

Measuring lepton number violation at colliders

Jan Hajer

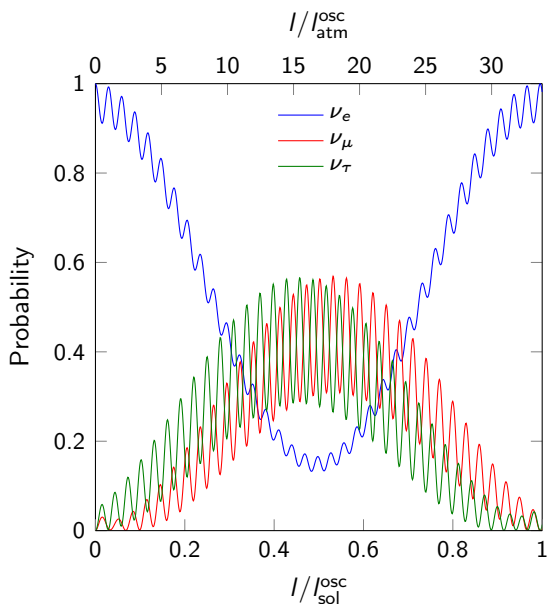
Centro de Física Teórica de Partículas, Instituto Superior Técnico, Universidade de Lisboa

Planck 2024

The 26th international conference from the Planck scale to the electroweak scale

Neutrino flavour oscillations and seesaw mechanism

Observed neutrino flavour oscillations



Can be explained by
at least to massive neutrinos

Right-handed Majorana neutrino N

$$\mathcal{L}_m = \begin{pmatrix} \vec{\nu} \\ N \end{pmatrix}^t \begin{pmatrix} 0 & \vec{m}_D \\ \vec{m}_D^T & m_M \end{pmatrix} \begin{pmatrix} \vec{\nu} \\ N \end{pmatrix}$$

Interaction governed by mixing parameter

$$\vec{\theta} = \frac{\vec{m}_D}{m_M} \quad \begin{array}{l} \text{Dirac mass} \\ \text{Majorana mass} \end{array}$$

Neutrino masses

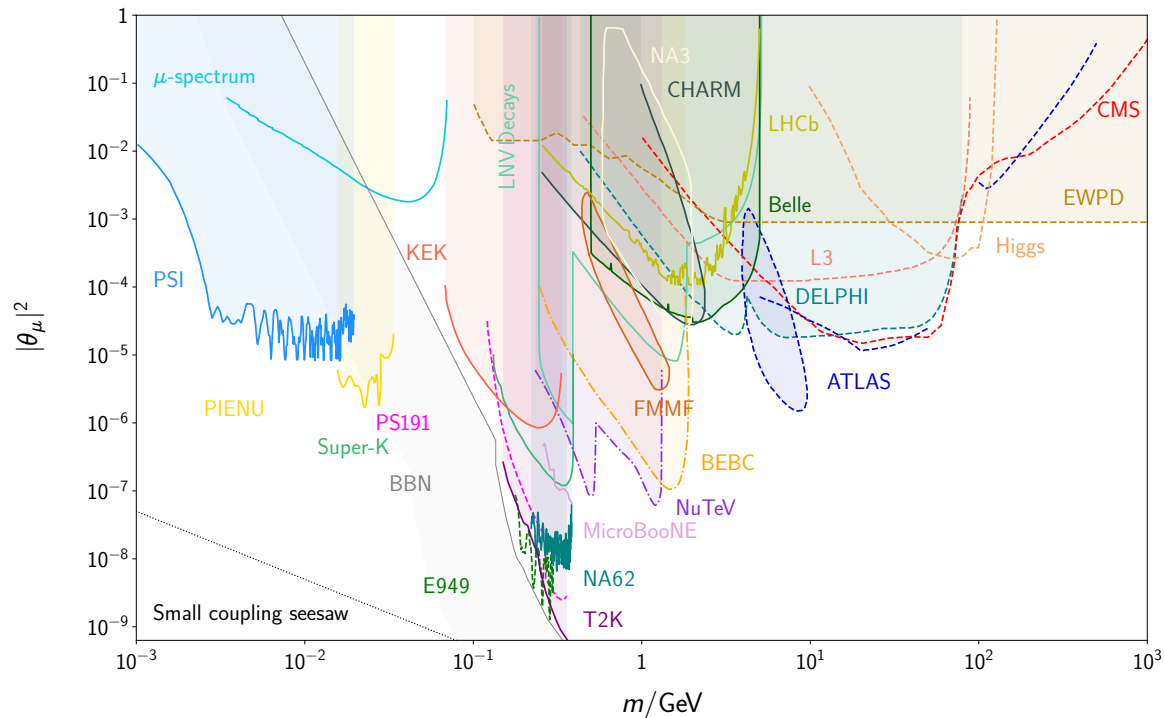
$$m_\nu = \frac{\vec{m}_D \vec{m}_D^T}{m_M} = \frac{\vec{\theta} \vec{\theta}^T}{m_M}$$

