#### Listening to Dark Sectors with Pulsar Timing Arrays

(work with Y. Cui, Y.-D. Tsai and Y. Tsai)

#### **Amitayus Banik**

Chungbuk National University

 $5^{\rm th}$  International Joint Workshop on the Standard Model and Beyond 2024 &  $3^{\rm rd}$  Gordon Godfrey Workshop on Astroparticle Physics

10<sup>th</sup> December 2024



• Measurements from Pulsar Timing Arrays (NANOGrav, EPTA, PPTA, CPTA, etc.) provide strong evidence for a stochastic GW background.

- Measurements from Pulsar Timing Arrays (NANOGrav, EPTA, PPTA, CPTA, etc.) provide strong evidence for a stochastic GW background.
- **Possible origins:** merging population of supermassive black holes (astrophysical) NANOGrav Collab. (2021, 2023), ...

- Measurements from **PTAs** (NANOGrav, EPTA, PPTA, CPTA) provide **strong** evidence for a stochastic GW background.
- **Possible origins:** merging population of supermassive black holes (astrophysical), **but** beyond-the-SM sources (phase transition, collapsing domain walls, cosmic string network, ...) also **allowed**.

..., NANOGrav Collab (2023), T. Bringmann et al. (2023), Y. Bai et al. (2023), J. Ellis et al. (2023), ...

- Measurements from **PTAs** (NANOGrav, EPTA, PPTA, CPTA) provide **strong** evidence for a stochastic GW background.
- Possible origins: merging population of supermassive black holes (astrophysical), but beyond-the-SM sources (phase transition, collapsing domain walls, cosmic string network, ...) also allowed.
   ..., NANOGray Collab (2023), T. Bringmann et al. (2023), Y. Bai et al. (2023), J. Ellis et al. (2023), ...
- Here: focus on first-order phase transitions occurring in a dark sector. ..., C. Han et al. (2023), S.-P. Li & K.-P. Xie (2023), ...

- Measurements from **PTAs** (NANOGrav, EPTA, PPTA, CPTA) provide **strong** evidence for a stochastic GW background.
- Possible origins: merging population of supermassive black holes (astrophysical), but beyond-the-SM sources (phase transition, collapsing domain walls, cosmic string network, ...) also allowed.
   ..., NANOGrav Collab (2023), T. Bringmann et al. (2023), Y. Bai et al. (2023), J. Ellis et al. (2023), ...
- Here: focus on first-order phase transitions occurring in a dark sector. ..., C. Han et al. (2023), S.-P. Li & K.-P. Xie (2023), ...
- In particular: connect observed PTA spectrum to microphysics model parameters as opposed to macroscopic phase transition (PT) parameters
   T. Bringmann et al. (2023), A. Addazi et al (2023), M. Winkler & K. Freese (2024), ...

• Need a scalar driving the PT

• Need a scalar driving the PT, additional bosonic degrees of freedom to **augment** PT.

- Need a scalar driving the PT, additional bosonic degrees of freedom to **augment** PT.
- Minimal setup: Gauged dark U(1) with complex scalar  $\Phi$

- Need a scalar driving the PT, additional bosonic degrees of freedom to **augment** PT.
- Minimal setup: Gauged dark U(1) with complex scalar  $\Phi$

$$\mathcal{L}_{\rm DSPT} \supset (D^{\mu}\Phi)^{\dagger} (D_{\mu}\Phi) - \mu^2 \, \Phi^{\dagger}\Phi + \frac{\lambda}{2} (\Phi^{\dagger}\Phi)^2$$

where  $D_{\mu}\Phi = \partial_{\mu}\Phi + ig_D A'_{\mu} \Phi$  and  $\mu^2, \lambda > 0$ 

- Need a scalar driving the PT, additional bosonic degrees of freedom to augment PT.
- Minimal setup: Gauged dark U(1) with complex scalar  $\Phi$

$$\mathcal{L}_{\text{DSPT}} \supset (D^{\mu}\Phi)^{\dagger} (D_{\mu}\Phi) - \mu^2 \, \Phi^{\dagger}\Phi + \frac{\lambda}{2} (\Phi^{\dagger}\Phi)^2$$

where  $D_{\mu}\Phi = \partial_{\mu}\Phi + ig_D A'_{\mu} \Phi$  and  $\mu^2, \lambda > 0$ 

• Symmetry breaking,

$$\langle\Phi
angle=rac{v_0+\phi}{\sqrt{2}}\,,\quad {
m with}\quad v_0=\sqrt{rac{\mu^2}{\lambda}}$$

