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NON-HOLOMORPHIC MODULAR A_4 SYMMETRIC SCOTOGENIC MODEL

Shenzhen MSU-BIT University

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Based on [Phys.Lett.B 860 \(2025\) 139171](#), [Phys.Lett.B 803 \(2020\) 135294](#), [Phys.Rev.D 109 \(2024\) 3, 035016](#)

In collaboration with Hiroshi Okada and Takaaki Nomura

OVERVIEW

- ▶ Motivation
- ▶ Origin of Neutrino Masses
- ▶ Modular Symmetries as Flavor Symmetries
- ▶ Neutrino Masses and \mathcal{A}_4
- ▶ Modular \mathcal{A}_4 Symmetric Scotogenic Model
- ▶ non-Holomorphic Modular \mathcal{A}_4 Symmetric Scotogenic Model
- ▶ Summary



MOTIVATION

Standard Model of Elementary Particles



- * Neutrino Mass Origin?
- * Source of Neutrino mixing? PMNS matrix
- * Dark Matter?

$$G_{SM} = SU(3)_c \times SU(2)_L \times U(1)_Y$$

$$\mathcal{L}^Y = \bar{u} Y_u Q H + \bar{Q} Y_d d H + \text{h.c.}$$

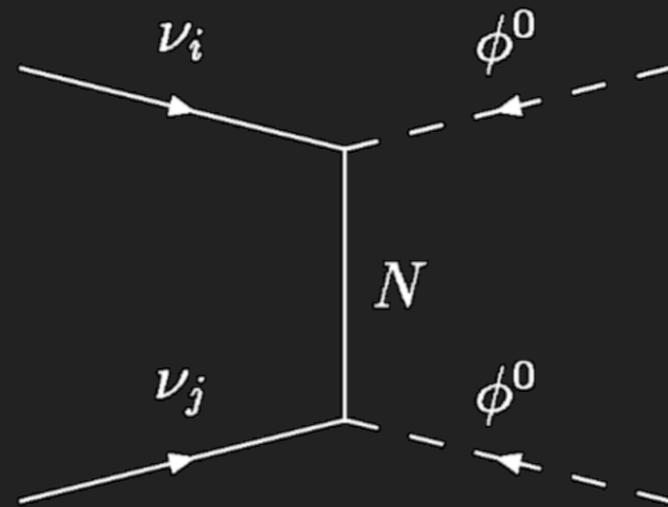
$$\langle h^0 \rangle = v \rightarrow m_f = Y_f v$$

ORIGIN OF NEUTRINO MASSES

Weinberg Operator 1979: $\frac{LHLH}{\Lambda} = \frac{(l^- H^+ - \nu H^0)(l^- H^+ - \nu H^0)}{\Lambda}$

Add $N_R \sim (1, 1, 0)$ under G_{SM}

$$\mathcal{L}_{new} = \bar{N} Y_D L H + m_N N_R N_R + \text{h.c.}$$



$$\begin{pmatrix} 0 & m_D \\ m_D & m_N \end{pmatrix} \Rightarrow m_\nu \simeq \frac{-m_D^2}{m_N}$$

$v \ll m_N \sim 10^{11} \text{ GeV} \rightarrow m_\nu \ll v$ with $Y_D \sim 1$
or $Y_D \ll 1 \rightarrow m_\nu \ll v$ with $m_N \sim O(10^{2-3} \text{ GeV})$

Seesaw-I

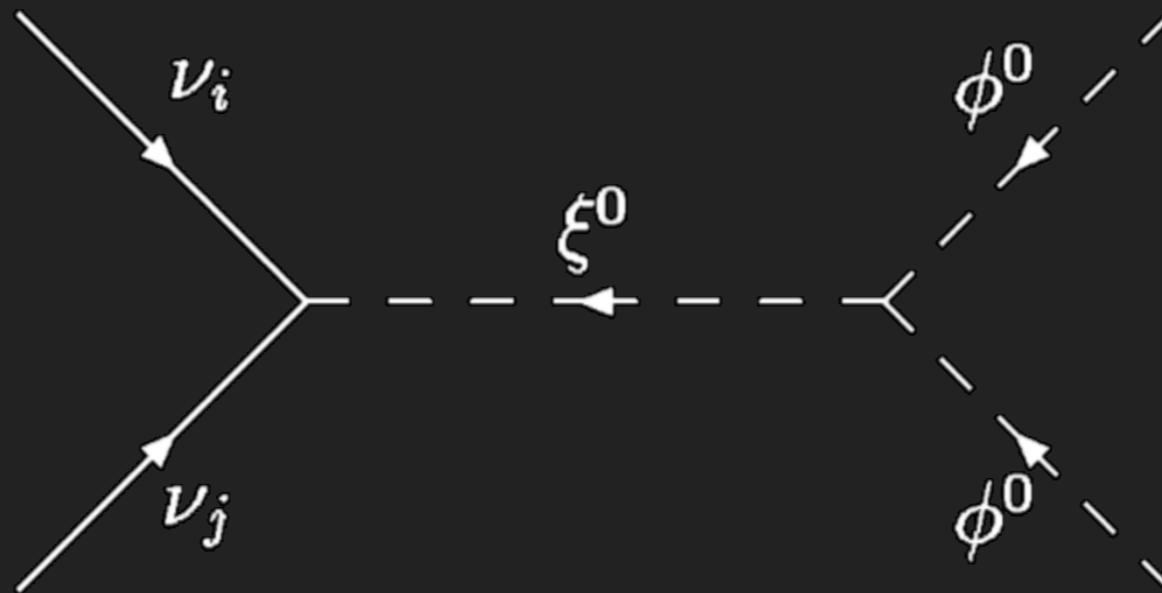
Seesaw-I

ORIGIN OF NEUTRINO MASSES

Seesaw-II

Add[1980] $\xi = (\xi^{++}, \xi^+, \xi^0) \sim (1, 3, 1)$ under G_{SM}

$$\mathcal{L}_{new} = YL\xi L - \mu H\xi H + \text{h.c.} \rightarrow m_\nu = Y \langle \xi^0 \rangle = -2 \frac{Y \mu v^2}{M_\xi}$$



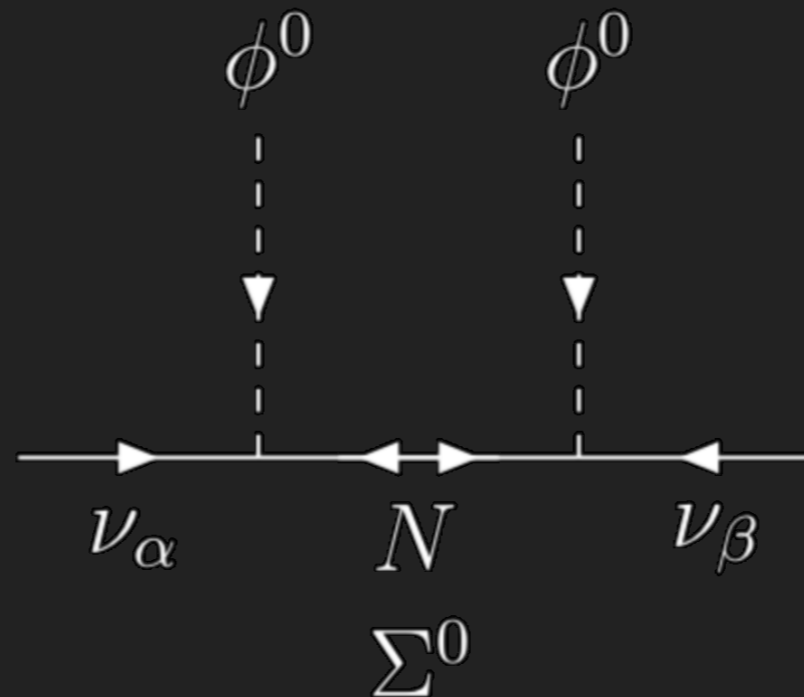
ORIGIN OF NEUTRINO MASSES

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Add $\Sigma_R \sim (1, 3, 0)$ under G_{SM}

