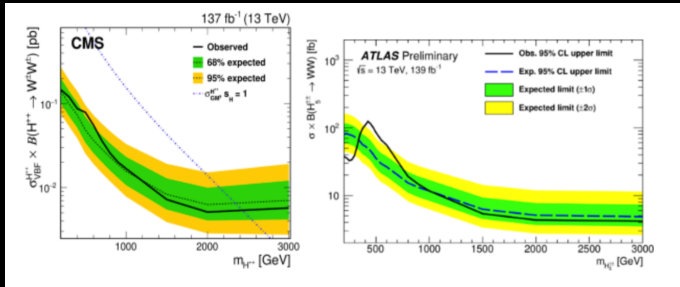
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Enter the charged scalars

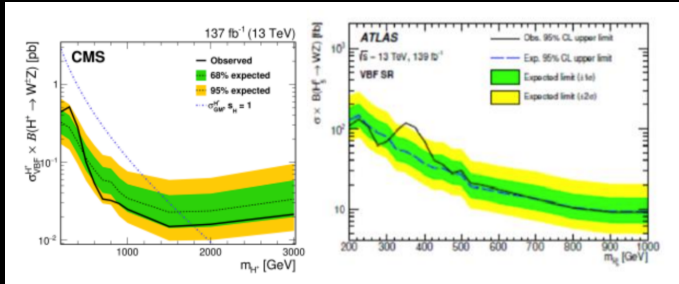
- Strong $H^{++}W^-W^- \implies M_{H^{++}} > 2 \text{ TeV}$ with 100% BR (CMS)
- Anything less means $H^{++} \rightarrow h_1^+h_2^+, h^+W^+$



[CMS, EPJC 2021, 2104.04762; ATLAS-CONF-23-023]

There is an indication of H^{++} at around 450 GeV:

Combined significance = 2.6σ



[CMS, EPJC 2021, 2104.04762; ATLAS, 2207.03925]

H^+ at around 375 GeV

Combined significance = 2.7σ

Final stock taking: What they can and cannot be

- h_{95} and H_{650} indications are strong and from more than one collab. There are several weak indications too, some of them we discussed
- H_{650} couples strongly to W^+W^-
Extensions with only singlets and/or doublets are out²

²As the data stands.

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2 custodial singlets, one triplet, one 5-plet, members of multiplets mass-degenerate in minimal version, only custodial singlets couple to fermions

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Also out if $H_{650}WW > H_{650}ZZ$, as indicated

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- *** After h_{125} , the combined significance of H_{650} passes 4σ
 - It has a strong coupling to WW and ZZ , indicates H^{++}
 - Whether it couples to fermions is not yet clear
- ** h_{95} is also close to 4σ
- * There are indications for H^+ and H^{++} too, none above 3σ at present. These are all global significances
 - We should wait for more data and an official combination, but if the trend continues ...
 - Extension by only $SU(2)$ singlets and doublets are ruled out, even if H^{++} hint goes away
 - There must be *at least* $SU(2)$ triplet scalars, but more than one such triplets to preserve custodial symmetry \Rightarrow something like Georgi-Machacek?

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Thank you for your patience

Backup slides

Backup: Combination of p -values

- p -value is *not* the probability that the null hypothesis H_0 is true. But it tells you at what CL you can reject H_0 .
- Be careful while combining two p -values, it is not just $p_* = p_1 p_2$, as p_* is not uniformly distributed over the unit interval (0:1).
- Suppose there are k independent tests with p_1, p_2, \dots, p_k , and let $p_* = p_1 p_2 \dots p_k$.
- Assuming H_0 , compute the probability $R_k(\rho)$ such that $p_* \leq \rho$.
- Quote $R_k(p_*)$ as the combined p -value.

Backup: Combination of p -values

- $k = 2$:

$$R_2(\rho) = \rho(1 - \ln \rho).$$

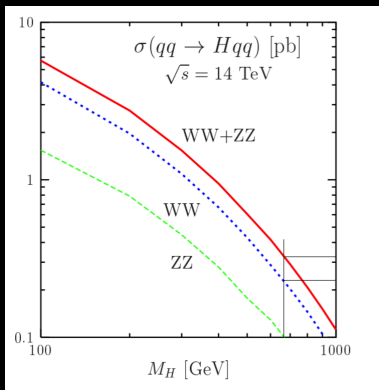
- For $k \geq 2$:

$$R_k(\rho) = \rho + \int_{\rho}^1 R_{k-1}(\rho/p) dp,$$

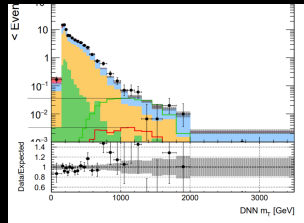
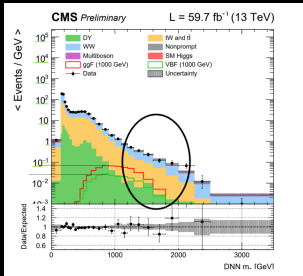
so that

$$R_3(\rho) = \rho \left\{ 1 - \ln \rho + \frac{1}{2} (\ln \rho)^2 \right\}.$$

Backup: H_{650} κ -values



Compare SM-like production \times -sec and the CMS number, goes as κ_W^2
 $\sqrt{160/220} = 0.85$



ggF has a large contamination from VBF: 7.0×10^{-2} and 2.5×10^{-2}

VBF is pure but the efficiency is less: 3.1×10^{-3} and 3.5×10^{-2}