

# INCLUSIVE DECAYS: SYNERGIES BETWEEN LATTICE AND CONTINUUM?

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# LATTICE VS HQE

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- Lattice can calculate  $\Gamma(B \rightarrow X_c l \bar{\nu}_l)$  and extract  $|V_{cb}|^{\text{inc}}$ .
- But before reaching 1-2% precision, **how do we use lattice** to improve the precision in the predictions obtained with the Heavy Quark Expansion (HQE)?
- The extraction of HQE parameters  $\mu_\pi^2, \mu_G^2, \rho_D^3, \rho_{LS}^3$  are relevant not only for  $\Gamma(B \rightarrow X_c l \bar{\nu}_l)$  and  $|V_{cb}|$  but also for  $|V_{ub}|, \tau_B, \Gamma(B \rightarrow X_s \gamma), \Gamma(B \rightarrow X_s l \bar{l})$  etc.
- Which quantities should be **calculated on the lattice** to improve the precision in  $\mu_\pi^2, \mu_G^2, \rho_D^3, \rho_{LS}^3, \dots?$

# POSSIBLE STRATEGIES

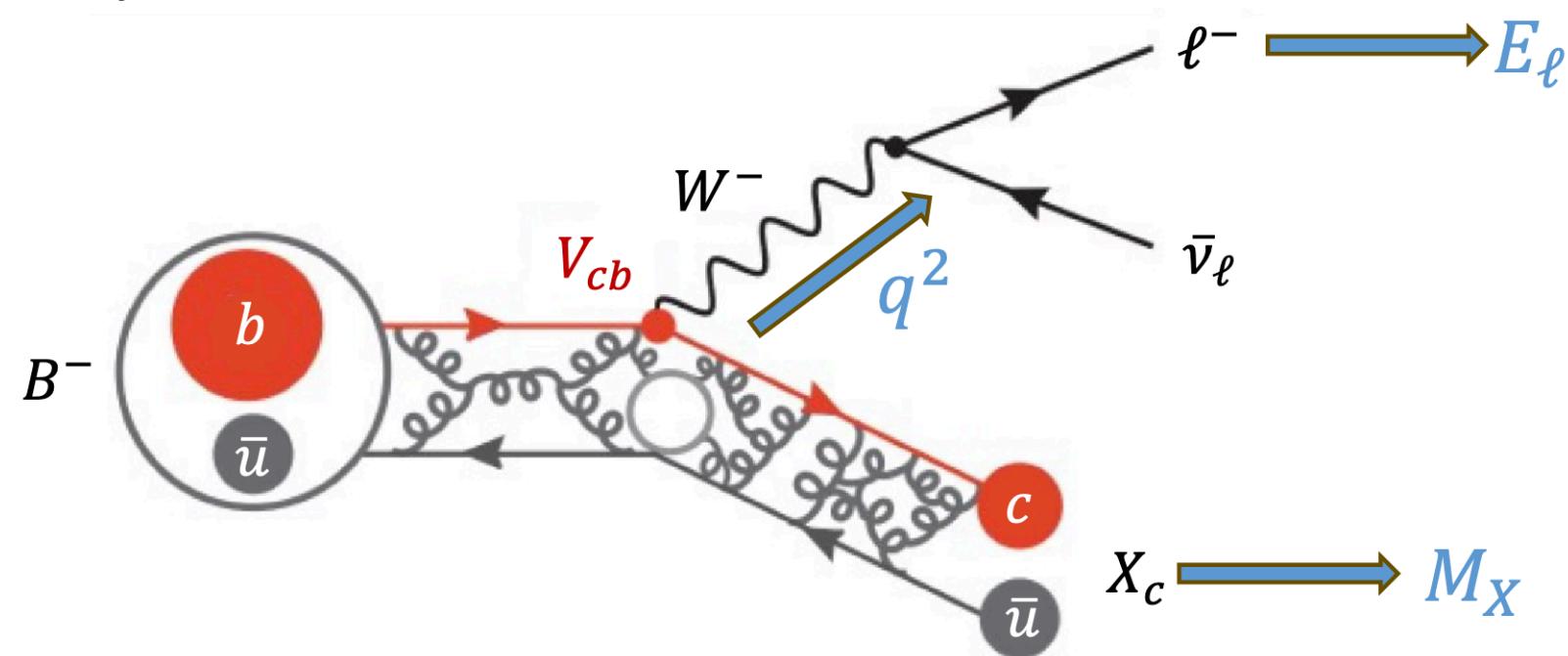
- In the OPE we usually consider moments of the differential distributions
- Perform global fits of semileptonic decays on lattice data

$$\langle (O)^n \rangle_{\text{cut}} = \int_{\text{cut}} (O)^n \frac{d\Gamma}{dO} dO \Bigg/ \int_{\text{cut}} \frac{d\Gamma}{dO} dO$$

$\langle E_l^n \rangle_{\text{cut}}$	$\langle (q^2)^n \rangle_{\text{cut}}$	$\langle (M_X^2)^n \rangle_{\text{cut}}$
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where we restrict the integration:  $E_l > E_{\text{cut}}$  or  $q^2 > q_{\text{cut}}^2$

$\mu_\pi^2, \mu_G^2, \rho_D^3, \rho_{LS}^3, \dots$
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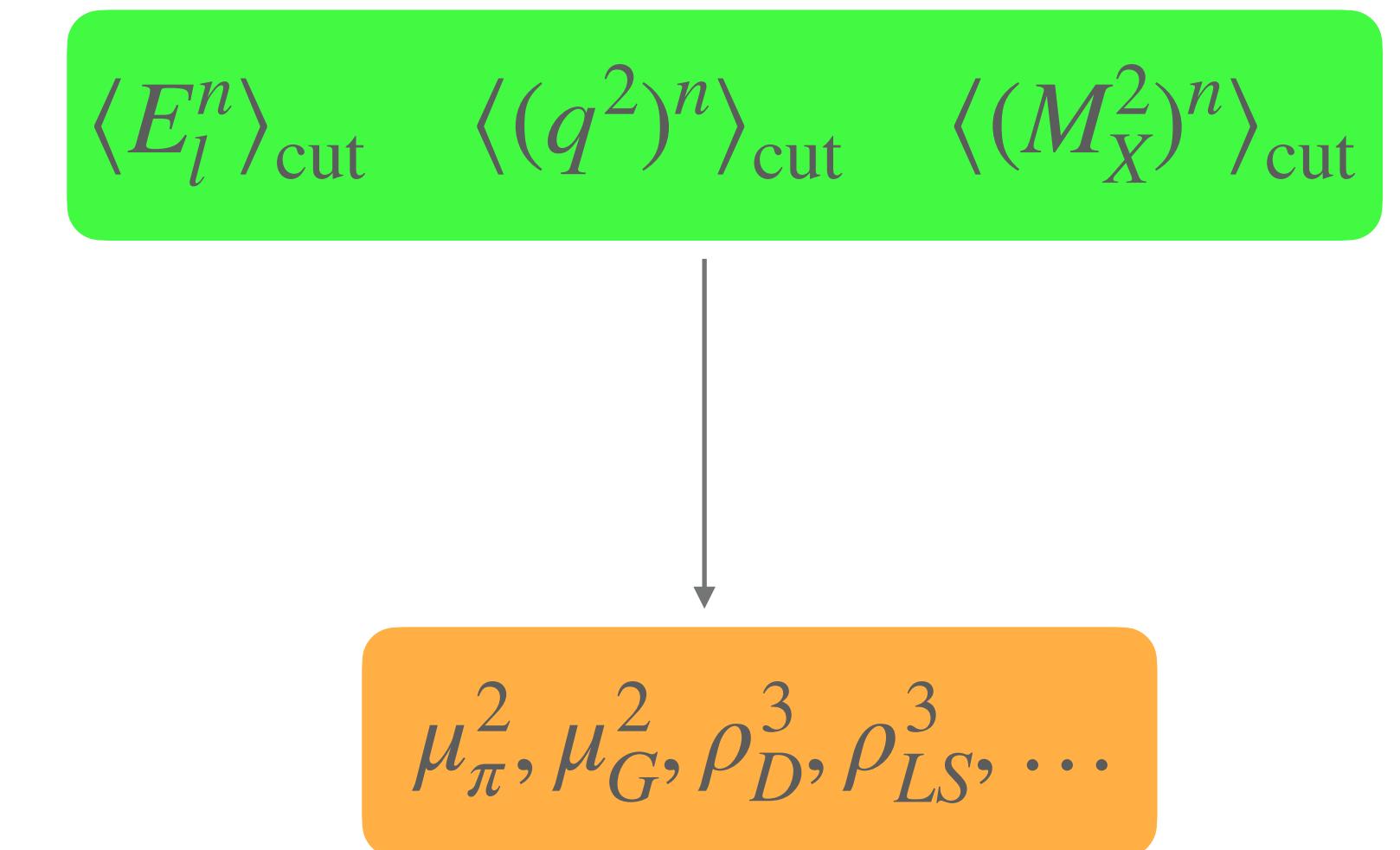


- $O = E_\ell$  : energy of the charged lepton
- $O = M_X^2$  : hadronic invariant mass
- $O = q^2$  : leptonic invariant mass

