

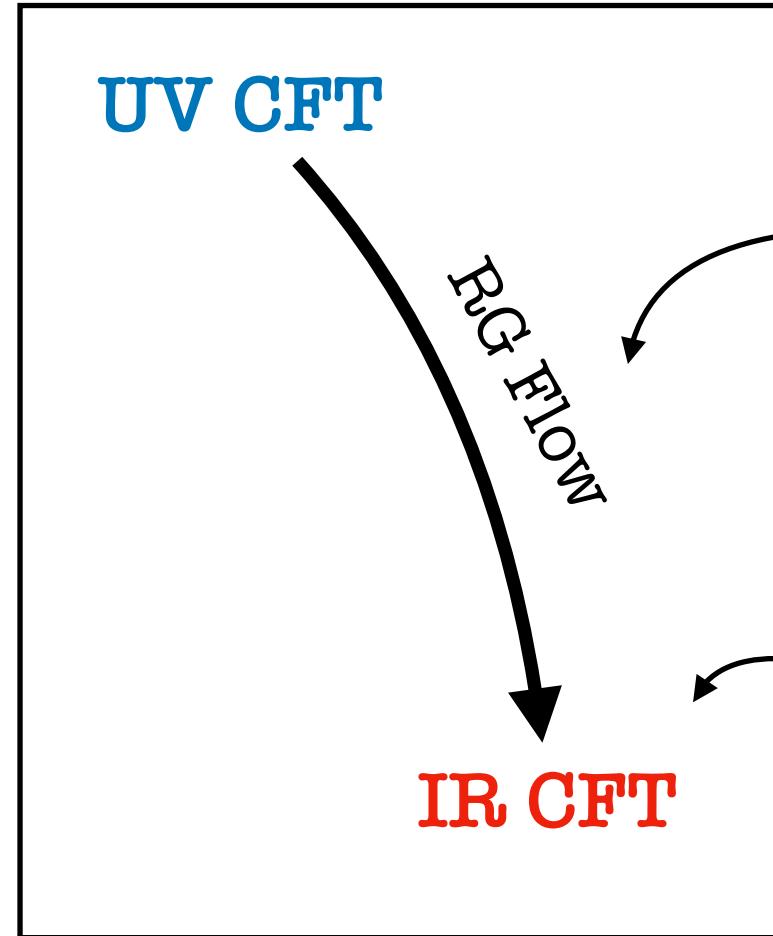
A Bootstrap Bridge between Gravity and QCD

Andrea Guerrieri

January 18, 2024



Bootstrap: What is (im)possible in the Space of QFTs



Example of a QFT

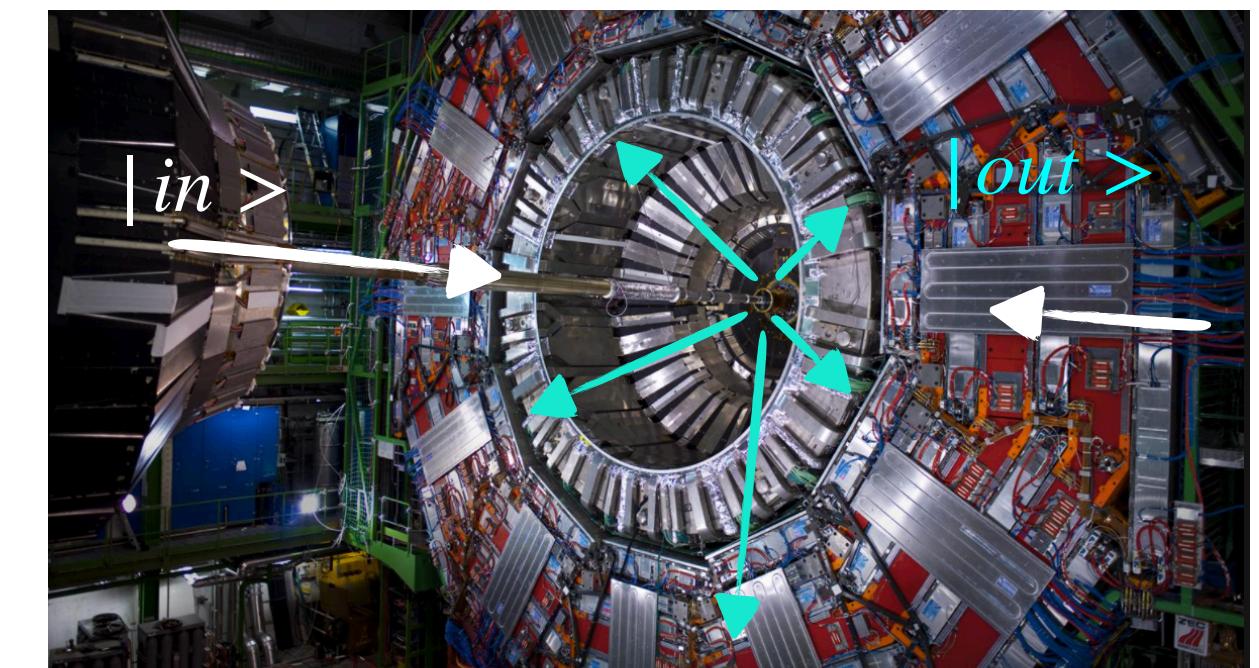
S-matrix Bootstrap:
study the possible RG-flows

Conformal Bootstrap:
study fixed points of the RG-flow

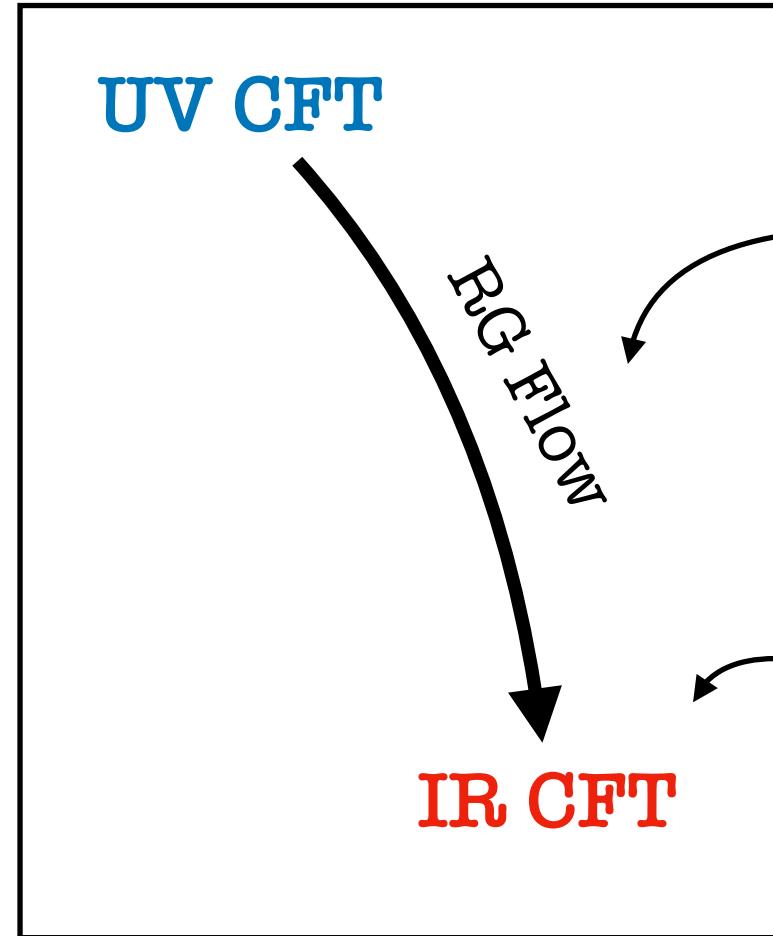
Bootstrap manifesto:
find the space of all possible physical observables assuming general principles

The S-matrix measures the probability of a scattering process

$$\mathcal{S}_{in \rightarrow out} \equiv \langle in | out \rangle$$



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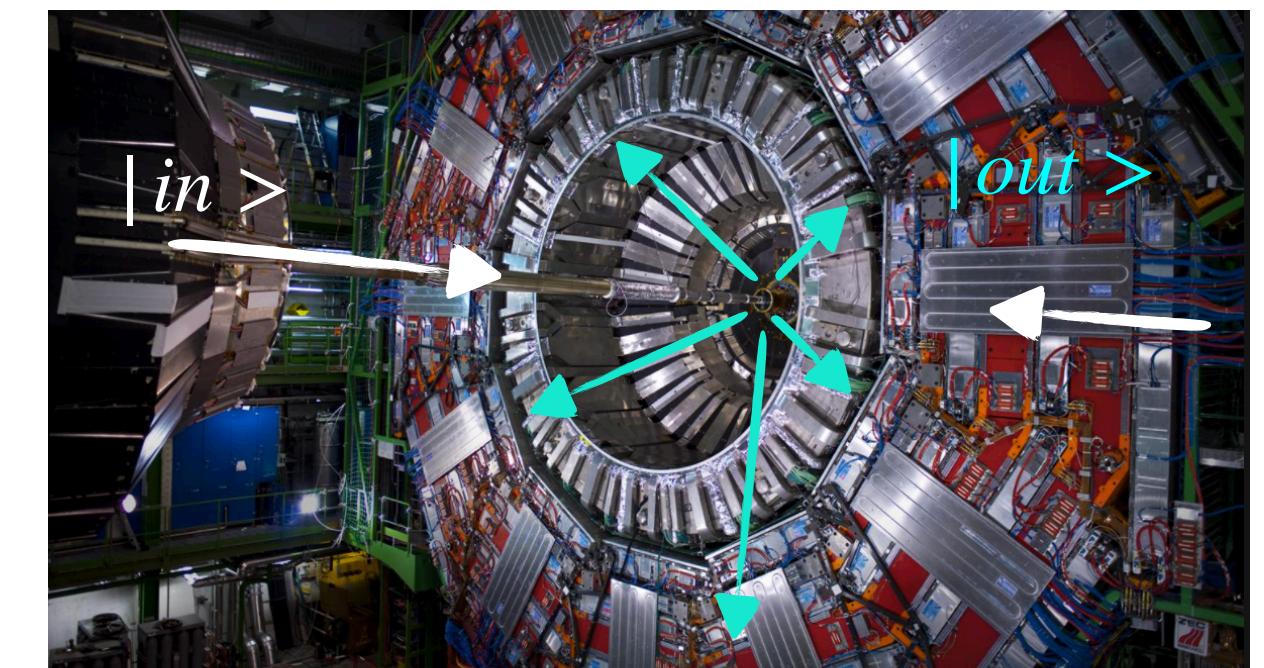
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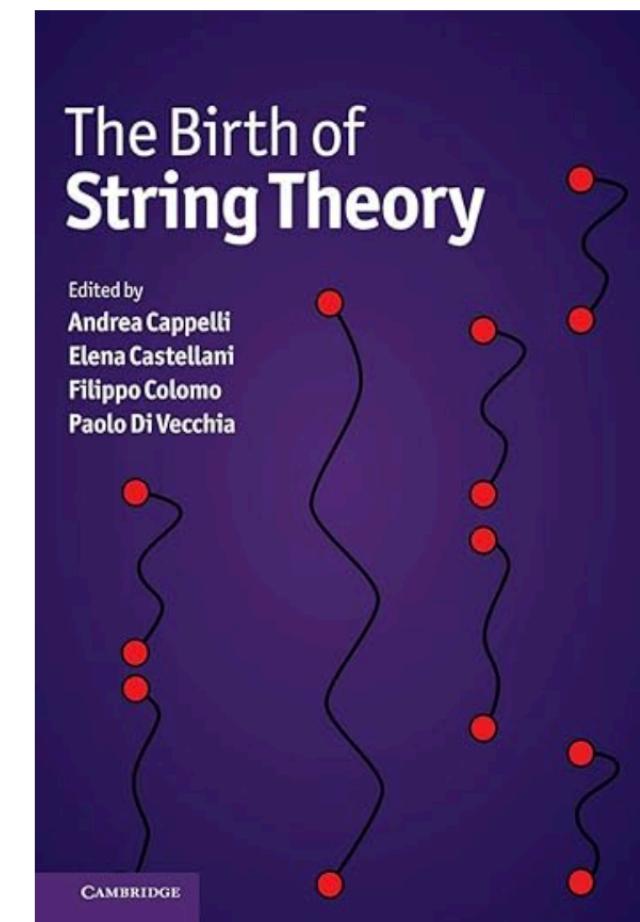
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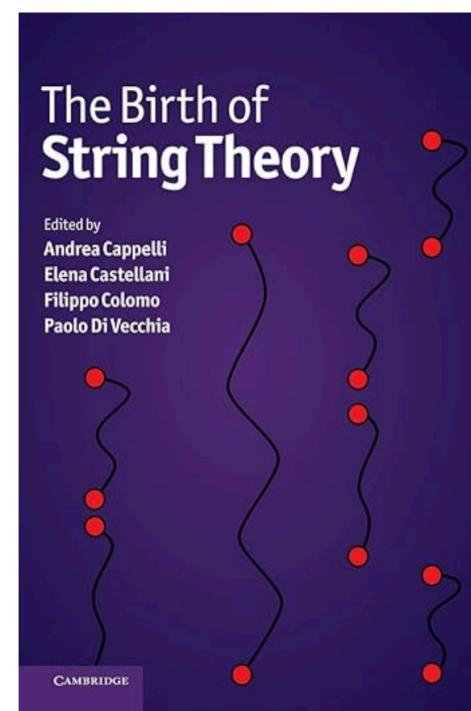
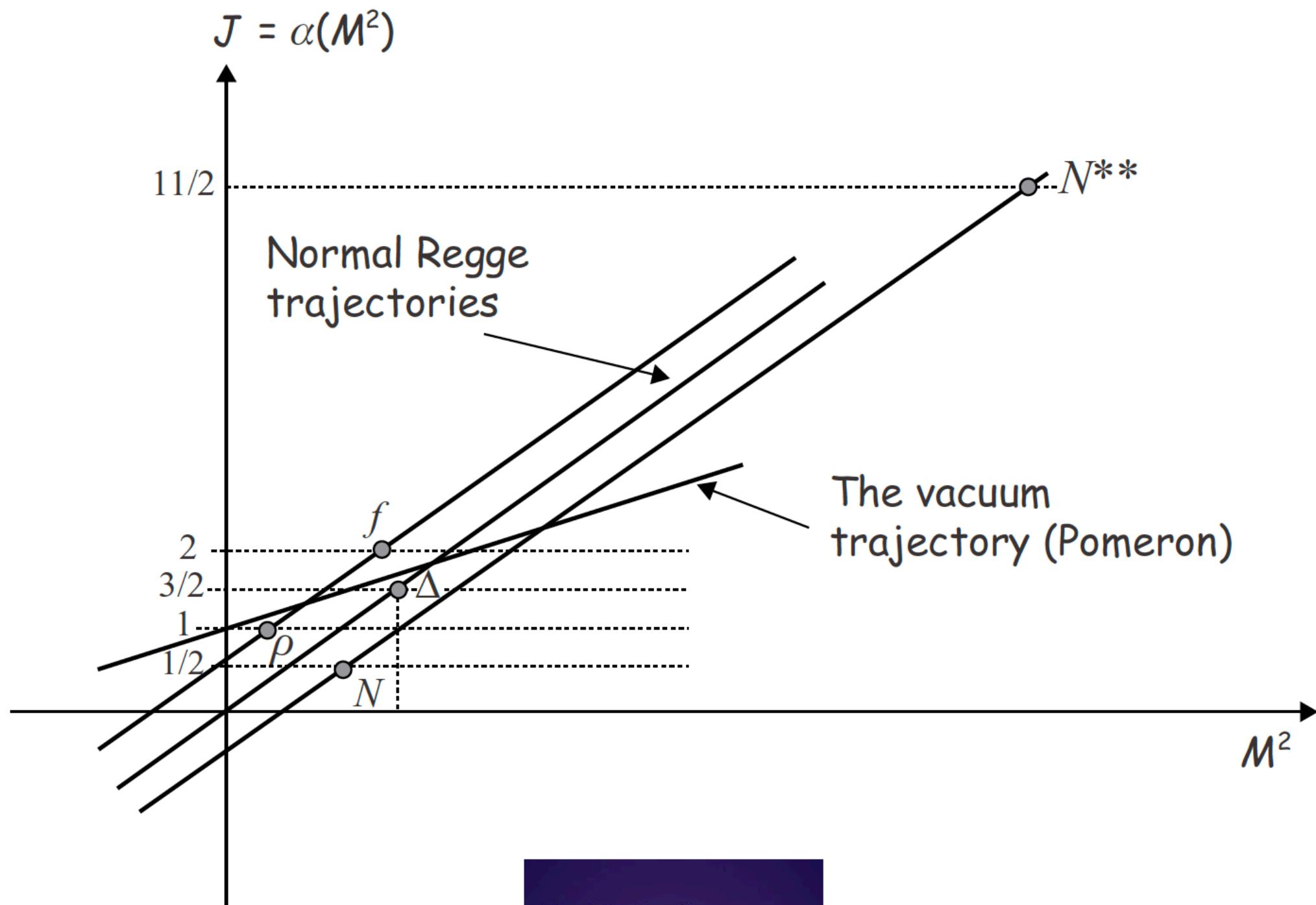


Point of the Talk: Strings populate the landscape of non-perturbative S-matrices



Strings and Regge trajectories

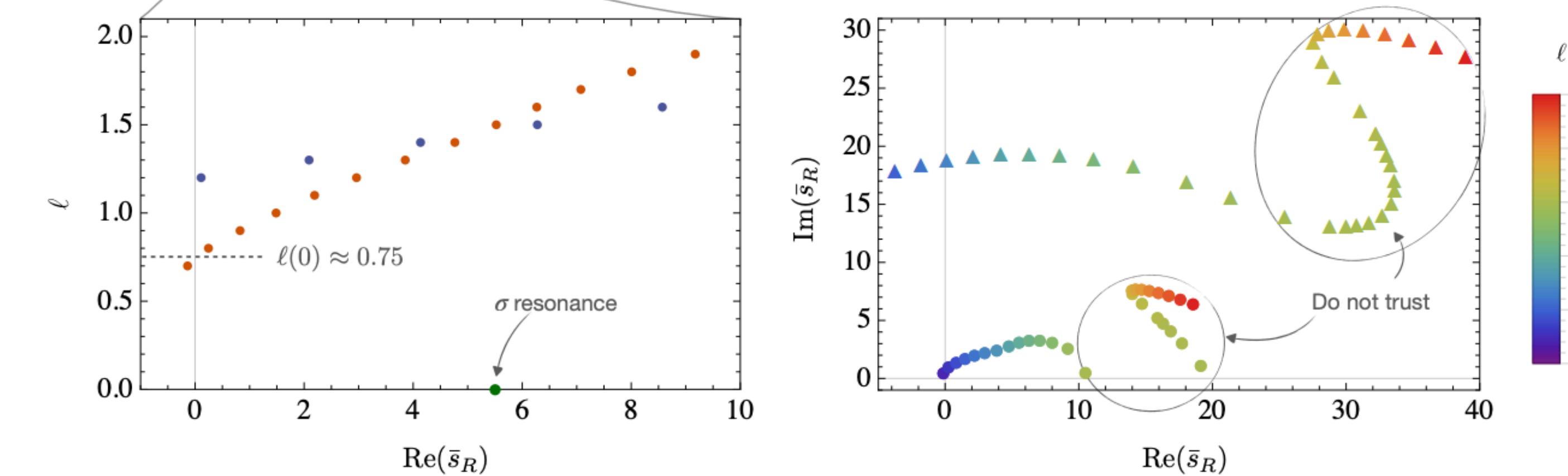
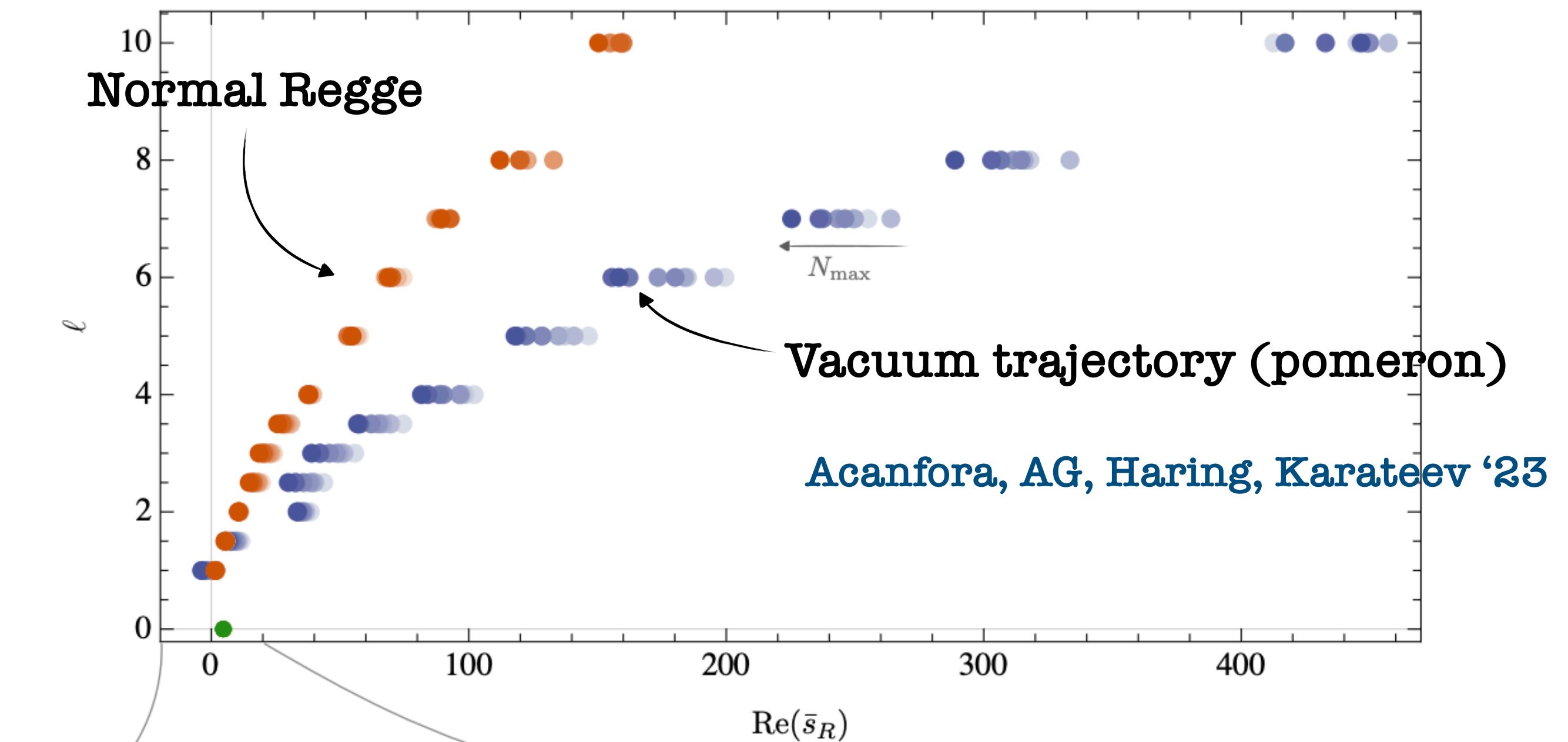
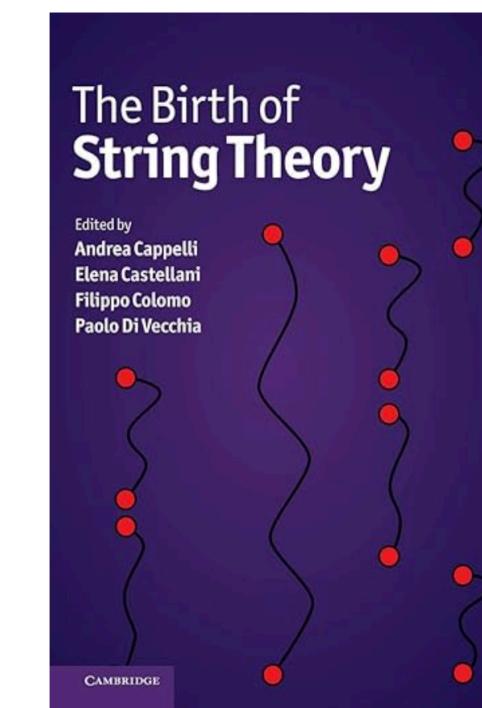
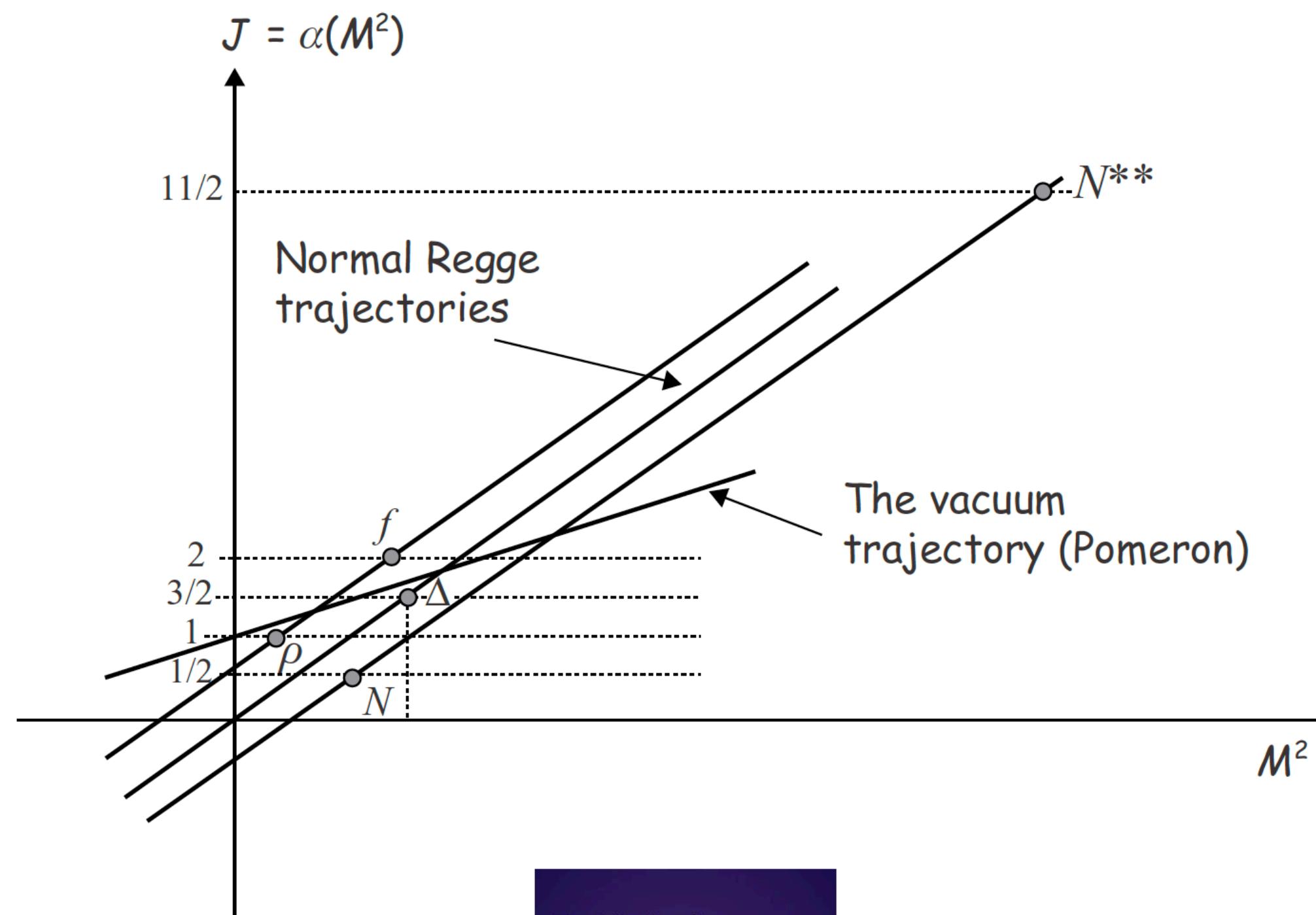
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Veneziano's picture!