SEESAW-III

$$\mathcal{L}_{new} = \bar{\Sigma} Y_D L H + m_N \Sigma_R \Sigma_R + h.c.$$

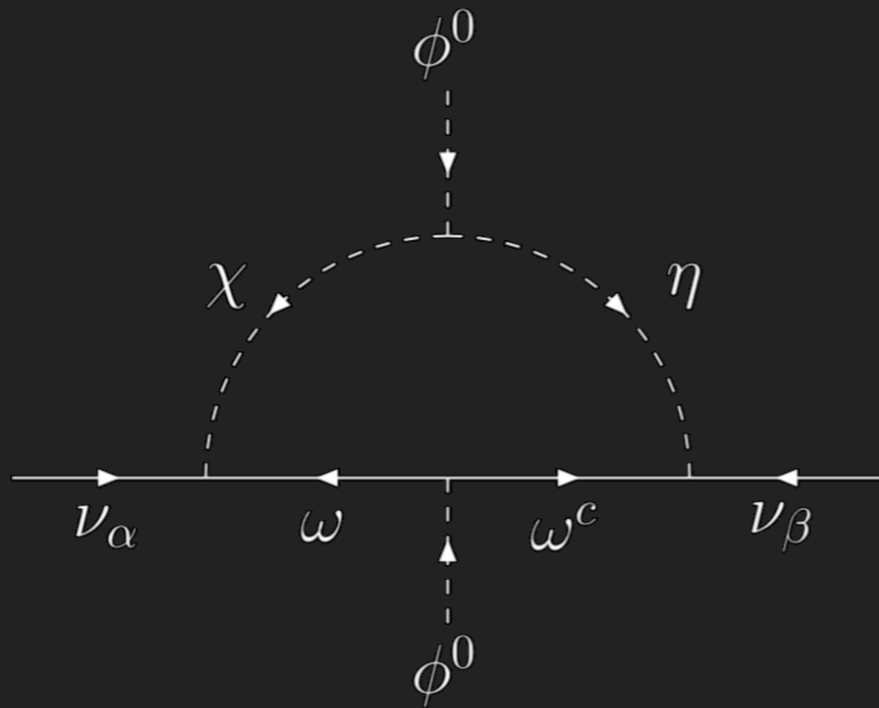


$$\begin{pmatrix} 0 & m_D \\ m_D & m_N \end{pmatrix} \Rightarrow m_\nu \simeq \frac{-m_D^2}{m_N}$$

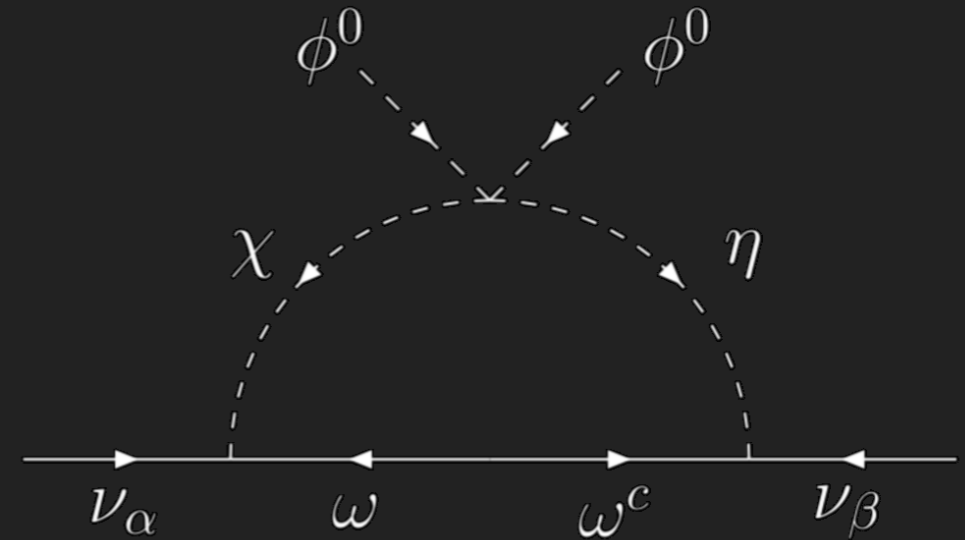
$\nu \ll m_N \sim 10^{11} \text{ GeV} \rightarrow m_\nu \ll \nu$ with $Y_D \sim 1$
or $Y_D \ll 1 \rightarrow m_\nu \ll \nu$ with $m_N \sim O(10^{2-3} \text{ GeV})$

ORIGIN OF NEUTRINO MASSES

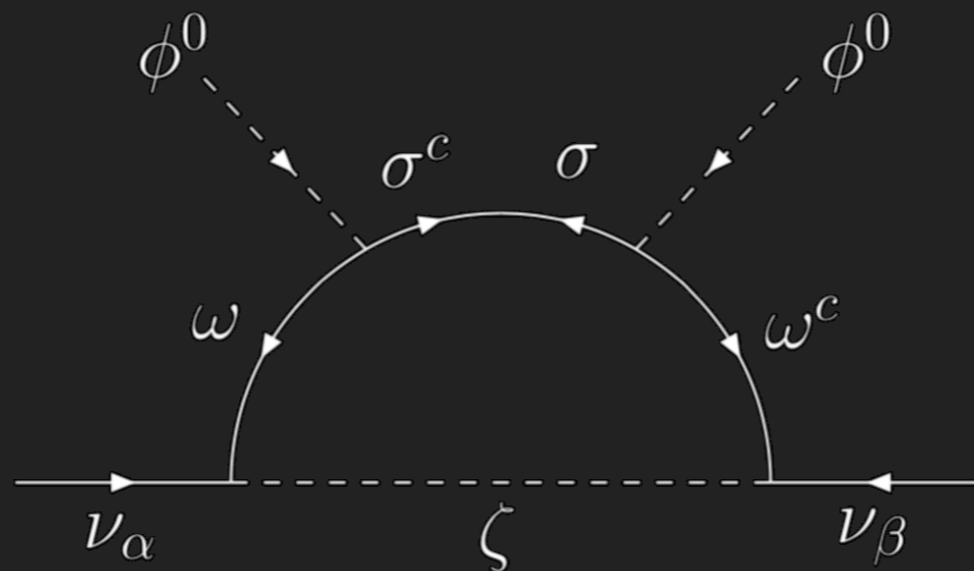
Radiative Majorana Neutrino Masses



Zee[1986]



Ma[2006]

