Flavour physics: status and prospects

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QCD Seminars

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Outline:

1. The problem of flavour

2. Open problems in semileptonic and rare B decays

3. A glance into BSM physics

Motivation

Despite the SM successes,
there are open problems:

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Hierarchy problem

dark matter/dark energy

flavour hierarchies

neutrino masses

gravity

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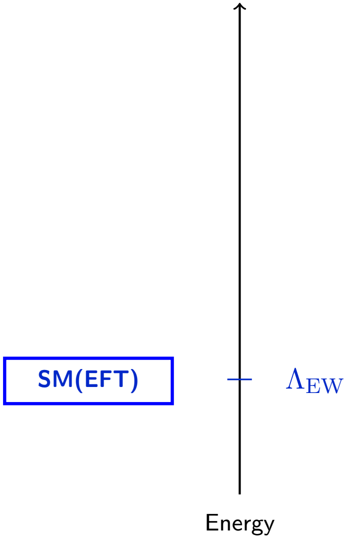

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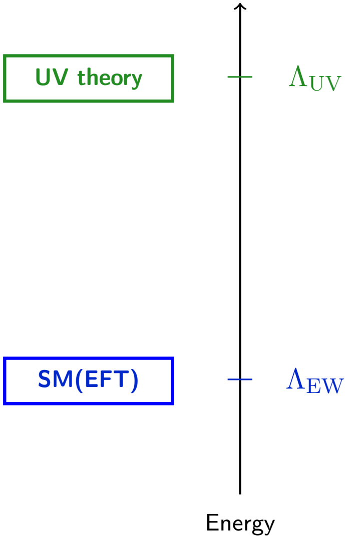

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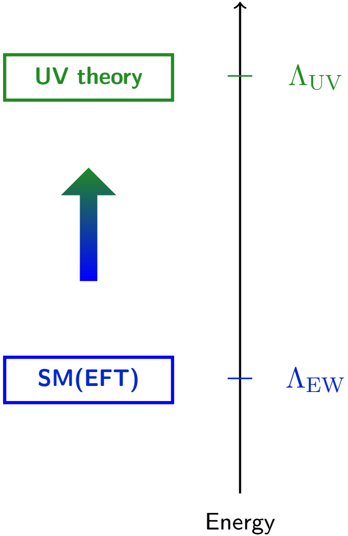


The (two) flavour problems

1. **The SM flavour problem**: The measured Yukawa pattern doesn't seem accidental

   $\Rightarrow$ Is there any deeper reason for that?

2. **The NP flavour problem**: If we regard the SM as an EFT valid below a certain energy cutoff Λ, why don't we see any deviations in flavour changing processes?

   $\Rightarrow$ Which is the flavour structure of BSM physics?

The SM flavour problem

$$\mathcal{L}_{\text{Yukawa}} \supset Y_u^{ij} \bar{Q}_L^i H u_R^j$$

$$Y_u \sim y_t \begin{pmatrix} & & 0.003 \\ & & 0.04 \\ & & 1 \end{pmatrix}$$

The SM flavour problem

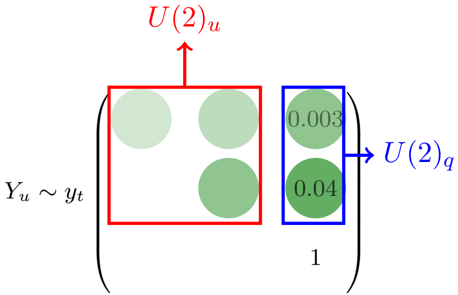
$$\mathcal{L}_{\text{Yukawa}} \supset Y_u^{ij} \bar{Q}_L^i H u_R^j$$

$$Y_u \sim y_t \begin{pmatrix} & & \\ & & \\ & & 1 \end{pmatrix}$$

Exact $U(2)^n$ limit

The SM flavour problem

$$\mathcal{L}_{\text{Yukawa}} \supset Y_u^{ij} \bar{Q}_L^i H u_R^j$$



An approximate $U(2)^n$ is acting on the light families!

The SM flavour problem

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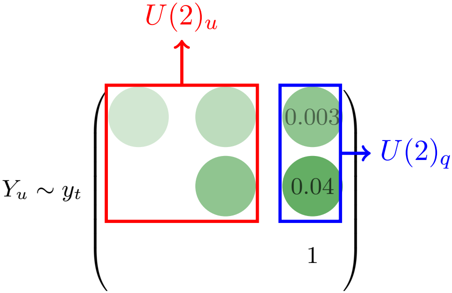


An approximate $U(2)^n$ is acting on the light families!

The NP flavour problem

$$\mathcal{L} = \boxed{\mathcal{L}_{\rm gauge}} + \boxed{\mathcal{L}_{\rm Higgs}}$$

Large Flavour symmetry

Three replica of the same fermion fields

$\boxed{U(3)^5 \text{ symmetry}}$

Flavour degeneracy is broken

The breaking is peculiar

- In the SM: accidental $U(3)^5 \to$ approx $U(2)^n$

The NP flavour problem

$$\mathcal{L} = \boxed{\mathcal{L}_{\text{gauge}}} + \boxed{\mathcal{L}_{\text{Higgs}}} + \boxed{\sum_{d,i} \frac{c_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^d}$$

Large Flavour symmetry

Three replica of the same fermion fields

$U(3)^5$ symmetry

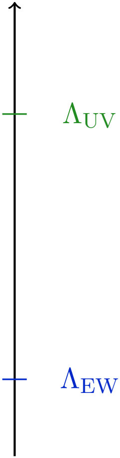
Flavour degeneracy is broken

The breaking is peculiar

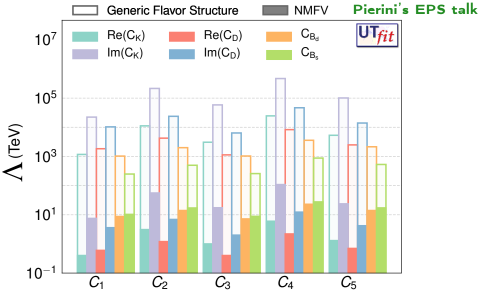
- In the SM: accidental $U(3)^5 \to$ approx $U(2)^n$
- **What happens when we switch on NP?**

The NP flavour problem

$$\mathcal{L} = \boxed{\mathcal{L}_{\text{gauge}}} + \boxed{\mathcal{L}_{\text{Higgs}}} + \boxed{\sum_{d,i} \frac{c_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^d}$$

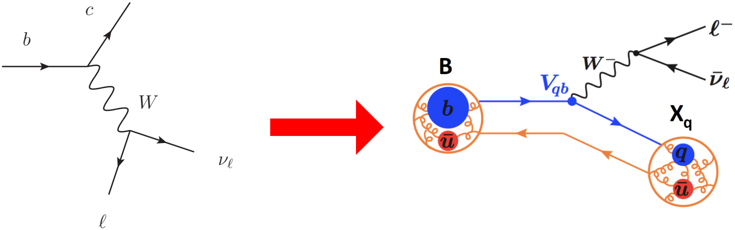


- What is the energy scale of NP?
- Why haven't observed any violation of accidental symmetries yet?



no breaking of the $U(2)^n$ flavour symmetry at low energies

Partonic vs Hadronic



$\mu_{\text{partonic}} = m_b$ $\qquad\qquad$ $\mu_{\text{hadronic}} = \Lambda_{\text{QCD}}$

Fundamental challenge to match partonic and hadronic descriptions

What's the problem for BSM?