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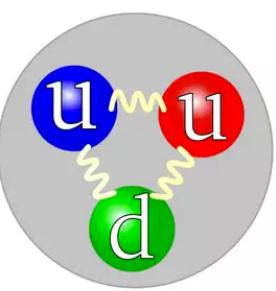
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Plan of the Talk

1) The Prehistory of String Theory: the Analytic S-matrix



2) Strings from Wightman axioms: SU(3) YM Glueball Scattering

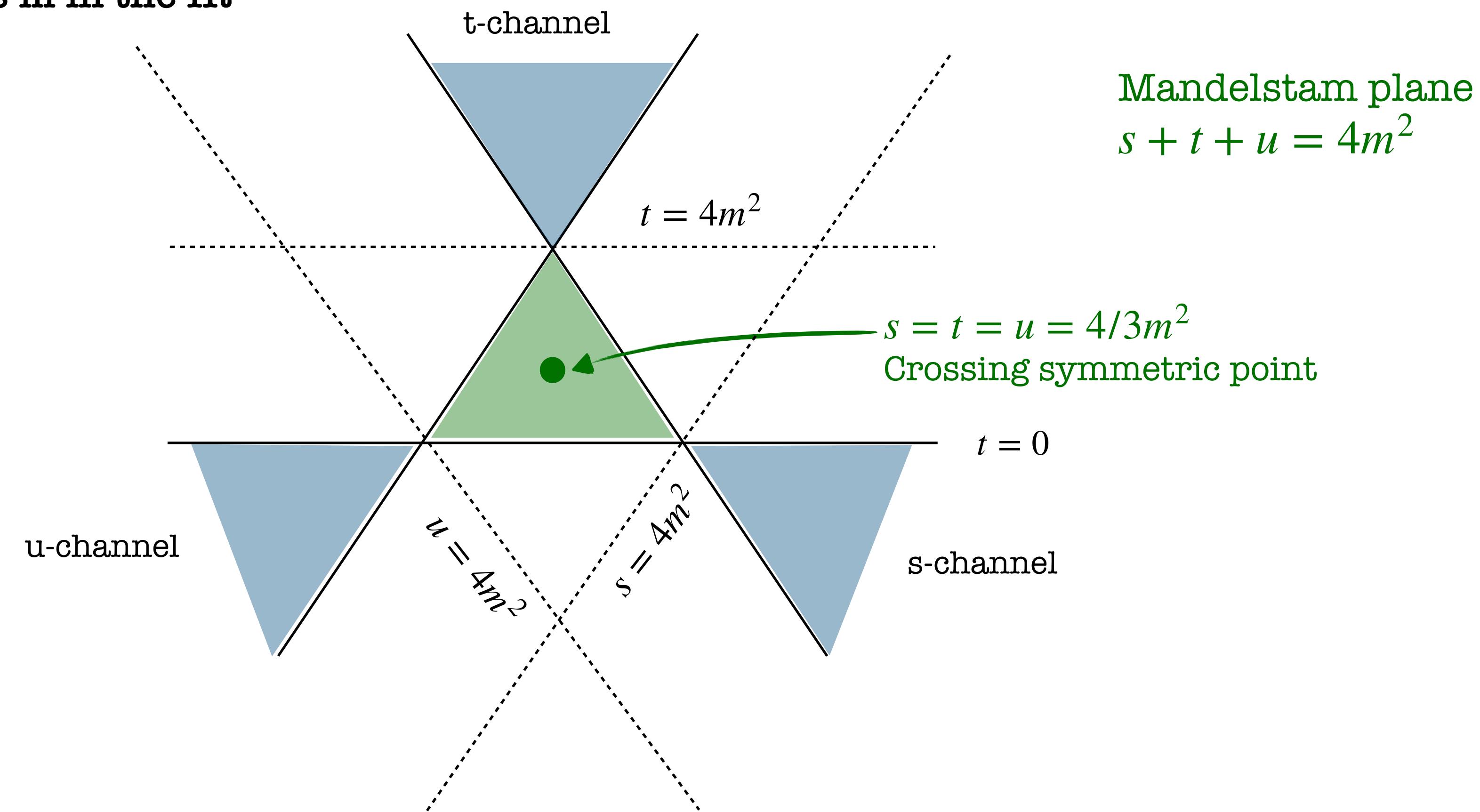


3) Strings from Gravity and no-go theorems

4) Bootstrap the world-sheet of confining Strings?

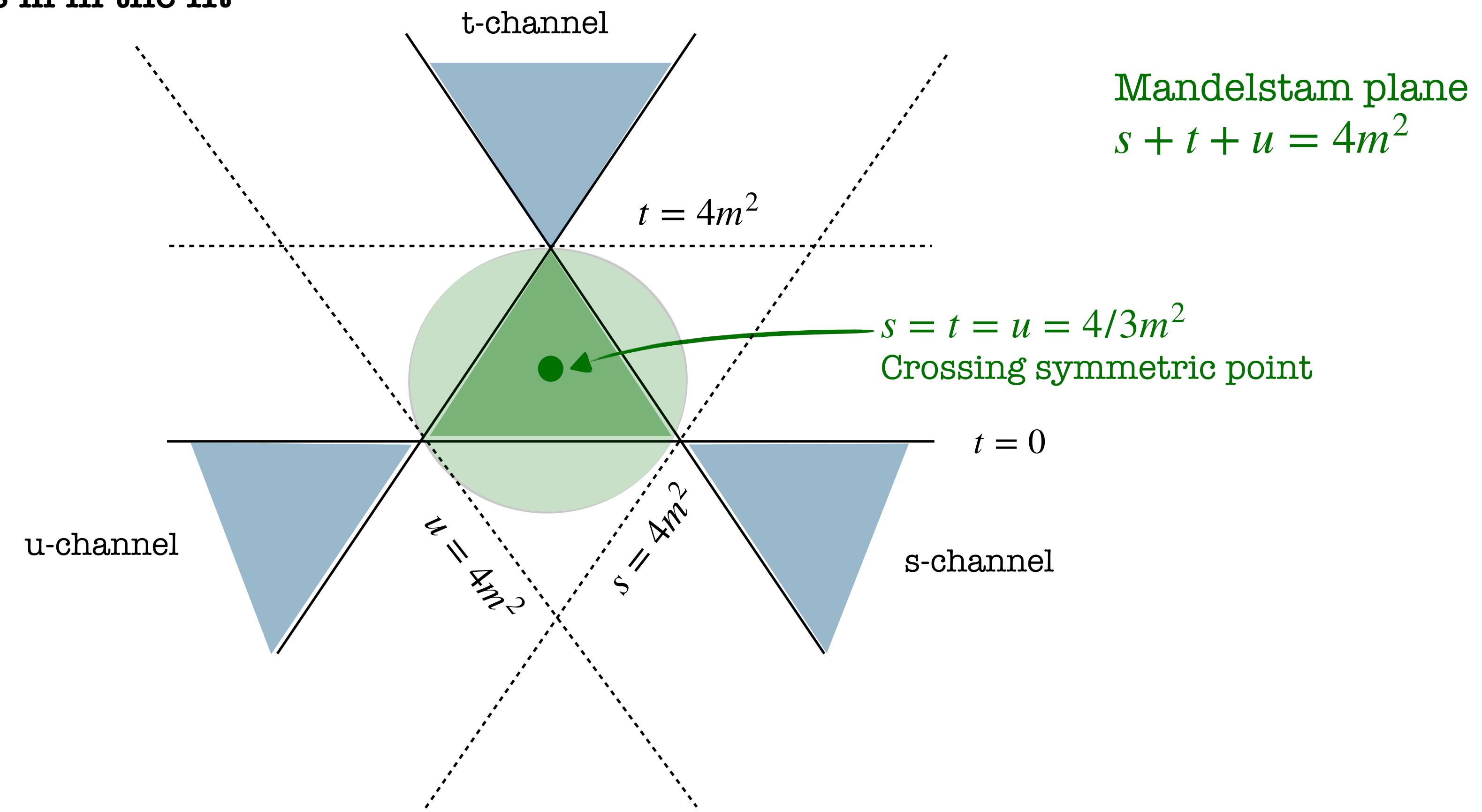
The Prehistory of String Theory: the analytic S-matrix 1

Example: 1 scalar field of mass m in the IR



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$$M(s, t, u) = -c_0 + c_2(\bar{s}^2 + \bar{t}^2 + \bar{u}^2) + c_3\bar{s}\bar{t}\bar{u} + \dots$$

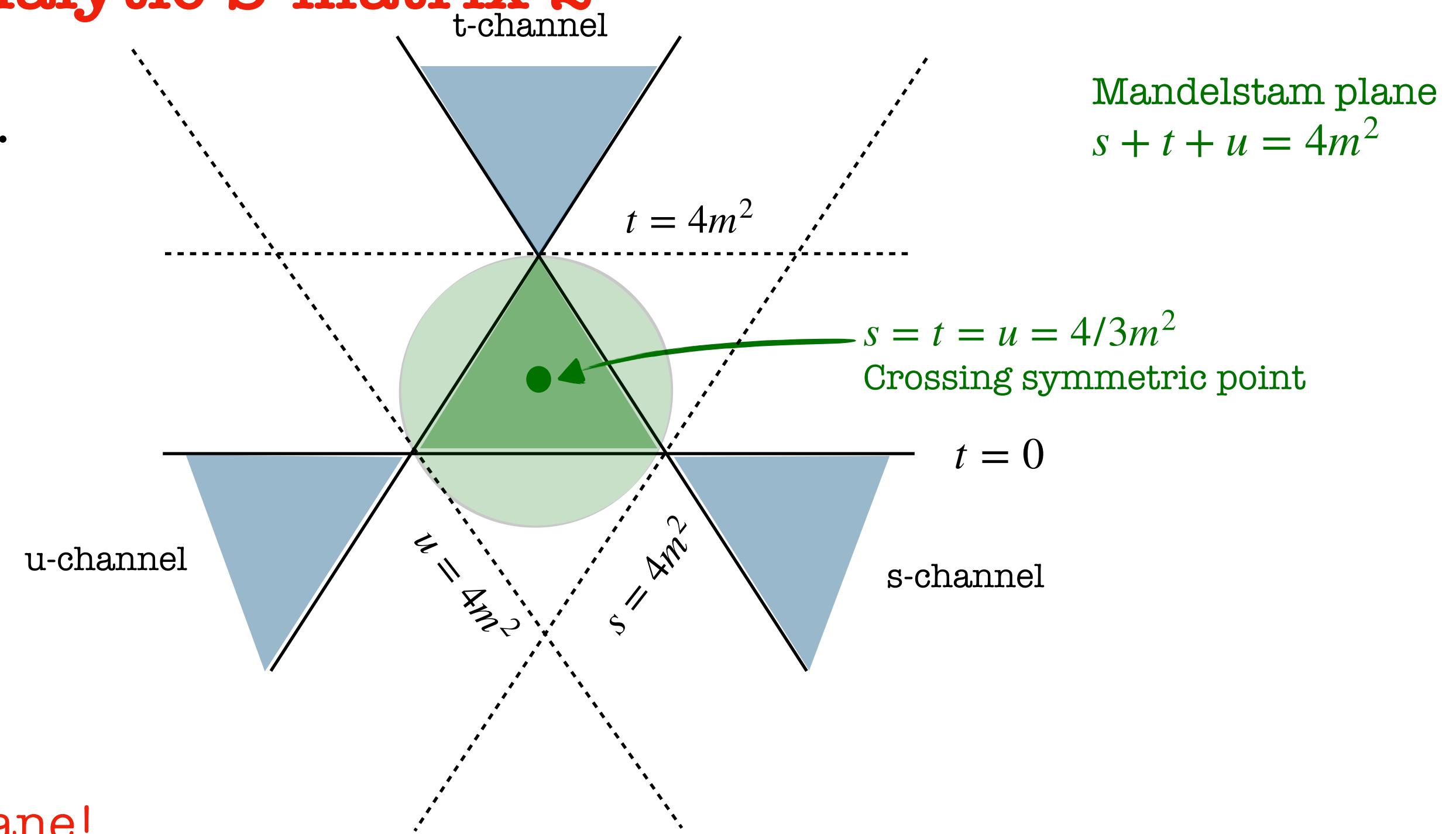
$$\bar{x} = x - \frac{4}{3}m^2$$

The set $\{c_0, c_2, c_3, \dots\}$ parametrizes the space of amplitudes

The Prehistory of String Theory: the analytic S-matrix 2

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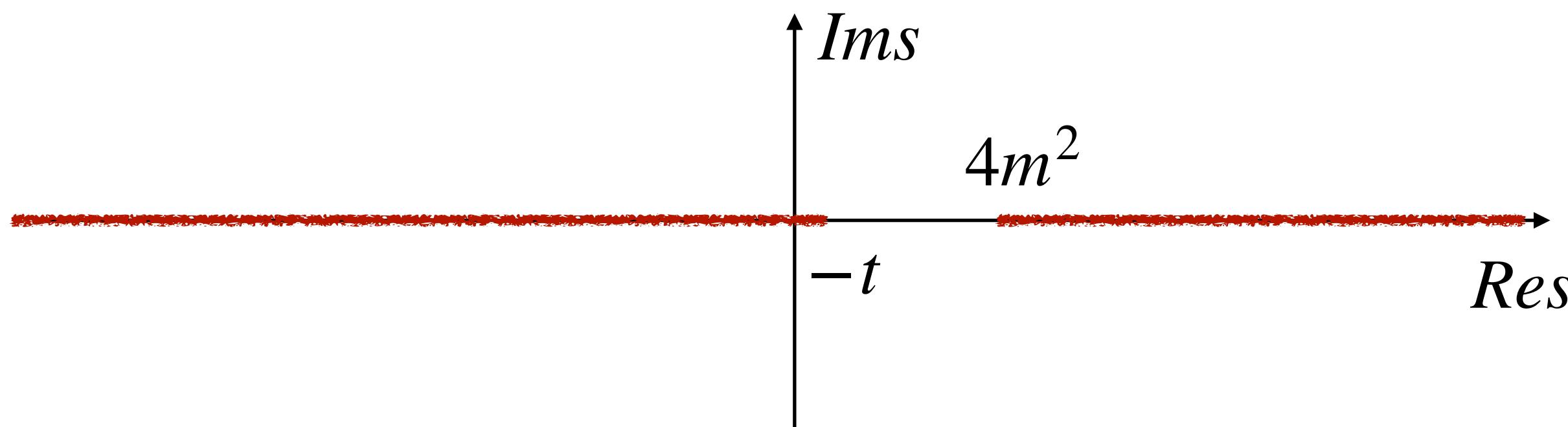
Space of amplitudes parametrized by $\{c_0, c_2, c_3, \dots\}$



Analyticity means we can go into the complex plane!

Analytic in the s-plane away from the cuts for all $-28m^2 < t < 4m^2$

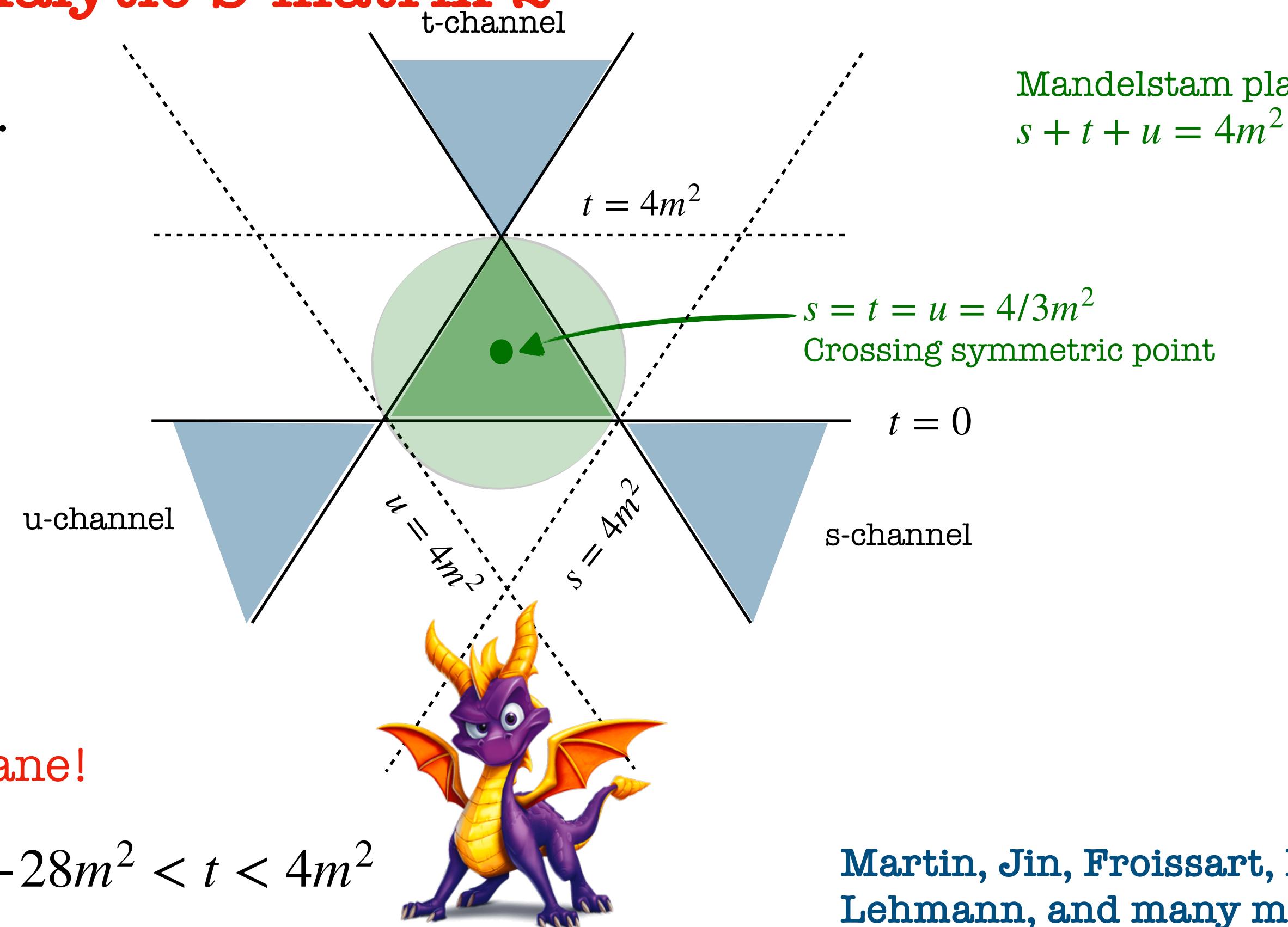
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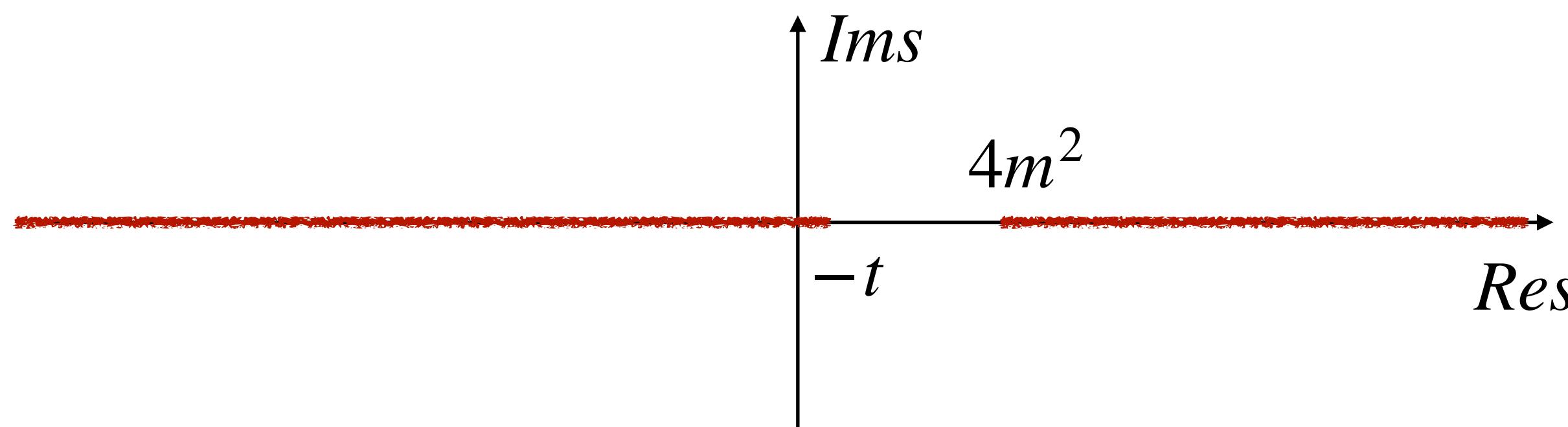
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$t < -28m^2$: double discontinuity region!

Correia, Sever, Zhiboedov 2111.12100
Tourkine, Zhiboedov 2303.08839

The beauty of Unitarity

$$M(s, t, u) = -c_0 + c_2(\bar{s}^2 + \bar{t}^2 + \bar{u}^2) + c_3\bar{s}\bar{t}\bar{u} + \dots$$

Froissart bound

Dispersive parameters \equiv operators of dimension ≥ 8

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$$M(s, t, u) = -\textcolor{red}{c}_0 + \boxed{c_2(\bar{s}^2 + \bar{t}^2 + \bar{u}^2) + c_3\bar{s}\bar{t}\bar{u} + \dots}$$

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$$t_0 = s_0 = 4/3m^2$$

Subtraction point

$$T_v(v, t_0) \equiv 16\pi \sum_{\ell=0}^{\infty} (2\ell+1) \textcolor{red}{Im}f_\ell(s) P_\ell(1 + 2t_0/(s-4)) \geq 0$$

Positivity

Legendre positivity

$$P_\ell(x) > 0, \quad x \geq 1$$

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NOT POSITIVE!

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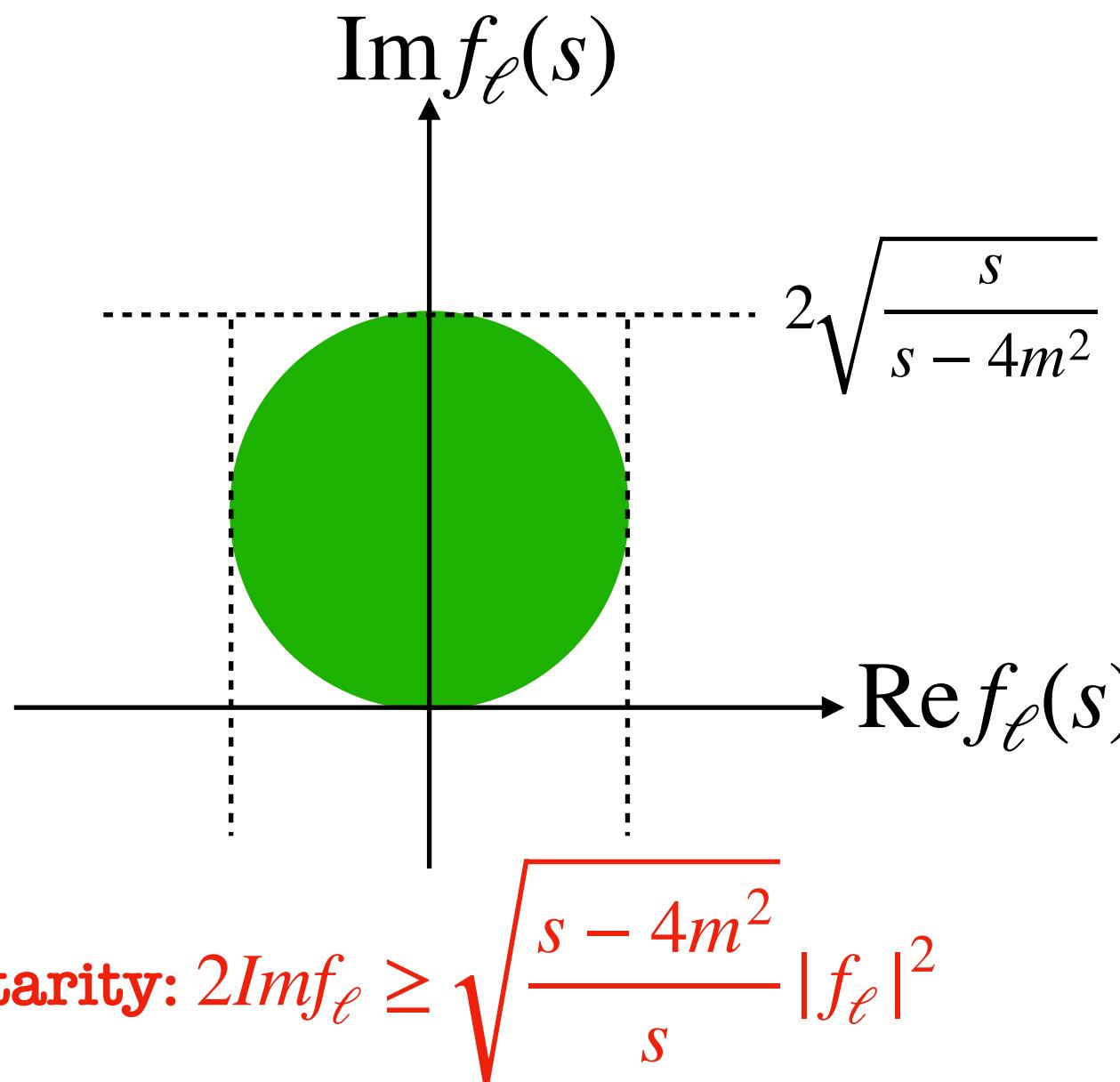
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Unitarity saves the day!



The Island of 4d scalar amplitudes

We bound c_0, c_2 using dispersion relations and unitarity!
It is an exercise in constrained optimization theory.

