

RG improvement and α_s determination from relativistic sum rules

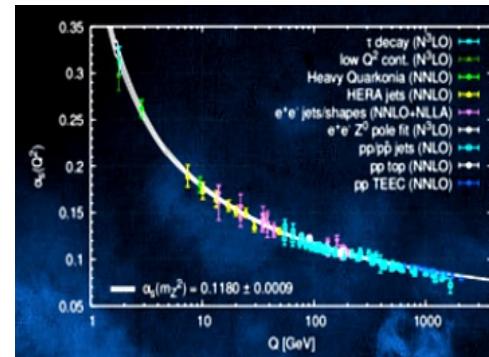
$$R_{\bar{q}q}(s) \equiv \frac{3s}{4\pi\alpha^2} \sigma(e^+e^- \rightarrow q\bar{q} + X)$$

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"Renormalization group improved determination of α_s , m_c , and m_b from the low energy moments of heavy quark current correlators" (Phys.Rev.D 108 (2023) 7, 074029)



Feb 05, 2024



$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^\alpha G_{\mu\nu}^\alpha + \sum_j \bar{q}_j (i \gamma^\mu D_\mu + m_j) q_j$$

where $G_{\mu\nu}^\alpha \equiv \partial_\mu A_\nu^\alpha - \partial_\nu A_\mu^\alpha + i f_{bc}^{~~a} A_\mu^b A_\nu^c$

and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$

alphas-2024: Workshop on precision measurements of the QCD coupling constant

Outline

- Motivation
- Definition of various quantities used.
- Prescriptions for the perturbative series.
 - a. FOPT and RGSPT
 - b. RG improvements.
- Issues with FOPT using charm moments
 - a) Charm mass determinations
 - b) Strong coupling determination
- Results

Motivation

- K. G. Chetyrkin, J. H. Kuhn, A. Maier, P. Maierhofer, P. Marquard, M. Steinhauser and C. Sturm,
``Addendum to Charm and bottom quark masses: An update [arXiv:1710.04249 [hep-ph]].

n	$m_c(3\text{GeV})$	exp	α_s	μ	np LO	total	$m_c(m_c)$
1	0.993	0.007	0.004	0.002	0.001	0.008	1.279
2	0.982	0.004	0.007	0.005	0.001	0.010	1.269
3	0.982	0.003	0.008	0.006	0.001	0.010	1.269
4	1.003	0.002	0.005	0.028	0.001	0.029	1.288

Quantities and their definitions

$$R_{\bar{q}q}(s) \equiv \frac{3s}{4\pi\alpha^2} \sigma(e^+e^- \rightarrow q\bar{q} + X) \simeq \frac{\sigma(e^+e^- \rightarrow q\bar{q} + X)}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

$q=\{c,b\}$

$$\mathcal{M}_n^{V,\text{th}} = \frac{12\pi^2 Q_q^2}{n!} \frac{d^n}{ds^n} \Pi^V(s) \Big|_{s=0}$$

$$\mathcal{M}_n^{P,\text{th}} = \frac{12\pi^2 Q_q^2}{n!} \frac{d^n}{ds^n} P(s) \Big|_{s=0}$$

$$P(s) = \frac{1}{s^2} \left(\Pi^P(s) - \Pi^P(0) - s \left[\frac{d}{ds} \Pi^P(s) \right]_{s=0} \right)$$

$$\Pi^P(q^2) \equiv i \int dx e^{iqx} \langle 0 | T\{j_P(x)j_P(0)\} | 0 \rangle$$

$$j_P = 2im_q \bar{q}(x)\gamma^5 q(x)$$

$$\mathcal{R}_n^X \equiv \frac{(\mathcal{M}_n^X)^{\frac{1}{n}}}{(\mathcal{M}_{n+1}^X)^{\frac{1}{n+1}}}$$

$$\begin{aligned} \mathcal{M}_q^{V,n} &= \int_0^\infty \frac{ds}{s^{n+1}} R_{q\bar{q}}(s) \\ &= \int_0^{s_{\text{th.}}} \frac{ds}{s^{n+1}} R_{q\bar{q}}^{\text{exp.}}(s) + \int_{s_{\text{th.}}}^\infty \frac{ds}{s^{n+1}} R_{q\bar{q}}^{\text{cont.}}(s) \end{aligned}$$

$$R_{\bar{q}q}^{\text{cont.}} \equiv 12\pi \text{Im} \Pi^V(s + i\epsilon) = 12\pi \times \frac{1}{2i} (\Pi^V(s + i\epsilon) - \Pi^V(s - i\epsilon))$$

$$(q^2 g_{\mu\nu} - q_\mu q_\nu) \Pi^V(q^2) = -i \int dx e^{iqx} \langle 0 | T\{j_\mu(x)j_\nu(0)\} | 0 \rangle$$

$$j_\mu(x) = \bar{q}(x)\gamma^\mu q(x)$$

$$\mathcal{M}_n^{X,\text{th.}} = \mathcal{M}_n^{X,\text{pert.}} + \mathcal{M}_n^{X,\text{n.p.}}$$

$$\mathcal{M}_n^{X,\text{pert.}} = m_q^{-2n} \sum_{i,j} T_{i,j}^{X,\text{pert.}} x^i L^j$$

$$\mathcal{M}_n^{X,\text{n.p.}} = \frac{\langle x G^2 \rangle_{i,j}^{\text{RGI}}}{(2m_q)^{4n+4}} \sum_{i,j} T_{i,j}^{X,\text{n.p.}} x^i L^j$$

$$x \equiv \alpha_s(\mu)/\pi$$

$$L \equiv \log(\mu^2/m_q^2)$$

$$m_q \equiv m_q(\mu)$$

Perturbative prescriptions

- Fixed order perturbation theory (FOPT)
- Renormalization group summed perturbation theory (RGSPt) N^n LL are summed using Renormalization group

$$\text{FOPT : } \mathcal{M}_n^{X, \text{pert}} = m_q^{-2n} \sum_{i,j} T_{i,j}^X x^i L^j \quad \beta(x) \equiv \mu^2 \frac{d}{d\mu^2} x(\mu) = - \sum_i \beta_i x^{i+2},$$

$$\text{RGSPt : } \mathcal{M}_n^{X, \Sigma} = m_q^{-2n} \sum_{i=0} x^i S_i(xL) \implies S_i(xL) = \sum_{n=i}^{\infty} T_{n,n-i}^X (xL)^{n-i} \quad \gamma_m(x) \equiv \mu^2 \frac{d}{d\mu^2} \log(m_q(\mu)) = - \sum_i \gamma_i x^{i+1}$$

$$\text{RGE : } \mu^2 \frac{d}{d\mu^2} \mathcal{M}_n^X = 0 \implies (\beta(x) \partial_x + \gamma_m(x) \partial_m + \partial_L) \mathcal{M}_n^X = 0$$

$$\sum_{i=0}^k \left[\beta_i (\delta_{i,0} + w - 1) S'_{k-i}(w) + (\beta_i(k-i) - 2n\gamma_i) S_{k-i}(w) \right] = 0$$

