

Ultrahigh-frequency GW backgrounds from “Cosmic Strings”

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with

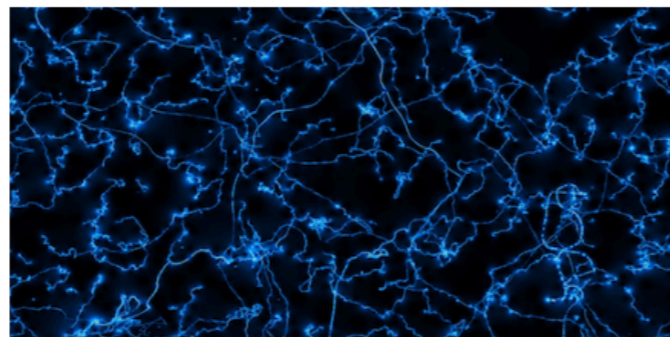
Géraldine Servant (U. Hamburg and DESY)

based on

Phys. Rev. D 109, 103538 [[arXiv: 2312.09281](https://arxiv.org/abs/2312.09281)]

Featured in Physics

Editors' Suggestion



C. Ringeval/Louvain University

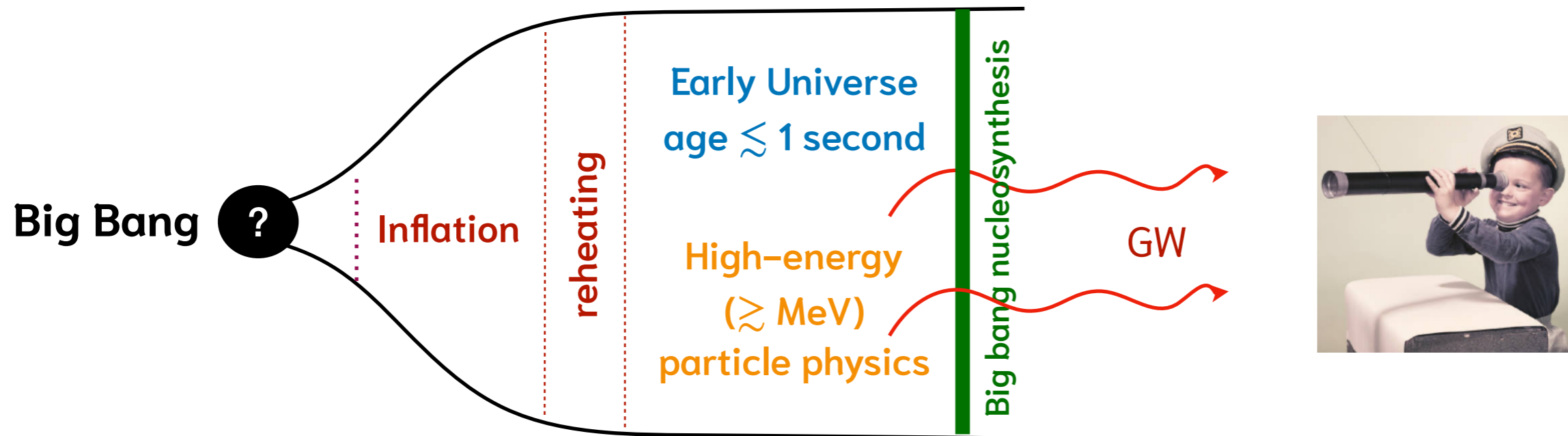
PhysiCS NEWS AND COMMENTARY

Cosmic Strings' Imprints in High-Frequency Gravitational Waves

May 23, 2024

Gravitational waves (GWs) as tools

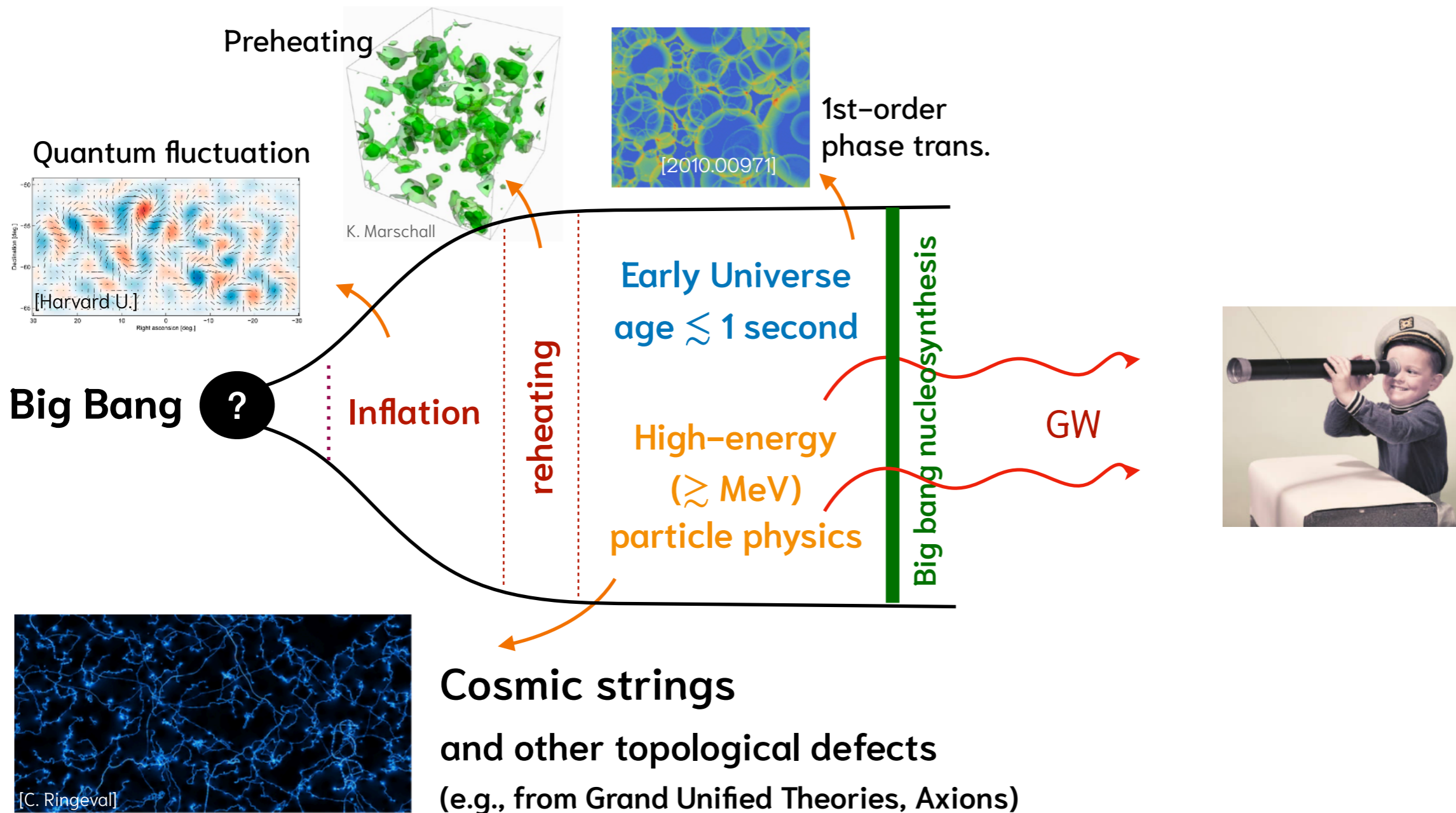
for charting **the early-Universe cosmology** and **high-energy particle physics**.



(See also e.g., talks by Stefan Antusch, Stephen King, and many others.)

Gravitational waves (GWs) as tools for charting the early-Universe cosmology and high-energy particle physics.

Many GW sources...(mostly BSM physics)

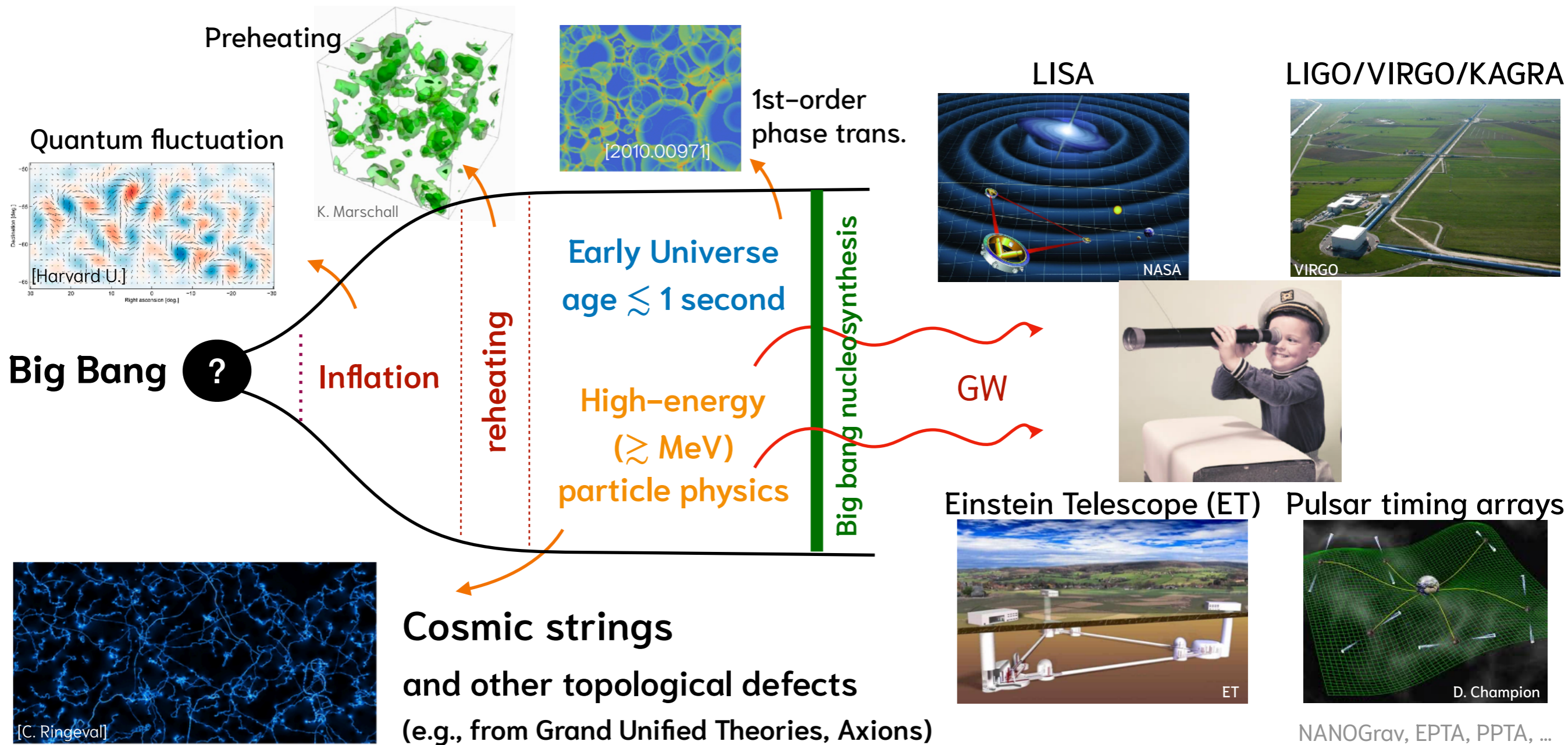


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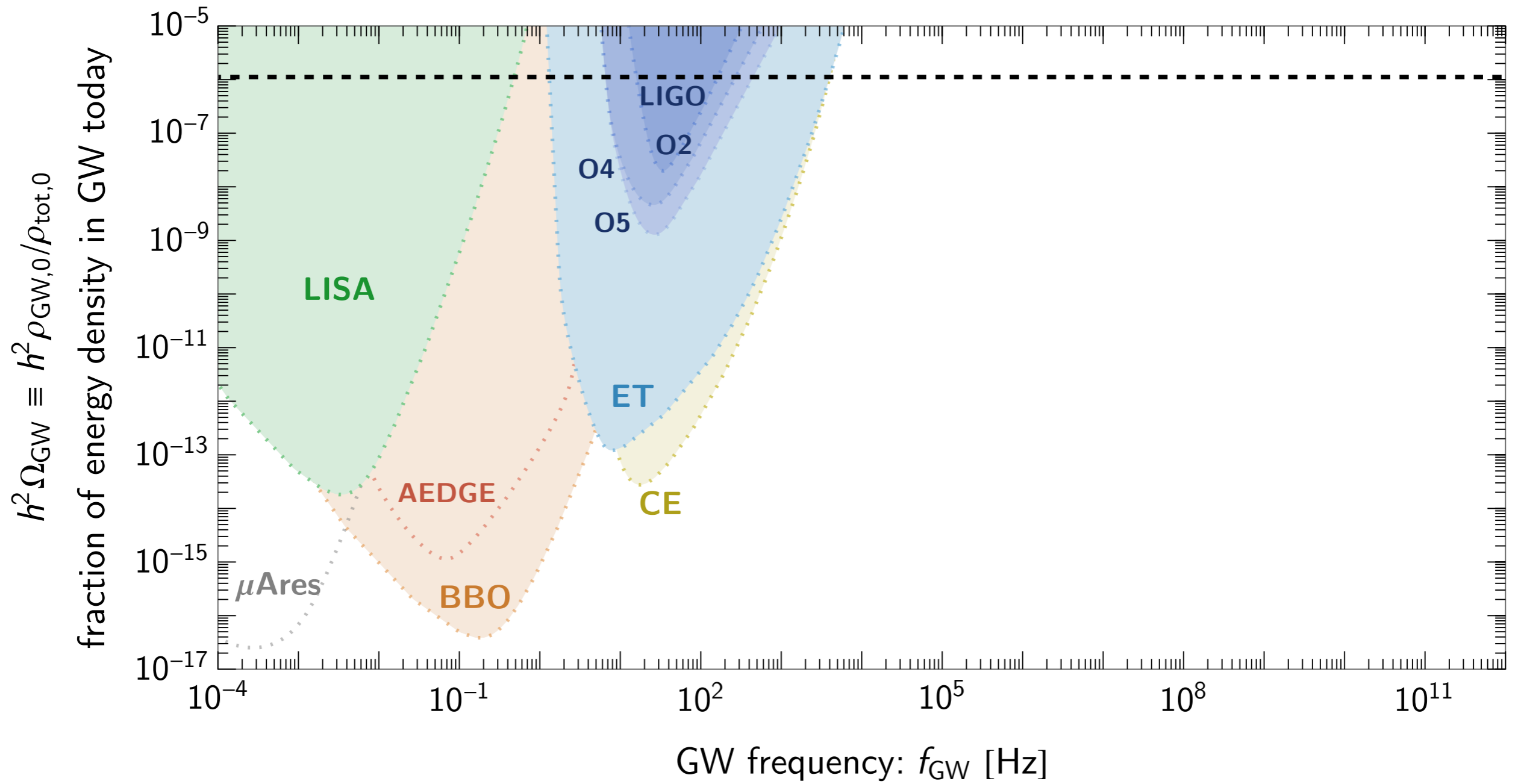
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Many ways to detect...

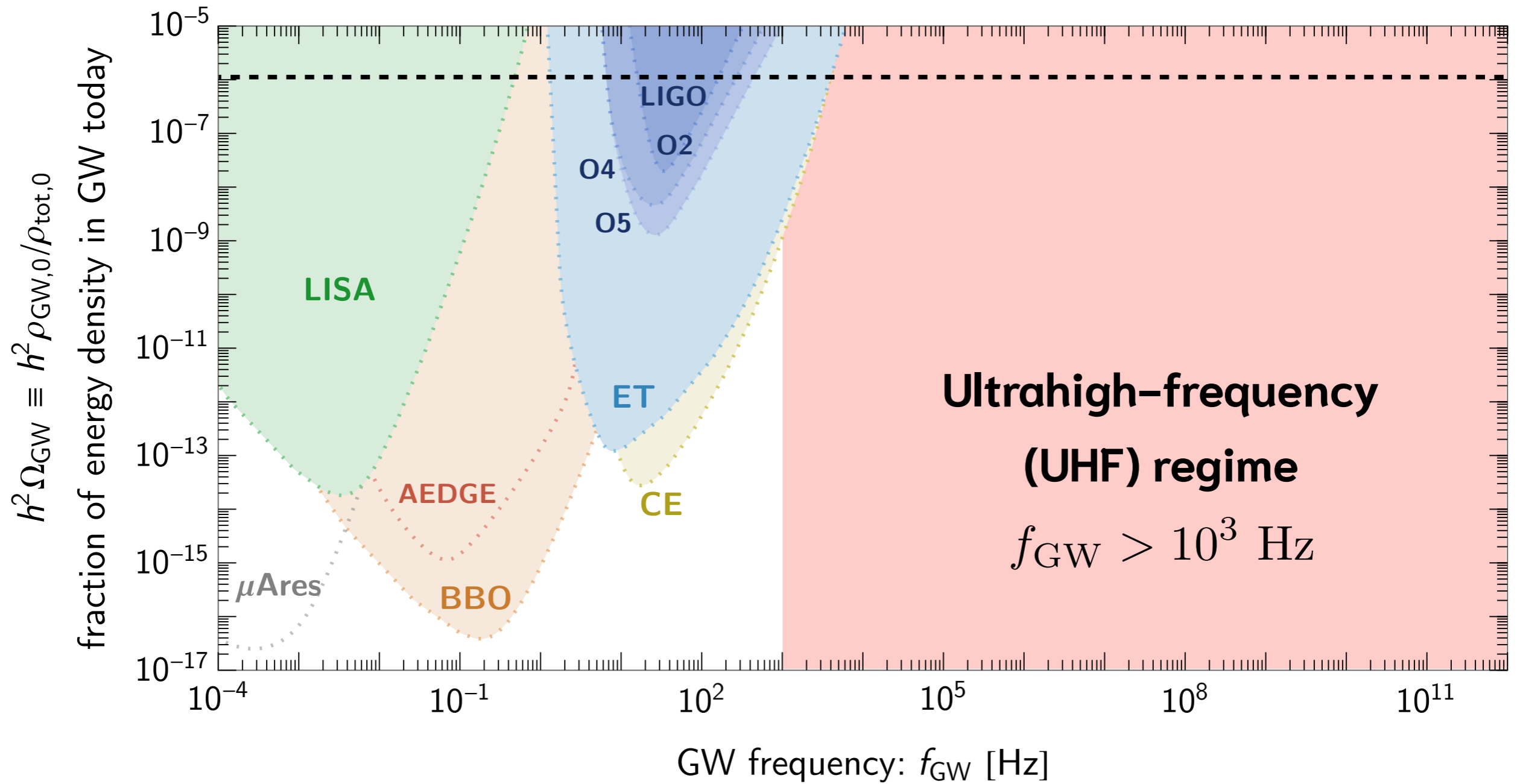


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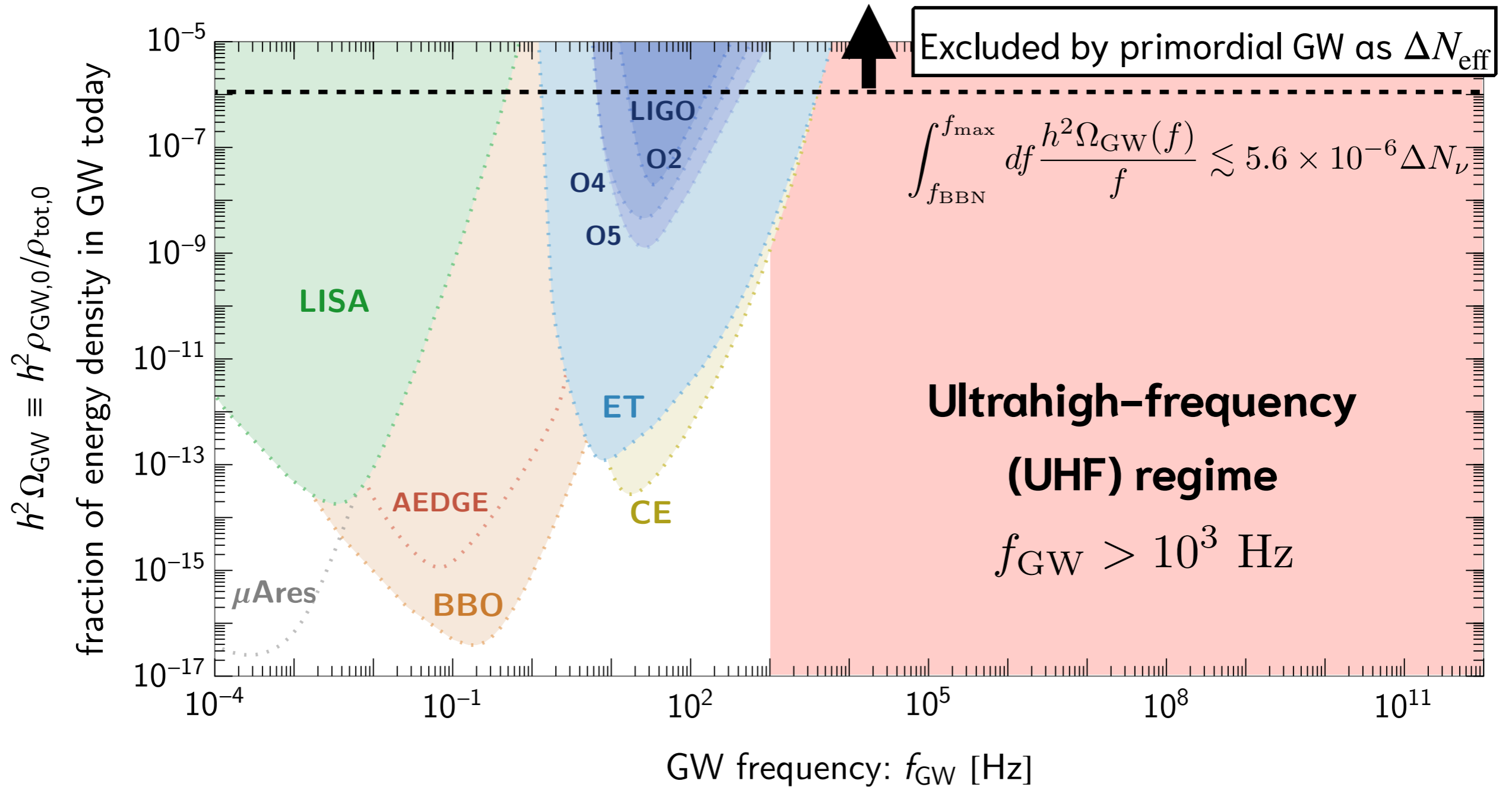
Gravitational-Wave Landscape