- Need a scalar driving the PT, additional bosonic degrees of freedom to augment PT.
- Minimal setup: Gauged dark U(1) with complex scalar  $\Phi$

$$\mathcal{L}_{\text{DSPT}} \supset (D^{\mu}\Phi)^{\dagger} (D_{\mu}\Phi) - \mu^2 \Phi^{\dagger}\Phi + \frac{\lambda}{2} (\Phi^{\dagger}\Phi)^2$$

where  $D_{\mu}\Phi = \partial_{\mu}\Phi + ig_D A'_{\mu} \Phi$  and  $\mu^2, \lambda > 0$ 

• Symmetry breaking,

$$\langle \Phi 
angle = rac{v_0 + \phi}{\sqrt{2}} \,, \quad {\rm with} \quad v_0 = \sqrt{rac{\mu^2}{\lambda}}$$

 Useful to also consider limit of SU(N), but need maximal breaking (avoid remnant massless degrees of freedom)

- Need a scalar driving the PT, additional bosonic degrees of freedom to augment PT.
- Minimal setup: Gauged dark U(1) with complex scalar  $\Phi$

$$\mathcal{L}_{\text{DSPT}} \supset (D^{\mu}\Phi)^{\dagger} (D_{\mu}\Phi) - \mu^2 \Phi^{\dagger}\Phi + \frac{\lambda}{2} (\Phi^{\dagger}\Phi)^2$$

where  $D_{\mu}\Phi = \partial_{\mu}\Phi + ig_D A'_{\mu} \Phi$  and  $\mu^2, \lambda > 0$ 

• Symmetry breaking,

$$\langle \Phi 
angle = rac{v_0 + \phi}{\sqrt{2}} \,, \quad {
m with} \quad v_0 = \sqrt{rac{\mu^2}{\lambda}}$$

- Useful to also consider limit of SU(N), but need maximal breaking (avoid remnant massless degrees of freedom)
- Can also consider a Yukawa term:  $y_D \bar{\Psi} \Phi \chi$ , with Dirac fermion  $\Psi$  and one singlet fermion  $\chi$ .

• Promote the vev  $v_0$  to background field  $\varphi \equiv \varphi(T)$ ; temperature evolution governed by finite-temperature effective potential

e.g. M. Laine and A. Vuorinen (2016)

$$V_{1-L}(\varphi,T) = V_0(\varphi) + V_{CW}(\varphi) + V_T(\varphi,T)$$

• Promote the vev  $v_0$  to background field  $\varphi \equiv \varphi(T)$ ; temperature evolution governed by **finite-temperature effective potential** 

e.g. M. Laine and A. Vuorinen (2016)

$$V_{1-L}(\varphi,T) = V_0(\varphi) + V_{CW}(\varphi) + V_T(\varphi,T)$$

• At high T,  $\varphi = 0$  and symmetry intact.

• Promote the vev  $v_0$  to background field  $\varphi \equiv \varphi(T)$ ; temperature evolution governed by **finite-temperature effective potential** e.g. M. Laire and A. Vuorinen (2016)

$$V_{1-L}(\varphi,T) = V_0(\varphi) + V_{CW}(\varphi) + V_T(\varphi,T)$$

- At high T,  $\varphi = 0$  and symmetry intact.
- At low T, end up at the true vacuum  $\varphi = v_0$  where the symmetry is broken.



• Promote the vev  $v_0$  to background field  $\varphi \equiv \varphi(T)$ ; temperature evolution governed by **finite-temperature effective potential** e.g. M. Laire and A. Vuorinen (2016)

$$V_{1-L}(\varphi,T) = V_0(\varphi) + V_{CW}(\varphi) + V_T(\varphi,T)$$

- At high T,  $\varphi = 0$  and symmetry intact.
- As low T,  $\varphi$  ends up at the true vacuum  $\varphi = v_0$  where the symmetry is broken.
- If a barrier forms (at critical temp. T<sub>c</sub>) separating the two vacua, the PT is said to be first-order (FOPT).