$w = 1 - \beta_0 x L$

$$S_0(w) = w^{\frac{2n\gamma_0}{\beta_0}}, \quad S_1(w) = \left(T_{1,0}^X + \frac{2nL_w\gamma_0(\beta_1 - 2\beta_0\gamma_0)}{\beta_0^2} + \frac{2n(\beta_1\gamma_0 - \beta_0\gamma_1)}{\beta_0^2} \right) w^{\frac{2n\gamma_0}{\beta_0}-1} + \frac{2n(\beta_0\gamma_1 - \beta_1\gamma_0) w^{\frac{2n\gamma_0}{\beta_0}}}{\beta_0^2}$$

$$\Omega_{n,a} \equiv \frac{\log^n(w)}{w^a}$$

positive integer

$a \propto \gamma_0 / \beta_0$

Other examples: strong coupling constant

FOPT:

$$x(Q) = x \left(1 + x\beta_0 L + x^2 (\beta_1 L + \beta_0^2 L^2) + x^3 \left(\beta_2 L + \frac{5}{2} \beta_1 \beta_0 L^2 + \beta_0^3 L^3 \right) \right) + \mathcal{O}(x^5)$$

RGSPT:

$$\begin{aligned} x(Q) = & \frac{x}{w} - \frac{x^2 \tilde{\beta}_1 L_w}{w^2} + x^3 \left(\frac{-\tilde{\beta}_1^2 + \tilde{\beta}_2 + \tilde{\beta}_1^2 L_w^2 - \tilde{\beta}_1^2 L_w}{w^3} + \frac{\tilde{\beta}_1^2 - \tilde{\beta}_2}{w^2} \right) + x^4 \left\{ \frac{-\frac{\tilde{\beta}_1^3}{2} + \tilde{\beta}_2 \tilde{\beta}_1 - \frac{\tilde{\beta}_3}{2}}{w^2} \right. \\ & + \frac{-\frac{\tilde{\beta}_1^3}{2} + \frac{\tilde{\beta}_3}{2} + \tilde{\beta}_1^3 (-L_w^3) + \frac{5}{2} \tilde{\beta}_1^3 L_w^2 + (2\tilde{\beta}_1^3 - 3\tilde{\beta}_1 \tilde{\beta}_2) L_w}{w^4} + \frac{\tilde{\beta}_1^3 - \tilde{\beta}_2 \tilde{\beta}_1 + (2\tilde{\beta}_1 \tilde{\beta}_2 - 2\tilde{\beta}_1^3) L_w}{w^3} \left. \right\} + \mathcal{O}(x^5) \end{aligned}$$

where: $w = 1 - \beta_0 x L$, $L_w \equiv \log(w)$ and $\tilde{\beta}_i = \beta_i / \beta_0$

Not only coupling and moments, even decoupling relations for quark masses and coupling can be summed in closed form!!!

Higher order results

$$\mathcal{M}_n^{X,\text{pert.}} = m_q^{-2n} \sum_{i=0,j=0}^{\infty,i} T_{i,j}^{X,\text{pert.}} x^i L^j \quad T_{i,0}^{X,\text{pert.}}$$

known to four loops ($\mathcal{O}(\alpha_s^3)$)

A. Maier, P. Maierhofer, P. Marquard and A. V. Smirnov,
 Nucl. Phys. B 824(2010), [arXiv:0907.2117 [hep-ph]].

$$\mathcal{M}_n^{X,\text{n.p.}} = \frac{\left\langle x G^2 \right\rangle_{\text{RGI}}}{(2m_q)^{4n+4}} \sum_{i,j} T_{i,j}^{X,\text{n.p.}} x^i L^j \quad T_{i,0}^{X,\text{n.p.}}$$

is known to two loops ($\mathcal{O}(\alpha_s)$)