Tiny neutrino masses are ensured for

- large m_M High scale seesaw
- small \vec{m}_D Small coupling seesaw

Sterile neutrinos/Heavy neutral leptons (HNLs)

- Inaccessibly heavy or
- Tiny interactions



Inaccessible: ■ Small coupling seesaw ■ High scale seesaw (at the GUT scale)

Symmetry-protected low-scale seesaw

Lepton number $L = n_\ell - n_{\bar{\ell}}$

Accidentally conserved in the Standard Model

Generalisation: 'Lepton number'-like symmetry

e.g. $U(1)_L$	$\vec{\nu}$	N_1	N_2
with charges	L	$+1$	-1

Symmetry breaking in the mass matrix

$$\mathcal{L}_m = \begin{pmatrix} \vec{\nu} \\ N_1 \\ N_2 \end{pmatrix}^t \begin{pmatrix} 0 & \vec{m}_D & \vec{\mu}_D \\ \vec{m}_D^T & \mu'_M & m_M \\ \vec{\mu}_D^T & m_M & \mu_M \end{pmatrix} \begin{pmatrix} \vec{\nu} \\ N_1 \\ N_2 \end{pmatrix}$$

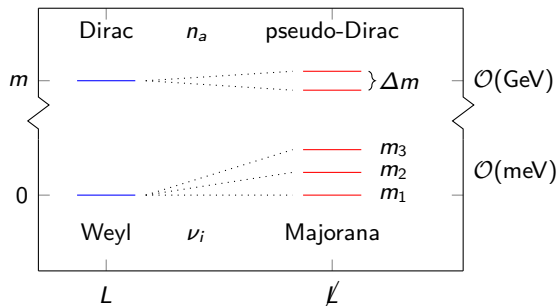
Symmetry L conserved

- Three massless neutrinos
 - Single Dirac heavy neutrino
- Corresponds to two degenerate Majoranas

Small symmetry breaking $\not\ll$

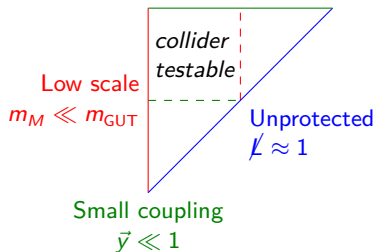
- Light neutrino masses $m_\nu \propto \not\ll$
- Heavy neutrino mass splitting $\Delta m \propto \not\ll$

Breaking induced neutrino mass splitting



Viable seesaw limits

Symmetry protected $\not\ll \ll 1$ Large coupling $\vec{y} \approx 1$ High scale $m_M \approx m_{\text{GUT}}$



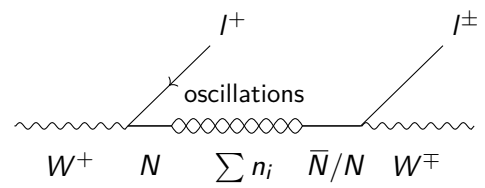
Heavy neutrino-antineutrino oscillations ($N\bar{N}$ Os)

[2210.10738]

Oscillations between events that have

- Lepton number conservation (LNC) $I^\pm I^\mp$
- Lepton number violation (LNV) $I^\pm I^\pm$