B-physics

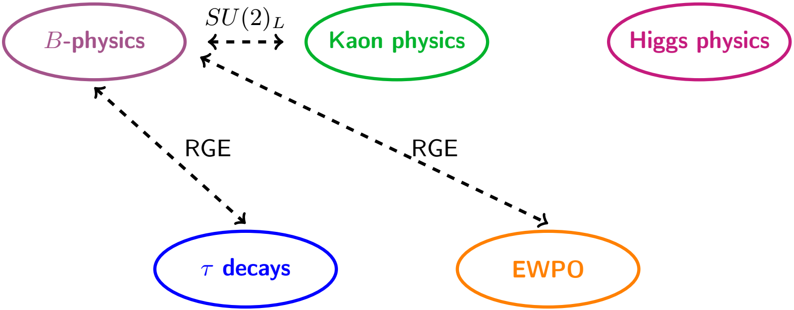
Kaon physics

Higgs physics

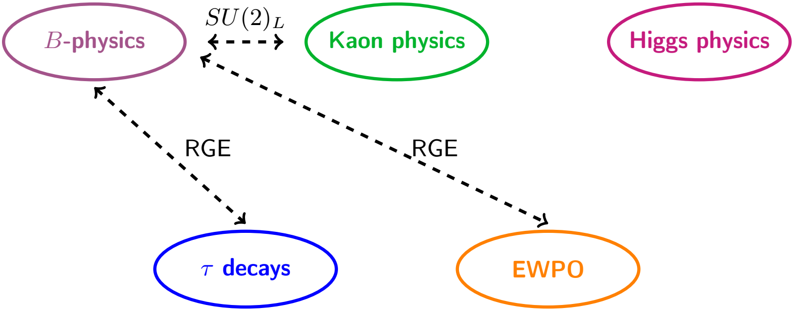
τ decays

EWPO

What's the problem for BSM?



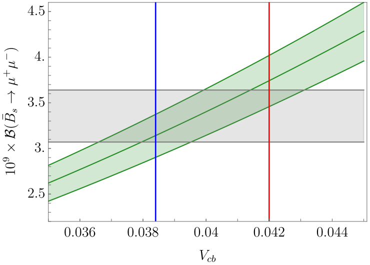
What's the problem for BSM?



How to satisfy all the constraints at the same time?

Semileptonic B decays

Long-standing puzzles in semileptonic decays



Inclusive vs Exclusive determination:

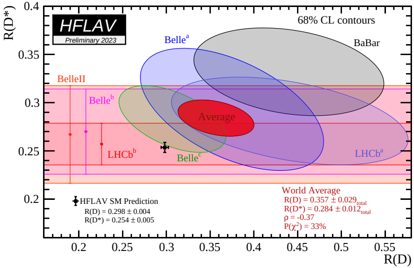
- Inclusive $B \to X_c \ell \bar{\nu}$ decays
- Exclusive decays

$$\Rightarrow \; B \to D^{(*)} \ell \bar{\nu}$$
$$\Rightarrow \; \Lambda_b \to \Lambda_c \mu \bar{\nu} / \Lambda_b \to p \mu \nu$$
$$\Rightarrow \; B_s \to D_s^{(*)} \ell \bar{\nu}$$
$$\Rightarrow \; B_s \to K \mu \nu / B_s \to D_s \mu \nu$$

Lepton flavour universality

$$R_{D^{(*)}} = \frac{\mathcal{B}(B \to D^{(*)} \tau \bar{\nu})}{\mathcal{B}(B \to D^{(*)} \ell \bar{\nu})}$$

- Current discrepancy at the order of 3.3σ
- Theory prediction is the arithmetic average of before 2021 estimates



Inclusive decays

Theory framework for $B \to X_c \ell \bar{\nu}$

Double expansion in $1/m$ and α_s

$$\Gamma_{sl} = \Gamma_0 f(\rho) \Big[1 + a_1 \left(\frac{\alpha_s}{\pi}\right) + a_2 \left(\frac{\alpha_s}{\pi}\right)^2 + a_3 \left(\frac{\alpha_s}{\pi}\right)^3 - \left(\frac{1}{2} - p_1 \left(\frac{\alpha_s}{\pi}\right)\right) \frac{\mu_\pi^2}{m_b^2}$$
$$+ \left(g_0 + g_1 \left(\frac{\alpha_s}{\pi}\right)\right) \frac{\mu_G^2(m_b)}{m_b^2} + d_0 \frac{\rho_D^3}{m_b^3} - g_0 \frac{\rho_{LS}^3}{m_b^3} + \ldots \Big]$$

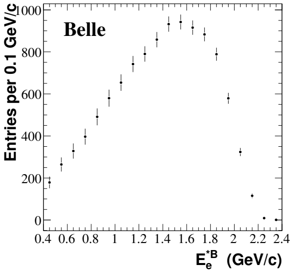
- The coefficients are known
- $\mu_\pi^2(\mu) = \frac{1}{2m_B} \langle B | \bar{b}_v (i\vec{D})^2 b_v | B \rangle_\mu \qquad \mu_G^2(\mu) = \frac{1}{2m_B} \langle B | \bar{b}_v \frac{i}{2} \sigma_{\mu\nu} G^{\mu\nu} b_v | B \rangle_\mu$

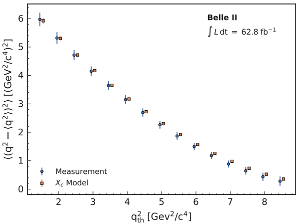
  $\Rightarrow$ No Lattice QCD determinations are available yet
- Use for the first time of α_s^3 corrections **[Fael, Schönwald, Steinhauser, '20]**
- Ellipses stands for higher orders

  $\Rightarrow$ proliferation of terms and loss of predictivity

How do we constrain the hadronic parameters?

We need information from kinematic distributions



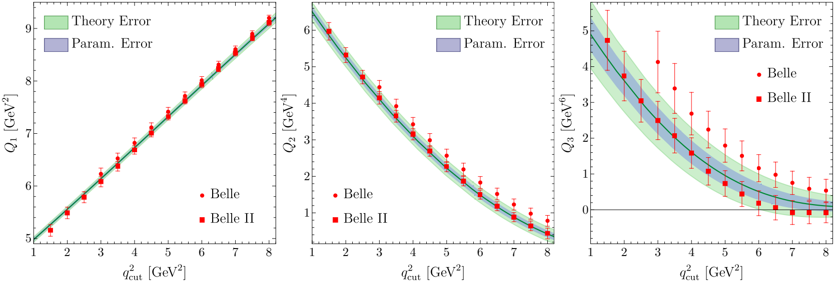


- Traditional method: Extract the hadronic parameters from moments of kinematic distributions in E_l and M_X
- New idea: Use q^2 moments to exploit the reduction of free parameters due to RPI **[Fael, Mannel, Vos, '18, Bernlochner et al, '22]**
- Measurements of branching fractions are needed and are at the moment quite old
- Can we do it on the lattice? **[Gambino, Hashimoto, '20, '23, Hashimoto, Jüttner, et al, '23]**

Global fit

[MB, Capdevila, Gambino, '21, Finauri, Gambino, '23]

	$m_b^{\rm kin}$	$\overline{m}_c$	μ_π^2	μ_G^2	ρ_D^3	ρ_{LS}^3	$10^2\mathrm{BR}_{c\ell\nu}$	$10^3\lvert V_{cb}\rvert$	$\chi^2_{\rm min}$(/dof)
without	4.573	1.092	0.477	0.306	0.185	-0.130	10.66	42.16	22.3
q^2-moments	0.012	0.008	0.056	0.050	0.031	0.092	0.15	0.51	0.474
Belle II	4.573	1.092	0.460	0.303	0.175	-0.118	10.65	42.08	26.4
	0.012	0.008	0.044	0.049	0.020	0.090	0.15	0.48	0.425
Belle	4.572	1.092	0.434	0.302	0.157	-0.100	10.64	41.96	28.1
	0.012	0.008	0.043	0.048	0.020	0.089	0.15	0.48	0.476
Belle &	4.572	1.092	0.449	0.301	0.167	-0.109	10.65	42.02	41.3
Belle II	0.012	0.008	0.042	0.048	0.018	0.089	0.15	0.48	0.559



About QED effects in inclusive decays

Why do we care about QED Effects?

