CLUSTER OF EXCELLENCE QUANTUM UNIVERSE



# Quadratic Coupling of Axions to Photons

based on arXiv: <u>2307.10362</u>

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# Outline of this talk

Axion-photon coupling:  $\mathcal{L} \supset -\frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$ 

#### Quadratic Coupling of Axions to Photons:

Which operator?  $(\partial a)^2 F_{\mu\nu} F^{\mu\nu}$  vs  $a^2 F_{\mu\nu} F^{\mu\nu}$ 

# Outline of this talk

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#### Quadratic Coupling of Axions to Photons:



# Outline of this talk

Axion-photon coupling:  $\mathcal{L} \supset -\frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$ 

#### Quadratic Coupling of Axions to Photons:



• No tree-level contribution from  $\mathcal{L} \supset -\frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$ 

$$\mathcal{M}_{t;+,-} = \frac{g_{a\gamma\gamma}^2}{t} \left[ \left( \varepsilon_1^+ \cdot \varepsilon_2^- \right) \left( p_1 \cdot q \right) - \left( \varepsilon_1^+ \cdot q \right) \left( \varepsilon_2^- \cdot p_1 \right) \right] \left( p_2 \cdot q \right) = 0$$

• No tree-level contribution from  $\mathcal{L} \supset -\frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$ 

◎ 1-loop contribution from QCD Chiral Lagrangian:



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● 1-loop contribution from QCD Chiral Lagrangian:



$$c_{F^2} = \pi m_u m_d / 3 (m_u + m_d)^2 \sim 0.2$$

The same result can be obtained from threshold corrections H. Kim, A. Lenoci, G. Perez, W. Ratzinger (2307.14962)

• No tree-level contribution from  $\mathcal{L} \supset -\frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$ 

◎ 1-loop contribution from QCD Chiral Lagrangian:

#### If axion make up dark matter:



See also H. Kim, A. Lenoci, G. Perez, W. Ratzinger (2307.14962)



Constraints from atomic clocks



Constraints from atomic interferometers (AION)



Constraints from MICROSCOPE mission (searching for violation of weak equivalent principle)



• Toy model 1: Shift symmetry-breaking EFT resulting from KSVZ-like setup

$$\mathcal{L}_{\rm UV} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \bar{\psi}_L i \not\!\!D \psi_L + \bar{\psi}_R i \not\!\!D \psi_R + \left( y \phi \bar{\psi}_L \psi_R + \text{h.c.} \right) + \partial_\mu \phi^\dagger \partial^\mu \phi - V(\phi^\dagger \phi)$$

Toy model 1: Shift symmetry-breaking EFT resulting from KSVZ-like setup

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Without explicit shift-symmetry-breaking terms:

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$$V(\phi^{\dagger}\phi) = \lambda \left(\phi^{\dagger}\phi - \frac{f_{a}^{2}}{2}\right)^{2} \text{ and } \phi = \frac{1}{\sqrt{2}}(\rho + f_{a})e^{ia/f_{a}}$$

$$\mathcal{L}_{a^{2}F^{2}}^{1-\text{loop}} = \underbrace{\frac{i^{2}}{16\pi^{2}}\frac{1}{3M_{\psi}^{2}}\left[M_{\psi}^{2}\frac{a^{2}}{f_{a}^{2}}\right](iQ_{\psi}e)^{2}F_{\mu\nu}F^{\mu\nu}}_{\psi} + \underbrace{\frac{i^{2}}{16\pi^{2}}\frac{2}{3M_{\psi}}\left[M_{\psi}\frac{(ia)^{2}}{2f_{a}^{2}}\right](iQ_{\psi}e)^{2}F_{\mu\nu}F^{\mu\nu}}_{\psi} = 0$$

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Toy model 1: Shift symmetry-breaking EFT resulting from KSVZ-like setup

$$\mathcal{L}_{\rm UV} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \bar{\psi}_L i D \psi_L + \bar{\psi}_R i D \psi_R + \left( y \phi \bar{\psi}_L \psi_R + \text{h.c.} \right) + \partial_\mu \phi^\dagger \partial^\mu \phi - V(\phi^\dagger \phi)$$

