Factorial Growth in pQCD and Implications for α_s

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Prototypical α_s Determination

- Consider an "effective charge" with a single hard scale:

$$\mathscr{R}(Q) = R(Q) + C_p \frac{\Lambda^p}{Q^p}$$

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power correction

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perturbative series in mass-independent scheme

$$R(Q) = \sum_{l=0} r_l(\mu/Q)\alpha_s(\mu)^{l+1}$$

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power p ↔ factorial growth

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Prototypical α_s Determination

- Consider an "effective charge" with a single hard scale:

physical quantity

power correction

$$\mathscr{R}(Q) = R(Q) + C_p \frac{\Lambda^p}{Q^p}$$

power $p \leftrightarrow$ factorial growth

"perturbative part"

perturbative series in mass-independent scheme

$$R(Q) = \sum_{l=0} r_l(\mu/Q)\alpha_s(\mu)^{l+1}$$

- Perturbative part and power correction inseparable.

Factorial Growth

- Even in quantum mechanics, high orders of perturbation theory grow factorially [*e.g.*, Bender & Wu 1971, 1973].
- Also in QFT [*e.g.*, Gross & Neveu 1974, Lautrup 1977].
- Static-energy r_l grow factorially (known for a long time):

$$r_l \sim R_0 (2\beta_0)^l \frac{\Gamma(l+1+b)}{\Gamma(1+b)} \equiv R_l$$

for $l \gg 1$. Here $b = \beta_1/2\beta_0^2 \overset{n_f=3}{=} 32/81 \approx 0.4$.

- Does $r_l = \{1, 1.38, 5.46, 26.7\}$ start growing by $l = 3$?

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Examples

- Static energy = energy between two static sources ($p = 1$):

$$E_0(r) = -\frac{C_F}{r} \sum_{l=0} v_l(\mu r) \alpha_s(\mu)^{l+1} + \Lambda_0$$

$$\mathscr{R}(1/r) = -r E_0(r)/C_F$$

Leino

- Its Fourier transform ($p > 1$)

$$\tilde{R}(q) = \sum_{l=0} a_l(\mu/q) \alpha_s(\mu)^{l+1}$$

- The "static force" ($p \geq 9$)

Mayer-Steudte

$$\mathfrak{F}(r) = -\frac{\mathrm{d}E_0}{\mathrm{d}r} \qquad \mathscr{F}(r) = F^{(1)}(1/r) = -r^2 \mathfrak{F}(r)/C_F$$

Outline

- Introduction
- Power Corrections and Factorial Growth
- New Approximation for Perturbative Series
- Borel Summation
- Worked Example: Static Energy
- Two or More Power Corrections
- Conclusions & Outlook

Power Corrections and Factorial Growth

μ Independence Yields Q Dependence

$$R(Q) = \sum_{l=0} r_l(\mu/Q)\alpha_s(\mu)^{l+1}$$

- Coefficients' μ dependence must cancel that of α_s, so RGE constrains Q dependence of $R(Q)$ (modulo massive loops).
- Use this observation to find the factorial growth of the r_l:
  - generalize Komijani [arXiv:1701.00347] study of pole mass;
  - new method [arXiv:2310.15137] simplifies and clarifies "minimal renormalon subtraction (MRS) [arXiv:1712.04983];
  - show factorial growth already at low orders.

Power-Term Removal

- Start with $\mathscr{R}(Q) = r_{-1} + R(Q) + C_p \frac{\Lambda^p}{Q^p}$.
- To eliminate Λ^p / Q^p, multiply by Q^p and differentiate:

$$r_{-1} + F(Q) \equiv \frac{1}{pQ^{p-1}} \frac{\mathrm{d}\, Q^p \mathscr{R}}{\mathrm{d}Q} \equiv \hat{Q}^{(p)} \mathscr{R}$$

- As a series $F(Q) = \sum_{k=0} f_k \alpha_s^{k+1}$.

$$f_k = r_k - \frac{2}{p} \sum_{j=0}^{k-1} (j+1) \beta_{k-1-j} r_j$$

- Differential equation $r(\alpha) + \frac{2}{p} \beta(\alpha) r'(\alpha) = f(\alpha)$.

Differential Equation

- Differential equation $r(\alpha) + \frac{2}{p}\beta(\alpha)r'(\alpha) = f(\alpha)$.
- Take $f(\alpha)$ as given and solve for $r(\alpha)$:
  - Komijani's solution reproduces R_l and yields R_0.
- Here, use only the elementary feature—
  - general solution is any particular solution plus a solution of the homogeneous equation (0 on RHS);
  - solution to homogeneous equation is $\propto \Lambda^p$.

My Solution

- The relation between the coefficients is a matrix equation

$$f_k^{(p)} = r_k - \frac{2}{p}\sum_{j=0}^{k-1}(j+1)\beta_{k-1-j}r_j$$

$$\boldsymbol{f}^{(p)} = \left[\mathbf{1} - \frac{2}{p}\mathbf{D}\right]\cdot\boldsymbol{r} \equiv \mathbf{Q}^{(p)}\cdot\boldsymbol{r}$$

and $\mathbf{D}$ is on the lower triangle.

- Matrix is infinite, but the lower triangular form makes a row-by-row solution straightforward.

- Notation to make the expressions compact: $\tau \equiv 2\beta_0/p$.

$$
\mathbf{Q}_{\mathrm{g}}^{(p)} = \begin{bmatrix}
1 & 0 & 0 & 0 & 0 & 0 & 0 & \cdots \\
-\tau & 1 & 0 & 0 & 0 & 0 & 0 & \cdots \\
-\tau^2 pb & -2\tau & 1 & 0 & 0 & 0 & 0 & \cdots \\
-\tau(\tau pb)^2 & -2\tau^2 pb & -3\tau & 1 & 0 & 0 & 0 & \cdots \\
-\tau(\tau pb)^3 & -2\tau(\tau pb)^2 & -3\tau^2 pb & -4\tau & 1 & 0 & 0 & \cdots \\
-\tau(\tau pb)^4 & -2\tau(\tau pb)^3 & -3\tau(\tau pb)^2 & -4\tau^2 pb & -5\tau & 1 & 0 & \cdots \\
-\tau(\tau pb)^5 & -2\tau(\tau pb)^4 & -3\tau(\tau pb)^3 & -4\tau(\tau pb)^2 & -5\tau^2 pb & -6\tau & 1 & \cdots \\
\vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \ddots & \ddots
\end{bmatrix}
$$

- As before $b = \beta_1/2\beta_0^2 \overset{n_f=3}{=} 32/81 \approx 0.4$.