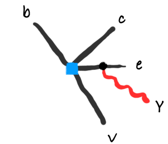
- We want to match the theory description with the experimental measurements that are always affected by photon emissions
- The MC PHOTOS accounts for QED effects, reporting results which can be compared with the non-radiative theory predictions
- PHOTOS knows only about real emission and obtains the virtual part by normalisation

$$\frac{d\Gamma}{dzdx} = \mathcal{F}^{(0)}(\omega_{\text{virtual}} + \omega_{\text{real}}) \Rightarrow \int dx(\omega_{\text{virtual}} + \omega_{\text{real}}) = 1$$

$\boxed{\text{Are virtual corrections under control?}}$

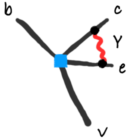
Leading contributions

1. **Collinear logs**: captured by splitting functions



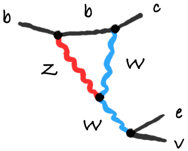
$$\sim \frac{\alpha_e}{\pi} \log^2 \left(\frac{m_b^2}{m_e^2} \right)$$

2. **Threshold effects** or Coulomb terms



$$\sim \frac{4\pi\alpha_e}{9}$$

3. **Wilson Coefficient**



$$\sim \frac{\alpha_e}{\pi} \left[\log \left(\frac{M_Z^2}{\mu^2} - \frac{11}{6} \right) \right]$$

Branching ratio

[Bigi, MB, Gambino, Haisch, Piccione, '23]

- The total branching ratio is not affected by large logs due to KLN theorem
- The large corrections are from the Wilson Coefficient and the threshold effects

$$\frac{\Gamma}{\Gamma^{(0)}g(\rho)} = 1 + \frac{\alpha}{\pi}\left[\ln\left(\frac{M_Z^2}{m_b^2}\right) - \frac{11}{6} + 5.516(14)\right]$$

$$= 1 + 1.43\% - 0.44\% + 1.32\% = 1 + 2.31\%$$

Wilson Coefficient

Threshold effects

- Large shift of the branching ratio of the same order of the current error on V_{cb}
- How do we incorporate in the current datasets?
- Moments are less sensitive because they are normalised

Global fit + QED

[Finauri, Gambino, '23]

- Implementation of QED corrections are analysis dependent
- BaBar provides branching fractions with and without radiation

$$R_{\mathrm{QCD}}^{\mathrm{new}} = \zeta_{\mathrm{QED}} R_{\mathrm{QCD}}^{\mathrm{Babar}}$$

$\Rightarrow$ ζ_{QED} accounts for the misalignment between the corrected BaBar results and the results from the full $\mathcal{O}(\alpha_e)$ computation

m_b^{kin}	$\overline{m}_c(2\,\mathrm{GeV})$	μ_π^2	$\mu_G^2(m_b)$	$\rho_D^3(m_b)$	ρ_{LS}^3	$\mathrm{BR}_{c\ell\nu}$	$10^3\lvert V_{cb}\rvert$
4.573	1.090	0.453	0.288	0.176	−0.113	10.62	41.95
0.012	0.010	0.043	0.049	0.019	0.090	0.15	0.48

- The central value shifts slightly
- Belle II data are needed to understand how to apply the correction
- Can we go beyond scalar QED?

Exclusive decays

Exclusive matrix elements

$$\langle H_c | J_\mu | H_b \rangle = \sum_i S_\mu^i \mathcal{F}_i$$

Exclusive matrix elements

$$\langle H_c | J_\mu | H_b \rangle = \sum_i S^i_\mu \mathcal{F}_i$$

form factor

scale Λ_{QCD}

independent

Lorentz structures

Exclusive matrix elements

$$\langle H_c | J_\mu | H_b \rangle = \sum_i S^i_\mu \mathcal{F}_i$$

form factor

scale Λ_{QCD}

independent Lorentz structures

Form factors determinations

- Lattice QCD
- QCD SR, LCSR

only points at specific kinematic points

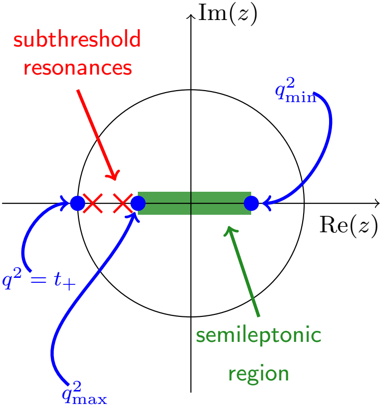
Form factors parametrisations

- HQET (CLN + improvements) $\Rightarrow$ reduce independent degrees of freedom
- Analytic properties $\rightarrow$ BGL

data points needed to fix the coefficients of the expansion

The z-expansion and unitarity

[Boyd, Grinstein, Lebed, '95, Caprini, Lellouch, Neubert, '98]



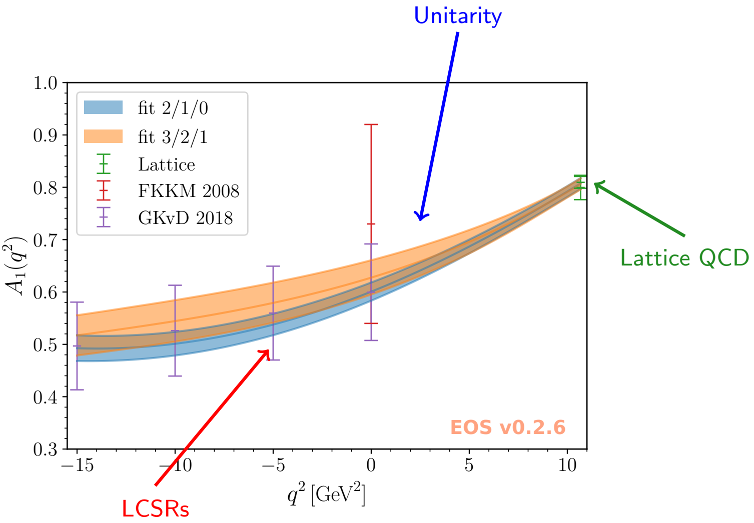
- in the complex plane form factors are real analytic functions
- q^2 is mapped onto the conformal complex variable z

$$z(q^2, t_0) = \frac{\sqrt{t_+ - q^2} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - q^2} + \sqrt{t_+ - t_0}}$$

- q^2 is mapped onto a disk in the complex z plane, where $|z(q^2, t_0)| < 1$

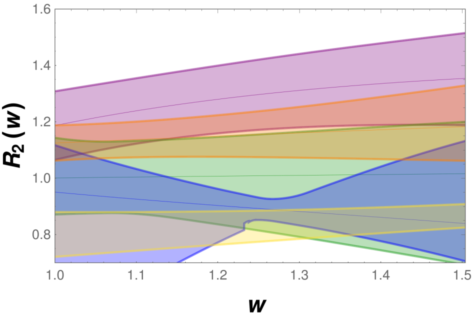
$$F_i = \frac{1}{P_i(z)\phi_i(z)} \sum_{k=0}^{n_i} a_k^i z^k$$

$$\sum_{k=0}^{n_i} |a_k^i|^2 < 1$$



Unitarity
1.0
0.9
0.8
0.7
0.6
0.5
0.4
0.3
$A_1(q^2)$
fit 2/1/0
fit 3/2/1
Lattice
FKKM 2008
GKvD 2018
Lattice QCD
EOS v0.2.6
−15
−10
−5
0
5
10
q^2 [GeV^2]
LCSRs

$B \to D^*$ after 2021

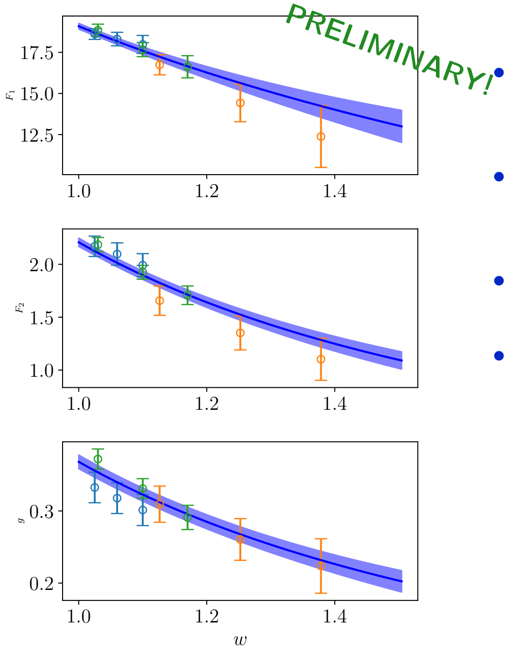


- FNAL/MILC '21
- HQE@$1/m_c^2$
- Exp data (BGL)
- JLQCD '23
- HPQCD '23

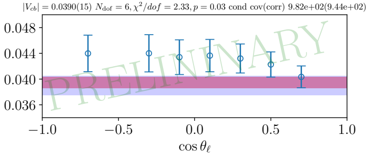
- Are the Lattice QCD datasets compatible?
- What's the source of the discrepancy with HQET? [MB, Harrison, Jung, ongoing]
- Why are experimental data so different?

What can we learn?

