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> Workshop on Tensions in Cosmology September 6-13 2023. Corfu, Greece

http://celrivera.wix.com/cosmology | celia.escamilla@nucleares.unam.mx

Setting the scene:

Cosmology Intertwined: the state-of-art

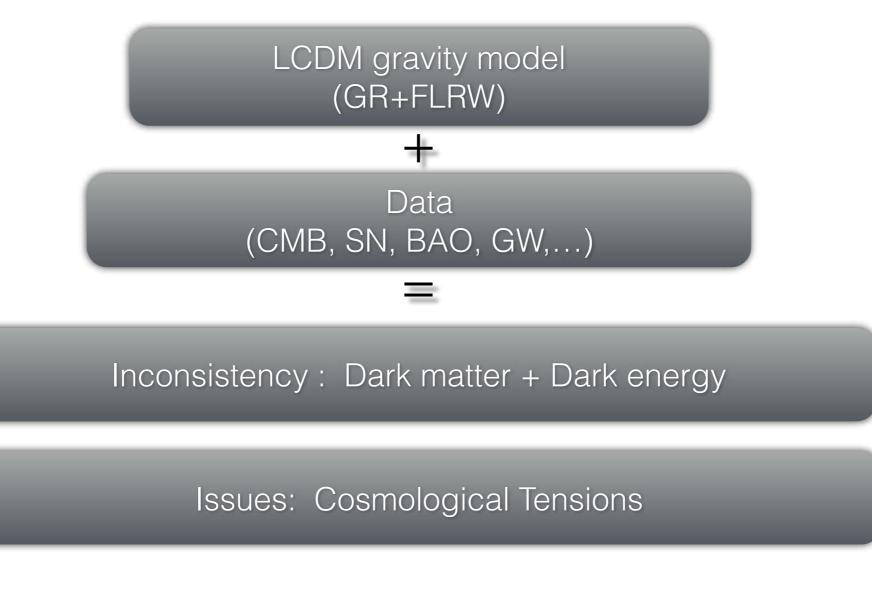
The precision problem:

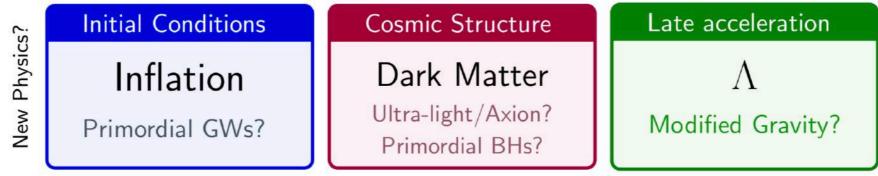
Beyond the standard cosmology

# The work:

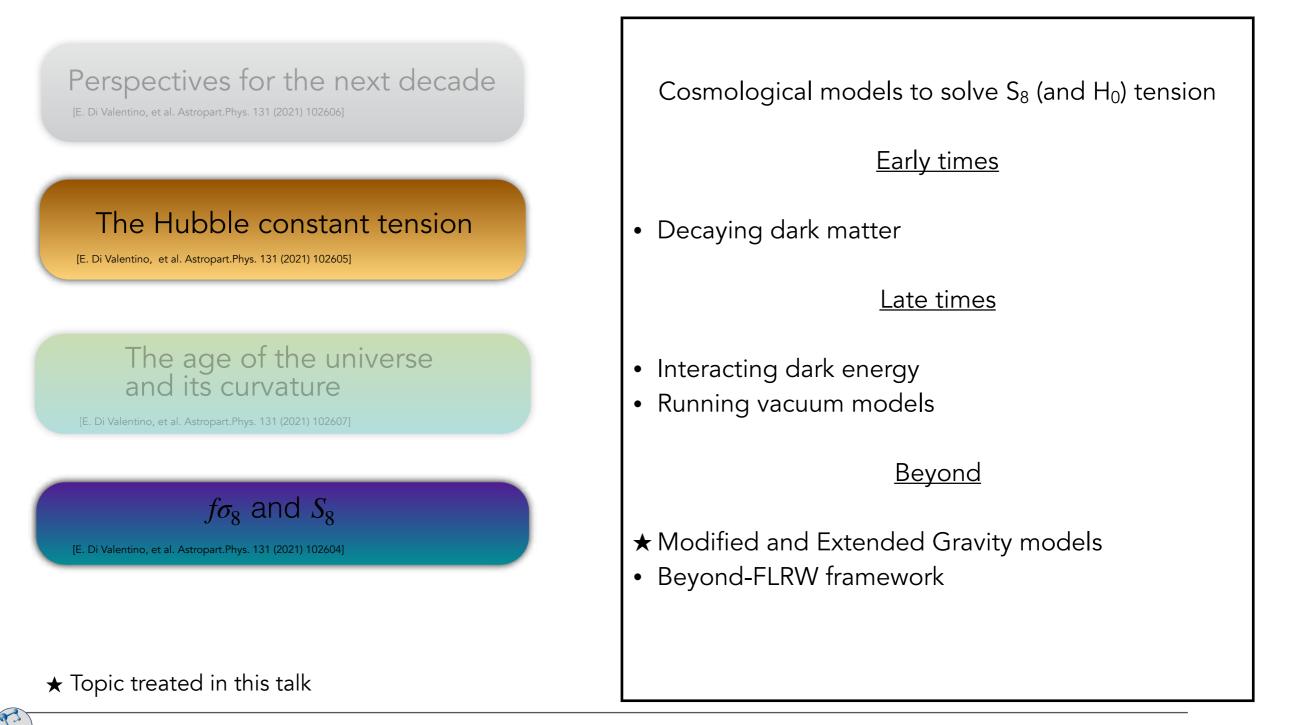
Beyond Einstein's gravity and its constraints

Outline





# Cosmology Intertwined: Snowmass Collaboration



# Cosmology Intertwined: Snowmass Collaboration and Cosmoverse

#### Perspectives for the next decade

[E. Di Valentino, et al. Astropart.Phys. 131 (2021) 102606]

#### The Hubble constant tension

[E. Di Valentino, et al. Astropart.Phys. 131 (2021) 102605]

# The age of the universe and its curvature

[E. Di Valentino, et al. Astropart.Phys. 131 (2021) 102607]



[E. Di Valentino, et al. Astropart.Phys. 131 (2021) 102604]

#### $\star$ Topic treated in this talk

- ★ Improve our understanding of systematic uncertainties
- ★ Maximize the amount of information that can be extracted from the data
- Improve our understanding of the physics on non-linear scales
- ★ De-standardize some of the LCDM assumptions.



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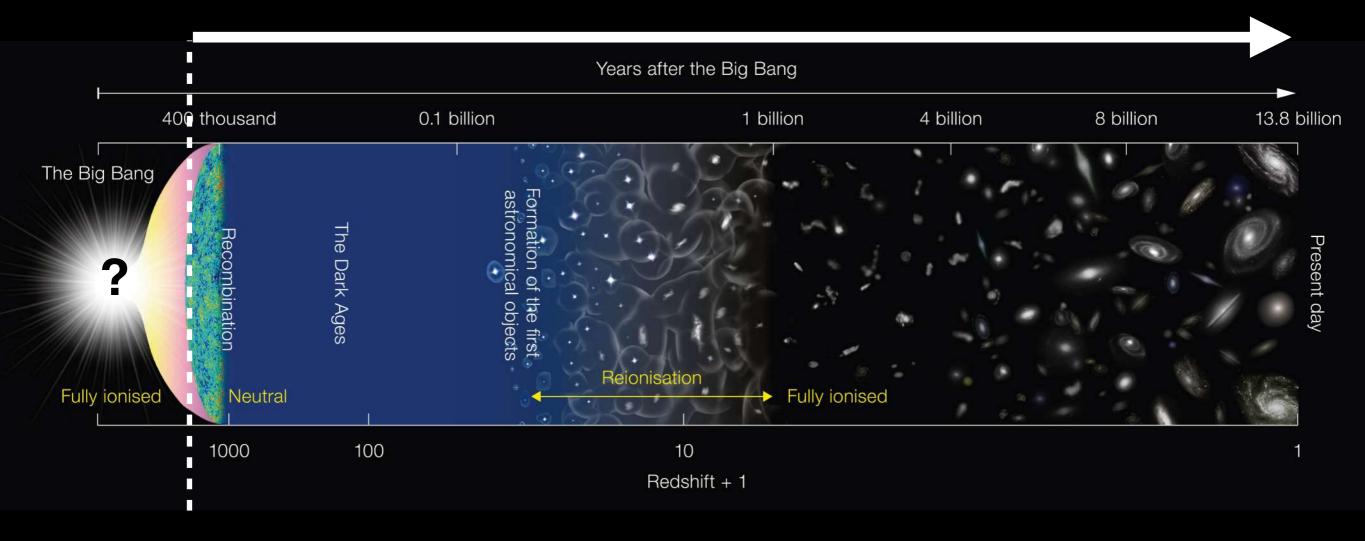
Two paths to infer/measure H<sub>0...</sub>