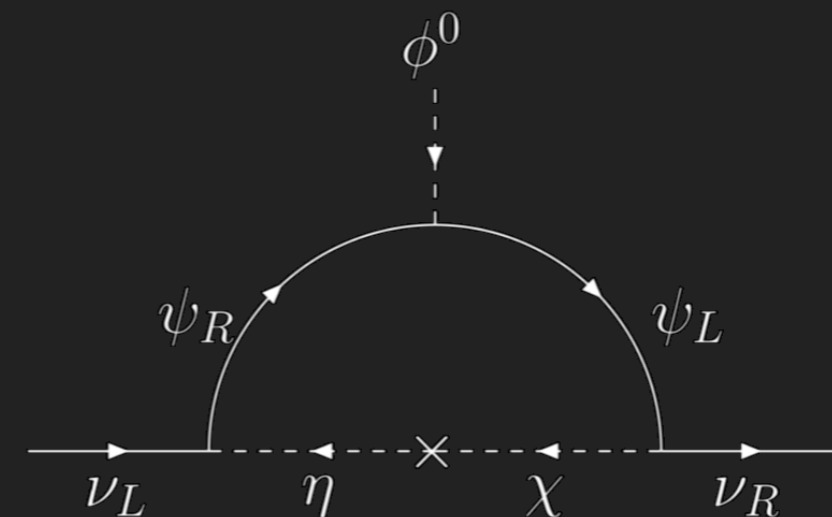
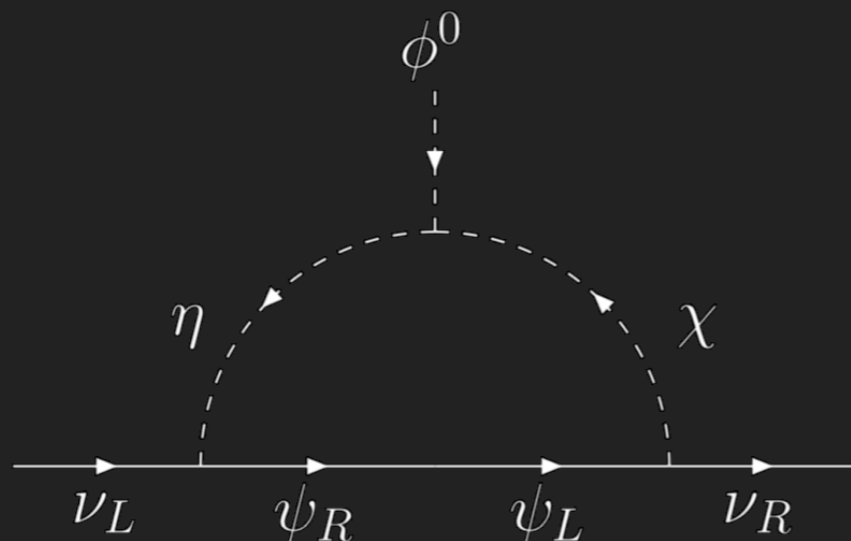
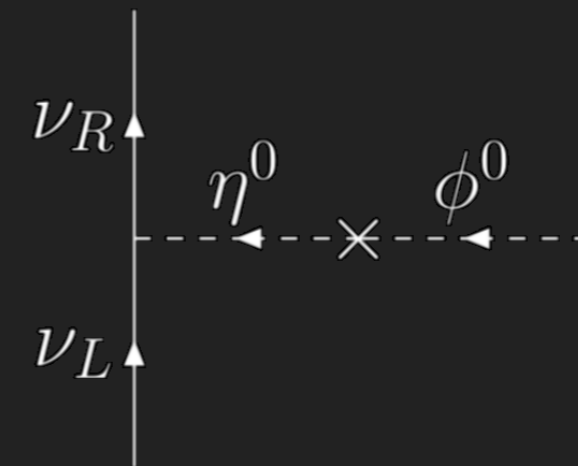
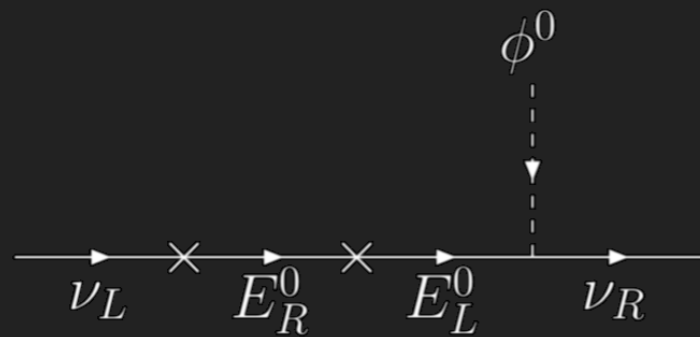
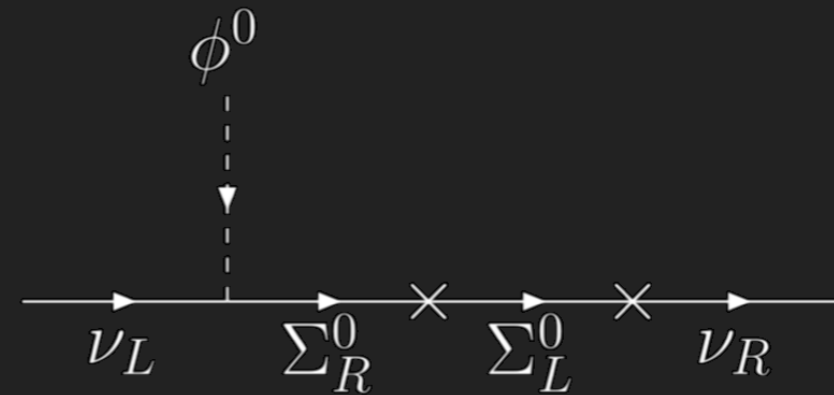
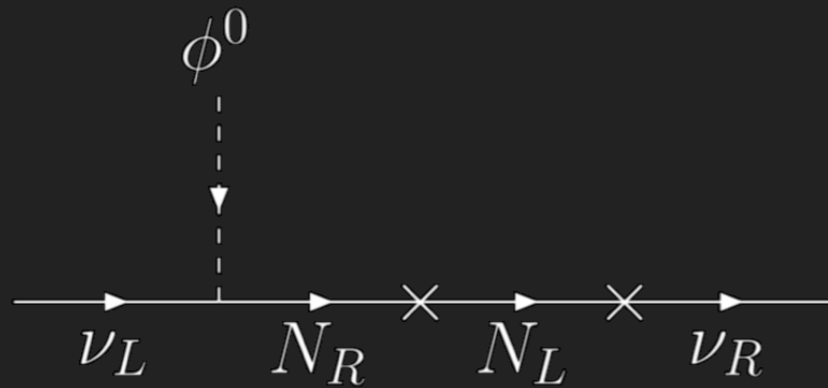


Fraser, Ma, OP[2014]

ORIGIN OF NEUTRINO MASSES

Dirac Neutrino Masses

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MODULAR SYMMETRY

$$\tau \longrightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d} \qquad S^2 = \mathcal{I}, \quad (ST)^3 = \mathcal{I}.$$

$$S = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad T = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, \quad S : \tau \longrightarrow -\frac{1}{\tau}, \quad T : \tau \longrightarrow \tau + 1.$$

$$\Gamma(N) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{Z}), \begin{pmatrix} a & b \\ c & d \end{pmatrix} \equiv \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \pmod{N} \right\} \qquad \Gamma_N \equiv \bar{\Gamma} / \bar{\Gamma}(N)$$

$$f_i(\gamma\tau) = (c\tau + d)^k \rho_{ij}(\gamma) f_j(\tau), \quad \gamma \in \Gamma(N) \qquad \Gamma_3 \simeq \mathcal{A}_4$$

$$y_1(\tau) = 1 + 12q + 36q^2 + 12q^3 + \dots$$

$$y_2(\tau) = -6q^{1/3}(1 + 7q + 8q^2 + \dots) \qquad q \equiv e^{i2\pi\tau}$$

$$y_3(\tau) = -18q^{2/3}(1 + 2q + 5q^2 + \dots)$$

[arXiv:[1706.08749](https://arxiv.org/abs/1706.08749)]

$$Y_3^{(4)} = (y_1^2 - y_2y_3, y_3^2 - y_1y_2, y_2^2 - y_1y_3) \qquad Y_1^{(6)} = y_1^3 + y_2^3 + y_3^3 - 3y_1y_2y_3$$

$$Y_1^{(4)} = y_1^2 + 2y_2y_3 \qquad Y_{3a}^{(6)} = (y_1^3 + 2y_1y_2y_3, y_1^2y_2 + 2y_2^2y_3, y_1^2y_3 + 2y_3^2y_2)$$

$$Y_{1'}^{(4)} = y_3^2 + 2y_1y_2 \qquad Y_{3b}^{(6)} = (y_3^3 + 2y_1y_2y_3, y_3^2y_1 + 2y_1^2y_2, y_3^2y_2 + 2y_2^2y_1)$$

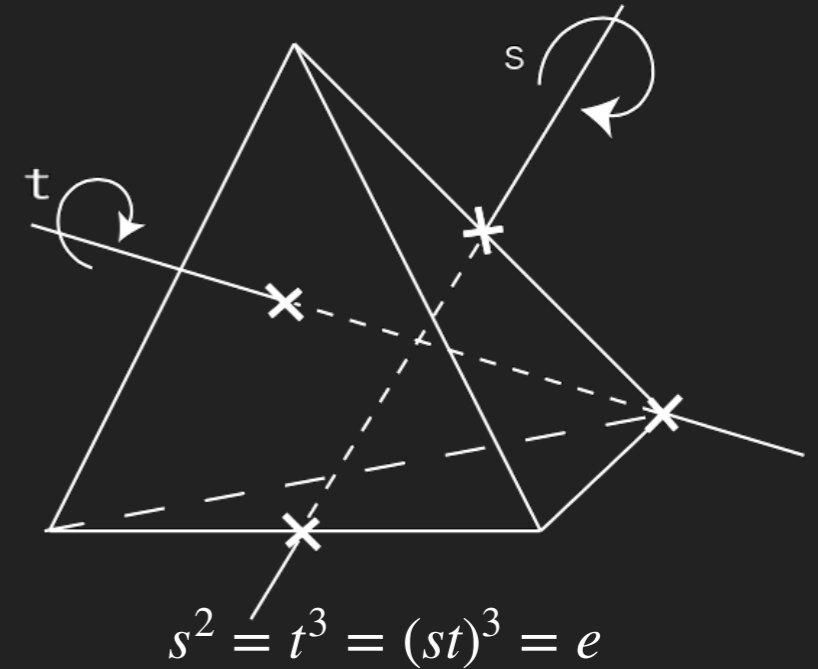
$$Y_{1''}^{(4)} = y_2^2 + 2y_1y_3 \qquad Y_{3c}^{(6)} = (y_2^3 + 2y_1y_2y_3, y_2^2y_3 + 2y_3^2y_1, y_2^2y_1 + 2y_1^2y_3)$$