- Scheme for α_s is chosen to simplify algebra (“geometric”):

$$
\beta(\alpha_{\mathrm{g}}) = -\frac{\beta_0 \alpha_{\mathrm{g}}^2}{1-(\beta_1/\beta_0)\alpha_{\mathrm{g}}}
$$

- Inverse reveals that factorial growth begins at low orders:

$$\mathbf{Q}_{\mathrm{g}}^{(p)\,-1} = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & \cdots \\ \tau & 1 & 0 & 0 & 0 & 0 & 0 & \cdots \\ \tau^2\frac{\Gamma(3+pb)}{\Gamma(2+pb)} & 2\tau & 1 & 0 & 0 & 0 & 0 & \cdots \\ \tau^3\frac{\Gamma(4+pb)}{\Gamma(2+pb)} & 2\tau^2\frac{\Gamma(4+pb)}{\Gamma(3+pb)} & 3\tau & 1 & 0 & 0 & 0 & \cdots \\ \tau^4\frac{\Gamma(5+pb)}{\Gamma(2+pb)} & 2\tau^3\frac{\Gamma(5+pb)}{\Gamma(3+pb)} & 3\tau^2\frac{\Gamma(5+pb)}{\Gamma(4+pb)} & 4\tau & 1 & 0 & 0 & \cdots \\ \tau^5\frac{\Gamma(6+pb)}{\Gamma(2+pb)} & 2\tau^4\frac{\Gamma(6+pb)}{\Gamma(3+pb)} & 3\tau^3\frac{\Gamma(6+pb)}{\Gamma(4+pb)} & 4\tau^2\frac{\Gamma(6+pb)}{\Gamma(5+pb)} & 5\tau & 1 & 0 & \cdots \\ \tau^6\frac{\Gamma(7+pb)}{\Gamma(2+pb)} & 2\tau^5\frac{\Gamma(7+pb)}{\Gamma(3+pb)} & 3\tau^4\frac{\Gamma(7+pb)}{\Gamma(4+pb)} & 4\tau^3\frac{\Gamma(7+pb)}{\Gamma(5+pb)} & 5\tau^2\frac{\Gamma(7+pb)}{\Gamma(6+pb)} & 6\tau & 1 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \ddots \end{bmatrix}$$

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well-known growth

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Komijani R_0 (truncated)

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\end{bmatrix}
$$

$$
\boldsymbol{r} = \mathbf{Q}_{\mathrm{g}}^{(p)\,-1} \cdot \boldsymbol{f}^{(p)}
$$

return

$$
r_l = \underbrace{\left(\frac{2\beta_0}{p}\right)^l \frac{\Gamma(l+1+pb)}{\Gamma(1+pb)}}_{\text{well-known growth}} \underbrace{\sum_{k=0}^{l-1} (k+1) \frac{\Gamma(1+pb)}{\Gamma(k+2+pb)} \left(\frac{p}{2\beta_0}\right)^k f_k^{(p)}}_{\text{Komijani } R_0 \text{ (truncated)}} \underbrace{+ f_l^{(p)}}_{\text{extra}}
$$

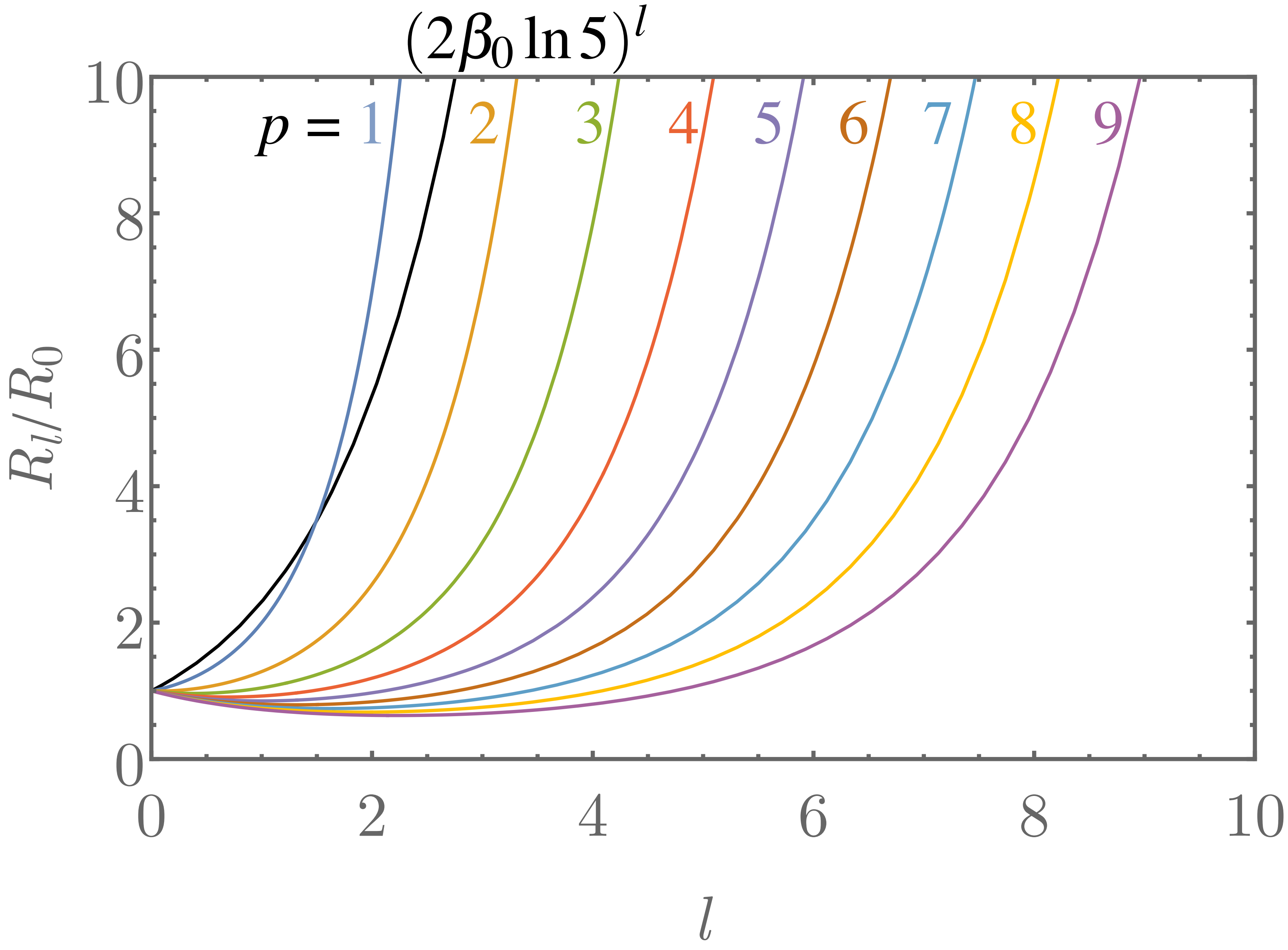
well-known growth

Komijani R_0 (truncated)

extra

Growth ↔ Power

- Larger $p \Rightarrow$ growth takes over at larger l.



