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Matthias Schott on behalf of the ATLAS Collaboration

# Measurement of the Strong Coupling Constant with ATLAS





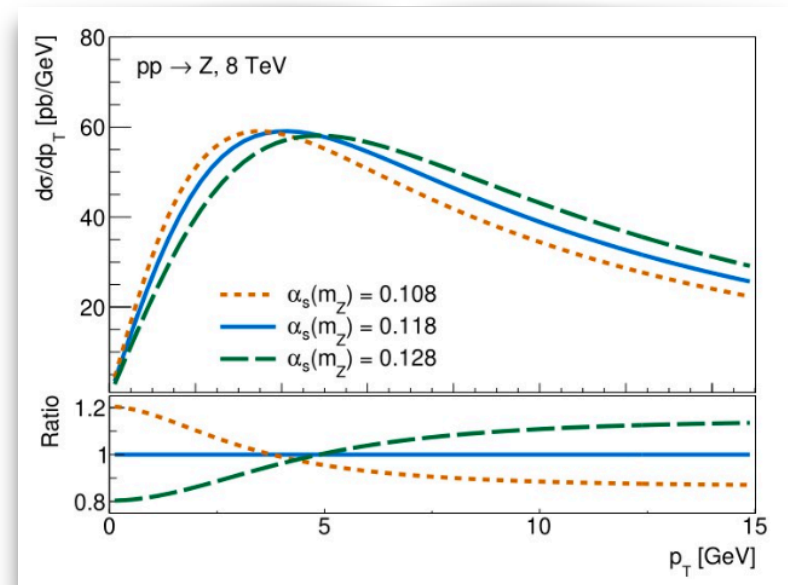
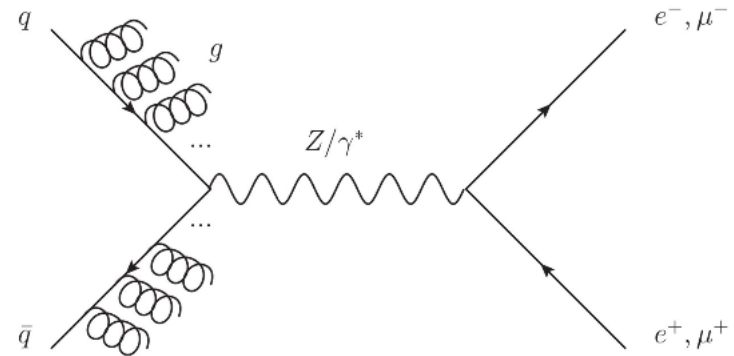
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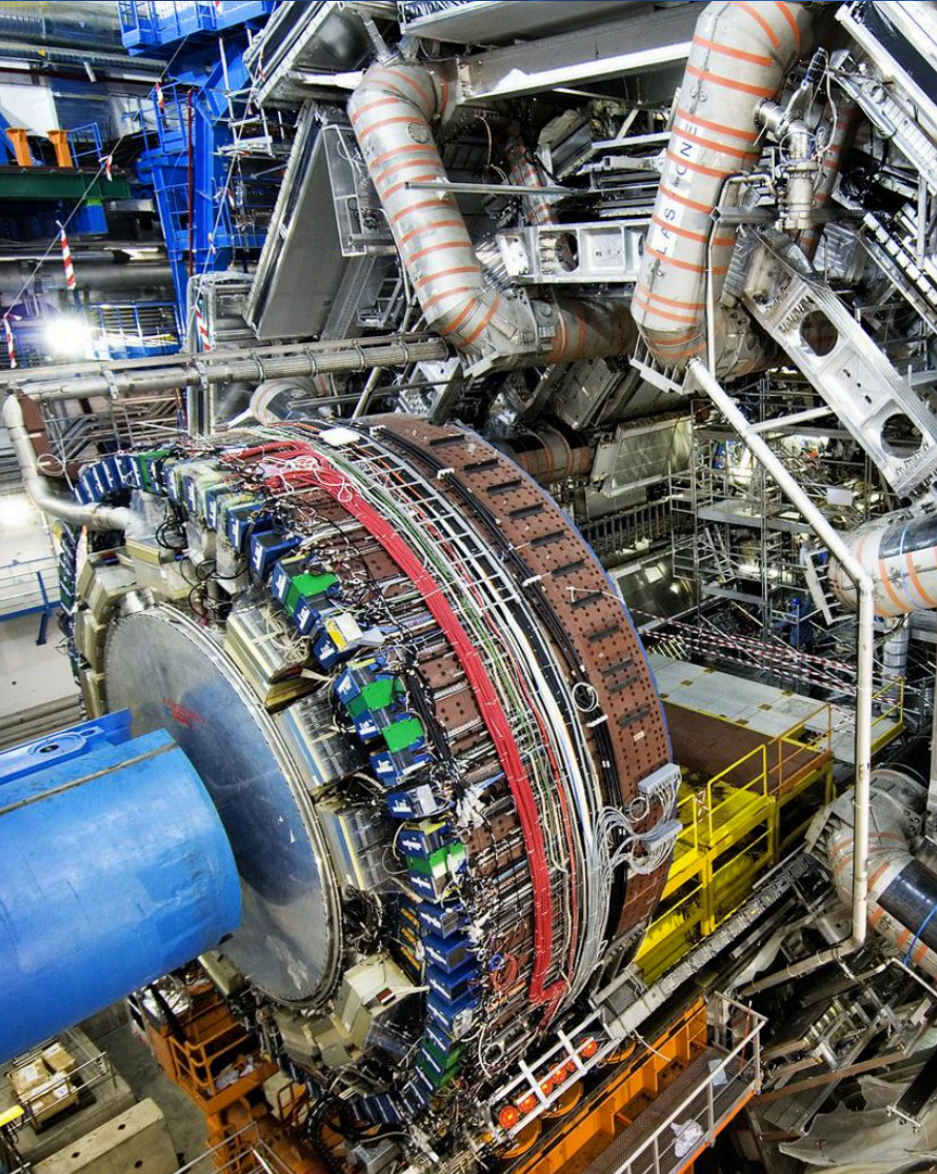
Drell-Yan Process

# Measure $\alpha_s(m_Z)$ from the $p_T(Z)$ distribution

- Z bosons produced in hadron collisions recoil against QCD initial-state radiation:
  - by momentum conservation, ISR gluons will boost the Z in the transverse plane
- The Sudakov factor is responsible for the existence of a peak in the Z-boson  $p_T$  distribution, at values of approximately 4 GeV
- The position of the peak is sensitive to  $\alpha_s(m_Z)$
- Semi-inclusive observable, which has advantages of
  - exclusive obs. (higher exp. sensitivity)
  - inclusive obs. (higher order theory, smaller non-pQCD effects)



# Anatomy of Drell-Yan differential cross section



- DY not only interesting physics wise, but required for all aspects of detector calibration
- DY cross section: factorize the production dynamic and the decay kinematic properties of the dilepton system

