

QMAP

EFT Measurements at the LHC

Markus Luty

UC Davis/QMAP

S. Chang, M. Chen, D. Liu, ML, arXiv:2212.06215
+ work in progress with S. Chang, T. Ma, A. Wulzer

Big Picture

- The Standard Model is the most general UV complete theory with the observed elementary particle content
- All parameters of the SM have been determined at the percent level or better
- Any observed deviation from the SM
 - ⇒ $\begin{cases} \text{new light particles} & \leftarrow \text{direct searches} \\ \text{new physics at high scales} & \leftarrow \text{precision measurement} \end{cases}$

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This talk:

- What are the observables that we want to measure?
- How should we measure/report them?

Effective Field Theory

- Effects of high scale new physics can be parameterized by adding new effective interactions to the SM

Effective Field Theory

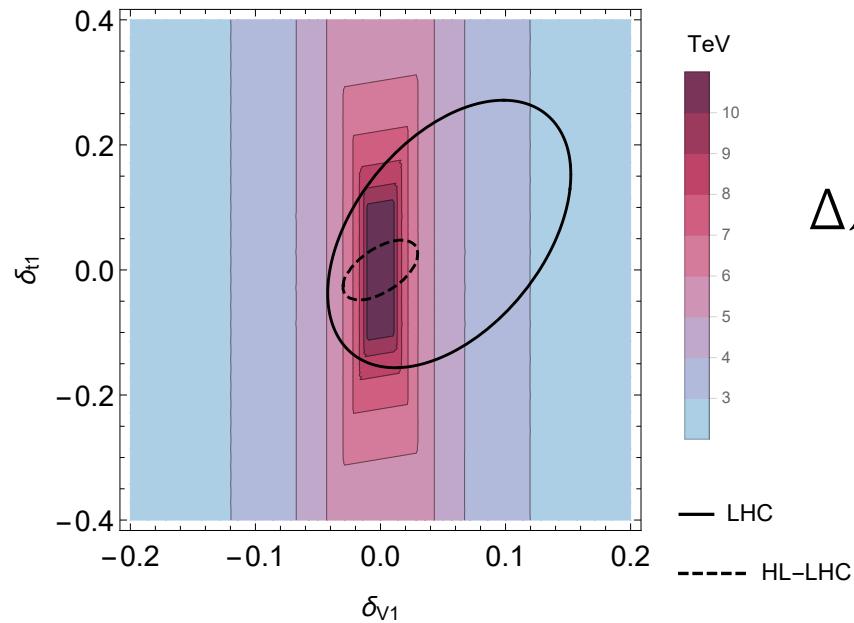
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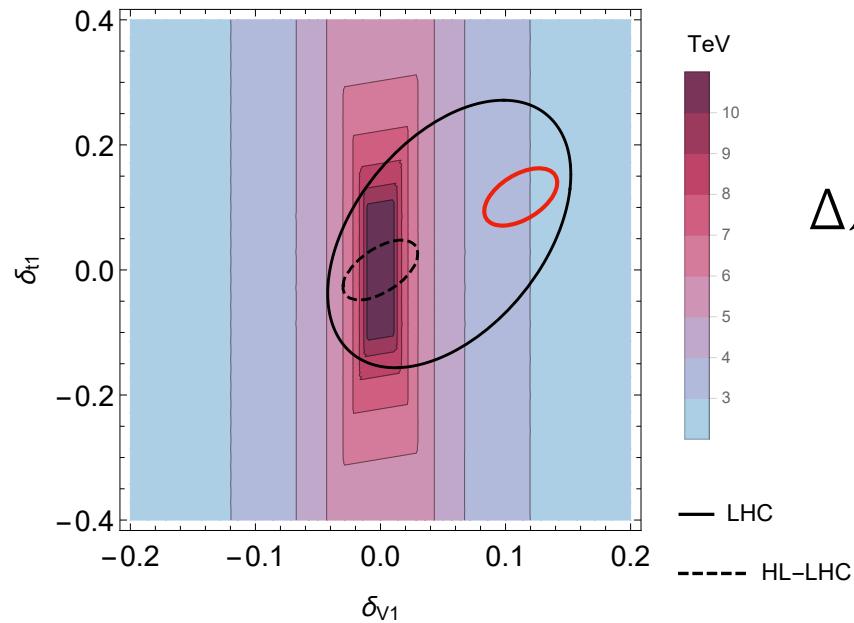
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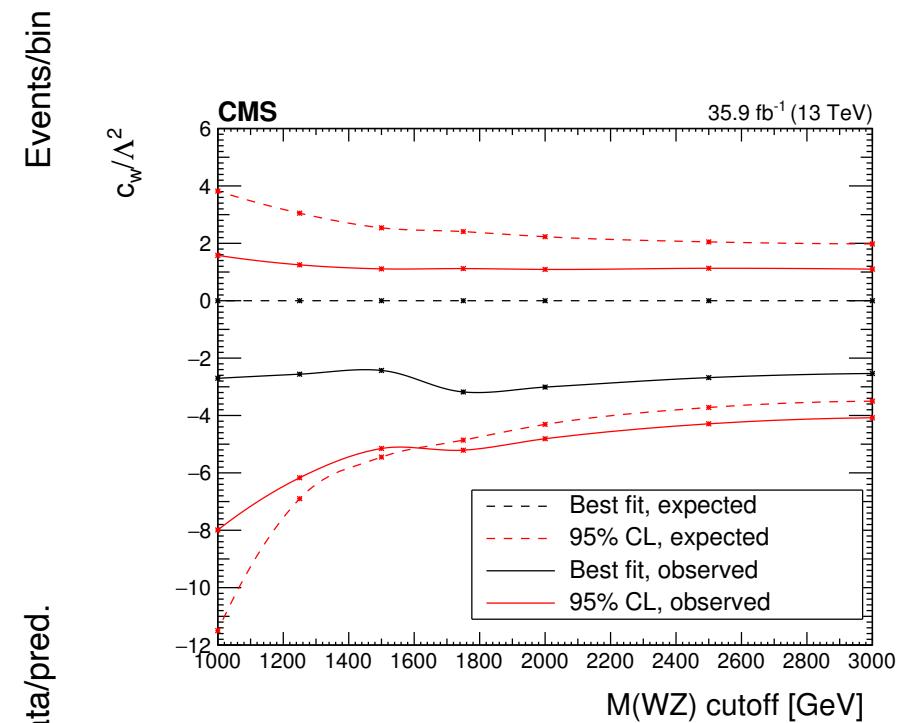
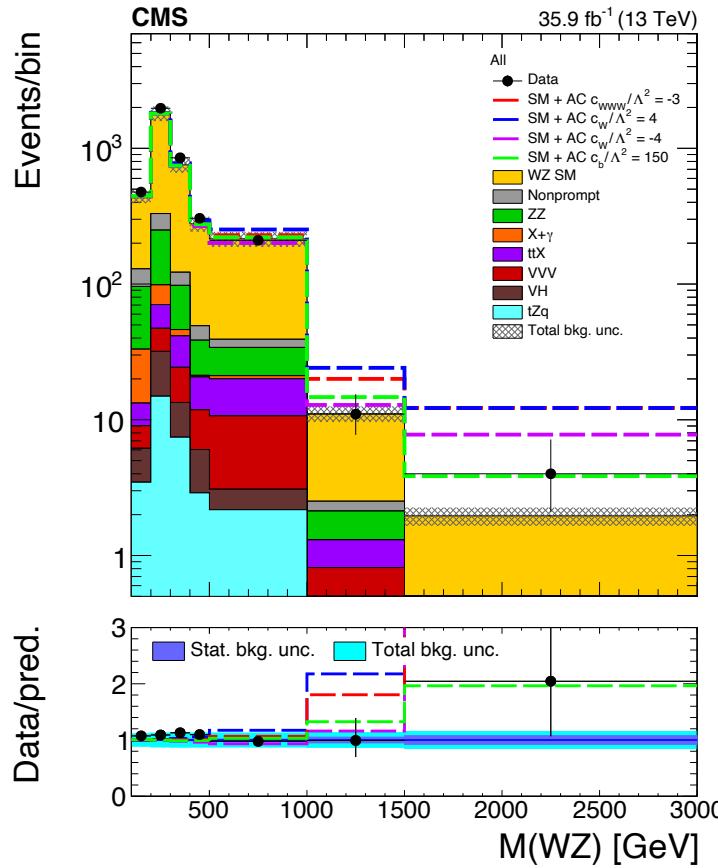
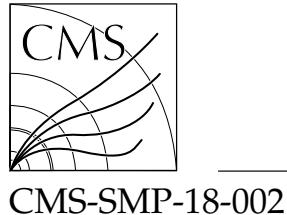
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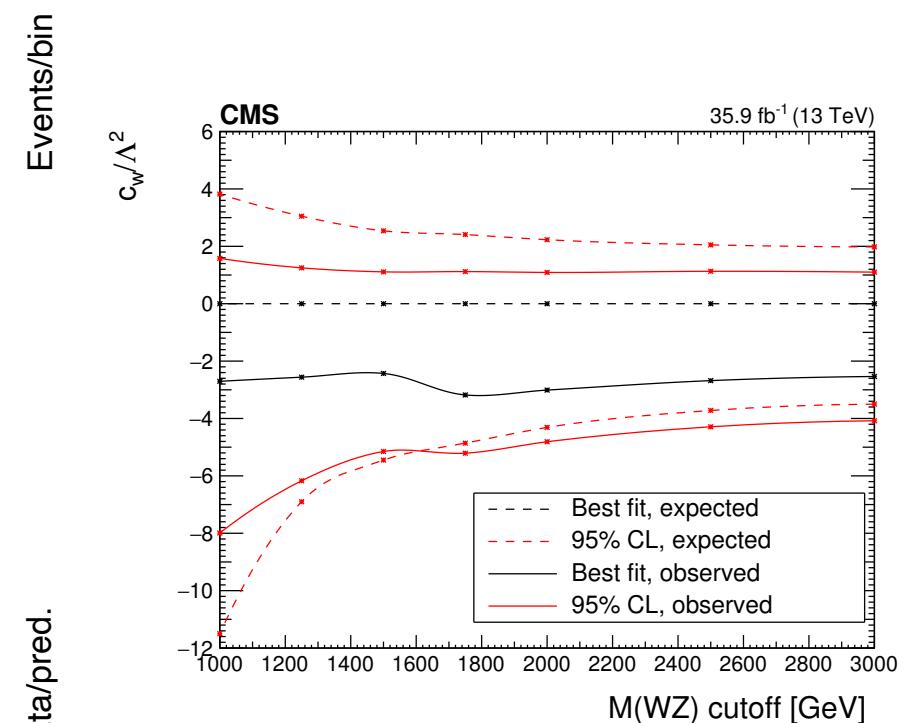
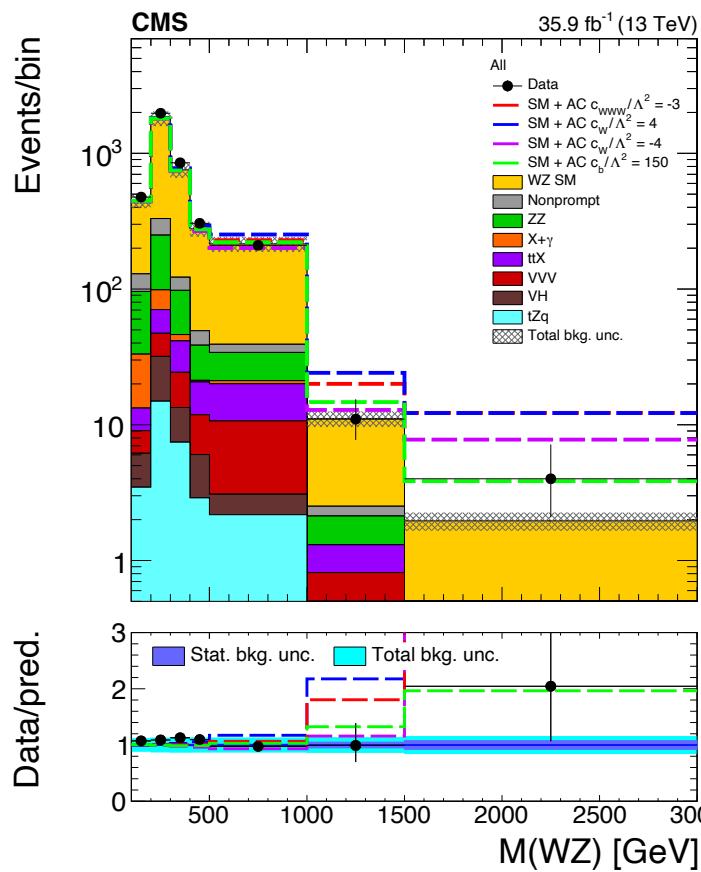
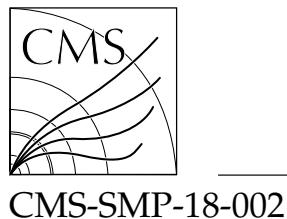
High Energy Signals