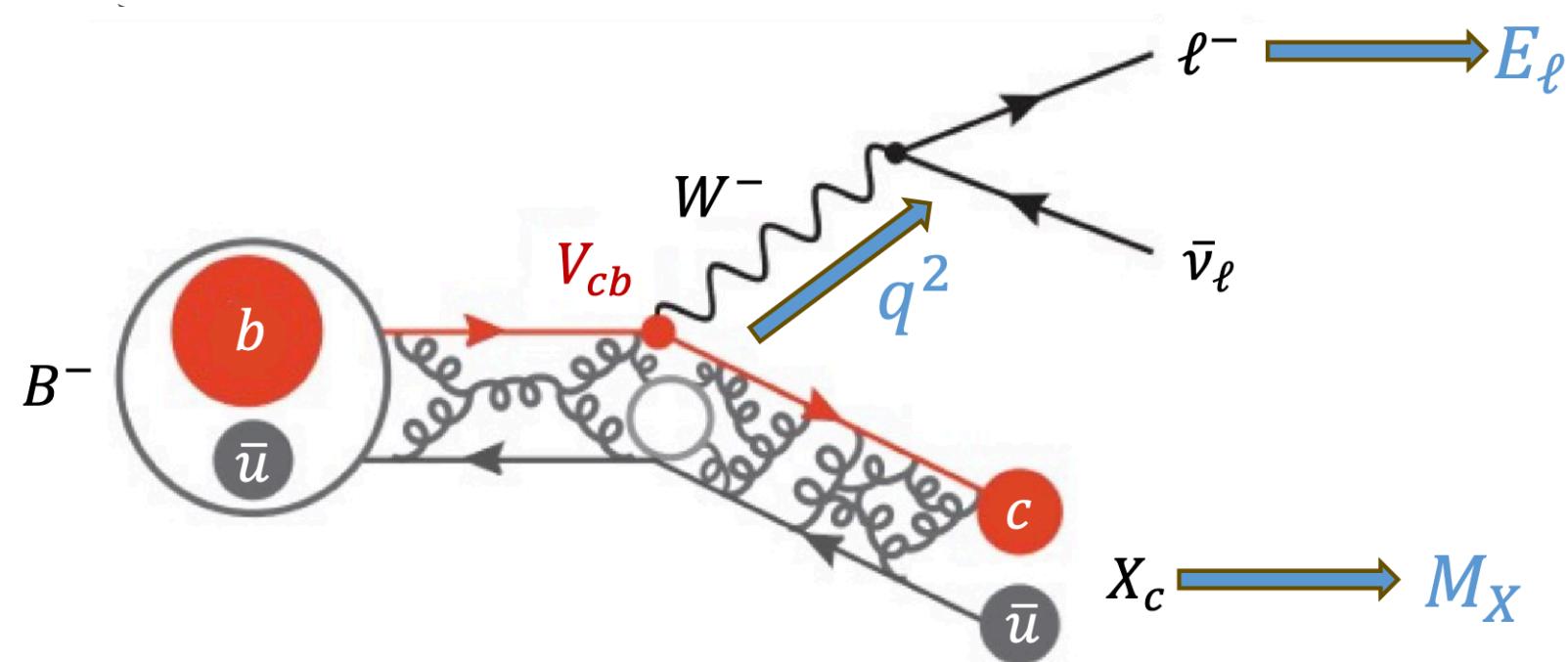
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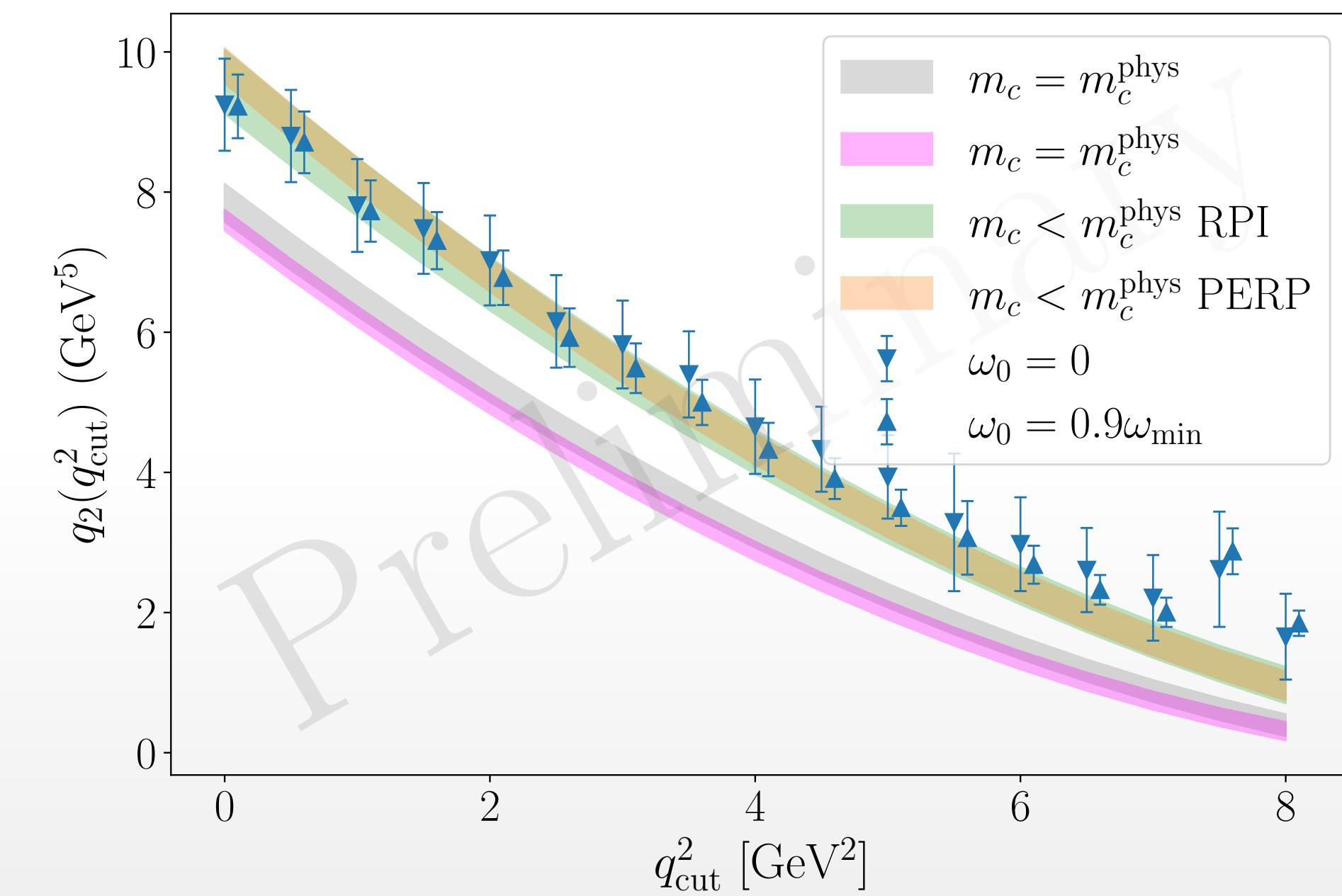
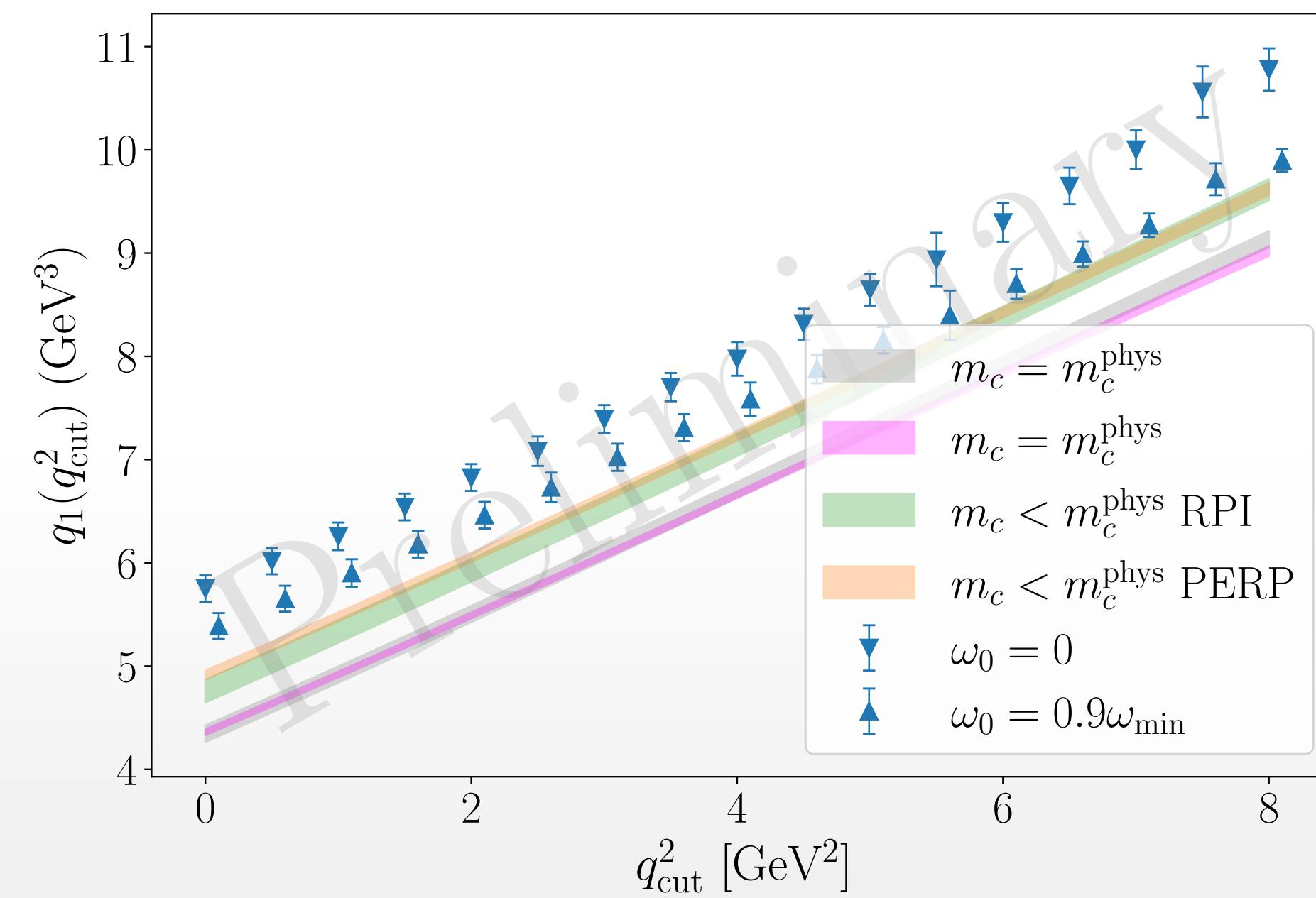
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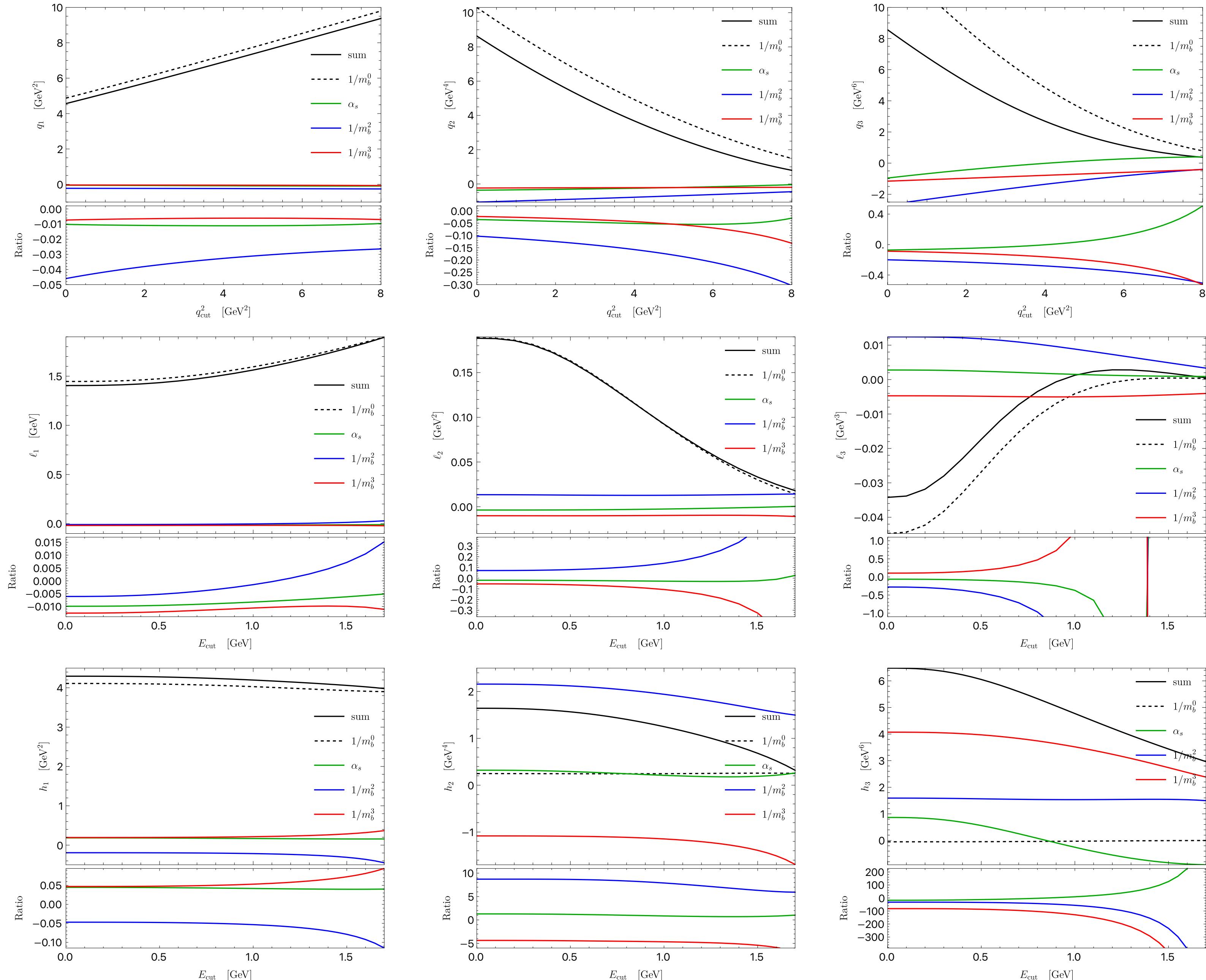
# Centralized moments: lattice “VS” continuum