NEUTRINO MASSES VIA \mathcal{A}_4 SYMMETRY

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$$\begin{aligned}
 C_1 &: \{e\}, & h &= 1, \\
 C_3 &: \{s, tst^2, t^2st\}, & h &= 2, \\
 C_4 &: \{t, ts, st, sts\}, & h &= 3, \\
 C_{4'} &: \{t^2, st^2, t^2s, tst\}, & h &= 3.
 \end{aligned}$$

	h	χ_1	$\chi_{1'}$	$\chi_{1''}$	χ_3
C_1	1	1	1	1	3
C_3	2	1	1	1	-1
C_4	3	1	ω	ω^2	0
$C_{4'}$	3	1	ω^2	ω	0



$$\underline{3} \times \underline{3} = \underline{1} + \underline{1'} + \underline{1''} + \underline{3} + \underline{3}.$$

$$\underline{1} = a_1 a_2 + b_1 b_2 + c_1 c_2,$$

$$\underline{1'} = a_1 a_2 + \omega^2 b_1 b_2 + \omega c_1 c_2,$$

$$\underline{1''} = a_1 a_2 + \omega b_1 b_2 + \omega^2 c_1 c_2.$$

$$\begin{aligned}
 (\mathbf{A})_3 \times (\mathbf{B})_3 &= (\mathbf{A} \cdot \mathbf{B})_1 + (\mathbf{A} \cdot \Sigma \cdot \mathbf{B})_{1'} + (\mathbf{A} \cdot \Sigma^* \cdot \mathbf{B})_{1''} \\
 &+ \begin{pmatrix} \{A_y B_z\} \\ \{A_z B_x\} \\ \{A_x B_y\} \end{pmatrix}_3 + \begin{pmatrix} [A_y B_z] \\ [A_z B_x] \\ [A_x B_y] \end{pmatrix}_3.
 \end{aligned}$$

[arXiv:[1003.3552](https://arxiv.org/abs/1003.3552)]

NEUTRINO MASSES VIA \mathcal{A}_4 SYMMETRY

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$$(\mathcal{A}_4, L) \quad \frac{1}{2} M N_{iR}^2 + f \bar{N}_{iR} (\nu_{iL} \eta^0 - l_{iL} \eta^+) + h_{ijk} \overline{(\nu_i, l_i)_L} l_{jR} \Phi_k + \text{H.c.},$$

$$(\nu_i, l_i)_L \sim (\underline{3}, 1),$$

$$l_{1R} \sim (\underline{1}, 1),$$

$$l_{2R} \sim (\underline{1}', 1),$$

$$l_{3R} \sim (\underline{1}'', 1),$$

$$N_{iR} \sim (\underline{3}, 0),$$

$$\Phi_i = (\phi_i^+, \phi_i^0) \sim (\underline{3}, 0),$$

$$\eta = (\eta^+, \eta^0) \sim (\underline{1}, -1).$$

$$U_L^\dagger \mathcal{M}_l U_R = \begin{bmatrix} \sqrt{3} h_1 v & 0 & 0 \\ 0 & \sqrt{3} h_2 v & 0 \\ 0 & 0 & \sqrt{3} h_3 v \end{bmatrix}$$

$$= \begin{bmatrix} m_e & 0 & 0 \\ 0 & m_\mu & 0 \\ 0 & 0 & m_\tau \end{bmatrix},$$

$$U_L = \frac{1}{\sqrt{3}} \begin{bmatrix} 1 & 1 & 1 \\ 1 & \omega & \omega^2 \\ 1 & \omega^2 & \omega \end{bmatrix}, \quad U_R = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

- * bimaximal and tribimaximal
- * Max $\sin(\theta_{23})$; $\sin(\theta_{13}) = 0$;
- * Degenerate neutrino masses
- * Soft breaking \rightarrow small mass splittings, $\sin(\theta_{13}) \neq 0$

$$\mathcal{M}_\nu = \frac{f^2 u^2}{M} U_L^T U_L = \frac{f^2 u^2}{M} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{bmatrix}$$

E. Ma, G. Rajasekaran, [PRD\(2001\)](#)
K.S. Babu, E. Ma, J.W.F. Valle, [PLB\(2002\)](#)

A MODULAR A_4 SYMMETRIC SCOTOGENIC MODEL

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	Fermions			Bosons	
	$(\bar{L}_{Le}, \bar{L}_{L\mu}, \bar{L}_{L\tau})$	$(e_{Re}, e_{R\mu}, e_{R\tau})$	N_R	H	η^*
$SU(2)_L$	2	1	1	2	2
$U(1)_Y$	$\frac{1}{2}$	-1	0	$\frac{1}{2}$	$-\frac{1}{2}$
A_4	$1, 1', 1''$	$1, 1'', 1'$	3	1	1
$-k$	0	0	-1	0	-3

	Couplings		
	$Y_1^{(6)}$	$Y_3^{(2)}$	$Y_3^{(4)}$
A_4	1	3	3
$-k$	6	2	4

$$\begin{aligned}
 -\mathcal{L}_{Lepton} = & \sum_{\ell=e,\mu,\tau} y_\ell \bar{L}_{L\ell} H e_{R\ell} \\
 & + \alpha_\nu \bar{L}_{Le} (Y_3^{(4)} \otimes N_R)_1 \tilde{\eta} + \beta_\nu \bar{L}_{L\mu} (Y_3^{(4)} \otimes N_R)_{1''} \tilde{\eta} \\
 & + \gamma_\nu \bar{L}_{L\tau} (Y_3^{(4)} \otimes N_R)_{1'} \tilde{\eta} \\
 & + M_1 (Y_3^{(2)} \otimes \bar{N}_R^C \otimes N_R) + \text{h.c.},
 \end{aligned}$$

$$Y_1^{(6)} = y_1^3 + y_2^3 + y_3^3 - 3y_1 y_2 y_3, \quad \rightarrow Y_1^{(6)} (H^\dagger \eta)^2$$

MASS MATRICES

$$M_N = \frac{M_1}{3} \begin{bmatrix} 2y_1 & -y_3 & -y_2 \\ -y_3 & 2y_2 & -y_1 \\ -y_2 & -y_{3,1} & 2y_3 \end{bmatrix}$$

$$Y_3^{(2)} = (y_1, y_2, y_3)$$

$$y_\eta = \begin{bmatrix} \alpha_\nu & 0 & 0 \\ 0 & \beta_\nu & 0 \\ 0 & 0 & \gamma_\nu \end{bmatrix} \begin{bmatrix} y'_1 & y'_3 & y'_2 \\ y'_3 & y'_2 & y'_1 \\ y'_2 & y'_1 & y'_3 \end{bmatrix} \quad Y_3^{(4)} \equiv (y'_1, y'_2, y'_3) = \begin{bmatrix} y_1^2 - y_2 y_3 \\ y_3^2 - y_1 y_2 \\ y_2^2 - y_1 y_3 \end{bmatrix}$$