Bonnier, Lopez, Mennessier, '70s
AG, Sever 2106.10257

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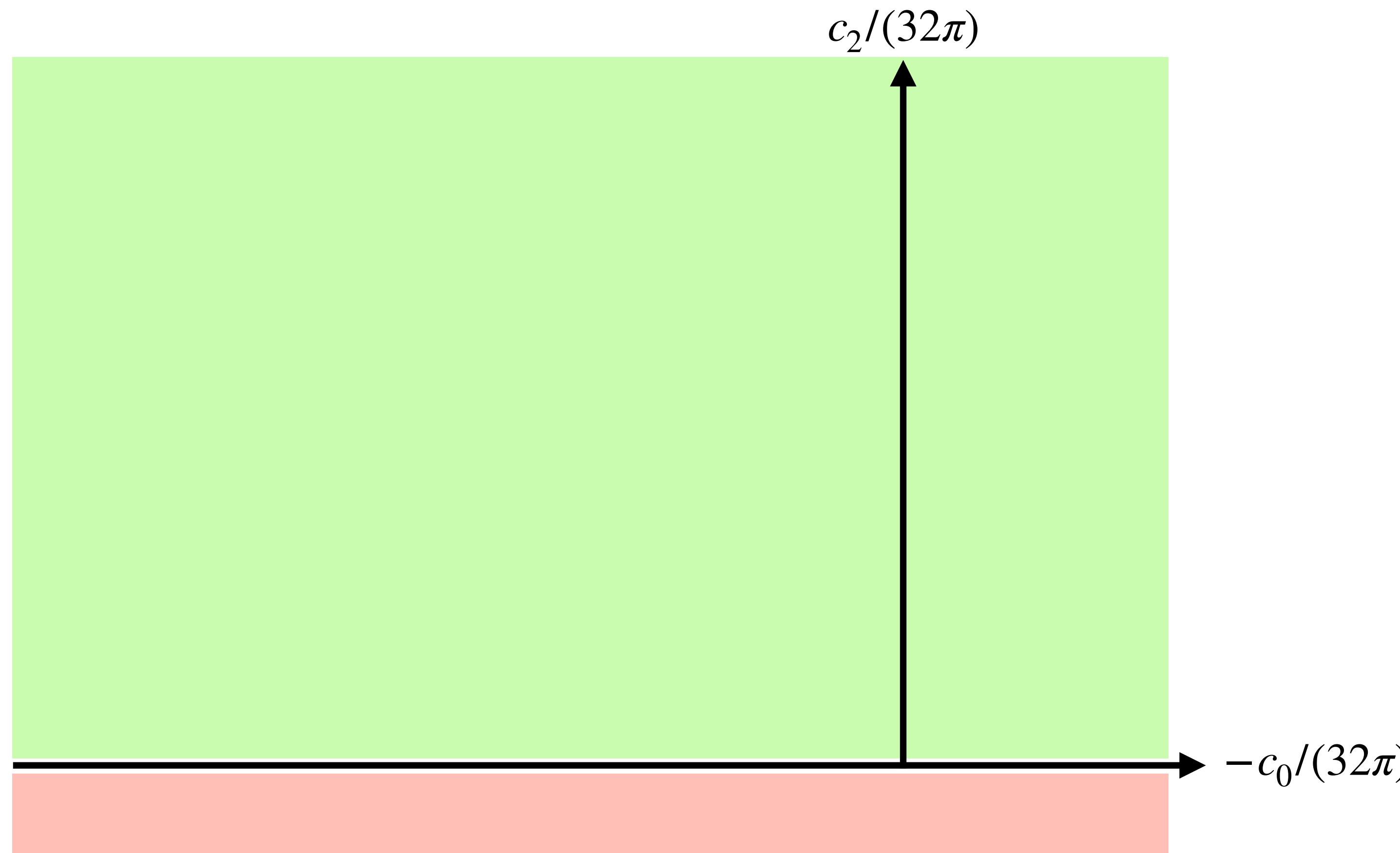
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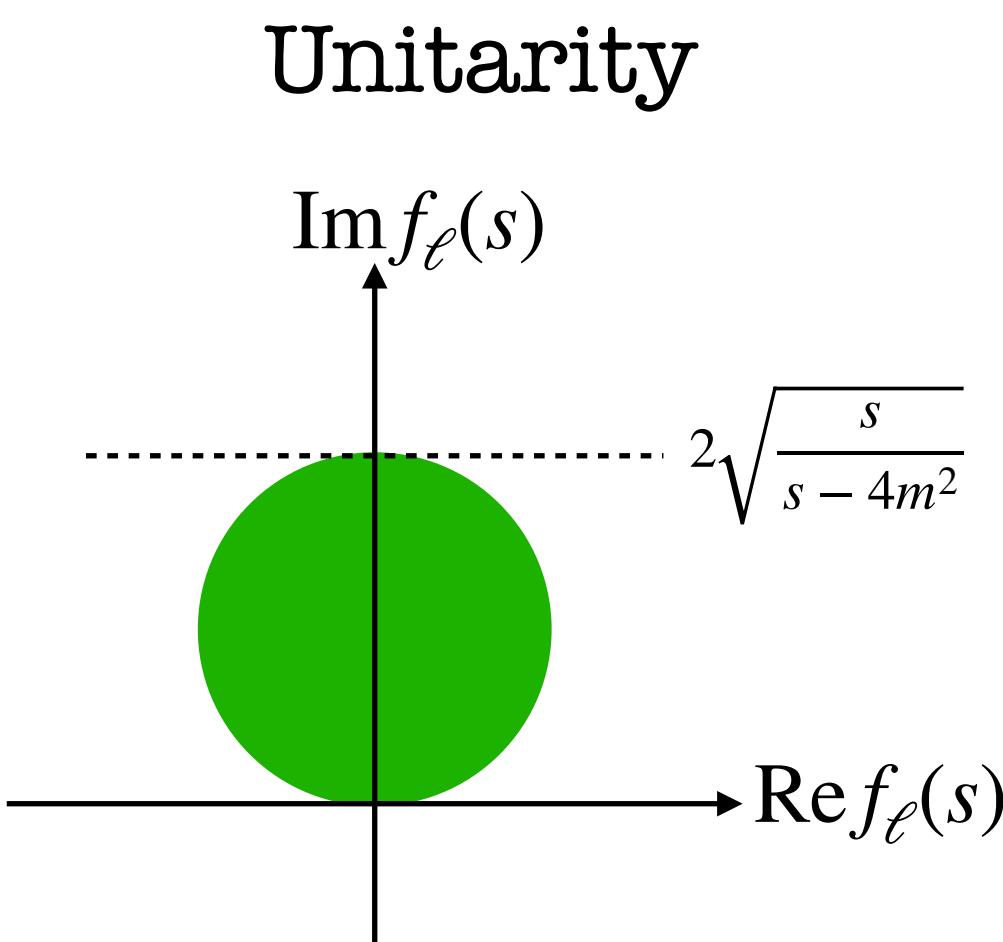
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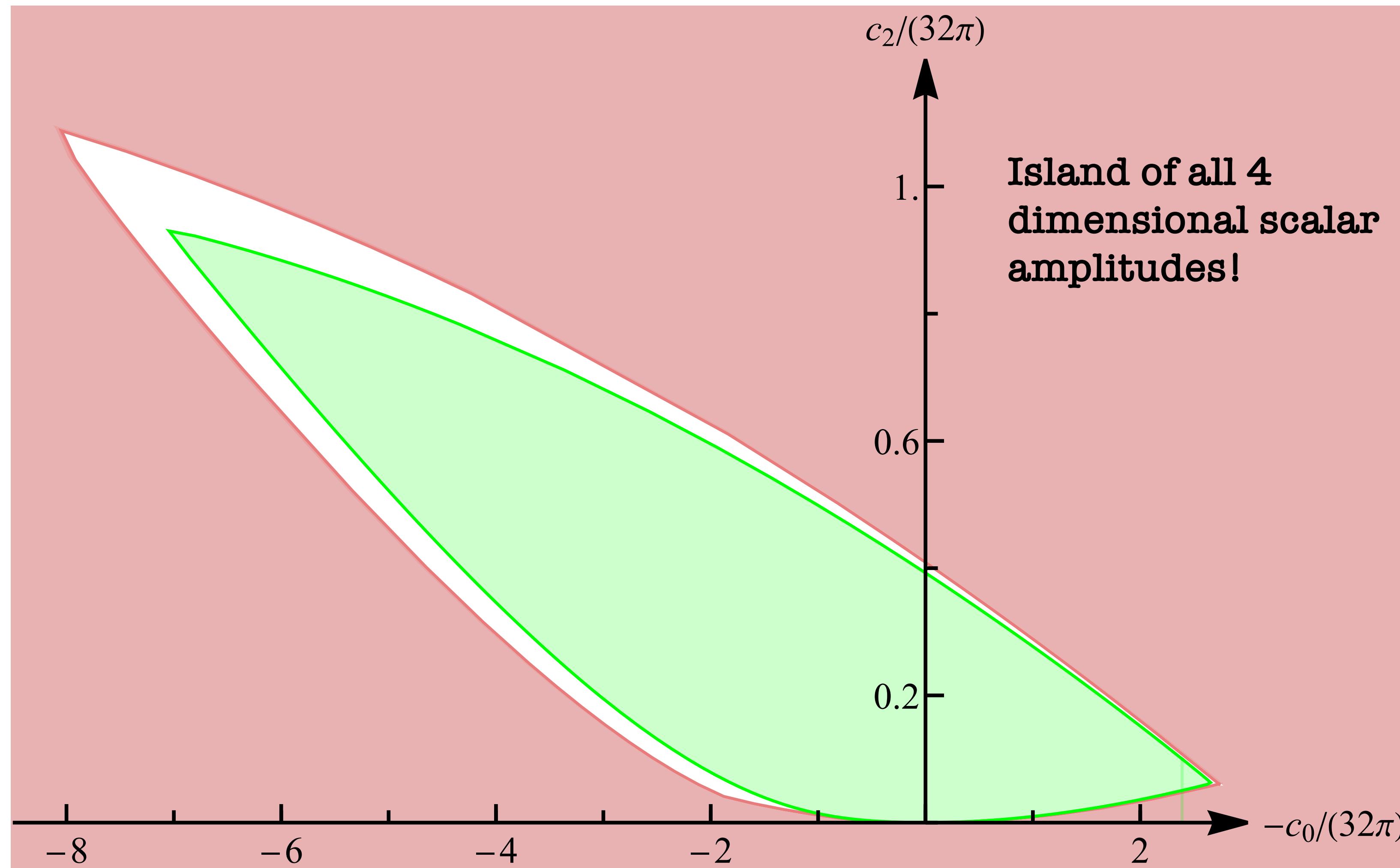
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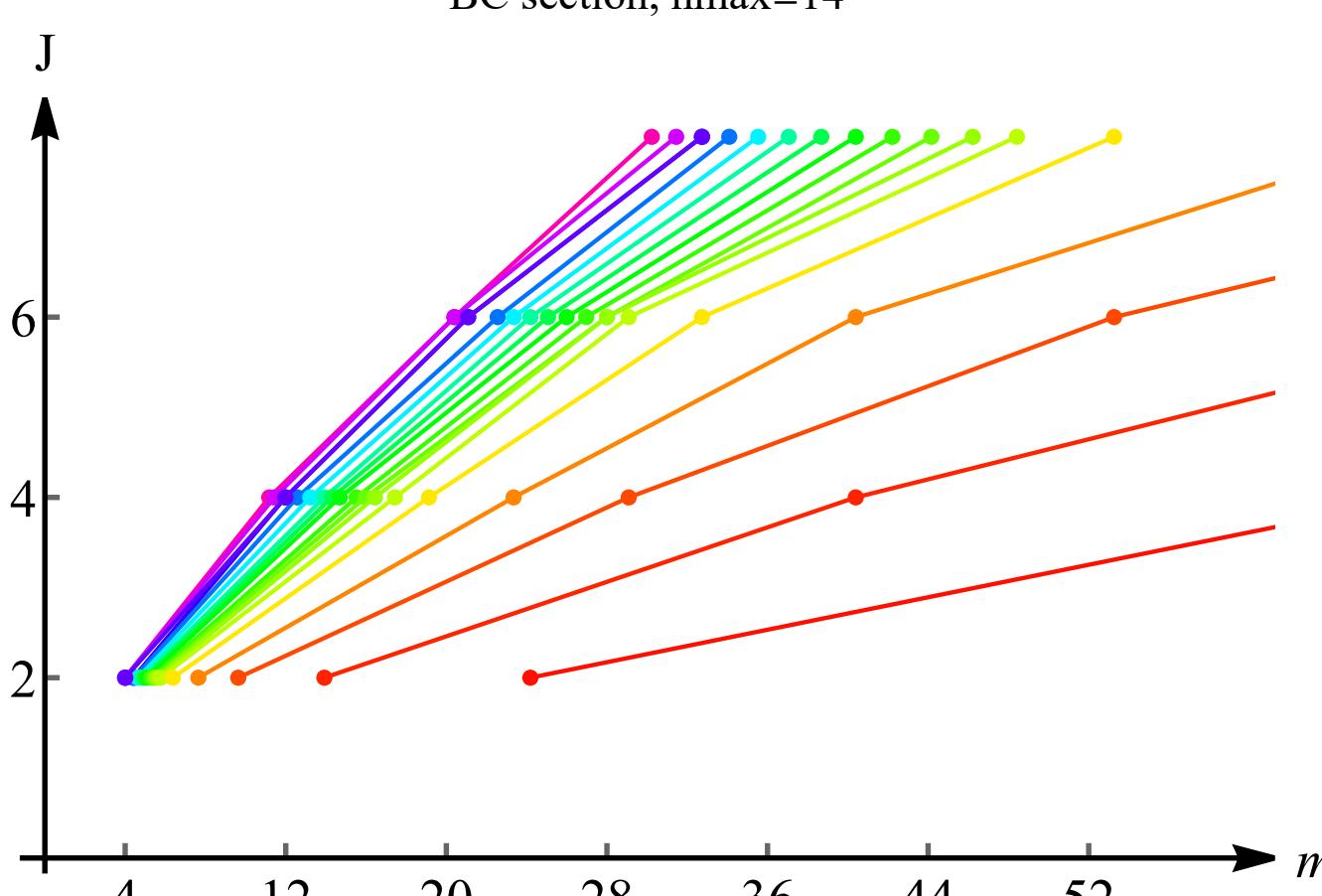
AG, Sever 2106.10257



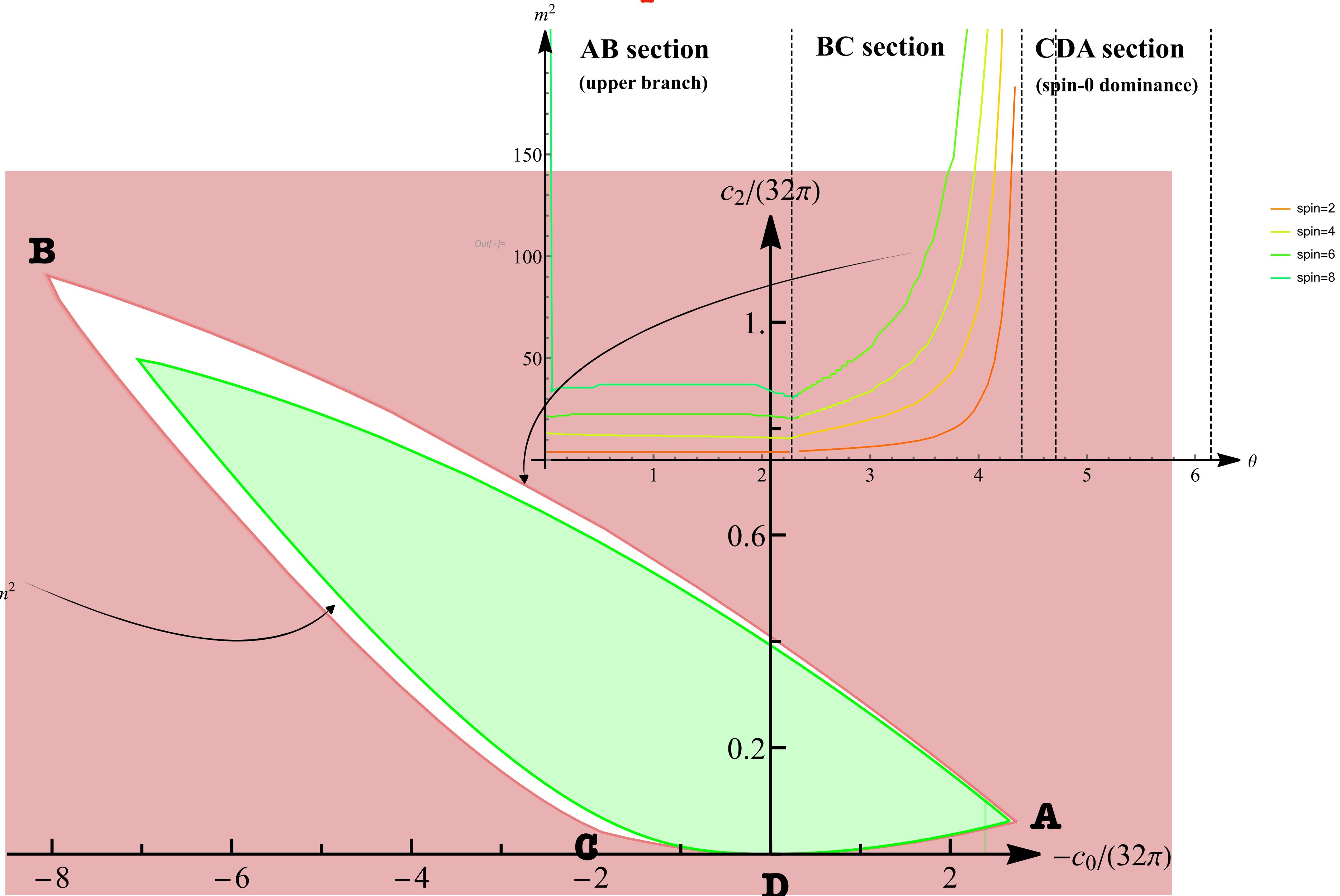
Elias-Miro, AG, Gümüş 2210.01502
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The Island of 4d scalar amplitudes



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Glueball Scattering in SU(3) pure YM

Regime in which the S-matrix Bootstrap shows its power: cutoff $\Lambda = 2m$, no small parameters.

(Still Hard after 50 years of QCD)

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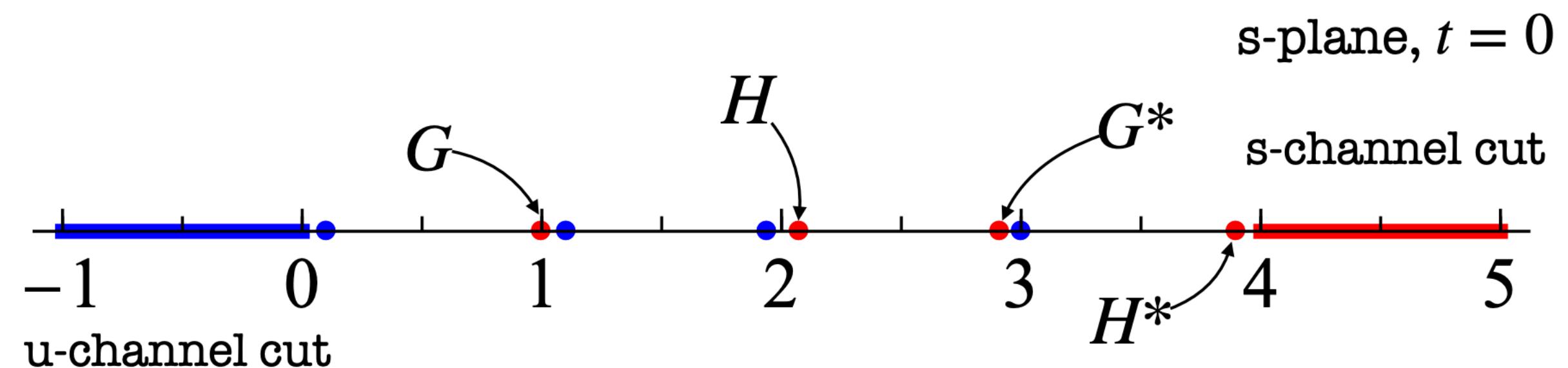
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Stable Glueballs spectrum

	J^{PC}	Mass
G	0^{++}	1
H	2^{++}	1.437 ± 0.006
G^*	0^{++}	1.72 ± 0.01
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Pole Structure in GG->GG scattering



Athenodorou, Teper 2007.06422, 2106.00364

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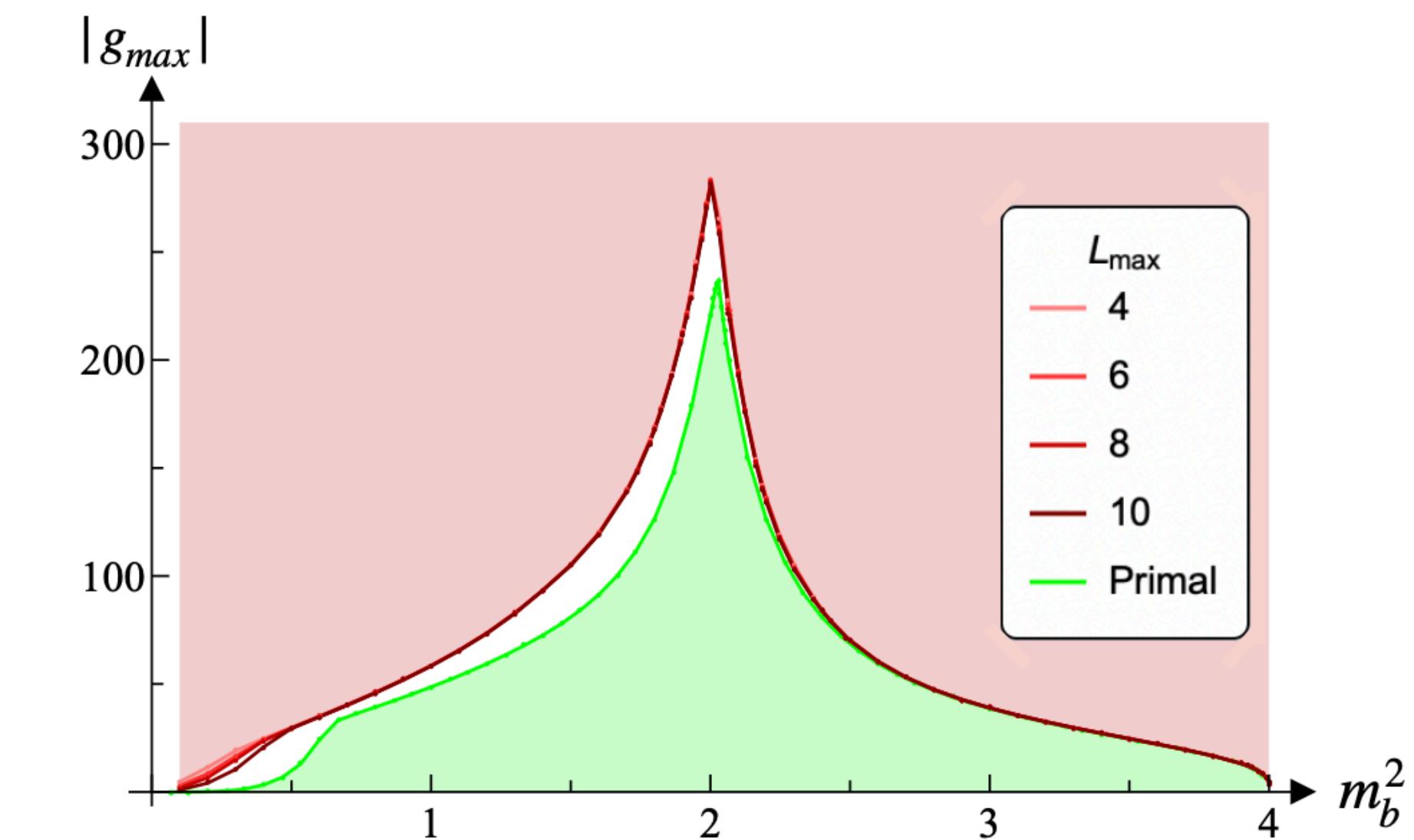
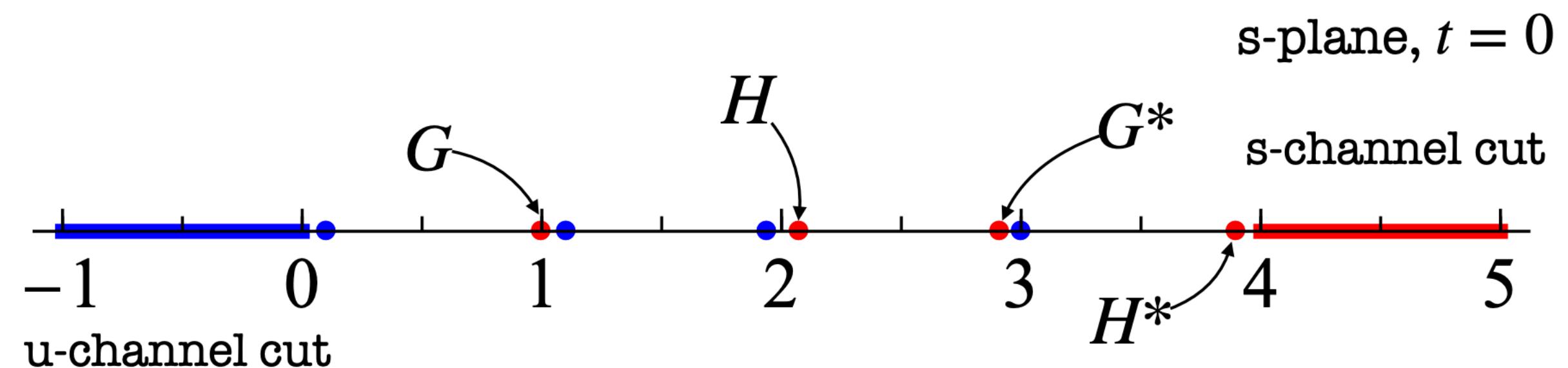
The maximum residue at the spin-0 pole is a simple problem

$$M \supset \frac{-g^2}{s - m_b^2} + \frac{-g^2}{t - m_b^2} + \frac{-g^2}{u - m_b^2} + \dots$$

Paulos, Penedones, Toledo, van Rees, Vieira
1708.06765

AG, Hebar, van Rees 2312.00127

Pole Structure in GG->GG scattering



Toy model: amplitudes maximizing the spin-2 coupling

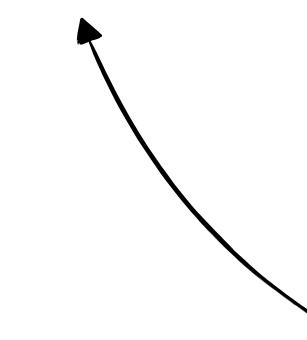
The maximum residue at the spin-2 pole is a hard problem (\mathbb{Z}_2 symmetry, no $\frac{1}{s - m^2}$ pole)

$$M \supset \frac{-g^2}{s - m_b^2} P_2 \left(1 + \frac{2t}{m_b^2 - 4m^2} \right) + \dots \sim t^2$$

AG, Hebbar, van Rees 2312.00127

Without Regge it would violate unitarity!