Preliminary results - feasibility study.

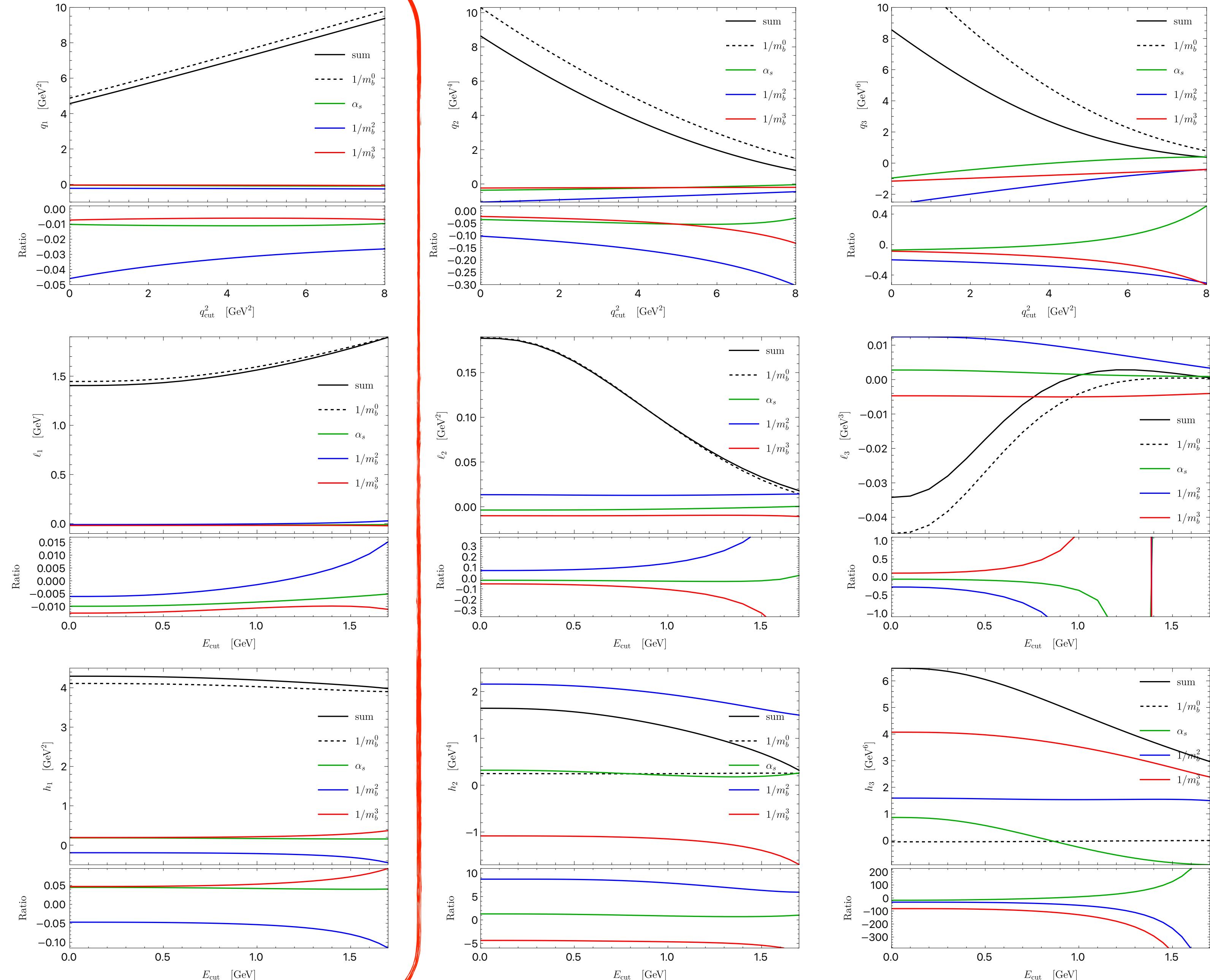


Lattice data (after extrapolation to the physical world) can be used to extract HQET parameters used in the OPE expansion

