

α_s FROM τ DECAY

Diogo Brito, Maarten Golterman, Kim Maltman, Santiago Peris
($\&$ Marcus Rodrigues, Wilder Schaaf)

- arXiv:2012.10440 (PRD)
- new arXiv to appear.

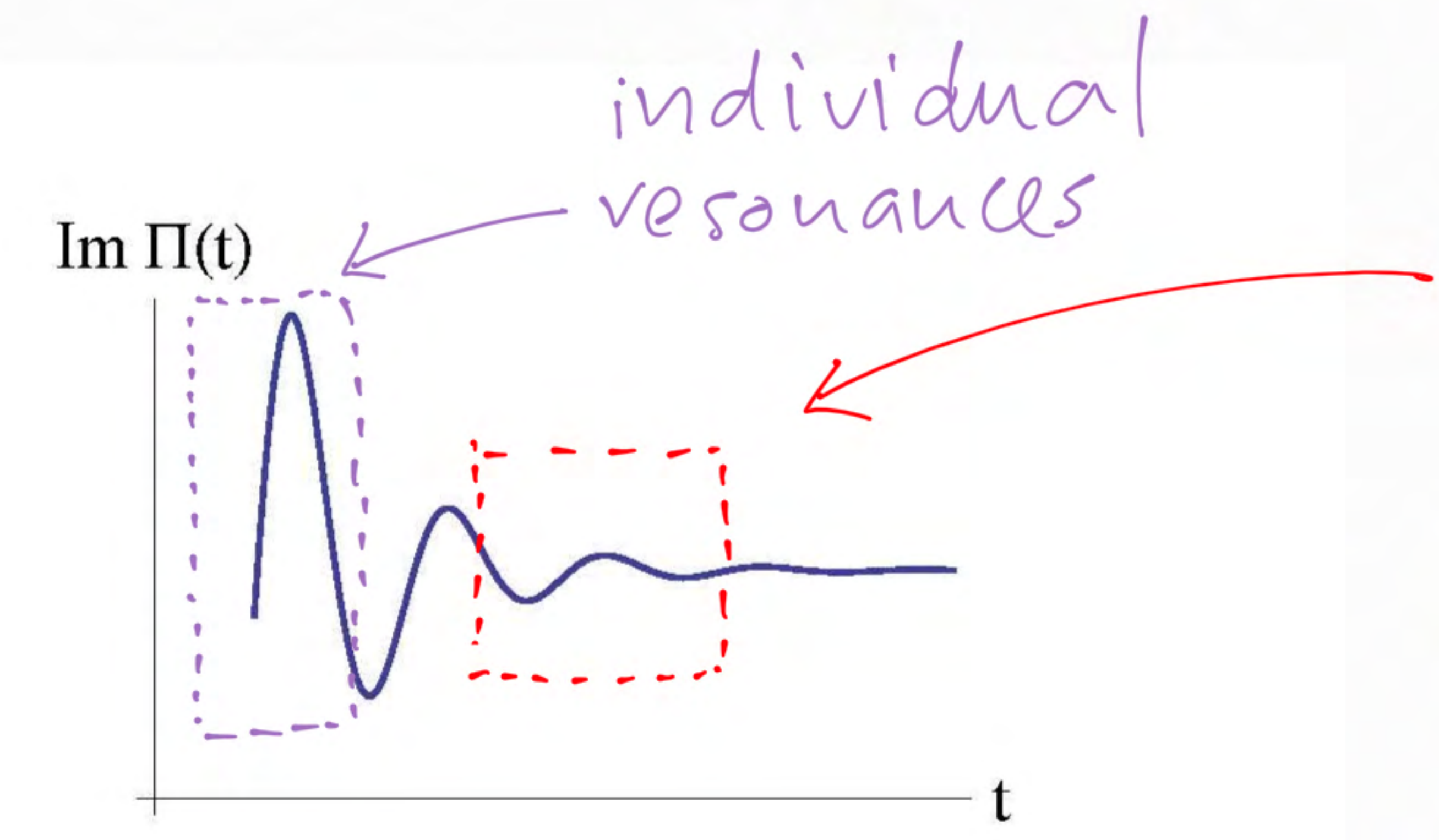
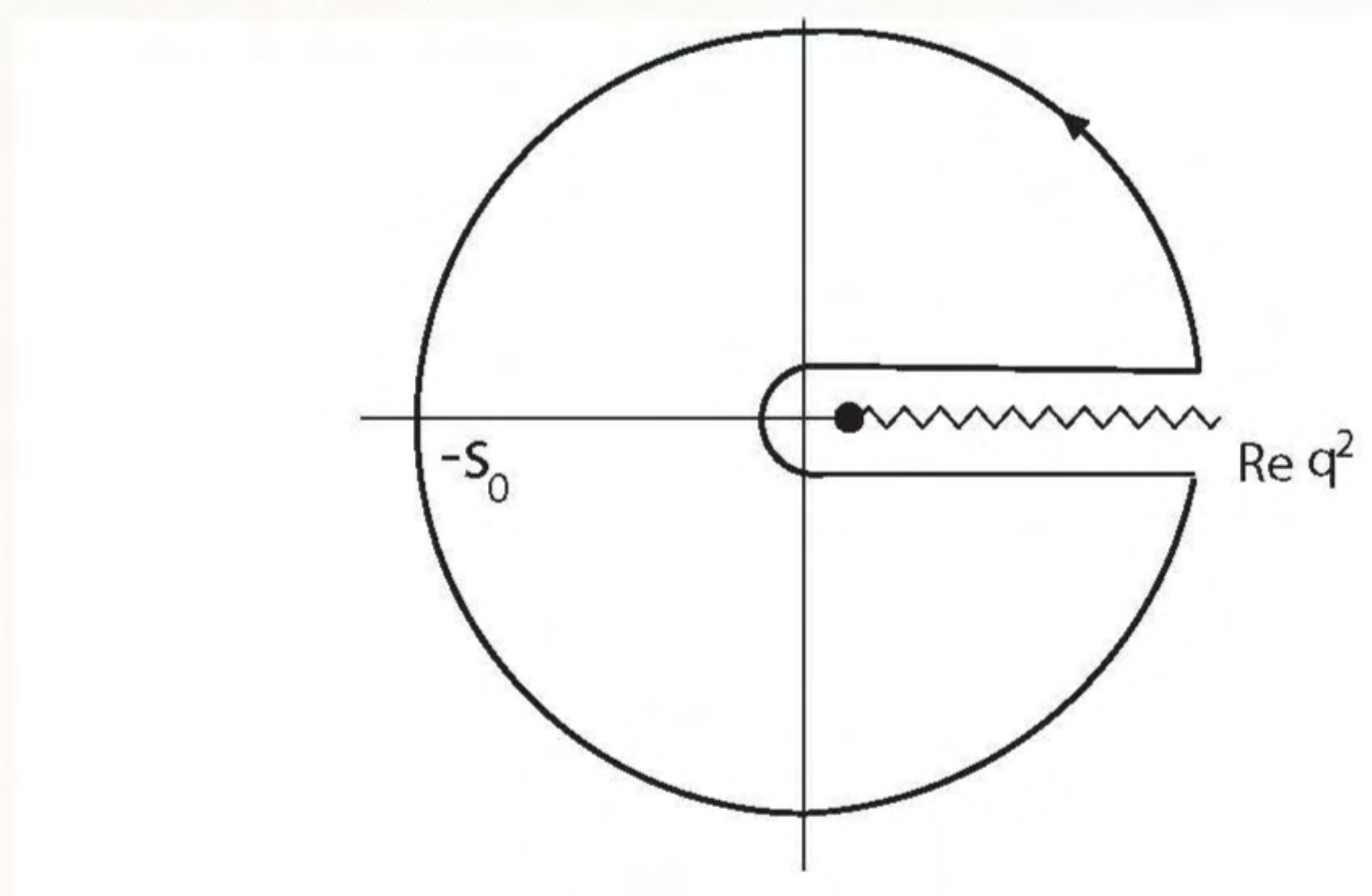
(see also arXiv:1611.03457 (PRD), 1907.03360 (PRD))

ALPHA_S (2024) @ TRENTO

FESRs, OPE & DVs

Shankar 77; Braaten, Narison, Pich '92

• $\Pi(q^2)$



wiggles!
(resonances overlap)
Poggio, Quinn, Weinberg '76

• "Cauchy's Theorem" ($z = q^2$, $w(z) = \text{polynomial}$):

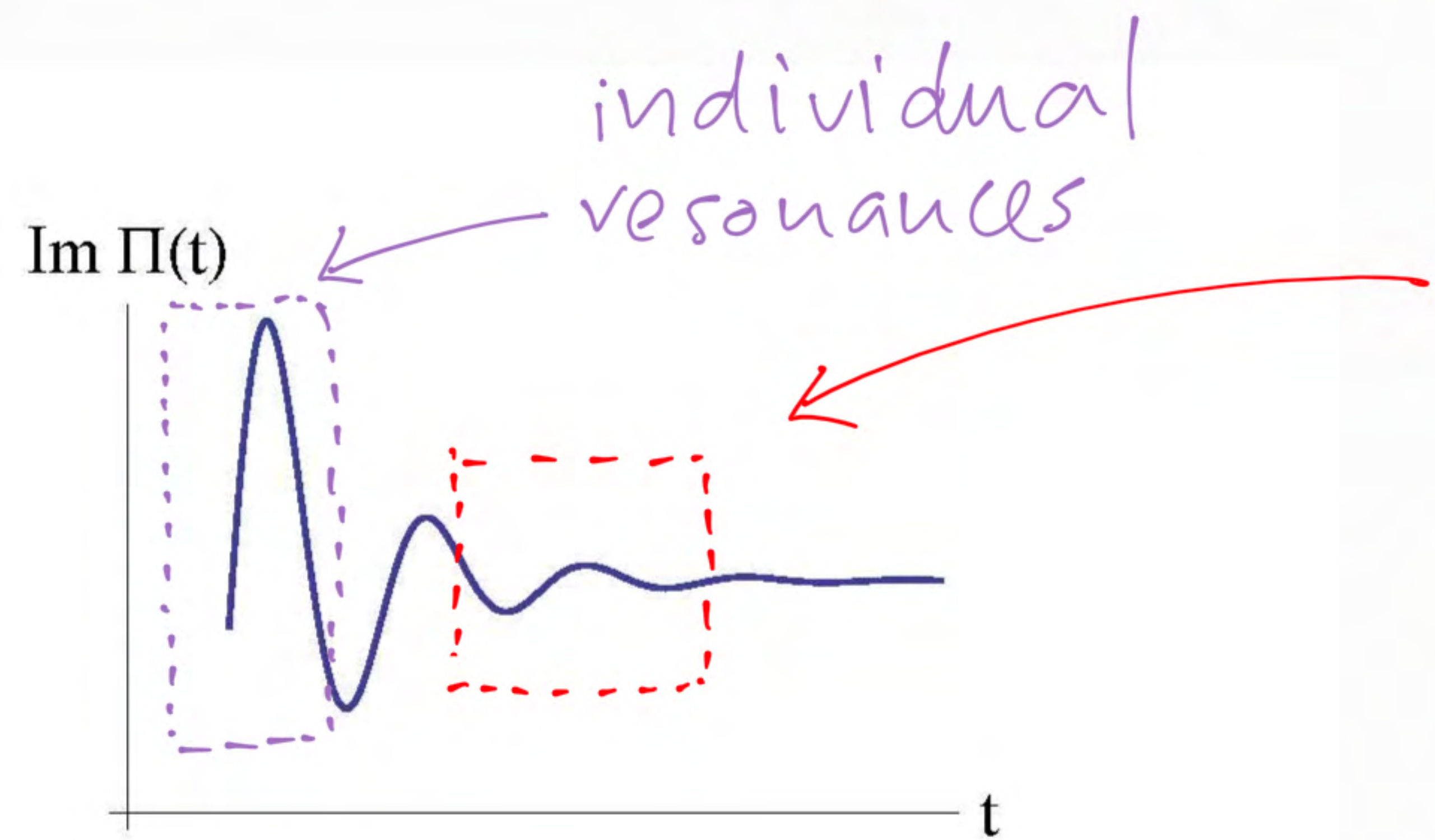
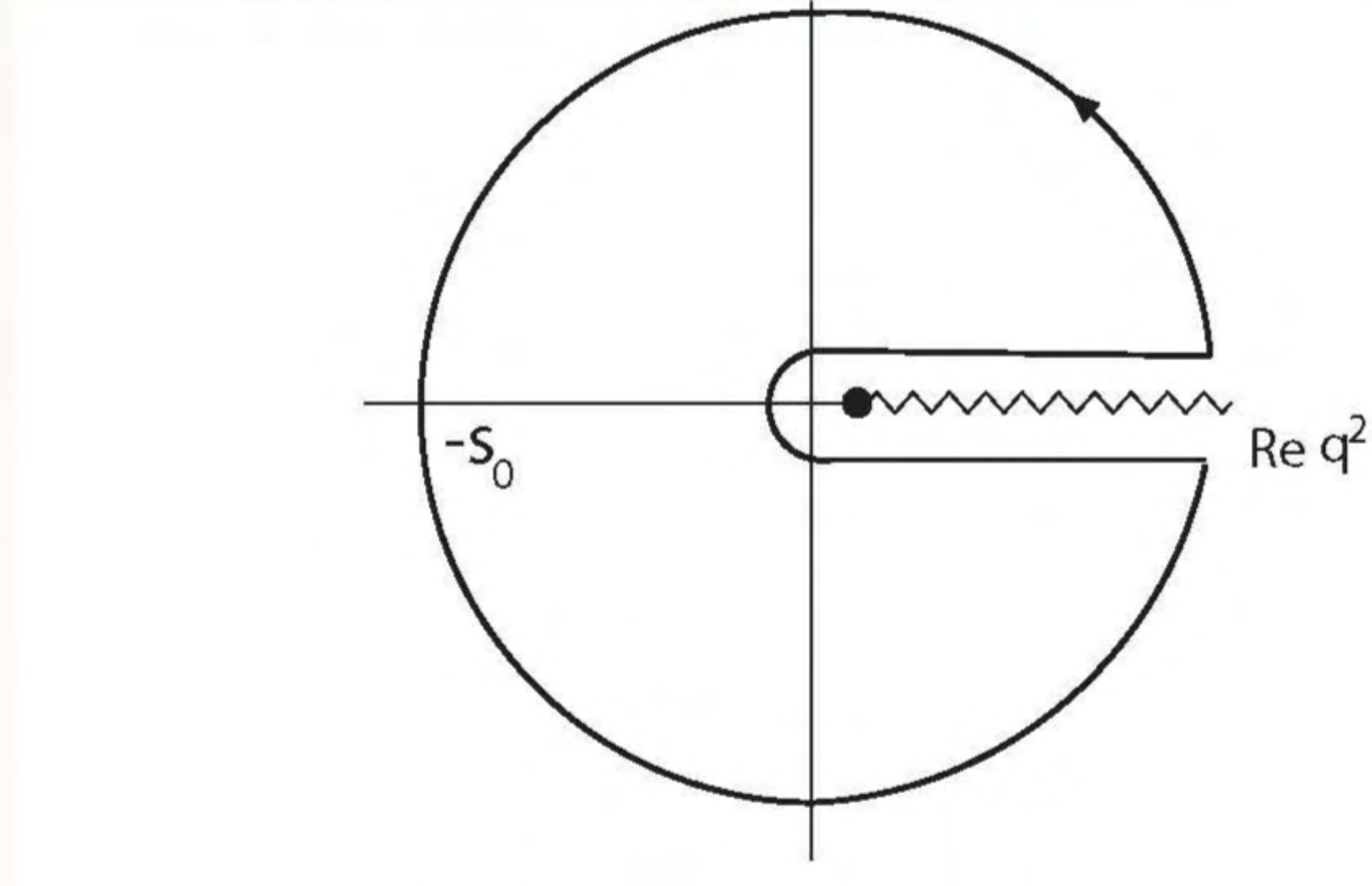
$$\begin{aligned}
 \int_0^{s_0} dt w(t) \underbrace{\frac{1}{\pi} \text{Im} \Pi(t)}_{\text{exp.}} &= -\frac{1}{2\pi} \oint_{|z|=s_0} dz w(z) \underbrace{\Pi(z)}_{\Pi_{\text{OPE}} + \Pi_{\text{DV}}} \\
 &= -\frac{1}{2\pi} \oint_{|z|=s_0} dz w(z) \underbrace{\Pi_{\text{OPE}}(z)}_{O(\alpha_s^5) + \text{Condensates}} - \int_{s_0}^{\infty} dt w(t) \frac{1}{\pi} \text{Im} \Pi_{\text{DV}}(t)
 \end{aligned}$$

