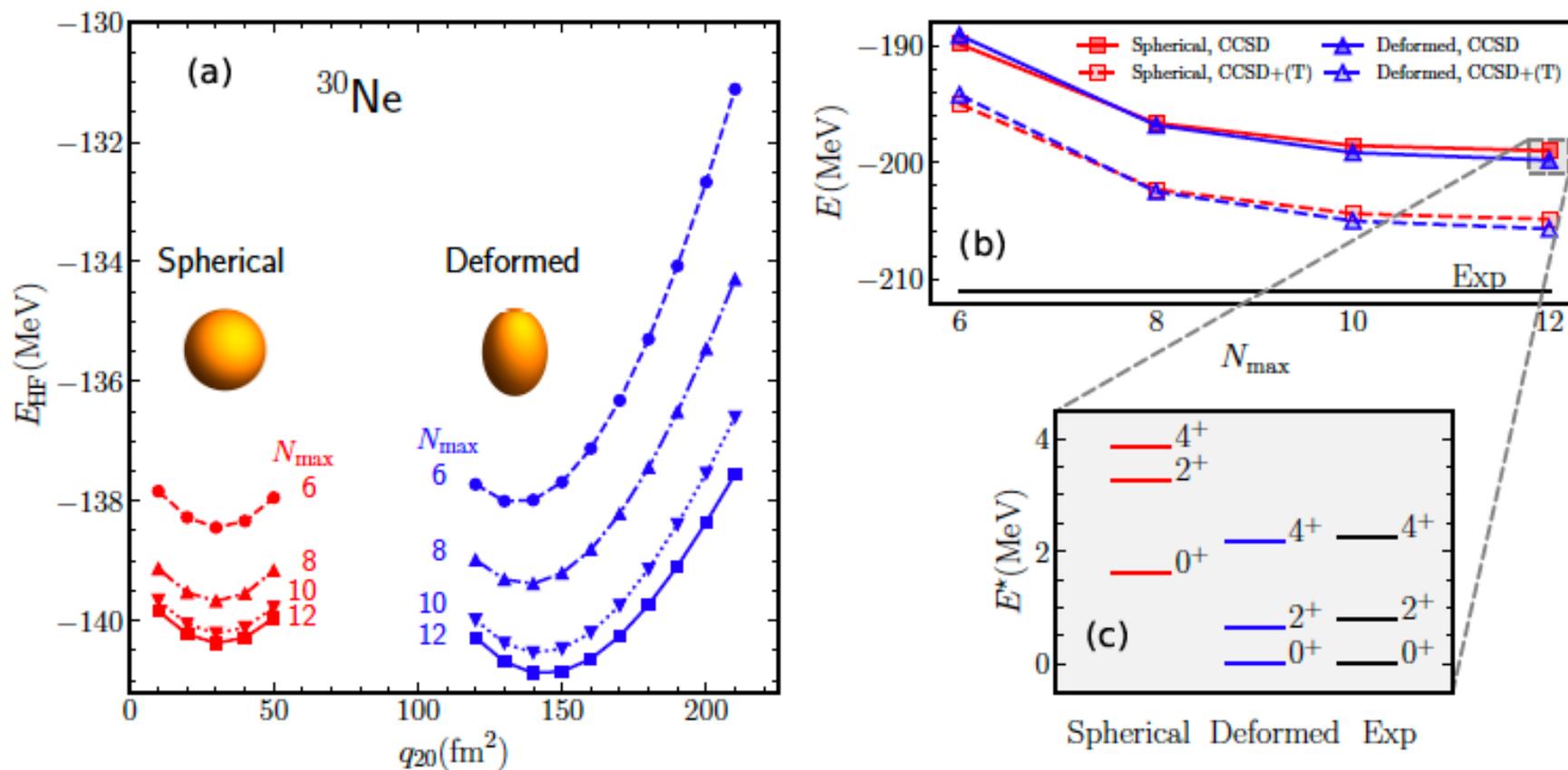


Deformed Nuclei at extreme isospin



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Halo Week

Gothenburg, June 10, 2024

Work supported by the US Department of Energy

Today's menu

- Baishan Hu, Zhonghao Sun, G. Hagen, TP, *Ab initio computations of strongly deformed nuclei around ^{80}Zr* , arXiv:2405.05052
- Zhonghao Sun, A. Ekström, C. Forssén, G. Hagen, G. R. Jansen, TP, *Multiscale physics of atomic nuclei from first principles*, arXiv:2404.00058
- B. Acharya, B. S. Hu, S. Bacca, G. Hagen, P. Navrátil, TP, *The magnetic dipole transition in ^{48}Ca* , Phys. Rev. Lett. 132, 232504 (2024).

Multiscale problem:

The bulk of the binding energy is from short-range correlations
Symmetry projection accounts for small details

1. Coester and Kümmel (1960), “Short-range correlations in nuclear wave functions”
2. Lipkin (1960), “Collective motion in many-particle systems: Part 1. the violation of conservation laws”

	E_{HF}	$E_{CCSD(T)}$	$E_{Proj.}$	$\langle J_{HF} \rangle$	$\langle J_{CCSD(T)} \rangle$
^8Be	-16.74	-50.24	-53.57	11.17	5.82
^{20}Ne	-59.62	-161.95	-164.21	21.26	12.09
^{34}Mg	-90.21	-264.34	-265.84	22.62	15.03

Data from Hagen et al., Phys. Rev. C 105, 064311 (2022)

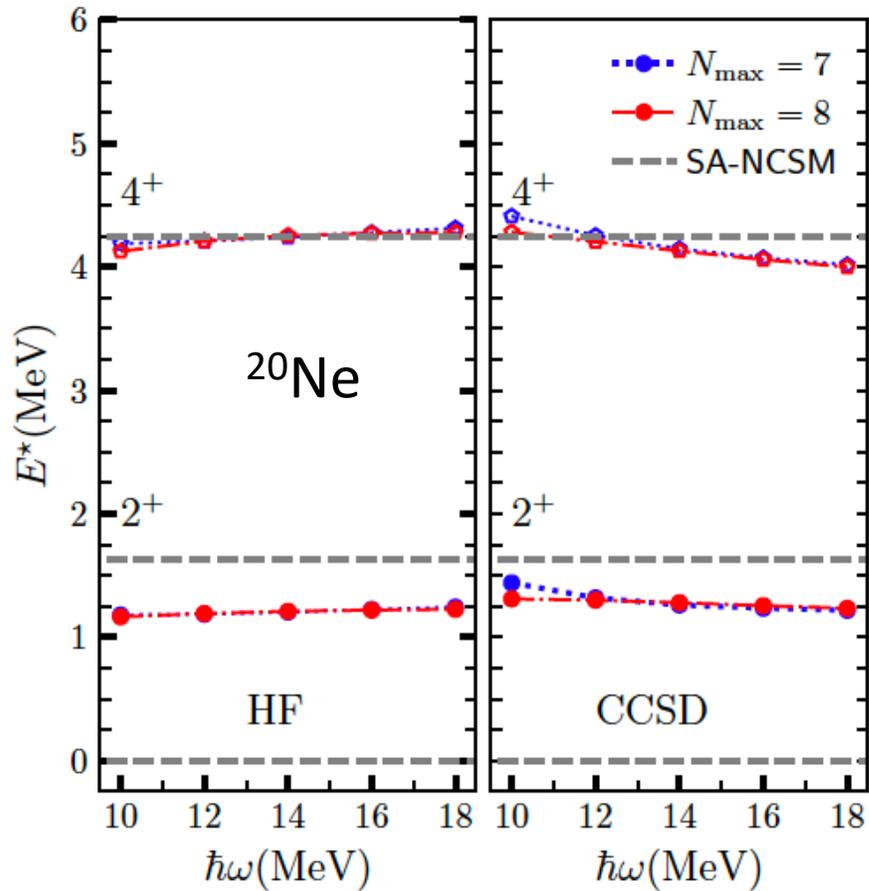
Energy gain from symmetry projection is small and not size extensive

Our approach

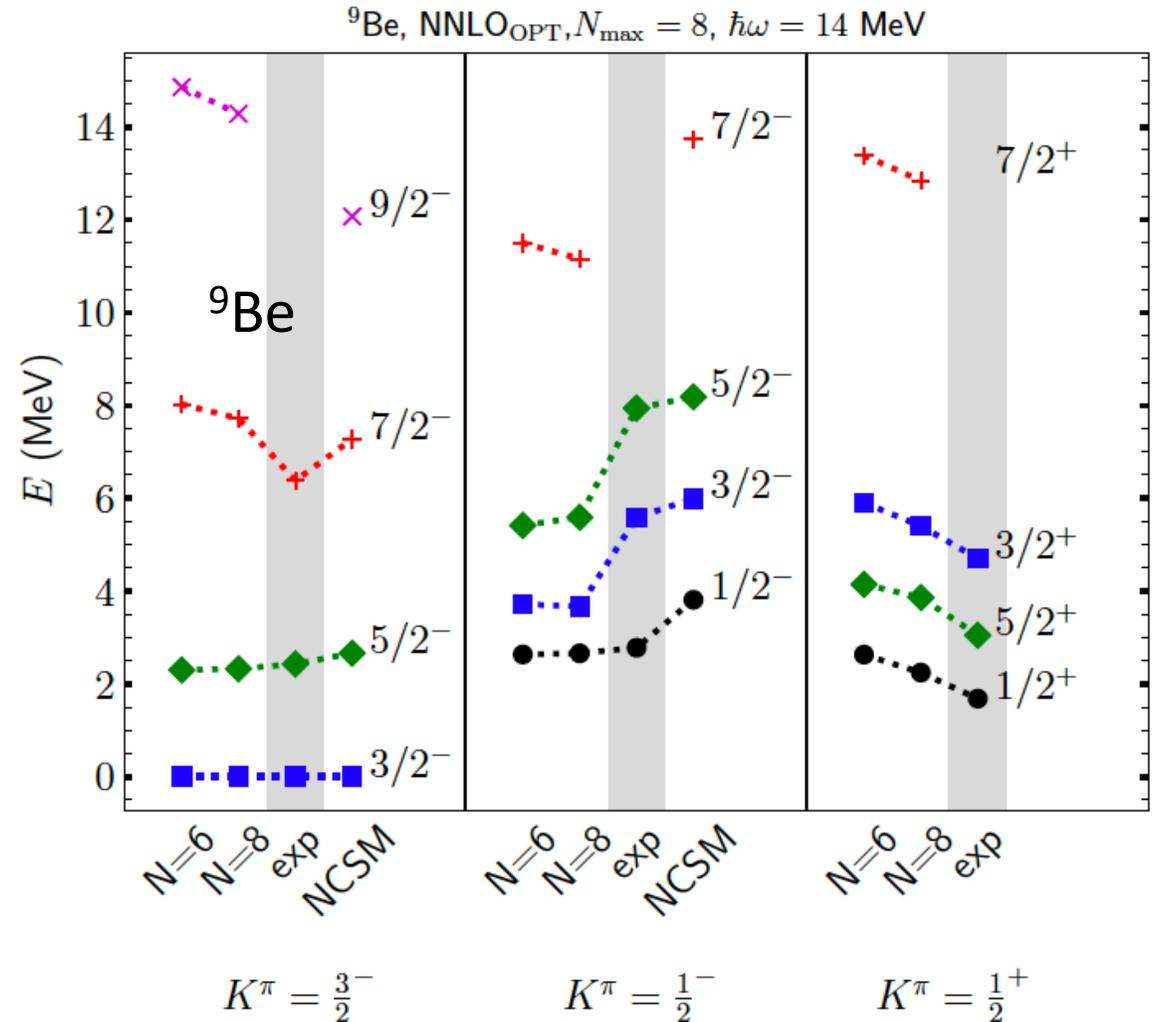
Include short-range correlations first, then long-range ones

1. Start from an axially symmetric reference state
2. Include short-range (“dynamical”) correlations via coupled cluster method
 - captures UV physics
3. Symmetry projection includes collective effects
 - captures IR physics

Meeting NCSM benchmarks



SA-NCSM: Launey, Dytrych, Sargsyan, Baker, Draayer, Eur. Phys. J. Special Topics 229, 2429 (2020).



NCSM: Caprio, Maris, Vary, Smith, Int. J. Mod. Phys. E 24, 1541002 (2015).

Shape coexistence

States with different shapes that are close in energy

Reviews: Heyde and Wood, *Rev. Mod. Phys.* 83, 1467 (2011); Gade and Liddick, *J. Phys. G* 43, 024001 (2016); Bonatsos, et al., *Atoms* 11, 117 (2023).

Observed in ^{30}Mg by Schwerdtfeger et al., *Phys. Rev. Lett.* 103, 012501 (2009) and in ^{32}Mg by Wimmer et al., *Phys. Rev. Lett.* 105, 252501 (2010).