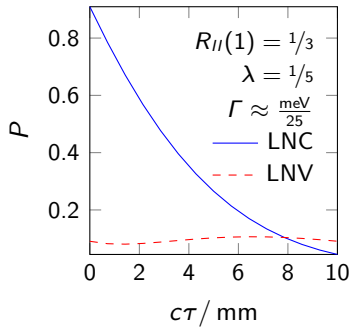
Oscillating mass eigenstates n_i



Oscillation frequency governed by Δm

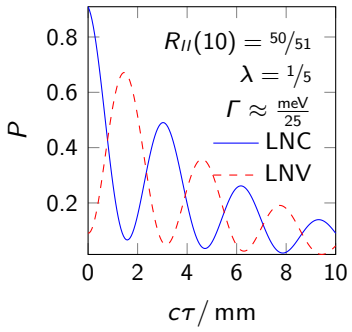
$$P_{\text{osc}}^{\text{LNC/LNV}}(\tau) = \frac{1 \pm \cos(\Delta m \tau)}{2}$$

Slow oscillation



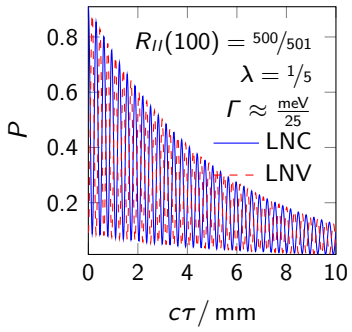
- Mostly LNC
- 'Dirac BM'-like

Intermediate oscillation

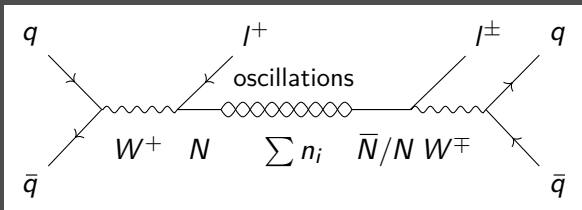


- Potentially resolvable
- $0 \leq R_{II} = \frac{P_{\text{LNC}}}{P_{\text{LNV}}} \leq 1$

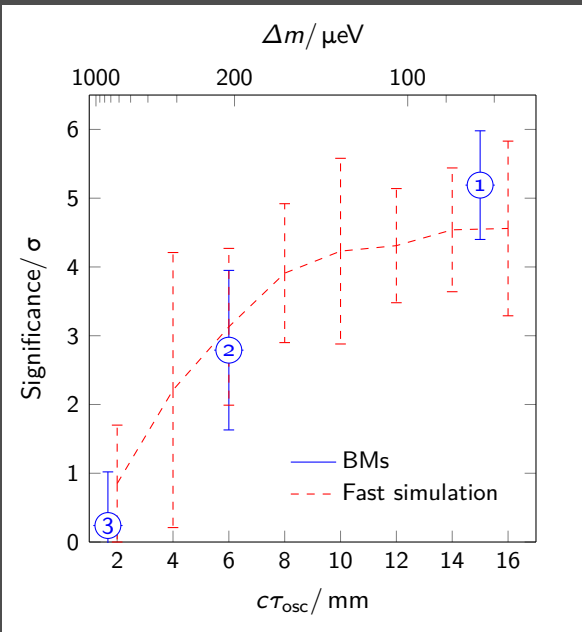
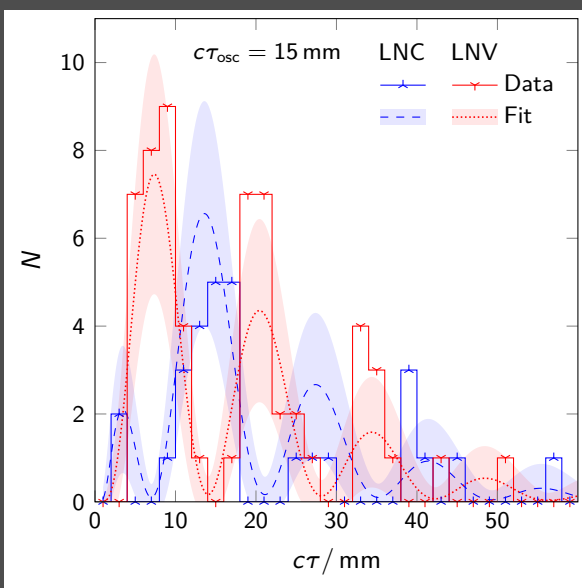
Fast oscillation