Cata, Golterman, S.P. '05, '08, '09

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Cata, Golterman, S.P. '05, '08, '09

➔

OPE
ASYMPTOTIC

A TALE OF TWO WORLDS

- Pert. Theory
 - CIPT ✗
 - FOPT (\overline{MS}) ✓ (Hoang's talk)

● Non Perturbative

* DV Approach:

$$\pi_{OPE} \sim N! \left[\frac{\Lambda_{QCD}^2}{m_c^2} \right]^N$$

(asymptotic)

$$\text{Im} \pi_{DV}(s) = e^{-\delta - \gamma s} \sin(\alpha + \beta s) \left(1 + \mathcal{O}\left(\frac{1}{s}, \frac{1}{N_c}, \frac{1}{\log s}\right) \right)$$

(Asymptotic Regge spectrum @ large N_c)

→ $\alpha_s, \delta, \gamma, \alpha, \beta$ extracted from fit to data in $s_0 \in [s_{min}, m_c^2]$

* tOPE:

$$\pi_{OPE} \sim \left[\frac{\Lambda_{QCD}^2}{m_c^2} \right]^N \ll 1$$

($N \gg 1$)

(convergent) FAPP

- "pinch": use polynomial w zero @ highest $s_0 \approx m_c^2$
- truncate OPE to lowest orders that allow a fit @ single $s_0 \approx m_c^2$
- set $DV = 0$

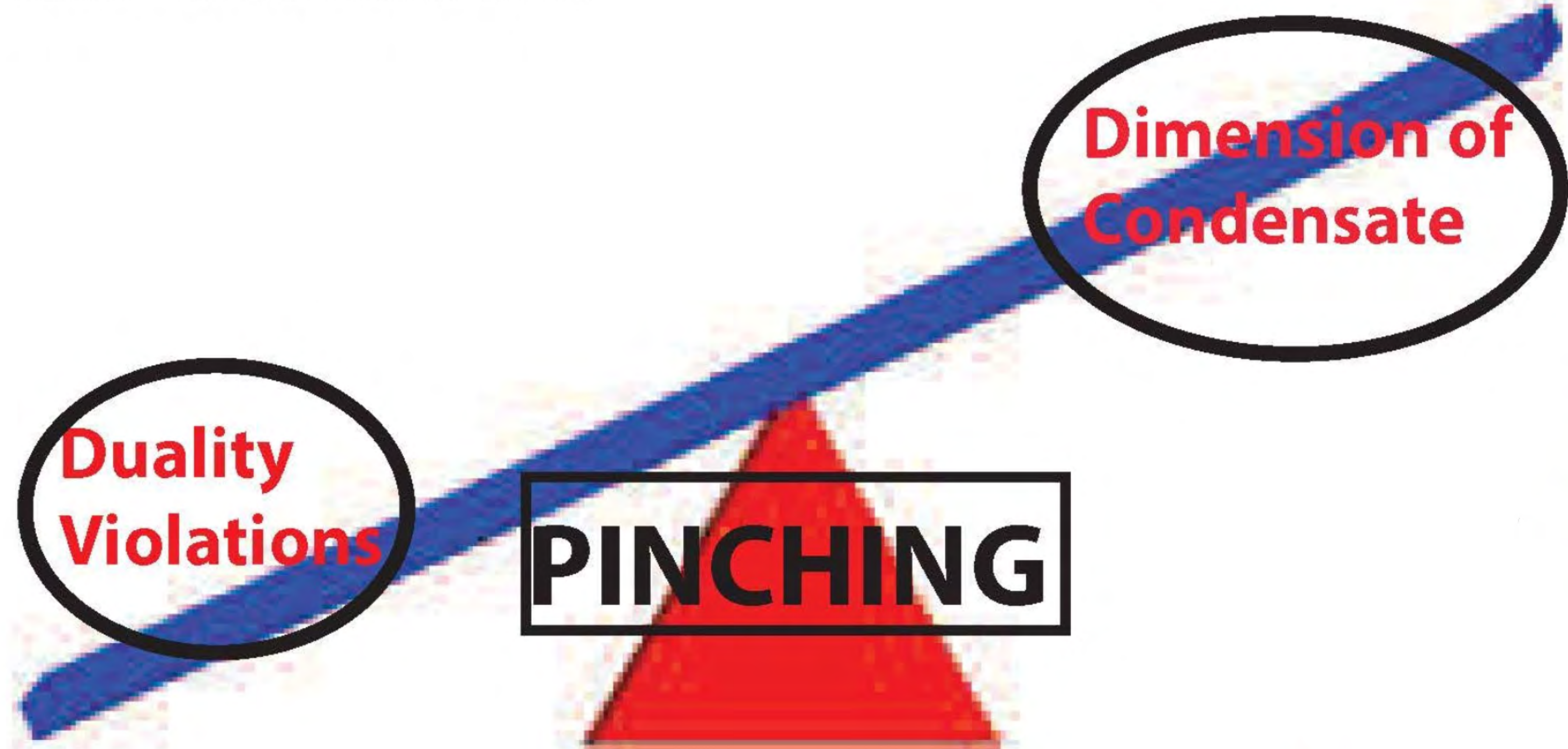
-3- BUT...

Main Theoretical Message

- with pinching \exists price to pay:

It's **NOT** possible to simultaneously suppress DVs and condensates

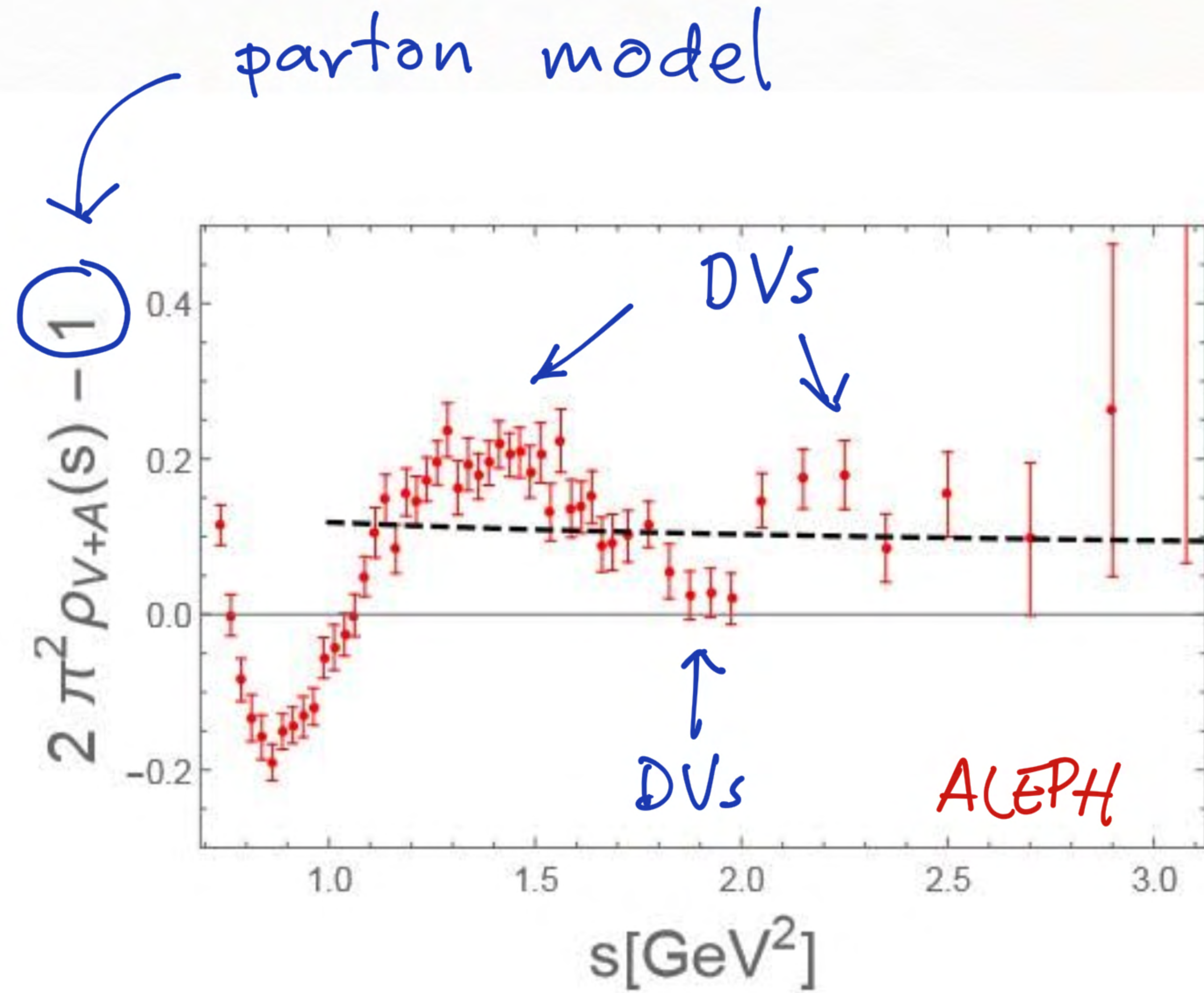
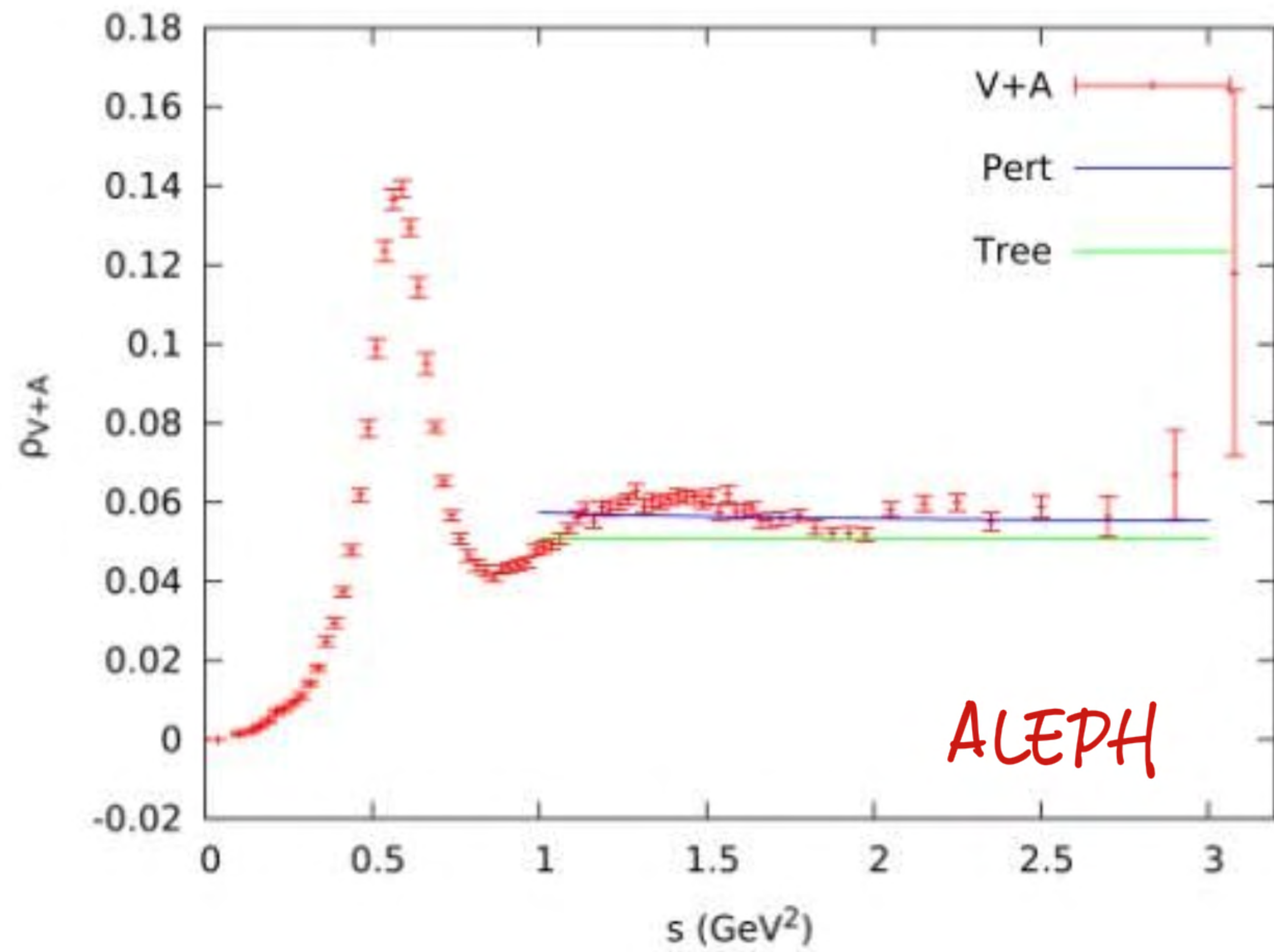
★ "Seesaw" mechanism at work:



OPE Asymptotic
 $\langle O_N \rangle \sim N!$
 \Downarrow
Restrict OPE to low orders consistently*

* $w(t) = t^N \leftrightarrow \text{OPE} = \frac{C_{2N+2}}{5^{N+1}}$

BEWARE!