## The work:

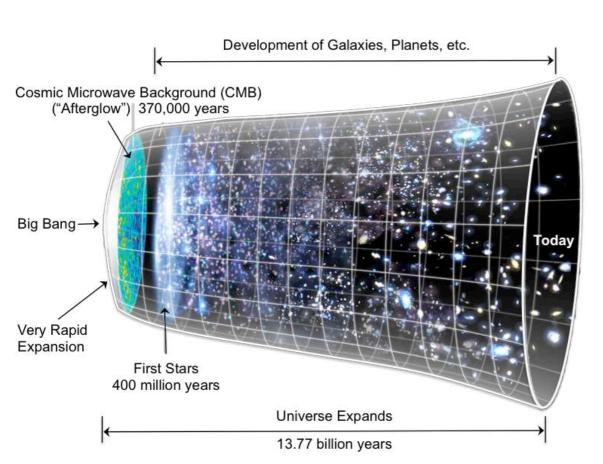
Beyond Einstein's gravity and its constraints



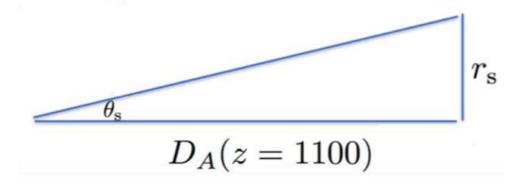
# From early cosmic times until today







#### The CMB as a (self-calibrated) standard ruler



Cosmic Microwave Background Radiation (CMB)

 $d_{\mathcal{A}}(z) = \frac{1}{1+z} \frac{1}{H_0 \sqrt{\Omega_{\mathcal{K}}}} \sinh \left[ H_0 \sqrt{\Omega_{\mathcal{K}}} \int_0^z \frac{dz'}{H(z')} \right]$ 



Nobel Prize in Physics 2019

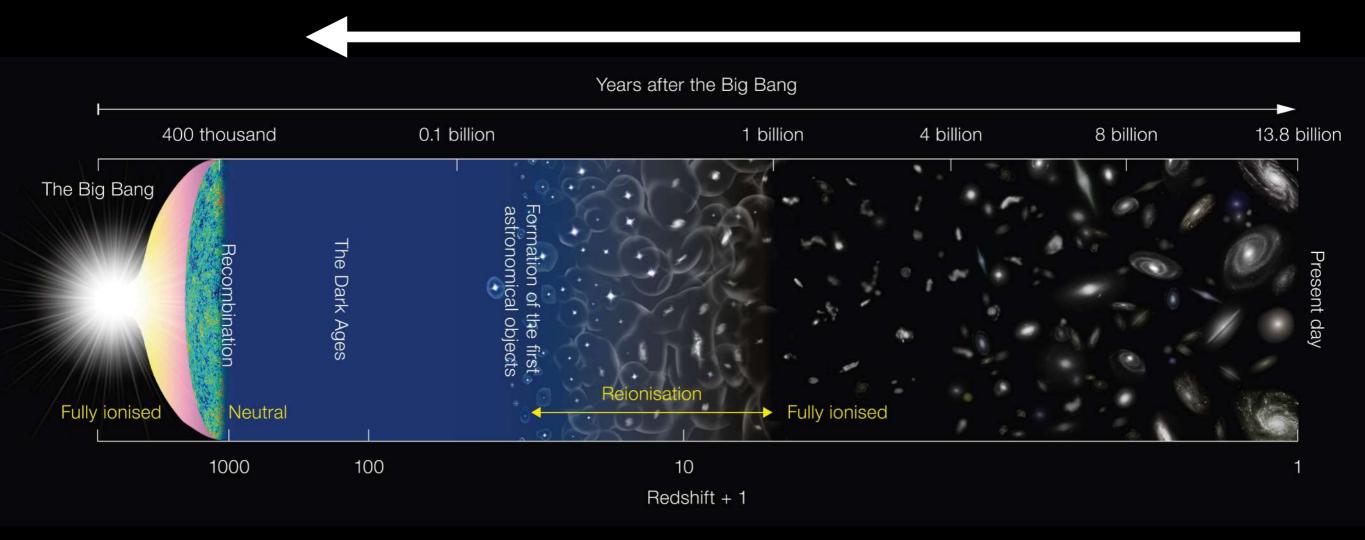
"for theoretical discoveries in cosmology" and "for the discovery of an exoplanet orbiting a solar-type star".



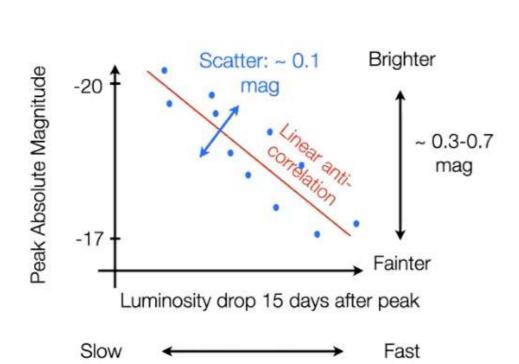
Media. James Peebles Prize share: 1/2 III. Niklas Elmehed. © Nob Media. Michel Mayor Prize share: 1/4

III. Niklas Elmehed. © Nobel Media. Didier Queloz Prize share: 1/4

# From today up to early cosmic times

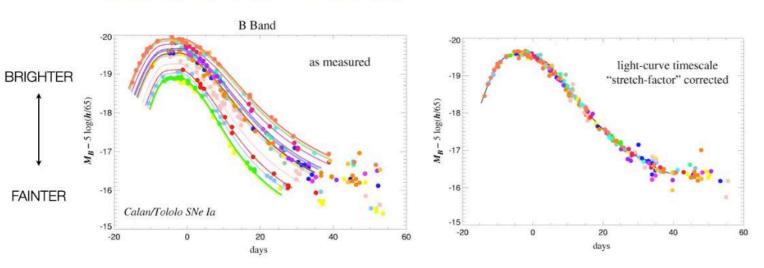


E



Supernovae Type la

# BEFORE CORRECTION AFTER CORRECTION



 $d_L(z) = (1+z) \frac{1}{H_0 \sqrt{\Omega_K}} \sinh \left[ H_0 \sqrt{\Omega_K} \int_0^z \frac{dz'}{H(z')} \right]$ 



Nobel Prize in Physics 2011

"for the discovery of the accelerated expansion of the Universe through observations of distant supernovae"