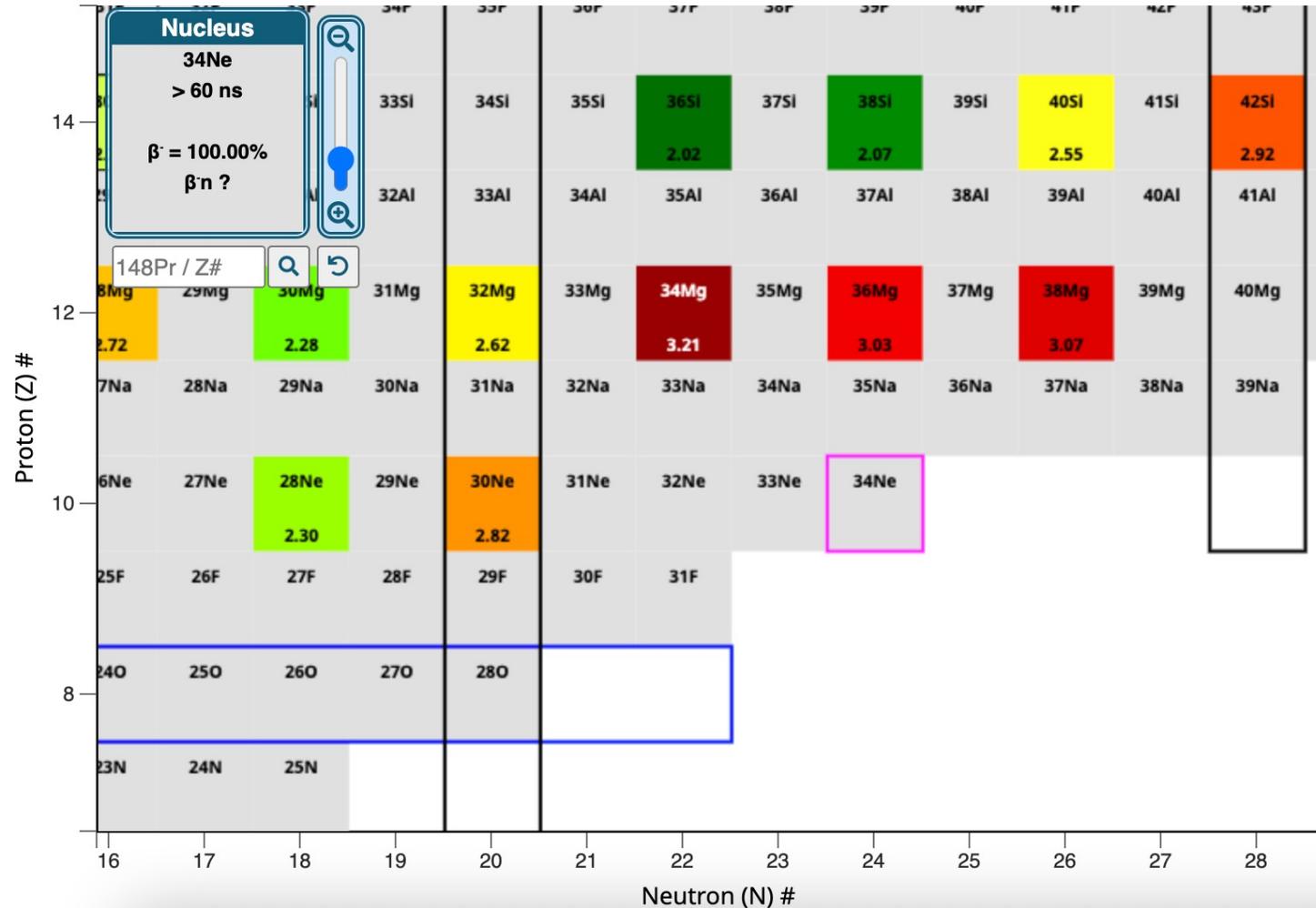
Theoretical descriptions: Reinhard et al., *Phys. Rev. C* 60, 014316 (1999); Rodríguez-Guzmán, Egido, and Robledo, *Nucl. Phys. A* 709, 201 (2002); Péru and Martini, *Eur. Phys. J. A* 50, 88 (2014); Caurier, Nowacki, and Poves, *Phys. Rev. C* 90, 014302 (2014); see also Tsunoda et al., *Nature* 587, 66 (2020).

Neutron-rich nuclei beyond $N \geq 20$ are deformed

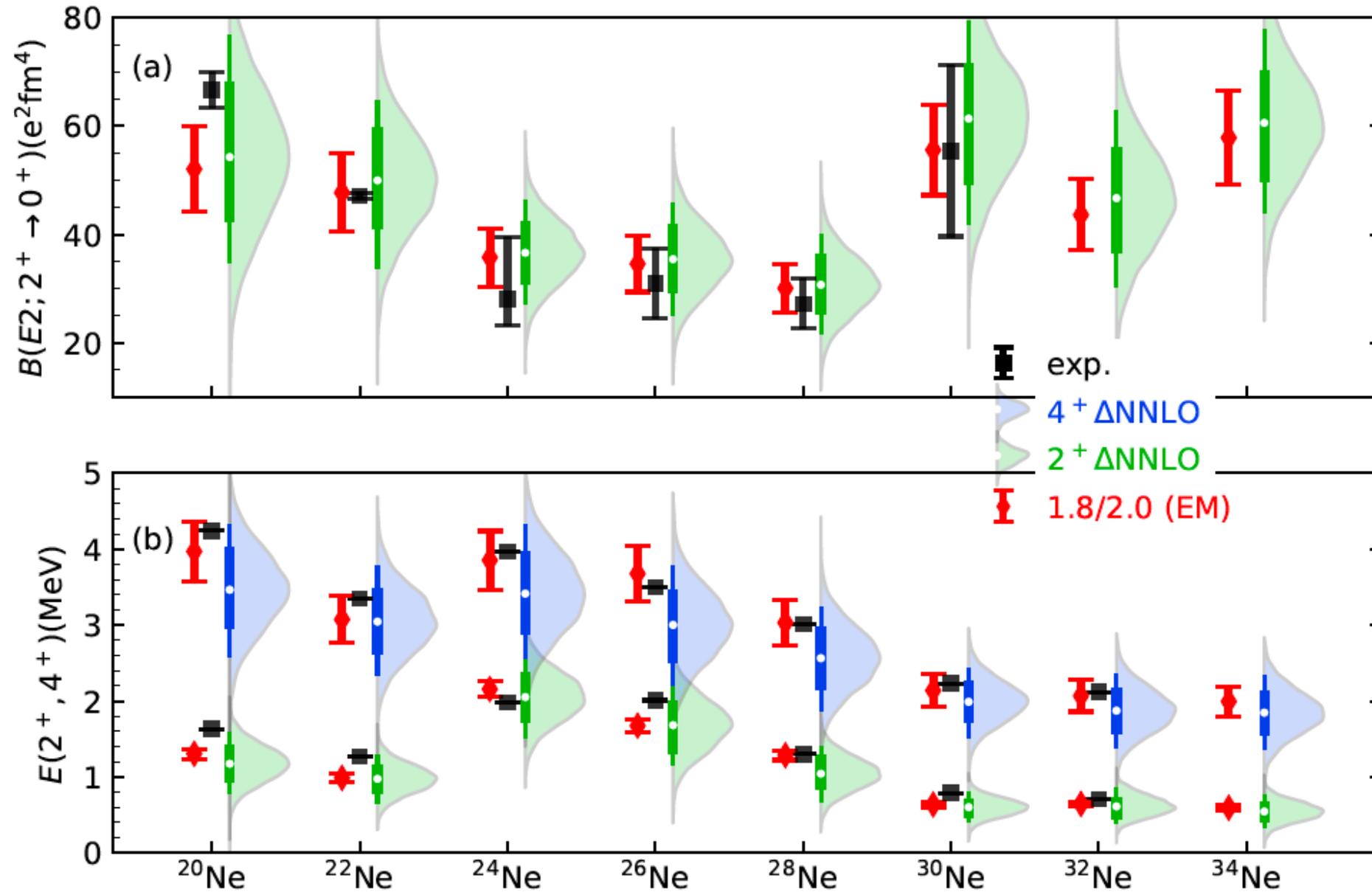
$$R_{4/2} \equiv \frac{E_{4^+}}{E_{2^+}}$$

$R_{4/2} = 10/3$ for a rigid rotor

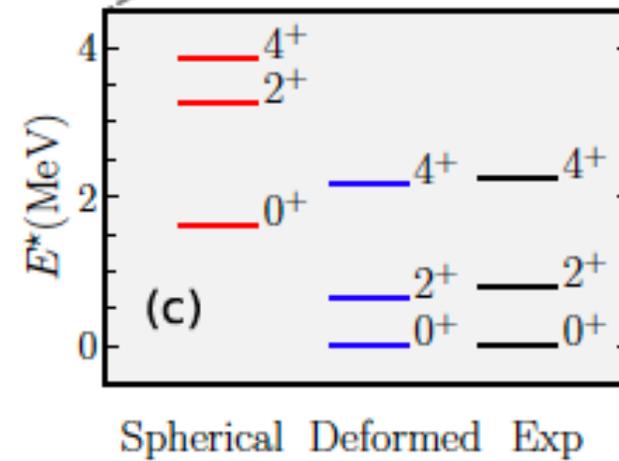
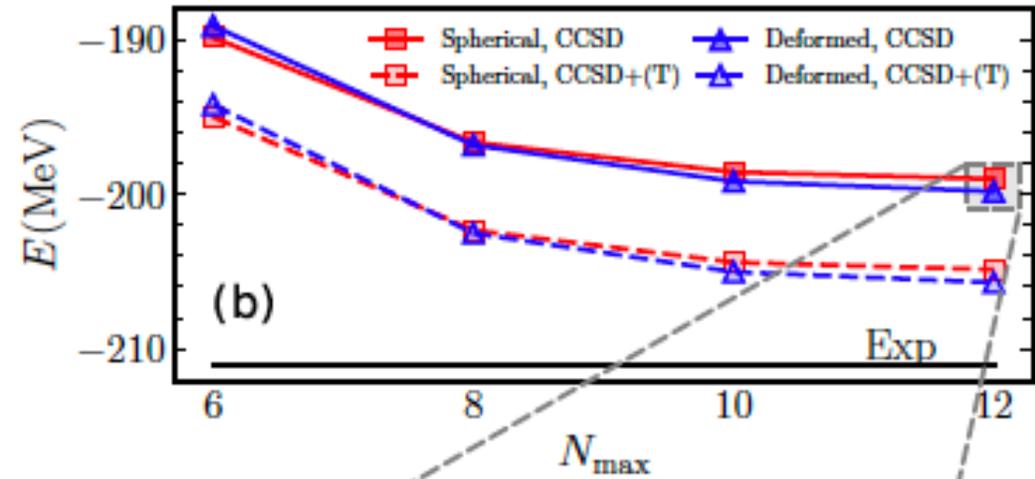
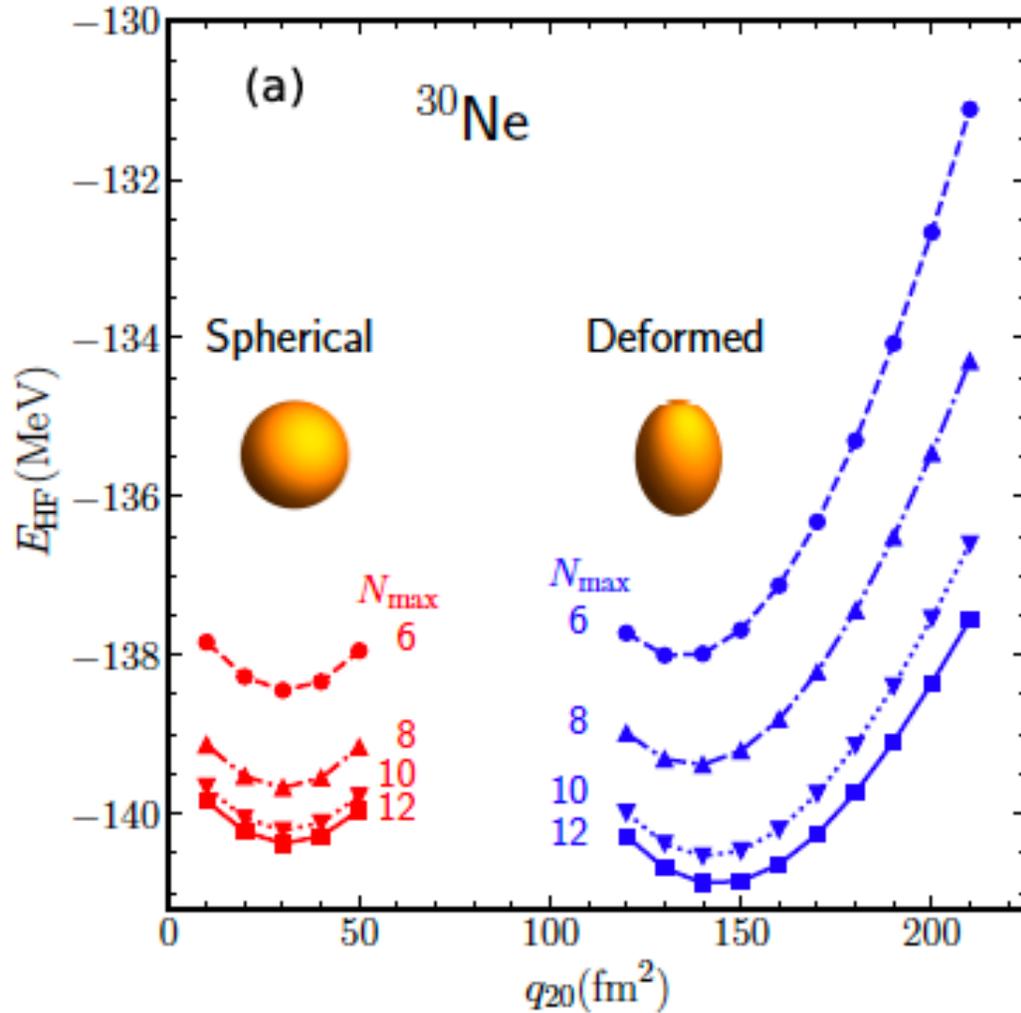
Simple picture: Spherical states (magic $N = 20$ number in the traditional shell model) coexist with deformed ground states