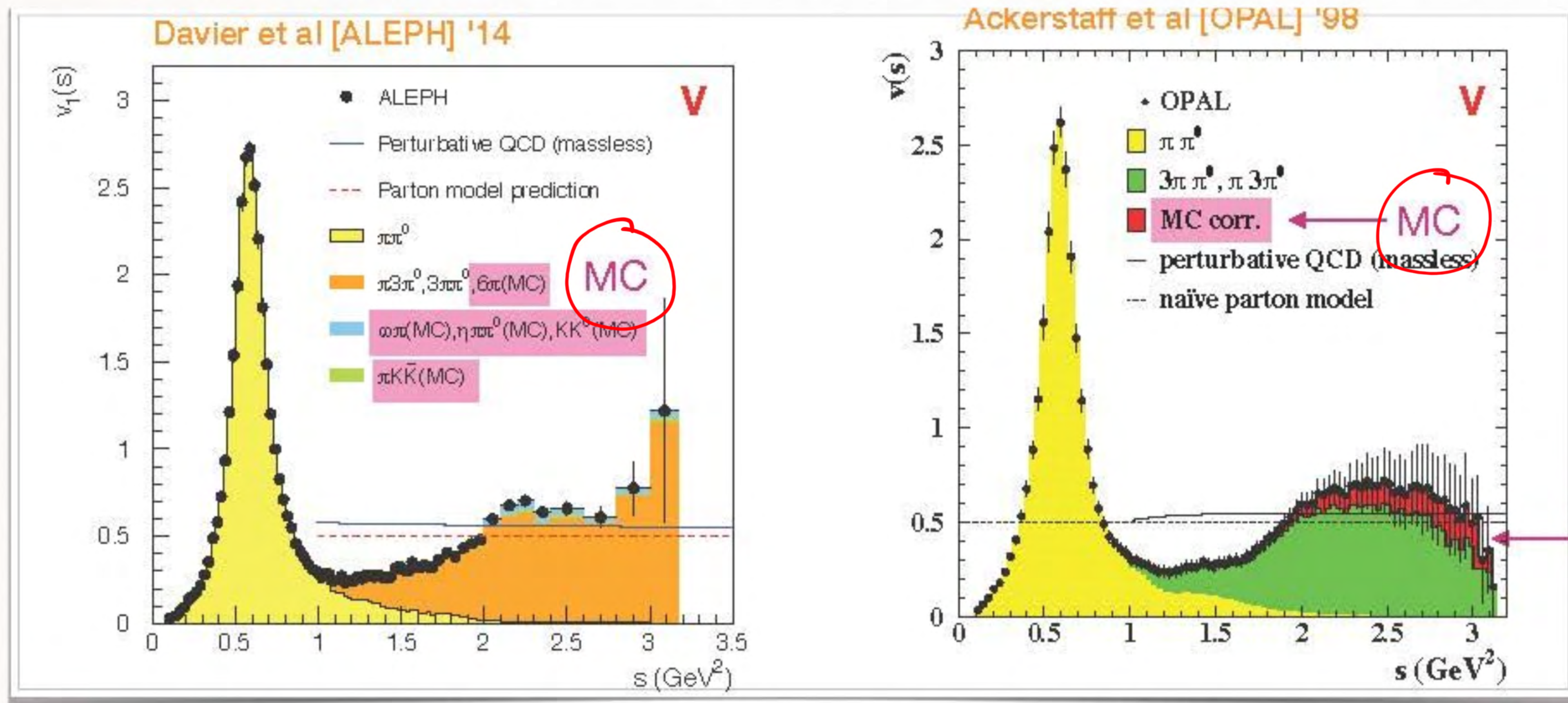
parton model
①

$\mathcal{O}(\alpha_s)$
perturbative

Pich, Rodriguez S. '22

ALEPH & OPAL Vector channel Data Sets

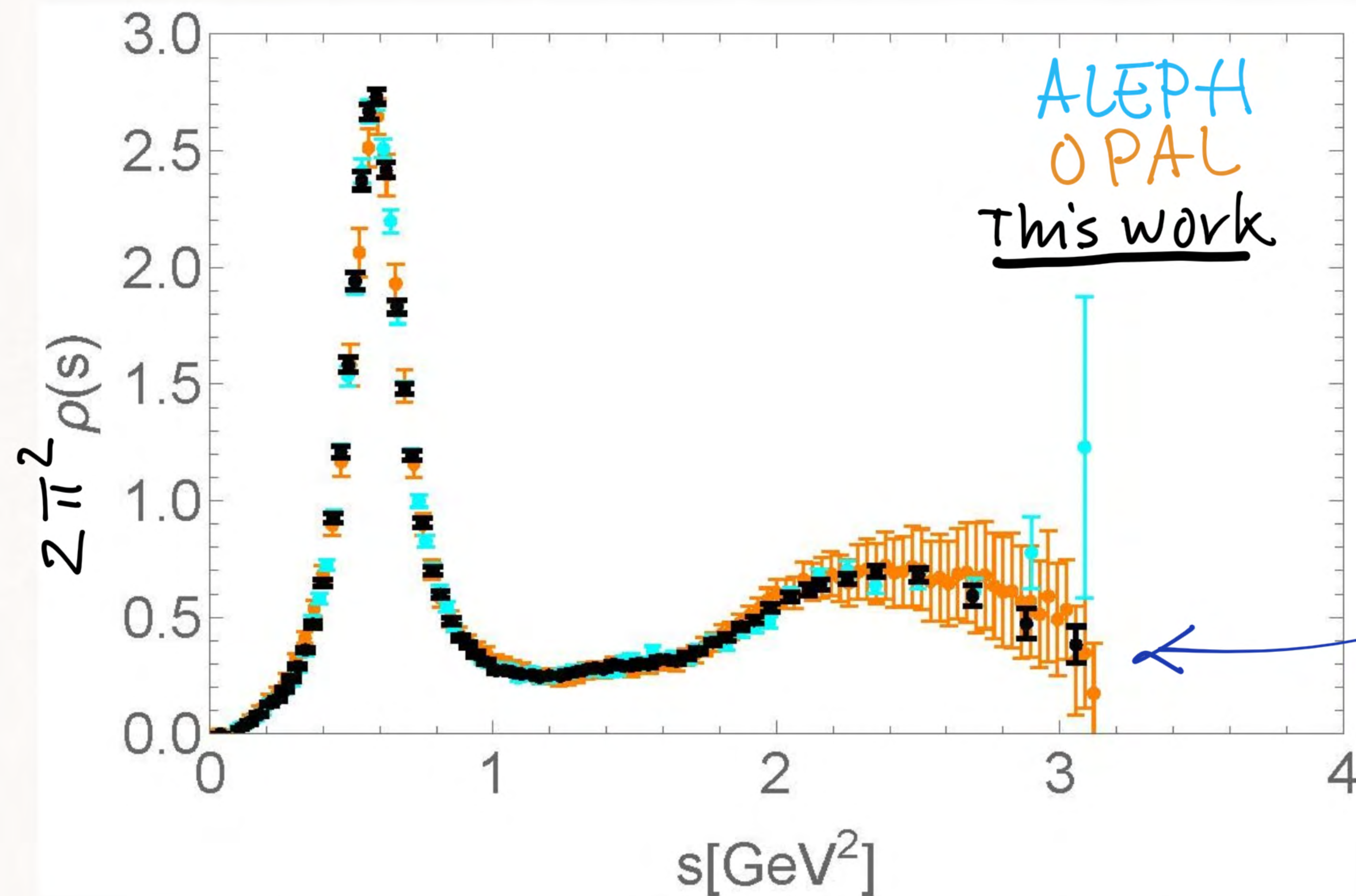
- V channel dominated by $\tau \rightarrow \nu_\tau \frac{2\pi}{4\pi}$
- Residual channels subdominant (but important for α_s !)
- \exists Monte Carlo (MC) input for several channels.



Recent measurements in e^+e^- allow to replace MC.

New vector spectral function

- Combined 2π & 4π channels from ALEPH & OPAL (à la Keshavarzi et al. '18 for $(g-2)_\mu$)
- No more Monte Carlo: 7 residual channels from e^+e^- using CVC (IB corr's small) and BABAR for $\tau \rightarrow K K_S \nu_\tau$ + updated BRs.

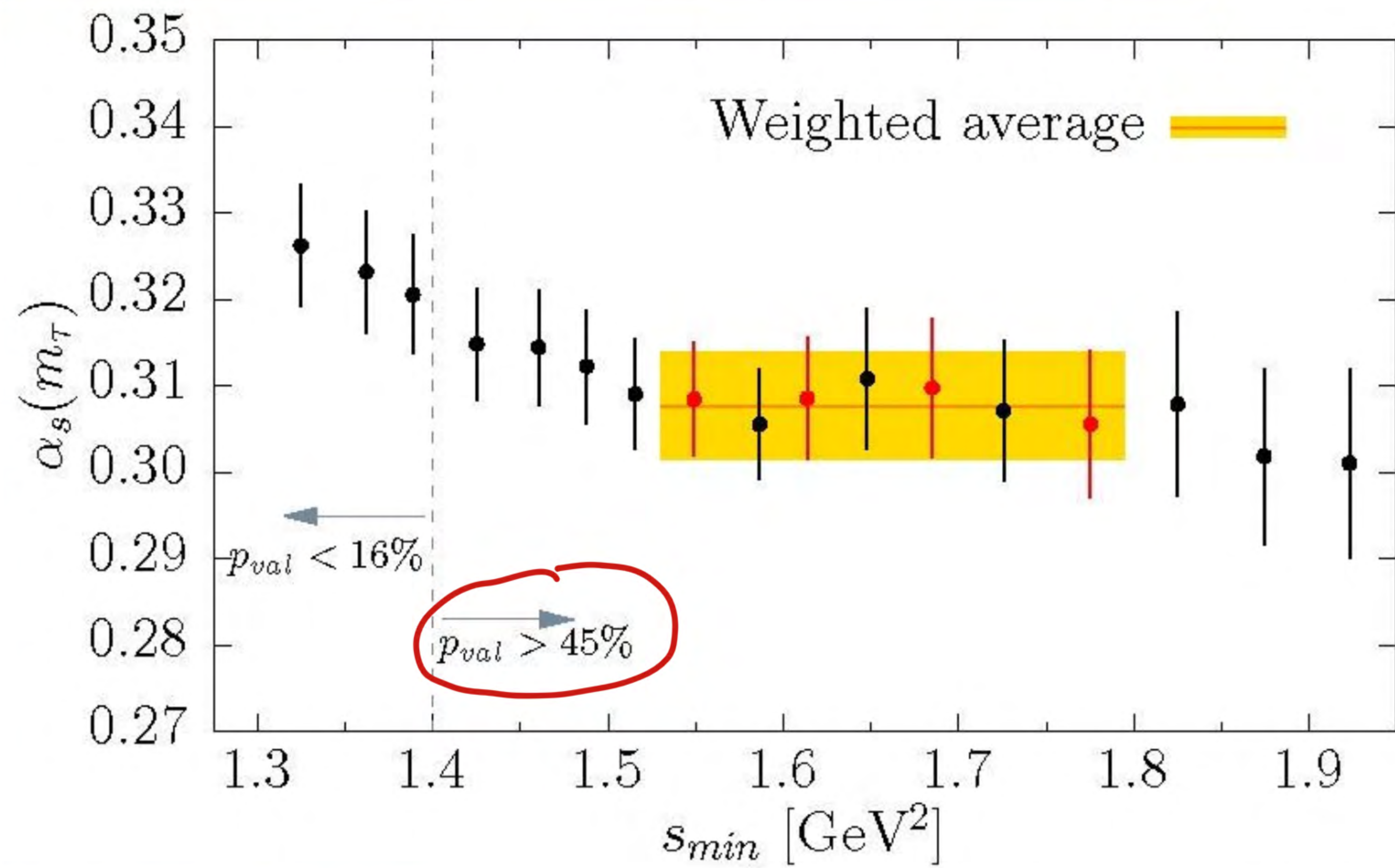


(all concentrated near end point)

- Dramatic error improvement near end point!