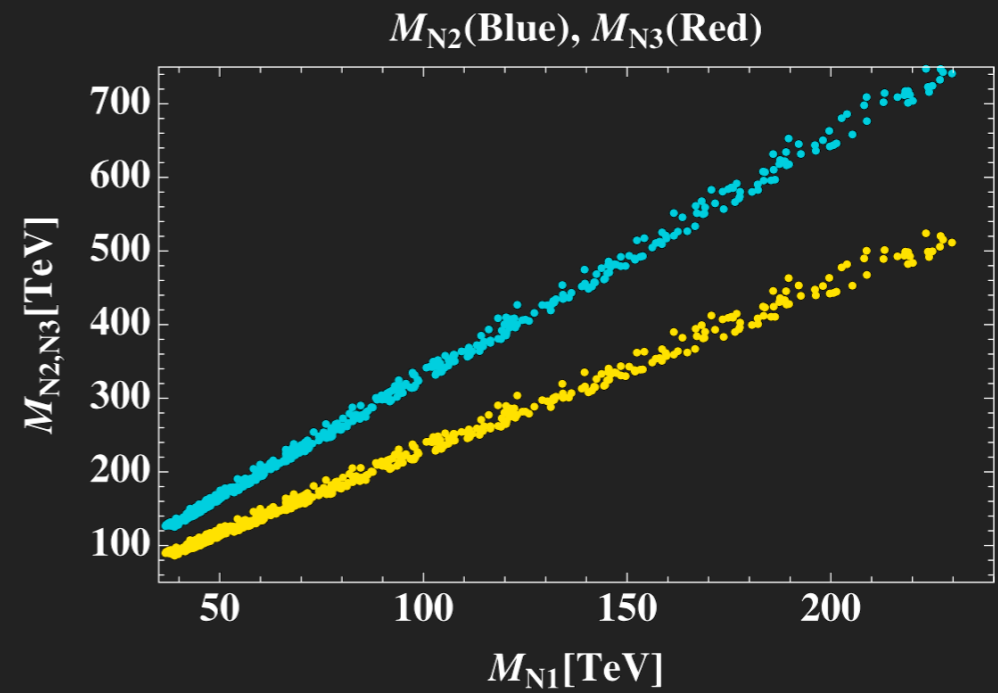
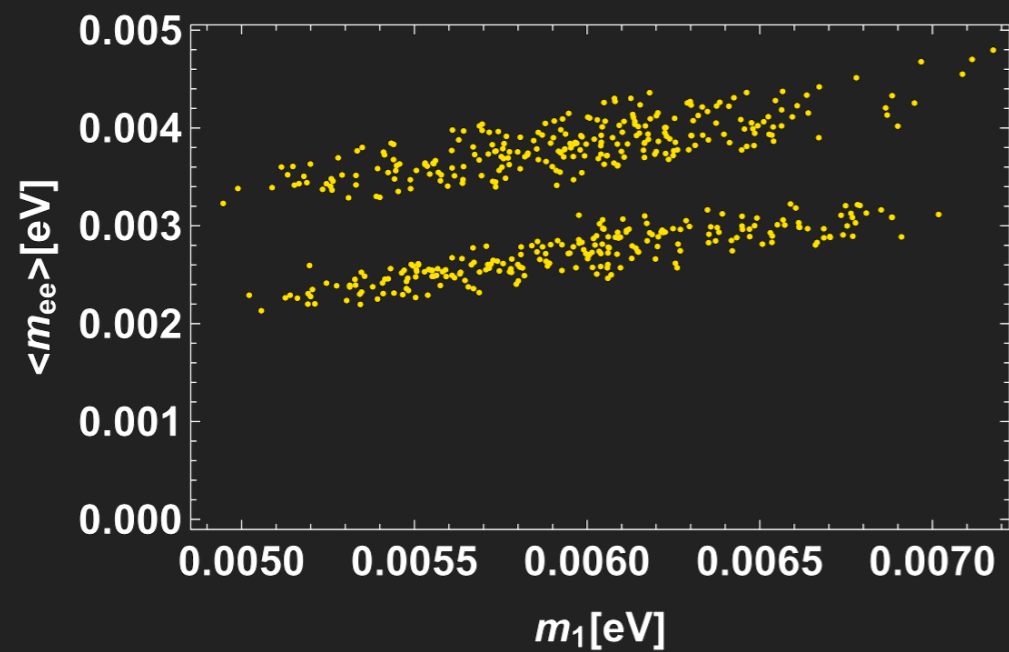
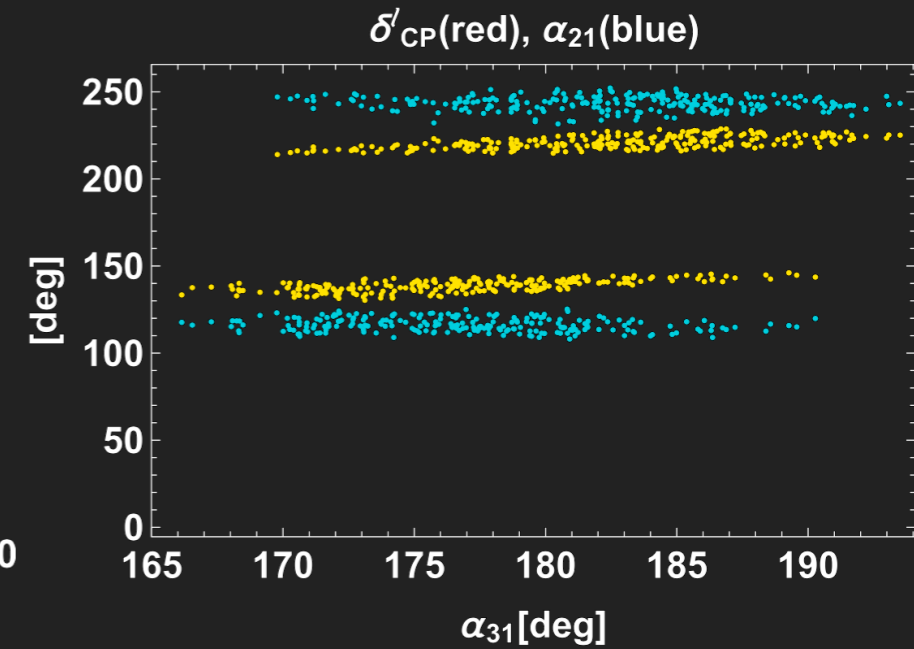
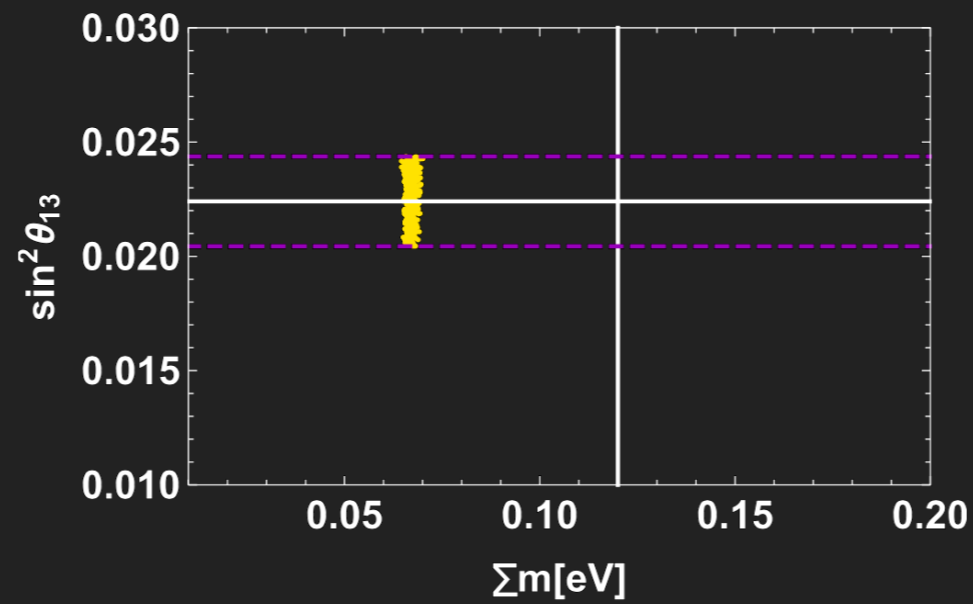
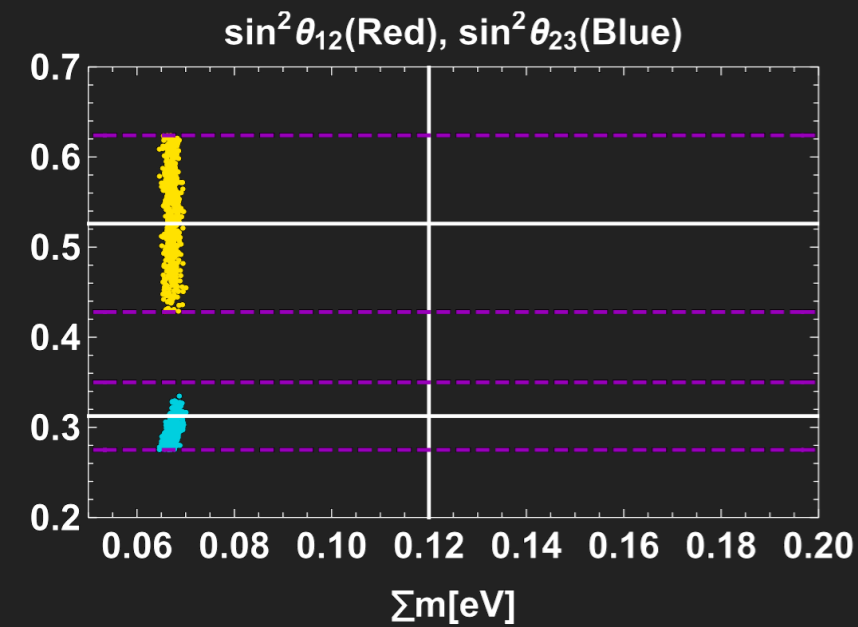
$$m_{\nu ij} \approx \sum_{\alpha=1-3} \frac{Y_{\eta i \alpha} D_{N \alpha} Y_{\eta \alpha j}^T}{(4\pi)^2} \times \left(\frac{m_R^2}{m_R^2 - D_{N \alpha}^2} \ln \left[\frac{m_R^2}{D_{N \alpha}^2} \right] - \frac{m_I^2}{m_I^2 - D_{N \alpha}^2} \ln \left[\frac{m_I^2}{D_{N \alpha}^2} \right] \right)$$