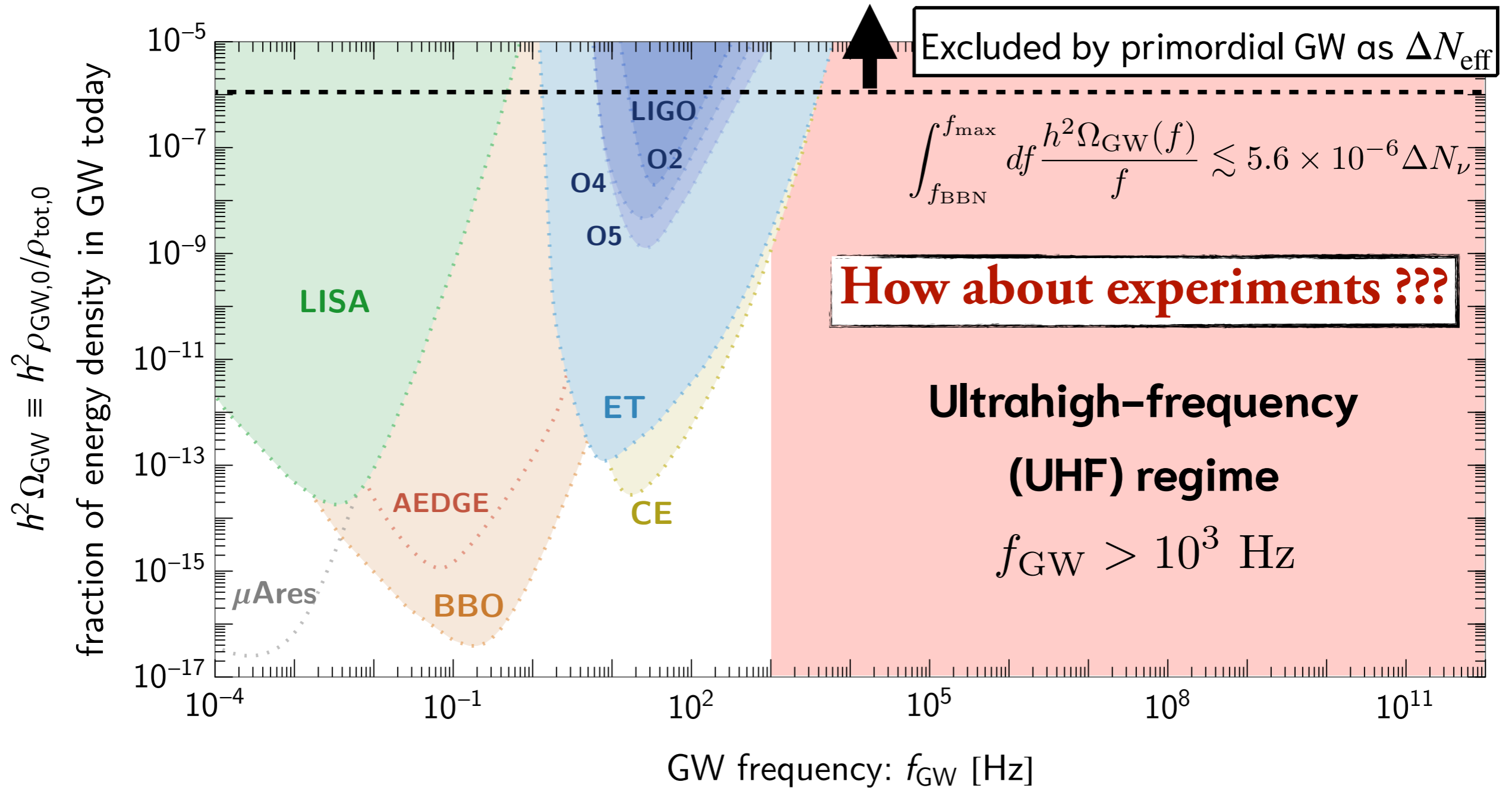
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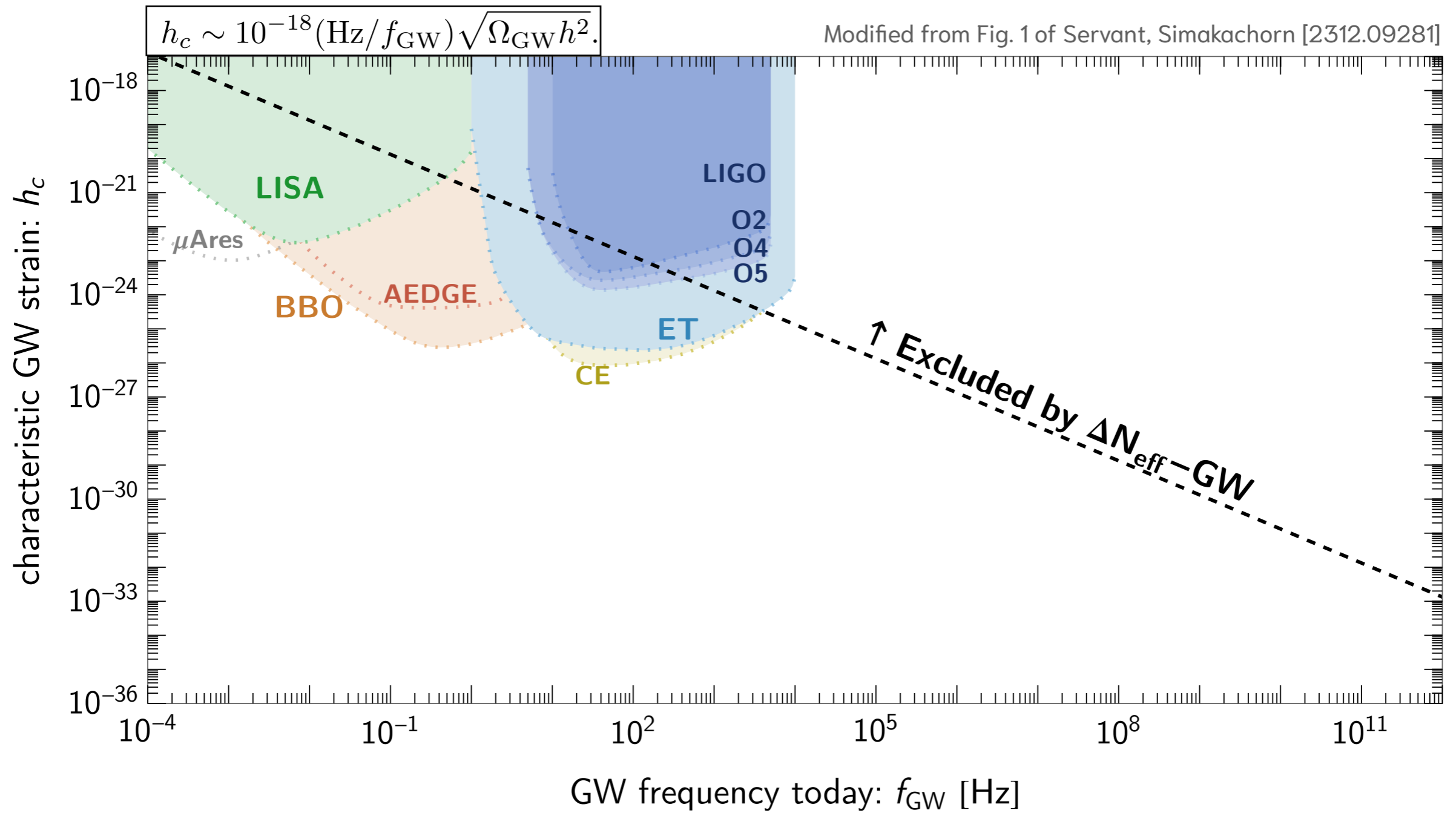
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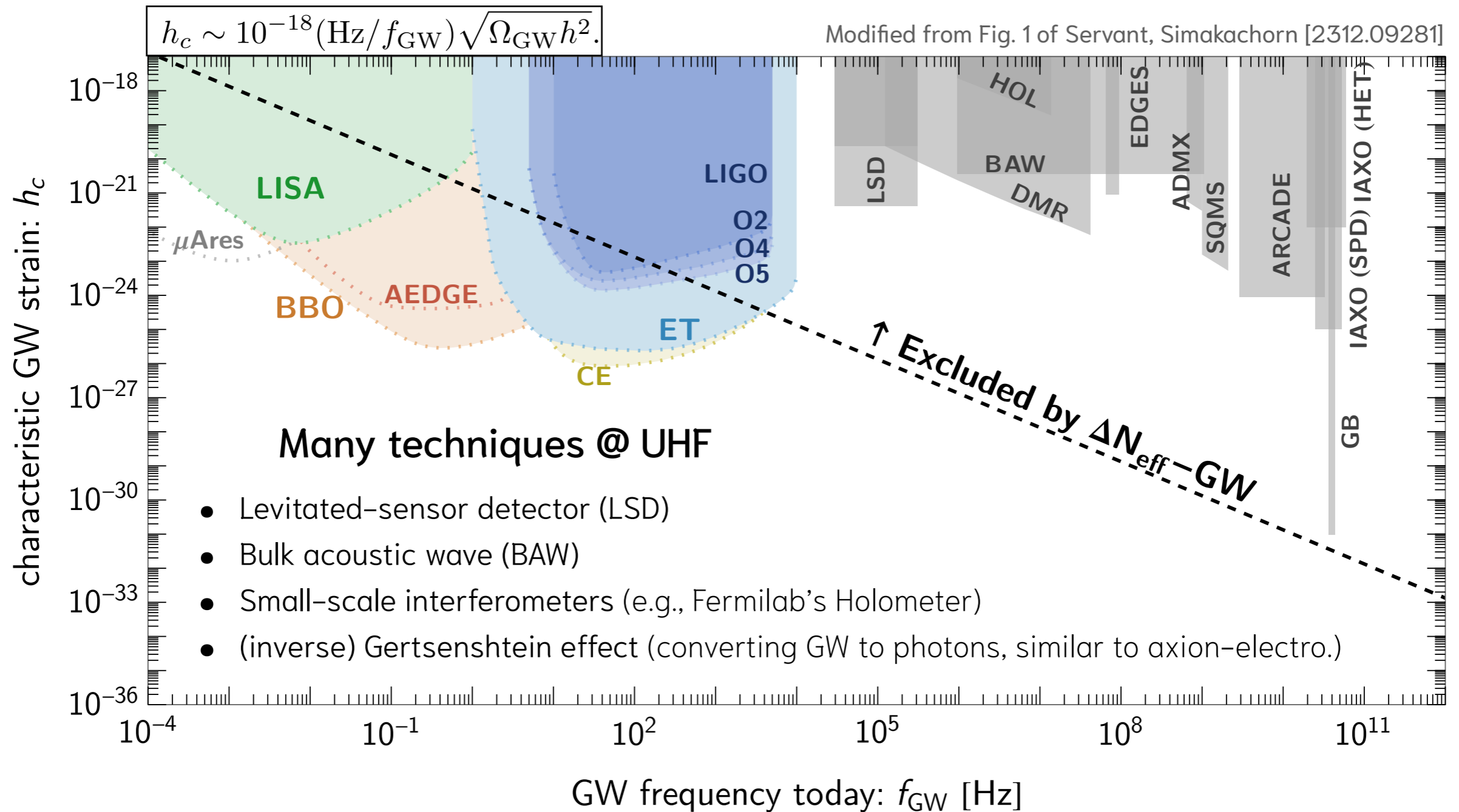
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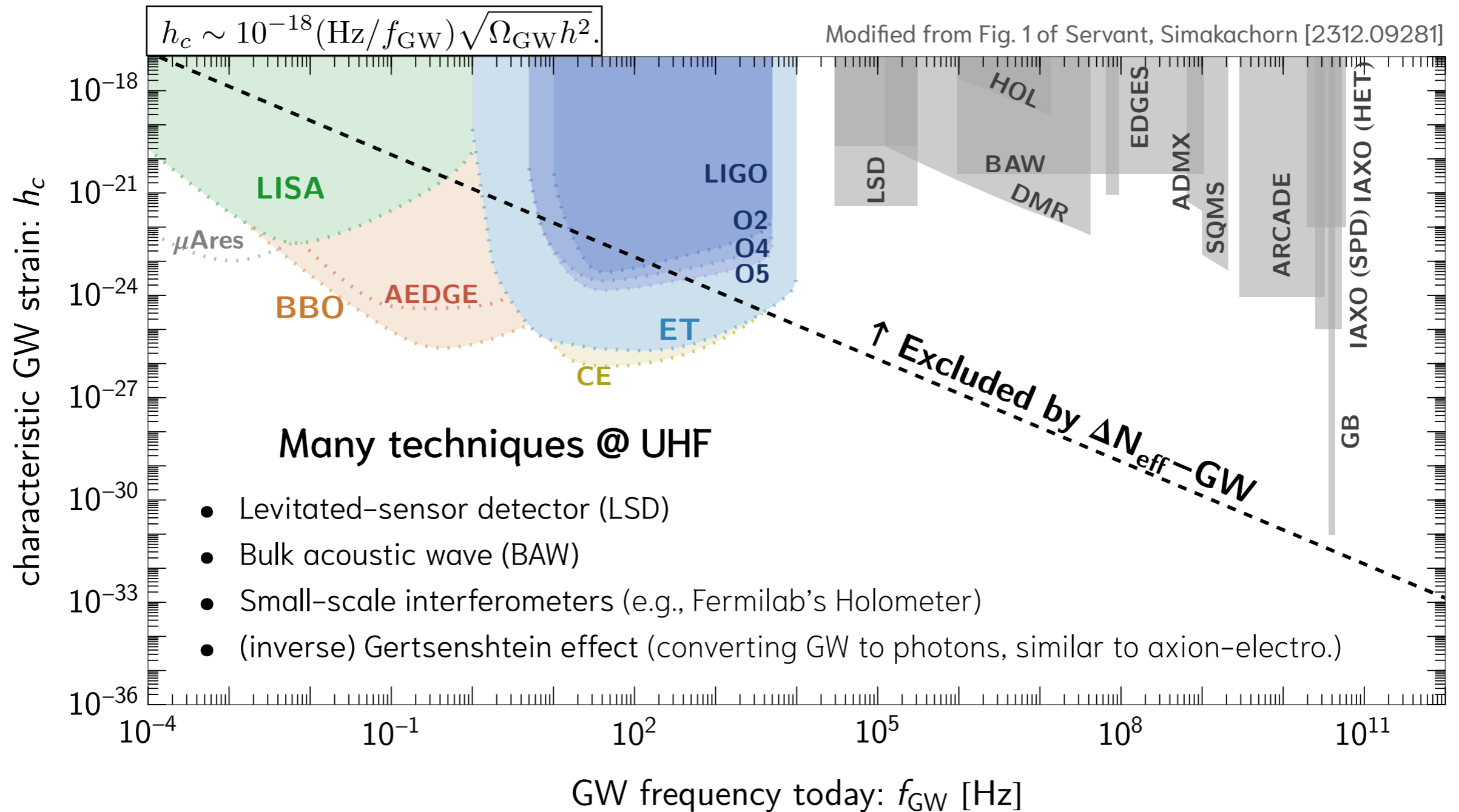
Current status of UHF GW experiments



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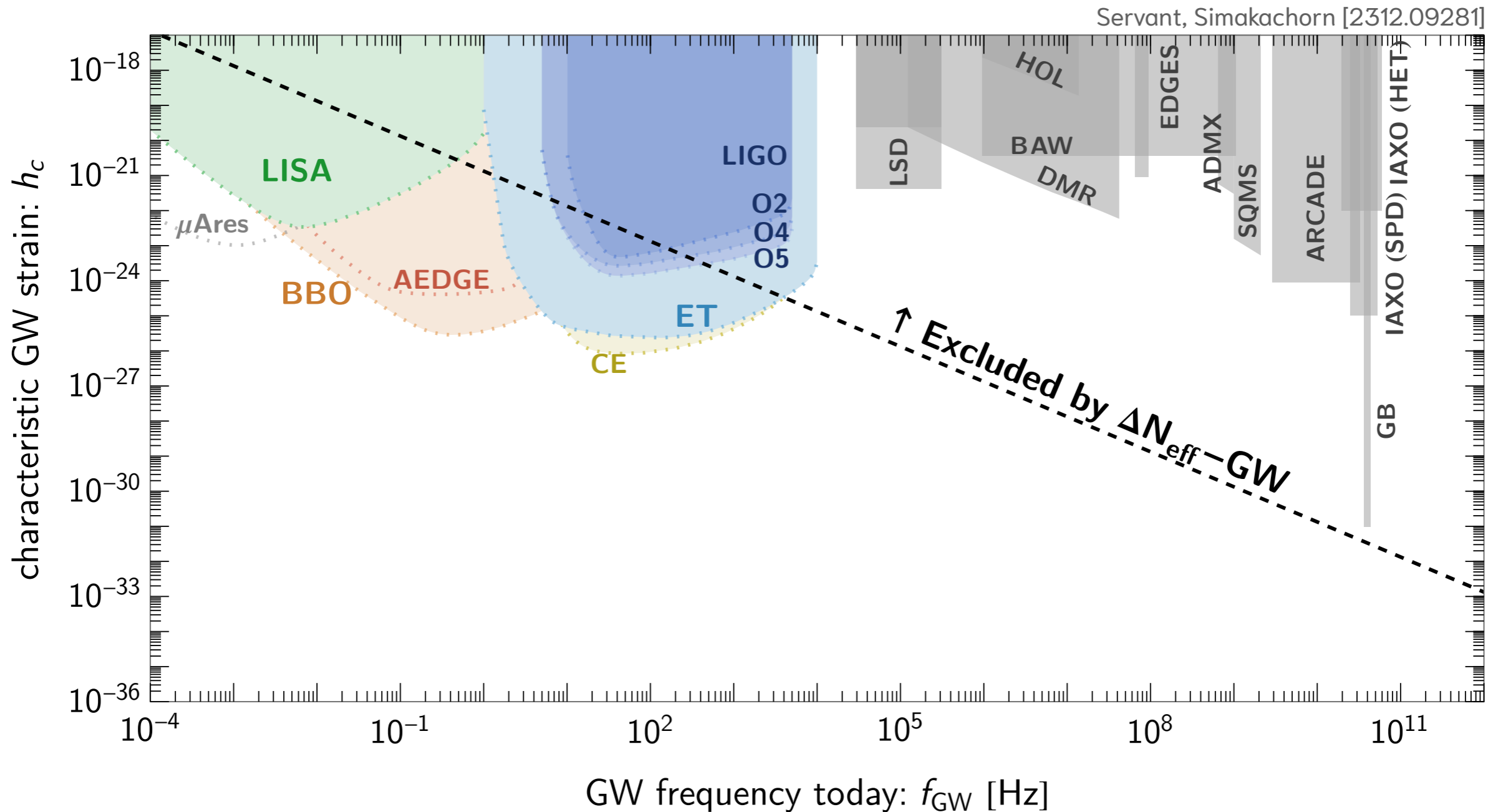


The ΔN_{eff} -GW bound is so far the strongest, but one should keep thinking because...

UHF regime = A golden window for new physics

Reasons:

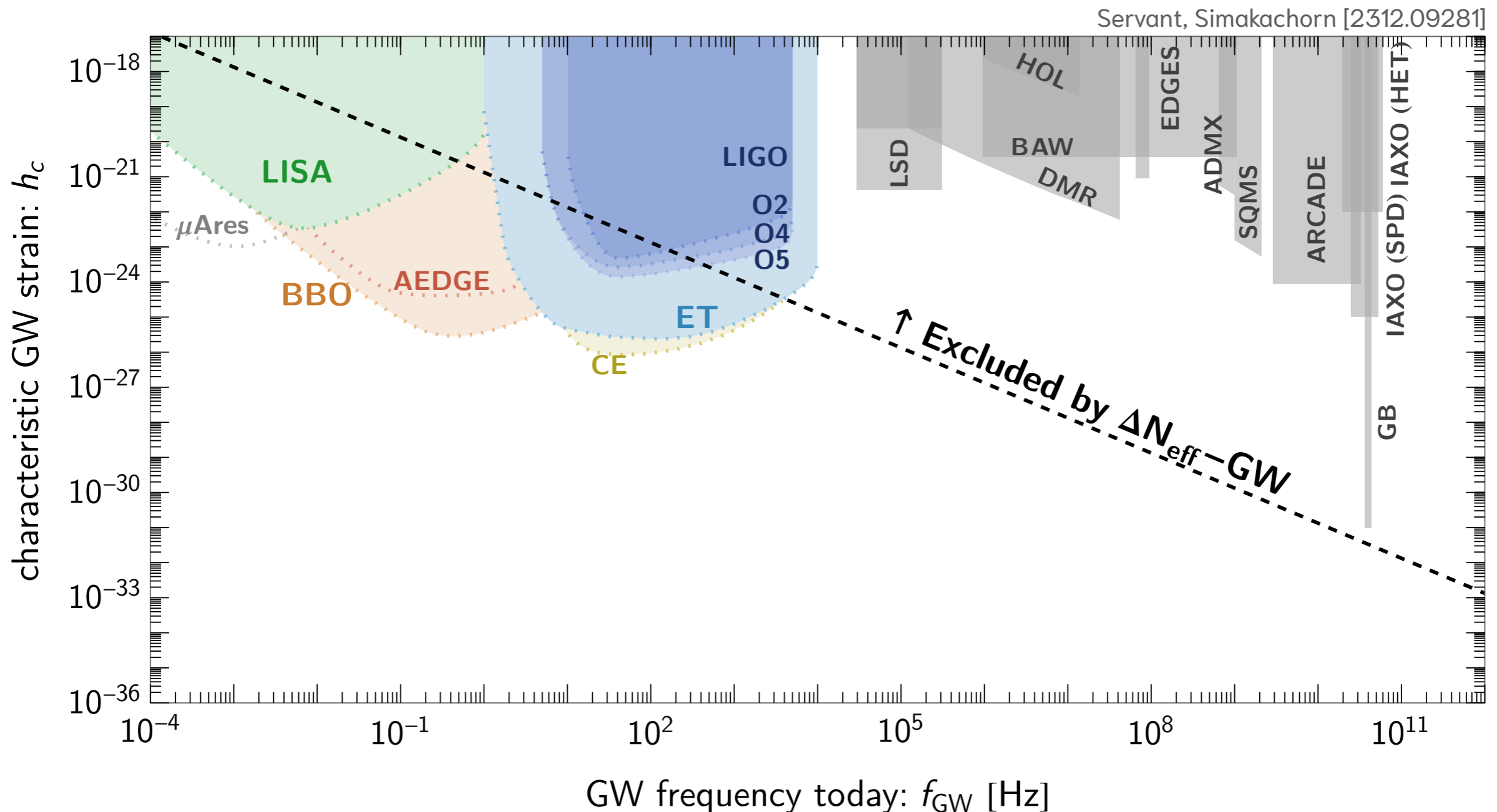
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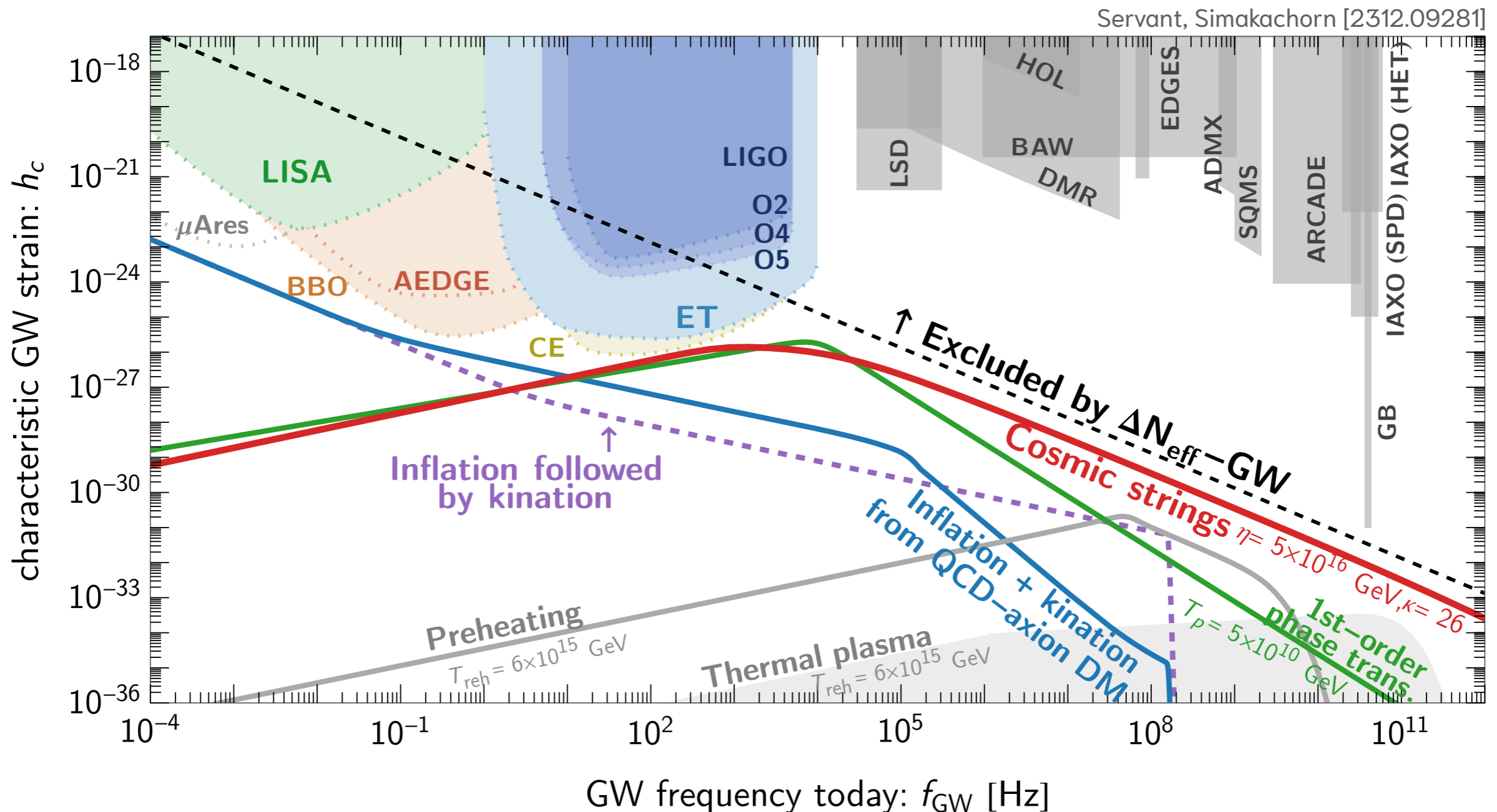


- Signals from well-motivated sources, $f_{\text{GW}} = f_{\text{GW},e} \left(\frac{a_e}{a_0} \right) \simeq 1 \text{ kHz} \left[\frac{H^{-1}(T)}{\lambda_{\text{GW}}} \right] \left(\frac{T}{10^{10} \text{ GeV}} \right)$,
both SM and BSM @ ultrahigh energies
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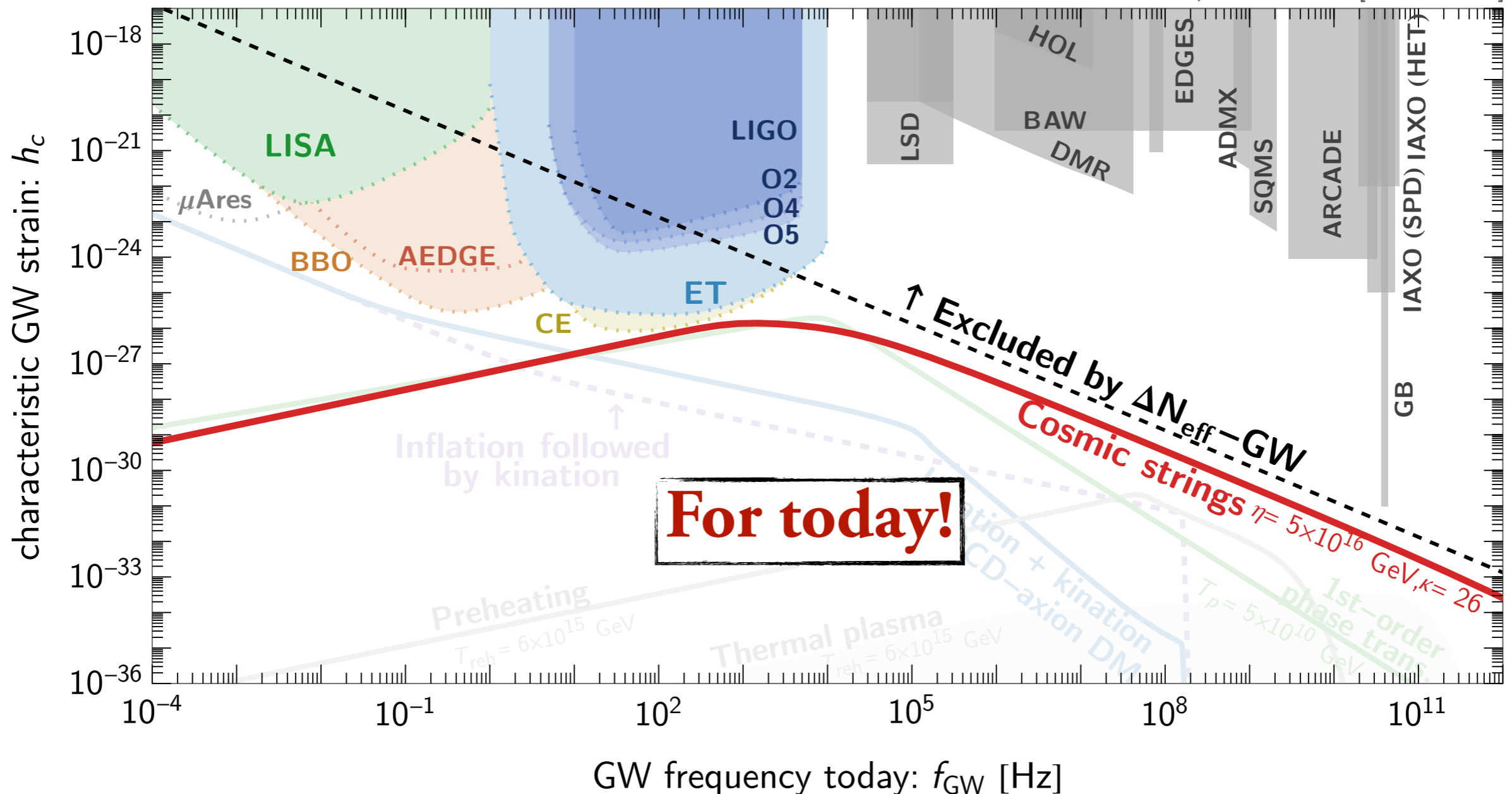
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Servant, Simakachorn [2312.09281]



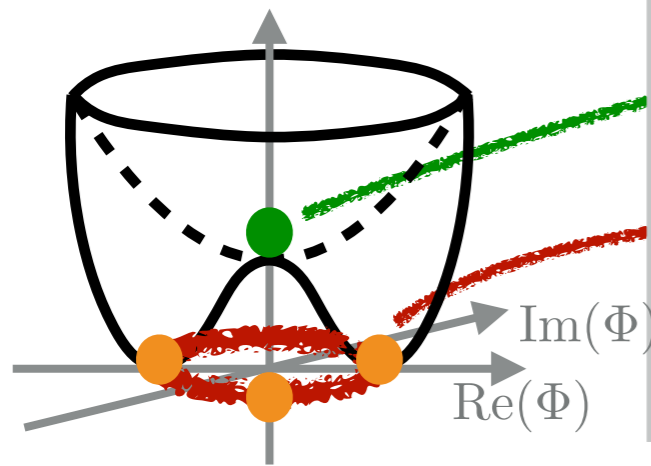
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Cosmic strings and their GW background (GWB)

Recent reviews, e.g., LISA cosmo [1909.00819], Gouttenoire, Servant, Simakachorn [1912.02569], Sousa '23

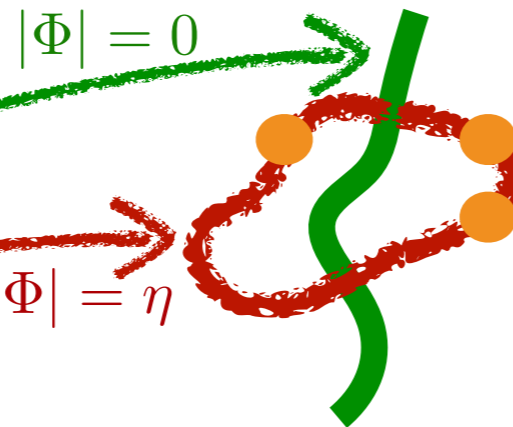
field space

e.g., $V(\Phi) \simeq \lambda(|\Phi|^2 - \eta^2)^2$

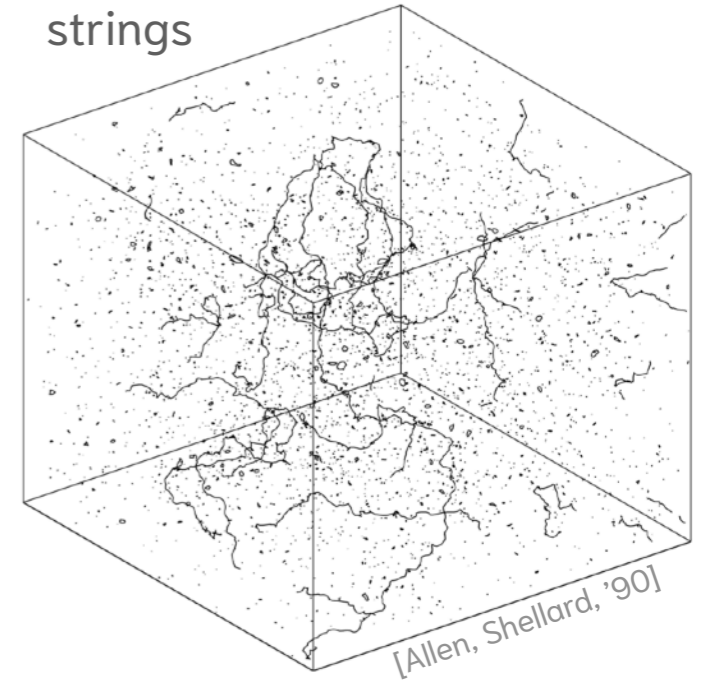


physical space

String's width $w \simeq m_{\Phi}^{-1} \simeq (\sqrt{\lambda}\eta)^{-1}$



Network of strings



[Allen, Shellard, '90]

A topological defect from spontaneous (global or local) symmetry breaking at the energy scale η
[Kibble, '76]

At late times, string's width \ll horizon size \Rightarrow **1D object**

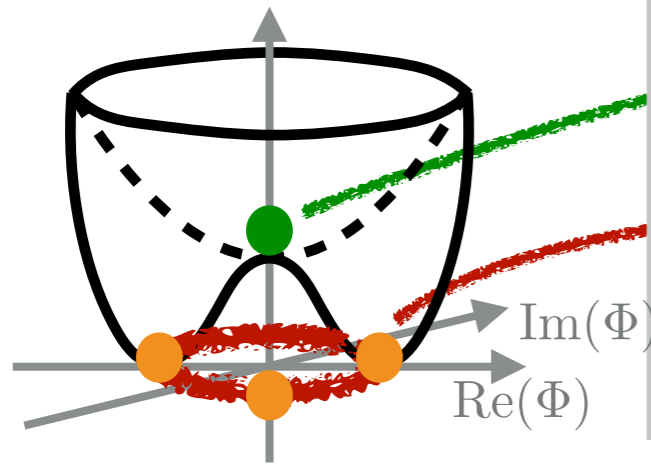
String tension: $G\mu \simeq \left(\frac{\eta}{m_{\text{Pl}}}\right)^2 \simeq 7 \times 10^{-7} \left(\frac{\eta}{10^{16} \text{ GeV}}\right)^2$

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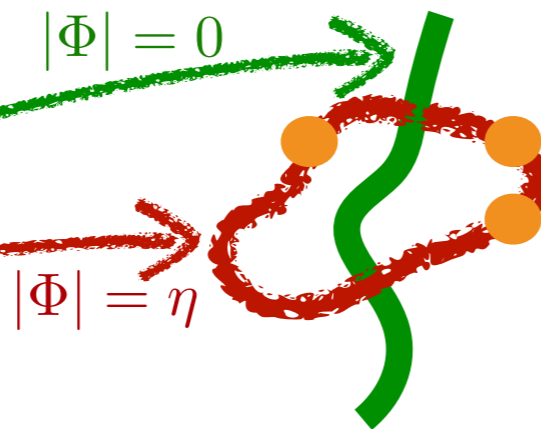
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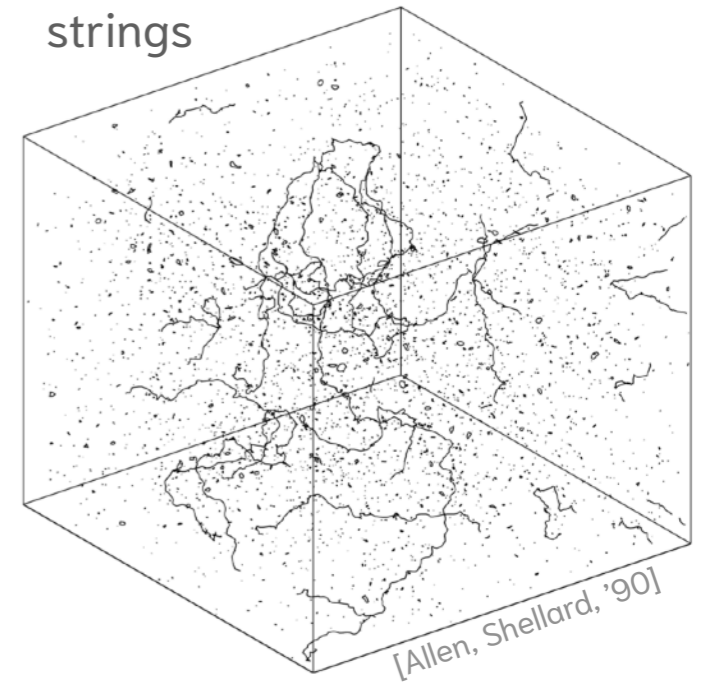


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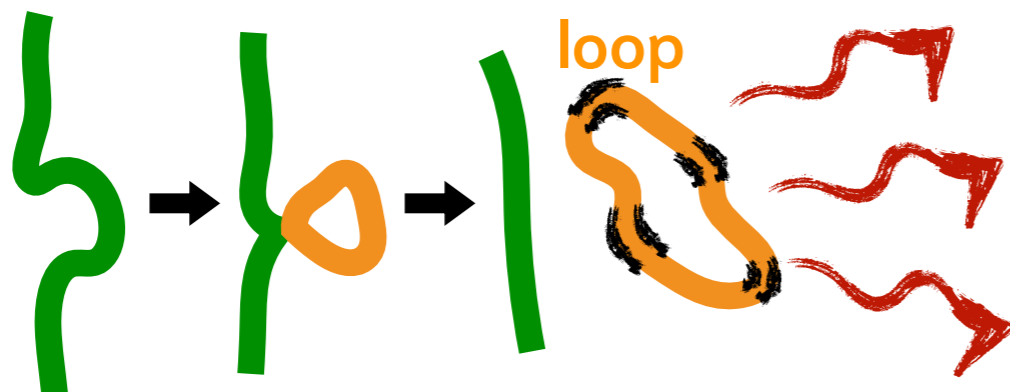
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GW emission from loops
(+ particle production for global strings)

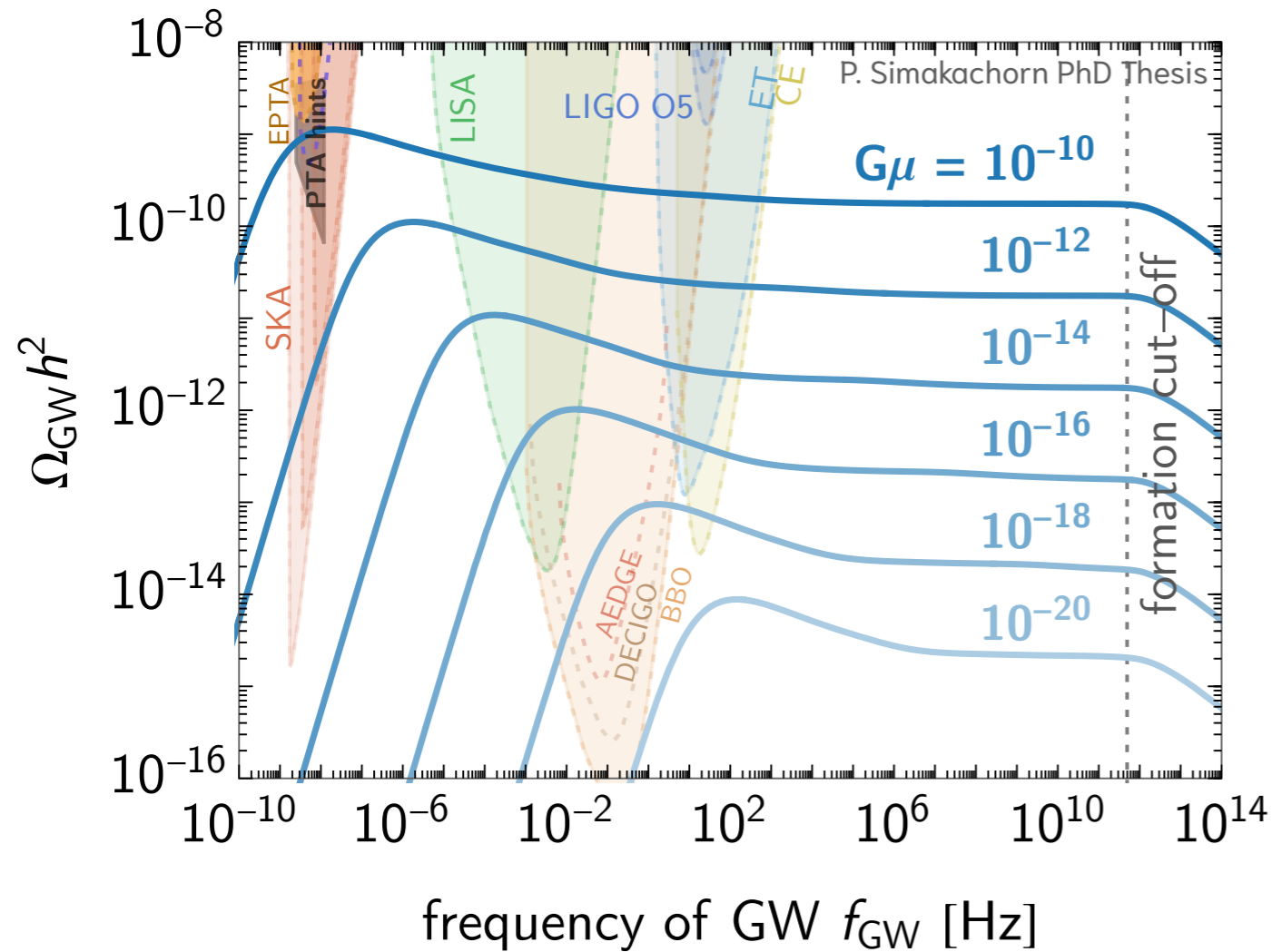
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GW from a loop

of loops produced along cosmic history
(from production time until today)