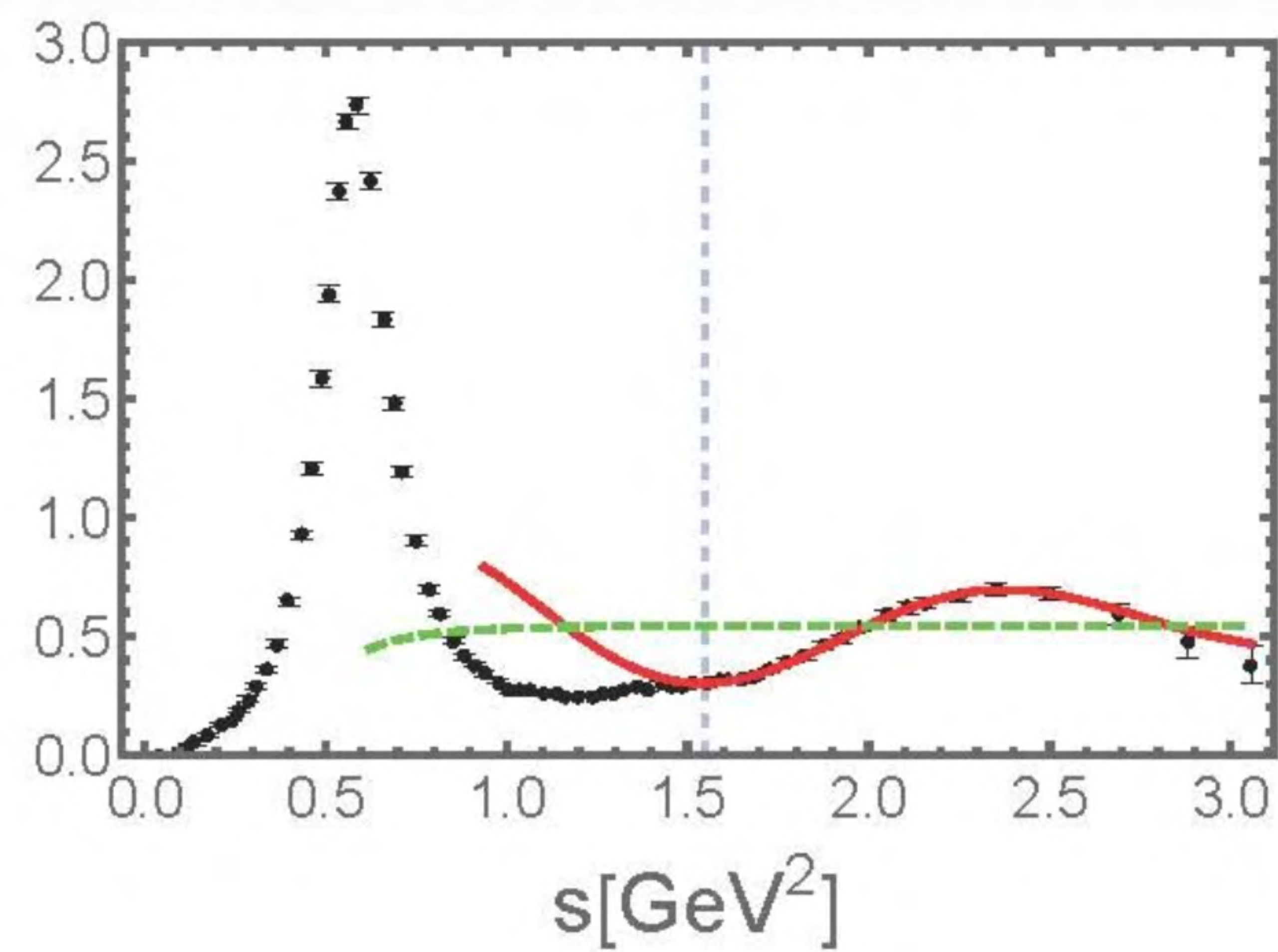
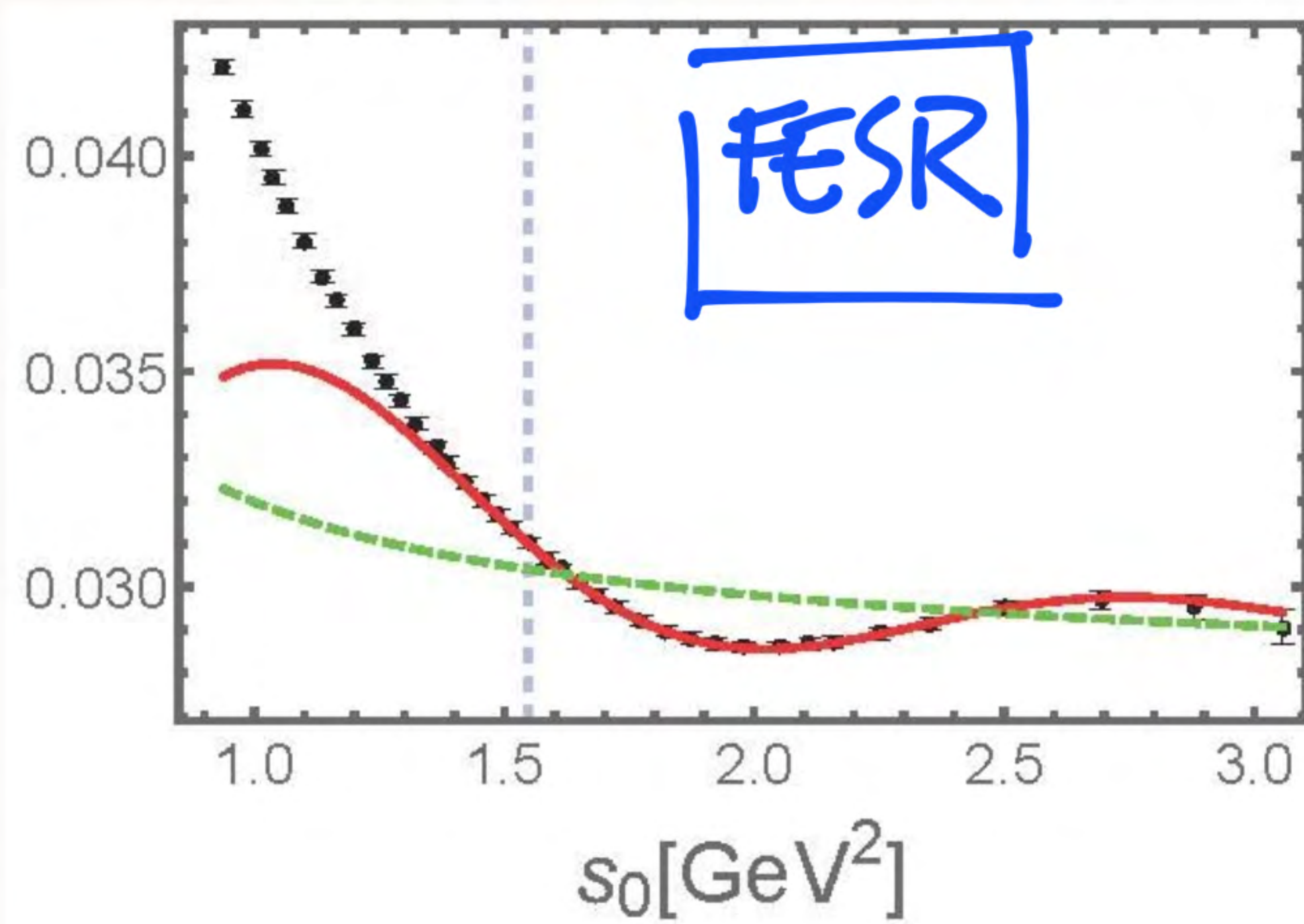
Extraction of α_s (I)

\bullet $w_0 = 1$



$s_0 \in [s_{min}, m_c^2]$ FESR

$$\int_0^{s_0} ds \frac{1}{\pi} \text{Im} \Pi_V(s) = -\frac{1}{2i\pi} \oint_{|z|=s_0} dz \Pi_V(z) = \int_{s_0}^{\infty} ds \frac{1}{\pi} \frac{\text{Im} \Pi(s)}{2V}$$



\oplus Data
 — OPE + DVs
 --- OPE w/o DVs
 $w_0 = 1 \leftrightarrow$ OPE = Pert. Theory

Extraction of α_s (II)

mom.	α_s	$c_6 [\text{GeV}^6]$
w_0	0.3077(65)	---
$w_0 \& w_2$	0.3091(69)	-0.0059(13)
$w_0 \& w_3$	0.3080(70)	-0.0070(12)
$w_0 \& w_4$	0.3079(70)	-0.0068(12)

$$w_0 = 1$$

$$w_2 = 1 - y^2$$

$$w_3 = (1 - y)^2 (1 + 2y)$$

$$w_4 = (1 - y^2)^2$$

- Several fits, single moments & combined
- Many fit windows $[s_{\min}, m_c^2]$
- Consistency: α_s , condensates & DV parameters.

⇓

$$\alpha_s(m_\tau) = 0.3077 \pm 0.0075 \quad (n_f=3, \text{FOPT})$$

$$\alpha_s(M_Z) = 0.1171 \pm 0.0010 \quad (n_f=5, \overline{MS})$$

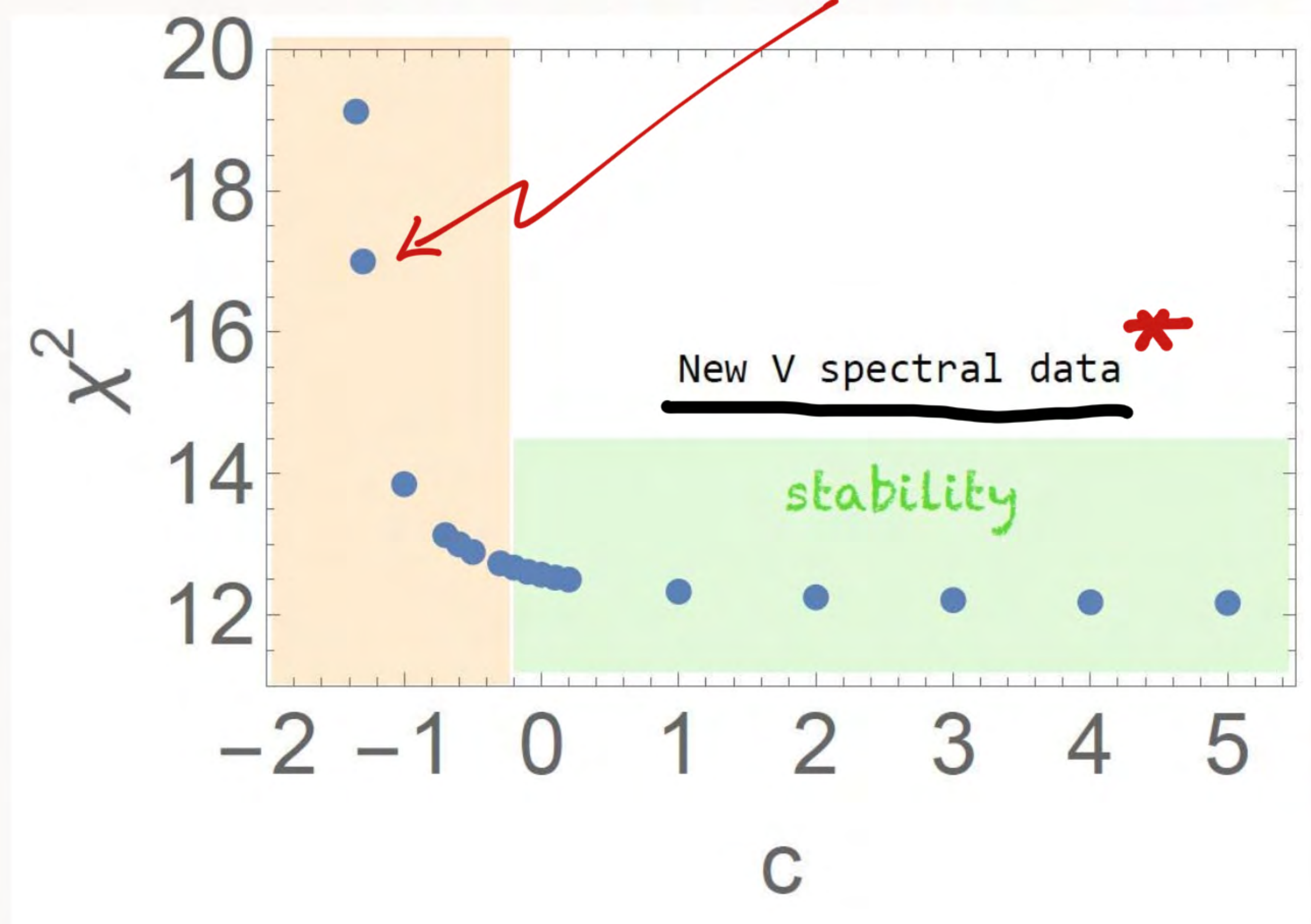
pert. series truncation
 &
 scale variation
 &
 DVs

Response to Criticisms from Pich & Rodriguez Sanchez (JHEP '22)

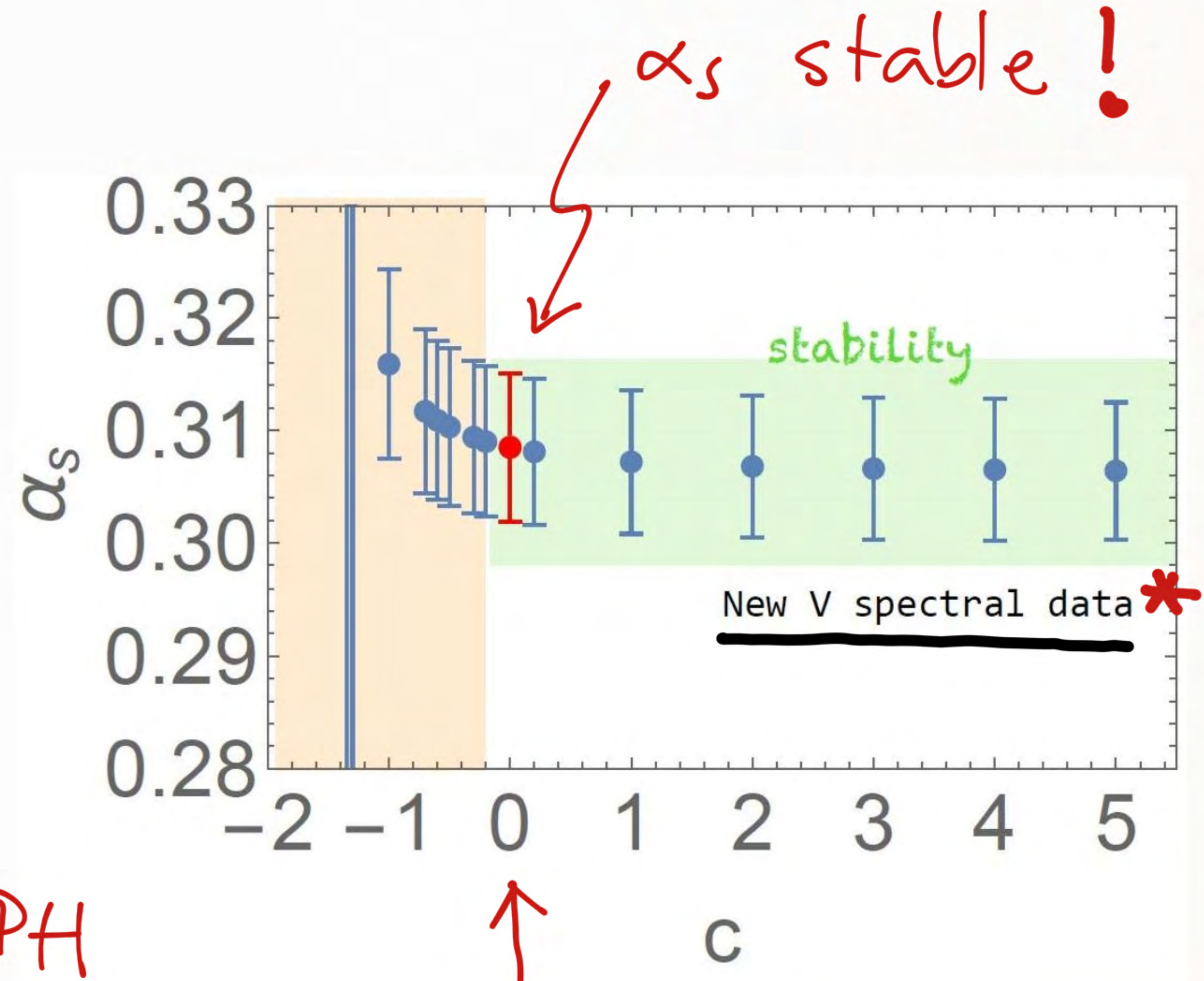
• "DVs unstable against variations"

$$P_{DV}(s) = e^{-\delta - \delta s} \sin(\alpha + \beta s) \left(1 + \frac{c}{s}\right) ; \quad 1.55 \text{ GeV}^2 \leq s_0 \leq 3.06 \text{ GeV}^2$$

choice of P-RS (w ALEPH data)



* better than ALEPH



PRS also consider: $\rho_{DV}(s) = s^\lambda e^{-\delta-\gamma s} \sin(\alpha+\beta s)$, $\lambda = 0, \dots, 8$

i) and the difference between $\lambda=0$ and $\lambda=8$ to claim an error in α_s of 0.016
However, a discrepancy in α_s of 0.034 (Table 2, PRS'22) is "amazing stability".

ii) totally ad-hoc: why not $\lambda=20$? Regge/Large N_c analysis leads to $\lambda=0$
+ possible $1/5$ corrections. Can you justify $\lambda=8$?

iii) Our improved V spectrum shows that p -value \downarrow when $\lambda \uparrow$.
So, $\lambda=8$ also disfavored by data.

we'll stick to $\lambda=0$.