New Approximation for Perturbative Series

Perturbative Series

- We must be back where we started, right? $\boldsymbol{r} = \mathbf{Q}^{-1} \cdot \boldsymbol{f} = \mathbf{Q}^{-1} \cdot \mathbf{Q} \cdot \boldsymbol{r}$
    - In practice, we know r_l and, hence, f_l for $l < L$.
      The formula returns these r_l (as it must).
- For $l \geq L$, the formula tells us (formally) the largest part.
- So truncate on f_l, not r_l. Evaluate $\sum_{l=0}^{\infty} r_l \alpha_s^{l+1}$ by—
    - taking exact r_l from the literature for $l < L$;
    - approximating $r_l \approx R_l$ for $l \geq L$.

Recap & Compendium

- That means $\sum_{l=0}^{\infty} r_l \alpha_s^{l+1} \to \sum_{l=0}^{L-1} r_l \alpha_s^{l+1} + \sum_{l=L}^{\infty} R_l^{(p)} \alpha_s^{l+1}$

  with

$$R_l^{(p)} \equiv R_0^{(p)} \left(\frac{2\beta_0}{p}\right)^l \frac{\Gamma(l+1+pb)}{\Gamma(1+pb)}$$

$$R_0^{(p)} \equiv \sum_{k=0}^{L-1} (k+1) \frac{\Gamma(1+pb)}{\Gamma(k+2+pb)} \left(\frac{p}{2\beta_0}\right)^k f_k^{(p)}$$

- Justified because the retained terms are formally larger than the ones omitted.

Borel Summation

Rearrange and React

- We have

$$R(Q) = \sum_{l=0}^{\infty} r_l \alpha_s^{l+1} \to \sum_{l=0}^{L-1} r_l \alpha_s^{l+1} + \sum_{l=L}^{\infty} R_l^{(p)} \alpha_s^{l+1}$$

$$= \underbrace{\sum_{l=0}^{L-1} \left(r_l - R_l^{(p)} \right) \alpha_s^{l+1}}_{R_{\mathrm{RS}}^{(p)}(Q)} + \underbrace{\sum_{l=0}^{\infty} R_l^{(p)} \alpha_s^{l+1}}_{R_{\mathrm{B}}^{(p)}(Q)}$$

- The "renormalon subtracted" part and the "Borel" part.

- The R_l from above yield divergent sum for R_{B}, but we're not done yet: use Borel summation to assign meaning.

Borel Summation

- Using the integral representation of $\Gamma(l+1)$:

$$R_{\mathrm{B}}^{(p)}(Q) = R_0^{(p)} \sum_{l=0}^{\infty} \left[\frac{\Gamma(l+1+pb)}{\Gamma(1+pb)\Gamma(l+1)} \int_0^{\infty} \left(\frac{2\beta_0 t}{p} \right)^l \mathrm{e}^{-t/\alpha_{\mathrm{g}}(Q)} \mathrm{d}t \right]$$

$$\to R_0^{(p)} \int_0^{\infty} \frac{\mathrm{e}^{-t/\alpha_{\mathrm{g}}(Q)}}{(1-2\beta_0 t/p)^{1+pb}} \mathrm{d}t$$

Mathematica knows the sum

where 2nd line comes from (illegally) swapping Σ and $\int$.

- Branch point in integrand at $t = p/2\beta_0$, dubbed "renormalon singularity" ['t Hooft 1979].

Borel Summation

- Split integration in two [BKKV, arXiv:1712.04983]:

$$R_{\mathrm{B}}^{(p)}(Q) = R_0^{(p)} \int_0^{p/2\beta_0} \frac{\mathrm{e}^{-t/\alpha_{\mathrm{g}}(Q)}}{(1-2\beta_0 t/p)^{1+pb}}\mathrm{d}t + R_0^{(p)} \int_{p/2\beta_0}^{\infty} \frac{\mathrm{e}^{-t/\alpha_{\mathrm{g}}(Q)}}{(1-2\beta_0 t/p)^{1+pb}}\mathrm{d}t$$

Mathematica knows the integrals

where ± on 2nd line comes from choice of contour.

- Without loss, absorb the second line into the power correction in $\mathscr{R}(Q)$.

Borel Summation

- Split integration in two [BKKV, arXiv:1712.04983]:

$$R_{\mathrm{B}}^{(p)}(Q) = R_0^{(p)} \frac{p}{2\beta_0} \mathscr{J}\left(pb, 1/2\beta_0\alpha_{\mathrm{g}}(Q)\right)$$

$$\pm R_0^{(p)} \mathrm{e}^{\pm \mathrm{i} pb\pi} \frac{p^{1+pb}}{2^{1+pb}\beta_0} \Gamma(-pb) \left[\frac{\mathrm{e}^{-1/[2\beta_0\alpha_{\mathrm{g}}(Q)]}}{[\beta_0\alpha_{\mathrm{g}}(Q)]^b}\right]^p$$

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$$\pm R_0^{(p)} \mathrm{e}^{\pm \mathrm{i} pb\pi} \frac{p^{1+pb}}{2^{1+pb}\beta_0} \Gamma(-pb) \left[\frac{\Lambda_{\overline{\mathrm{MS}}}}{Q}\right]^p$$

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Definition and Properties of $\mathscr{J}$

- Thus, we now define

$$R_{\mathrm{B}}^{(p)}(Q) = R_0^{(p)} \frac{p}{2\beta_0} \mathscr{J}(pb, 1/2\beta_0\alpha_{\mathrm{g}}(Q))$$