[MB, Jüttner, Tsang, in preparation]



- Combining Lattice QCD results in a BGL fit is possible
- Unitarity is essential to contain uncertainties [Flynn, Jüttner, Tsang, '23]
- Difference in slope is the real issue
- Pheno still ongoing, not all kinematic distribution yield a good fit for V_{cb}

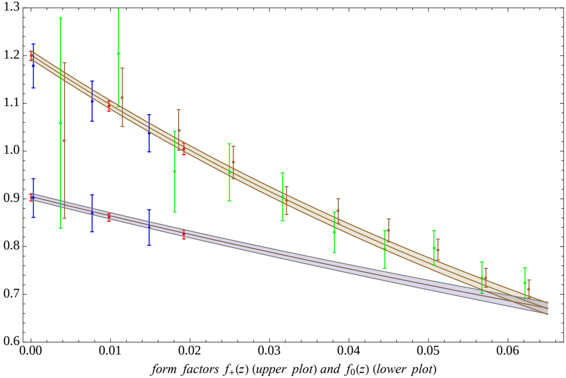


[See also:2310.03680]

$B \to D$

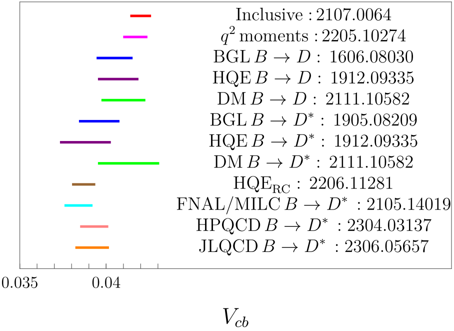
[Bigi, Gambino, '16]

- Belle+Babar data and HPQCD+FNAL/MILC Lattice points



$$|V_{cb}| = (40.49 \pm 0.97) \times 10^{-3}$$

Pheno Status 1

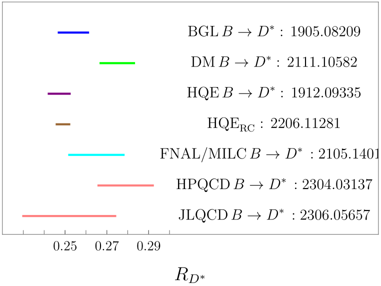


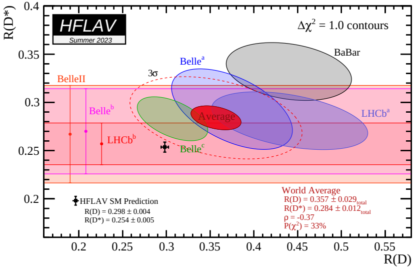
- The inclusive determination is solid
- No evident issues for $B \to D$
- Spread between inclusive and exclusive up to $3 - 4\sigma$
- Work in progress for the theory predictions of $B \to D^*$ to understand the various tensions

  $\Rightarrow$ Do we have to correct for QED?

- New experimental data are available are under scrutiny

Pheno status 2





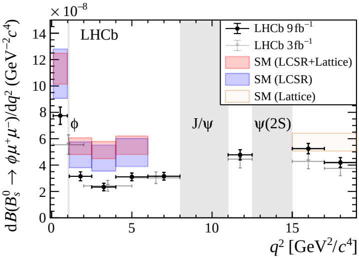
- New Lattice QCD results point to larger values for R_{D^*}

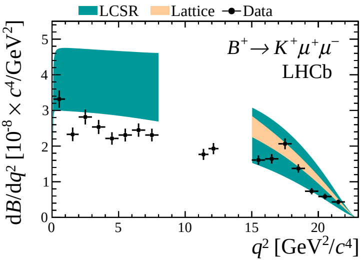
  $\Rightarrow$ Difference in the slopes is crucial and has to be understood

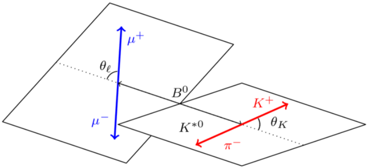
- No change in R_D, where Lattice QCD results, LCSRs, HQET and experimental data agree very well with each other

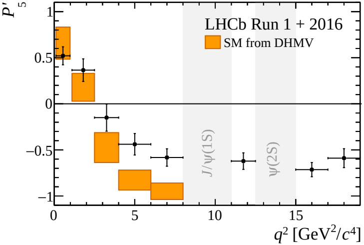
Rare B decays

“Anomalies” in $b \to s\mu^+\mu^-$ transitions



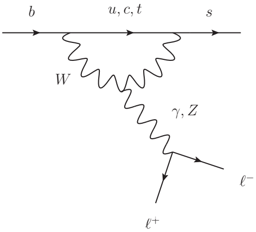






EFT for b decays

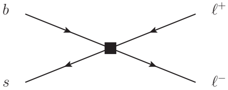
$\Lambda \sim m_t$



$$\mathcal{A}(H_b \to H_s) = \langle H_s | \mathcal{L}_{\text{SM}} | H_b \rangle$$

matching
and
running

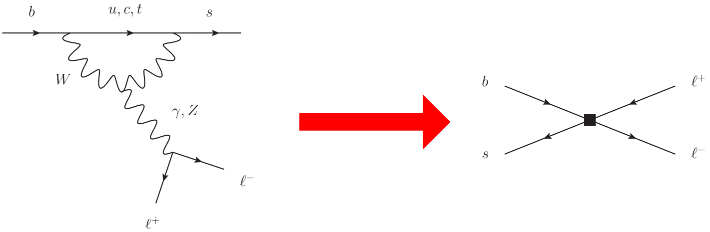
$\Lambda \sim m_b$



$$\mathcal{A}(H_b \to H_s) = \frac{4G_F}{\sqrt{2}} \sum \mathcal{C}_i(\mu) \langle H_s | \mathcal{O}_i | H_b \rangle$$

Energy (Λ)

$b \to s\ell\ell$



$$\mathcal{H}_{\text{eff}} = -4\frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \left[-\mathcal{C}_1 \mathcal{O}_1 - \mathcal{C}_2 \mathcal{O}_2 + \mathcal{C}_7 \mathcal{O}_7 + \mathcal{C}_9 \mathcal{O}_9 + \mathcal{C}_{10} \mathcal{O}_{10} \right]$$

$$\mathcal{O}_1 = (\bar{s}\gamma^\mu P_L b)\,(\bar{c}\gamma_\mu c) \qquad \mathcal{O}_2 = (\bar{s}\gamma^\mu T^a P_L b)\,(\bar{c}\gamma_\mu T^a c)$$

$$\mathcal{O}_9 = (\bar{s}\gamma^\mu P_L b)\,(\bar{\ell}\gamma_\mu \ell) \qquad \mathcal{O}_{10} = (\bar{s}\gamma^\mu P_L b)\,(\bar{\ell}\gamma_\mu \gamma_5 \ell)$$

$$\mathcal{O}_7 = (\bar{s}\sigma^{\mu\nu} P_R b)\,F_{\mu\nu}$$

- Wilson coefficients are calculated at NNLO

Gorbahn, Haisch, '04, Bobeth, Gambino, Gorbahn, Haisch, '11

- The running to $\mu = m_b$ is known

$B \to K^{(*)}\ell^+\ell^-$

$$\mathcal{A}_\lambda^{L,R} = \mathcal{N}_\lambda \left\{ (\mathcal{C}_9 \mp \mathcal{C}_{10})\mathcal{F}_\lambda(q^2) + \frac{2m_b M_B}{q^2} \left[\mathcal{C}_7 \mathcal{F}_\lambda^T(q^2) - 16\pi^2 \frac{M_B}{m_b} \mathcal{H}_\lambda(q^2) \right] \right\}$$

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local: $\mathcal{O}_9, \mathcal{O}_{10}, \mathcal{O}_7$

$B \to K^{(*)}\ell^+\ell^-$

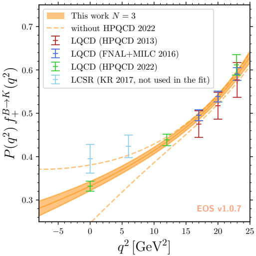
non-local: $\mathcal{O}_1, \mathcal{O}_2$