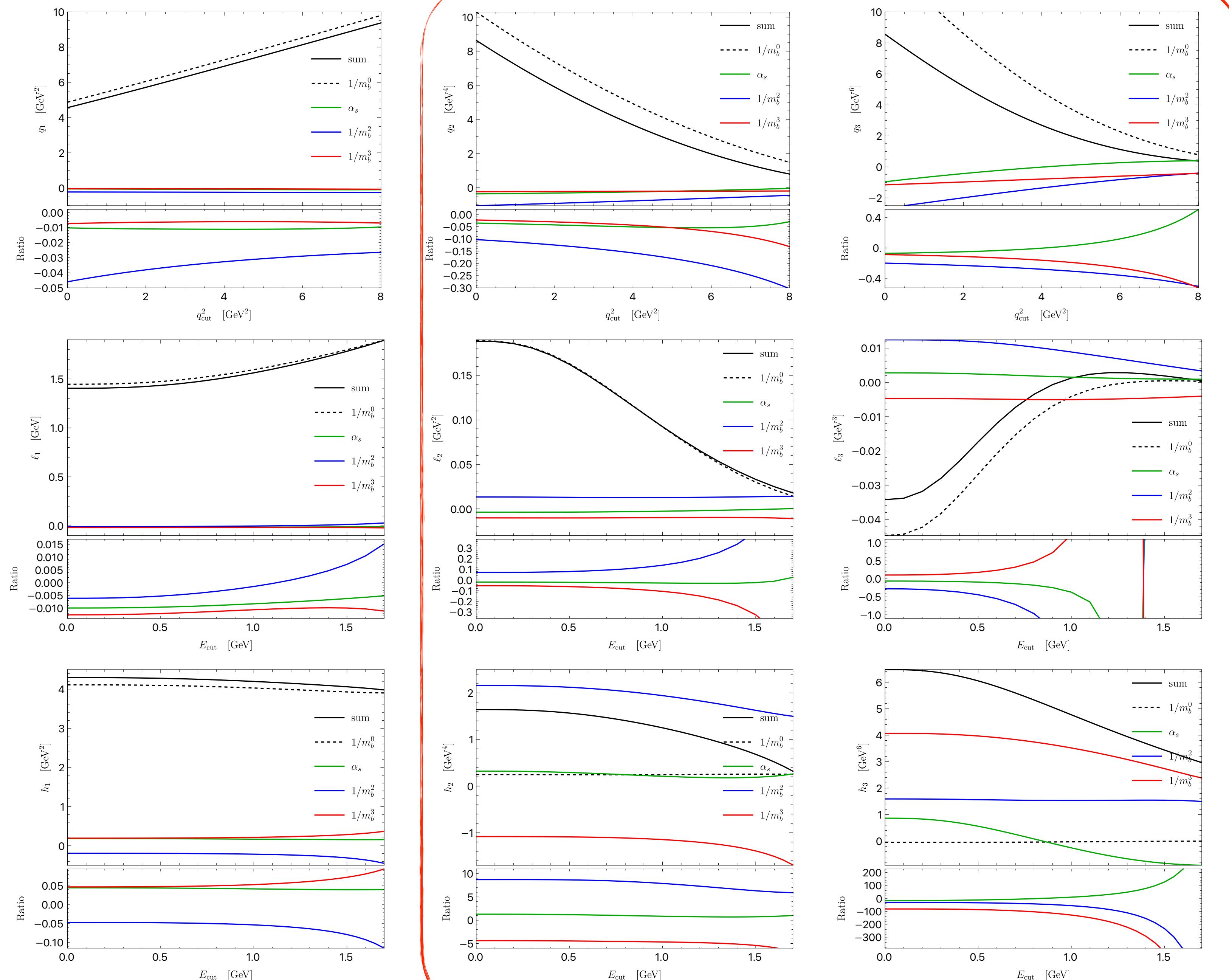
► Different moments have different sensitivities to the power corrections