- "logarithms in condensates invalidate analysis"

$$C_D^L(s) = C_D(\mu^2) \left(1 + L_D \log \frac{-s}{\mu^2} \right)$$

PRS choose $L_D \sim 0.2$

However $L_4 = \frac{11}{8} \left(\frac{\alpha_s}{\pi} \right)^2 \approx 0.012$ & $L_6 = -\frac{19}{63} \left(\frac{\alpha_s}{\pi} \right) \approx 0.03 \ll 0.2$

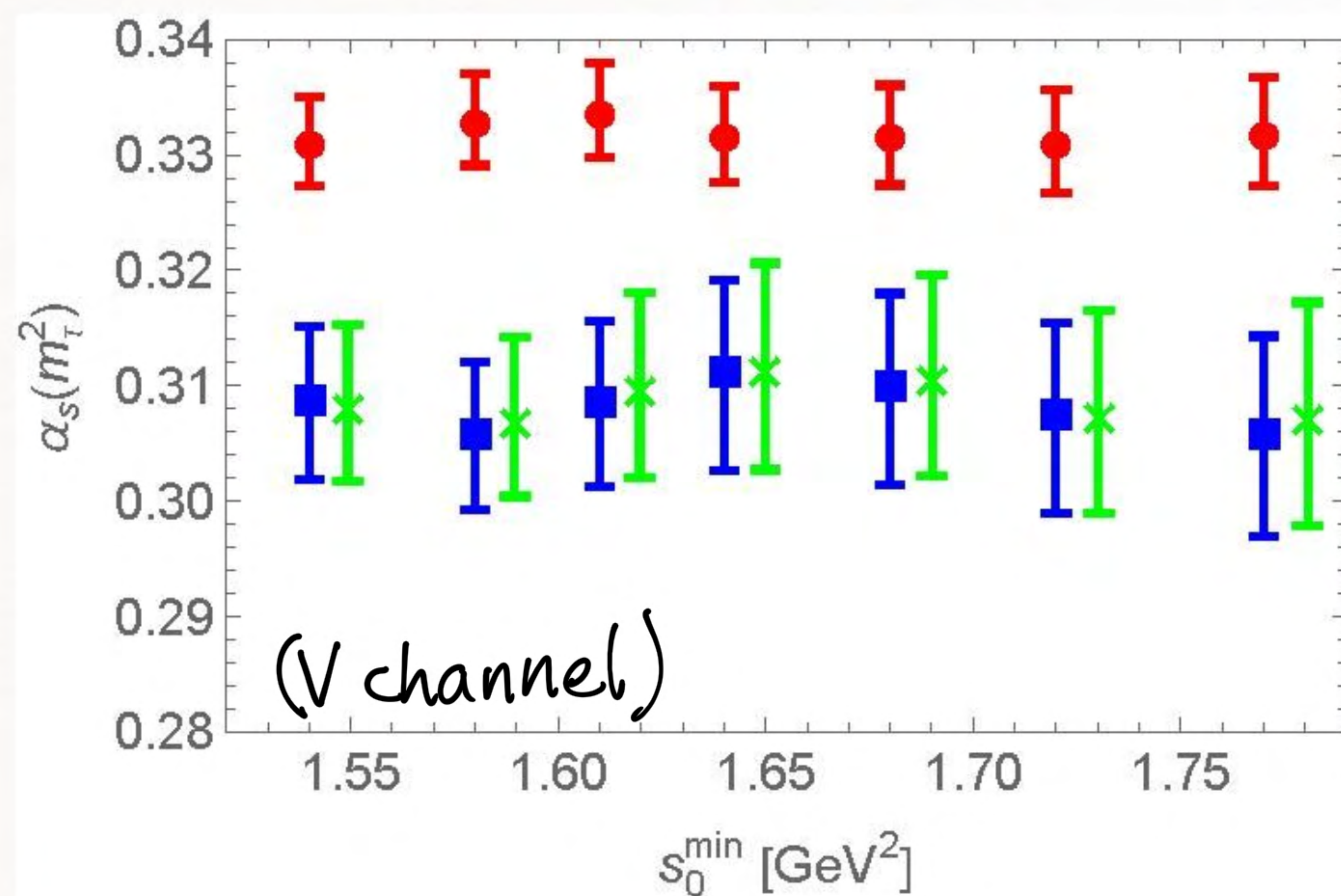
(large N_c factorization)

These values of C_4^L & C_6^L do not affect our analysis.

In fact $L_D \neq 0$ would affect FESR s_0 -scaling: not seen in our fits.

(And recall OPE asymptotic: unreliable at large condensate dimension.)

- Tautology: " α_s determined by $w_0=1$, other weights redundant"



- Fit to $w_0=1$
- Fit to $w_0=1, w_2=1-x^2, OPE \sim 1/s_0^3$ (right)
- Fit " " " but $OPE \sim 1/s_0^5$ (wrong!)



NO REDUNDANCY! (math error in PRS)
cf. Backup slides

True Redundancy in tOPE

and, in general, in single- s_0 analyses.

- Take, e.g., the optimal weights: $w_{2m}(x) = 1 - (m+2)x^{m+1} + (m+1)x^{m+2}$
 $(m=1, \dots, 5)$ $(Pich, Rodriguez S. '16)$

OPE balance :

$w_{21} \rightarrow C_6, C_8$	$w_{23} \rightarrow C_{10}, C_{12}$	$w_{25} \rightarrow C_{14}, C_{16}$
$w_{22} \rightarrow C_8, C_{10}$	$w_{24} \rightarrow C_{12}, C_{14}$	

Following ALEPH (OPAL) analysis, set $C_{10} = C_{12} = C_{14} = C_{16} = 0$ $\underbrace{V+A}_{s_0 = 2.8 \text{ GeV}^2}$

- Fits :
- $w_{23}, w_{24}, w_{25} \rightarrow \alpha_s = 0.3125(23)$, NO OPE AT ALL! ($\chi^2 = 11.6/2$)
 - add $w_{22} \rightarrow \alpha_s = 0.3125(23)$, determines C_8 ($\chi^2 = 11.6/2$)
 - add $w_{21} \rightarrow \alpha_s = 0.3125(23)$, determines C_6 (same C_8) ($\chi^2 = 11.6/2$)

These are not good fits & C_8, C_6 very unstable wrt \otimes

Therefore

i) value of α_s determined by Pert. Theory only
would get same value even if, e.g., $\pi_{\text{OPE}}(q^2) \sim (\log q^2)^{100}$ "
(or any other behavior)

ii) i.e. same as in QED
might as well calculate the integral of the spectral function directly.
Then, why bother about FESRs, Cauchy contours, etc...?

$$\int_0^{s_0} dt w(t) \rho(t)_{\text{EXP}} = \int_0^{s_0} dt w(t) \rho(t)_{\text{Pert. theory}}$$

let's focus on this last example on V+A (ALEPH):

$s_0 = 2.8 \text{ GeV}^2$
 $\alpha_s(m_\tau)_{\textcircled{1}} - \alpha_s(m_\tau)_{\textcircled{2}} = 0.0103 (38)_{\text{exp \& th}}$, 2.75

(-4 bin) \downarrow (-60 MeV)

$s_0 = 2.6 \text{ GeV}^2$
 $\alpha_s(m_\tau)_{\textcircled{1}} - \alpha_s(m_\tau)_{\textcircled{2}} = 0.0102 (29)_{\text{exp \& th}}$, 3.65

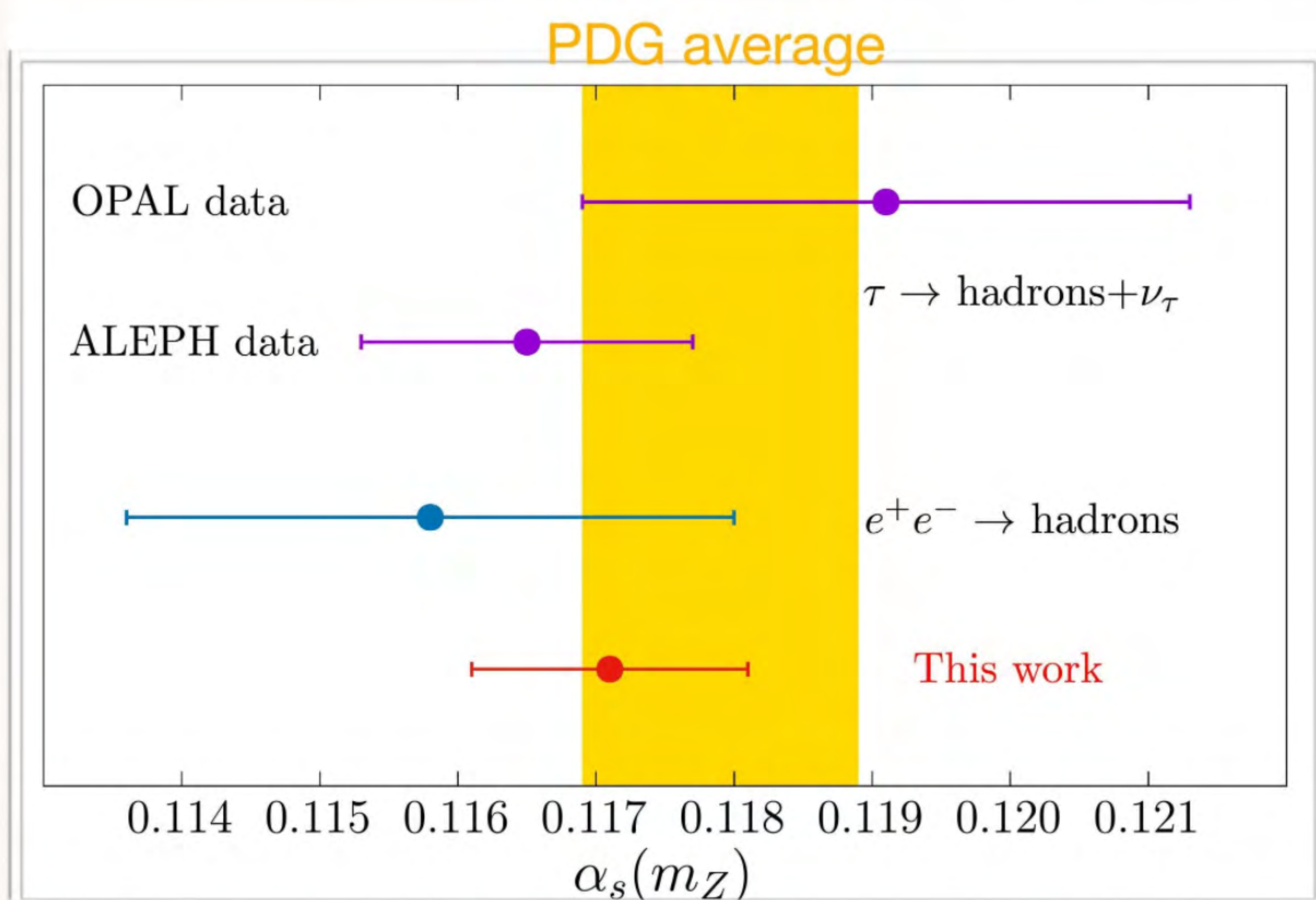
(includes correlations)

} $\textcircled{1}$ W_{23}, W_{24}, W_{25}
 } $\textcircled{2}$ W_{25}

*** SIGNIFICANT DISCREPANCY ***

(Even worse problems using our improved V spectral function: \sim 8-10 σ)

CONCLUSIONS



DB, Golterman, Jamin, Mahdavi, Maltman, Osborne, Peris, '12

DB, Golterman, Maltman, Osborne, Peris, '15

DB, Golterman, Keshavarzi, Maltman, Nomura, Peris, Teubner '18

DB, Golterman, Maltman, Peris, Rodrigues, Schaaf, '21

$$\alpha_s(m_\tau) = 0.3077 \pm 0.0075$$



$$\alpha_s(M_Z) = 0.1171 \pm 0.0010 \quad (n_f=5)$$

$$\alpha_s(M_Z) = 0.1180 \pm 0.0009$$

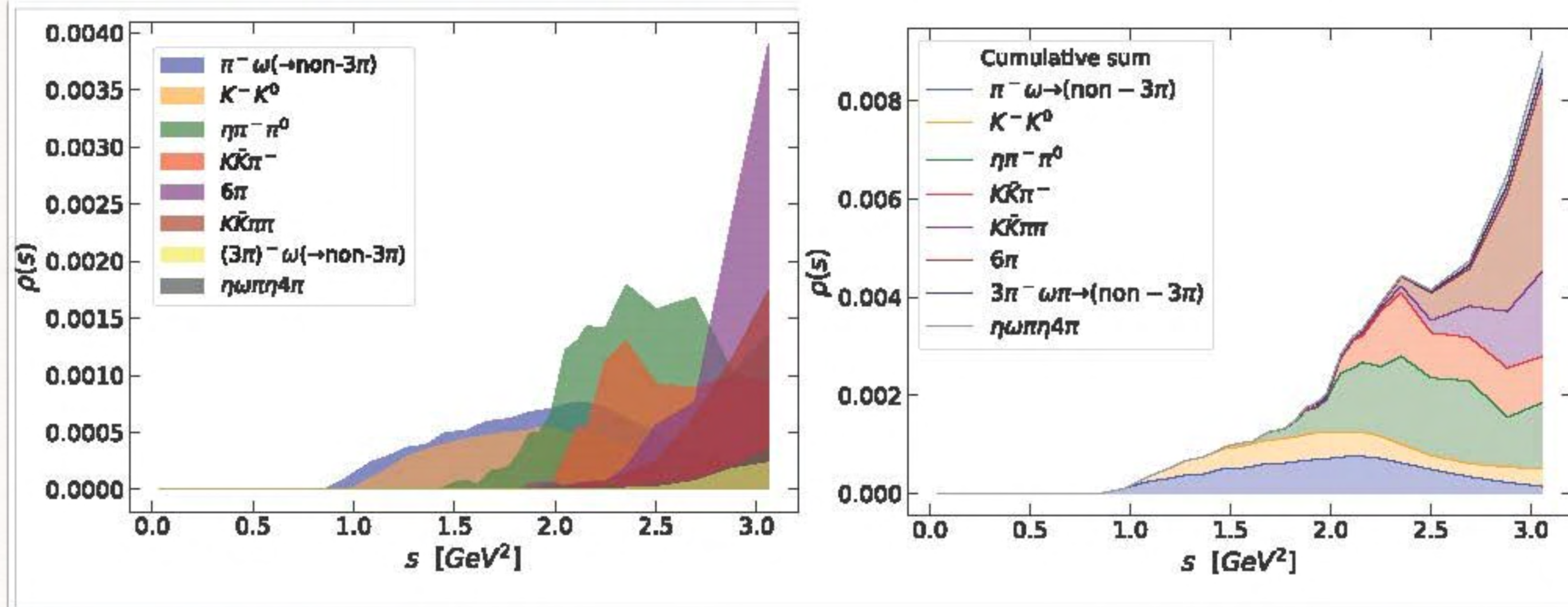
PDG

- Based on new V isovector spect. function w/o Monte Carlo input.
- PRS criticisms *mistaken or unsupported by data.*
- single- s_0 analysis of α_s in $tOPE$ purely perturbative, redundant across set of weights and showing inconsistency in α_s values: *Unreliable*