Saul Perlmutter

Prize share: 1/2





© The Nobel Foundation. Photo: U Montan Adam G. Riess Prize share: 1/4

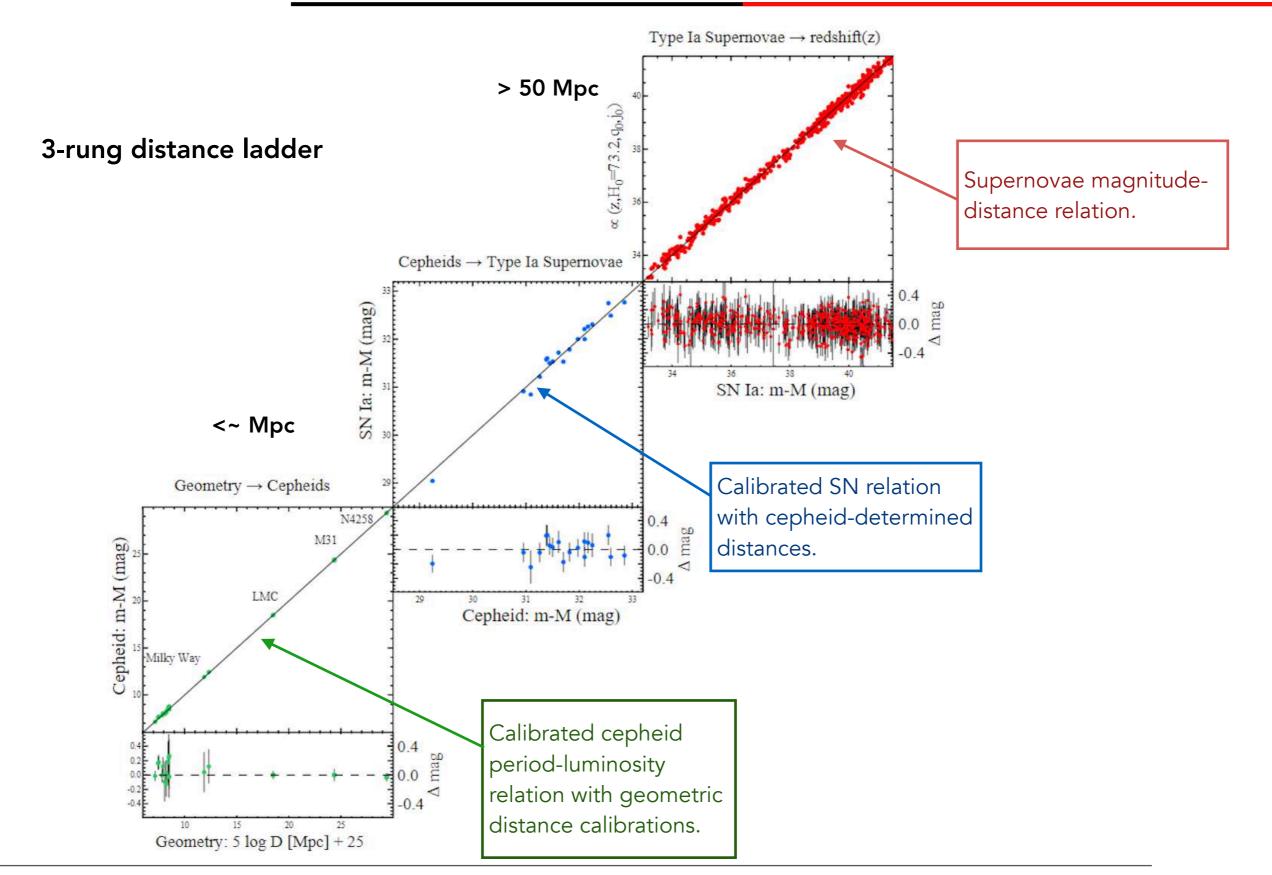
 $\mu = m_B - M - \alpha x_1 - \beta c$ 

Brian P. Schmidt

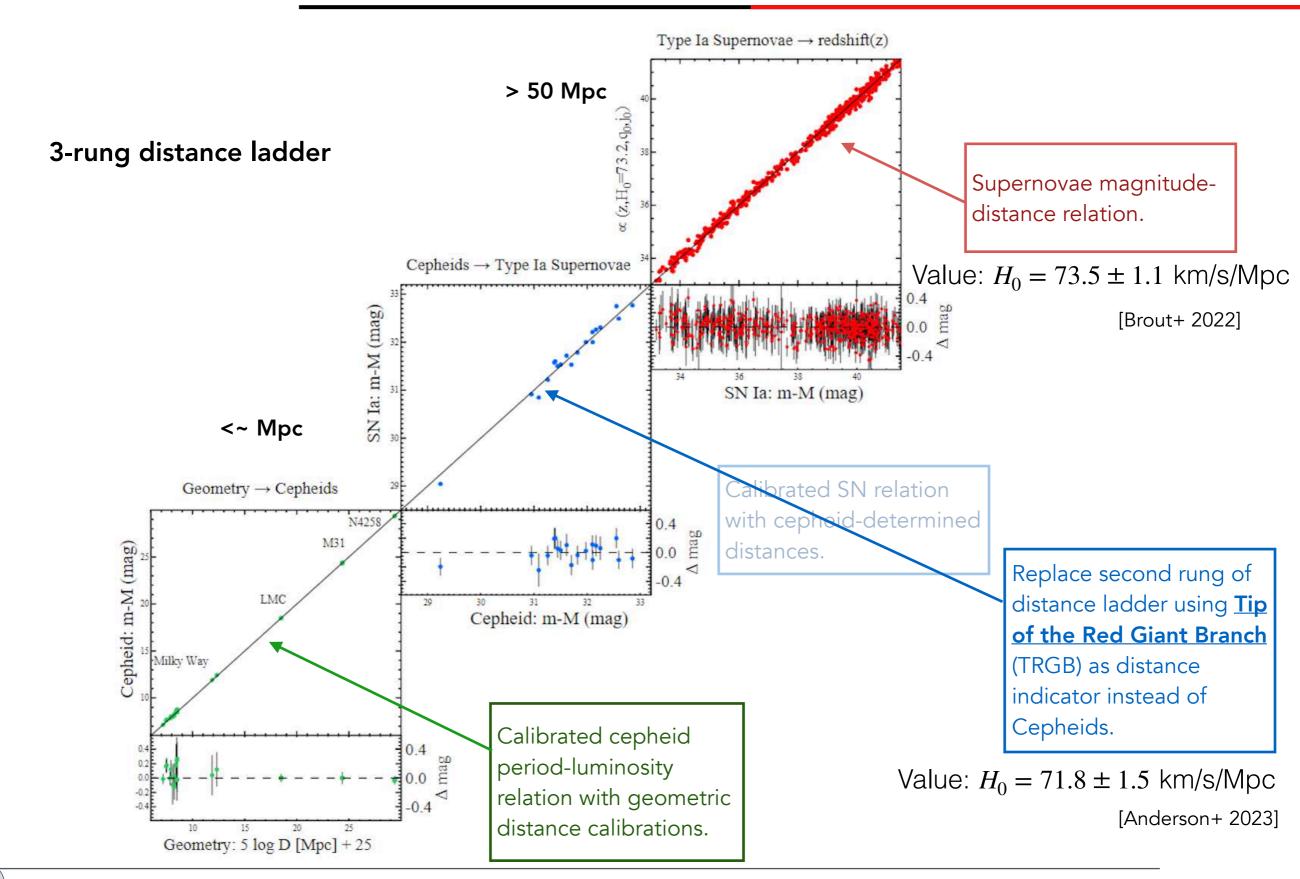
Prize share: 1/4

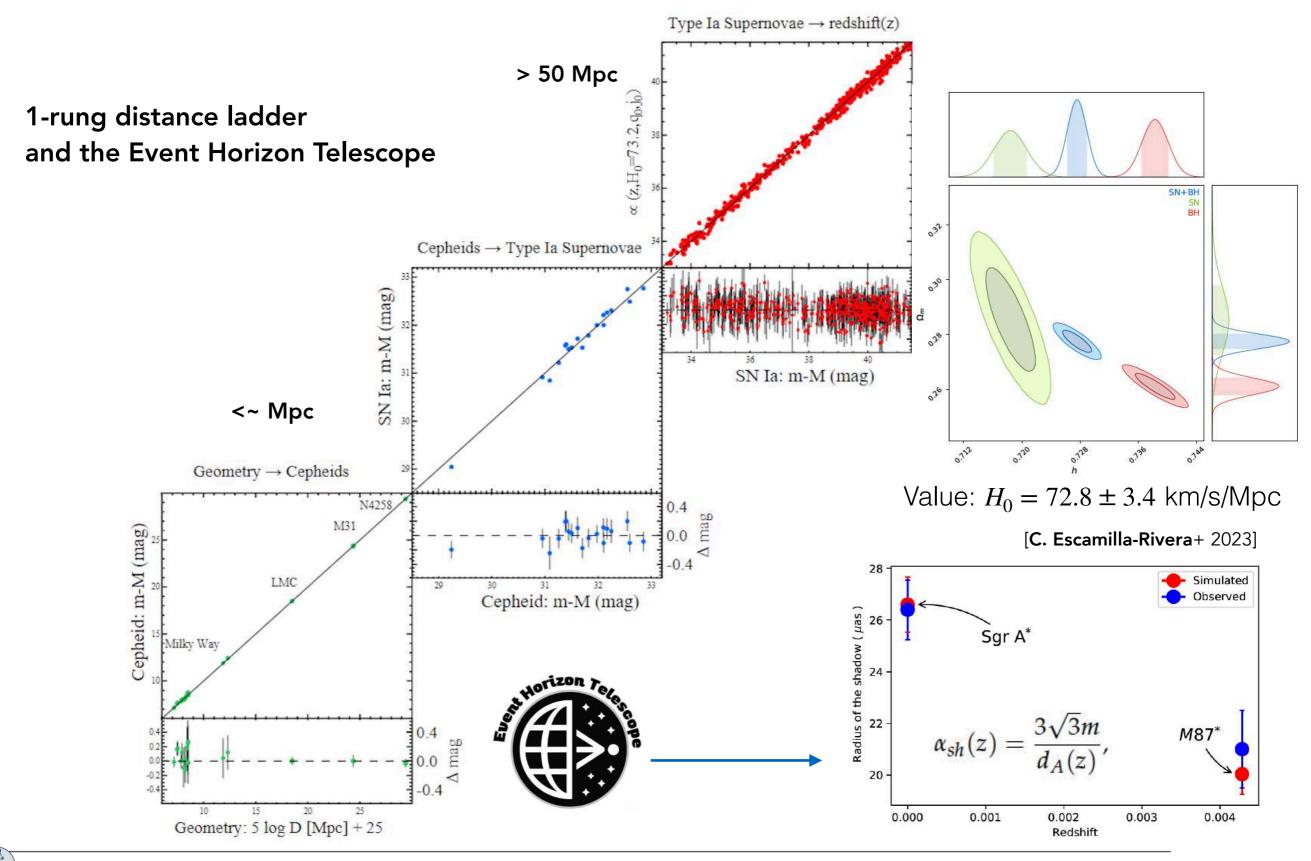
[Philips, ApJ 413 (1993)]











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# The precision problem:

Beyond the standard cosmology

Two paths to infer/measure H<sub>0</sub>... we have work to do!

### The work:

Beyond Einstein's gravity and its constraints

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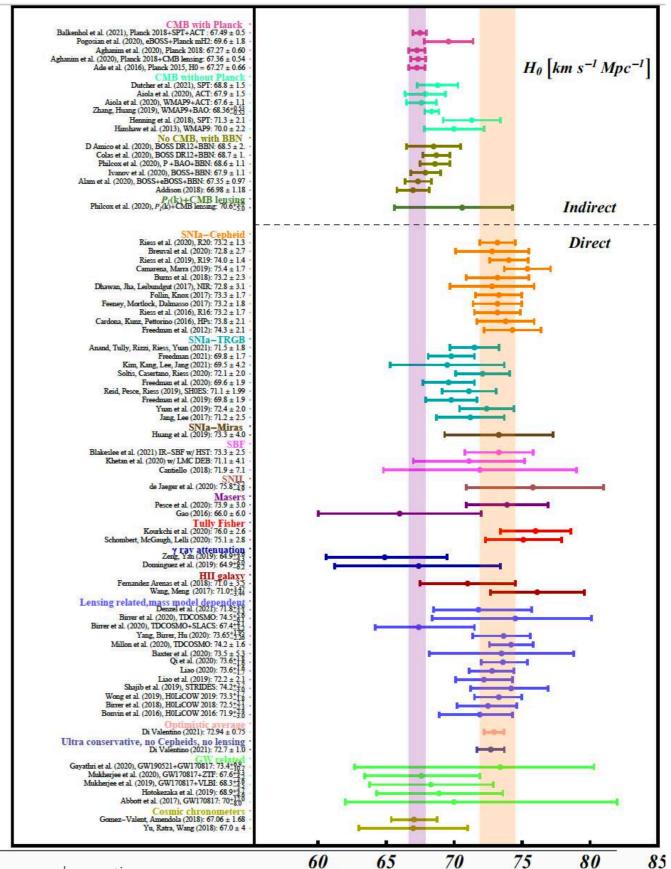
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#### Beyond Einstein's gravity and its constraints



#### Possible solutions of $H_0$

[Cosmology Intertwined: A Review of the Particle Physics, Astrophysics, and Cosmology Associated with the Cosmological Tensions and Anomalies. J. High En. Astrophys.(2022)] Lovelock's theorem (1971) [version from Clifton+ Phys.Rept. 513 (2012)]

"The only second-order, local gravitational field equations derivable from an action containing solely the 4D metric tensor (plus related tensors) are the Einstein field equations with a cosmological constant"

$$S_{grav} = \frac{M_{\rm Pl}^2}{2} \int \sqrt{-g} d^4 x[R]$$

