# Soft OCD Review - with some emphasis on string models, LHC, and PYTHIA 

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## Soft QCD — Theory Models

A

## Regge Theory


(a)

(b)

Optical Theorem

+ Eikonal multi-Pomeron exchanges

$$
\sigma_{\text {tot,inel }} \propto S^{\varepsilon} \text { or } \log ^{2}(s)
$$

Cut Pomerons $\rightarrow$ Flux Tubes (strings)
Uncut Pomerons $\rightarrow$ Elastic (\& eikonalization)
Cuts unify treatment of all soft processes
$E L, S D, D D, \ldots, N D$
Perturbative contributions added above $\mathrm{Q}_{0}$

B

## pOCD-Based



+ Unitarity \& IR Regularisation
$\rightarrow$ Multi-parton interactions (MPI)
+ Parton Showers \& Hadronization
Regulate $\mathrm{d} \sigma$ at low рто $\sim$ few GeV
Screening/Saturation $\rightarrow \sqrt{s}$-dependent $\mathrm{p}_{\text {то }}$
Total cross sections from Regge Theory
(Donnachie-Landshoff + Parametrizations)


## The Structure of an LHC pp Collision



## Linear Confinement

On lattice, compute potential energy of a colour-singlet $q \bar{q}$ state, as function of the distance, $R$, between the $q$ and $\bar{q}$ :


Linear Term $=$ Model as strings (Lund Model)

## String Fragmentation in One Slide

## The string model provides a mapping:

Quarks > String endpoints
Gluons > Kinks on strings
Further evolution then governed by string world sheet (area law)

+ string breaks by tunnelling
By analogy with "Schwinger mechanism" in QED (electronpositron pair production in strong electric field)
> (Jets of) Hadrons!


String breaks by quark pair production
$\Longrightarrow$ strangeness suppression

$$
\propto \frac{\exp \left(\frac{-\pi m_{s}^{2}}{\kappa}\right)}{\exp \left(\frac{-\pi m_{u, d}^{2}}{\kappa}\right)}
$$

## Alternative: The Cluster Model — Used in HERWIG \& SHERPA

## In "unquenched" QCD

$g \rightarrow q \bar{q} \Longrightarrow$ The strings will "break"
Non-perturbative so can't use $P_{g \rightarrow q \bar{q}}(z)$
Alternative: force $g \rightarrow q \bar{q}$ at end of shower


## New Directions in String Fragmentation

Regard tension $\kappa$ as an emergent quantity

Cyclonic and Anticyclonic Winds


May depend on spatial coordinate $\sigma$ ?
Work in progress with E. Carragher \& J. March-Russell (Oxford).
May depend on environment? (e.g., other strings nearby)
Two approaches (so far) within Lund string-model context:
Colour Ropes [Bierlich et al. 2015] + several more recent
Close-Packing [Fischer \& Sjöstrand 2017] + Work in progress with L. Bernardinis \& V. Zaccolo (Trieste)

## The Environment — in Hadronic Collisions

In hadronic collisions, we are not hadronizing a simple $q-g-\ldots-g-\bar{q}$ string
Coloured initial states + gluon exchanges
$\Longrightarrow$ more complicated colour flows
Also: Protons are composite
One proton = beam of partons

+ QCD $2 \rightarrow 2$ scattering diverges at low PT
$\Longrightarrow \sigma_{\text {parton-parton }}\left(\hat{p}_{\perp}\right)>\sigma_{\text {proton-proton }}$
Interpretation: $\frac{\sigma_{\text {parton-parton }}\left(\hat{p}_{\perp}\right)}{\sigma_{\text {hadron-hadron }}} \sim\langle n\rangle_{\text {parton-parton }}\left(\hat{p}_{\perp}\right)$
(Regulated at low $\hat{p}_{\perp}$ by IR cutoff $\sim$ colour screening)
Multiple Parton-Parton Interactions (MPI)
$\rightarrow$ Additional colour exchanges



## MPI \& Confinement

## MPI / cut pomerons $\Rightarrow$ lots of coloured partons scattered into final state

With significant overlaps in phase space Who gets confined with whom?