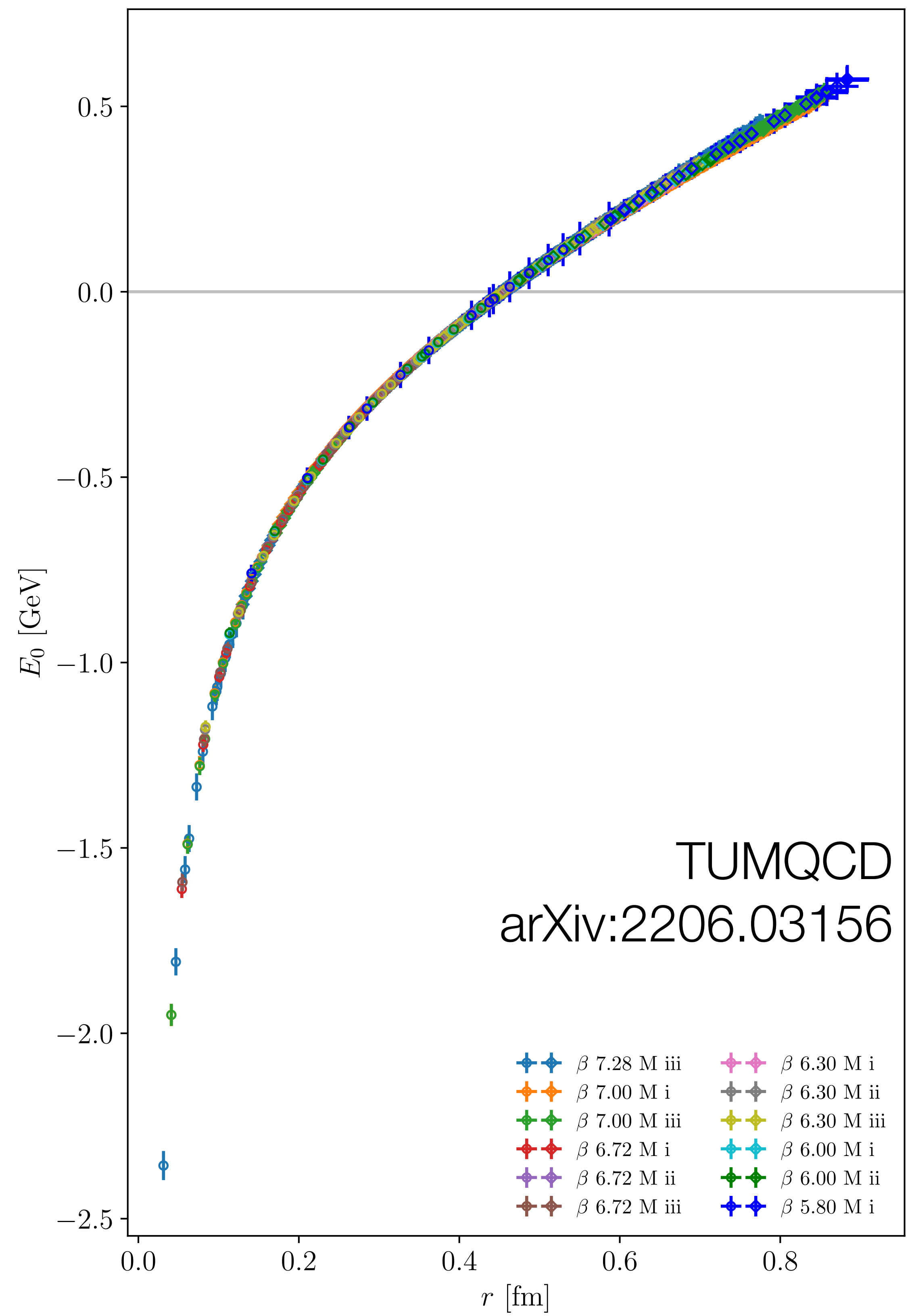
$$\mathscr{J}(c,y) = \mathrm{e}^{-y}\Gamma(-c)\gamma^{\star}(-c,-y)$$

where $\gamma^{\star}(a,x)$ is an analytic function of both a and x:

limiting function of the incomplete gamma function

- convergent expansion in $x = -1/2\beta_0\alpha_{\mathrm{g}}$;
- asymptotic expansion in α_{g} regenerates the starting point; the dropped term is $\mathrm{O}(\mathrm{e}^{-p/2\beta_0\alpha_{\mathrm{g}}})$.

Static Energy



Static Energy

- Quantity extracted from oblong Wilson loops:
  - perturbative potential has IR divergences starting at 3 loops [Appelquist, Dine, Muzinich 1978];
  - compensated by multipole (retardation) term [Brambilla, Pineda, Soto, Vairo 1999, 2000].
- Perturbative series:

$$E_0(r) = -\frac{C_F}{r}\sum_{l=0} v_l(\mu r)\alpha_s(\mu)^{l+1} + \Lambda_0$$

- In notation used above, $Q \to 1/r$, $\mathscr{R}(1/r) = -rE_0(r)/C_F$.

Related Quantities

- Perturbation theory carried out in momentum space:

$$\tilde{R}(q) = \sum_{l=0} a_l(\mu/q)\alpha_s(\mu)^{l+1}$$

- Leading power/factorial comes from Fourier transform, so $\tilde{R}(q)$ has $p > 1$.

- The “static force”

$$\mathfrak{F}(r) = -\frac{\mathrm{d}E_0}{\mathrm{d}r} \qquad \mathscr{F}(r) = F^{(1)}(1/r) = -r^2\mathfrak{F}(r)/C_F$$

has no power corrections (until instantons at $p \geq 9$).