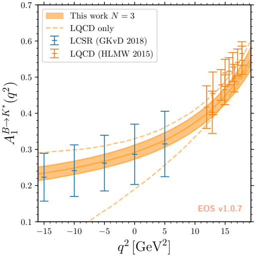
$$\mathcal{A}_\lambda^{L,R} = \mathcal{N}_\lambda \left\{ (\mathcal{C}_9 \mp \mathcal{C}_{10}) \mathcal{F}_\lambda(q^2) + \frac{2 m_b M_B}{q^2} \left[\mathcal{C}_7 \mathcal{F}_\lambda^T(q^2) - 16\pi^2 \frac{M_B}{m_b} \mathcal{H}_\lambda(q^2) \right] \right\}$$

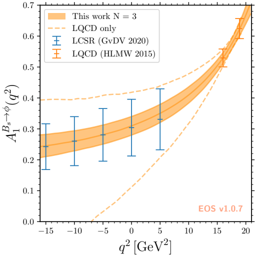
local: $\mathcal{O}_9, \mathcal{O}_{10}, \mathcal{O}_7$

Local form factors

$$\mathcal{F}_\lambda^{(T)} = \langle K^{(*)}(k)|\bar{s}\Gamma_\lambda^{(T)} b|\bar{B}(k+q)\rangle$$



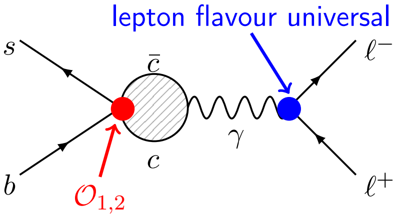




See 2305.06301 for details
Other references: 1503.05534, 1811.00983

Non-local form factors

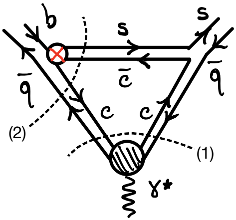
$$\mathcal{H}_\lambda(q^2) = iP_\mu^\lambda \int d^4x\, e^{iqx} \langle K^{(*)}(k)| T\left\{\mathcal{J}^\mu_{\mathsf{em}}, C_i\mathcal{O}_i(0)\right\}|\bar{B}(k+q)\rangle$$



$$C_9 \to C_9^{\text{eff}}(q^2) = C_9 + C_9^{\text{LD}}(q^2)$$

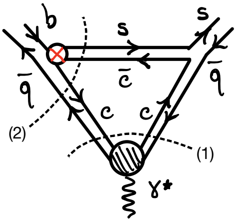
How do we parametrise these long-distance effects?

Charm-loop effects in $b \to s\ell^+\ell^-$



- Conformal transformation $q^2 \mapsto z(q^2)$, with $|z| < 1$
- $C_9^{\text{LD}} \propto \alpha_n z^n$ [1707.07305]
- Dispersive analysis allow to determine the truncation order of the series [2011.09813,2206.03797]
- Effects are **small**

Charm-loop effects in $b \to s\ell^+\ell^-$



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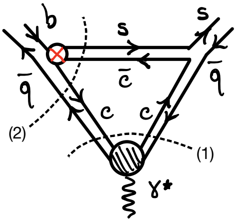
Is this all?

Charm-loop effects in $b \to s\ell^+\ell^-$



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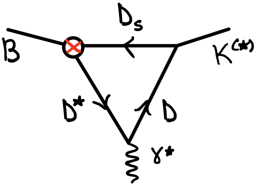


[2212.10516]

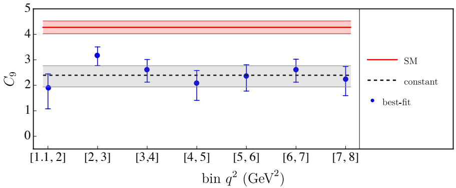
- Are these contributions included?
- Are they large that they can reconcile the tension in $B \to K^*\mu\mu$?

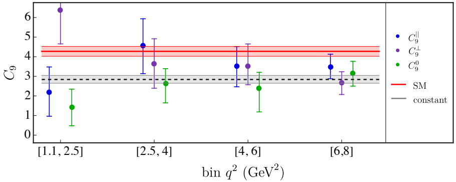
Charm loop effects in $B \to K^{(*)}\mu^+\mu^-$

MB, Isidori, Maechler, Tinari, 2401.18007

- Can we extract some hints of the shape of $C_9^{\rm LD}(q^2)$ from data?
  - $\Rightarrow$ NP yields a **constant** effect in the whole kinematic region
- Is the current sensitivity enough to claim anything?

$$C_9^{\rm eff} = C_9 + \sum_V \eta_V^\lambda e^{i\delta_V^\lambda} \frac{q^2}{(m_V^2)} \frac{m_V \Gamma_V}{m_V^2 - q^2 - i m_V \Gamma_V}$$

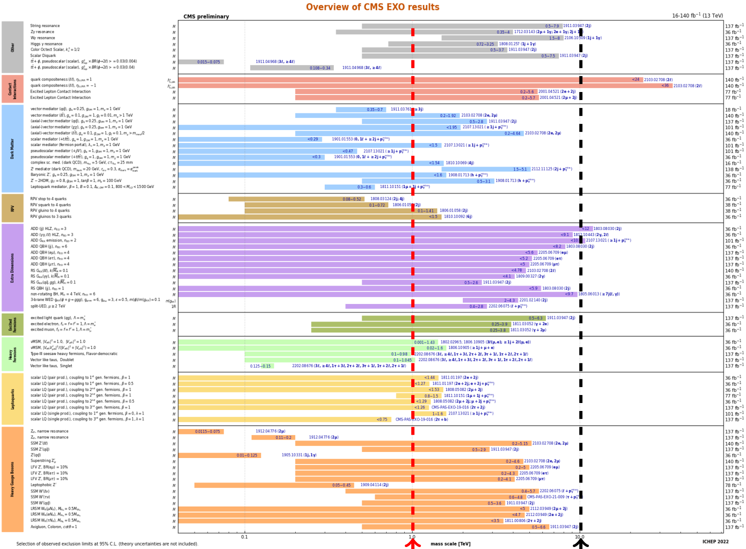




No evidence for q^2 dependence

A glance into BSM physics

Status of high energy bounds

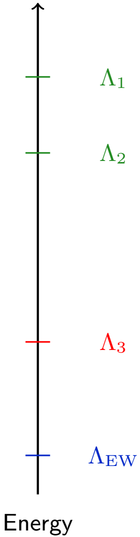


3rd generation

universal new physics

Flavour Non-Universal New Physics

Dvali, Shifman, '00
Panico, Pomarol, '16
MB, Cornella, Fuentes-Martin, Isidori '17
Allwicher, Isidori, Thomsen '20
Barbieri, Cornella, Isidori, '21
Davighi, Isidori '21



Basic idea:

- 1st and 2nd have small masses and small couplings to NP because they are generated by dynamics at a heavier scale
- 3rd generation is linked to dynamics at lower scales and has stronger couplings

Flavour deconstruction:
fermion families interact with different gauge groups and flavour hierarchies emerge as accidental symmetries

Flavour Non-Universal New Physics

Dvali, Shifman, '00
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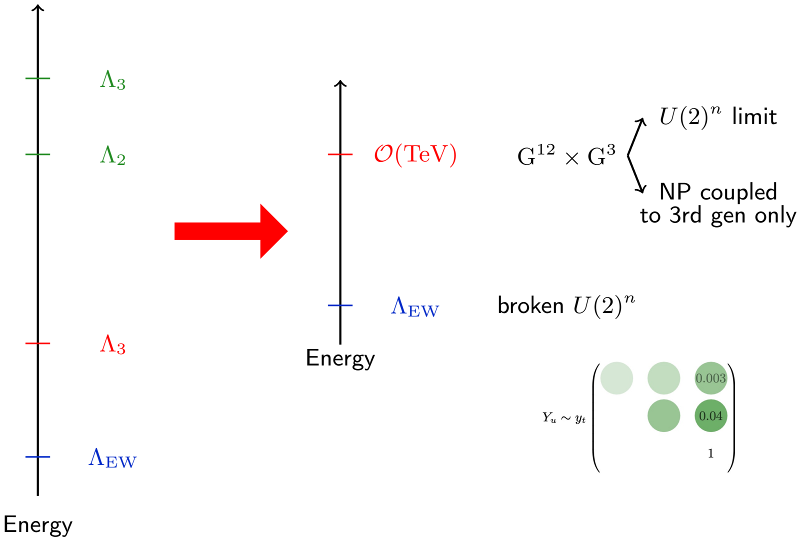
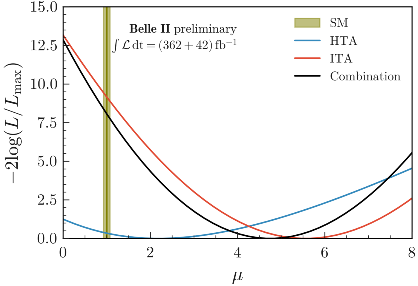
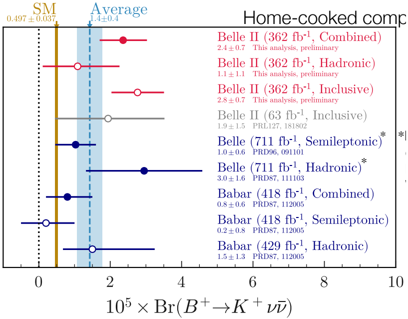


