

A model of pseudo-Nambu-Goldstone dark matter with two complex scalars

Yu Hamada (DESY)

based on

[arXiv:2401.02397](https://arxiv.org/abs/2401.02397), JHEP 05 (2024) 076

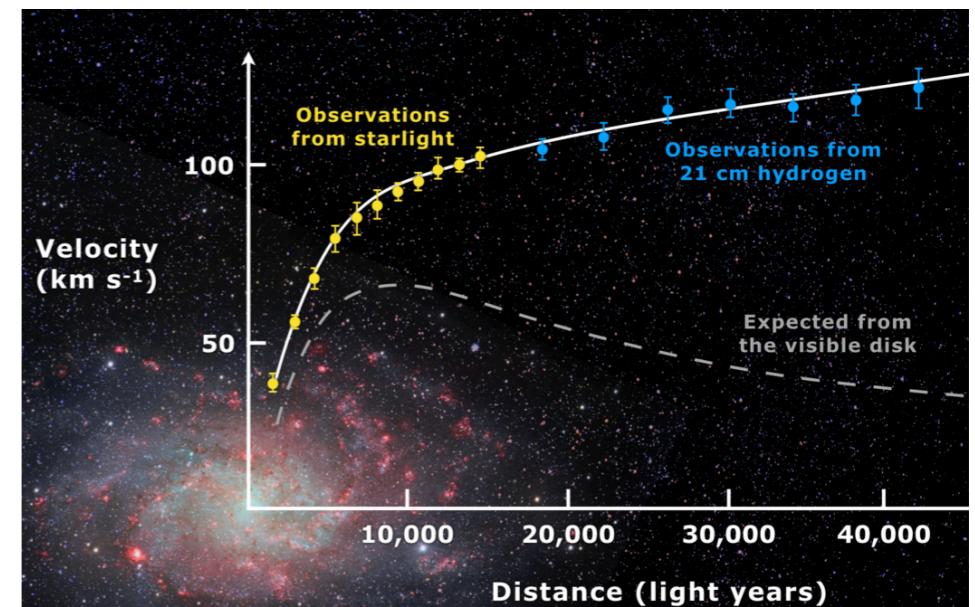
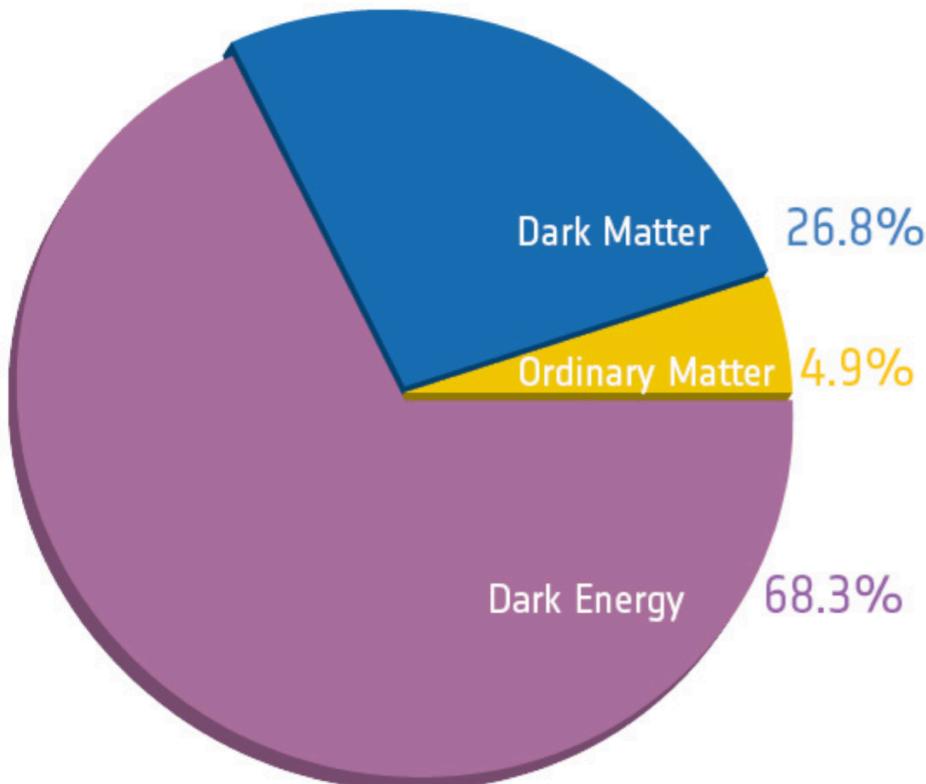
in collaboration with

T. Abe (Tokyo U. of Science) and K. Tsumura (Kyushu U.)

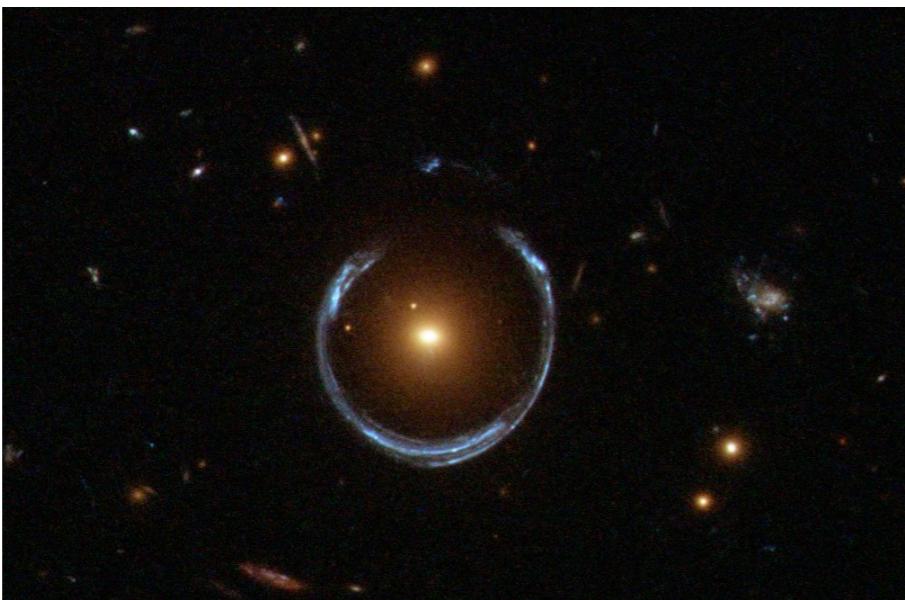


Dark matter

- There are many evidences for dark matter



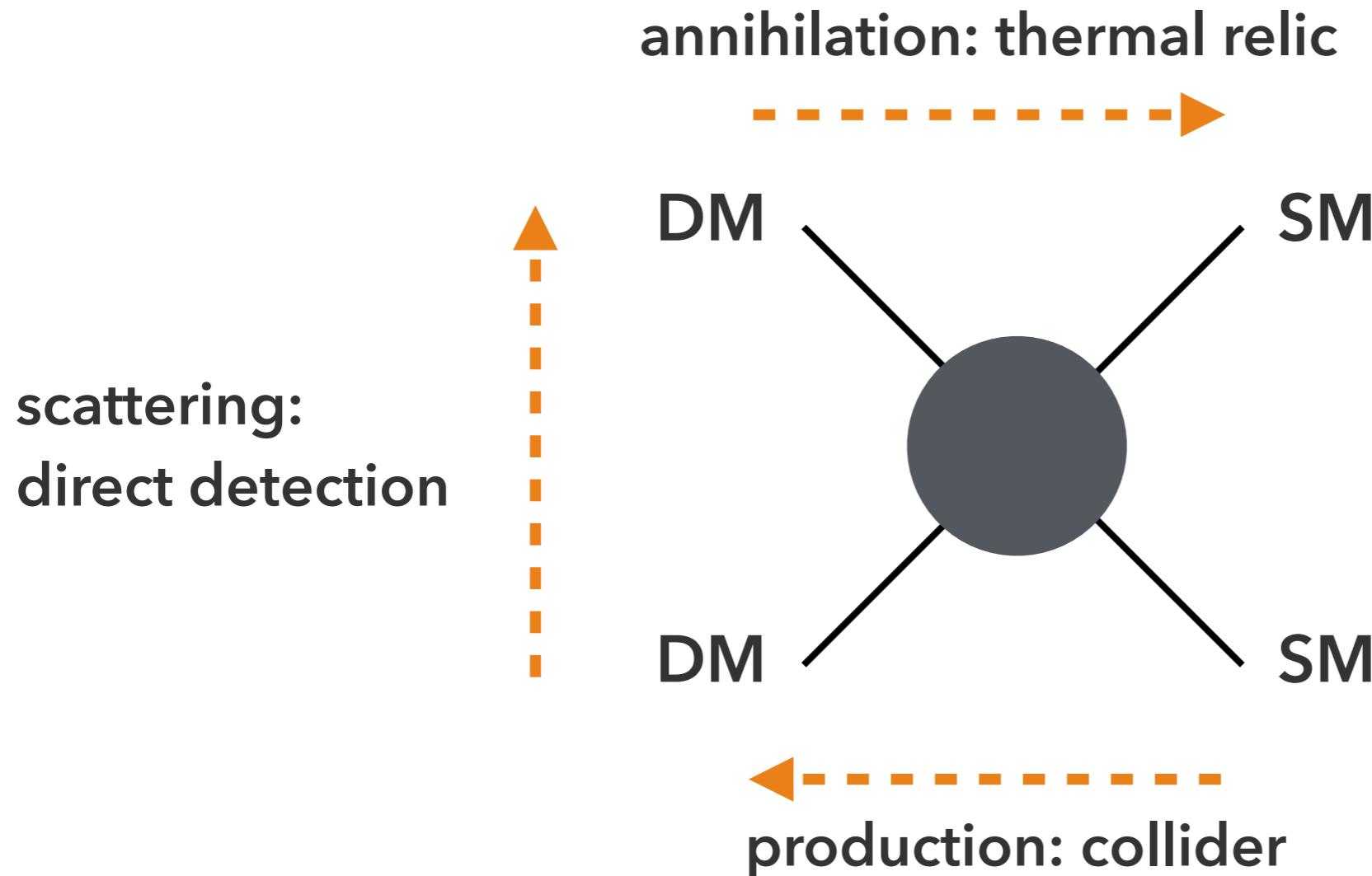
Galaxy rotation curve (from Wikipedia)