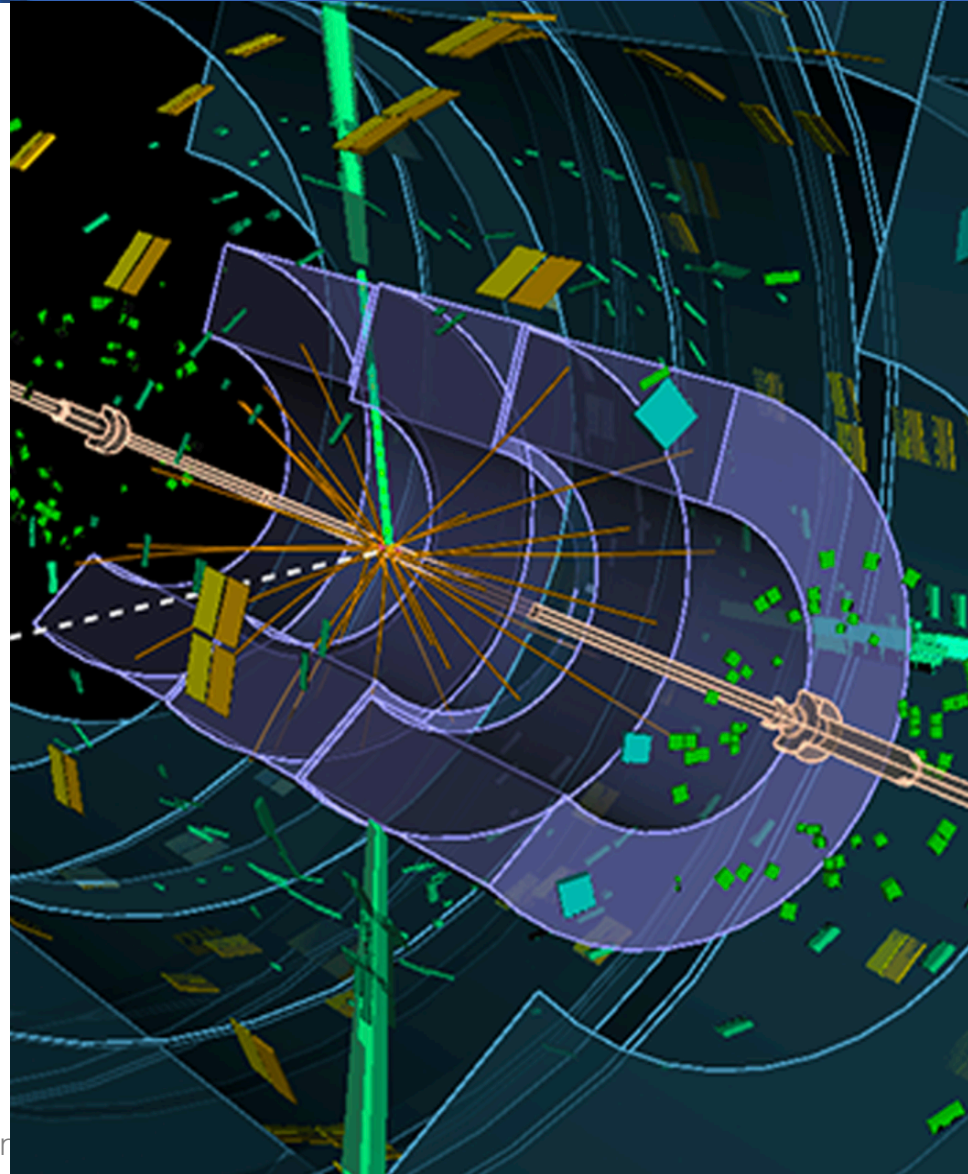
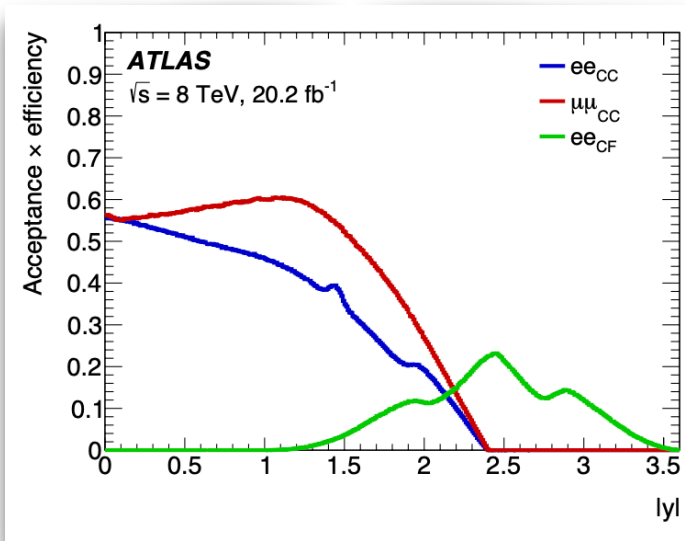
$$\frac{d\sigma}{dpdq} = \frac{d^3\sigma^{U+L}}{dp_T dy dm} \left( 1 + \cos^2 \theta + \sum_{i=0}^7 A_i(y, p_T, m) P_i(\cos \theta, \phi) \right)$$

- $ds/dp_T$ : transverse dynamics
- $ds/dy$ : longitudinal dynamics (PDFs)
- decomposition of  $(\cos \theta, \Phi)$  into 9 helicity cross sections
- basis of spherical harmonics

# Event selection

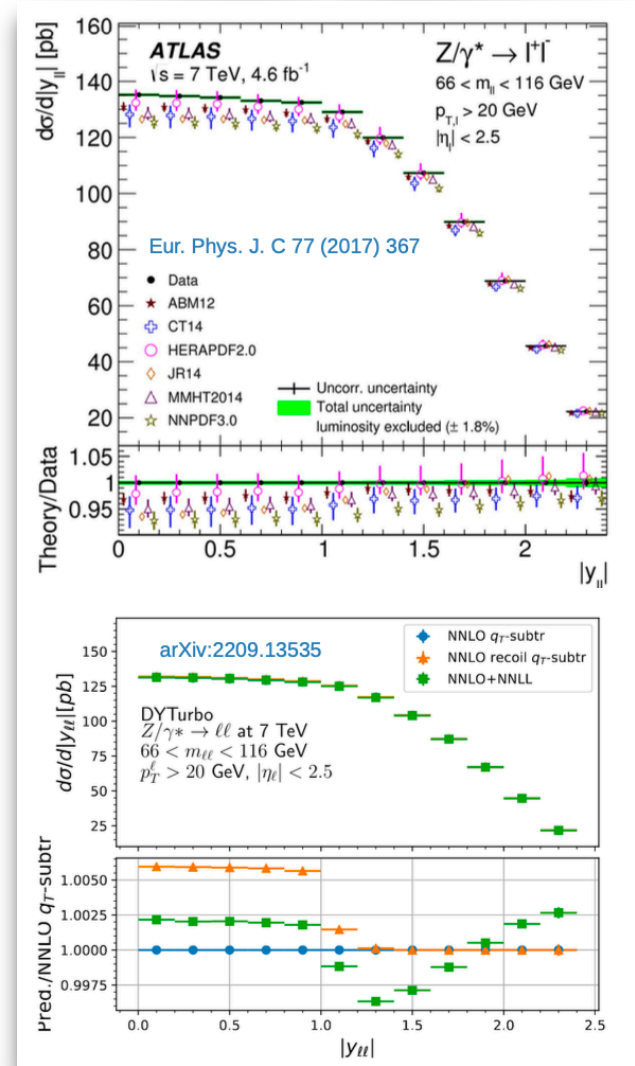
## ■ Three Channels

- eeCC: two electrons (6.2M events)  
 $p_T > 20$  GeV,  $|\eta| < 2.4$
- $\mu\mu$ CC: two muons (7.8M events)  
 $p_T > 20$  GeV,  $|\eta| < 2.4$
- eeCF: central electron with  $p_T > 25$  GeV,  
 $|\eta| < 2.4$ , forward electron with  $p_T > 20$   
GeV,  $2.5 < |\eta| < 4.9$  (1.2M events)
- $80 < m_{ll} < 100$  GeV



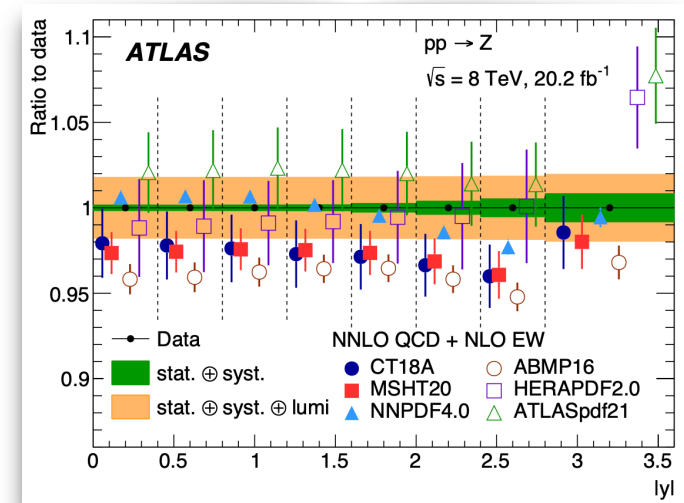
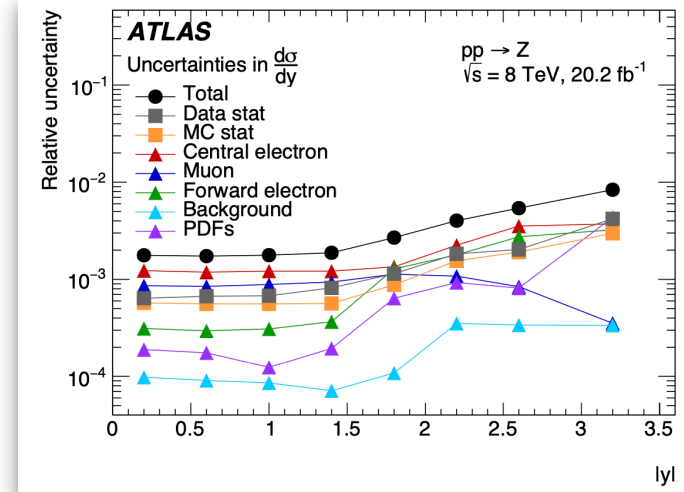
# Full-lepton phase space rapidity cross section