Coupled cluster theory: collectivity of neon nuclei

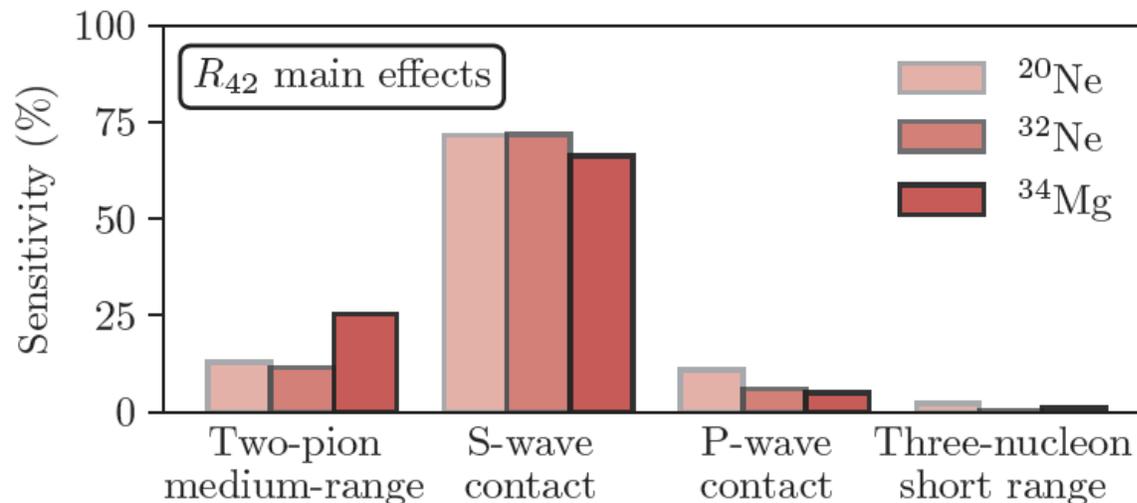


Prediction: Shape coexistence in ^{30}Ne

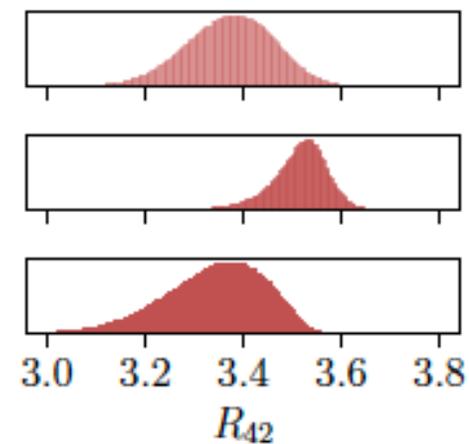
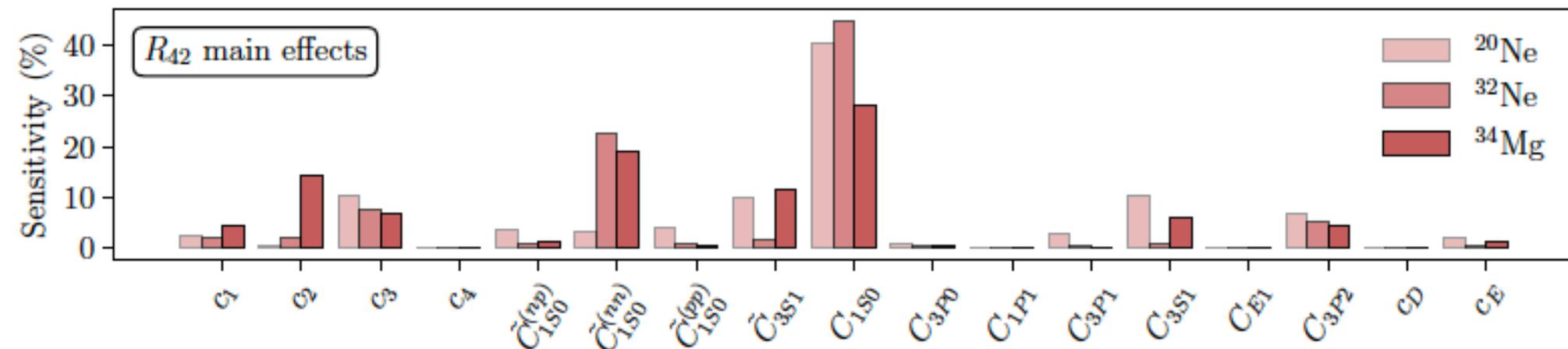


What drives nuclear deformation in chiral EFT?