Gravitational lens (from Wikipedia)

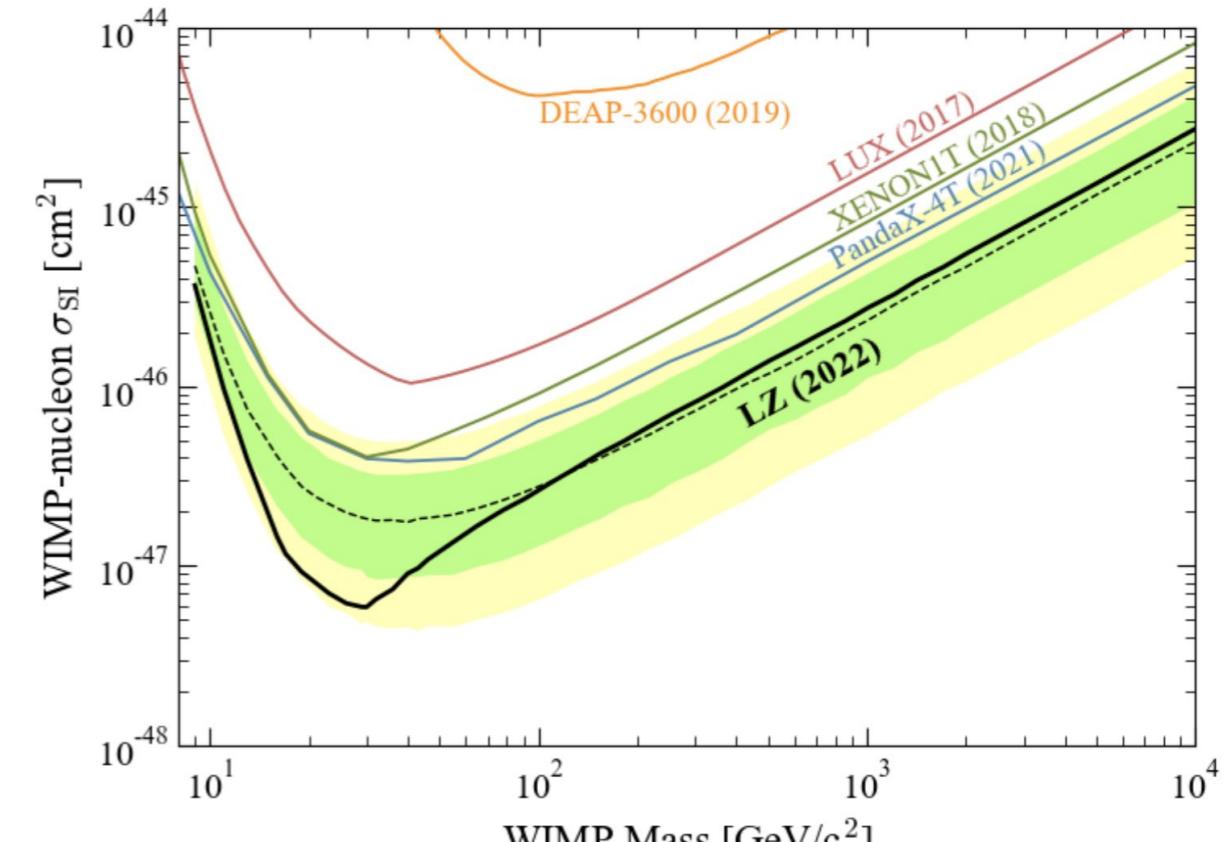
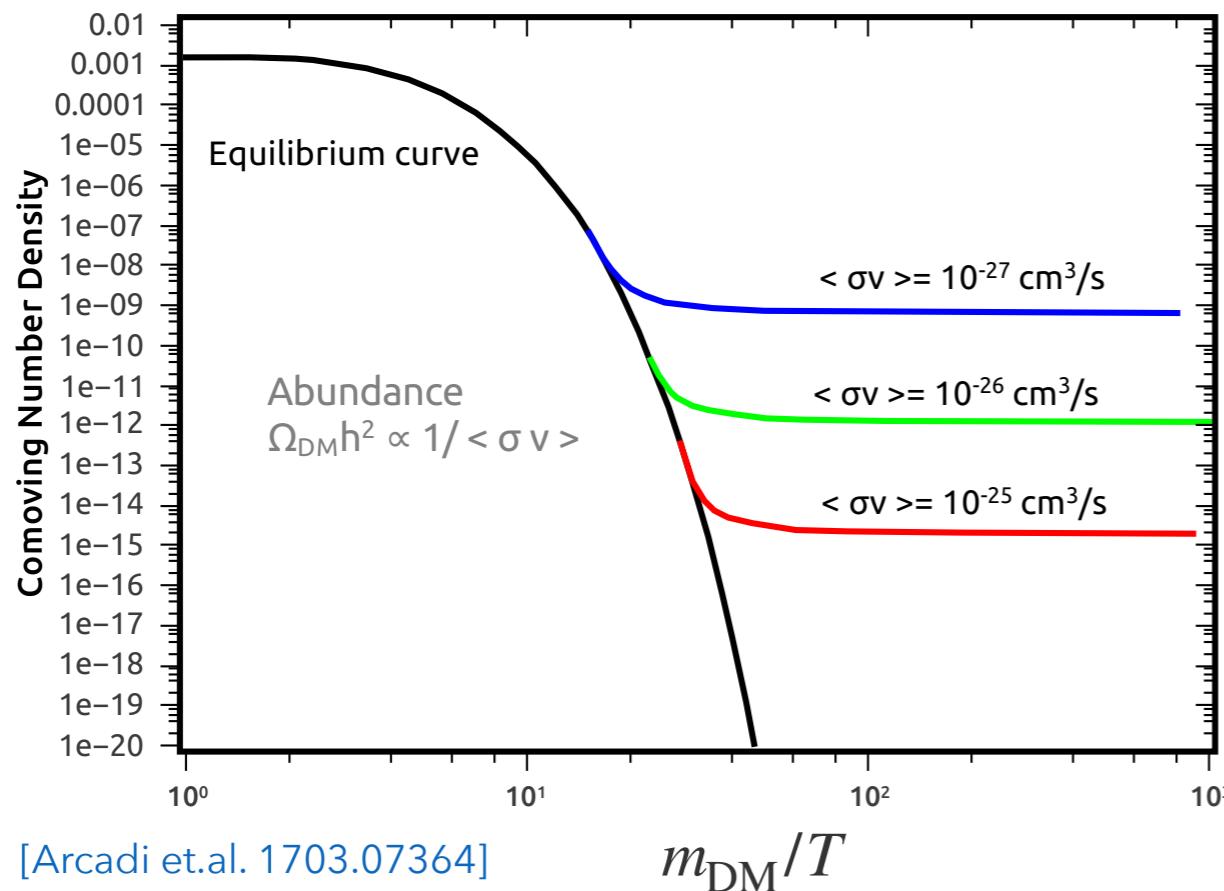
WIMP (Weakly Interacting Massive Particle)

- "weakly" interacting with SM sector
- explain DM relic with thermal freeze-out mechanism
- correlated processes → nice testability



DM relic vs Direct detection

- must keep DM annihilation rate (DM relic)
- severe constraint on DM-nucleon σ_{SI} (direct detection)



naively incompatible → needs some trick!

pseudo-Nambu-Goldston DM

[Gross-Lebedev-Toma '17]

$$V(\phi, H) = V_{\text{SM}}(H) - \mu_\phi^2 |\phi|^2 + \lambda_\phi |\phi|^4 + \lambda_{\phi H} |\phi|^2 |H|^2 + (\mu_{\text{break}}^2 \phi^2 + \text{h.c.})$$

$$\downarrow \\ m_{\text{DM}}^2$$

- SM Higgs H + singlet complex scalar ϕ
- global $U(1)$: $\phi \rightarrow e^{i\theta} \phi$, softly broken by μ_{break}
- VEV $\langle \phi \rangle \neq 0 \rightarrow \text{pseudo NG boson w/ derivative interaction}$

[cf. soft pion theorem]

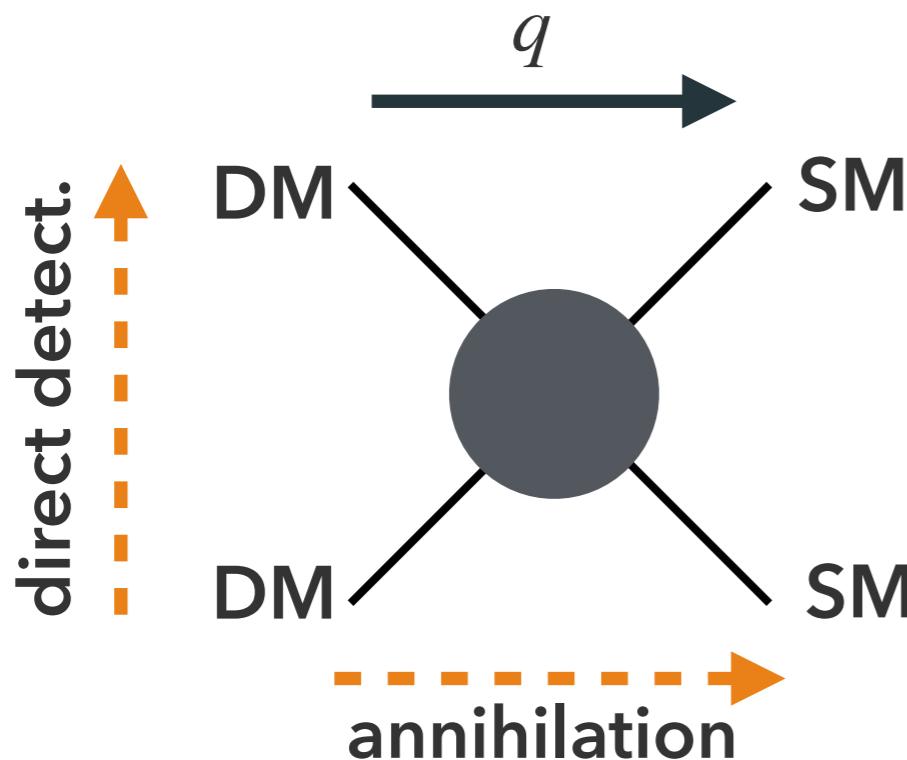
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[cf. soft pion theorem]

$$\alpha q^2 = \begin{cases} t \simeq 0 & @ \text{direct detection} \\ s \simeq 4m_{\text{DM}}^2 & @ \text{annihilation} \end{cases}$$

e evade the dilemma elegantly!