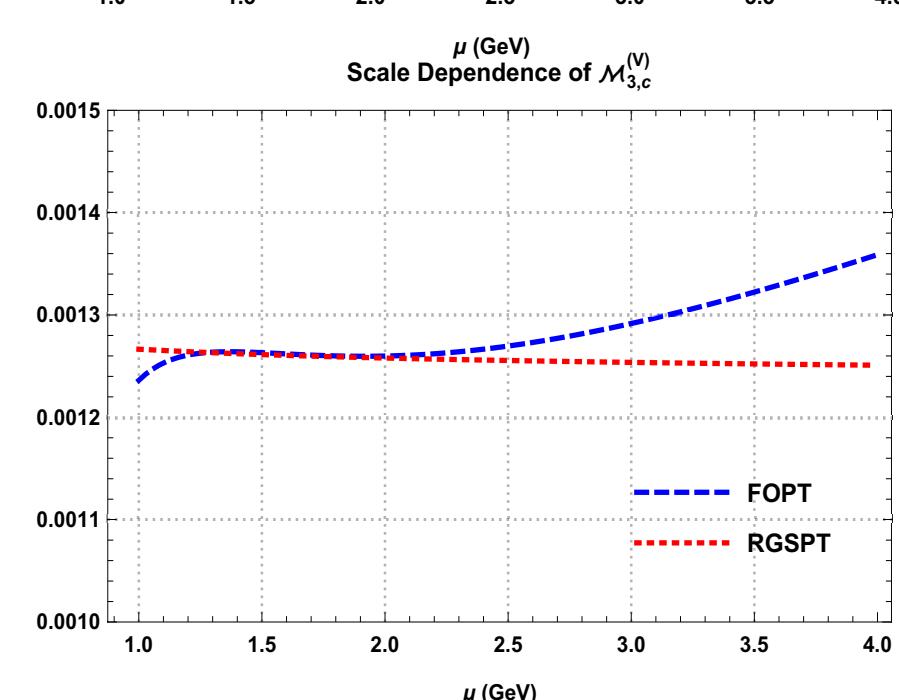
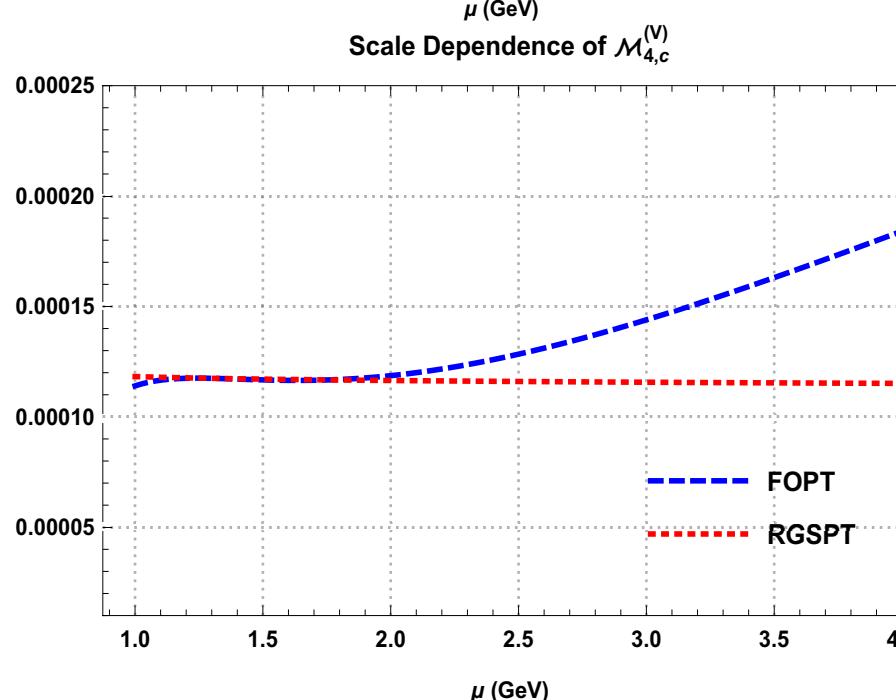
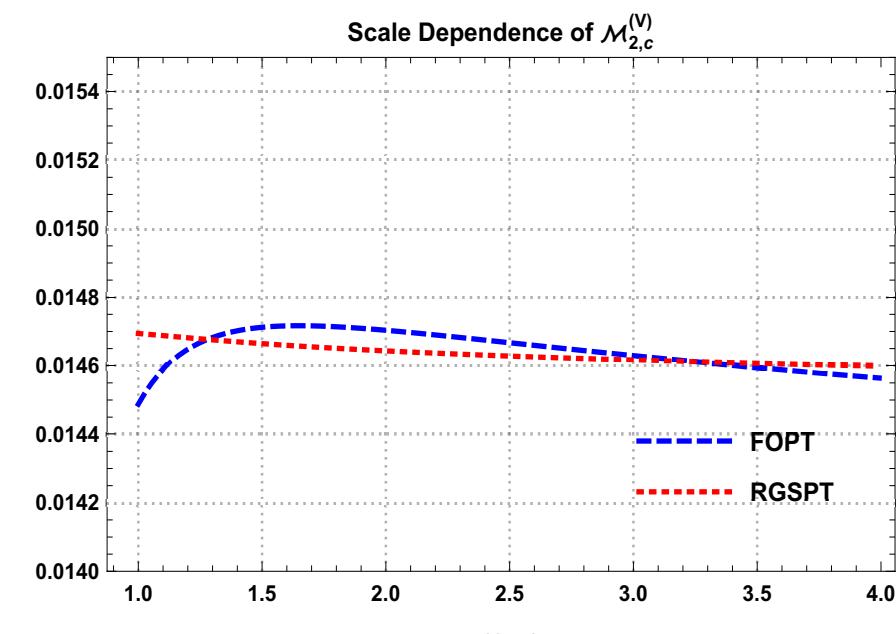
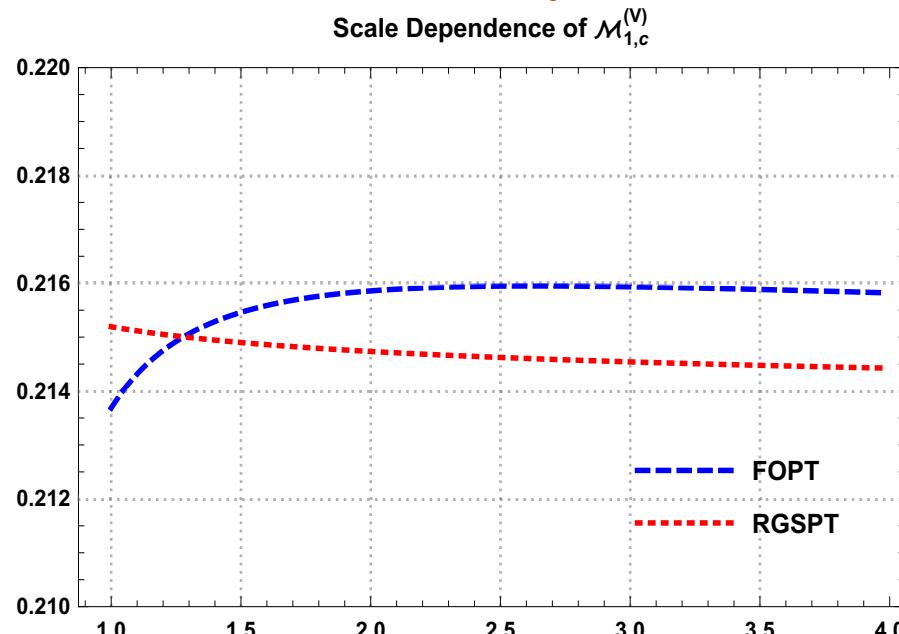
D. J. Broadhurst, et al., Phys. Lett. B 329 (1994),
 [arXiv:hep-ph/9403274 [hep-ph]]

$$\begin{aligned} x &\equiv \alpha_s(\mu)/\pi \\ m_q &\equiv m_q(\mu) \\ L &\equiv \log(\mu^2/m_q^2) \end{aligned} \quad \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle_{\text{RGI}} = 0.006 \pm 0.012 \text{ GeV}^4$$

B. L. Ioffe, Prog. Part. Nucl. Phys. 56 (2006),
 [arXiv:hep-ph/0502148 [hep-ph]].

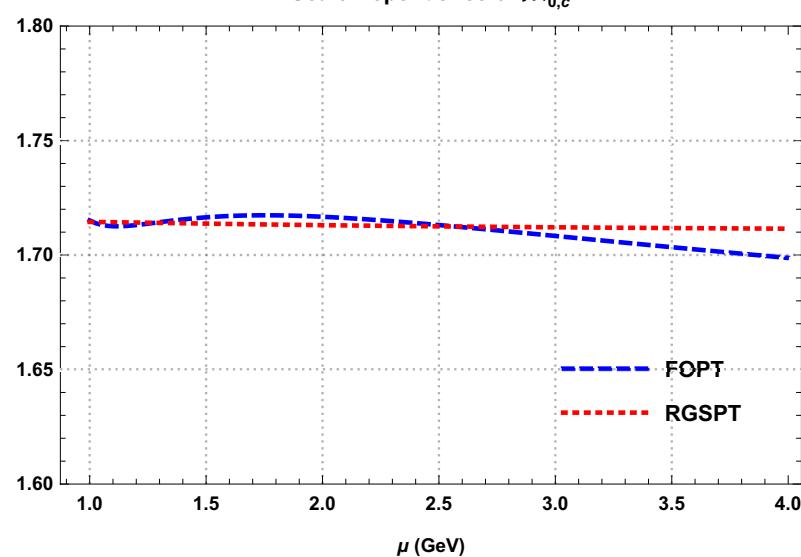
$$m_q(\mu) = M_q \left(1 - x(\mu) \left(\frac{4}{3} + \log \left(\frac{\mu^2}{M_q^2} \right) \right) \right) + \mathcal{O}(x^2) \quad m_c(q) = \overline{m}_c \int_{x(\overline{m}_c)}^{x(q)} dx e^{\left(\frac{\gamma(x)}{\beta(x)} \right)}$$

