

Precise characterization of the β^+ decay of ^8B to ^8Be



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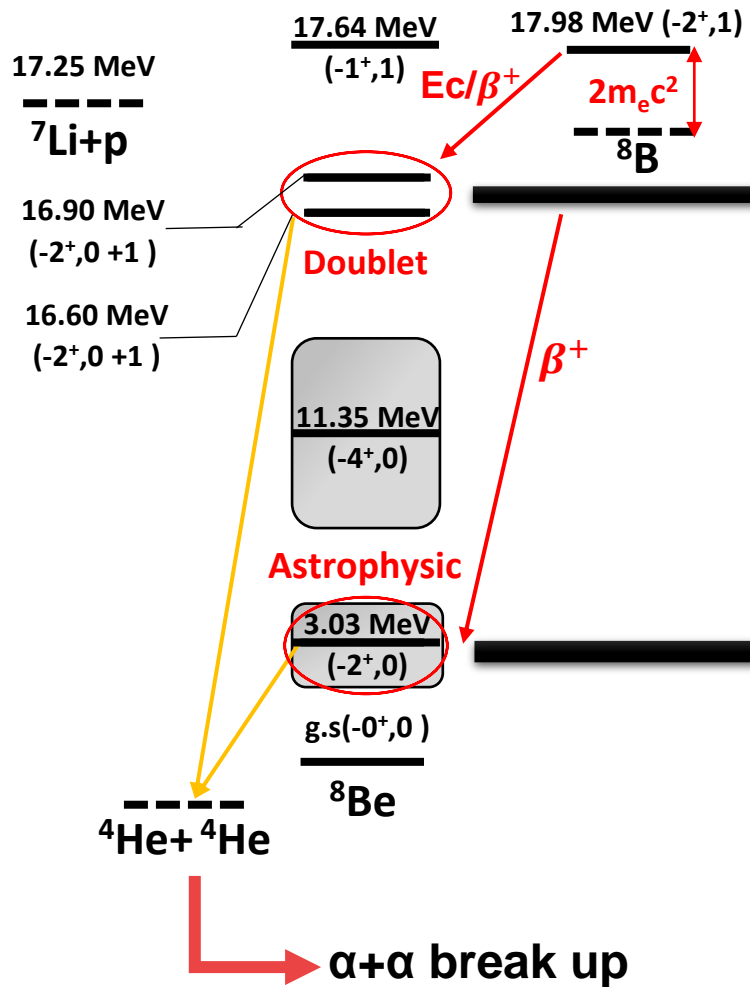


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The β^+ decay of ^8B to ^8Be is of interest for **astrophysics** and **nuclear structure**

Decay scheme



The 16.6 and 16.9 MeV levels of ^8Be are considered to form a **fully mixed 2^+ isospin doublet** ($^7\text{Li} \otimes p ; ^7\text{Be} \otimes n$)

E. Matt, et al. Phys. Rev. Lett. 9(2):17...

Main FOCUS

...nD Thesis (UCM, Department of Physics, Sep. 2020)

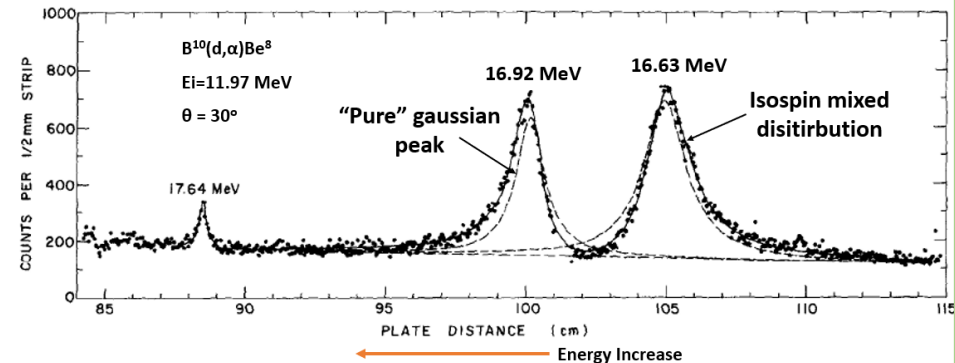
Source of high-energy solar neutrinos above 2 MeV

Main contributor to what was known as the **“solar neutrino problem”**

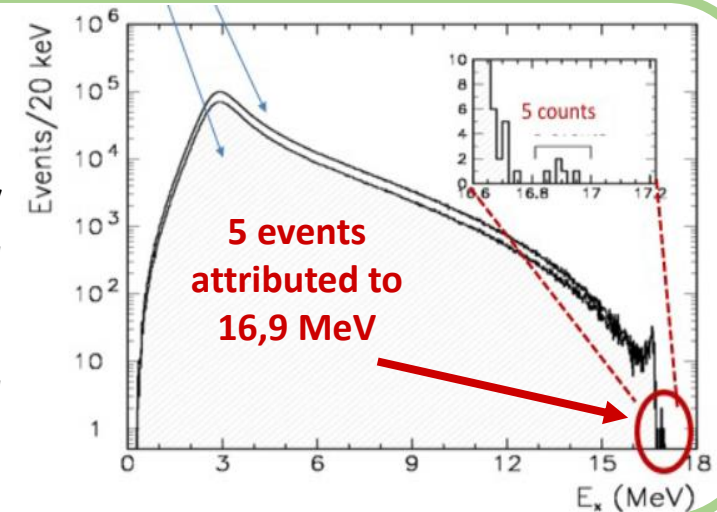
JYFL08

O Kirsebom. Phys. Rev. C, 83(6):065802–065822, 2011.