They must restore $M(s \rightarrow \infty, t \leq 0) < s \log^2 s$



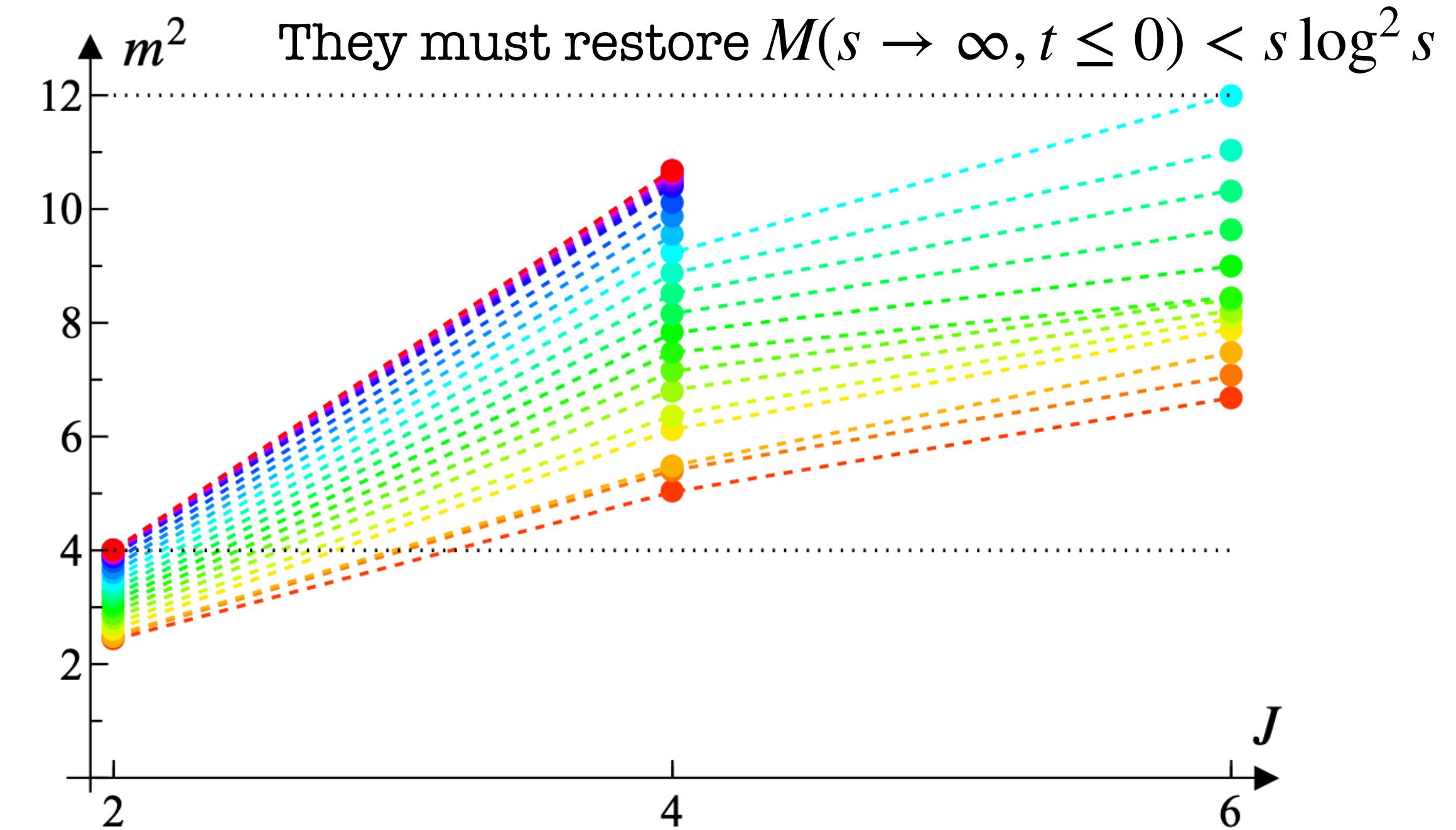
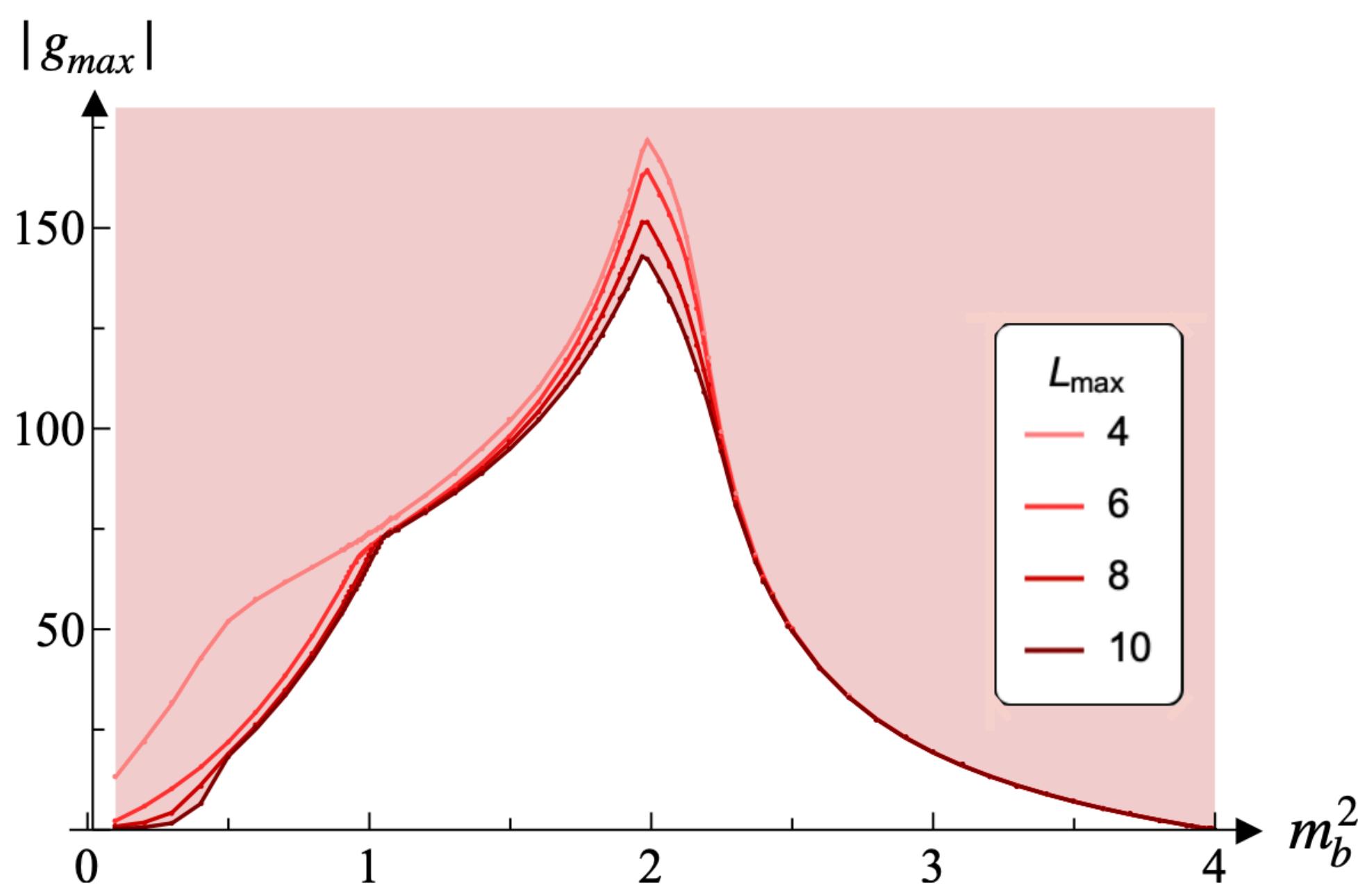
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The Glue-Hedron

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AG, Hebbar, van Rees 2312.00127

$\max g_G $	$\max g_H $	$\max g_{G^*} $	$\max g_{H^*} $
213	158	224	2.15
206	156	217	—

SU(3) YM Lattice $g_G \approx 50 \pm 7$

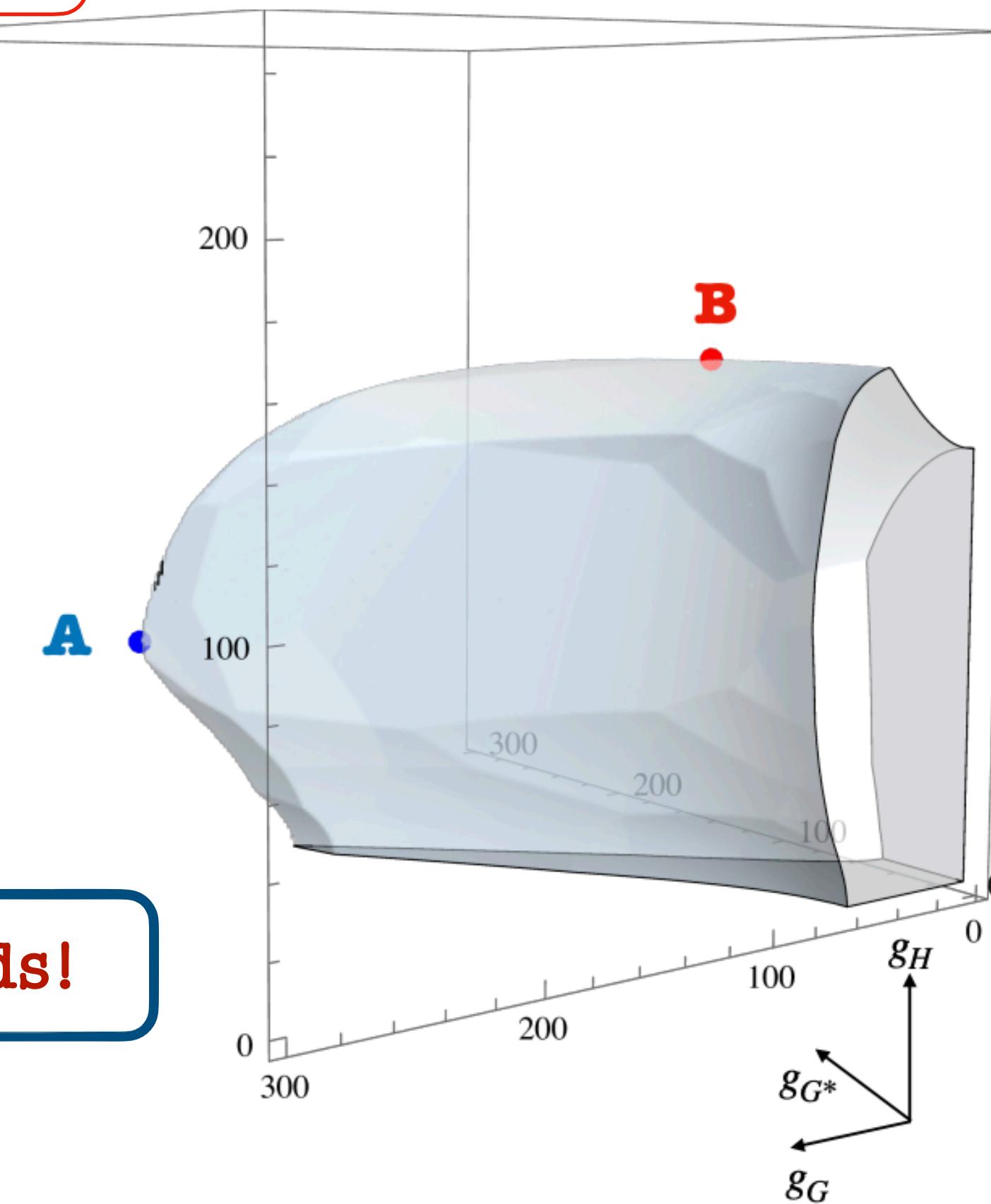
De Forcrand, Schierloz, Schneider, Teper '85

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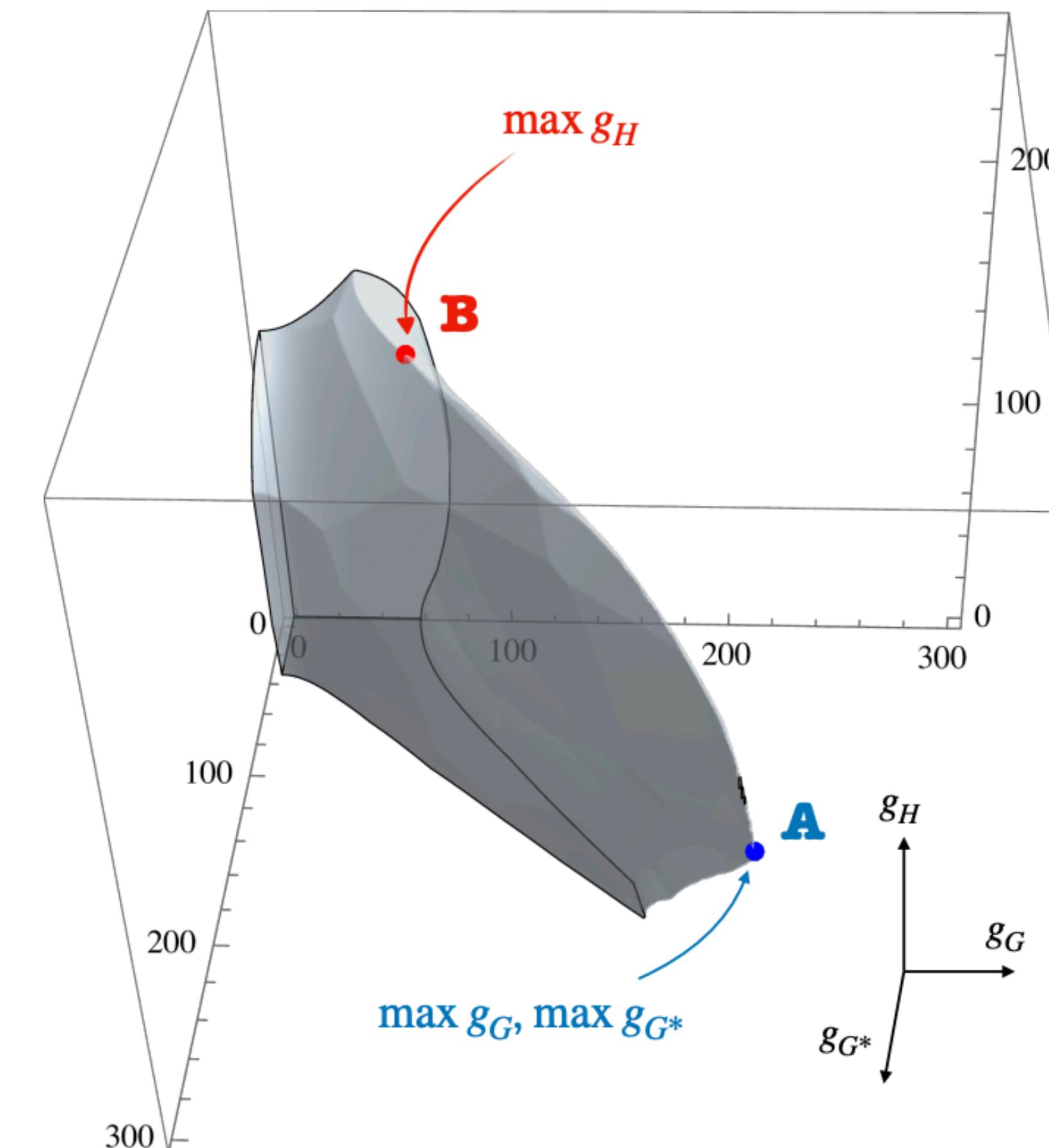
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Rigorous Bounds!





QG Bootstrap and No Go Theorems

D=10, Maximal Susy, turn off all couplings except G_D

$$A_{QG} = \int \sqrt{-g} (R + 0 \times R^2 + 0 \times R^3 + \alpha_D R^4 + \dots)$$

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Amplitude for graviton scattering in the EFT

$$A(s, t, u) = 8\pi G_D \left(\underbrace{\frac{1}{stu}}_{\text{Sugra}} + \alpha_D \ell_P^6 + \mathcal{O}(s \log s) \right)$$

First quantum correction $\alpha_D R^4$

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1) **No-go Theorem:** If gravity is quantum, it must be corrected at leading order!

Can be violated only if there are new physical principles!

- AG, J. Penedones and P. Vieira, [2102.02847](#)
- AG, H. Murali, J. Penedones and P. Vieira, [2212.00151](#)

E.g. D=10

$$\alpha(\text{Bootstrap}) \geq 0.126 \pm 0.006$$

$$\alpha(\text{String Theory}) \geq 3^{1/4} \zeta\left(\frac{3}{2}\right) \left(\zeta\left(\frac{3}{2}, \frac{1}{2}\right) - \zeta\left(\frac{3}{2}, \frac{2}{3}\right) \right) / \sqrt{2} \simeq 0.1389\dots$$

The space of deformations compatible with String Theory!

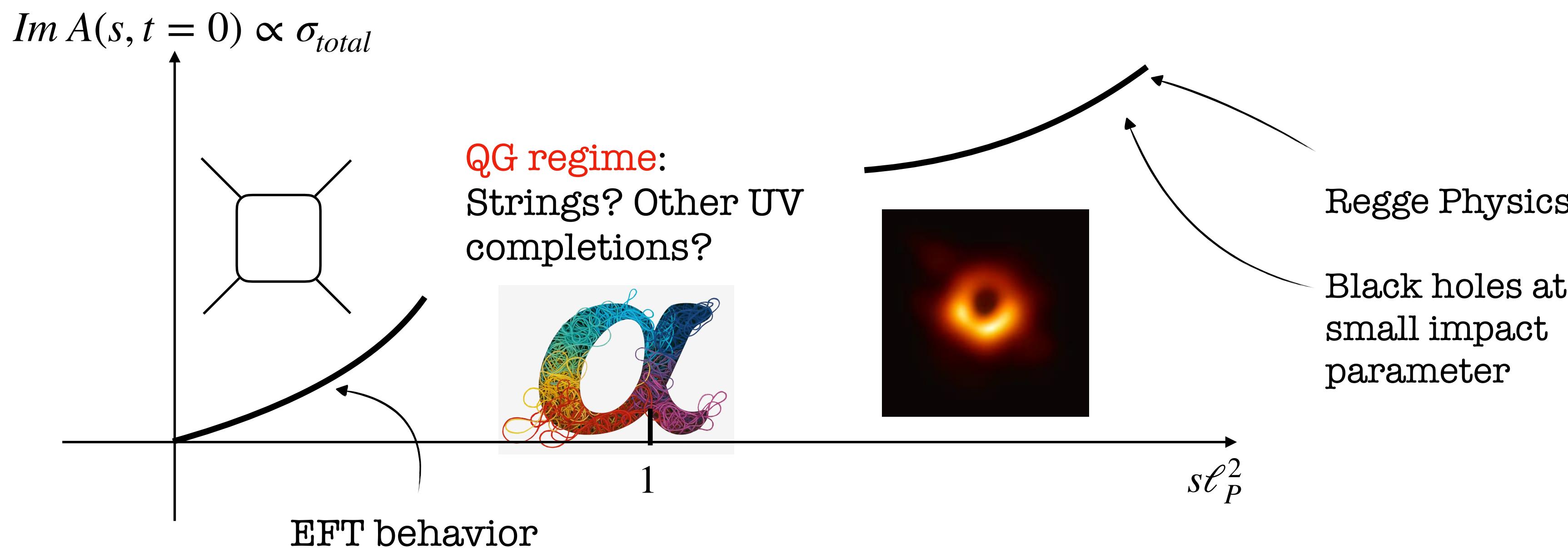
Wilson coefficients as probes of gravity UV completions

Can α_D take any value?

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Can α_D take any value?

α_D knows about the theory at all scales! $\alpha_D \propto \int_0^\infty \frac{\text{Im}A(s, t=0)}{s} ds \geq 0$



QG Bootstrap: Methodology

- There exist a parametrization of the non-perturbative amplitude manifestly **crossing symmetric** and **analytic**

$$A(s, t, u) = \underbrace{\frac{8\pi G_D}{stu}}_{\text{Sugra}} + \prod_{X=s,t,u} (\rho_X + 1)^2 \sum_{a+b+c \leq N} \nu_{(a,b,c)} \rho_s^a \rho_t^b \rho_u^c$$

UV completion

Crossing

Analyticity

$$\rho(s) = \frac{1 - \sqrt{-s}}{1 + \sqrt{-s}}$$

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- Impose **unitarity** numerically (the $\nu_{(a,b,c)}$ cannot vary arbitrarily)

Linear operation: $T(s, x) \rightarrow S_\ell(s) \quad |S_\ell(s)|^2 \leq 1, \quad s > 0, \quad \ell = 0, 2, \dots, \infty$ (10D)

$$S_\ell(s) = 1 + i \frac{s^3}{2^{18} 3 \pi^4} \int_{-1}^1 (1 - x^2)^3 \frac{C_\ell^{7/2}(x)}{C_\ell^{7/2}(1)} T(s, x)$$

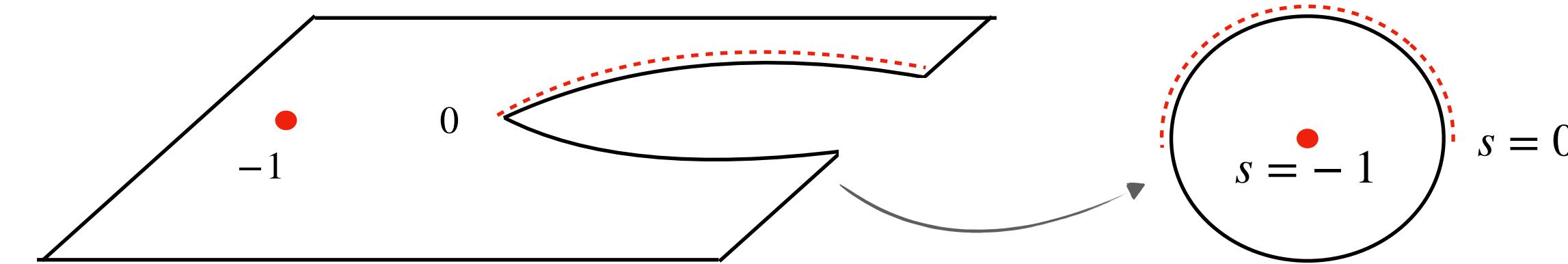
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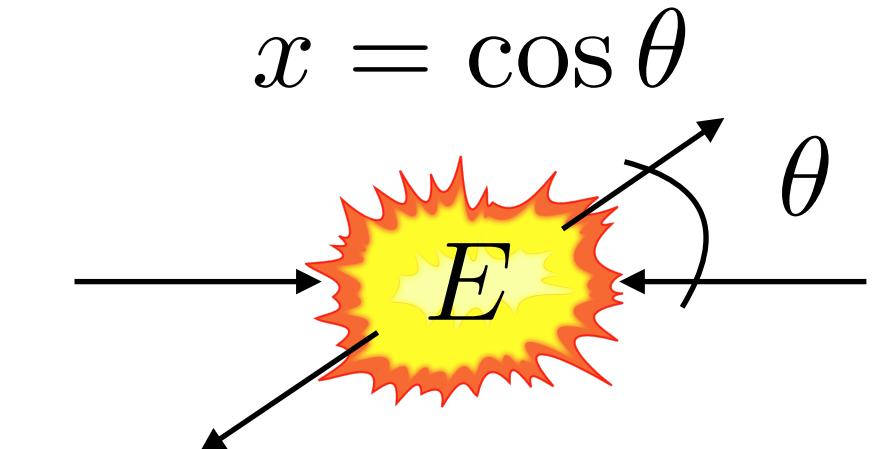
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(10D)

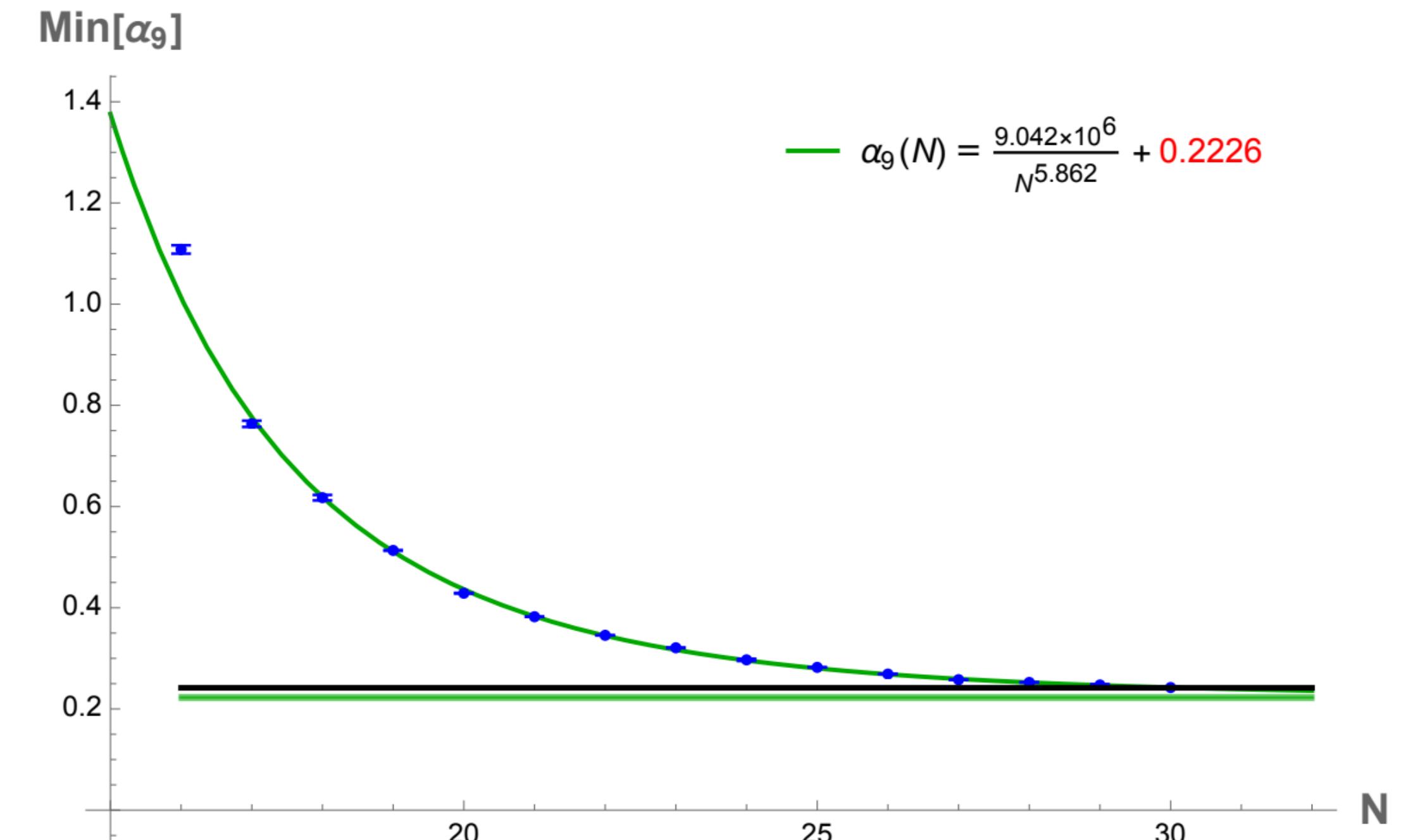
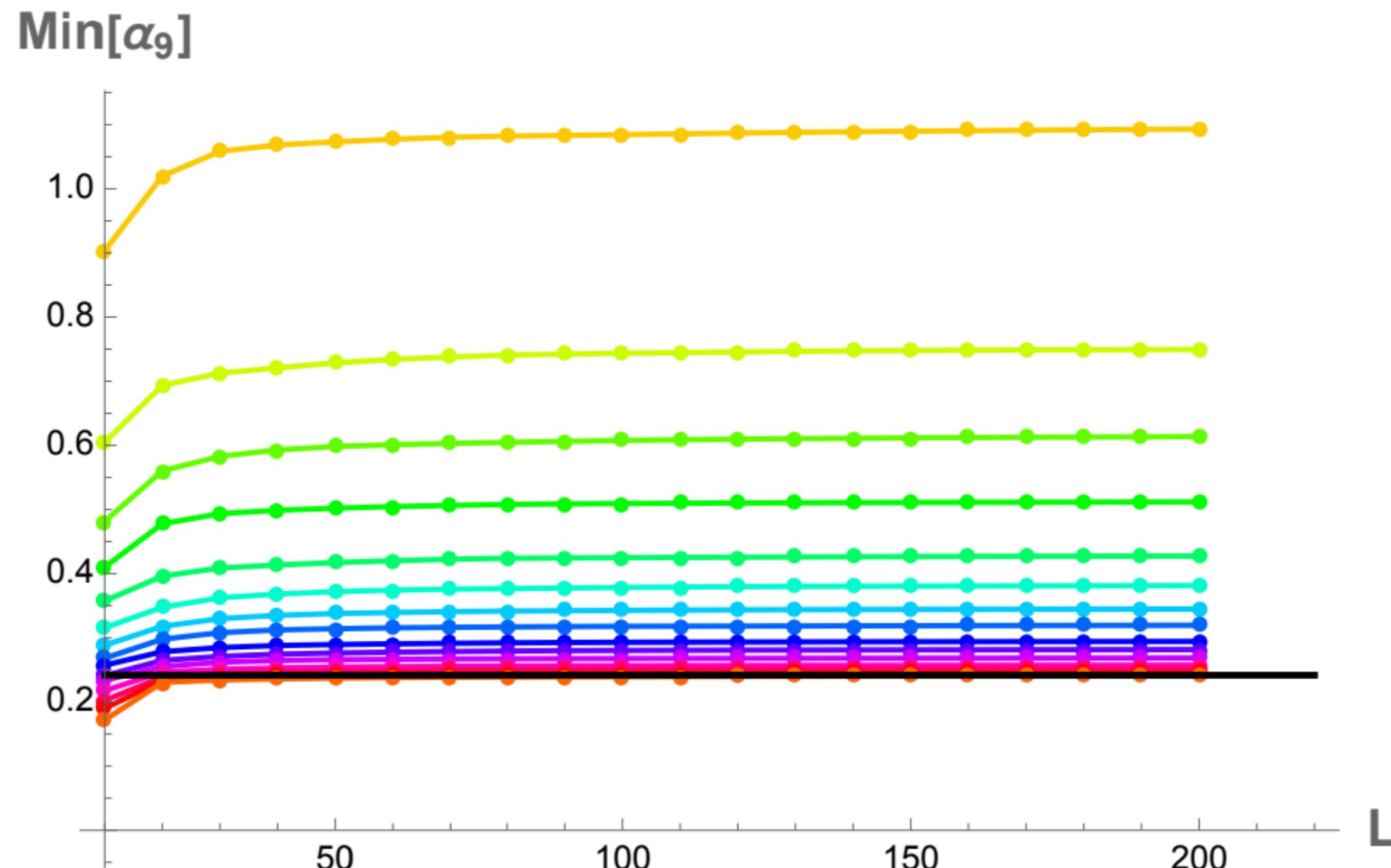
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- **Bootstrap as a semi-definite optimization problem**