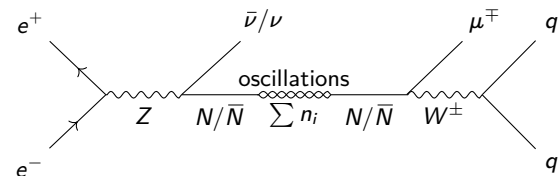
- Unresolvable
- LNV as frequent as LNC
- 'Majorana BM'-like



LNV is measured by comparing the charges of the two leptons



Single charged lepton



Measurement

- LNV cannot be measured using two charges
- One can still measure angular distributions

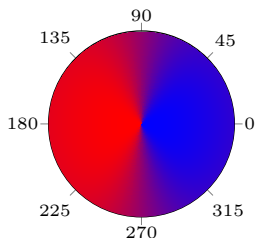
Angular dependent probability

$$P_{l^\mp}(\cos\theta, \tau) := \frac{1}{\sigma} \frac{d\sigma(\cos\theta)}{d\cos\theta} P_{\text{osc}}^{\text{LNC/LNV}}(\tau)$$

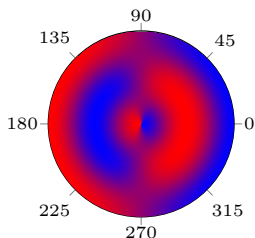
Probability of measuring charged leptons

- linked to forward backward asymmetry of neutrino production (see 'Dirac BM'-like)
- l^- from non-oscillating N or from oscillating \bar{N} (similar for l^+)