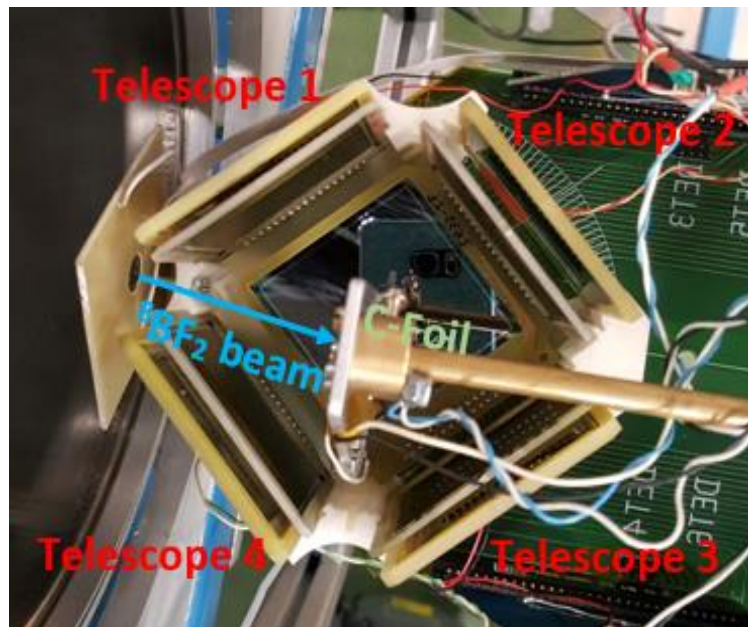
- ✓ The 16,6 MeV and 16,9 MeV have been observed through **reaction experiments**.
- ✓ The line shape can only be explained by **assuming an equal isospin mixture**. (W.D. Callender and C.P. Brown, Physical Review 2 (1970))
- ✓ Isospin mixture can be determined through the Fermi (F) and Gamow-Teller (GT) feeding to the levels.
- ✓ The β^+ decay is sensible to B_F and B_{GT} , but has a problem: **statistics**.



- The 3 MeV level takes most of the β -feeding, (>88%).
- 16,6 MeV : observed by several groups in EC/ β^+ -decay studies (E. Matt et al., Physics Letters 9 (1964) 174. +C.P. Browne et al., Physics letters 23 (1966) 371).
- 16,9 MeV : first hinted at the JYFL08 experiment (O. Kirsebom et al., Phys. Rev. C 83 (2011) 065802.),



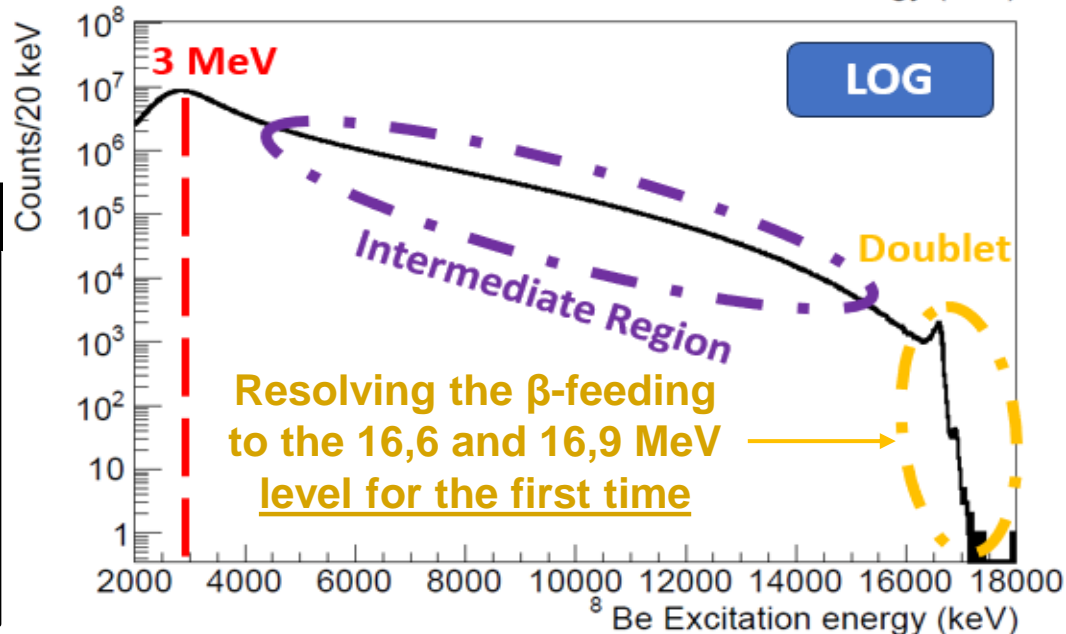
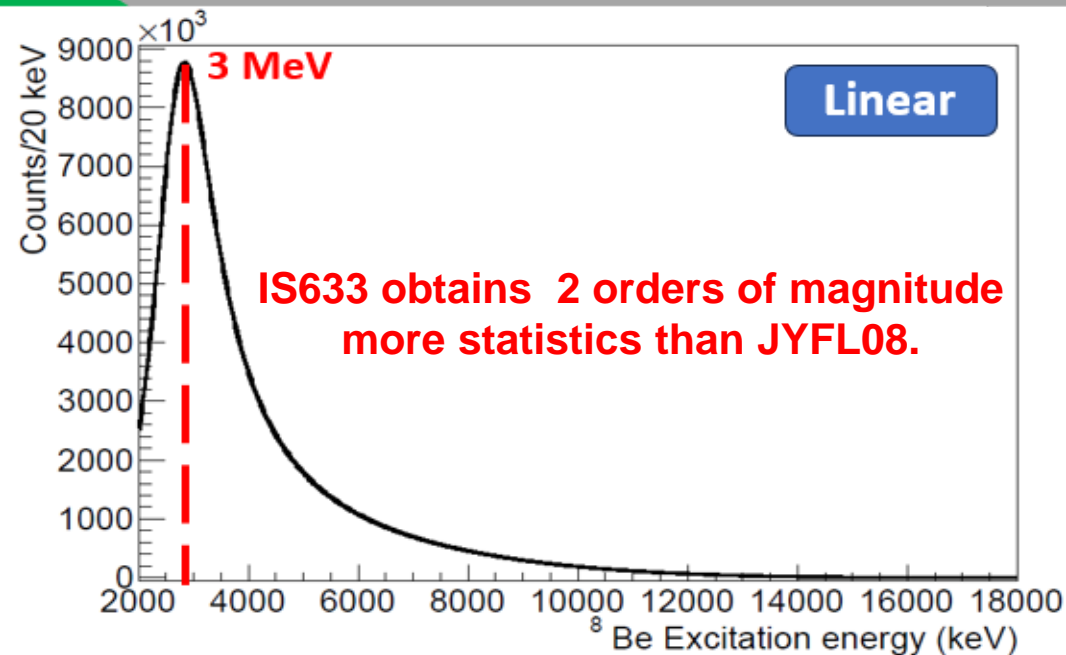
Experiment IS633 was designed to improve on the previous result