Exploit hardness of BSM contribution...
...but beware of breakdown of EFT



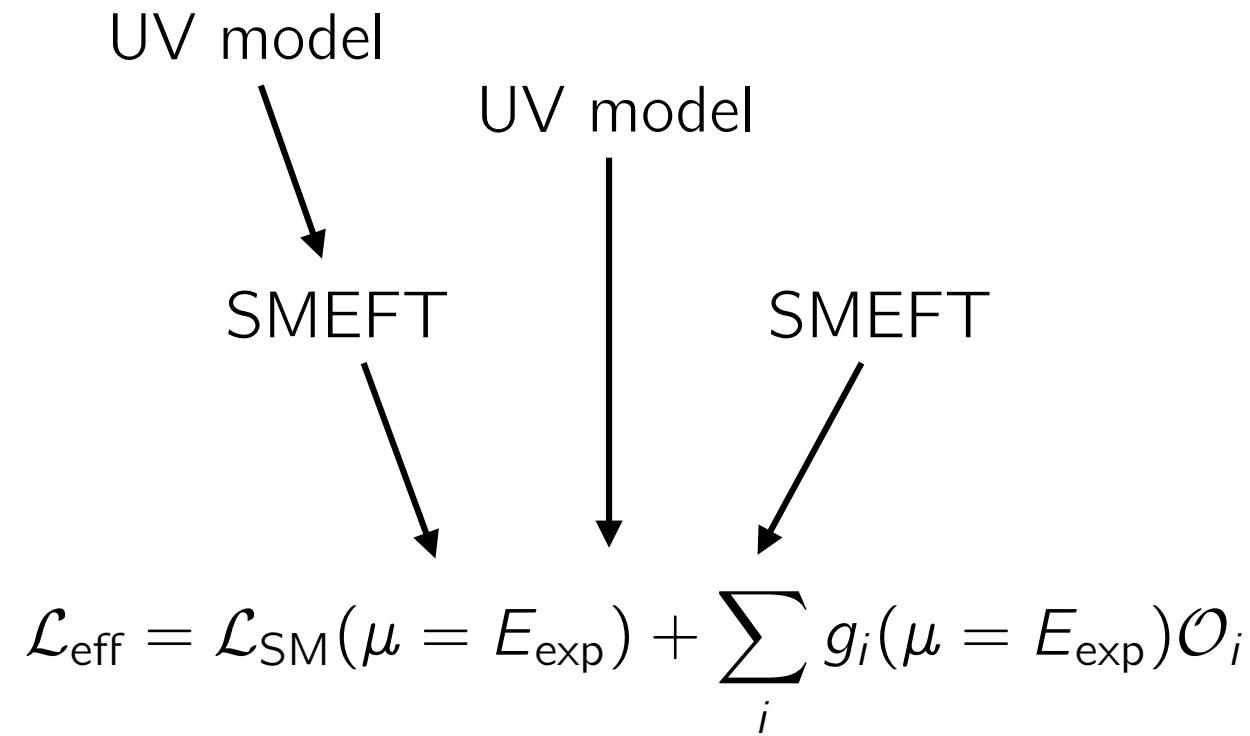
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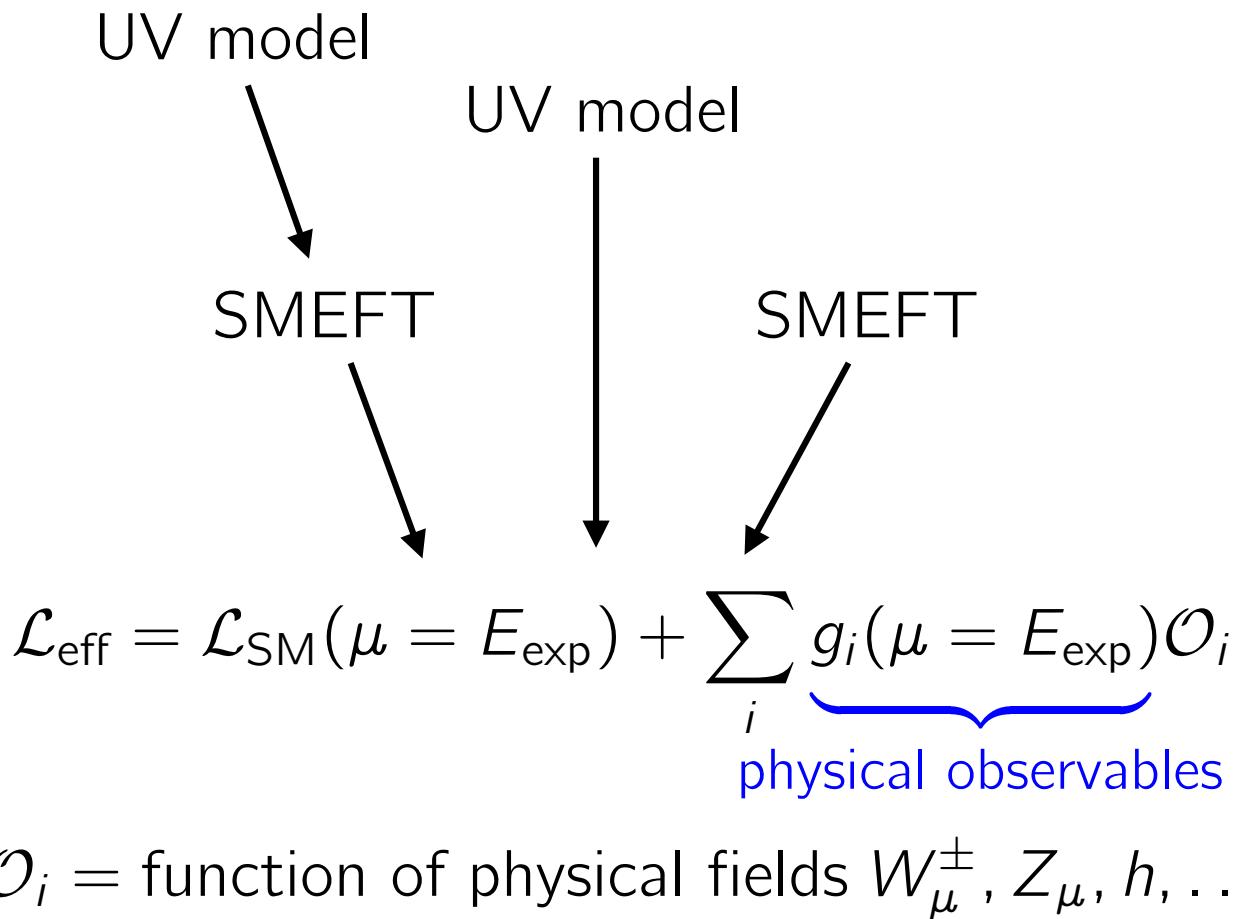
Can we be more systematic?