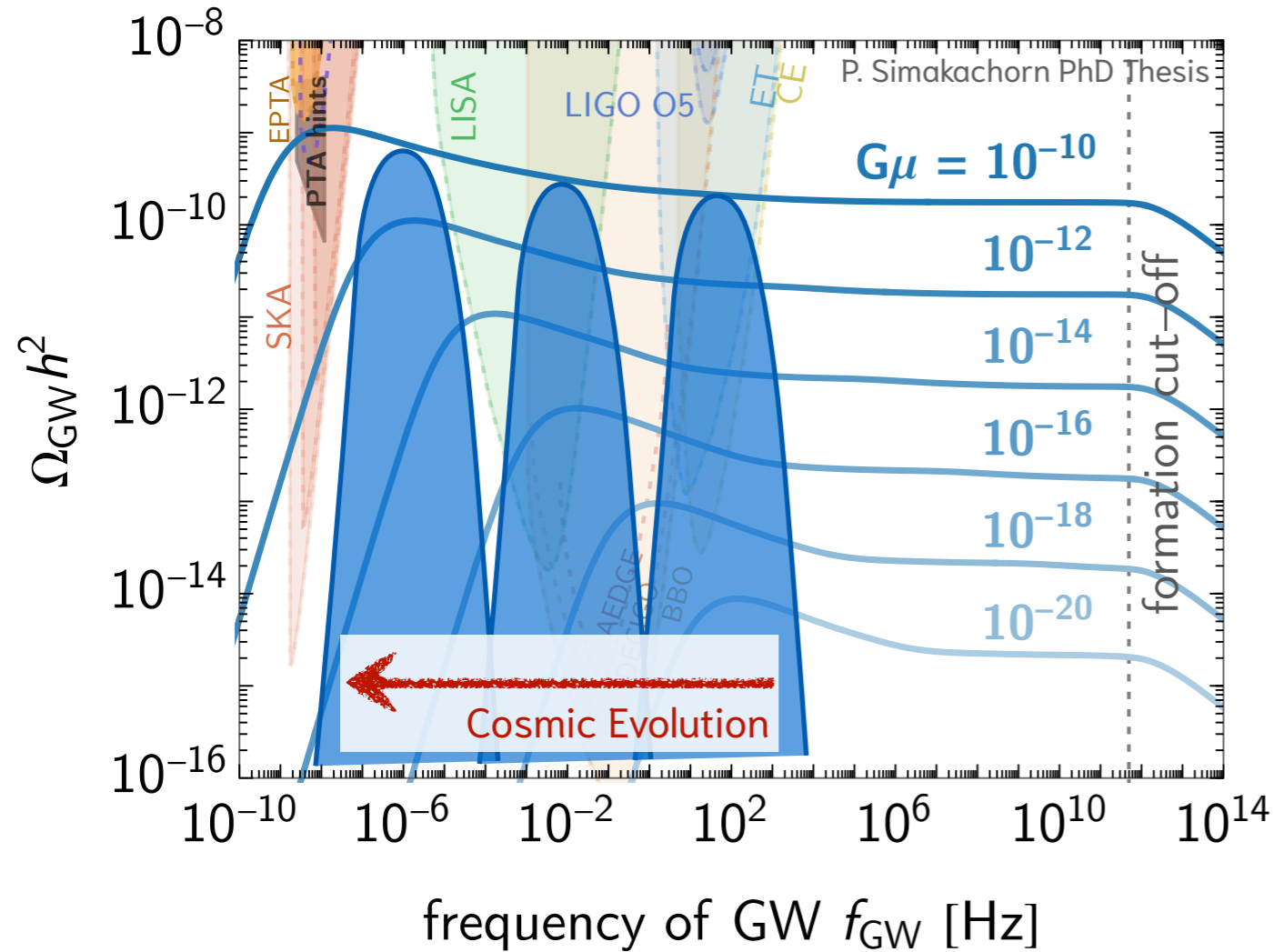
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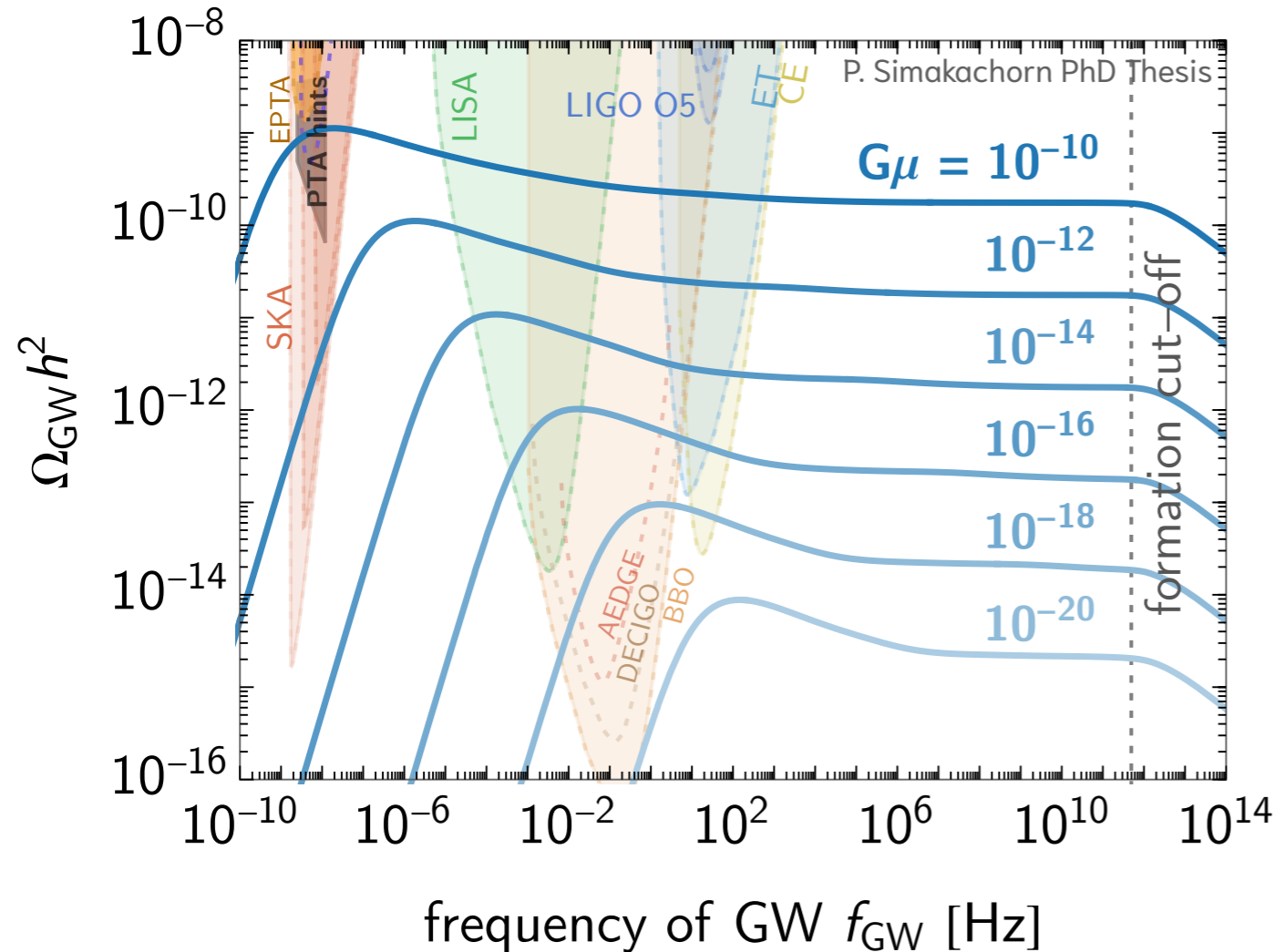
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Not so large UHF signals due to observations @ low-frequency.

LVK (LIGO-VIRGO-KAGRA) @ ~10 Hz

$$\Omega_{\text{GW}} h^2 \lesssim 10^{-8} \Rightarrow G\mu \lesssim 10^{-7}$$

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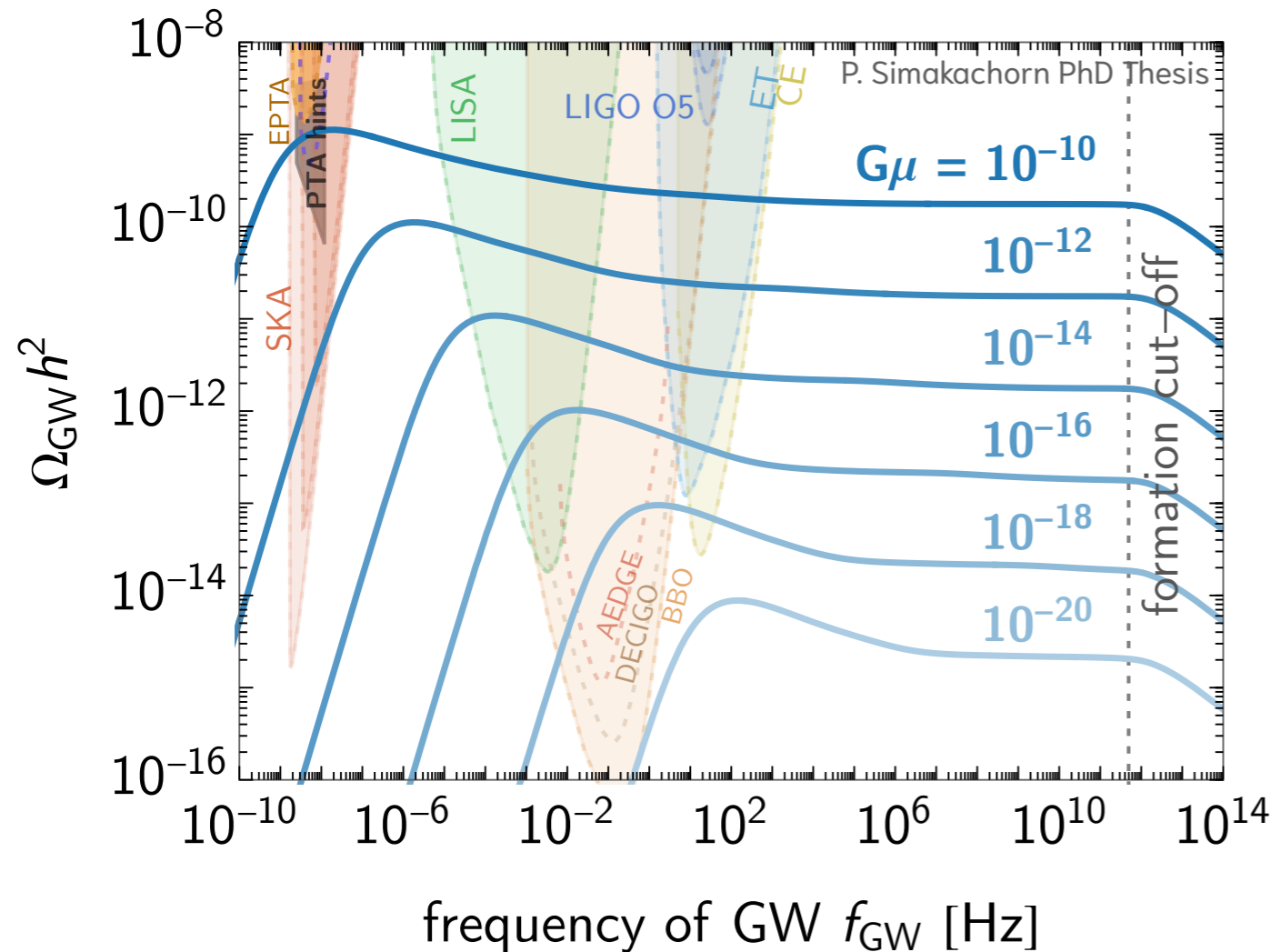
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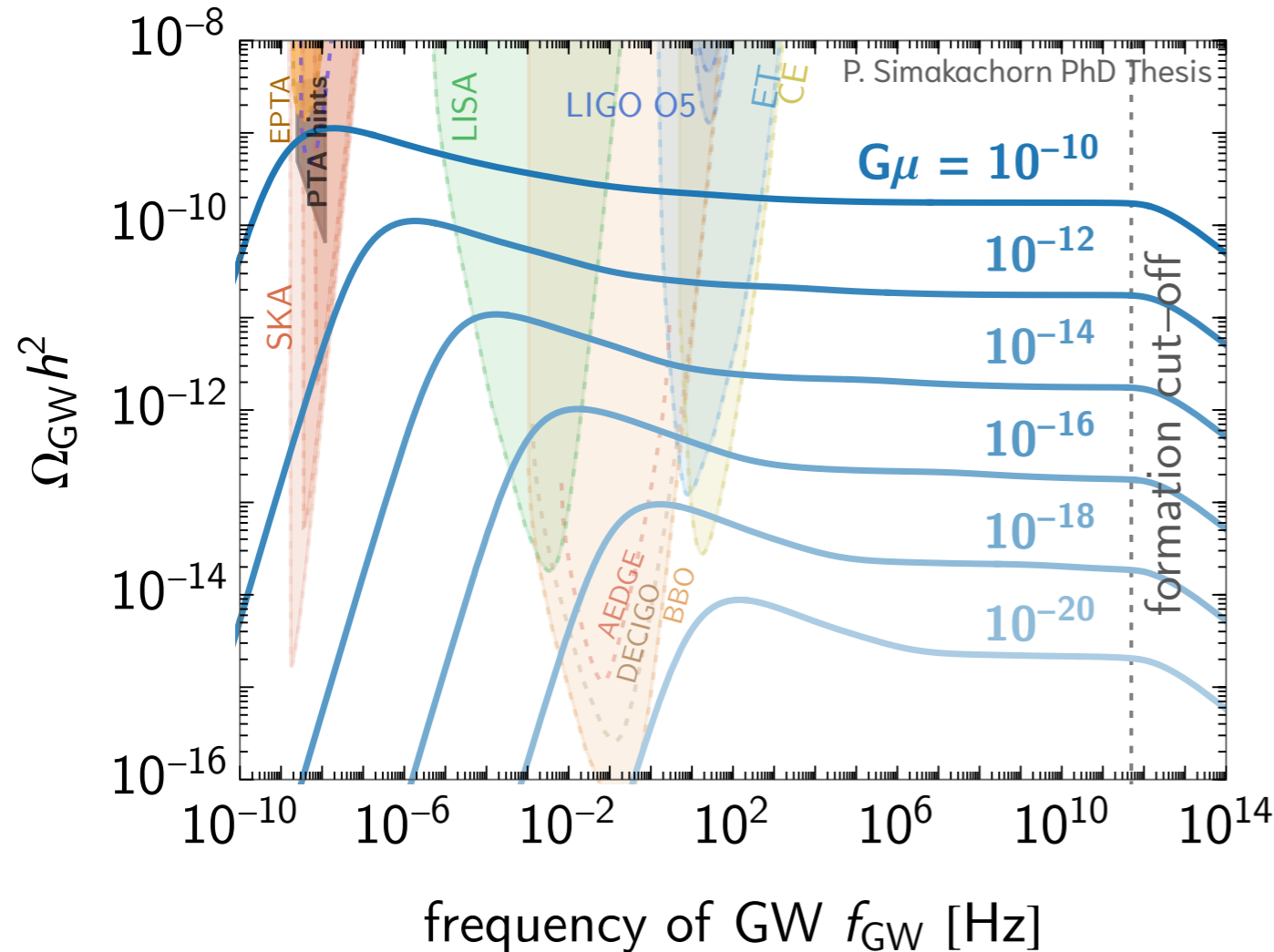
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String network decays!

Metastable cosmic strings

(Also see talks of S. Antusch and S. King)

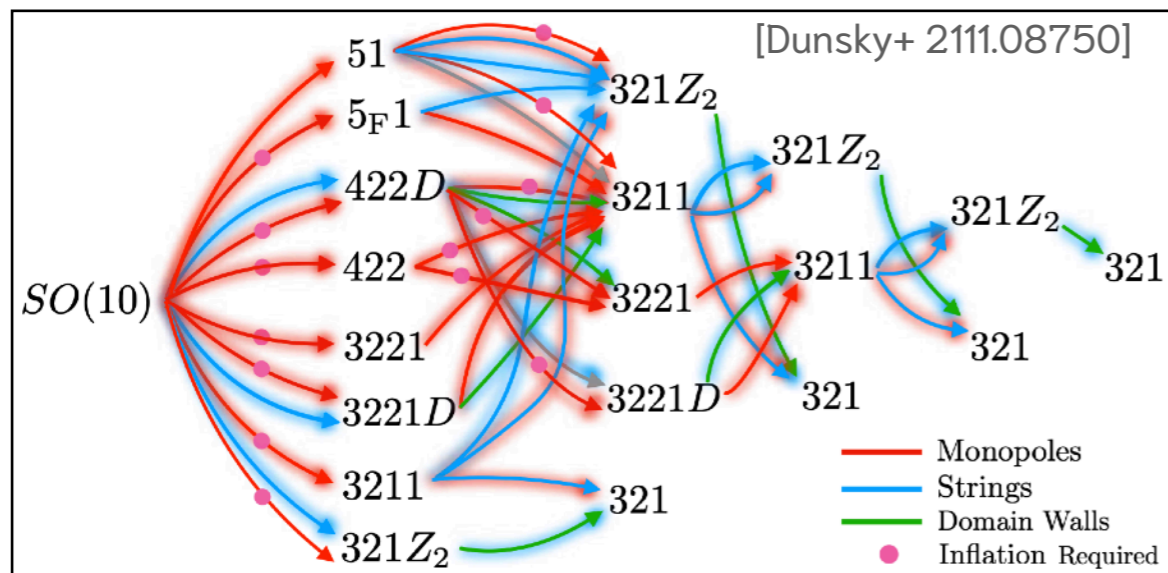
Multiple-step phase transitions \Rightarrow Hybrid defects

E.g., Grand unified theories, hybrid inflation models, post-inflationary axion

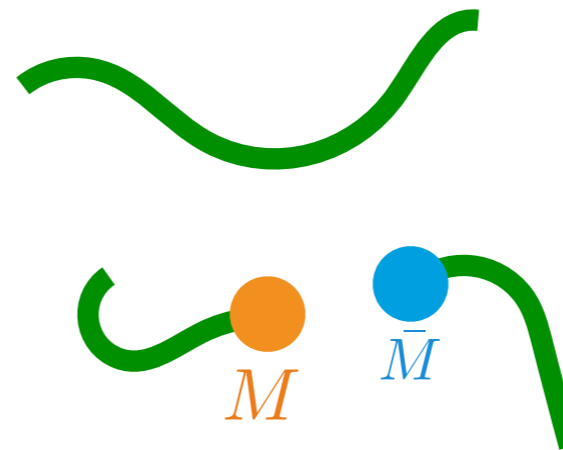
Buchmüller, Domcke, Schmitz+
 [1912.03695, 2102.08923, 2107.04578, 2307.04691]
 King, Pascoli, Turner, Zhou [2005.13549, 2106.15634]
 Murayama+ [1908.03227, 2111.08750]

Vilenkin, Everett '82
 Hiramatsu+ [1012.4558, 1202.5851]

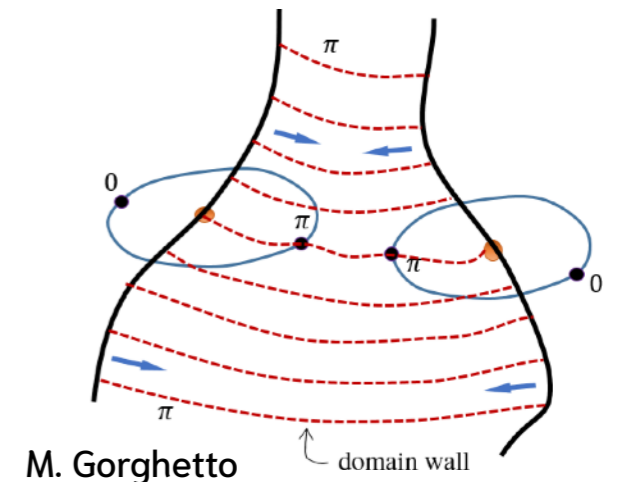
and many others...



string-monopole



String-domain wall



See also decaying defects

via Hubble-induced phase transition [see the talk by Javier Rubio]

Kamada, Yamada [1407.2882] Bettoni, Domènech, Rubio [1810.11117]

Example of metastable cosmic strings

energy
scale

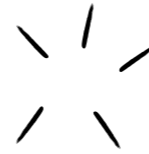


Symmetry breaking #1



Monopole formation @ energy scale m_M

Inflation



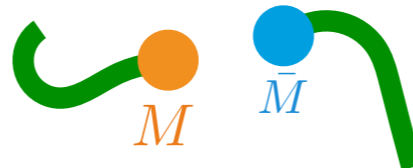
Diluting monopoles,
otherwise monopole domination.