- 4 Si $\Delta E-E$ telescopes DSSD + PAD
- stop α : 1-9 MeV -- 60 μ m (U6 & U2)
- stop α : 1-7 MeV -- 40 μ m-- 40 μ m (U3 & U4)
- (low β -response)

Our experiment follows a **four-step technique**

1. ^8B nucleus is implanted in a ^{12}C foil
2. The β^+ decay of ^8B ($T_{1/2} = 771.17(94)$ ms) populates states of ^8Be
3. ^8Be is unbound $\rightarrow \alpha-\alpha$ break up
4. Reconstruction of the $\alpha-\alpha$ coincidence spectrum through a system of four telescopes



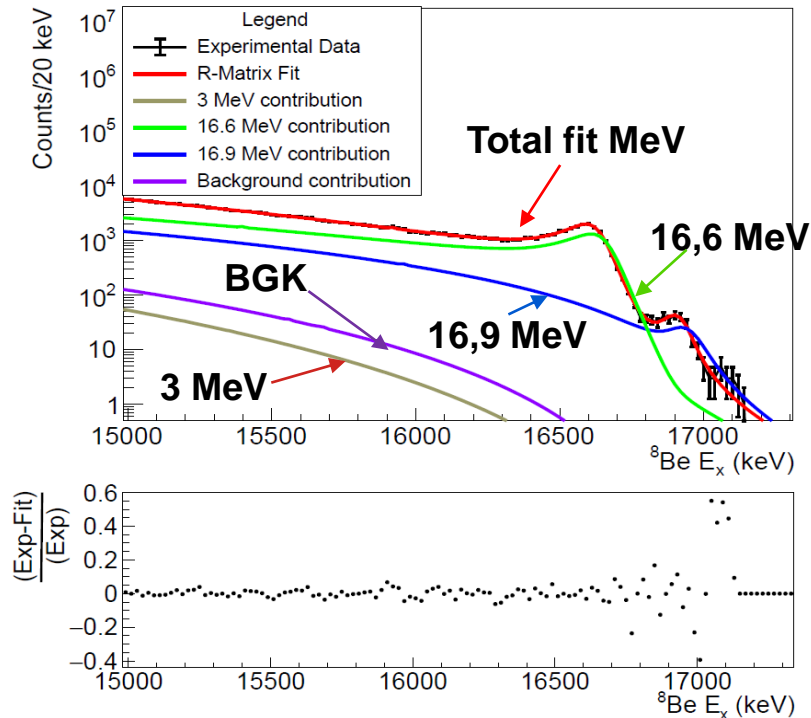
Our objective is to extract B_F , B_{GT} and Γ from the spectrum.

Levels fed through β^+ /EC are **to broad to be fitted** with a simple function (Landau, Gauss, ...)

Instead a **4 level R-Matrix** algorithm was used $R_{c'c} = \sum_{\lambda} \frac{\gamma_{\lambda c'} \gamma_{\lambda c}}{E_{\lambda} - E}$

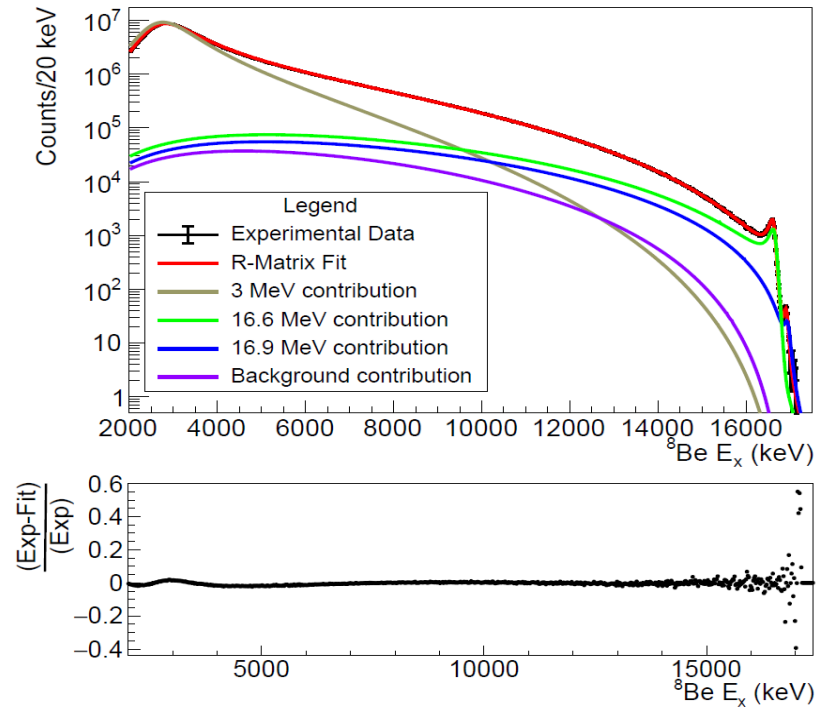
3 MeV → Dominant
 16,6 MeV } Main decay
 16,9 MeV }
 BKG → Level tails

Local 2+ doublet fit



The contribution of the broad 3 MeV state does not influence in the region of the 2+ doublet.