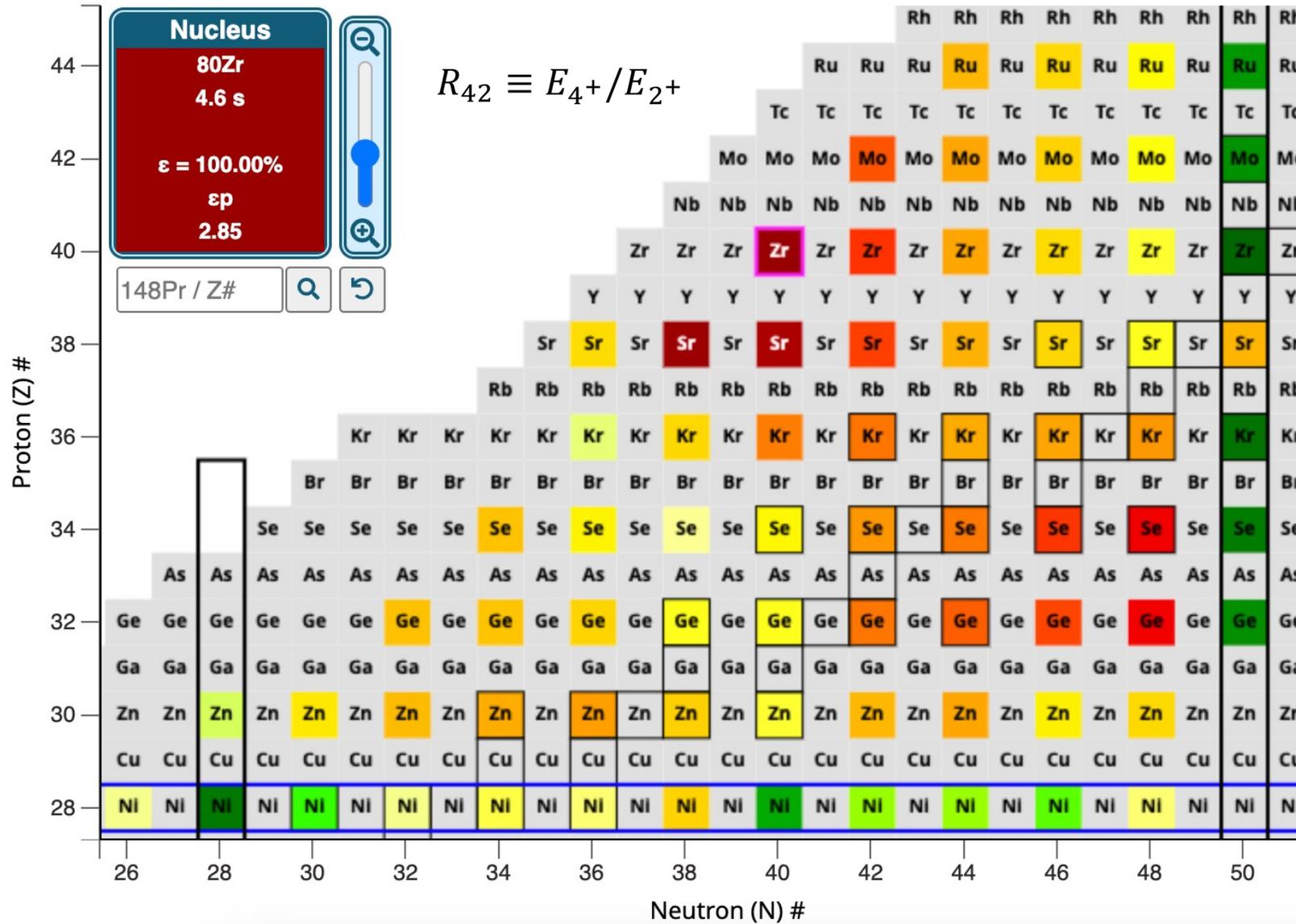
Executive summary



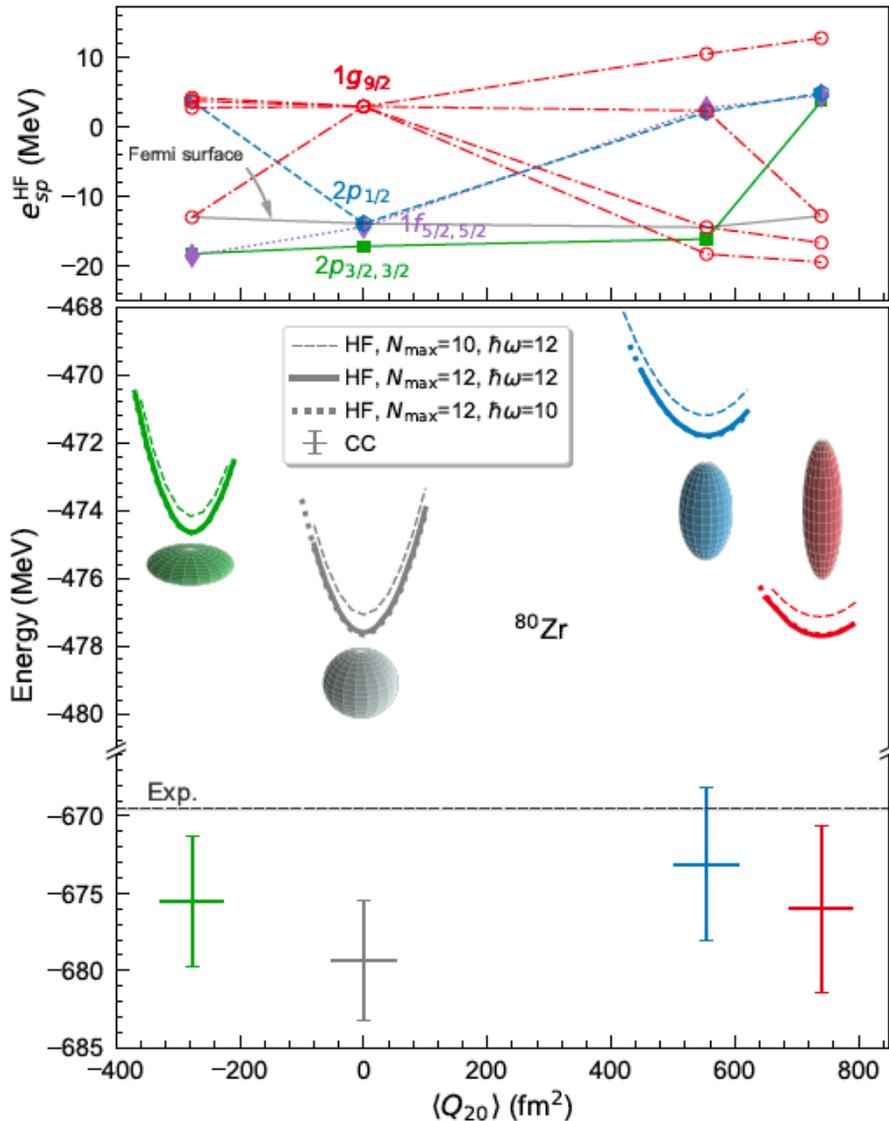
More detailed view



The region around ^{80}Zr



Shapes of ^{80}Zr



Quadrupole constrained HF computations

- several minima identified
- angular momentum projected

Shape coexistence identified

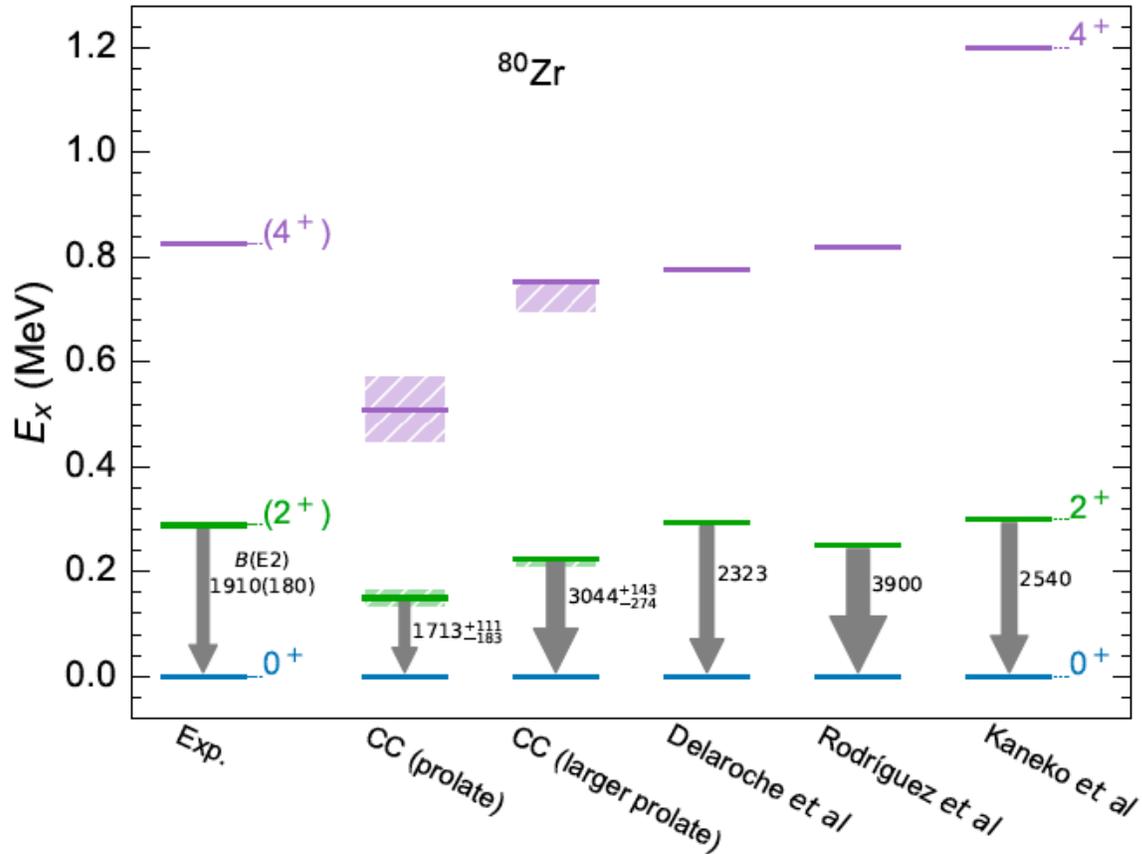
- coupled-cluster computations too uncertain to predict shape of ground state

Used Miyagi (2023) for 3NFs in large model spaces

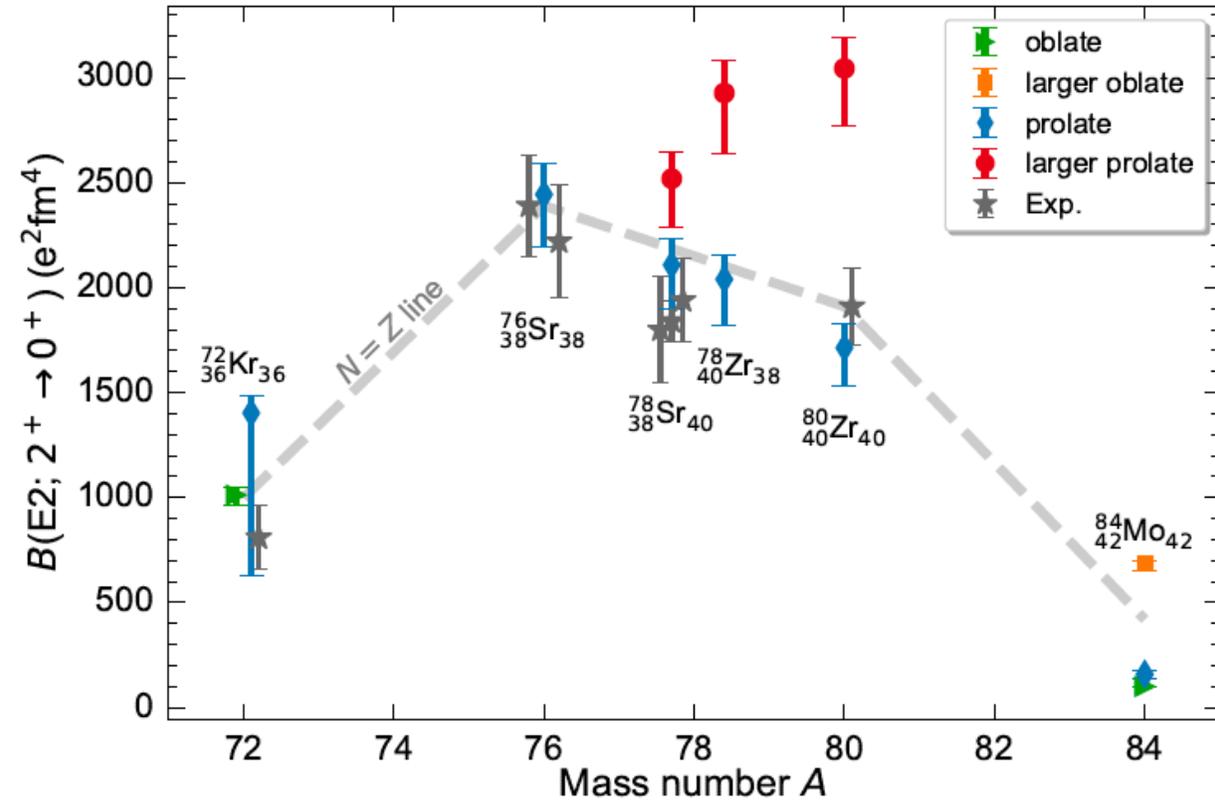
Fun fact: ^{80}Zr has higher energy than two ^{40}Ca nuclei

The region around ^{80}Zr

Spectrum prefers larger prolate shape



$B(E2)$ prefers prolate band



Why do people care about M1 transitions?

Supernova 1987A



February 24, 1987
Las Campanas Observatory

M1 spin excitations are dominated by isovector contributions.

The isovector-0 component of the Gamow-Teller operator translates to inelastic neutral-current neutrino-nucleus reactions at energies relevant for supernovae.

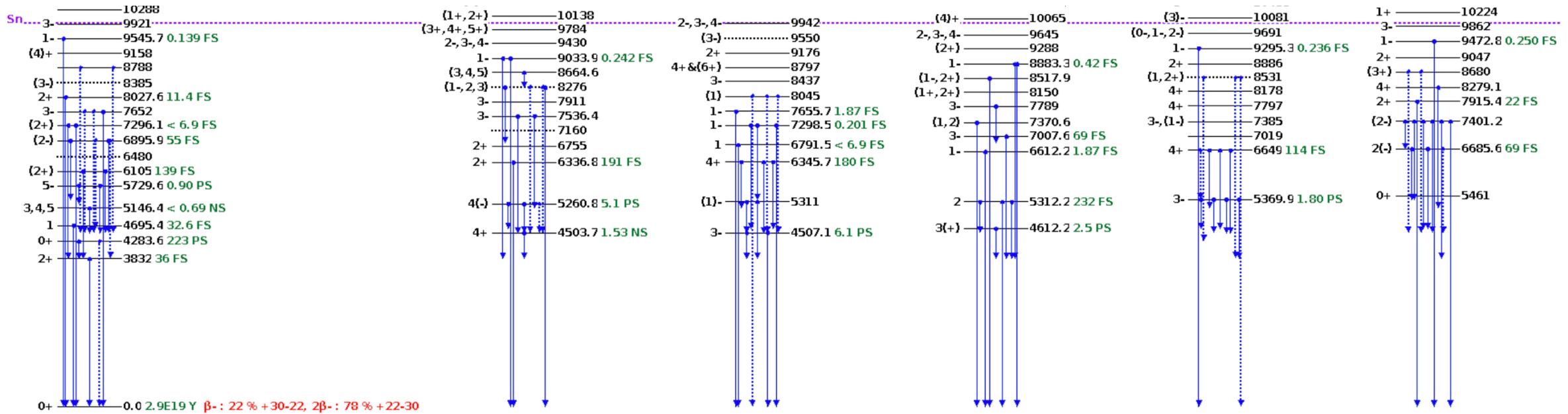
Our understanding of M1 impacts supernovae signals and dynamics.

Lüttge, von Neumann-Cosel, Neumeyer, Richter, Nucl Phys A (1996);
Langanke, Martinez-Pinedo, von Neumann-Cosel, Richter, Phys Rev Lett (2004);
Loens, Langanke, Martinez-Pinedo, Sieja, EPJA (2012);
Tornow et al, Phys Letts B (2022).

Review on *M1*:

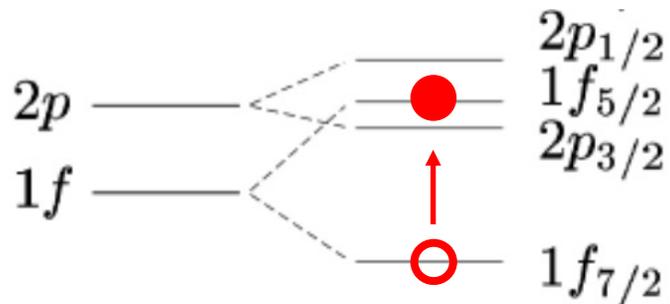
K. Heyde, P. von Neumann-Cosel, A. Richter, Rev. Mod. Phys. 82, 2365 (2010).

The resonant 1^+ HALO state in ^{48}Ca at 10.224 MeV



Scattering / reactions that probe the 1^+ state: (e, e') , (p, p') , (p, n) , or (γ, n)

Simple picture of the 1^+ state: neutron $1p$ - $1h$ excitation; extreme single-particle model: $B(M1) = 12 \mu_N^2$



	S_n (MeV)	ΔE (MeV)	Γ (keV)	$1p$ - $1h$
$\Delta\text{NNLO}_{\text{GO}}(394)$	9.74	-0.44	0	91%
$\Delta\text{NNLO}_{\text{GO}}(450)$	9.38	-1.26	0	91%
NNLO_{sat}	9.34	-0.23	0	91%
1.8/2.0(EM)	10.00	0.55	4	92%
Experiment	9.95	0.28	≤ 17	

The status

(e, e') scattering: $B(M1) = 4.0 \pm 0.3 \mu_N^2$ [Steffen et al 1980; 1983]

(γ, n) scattering: $B(M1) = 6.8 \pm 0.5 \mu_N^2$ [Tompkin et al 2011]

(p, p') scattering: $B(M1) = 3.85(32) - 4.63(38) \mu_N^2$ [Birkhan et al 2016]

Extreme s.p. model: $B(M1) = 12 \mu_N^2$

Theory has a hard time to reproduce a large amount of quenching

A. Harting, W. Weise, H. Toki, and A. Richter, Physics Letters B 104, 261 (1981).

J. B. McGrory and B. H. Wildenthal, Phys. Lett. B 103, 173 (1981).

Toru Suzuki, S. Krewald, and J. Speth, Phys. Lett. B 107, 9 (1981).