- Interpretation of fiducial cross sections hampered by breakdown of fixed order perturbation theory
  - Problem: low  $p_T(Z)$  spectrum impacts  $p_T$  lepton spectrum
- Proposed solutions:
  - Change the definition of fiducial cuts [arXiv:2106.08329](https://arxiv.org/abs/2106.08329)
  - Use Ai theory predictions to extrapolate the measured cross sections [arXiv:2001.02933](https://arxiv.org/abs/2001.02933)
  - Include resummation corrections into predictions [arXiv:2209.13535](https://arxiv.org/abs/2209.13535) Amoroso et al.
  - All above solutions introduce either experimental or theoretical uncertainties/problems
- Ai-based elegant solution:
  - Fiducial cuts removed by analytic integration of  $(\cos\theta, \Phi)$  in the full phase space of the decay leptons through the measured Ai coefficients
  - few permille total uncertainties for  $ds/dy$  and negligible theoretical uncertainties



# Full-lepton phase space rapidity cross section

- Exquisite permille level precision in the central region
- Subpercent uncertainties up to  $|y| < 3.6$  thanks to dedicated forward electron calibration
- First comparison to N3LO QCD predictions
- Enables precise and unambiguous PDF interpretation with QCD scale variations now smaller than PDF uncertainties



# Measurement methodology

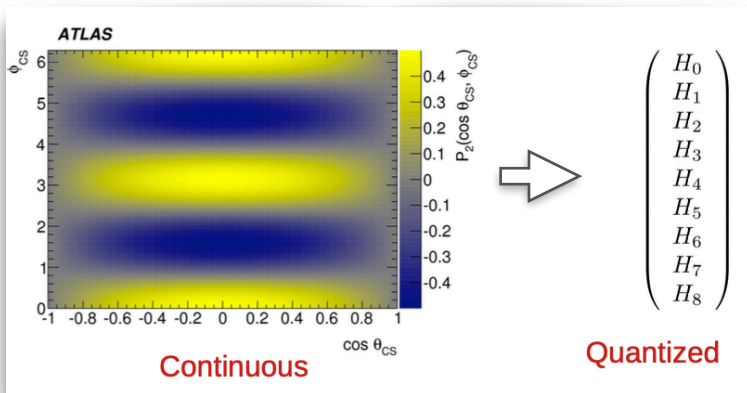
## Expected Yield

Reco  $(p_T^z, y^z, m^z, \cos\theta, \phi)$  bin

$$N_{\text{exp}}^n(A, \sigma, \theta) = \left\{ \sum_{j=1}^{N_{\text{bins}}^{\text{ana}}} \mathcal{L} \sigma_j \left[ t_{8j}^n(\beta) + \sum_{i=0}^7 A_{ij} t_{ij}^n(\beta) \right] \right\} \gamma^n + \sum_B T_B^n(\beta)$$

Truth  $(p_T^z, y^z, m^z)$  bin      Cross section      Angular coefficient      Templated polynomial      Background template

- Likelihood defined in 22528  $(\cos\theta, \Phi, p_T, y)$  bins
- Parameters of interests are the 8  $A_i$  + 1 cross section in  $p_T$ - $y$  bins: 9 parameters in 176 bins

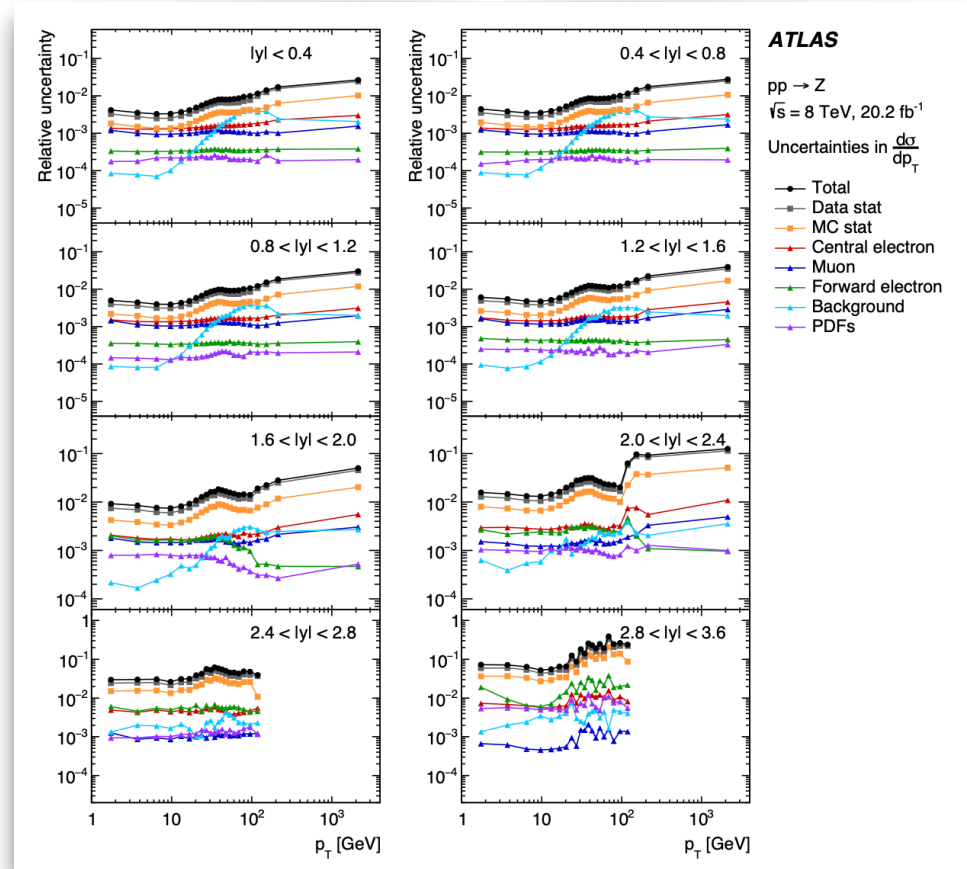
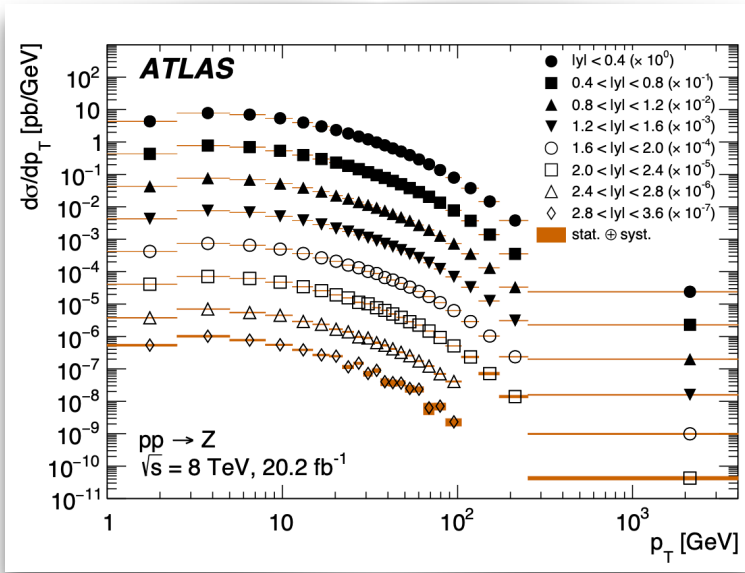


- Measuring the angular coefficients corresponds to building a synthetic “quantized” representation of the  $(\cos\theta, \Phi)$  kinematic space
- Trade systematics for statistics
- Very powerful: avoids theoretical extrapolation of fiducial lepton cuts to full phase space and thereby opens the door to a rich field of precise interpretations