Global fit



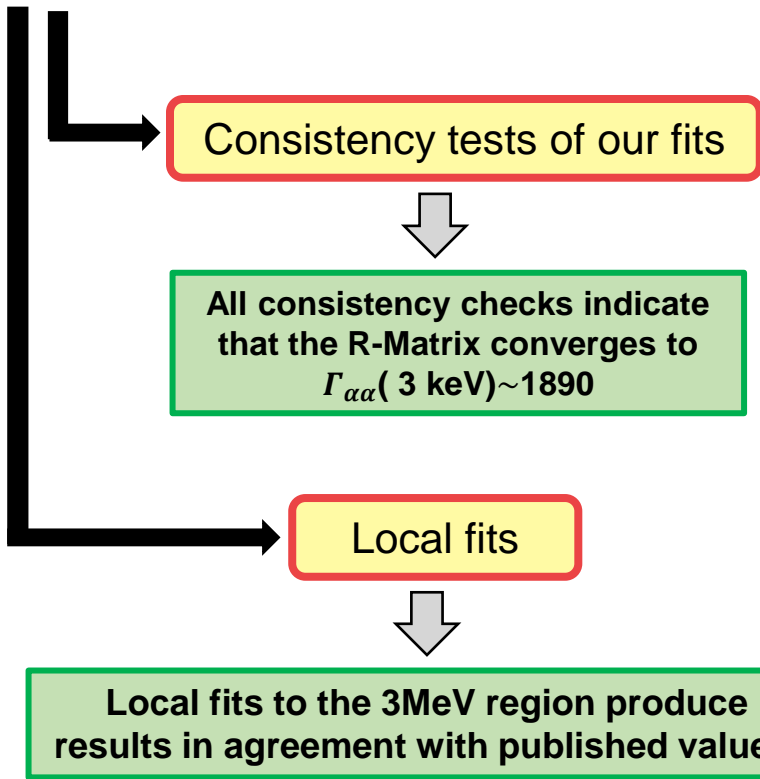
The global fit includes the contributions of the 3 MeV and Intermediate region.

There is a **discrepancy** between the Decay Width ($\Gamma_{\alpha\alpha}$) obtained through our R-Matrix and those of the **adopted published values [Tilley (2004)]**



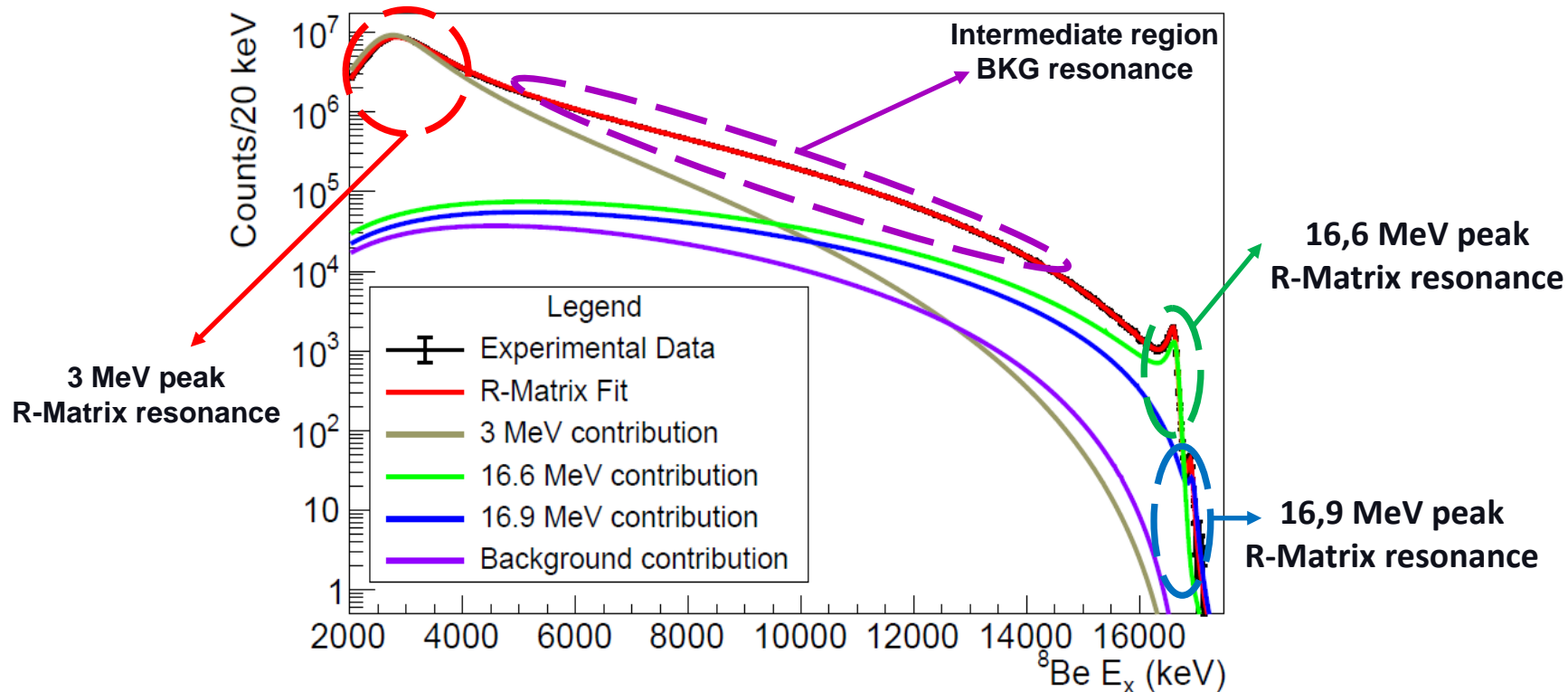
We have performed **two types of cross-checks** to find the reason for this discrepancy