• FOPTs proceed through nucleation of **bubbles** of the true vacuum;

 FOPTs proceed through nucleation of **bubbles** of the true vacuum; rate per Hubble volume given by

S.R. Coleman (1977), C. G. Callan, Jr. and S. R. Coleman (1977), A. D. Linde (1981, 1983)

$$\Gamma(T) = T^4 \left(\frac{S_3}{2\pi T}\right)^{\frac{3}{2}} \exp\left(-\frac{S_3}{T}\right)$$



• FOPTs proceed through nucleation of **bubbles** of the true vacuum; rate per Hubble volume given by

S.R. Coleman (1977), C. G. Callan, Jr. and S. R. Coleman (1977), A. D. Linde (1981, 1983)

$$\Gamma(T) = T^4 \left(\frac{S_3}{2\pi T}\right)^{\frac{3}{2}} \exp\left(-\frac{S_3}{T}\right)$$



• Bubble nucleation occurs at  $T_*$  determined by

$$\Gamma(T_*) \sim H^4(T_*) \,,$$

• FOPTs proceed through nucleation of **bubbles** of the true vacuum; rate per Hubble volume given by

S.R. Coleman (1977), C. G. Callan, Jr. and S. R. Coleman (1977), A. D. Linde (1981, 1983)

$$\Gamma(T) = T^4 \left(\frac{S_3}{2\pi T}\right)^{\frac{3}{2}} \exp\left(-\frac{S_3}{T}\right)$$



• Bubble nucleation occurs at  $T_*$  determined by

$$\Gamma(T_*) \sim H^4(T_*)$$
, where  $H^2(T) = \frac{8\pi G_N}{3} \rho_R(T) \equiv \frac{8\pi G_N}{3} \left(\frac{\pi^2}{30}g_{*, \text{ tot}} T^4\right)$ 

#### Phase Transition Parameters

Obtain

$$S_3 = 4\pi \int_0^\infty dr \, r^2 \left[ \frac{1}{2} \left( \frac{d\varphi_b}{dr} \right)^2 + V_{1-\mathrm{L}}(\varphi_b, T) \right] \,,$$

the 3-D Euclidean action for bounce configuration, using FindBounce. v.  $_{\hbox{\scriptsize V. Guada, et. al. (2020)}}$ 

#### Phase Transition Parameters

Obtain

$$S_3 = 4\pi \int_0^\infty dr \, r^2 \left[ \frac{1}{2} \left( \frac{d\varphi_b}{dr} \right)^2 + V_{1-L}(\varphi_b, T) \right] \,,$$

the 3-D Euclidean action for bounce configuration, using FindBounce. v.  $_{\hbox{\scriptsize Guada, et. al. (2020)}}$ 

• PT characterized by its inverse duration  $\beta$  and strength  $\alpha_*$  (~ latent heat)

$$\frac{\beta}{H_{\star}} \equiv T_{\star} \left[ \frac{d}{dT} \left( \frac{S_3}{T} \right) \right] \Big|_{T_{\star}}, \quad \alpha_{\star} \equiv \frac{1}{\rho_R(T_{\star})} \left( \Delta V \Big|_{T_{\star}} - T_{\star} \left. \frac{\partial \Delta V}{\partial T} \Big|_{T_{\star}} \right),$$

#### Phase Transition Parameters

Obtain

$$S_3 = 4\pi \int_0^\infty dr \, r^2 \left[ \frac{1}{2} \left( \frac{d\varphi_b}{dr} \right)^2 + V_{1-\mathbf{L}}(\varphi_b, T) \right] \,,$$

the 3-D Euclidean action for bounce configuration, using FindBounce. v.  $_{\text{Guada, et. al. (2020)}}$ 

• PT characterized by its inverse duration  $\beta$  and strength  $\alpha_*$  ( $\sim$  latent heat)

$$\frac{\beta}{H_{\star}} \equiv T_{\star} \left[ \frac{d}{dT} \left( \frac{S_3}{T} \right) \right] \Big|_{T_{\star}}, \quad \alpha_{\star} \equiv \frac{1}{\rho_R(T_{\star})} \left( \Delta V \Big|_{T_{\star}} - T_{\star} \left. \frac{\partial \Delta V}{\partial T} \Big|_{T_{\star}} \right),$$

• Model dependence comes from  $V_{1-L}(\varphi_b, T) \longrightarrow \mathsf{PT}$  parameters end up depending on the model parameters  $[\mu, \lambda, g_D, (y_D, N)]$ 

## U(1) Example: $T_*$

Work with pure U(1) group,  $y_D = 0$ . Fix  $\lambda = 0.05$ , study trend of  $T_*$  in  $(\mu, g_D)$  plane.