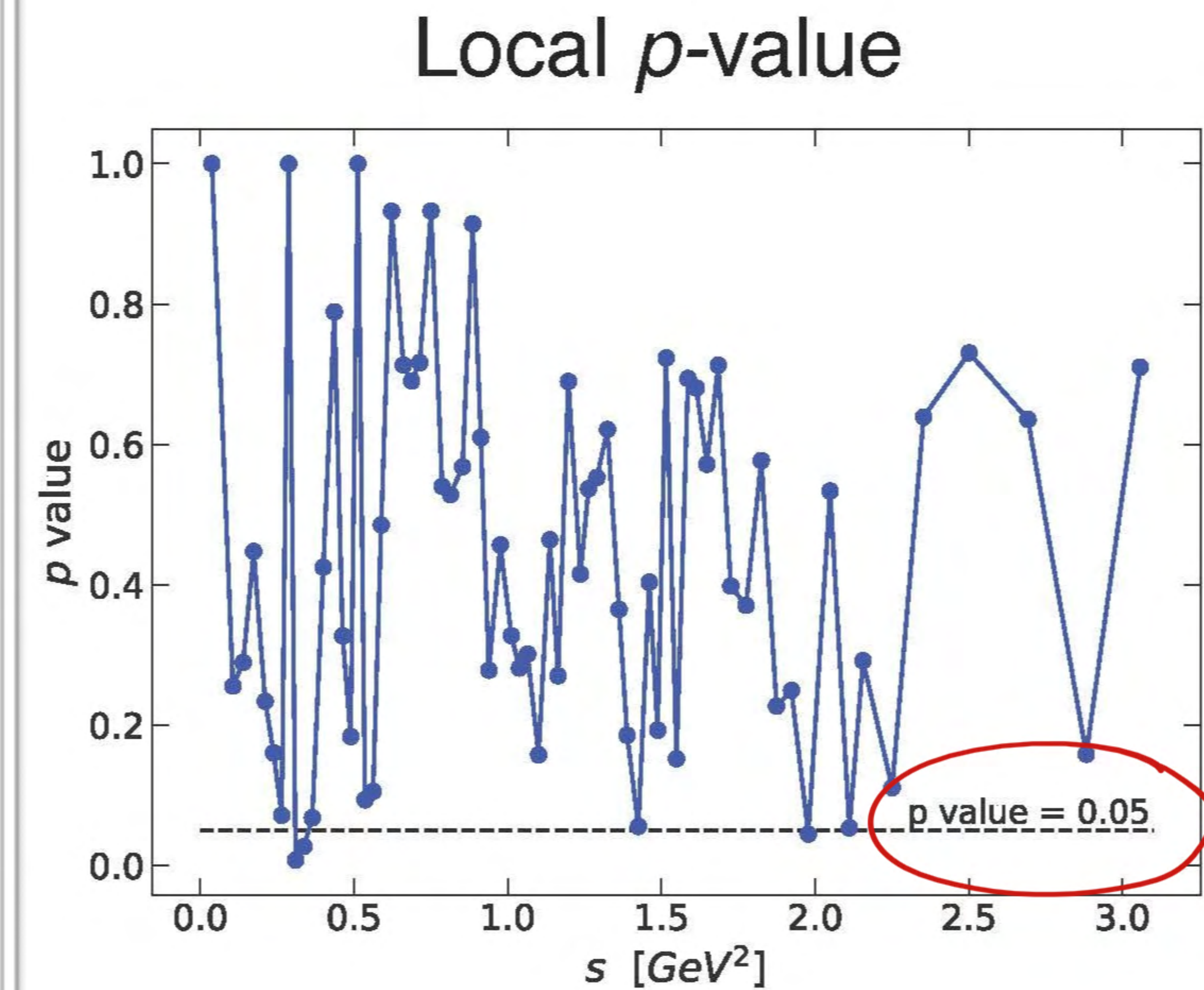
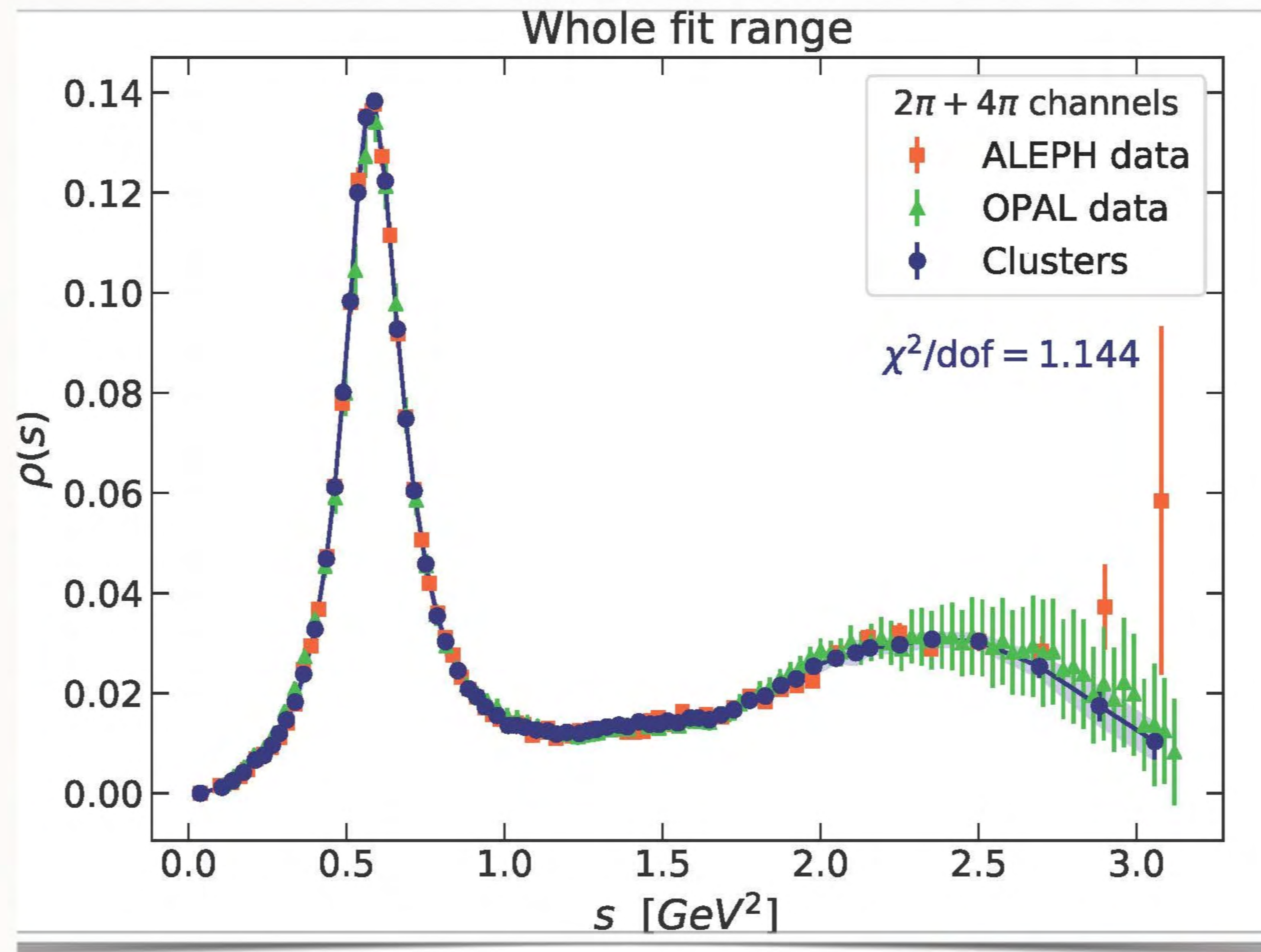
BACK-UP
SLIDES

Improved V spectrum : residual channels ($\sim 2\%$ of total).