Parameters		Tilley (2004)	Global Fit	Local 2+ Fit
r_0 (fm)		1.35	1.35	1.35
2_0^+	E(keV)	3030 (30)	3036	
	B_F	Beyond Wigner limit	0	
	B_{GT}		0,011485	
	$\Gamma_{\alpha\alpha}$ (keV)	1513(15)	1883,1	
2_1^+	E(keV)	16626(3)	16623	16632(70)
	B_F		1,282	0.32
	B_{GT}		0,572277	1.17
	$\Gamma_{\alpha\alpha}$ (keV)	108.1(5)	114,67	129.5
2_2^+	E(keV)	16922	16913	16919(20)
	B_F		0,39353	1.44(79)
	B_{GT}		1,046	0.35(49)
	$\Gamma_{\alpha\alpha}$ (keV)	74.0(4)	99,179	108(13)
2_{BKG}^+	E(keV)		37000	
	B_F		0	
	B_{GT}		0,093286	
	$\Gamma_{\alpha\alpha}$ (keV)		12116	



Local fits produce results in agreement with the literature → global fits don't

The problem appears in the intermediate region (BKG level) → distortion of the 3 MeV resonance



- R-Matrix decomposes the spectrum in resonant levels.
- For an excitation to continuum, a virtual resonance must be used.
- This only works if the continuum is close to the resonant levels.
- But if that is not the case R-Matrix will not work.

R-Matrix can not fit the whole spectrum due to the intermediate (not resonant) region.

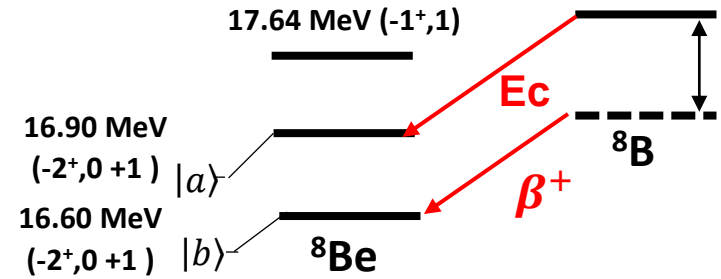
Assuming 16,9 and 16,6 can be expressed as a mixture of two states with well defined isospin.

$$|16,6\rangle = \alpha|T=0\rangle + \beta|T=1\rangle \quad \alpha^2 + \beta^2 = 1$$

$$|16,9\rangle = \beta|T=0\rangle - \alpha|T=1\rangle \quad \text{mixing coefficients}$$

If the states are completely mixed : $\alpha^2/\beta^2 = 1$

$M_{0,GT} \gg M_{1,GT} = 0$ Assumed due to S.M calculations



How can we find α^2/β^2 ?

Method 1

$$\frac{\alpha^2}{\beta^2} = \frac{B_{16.9,F}}{B_{16.6,F}}$$

Method 2

$$\frac{\alpha^2}{\beta^2} = \frac{B_{16.6,GT}}{B_{16.9,GT}}$$

Method 3

$$\alpha^2 = \frac{\Gamma_{16.6}}{\Gamma_0} = \frac{\Gamma_{16.6}}{\Gamma_{16.6} + \Gamma_{16.9}}$$

$$\beta^2 = \frac{\Gamma_{16.9}}{\Gamma_0} = \frac{\Gamma_{16.9}}{\Gamma_{16.6} + \Gamma_{16.9}}$$

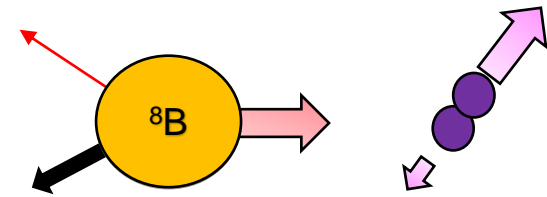
Method	Isospin coefficient ratio (α^2/β^2)	
	Local 2 ⁺ doublet Fit	Global Fit
B_F	5(12)	1,71(48)
B_{GT}	3(5)	1,73(42)
Γ	1,20(14)	1,22(13)

The results from the Fermi and GT feeding are inconsistent between local and global fits.

The Isospin coefficient ratio obtained from the decay width is in accordance with expectations.

An alternative method has been used : the **β-Recoil**

- The ${}^8\text{B}$ suffers a β^+ decay emitting an e^+ and ν_e
- Due to **momentum conservation** the ${}^8\text{Be}$ has a recoil
- This additional momentum is transferred (not always equally) to the emitted α particles



The β -recoil distribution is dependent on the ${}^8\text{Be}$ energy.
(D. Schardt and K. Riisager, Z. Phys. A 345 (1997))

$$w(t) = N_\beta \int_{W_{min}}^{W_0} F(Z, W) \frac{p_\beta}{2k} \left[(c_1 - c_2) W (W_0 - W) - A \frac{c_2^2 - c_1^2}{2} p_\beta \frac{t}{k} - A \frac{c_2^3 - c_1^3}{3} p_\beta^2 \right] dW + N_{EC} \frac{f_{BK} + f_{BL1}}{2T_{max}}$$

Recoil Distribution $\underbrace{\hspace{10em}}$ Fermi function $\underbrace{\hspace{10em}}$ Asymmetry parameter $\underbrace{\hspace{10em}}$ Ec-component