$$y_1(\tau) = \frac{i}{2\pi} \left(\frac{\eta'(\tau/3)}{\eta(\tau/3)} + \frac{\eta'((\tau+1)/3)}{\eta((\tau+1)/3)} + \frac{\eta'((\tau+2)/3)}{\eta((\tau+2)/3)} - \frac{27\eta'(3\tau)}{\eta(3\tau)} \right)$$

$$y_2(\tau) = \frac{-i}{\pi} \left(\frac{\eta'(\tau/3)}{\eta(\tau/3)} + \omega^2 \frac{\eta'((\tau+1)/3)}{\eta((\tau+1)/3)} + \omega \frac{\eta'((\tau+2)/3)}{\eta((\tau+2)/3)} \right)$$

$$y_3(\tau) = \frac{-i}{\pi} \left(\frac{\eta'(\tau/3)}{\eta(\tau/3)} + \omega \frac{\eta'((\tau+1)/3)}{\eta((\tau+1)/3)} + \omega^2 \frac{\eta'((\tau+2)/3)}{\eta((\tau+2)/3)} \right)$$

RESULT PLOTS



NON-HOLOMORPHIC MODULAR \mathcal{A}_4 SYMMETRIC SCOTOGENIC MODEL¹⁵

Fields	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	\mathcal{A}_4	k
\overline{L}_L	1	2	$\frac{1}{2}$	3	-1
e_R	1	1	-1	3	1
N_R	1	1	0	3	0
H	1	2	$\frac{1}{2}$	1	0
η	1	2	$-\frac{1}{2}$	1	1

$$\begin{aligned}
 \mathcal{L}_{Y_\ell} = & M_0 [\overline{N}_R^c N_R]_1 + M_3 Y_3^{(0)} [\overline{N}_R^c N_R]_{3_s} \\
 & + y_\ell [\overline{L}_L e_R]_1 H + y_{\ell_3} Y_3^{(0)} [\overline{L}_L e_R]_{3_s} H \\
 & + y'_{\ell_3} Y_3^{(0)} [\overline{L}_L e_R]_{3_a} H + y_D [\overline{L}_L N_R]_1 \eta \\
 & + y_{D_3} Y_3^{(0)} [\overline{L}_L N_R]_{3_s} \eta + y'_{D_3} Y_3^{(0)} [\overline{L}_L N_R]_{3_a} \eta \\
 & + \text{h.c.},
 \end{aligned}$$

$$\begin{aligned}
 V = & \mu_\eta^2 \eta^\dagger \eta + \tilde{\lambda} Y_1^{(-2)} (H \eta)^2 + \lambda_{H\eta} (H^\dagger H) (\eta^\dagger \eta) \\
 & + \lambda'_{H\eta} |H \eta|^2 + \text{other trivial terms},
 \end{aligned}$$

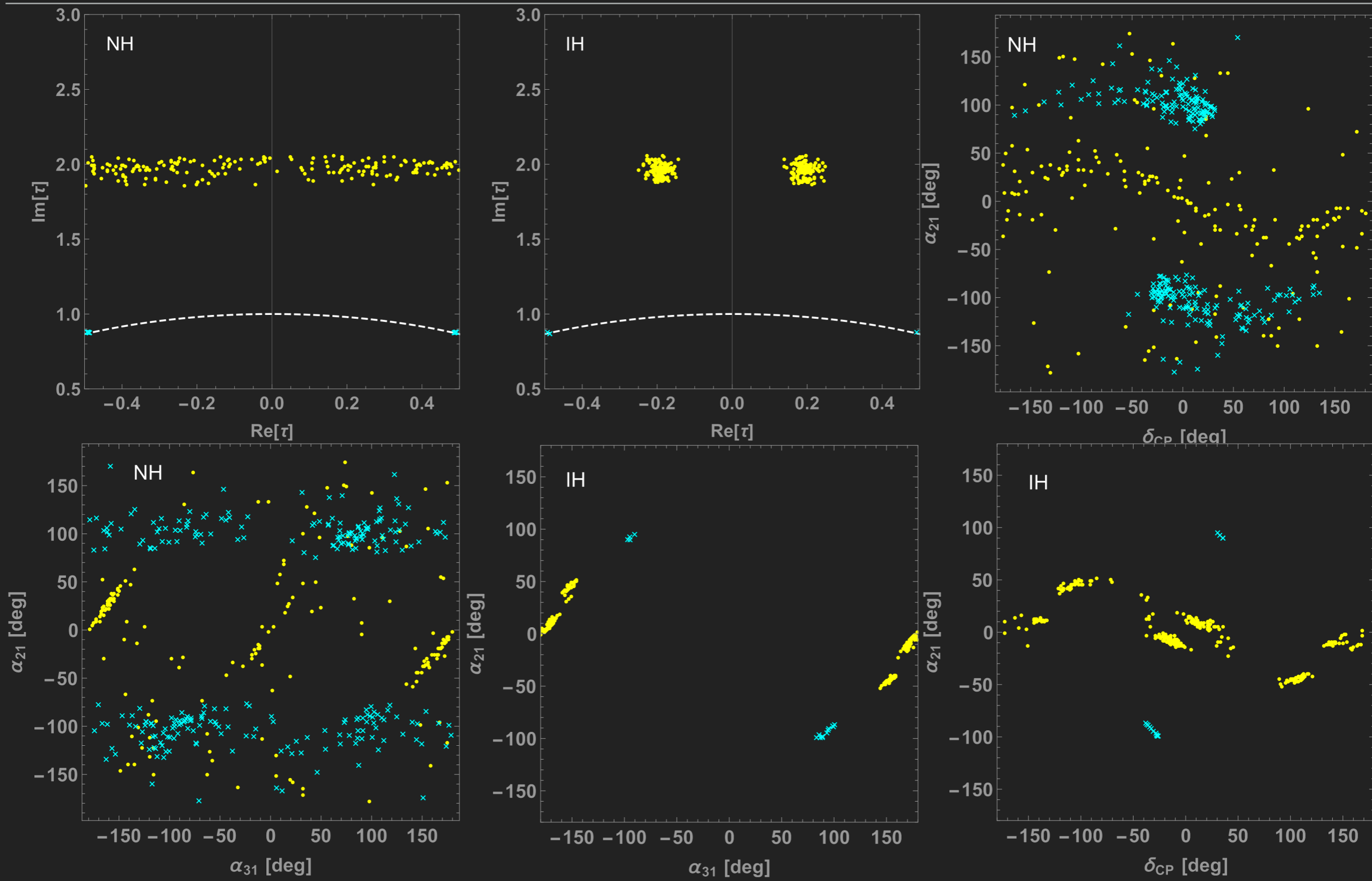
$$M_N = M_3 \begin{bmatrix} \tilde{M}_0 + 2y_1 & -y_3 & -y_2 \\ -y_3 & 2y_2 & \tilde{M}_0 - y_1 \\ -y_2 & \tilde{M}_0 - y_1 & 2y_3 \end{bmatrix}$$