Symmetry breaking #2



String formation @ energy scale η

Decay of strings



Monopole nucleation and strings decay

Example of metastable cosmic strings

energy
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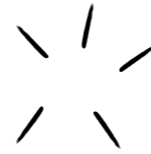


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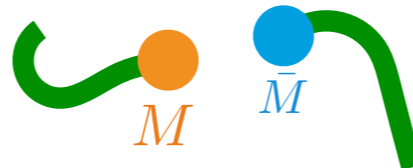
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Monopole nucleation and strings decay

Nucleation rate per unit length:

(assuming thin strings and point-like monopole)

Vilenkin '82, Preskill, Vilenkin '93 Monin,+ '08

Beyond thin soliton approx. :

Dvali+ 2210.14947, Chitose+ 2312.15662

$$\Gamma_d = \frac{\mu}{2\pi} e^{-\pi\kappa} \quad \text{with} \quad \sqrt{\kappa} \equiv \frac{m_M}{\eta} \left[\frac{\text{monopole scale}}{\text{string scale}} \right] > 1$$

happens when $\Gamma_d l(t_{\text{brk}}) = H(t_{\text{brk}})$

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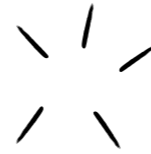


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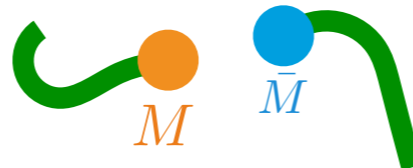
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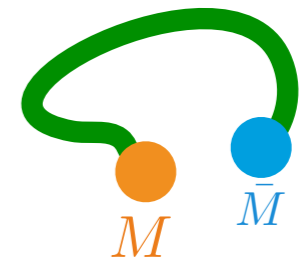
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Two effects on loops:

$$n_{\text{loop}}^{\text{brk}} = n_{\text{loop}}^{\text{stable}} \Theta(t_{\text{brk}} - t_i) e^{-\Gamma_d [l(t-t_{\text{brk}}) \cdots]}$$



1. Shut off loops production

2. Depleting loop number density



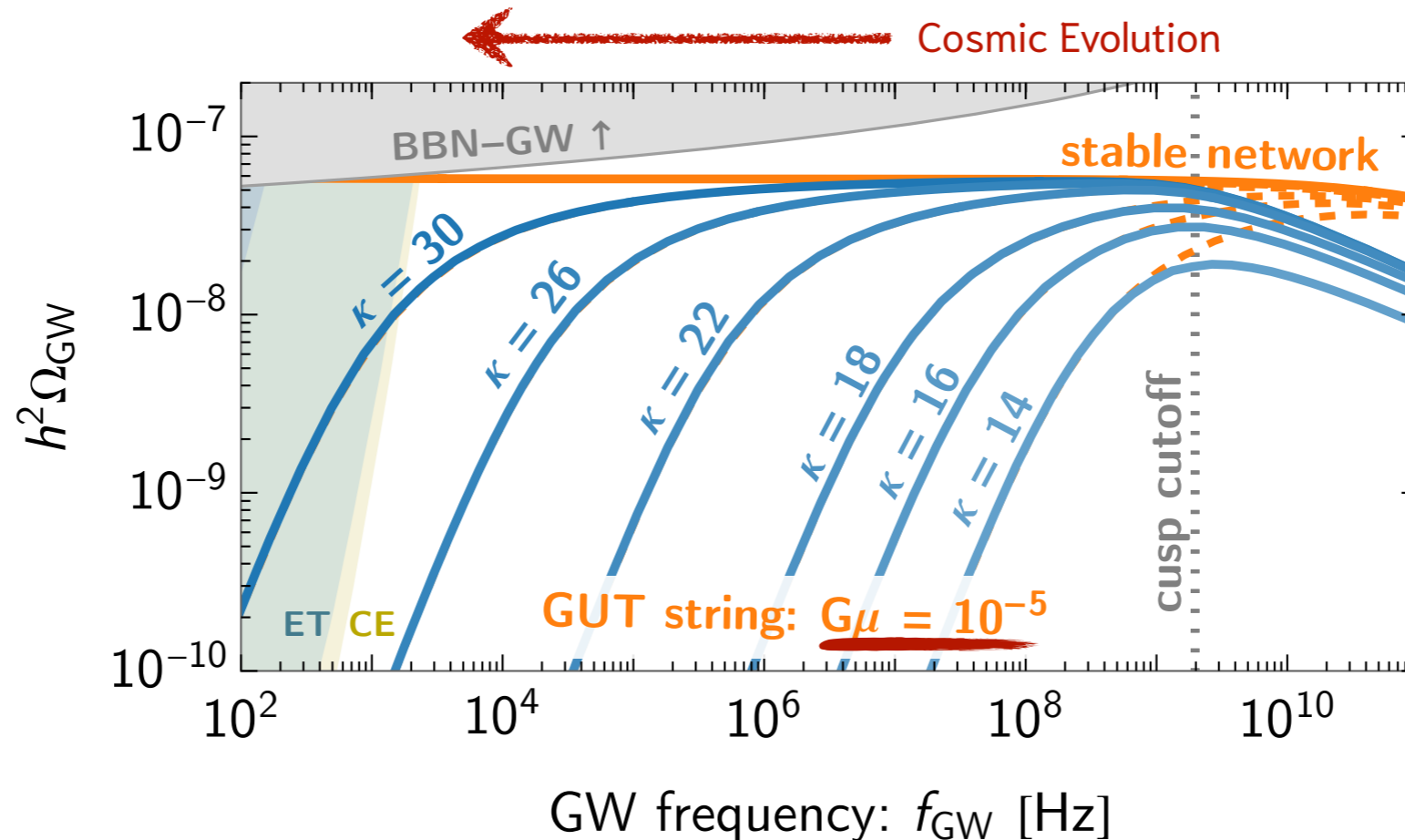
dominant factor to suppress GWB

Large GWB @ UHF from metastable cosmic strings

Servant, Simakachorn [2312.09281]

Strings get cut @ late times \Rightarrow
suppressed GW at low-frequencies

Infrared (IR) cutoff:
 $f_{\text{GW}}^{\text{meta}} \simeq \text{MHz} \left(10^{-5} / G\mu\right)^{\frac{1}{2}} e^{-\frac{\pi}{4}(\kappa-25)}$



The GWB can saturate the ΔN_{eff} -GW bound,
even if no signal is observed at sub-kHz experiments.

Any UHF experiment with sensitivity
slightly better than ΔN_{eff} -GW bound probes into GUT-scale physics.

Other signatures of particle physics @ UHF: UV cutoffs

UHF signal from smaller loops at earlier times.

$$f_{\text{GW}} \simeq 50 \text{ kHz} \left[\frac{T}{10^{10} \text{ GeV}} \right] \left[\frac{10^{-5}}{G\mu} \right]^{\frac{1}{2}} \quad \left| \quad l_{\text{loop}} \sim H^{-1} \sim 10^{-10} \text{ GeV} \left[\frac{10^{14} \text{ GeV}}{T} \right]^2 \right.$$

1D description breaks e.g.,
small-scale structures (e.g., cusp and kink).

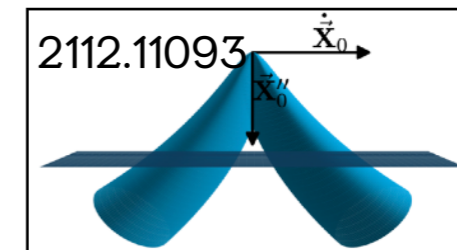
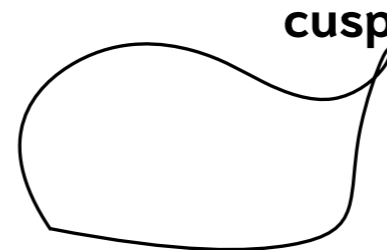
E.g. GW emission is inefficiently
due to particle production from **cusps** if

$$f_{\text{GW}} > f_{\text{GW}}^{\text{cusp}} \simeq \text{GHz} \lambda^{1/8} \left(\frac{G\mu}{10^{-5}} \right)^{3/4}$$

scalar
self-coupling

$$V(\Phi) \simeq \lambda(|\Phi|^2 - \eta^2)^2$$

$$w \simeq m_{\Phi}^{-1} \simeq (\sqrt{\lambda}\eta)^{-1}$$



$$\left. \frac{dE}{dt} \right|_{\text{cusp}} \sim \lambda^{-\frac{1}{4}} \frac{\mu^{3/4}}{l_{\text{loop}}^{1/2}}$$

Other UV cutoffs: the network formation, thermal friction, and particle production from kinks.

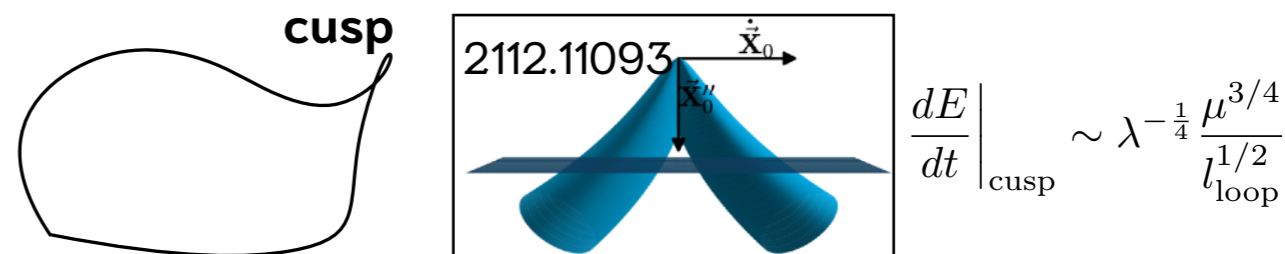
Auclair, Steer, Vachaspati [1911.12066], Gouttenoire, Servant, Simakachorn [1912.02569], Servant, Simakachorn [2312.09281]

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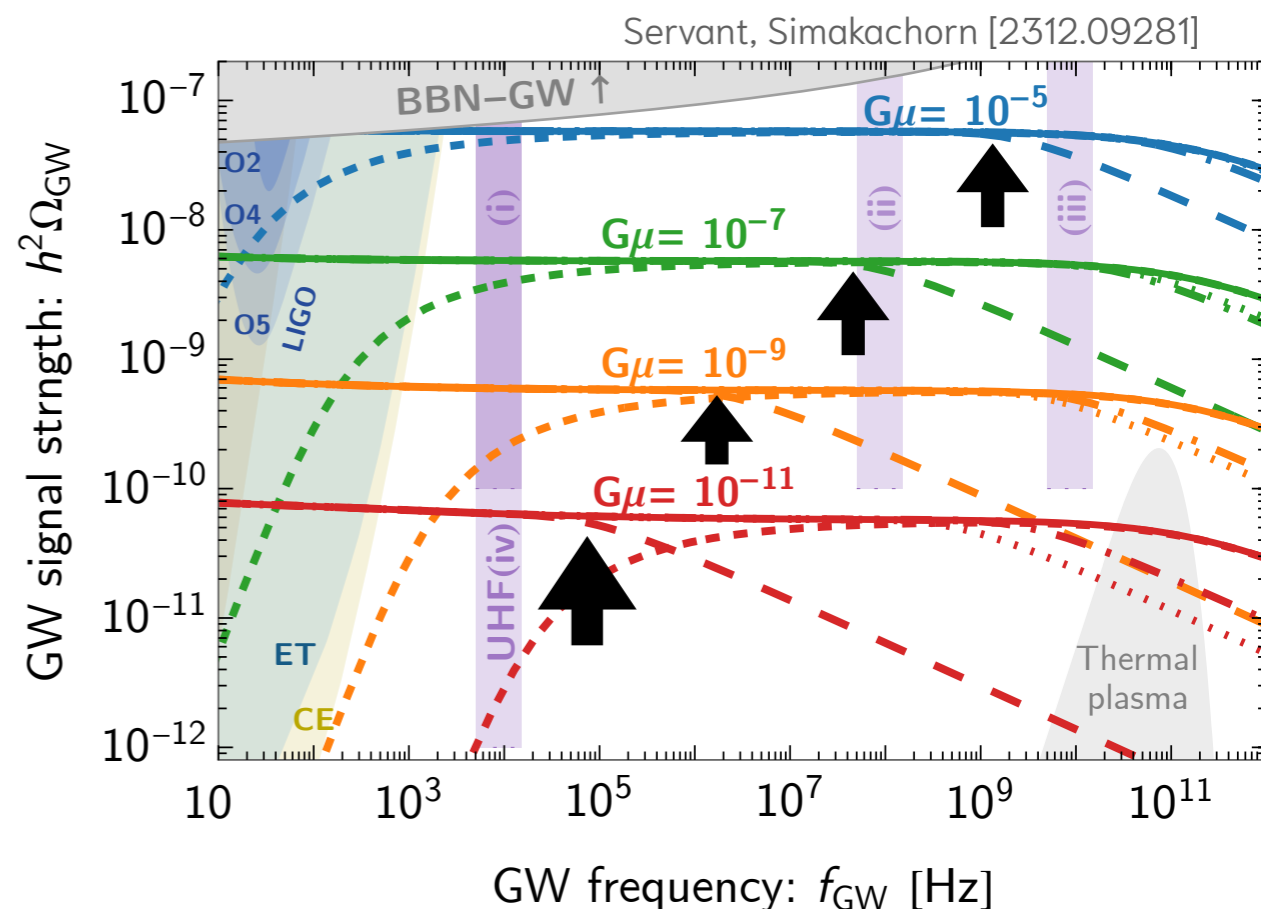
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UV cutoff in the GW spectrum.



Other UV cutoffs: the network formation, thermal friction, and particle production from kinks.

Auclair, Steer, Vachaspati [1911.12066], Gouttenoire, Servant, Simakachorn [1912.02569], Servant, Simakachorn [2312.09281]

Reconstruction of the scalar potential via GW

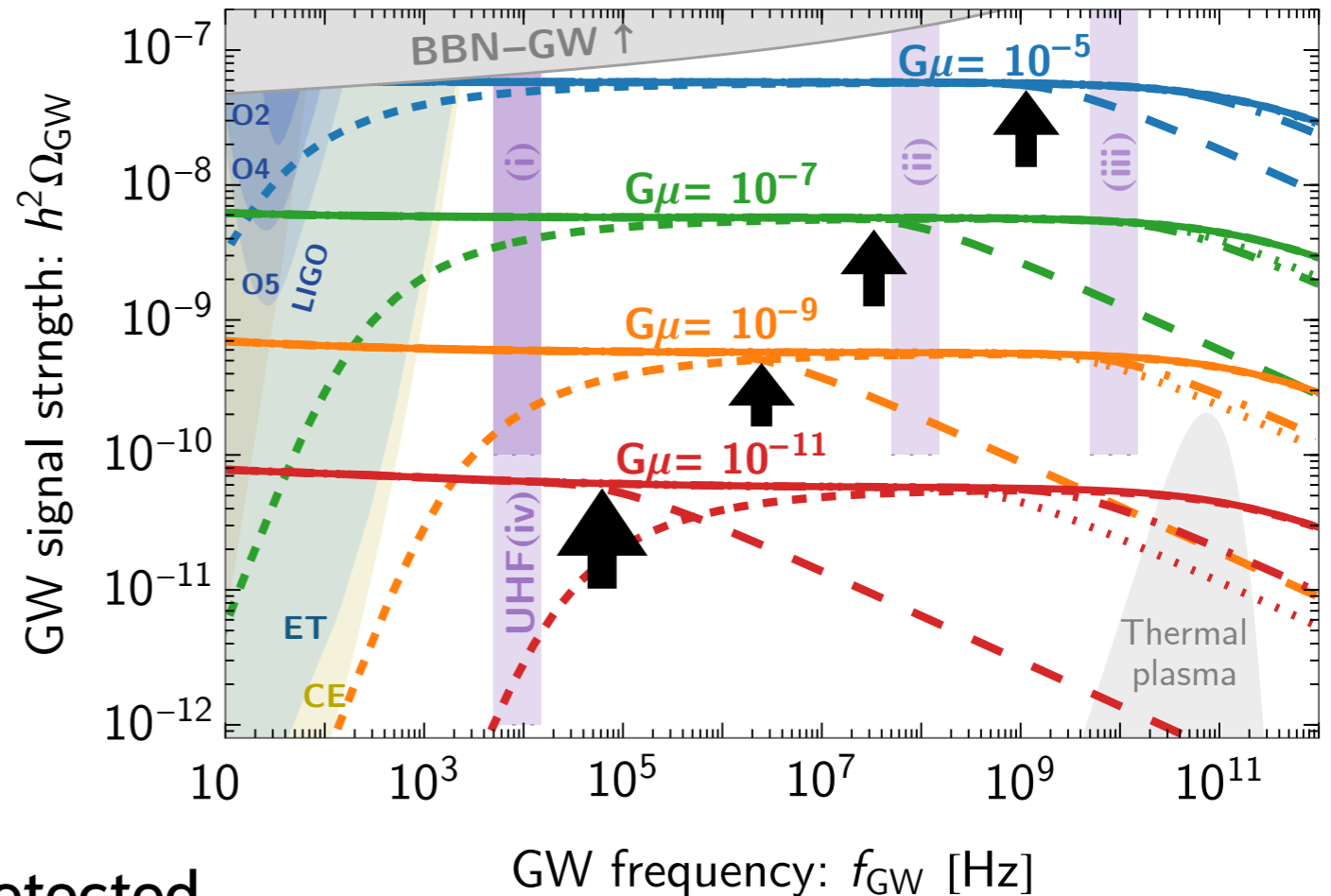
Servant, Simakachorn [2312.09281]

UV cutoff in the GW spectrum
e.g., from cusp,
moves with the string's fatness

$$w \simeq m_{\Phi}^{-1} \simeq (\sqrt{\lambda}\eta)^{-1}$$

$$f_{\text{GW}}^{\text{cusp}} \simeq \text{GHz} \lambda^{1/8} \left(\frac{G\mu}{10^{-5}} \right)^{3/4}$$

$$V(\Phi) \simeq \lambda(|\Phi|^2 - \eta^2)^2$$



How to extract the UV cutoff if GWB is detected.

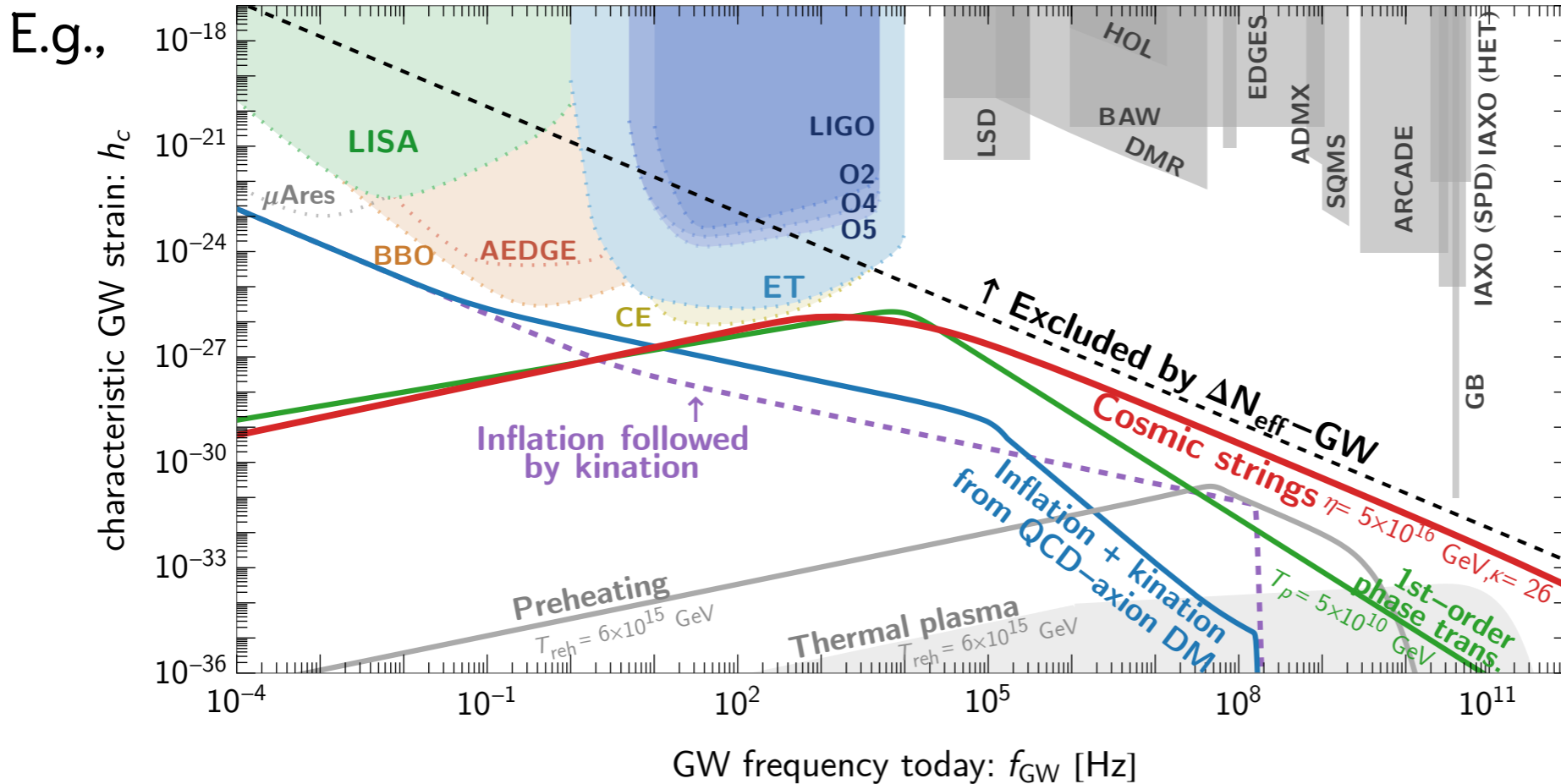
- Detect directly the cutoff (need some luck)
- Several detectors at different frequencies.
Detect the flat part and the UV slope, \Rightarrow UV cutoff at the intersection (more generic)

Conclusion

**UHF regime (beyond kHz) =
A golden window for new physics (above 10^{10} GeV)**

So far...the theoretical ΔN_{eff} -GW bound is the strongest.

But any experiment slightly more sensitive than this bound will probe into the Universe at ultrahigh energy scales.



**Large UHF signal
from metastable strings.**

**The reconstruction
of scalar potential
(e.g., self-coupling λ)**

**is possible via
the observed UV cutoff
in the spectrum!**

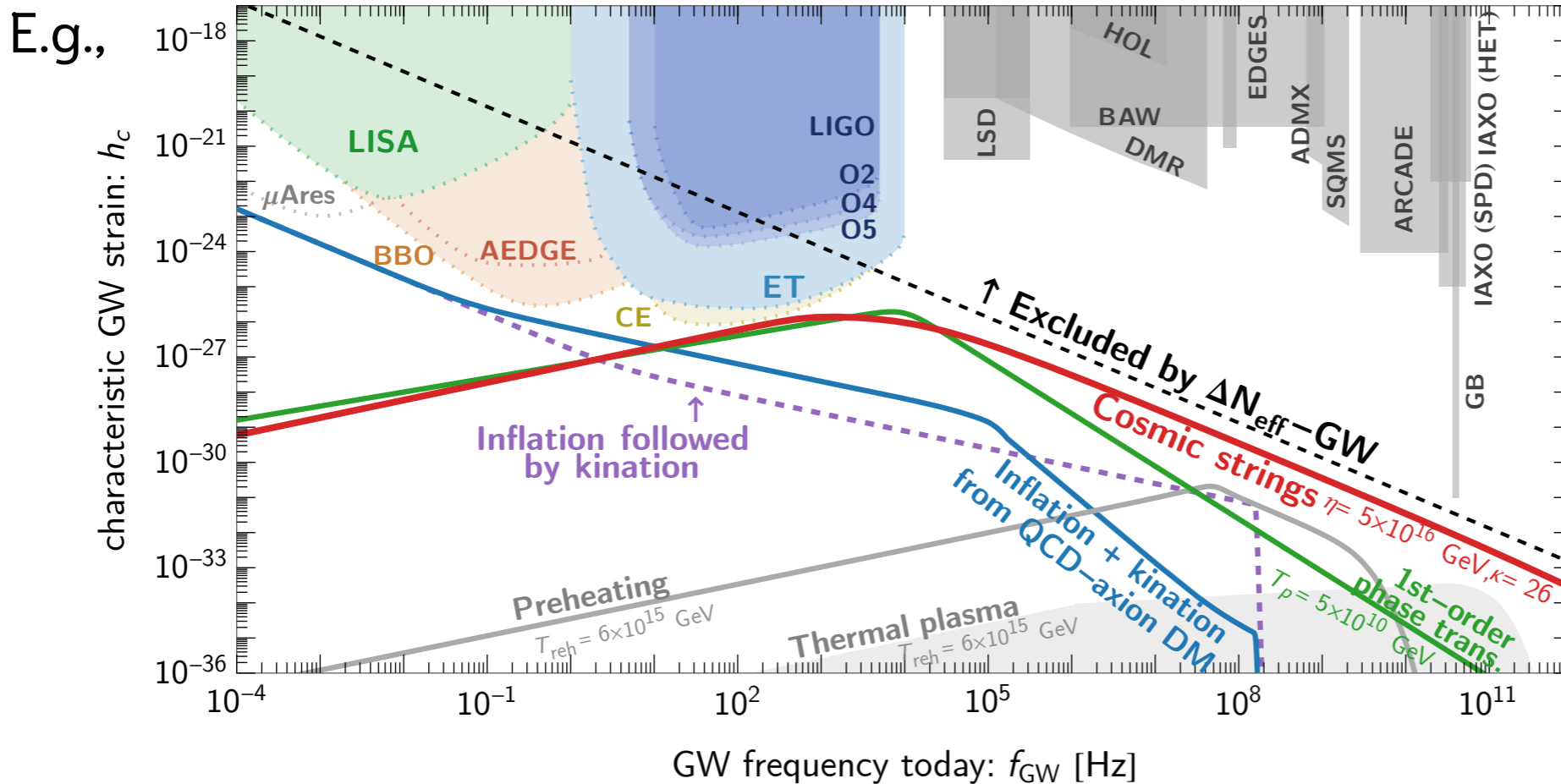
Servant, Simakachorn [2312.09281]

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Servant, Simakachorn [2312.09281]

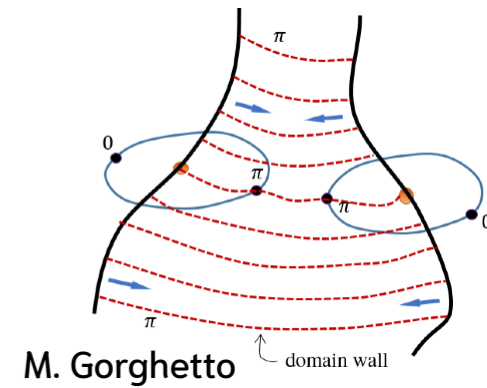
We also look at global strings or axionic strings...