β-component

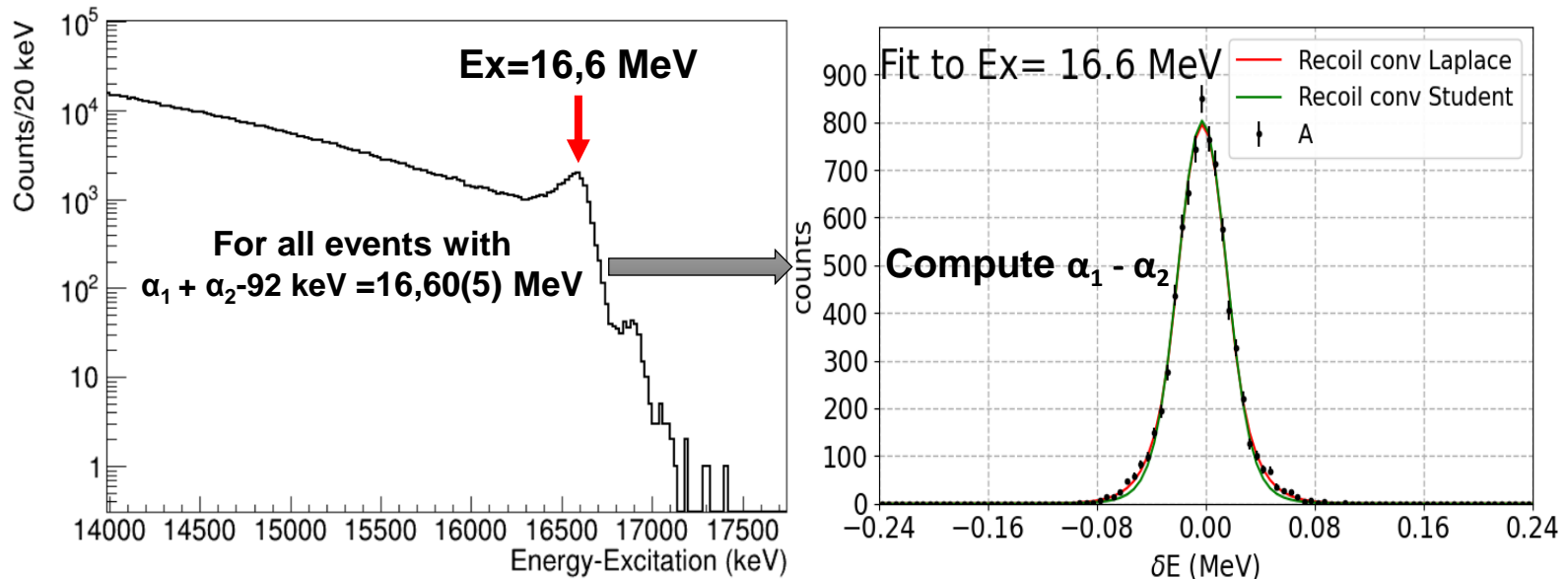
For ${}^8\text{B} \rightarrow {}^8\text{Be}$ the asymmetry parameter is directly related with B_F and B_{GT}

$$\frac{B_F}{B_{GT}} \left(\frac{g_V^2}{g_A^2} \right) = - \frac{(A+1)}{(A-1)} \begin{cases} A=1 \rightarrow B_{GT}=0 \\ A=0 \rightarrow B_F=B_{GT} \\ A=-1 \rightarrow B_F=0 \end{cases}$$

For a fixed value of E_x (i.e $\alpha_1 + \alpha_2 - 92$ keV) the β -recoil spectrum can be used to obtain B_F and B_{GT}



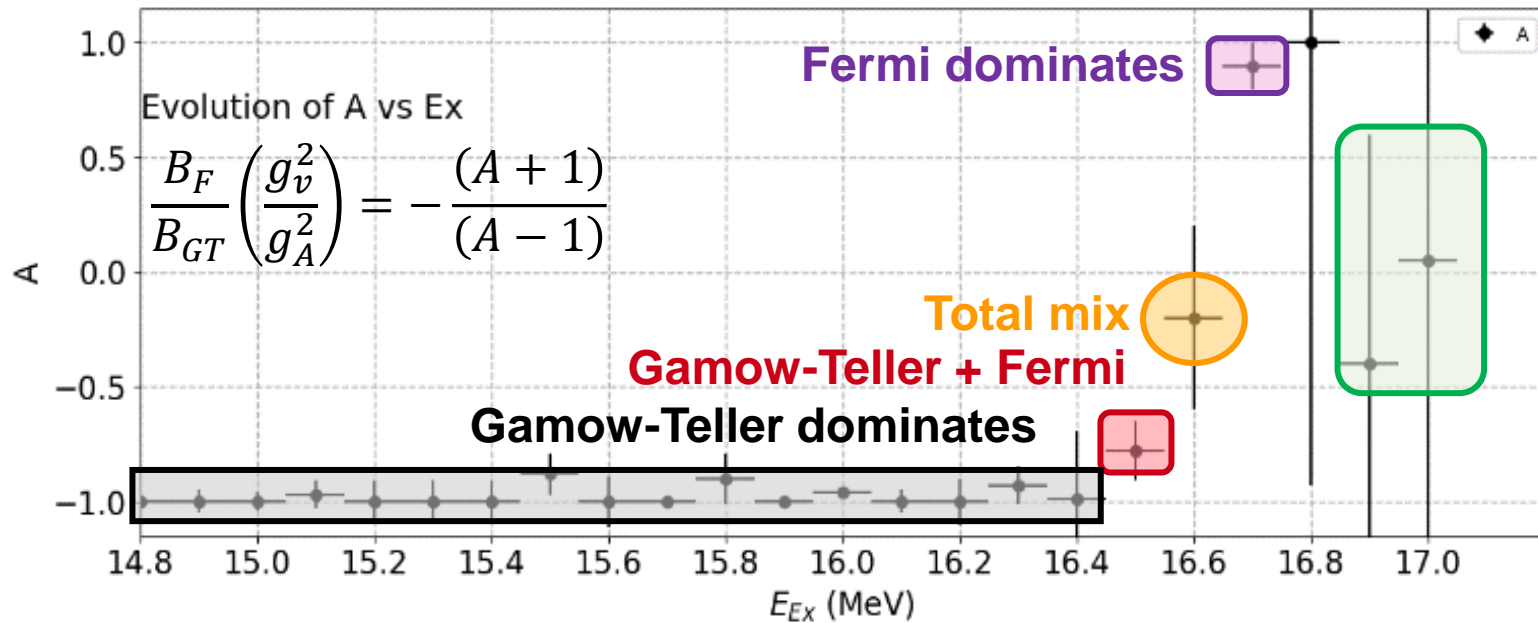
The β -Recoil spectrum is equivalent to the $\alpha_1 - \alpha_2$ spectrum