FindMinimum $\alpha_D[\nu_{(abc)}]$ subject to unitarity

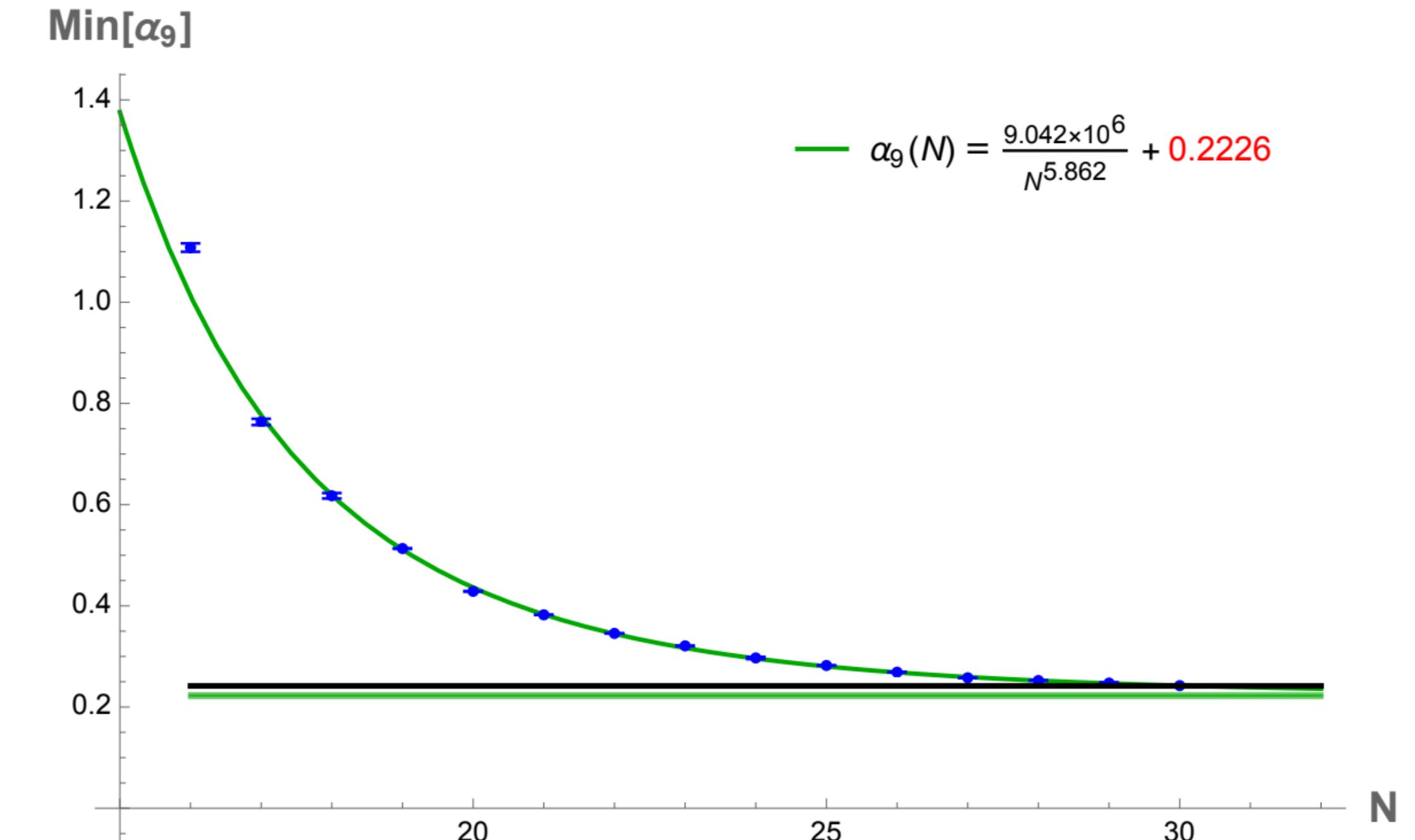
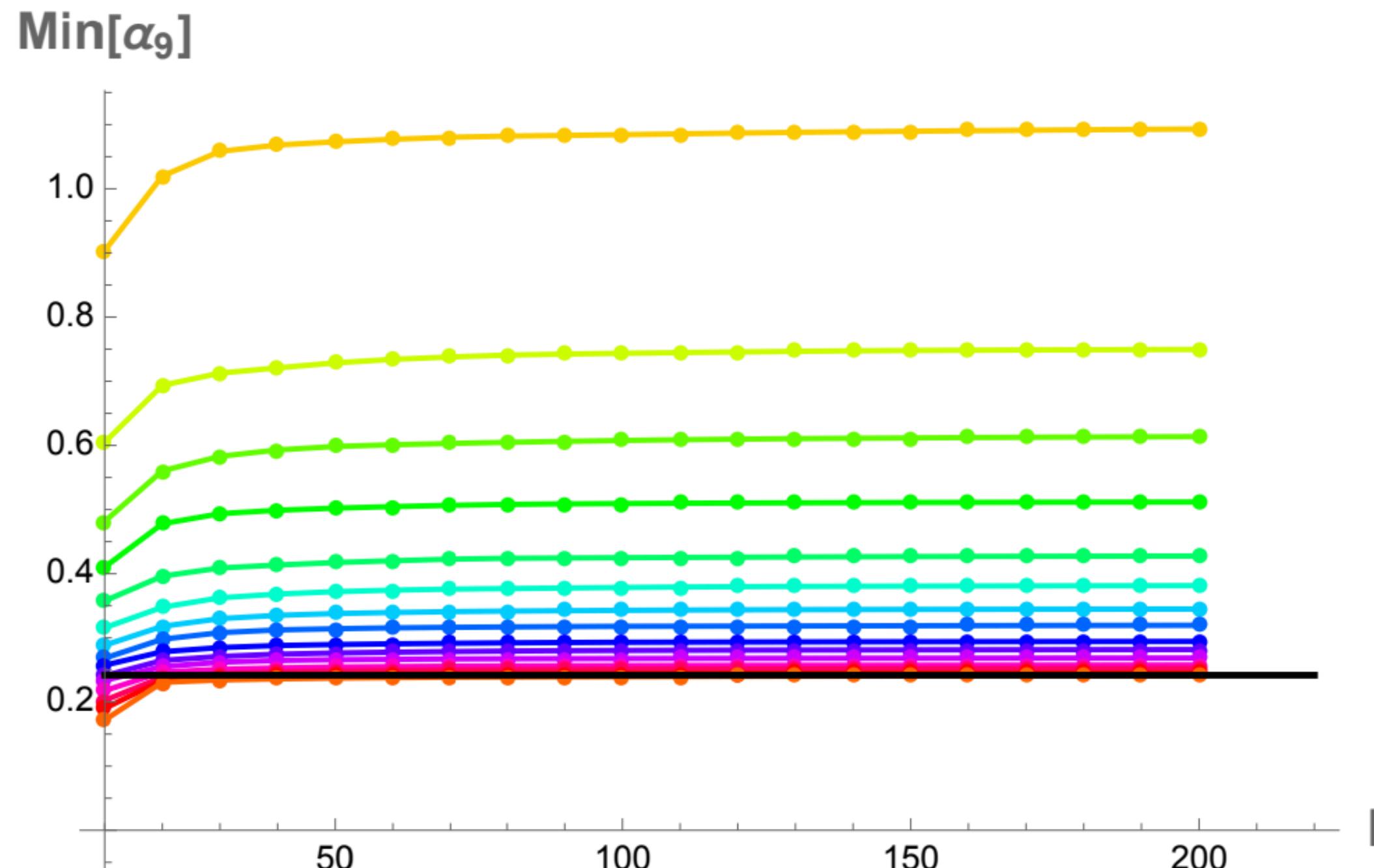


QG Bootstrap: What we learn 1



Existence of Universal lower bound depending on low energy SUGRA, analyticity, crossing, and unitarity

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1. The Bound on α_D

$$\alpha_D^{\min} < \alpha_D < \infty$$

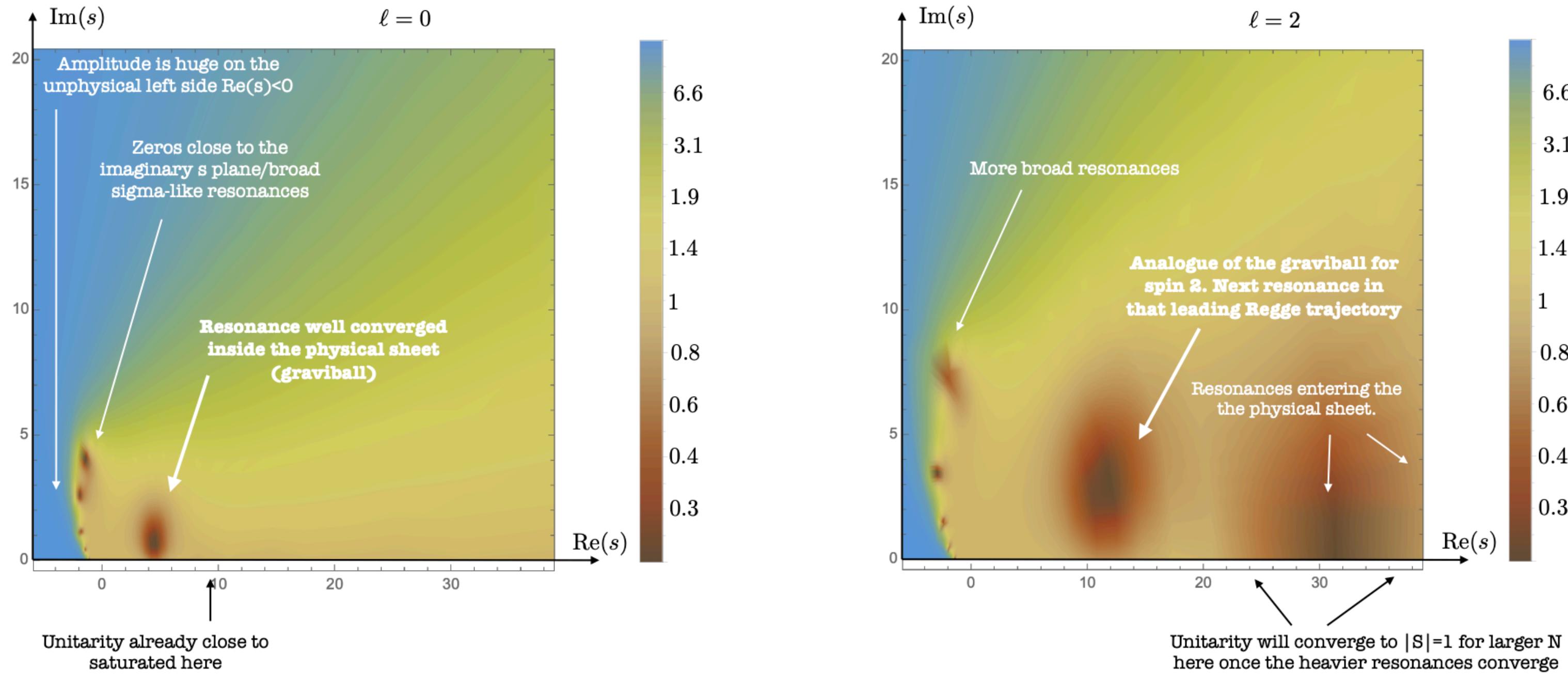
Dimension	String/M theory	Bootstrap α_D^{\min}
9	≥ 0.2411	0.223 ± 0.002
10	≥ 0.1389	0.124 ± 0.003
11	0.1304	0.101 ± 0.005

- a) String Theory in 9, and 10 dimensions almost saturates the allowed region for α
- b) α for M-theory is close to the boundary of the allowed region

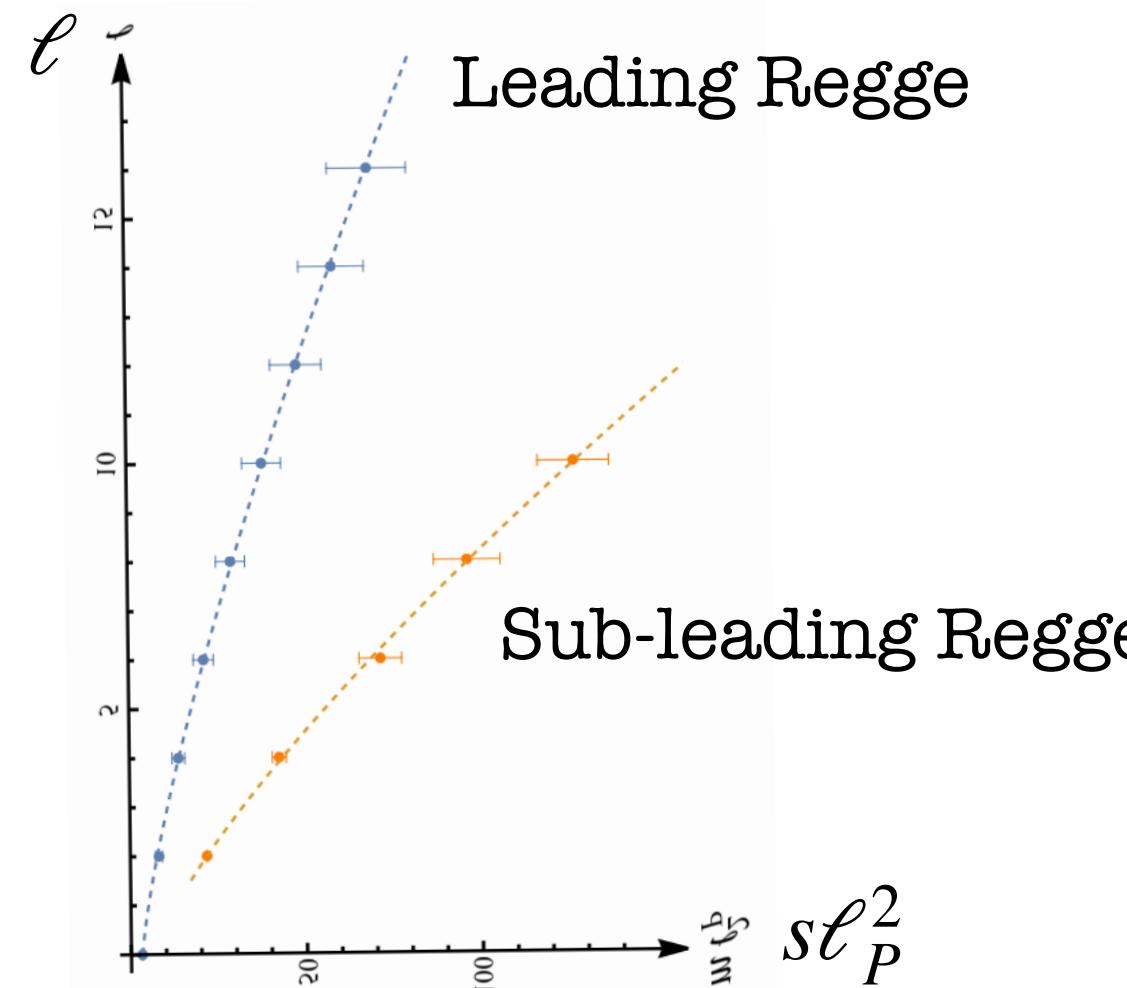
Green, Gutperle
Green, Vanhove
Green, Russo, Vanhove

hep-th/9701093
hep-th/9704145
hep-th/0610299

QG Bootstrap: What we learn 2



We can reconstruct the solution that minimizes α_D and study this non-perturbative amplitude



Resonance spectrum organizes in (curved) Regge trajectories

Stringy Spectrum although there is no assumption about the UV completion

Can we directly Bootstrap the world-sheet of the Hadronic String?

String Theory from Gravity



What about the Hadronic String Theory?

Can we directly Bootstrap the world-sheet of the Hadronic String?

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What about the Hadronic String Theory?

Simplest case: massless modes of long Strings in 3D

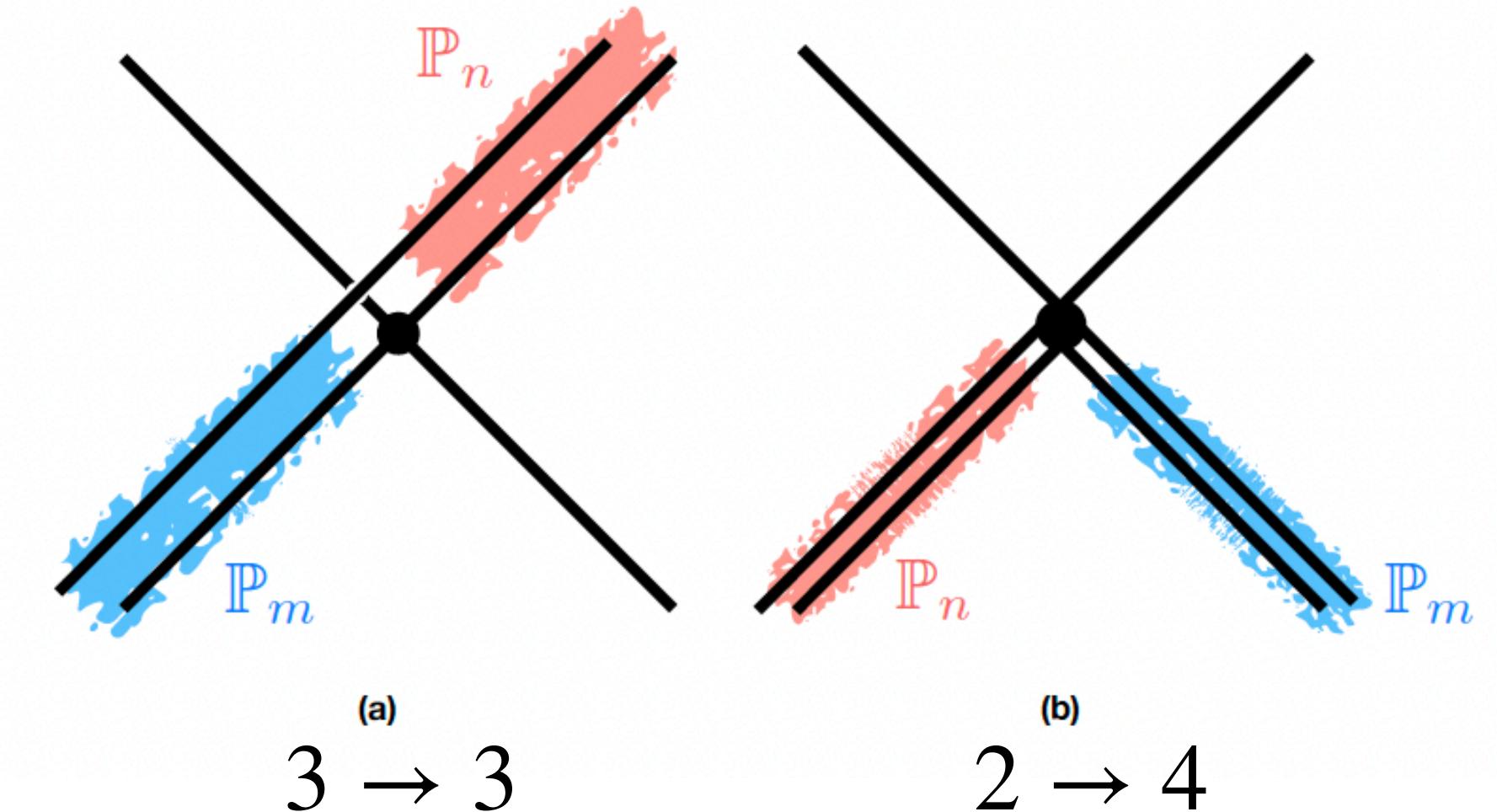
- AG, A. Homrich, J. Penedones and P. Vieira, to appear

Idea: project multi-particle states into jet states

Problem decomposes into a bunch of 2->2 processes

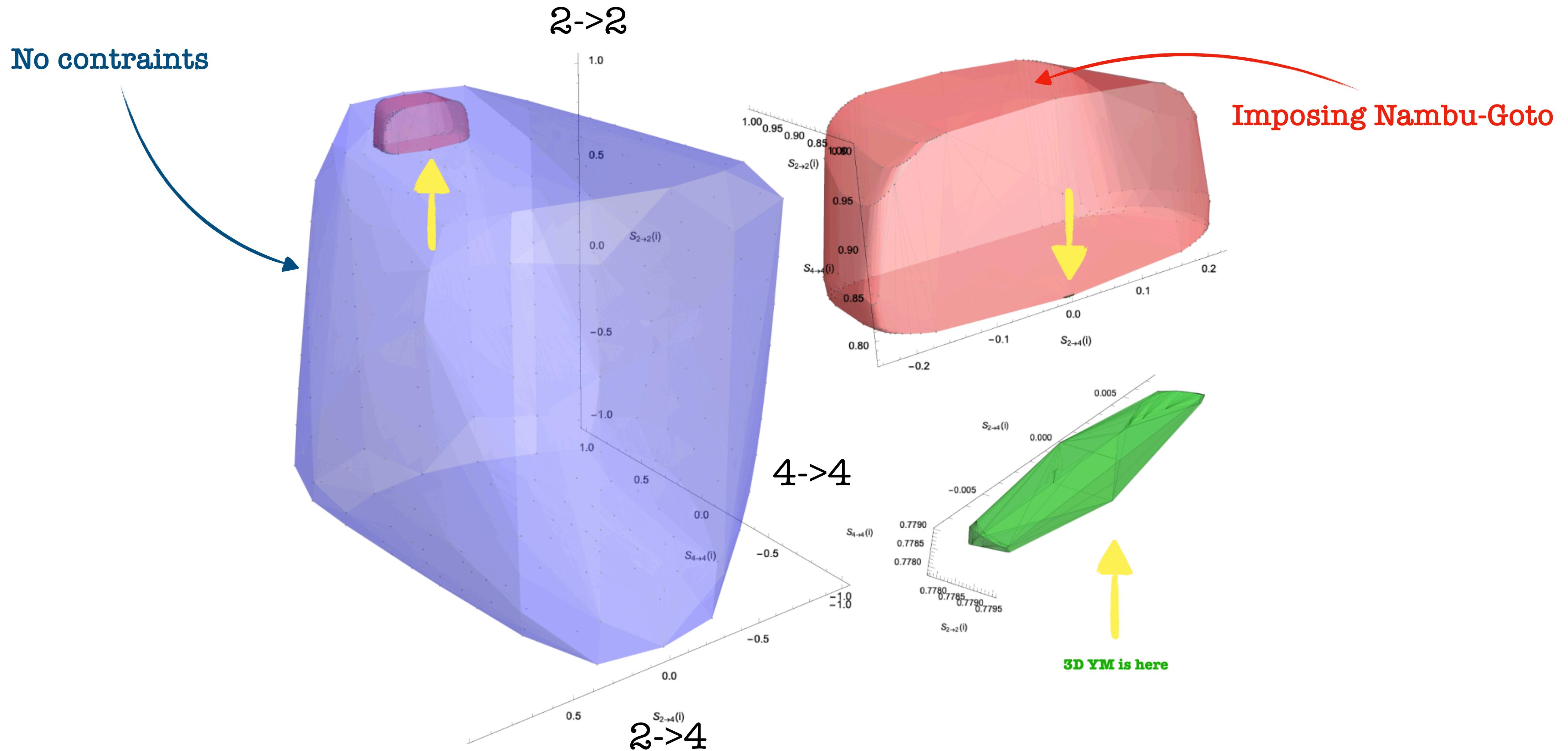
2-particle Jet State

$$|n, P\rangle \equiv \sqrt{2n+1} \int_0^1 d\alpha \frac{P_n(2\alpha - 1)}{\sqrt{8\pi\alpha(1-\alpha)}} |\alpha, (1-\alpha), P\rangle_2$$



$$\begin{aligned} S_{11 \rightarrow 11} &= \text{Diagram } 1, \quad S_{1n \rightarrow 1m} = \text{Diagram } 2, \quad S_{n1 \rightarrow m1} = \text{Diagram } 3, \\ S_{n1 \rightarrow 1m} &= \text{Diagram } 4, \quad S_{1n \rightarrow m1} = \text{Diagram } 5, \quad S_{11 \rightarrow nm} = \text{Diagram } 6, \\ S_{nm \rightarrow 11} &= \text{Diagram } 7 \quad \text{and finally } S_{pn \rightarrow rm} = \text{Diagram } 8. \end{aligned} \quad (6)$$

The Multi-Particle Matrioska coming soon...