RG behavior of perturbative moments in FOPT and RGSPT

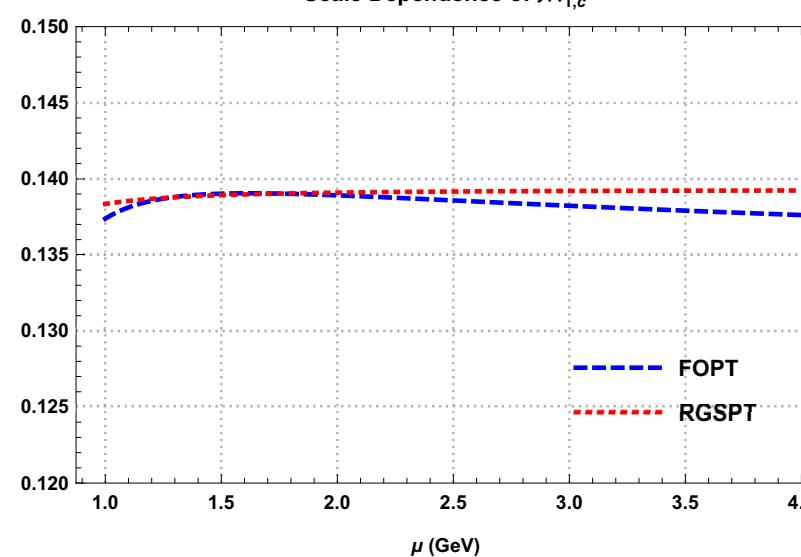


RG behavior of moments contd...

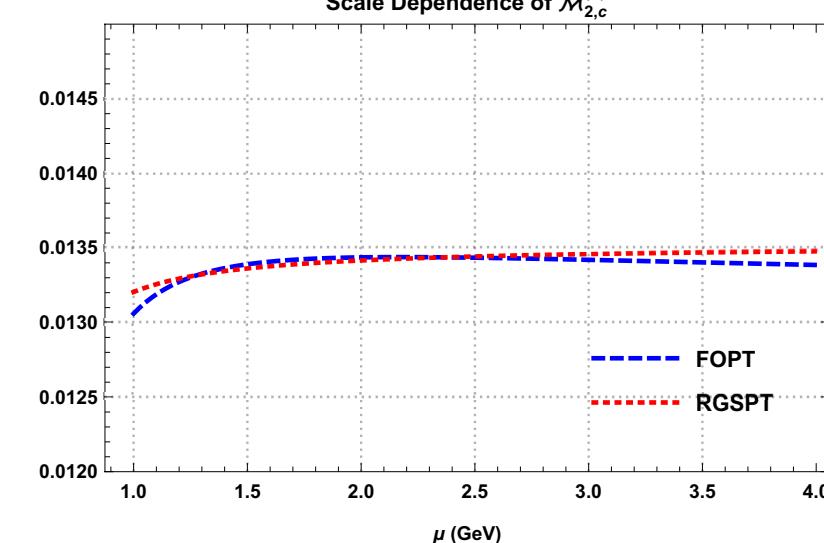
Scale Dependence of $\mathcal{M}_{0,c}^{(P)}$



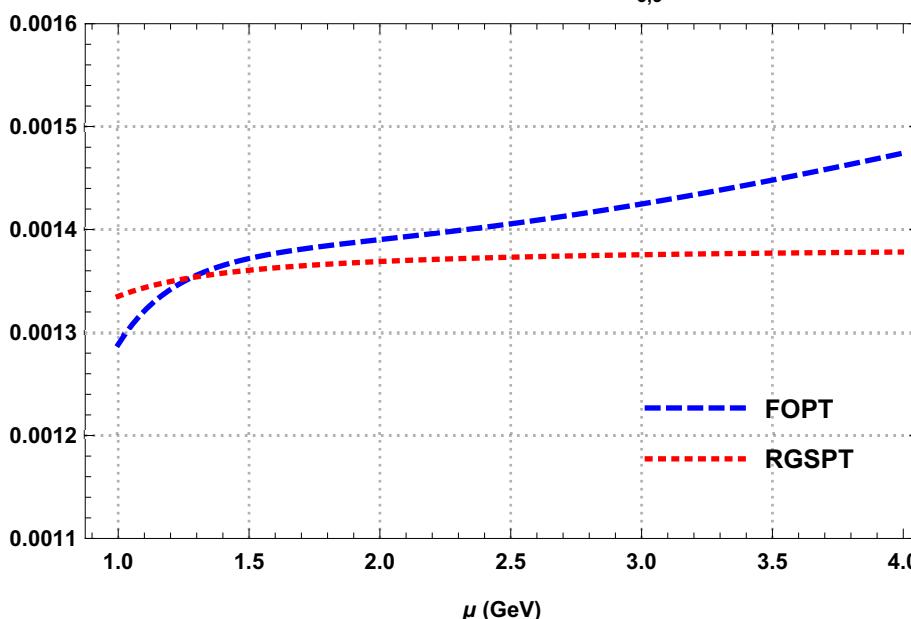
Scale Dependence of $\mathcal{M}_{1,c}^{(P)}$



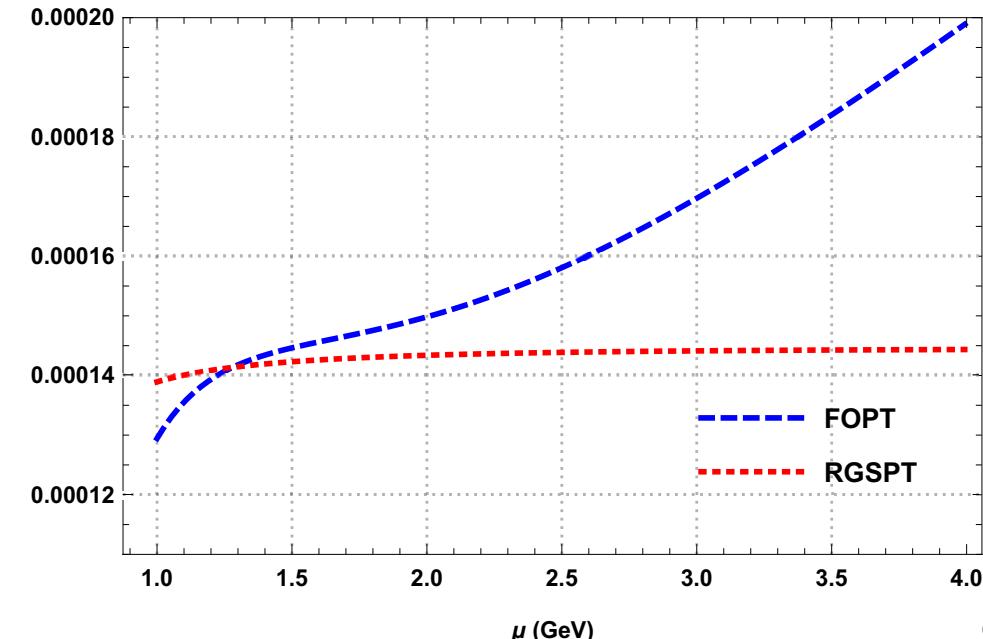
Scale Dependence of $\mathcal{M}_{2,c}^{(P)}$