Cons of original pNG DM model

$$V(\phi, H) = V_{\text{SM}}(H) - \mu_\phi^2 |\phi|^2 + \lambda_\phi |\phi|^4 + \lambda_{\phi H} |\phi|^2 |H|^2 + (\mu_{\text{break}}^2 \phi^2 + \text{h.c.})$$

- SSB of $\mathbb{Z}_2 \rightarrow$ **domain wall problem**
- other soft-breaking terms ϕ, ϕ^3 are dropped by hand

→ **needs to be improved**

- other pNG DM models in literature:

[Abe-Toma-Tsumura 2001.03954] [Okada-Raut-Shafi, 2001.05910]

[Abe et. al., 2104.13523] [Okada et. al., 2105.03419]

[Abe-Hamada, 2205.11919] [Liu et. al., 2208.06653]

[Otsuka- et. al. 2210.08696] [Chiang et. al., 2311.13753]

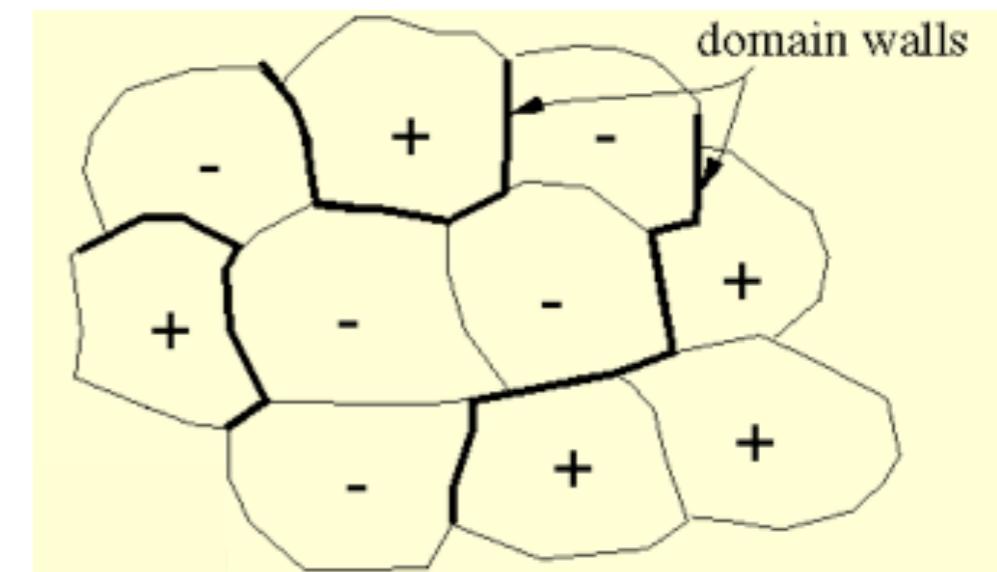


fig. from https://www.ctc.cam.ac.uk/outreach/origins/cosmic_structures_two.php

We propose an alternative pNG DM model

Plan of talk

- Introduction
- Our model
- Prediction and constraints
- Summary

Our model

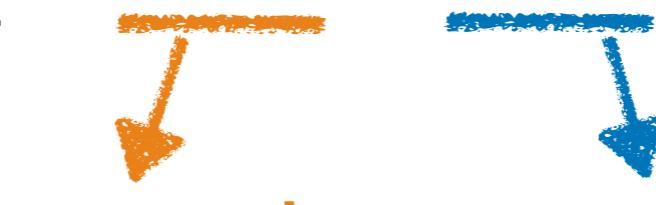
Model

[Abe-YH-Tsumura 2401.02397]

- scalar extension w/ two singlet complex scalars ϕ_1, ϕ_2
- impose
$$\frac{U(1)_1 \times U(1)_2}{\text{(phase rot. of } \phi_1, \phi_2)}$$
$$= U(1)_V \times U(1)_A$$

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gauged **global, softly broken**

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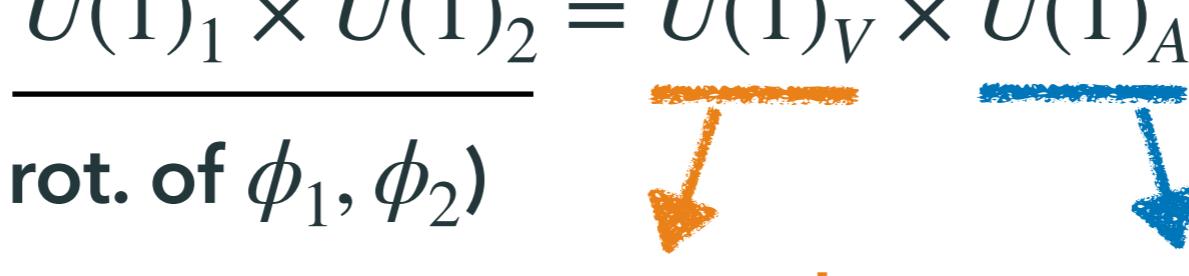
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- exchange \mathbb{Z}_2 symmetry: $\phi_1 \leftrightarrow \phi_2$
unique soft-breaking term

$$\begin{aligned} V(H, \phi_1, \phi_2) = & m_1^2 (|\phi_1|^2 + |\phi_2|^2) - (m_{12}^2 \phi_1^* \phi_2 + \text{h.c.}) \\ & + \frac{\lambda_1}{2} (|\phi_1|^4 + |\phi_2|^4) + \lambda_3 |\phi_1|^2 |\phi_2|^2 \\ & - m_H^2 |H|^2 + \lambda_H |H|^4 + \lambda_{H1} |H|^2 (|\phi_1|^2 + |\phi_2|^2) \end{aligned}$$

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gauged global, softly broken
 - exchange \mathbb{Z}_2 symmetry: $\phi_1 \leftrightarrow \phi_2$
unique soft-breaking term
 - $\langle \phi_1 \rangle = \langle \phi_2 \rangle = v_s/2 \rightarrow$ spontaneously breaks **only** two $U(1)$'s
→ no DW problem
- $$V(H, \phi_1, \phi_2) = m_1^2 (|\phi_1|^2 + |\phi_2|^2) - (m_{12}^2 \phi_1^* \phi_2 + \text{h.c.})$$
$$+ \frac{\lambda_1}{2} (|\phi_1|^4 + |\phi_2|^4) + \lambda_3 |\phi_1|^2 |\phi_2|^2$$
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Particle contents

[Abe-YH-Tsumura 2401.02397]

- convenient to introduce "Higgs basis"

$$\phi_+ \equiv \frac{1}{\sqrt{2}} (\phi_1 + \phi_2) \quad \phi_- \equiv \frac{1}{\sqrt{2}} (\phi_1 - \phi_2)$$

\mathbb{Z}_2 even \mathbb{Z}_2 odd, $\langle \phi_- \rangle = 0$

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\mathbb{Z}_2 odd, $\langle \phi_- \rangle = 0$

- expand around VEVs:

$$H = \begin{pmatrix} i\pi_W^+ \\ \frac{1}{\sqrt{2}} (v_h + \sigma + i\pi_Z) \end{pmatrix} \quad \phi_+ = \frac{1}{\sqrt{2}} (v_s + s_+ + i\pi_V) \quad \phi_- = \frac{1}{\sqrt{2}} (s_- + ia)$$

mix to be h, h'

**stable pNG DM
coming from $U(1)_A$**

The diagram illustrates the mixing of the Higgs field H into the physical fields ϕ_+ and ϕ_- . The Higgs field H is represented by a column vector with components $i\pi_W^+$ and $\frac{1}{\sqrt{2}} (v_h + \sigma + i\pi_Z)$. The mixing is shown as a curved arrow pointing from H to the mixing terms in the definitions of ϕ_+ and ϕ_- . The mixing terms are $\frac{1}{\sqrt{2}} (v_s + s_+ + i\pi_V)$ for ϕ_+ and $\frac{1}{\sqrt{2}} (s_- + ia)$ for ϕ_- . An orange arrow points from the mixing term in ϕ_- to the text "stable pNG DM coming from $U(1)_A$ ".