Axionic (or global) strings

$$\Omega_{\text{GW}} \propto \eta^4 \text{ or } f_a^4 \quad \text{with } f_a: \text{Peccei-Quinn symmetry-breaking scale}$$

Strings attach to domain walls and collapse: $T_{\text{dec}} \sim 10^9 \text{ GeV} \sqrt{m_a/\text{GeV}}$

String-domain wall



Light axion ($m_a \lesssim 10^{-22} \text{ eV}$)

\Rightarrow \sim stable strings

Small UHF signal

- ΔN_{eff} -Goldstone bound

$$f_a \lesssim \mathcal{O}(1 - 3) \times 10^{15} \text{ GeV}$$

Cui, Chang '21, Hardy, Nicoleuscu, Gorghetto '21

- Pulsar-timing arrays

$$f_a \lesssim 2.8 \times 10^{15} \text{ GeV}$$

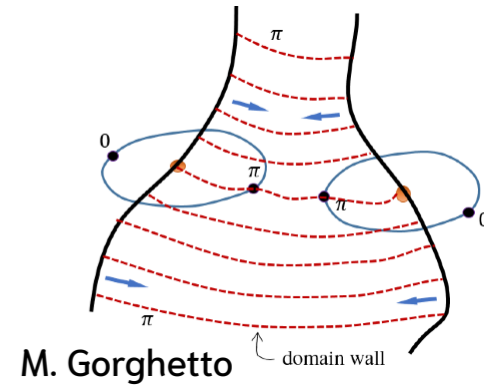
Servant, Simakachorn [2307.03121]

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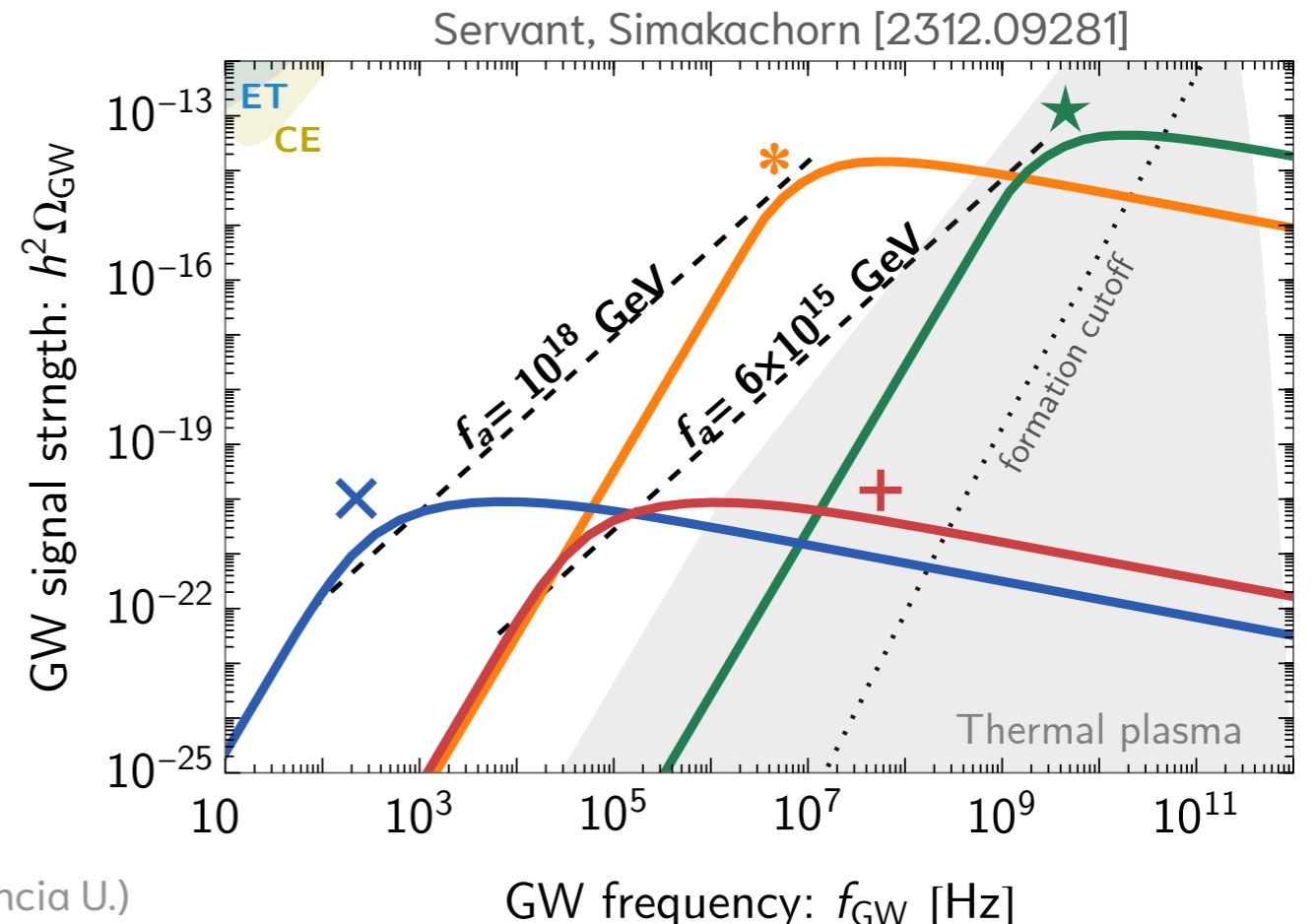
Servant, Simakachorn [2307.03121]

Heavy axion ($m_a \gtrsim \text{GeV}$)

\Rightarrow IR cutoff in UHF

Small signal, even for **large** f_a .

GWB is diluted by matter domination from axions produced from string collapse.

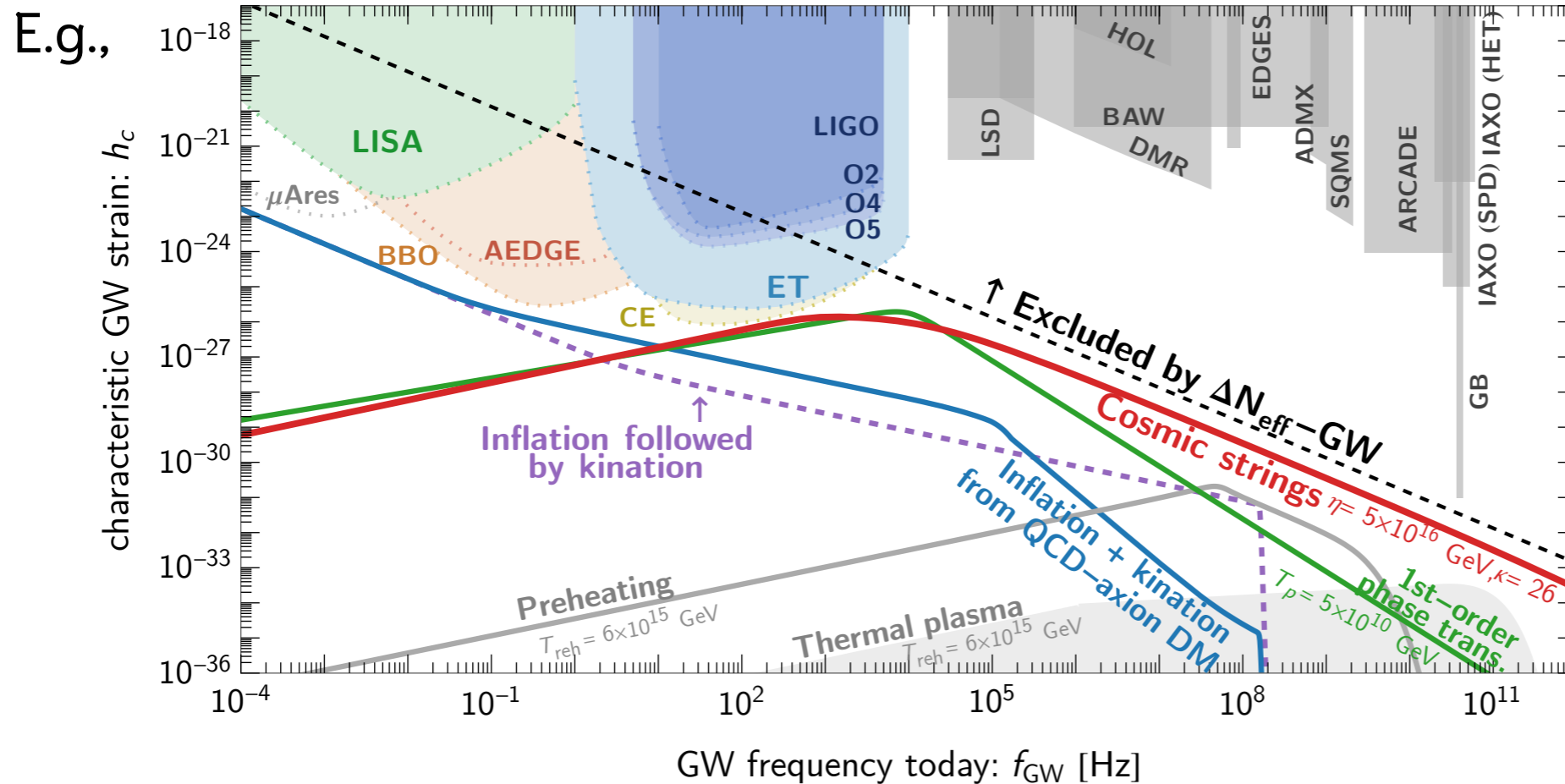


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Servant, Simakachorn [2312.09281]

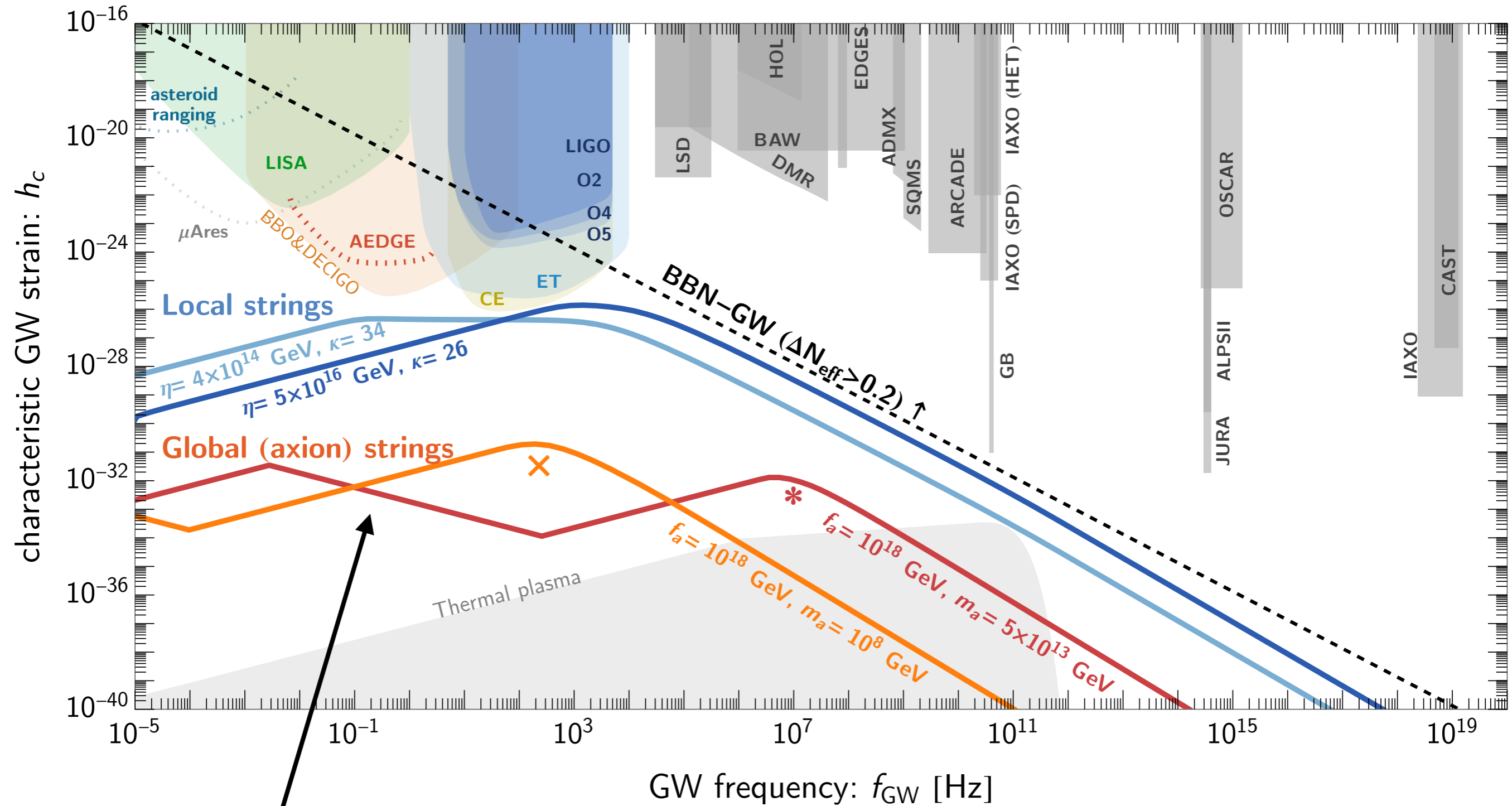
GW signal in UHF from axion strings is suppressed,
i.e., requiring at least sensitivity $h_c \simeq 10^{-31}$.

Thank You!

Back-up

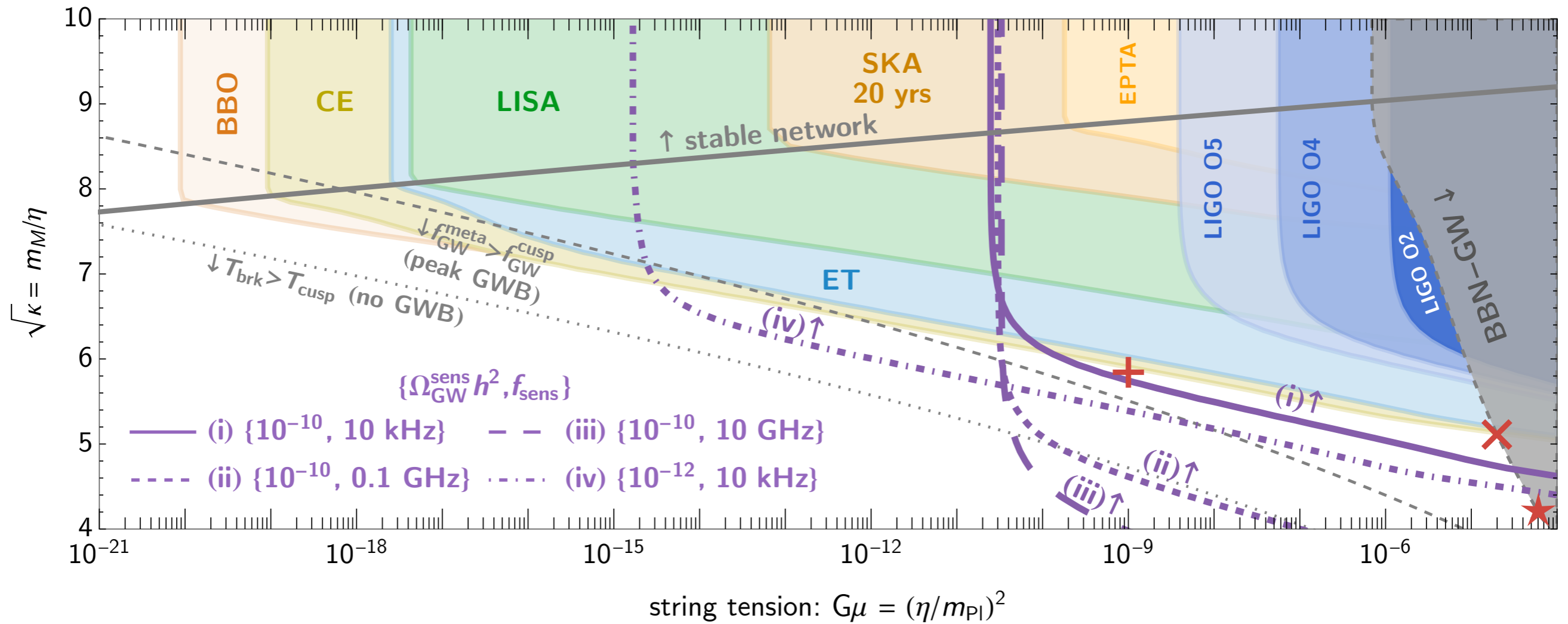
UHF GWB from local and global (axionic) strings (Best cases)

Servant, Simakachorn [2312.09281]

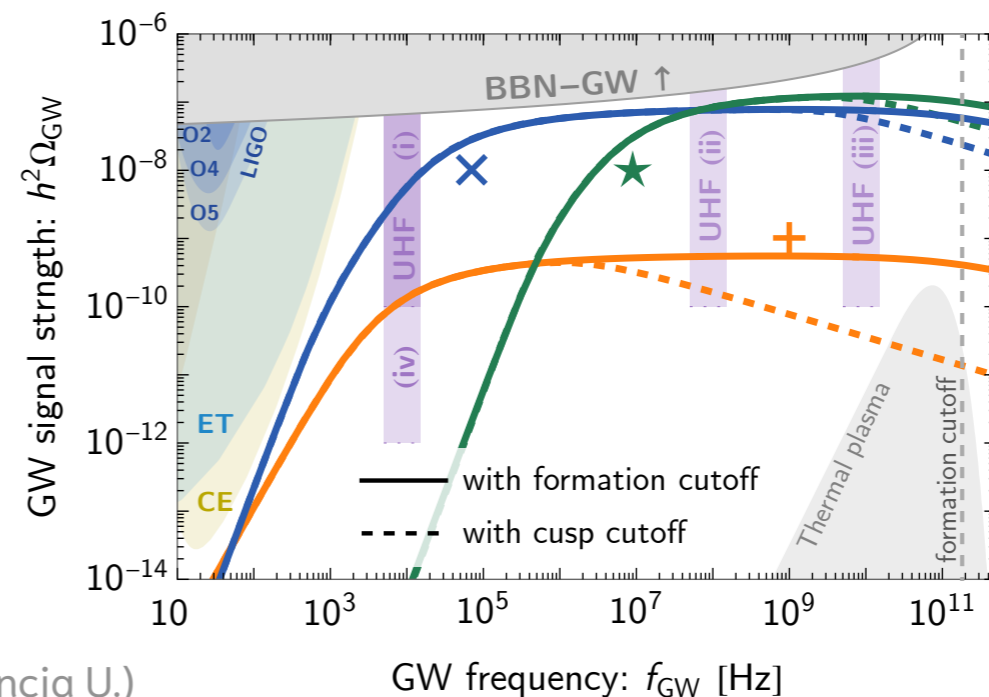


Low-frequency slope is changed by the modified causality tail during the axion matter domination.