► First moments are rather insensitive on power corrections

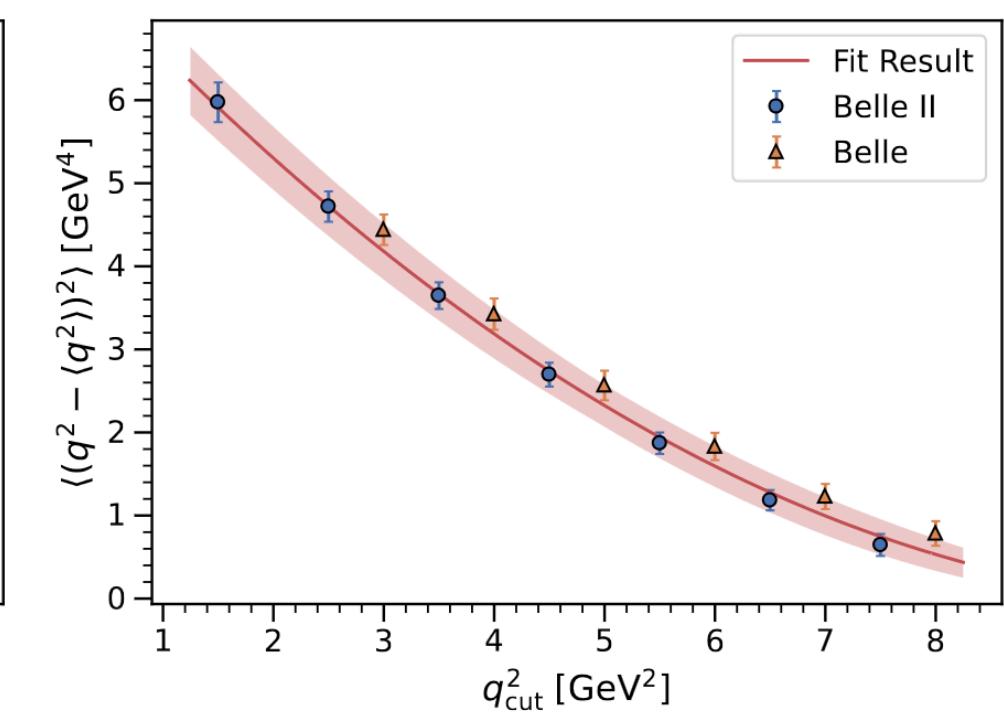
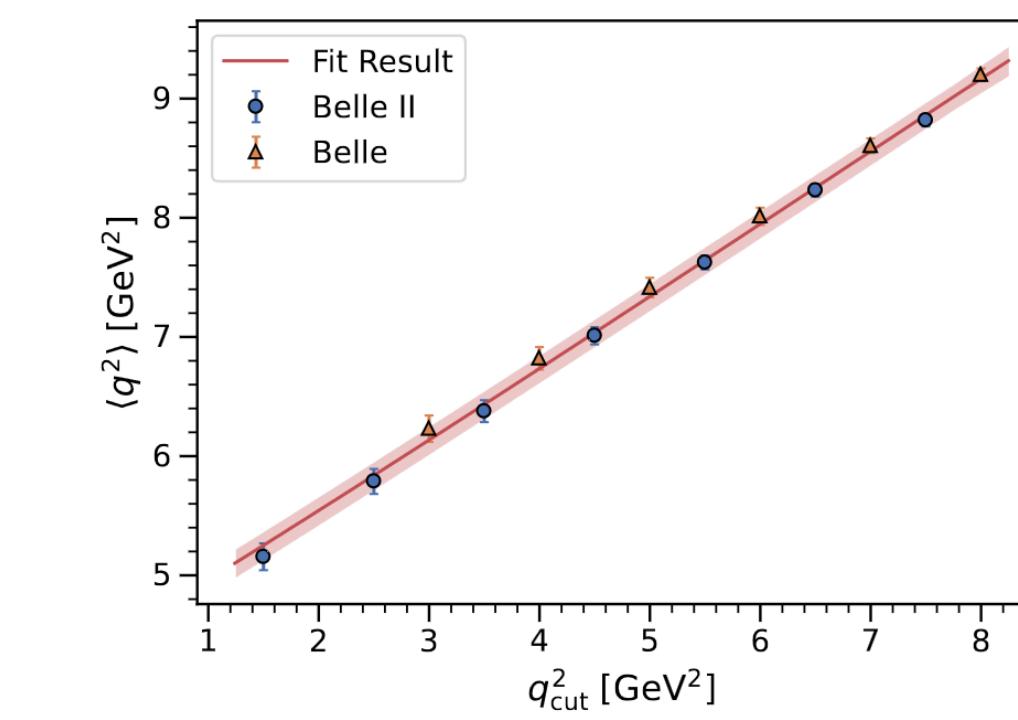
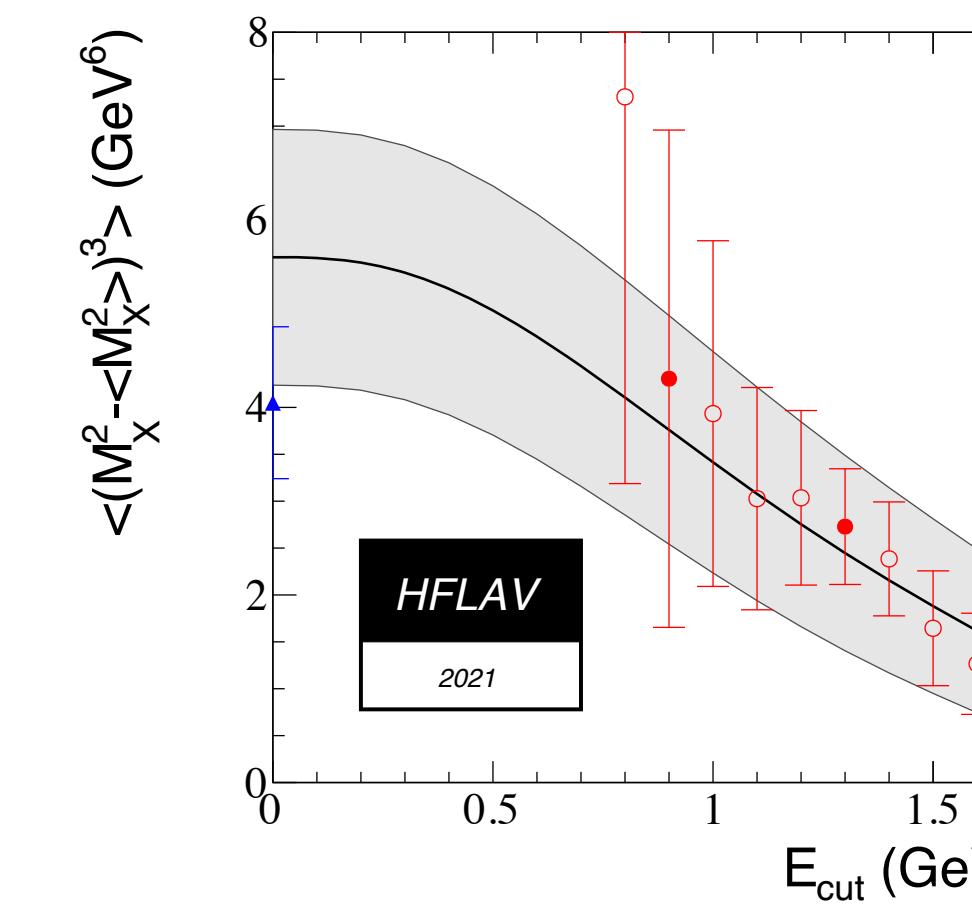
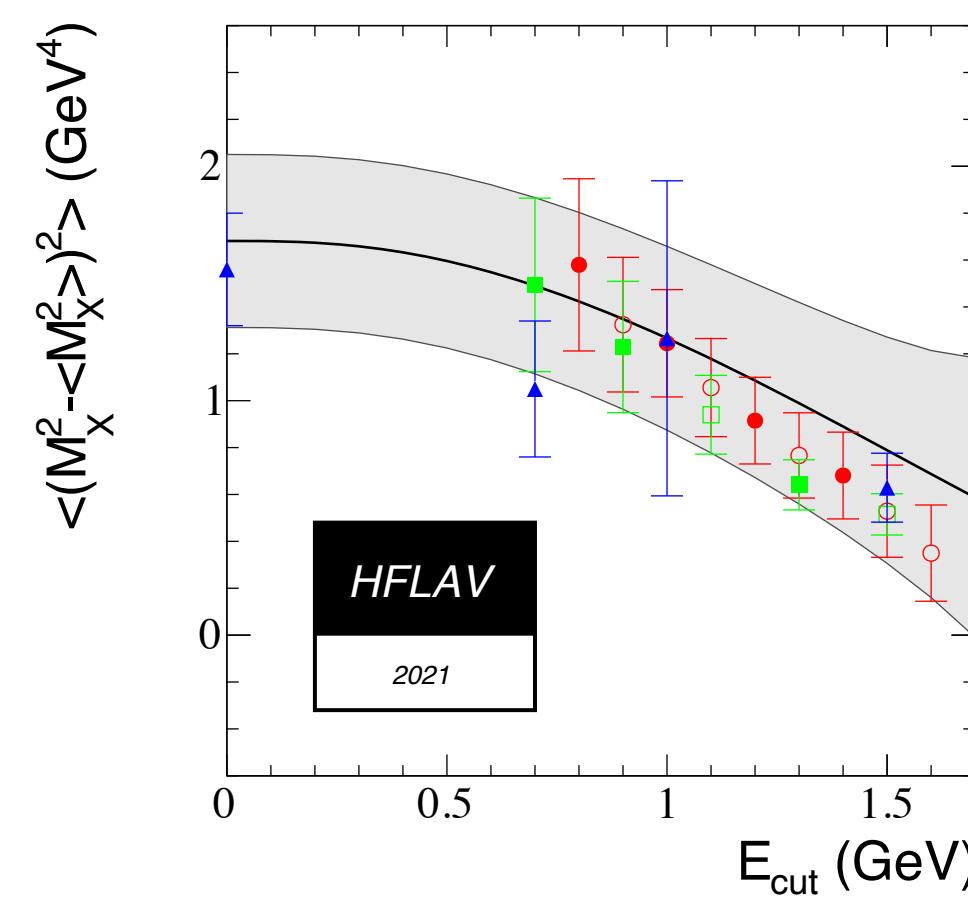
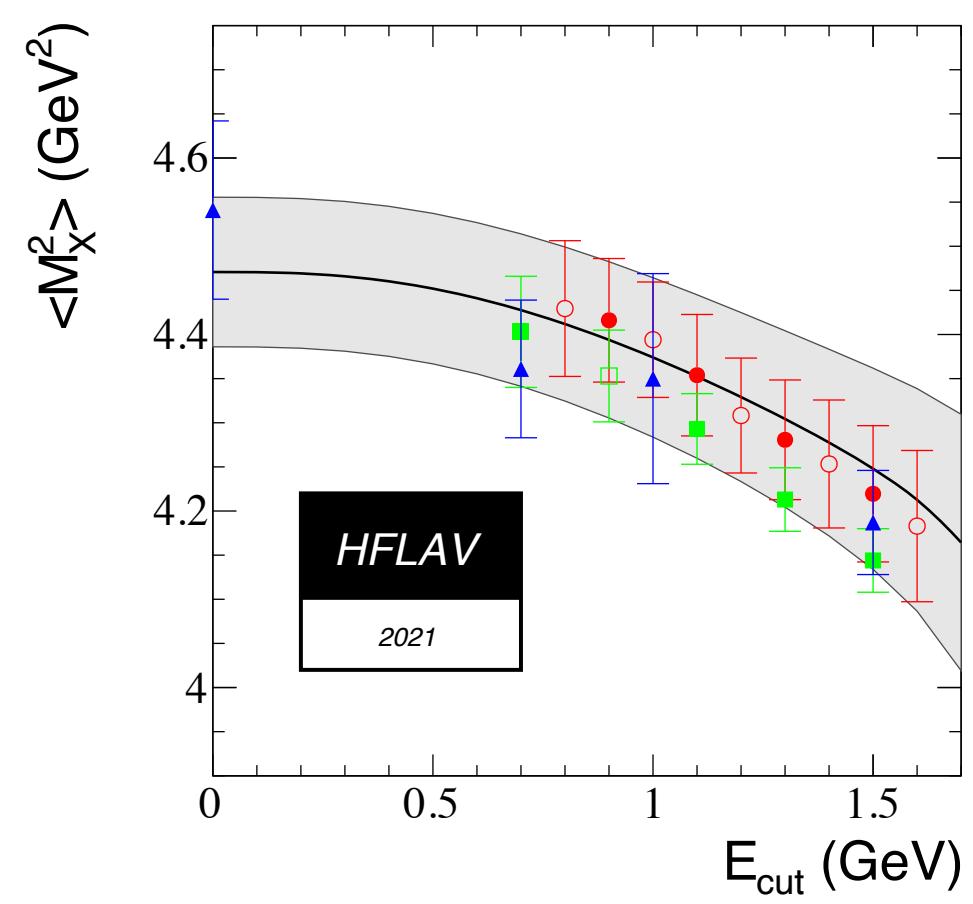


- First moments are rather insensitive on power corrections
- Higher moments are sensitive to power corrections