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\mathbb{Z}_2 even

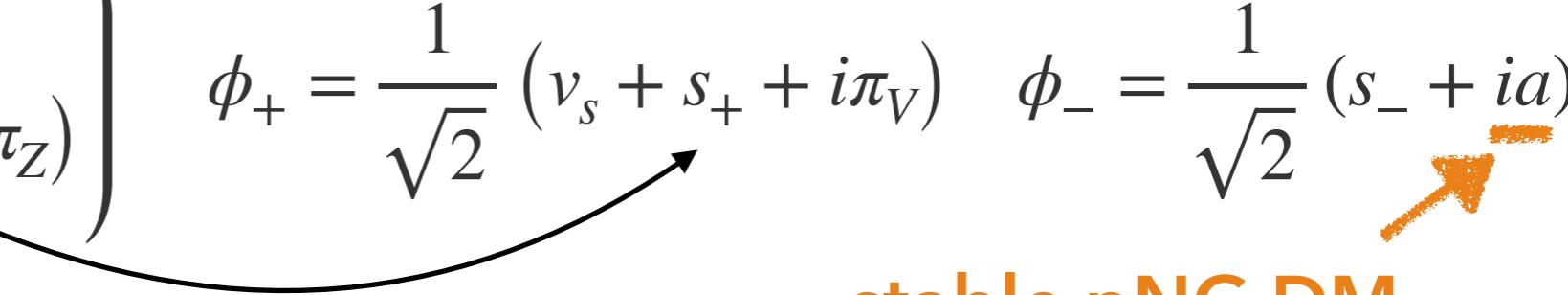
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**stable pNG DM
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- s_- is also \mathbb{Z}_2 odd \rightarrow assume $m_{s_-} > m_{\text{DM}}$

Gauge kinetic mixing

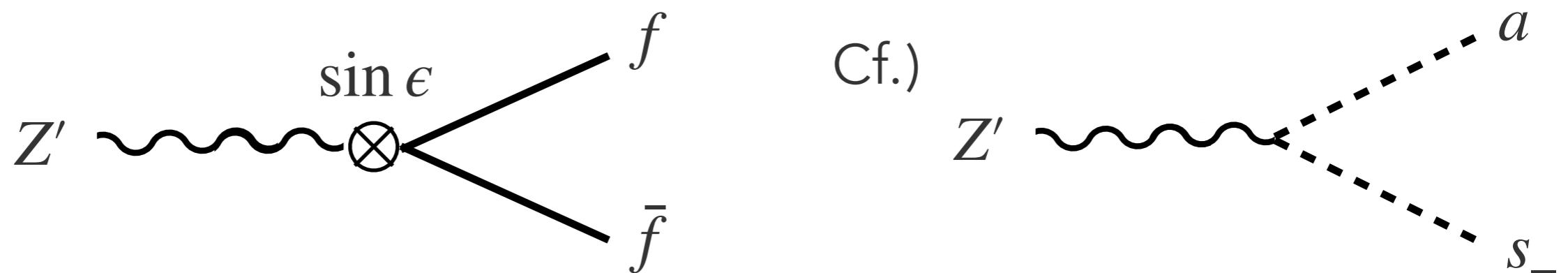
[Abe-YH-Tsumura 2401.02397]

- the new gauge field V_μ mixes with SM Y_μ : $\mathcal{L}_{\text{mix}} = -\frac{\sin \epsilon}{2} Y^{\mu\nu} V_{\mu\nu}$
 - obtain masses from VEVs of SM Higgs H and new scalar ϕ_+
 - mass eigenstates: Z_μ A_μ $Z'_\mu \leftarrow \mathbf{new massive gauge boson}$

Gauge kinetic mixing

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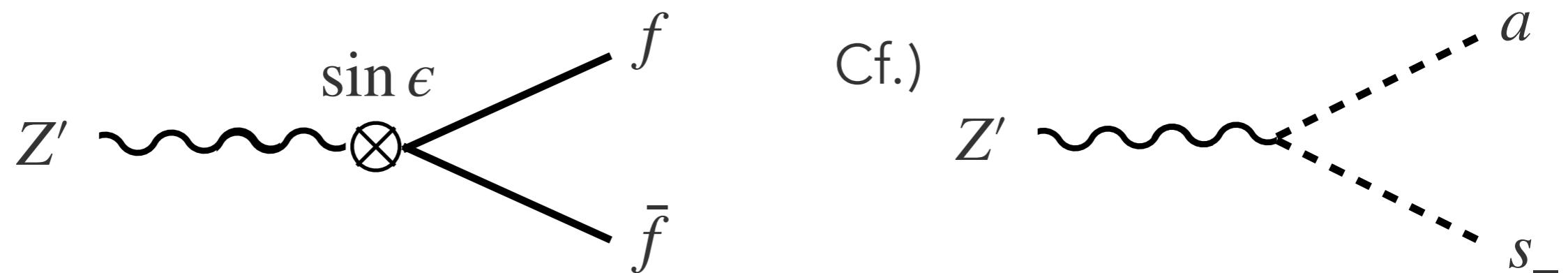
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bonus: Z' can be taken light, say, $m_{Z'} \sim \mathcal{O}(100) \text{ GeV}$

→ its coupling constant remains sufficiently perturbative

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Prediction and constraints

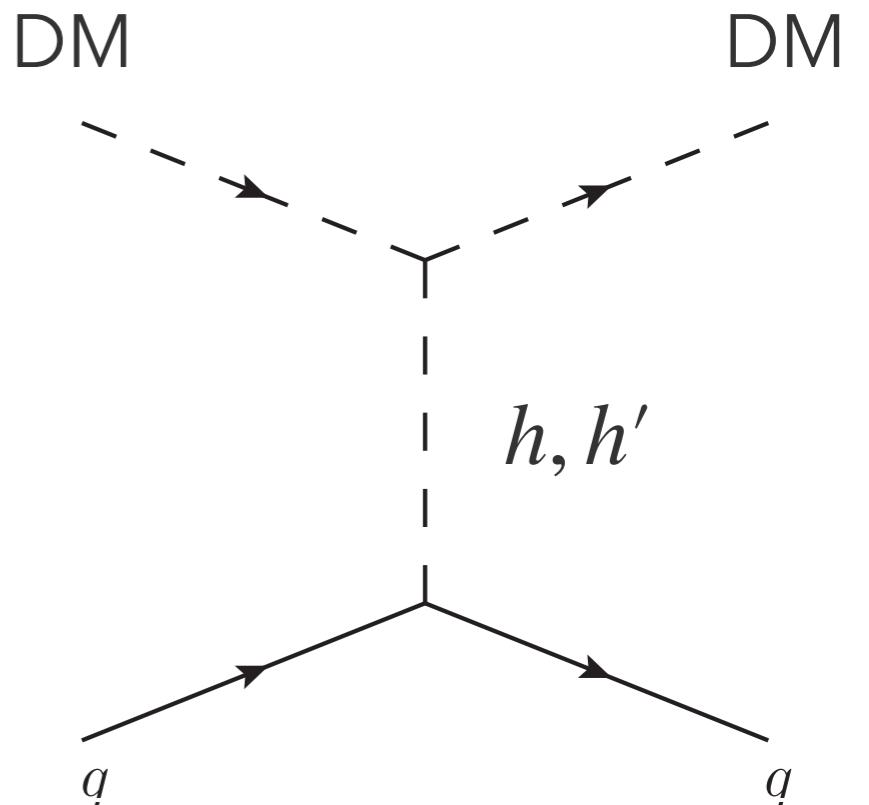
Direct detection @ tree level

[Abe-YH-Tsumura 2401.02397]

- easily shown in non-linear rep.

$$\phi_1 = \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} v_s + s_1 \right) \exp \left[+i \frac{\pi_a}{v_s} \right]$$

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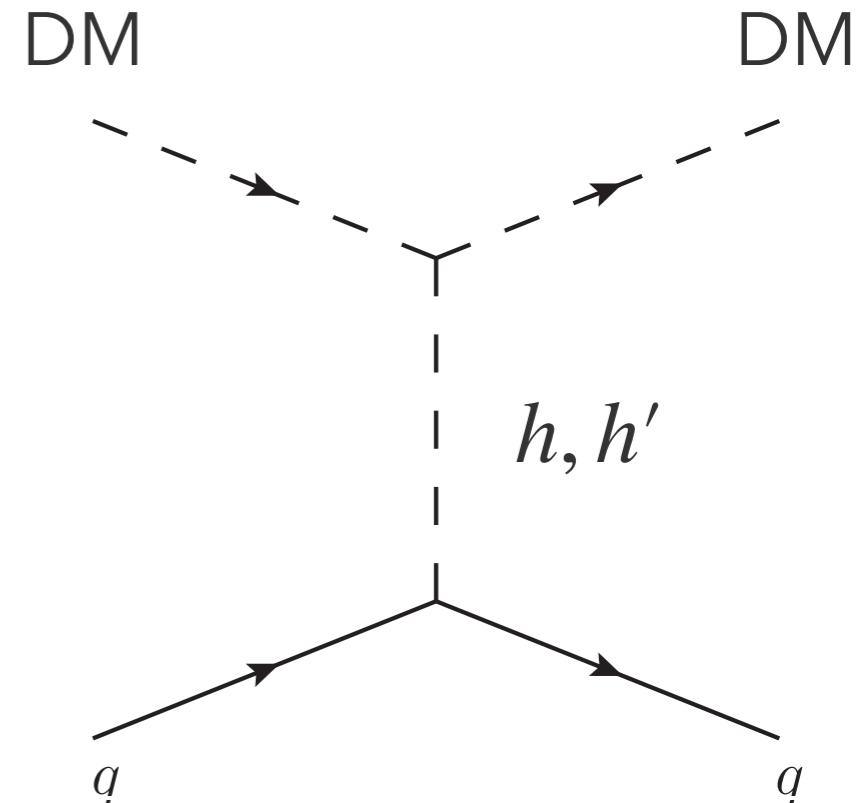
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- relevant cubic coupling

$$s_+ = (s_1 + s_2)/\sqrt{2} = (\text{linear comb of } h h')$$

$$\mathcal{L} \Big|_{s\pi_a\pi_a} = -\frac{2}{v_s} \left[s_+ \pi_a \underbrace{\left(\partial^2 + m_{\text{DM}}^2 \right) \pi_a}_{\rightarrow 0} + \pi_a \partial_\mu \pi_a \underbrace{\partial^\mu s_+}_{\text{on-shell cond.}} \right]$$

$\rightarrow 0$
on-shell cond.