$B^+ \to K^+\nu\bar{\nu}$ from Belle II





- First evidence of the $B^+ \to K^+\nu\bar{\nu}$ process at 3.6σ with

$$\mathcal{B}(B^+ \to K^+\nu\bar{\nu}) = (2.4 \pm 0.5(\mathrm{stat})^{+0.5}_{-0.4}(\mathrm{syst})) \times 10^{-5}$$

- Tension with the SM of $\sim 2.8\sigma$

What do we expect in the SMEFT?

$$\mathcal{L}_{\rm EFT} \supset \boxed{\frac{C_{bc\tau\tau}}{\Lambda^2}}(\bar{b}^i_L \gamma_\nu c^j_L)(\bar{\nu}_\tau \gamma^\mu \tau_L)$$

From $U(2)^n \Rightarrow C_{bc\tau\tau} \sim V_{cb}\mathcal{O}(1)$

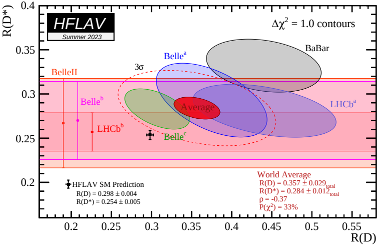
From $R_{D^{(*)}} \Rightarrow \Lambda \sim \mathcal{O}(\mathrm{TeV})$

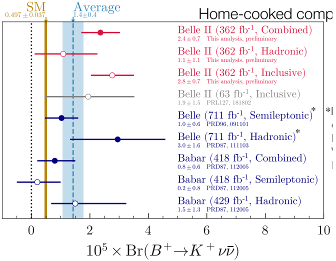
Using $SU(2)_L$ invariance, we have

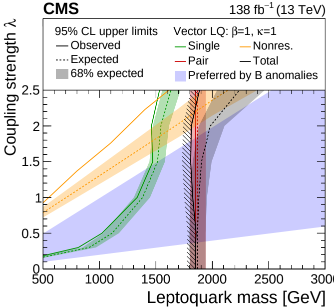
$$\mathcal{L}_{\rm EFT} \supset \boxed{\frac{C_{bs\tau\tau}}{\Lambda^2}}(\bar{b}^i_L \gamma_\nu s^j_L)(\bar{\nu}_\tau \gamma^\mu \nu_\tau)$$

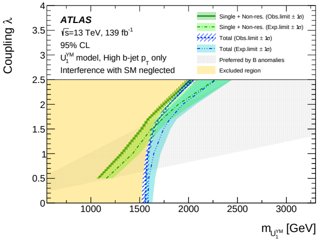
From $U(2)^n \Rightarrow C_{bs\tau\tau} \sim V_{cb}\mathcal{O}(1)$

Belle II measurement of $B \to K\nu\bar{\nu}$
in agreement with $U(2)^n$









The present hints align well together, but it is too soon to claim victory...

Conclusions

- Flavour physics is a powerful test for new physics living at different energy scales

- At the current status, we haven't observed any clear sign of new physics

- No clear sign of new physics can hint to a peculiar structure for the flavour structure of NP and to flavour deconstruction

  $\Rightarrow$ Theoretical and Experimental efforts will shed light on puzzles in hadronic predictions, aiming to a deeper understanding of the SM

  $\Rightarrow$ From the phenomenological point of view, a few hints point to a strong link between new physics and the third generations, with possible new physics reach close to the current searches

Appendix

Measuring V_{cb}

Interaction basis

$$-\mathcal{L}_{\mathrm{Y}} = Y_d^{ij} \bar{Q}_L^i H d_R^j + Y_u^{ij} \bar{Q}_L^i \tilde{H} u_R^j + \mathrm{h.c.}$$

Non-diagonal Yukawa

Mass basis

$$\mathcal{L}_{cc} \propto \bar{u}_L^i \gamma^\mu d_L^j W_\mu^+ V_{ij}$$

CKM matrix

V_{cb} extraction

$$\mathcal{O}_{\mathrm{theory}}(V_{cb}, \vec{\mu}) = \mathcal{O}_{\mathrm{exp}}$$

theory inputs needed

$B \to D^{(*)}$ form factors

- 7 (SM) + 3 (NP) form factors
- Lattice computation for $q^2 \neq q^2_{\text{max}}$ only for $B \to D$
- Calculation usually give only a few points
- q^2 dependence must be inferred
- Conformal variable z

$$z(q^2, t_0) = \frac{\sqrt{t_+ - q^2} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - q^2} + \sqrt{t_+ - t_0}}$$

- $t_+ = (m_B + m_{D^{(*)}})^2$ pair production threshold
- $t_0 < t_+$ free parameter that can be used to minimise $|z_{\text{max}}|$
- $|z| \ll 1$, in the $B \to D$ case $|z| < 0.06$

The HQE parametrisation 1

- Expansion of QCD Lagrangian in $1/m_{b,c} + \alpha_s$ corrections

[Caprini, Lellouch, Neubert, '97]

- In the limit $m_{b,c} \to \infty$: all $B \to D^{(*)}$ form factors are given by a **single** Isgur-Wise function

$$F_i \sim \xi$$

- at higher orders the form factors are still related $\Rightarrow$ **reduction** of free parameters

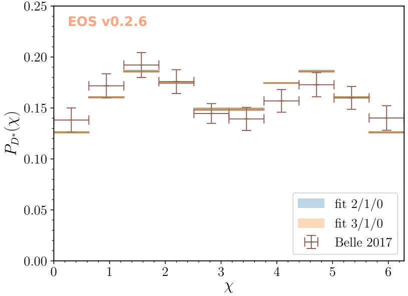
$$F_i \sim \left(1 + \frac{\alpha_s}{\pi}\right) \xi + \frac{\Lambda_{\mathsf{QCD}}}{2m_b} \xi^i_{\mathsf{SL}} + \frac{\Lambda_{\mathsf{QCD}}}{2m_c} \xi^i_{\mathsf{SL}}$$

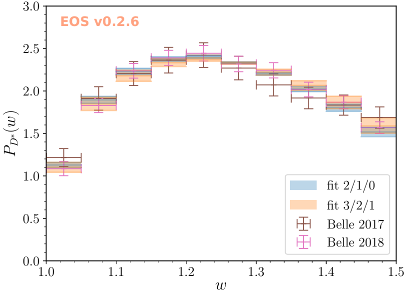
  - at this order 1 leading and 3 subleading functions enter
  - ξ^i are not predicted by HQE, they have to be determined using some other information

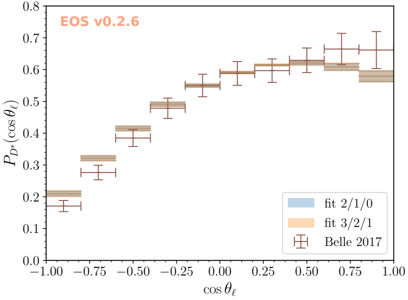
The HQE parametrisation 2

- Important point in the HQE expansion: $q^2 = q^2_{\text{max}}$
- At this point Luke's Theorem applies: the subleading corrections vanish for some form factors
- The leading Isgur-Wise function is normalised: $\xi(q^2 = q^2_{\text{max}}) = 1$
- **Problem**: contradiction with lattice data!
- $1/m_c^2$ corrections **have to be systematically included** **[Jung, Straub, '18, MB, M.Jung, D.van Dyk, '19]**
  - well motivated also since $\alpha_s/\pi \sim 1/m_b \sim 1/m_c^2$