$$S_{grav} = \frac{M_{\text{Pl}}^2}{2} \int \sqrt{-g} d^4 x \Big[ \phi R - \frac{\omega(\phi)}{\phi} (\nabla \phi^2) - 2V(\phi) \Big]$$

$$S_{grav} = \frac{M_{\text{D}}^2}{2} \int \sqrt{-\gamma} d^D x [\mathcal{R} + \alpha \mathcal{G}]$$

$$S_{grav} = \frac{M_{\text{Pl}}^2}{2} \int \sqrt{-g} d^4 x [R + \beta_1 R \nabla_\mu \nabla^\mu R + \beta_2 \nabla_\mu R_{\beta\gamma} \nabla^\mu R^{\beta\gamma}]$$

$$S_{grav} = \frac{M_{\text{Pl}}^2}{2} \int \sqrt{-g} d^4 x \Big[ R + f \Big( \frac{1}{\Box} R \Big) \Big]$$

$$S_{grav} = ?!$$

#### Hubble tension from the Lovelock's theorem break

"The only second-order, local gravitational field equations derivable from an action containing solely the 4D metric tensor (plus related tensors) are the Einstein field equations with a cosmological constant"

$$S_{grav} = \frac{M_{\text{Pl}}^2}{2} \int \sqrt{-g} d^4 x \Big[ \phi R - \frac{\omega(\phi)}{\phi} (\nabla \phi^2) - 2V(\phi) \Big]$$
  

$$S_{grav} = \frac{M_{\text{D}}^2}{2} \int \sqrt{-\gamma} d^D x [\mathcal{R} + \alpha \mathcal{G}]$$
  

$$S_{grav} = \frac{M_{\text{Pl}}^2}{2} \int \sqrt{-g} d^4 x [R + \beta_1 R \nabla_\mu \nabla^\mu R + \beta_2 \nabla_\mu R_{\beta\gamma} \nabla^\mu R^{\beta\gamma}]$$
  

$$S_{grav} = \frac{M_{\text{Pl}}^2}{2} \int \sqrt{-g} d^4 x \Big[ R + f\Big(\frac{1}{\Box}R\Big) \Big]$$

Value:  $H_0 = 69.9^{+0.84}_{-0.86}$  km/s/Mpc [M. Gonzalez et al. JCAP 10 (2021) 028]