What's next?

Q0: Can we construct non-perturbative scattering amplitudes and understand their properties?

Q1: Is it String Theory the unique UV completion of Sugra?

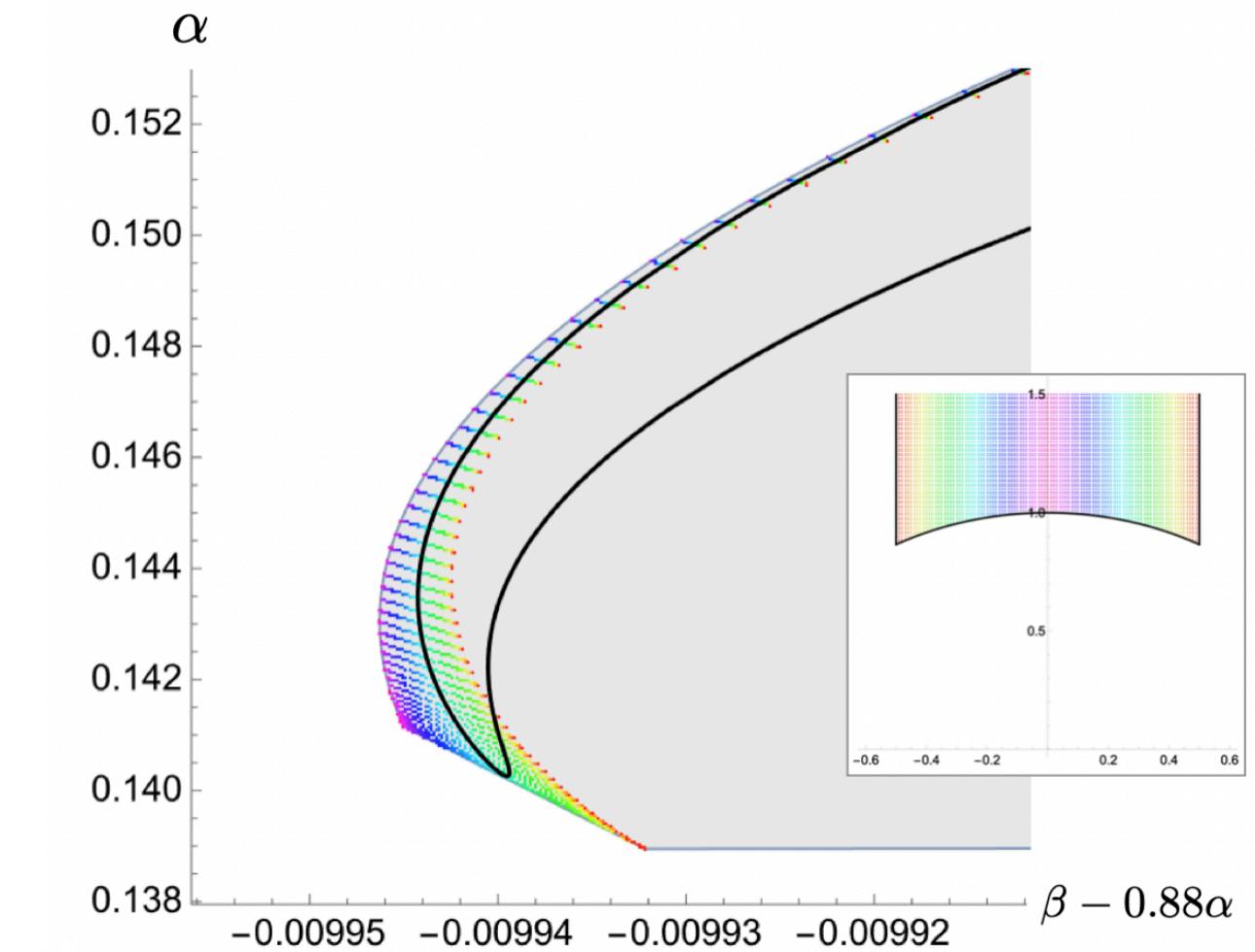
Q1.1: What bounds on D^4R^4 , D^6R^4 , ... operators?
What happens in lower dimensions?

Q1.2: Is the discrepancy compatible with black hole production?

$$\alpha \geq \frac{16}{3\pi^4 \ell_P^{14}} \sum_l (l+1)_6 (2l+7) \int_0^\infty ds \frac{\eta_l(s)}{s^8}$$

The value of α increases as soon as we have inelasticity

Rotating black hole in 10D



We need a model: $Prob_{2 \rightarrow 2} \sim \text{Exp}(-S_{BH}(\text{Area}))$

Q2: How the UV completion of pure Einstein Gravity in $D \geq 5$ looks like?

Q3: Can we perform a non-perturbative gravity Bootstrap in 4D?

Q4: Can we put QCD phenomenology on a more rigorous footing?

$$A^{QG} = 8\pi G_N^d \int d^d x \sqrt{-g} (R + a_2 R^2 + a_4 R^3 + \alpha R^4 + \dots)$$

After dinner suggestion

Cachaça Degustation at **Chachaçaria.bar**



See you there, or look for me or these gentlemen after dinner



This is a great caipirinha

Backup Slides

An analytic bound on scattering

Goal: we bound $c_4 \iff$ we bound Δ_3

What are the non-perturbative properties of the branons scattering amplitude?

Unitarity: define $S(s) = 1 + \frac{i}{2s} T_{2 \rightarrow 2}(s)$ then $|S(s)|^2 \leq 1$ for $s > 0$

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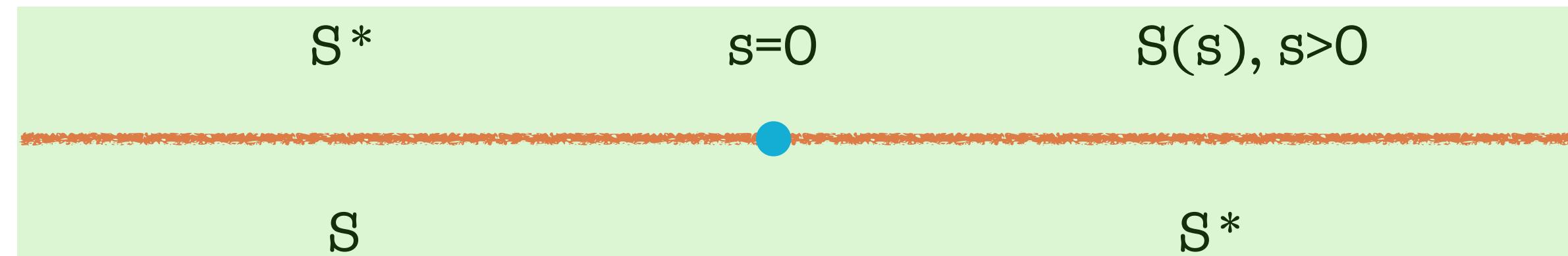
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Analytic away from the real axis

Analyticity



Low Energy Constraints: $S(s) = 1 + i\frac{s}{4} - \frac{s^2}{32} + i(\gamma_3 - \frac{1}{384})s^3 + \dots$

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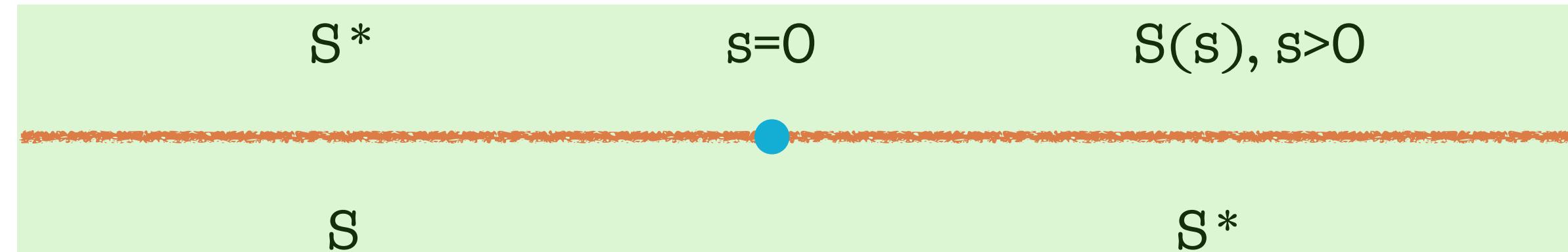
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Caristo, Caselle, Magnoli, Nada, Panero '21

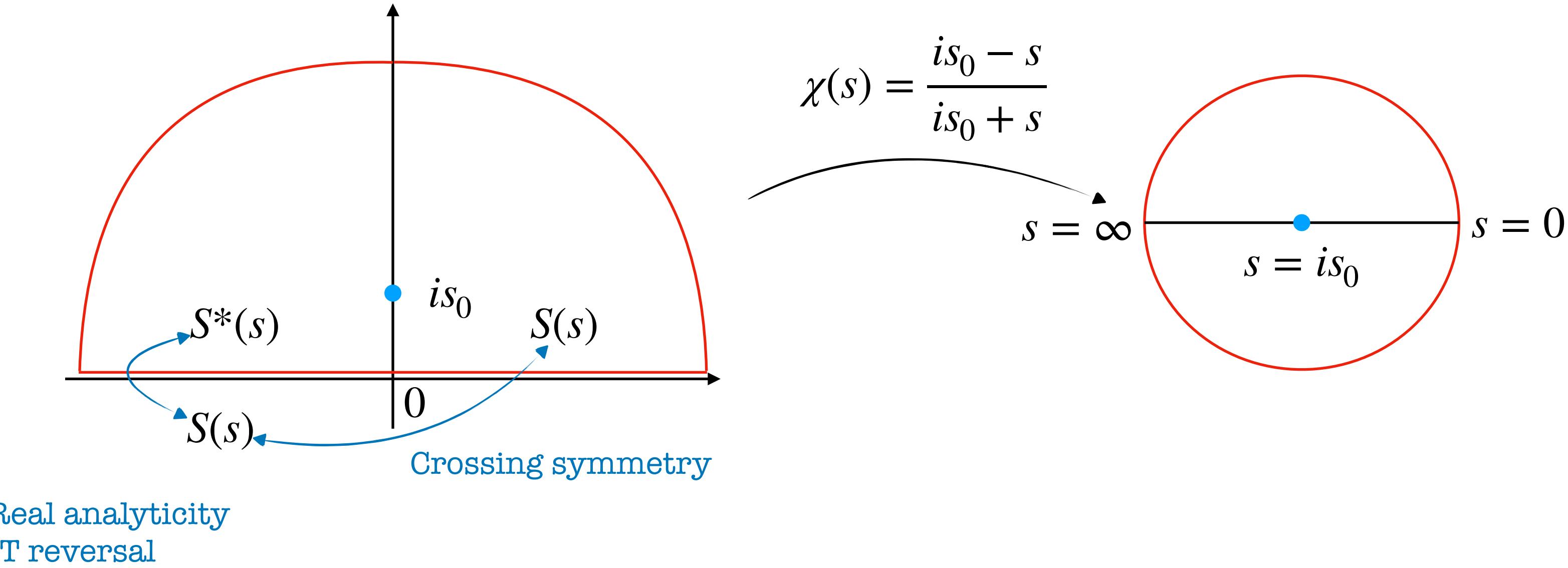
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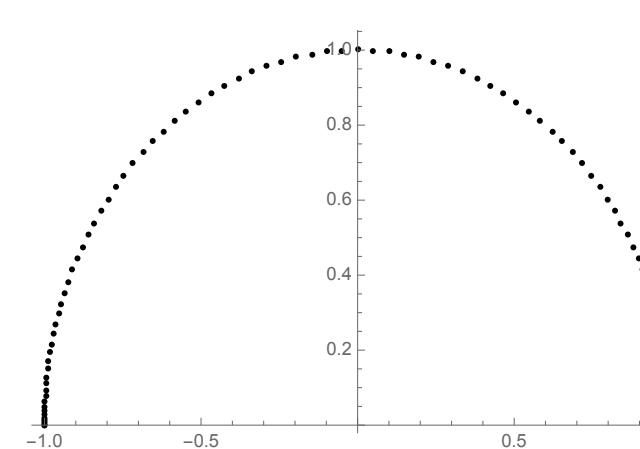
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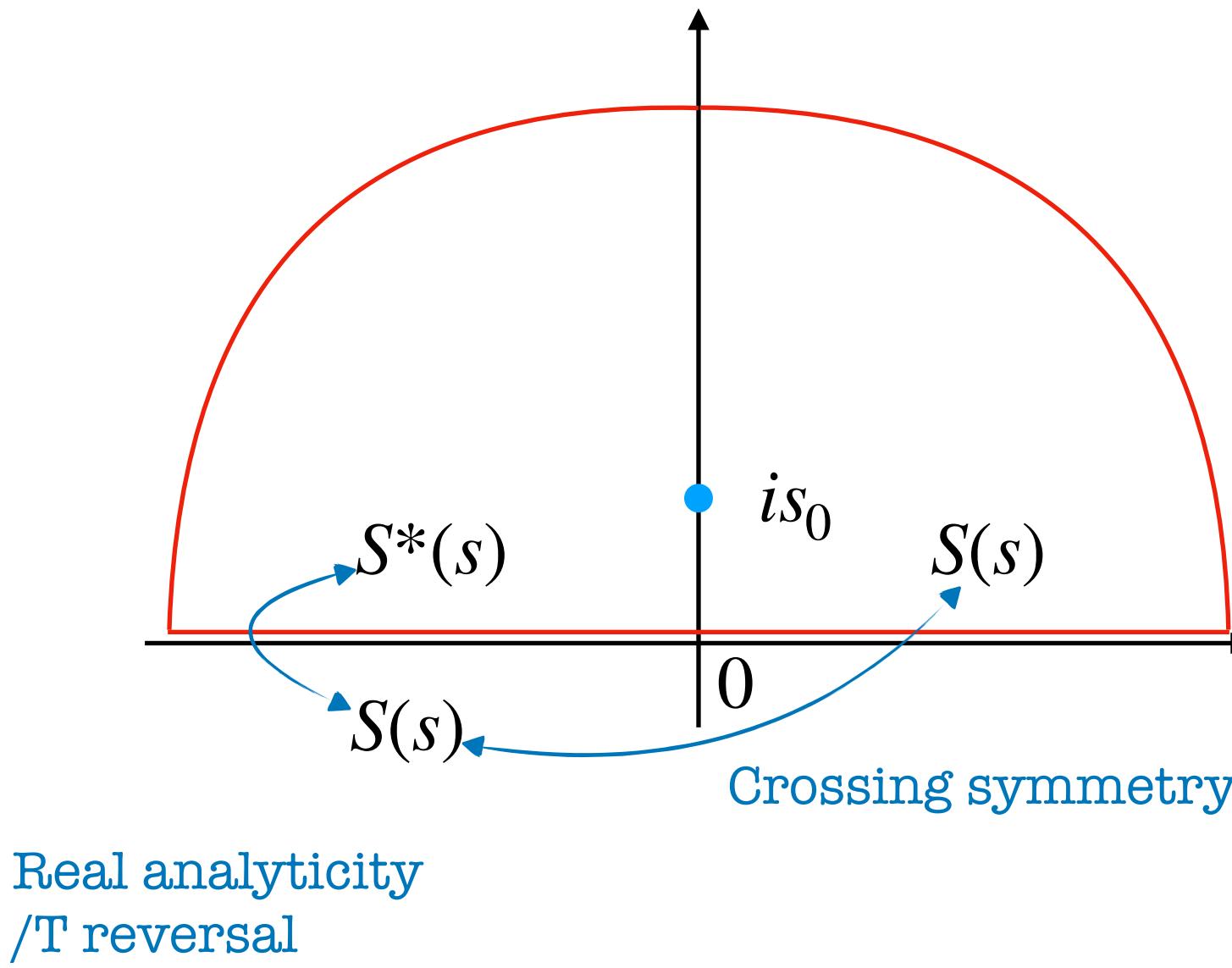
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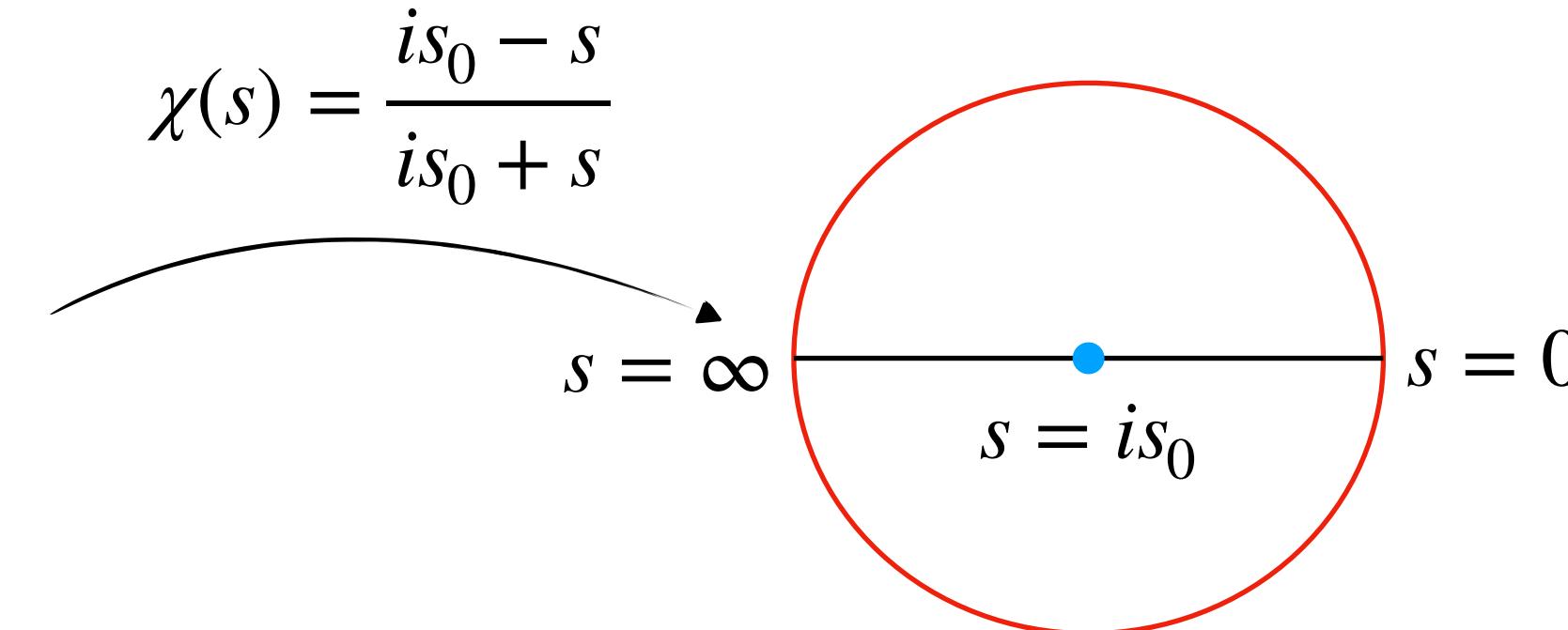
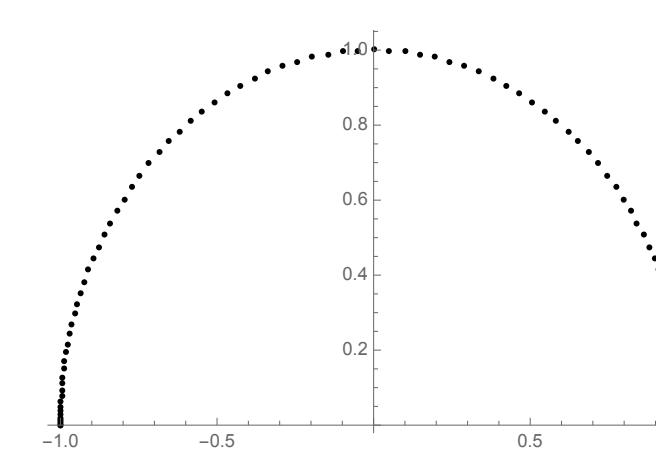
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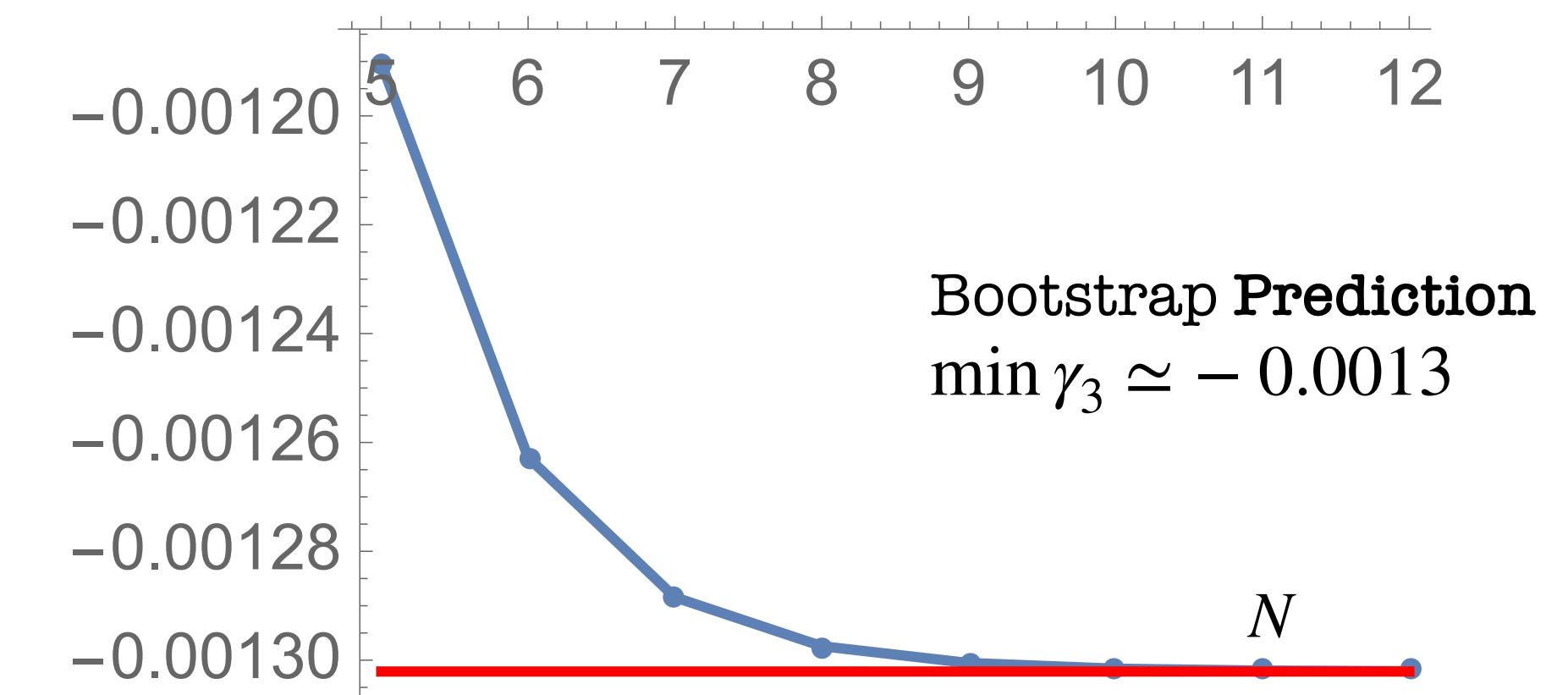
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Naively: Stronger constraints!

$$\sum_n P_{2 \rightarrow n} = 1 \implies P_{2 \rightarrow 2} + P_{2 \rightarrow 4} + \dots \leq 1$$