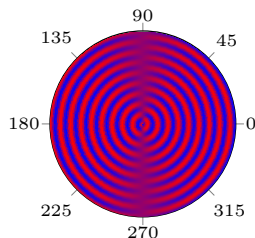
'Dirac BM'-like



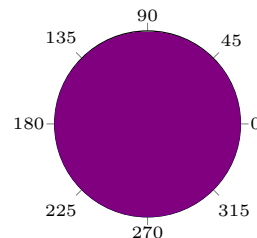
Slow oscillation



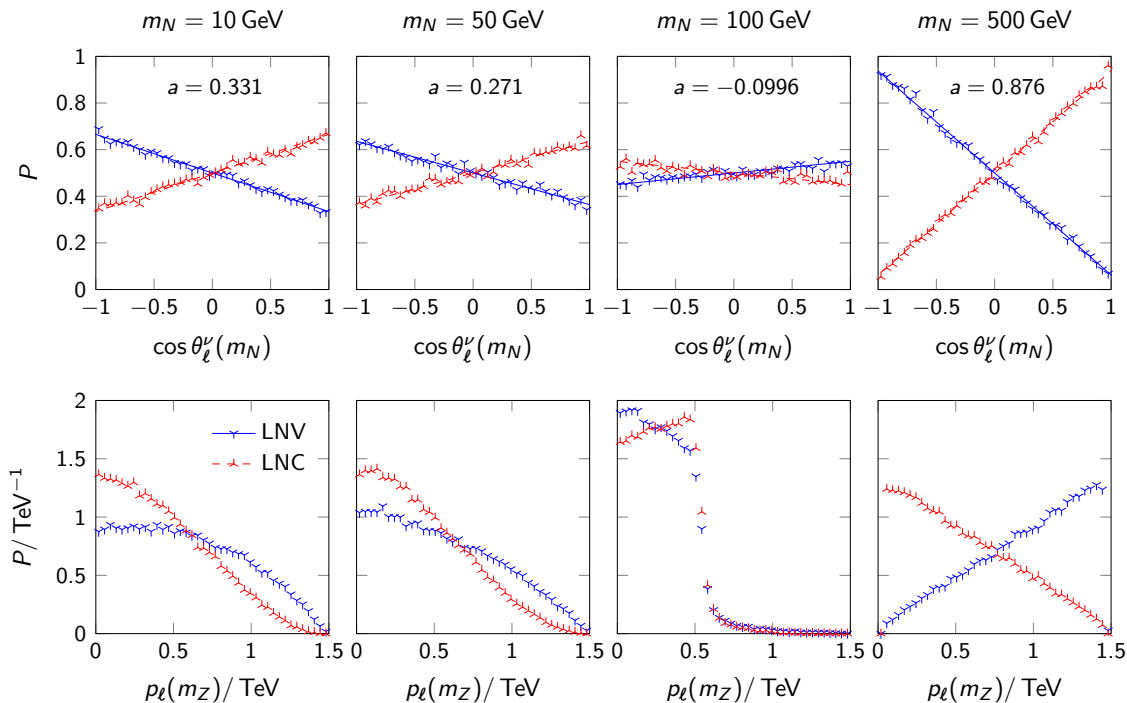
Fast oscillation



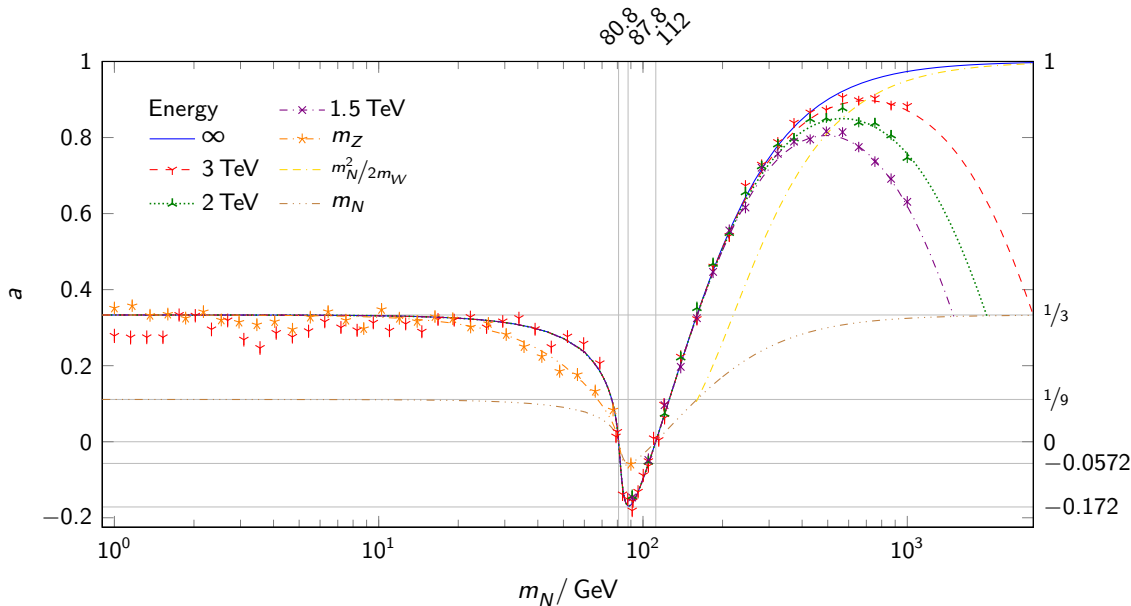
'Majorana BM'-like



$$d\sigma_{c/\nu}(\cos\alpha) \propto (\sigma_0 \mp \sigma_1 \cos\alpha) d\cos\alpha, \quad P_{c/\nu}^M(\cos\alpha) = (1 \mp a \cos\alpha)/2, \quad a = \sigma_1/\sigma_0,$$



Opening angle asymmetry is sensitive to LNV



For the achievable sensitivity see talk by Bruno Oliveira

at Thursday after lunch in parallel session PII.4 in Room QA 1.4

Inadequate frameworks for oscillating relativistic particles

- Quantum mechanics
- Plane-wave QFT

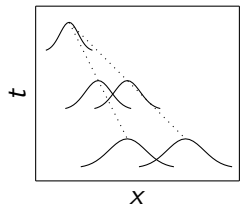
QFT with external wave packets

- Gaussian wave packets with width σ
- External widths are experiment depended parameters
- Internal widths are calculated

Transition amplitude in QFT with external wave packets Φ

$$\mathcal{A}(x) = \left\langle \Phi(x'') \left| \mathcal{T} \exp \left[-i \int \mathcal{H}(x') d^4x' \right] - \mathbb{1} \right| \Phi(x') \right\rangle$$

