Determination of α_S from the Z-boson transverse-momentum distribution at the Tevatron

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Based on:

Stefano Camarda, G. F., Matthias Schott

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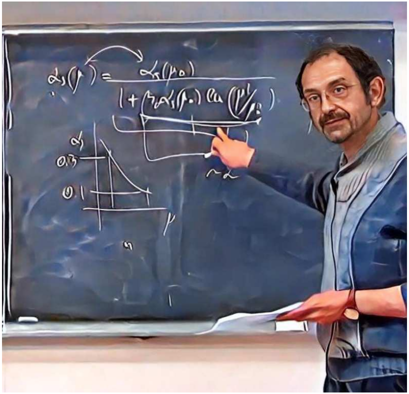
Workshop α_S-2024

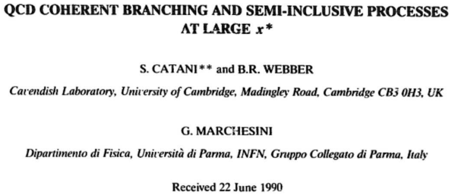
ECT* – Trento – 7/2/2024

Stefano Catani (1958-2024)



Wonderful person, outstanding physicist

The idea: α_S from semi-inclusive processes


QCD COHERENT BRANCHING AND SEMI-INCLUSIVE PROCESSES AT LARGE x*

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$$\alpha_s^{(MC)} = \alpha_s^{(\overline{MS})}\left(1 + K\frac{\alpha_s^{(\overline{MS})}}{2\pi}\right),$$

$$\Lambda_{MC} = \Lambda_{\overline{MS}} \exp(K/4\pi\beta_0)$$
$$\simeq 1.569\,\Lambda_{\overline{MS}} \quad \text{for } N_f = 5.$$

In this paper we have studied [...] the next-to-leading logarithmic terms in **semi-inclusive hard processes such as the DIS and DY processes** *at large x. Since the Monte Carlo algorithm with these improvements is accurate to next-to-leading order in the large-x region,* **it can be used to determine the fundamental QCD scale $\Lambda_{\overline{MS}}$**

The idea: α_S from semi-inclusive processes

Advantages:

- **higher sensitivity to α_S** w.r.t. *inclusive* observables;
- calculable at **higher theoretical accuracy** w.r.t. *exclusive* observables.

Challenges:

- sensitivity to **infrared (Sudakov) logs**;
- sensitivity **non perturbative QCD** effects.

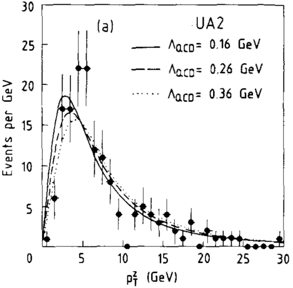
Classical semi-inclusive obs. at hadron colliders:
high invariant-mass Drell–Yan lepton pair
at small transverse-momentum (q_T).

α_S from Z-boson q_T distribution

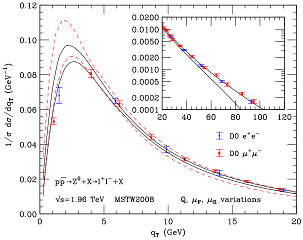
this talk

M. Schott talk

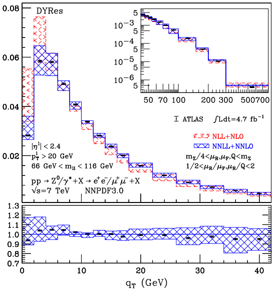
$Sp\bar{p}S$ ($\sqrt{s} = 0.63$ TeV) **Tevatron ($\sqrt{s} = 1.96$ TeV)** **LHC ($\sqrt{s} = 7 - 8$ TeV)**



[UA2 Coll.('92)]
compared with
[Altarelli et al.('84)]

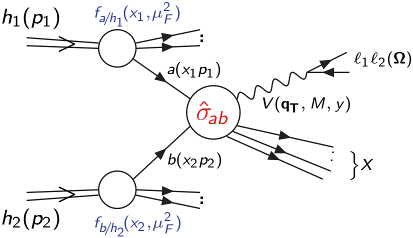


[D0 Coll.('08,'10)]
compared with
[Catani et al.('10)]



[ATLAS Coll.('14)]
compared with
[Catani et al.('15)]