On-Shell Approach

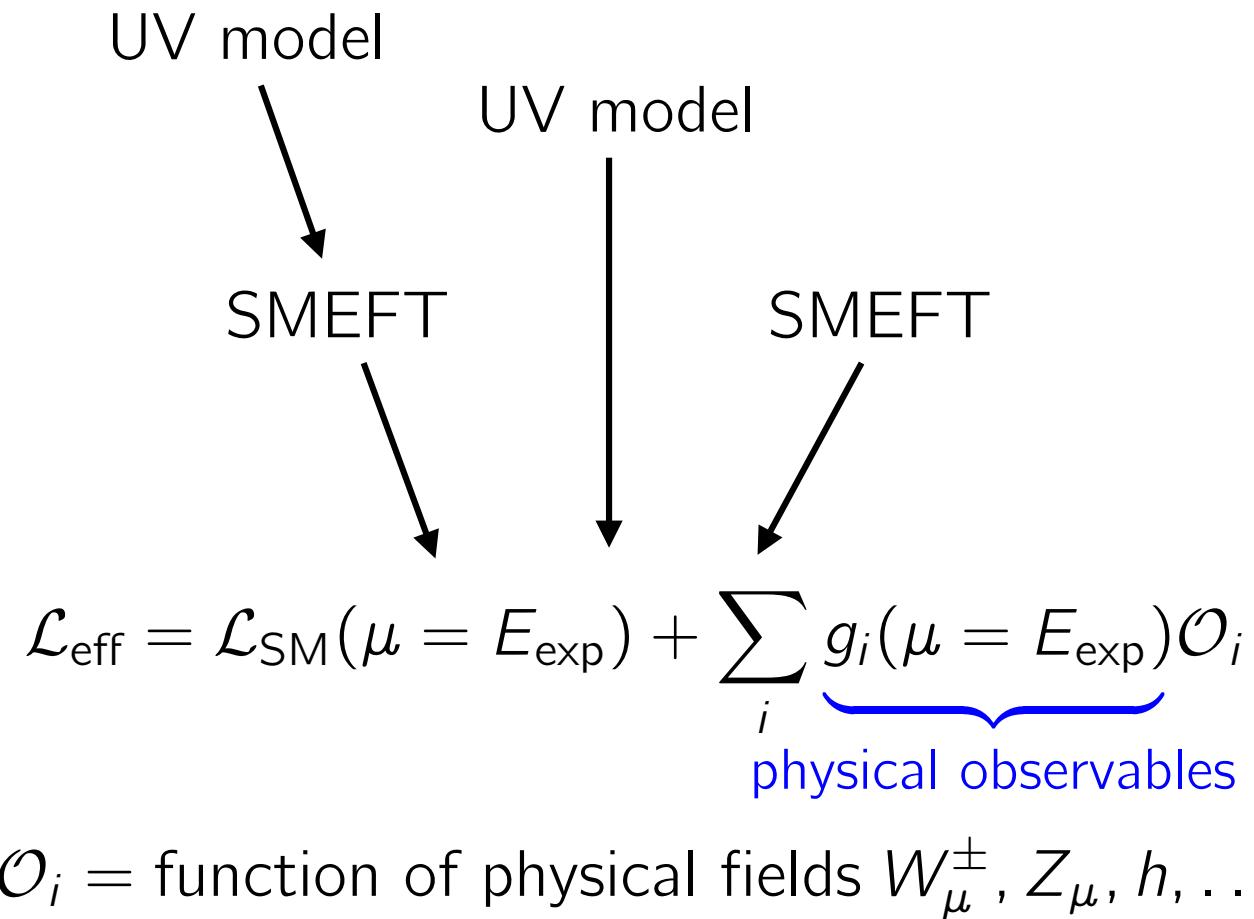


\mathcal{O}_i = function of physical fields $W_\mu^\pm, Z_\mu, h, \dots$

On-Shell Approach

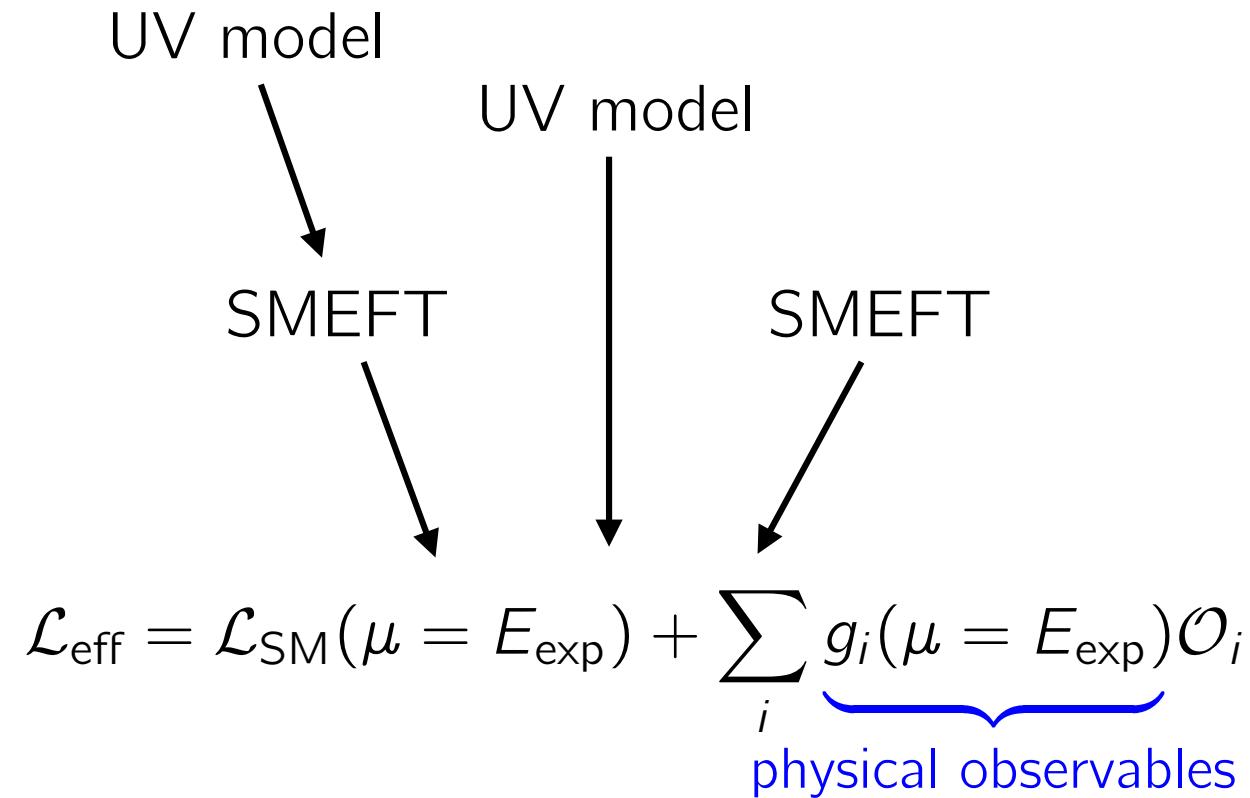


On-Shell Approach



Bottom-up approach: experimentally constrain $g_i(\mu = E_{\text{exp}})$

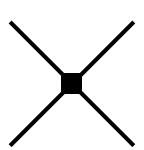
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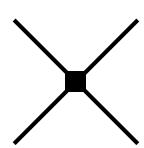
Classify: $\mathcal{O}_i \leftrightarrow$ Feynman rules for on-shell particles

Primary Observables



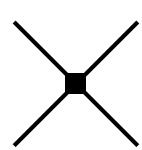
$$\sum_i g_i \mathcal{O}_i^{(4)} = \sum_{\alpha} g_{\alpha} \mathcal{O}_{\alpha}^{(4)} \left[1 + \frac{c_1}{M^2} s + \frac{c_2}{M^2} t + \frac{c_3}{M^4} st + \dots \right]$$

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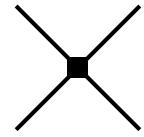
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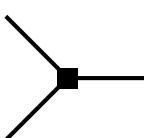


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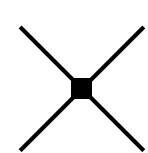


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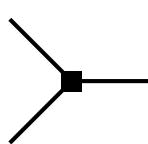


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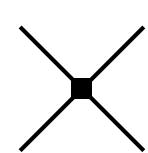
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- There are finitely many primary operators

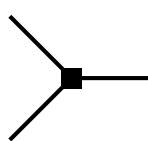
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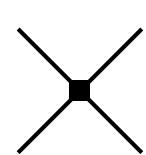
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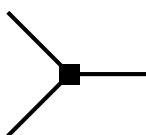
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- Generically descendants are subleading if EFT is valid

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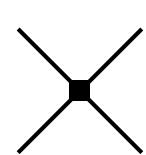
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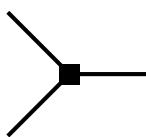
Exceptions:

- Primary operators may be suppressed in certain UV models
e.g. PNGB models $\Rightarrow h$ shift symmetry
- Accidental cancelations

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In such cases, define “physical primary” to be leading operator

Primary Observables

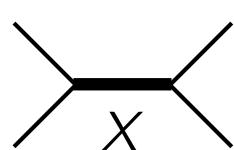
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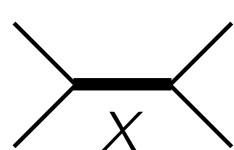
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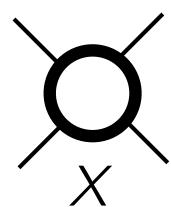
A Feynman diagram showing a central horizontal black line labeled 'X' with two diagonal lines extending from its ends, representing a resonance particle.
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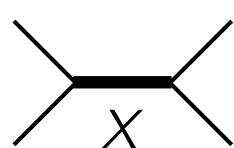