The $\alpha_1 - \alpha_2$ spectrum is fitted using the recoil function folded with a response function (Laplace or T-student) with A as a fit parameter

To get a satisfactory fit; the N_β and N_{EC} parameters must be theoretically computed and fixed.

The Evolution of A can be used to deduce the evolution of B_F and B_{GT}



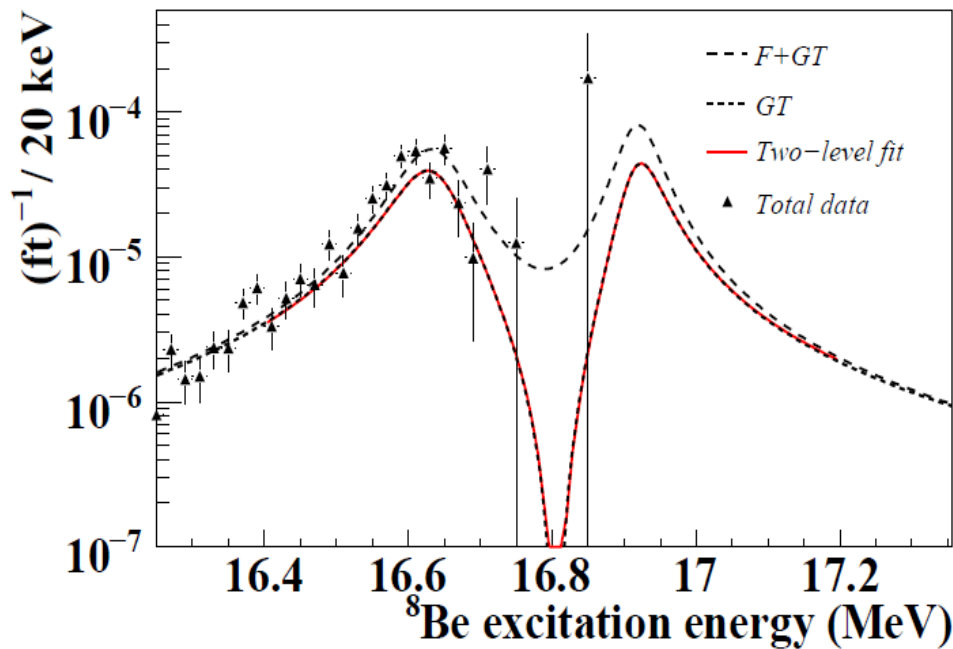
- 1) From 3 MeV to 16,2 MeV Gamow-Teller dominates the decay. ($A=-1,00(1) \rightarrow B_F=0 ; B_{GT}=1$).
- 2) Penetrating the doublet region, the influence of the Fermi contribution becomes noticeable.
- 3) In the 16,60 (5) MeV region $A=0,10(8) \rightarrow B_F= B_{GT}=0.5$, there is a total Fermi and GT coexistence
- 4) Between the resonances, the Fermi contribution dominates ($A=1,0(4) \rightarrow B_F=1 ; B_{GT}=0$)
- 5) In 16,90(5) MeV there is a very high uncertainty!!

The Gamow Teller and Fermi contributions can be distinguished experimentally (for the first time!!) and three observations can be made

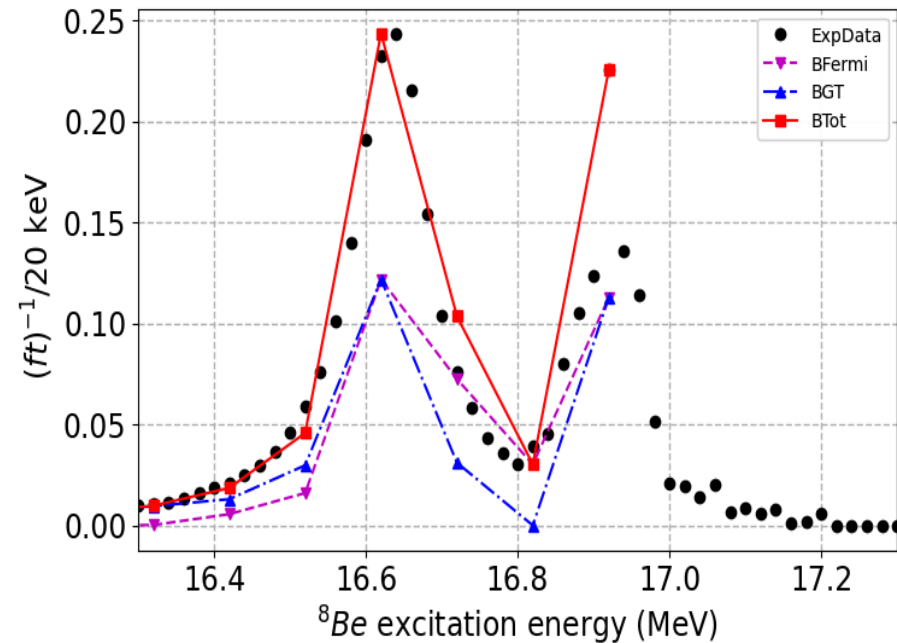
G-T dominates up to 16,3 MeV;
F and GT coexist in the resonances;
Fermi dominates between resonances