$\rightarrow 0$
low q

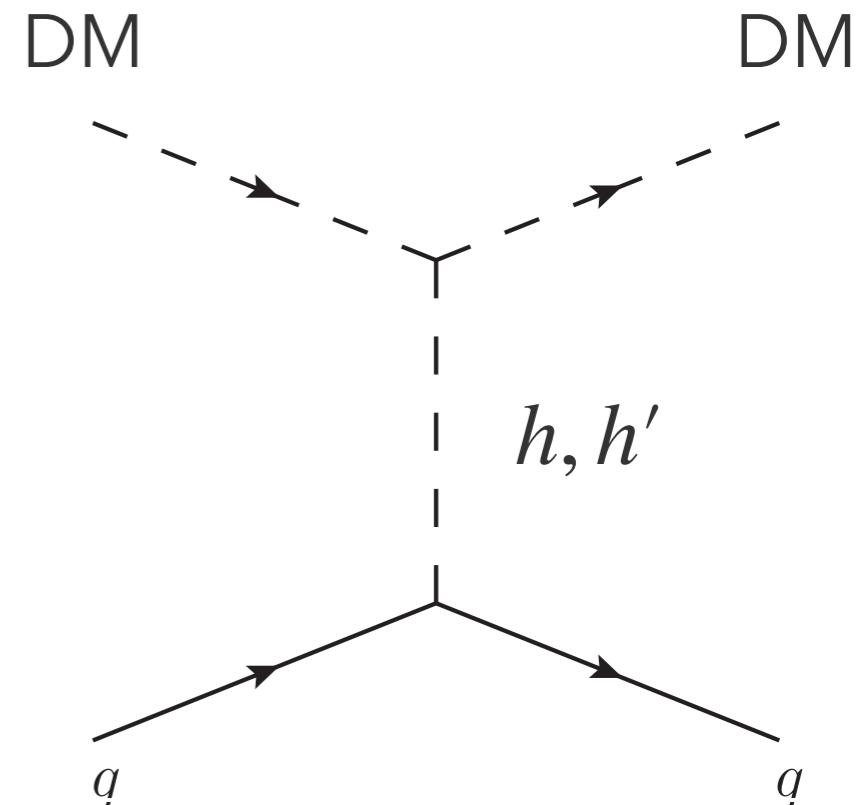
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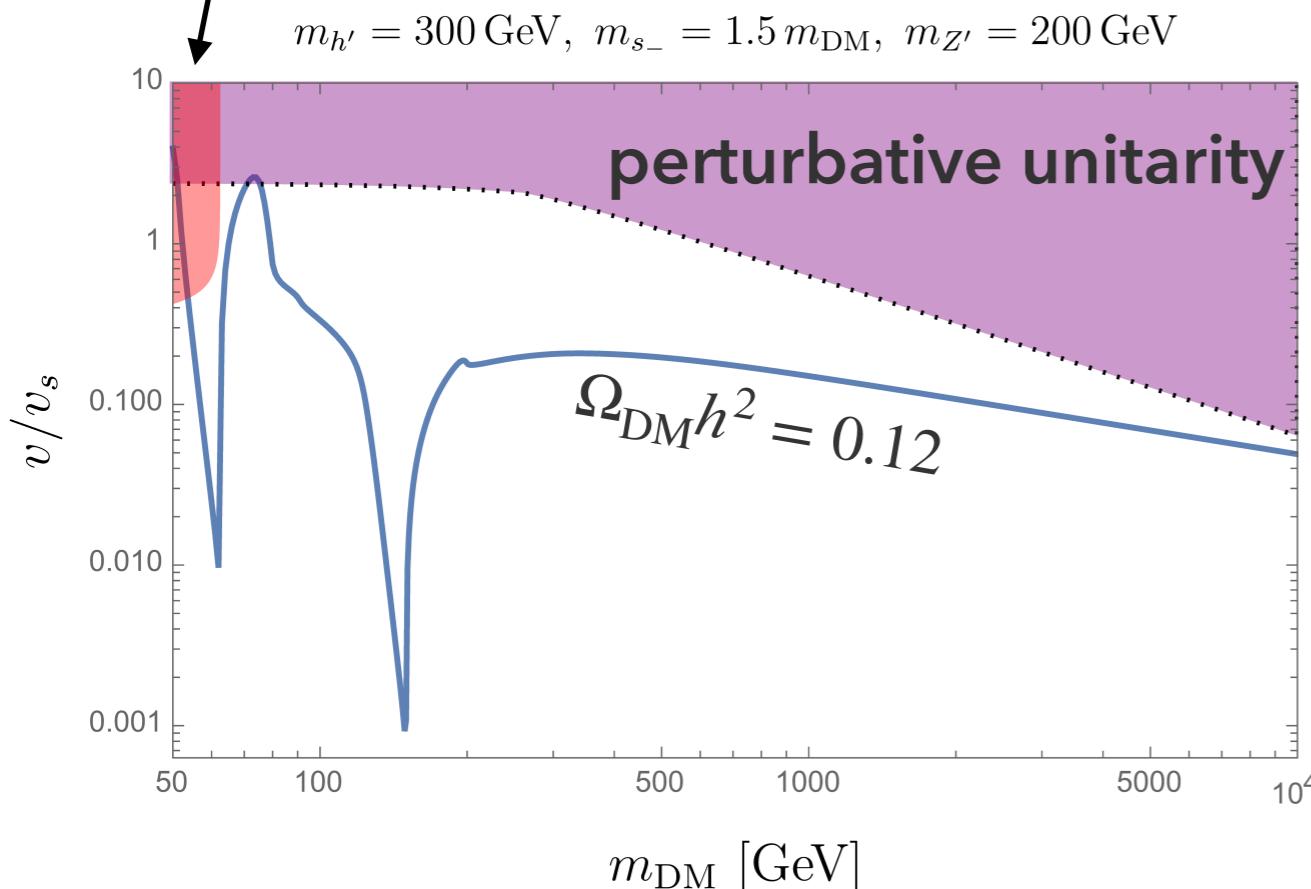
low q

vanishes!

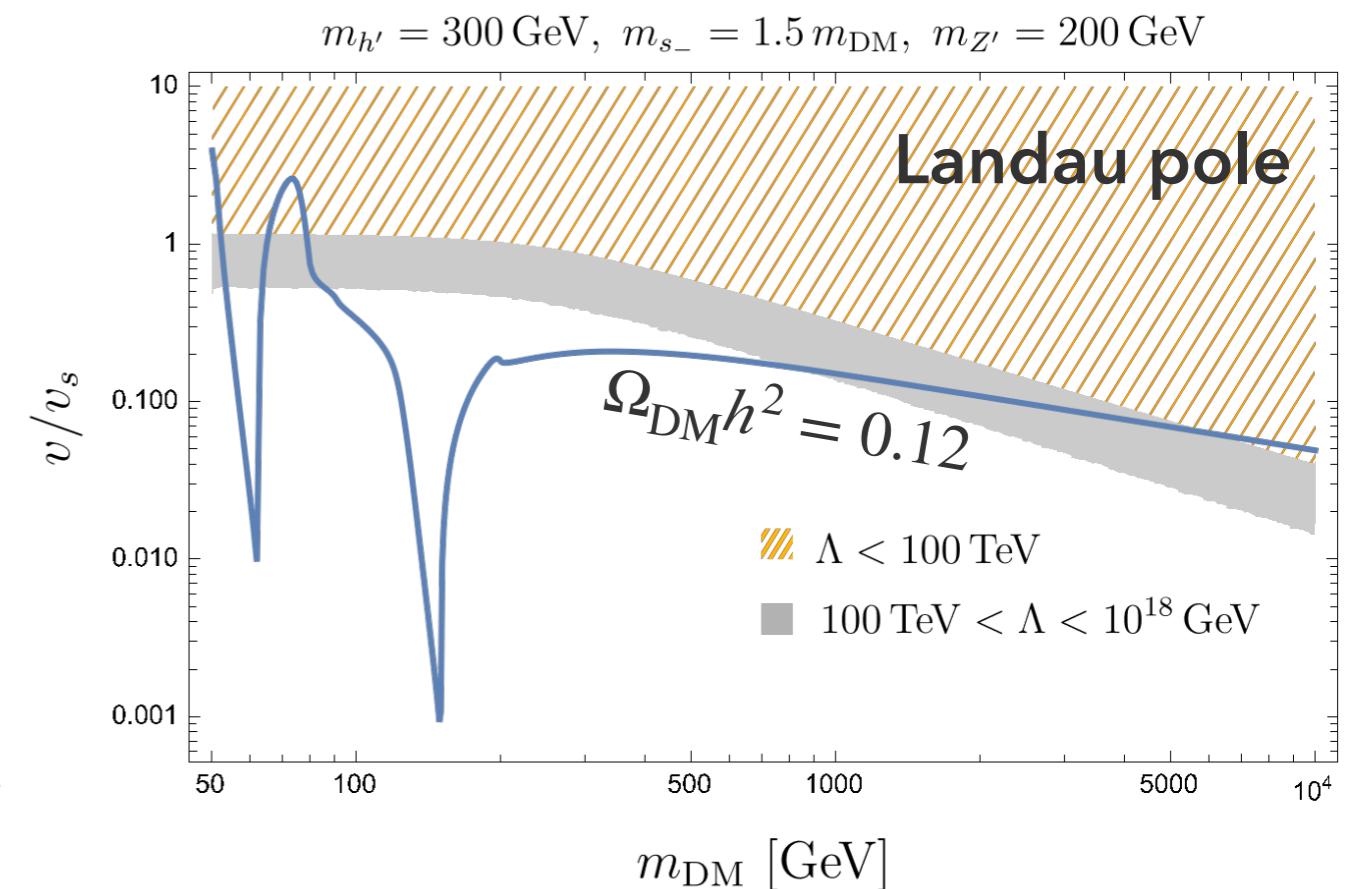
DM thermal relic

[Abe-YH-Tsumura 2401.02397]