# HOW TO IMPROVE

- Some experimental data are very precise, some are not



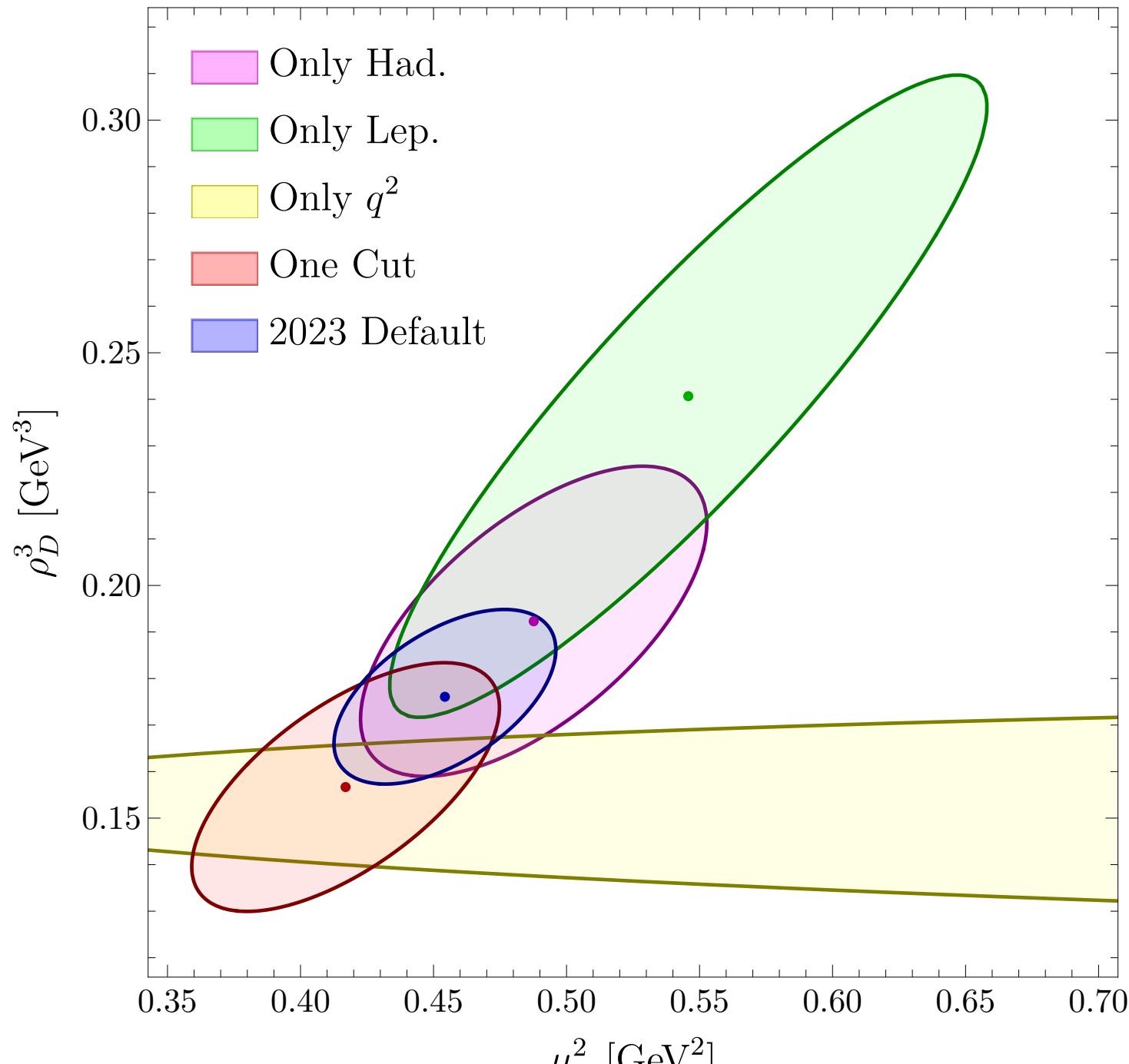
Bernlochner, et al, JHEP 10 (2022) 068

based on Gambino, Schwanda Phys.Rev.D 89 (2014) 1, 014022

- Can lattice substantially improve the precision?

# COMPLEMENTARITY BETWEEN OBSERVABLES

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Gambino, Finauri, JHEP 02 (2024) 206

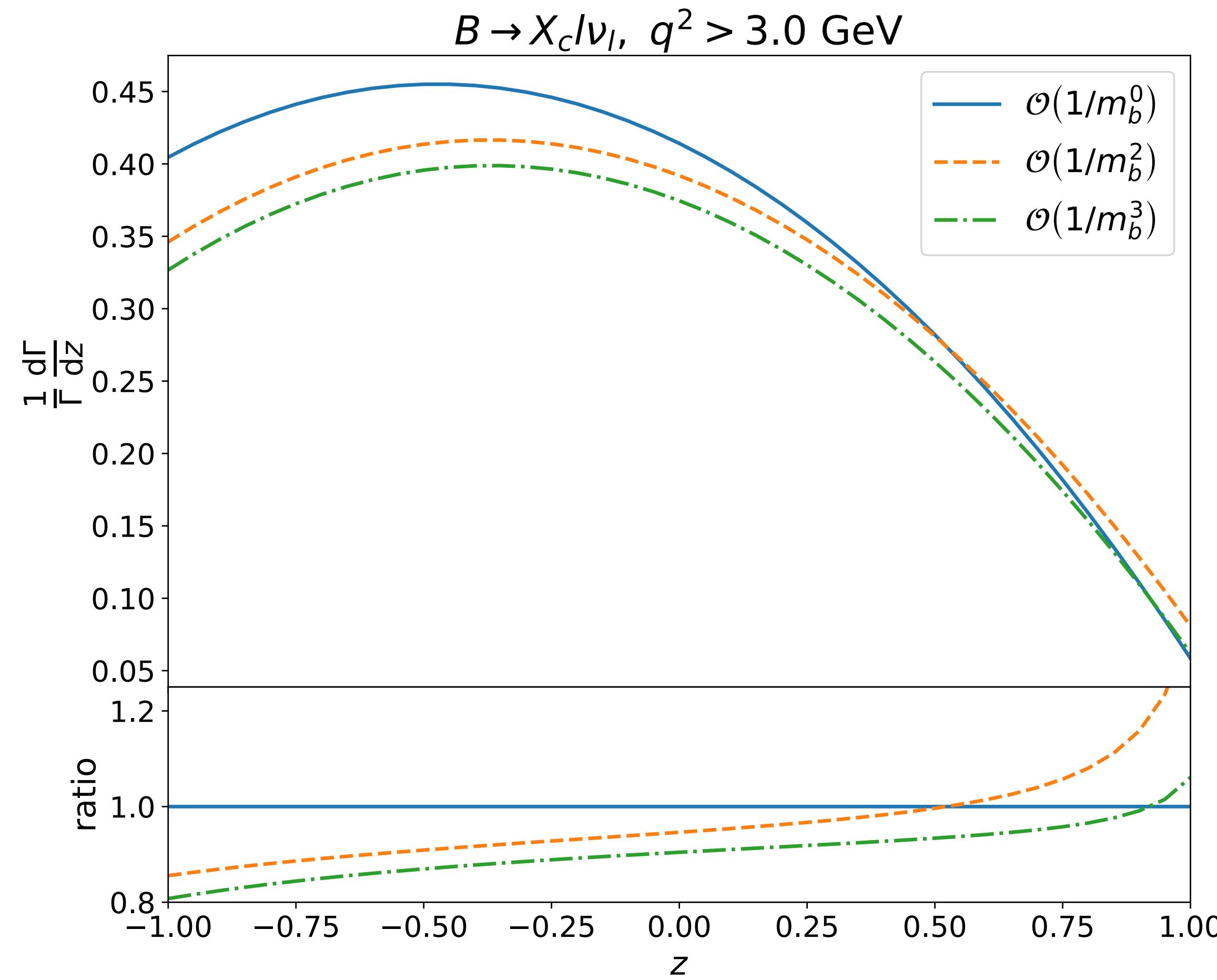
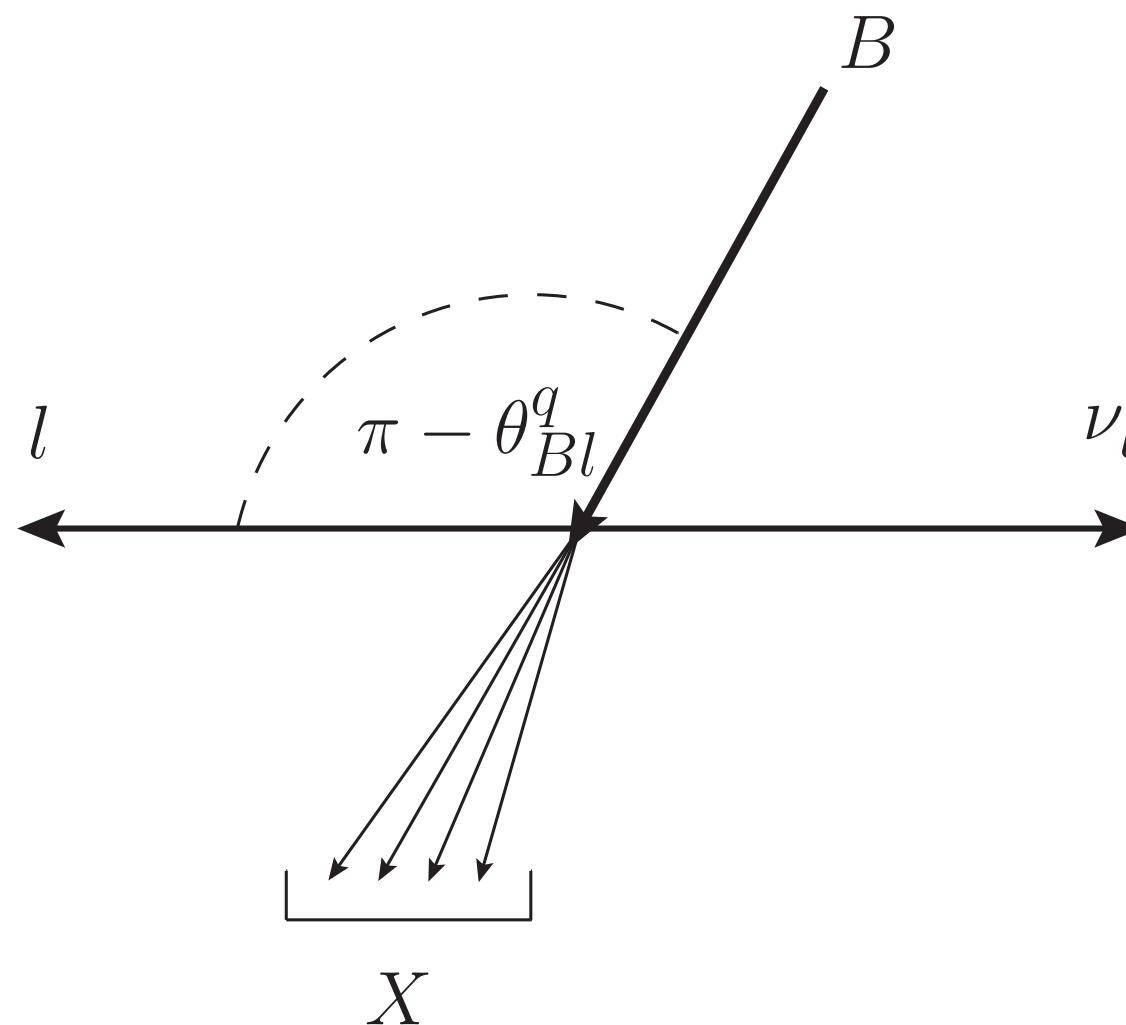
- Moments exhibit “flat directions” in the space of HQE parameters.
- Use “orthogonal observables” to break the degeneracy and improve the sensitivity of the fits on the HQE
- It can be beneficial to keep separated the V and A currents for all moments.
- Consider also scalar and pseudo-scalar currents.
- Expressions available up to  $1/m_b^3$  and  $O(\alpha_s)$  at  $1/m_b^0$