## U(1) Example: $T_*$

Work with pure U(1) group,  $y_D = 0$ . Fix  $\lambda = 0.05$ , study trend of  $T_*$  in  $(\mu, g_D)$  plane.



## U(1) Example: $T_*$

Work with pure U(1) group,  $y_D = 0$ . Fix  $\lambda = 0.05$ , study trend of  $T_*$  in  $(\mu, g_D)$  plane.



#### $\mu$ fixes the **relative scale** of $T_*!$

## U(1) Example: $\beta/H_*$ and $\alpha_*$

Set  $\mu = 2$  MeV and  $y_D = 0$ . Study trend of  $\beta/H_*$  and  $\alpha_*$  in  $(g_D, \lambda)$  plane.

## U(1) Example: $\beta/H_*$ and $\alpha_*$

Set  $\mu = 2$  MeV and  $y_D = 0$ . Study trend of  $\beta/H_*$  and  $\alpha_*$  in  $(g_D, \lambda)$  plane.



## U(1) Example: $\beta/H_*$ and $\alpha_*$

Set  $\mu = 2$  MeV and  $y_D = 0$ . Study trend of  $\beta/H_*$  and  $\alpha_*$  in  $(g_D, \lambda)$  plane.



Larger couplings: PT lasts **longer**; increasing  $g_D$  (for fixed  $\lambda$ ): PT is **stronger**.

Sources: bubble collisions, sound waves, and turbulence
 [A. Kosowsky et al. (1992, 1993), S. J. Huber and T. Konstandin (2008), D. J. Weir (2016)], [Hindmarsh et al. (2013, 2015, 2017)],
 [Kosowsky et al. (2002), Dolgov et al. (2002), Caprini et al. (2009)]



Sources: bubble collisions, sound waves, and turbulence
 [A. Kosowsky et al. (1992, 1993), S. J. Huber and T. Konstandin (2008), D. J. Weir (2016)], [Hindmarsh et al. (2013, 2015, 2017)],
 [Kosowsky et al. (2002), Dolgov et al. (2002), Caprini et al. (2009)]



• Depending on microphysics, different contributions dominate

• Sources: bubble collisions, sound waves, and turbulence [A. Kosowsky et al. (1992, 1993), S. J. Huber and T. Konstandin (2008), D. J. Weir (2016)], [Hindmarsh et al. (2013, 2015, 2017)], [Kosowsky et al. (2002), Dolgov et al. (2002), Caprini et al. (2009)]



 Depending on microphysics, different contributions dominate → in our case, sound waves dominant due to strong interactions of scalar with plasma

Sources: bubble collisions, sound waves, and turbulence
 [A. Kosowsky et al. (1992, 1993), S. J. Huber and T. Konstandin (2008), D. J. Weir (2016)], [Hindmarsh et al. (2013, 2015, 2017)],
 [Kosowsky et al. (2002), Dolgov et al. (2002), Caprini et al. (2009)]



• Depending on microphysics, different contributions **dominate**  $\rightarrow$  in our case, sound waves dominant due to strong interactions of scalar with plasma (sub-leading contribution from turbulence) Rev. by Caprini et al. (2015, 2018, 2019)

$$\Omega_{\rm GW}^{\rm sw} h^2(f) = \Omega_{\rm GW}^{\rm sw, peak} h^2 \left(\frac{f}{f_{\rm peak}^{\rm sw}}\right)^3 \left(\frac{7}{4+3\left(\frac{f}{f_{\rm peak}^{\rm sw}}\right)^2}\right)^{\frac{7}{2}}$$

Sources: bubble collisions, sound waves, and turbulence
 [A. Kosowsky et al. (1992, 1993), S. J. Huber and T. Konstandin (2008), D. J. Weir (2016)], [Hindmarsh et al. (2013, 2015, 2017)],
 [Kosowsky et al. (2002), Dolgov et al. (2002), Caprini et al. (2009)]



• Depending on microphysics, different contributions **dominate**  $\rightarrow$  in our case, sound waves dominant due to strong interactions of scalar with plasma (sub-leading contribution from turbulence) Rev. by Caprini et al. (2015, 2018, 2019)

$$\Omega_{\rm GW}^{\rm sw} h^2(f) = \Omega_{\rm GW}^{\rm sw, peak} h^2 \left(\frac{f}{f_{\rm peak}^{\rm sw}}\right)^3 \left(\frac{7}{4+3\left(\frac{f}{f_{\rm peak}^{\rm sw}}\right)^2}\right)^{\frac{7}{2}}$$

Improvements: T. Ghosh et al. (2023); H.-K. Guo et al. (2024), also see his plenary talk tomorrow!