Drell–Yan q_T distribution



$$\mathbf{h_1(p_1) + h_2(p_2) \;\rightarrow\; V + X \;\rightarrow\; \ell_1 + \ell_2 + X}$$

where $\quad V = Z^0/\gamma^*, W^\pm$

QCD factorization formula:

$$\frac{d\sigma}{dq_T^2} = \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 \, f_{a/h_1}(x_1, \mu_F^2) \, f_{b/h_2}(x_2, \mu_F^2) \, \frac{d\hat{\sigma}_{ab}}{dq_T^2}(\alpha_S(\mu_R^2), \mu_R^2, \mu_F^2).$$

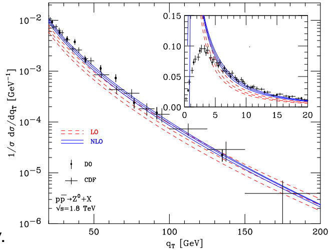
Fixed-order perturbative expansion reliable only for $q_T \sim M$. When $q_T \ll M$:

$$\int_0^{q_T^2} d\bar{q}_T^2 \frac{d\hat{\sigma}_{q\bar{q}}}{d\bar{q}_T^2} \sim 1 + \alpha_S \left[c_{12} L_{q_T}^2 + c_{11} L_{q_T} + \cdots \right]$$

$$+ \alpha_S^2 \left[c_{24} L_{q_T}^4 + \cdots + c_{21} L_{q_T} + \cdots \right] + \mathcal{O}(\alpha_S^3)$$

with $\alpha_S^n L_{q_T}^m \equiv \alpha_S^n \log^m(M^2/q_T^2) \gtrsim 1$.

Resummation of logarithmic corrections mandatory.



$\mathbf{q_T}$ resummation in QCD

[Catani,de Florian,Grazzini('01)] [Bozzi,Catani,de Florian,Grazzini('03,'06)]

$$\frac{d\hat{\sigma}}{dq_T^2} = \frac{d\hat{\sigma}^{(res)}}{dq_T^2} + \frac{d\hat{\sigma}^{(fin)}}{dq_T^2};$$

In the impact parameter space: $q_T \ll M \Leftrightarrow Mb \gg 1, \quad \log M/q_T \gg 1 \Leftrightarrow \log Mb \gg 1$

$$\frac{d\hat{\sigma}^{(res)}}{dq_T^2} = \frac{M^2}{\hat{s}} \int \frac{d^2\mathbf{b}}{4\pi} e^{i\mathbf{b}\cdot\mathbf{q_T}} \mathcal{W}(b, M),$$

In the Mellin space (with respect to $z = M^2/\hat{s}$) we have:

$$\mathcal{W}_N(b,M) = \mathcal{H}_N(\alpha_S) \times \exp\left\{\mathcal{G}_N(\alpha_S, L)\right\}$$

with $\quad L \equiv \log(M^2 b^2) \quad$ and $\quad \alpha_S L \sim 1$

$$\mathcal{G}(\alpha_S, L) = L\, g^{(1)}(\alpha_S L) + g^{(2)}(\alpha_S L) + \frac{\alpha_S}{\pi} g^{(3)}(\alpha_S L) + \cdots \qquad \mathcal{H}(\alpha_S) = \hat{\sigma}^{(0)} \left(1 + \frac{\alpha_S}{\pi} \mathcal{H}^{(1)} + \left(\frac{\alpha_S}{\pi}\right)^2 \mathcal{H}^{(2)} + \cdots\right)$$

LL ($\sim \alpha_S^n L^{n+1}$): $g^{(1)}$, $(\hat{\sigma}^{(0)})$; NLL ($\sim \alpha_S^n L^n$): $g^{(2)}$, $\mathcal{H}^{(1)}$; $\cdots$ N^{k}LL ($\sim \alpha_S^n L^{n+k-1}$): $g^{(k+1)}$, $\mathcal{H}^{(k)}$;

Resummed result at small q_T *matched* with corresponding fixed "finite" part at large q_T: *uniform accuracy* for $q_T \ll M$ and $q_T \sim M$.

- Resummed effects exponentiated in a universal of Sudakov form factor, process-dependence factorized in the hard-virtual factor $H_c^F(\alpha_S)$ via all-order formula [Catani, Cieri, de Florian, G.F., Grazzini('14)].
- Resummation performed at partonic cross section level: (collinear) PDF evaluated at $\mu_F \sim M$, $f_N(b_0^2/b^2) = \exp\left\{-\int_{b_0^2/b^2}^{\mu_F^2} \frac{dq^2}{q^2}\,\gamma_N(\alpha_S(q^2))\right\} f_N(\mu_F^2)$: no PDF extrapolation in the non perturbative region, study of μ_R and μ_F dependence as in fixed-order calculations.
- No need for NP models: Landau singularity of α_S regularized using a *Minimal Prescription* without power-suppressed corrections [Laenen et al.('00)],[Catani et al.('96)].
- Introduction of resummation scale $Q \sim M$: variations give an estimate of the uncertainty from uncalculated logarithmic corrections.