G. F. Bertsch, Nucl. Phys. A 354, 157 (1981).

M. Kohno and D. W. L. Sprung, Phys. Rev. C 26, 297 (1982).

K. Takayanagi, K. Shimizu, and A. Arima, Nucl. Phys. A 481, 313 (1988).

M. G. E. Brand, K. Allaart, and W. H. Dickhoff, Nucl. Phys. A 509, 1 (1990).

B. A. Brown and W. A. Richter, Phys. Rev. C 58, 2099 (1998).

J. D. Holt, J. Menendez, J. Simonis, and A. Schwenk, Phys. Rev. C 90, 024312 (2014).

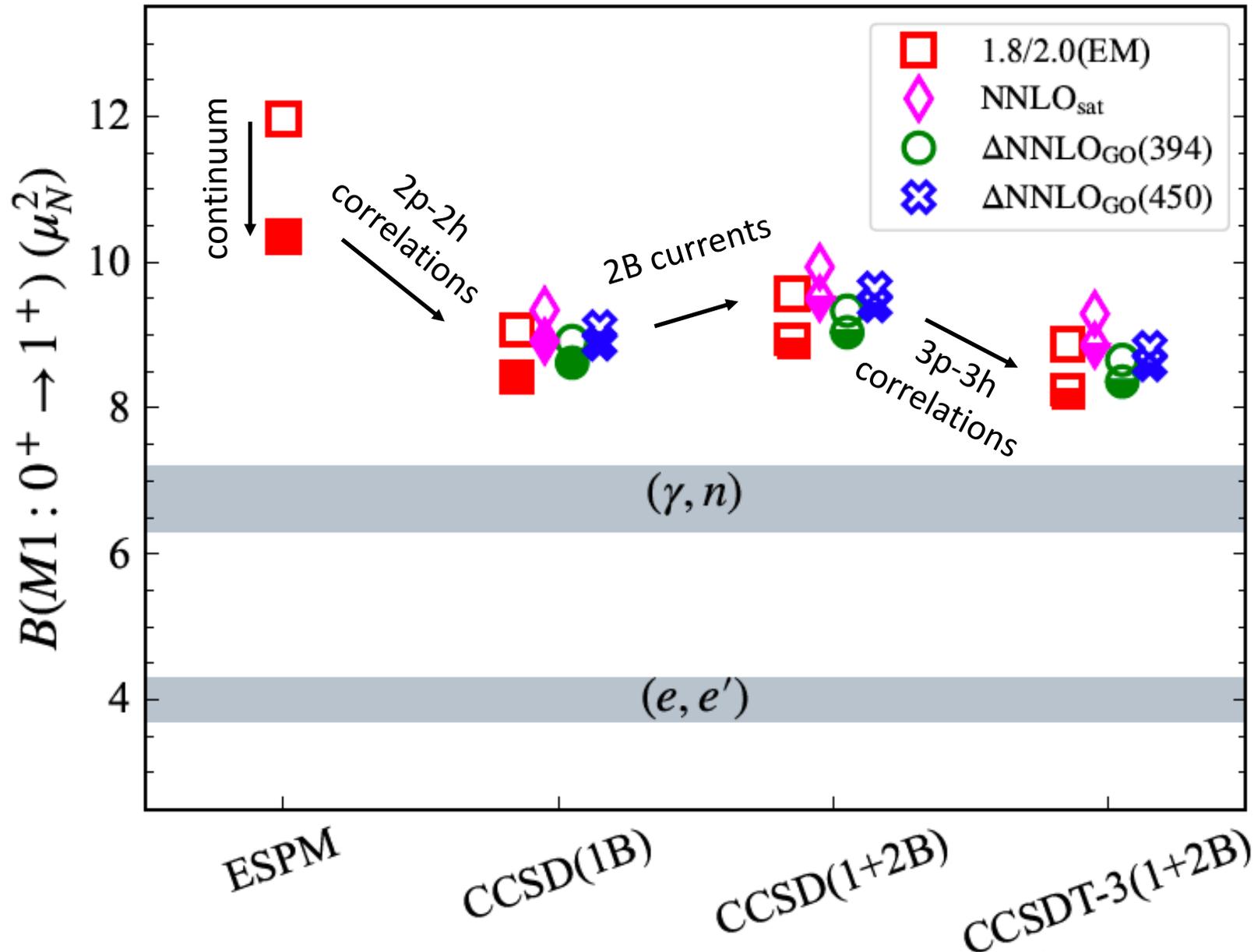
J. Wilhelmy, et al., Phys. Rev. C 98, 034315 (2018).

} Meson-exchange currents claimed to explain small $B(M1)$

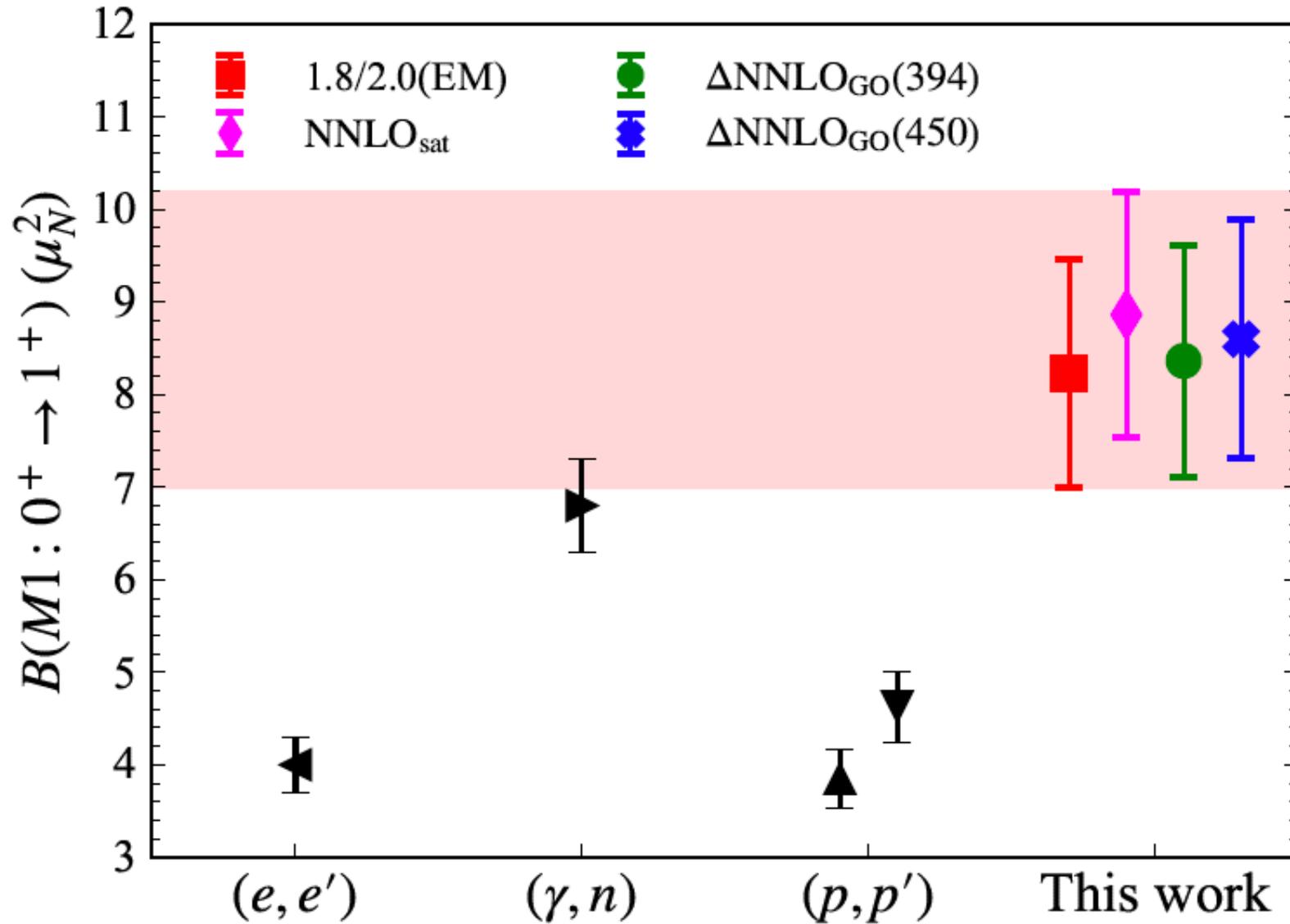
} All too high $B(M1)$;
 $B(M1) = 7 - 8 \mu_N^2$;
 $B(M1) > 5.1 \mu_N^2$;

} Reproduce (e, e') $B(M1)$ if quenched

Contributions to $B(M1)$



Final result



Summary

Breaking and restoring symmetries

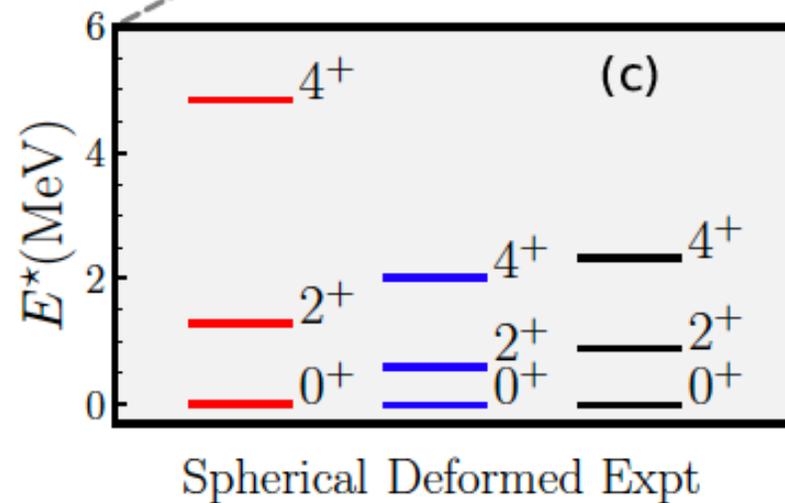
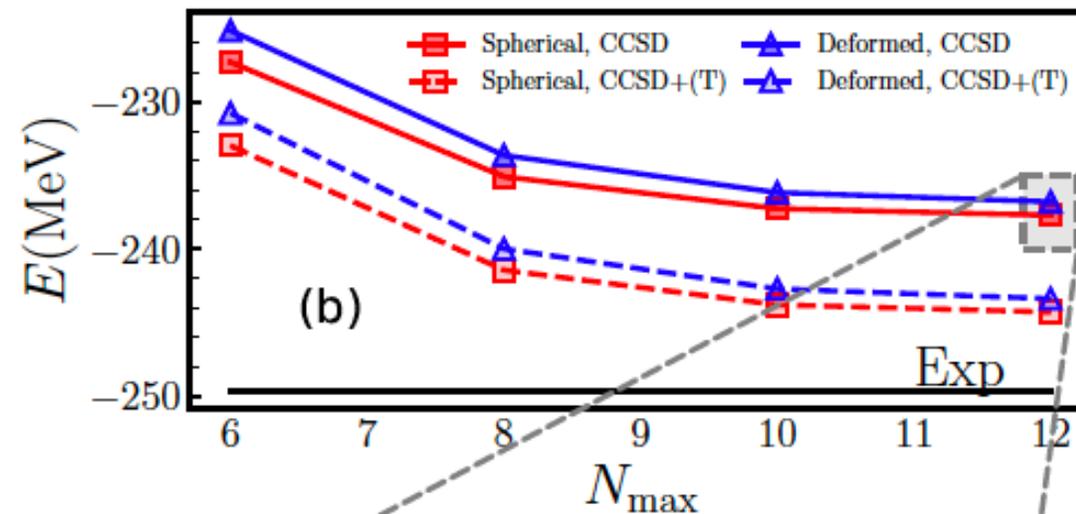
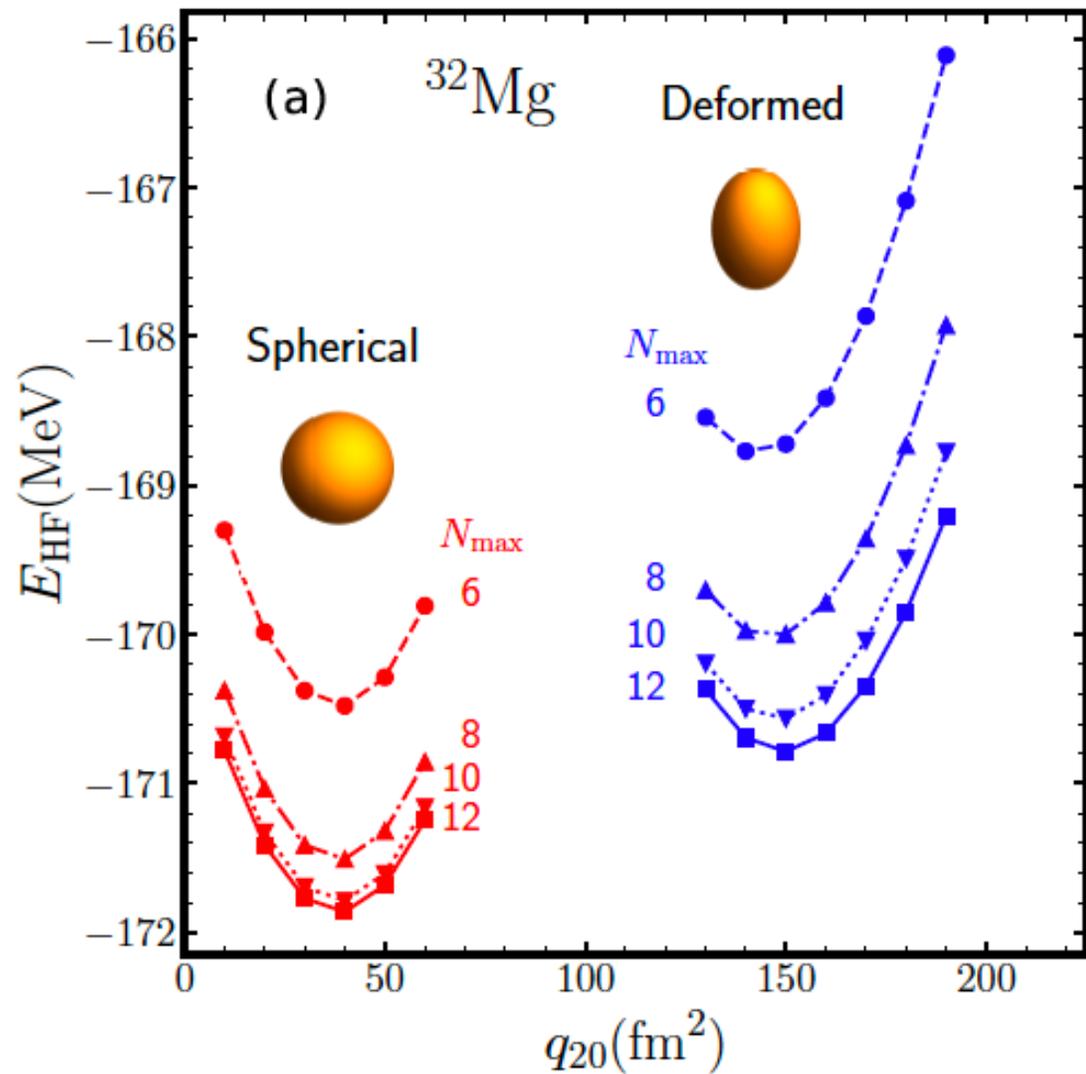
Exploits separation of scale between universal collective and specific UV physics