## Each has a colour ambiguity $\sim 1 / N_{C}^{2} \sim 10 \%$

E.g.: random triplet charge has $1 / 9$ chance to be in singlet state with random antitriplet:

$$
\begin{aligned}
& 3 \otimes \overline{3}=8 \oplus 1 \\
& 3 \otimes 8=15+6+3, \text { etc. }
\end{aligned}
$$

Many charges $\rightarrow$ Colour Reconnections* (CR) more likely than not

$$
\text { Expect Prob(no CR) } \propto\left(1-\frac{1}{N_{C}^{2}}\right)^{n_{\mathrm{MPI}}}
$$


$\leftrightarrow$ related to coalescence models

> "Parton Level"
(Event structure before confinement)

## What about Baryon Number?

Types of string topologies:

## Open Strings


$(3 \otimes \overline{3})_{\text {singlet }}=\frac{1}{9}$

## SU(3) String Junction

## Closed Strings


$(8 \otimes \overline{8})_{\text {singlet }}=\frac{1}{64}$

$(3 \otimes 3 \otimes 3)_{\text {singlet }}=\frac{1}{27}$
Could we get these at LHC?

## Stochastic sampling of $\operatorname{SU}(3)$ group probabilities (e.g., $3 \otimes 3=6 \oplus \overline{3}$ )

$\Longrightarrow$ Random (re)connections in colour space (weighted by group weights)
"QCD Colour Reconnections"

Example of possible colour configuration
正 Choose this string configuration instead if "string length" ~ total potential energy is lower



New source of baryon-


[^0]ALICE 2021


## Non-Linear String Dynamics?

Count \# of (oriented) flux lines crossing $y=0$ in pp collisions at LHC (according to PYTHIA) - And classify by SU(3) multiplet:


Confining fields may be reaching higher effective representations than simple $q \bar{q}$ (3) ones.
E.g.: 27

$\rightarrow$ Is "emergent tension" driving strangeness enhancement in pp?

Colour Ropes (Bierlich et al.),

+ Close-Packing: Altmann, Bernardinis,
Jueid, PS, Zaccolo (in progress)


## Work in Progress: Strangeness Enhancement from Close-Packing

## Idea: each string exists in an effective background produced by the others

Close-packing


Dense string environments
$\rightarrow$ Casimir scaling of effective string tension
$\rightarrow$ Higher probability of strange quarks

## Strange Junctions



String breaks VS.

String tension could be different from the vacuum case compared to near a junction
Altmann, Bernardinis, Jueid, PS, Zaccolo (in progress)


$\leftrightarrow$ Impact on EAS muon rates?

## Beyond Strings — QGP?

## Currently most realistic complete approach for $\mathrm{pp} \leftrightarrow \mathrm{pA} \leftrightarrow \mathrm{AA}$ ?

The core-corona solution [Werner 2007]: mix discrete strings with continuous QGP


Allows smooth transition between string and hydro descriptions. Implemented in EPOS MC Qualitatively agrees with ALICE strangeness data (but too steep rise with multiplicity?)

## Thorny Issue ! The Proton-to-Pion Ratio

## Note:

Observed $p / \pi$ in pp collisions at LHC is lower than in $e^{+} e^{-}$ones (LEP).

I think this is now the main challenge for strangeness-enhancement models

Interactions?
Upscattering/Annihilation?
Octet vs Triplet fragmentation?...?


Forward region important for cosmic-ray physics $\Rightarrow$ LHCf.

Also for FASER/. . . and the Forward Physics Facility.

Wide spread of predictions; no generator perfect. PYTHIA: $\pi^{0}$ too hard, n too soft.

May require improved modelling of

- beam remnant,
- diffraction, and
- $\mathrm{c} / \mathrm{b} / \tau$ production.