$$\ln\left(M^2 b^2\right) \rightarrow \ln\left(Q^2 b^2\right) + \ln\left(M^2/Q^2\right)$$

- Perturbative unitarity constraint: recover *exactly* the total cross-section (upon integration on q_T)

$$\ln\left(Q^2 b^2\right) \rightarrow \widetilde{L} \equiv \ln\left(Q^2 b^2 + 1\right) \quad \Rightarrow \quad \exp\left\{\alpha_S^n \widetilde{L}^k\right\}\Big|_{b=0} = 1 \ \Rightarrow \ \int_0^\infty dq_T^2 \left(\frac{d\hat{\sigma}}{dq_T^2}\right) = \hat{\sigma}^{(tot)};$$

- General procedure to treat the q_T recoil [Catani, de Florian, G.F., Grazzini('15)]:

$$\frac{d\hat{\sigma}^{(0)}}{d\mathbf{\Omega}} = \hat{\sigma}^{(0)}(M^2)\, F(\mathbf{q_T}; M^2, \mathbf{\Omega}) \ \ \textit{with} \ \ F(\mathbf{q_T}; M^2, \mathbf{\Omega}) = F(\mathbf{0}; M^2, \mathbf{\Omega}) + \mathcal{O}(\mathbf{q_T^2}/M^2)$$

q_T resummation: perturbative accuracy

- Formalism implemented in numerically efficient and publicly available code:

  DYTurbo: computes resummed and fixed-order fiducial cross section and related distributions it retains full kinematics of the vector boson and of its leptonic decay products [Camarda,Boonekamp,Bozzi,Catani,Cieri,Cuth,G.F.,de Florian,Glazov, Grazzini,Vincter,Schott('20)]

  https://dyturbo.hepforge.org.

- We have explicitly included in DYTurbo up to:
  - N^{4}LL logarithmic contributions to all orders (i.e. up to $\exp(\sim\alpha_S^n L^{n-3})$);
  - Approximated N^{4}LO corrections (i.e. up to $\mathcal{O}(\alpha_S^4)$) at small q_T;
  - NLO corrections (i.e. up to $\mathcal{O}(\alpha_S^2)$) at large q_T;
- Matching with NNLO corrections (i.e. up to $\mathcal{O}(\alpha_S^3)$) at large q_T from results in [Boughezal et al.('16)], [Gehrmann-De Ridder et al.('16)], [MCFM ('23)];
- Results up to N^{3}LO (i.e. up to $\mathcal{O}(\alpha_S^3)$) recovered for the total cross section (from unitarity).

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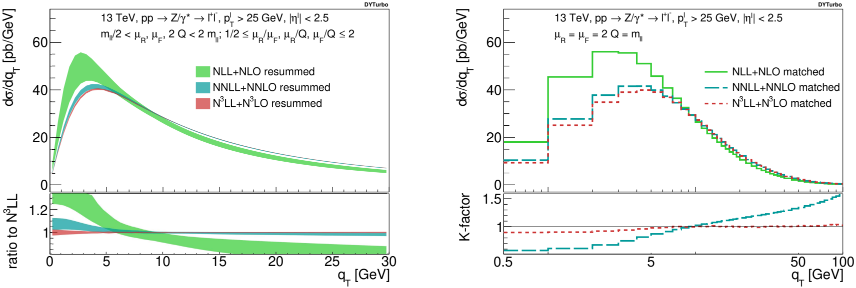
- Results up to N^3LO (i.e. up to $\mathcal{O}(\alpha_S^3)$) recovered for the total cross section (from unitarity).

Z/γ^* production at N^{3}LL+N^{3}LO (resummed and matched)

[Camarda,Cieri,G.F.('21)]

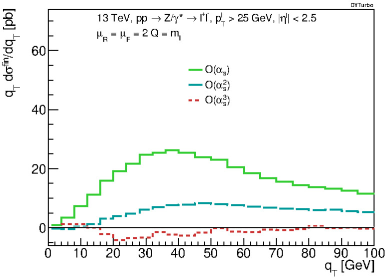


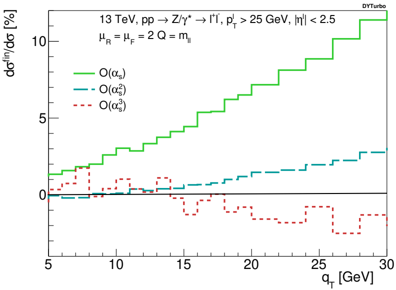
DYTurbo results. Resummed (left) and matched (right) NLL, NNLL and N^{3}LL bands for Z/γ^* q_T spectrum.