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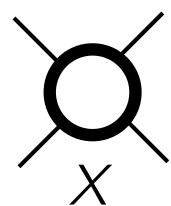

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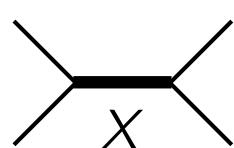

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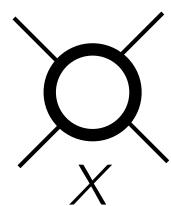
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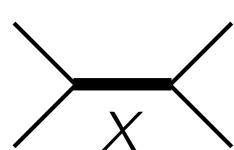

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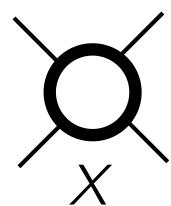
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$$M \lesssim E_{\max}(g_{\alpha})$$

Primary Observables

i	$\mathcal{O}_i^{hZ\bar{f}f}$	CP	$d_{\mathcal{O}_i}$	SMEFT Operator	c Unitarity Bound
1	$hZ^\mu \bar{\psi}_L \gamma_\mu \psi_L$	+	5	$i(H^\dagger \overset{\leftrightarrow}{D}_\mu H) \bar{Q}_L \gamma^\mu Q_L$	$\frac{0.6}{E_{\text{TeV}}^2}, \frac{5}{E_{\text{TeV}}^4}$
2	$hZ^\mu \bar{\psi}_R \gamma_\mu \psi_R$	+	5	$i(H^\dagger \overset{\leftrightarrow}{D}_\mu H) \bar{u}_R \gamma^\mu u_R$	$\frac{0.6}{E_{\text{TeV}}^2}, \frac{5}{E_{\text{TeV}}^4}$
3	$hZ^{\mu\nu} \bar{\psi}_L \sigma_{\mu\nu} \psi_R + \text{h.c.}$	+	6	$\bar{Q}_L \sigma^{\mu\nu} u_R \sigma^a \tilde{H} W_{\mu\nu}^a + \text{h.c.}$	$\frac{2}{E_{\text{TeV}}^2}, \frac{10}{E_{\text{TeV}}^4}$
4	$ih\tilde{Z}_{\mu\nu} \bar{\psi}_L \sigma^{\mu\nu} \psi_R + \text{h.c.}$	-	6	$i\bar{Q}_L \sigma^{\mu\nu} u_R \sigma^a \tilde{H} \tilde{W}_{\mu\nu}^a + \text{h.c.}$	$\frac{2}{E_{\text{TeV}}^2}, \frac{10}{E_{\text{TeV}}^4}$
5	$ihZ^\mu (\bar{\psi}_L \overset{\leftrightarrow}{\partial}_\mu \psi_R) + \text{h.c.}$	+	6	$(H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \overset{\leftrightarrow}{D}^\mu u_R) \tilde{H} + \text{h.c.}$	
6	$hZ^\mu \partial_\mu (\bar{\psi}_L \psi_R) + \text{h.c.}$	-	6	$i(H^\dagger \overset{\leftrightarrow}{D}_\mu H) D^\mu (\bar{Q}_L u_R) \tilde{H} + \text{h.c.}$	$\frac{0.1}{E_{\text{TeV}}^3}, \frac{4}{E_{\text{TeV}}^6}$
7	$ihZ^\mu \partial_\mu (\bar{\psi}_L \psi_R) + \text{h.c.}$	+	6	$(H^\dagger \overset{\leftrightarrow}{D}_\mu H) D^\mu (\bar{Q}_L u_R) \tilde{H} + \text{h.c.}$	
8	$hZ^\mu (\bar{\psi}_L \overset{\leftrightarrow}{\partial}_\mu \psi_R) + \text{h.c.}$	-	6	$i(H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \overset{\leftrightarrow}{D}^\mu u_R) \tilde{H} + \text{h.c.}$	
9	$ih\tilde{Z}_{\mu\nu} (\bar{\psi}_L \gamma^\mu \overset{\leftrightarrow}{\partial}^\nu \psi_L)$	+	7	$i H ^2 \tilde{W}^{a\mu\nu} (\bar{Q}_L \gamma_\mu \sigma^a \overset{\leftrightarrow}{D}_\nu Q_L)$	
10	$h\tilde{Z}_{\mu\nu} \partial^\mu (\bar{\psi}_L \gamma^\nu \psi_L)$	-	7	$ H ^2 \tilde{W}^{a\mu\nu} D_\mu (\bar{Q}_L \gamma_\nu \sigma^a Q_L)$	$\frac{0.4}{E_{\text{TeV}}^3}, \frac{1}{E_{\text{TeV}}^4}$
11	$ih\tilde{Z}_{\mu\nu} (\bar{\psi}_R \gamma^\mu \overset{\leftrightarrow}{\partial}^\nu \psi_R)$	+	7	$i H ^2 \tilde{B}^{\mu\nu} (\bar{u}_R \gamma_\mu \overset{\leftrightarrow}{D}_\nu u_R)$	
12	$h\tilde{Z}_{\mu\nu} \partial^\mu (\bar{\psi}_R \gamma^\nu \psi_R)$	-	7	$ H ^2 \tilde{B}^{\mu\nu} D_\mu (\bar{u}_R \gamma_\nu u_R)$	

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$$f_{\alpha}(x, y) \xrightarrow{|x|, |y| \ll 1} A_{\alpha}x + B_{\alpha}y + O(x^2, y^2, xy) \quad A_{\alpha}, B_{\alpha} \sim 1$$

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- Statistical error model

$$f_{\alpha}(x, y) = \sum_{i=1}^N c_{\alpha i} e_i(x, y)$$

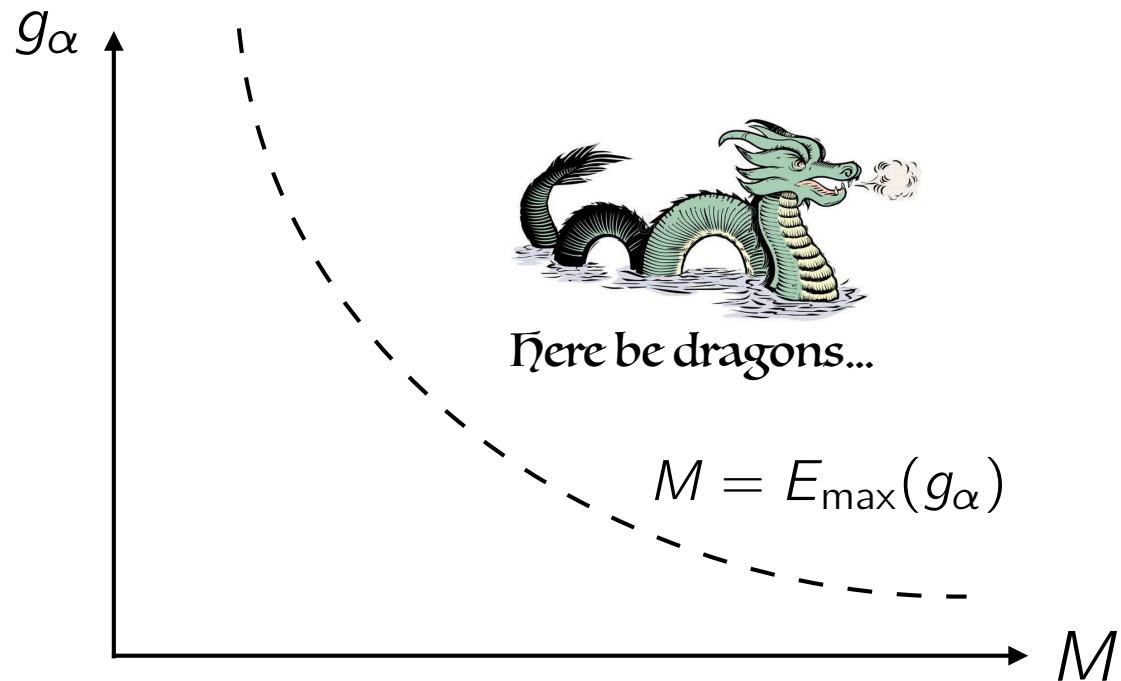
$c_{\alpha i}$ = nuisance parameters $\langle c_{\alpha i} \rangle = 0, \delta c_{\alpha i} = 1$