#### Model Parameters and the Spectrum

$$\begin{split} \Omega_{\rm GW}^{\rm sw,peak} h^2 &= 5.71 \times 10^{-8} \, v_w \, \left(\frac{10}{g_{*,\rm tot}}\right)^{\frac{1}{3}} \left(\frac{\beta/H_*}{100}\right)^{-1} \left(\frac{\kappa_{\rm sw} \, \alpha_*}{1+\alpha_*}\right)^2 \, \Upsilon \\ f_{\rm peak}^{\rm sw} &= \frac{1.3 \times 10^{-8} \, {\rm Hz}}{v_w} \left(\frac{\beta/H_*}{100}\right) \left(\frac{g_{*,\rm tot}}{10}\right)^{\frac{1}{6}} \left(\frac{T_*}{1 \, {\rm MeV}}\right) \end{split}$$

#### Model Parameters and the Spectrum

$$\begin{split} \Omega_{\rm GW}^{\rm sw,peak} h^2 &= 5.71 \times 10^{-8} \, v_w \, \left(\frac{10}{g_{*,\rm tot}}\right)^{\frac{1}{3}} \left(\frac{\beta/H_*}{100}\right)^{-1} \left(\frac{\kappa_{\rm sw} \, \alpha_*}{1+\alpha_*}\right)^2 \, \Upsilon \\ f_{\rm peak}^{\rm sw} &= \frac{1.3 \times 10^{-8} \, {\rm Hz}}{v_w} \left(\frac{\beta/H_*}{100}\right) \left(\frac{g_{*,\rm tot}}{10}\right)^{\frac{1}{6}} \left(\frac{T_*}{1 \, {\rm MeV}}\right) \end{split}$$



#### Model Parameters and the Spectrum

$$\Omega_{\rm GW}^{\rm sw,peak} h^2 = 5.71 \times 10^{-8} v_w \left(\frac{10}{g_{*,\rm tot}}\right)^{\frac{1}{3}} \left(\frac{\beta/H_*}{100}\right)^{-1} \left(\frac{\kappa_{\rm sw} \alpha_*}{1+\alpha_*}\right)^2 \Upsilon$$
$$f_{\rm peak}^{\rm sw} = \frac{1.3 \times 10^{-8} \,\mathrm{Hz}}{v_w} \left(\frac{\beta/H_*}{100}\right) \left(\frac{g_{*,\rm tot}}{10}\right)^{\frac{1}{6}} \left(\frac{T_*}{1\,\mathrm{MeV}}\right)$$



#### Fitting to the PTA Data

#### Fitting to the PTA Data

• To study the parameter space **consistent** with the PTA data, need mean values and  $1\sigma$  error bars.

#### Fitting to the PTA Data

- To study the parameter space **consistent** with the PTA data, need mean values and  $1\sigma$  error bars.
- Test compatibility using

$$\chi^2 = \sum_{i=1}^{41} \left( \frac{\log \Omega_{\mathrm{GW},i}^{\mathrm{model}} - \log \Omega_{\mathrm{GW},i}^{\mathrm{data}}}{\Delta \log \Omega_{\mathrm{GW},i}^{\mathrm{data}}} \right)^2$$

 $\Omega^{\text{model}}_{\text{GW},i} \equiv \text{model prediction},$  $\Omega^{\text{data}}_{\text{GW},i} \equiv \text{mean values of PTA data},$  $\Delta = \text{error bars of PTA data}$ 

• Need the degrees of freedom to determine confidence levels,

 $n_{\rm d.o.f.} = 41 - n_{\rm params}$ 



M. Winkler and K. Freese (2024)







• Extremely narrow parameter space for which a DS PT can account for the PTA data.



- Extremely narrow parameter space for which a DS PT can account for the PTA data.
- Expected to be a generic problem for most models.