Lower panel: ratio with respect to the N^{3}LL central value.

Z/γ^* production: finite part

[Camarda,Cieri,G.F.('21)]

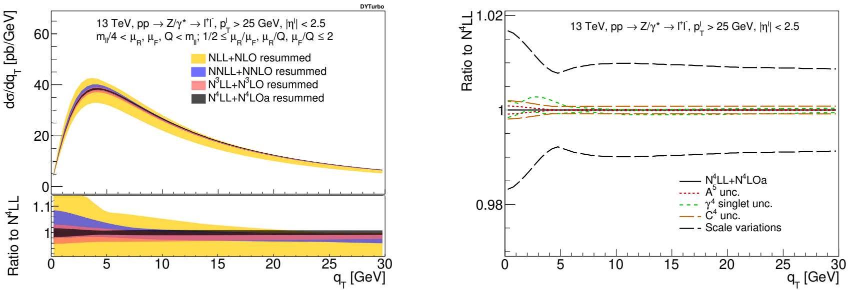




Finite part at $\mathcal{O}(\alpha_S)$, $\mathcal{O}(\alpha_S^2)$ and $\mathcal{O}(\alpha_S^3)$ (left) and ratio wrt matched results (right).

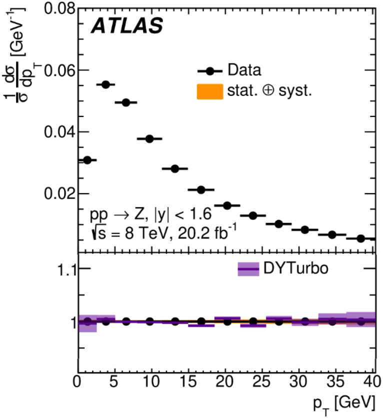
Z/γ^* production at N^{4}LL+N^{4}LOa resummed

[Camarda,Cieri,G.F.('23)]



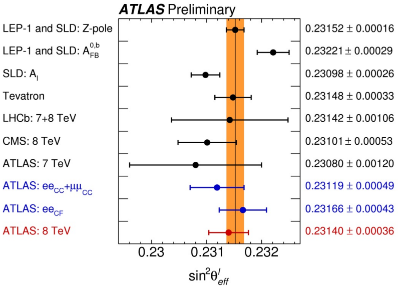
DYTurbo results. Left: Resummed NLL, NNLL, N^{3}LL and N^{4}LLa bands for Z/γ^* (left). Right: Uncertainties from approximations of the perturbative coefficients at N4LL+N4LOa compared to scale variations.

Z/γ^* production theory vs data

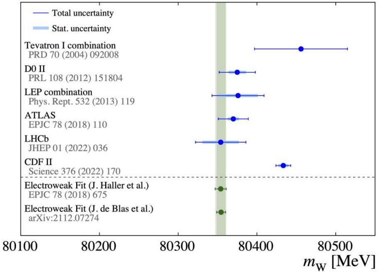


DYTurbo results at N^{4}LLa accuracy compared with data [ATLAS Coll.('23)].
Time performance of $\mathcal{O}$(*seconds*): (with exception of V+jet term with fiducial lepton cuts).

Modelling Z (and W) production for $\sin^2\theta^l_{eff}$ and M_W determinations



Comparison of the measurements of the $\sin^2\theta^l_{eff}$.



Measured values of M_W compared with the prediction of from the global electroweak fit

Combining QED and QCD q_T resummation

[Cieri,G.F.,Sborlini('18)]

We start considering QED contributions to the q_T spectrum in the case of colourless and **neutral** high mass systems, e.g. on-shell Z boson production

$$h_1 + h_2 \;\rightarrow\; Z^0 + X$$

In the impact parameter and Mellin spaces resummed partonic cross section reads:

$$\mathcal{W}_N(b,M) = \hat{\sigma}^{(0)}\,\mathcal{H}'_N(\alpha_S,\alpha) \times \exp\left\{\mathcal{G}'_N(\alpha_S,\alpha,L)\right\}$$