Value:  $H_0 = 68.8 \pm 0.9$  km/s/Mpc [D. Wang and D. Mota. Phys.Dark Univ. 32 (2021)]

Value:  $H_0 = 69.22^{+0.66}_{-0.73}$  km/s/Mpc [S. Odintsov et al. Nucl.Phys.B 966 (2021)]

Value:  $H_0 = 68.74^{+0.59}_{-0.51}$  km/s/Mpc [E. Belgacem et al. JCAP 04 010 (2020)]

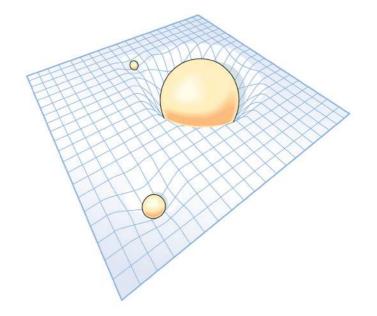
Value:  $H_0 = 69.8 \pm 2.0 \text{ km/s/Mpc}$ 

[C. Escamilla-Rivera and J. Fabris Universe 2021]

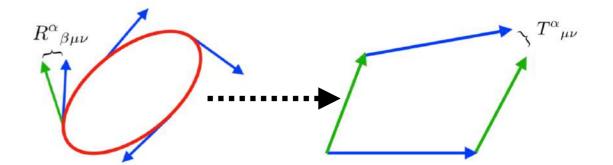
 $S_{grav} = ?!$ 

#### Rethinking the connection...

Space-time tells matter how to move; matter tells space-time how to curve. — J. Wheeler



Curvature is a property of the connection not of the space-time



If a manifold is differentiable ..... We can define **tetrads** (gives basis for vector on the tangent space)

Two "completely" equivalent ways of understading gravity:

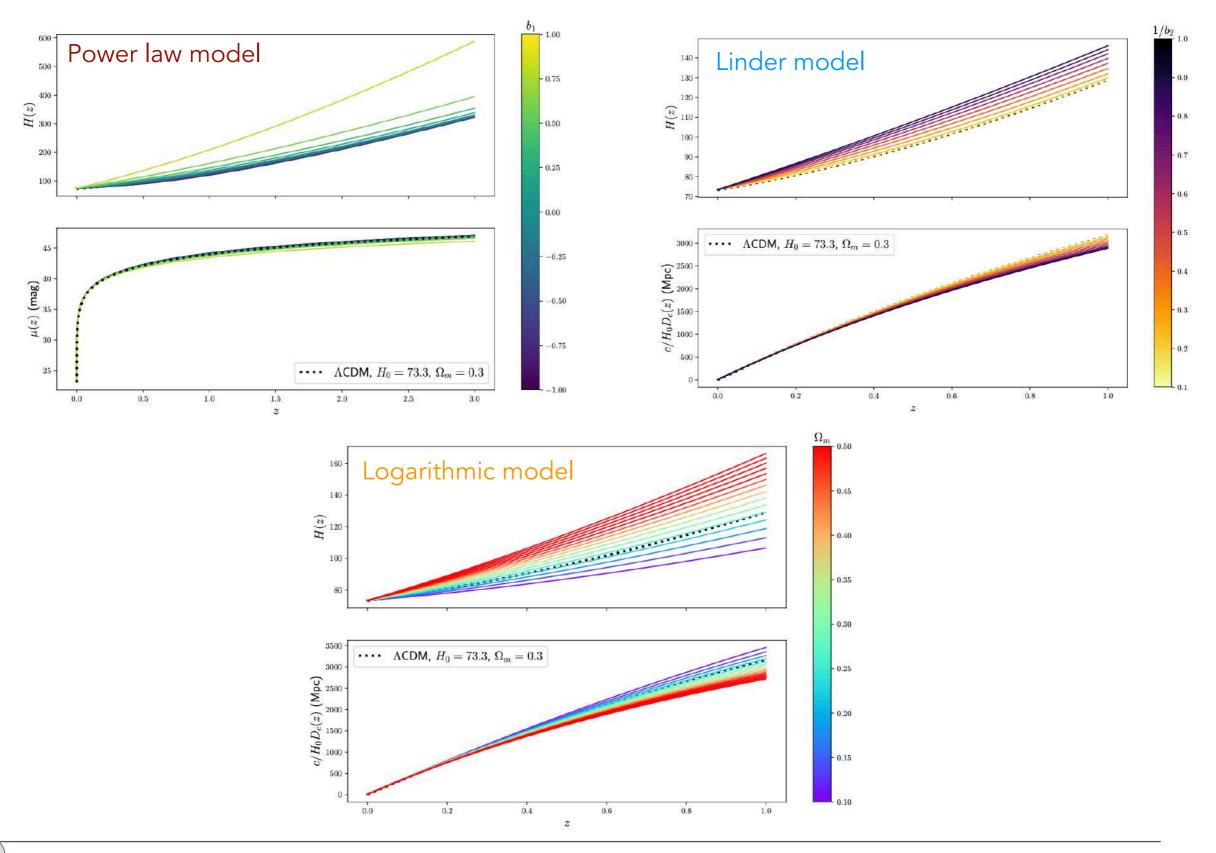


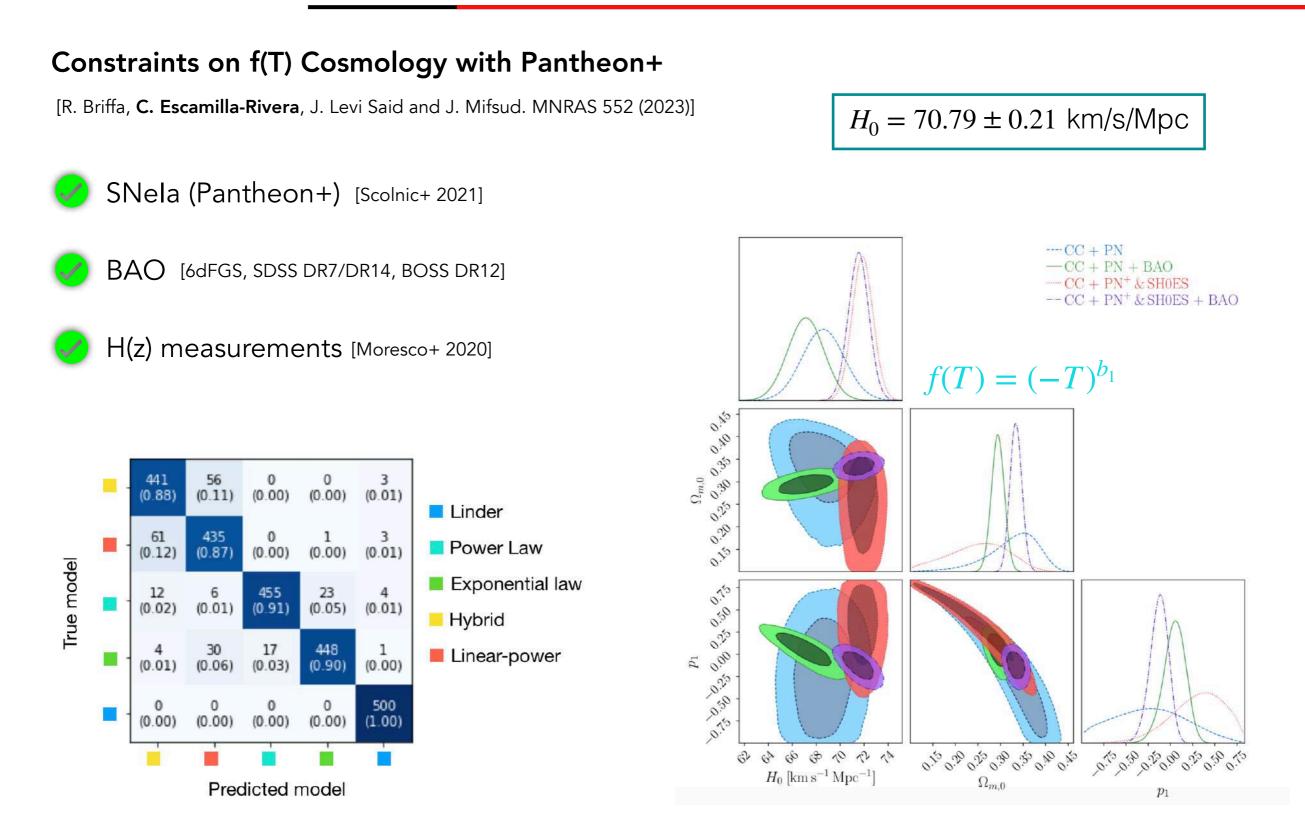
**TEGR** —-> Weitzenböck connection —-> Torsion with vanishing curvature

$$\tilde{\Gamma}^{\rho}_{\mu\nu} = E^{\rho}_{a} \left( \partial_{\mu} e^{a}_{\nu} + w^{a}_{b\mu} e^{b}_{\nu} \right) \qquad T^{\rho}_{\mu\nu} = \tilde{\Gamma}^{\rho}_{\nu\mu} - \tilde{\Gamma}^{\rho}_{\mu\nu}$$