Conceptually simple & computationally affordable

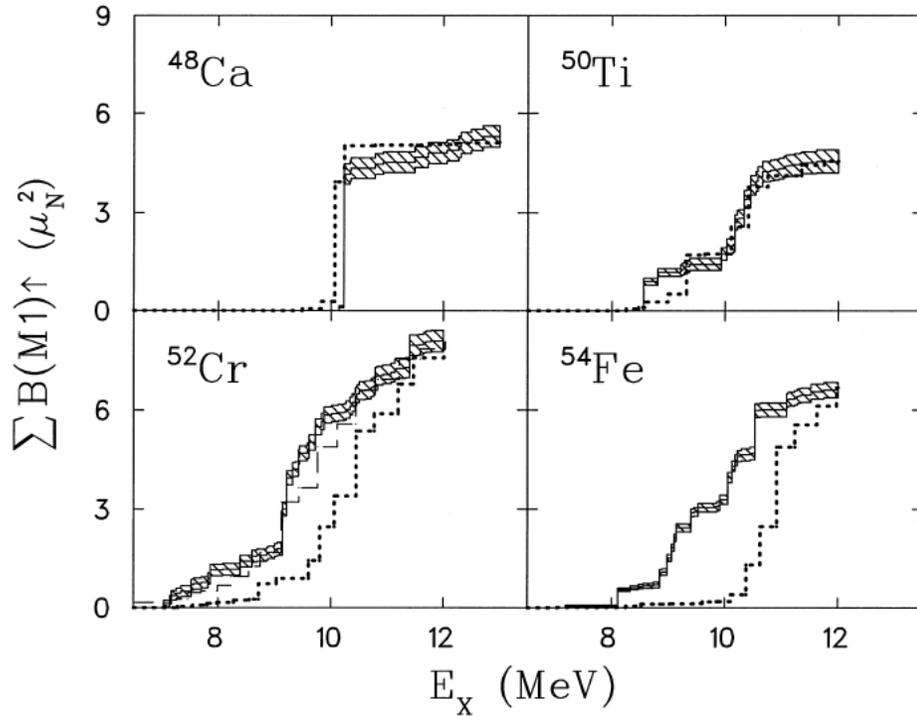
- Shape coexistence in ^{30}Ne
- Connected deformation to microscopic forces
 - Much improved $B(E2)$ values
 - ^{3x}Ne , ^{3x}Mg , ^{80}Zr
- $B(M1)$ in ^{48}Ca larger than (e,e') but in agreement with (n,γ)
 - Two-body currents do not quench $B(M1)$

Thank you!

Confirmation: Shape coexistence in ^{32}Mg

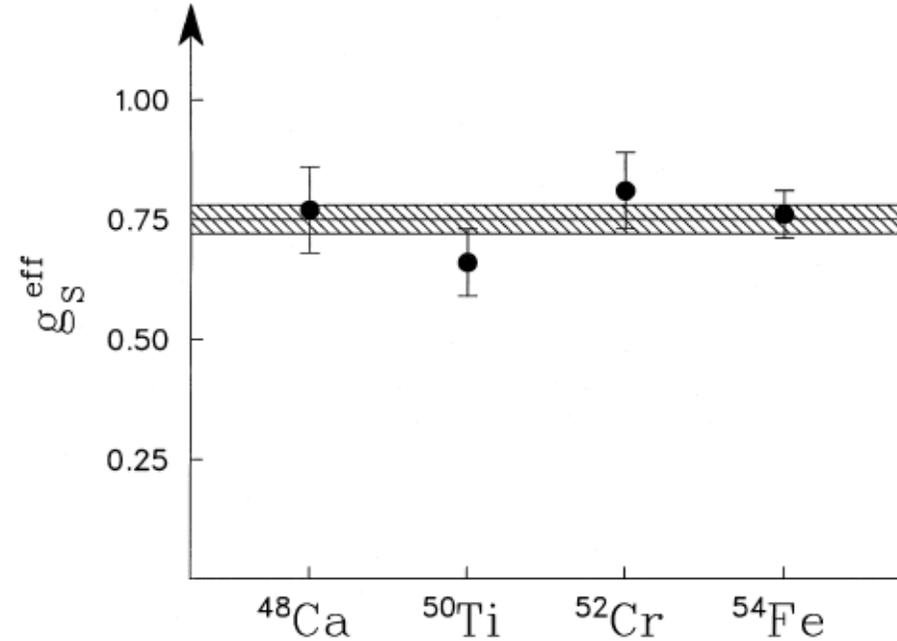


Why could/should there be quenching?



Results from (e, e') scattering match quenched shell-model results

Von Neumann-Cosel, Poves, Retamosa, Richter, Phys Letts B (1998)



Proposed: $B(M1)$ is quenched similarly to $B(GT)$ in pf shell nuclei

→ Impacts (re)analyses of (p, p') experiments using the “unit cross section” method

Two-body currents do not quench M1 transitions in light nuclei

$J_i^\pi \rightarrow J_f^\pi$	Method	IA	MEC				Total
			$\pi + \rho$	MS	MD	Δ	
${}^6\text{Li}(0^+; 1) \rightarrow {}^6\text{Li}(1^+; 0)$	VMC	3.683(14)	0.307	0.003	0.010	-0.053	3.950(14)
${}^6\text{Li}(0^+; 1) \rightarrow {}^6\text{Li}(1^+; 0)$	GFMC	3.587(16)	0.323	0.002	0.012	-0.048	3.876(14)
${}^7\text{Li}(\frac{1}{2}^-) \rightarrow {}^7\text{Li}(\frac{3}{2}^-)$	VMC	2.743(17)	0.396	0.006	-0.017	-0.034	3.162(22)
${}^7\text{Li}(\frac{1}{2}^-) \rightarrow {}^7\text{Li}(\frac{3}{2}^-)$	GFMC	2.677(19)	0.395	0.011	-0.017	0.072	3.138(22)
${}^7\text{Be}(\frac{1}{2}^-) \rightarrow {}^7\text{Be}(\frac{3}{2}^-)$	VMC	2.420(30)	0.390	-0.005	0.010	-0.024	2.791(36)
${}^7\text{Be}(\frac{1}{2}^-) \rightarrow {}^7\text{Be}(\frac{3}{2}^-)$	GFMC	2.374(31)	0.394	-0.010	0.010	-0.002	2.766(36)

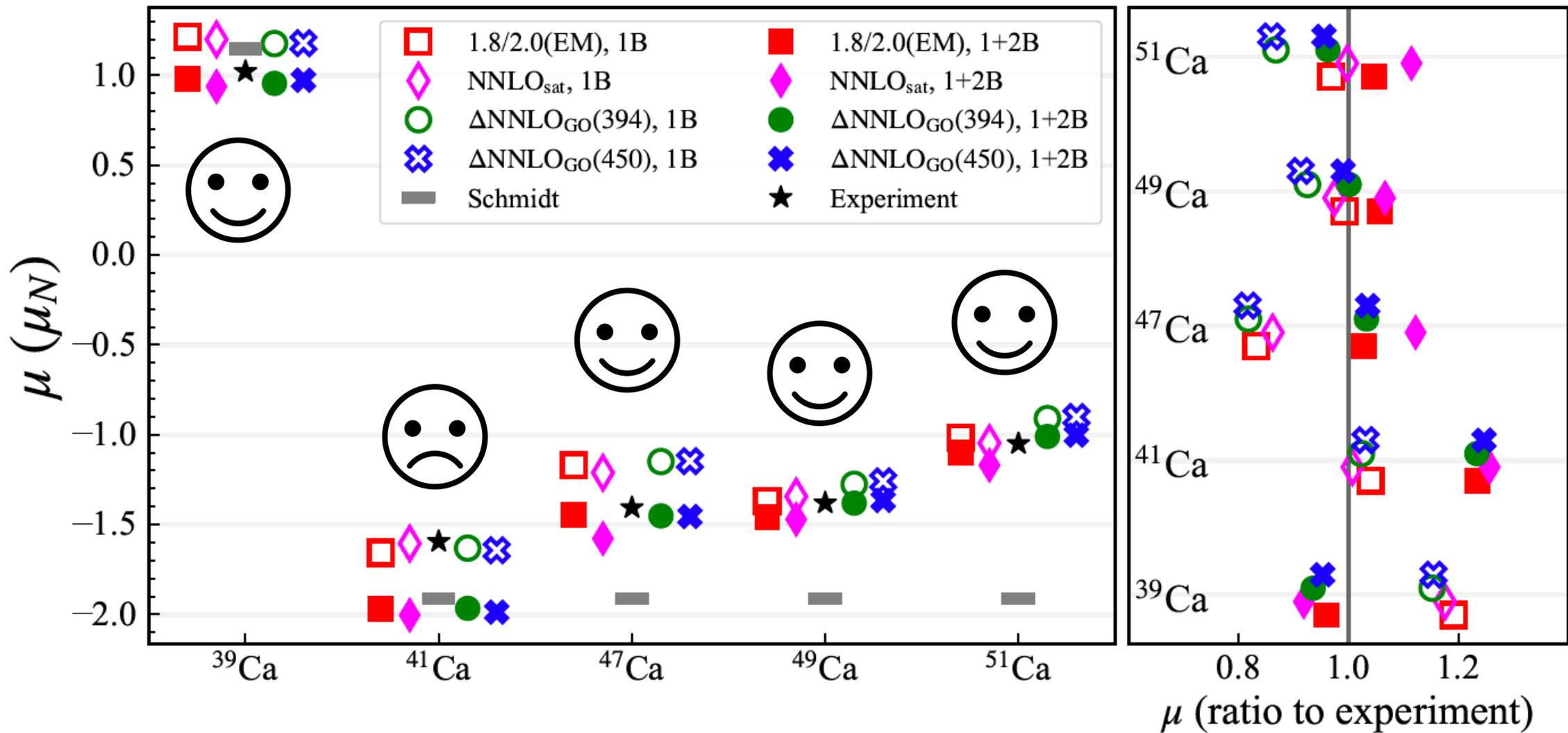
Marcucci, Muslema Pervin, Pieper, Schiavilla,
Wiringa, Phys Rev C 78, 065501 (2008)

This is similar to
what we will use

This is perhaps
similar to what
people used in
the 1980s

Two-body currents for $M1$ transitions differ from those for Gamow-Teller transitions