Original data sets from: BABAR, SND and CMD-3

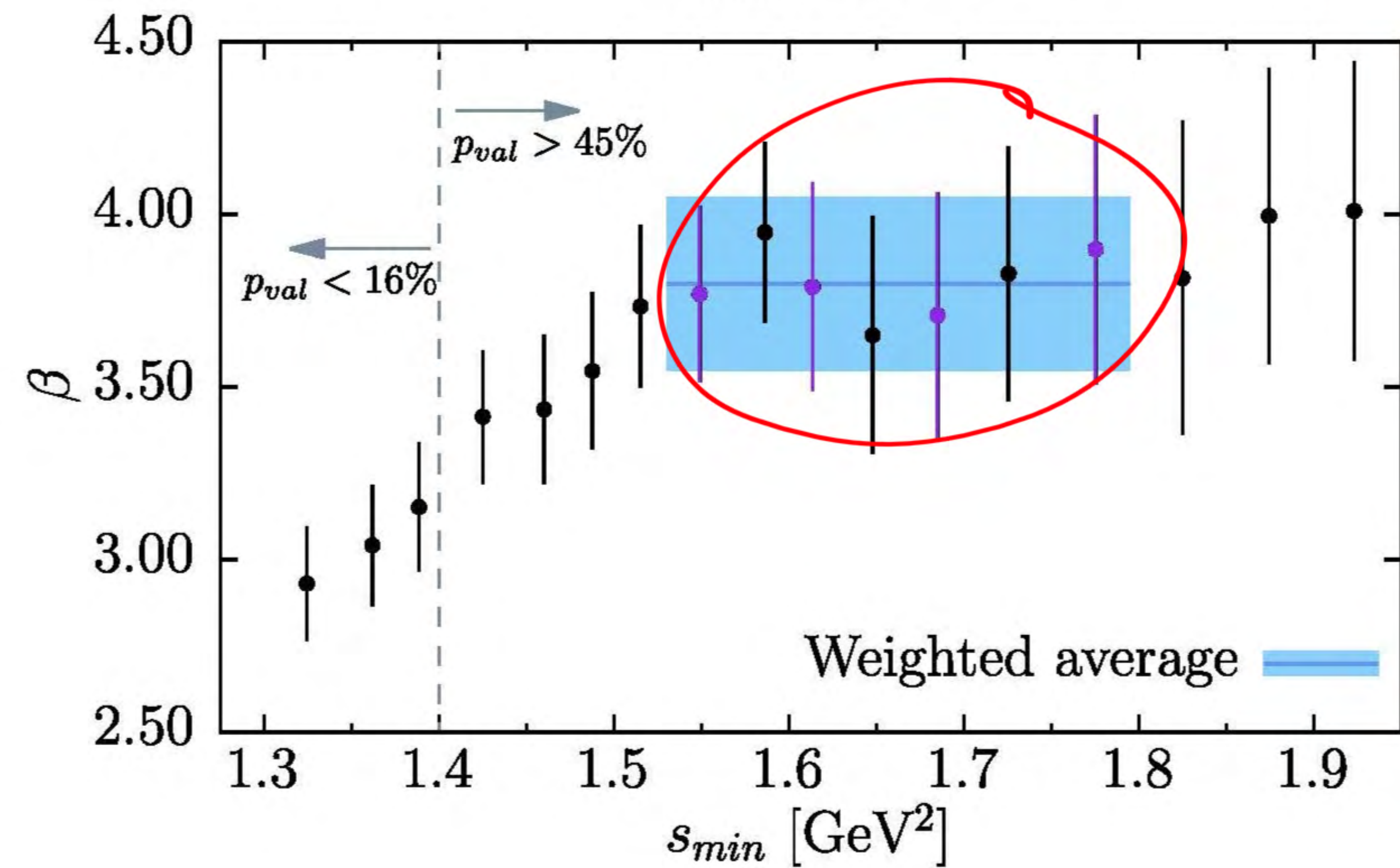
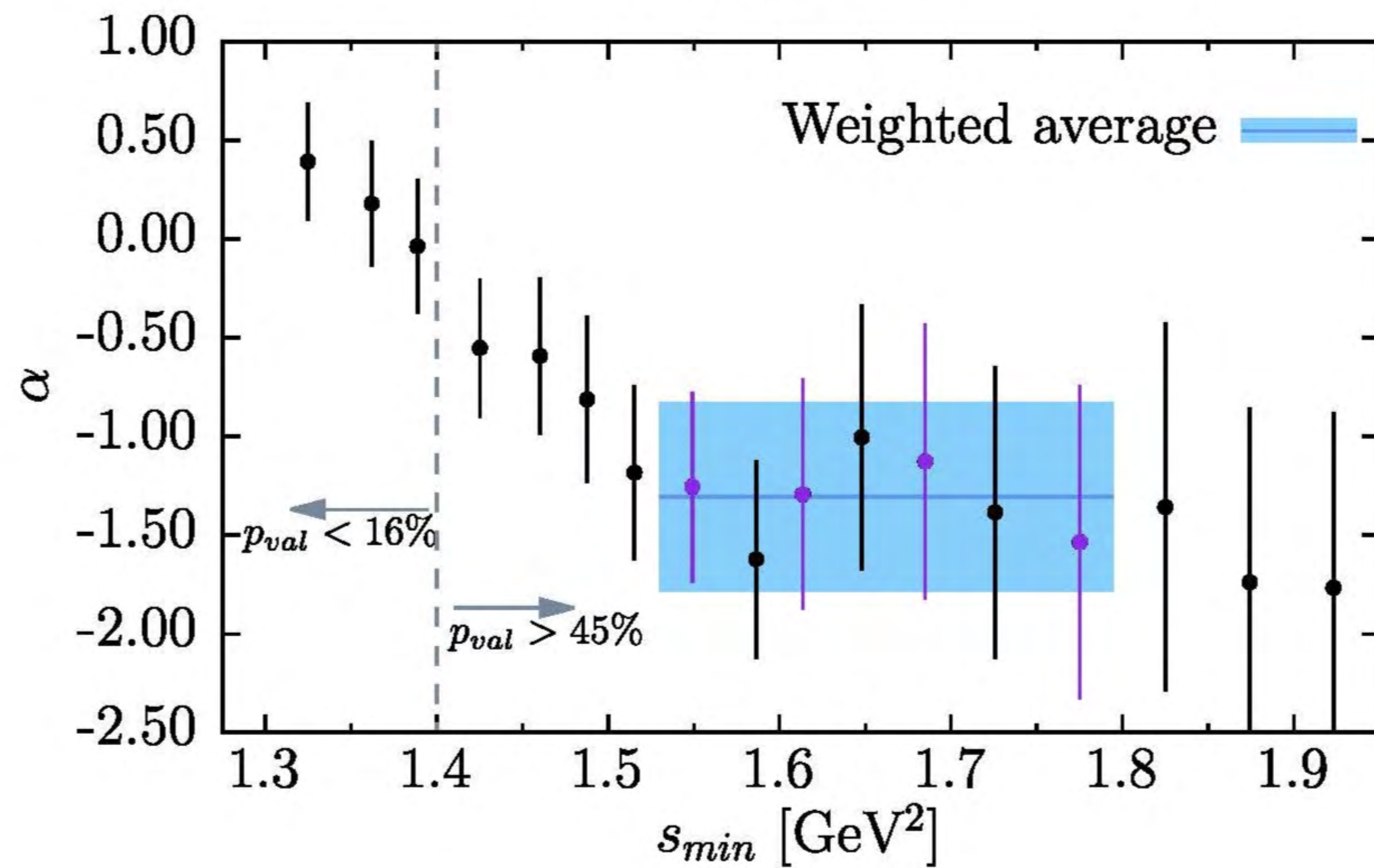
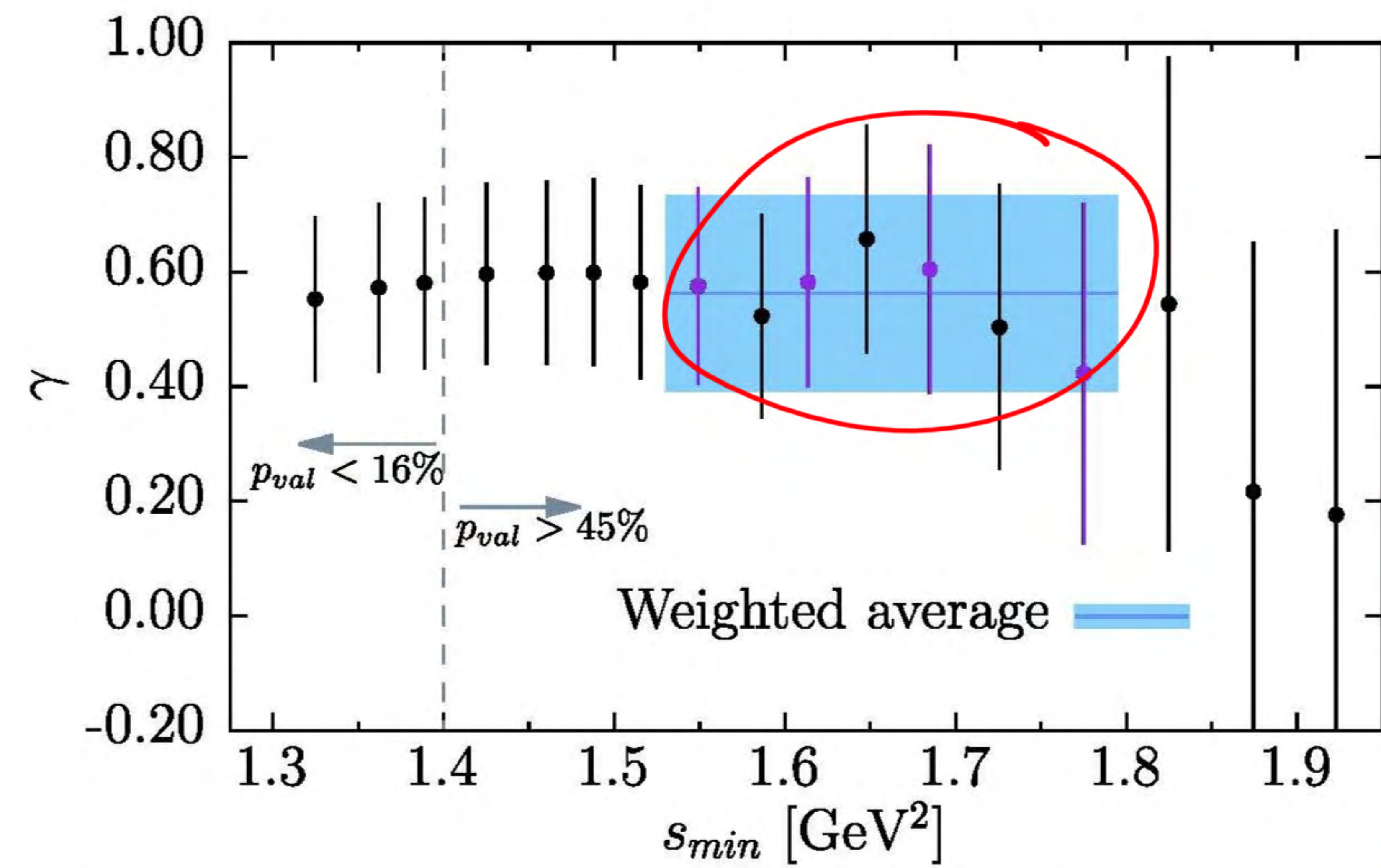
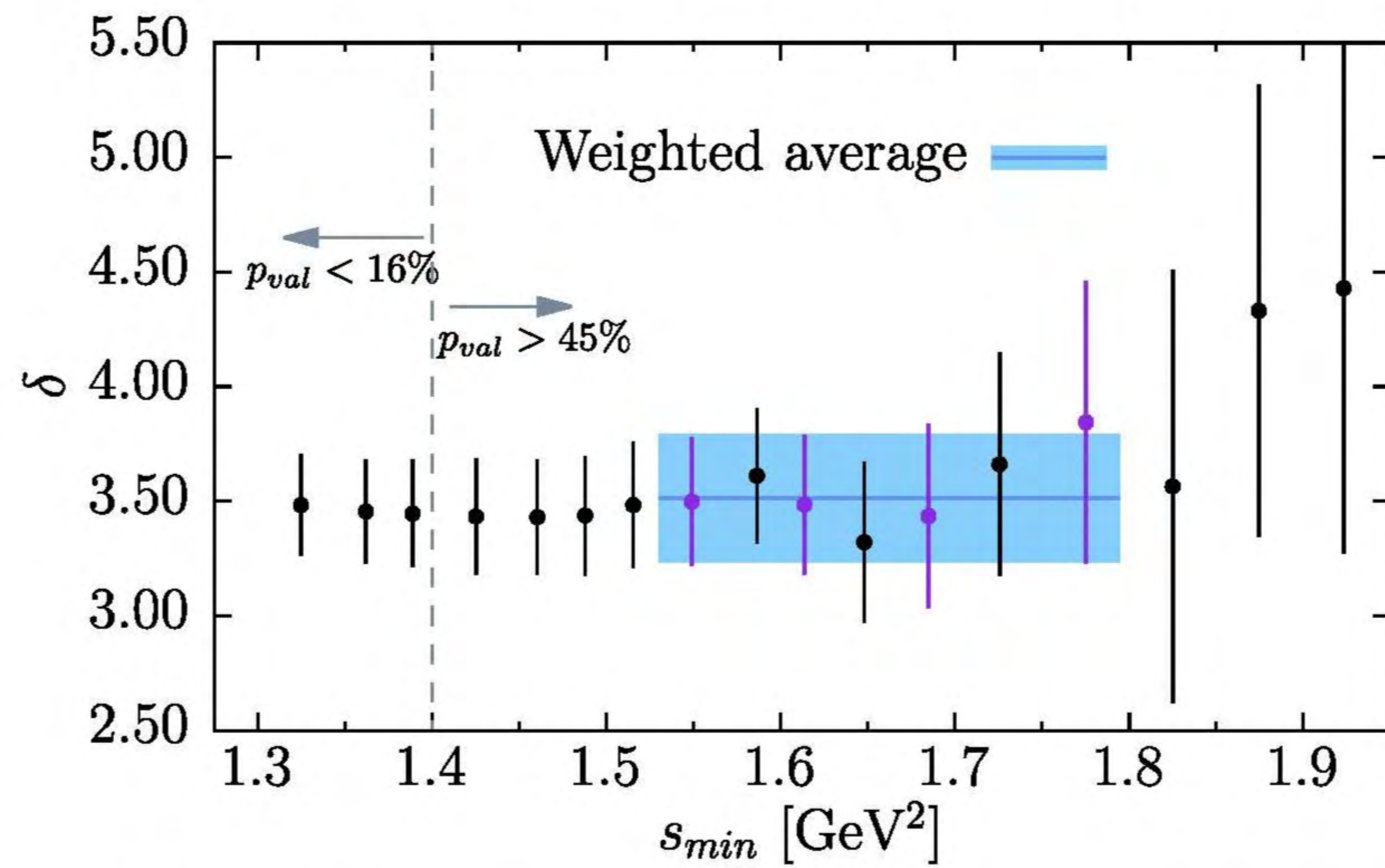


Improved V spectrum



DB, Golterman, Maltman, Peris, Rodrigues and Schaaf, arXiv:2012.10440

DV parameters

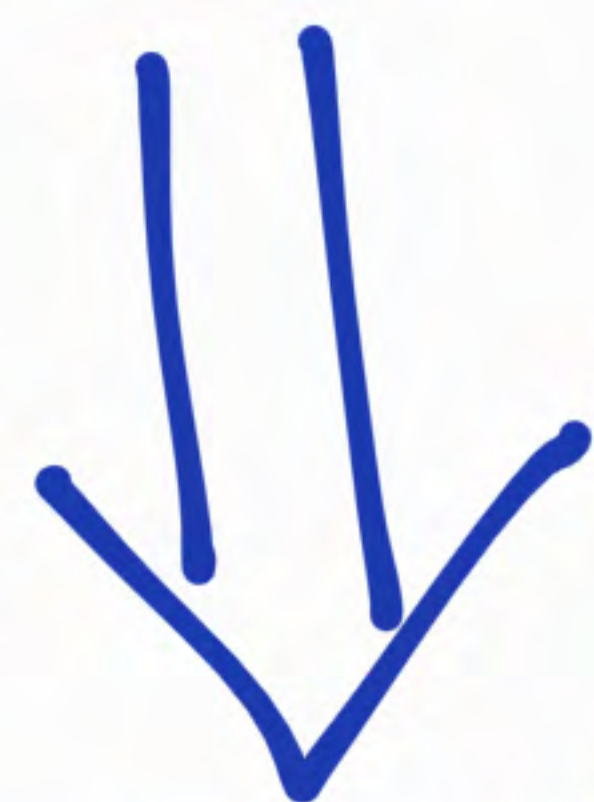
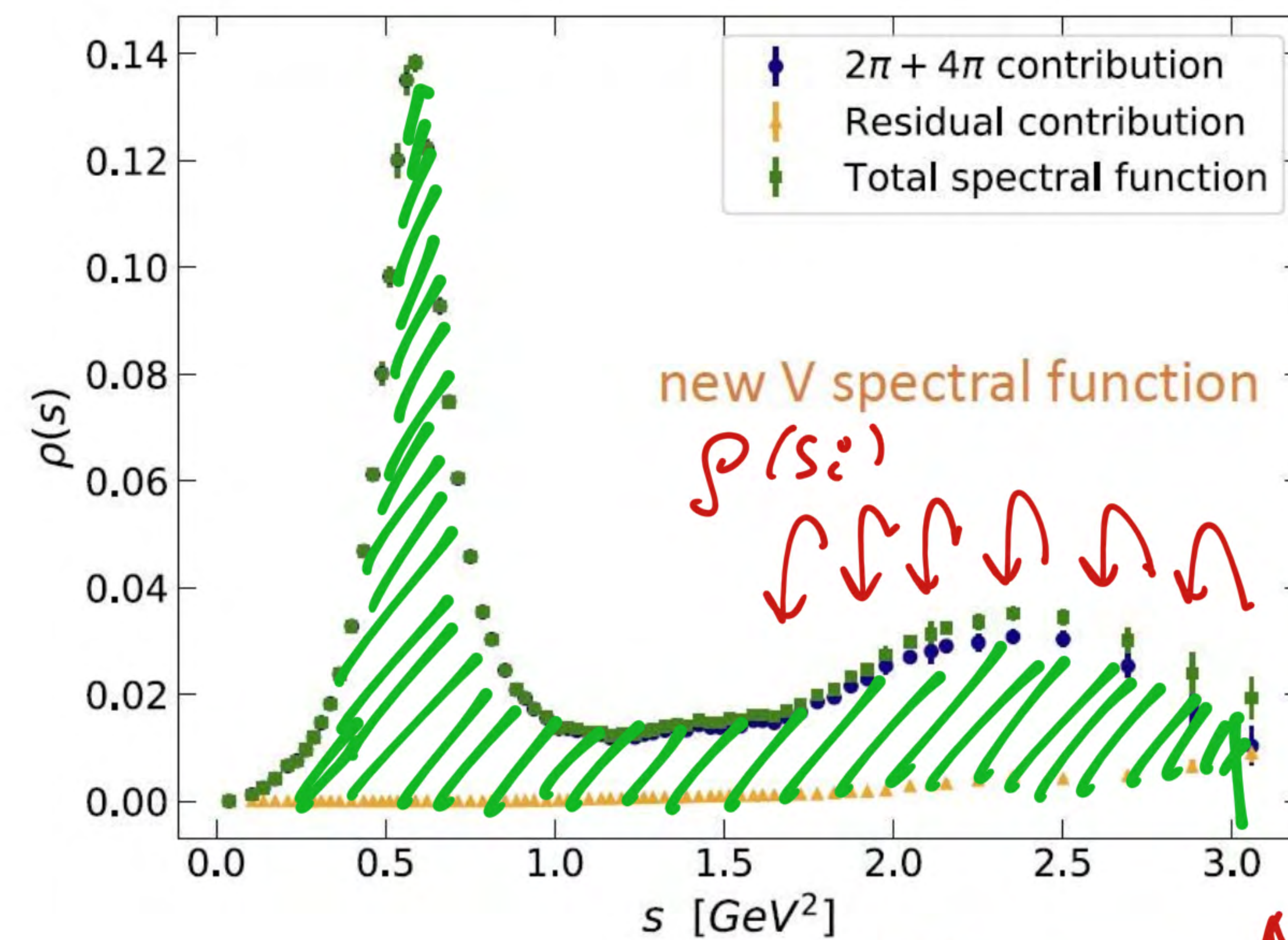
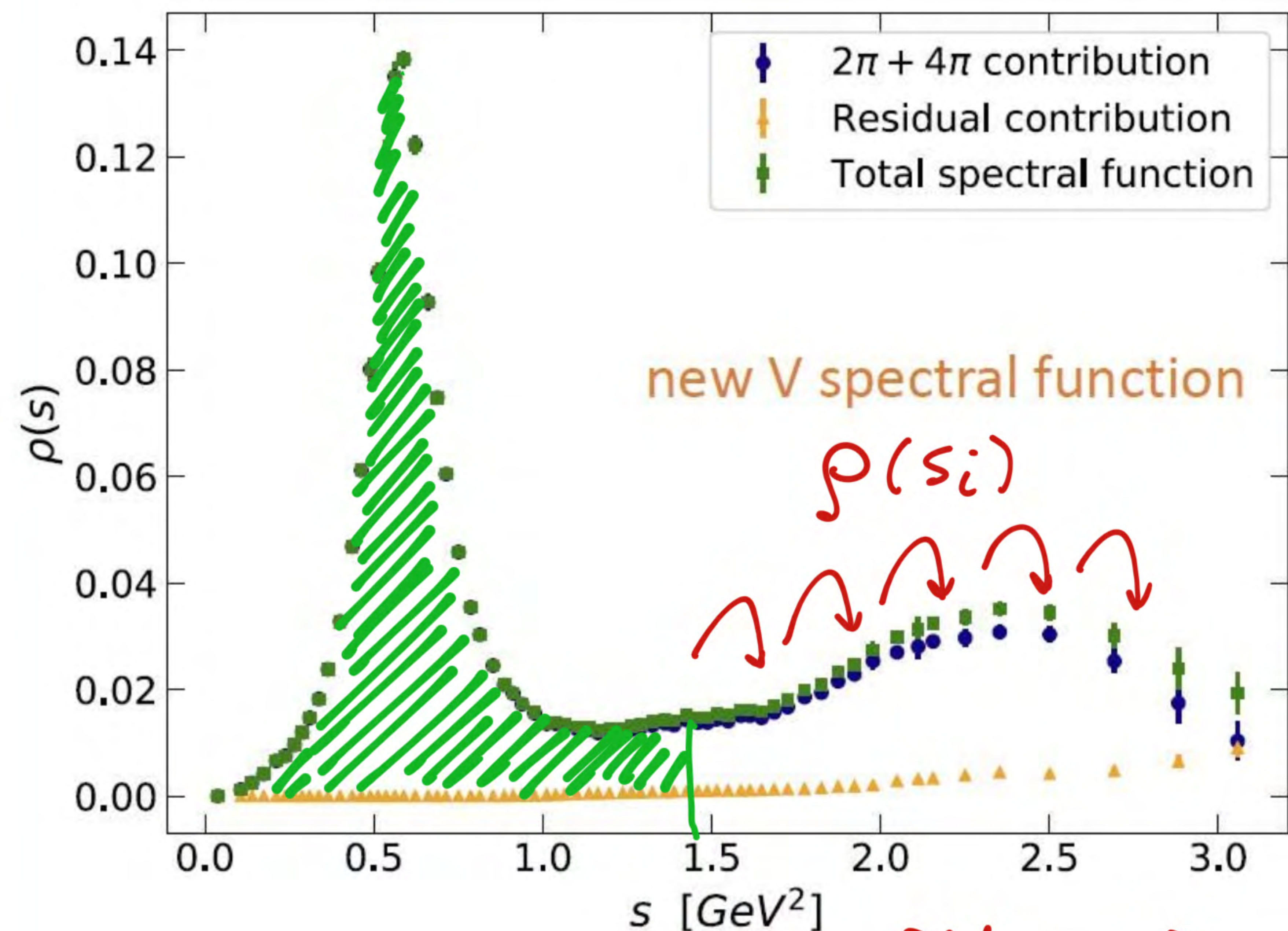