higgs invisible decay



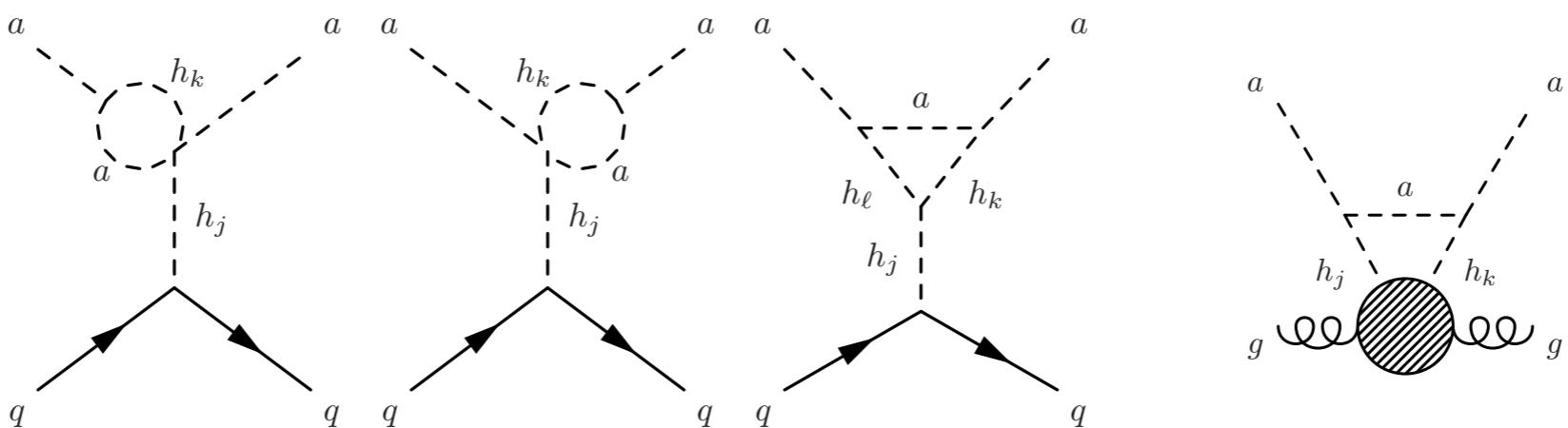
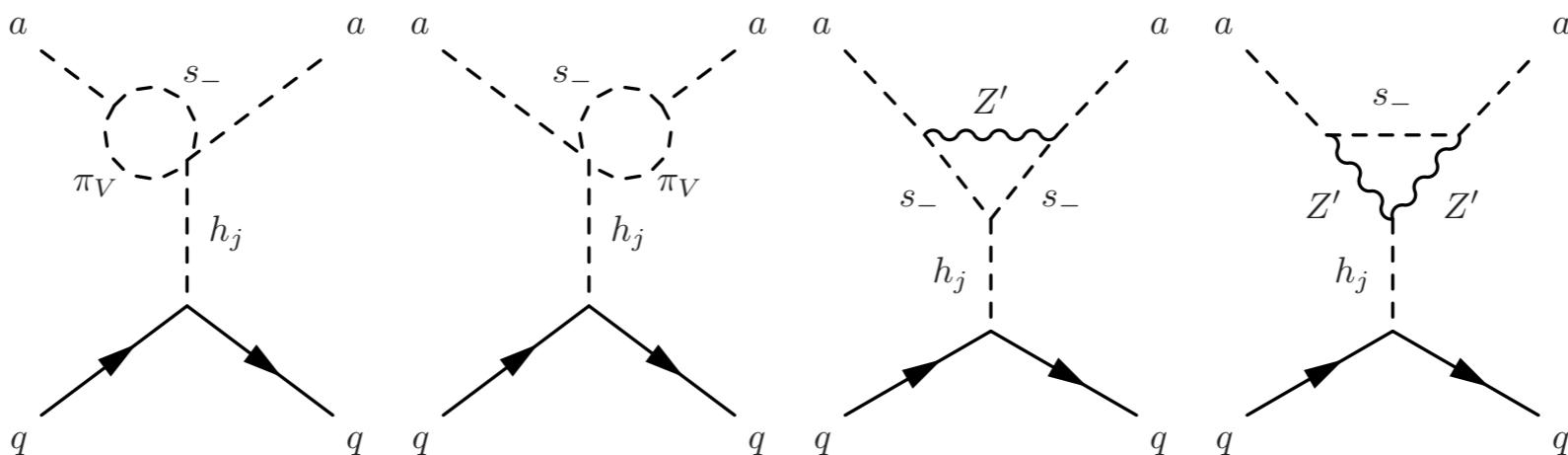
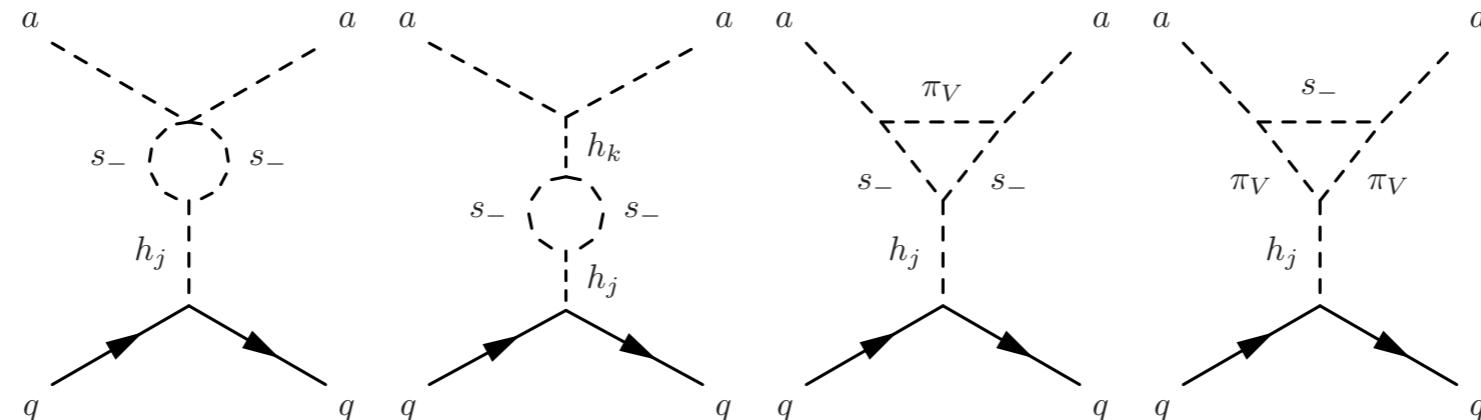
$$\sin \xi = 0.1, \sin \epsilon = 10^{-4}$$



**less constrained by perturbativity and Landau pole,
thanks to the light Z' mass**

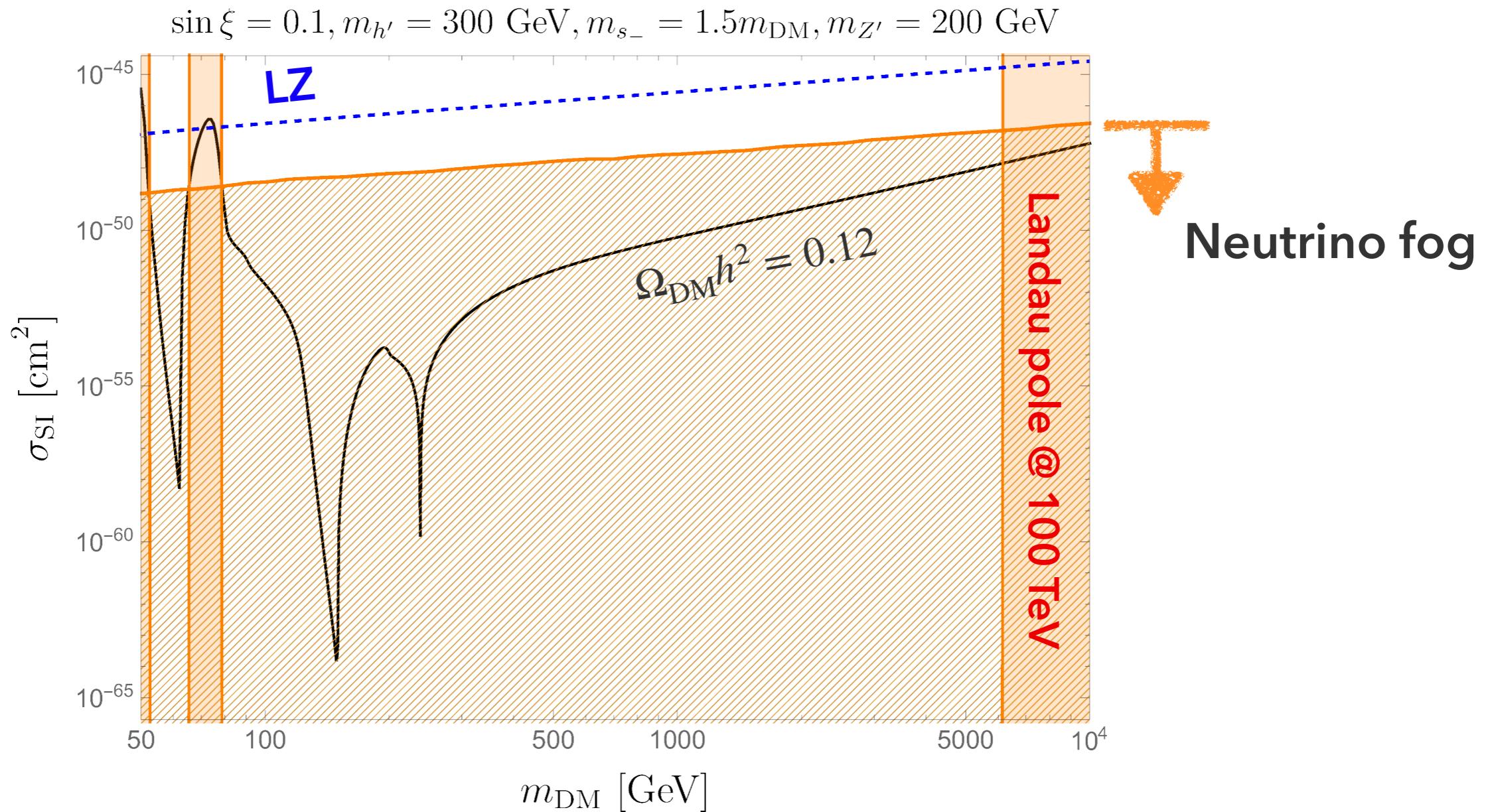
Direct detection @ loop level

[Abe-YH-Tsumura 2401.02397]



Direct detection @ loop level

[Abe-YH-Tsumura 2401.02397]



To test the model, we should go into neutrino fog!

Summary

- pNG DM models evade direct detection constraint keeping the thermal DM relic.
- The original model suffers from DW problem and assumes the other soft-breaking terms to vanish.
- We proposed a new pNG DM model in which dark sector consists of two scalars ϕ_1, ϕ_2 and gauge boson Z'_μ
- pNG DM comes from SSB of $U(1)_A$ and is stabilized by exchange symmetry $\phi_1 \leftrightarrow \phi_2$.
- Direct detection signal is covered by neutrino fog
→ experiments must go further!