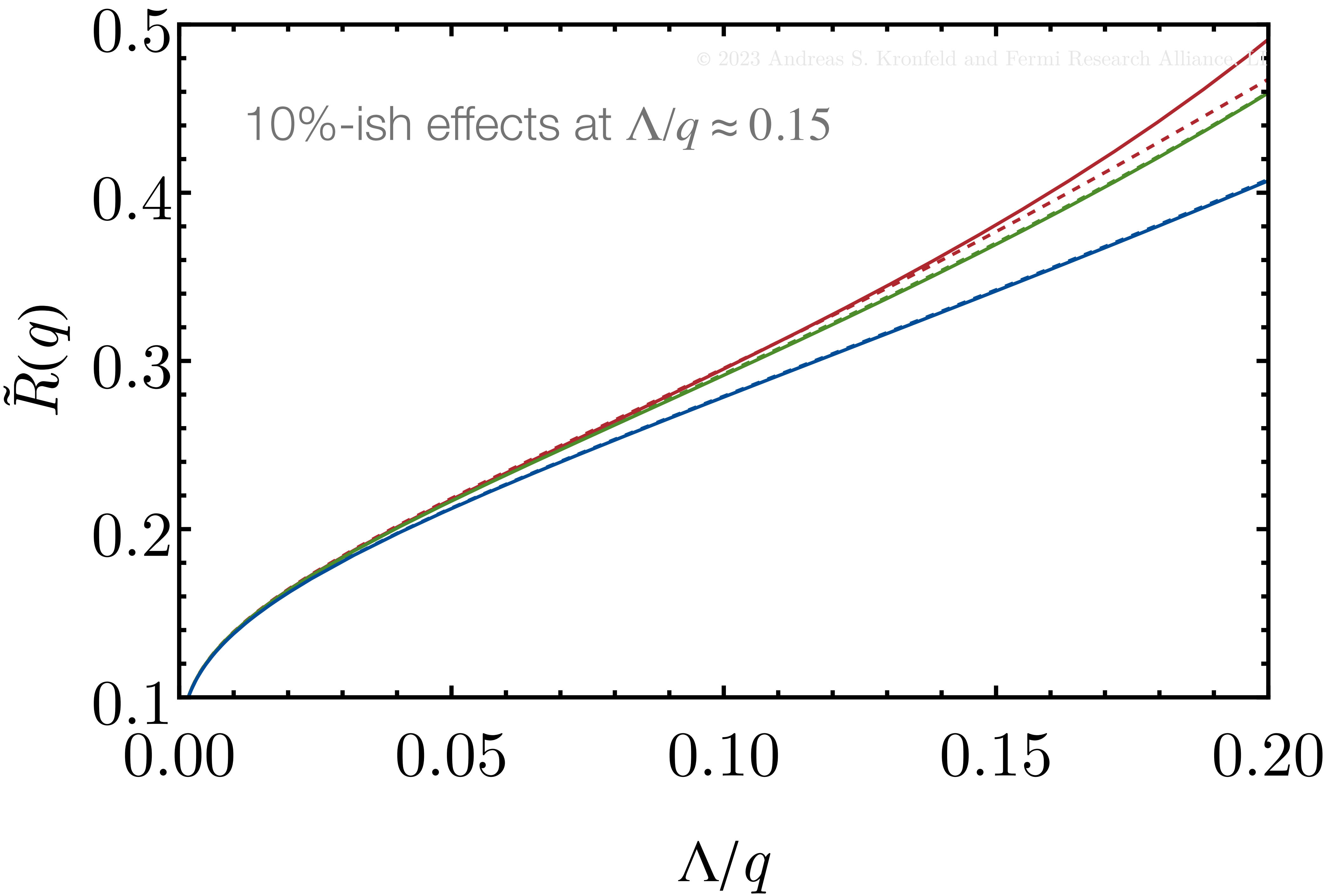
Coefficients at $\mu = 1/r$ or $\mu = q$

l	$\overline{\text{MS}}$ $a_l(1)$	$\overline{\text{MS}}$ $f_l(1)$	geometric $a_l(1)$	geometric $f_l(1)$	α_2 $a_l(1)$	α_2 $f_l(1)$
0	1	1	1	1	1	1
1	0.557 042	−0.048 552	0.557 042	−0.048 552	0.557 042	−0.048 552
2	1.702 18	0.687 291	1.834 97	0.820 079	1.834 97	0.820 079
3	2.436 87	0.323 257	2.832 68	0.558 242	3.013 89	0.739 452

l	$\overline{\text{MS}}$ $v_l(1)$	$\overline{\text{MS}}$ $v_l(1)-V_l(1)$	geometric $v_l(1)$	geometric $v_l(1)-V_l(1)$	α_2 $v_l(1)$	α_2 $v_l(1)-V_l(1)$
0	1	0.206 061	1	0.182 531	1	0.177 584
1	1.383 84	−0.202 668	1.383 84	−0.249 689	1.383 84	−0.259 574
2	5.462 28	0.019 479	5.595 07	−0.009 046	5.595 07	−0.042 959
3	26.6880	0.219 262	27.3034	0.050 179	27.4846	0.066 468

Good Series (at most $p > 1$ growth)