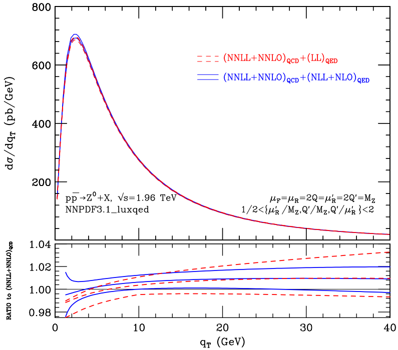
$$\mathcal{G}'(\alpha_S,\alpha,L) = \mathcal{G}(\alpha_S,L) + L\,g'^{(1)}(\alpha L) + g'^{(2)}(\alpha L) + \sum_{n=3}^{\infty}\left(\frac{\alpha}{\pi}\right)^{n-2} g'^{(n)}(\alpha L)$$

$$+\, g'^{(1,1)}(\alpha_S L,\alpha L) + \sum_{\substack{n,m=1\\ n+m\neq 2}}^{\infty}\left(\frac{\alpha_S}{\pi}\right)^{n-1}\left(\frac{\alpha}{\pi}\right)^{m-1} g_N'^{(n,m)}(\alpha_S L,\alpha L)$$

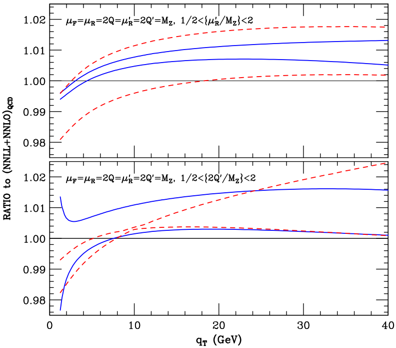
$$\mathcal{H}'(\alpha_S,\alpha) \;=\; \mathcal{H}(\alpha_S) + \frac{\alpha}{\pi}\mathcal{H}'^{(1)} + \sum_{n=2}^{\infty}\left(\frac{\alpha}{\pi}\right)^{n}\mathcal{H}_N'^{(n)} + \sum_{n,m=1}^{\infty}\left(\frac{\alpha_S}{\pi}\right)^{n}\left(\frac{\alpha}{\pi}\right)^{m}\mathcal{H}_N'^{F\,(n,m)}$$

LL QED ($\sim\alpha^n L^{n+1}$): $g'^{(1)}$; NLL QED ($\sim\alpha^n L^n$): $g'^{(2)}$, $\mathcal{H}'^{(1)}$;
LL mixed QCD-QED ($\sim\alpha_S^n\alpha^n L^{2n}$): $g'^{(1,1)}$;

Combined QED and QCD q_T resummation for Z production at the Tevatron [Cieri,G.F.,Sborlini ('18)]

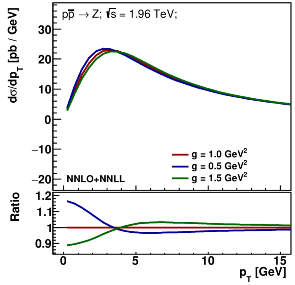


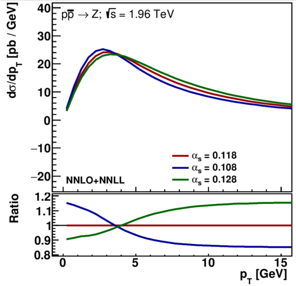
Z qT spectrum at the LHC. NNLL+NNLO QCD combined with the LL (red dashed) and NLL+NLO (blue solid) QED with corresponding QED uncertainty bands.



Ratio of the resummation (upper panel) and renormalization (lower panel) QED scale-dependent results with respect to the central value NNLL+NNLO QCD result.

Non perturbative effects



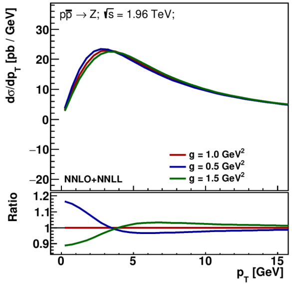


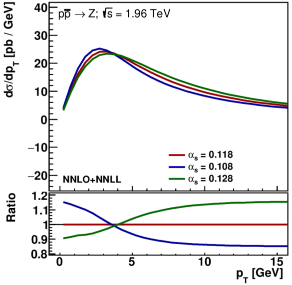
- Up to now discussed result in a complete perturbative framework (except for PDFs).
- Non perturbative *intrinsic* k_T effects parametrized by a NP form factor $S_{NP} = \exp\{-gb^2\}$ with $0 < g < 1.2\, GeV^2$:

$$\exp\{\mathcal{G}_N(\alpha_S, \widetilde{L})\} \quad \to \quad \exp\{\mathcal{G}_N(\alpha_S, \widetilde{L})\}\, S_{NP}$$

- NP effects increase the hardness of the q_T spectrum at small values of q_T. Non trivial interplay of perturbative and NP effects.
- However possible to disentangle the effects: scale of the NP effects is $\langle q_T \rangle \sim 1\, GeV$ ($g \sim 0.5 GeV^2$), scale of "soft gluon" recoil is $\langle q_T \rangle \sim 10\, GeV$.