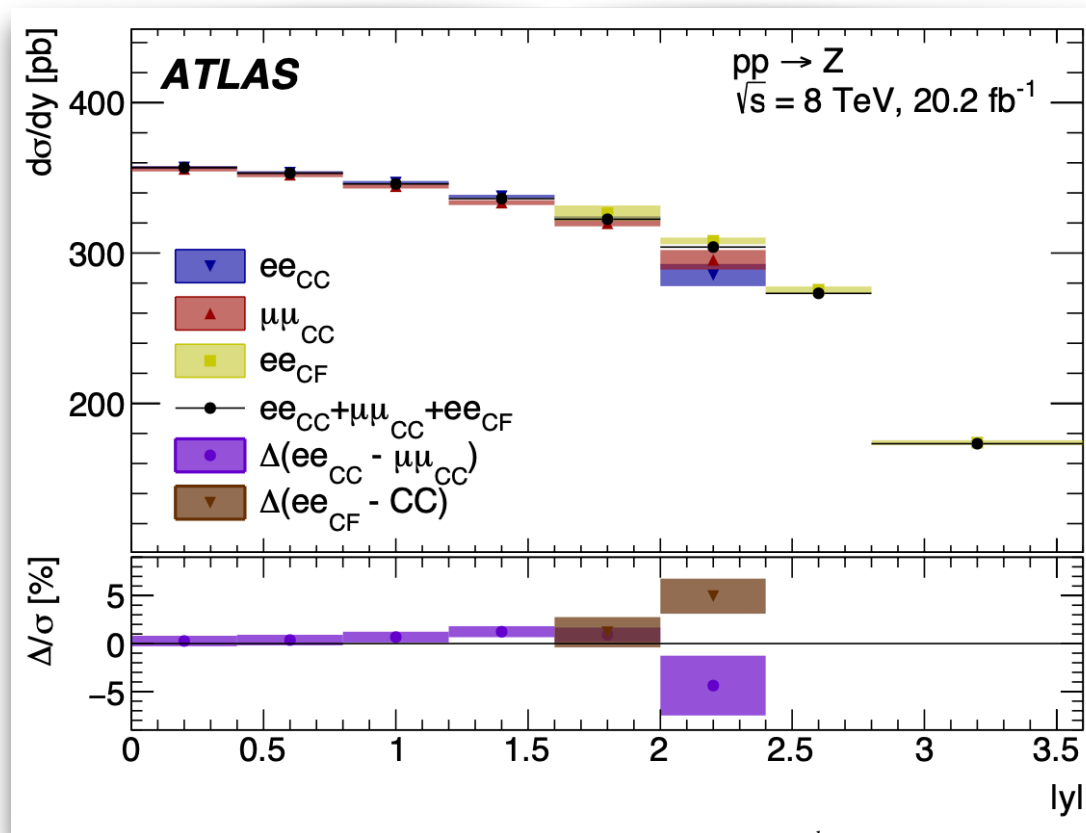
# $d\sigma/dp_T dy$ measurement uncertainties

- First measurement at the LHC of full-lepton phase space cross sections
- Statistically dominated measurement
- Negligible theory uncertainties: cross sections are parameters of the fit, and not the result of an extrapolation



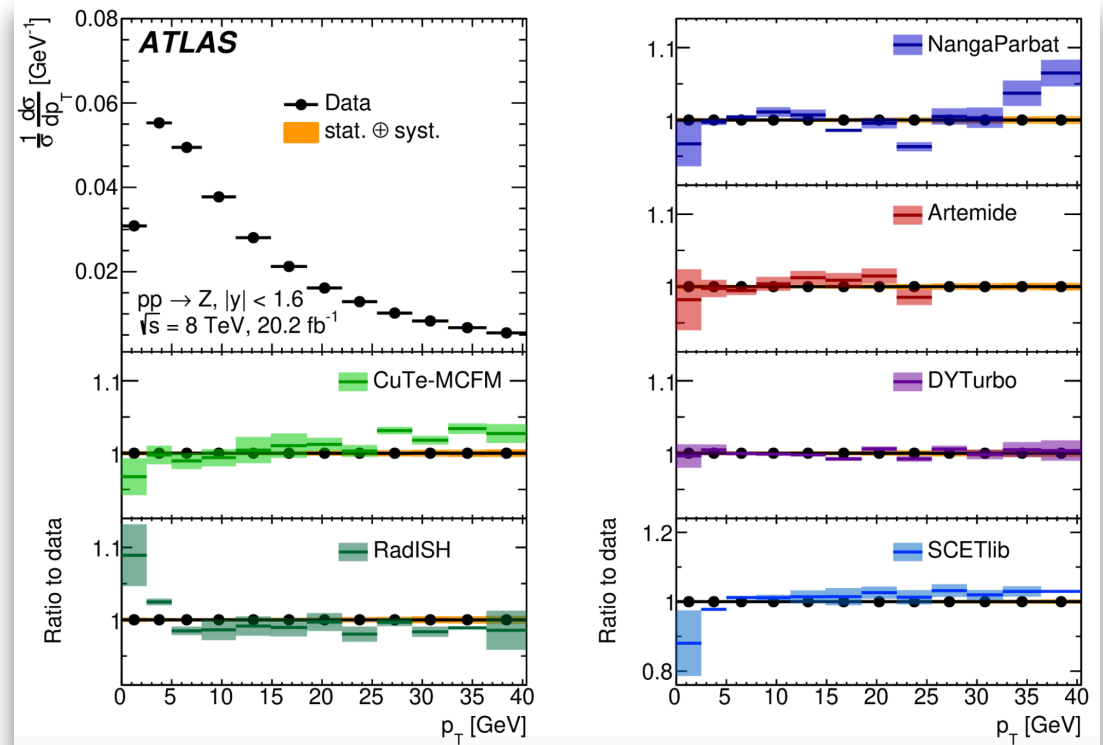
# Channels compatibility

- All three channels ( $ee_{CC}$ ,  $\mu\mu_{CC}$ ,  $ee_{CF}$ ) yield compatible results
  - Important cross-check of detector calibration
  - Forward electrons allow to minimize PDF dependence



# $p_T$ cross section measurement

- Measurement compared to six predictions currently involved in the at N<sup>3</sup>LL/ N<sup>4</sup>LL logarithmic accuracy
  - including O( $a_s^3$ ) matching from MCFM/NNLOJET
- Excellent agreement between data and predictions
  - impressive progress in the understanding of the boson  $p_T$  modelling from the experimental and theoretical points of view





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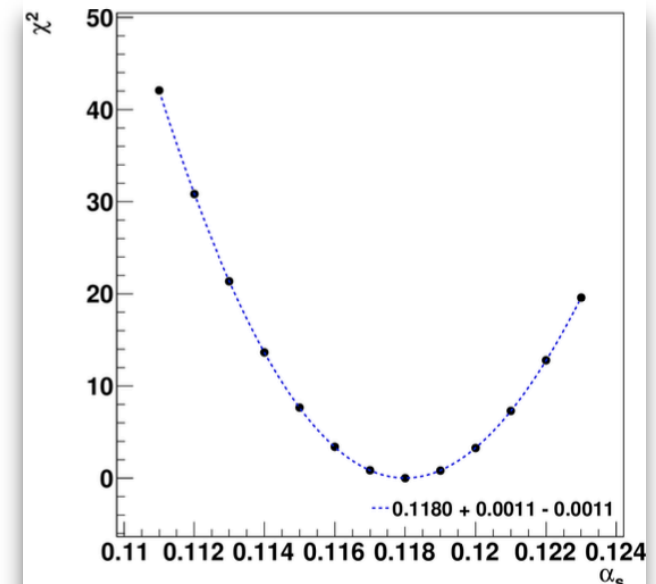


Extraction of the strong  
coupling constant

- DYTurbo interfaced to xFitter arXiv:1410.4412
- Evaluate  $\chi^2(\alpha_s)$  with as variations as provided in LHAPDF
- Include experimental ( $\beta_{j,\text{exp}}$ ) and PDF ( $\beta_{k,\text{th}}$ ) uncertainties in the  $\chi^2$