Comparison with kinematical distributions



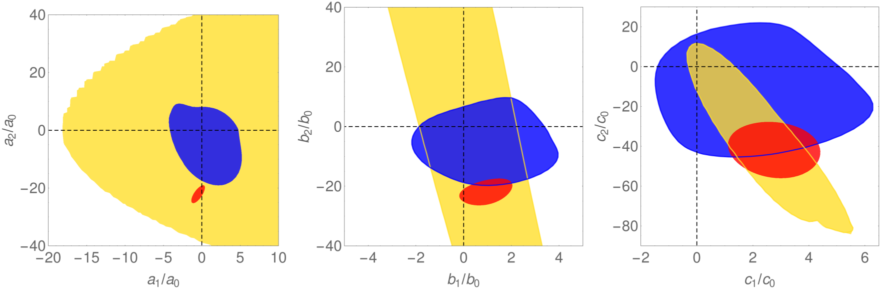




good agreement with kinematical distributions

Fit stability

- BGL fit to Belle 2017 and 2018 data (yellow)
- HQE fit 2/1/0 (red)
- HQE fit 3/2/1 (blue)



- compatibily of HQE fit with data driven one
- 2/1/0 underestimates massively uncertainties

$\boxed{3/2/1 \text{ is our nominal fit}}$

HQET in a nutshell

- In HQET it is convenient to work with **velocities** instead of momenta
- Instead of q^2 we use the dimensionless variable $w = v_B \cdot v_{D^*}$
- When the $B(b)$ decays such that the $D^*(c)$ is at rest in the $B(b)$ frame

$$v_B = v_{D^*} \qquad \Rightarrow \qquad w = 1$$

- The brown muck doesn't realise that anything changed
- At zero recoil, the leading IW function is **normalized**

$$\xi(w = 1) = 1$$

BGL vs CLN parametrisations

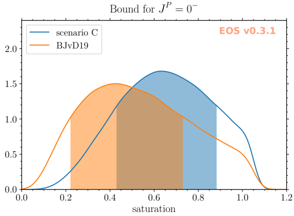
CLN [Caprini, Lellouch, Neubert, '97]

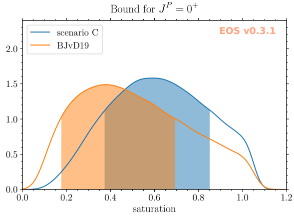
- Expansion of FFs using HQET
- $1/m_{b,c}$ corrections included
- Expansion of leading IW function up to 2nd order in $(w-1)$

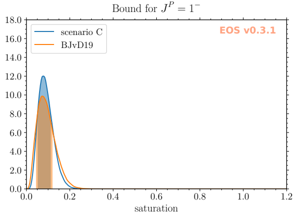
BGL [Boyd, Grinstein, Lebed, '95]

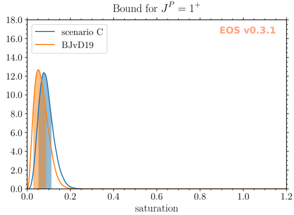
- Based on analyticity of the form factors
- Expansion of FFs using the conformal variable z
- Large number of free parameters

Results: unitary bounds

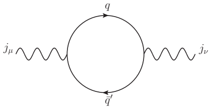








Unitarity Bounds



$$= i \int d^4x\, e^{iqx} \langle 0 | T\left\{ j_\mu(x), j_\nu^\dagger(0) \right\} |0\rangle = (g_{\mu\nu} - q_\mu q_\nu)\Pi(q^2)$$

- If $q^2 \ll m_b^2$ we can calculate $\Pi(q^2)$ via perturbative techniques $\Rightarrow \chi(0)$
- Dispersion relations link $\mathrm{Im}\left(\Pi(q^2)\right)$ to sum over matrix elements

$$\sum_i |F_i(0)|^2 < \chi(0)$$

[Boyd, Grinstein,Lebed, '95
Caprini, Lellouch, Neubert, '97]

- The sum runs over <u>all</u> possible states hadronic decays mediated by a current $\bar{c}\Gamma_\mu b$
  - The unitarity bounds are more effective the most states are included in the sum
  - The unitarity bounds introduce correlations between FFs of different decays
  - $B_s \to D_s^{(*)}$ decays are expected to be of the same order of $B_{u,d} \to D_{u,d}^{(*)}$ decays due to $SU(3)_F$ simmetry

Theory framework

$$\Gamma = \frac{1}{m_B}\mathrm{Im}\int d^4x \langle B(p)|T\left\{\mathcal{H}^\dagger_{\mathrm{eff}}(x)\mathcal{H}_{\mathrm{eff}}(0)\right\}|B(p)\rangle$$

Theory framework

$$\Gamma = \frac{1}{m_B} \mathrm{Im} \int d^4x \langle B(p) | T \left\{ \mathcal{H}^\dagger_{\mathrm{eff}}(x) \mathcal{H}_{\mathrm{eff}}(0) \right\} | B(p) \rangle$$

$$\uparrow$$

$$\sum_{n,i} \frac{1}{m_b^n} \mathcal{C}_{n,i} \mathcal{O}_{n+3,i}$$

Theory framework

$$\Gamma = \frac{1}{m_B}\mathrm{Im}\int d^4x \langle B(p)|T\left\{\mathcal{H}^\dagger_{\mathrm{eff}}(x)\mathcal{H}_{\mathrm{eff}}(0)\right\}|B(p)\rangle$$

$$\sum_{n,i} \frac{1}{m_b^n}\mathcal{C}_{n,i}\mathcal{O}_{n+3,i}$$

- The Wilson coefficients are calculated perturbatively
- The matrix elements $\langle B(p)|\mathcal{O}_{n+3,i}|B(p)\rangle$ are non perturbative
  - $\Rightarrow$ They need to be determined with non-perturbative methods, e.g. Lattice QCD
  - $\Rightarrow$ They can be extracted from data
  - $\Rightarrow$ With large n, large number of operators

Theory framework

$$\Gamma = \frac{1}{m_B} \mathrm{Im} \int d^4x \langle B(p)| T\left\{\mathcal{H}_{\mathrm{eff}}^\dagger(x)\mathcal{H}_{\mathrm{eff}}(0)\right\}|B(p)\rangle$$

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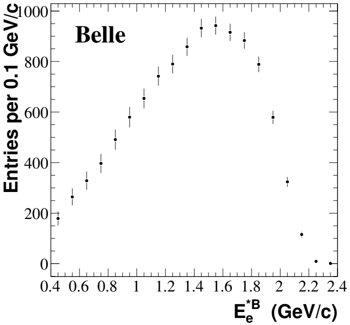
loss of predictivity

How do we constrain the OPE parameters?



- Lepton energy and hadronic invariant mass distributions can be used to extract non perturbative information
- Moments of the kinematic distributions

$$\langle E_\ell^n \rangle = \frac{\int_{E_\ell > E_{\ell,\mathrm{cut}}} dE_\ell E_\ell^n \frac{d\Gamma}{dE_\ell}}{\Gamma_{E_\ell > E_{\ell,\mathrm{cut}}}}$$

$$R^* = \frac{\int_{E_\ell > E_{\ell,\mathrm{cut}}} dE_\ell \frac{d\Gamma}{dE_\ell}}{\int dE_\ell \frac{d\Gamma}{dE_\ell}}$$

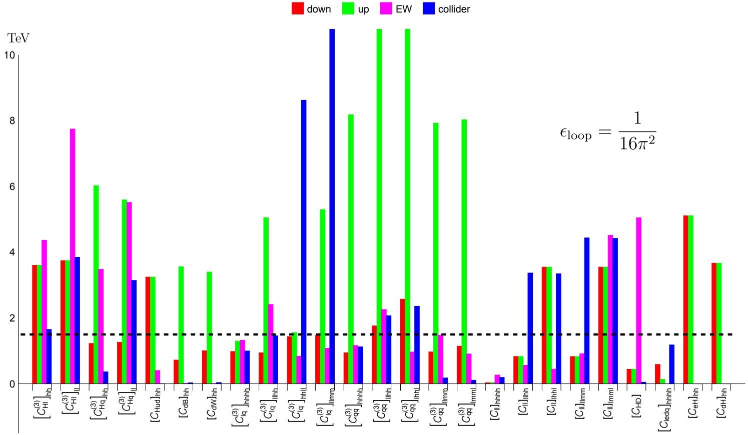
- Similar definition for hadronic mass moments