Non perturbative effects





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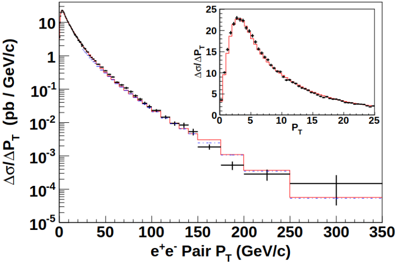
Z-boson q_T measurement at CDF

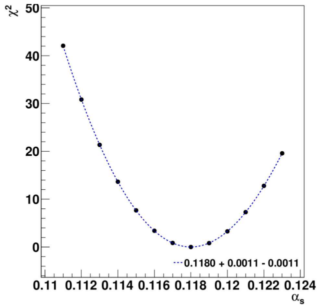
The CDF measurement of $Z/\gamma^* \to e^+e^-$ ($\sqrt{s} = 1.96\,TeV$ with $\int \mathcal{L} = 2.1 fb^{-1}$) [CDF Coll.('10)] is ideal for $\alpha_S(m_Z)$ determination.



[CDF Coll.('10)]

- Measurement in full-lepton phase space with small extrapolation using angular coefficients method ⇒ allows fast analytic predictions with DYTurbo.
- $p\bar{p}$ collisions: small contribution from heavy-flavour in initial state (0.4% $b\bar{b} \to Z$, 1.3% $c\bar{c} \to Z$). Quark mass effects negligible.
- Low pile-up and good electron resolution. Fine q_T bins ($0.5\,GeV$) with relatively small bin-to-bin correlations.

Methodology



- DYTurbo interfaced to xFitter.
  Fit region: $Z\ q_T < 30\,GeV$, predictions at $N^3LL+\mathcal{O}(\alpha_S^3)$ (i.e. N^3LL+N^3LO at low q_T) with NNPDF4.0 PDF at NNLO.
- Defined χ^2 with experimental (β_{exp}) and PDFs (β_{th}) uncertainties (equivalent to including the new dataset in the PDF using profiling/reweighting).
- The non-perturbative form factor is $S_{NP} = \exp\{-gb^2\}$ with g left free in the fit.

$$\chi^2(\beta_{\rm exp},\beta_{\rm th}) = \sum_{i=1}^{N_{\rm data}} \frac{\left(\sigma_i^{\rm exp} + \sum_j \Gamma_{ij}^{\rm exp}\beta_{j,\rm exp} - \sigma_i^{\rm th} - \sum_k \Gamma_{ik}^{\rm th}\beta_{k,\rm th}\right)^2}{\Delta_i^2} + \sum_j \beta_{j,\rm exp}^2 + \sum_k \beta_{k,\rm th}^2 \,.$$

Theory uncertainties

	PDF fit	Hessian profiling
$\alpha_S(m_Z)$	0.1188 ± 0.0008	0.1184 ± 0.0006
g [GeV^2]	0.69 ± 0.05	0.71 ± 0.05
Dataset	χ^2/points	χ^2/points
NC DIS H1-ZEUS e^+p	955/905	
CC DIS H1-ZEUS e^+p	46/39	
NC DIS H1-ZEUS e^-p	219/159	
CC DIS H1-ZEUS e^-p	53/42	
H1-ZEUS correlated χ^2	91	
CDF Z p_T	41/55	40/55
Total	1405 / 1184	

- Bias from α_S-PDFs correlations [Forte,Kassabov('20)] $\rightarrow$ PDFs refitted.
- Other PDF sets considered. CT18, CT18Z, MSHT20, HERAPDF2.0, ABMP16 (NNLO) and MSHT20an3lo. The midpoint value is the nominal result and the PDF envelope as an additional uncertainty.
- Uncertainty from missing higher orders: $m_{ll}/2 < \{\mu_R, \mu_F, Q\} < 2m_{ll}$ with $0.5 < \{\mu_R/\mu_F, \mu_R/Q, \mu_F/Q\} < 2$.
- NP effects: b_*-pr. $b_* = b/\sqrt{1 + b^2/b^2_{lim}}$ ($b_{lim} = 2 - 3$ GeV^{-1}) and *minimal* pr. ($b_{lim} \rightarrow \infty$); quartic term $\exp(-qb^4)$ and different parametrization of S_{NP} [Collins,Rogers('15)].
- Uncerainty from finite component at $\mathcal{O}(\alpha_S^3)$.
- Check with D0 data and fit boundaries.