Decoherence



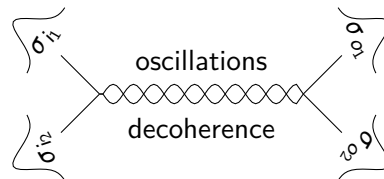
Result can be expressed with effective damping parameter λ

Damped oscillations

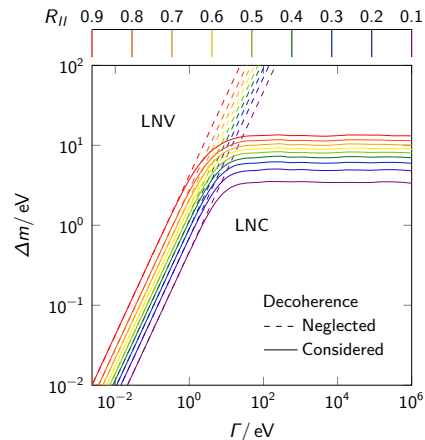
$$P_{\text{osc}}^{\text{LNC/LNV}}(\tau) = \frac{1 \pm \cos(\Delta m \tau) e^{-\lambda}}{2}$$

LNV can be drastically enhanced

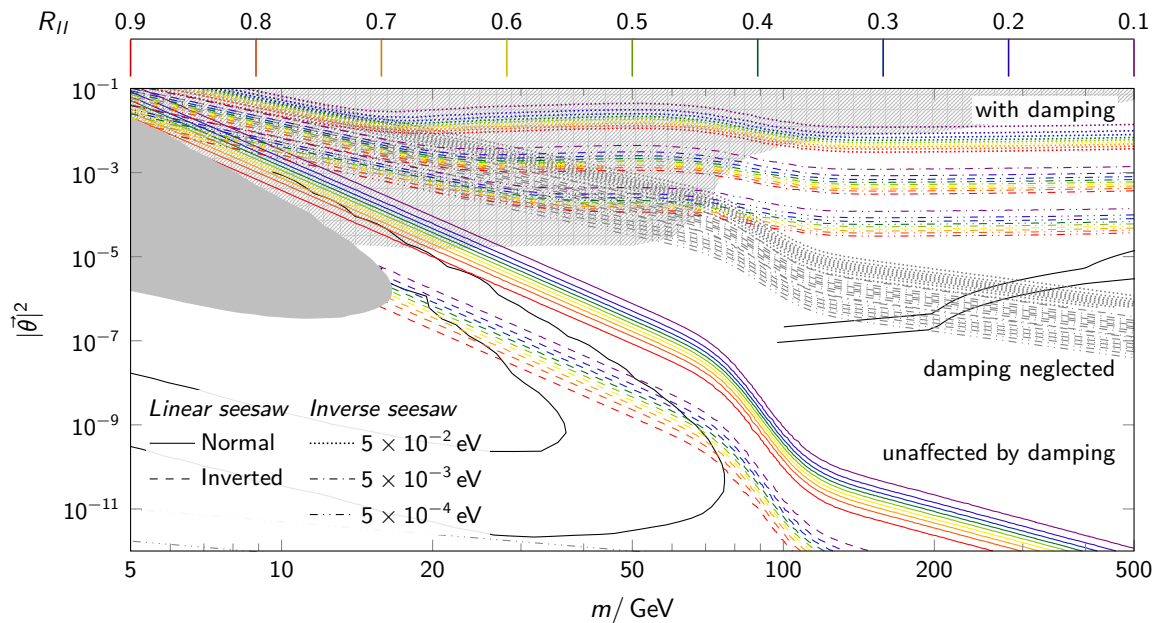
Width of external wave packets σ



Impact on $N\bar{N}O$ s



Decoherence at the LHC



Minimal linear seesaw

Not affected by decoherence

Inverse seesaw

LNV significantly increased

- Collider testable Type I seesaw models predict pseudo-Dirac HNLs
- Collider testable single Majorana or Dirac HNLs cannot explain neutrino masses
- Pseudo-Dirac HNLs can oscillate between LNC and LNV events
- In the absence of countable LNV these $N\bar{N}O$ s are the only unambiguous measurement of LNV
- Theses $N\bar{N}O$ s are detectable at the HL-LHC and future lepton colliders
- Decoherence of $N\bar{N}O$ s are extremely relevant

References

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