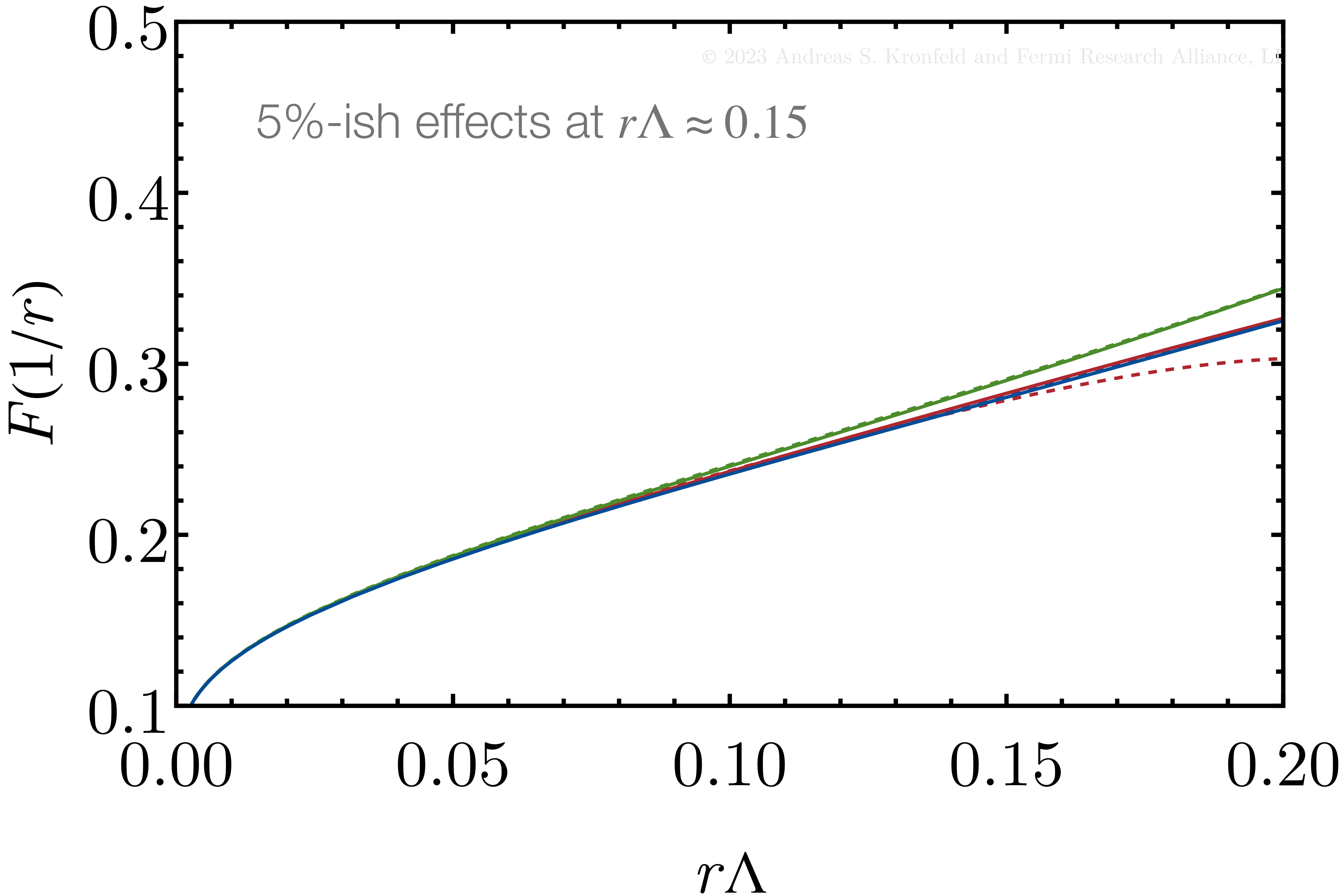
$\mu = sq$



$s = ½$
$s = 1$
$s = 2$

Great Series (instanton power $p \geq 9$)

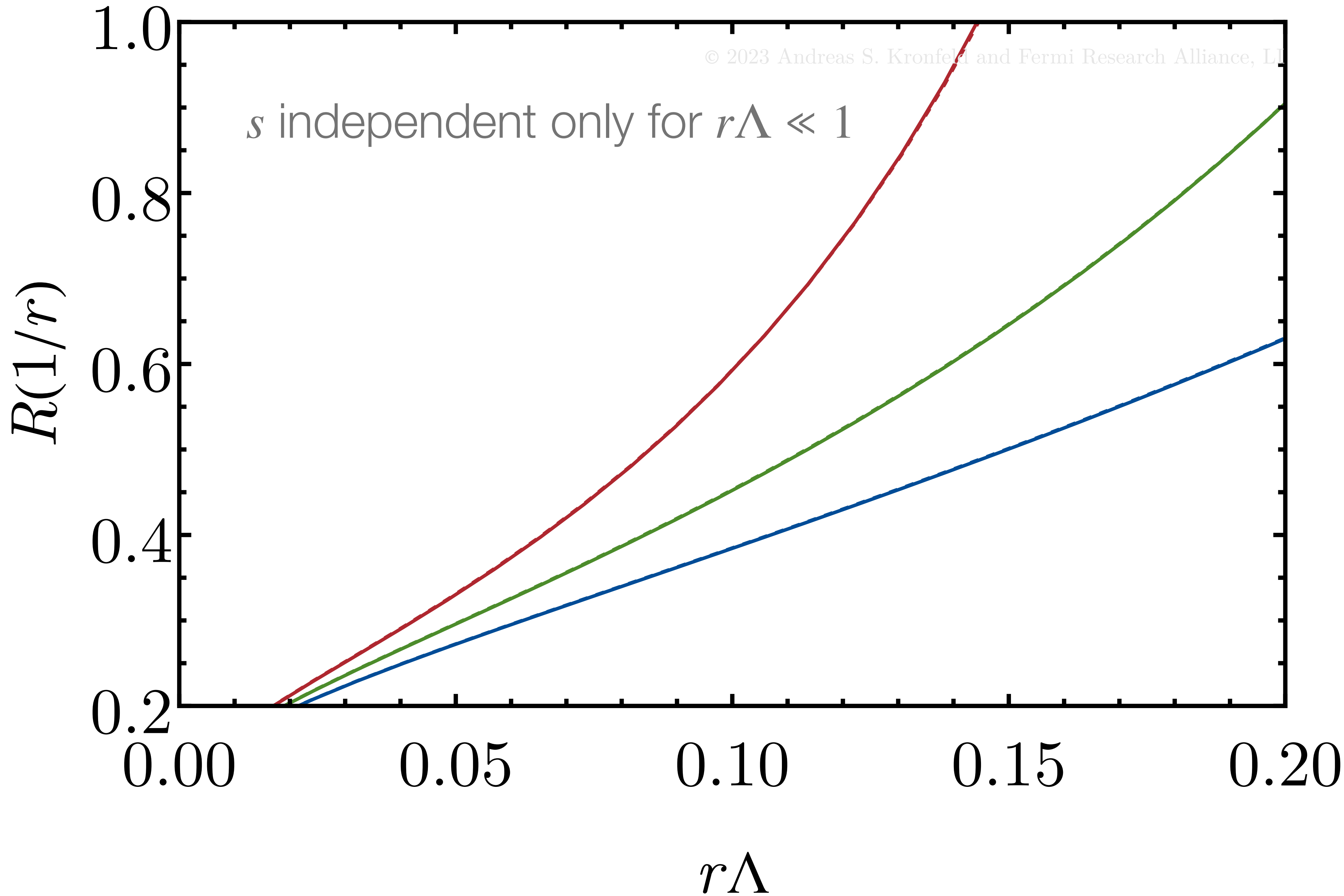
$\mu = s/r$



$s = ½$
$s = 1$
$s = 2$

Horrible Series ($p = 1$)