MF, Rahimi, Vos, 2208.04282 [hep-ph]

# FORWARD-BACKWARD ASYMMETRIES

$$A_{FB} = \frac{\int_{-1}^0 \frac{d\Gamma}{dz} - \int_{-1}^0 \frac{d\Gamma}{dz}}{\int_{-1}^0 \frac{d\Gamma}{dz} + \int_{-1}^0 \frac{d\Gamma}{dz}}$$

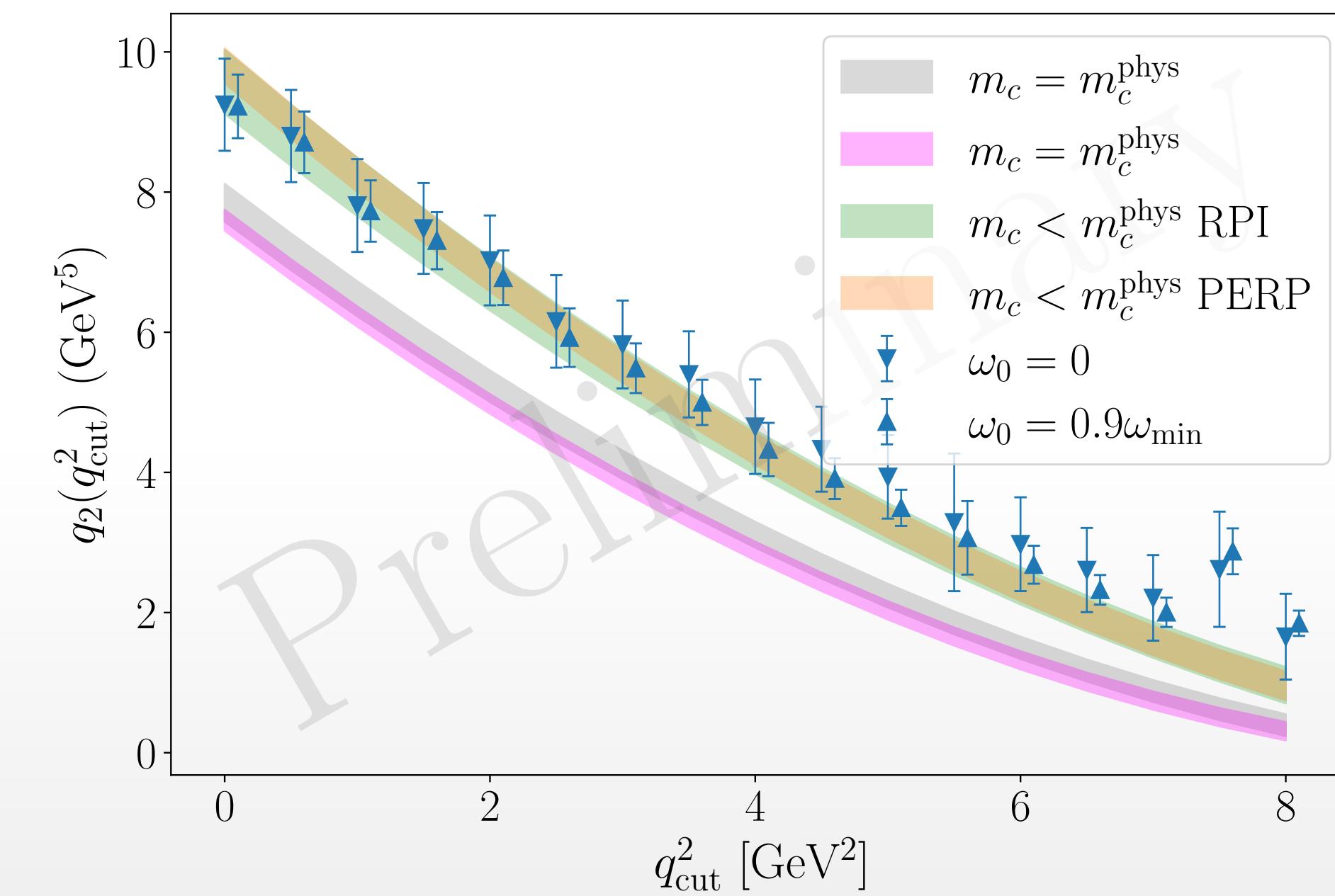
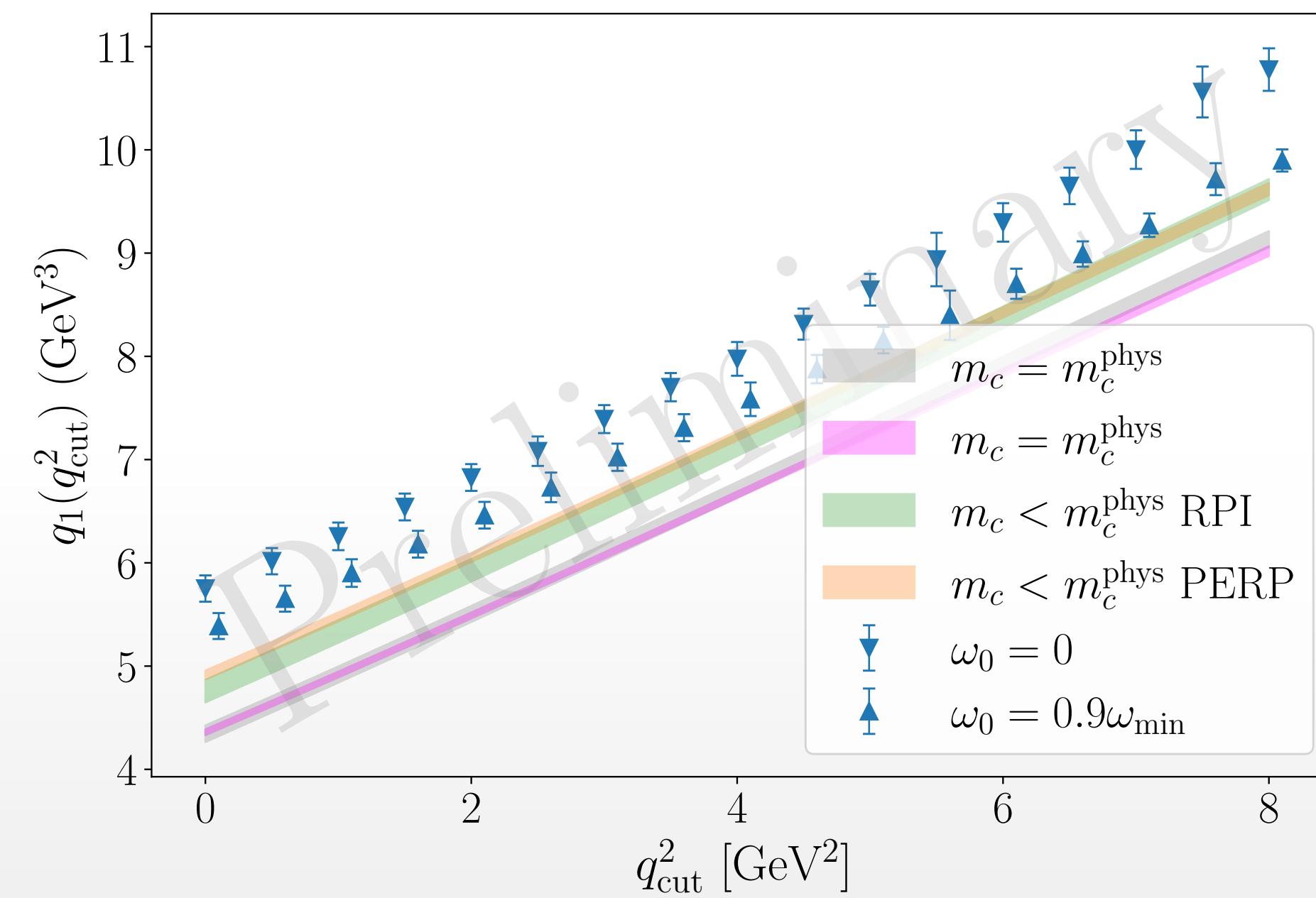
$$z = \cos \theta = \frac{\nu \cdot p_\nu - \nu \cdot p_\ell}{\sqrt{(\nu \cdot q)^2 - q^2}}$$