Scale Dependence of $\mathcal{M}_{3,c}^{(P)}$



Scale Dependence of $\mathcal{M}_{4,c}^{(P)}$



Experimental and lattice moments

Moments	Ref. [1]	Ref. [2]
\mathcal{M}_1^V	2.121 ± 0.036	2.154 ± 0.023
\mathcal{M}_2^V	1.478 ± 0.028	1.490 ± 0.017
\mathcal{M}_3^V	1.302 ± 0.027	1.308 ± 0.016
\mathcal{M}_4^V	1.243 ± 0.028	1.248 ± 0.016

Moments for the vector channel $10^{-n} GeV^{-2n}$.

Moments	Ref. [3]	Ref. [4]	Ref. [5]	Ref. [6]	Ref. [7]
\mathcal{M}_1^P	1.404 ± 0.019	1.395 ± 0.005	1.385 ± 0.007	1.386 ± 0.005	1.387 ± 0.004
\mathcal{M}_2^P	1.359 ± 0.041	1.365 ± 0.012	1.345 ± 0.032	1.349 ± 0.012	1.344 ± 0.010
\mathcal{M}_3^P	1.425 ± 0.059	1.415 ± 0.010	1.406 ± 0.048	1.461 ± 0.050	1.395 ± 0.022

Pseudoscalar moment from lattice QCD (units of $10^{-n} GeV^{-2n}$).}

1. B. Dehnadi, A. H. Hoang, V. Mateu and S. M. Zebarjad, JHEP **09** (2013), 103
2. K. G. Chetyrkin, J. H. Kuhn, A. Maier, P. Maierhofer, P. Marquard, M. Steinhauser and C. Sturm, Phys-RevD.96.116007
3. I. Allison *et al.* [HPQCD], Phys. Rev. D **78** (2008), 054513
4. C. McNeile *et al.*, Phys. Rev. D **82** (2010), 034512
5. K. Nakayama, B. Fahy and S. Hashimoto, Phys. Rev. D **94** (2016) no.5, 054507
6. Y. Maezawa and P. Petreczky, Phys. Rev. D **94** (2016) no.3, 034507
7. P. Petreczky and J. H. Weber, Phys. Rev. D **100** (2019) no.3, 034519

Vector
moments

Pseudoscalar
Moments

Application to charm moments: Vector moments

Sources	Moments	FOPT						RGSPT					
		$m_c(3\text{ GeV})$	Theo. Unc.			total	Exp. Unc.	$m_c(3\text{ GeV})$	Theo. Unc.			total	Exp. Unc.
			α_s	μ	n.p.				μ	n.p.	total		
Ref. [1]	\mathcal{M}_1^V	1005.4(13.9)	3.2	7.6	0.2	8.3	11.2	1000.2(12.3)	3.7	1.9	2.3	4.8	11.3
	\mathcal{M}_2^V	997.2(19.8)	4.7	11.3	14.4	18.9	6.1	988.5(9.2)	5.4	1.6	3.5	6.6	6.3
	\mathcal{M}_3^V	1022.1(127.8)	3.4	41.8	120.6	127.7	4.0	983.4(9.2)	6.7	1.4	3.9	7.8	4.9
	\mathcal{M}_4^V	1077.3(113.6)	1.0	100.5	52.9	113.6	2.8	980.5(8.9)	7.7	0.9	1.8	8.0	3.9
Ref. [2]	\mathcal{M}_1^V	995.4(10.8)	3.3	7.6	0.0	8.3	6.9	990.1(8.5)	3.6	2.0	2.3	4.8	7.0
	\mathcal{M}_2^V	994.6(19.5)	4.7	11.3	14.7	19.1	3.7	985.8(7.7)	5.4	1.6	3.5	6.6	3.8
	\mathcal{M}_3^V	1021.3(126.5)	3.4	41.9	126.5	133.3	2.3	982.3(8.3)	6.7	1.4	3.9	7.8	2.8
	\mathcal{M}_4^V	1076.8(113.8)	1.0	100.7	52.9	113.8	1.6	979.8(8.3)	7.7	0.9	1.8	8.0	2.2

The pole mass of the charm quark is used as input in the non-perturbative condensate terms. $\mu \in [1, 4] \text{ GeV}$

Sources	Moments	FOPT						RGSPT					
		$m_c(3\text{ GeV})$	Theo. Unc.			total	Exp. Unc.	$m_c(3\text{ GeV})$	Theo. Unc.			total	Exp. Unc.
			α_s	μ	n.p.				μ	n.p.	total		
Ref. [1]	\mathcal{M}_1^V	1004.8(13.7)	3.3	7.1	1.4	7.9	11.2	1000.9(12.1)	3.7	1.9	0.9	4.3	11.3
	\mathcal{M}_2^V	989.4(9.2)	5.4	3.5	2.0	6.7	6.3	989.6(8.6)	5.5	1.8	1.0	5.8	6.3
	\mathcal{M}_3^V	990.9(13.5)	5.4	10.9	2.5	12.6	4.7	984.8(8.7)	6.8	2.4	1.0	7.3	4.8
	\mathcal{M}_4^V	1014.5(37.7)	3.8	37.2	2.9	37.5	3.4	980.9(9.8)	8.0	3.9	0.9	9.0	3.9
Ref. [2]	\mathcal{M}_1^V	995.1(10.3)	3.3	6.8	0.7	7.6	7.0	990.8(8.2)	3.8	2.0	0.9	4.3	7.0
	\mathcal{M}_2^V	987.3(7.4)	5.4	3.2	0.8	6.3	3.8	986.9(7.0)	5.5	1.8	1.0	5.9	3.8
	\mathcal{M}_3^V	990.7(12.7)	5.8	10.9	2.7	12.4	2.7	983.7(7.8)	6.8	2.4	1.1	7.3	2.8
	\mathcal{M}_4^V	1014.9(37.0)	3.8	36.7	0.8	36.9	2.0	980.2(9.3)	8.1	3.9	1.0	9.0	2.2