Adding a shift-symmetry-breaking terms to the potential:

$$V(\phi^{\dagger}\phi) \supset \lambda \left(\phi^{\dagger}\phi - \frac{f_a^2}{2}\right)^2 + g^2 \left(\phi^{\dagger}\phi - \frac{f_a^2}{2}\right) \left(1 - \cos\left(\frac{a}{f_a}\right)\right) \quad \text{and} \quad \phi = \frac{1}{\sqrt{2}}(\rho + f_a)e^{ia/f_a}$$

Preserves CP,  $\mathbb{Z}_n$  symmetry, no mass term of an ALP at tree-level

Toy model 1: Shift symmetry-breaking EFT resulting from KSVZ-like setup

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Adding a shift-symmetry-breaking terms to the potential:

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Integrating out:  $\rho$ ,  $\psi$   

$$\mathcal{L}_{a^{2}F^{2}}^{1-\text{loop}} \supset \frac{i^{2}}{16\pi^{2}} \frac{2}{3M_{\psi}} \left[M_{\psi}\frac{\rho}{f_{a}}\right] (iQ_{\psi}e)^{2}F_{\mu\nu}F^{\mu\nu}\Big|_{\rho=\rho_{c}(a)}$$

$$\supset \frac{1}{48\pi^{2}}(Q_{\psi}e)^{2}\frac{g^{2}}{f_{a}^{2}M_{\rho}^{2}}a^{2}F_{\mu\nu}F^{\mu\nu}$$

One can derive the variation of  $\alpha(a)$ :

$$c_{_{F^2}} = \frac{4\pi}{3} Q_{\psi}^2 \frac{g^2}{M_{\rho}^2} \,, \quad \alpha(a) = \alpha \left( 1 + \frac{Q_{\psi}^2 \alpha}{3\pi} \frac{g^2 a^2}{M_{\rho}^2 f_a^2} \right)$$

#### • Toy model 2: QCD-like dynamics for an ALP

Main ingredients: Dark QCD-like & Dark photon sectors:  $SU(N)' \otimes U(1)'$ Chiral fermions charged under the symmetries of the dark sectors

Using SU(N)' instanton to break the shift-symmetry of an ALP:

$$V(a) = -m_{\pi'}^2 f_{\pi'}^2 \cos\left(\frac{a}{2f_a}\right) \simeq -\frac{1}{2}m_a^2 a^2 \qquad \text{(assuming: } m_{u'} = m_{d'}\text{)}$$

Analogously to QCD axion computations, the fine structure constant in the dark sector is also modified:

$$\alpha' \simeq \alpha'_0 \left( 1 + \frac{\alpha'_0 a^2}{48 \pi f_a^2} \right)$$

Using kinetic mixing  $\chi$  between photon & dark photon:

$$\alpha \simeq \alpha_0 \left( 1 + \chi^2 \alpha'_0 \frac{a^2}{48 \pi f_a^2} \right)$$

## Bounds in axion-photon couplings parameter space: ALPs



## Sensitivity of the quadratic coupling: ALPs



## Quadratic coupling of Axions to Photons: Summary

#### • For the QCD axion:

1.) Dynamics endowing axion with a mass can also lead to a quadratic coupling of axions to photons

2.) Other pre-existing constraints are stronger, but quadratic coupling offers new ways to probe parts of parameter space and exploit the precision of table-top experiments of fundamental constants.

#### For Axion-like-Particles:

1.) Resulting constrain may be the strongest bounds in large regions of parameter space

2.) Constructing a UV completion for the shift-breaking coupling with less finetuning of the ALP mass is not easy

# **Backup slides**

## **Backup slides:** One-Loop Effective Action

Path integral formalism: 
$$e^{iS_{eff}[\psi_{SM}^{L}](\mu)} = \int \mathcal{D}\psi_{BSM}^{H} e^{iS[\psi_{BSM}^{H},\psi_{SM}^{L}](\mu)}$$