$\mu = s/r$



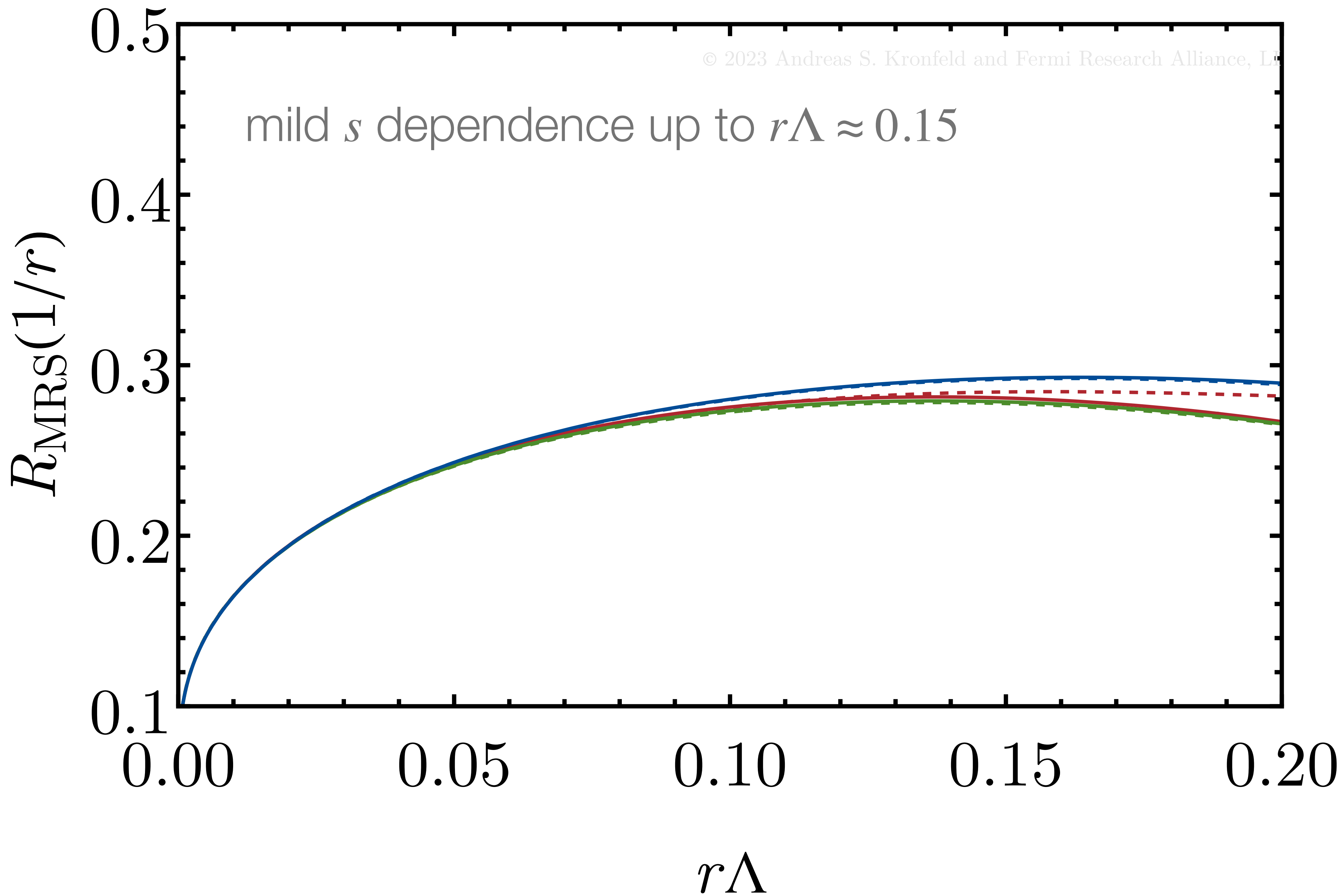
— $s = ½$

— $s = 1$

— $s = 2$

MRS Series

$\mu = s/r$



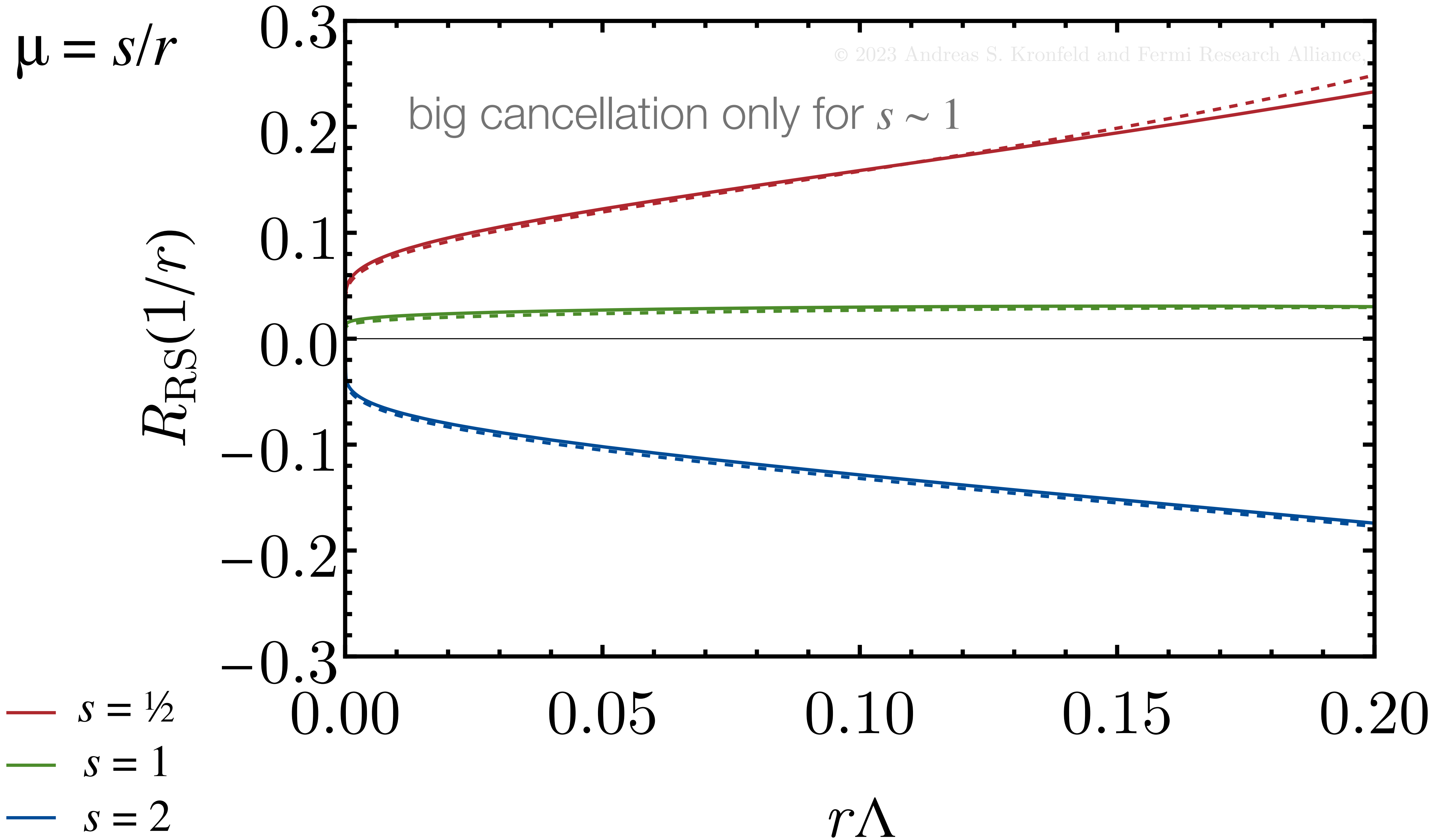
$s = ½$

$s = 1$

$s = 2$

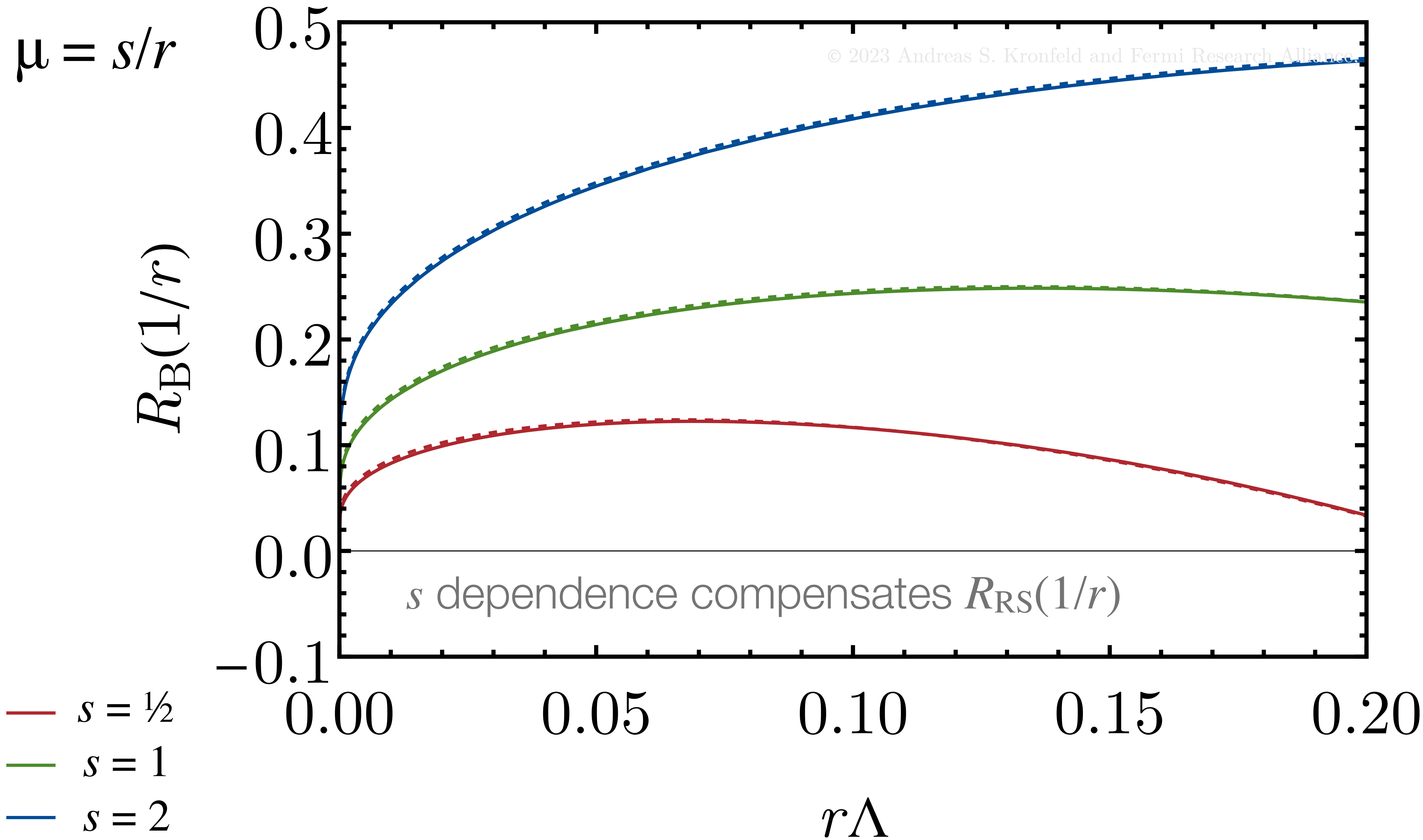
Renormalon Subtracted Series

$\mu = s/r$



The part that is a convergent series in $1/\alpha_s$

$\mu = s/r$



Fitting with Power Corrections

- The Λ on the horizontal axis is $\Lambda_{\overline{\text{MS}}}$—
  - fits to data will have this as free parameter, i.e., optimization will stretch/shrink the curves to fit.
- Let's go back to the plots and get a feel for adding small amounts of order $(\Lambda/q)^{2 \text{ or } 3 \text{ or } 4}$, $(\Lambda r)^9$, or Λr.
- Disentangling power-law and logarithmic dependence seems hard for $\tilde{R}(q)$ and $R(1/r)$, but not for $F(1/r)$ and $R_{\text{MRS}}(1/r)$.

★

Two or More Power Corrections

Next Approximation