Application to charm moments: pseudoscalar moments

Sources	Moments	FOPT						RGSPt					
		$m_c(3\text{ GeV})$	Theo. Unc.			Exp. Unc.	$m_c(3\text{ GeV})$	Theo. Unc.			Exp. Unc.		
			α_s	μ	n.p.			μ	n.p.	total			
Ref. [3]	\mathcal{M}_1^P	983.6(10.0)	1.1	5.0	2.4	5.7	8.2	989.3(9.0)	1.4	3.5	0.7	3.8	8.1
	\mathcal{M}_2^P	988.3(12.5)	1.7	6.8	3.6	12.5	9.8	990.5(11.4)	1.5	5.8	0.9	6.0	9.7
	\mathcal{M}_3^P	998.9(29.9)	2.2	26.6	10.5	28.6	8.5	985.4(11.6)	3.2	6.4	2.0	7.4	9.0
Ref. [4]	\mathcal{M}_1^P	987.1(6.1)	1.1	5.0	2.3	5.5	2.4	992.8(4.5)	1.4	3.5	3.5	0.7	2.3
	\mathcal{M}_2^P	986.9(8.3)	1.7	6.7	3.6	7.8	2.7	989.1(6.6)	1.5	5.8	0.9	6.0	2.7
	\mathcal{M}_3^P	1000.2(28.6)	2.2	26.5	10.4	28.6	1.4	986.9(7.6)	3.2	6.4	2.0	7.4	1.5
Ref. [5]	\mathcal{M}_1^P	991.7(6.4)	1.1	4.9	2.2	5.5	3.2	997.3(5.0)	1.4	3.5	0.7	3.8	3.2
	\mathcal{M}_2^P	991.5(10.9)	1.6	6.8	3.5	7.9	7.5	993.6(9.6)	1.5	5.8	0.9	6.1	7.5
	\mathcal{M}_3^P	1001.5(29.4)	2.2	26.5	10.3	28.5	7.1	988.2(10.5)	3.2	6.4	2.0	7.4	7.5
Ref. [6]	\mathcal{M}_1^P	991.2(6.0)	1.1	4.9	2.3	5.5	2.4	996.8(4.5)	1.4	3.5	0.7	3.8	2.3
	\mathcal{M}_2^P	990.5(8.3)	1.6	6.8	3.5	7.9	2.8	992.7(6.6)	1.5	5.8	0.9	6.0	2.7
	\mathcal{M}_3^P	993.8(29.7)	2.2	26.7	10.9	28.9	7.0	980.0(10.5)	3.3	6.3	2.0	7.4	7.4
Ref. [7]	\mathcal{M}_1^P	990.6(5.9)	1.1	4.9	2.3	5.5	1.9	996.2(4.2)	1.4	3.5	0.7	3.8	1.9
	\mathcal{M}_2^P	991.6(8.2)	1.6	3.5	7.9	2.3	9.8	993.7(6.5)	1.5	5.8	0.9	6.1	2.3
	\mathcal{M}_3^P	1003.2(28.6)	2.2	26.5	10.2	28.6	3.2	989.9(8.2)	3.2	6.4	2.0	7.4	3.4

Final Value: $m_c(3\text{ GeV}) = 0.9962(42) \text{ GeV}$

$$\implies m_c(m_c) = 1.2811(38) \text{ GeV}$$

PDG Value: $m_c(m_c) = 1.27 \pm 0.02 \text{ GeV}$

REvolver:

A. H. Hoang, C. Lepenik and V. Mateu,
 Comput. Phys. Commun. 270 (2022), 108145
 [arXiv:2102.01085 [hep-ph]].

Strong coupling determination

Vector moments

n	\mathcal{R}_n^V	Refs. [1,2]
1	1.770 ± 0.017	
2	1.1173 ± 0.0023	
3	1.03536 ± 0.00084	

Pseudoscalar moments

Moments	Ref. [3]	Ref. [4]	Ref. [5]	Ref. [6]	Ref. [7]	Ref. [8]
\mathcal{M}_0^P	1.708 ± 0.007	1.708 ± 0.005	–	1.699 ± 0.008	1.705 ± 0.005	1.7037 ± 0.0027
\mathcal{R}_1^P	1.197 ± 0.004	–	1.188 ± 0.004	1.199 ± 0.004	1.1886 ± 0.013	1.1881 ± 0.0007
\mathcal{R}_2^P	1.033 ± 0.004	–	1.0341 ± 0.0018	1.0344 ± 0.0013	1.0324 ± 0.0016	–

These \mathcal{M}_n^P is in the units of $10^{-n} \text{ GeV}^{-2n}$

Refs.:

1. D. Boito and V. Mateu, Phys. Lett. B **806** (2020), 135482
2. D. Boito and V. Mateu, JHEP **03** (2020), 094

3. I. Allison *et al.* [HPQCD], Phys. Rev. D **78** (2008), 054513
4. C. McNeile *et al.*, Phys. Rev. D **82** (2010), 034512
5. Y. Maezawa and P. Petreczky, Phys. Rev. D **94** (2016) no.3, 034507
6. K. Nakayama, B. Fahy and S. Hashimoto, Phys. Rev. D **94** (2016) no.5, 054507
7. P. Petreczky and J. H. Weber, Phys. Rev. D **100** (2019) no.3, 034519
8. P. Petreczky and J. H. Weber, Eur. Phys. J. C **82** (2022) no.1, 64

Vector moments

Pseudoscalar
Moments

Strong coupling determination: Vector moments

$\overline{\text{MS}}$ value of the quark mass in the condensate terms:

$$m_c(q) = \overline{m}_c \int_{x(\overline{m}_c)}^{x(q)} dx e^{\left(\frac{\gamma(x)}{\beta(x)} \right)}$$