String Theory and M-theory Expectations

α_D is 1-loop exact up to non-perturbative corrections

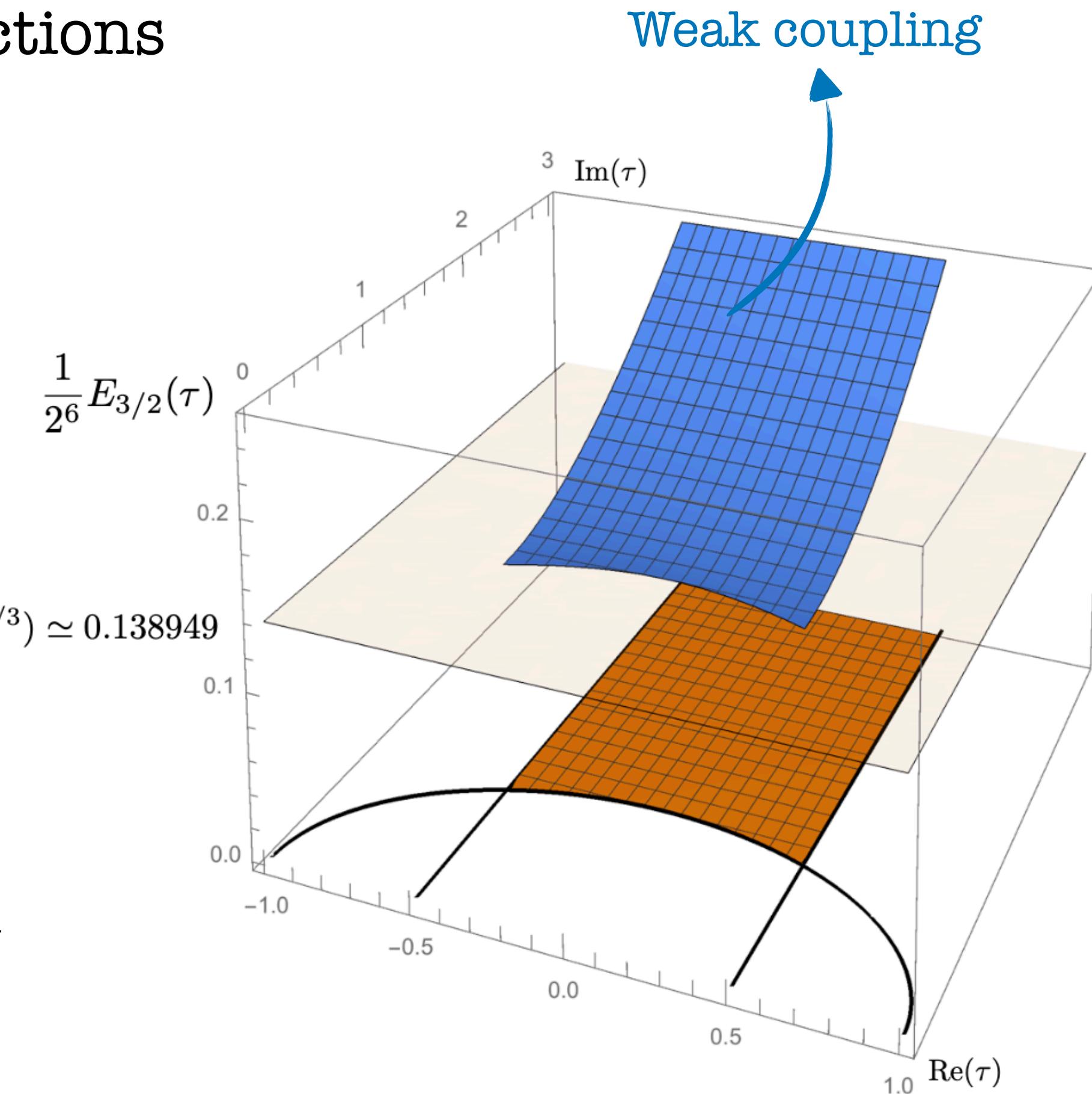
Min α_{10} realized in type IIB

$$\text{D=10 Type IIB: } \alpha_{10}^{IIB} = \frac{1}{2^6} E_{3/2}(\tau, \bar{\tau}) \geq 0.139\dots$$

$$\text{D=9: } \alpha_9(\tau, \nu) = \frac{1}{2^6} \left(\nu^{-3/7} E_{3/2}(\tau, \bar{\tau}) + \frac{2\pi^2}{3} \nu^{4/7} \right) \geq 0.2417\dots$$

$$\nu = \left(\frac{r}{\ell_s} \right)^{7/4} \sqrt{g_9} = \left(\frac{\ell_P}{\tilde{r}} \right)^{7/4}$$

$$\text{D=11: } \alpha_{11} = \frac{(2\pi)^2}{3 \times 2^7} = 0.1028\dots$$



Green, Gutperle
Green, Vanhove
Green, Russo, Vanhove

hep-th/9701093
hep-th/9704145
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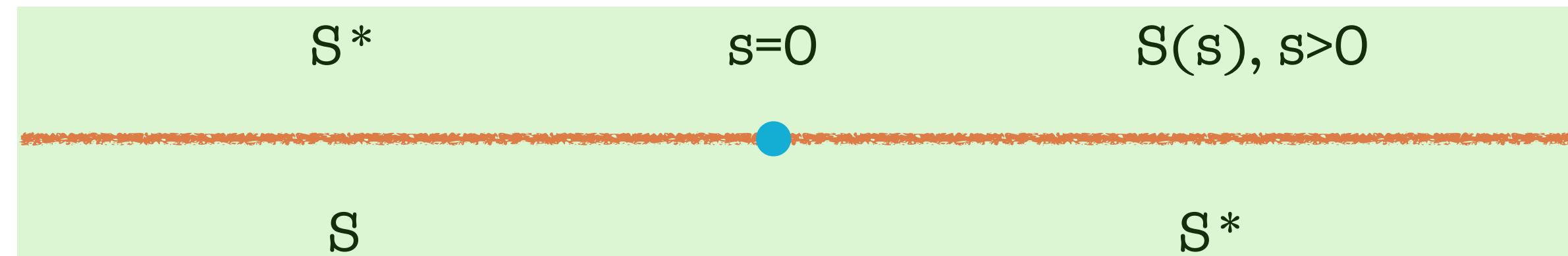
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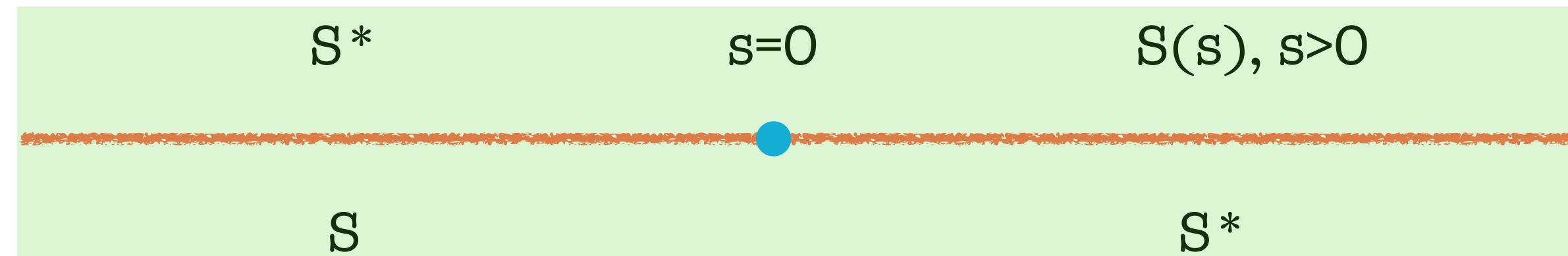
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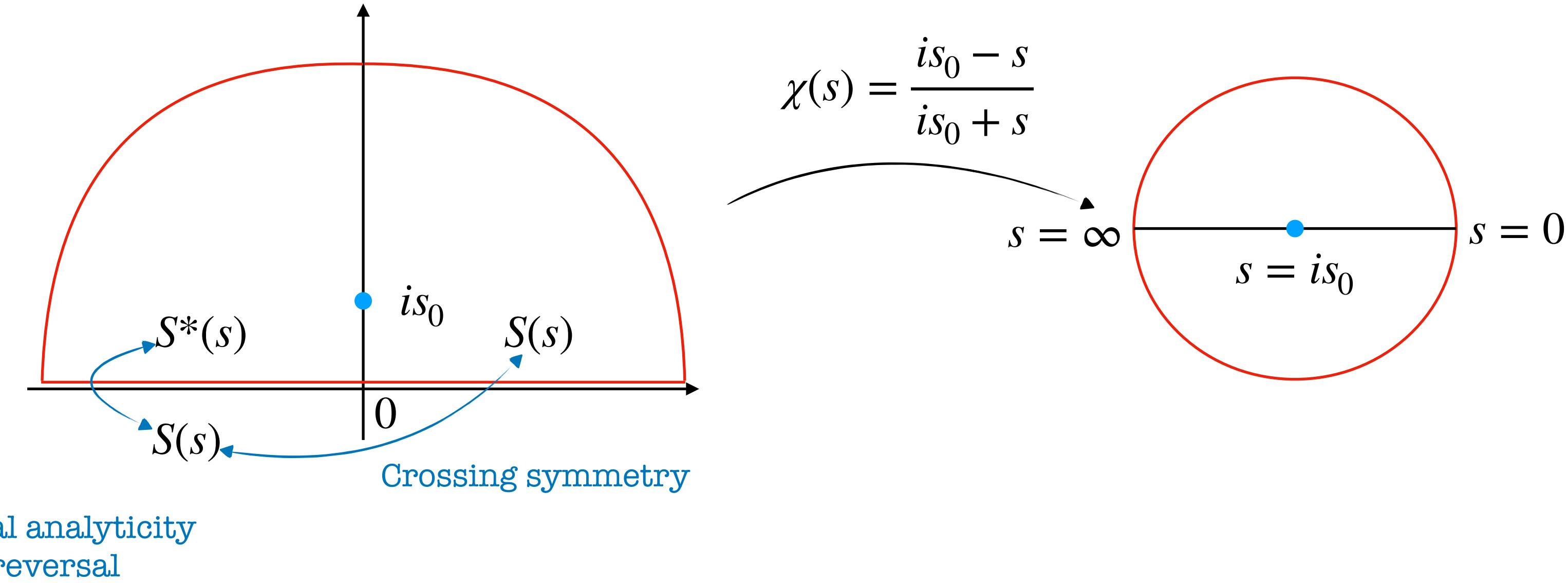
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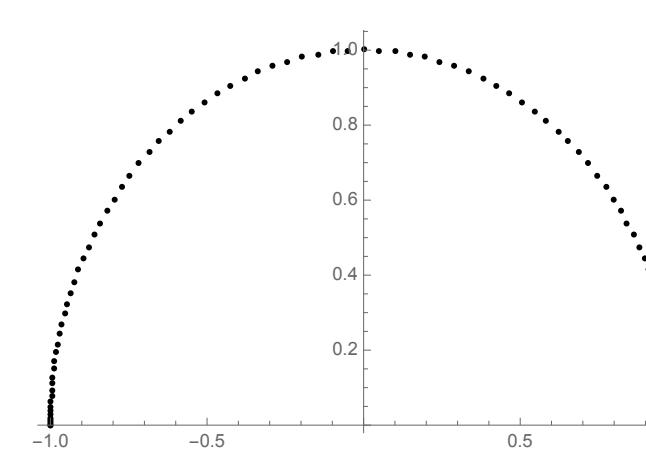
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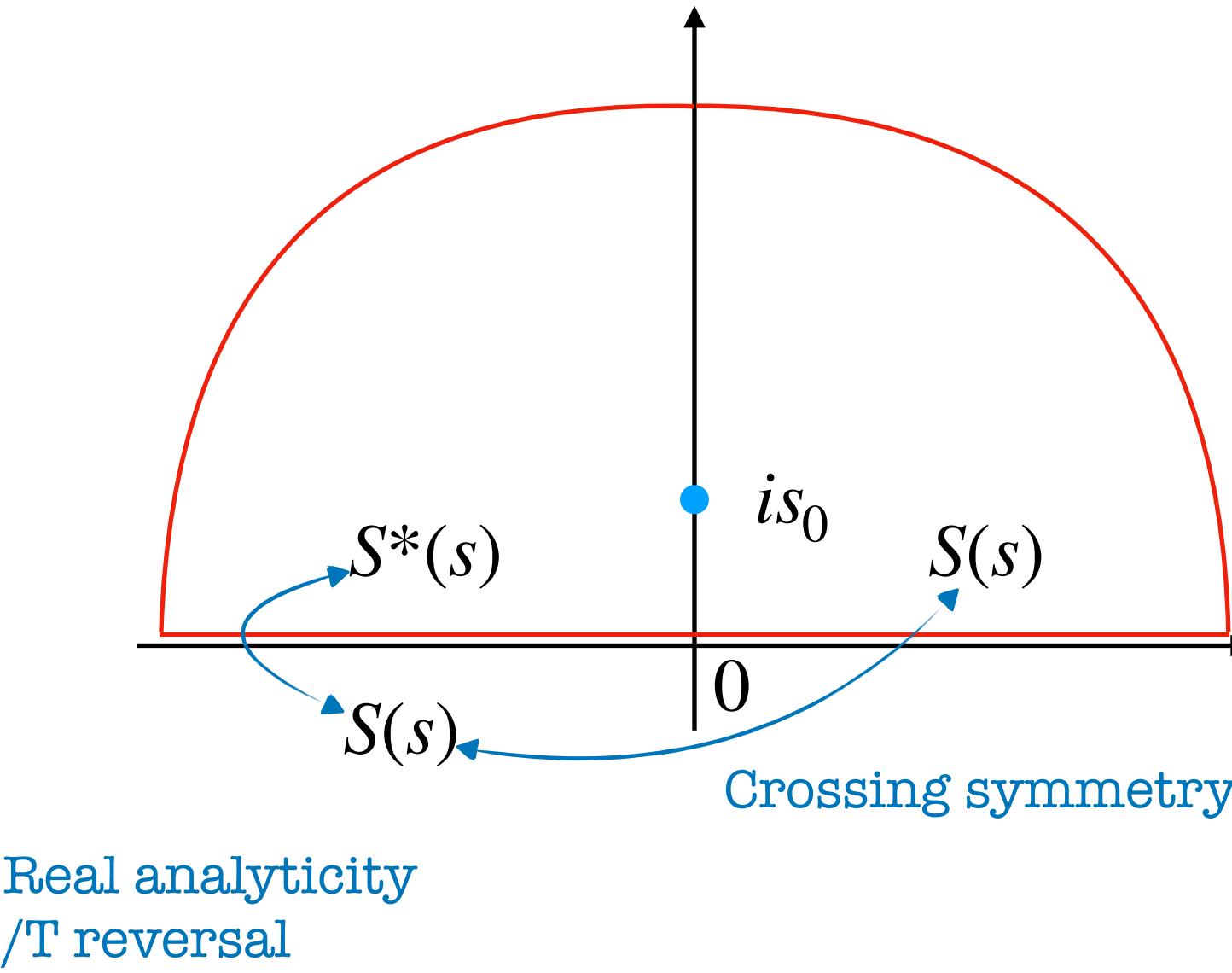
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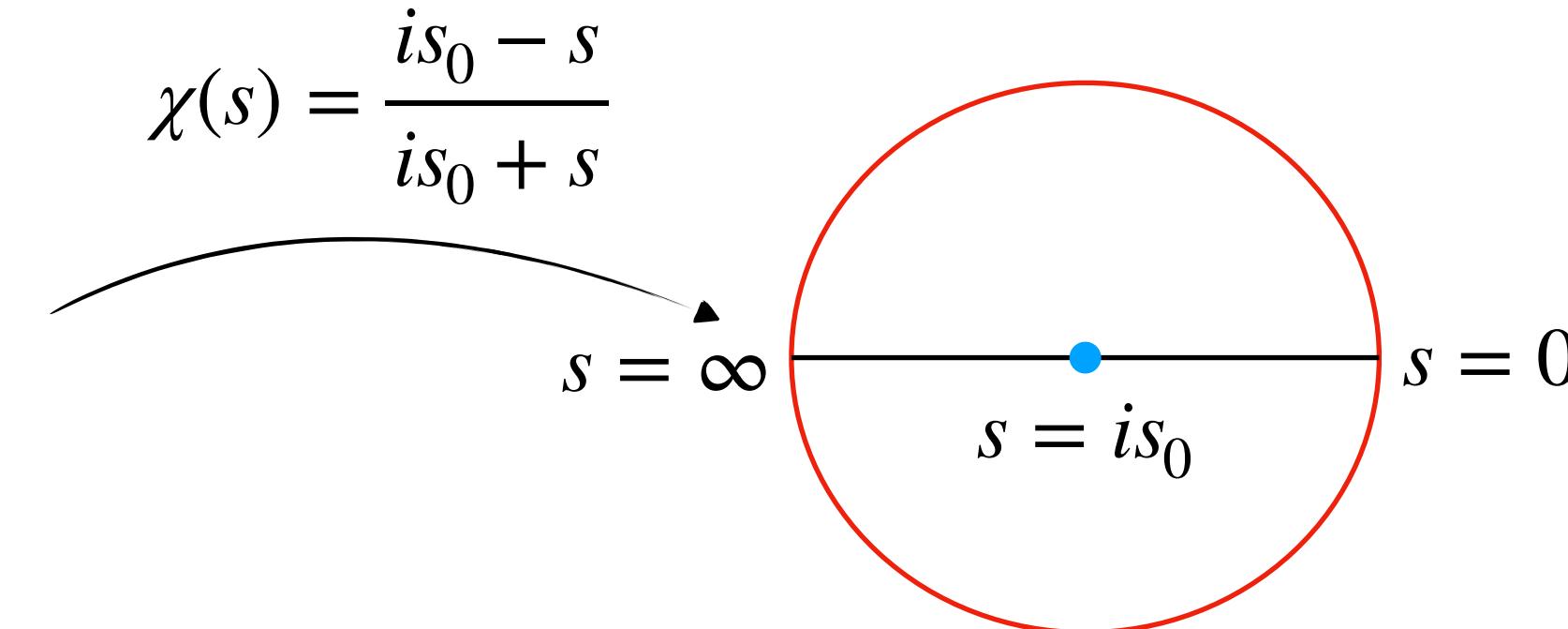
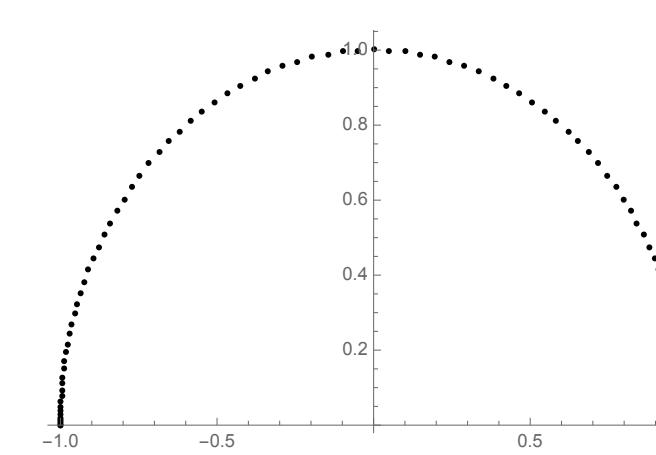
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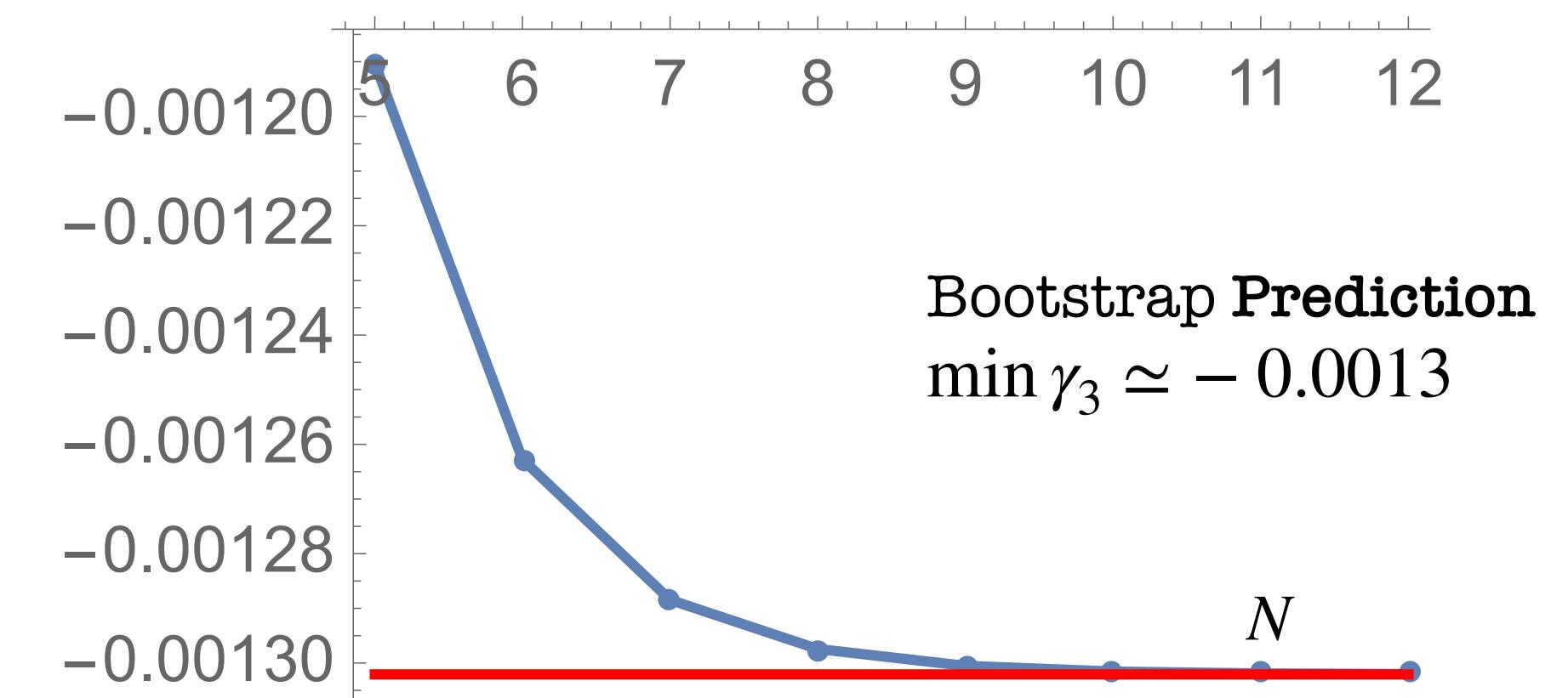
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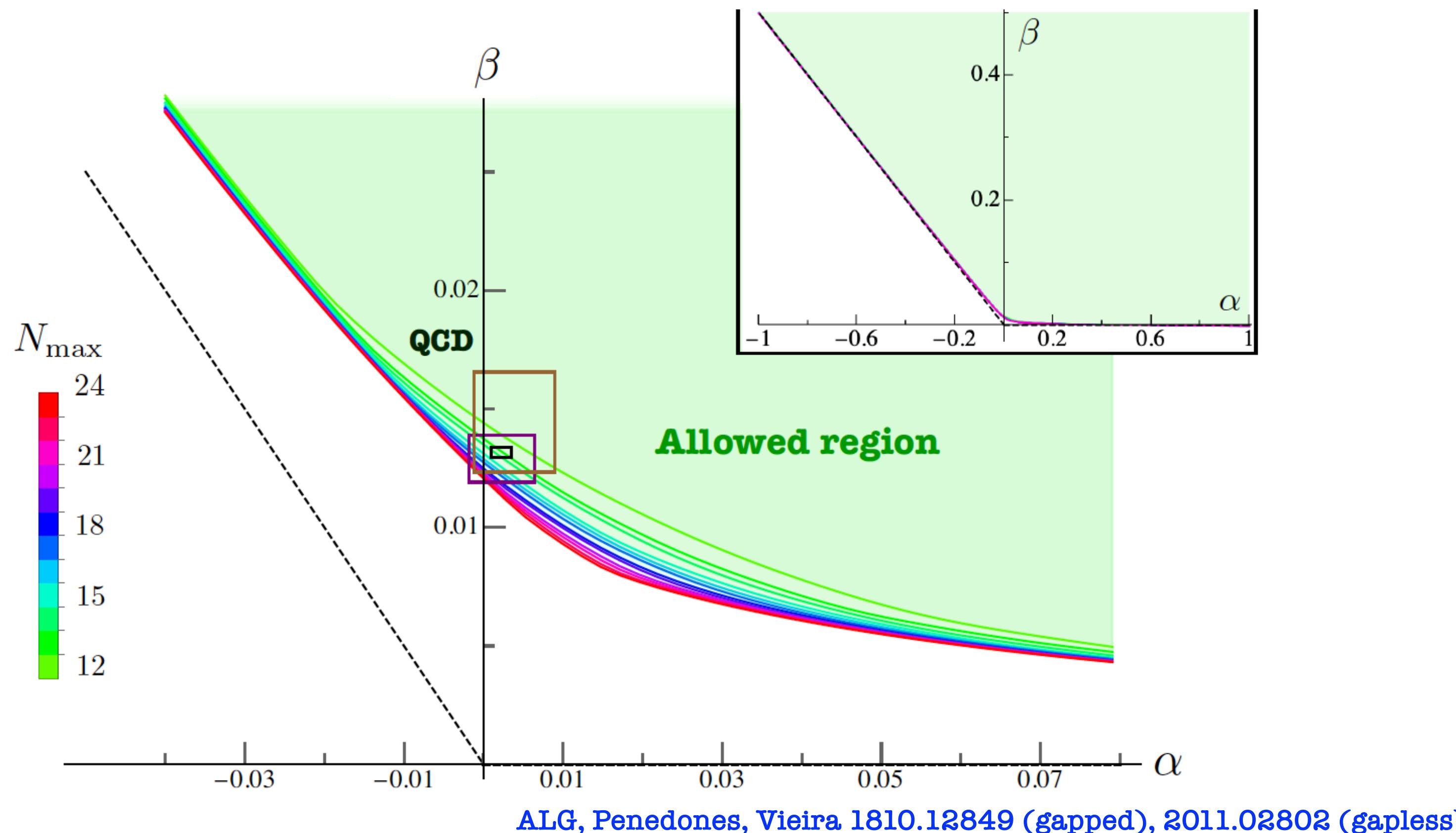
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Low energy QCD

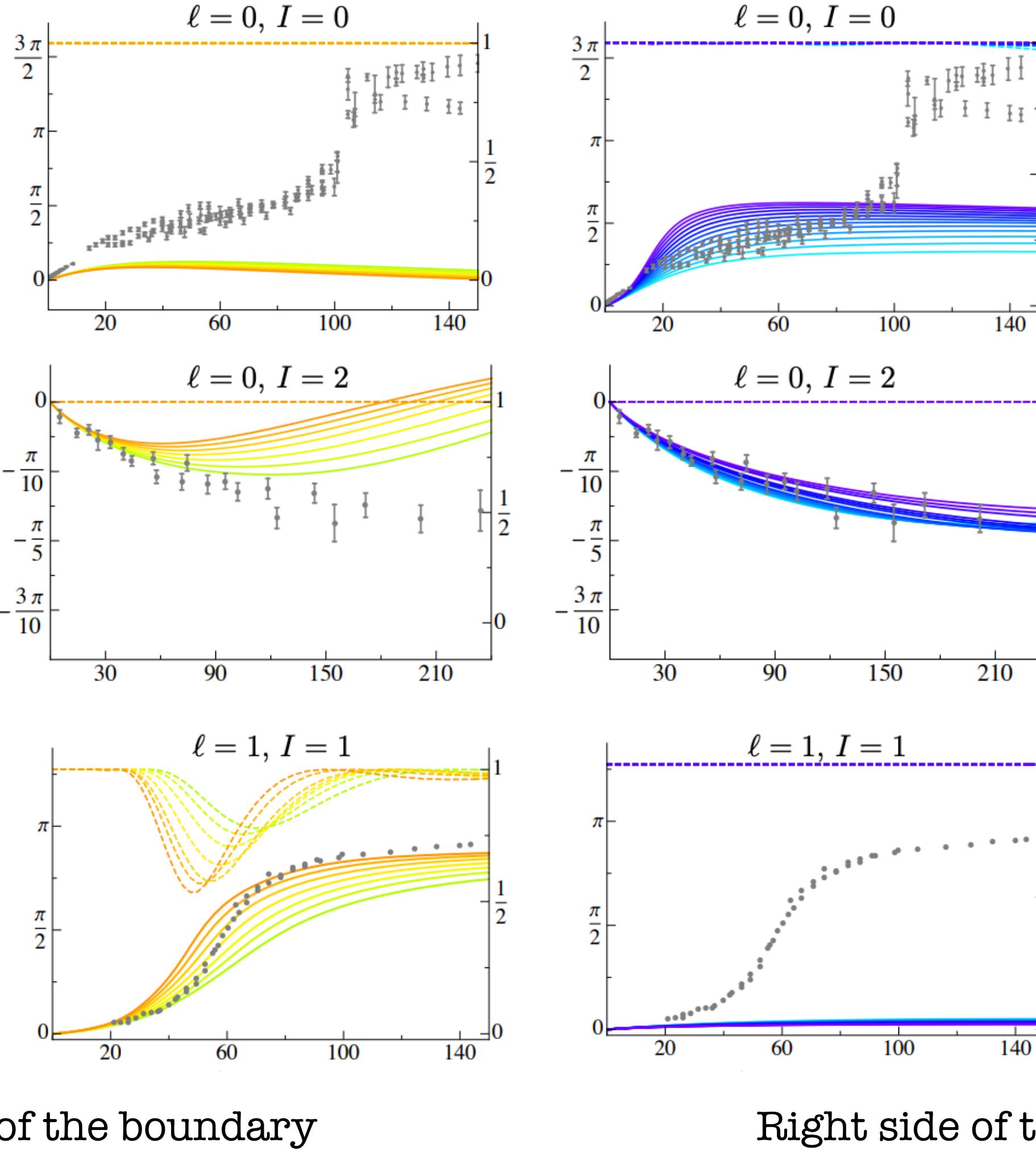
In QCD dynamical mass generation, non-perturbative RG flow

$$\text{Amplitude} = \frac{s}{f^2} + \alpha \frac{s^2}{f^4} + \beta \frac{t^2 + u^2}{f^4} + \log s + \text{UV completion}$$



α, β can be only computed using lattice QCD today or extracted from data!!!

Non perturbative S-matrices from Bootstrap



Left side of the boundary

Right side of the boundary

What can we add to nail down QCD?