[S.Bahamonde, K. Dialektopoulos, **C. Escamilla-Rivera**, G. Farrugia, V. Gakis, M. Hendry, M. Hohmann, J. Levi Said, J. Mifsud and E. Di Valentino. Rept.Prog.Phys. 86 (2023)]

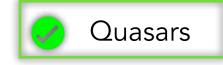
			Test					
Popular models of f(T)*	Cosmological behaviour	Maturity	Late-time	Early-time				
$f(T) = (-T)^{b_1}$	Stability 🥑	Reproduces	SNela 🥏	[Aguilar, <b>C. Escamilla-</b> <b>Rivera</b> , J. Levi Said, J. Mifsud. In preparation				
$f(T) = T_0 \left[ 1 - Exp \left( -b_2 \sqrt{T/T_0} \right) \right]$	Fine-tuning 🥏	Reproduces	BAO 🥏	(2023)]				
$f(T) = T_0 \sqrt{T/b_4 T_0} \log(b_4 T_0/T)$	Cosmic 🥏 acceleration	w(z)CDM	H(z) 🥏					
$S = \frac{1}{2\kappa^2} \int d^4x e[-T + f(T)] + \int d^4x e \mathscr{L}_m$								
* Power law model — Linder model — Logarithmic model $H^2 + \frac{T}{3}f_T - \frac{f}{6} = \frac{\kappa^2}{3}\rho$ ,								
		$\dot{H}(1 -$	$f_T - 2Tf_{TT})$	$= -\frac{\kappa^2}{2} \left(\rho + p\right) ,$				