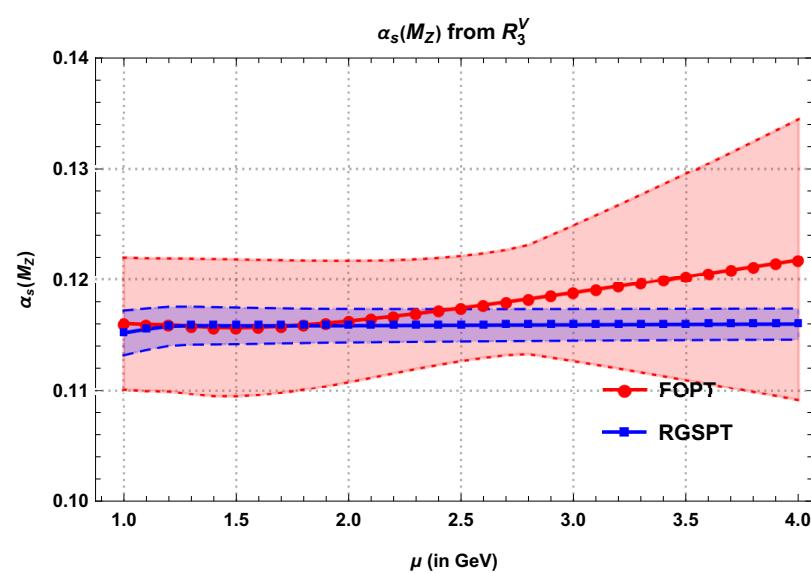
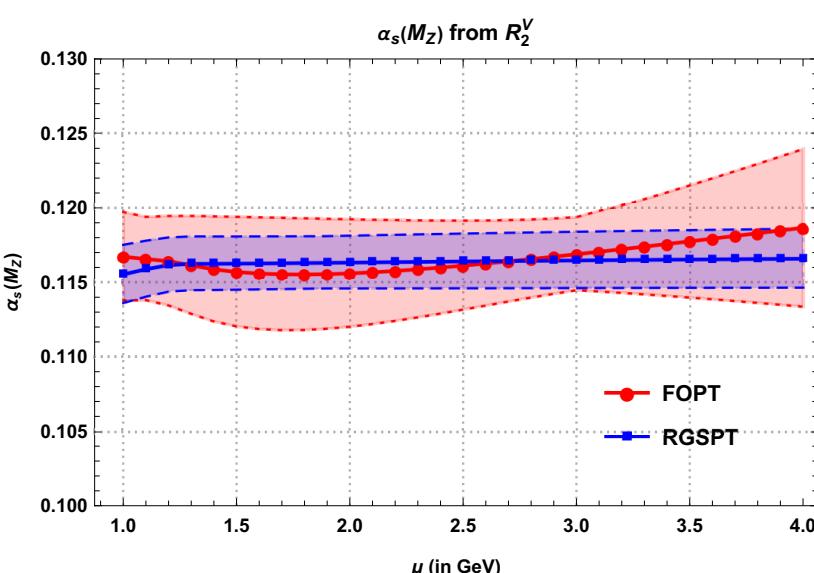
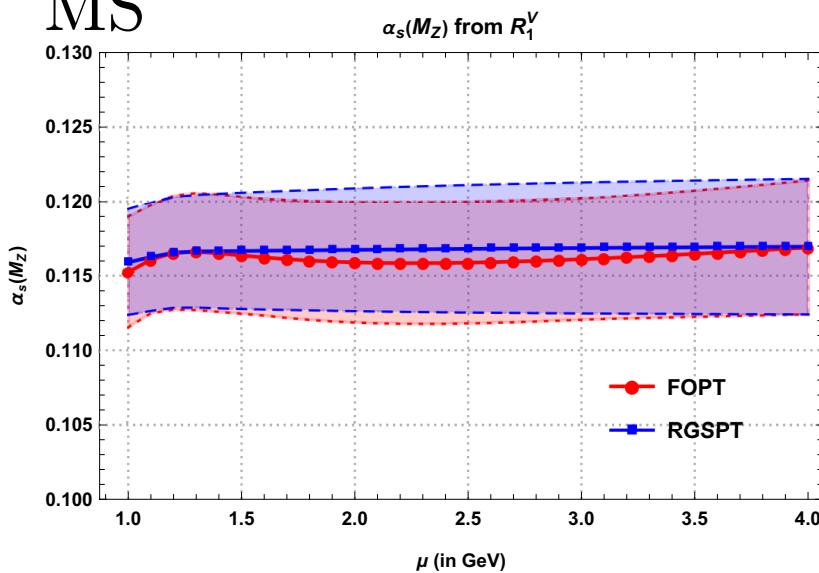
Moment	FOPT					RGSPT						
	$\alpha_s(M_Z)$	Theo. Unc.				Exp. Unc.	$\alpha_s(M_Z)$	Theo. Unc.				Exp. Unc.
		m_c	μ	n.p.	total			m_c	μ	n.p.	total	
\mathcal{R}_1^V	0.1167(39)	3	13	8	16	36	0.1167(38)	3	7	8	11	36
\mathcal{R}_2^V	0.1163(31)	4	27	12	29	11	0.1163(18)	3	8	12	15	11
\mathcal{R}_3^V	0.1159(60)	4	58	14	60	5	0.1159(17)	3	6	14	16	5

Pole value of the quark mass in the condensate terms:

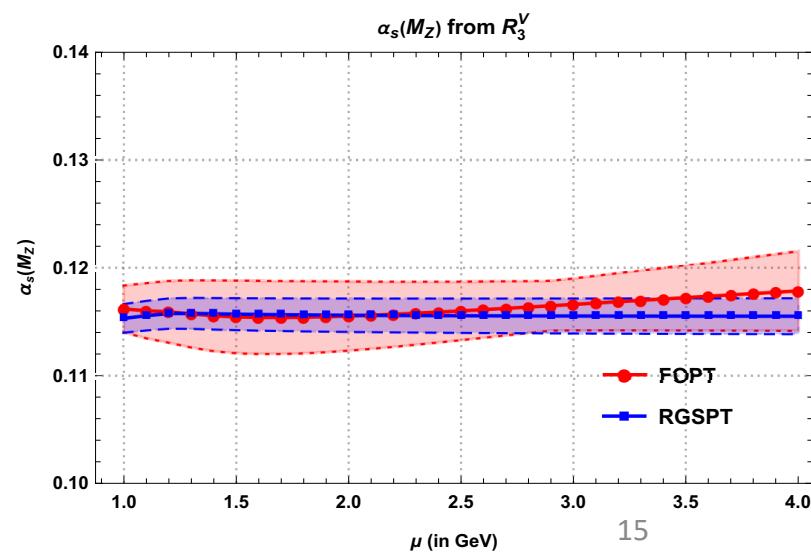
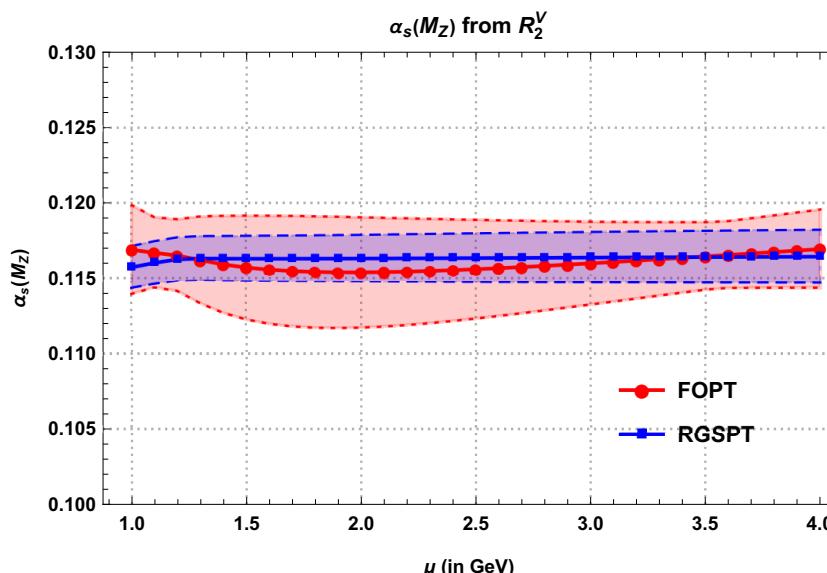
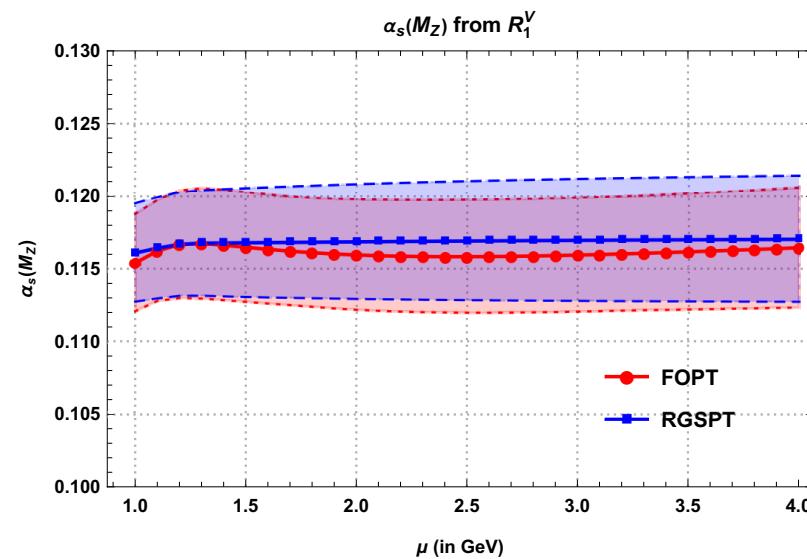
Moment	FOPT					RGSPT						
	$\alpha_s(M_Z)$	Theo. Unc.				Exp. Unc.	$\alpha_s(M_Z)$	Theo. Unc.				Exp. Unc.
		m_c	μ	n.p.	total			m_c	μ	n.p.	total	
\mathcal{R}_1^V	0.1169(38)	3	13	5	15	35	0.1169(36)	2	6	5	8	35
\mathcal{R}_2^V	0.1164(28)	4	24	9	26	10	0.1164(15)	2	5	6	8	10
\mathcal{R}_3^V	0.1159(30)	3	27	13	30	5	0.1159(14)	2	2	6	6	5