Detectability of GWB from metastable local strings



Servant, Simakachorn [2312.09281]



UV cutoffs in cosmic-string GWB

Servant, Simakachorn [2312.09281]

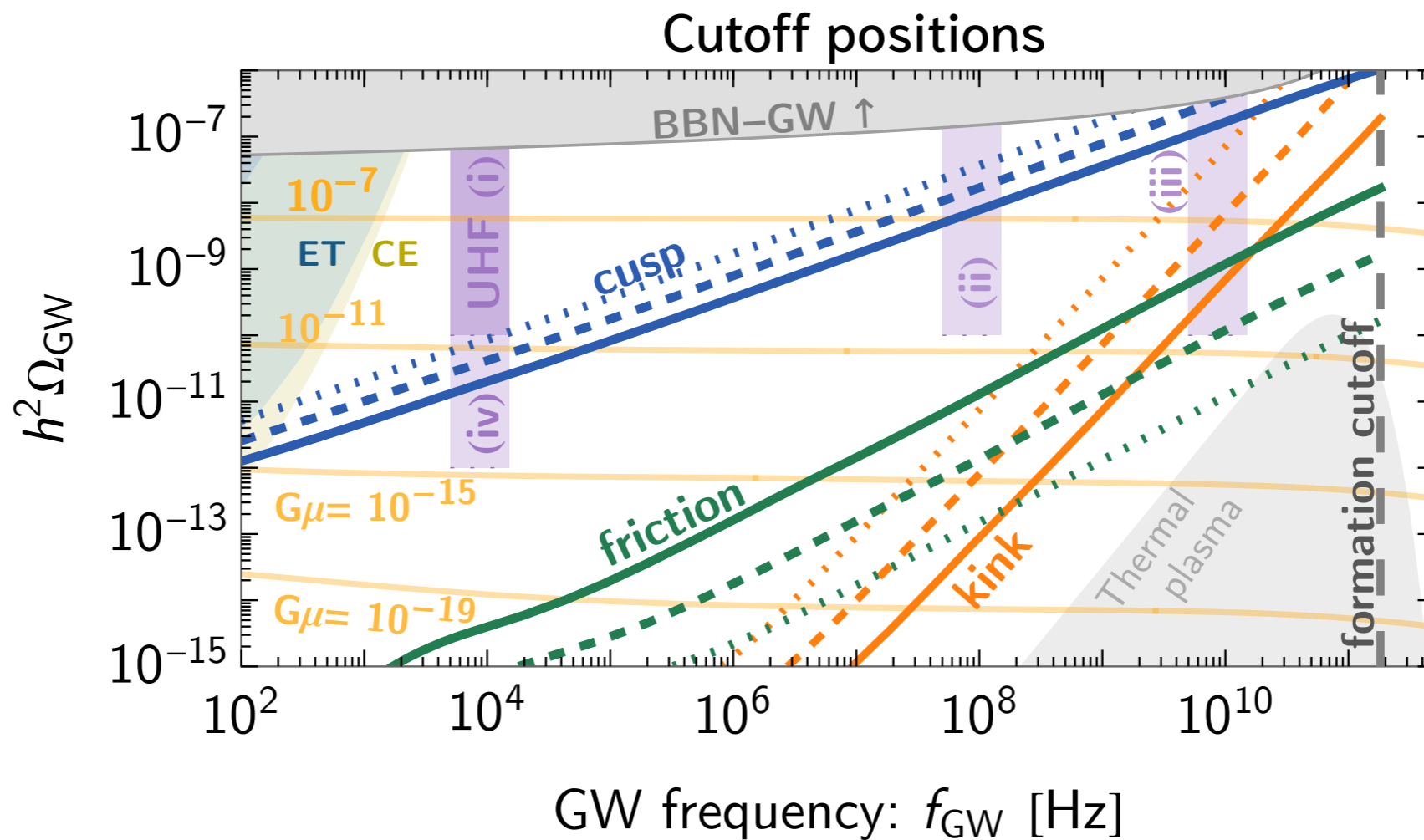
$$f_{\text{GW}}^{\text{form}} \simeq 182 \text{ GHz} \sqrt{\lambda} \left[\frac{g_*(T_{\text{form}})}{106.75} \right]^{\frac{1}{4}}$$

$$f_{\text{GW}}^{\text{cusp}} \simeq 62.3 \text{ kHz} \sqrt{\frac{1}{\beta_c}} \left(\frac{G\mu}{10^{-11}} \right)^{3/4}$$

$$f_{\text{GW}}^{\text{fric}} \simeq 0.45 \text{ GHz} \beta_{\text{fric}}^{-1} \left(\frac{G\mu}{10^{-11}} \right)^{\frac{1}{2}} \left[\frac{g_*(T)}{106.75} \right]^{\frac{3}{4}}$$

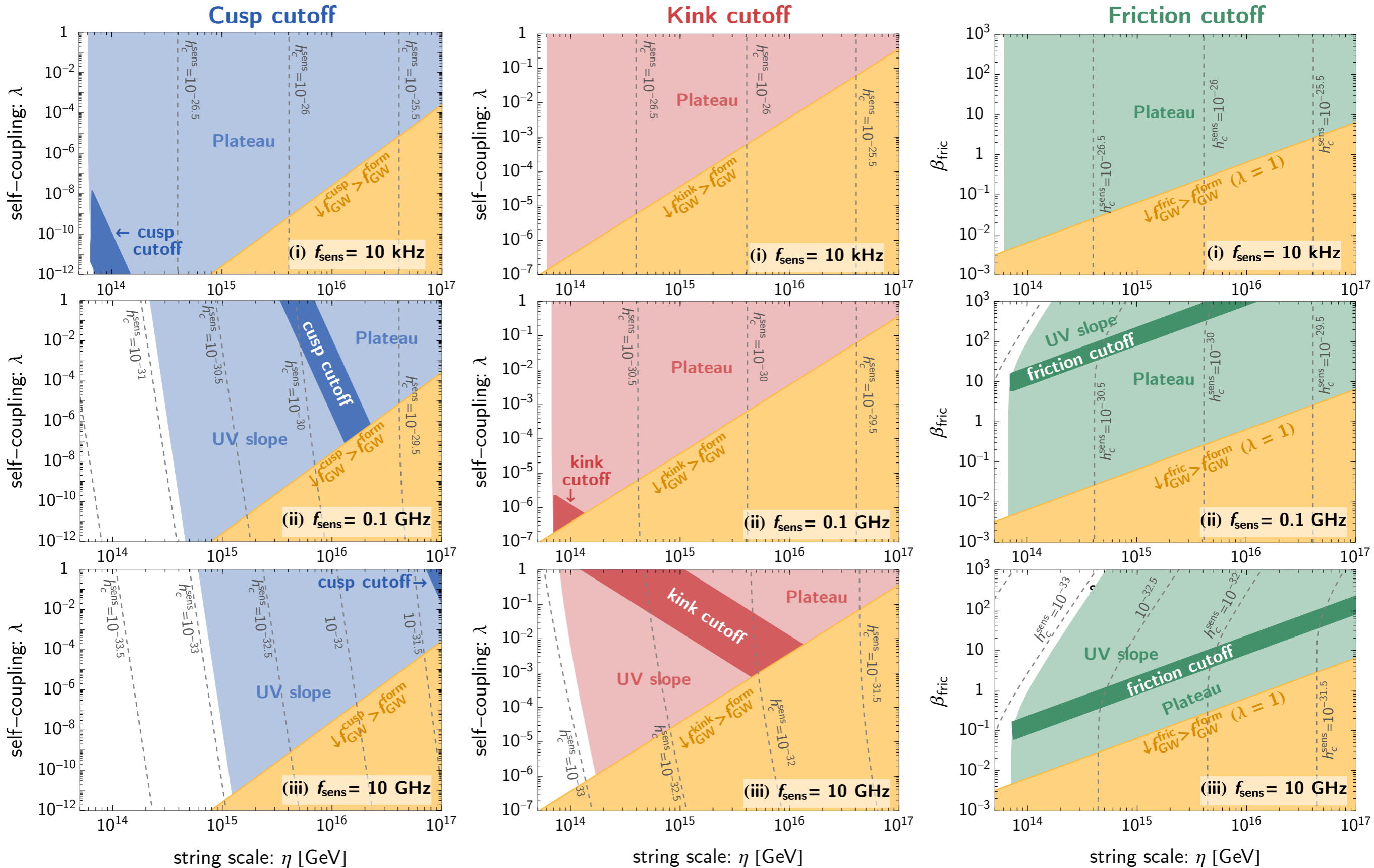
$$f_{\text{GW}}^{\text{kink}} \simeq 2.79 \text{ GHz} \sqrt{\frac{1}{\beta_k}} \left(\frac{G\mu}{10^{-11}} \right)^{1/4}$$

$$\beta_c \simeq N_c \lambda^{-1/4}, \text{ and } \beta_k \simeq N_{kk} \lambda^{-1/2}$$



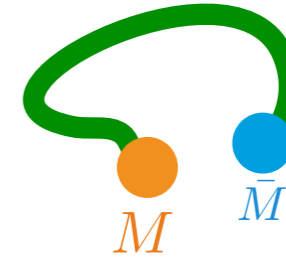
Reconstruction of scalar potential with UHF GWB

Servant, Simakachorn [2312.09281]

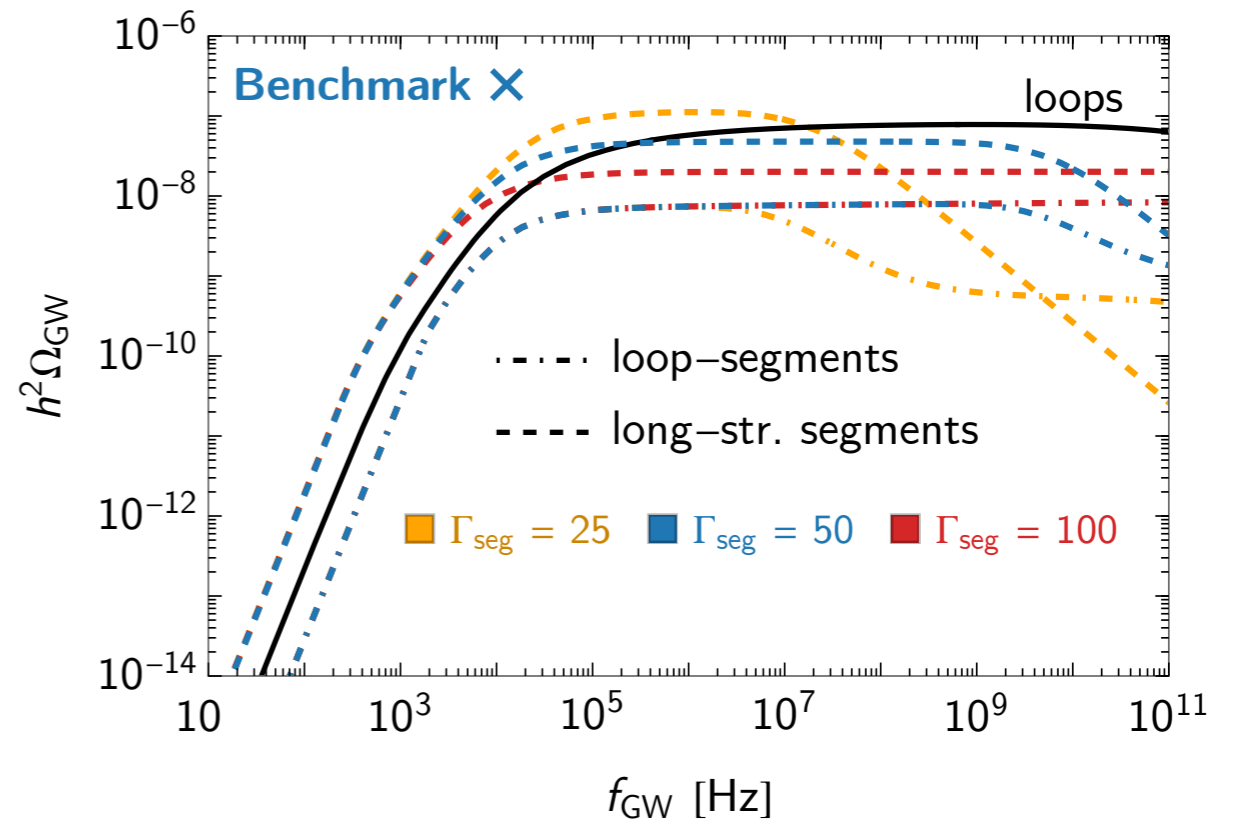
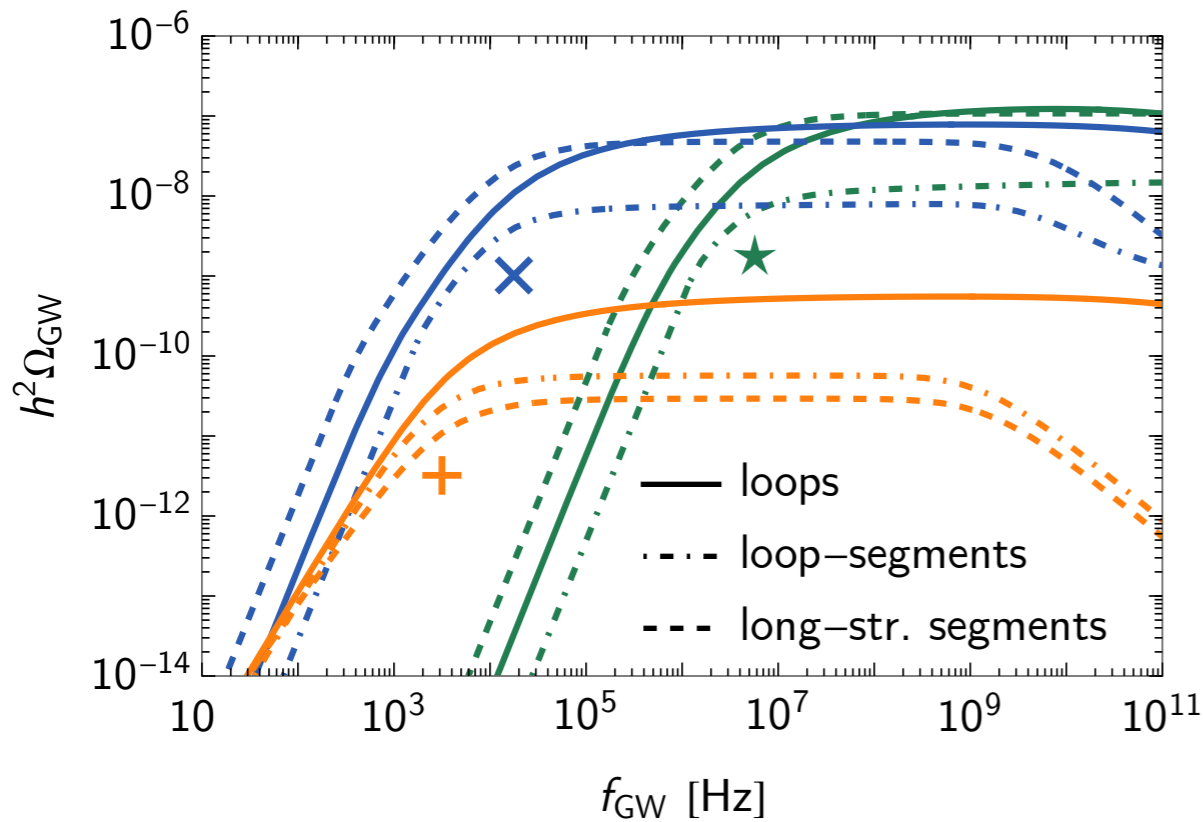


GWB from string segments

Servant, Simakachorn [2312.09281]



Currently, there are several uncertainties in calculating the GW from segments.



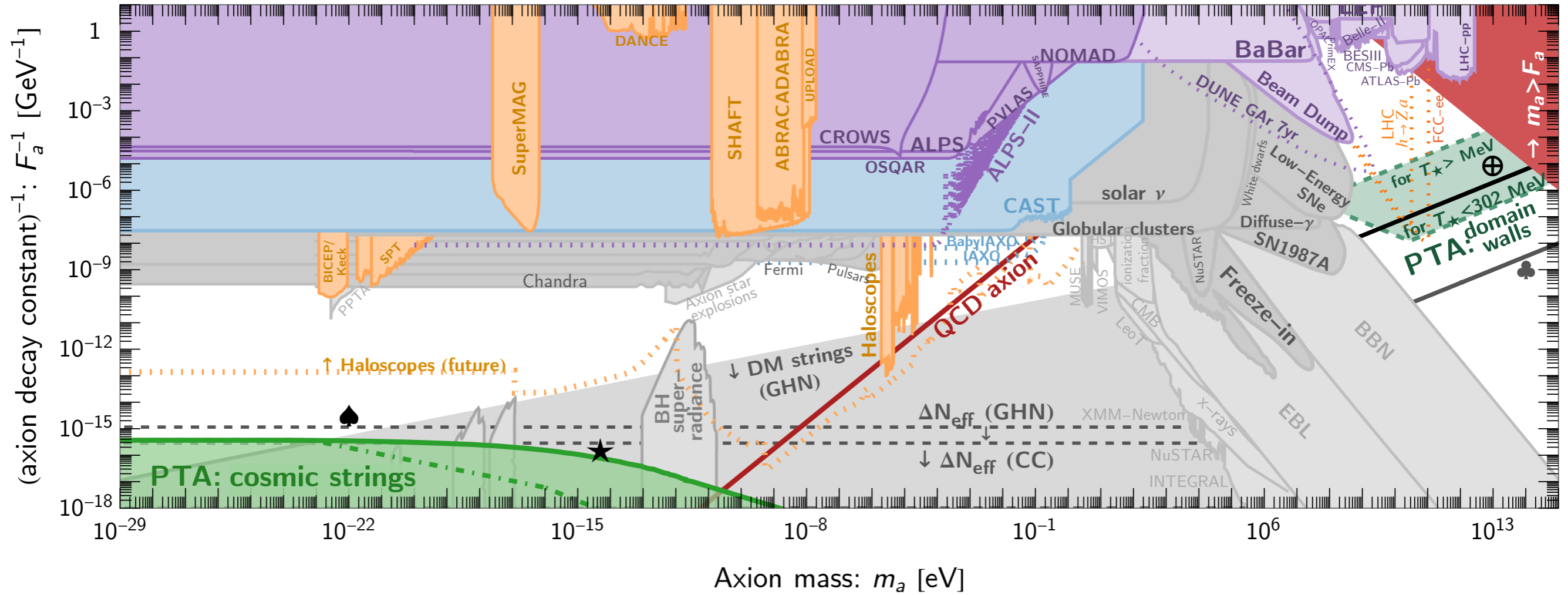
See more details also in Buchmüller, Domcke, Schmitz [2107.04578]

For going beyond the thin soliton approximation
see Dvali, Valbuena-Bermudez, Zantedeschi [2210.14947].

Pulsar timing array constraints on postinflationary axion

Servant, Simakachorn [2307.03121]

Using NANOGrav 15-year data



Axion matter domination from axionic string decay

Servant, Simakachorn [2312.09281]

Axion string-wall system decays.

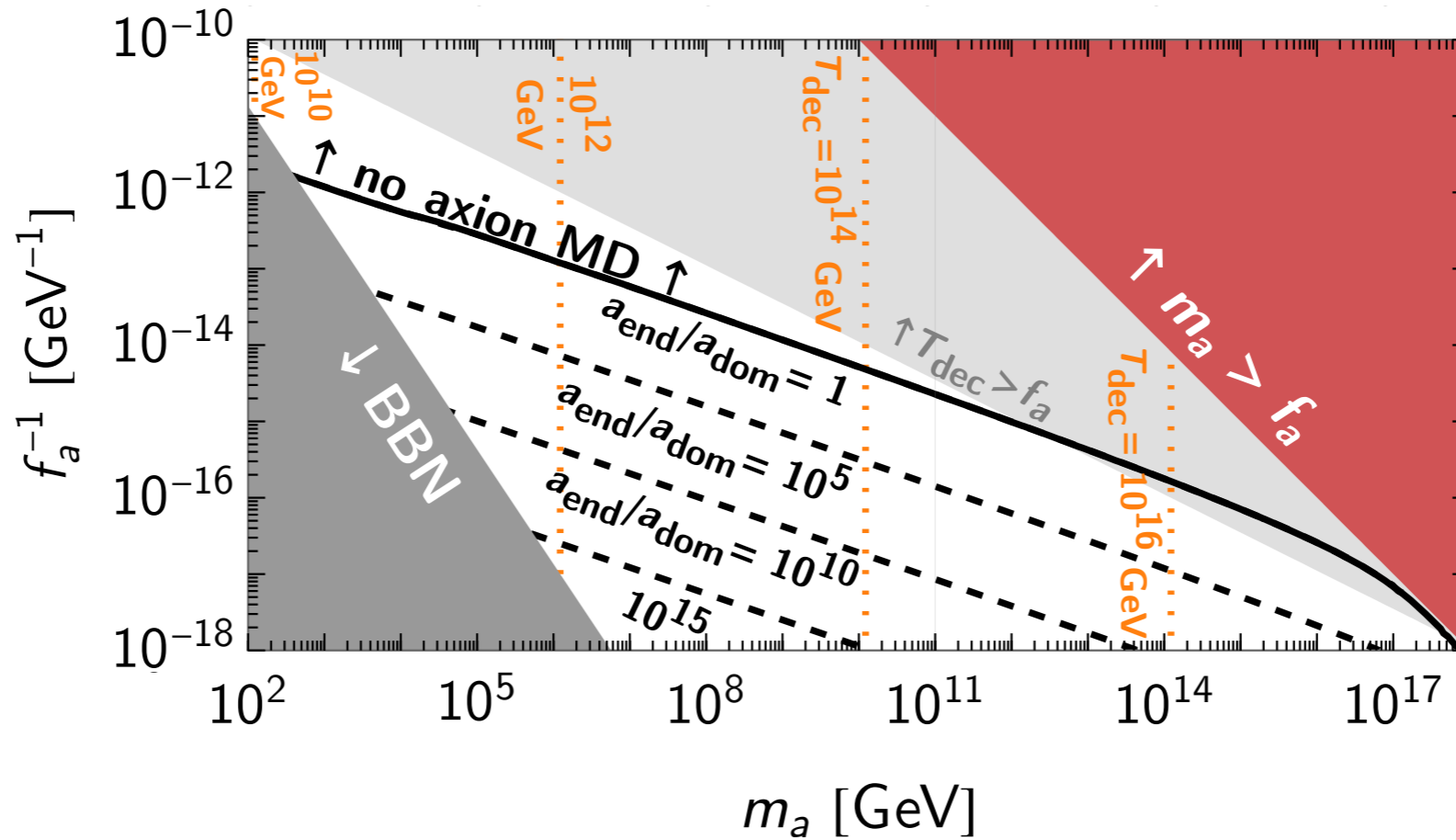
Axion-matter domination

Axions decay into photons

$$T_{\text{dec}} \sim 10^9 \text{ GeV} \sqrt{m_a/\text{GeV}}$$

$$T_{\text{dom}} \simeq T_{\text{dec}} G\mu(T_{\text{dec}})$$

$$T_{a\gamma} \simeq 4.2 \text{ MeV} \left[\frac{106.75}{g_*(T_{a\gamma})} \right]^{\frac{1}{4}} \left(\frac{m_a}{\text{TeV}} \right)^{\frac{3}{2}} \left[\frac{10^{12} \text{ GeV}}{f_a} \right]$$