Turczyk, JHEP 04 (2016) 131  
Herren, SciPost Phys. 14 (2023) 020

# Centralized moments: lattice “VS” continuum

Preliminary results - feasibility study.



Lattice data (after extrapolation to the physical world) can be used to extract HQET parameters used in the OPE expansion

# BOTTOM AND CHARM MASSES

e.g.

$$\mu_\pi^2 = \frac{1}{2M_B} \langle B | \bar{b}_\nu (iD)^2 b_\nu | B \rangle$$

- $m_b$  must be physical to extract the HQE parameters
- Physical or unphysical charm mass?

$\mu_\pi^2, \mu_G^2, \dots$  are independent on  $m_c$   
however they depend on  $M_B$

$$\begin{aligned} \langle O^n \rangle_{\text{cut}} = (m_b)^{mn} & \left[ X_0^{(O,n)} + \frac{\alpha_s}{\pi} X_1^{(O,n)} + \left( \frac{\alpha_s}{\pi} \right)^2 X_2^{(O,n)} + \frac{\mu_\pi^2}{m_b^2} \left( p_0^{(O,n)} + \frac{\alpha_s}{\pi} p_1^{(O,n)} + \dots \right) \right. \\ & + \frac{\mu_G^2}{m_b^2} \left( g_0^{(O,n)} + \frac{\alpha_s}{\pi} g_1^{(O,n)} + \dots \right) + \frac{\rho_D^3}{m_b^3} \left( d_0^{(O,n)} + \frac{\alpha_s}{\pi} d_1^{(O,n)} + \dots \right) + \frac{\rho_{LS}}{m_b^2} \left( l_0^{(O,n)} + \frac{\alpha_s}{\pi} l_1^{(O,n)} + \dots \right) + O\left(\frac{1}{m_b^4}\right) \end{aligned}$$

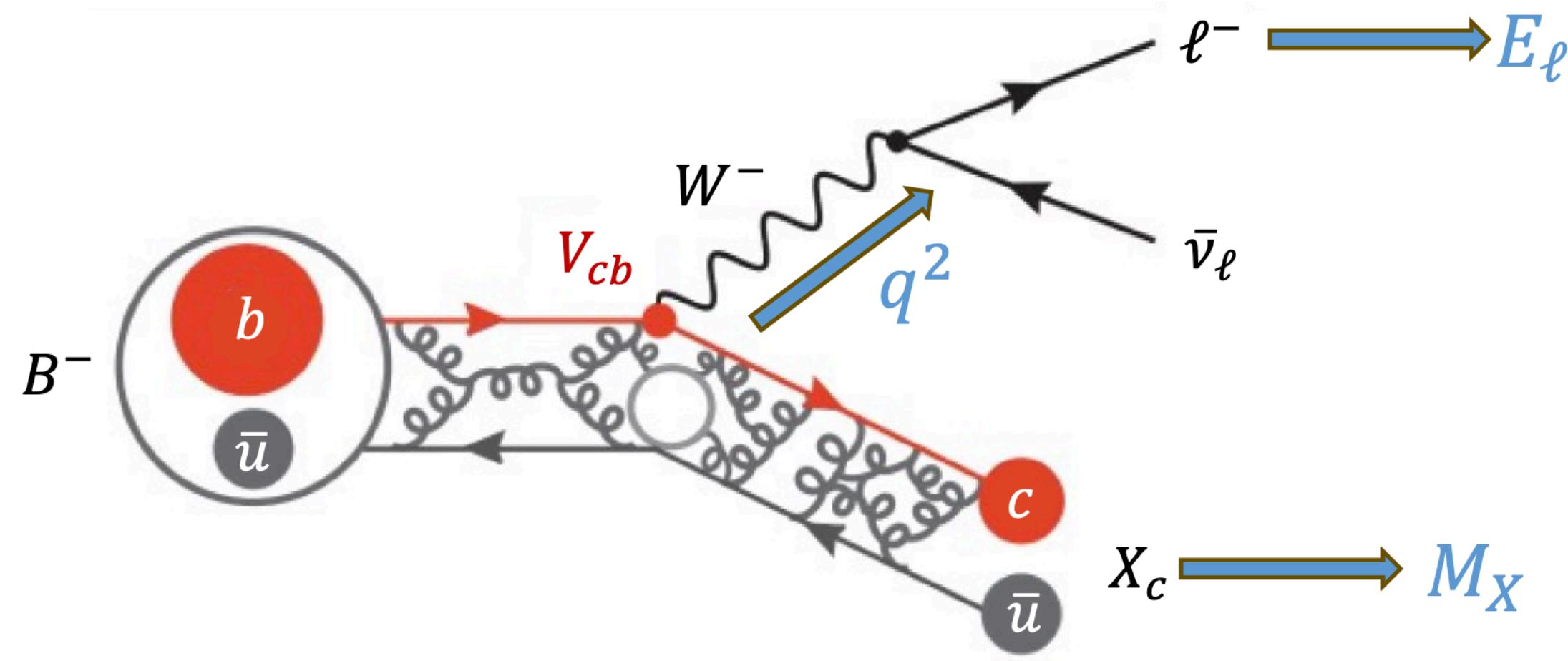
➤ HQET estimate:

$$M_D = m_c + \overline{\Lambda} + \frac{\mu_\pi^2 - \frac{d_H}{3} \mu_G^2}{2m_c} + O\left(\frac{1}{m_c^2}\right)$$

- With  $m_c < m_c^{\text{phys}}$ , lattice should provide reference value for  $\overline{m}_c$ !

# SMALL VELOCITY SUM RULES

Bigi, Shifman, Uraltsev, Vainshtein, *Phys.Rev.D* 56 (1997) 4017



Excitation energy of the  $X_c$  system

$$\omega_X = M_B - q_0 - \sqrt{M_D^2 + \mathbf{q}^2}$$

Maximum energy of the dilepton system

$$q_0^{\max} = M_B - \sqrt{M_D^2 + \mathbf{q}^2}$$

$$I_n(\mathbf{q}^2, \mu) = \int_{q_0^{\max} - \mu}^{q_0^{\max}} dq_0 \omega_X^n \frac{d^2 \Gamma}{dq_0 d\mathbf{q}^2}$$

$\mu \simeq 1 \text{ GeV}$

# SMALL VELOCITY SUM RULES

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$$I_1(\mathbf{q}^2, \mu) = I_0(\mathbf{q}^2) \frac{\mathbf{q}^2}{2m_c^2} \bar{\Lambda}(\mu) + O(|\mathbf{q}|^3, \Lambda_{\text{QCD}}^2)$$

$$I_2(\mathbf{q}^2, \mu) = I_0(\mathbf{q}^2) \frac{\mathbf{q}^2}{3m_c^2} \mu_\pi^2(\mu) + O(|\mathbf{q}|^3, \Lambda_{\text{QCD}}^3)$$

$$I_3(\mathbf{q}^2, \mu) = I_0(\mathbf{q}^2) \frac{\mathbf{q}^2}{2m_c^2} \rho_D^3(\mu) + O(|\mathbf{q}|^3, \Lambda_{\text{QCD}}^4)$$

- SV sum rules are independent on the current (S,V,A)
- Can lattice verify these relations?
- Can one extract  $\mu_\pi^2, \rho_D^3$  by studying the limit  $|\mathbf{q}| \rightarrow 0$  and  $m_Q \rightarrow \infty$

Note: the HQE parameters extracted here are in  $m_Q \rightarrow \infty$  limit



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