Work in progress with H. Murali

D=4 Strings and the Axion

In D=4 two leading deviations from Nambu-Goto α_3, β_3

$$\mathcal{A}_{EFF} = \int d^2\sigma \sqrt{-h} \left(\frac{1}{\ell_s^2} + \dots + \color{red}\alpha_3\color{black} \ell_s^6 K^4 + \color{red}\beta_3\color{black} \ell_s^6 R^2 + \dots \right)$$
$$\gamma_3 = \alpha_3 - \beta_3$$

New Effect in the amplitude: Polchinski-Strominger term $\propto \alpha_2 = \frac{D-26}{384\pi}$

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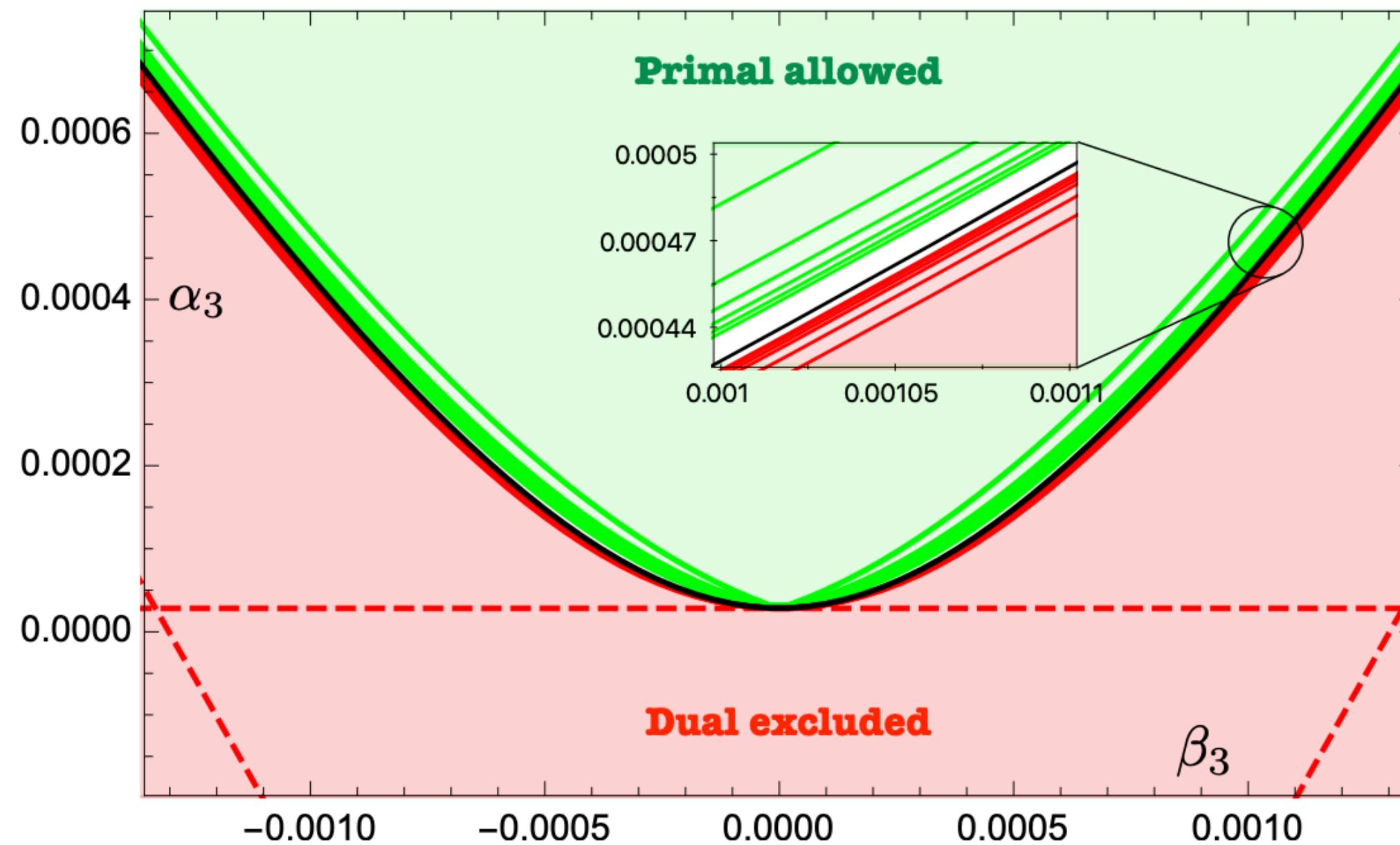
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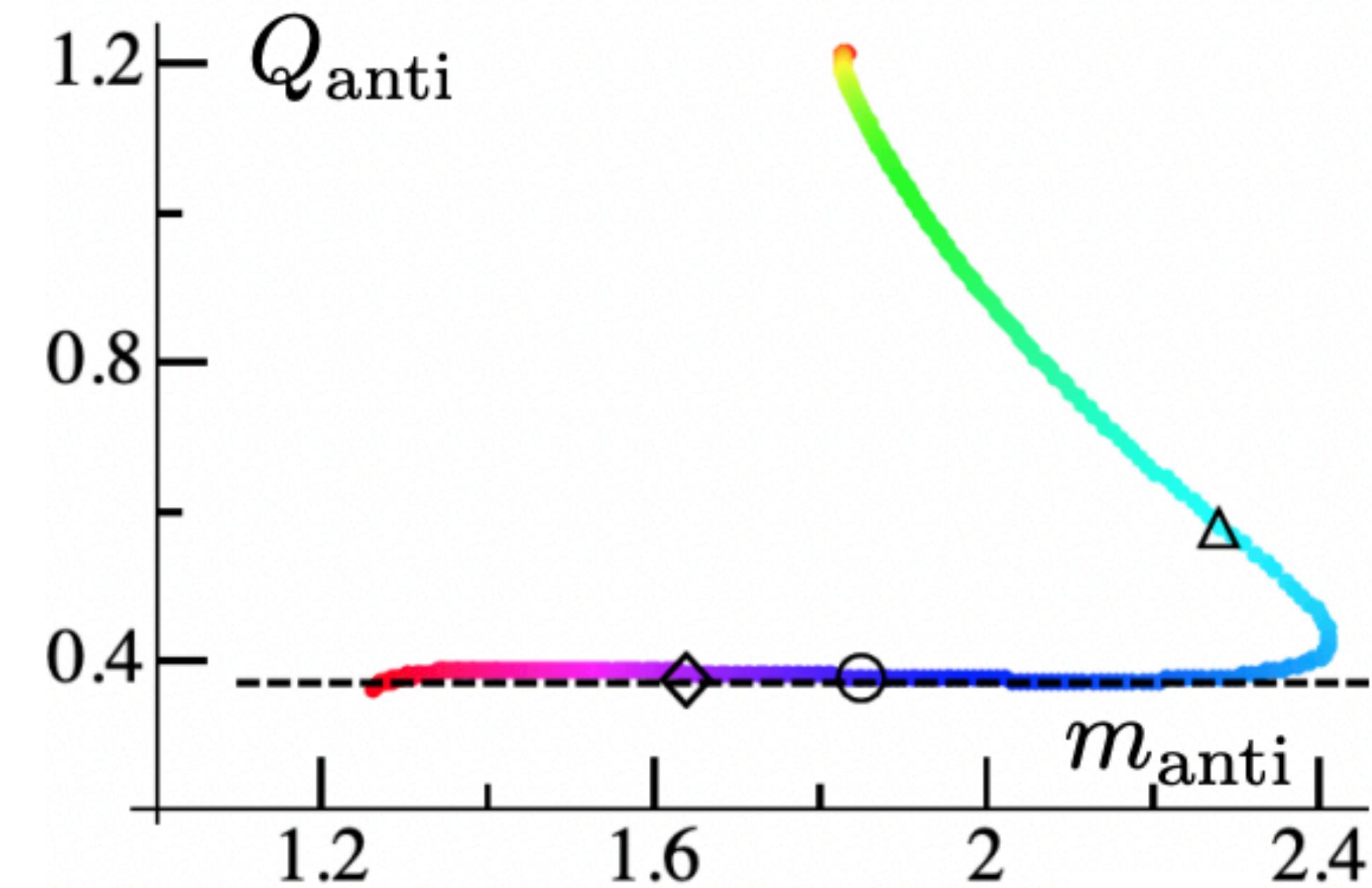
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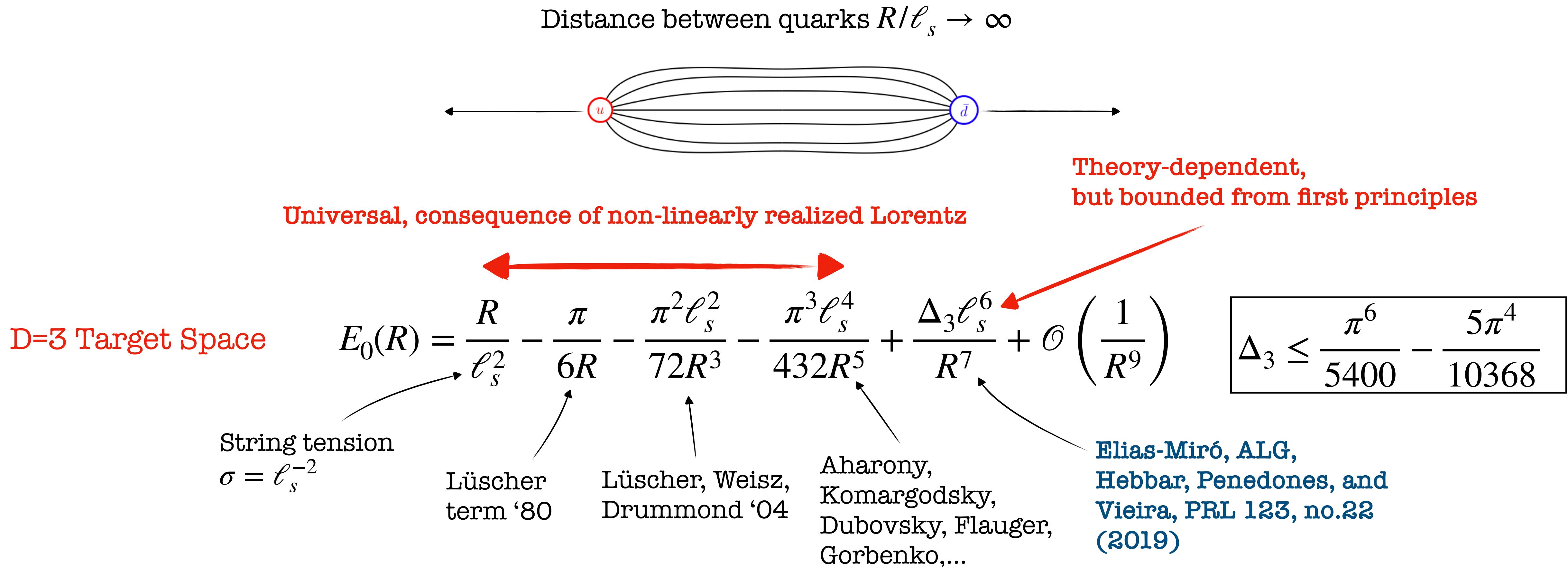


The extremal amplitude has an axion resonance

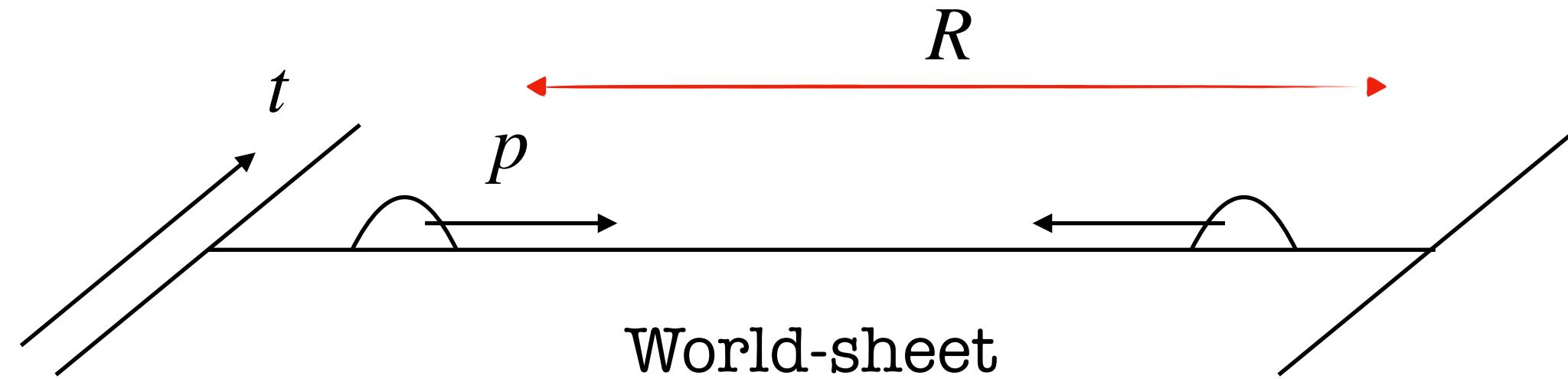


Bounds on the static $q\bar{q}$ potential (toy for quantum gravity)

Application of the S-matrix Bootstrap to fundamental questions about confinement



Effective String Theory



Physical Degrees of freedom: X^i with $i=1,\dots,D-2$, massless Goldstones

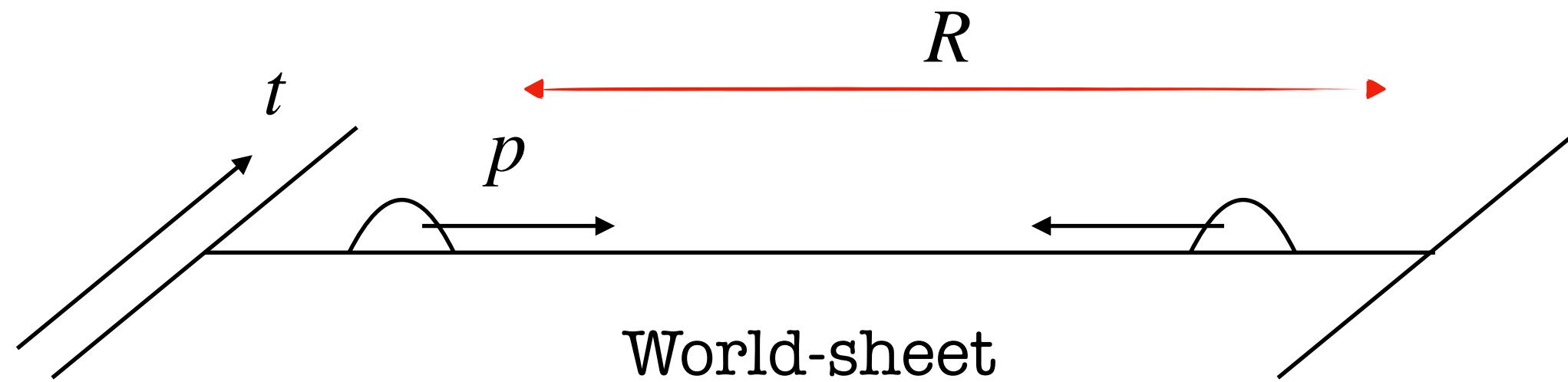
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We have an action, we can compute the S-matrix, but is this useful?

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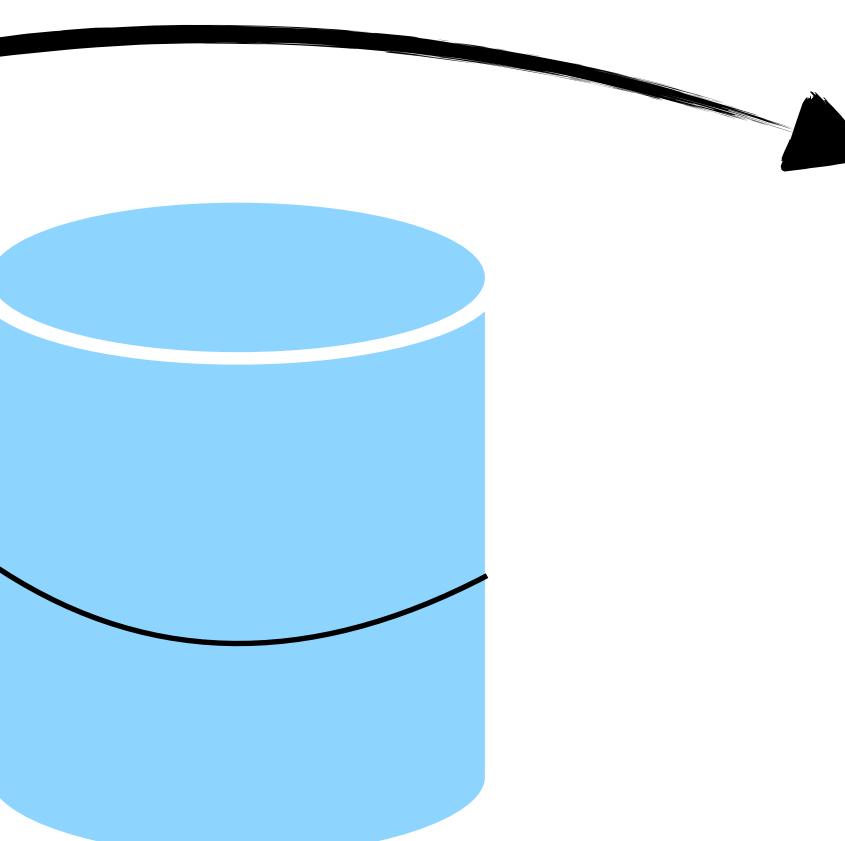
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Thermodynamic Bethe Ansatz

$$T_{2 \rightarrow 2}(s) = \ell_s^2 \frac{s^2}{2} + i \ell_s^4 \frac{s^3}{16} + \ell_s^6 \left(2\gamma_3 - \frac{1}{192} \right) s^4 + \dots$$



Finite Volume Energy Levels from Infinite Space Scattering

$$E_0(R) = R + \frac{1}{\pi R} \int_0^\infty dq \log(1 - e^{-\epsilon(q)})$$

$$\epsilon(p) = pR + \frac{1}{2\pi} \int_0^\infty dq \frac{\partial}{\partial q} \delta(4pq) \log(1 - e^{-\epsilon(q)})$$

$$\Delta_3 = -\frac{32\gamma_3\pi^6}{225} - \frac{5\pi^4}{10368}$$

We solve them in the $1/R$ expansion

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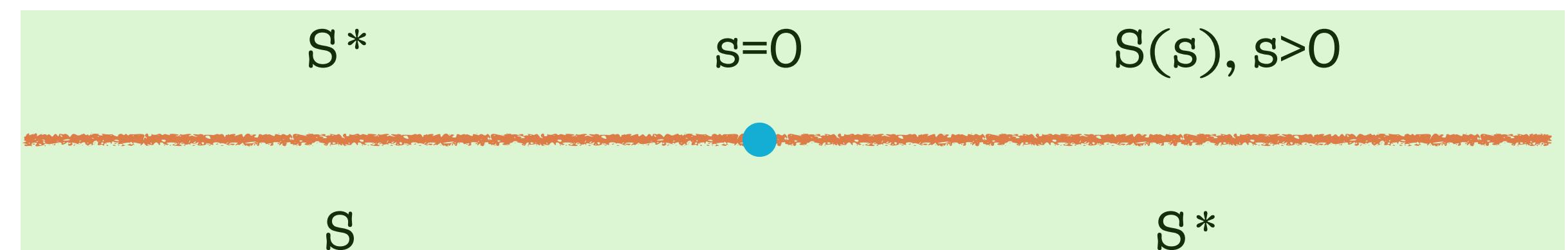
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gauge group	\mathbb{Z}_2	$SU(2)$	$SU(6)$	$SU(\infty)$
$\gamma_3 \times 768$	-0.4	[4]	-0.3	[5]

[4] Baffigo, Caselle '23

[5] Caristo, Caselle, Magnoli, Nada, Panero '21

[1,6] Dubovsky, Gorbenko, et al

The S-matrix Bootstrap represent a novel framework connecting different fields in physics and mathematics

Optimization theory

Development of new algorithms and strategies

Analytic solutions and Geometric Function Theory

1+1 QFTs and Integrability

Integrable Models as application and testing ground for the Bootstrap

Approximate integrability to study the QCD String

String and M-theory

Non-perturbative Dualities

Nonperturbative properties of String scattering Amplitudes

Perturbative regimes

Large Spin expansion, and Spin analyticity

Loop expansions and match with QFT

S-matrix Bootstrap



CFT Bootstrap

Celestial Holography

Matrix Models Bootstrap

Q: Can we get rid of Lagrangians and reformulate QFTs and Strings using more general physical principles?

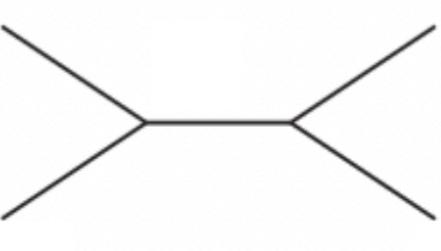
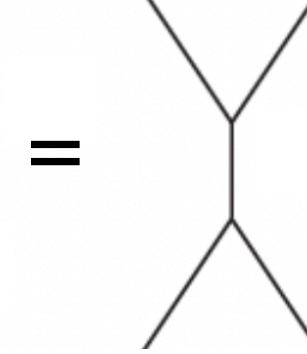
Bootstrap as an Optimization Problem

Goal: Find the optimal value of a physical observables constrained by the laws of nature

Math + Physics

Causality, crossing, and global symmetries

Linear Constraints: $M_{2 \rightarrow 2}(s | t) = c + \int_{4m^2}^{\infty} \text{Disc}_z M_{2 \rightarrow 2}(z | t) K(z, s | t) dz$

 $=$ 
$$M_{2 \rightarrow 2}(s | t) = M_{2 \rightarrow 2}(s | 4m^2 - s - t)$$

Quantum mechanics

Unitarity Constraints: $\mathbb{I} - \mathbb{S}\mathbb{S}^\dagger \succeq 0$

$\mathbb{S} \supset M_{n \rightarrow m}$ all possible processes

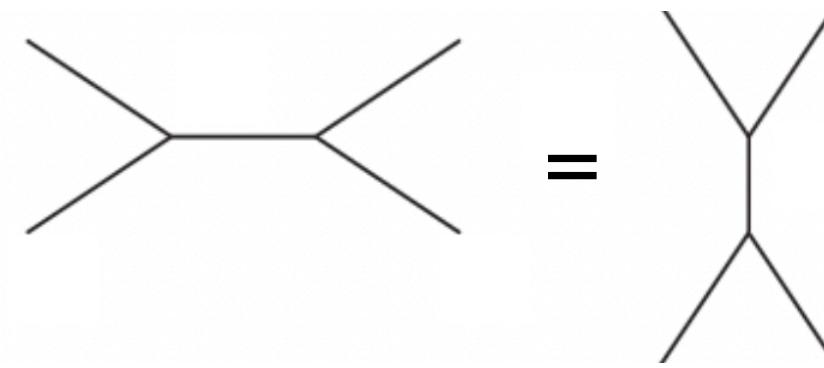
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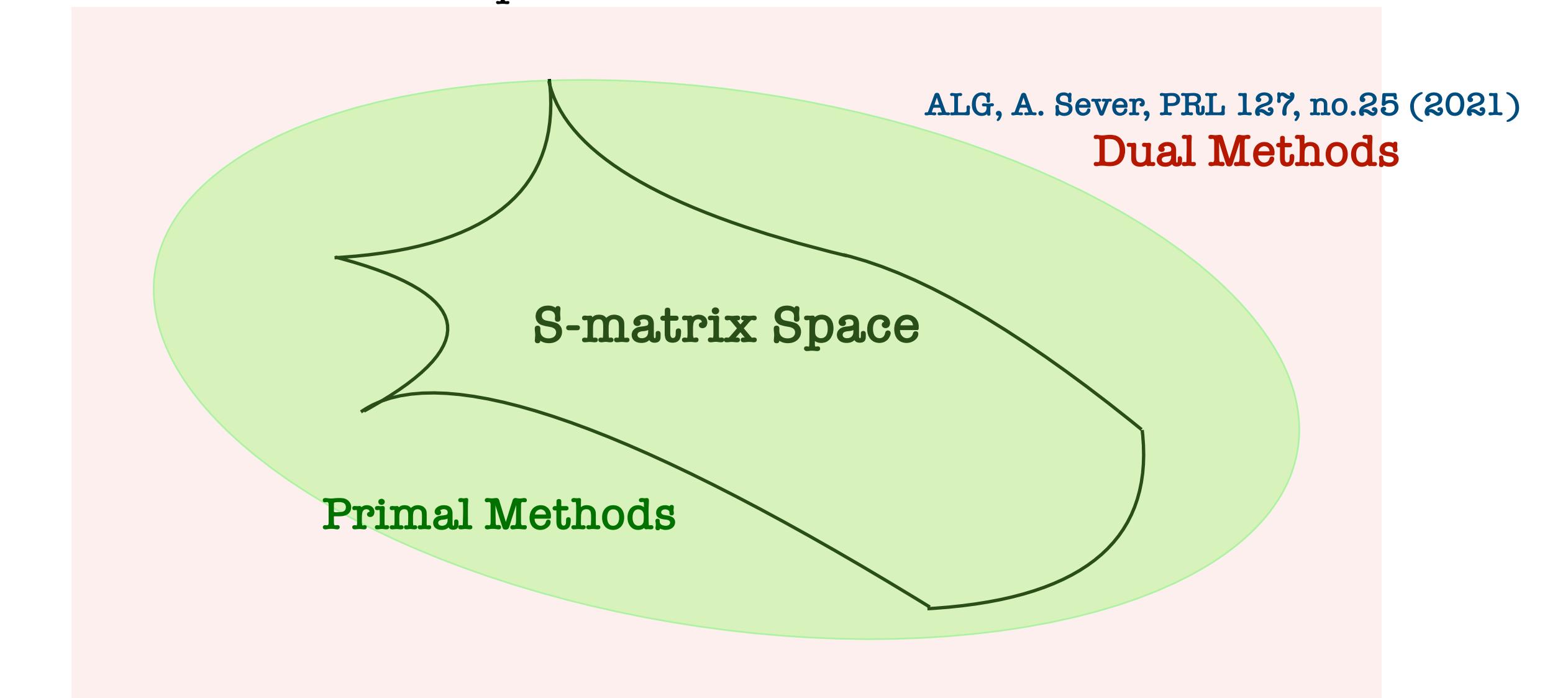
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Computer Science

With Standard Optimization Tools (SDPB), we only explored the Convex Hull



Use AI and Deep Learning Algorithms

Unsupervised Reinforcement Learning to explore
the S-matrix Space

We relax Unitarity

$$\sum_n P_{2 \rightarrow n} = 1 \implies P_{2 \rightarrow 2} \leq 1$$

Standard Search Algorithms

Combine Gradient Methods with SDPB
(software created by the Bootstrap Collaboration)

Why we need for a theoretical collider



Why we need for a theoretical collider



Theoretical physicists are impatient!

The New Bootstrap Manifesto:
find the space of all possible
physical observables assuming
general principles

$$\mathcal{S}_{in \rightarrow out} \equiv \langle in | out \rangle$$

