

# Astrophysical Phenomenology of Light Colored Higgs Triplet

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# Astrophysical Phenomenology of Light Colored Higgs Triplet

- Introduction
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- Conclusions

# Introduction: Who is T and why is it interesting

**Triplet**  $(3, 1)_{+\frac{1}{3}}$

GUTs:

$$SM : SU(3) \times SU(2) \times U(1) \subset SU(5) \begin{cases} \subset SO_{10} \\ \dots \\ \subset SU_6 \end{cases}$$

Doublet Triplet Splitting Problem:

$$\begin{pmatrix} T^1 \\ T^2 \\ T^3 \\ H^- \\ H^0 \end{pmatrix} \left. \vphantom{\begin{pmatrix} T^1 \\ T^2 \\ T^3 \\ H^- \\ H^0 \end{pmatrix}} \right\} \text{colored} \Rightarrow T - - - \begin{cases} q, q \\ q, l \end{cases} \Rightarrow \text{mediates proton decay} \Rightarrow \text{Yukawas} \text{ ⚡}$$

# Introduction: Who is T and why is it interesting

## Solutions:

- **Standard approach:**

$$\Rightarrow m_H \text{ light and } m_T = \mathcal{O}(M_{GUT})$$

- **Decoupled T:**  $m_H$  and  $m_T$  light,  $\Gamma \sim m_T \left(\frac{m_{3/2}}{M_{GUT}}\right)^2$

$\Rightarrow \tau > \text{seconds}$ : Potentially interesting phenomenology **Dvali 1992**

$$H_u Q u^c + H_d Q d^c$$

by GUT-breaking couplings vanish



$$\cancel{T_u Q L} + \cancel{T_d d_c u^c}$$

# Mass Range of the Light Triplet

$$m_T \sim [0.3, 7] \text{TeV}$$

Mass calculated in  $MSSM_T$ : T included in MSSM  $\Rightarrow$  modified RGEs

$$\mu \frac{\partial}{\partial \mu} g = -\frac{g^3}{(4\pi)^2} b_i$$

with adjusted coefficients  $b_i$ :

$$b_i^{MSSM_T} = b_i^{MSSM} - T_F(\text{Triplet})$$

Reminder:

$$b_i^{MSSM} = \left(\frac{11}{3} - \frac{2}{3}\right) T_G - \left(\frac{2}{3} + \frac{1}{3}\right) T_F - \left(\frac{1}{3} + \frac{2}{3}\right) T_S$$

$$\text{with } T_G = \begin{cases} 0 & \text{for } U(1) \\ N & \text{for } SU(N) \end{cases} \quad \text{and } T_{S,F} = \begin{cases} \frac{3}{5} Y_{\Phi_i}^2 & \text{for } U(1) \\ \frac{1}{2} & \text{for } SU(N) \end{cases}$$

# Cosmological Production - Abundance

Thermal production with freeze-out at ( $T_* = T_{QCD}$ ):

$$x_* = \frac{m_T}{T_*} \sim [10^3, 10^4]$$

$$\Omega_T = \frac{\tilde{g}_*^{\frac{1}{2}}}{g_* h_{75}^2} x_*^{\frac{3}{2}} \left( \frac{3 \times 10^{-38} \text{cm}^2}{\sigma_*} \right) \sim [10^{-7}, 10^{-6}]$$

- ▶  $\sigma_* = 10^{-26} \text{cm}^2$
- ▶  $\tilde{g}_* = 19.5$  and  $g_* = 3$  are the total number of effectively massless d.o.f. for all particles and for the relic respectively
- ▶  $h_{75} \sim \mathcal{O}(1)$

Too small to constitute 100% of DM.. BUT...

SUSY: consider also boundstates with  $\tilde{T}^c$  antiparticle of its superpartner

# Cosmological Production - T-Mesons

T-Meson: boundstate of Triplet with another colored particle:  $Td$ ,  $T\tilde{T}^c$

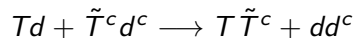
- **$T\tilde{T}^c$ -production:** cf. Redi et al. 2018

At  $T \sim T_{QCD}$ :

$T$  and  $\tilde{T}$  confine by picking up a d-quark ( $\Omega_{Td} \sim 1$ )  
or with each other (small abundance  $\Omega_{T\tilde{T}^c} \sim 10^{-14}$ ).

Later, at  $T \lesssim T_{QCD}$ :

another process is taking place, yielding a much higher abundance of  
 $T\tilde{T}^c$  ( $\Omega_{T\tilde{T}^c} \sim 10^{-1} = \Omega_{DM}$ ):



- **$T\tilde{T}^c$ -decay:**



# Cosmological Production - T-Mesons

$\sigma_{T\bar{T}c} \sim R_{Bohr}^2$ ?  $R_{Bohr} \sim \frac{1}{\alpha_{QCD} m_T}$  cf. Redi et al. 2018

here: both constituents charged  $\rightarrow$  el. dipole moment  $\mathcal{D} = \frac{e}{m_T}$

$\Rightarrow \sigma \sim 6.4 \times 10^{-32} Z^2 (\mathcal{D} \times 10^{17})^2 \text{cm}^2$  Sigurdson et al. 2004

equating this with the extrapolated linear  $\sigma(m)$ -function of:

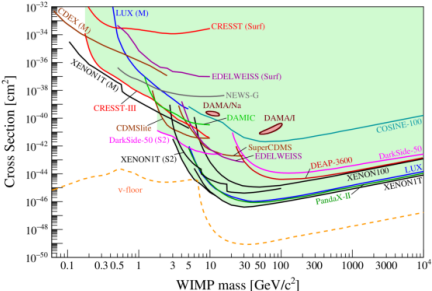


Fig.: Direct detection bounds on dark matter cross sections. Billard et al. 2022



# Cosmological Production - T-Mesons

bound on Triplet-Mass:  $m_T > 10^8 \text{ GeV} \ll 10^{16} \text{ GeV}$

Larger than calculated in  $MSSM_T$ : We don't have to stick to the minimal case - mass is a free parameter

$\Rightarrow$  Appropriate DM-candidate!