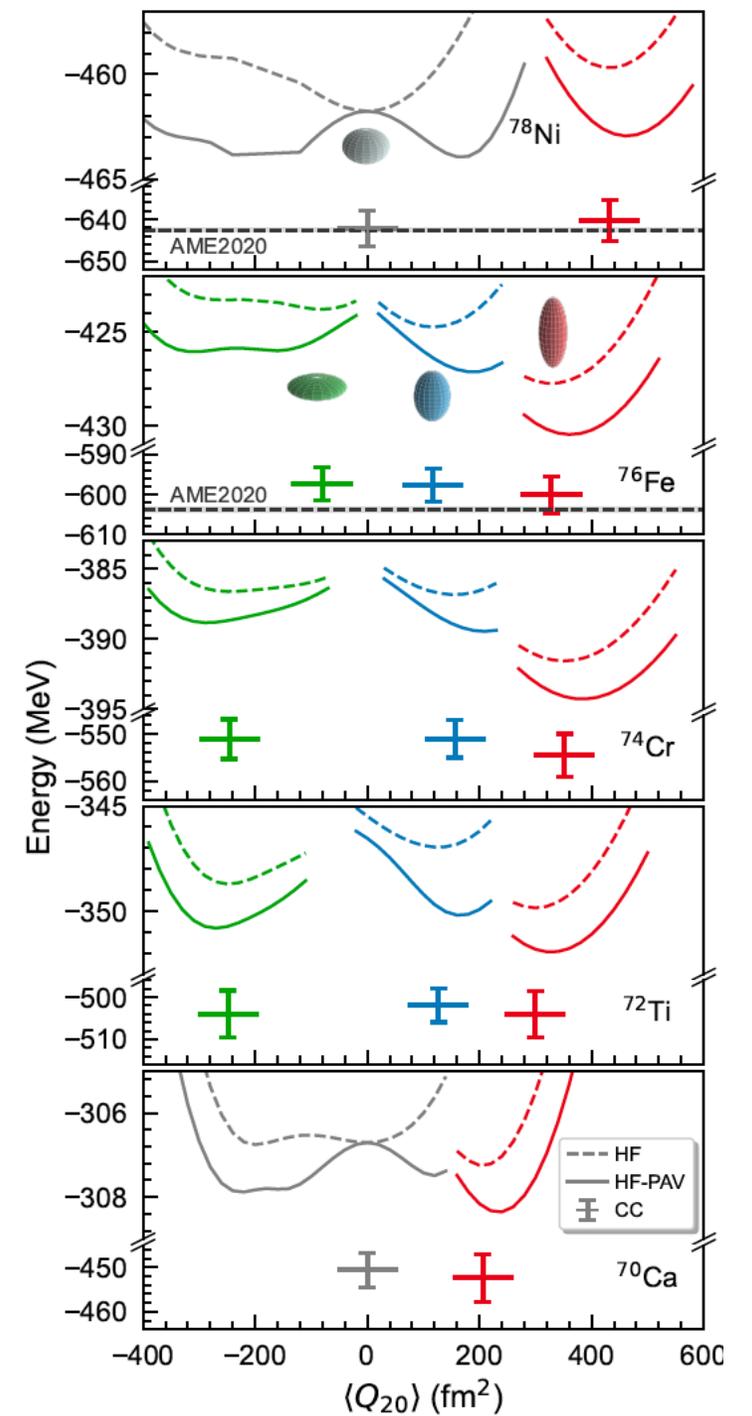
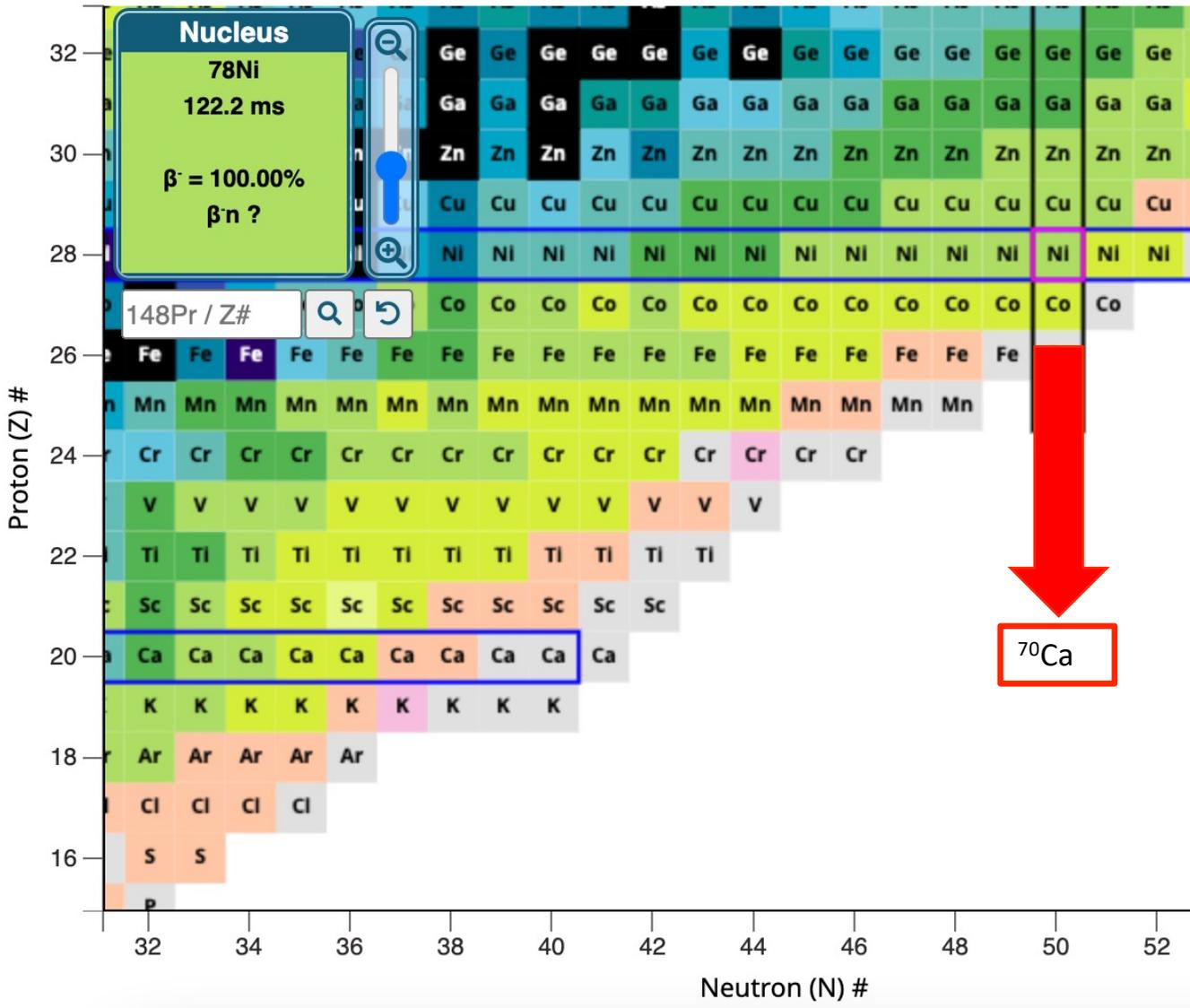
Magnetic moments



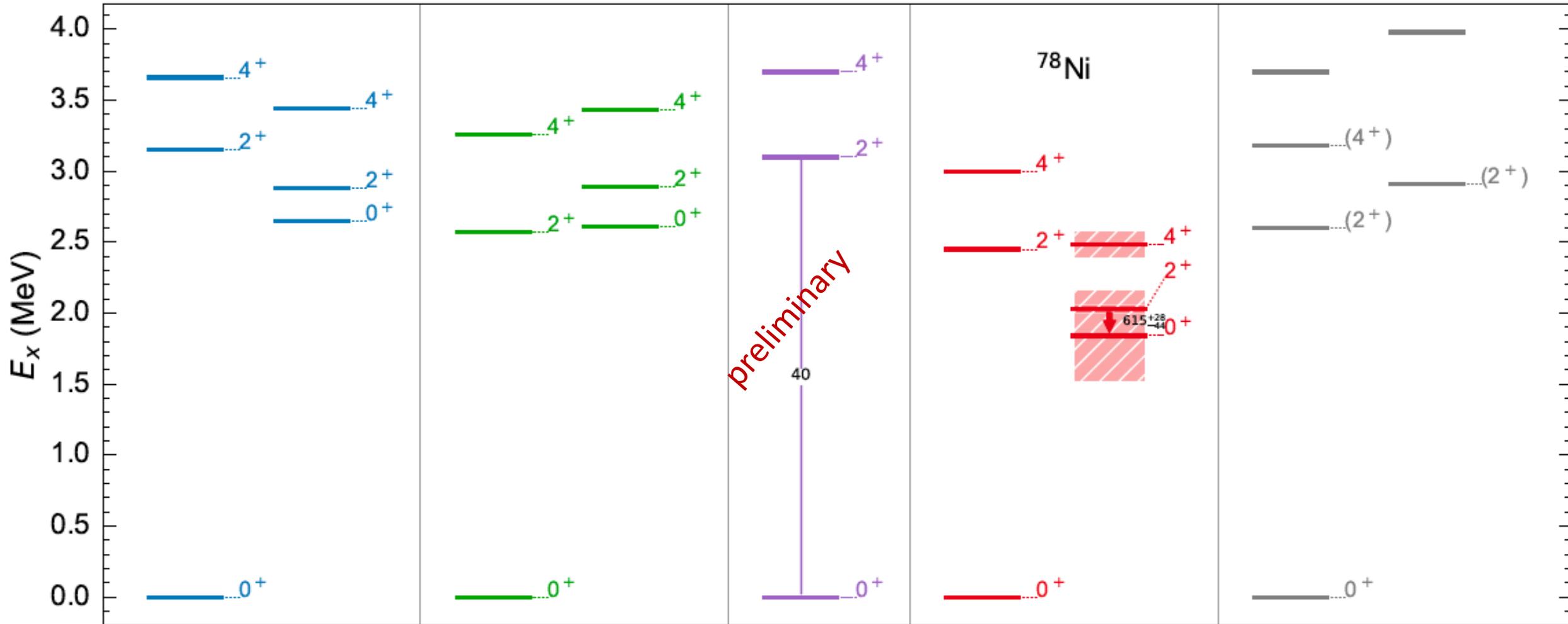
Bijaya Acharya et al., Phys. Rev. Lett. 132, 232504 (2024)

Takayuki Miyagi et al., Phys. Rev. Lett. 132, 232503 (2024): multi-shell VS-IMSRG calculation accurate for ^{41}Ca .²⁴

South of ^{78}Ni



Structure of ^{78}Ni



LSSM

Nowacki, Poves, Caurier,
Bounthong (2016)

MCSM

Taniuchi et al. (2019)

IMSRG

Tichai, Kapás, Miyagi,
Werner, Legeza,
Schwenk, Zarand (2024)

CC

Expt.

Taniuchi et al. (2019)

Ab initio comparable to mean field computations

Nucleus	Exp.	This work	Ref. [29]	Other
^{80}Zr	1910(180) ^a	1713 ⁺¹¹¹ ₋₁₈₃	2323	3900 ^b
		3044 ⁺¹⁴³ ₋₂₇₄		2540 ^f
^{78}Zr	not known	2040 ⁺¹¹⁸ ₋₂₂₀	2504	
		2927 ⁺¹⁵⁵ ₋₂₈₈		
^{78}Sr	1840(100) ^a	2108 ⁺¹²¹ ₋₂₁₁	1989	2291 ^f
		2519 ⁺¹²⁵ ₋₂₂₈		
^{76}Sr	2390(240) ^a	2444 ⁺¹⁴⁵ ₋₂₄₈	2350	2175 ^f
^{72}Kr	810(150) ^c	1012 ⁺³⁶ ₋₅₀	819	763 ^d
	999(129) ^e	1403 ⁺⁸⁴ ₋₇₇₅		1097 ^f