[R. Sandoval, C. Escamilla-Rivera, R. Briffa, and J. Levi Said, 2309.03675 (2023)]

```
SNela (Pantheon+) + BAO + H(z)
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XA sample [Negrete+ 2014-2018, Marziani+ 2018, ]

- ~ 250 objects
- 1. Radiate near the Eddington limit (relation between this limit and the BH mass)
- 2. BH mass can be obtained through the virialized relation
- 3. Using an ionization parameter we can reach an expression for the luminosity

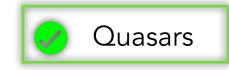
 $L(\text{FWHM}) = 7.88 \times 10^{44} (\text{FWHM})_{1000}^4,$ 

$$\mu = 2.5[\log L - \log(f_{\lambda}\lambda)] - 100.19 + 5\log(1+z),$$

$$\chi_{\mathrm{xA}}^2 = -\frac{1}{2} \sum_{i} \left[ \frac{(\mu_i - \mu(z_i, \Theta))^2}{\delta \mu_i^2} + \ln(\delta \mu_i^2) \right]$$



[R. Sandoval, C. Escamilla-Rivera, R. Briffa, and J. Levi Said, 2309.03675 (2023)]



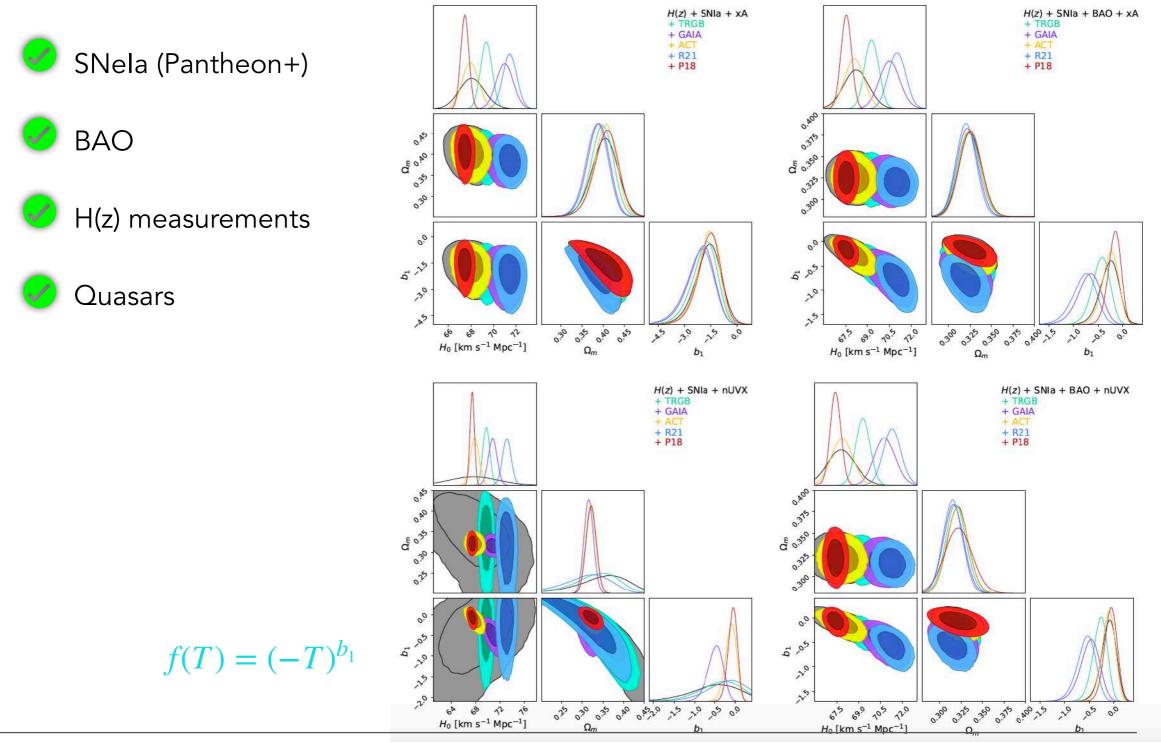
#### X-UV (nUVX) sample [Lusso+ 2020]

- 2421 objects (0 < z < 7.54)
- Cosmological candle:  $F_{UV} F_X$

$$\log(d_L) = \frac{\left[\log F_{\rm X} - \gamma F_{\rm UV}\right]}{2(\gamma - 1)} + \beta',$$
$$\mu = \frac{5}{2(\gamma - 1)}(\log F_{\rm X} - \gamma F_{\rm UV}) + 5\beta',$$

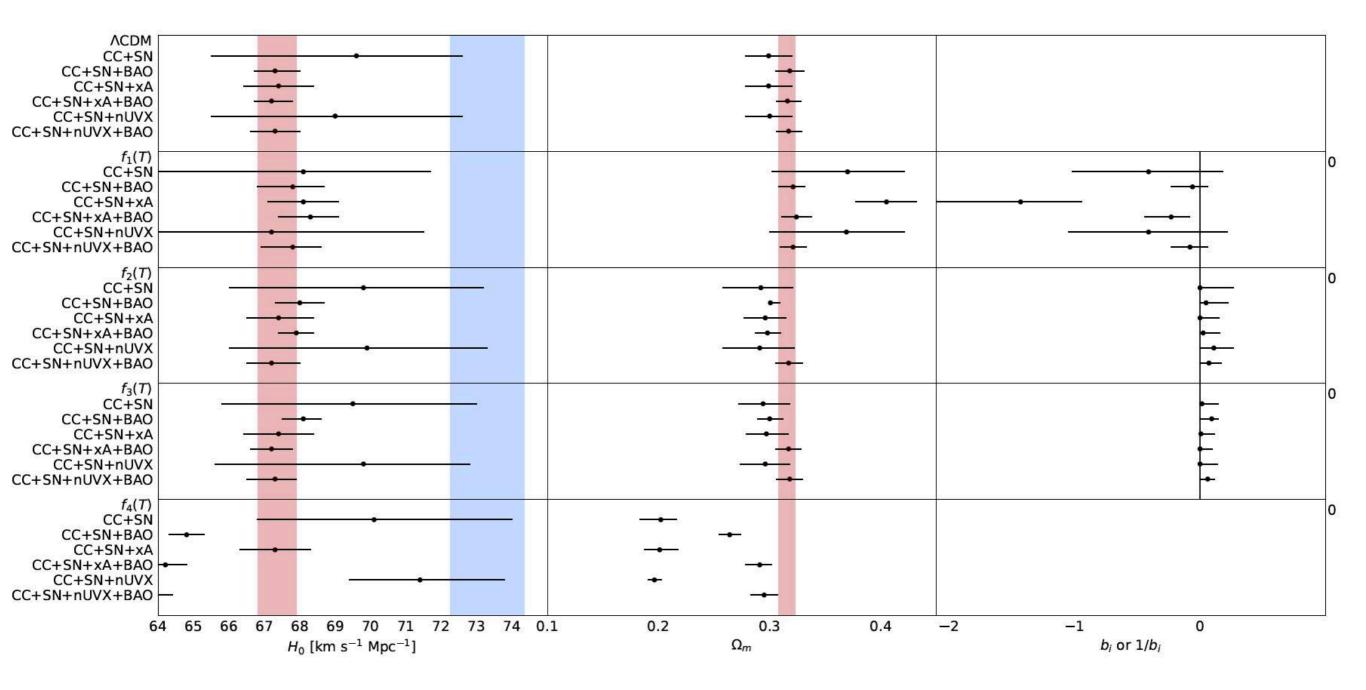
$$\mu = \frac{5}{2(\gamma - 1)} (\log F_{\rm X} - \gamma F_{\rm UV}) + 5\beta',$$

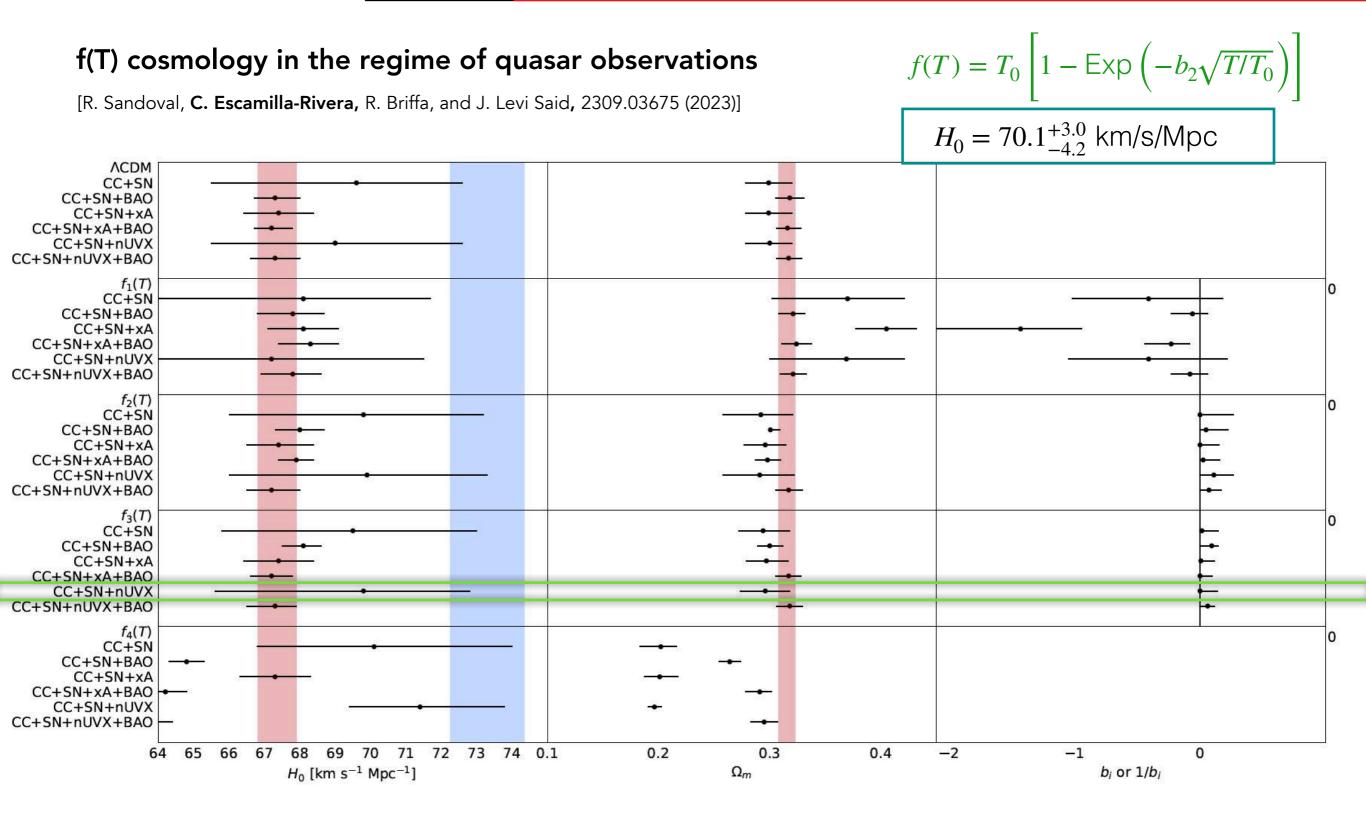
[R. Sandoval, C. Escamilla-Rivera, R. Briffa, and J. Levi Said, 2309.03675 (2023)]



© C. Escamilla-Rivera. ICN-UNAM || f(T) cosmology in the regime of quasar observations

[R. Sandoval, C. Escamilla-Rivera, R. Briffa, and J. Levi Said, 2309.03675 (2023)]





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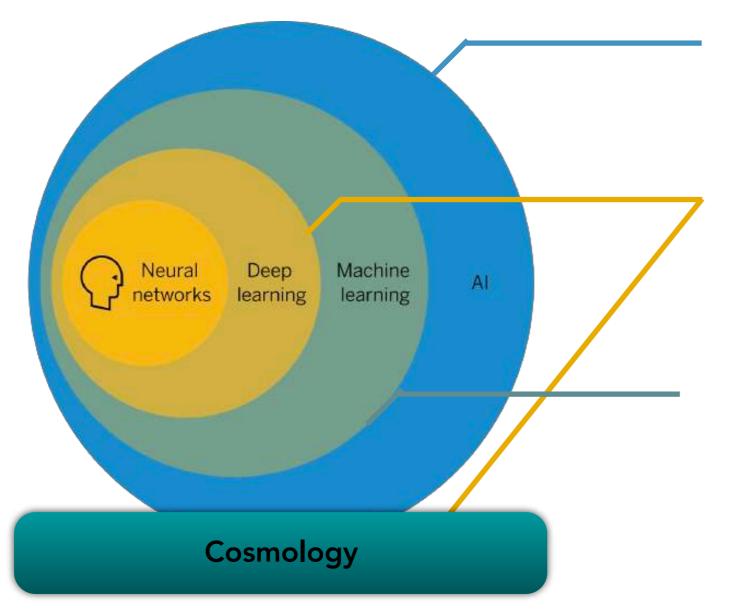
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#### and beyond: Machine Learning

Nonparametric and model independent reconstructions approaches							
Method	Assumption of	Binned	Low efficiency at	Underestimation	High computa-		
	prior/models		high $z$	of errors $z$	tional cost		