# Cosmological Production - DM-candidate:

- ▶ not yet excluded by experimental bounds
- ▶ long-lived
- $T\tilde{T}^c$  stable?
  - ▶ **no** if  $m_T \neq m_{\tilde{T}}$ ,  
i.e model-dependent LSP (e.g. Neutralino),  
 $T$  and  $\tilde{T}$  are N<sub>ext</sub>-to-LSP and NNLSP: LSP-DM
  - ▶ **yes** for  $m_T = m_{\tilde{T}}$ :  $T\tilde{T}^c$ -DM if  $m_T > 10^8$  GeV
- $Td$  stable?
  - ▶  $T$ : yes (assumption in our project)
  - ▶  $d$ : could decay via weak i.a.:  $Td \rightarrow Tu + e^- + \nu_e$   
lifetime:  $\tau_T \sim 15\text{min} \sim 10^3\text{sec}$

No DM but interesting collider phenomenology (another project..)

# $Td$ -Mesons as heavy Neutrons

$Td$ -Meson: same quantum numbers as neutron but

$$\frac{10^2 \text{ GeV}}{1 \text{ GeV}} \lesssim \frac{m_{Td}}{m_{neutron}} \lesssim \frac{10^{16} \text{ GeV}}{1 \text{ GeV}}$$

$Td$  is a BBN-particle cf. Kusakabe et al. 2009

→ if it exists around us - can it be inside the atoms?

Or will it be stuck inside center of the stars/planets

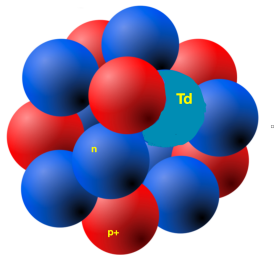


Fig.: Nucleus with  $Td$

# Conclusions

- Colored Higgs Triplet is universally predicted by Grand Unification
- This particle can be long-lived and therefore has an interesting phenomenology and cosmology

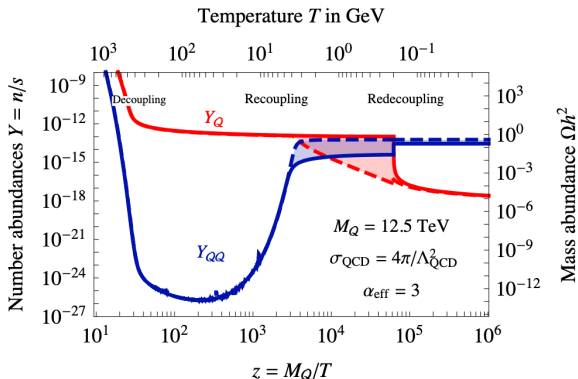
cf. Dvali 1992; Dvali 1995; Dvali,Shifman 1997; Dvali,AJ 2020

- Colored Higgs Triplet is not yet excluded by experimental bounds - lead to serious Dark Matter candidates (T-Meson or LSP-DM)
- Open questions: Td-Meson is BBN particle. Can the Td-Meson hide inside our daily life matter? Or accumulate inside the core of our planet?

A photograph of a bright blue sky filled with scattered white clouds. The word "THANKS" is written across the center of the image in a yellow, hand-drawn, cursive font. The letters are slightly tilted upwards from left to right. The background consists of a clear blue sky with various cloud formations, including some wispy clouds and larger, more puffy ones.

THANKS

# BackUp: Abundance of $Td$ - vs. $T\tilde{T}^c$ -Meson



**Fig.:** Cosmological evolution of the abundances of Q states and of QQ DM states for  $M_Q = 12.5 \text{ TeV}$ . The uncertain phase at  $T \geq \Lambda_{\text{QCD}}$  negligibly affects the final relic abundances: the dashed curves assume non-perturbative effects before confinement estimated as  $\sigma = \sigma_{\text{QCD}} \frac{\Lambda_{\text{QCD}}^2}{T}$ ; the solid curves neglect such effects. The mass abundance on the right axis is computed assuming QQ particles with mass  $2M_Q$ . [Redi et al. 2018](#)

# BackUp: SO(10)-example for Decoupled Triplet Dvali 1992

To generate quark and lepton masses:

- standard version of SO(10):  $g_{10} \otimes 16 \otimes 16$  ⚡
- pair of hypothetical heavy supermultiplets  $144 + \overline{144}$  which couple to 45, 10 and 16:

$$W = g_{10} 16 \gamma_i \overline{144}_k \otimes 45_{ik} + M_{144} \overline{144}_k \otimes 144_k + g'_{10} 16 \otimes 144_k \otimes 10_k$$

with  $M_{144} \sim M_{GUT}$

Effective operator for energies below  $M_{GUT}$ :

$$\frac{gg'}{M_{144}} 45_{ik} 10_k \otimes 16 \gamma_i 16 \Rightarrow \text{Yukawa with Higgs Triplet suppressed}$$