$e_i(x, y)$ = basis of shape functions

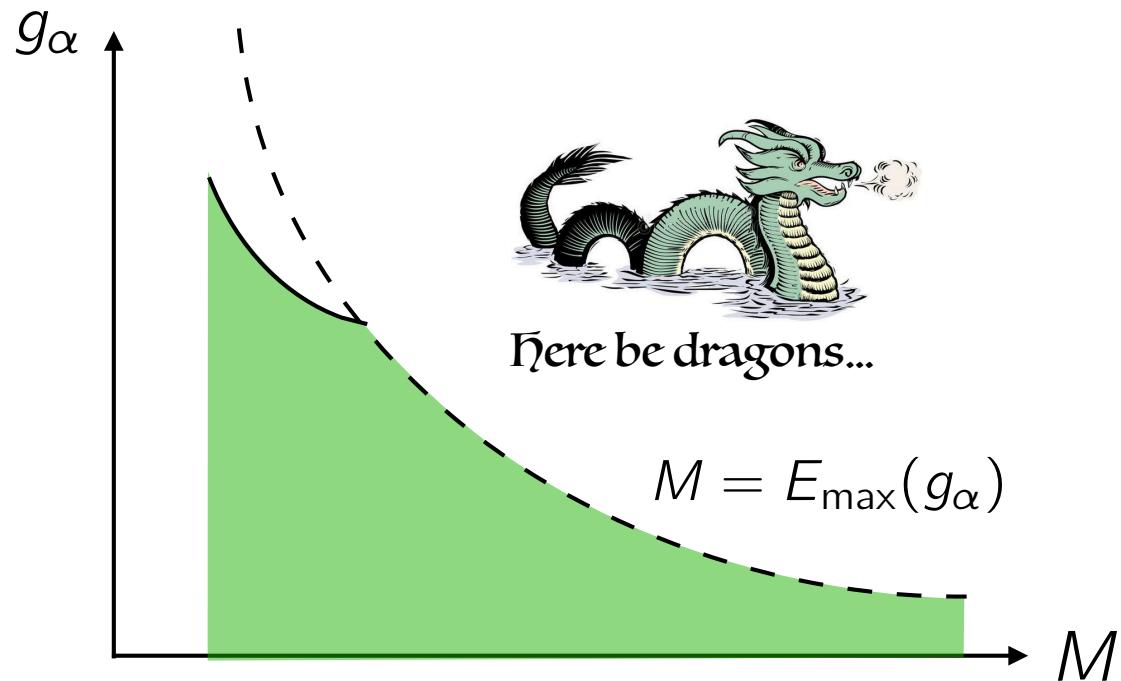
Constraining Primary Observables



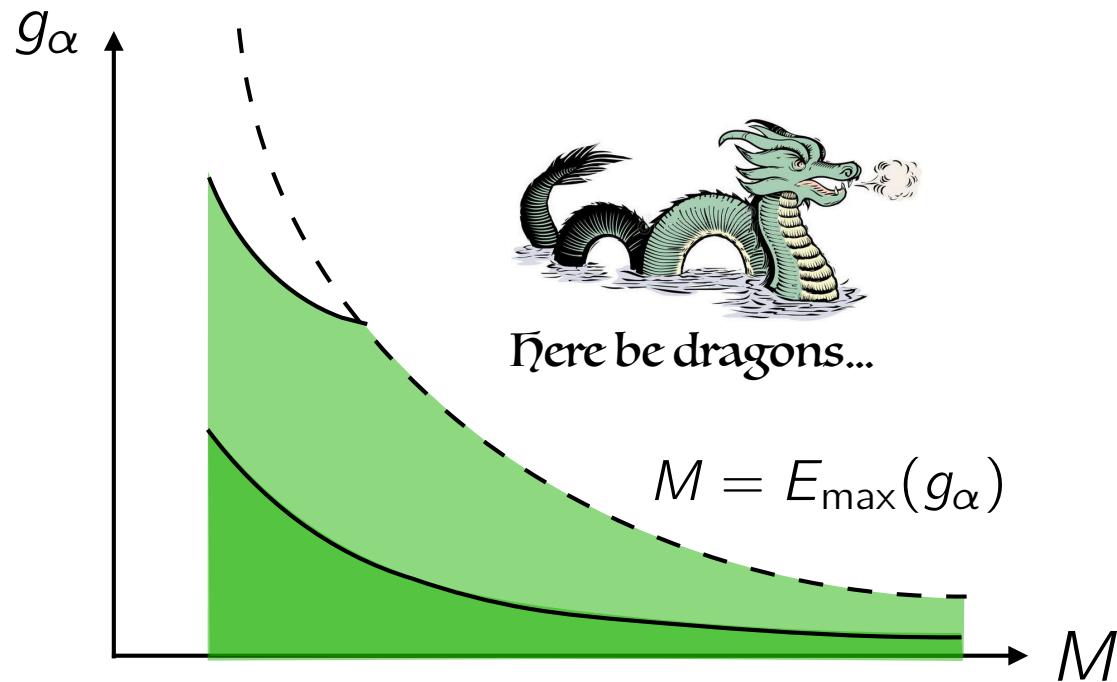
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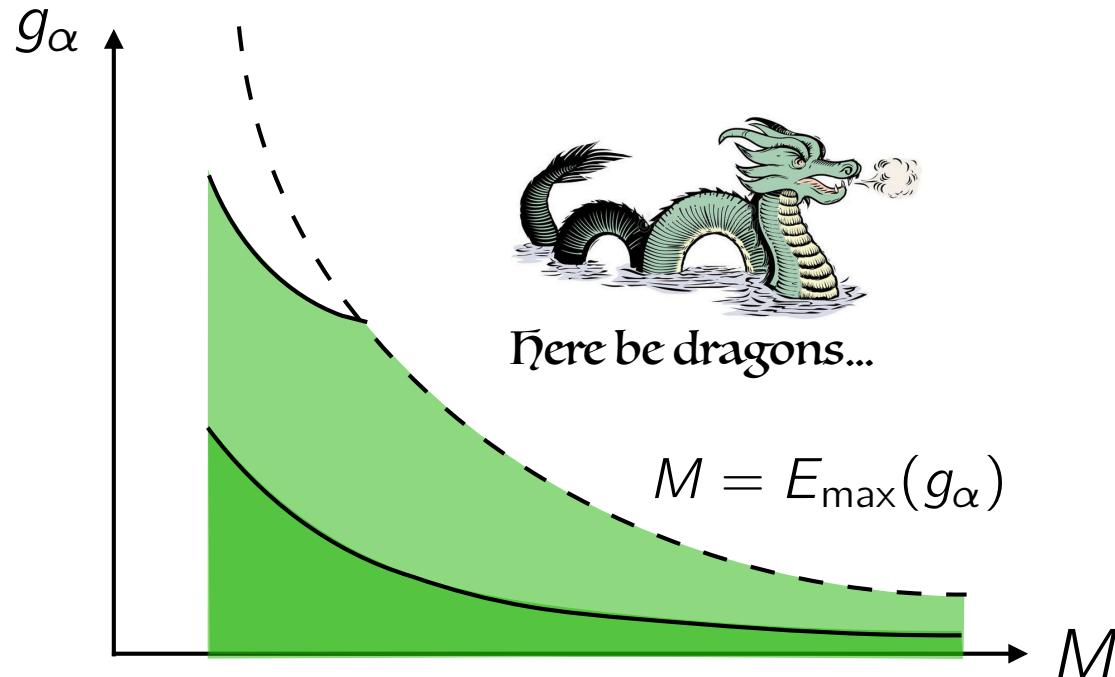
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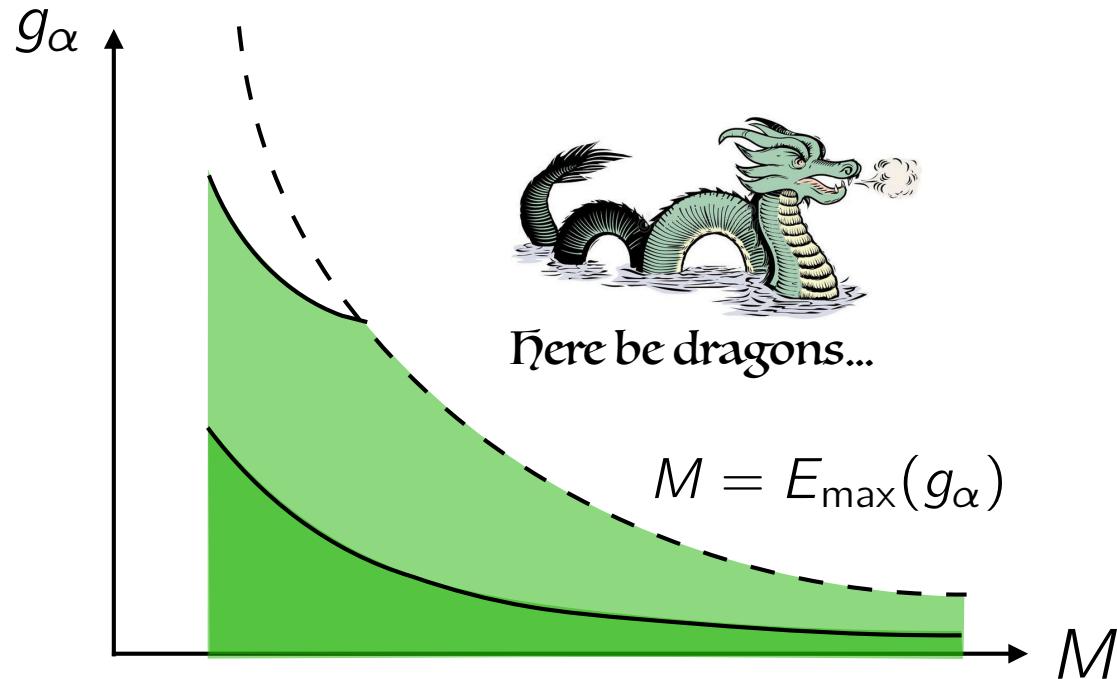