- If there is another power correction with $p_2 > p_1 = p$, then f_k will grow in a similar but slower fashion.

- Apply previous procedure with p_1; then repeat with p_2:

$$\boldsymbol{f}^{\{p_1,p_2\}} \equiv \mathbf{Q}^{(p_2)} \cdot \mathbf{Q}^{(p_1)} \cdot \boldsymbol{r}$$

$$\Rightarrow \boldsymbol{r} = \mathbf{Q}^{(p_1)^{-1}} \cdot \mathbf{Q}^{(p_2)^{-1}} \cdot \boldsymbol{f}^{\{p_1,p_2\}}$$

$$= \left[\frac{p_2}{p_2 - p_1} \mathbf{Q}^{(p_1)^{-1}} + \frac{p_1}{p_1 - p_2} \mathbf{Q}^{(p_2)^{-1}} \right] \cdot \boldsymbol{f}^{\{p_1,p_2\}}$$

- Extension to any sequence of higher powers by induction.

Summary

Summary

- MRS revisited for any sequence of power corrections ↔ dominant, subdominant, sub-subdominant, … growth.
- Formulas for growth and normalization both follow from RGE and hold exactly at low orders.
- Cancellation scale dependent, but total is not.
- Scale dependence is mild.
- Standard to sum logarithms; let’s sum factorials too!

Thank you for your attention

Questions?