Find classical solution by solving EOM:

$$\frac{\delta S\left[\psi_{BSM}^{H},\psi_{SM}^{L}\right]}{\delta\psi_{BSM}^{H}}\bigg|_{\psi_{BSM}^{H}=\psi_{BSM,c}^{H}}=0 \Rightarrow \psi_{BSM,c}^{H}(\psi_{SM}^{L})$$

Expand action around minimum:

$$S\left[\psi_{BSM}^{\boldsymbol{H}}\right] = S\left[\psi_{BSM,c}^{\boldsymbol{H}} + \eta\right] = S\left[\psi_{BSM,c}^{\boldsymbol{H}}\right] + \frac{1}{2} \left.\frac{\delta^2 S}{\delta(\psi_{BSM}^{\boldsymbol{H}})^2}\right|_{\psi_{BSM,c}^{\boldsymbol{H}}} \eta^2 + \mathcal{O}(\eta^3)$$

Integrate over quantum fluctuation  $\eta$  :

$$e^{iS_{eff}\left[\psi_{SM}^{L}\right]} = e^{iS\left[\psi_{BSM,c}^{H}\right]} \left[ \det\left(-\frac{\delta^{2}S}{\delta(\psi_{BSM}^{H})^{2}}\Big|_{\psi_{BSM,c}^{H}}\right) \right]^{-c_{s}}$$

 $c_s$  is spin factor ( $c_s = +1/2$  for real scalar, -1 for Dirac fermion)

Re-write the determinant,  $\det(A) = e^{\operatorname{Tr} \log A}$  :

$$S_{eff} \left[ \psi_{SM}^{L} \right] = S \left[ \psi_{BSM,c}^{H} \left( \psi_{SM}^{L} \right) , \psi_{SM}^{L} \right] + ic_s \operatorname{Tr} \log \left( - \left. \frac{\delta^2 S}{\delta (\psi_{BSM}^{H})^2} \right|_{\psi_{BSM,c}^{H}} \right)$$

Tree-level

One-loop level

#### Backup slides—Building axion EFTs: One-loop matching using functional method

We parametrise the shape of UV Lagrangian as follows:

$$\mathcal{L}_{\mathrm{UV}}^{\mathrm{fermion}} \left[ \Psi_H, \phi \right] \supset \bar{\Psi}_H \left[ i D_\mu \gamma^\mu - M + X[\phi] \right] \Psi_H$$

General coupling with background fields

Example:  $X[\phi] = V_{\mu}[\phi]\gamma^{\mu} - A_{\mu}[\phi]\gamma^{\mu}\gamma^{5} - W_{1}[\phi]i\gamma^{5}$ 

Path Integral: extract the one-loop (heavy-only) piece:  $e^{iS_{eff}[\phi_L]} = \int \mathcal{D}\bar{\Psi}_H \mathcal{D}\Psi_H e^{iS_{UV}[\Psi_H, \phi_L]}$ 

$$S_{eff}^{1-loop} = -i \operatorname{Tr} \log \left( \left. - \frac{\delta^2 S}{\delta \Psi_H^2} \right|_{\Psi_{H,c}} \right) = -i \operatorname{Tr} \log \left( i D_\mu \gamma^\mu - M + X[\phi] \right)$$

Evaluating the functional trace:  $\operatorname{Tr} \mathcal{O}(i \not \!\!D, X) = \int d^4x \int \frac{d^4q}{(2\pi)^4} \operatorname{tr} \mathcal{O}(i \not \!\!D - \not \!\!\!q, X)$ 

$$\mathcal{L}_{\rm EFT}^{\rm 1loop} = i \operatorname{tr} \sum_{n=1}^{\infty} \frac{1}{n} \int \frac{d^4q}{(2\pi)^4} \left[ \frac{-1}{\not(2\pi)^4} \left( -iD_{\mu}\gamma^{\mu} - V_{\mu}[\phi]\gamma^{\mu} + A_{\mu}[\phi]\gamma^{\mu}\gamma^5 + W_1[\phi]i\gamma^5 \right) \right]^n$$
  
Encapsulate axion derivative couplings

Expanding order by order (ex: up to n=6)

- Integrating over momentum q (use Dimensional Regularisation for divergence integrals)
- ullet Evaluating the Dirac traces (careful with  $\gamma^5$  )