Theory uncertainties

	$\alpha_S(m_Z)$	g [GeV^2]	χ^2/dof
NNPDF4.0	0.1192 ± 0.0008	0.66 ± 0.05	41/53
CT18	0.1189 ± 0.0010	0.67 ± 0.05	40/53
CT18Z	0.1198 ± 0.0009	0.62 ± 0.05	41/53
MSHT20	0.1185 ± 0.0009	0.72 ± 0.05	40/53
HERAPDF2.0	0.1188 ± 0.0008	0.69 ± 0.05	40/53
ABMP16	0.1185 ± 0.0007	0.62 ± 0.05	42/53
MSHT20an3lo (N^{4}LL)	0.1184 ± 0.0009	0.73 ± 0.05	40/53
PDF fit	0.1184 ± 0.0006	0.71 ± 0.05	1405/1184

- Bias from α_S-PDFs correlations [Forte,Kassabov('20)] $\to$ PDFs refitted.
- Other PDF sets considered. CT18, CT18Z, MSHT20, HERAPDF2.0, ABMP16 (NNLO) and MSHT20an3lo. The midpoint value is the nominal result and the PDF envelope as an additional uncertainty.
- Uncertainty from missing higher orders: $m_{ll}/2 < \{\mu_R, \mu_F, Q\} < 2m_{ll}$ with $0.5 < \{\mu_R/\mu_F, \mu_R/Q, \mu_F/Q\} < 2$.
- NP effects: b_*-pr. $b_* = b/\sqrt{1 + b^2/b_{lim}^2}$ ($b_{lim} = 2 - 3$ GeV^{-1}) and *minimal* pr. ($b_{lim} \to \infty$); quartic term $\exp(-qb^4)$ and different parametrization of S_{NP} [Collins,Rogers('15)].
- Unceraintny from finite component at $\mathcal{O}(\alpha_S^3)$.
- Check with D0 data and fit boundaries.

Theory uncertainties

$\mu_R/m_{\ell\ell}$	$\mu_F/m_{\ell\ell}$	$Q/m_{\ell\ell}$	$\alpha_S(m_Z)$	g [GeV^2]	χ^2/dof
1	1	1	0.1192 ± 0.0008	0.66 ± 0.05	41/53
1	1	2	0.1183 ± 0.0007	0.77 ± 0.05	40/53
1	1	0.5	0.1196 ± 0.0008	0.57 ± 0.05	42/53
1	2	1	0.1194 ± 0.0008	0.66 ± 0.05	41/53
1	2	2	0.1183 ± 0.0007	0.77 ± 0.05	41/53
1	0.5	1	0.1193 ± 0.0008	0.68 ± 0.05	42/53
1	0.5	0.5	0.1196 ± 0.0008	0.59 ± 0.05	42/53
2	1	1	0.1193 ± 0.0008	0.67 ± 0.05	42/53
2	1	2	0.1194 ± 0.0008	0.70 ± 0.05	41/53
2	2	1	0.1192 ± 0.0008	0.65 ± 0.05	42/53
2	2	2	0.1192 ± 0.0008	0.67 ± 0.05	41/53
0.5	1	1	0.1184 ± 0.0007	0.75 ± 0.05	42/53
0.5	1	0.5	0.1192 ± 0.0007	0.64 ± 0.05	41/53
0.5	0.5	1	0.1183 ± 0.0007	0.75 ± 0.05	42/53
0.5	0.5	0.5	0.1192 ± 0.0007	0.64 ± 0.05	42/53

- Bias from α_S-PDFs correlations [Forte,Kassabov('20)] $\rightarrow$ PDFs refitted.
- Other PDF sets considered. CT18, CT18Z, MSHT20, HERAPDF2.0, ABMP16 (NNLO) and MSHT20an3lo. The midpoint value is the nominal result and the PDF envelope as an additional uncertainty.
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- NP effects: b_*-pr. $b_* = b/\sqrt{1 + b^2/b_{lim}^2}$ ($b_{lim} = 2 - 3$ GeV^{-1}) and *minimal* pr. ($b_{lim} \rightarrow \infty$); quartic term $\exp(-qb^4)$ and different parametrization of S_{NP} [Collins,Rogers('15)].
- Unceraint from finite component at $\mathcal{O}(\alpha_S^3)$.
- Check with D0 data and fit boundaries.