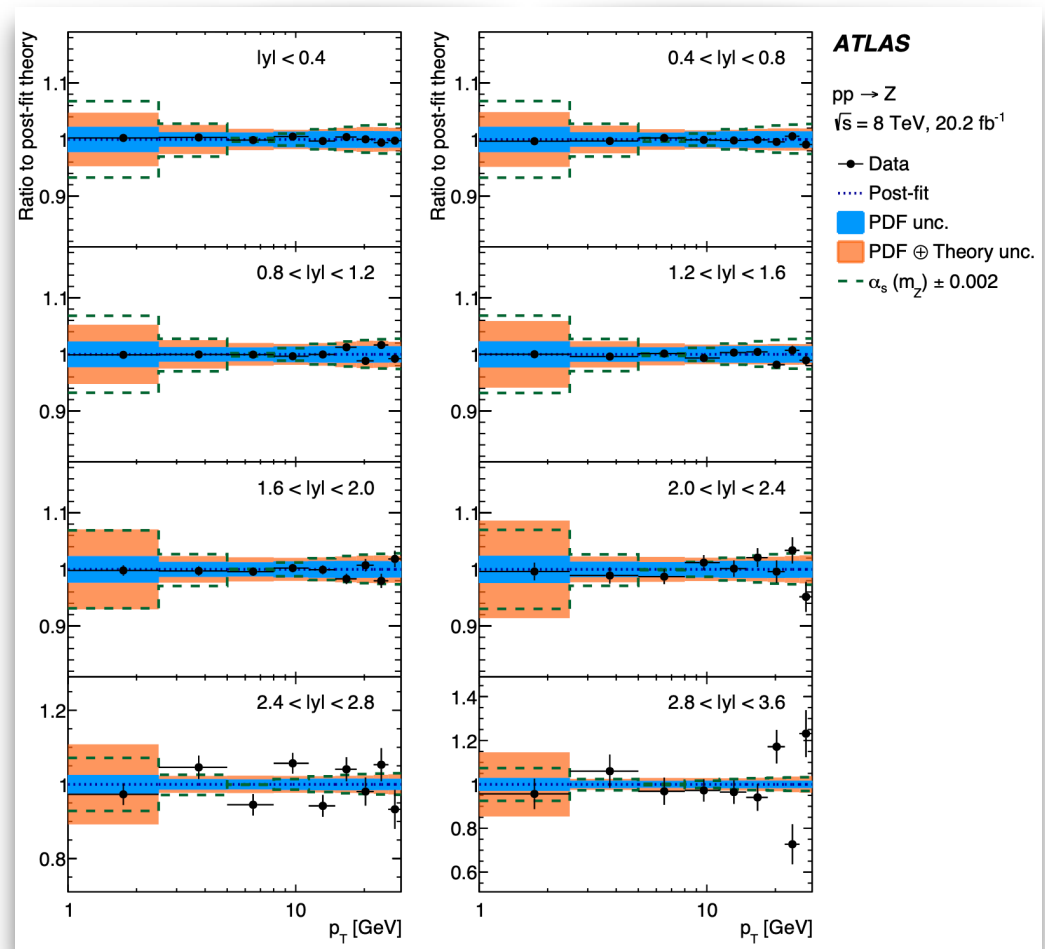
$$\chi^2(\beta_{\text{exp}}, \beta_{\text{th}}) = \sum_{i=1}^{N_{\text{data}}} \frac{\left( \sigma_i^{\text{exp}} + \sum_j \Gamma_{ij}^{\text{exp}} \beta_{j,\text{exp}} - \sigma_i^{\text{th}} - \sum_k \Gamma_{ik}^{\text{th}} \beta_{k,\text{th}} \right)^2}{\Delta_i^2} + \sum_j \beta_{j,\text{exp}}^2 + \sum_k \beta_{k,\text{th}}^2$$

- At each value of  $\alpha_s(m_Z)$  the  $\beta_{k,\text{th}}$  terms explore the PDF space to find the best fit to the  $p_{\text{T}}(Z)$  data
  - $\rightarrow$  equivalent to including the new dataset in the PDF without refitting, using profiling/reweighting Eur.Phys.J.C 75 (2015) 9, 458
- The non-perturbative form factor is added with unconstrained nuisance parameters ( $b = 0$ ) i.e. left free in the fit
- Fit the region of  $p_{\text{T}}(Z) < 29\text{GeV}$



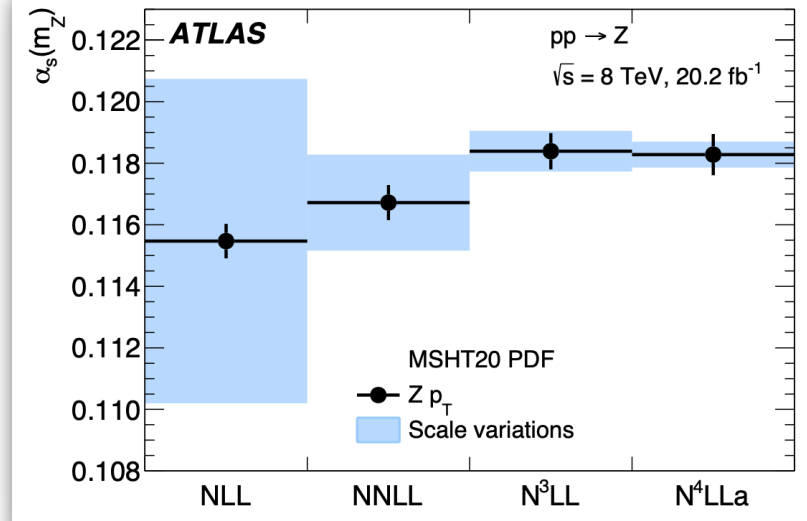
# Determination of $a_s(m_Z)$ from $p_T(Z)$ at 8 TeV

- $a_s(m_Z)$  from a fit to the double-differential  $p_T$ - $YZ$  cross section measured in full-lepton phase space
- Postfit  $\chi^2/\text{dof} = 82/72$
- Determination performed at lower orders, demonstrating good convergence of the perturbative series



# Uncertainties

- Use MSHT20 PDF (only PDF Set which is available at N3LO order)
- Repeat Fit using lower orders (also with MSHT20)
  - $\alpha_S$  at higher orders is always within uncertainties of lower orders
- Scale Variations: independent  $u_R$ ,  $u_F$  and  $Q$  variations
- Variations of the lower fit range (0 GeV  $\rightarrow$  5 GeV) and the upper fit range (29 GeV  $\rightarrow$  22 GeV) of the fit range, are performed to test the stability
  - test for non-perturbative and quark flavor effects



|                               |   |
|-------------------------------|---|
| Experimental uncertainty      | $\pm 0.44$                                      |
| PDF uncertainty               | $\pm 0.51$                                      |
| Scale variation uncertainties | $\pm 0.42$                                      |
| Matching to fixed order       | 0 $-0.08$                                       |
| Non-perturbative model        | $+0.12$ $-0.20$                                 |
| Flavour model                 | $+0.40$ $-0.29$                                 |
| QED ISR                       | $\pm 0.14$                                      |
| $N^4LL$ approximation         | $\pm 0.04$                                      |
| <b>Total</b>                  | <b><math>+0.91</math>    <math>-0.88</math></b> |

# NNLO PDF Sets (1/2)

- At order N4LLa+N3LO, only one N3LO PDF set available: MSHT20aN3LO
  - Study the dependence of the results on the choice of PDF set by fitting one order lower, i.e. N3LLa+N3LO using NNLO PDFs.