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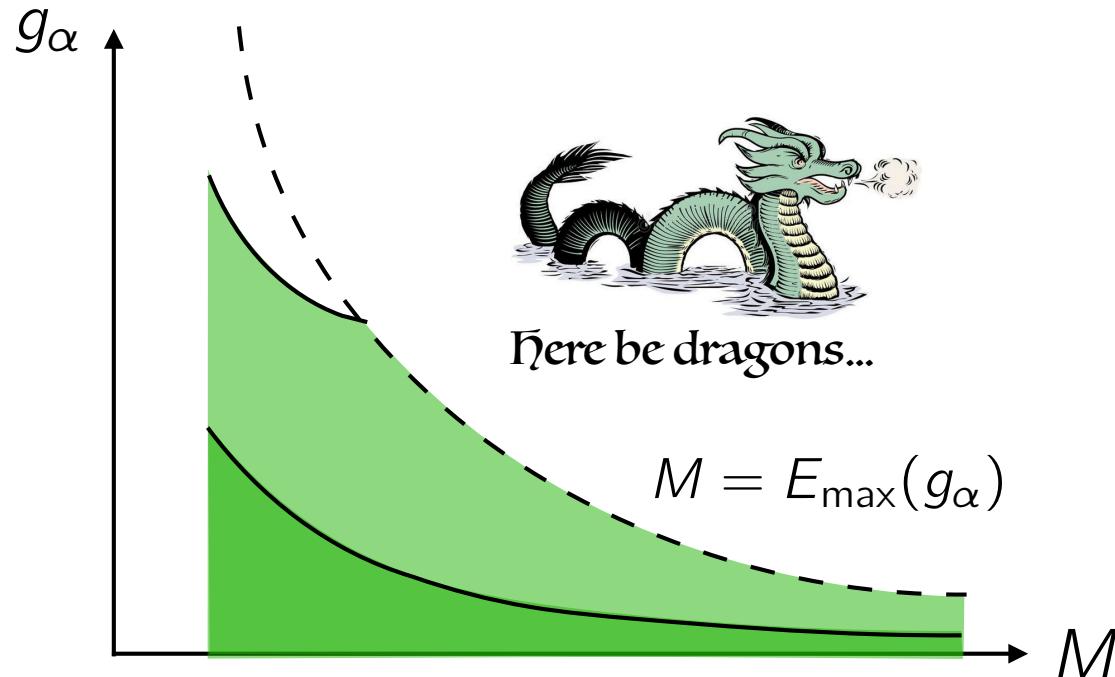
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Constraining Primary Observables



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- Can compare/combine different channels
- Constraints can be directly compared to UV models

A Simple Analysis

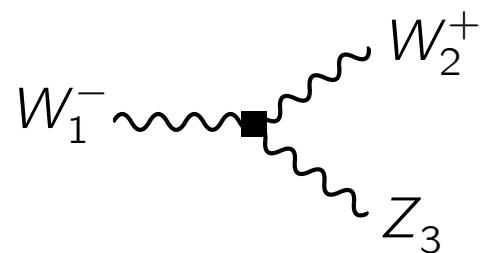
Work in progress...

$$\Delta \mathcal{L} = i\delta g W_\mu^- \overset{\leftrightarrow}{\partial}_\nu W_\mu^+ Z_\nu$$

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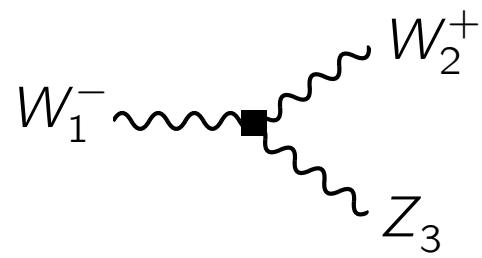
Feynman diagram showing a vertex correction for a Z_3 boson line. A horizontal wavy line labeled W_1^- enters from the left and meets a vertical wavy line labeled W_2^+ at a vertex. From this vertex, another vertical wavy line labeled Z_3 continues downwards. The vertex is marked with a small black square.