Some possible actions for harder baryons and softer mesons:

- Use QCDCR for better central baryon production. [Christiansen \& PS, 1505.01681]
- Make diquark remnant take more than twice quark ditto: (already default) helps some.
- In string diquark picture B and $\overline{\mathrm{B}}$ are nearest neighbours, but with popcorn allow intermediate meson: ... BM $\overline{\mathrm{B}} \ldots$ Thus leading diquark either BMM . . . or MBM.... New: forbid latter possibility (or only suppress it).

- Normal fragmentation function

$$
f(z) \propto \frac{1}{z}(1-z)^{a} \exp \left(-\frac{b m_{\perp}^{2}}{z}\right), \quad z=\frac{\left(E+p_{z}\right)_{\text {hadron }}}{\left(E+p_{z}\right)_{\text {left in string }}}
$$

modified with separately tuned ( $a$ and) $b$ for leading diquark.

- Reduce primordial $k_{\perp}$ in remnant for soft collisions.






## A New Framework for Hadronic Collisions ( $\rightarrow$ Cosmic Rays)

Based on 2 articles by Marius Utheim \& TS: "A Framework for Hadronic Rescattering in pp Collisions", Eur. Phys. J. C80 (2020) 907, arXiv:2005.05658 "Hadron Interactions for Arbitrary Energies and Species, with Applications to Cosmic Rays", Eur. Phys. J. C82 (2022) 21, arXiv:2108. 03481

- Models arbitrary hadron-hadron collisions at low energies.
- Models arbitrary hadron-p/n collisions at any energy.
- Initialization slow, $\sim 15$ minutes, * but thereafter works for any hadron-p/n at any energy, and $\star$ initialization data can be saved, so only need to do once.
- The Angantyr nuclear geometry part used to extend to hadron-nucleus at any energy.
- Native C++ simplifies interfacing Pythia $8 \leftrightarrow$ Corsika 8 .
- So far limited comparisons with data.


## Comparison to Other Models -

Maximilian Reininghaus, TS, M. Utheim, arXiv:2303:02792


Additive quark rule $\sigma_{\pi \mathrm{p}} \approx(2 / 3) \sigma_{\mathrm{pp}}$ at high energies.
Simple extension to pA: $\sigma_{h \mathrm{~A}}=\frac{A}{\left\langle n_{\text {coll }}\right\rangle} \sigma_{h \mathrm{p}} \quad$ where $\left\langle n_{\text {coll }}\right\rangle$ comes from Angantyr

## Comparison to Other Models — 2

Hadronic
cascades quite different




EM cascades quickly decouple from hadronic ones

## Some Further PYTHIA aspects:

Includes charm and bottom (and jets)
Native C++ $\rightarrow$ multithreading
Users can do tunings themselves
$\rightarrow$ study air-shower / accelerator interplay

Thank you

## Extra Slides

## Returning to Strings: the String Fragmentation Function

Schwinger $\Longrightarrow$ Gaussian $p_{\perp}$ spectrum (transverse to string axis) \& Prob(d:u:s) $\approx 1: 1: 0.2$
The meson $M$ takes a fraction $z$ of the quark momentum, Probability distribution in $z \in[0,1]$ parametrised by Fragmentation Function, $f\left(z, Q_{\mathrm{HAD}}^{2}\right)$


## (Note on the Length of Strings)

## In Spacetime:

String tension $\approx 1 \mathrm{GeV} / \mathrm{fm} \rightarrow$ a $50-\mathrm{GeV}$ quark can travel 50 fm before all its kinetic energy is transformed to potential energy in the string. Then it must start moving the other way. ( $\rightarrow$ "yo-yo" model of mesons. Note: string breaks $\rightarrow$ several mesons)

## The MC implementation is formulated in momentum space

Lightcone momenta $p_{ \pm}=E \pm p_{z}$ along string axis
$\rightarrow$ Rapidity (along string axis) and $p_{\perp}$ transverse to it