Backup

Comparison w/ other pNG DM models

	Gauge	Global sym.	w/ soft breaking	DM stability	DM d.o.f.	DW
original	---	$[U(1) \times Z_2]$	$\text{dim } 2 \rightarrow Z_2$	Z_2	real	x
B-L (GUT)	$U(1)_{B-L}$	$U(1) \times U(1)$	$\text{dim } 3 \rightarrow U(1)_{B-L}$	decaying	real	○
Dark $U(1)_D$	$U(1)_D$	$U(1) \times U(1) [\times C_D]$	$\text{dim } 3 \rightarrow U(1)_D$	decaying	real	○
Abe-Hamada	$U(1)_D$	$U(1) \times SU(2) [\times C_D]$	$\text{dim } 2 \rightarrow U(1)_D \times U(1)$	$U(1)$	complex	○
Dark $SU(2)_D$	$SU(2)_D$	$SO(4)$	$\text{dim } 3 \rightarrow SU(2)_D \times U(1)$	$U(1)$	complex	○
our work	$U(1)_D$	$U(1) \times U(1) [\times S_2]$	$\text{dim } 2 \rightarrow U(1)_D \times S_2$	Z_2	real	○

Constraints

- perturbative unitarity

$$|\lambda_H| < 4\pi,$$

$$|2\lambda_1 - \lambda_3| < 8\pi,$$

$$|\lambda_3| < 8\pi,$$

$$\left| 2\lambda_1 + \lambda_3 + 6\lambda_H \pm \sqrt{(2\lambda_1 + \lambda_3 - 6\lambda_H)^2 + 16\lambda_{H1}^2} \right| < 16\pi$$

- potential boundedness from below

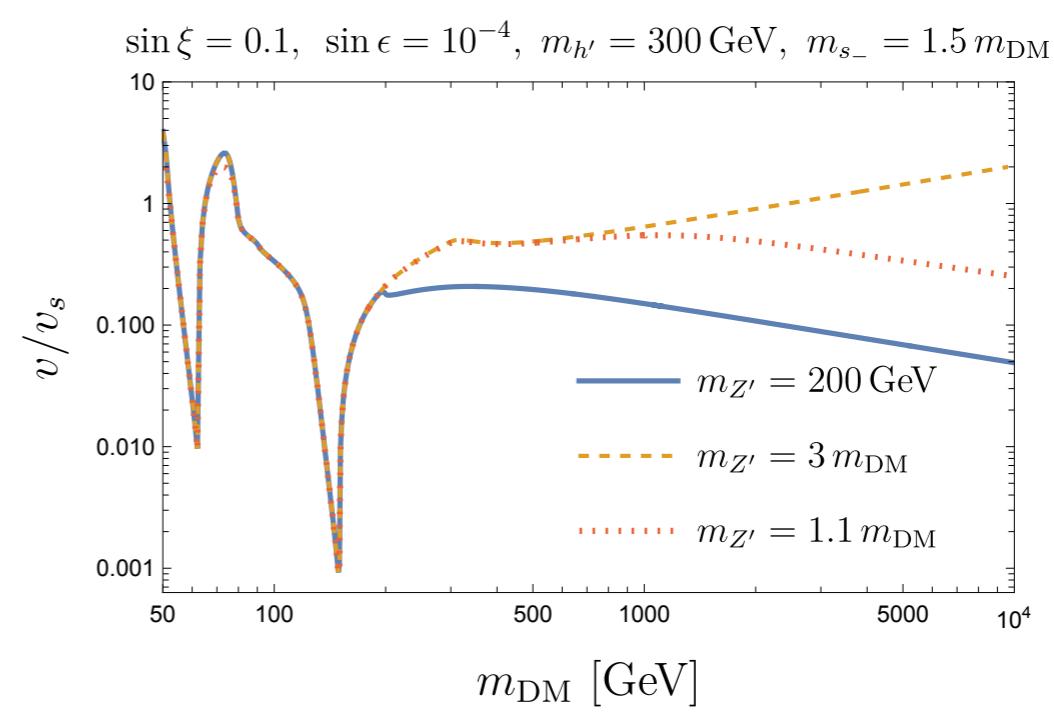
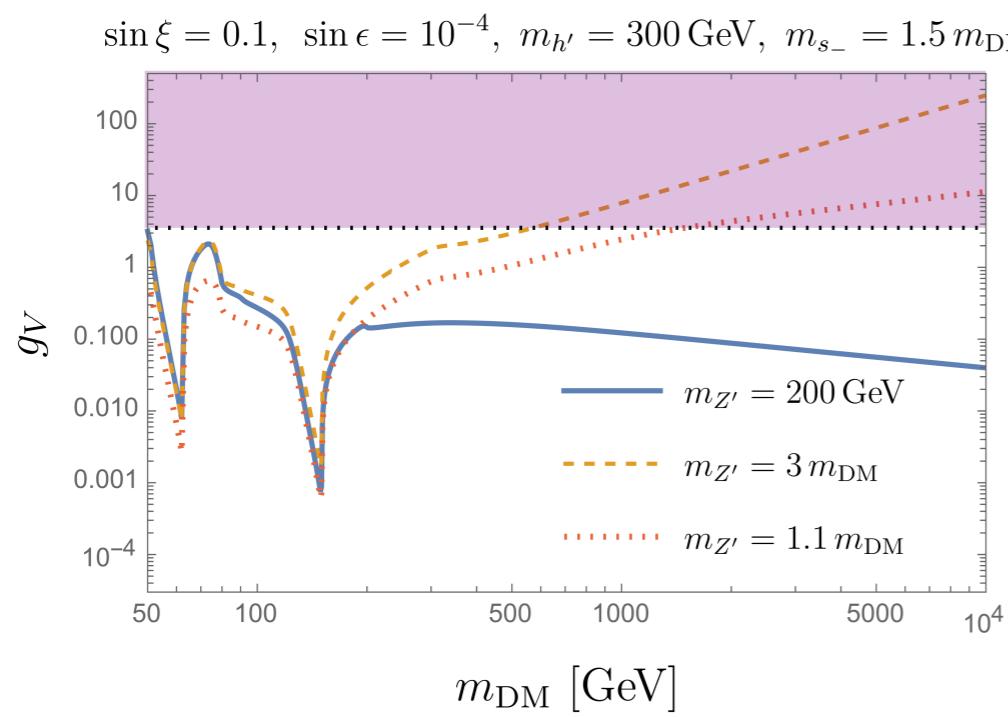
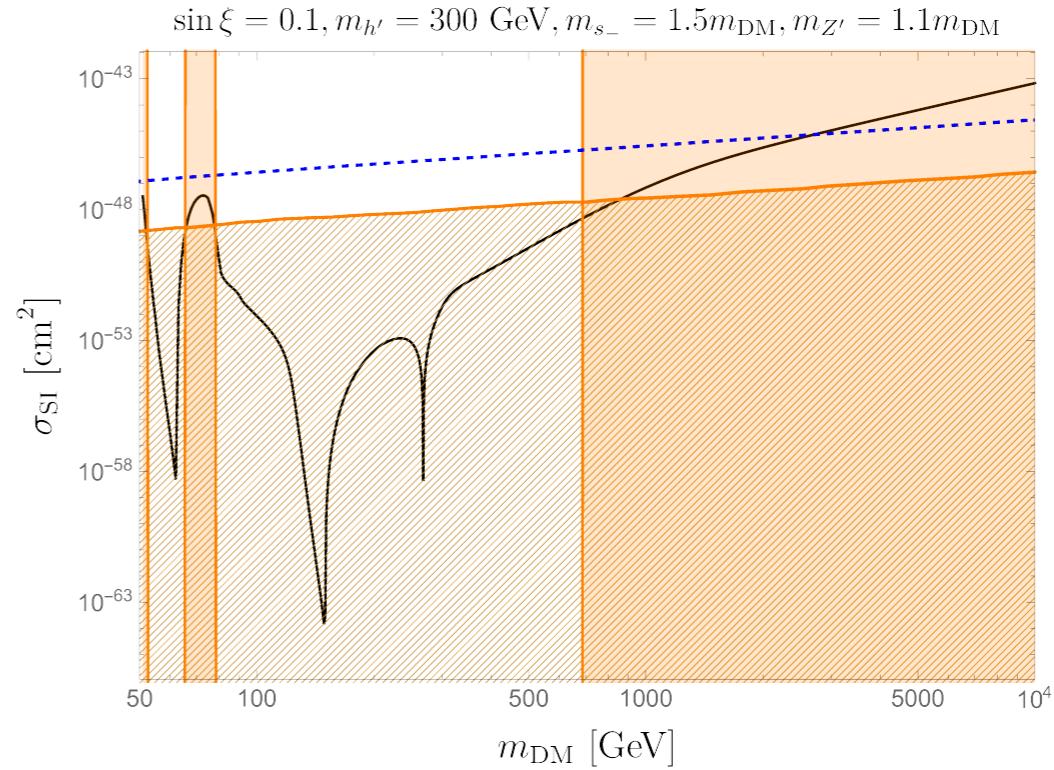
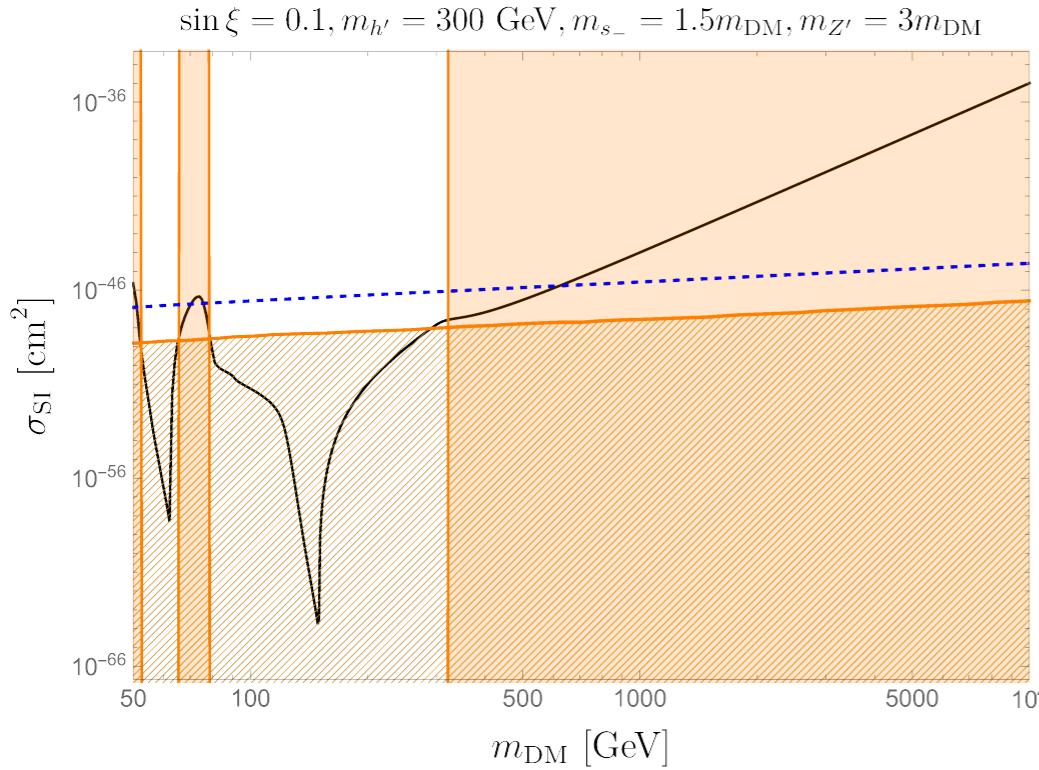
$$\lambda_1 > 0, \quad \lambda_H > 0, \quad \lambda_1 + \lambda_3 > 0,$$