This result is consistent with β decay calculations (S.Hyldegaard, PhD thesis. Aarhus University (2010)) and hinted in reaction experiments (W.D. Callender and C.P.Brown, Physical Review 2 (1970))

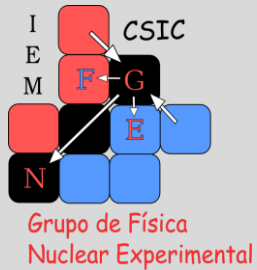
(S.Hyldegaard-JYFL08)



(Experimental)



- IS633 is the **first experiment** that enables to study the 2^+ doublet of by β /EC where Fermi and Gamow-Teller contributions could be separated.
- **R-matrix local fits** to the ^8Be excitation spectrum to either low or high energy-region produces **good results** (Comparison with JYFL08 experiment assure that **IS633 is consistent with previous results** and no extra effect from high statistics were found).
- **A R-matrix fit to the full ^8Be excitation spectrum** produces Γ values for the 3 MeV state that **differs from the the literature** (This could be due to a very broad intermediate level or caused by coupling to the “non-resonant” continuum).
- The uncertainties obtained for the B_F and B_{GT} were rather large.
- An additional analysis method was used employing the β -Recoil spectrum for each energy bin.
- The β -Recoil analysis confirms that GT transitions dominates up to 16,3 MeV ^8Be excitation energy; **F and GT contributions coexist in the resonances and the Fermi component dominates between resonances** (first experimental confirmation) .



Special thanks to the members of the **MAGISOL/IS633** collaboration

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***And you for your attention !!!
Questions?***

The β -recoil distribution is dependent on the ^8Be energy. (D. Schardt and K. Riisager, Z. Phys. A 345 (1997))

$$w(t) = N_\beta \int_{W_{min}}^{W_0} \underbrace{F(Z, W)}_{\text{Fermi function}} \frac{p_\beta}{2k} \left[\underbrace{(c_1 - c_2)W(W_0 - W)}_{\beta\text{-component}} - \underbrace{A \frac{c_2^2 - c_1^2}{2} p_\beta \frac{t}{k}}_{\text{Asymmetry parameter}} - \underbrace{A \frac{c_2^3 - c_1^3}{3} p_\beta^2}_{\text{Ec-component}} \right] dW + N_{EC} \frac{f_{BK} + f_{BL1}}{2T_{max}}$$

Recoil Distribution

$$W_{min} = \begin{cases} \frac{m^2 + (W_0 + t/k)^2}{2(W_0 + t/k)} & \text{for } \frac{t}{k} < -(W_0 - m) \\ m & \text{for } \left| \frac{t}{k} \right| \leq (W_0 - m) \\ \frac{m^2 + (W_0 - \beta t/k)^2}{2(W_0 - t/k)} & \text{for } \frac{t}{k} > (W_0 - m) \end{cases}$$

$$c_1 = \min \left\{ \cos \psi_{max}, \max \left[\cos \psi_{min}, \frac{W - W_0 - t/k}{p_\beta} \right] \right\}$$

$$c_2 = \max \left\{ \cos \psi_{min}, \min \left[\cos \psi_{max}, \frac{W - W_0 - t/k}{p_\beta} \right] \right\}$$

$$A = \frac{g_v^2 B_F - (1/3 + 2/30 \tau \theta) g_A^2 B_{GT}}{g_v^2 B_F + g_A^2 B_{GT}}$$

$$\theta = \begin{cases} -\frac{J' + 1}{2J' - 1} & J' = J + 1 \\ 1 & J' = J \\ -\frac{J'}{2J' + 3} & J' = J - 1 \end{cases}$$

$$\tau = 10 \left[\frac{L(L+1)(2L+1)}{(2L-1)(2L+3)} \right]^{1/2} \left[\frac{(2J' - 1)(2J' + 1)(2J' + 3)}{J'(J' + 1)} \right]^{1/2} W(2; J' L J'' J' L)$$

For $^8\text{B} \xrightarrow{+\beta} ^8\text{Be}$ $\Delta J=0$ and $\Delta L=2$ therefore $\tau\theta = 10$ and B_F and B_{GT} can be deduced from A

$$A = \frac{g_v^2 B_F - (1/3 + 2/30 \tau \theta) g_A^2 B_{GT}}{g_v^2 B_F + g_A^2 B_{GT}} \Rightarrow \frac{B_F}{B_{GT}} \left(\frac{g_v^2}{g_A^2} \right) = -\frac{(A + 1)}{(A - 1)} \begin{cases} A=1 \rightarrow B_{GT}=0 \\ A=0 \rightarrow B_F=B_{GT} \\ A=-1 \rightarrow B_F=0 \end{cases}$$

A series of cross-checks were performed to ensure that that results are consistent and do not suffer from systematic errors such as summing or piled-up.

Check I: $\Gamma_{\alpha\alpha}$ = Fixed to the literature values

The R-Matrix does not generate good results ($\chi^2= 3991$; $\Gamma_{\alpha\alpha}(3 \text{ keV})= 1415.7$).

Check II: $\Gamma_{\alpha\alpha}$ = Set to the literature values (but allowed to change)

The R-Matrix parameters converge to the same values as before.
($\chi^2= 14$; $\Gamma_{\alpha\alpha}(3 \text{ keV})= 1949.8$)

Check III: Set the energy of the BKG level to 37 MeV

R-Matrix generates the best fit, the results also diverges form the literature
($\chi^2= 12$; $\Gamma_{\alpha\alpha}(3 \text{ keV})= 1883.1$)

All consistency checks indicate that the R-Matrix converges to $\Gamma_{\alpha\alpha}(3 \text{ keV})\sim 1890$