agreement
with
Regge/large N_c
analysis
(not obvious)

Is α_s fixed at $s_0 = 1.55 \text{ GeV}^2$

OR

α_s fixed at $s_0 = m_\tau^2$



NEITHER!

(truth: both α_s & DVs fixed in window $s_0 \in [1.55, m_\tau^2] \text{ GeV}^2$)

Leading dim-6 contribution
to $\delta\mathcal{P}\bar{\tau}$ are 4-quark
condensates:

In the large- N_c approximation, these operators can be expressed, after Fierzing, in terms of $\langle\bar{q}q\rangle$ condensates with the result:

$$\begin{aligned}\langle O_1\rangle &= \frac{1}{2}\langle\bar{u}u\rangle\langle\bar{d}d\rangle, \quad \langle O_2\rangle = -\frac{1}{2}\langle\bar{u}u\rangle\langle\bar{d}d\rangle, \quad \langle O_4\rangle = -\frac{1}{2}\langle\bar{u}u\rangle^2 - \frac{1}{2}\langle\bar{d}d\rangle^2 \\ \langle O_5\rangle &= \frac{1}{2}\langle\bar{u}u\rangle^2 + \frac{1}{2}\langle\bar{d}d\rangle^2, \quad \langle O_7\rangle = -\frac{1}{2}\langle\bar{u}u\rangle^2 - \frac{1}{2}\langle\bar{d}d\rangle^2 - \frac{1}{2}\langle\bar{s}s\rangle^2,\end{aligned}\quad (44)$$

with $\langle O_3\rangle = \langle O_6\rangle = 0$. These condensates give the following contribution to the vector two-point correlator

$$\begin{aligned}(-s)^3\Pi_{V;ud}^{(0+1)}(s) &= -8\pi^2\left[1 - \frac{9}{8}\left(\frac{\alpha_s}{\pi}\right)L\right]\left(\frac{\alpha_s}{\pi}\right)\langle O_1\rangle + 5\pi^2L\left(\frac{\alpha_s}{\pi}\right)^2\langle O_2\rangle \\ &\quad - \frac{8\pi^2}{9}\left[1 - \frac{95}{72}L\left(\frac{\alpha_s}{\pi}\right)\right]\left(\frac{\alpha_s}{\pi}\right)\langle O_4\rangle + \frac{5\pi^2}{9}L\left(\frac{\alpha_s}{\pi}\right)^2\langle O_5\rangle \\ &\quad + \frac{24\pi^2}{81}L\left(\frac{\alpha_s}{\pi}\right)^2\langle O_7\rangle,\end{aligned}\quad (45)$$

with $L = \log(-s/s_0)$, $\alpha_s \equiv \alpha_s(s_0)$ and $\langle O_i\rangle \equiv \langle O_i\rangle(s_0)$. A further simplification may be obtained in the $SU(3)$ limit where $\langle\bar{u}u\rangle = \langle\bar{d}d\rangle = \langle\bar{s}s\rangle \equiv \langle\bar{q}q\rangle$,

$$(-s)^3\Pi_{V;ud}^{(0+1)}(s) = -\frac{28}{9}\pi\rho\alpha_s\langle\bar{q}q\rangle^2\left(1 - \frac{19}{63}\alpha_sL\right),\quad (46)$$

where ρ is a factor parametrizing the possible deviations from the large- N_c and $SU(3)$ limits.

spectral moments: $I_n(s_k) = \sum_{i=1}^k \rho(s_i) s_i^n \Delta$, $s_k \in \{s_{N_0}, \dots, s_{N_0+N}\}$
 bin width

theory representation: $t_n(s_k)$ weight

$$\begin{pmatrix} I_n(s_{N_0}) - t_n(s_{N_0}) \\ I_n(s_{N_0} + \Delta) - t_n(s_{N_0} + \Delta) \\ I_n(s_{N_0} + 2\Delta) - t_n(s_{N_0} + 2\Delta) \\ \vdots \\ I_n(s_{N_0+N}) - t_n(s_{N_0+N}) \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 & 0 & \dots \\ 1 & (s_{N_0} + \Delta)^n \Delta & 0 & 0 & \dots \\ 1 & (s_{N_0} + \Delta)^n \Delta & (s_{N_0} + 2\Delta)^n \Delta & 0 & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 1 & (s_{N_0} + \Delta)^n \Delta & \dots & s_{N_0+N}^n \Delta & \dots \end{pmatrix} \begin{pmatrix} \rho(s_{N_0}) - \frac{I_n(s_{N_0}) - t_n(s_{N_0})}{(s_{N_0})^n \Delta} \\ \rho(s_{N_0} + \Delta) - \frac{t_n(s_{N_0} + \Delta) - t_n(s_{N_0})}{(s_{N_0} + \Delta)^n \Delta} \\ \rho(s_{N_0} + 2\Delta) - \frac{t_n(s_{N_0} + 2\Delta) - t_n(s_{N_0} + \Delta)}{(s_{N_0} + 2\Delta)^n \Delta} \\ \vdots \\ \rho(s_{N_0+N}) - \frac{t_n(s_{N_0+N}) - t_n(s_{N_0+N} - \Delta)}{s_{N_0+N}^n \Delta} \end{pmatrix}$$

If $t_n(s_k)$ neglected, looks like a fit to $\rho(s_k)$ independent of weight "n".
 However, when $t_n(s_k)$ included, there is nontrivial constraint on t_n & t_m .

$$\frac{t_n(s_{N_0} + k\Delta) - t_n(s_{N_0} + (k-1)\Delta)}{(s_{N_0} + k\Delta)^n} = \frac{t_m(s_{N_0} + k\Delta) - t_m(s_{N_0} + (k-1)\Delta)}{(s_{N_0} + k\Delta)^m}$$

ALPHA V+A $s_0 = 2.8 \text{ GeV}^2$

Redundancy in tOPE strategy :

Assumption $C_{12} = C_{14} = C_{16} = 0$

"Optimal weights" (PRS'16):

$$w_{21} = 1 - 3y^2 + 2y^3; \quad w_{22} = 1 - 4y^3 + 3y^4$$

$$w_{23} = 1 - 5y^4 + 4y^5$$

$$\begin{aligned} w_{24} &= 1 - 6y^5 + 5y^6 \\ w_{25} &= 1 - 7y^6 + 6y^7 \end{aligned}$$

\Rightarrow fix α_s
w/o OPE!

Results

\Rightarrow REDUNDANCY

$$w_{24} \oplus w_{25}: \quad \alpha_s = 0.3168(27), \quad \chi^2 = 3.06933$$

$$\oplus w_{23}: \quad \alpha_s = 0.3168(27), \quad \chi^2 = 3.06933$$

$$\oplus w_{22}: \quad \alpha_s = 0.3168(27), \quad \chi^2 = 3.06933$$

$$\oplus w_{21}: \quad \alpha_s = 0.3168(27), \quad \chi^2 = 3.06933$$

$$C_{10} = -41(25) \times 10^{-4} \text{ GeV}^{10}$$

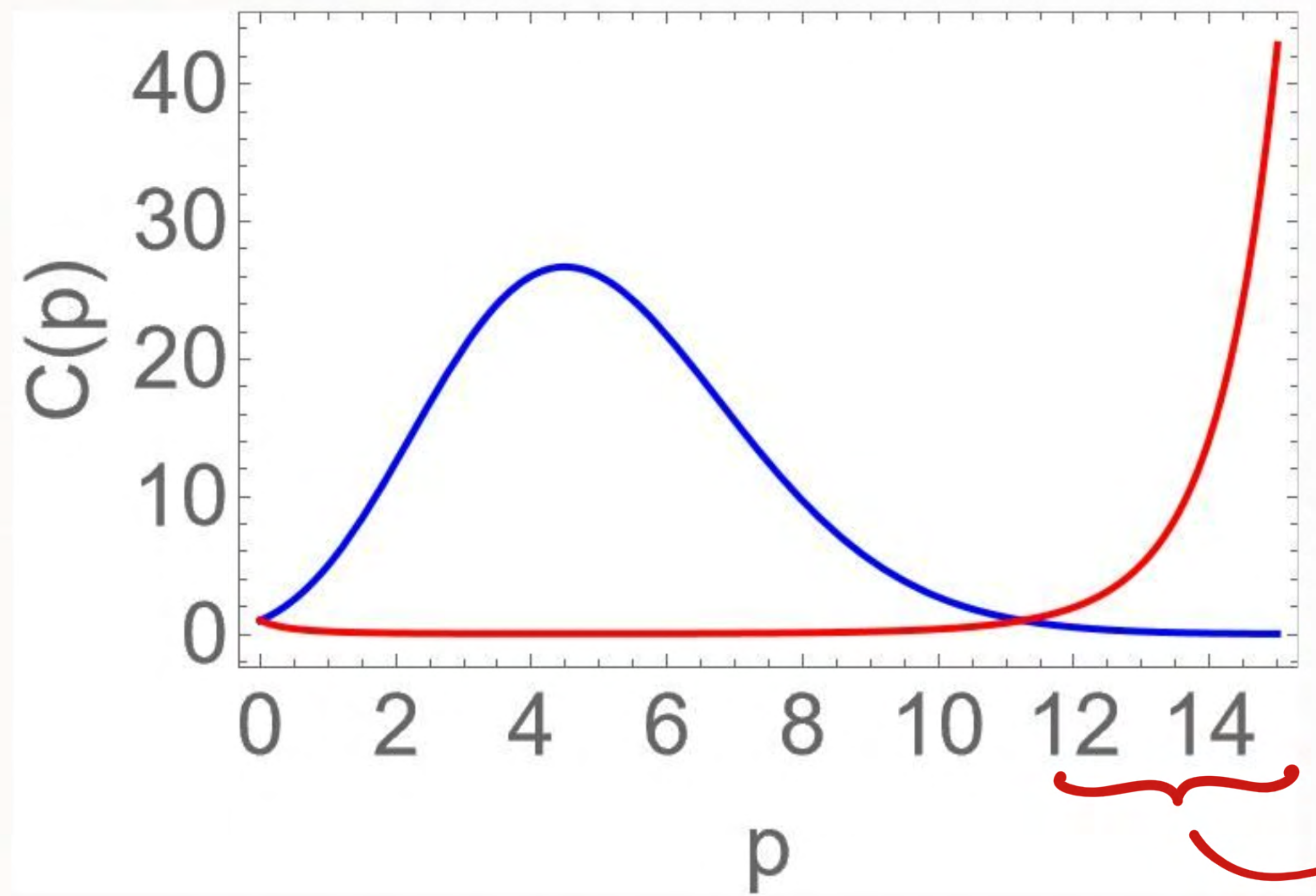
$$C_{10} = \text{"}, \quad C_8 = 16(14) \times 10^{-4} \text{ GeV}^8$$

$$C_{10} = \text{"}, \quad C_8 = \text{"}, \quad C_6 = 5(5) \times 10^{-4} \text{ GeV}^6$$

$$w_{25} \text{ only}: \quad \alpha_s = 0.3202(24)$$

Test Pert. Theory, not the OPE!

Asymptotic vs. Convergent



$$f(z) = \sum_{p=1}^N C(p) z^p \quad \text{with } N \rightarrow \infty$$

$$C(p) = \left(\frac{5^p}{p!} \right), \text{ convergent}$$

$$C(p) = \left(\frac{p!}{5^p} \right), \text{ Asymptotic}$$

Must avoid being sensitive to high-order contributions

[Poincare (Circa 1893): discussion Geometers vs. Astronomers]