- The moments give access to the distribution, but not to the normalisation
- They admit an HQE as the rate

$\Rightarrow$ No $\mathcal{O}(\alpha_s^3)$ terms are known yet

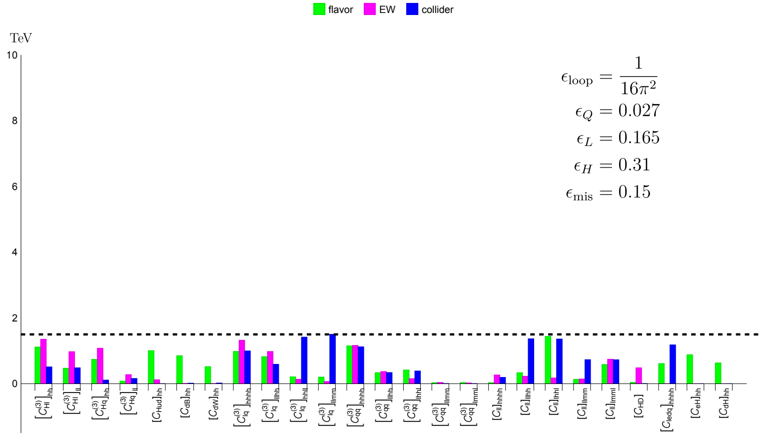
SMEFT with Flavour 1

[Allwicher, Cornella, Isidori, Stefanek, in preparation]



SMEFT with Flavour 2

[Allwicher, Cornella, Isidori, Stefanek, in preparation]

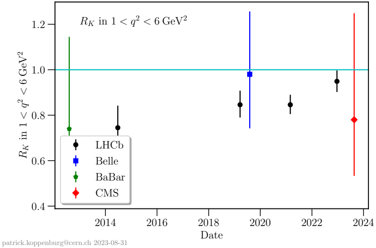


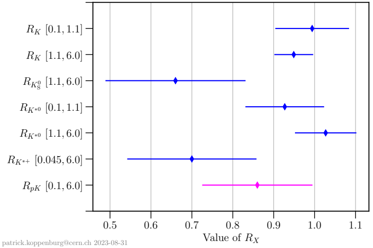
Lepton Flavour Universality violation

$$R_X = \frac{\mathcal{B}(H_b \to X\mu^+\mu^-)}{\mathcal{B}(H_b \to Xe^+e^-)}$$

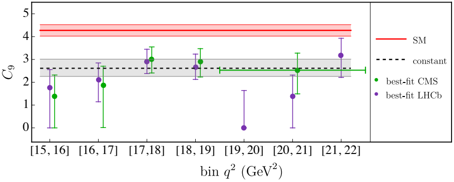
- Test of Lepton Flavour Universality, which is one of the building principles of the SM
- With ratios, we reduce hadronic uncertainties at large extent
- For $q^2 \gg m_\ell^2 \to R_X = 1$
- Leading theoretical uncertainty coming from QED effects $\sim 1\%$

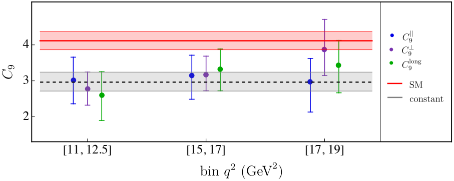
MB, Isidori, Pattori, '16
Isidori, Lancerini, Nabeebaccus, Zwicky, '22

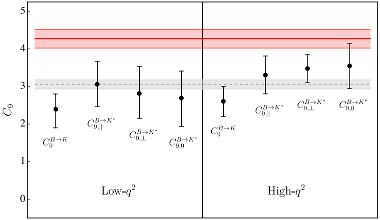




C_9 from $B \to K^{(*)}\mu^+\mu^-$ data





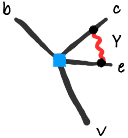


The inclusive case

- If wrt QCD the hadronic and leptonic system are separated, QED corrections mix them

  $\Rightarrow$ Defining fully inclusive observables is harder

  $\Rightarrow$ Analogy with experiments is essential



- The OPE is still valid for the total decay width

- At the differential level, this is generally not true

  $\Rightarrow$ Large contributions factorise wrt to tree-level

  $\Rightarrow$ Useful to go beyond NLO

Two calculation approaches

1. Splitting Functions

$$\left(\frac{d\Gamma}{dy}\right)^{(1)} = \frac{\alpha}{2\pi}\,\underbrace{\bar{L}_{b/e}}_{\log(m_b^2/m_e^2)} \int_y^{1-\rho} \frac{dx}{x}\,\underbrace{P_{ee}^{(0)}}_{\text{plus distribution}}\left(\frac{y}{x}\right)\left(\frac{d\Gamma}{dx}\right)^{(0)}$$

- Correction vanishes for the inclusive branching fraction
- Suitable for evaluating $\mathcal{O}(\alpha^2)$ and $\mathcal{O}(\alpha/m_b^n)$ corrections

2. Full $\mathcal{O}(\alpha)$ corrections

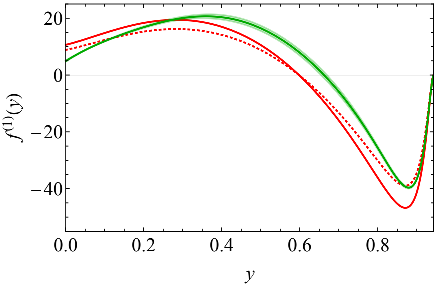
- Access all corrections, not only the one that factorise
- Real corrections are computationally expensive
  - $\Rightarrow$ Cuba library employed to carry out the 4-body integration
  - $\Rightarrow$ Phase space splitting used to reduce the size of the integrands

Lepton Energy spectrum

[Bigi, MB, Gambino, Haisch, Piccione, '23]

- We compute bins in the lepton energy using the full $\mathcal{O}(\alpha)$ calculation
- We compare them to the results given by the splitting functions
- The difference the two calculations for the lepton energy spectrum and obtain a full analytic formula for the radiative corrections

  $\Rightarrow$ Relatively small, easy-to-use formula to obtain branching fractions, lepton energy moments w/o cuts



$$f^{(1)}(y) = \frac{\bar{L}_{b/e}}{2} f^{(1)}_{\mathrm{LL}}(y) + \Delta f^{(1)}(y)$$

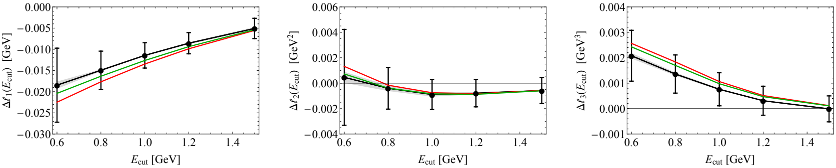
Comparison with data

[Bigi, MB, Gambino, Haisch, Piccione, '23]

- Babar provides data with and without applying PHOTOS to subtract QED effects

  $\Rightarrow$ Perfect ground to test our calculations

  $\Rightarrow$ Not the same for Belle at the moment, could be possible for future analysis?



- The moments, since they are normalised, are not affected by the large threshold corrections
- The agreement with BaBar is very good

$$\langle E_\ell^n \rangle = \frac{\int_{E_\ell > E_{\ell,\rm cut}} dE_\ell E_\ell^n \frac{d\Gamma}{dE_\ell}}{\Gamma_{E_\ell > E_{\ell,\rm cut}}}$$

QED for exclusive decays

- For $B^0 \to D^+ \ell \bar{\nu}$, the threshold effects were calculated and are $1 + \alpha\pi$

  [Ginsberg, '66, De Boer, Kitahara, Nisandzic, '18]

- For $B^0 \to D^{*+} \ell \bar{\nu}$, the threshold effects might have a different structure because the hadronic matrix element is different

  $\Rightarrow$ To verify explicitly

- Structure-dependent terms are unknown, but maybe something is doable in the HQE?

- How do we reconcile the threshold effects between the exclusive and the inclusive?

$$\mathcal{B}(B \to X_c \ell \nu) = \mathcal{B}(B \to D \ell \nu) + \mathcal{B}(B \to D^* \ell \nu) + \mathcal{B}(B \to D^{**} \ell \nu) + \ldots$$