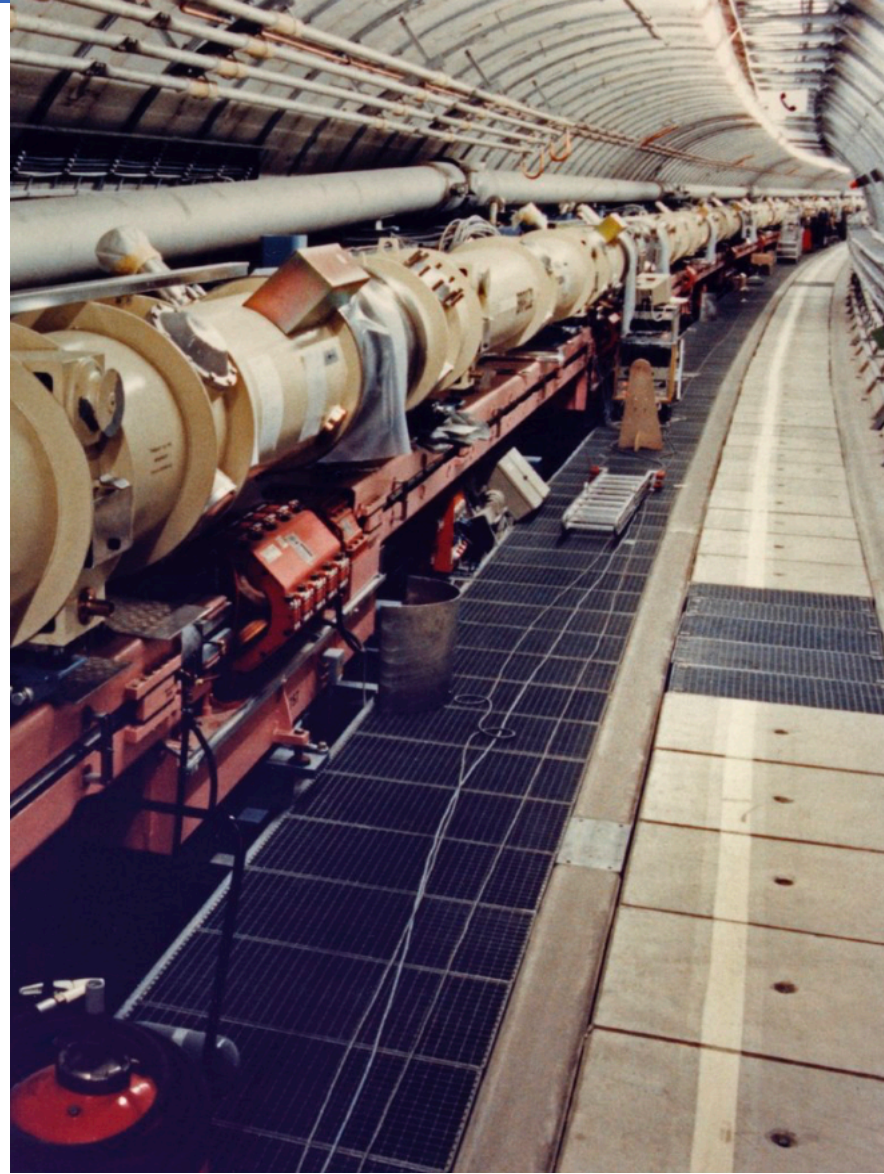
| PDF set        | $\alpha_s(m_Z)$ | PDF uncertainty | $g$ [GeV <sup>2</sup> ] | $q$ [GeV <sup>4</sup> ] | $\chi^2/\text{dof}$ |
|----------------|-----------------|-----------------|-------------------------|-------------------------|---------------------|
| MSHT20 [32]    | 0.11839         | 0.00040         | 0.44                    | -0.07                   | 96.0 /69            |
| NNPDF40 [78]   | 0.11779         | 0.00024         | 0.50                    | -0.08                   | 116.0/69            |
| CT18A [79]     | 0.11982         | 0.00050         | 0.36                    | -0.03                   | 97.7 /69            |
| HERAPDF20 [63] | 0.11890         | 0.00027         | 0.40                    | -0.04                   | 132.3/69            |

- At this order, the spread observed is  $\pm 0.00102$ 
  - driven by the difference between the NNPDF4.0 and CT18A PDF sets
- What causes this difference?



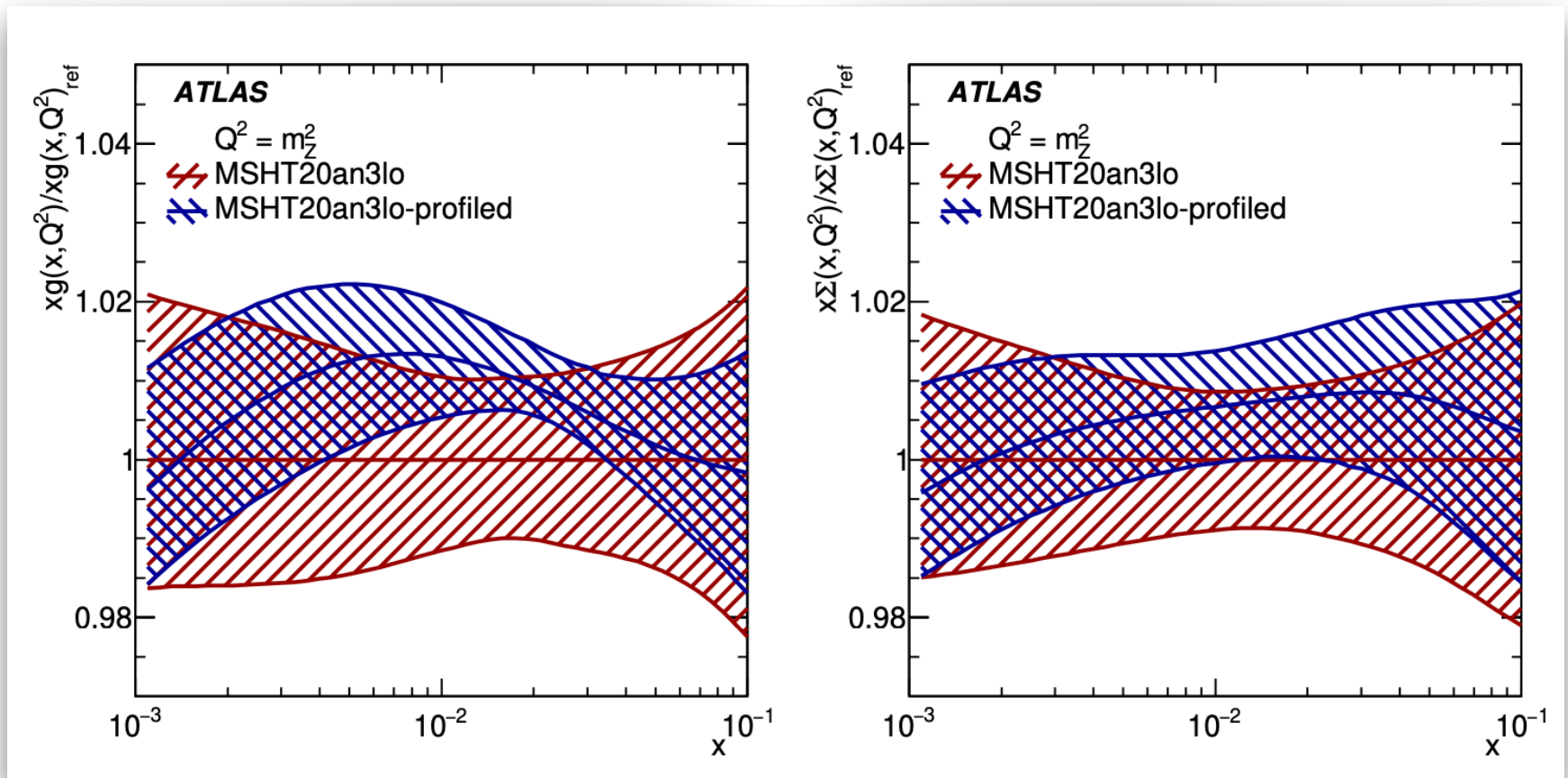
# NNLO PDF Sets (2/2)

- $p_T(Z)$  very sensitive to gluon PDF
  - PDF determinations at NNLO are affected by significant tension between low- $x$  (from HERA) and high- $x$  gluon PDFs (Hadron Collider Jets)
- Effect of this tension has been studied at N<sup>3</sup>LL
  - Refit when including DIS cross-section data from HERA with  $Q^2 > 10 \text{ GeV}^2$  plus  $p_T(Z)$ 
    - → Double Count HERA to reduce impact of other data-sets. CT18A is shifted downwards and the spread is reduced to  $\pm 0.00016$ .
- Using MSHT20an<sup>3</sup>lo largely removes the tension in the gluon PDF
  - indicated by the significant improvement in the  $\chi^2$  of the  $p_T(Z)$  and DIS Data.
    - → indication that the spread of PDFs at NNLO is not representative of the true PDF uncertainty on N<sup>3</sup>LO.