$$= \delta g (\epsilon_1 \cdot \epsilon_2) (p_1 - p_2) \cdot \epsilon_3$$

A Simple Analysis

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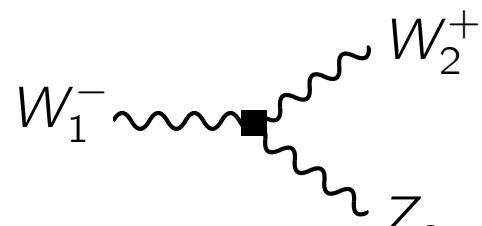


$$\begin{aligned} &= \delta g (\epsilon_1 \cdot \epsilon_2) (p_1 - p_2) \cdot \epsilon_3 \\ &\quad \times \left[1 + f \left(\frac{p_1^2}{M^2}, \frac{p_2^2}{M^2}, \frac{p_3^2}{M^2} \right) \right] \end{aligned}$$

A Simple Analysis

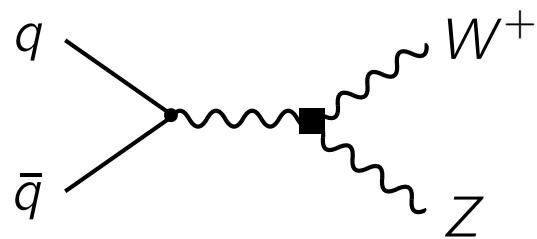
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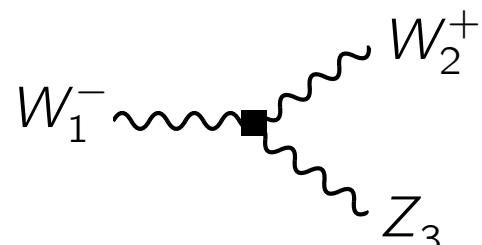
$$\sim \sqrt{s}$$

$$s \gg m_W^2$$

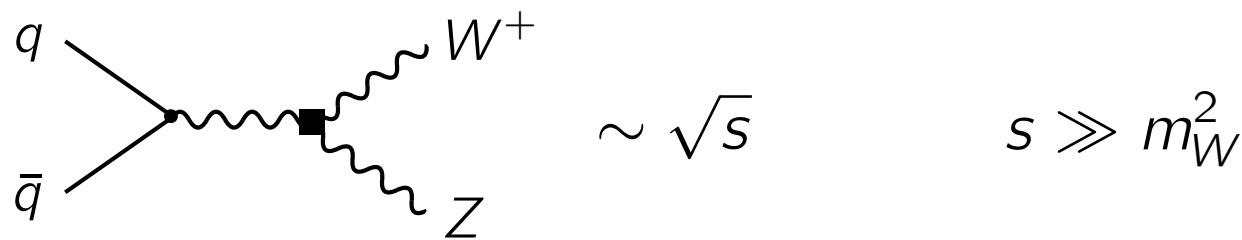
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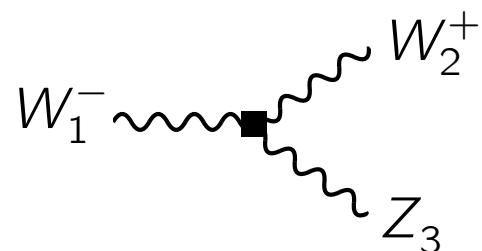


$$\mathcal{M} = \mathcal{M}_{\text{SM}} + \delta g \left[1 + f \left(\frac{s}{M^2}, \frac{m_W^2}{M^2}, \frac{m_Z^2}{M^2} \right) \right] \delta \mathcal{M}$$

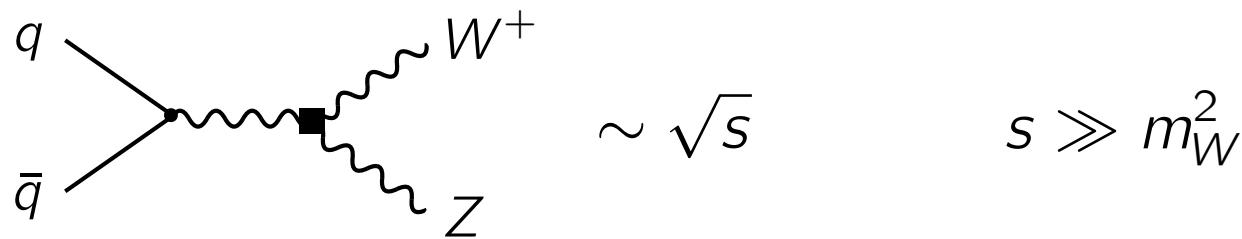
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Monte Carlo data can be generated by reweighting

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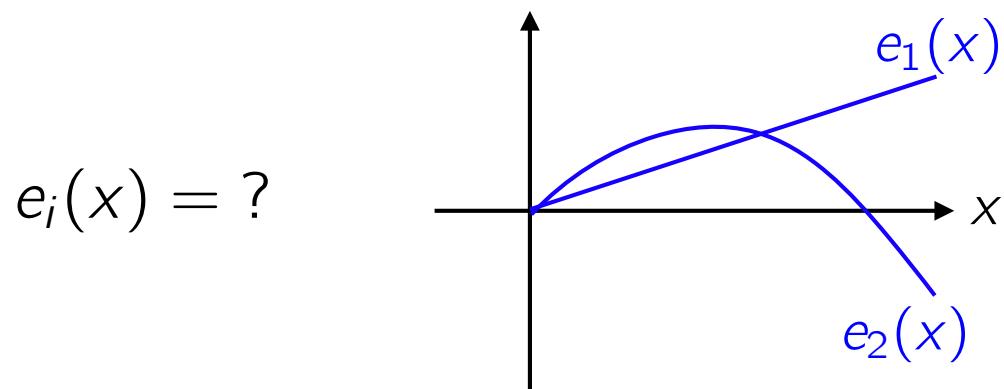
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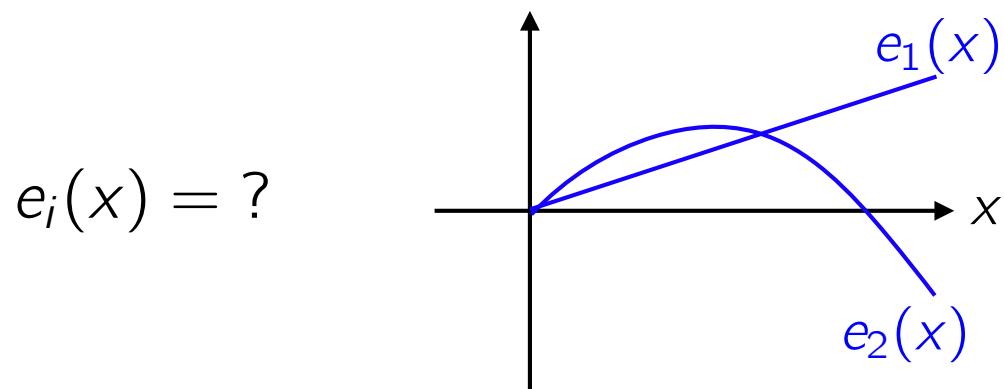
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c_i = nuisance parameters

$$L(\delta g, c) = L_{\text{data}}(\delta g) L_{\text{prior}}(c)$$

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- Contribution to LHC legacy... 