Theory uncertainties

	$\alpha_S(m_Z)$	g [GeV^2]
$b_{\lim} = 2$ GeV^{-1}	0.1187 ± 0.0007	0.83 ± 0.05
$b_{\lim} \to \infty$	0.1199 ± 0.0008	0.42 ± 0.05
g_k	0.1186 ± 0.0008	0.65 ± 0.05
$q = 0.1$ GeV^4	0.1197 ± 0.0008	0.51 ± 0.05
VFN PDF evolution	0.1190 ± 0.0007	0.71 ± 0.05

- Bias from α_S-PDFs correlations [Forte,Kassabov('20)] $\to$ PDFs refitted.
- Other PDF sets considered. CT18, CT18Z, MSHT20, HERAPDF2.0, ABMP16 (NNLO) and MSHT20an3lo. The midpoint value is the nominal result and the PDF envelope as an additional uncertainty.
- Uncertainty from missing higher orders: $m_{ll}/2 < \{\mu_R, \mu_F, Q\} < 2m_{ll}$ with $0.5 < \{\mu_R/\mu_F, \mu_R/Q, \mu_F/Q\} < 2$.
- NP effects: b_*-pr. $b_* = b/\sqrt{1 + b^2/b_{lim}^2}$ ($b_{lim} = 2 - 3$ GeV^{-1}) and *minimal* pr. ($b_{lim} \to \infty$); quartic term $\exp(-qb^4)$ and different parametrization of S_{NP} [Collins,Rogers('15)].
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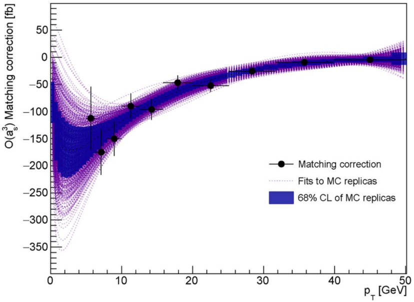
Theory uncertainties



- Bias from α_S-PDFs correlations [Forte,Kassabov('20)] $\rightarrow$ PDFs refitted.
- Other PDF sets considered. CT18, CT18Z, MSHT20, HERAPDF2.0, ABMP16 (NNLO) and MSHT20an3lo. The midpoint value is the nominal result and the PDF envelope as an additional uncertainty.
- Uncertainty from missing higher orders: $m_{ll}/2 < \{\mu_R, \mu_F, Q\} < 2m_{ll}$ with $0.5 < \{\mu_R/\mu_F, \mu_R/Q, \mu_F/Q\} < 2$.
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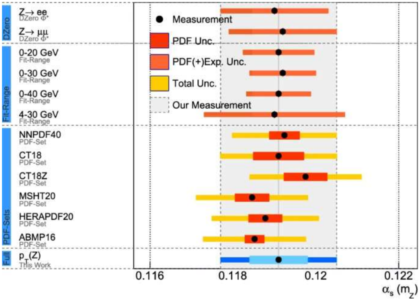

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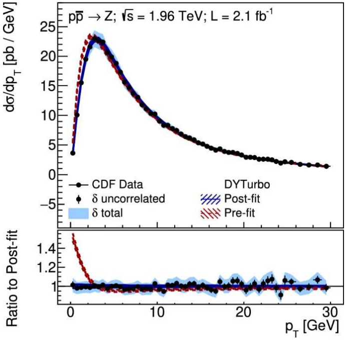
Fit results



Statistical uncertainty		± 0.7	
Experimental systematic uncertainty		± 0.1	
PDF uncertainty (NNPDF4.0)		± 0.4	
PDF uncertainty (envelope of PDFs)		± 0.7	
Scale variations uncertainties	$+0.4$		-0.9
Matching at $\mathcal{O}(\alpha_S^3)$		± 0.1	
Non-perturbative model		± 0.7	
Flavour model	0		-0.3
QED ISR		$< \pm 0.1$	
Lower limit of fit range		± 0.2	
Total	$+1.3$		-1.6

Simultaneous fit of $\alpha_S(m_Z)$ and g at N^{3}LL+$\mathcal{O}(\alpha_S^3)$ (N^{3}LL+N^{3}LO):

$$\alpha_S(m_Z) = 0.1191^{+0.0013}_{-0.0016}$$

$$g = 0.66 \pm 0.05\, GeV^2$$

Conclusions

- Novel methodology for determination of $\alpha_S(m_Z)$ based on Z-boson small-q_T distribution.
- Based on N^3LL+N^3LO resummed QCD predictions.
- Result in agreement with the world average. Uncertainty comparable to other determinations.
- Precise collider determination: 1.2% relative uncerainty.
- Crucial development of `DYTurbo` program to compute fast and accurate theoretical predictions:
  `https://dyturbo.hepforge.org`