Principal Components	1	X	1	1	1		
Analysis (PCA)							
Nonlinear Inverse Ap-	1	X	1	X	1		
proach (NIA)							
Dipole of the Lumi-	X	×	1	X	X		
nosity Distance method							
(DLD)							
Nodal Reconstruction	1	X	×	X	X		
(NR)							
Genetic Algorithms	X	X	X	1	X		
(GA)							
Reconstructions of	X	×	×	1	1		
the Expansion History							
(MIR-I,II,III)							
Gaussian Processes	1	X	X	1	1		
(GP)							

[C. Escamilla-Rivera+ JCAP 10 (2021) 016]

# In which stage is "NN" Cosmology?



**Artificial Intelligence:** perceive their environment and define a course of action.

**Deep learning:** tasks are organised in consecutive layers, builded on the output of previous ones. Mimics the distributed approach to problem-solving.

Machine learning: tasks are complete without being explicitly programmed to do so.

Maping from observations to theory

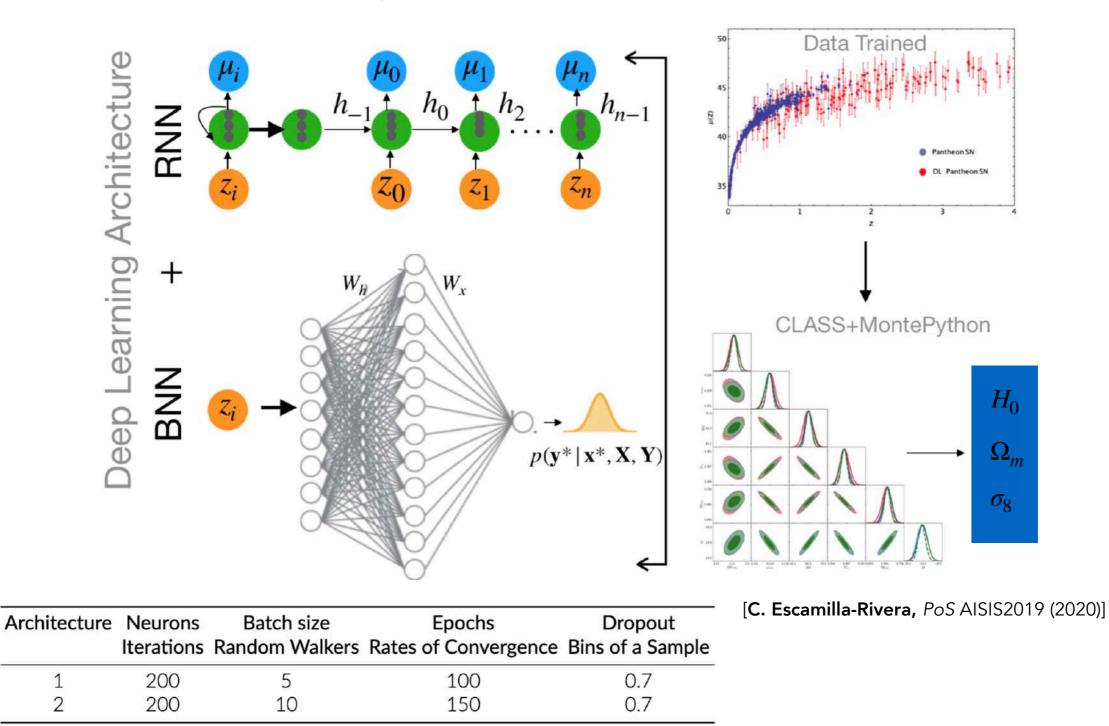
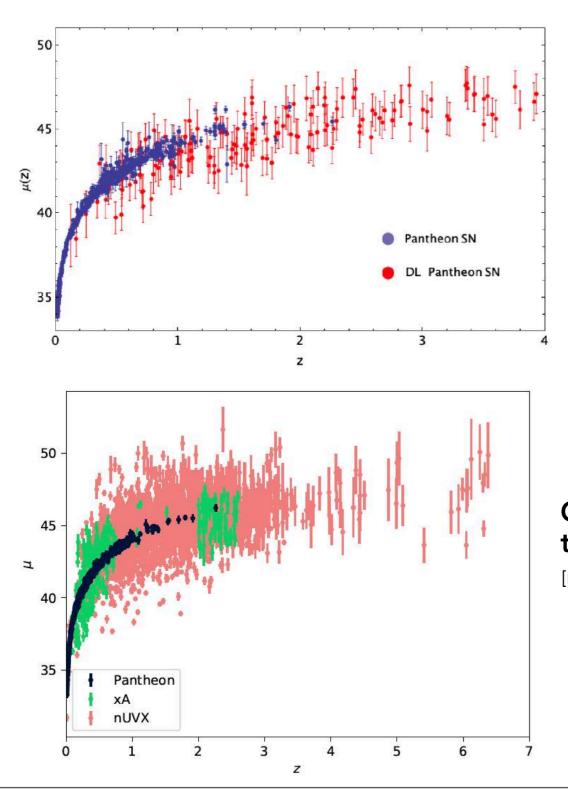


Table 1. Choice of hyperparameters in both RNN architectures and their association with Bayesian cosmology processing language.

#### Observations data mining results



# A deep learning approach to cosmological dark energy models

- [C. Escamilla-Rivera, M.Carvajal and S. Capozziello. JCAP (2020)]
- [C. Escamilla-Rivera and C. Zamora. JCAP (2020)]
- [C. Escamilla-Rivera, M.Carvajal, C. Zamora and M. Hendry. JCAP (2022)]

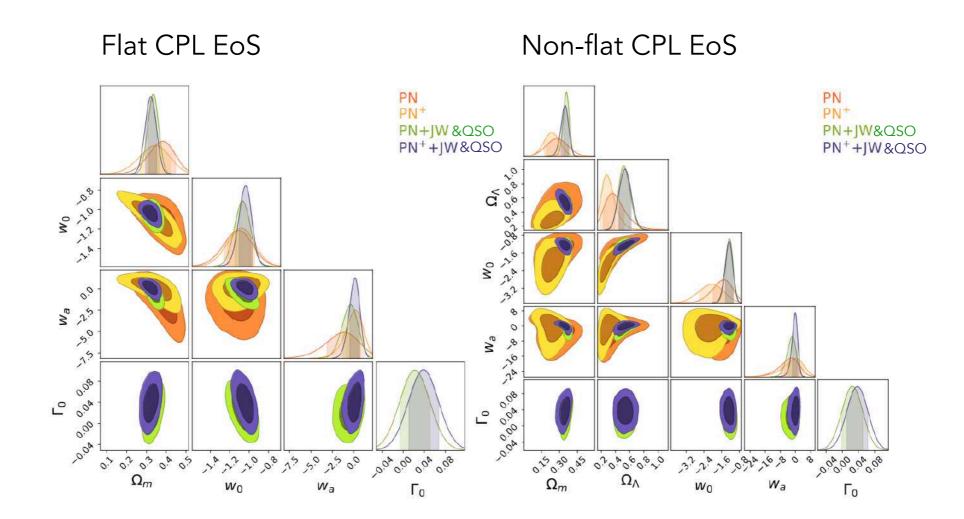
# Constraints on dark energy using deep learning training for quasars and JWST forecasting

[P. Maldonado, C. Escamilla-Rivera and R. Sandoval, in preparation (2023)]

#### Constraints on dark energy using deep learning training for quasars

[P. Maldonado, C. Escamilla-Rivera and R. Sandoval, in preparation (2023)]





## Conclusions

- Teleparallel Cosmology is maturing fast
- Quasars can help to relax local tensions (!)
- We can now classify models (likelihood free inference) using training data

 Data + NN can improve constraints on cosmological parameters, e.g H<sub>0</sub>

# Thanks for your kind attention