# PDF profiling

- PDF profiling at the best  $\alpha_s(m_Z)$  shows reduction of gluon and sea quark PDF uncertainties



# Non perturbative QCD model

- NP model is generally determined from the data, parameters values depend on the chosen prescription to avoid the Landau pole in b-space

$$S_{NP}(b) = \exp \left[ -g_j(b) - g_K(b) \log \frac{m_{\ell\ell}^2}{Q_0} \right]$$

$$g_j(b) = \frac{g b^2}{\sqrt{1 + \lambda b^2}} + \text{sign}(q) \left( 1 - \exp[-|q| b^4] \right)$$

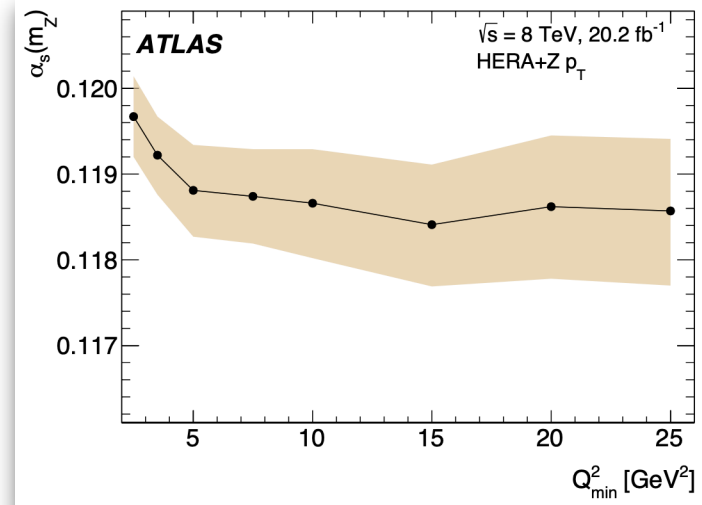
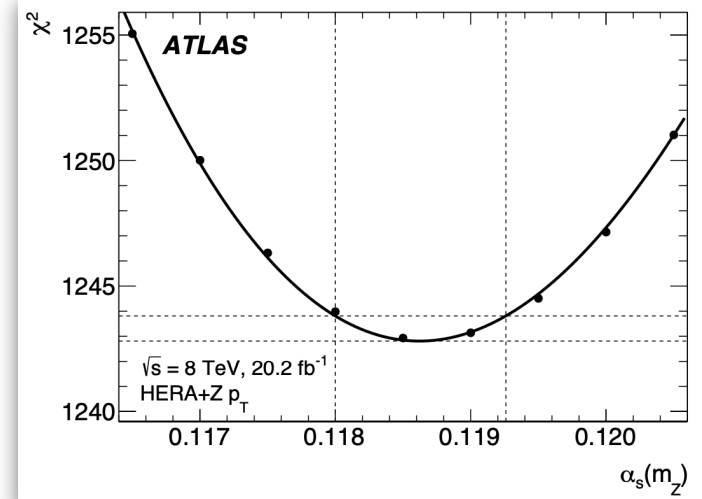
$$g_K(b) = g_0 \left( 1 - \exp \left[ -\frac{C_F \alpha_s(b_0/b_*) b^2}{\pi g_0 b_{lim}^2} \right] \right)$$

$$b_* = \frac{b}{1 + b^2/b_{lim}^2}$$

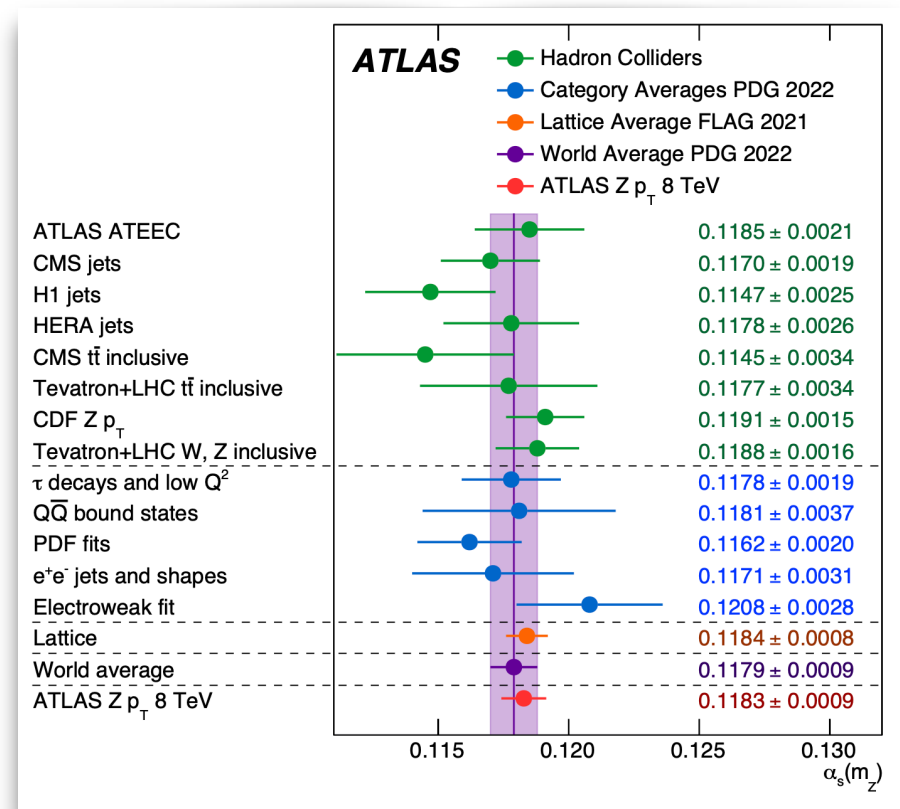
- $g_j$  functions include a quadratic/quartic term:  $g$  and  $q$  free parameters of the fit
  - The theory should not depend on  $b_{lim}$  (freezing scale) and  $Q_0$  (starting scale), provided SNP is flexible enough.  $Q_0$  and  $b_{lim}$  estimated as parameterisation unc.
- $g_0$  controls the very high  $b$  (very small  $p_T$ ) behaviour, should be fitted to data, but we have no sensitivity to it, so it is varied
- Lambda controls transition from Gaussian to exponential: varied between 0.5-2
  - Fits excluding 0-5 GeV yields  $\alpha_s(m_Z)$  with a spread of + 0.00017 – 0.00010
  - Fit uncertainty increased from 0.00067 to 0.00071
  - Correlation between  $\alpha_s(m_Z)$  and  $g$  largely reduced

# More checks

- Simultaneous determination of  $\alpha_s$ , the PDFs, and the non-perturbative parameters
  - N3LL+N3LO, with PDFs evolved at NNLO.
- The light-quark coefficient functions of the DIS cross sections are calculated in the  $\overline{MS}$  scheme.
- The heavy quarks ( $c$ ,  $b$ ) generated dynamically
  - using general-mass variable-flavour-number scheme, with up to five active quark flavours.
- Fits performed at fixed values of  $\alpha_s$  via a quadratic interpolation of the  $\chi^2$  function
  - $0.11866 \pm 0.00064$
- The dependence of  $\alpha_s$  on the minimum squared four-momentum transfer  $Q^2$  of the HERA data is studied in the range from 2.5 GeV to 25 GeV
  - No sign. dependence is observed for  $>5$  GeV.



- $\alpha_s(m_Z) = 0.11828 \pm 0.00084 \mp 0.00088$
- Most precise experimental determination of  $\alpha_s(m_Z)$ , as precise as the PDG and Lattice world averages
- First  $\alpha_s(m_Z)$  determination at N3LO+N4LL
- Clean experimental signature (leptons) with highest exp sensitivity
- Determination focusing on the Sudakov region (usually avoided to determine  $\alpha_s$ )
- Observable not suitable for inclusion in PDF fits  $\rightarrow$  no correlation with  $\alpha_s(m_Z)$  determinations from PDF fits





# Summary

- New window for the determination of the strong coupling using the transverse momentum of  $Z$  bosons
- New measurements might reduce further PDF uncertainties
- New measurements required to constrain further non-perturbative effects



# Orders

|            | Virtual            |      | Sudakov |                |        | Real         |
|------------|--------------------|------|---------|----------------|--------|--------------|
|            | H[ $\delta(1-z)$ ] | H[z] | Cusp AD | Collinear, RAD | PDF    | CT,V+jet     |
| LL+LO      | 1                  | 1    | 1-loop  | 0              | const. | 1            |
| NLL+NLO    | $\alpha_s$         | C1   | 2-loop  | 1-loop         | LO     | $\alpha_s$   |
| NLL*+NLO   | $\alpha_s$         | C1   | 2-loop  | 1-loop         | NLO    | $\alpha_s$   |
| NNLL+NNLO  | $\alpha_s^2$       | C2   | 3-loop  | 2-loop         | NLO    | $\alpha_s^2$ |
| N3LL+N3LO  | $\alpha_s^3$       | C3   | 4-loop  | 3-loop         | NNLO   | $\alpha_s^3$ |
| N4LLa+N3LO | $\alpha_s^4$       | C4   | 5-loop  | 4-loop         | N3LO   | $\alpha_s^4$ |