If the quark gives all its energy to a single pion traveling along the $z$ axis

$$
y=\frac{1}{2} \ln \left(\frac{E+p_{z}}{E-p_{z}}\right)=\frac{1}{2} \ln \left(\frac{\left(E+p_{z}\right)^{2}}{E^{2}-p_{z}^{2}}\right) \quad \rightarrow \quad y_{\max } \sim \ln \left(\frac{2 E_{q}}{m_{\pi}}\right) \underset{\substack{\text { Increasing } E_{q} \rightarrow \text { logarithmic } \\ \text { growth in rapidity range }}}{\substack{\text { and }}}
$$

## Particle Production:

Scaling in $z \Longrightarrow$ flat in rapidity (long. boost invariance)

"Lightcone scaling"

$$
\left\langle n_{\mathrm{ch}}\right\rangle \approx c_{0}+c_{1} \ln E_{\mathrm{cm}}, \sim \text { Poissonian multiplicity distribution }
$$

## A Brief History of MPI (in PYTHIA)

## 1987 [Sjöstrand \& van Zijl, Phys.Rev.D 36 (1987) 2019]

Cast MPI as Sudakov-style evolution equation Analogous to $\sigma_{\mathrm{X}+\mathrm{jet}}\left(p_{\perp}\right) / \sigma_{\mathrm{X}}$ for parton showers
$\frac{d P_{\text {hardest }}}{d^{2} b d x_{T 1}}=p\left(x_{T 1}, b\right) \exp \left\{-\int_{x_{T 1}}^{1} p\left(x_{T}^{\prime}, b\right) d x_{T}^{\prime}\right\}$
$\mathrm{p} \propto \sigma_{2 \rightarrow 2}\left(x_{T}, b\right) / \sigma_{p p} \quad ; \quad x_{T}=2 \hat{p}_{\perp} / \sqrt{s}$
with Impact-parameter dependence

b
Crucial to describe "Underlying Event"
a.k.a. "Jet Pedestal": hard jets are accompanied by — and sit on top of - higher-than average particle densities
"Outside the [jet], a constant $\mathrm{E}_{\mathrm{T}}$ plateau is observed, whose height is independent of the jet $\mathrm{E}_{\mathrm{T}}$. Its value is substantially higher than the one observed for minimum bias events." (compared with the average $=$ minimum-bias pp collision)

## Interleaved Evolution in PYTHIA

## Interleaved Evolution

2005
[Sjöstrand \& PS, Eur.Phys.J.C 39 (2005) 129] Interleave MPI \& ISR evolutions in one common sequence of $p$ T
$\rightarrow$ ISR \& MPI "compete" for the available $x$ in the proton remnant.

2011 [Corke \& Sjöstrand. JHEP 03 (2011) 3 32] Also include FSR in interleaving


## Example for pp collisions at 13 TeV — PYTHIA's default MPI model


*note: can be arbitrarily soft

## Strings should push each other transversely

 Colour-electric fields $\rightarrow$ Classical forceModel string radial shape \& shoving physics
$\Longrightarrow$ force $f\left(d_{\perp}\right)=\frac{g \kappa d_{\perp}}{R^{2}} \exp \left(-\frac{d_{\perp}^{2}}{4 R^{2}}\right)$

$g$ : fraction of energy in chromo-electric field (as opposed to in condensate or magnetic flux)
$d_{\perp}$ : transverse distance (in string-string "shoving frame")
$R$ : string radius
$\kappa$ : string tension $\sim 1 \mathrm{GeV} / \mathrm{fm}$


## What a strange world we live in, said Alice

We know ratios of strange hadrons to pions strongly increase with event activity Landmark measurement by ALICE (2017)


## What could be driving this?



## Confront with Measurements

## LHC experiments report very large (factor-10) enhancements in heavy-

 flavour baryon-to-meson ratios at low $\mathrm{p}_{\mathrm{T}}$ !

## Confront with Measurements: Strangeness

## What about Strange heavy-flavour baryons?





String Formation Beyond Leading Colour, Christiansen \& PZS, 1505.01681

New: String Junctions Revisited, Altmann \& PZS, 2404.12040

Also: baryon asymmetry diluted by extra baryon pairs



[^0]:    Illustration by J. Altmann