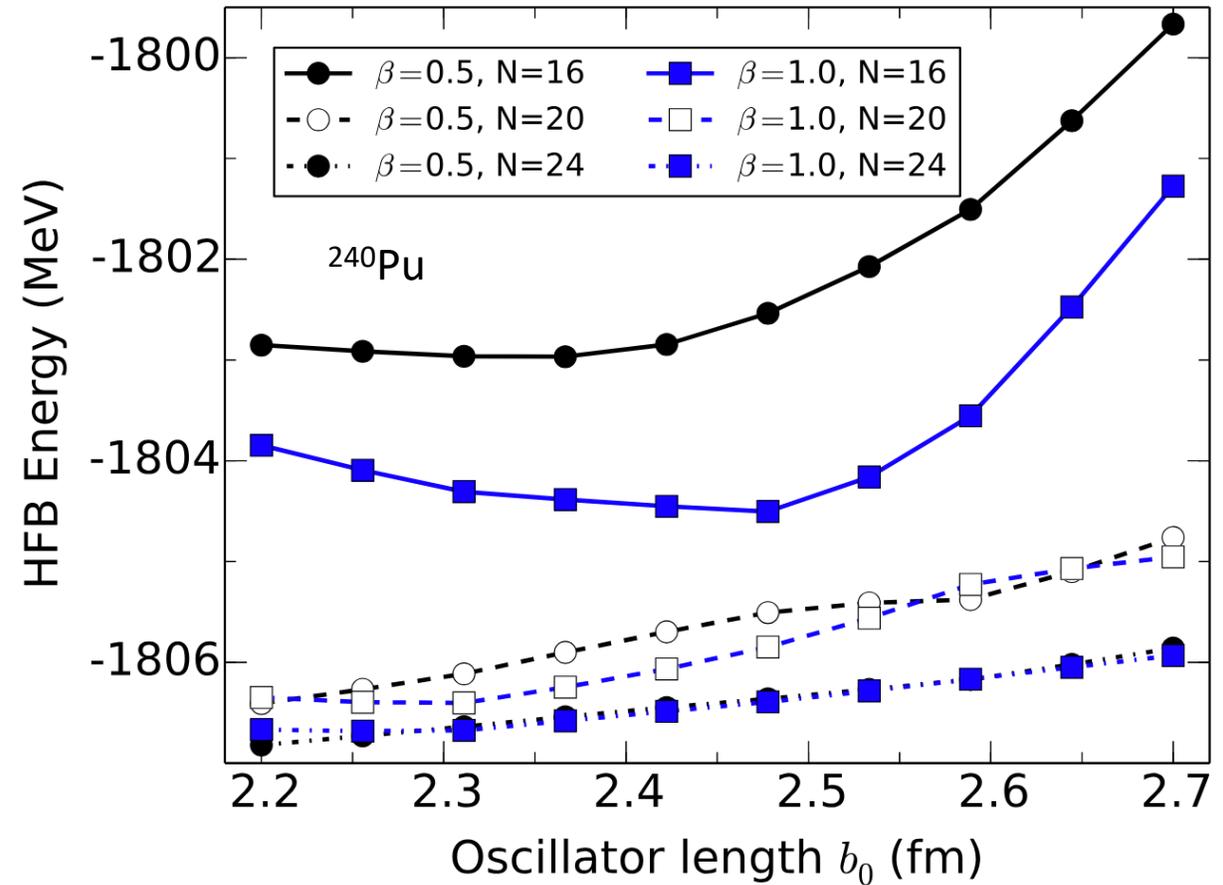
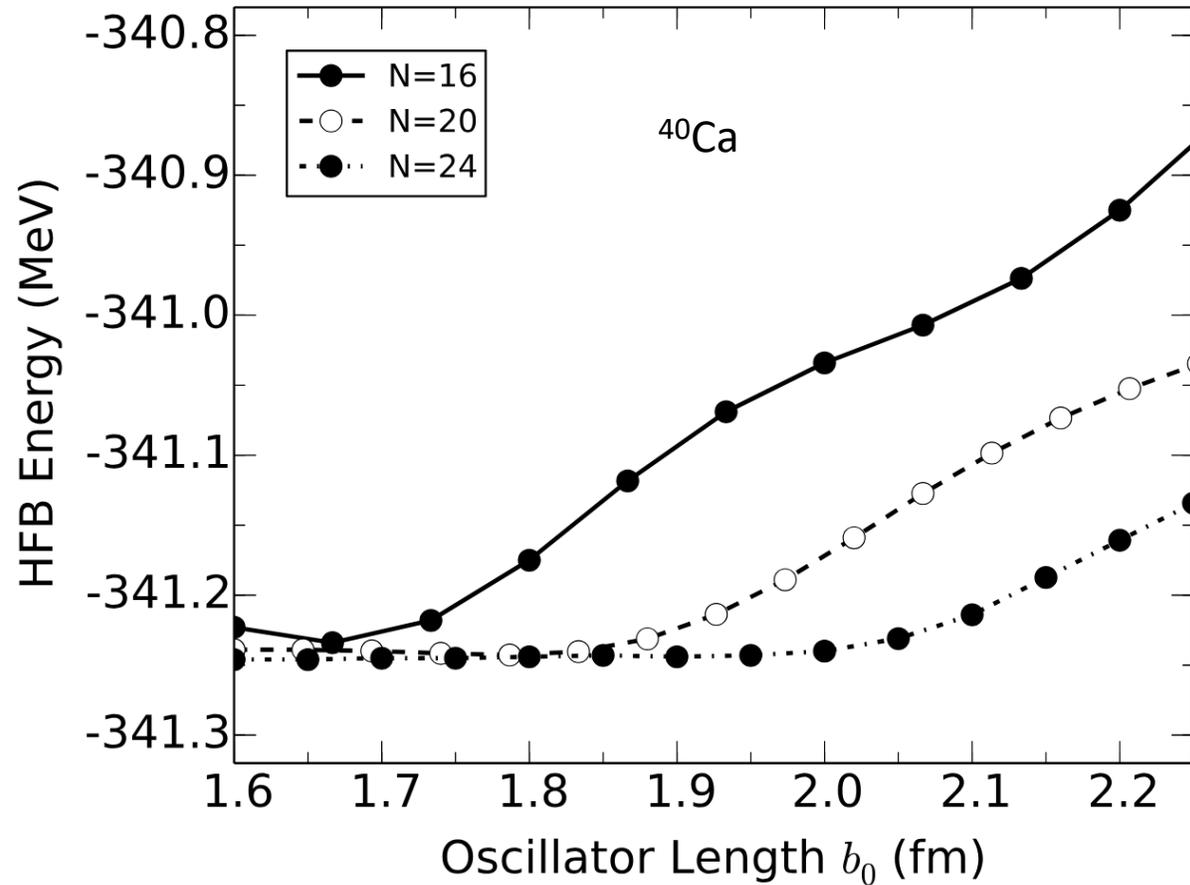
[29] = Delaroche et al, PRC (2010)

b = Rodriguez & Egido, Phys Lett B (2011)

d = Bender, Bonche, Heenen, PRC (2006)

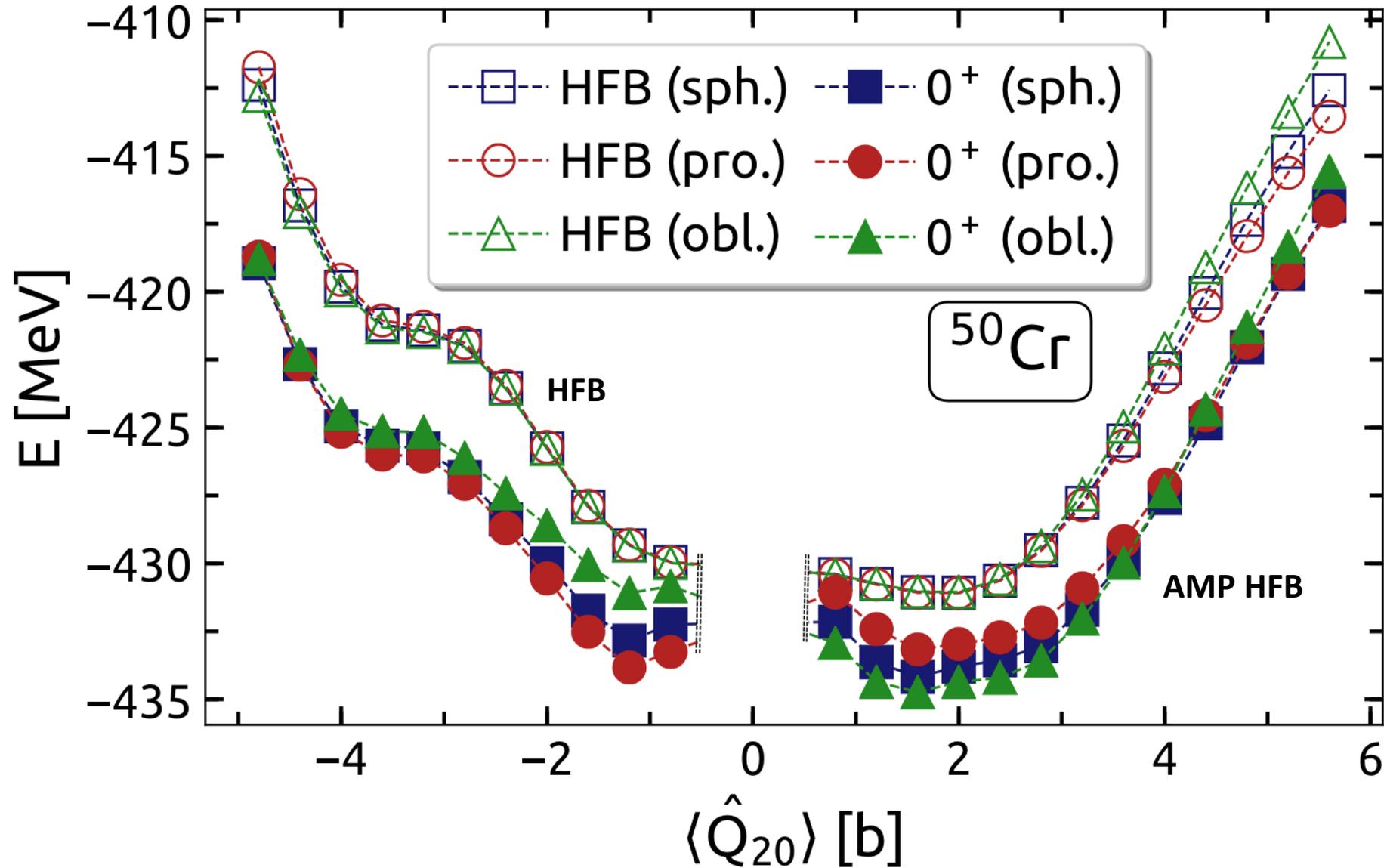
f = Kaneko, Shimizu, Mizusaki, Phys Lett B (2021)

Convergence of mean-field computations

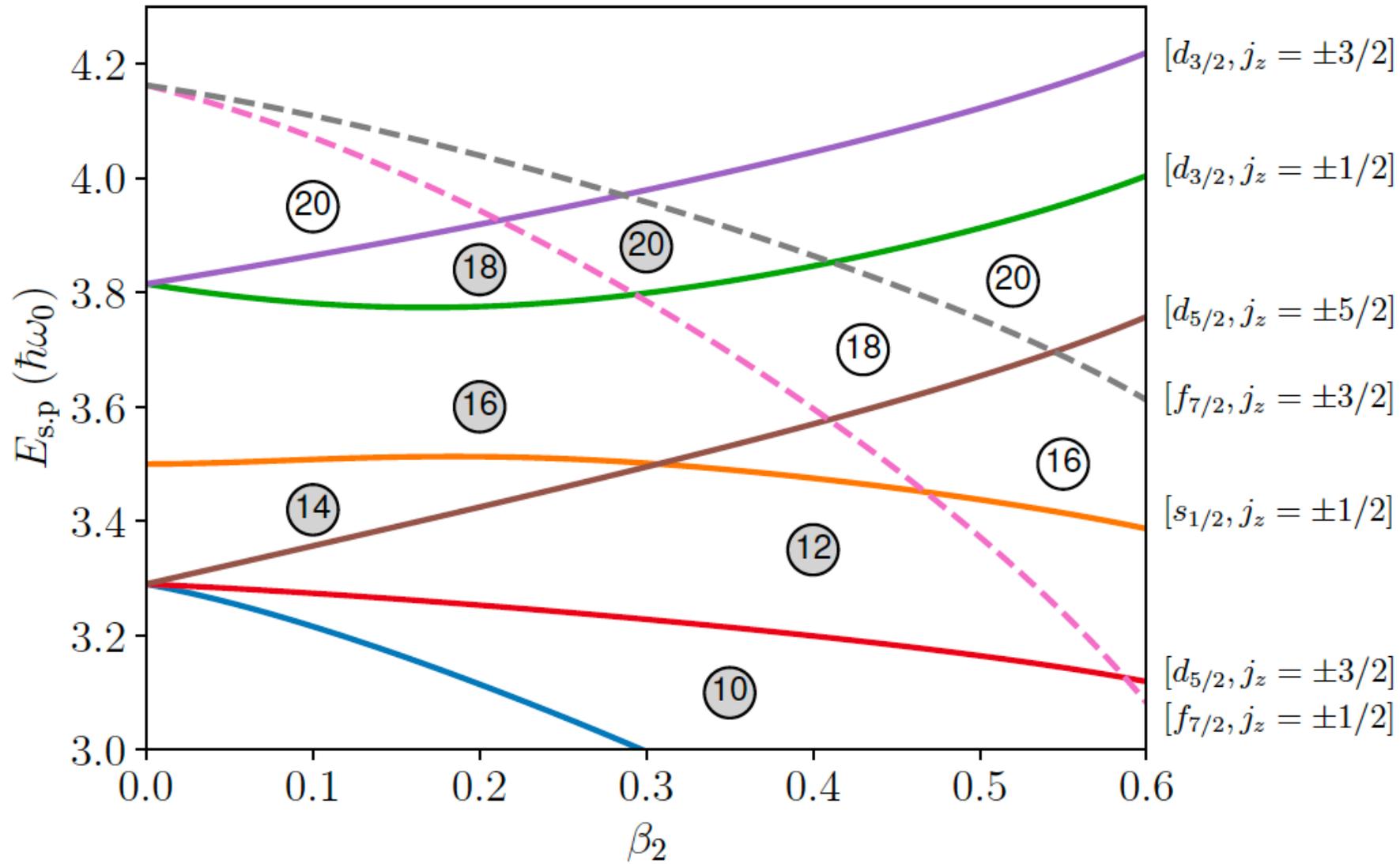


Schunck, McDonnell, Sarich, Wild, Higdon, J Phys G (2015)

Convergence of mean-field computations



Nilsson model reminder



Towards coupled cluster computations of Schiff moments in radium nuclei