$$y_\eta = \begin{bmatrix} y_D + 2y_{D_3}y_1 & -(y_{D_3} - y'_{D_3})y_3 & -(y_{D_3} + y'_{D_3})y_2 \\ -(y_{D_3} + y'_{D_3})y_3 & 2y_{D_3}y_2 & y_D - (y_{D_3} - y'_{D_3})y_1 \\ -(y_{D_3} - y'_{D_3})y_2 & y_D - (y_{D_3} + y'_{D_3})y_1 & 2y_{D_3}y_3 \end{bmatrix} (\overline{L}_L y_\eta N_R) \eta.$$

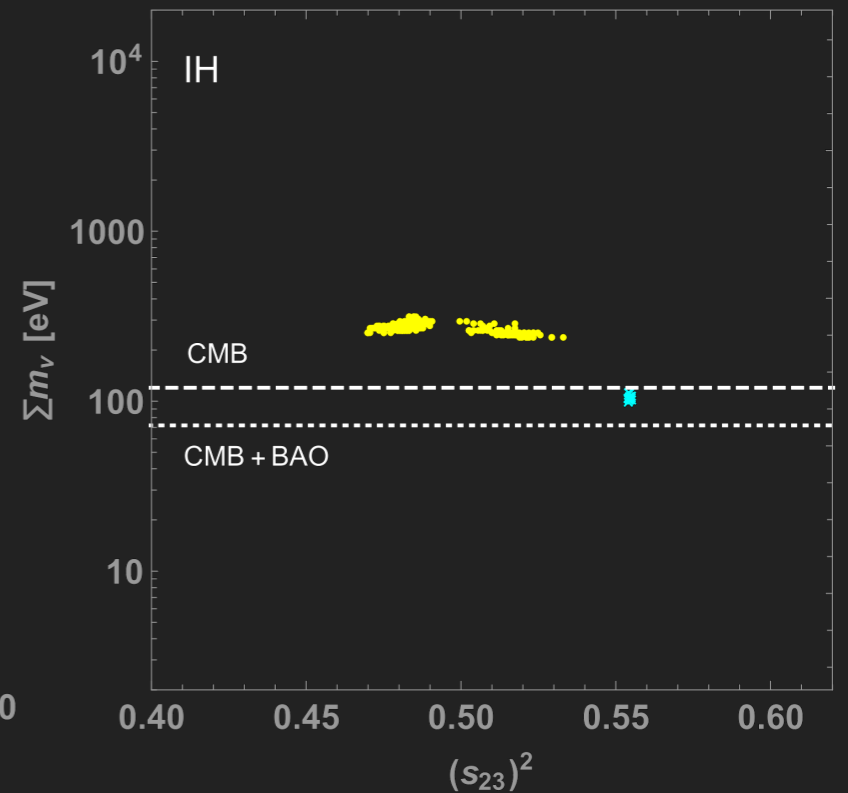
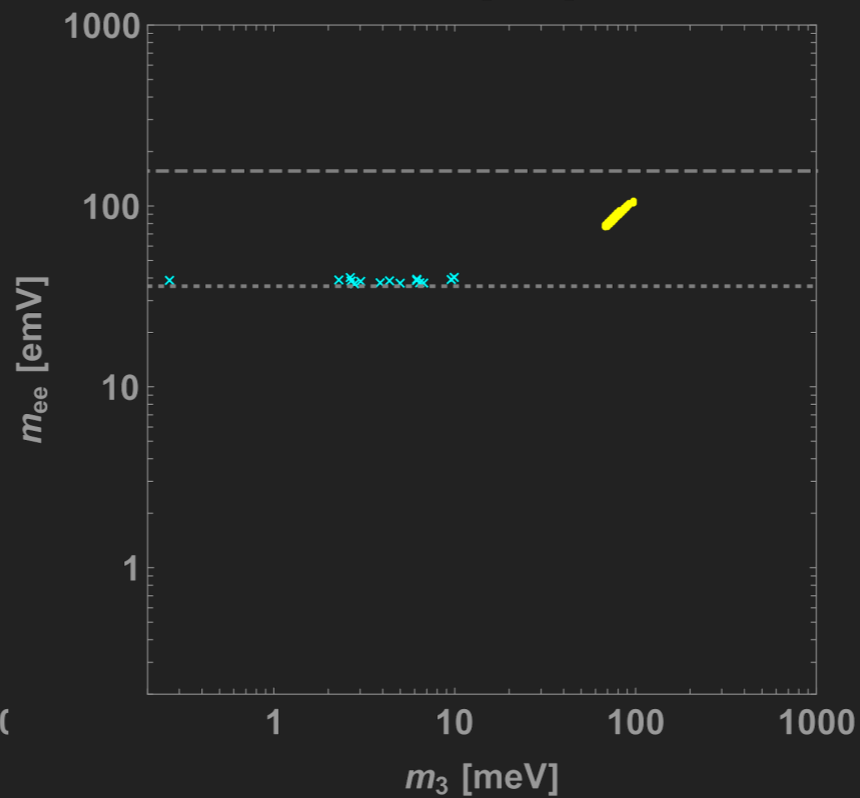
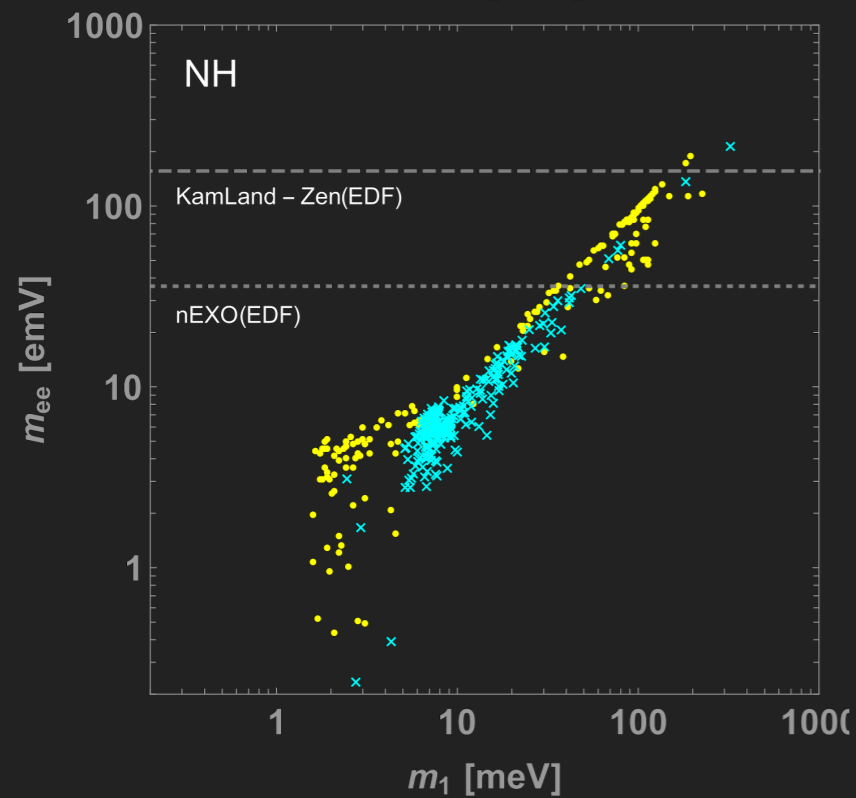
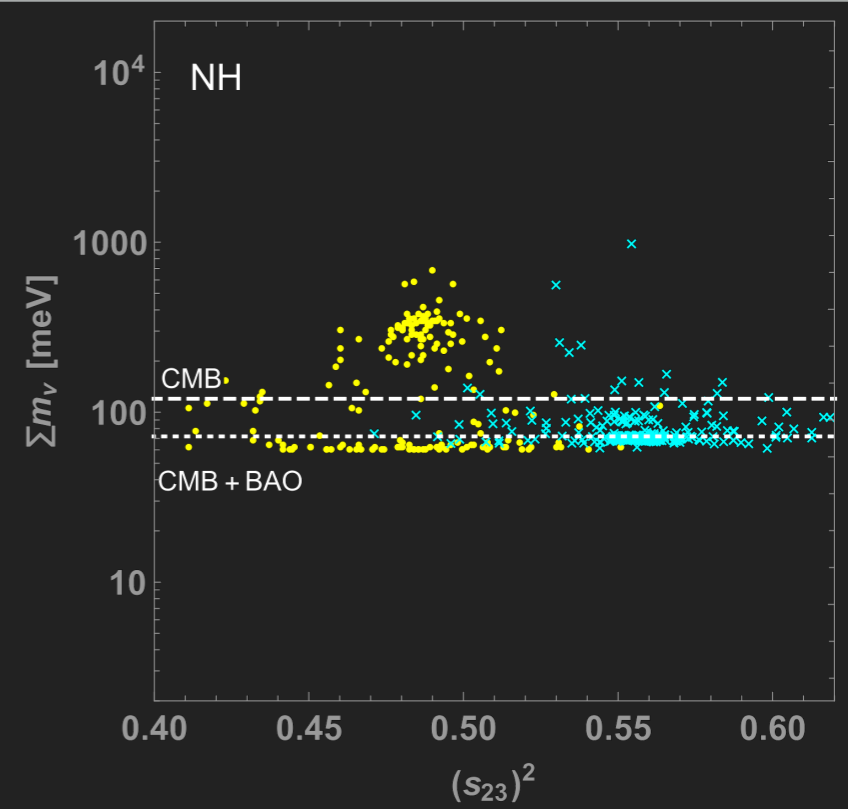
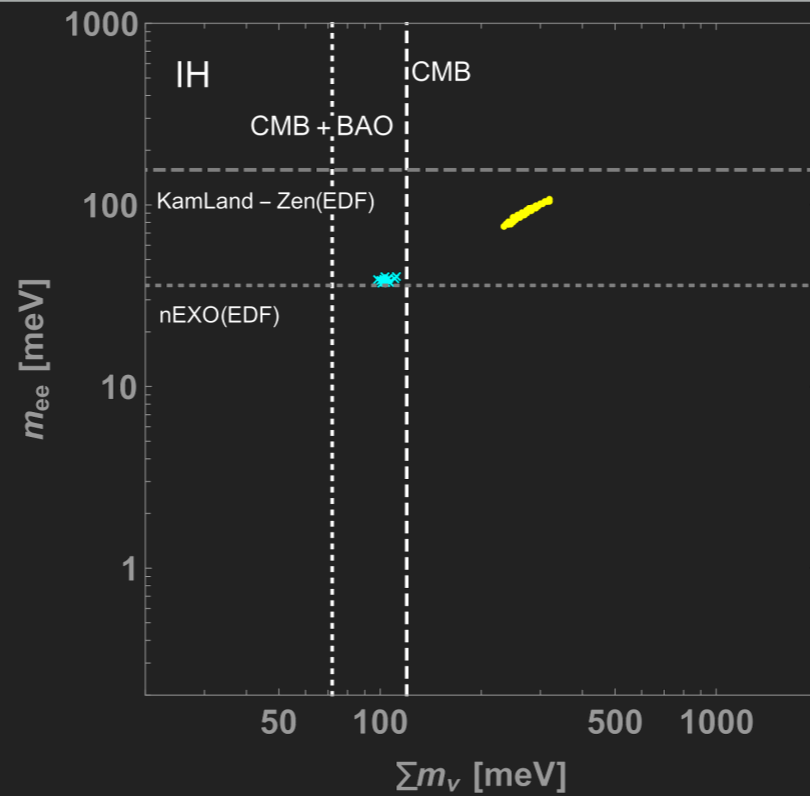
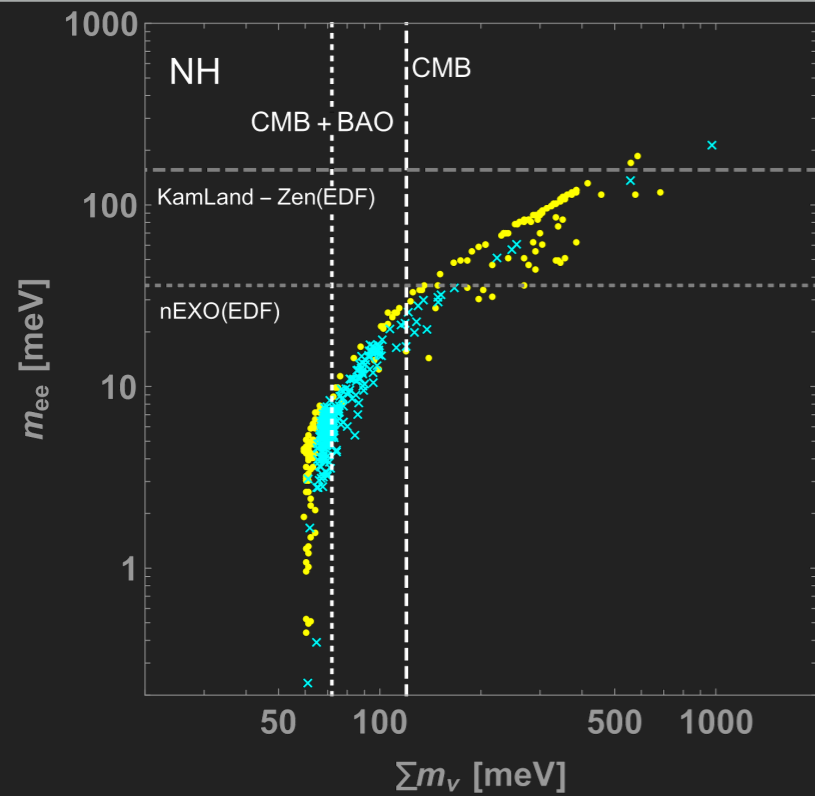
$$\frac{m_\ell}{v} = \begin{bmatrix} y_\ell + 2y_{\ell_3}y_1 & -(y_{\ell_3} - y'_{\ell_3})y_3 & -(y_{\ell_3} + y'_{\ell_3})y_2 \\ -(y_{\ell_3} + y'_{\ell_3})y_3 & 2y_{\ell_3}y_2 & y_\ell - (y_{\ell_3} - y'_{\ell_3})y_1 \\ -(y_{\ell_3} - y'_{\ell_3})y_2 & y_\ell - (y_{\ell_3} + y'_{\ell_3})y_1 & 2y_{\ell_3}y_3 \end{bmatrix}$$

$$m_{\nu_{ij}} \approx \sum_{\alpha=1-3} \frac{Y_{\eta i \alpha} D_{N \alpha} Y_{\eta \alpha j}^T}{(4\pi)^2} \times \left(\frac{m_R^2}{m_R^2 - D_{N \alpha}^2} \ln \left[\frac{m_R^2}{D_{N \alpha}^2} \right] - \frac{m_I^2}{m_I^2 - D_{N \alpha}^2} \ln \left[\frac{m_I^2}{D_{N \alpha}^2} \right] \right)$$

RESULT PLOTS



RESULT PLOTS



SUMMARY

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- ▶ Naturally small neutrino masses (radiative model)
- ▶ Minimal field content (\mathcal{A}_4)
- ▶ Candidate for dark matter (scotogenic = darkness)
- ▶ Prediction for PMNS matrix/neutrino mixing (modular)
- ▶ Correlation between Majorana and Dirac phases
- ▶ Model is testable by current and future neutrino experiments (KATRIN, $0\nu\beta\beta$) and cosmological observations (DESI2024)

The end

THANK YOU!