Behavior of strong coupling determination from moments

MS



Pole mass in the condensate terms



Strong coupling determination from lattice moments

Sources	Moments	FOPT						RGSPt					
		$\alpha_s(M_Z)$	Theo. Unc.				Exp. Unc.	$\alpha_s(M_Z)$	Theo. Unc.				Exp. Unc.
			m_c	μ	n.p.	total			m_c	μ	n.p.	total	
Ref. [3]	\mathcal{M}_0^P	0.1172(20)	3	19	3	19	6	0.1172(8)	3	3	3	5	6
	\mathcal{R}_1^P	0.1182(43)	4	42	5	43	6	0.1181(15)	3	12	5	13	6
	\mathcal{R}_2^P	0.1150(53)	4	50	9	51	15	0.1149(18)	3	7	9	11	15
Ref. [4]	\mathcal{M}_0^P	0.1172(20)	3	19	3	19	5	0.1172(7)	3	3	3	5	5
Ref. [5]	\mathcal{M}_0^P	0.1168(48)	3	8	6	47	7	0.1168(13)	3	9	6	11	7
	\mathcal{R}_1^P	0.1152(51)	4	50	8	50	6	0.1152(13)	3	7	8	11	6
Ref. [6]	\mathcal{M}_0^P	0.1164(20)	3	18	4	19	7	0.1164(9)	3	3	4	5	7
	\mathcal{R}_1^P	0.1182(43)	4	42	5	43	6	0.1184(15)	3	13	5	14	6
	\mathcal{R}_2^P	0.1153(50)	4	49	8	50	5	0.1153(12)	3	7	8	7	5
Ref. [7]	\mathcal{M}_0^P	0.1169(20)	3	19	3	19	5	0.1169(7)	3	13	3	5	5
	\mathcal{R}_1^P	0.1169(47)	4	47	6	47	2	0.1169(12)	3	10	6	12	2
	\mathcal{R}_2^P	0.1146(53)	4	52	9	53	6	0.1146(13)	3	6	9	11	6
Ref. [8]	\mathcal{M}_0^P	0.1168(19)	3	19	3	19	2	0.1168(13)	3	9	6	11	7
	\mathcal{R}_1^P	0.1168(47)	4	47	6	47	1	0.1168(12)	3	1	6	11	1

Refs.:

3 . I. Allison *et al.* [HPQCD], Phys. Rev. D **78** (2008), 054513

4 . C. McNeile, C. T. H. Davies, E. Follana, K. Hornbostel and G. P. Lepage, Phys. Rev. D **82** (2010), 034512

5 . Y. Maezawa and P. Petreczky, Phys. Rev. D **94** (2016) no.3, 034507

6 . K. Nakayama, B. Fahy and S. Hashimoto, Phys. Rev. D **94** (2016) no.5, 054507

7 . P. Petreczky and J. H. Weber, Phys. Rev. D **100** (2019) no.3, 034519

8 . P. Petreczky and J. H. Weber, Eur. Phys. J. C **82** (2022) no.1, 64

$$\text{Final value: } \alpha_s(M_Z) = \{0.1172(7), 0.1169(7)\}$$

$$\implies \alpha_s(M_Z) = 0.1171(7)$$

$$\text{PDG value: } \alpha_s(M_Z) = 0.1179(9)$$

Future directions

1. Proper correlation of different moments.
2. Truncation uncertainty.
3. Improved continuum contributions.

$$\mathcal{M}_q^{V,n} = \int_0^\infty \frac{ds}{s^{n+1}} R_{q\bar{q}}(s) = \int_0^{s_{\text{th.}}} \frac{ds}{s^{n+1}} R_{q\bar{q}}^{\text{exp.}}(s) + \int_{s_{\text{th.}}}^\infty \frac{ds}{s^{n+1}} R_{q\bar{q}}^{\text{cont.}}(s)$$

$$R_{\bar{q}q}^{\text{cont.}} \equiv 12\pi \text{Im} \Pi^V(s + i\epsilon) = 12\pi \times \frac{1}{2i} (\Pi^V(s + i\epsilon) - \Pi^V(s - i\epsilon))$$

Analytic continuation: M. S. A. Alam Khan, “Renormalization group summation and analytic continuation from spacelike to timeline regions,” Phys. Rev. D 108 (2023) no.1, 014028 [arXiv:2306.10262 [hep-ph]].

Other applications in the determination of α_S using finite energy sum rules as well as Borel-Laplace and Laplace sum rules for $e^+e^- \rightarrow \text{hadrons}$

Thank you