#### Best Fit to the PTA Data

MODEL	$(\mu, \lambda, g_D, y_D)$	$(T_*, \alpha_*, \beta/H_*)$	$\chi^2$
U(1)	(2.4 MeV, 0.034, 1.2, 0)	(2.35 MeV, 0.96, 47)	36
U(1) w/ $\Psi$ and $\chi$	(2.7 MeV, 0.0575, 1.5, 0.8)	(2.5 MeV, 0.76, 38)	29
SU(2)	(2.0 MeV, 0.052, 1.14, 0)	(1.9 MeV, 0.79, 58)	30
SU(3)	(1.8 MeV, 0.029, 0.68, 0)	(1.8 MeV, 1.06, 70)	40



• With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)

• With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$ 

- With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$
- Physical dark Higgs  $\phi$  emerges as lightest stable particle  $\rightarrow$  can be a **dark matter candidate**

- With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$
- Physical dark Higgs  $\phi$  emerges as lightest stable particle  $\rightarrow$  can be a **dark matter candidate**, but need to deplete its energy density through decay/annihilation to SM particles.

- With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$
- Physical dark Higgs  $\phi$  emerges as lightest stable particle  $\rightarrow$  can be a **dark matter candidate**, but need to deplete its energy density through decay/annihilation to SM particles.
- Higgs mixing?

- With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$
- Physical dark Higgs  $\phi$  emerges as lightest stable particle  $\rightarrow$  can be a **dark matter candidate**, but need to deplete its energy density through decay/annihilation to SM particles.
- Higgs mixing?  $\times$  strongly constrained and  $\phi$  too long-lived  $\tau>0.1~{\rm s}~({\rm BBN}~{\rm bound})$ 
  - A. M. Sirunyan et al. (2018)

- With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$
- Physical dark Higgs  $\phi$  emerges as lightest stable particle  $\rightarrow$  can be a **dark matter candidate**, but need to deplete its energy density through decay/annihilation to SM particles.
- Higgs mixing?  $\times$  strongly constrained and  $\phi$  too long-lived  $\tau>0.1~{\rm s}~({\rm BBN}~{\rm bound})$

A. M. Sirunyan et al. (2018)

• Kinetic mixing? can work

- With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$
- Physical dark Higgs  $\phi$  emerges as lightest stable particle  $\rightarrow$  can be a **dark matter candidate**, but need to deplete its energy density through decay/annihilation to SM particles.
- Higgs mixing?  $\times$  strongly constrained and  $\phi$  too long-lived  $\tau>0.1~{\rm s}~({\rm BBN}~{\rm bound})$

A. M. Sirunyan et al. (2018)

• Kinetic mixing? can work, proceeds through *forbidden* annihilation channel  $\phi \phi \rightarrow A'A'$  followed by  $A' \rightarrow e^+e^-$ .

R. T. D'Agnolo and J. T. Ruderman (2015)

- With a thermalized dark sector, preferred parameters lead to  $T_*\sim 2-2.5$  MeV  $\longrightarrow$  BBN constraints Y. Bai & M. Korwal (2021)  $\checkmark$
- Physical dark Higgs  $\phi$  emerges as lightest stable particle  $\rightarrow$  can be a **dark matter candidate**, but need to deplete its energy density through decay/annihilation to SM particles.
- Higgs mixing?  $\times$  strongly constrained and  $\phi$  too long-lived  $\tau>0.1~{\rm s}~({\rm BBN}~{\rm bound})$

A. M. Sirunyan et al. (2018)

• Kinetic mixing? can work, proceeds through *forbidden* annihilation channel  $\phi \phi \rightarrow A'A'$  followed by  $A' \rightarrow e^+e^-$ .

R. T. D'Agnolo and J. T. Ruderman (2015)

$$\Omega_{\phi}h^2 \sim 10^{-13} \left(\frac{m_{\phi}}{3\,\mathrm{MeV}}\right)^2 e^{x_f\,\Delta}, \quad \Delta \equiv \frac{2(m_{A'}-m_{\phi})}{m_{\phi}}$$

## Summary and Outlook

- A dark sector phase transition can account for the **whole PTA signal**, but parameter space is extremely constrained.
- Leads to predictive phenomenology for the underlying model.
- Interesting implications for the DS cosmology ("can" get right DM).

# Thank you for your attention!

# Backup

## $\beta/H_*$ and $\alpha_*$ for SU(N)

Set  $\mu = 1$  MeV and  $y_D = 0$ . Study trend of  $\beta/H_*$  and  $\alpha_*$  in  $(g_D, \lambda)$  plane.



Enlarging the gauge group (N = 2: solid lines, N = 3: dashed lines) 'tilts' the parameter space!

### Yukawa Coupling

Work with U(1) group with two fermions. Fix  $\lambda = 0.05$  and  $\mu = 1$  MeV to study trends in  $(g_D, y_D)$  plane.



Can increase  $g_D$  to compensate for negative fermionic contribution!