$$\lambda_{H1} + \sqrt{2\lambda_1 \lambda_H} > 0, \quad \lambda_{H1} + \sqrt{(\lambda_1 + \lambda_3) \lambda_H} > 0$$

1-loop RGE

$$(4\pi)^2 \beta_{\lambda_1} \equiv (4\pi)^2 \frac{d\lambda_1(\mu)}{d \ln \mu} = 10\lambda_1^2 + 2\lambda_3^2 + 12g_V^4 - 12g_V^2\lambda_1,$$
$$(4\pi)^2 \beta_{\lambda_3} \equiv (4\pi)^2 \frac{d\lambda_2(\mu)}{d \ln \mu} = 4\lambda_3^2 + 8\lambda_1\lambda_3 + 12g_V^4 - 12g_V^2\lambda_3,$$
$$(4\pi)^2 \beta_{g_V} \equiv (4\pi)^2 \frac{dg_V(\mu)}{d \ln \mu} = \frac{2}{3}g_V^3$$

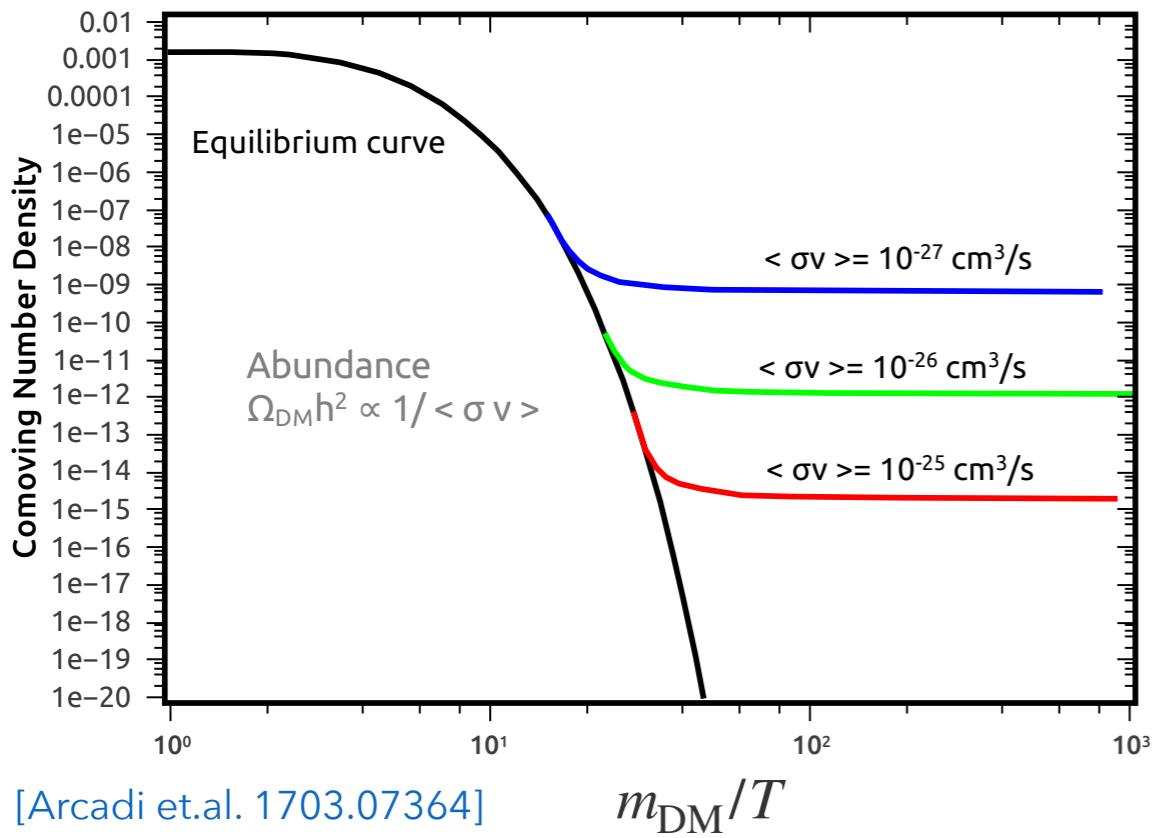
Other benchmarks



DM relic vs Direct detection

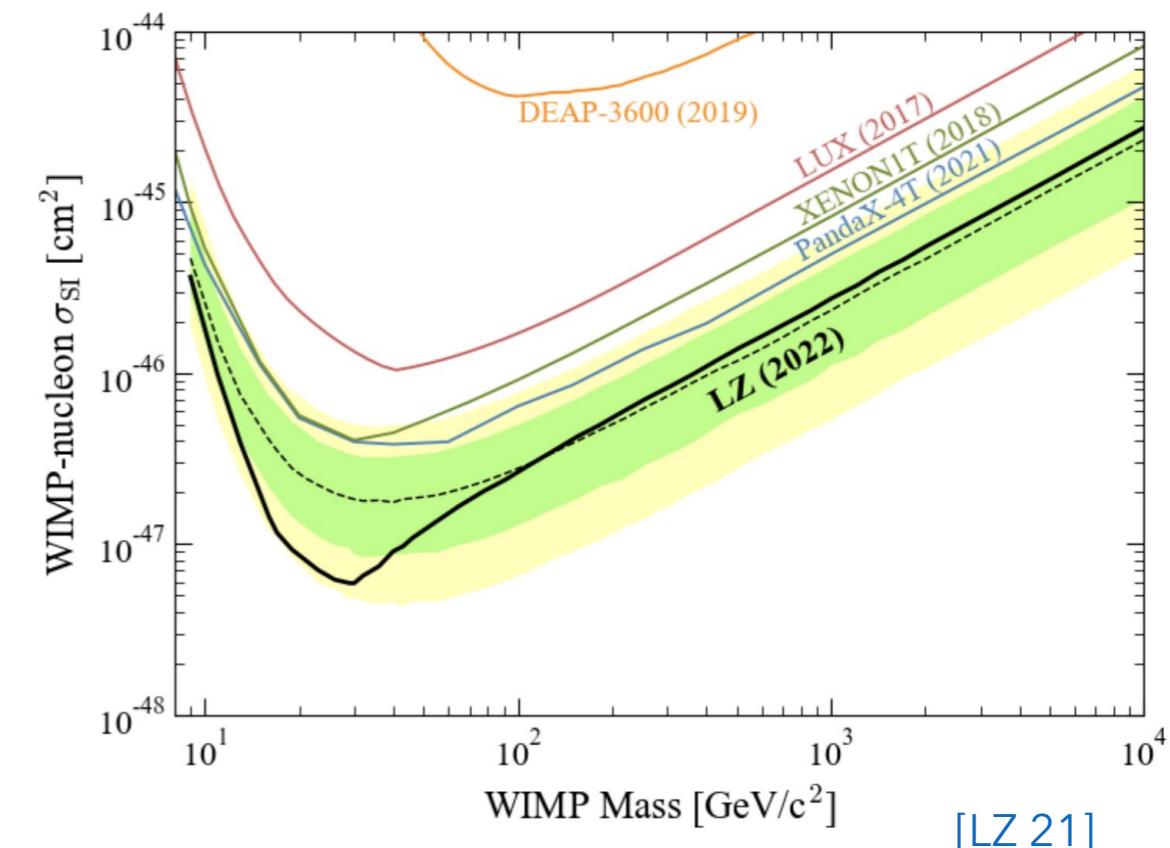
Keep annihilation rate
(DM relic)

$$\langle \sigma v \rangle \sim 10^{-26} \text{ cm}^3/\text{s} \sim 10^{-19} \text{ cm}^2$$



Suppress SI cross section
(direct detection)

$$\sigma_{\text{SI}} \lesssim 10^{-47} \text{ cm}^2$$



seems incompatible \rightarrow needs some trick!

$\mathcal{O}(4)$ symmetric limit

- assuming $\lambda_1 = \lambda_3$,

$$\begin{aligned} V(H, \phi_+, \phi_-) = & m_1^2(|\phi_+|^2 + |\phi_-|^2) - m_{12}^2(|\phi_+|^2 - |\phi_-|^2) \\ & + \frac{\lambda_1}{2} (|\phi_+|^2 + |\phi_-|^2)^2 \\ & - m_H^2 |H|^2 + \lambda_H |H|^4 + \lambda_{H1} |H|^2 (|\phi_+|^2 + |\phi_-|^2) \end{aligned}$$

- emergent $\mathcal{O}(4)$ symmetry (softly broken by m_{12}^2)
→ degeneracy $m_{s_-} = m_{\text{DM}}$

- stable DM: $\phi_- = \frac{1}{\sqrt{2}}(s_- + ia)$

→ equivalent to previous model (Abe-Hamada)