Known analytically

Approximated numerically

Unknown, estimated with series acceleration

Not included



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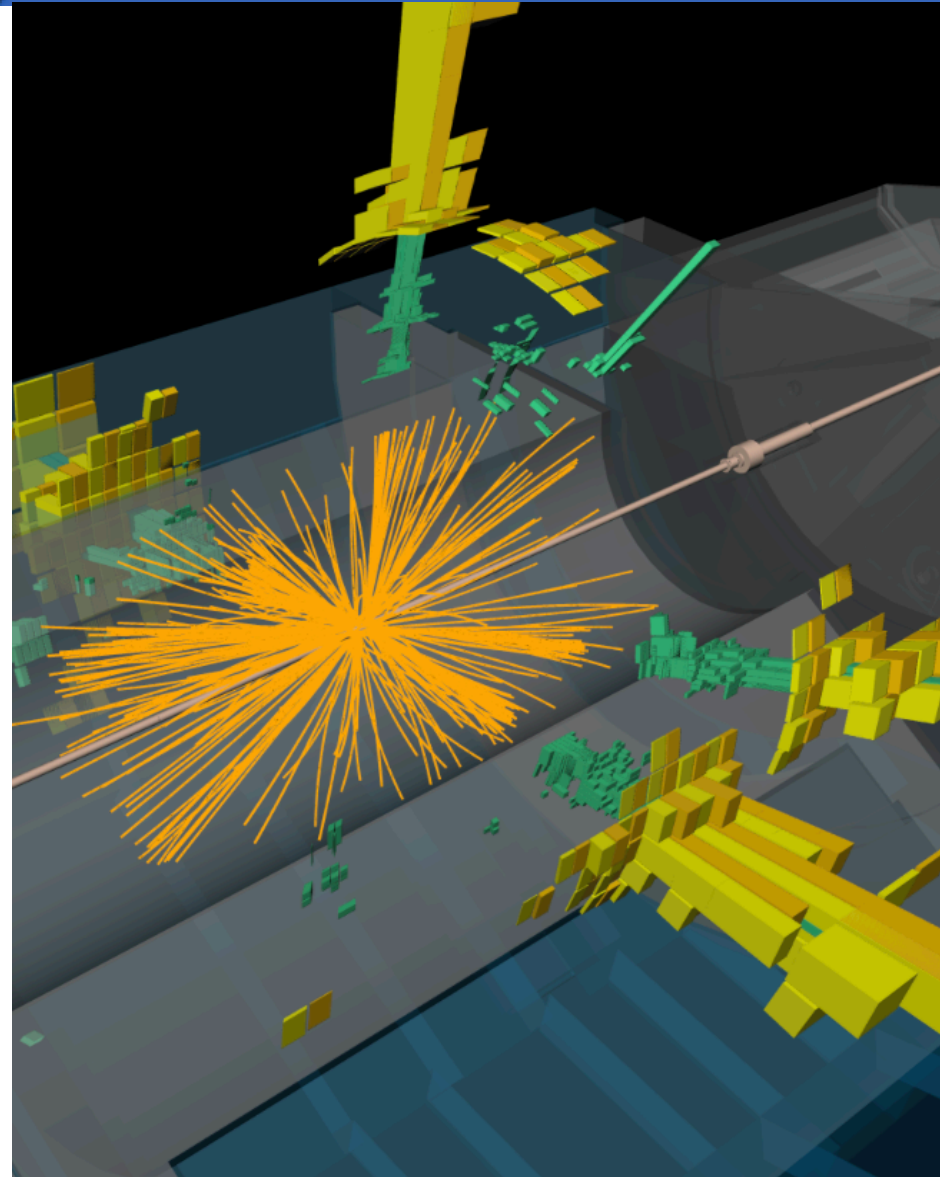
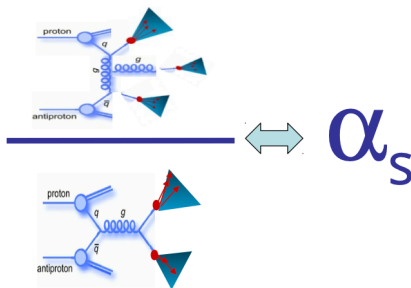


# Strong Coupling in Jet Correlations



# Jet Measurements

- Use Transverse energy-energy correlations (TEEC) and Associated asymmetry (ATEEC)
- TEEC: defined as the transverse-energy-weighted distribution of the azimuthal differences between jet pairs in the final state
- ATEEC: defined as the difference between the forward ( $\cos\phi > 0$ ) and the backward ( $\cos\phi < 0$ ) part of TEEC
- TEEC and ATEEC functions are sensitive to gluon radiation and show a clear dependence on the strong coupling

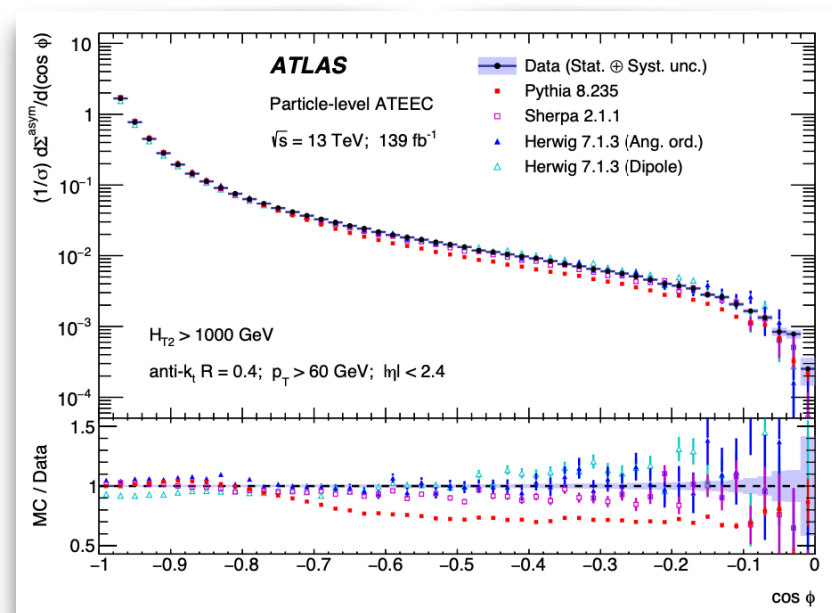
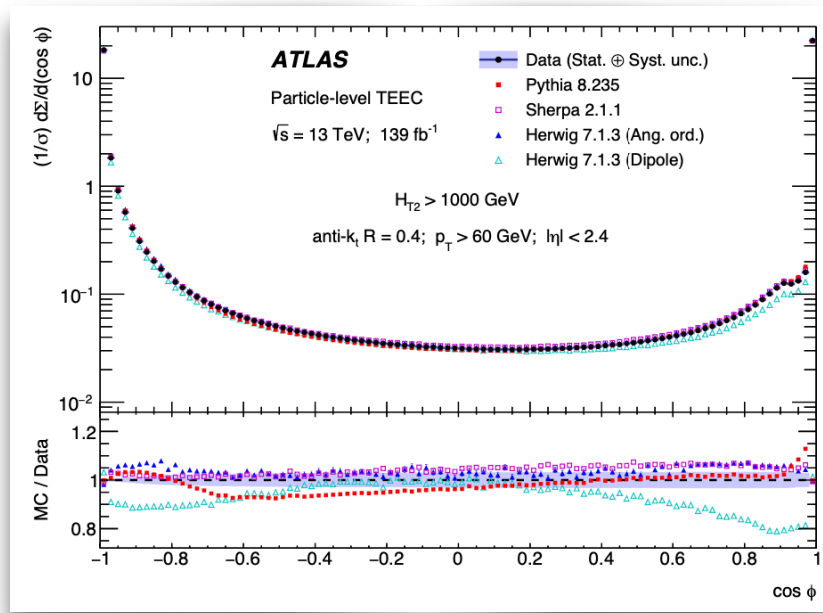


# Event Selection and Results

ATLAS Collab. Eur.Phys.J. C77 (2017) 12, 872

- Selection: Multijet events
  - Leading Jet  $p_T > 460$  GeV
  - Leading Two Jets with  $p_{T1} + p_{T2} > 1$  TeV
  - All Jet  $p_T$ 's  $> 60$  GeV
  - 57.6M events at 8 TeV

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# Extraction of $\alpha_s(m_Z)$

- Determination of  $\alpha_s$  with a  $\chi^2$  fit
  - Predictions based in NNLO
  - Dominant uncertainties from PDFs

$$\alpha_s(m_Z) = 0.1175 \pm 0.0006 \text{ (exp.)}_{-0.0017}^{+0.0034} \text{ (theo.)}$$

$$\alpha_s(m_Z) = 0.1185 \pm 0.0009 \text{ (exp.)}_{-0.0012}^{+0.0025} \text{ (theo.)}$$