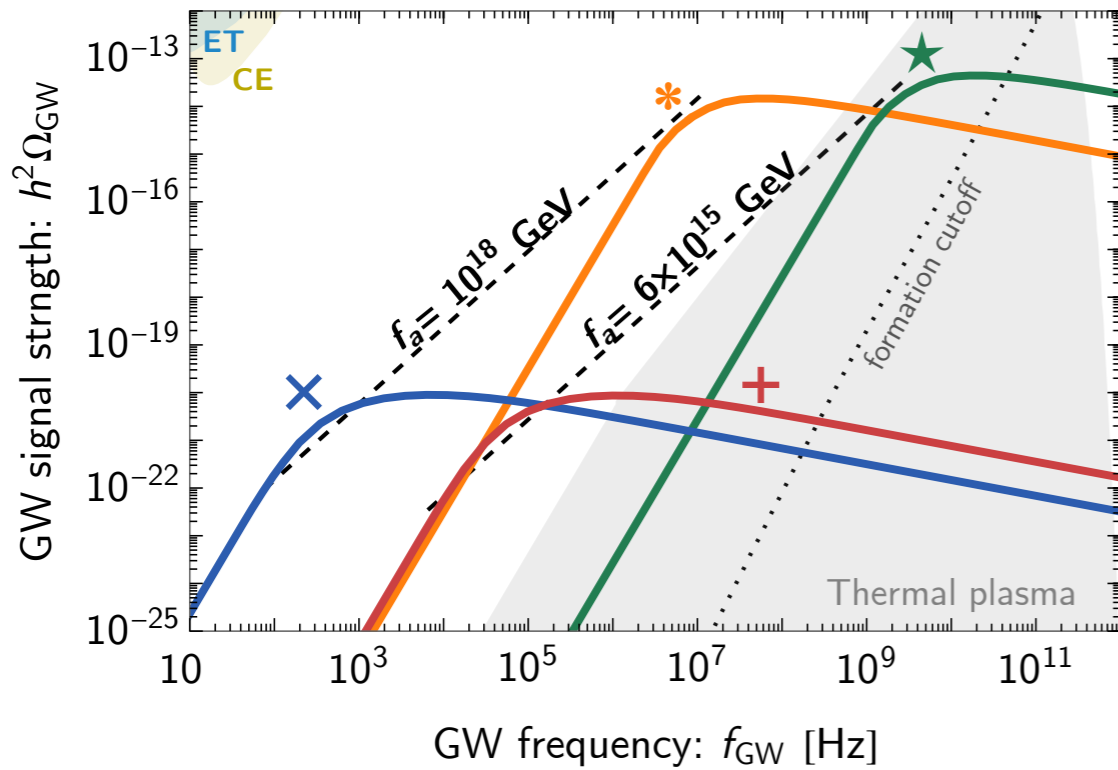
Duration of Axion matter domination

$$\mathcal{B} \equiv \frac{a_{\text{dom}}}{a_{\text{end}}} = \left[\left(\frac{3\sqrt{10}}{64\pi^2} \right) \frac{m_a^3 g_{a\gamma}^2 M_{\text{Pl}}}{g_*^{1/2}(T_{\text{dom}}) T_{\text{dom}}^2} \right]^{\frac{2}{3}} \leq 1,$$

$$g_{a\gamma} = 1.92\alpha_{\text{em}}/(2\pi f_a)$$

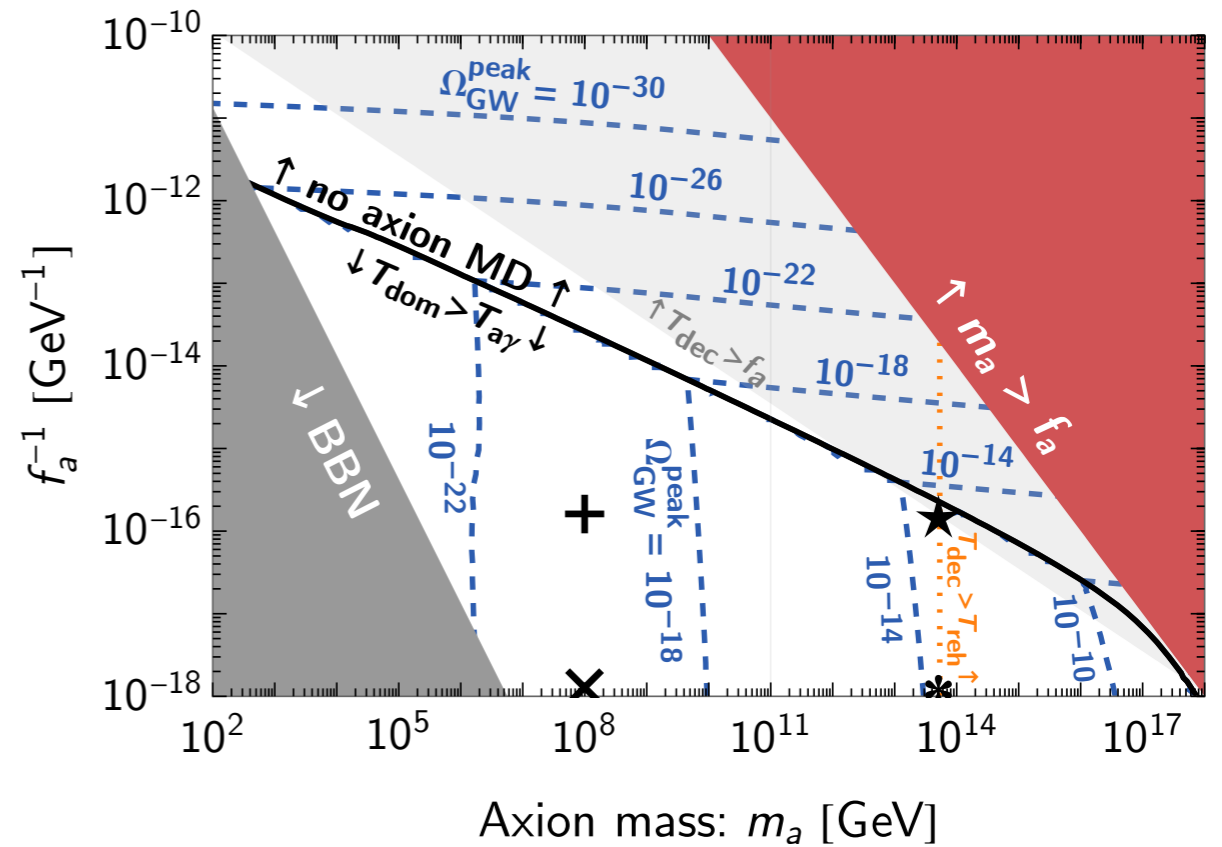
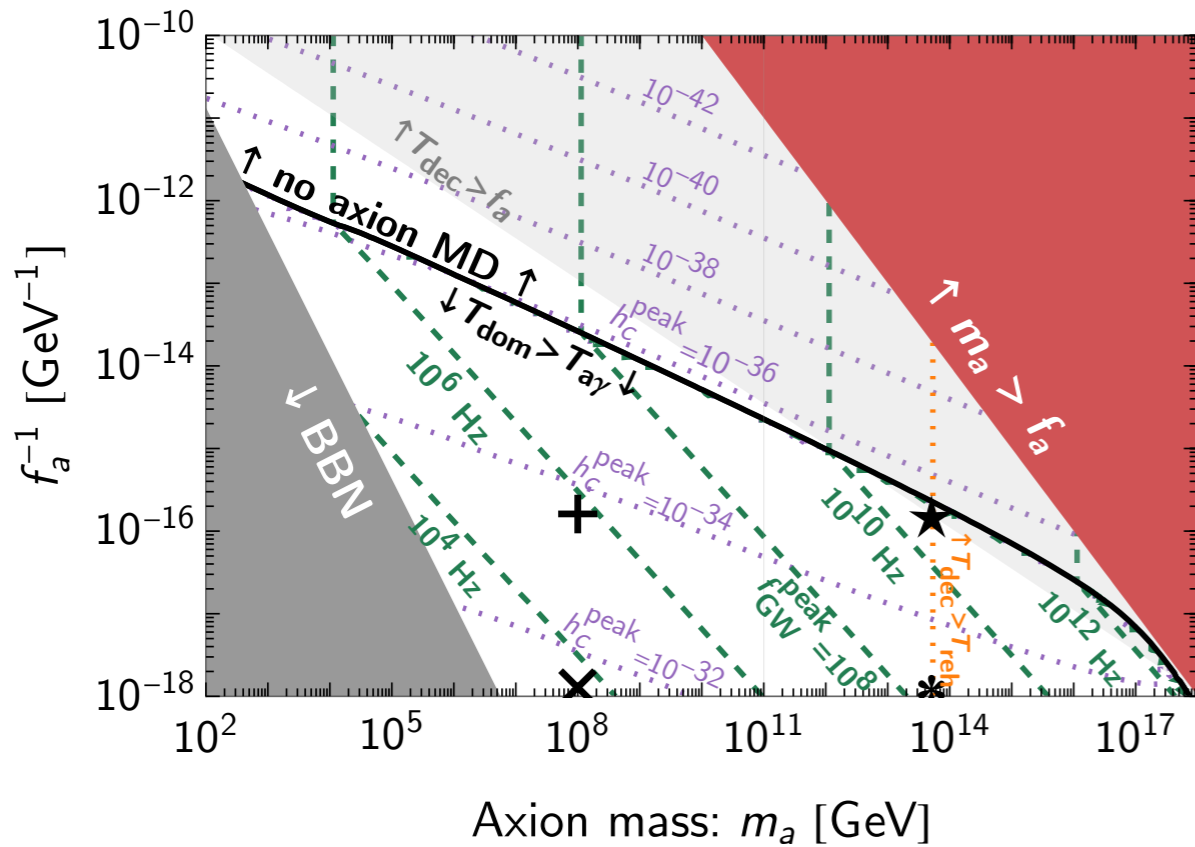
Suppressed UHF GWB from axion strings

Servant, Simakachorn [2312.09281]



$$\Omega_{\text{GW}}(f_{\text{GW}}) = \Omega_{\text{GW}}^{\text{RD}}[f_{\text{GW}}^{\text{RD}}(f_{\text{GW}})] \frac{\mathcal{G}(T_{\text{end}})}{\mathcal{G}(T_{\text{dom}})} \mathcal{B}.$$

$$f_{\text{GW}} = f_{\text{GW}}^{\text{RD}} \left[\frac{\mathcal{G}(T_{\text{end}})}{\mathcal{G}(T_{\text{dom}})} \right]^{\frac{1}{4}} \mathcal{B}^{\frac{1}{4}}.$$



The scale-invariant local-string GWB during radiation-domination (simple argument)

Fraction of energy density
in GW today

$$\Omega_{\text{GW},0} = \left(\frac{\rho_{\text{GW,prod}}}{\rho_{\text{tot},0}} \right) \left(\frac{a_{\text{prod}}}{a_0} \right)^4 = \left(\frac{\rho_{\text{GW,prod}}}{\rho_{\text{tot,prod}}} \right) \left(\frac{\rho_{\text{tot,prod}}}{\rho_{\text{tot},0}} \right) \left(\frac{a_{\text{prod}}}{a_0} \right)^4$$

constant

Long-lasting

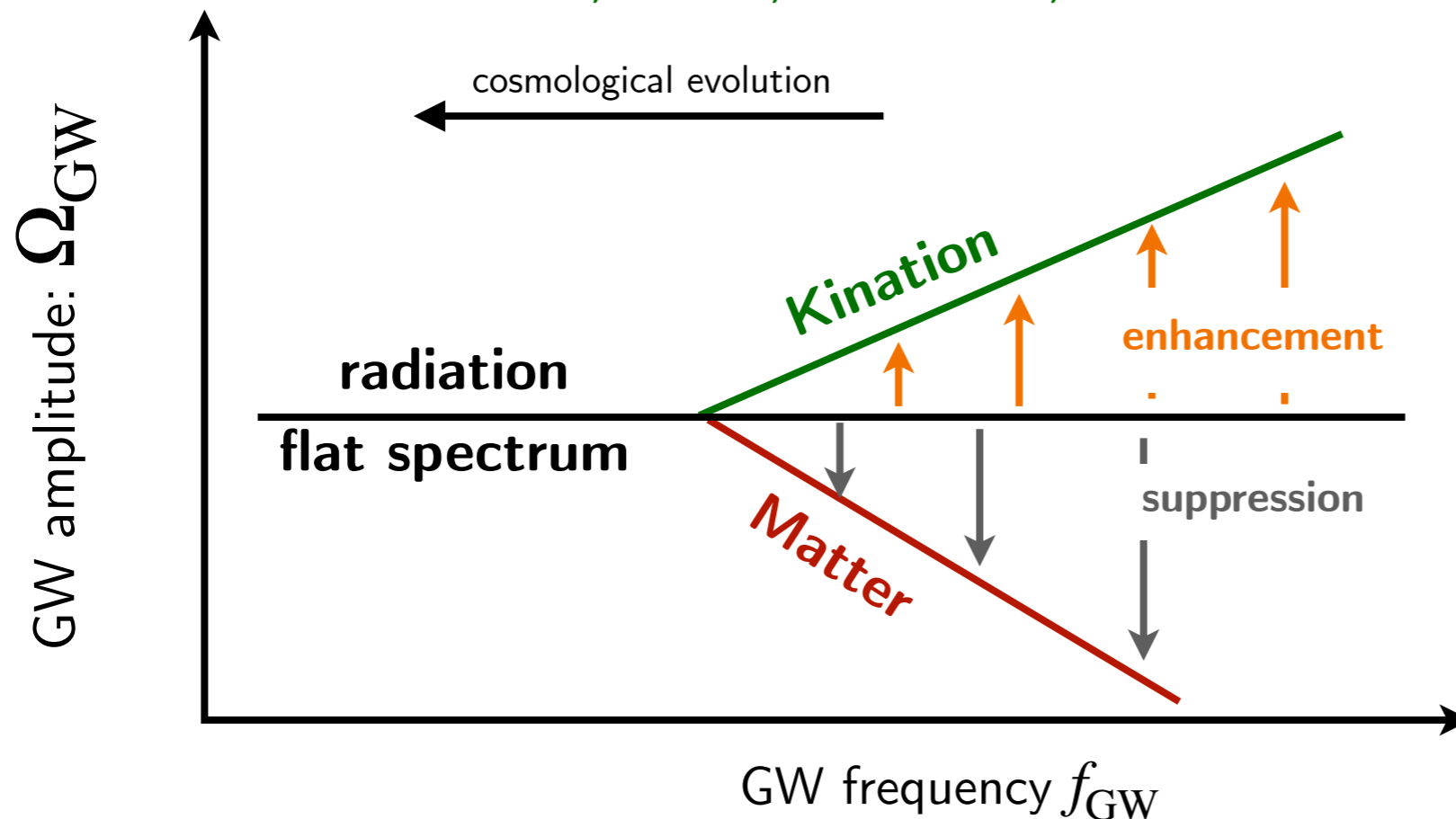
Inflationary GW

(scale-invariant) tensor perturbation: $\Delta_h^2 \simeq (H/M_{\text{Pl}})^2$
 $(\rho_{\text{GW}}/\rho_{\text{tot}})_{\text{prod}} = \text{constant}$

Cosmic-string GW:

$\rho_{\text{GW}} \propto \rho_{\text{string-network}} \propto \rho_{\text{tot}}$
 in the so-called “scaling regime”

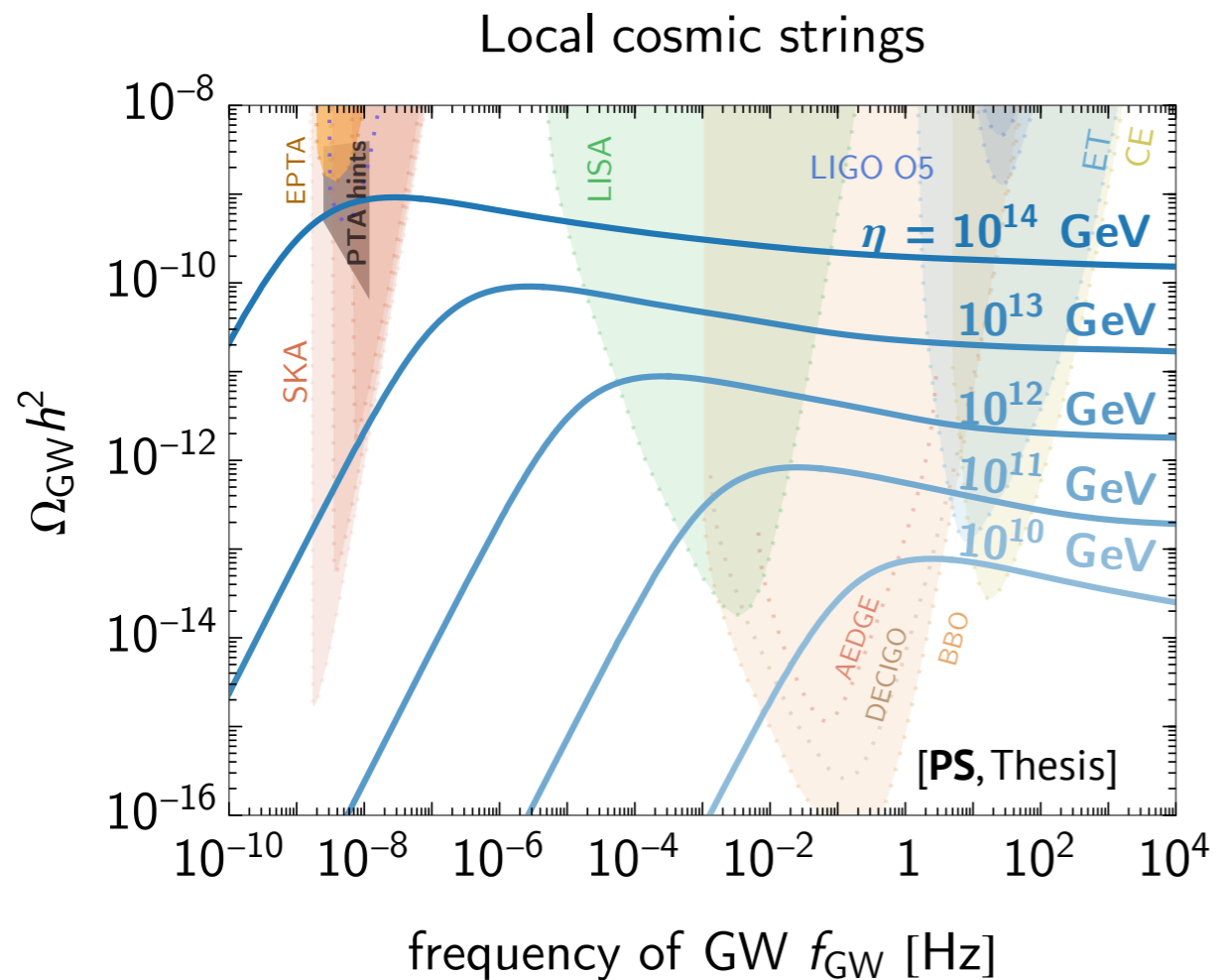
Gouttenoire, Servant, Simakachorn, 2111.01150



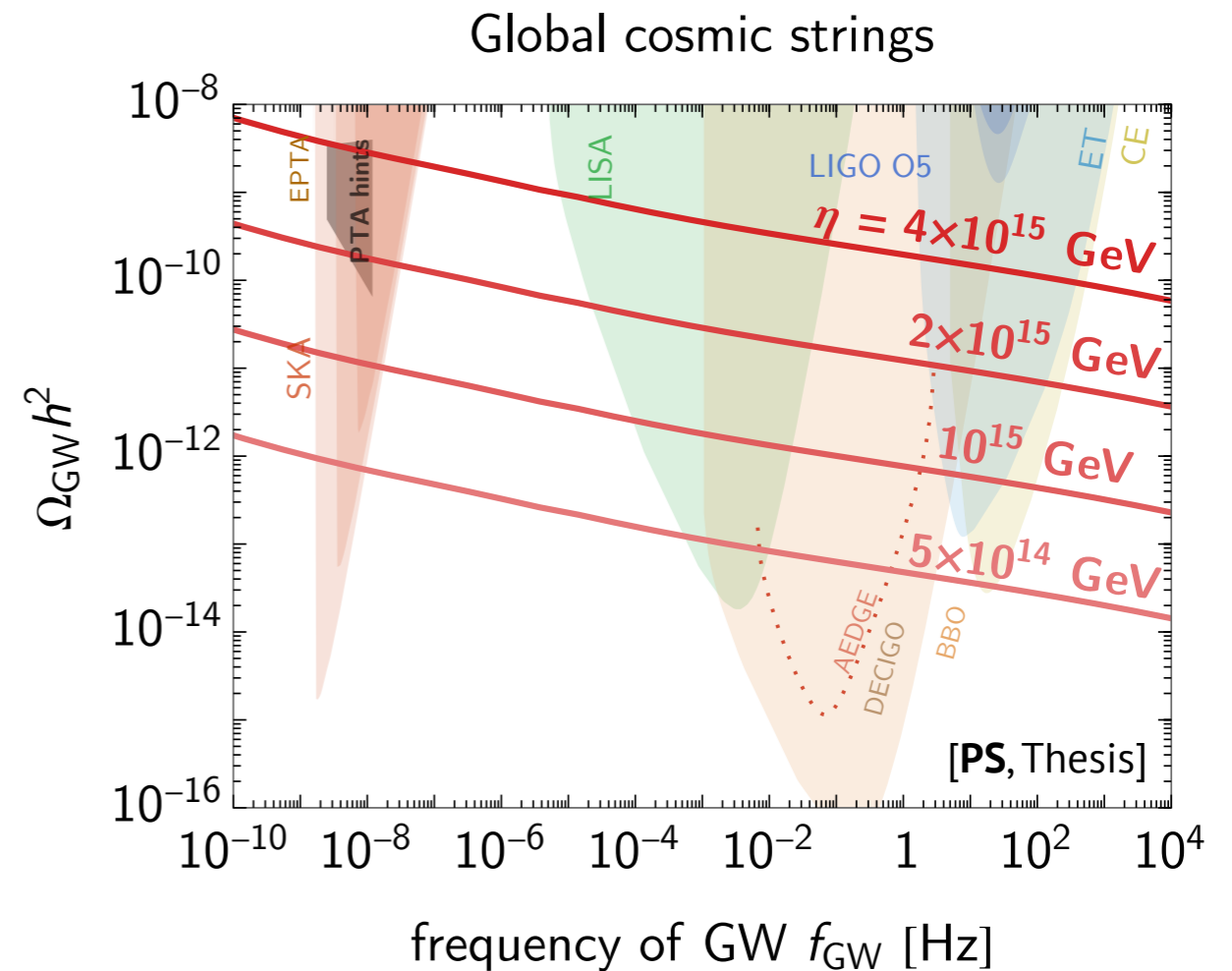
$$\rho_{\text{tot}} \propto \begin{cases} a^{-3} & \text{for matter era,} \\ a^{-4} & \text{for radiation era,} \\ a^{-6} & \text{for kination era} \end{cases}$$

GW from cosmic strings

generated from spontaneous symmetry breaking at an energy scale η



$$\Omega_{\text{GW}}^{\text{local}} \propto \sqrt{G\mu} \propto \eta$$



$$\Omega_{\text{GW}}^{\text{global}} \propto \eta^4$$

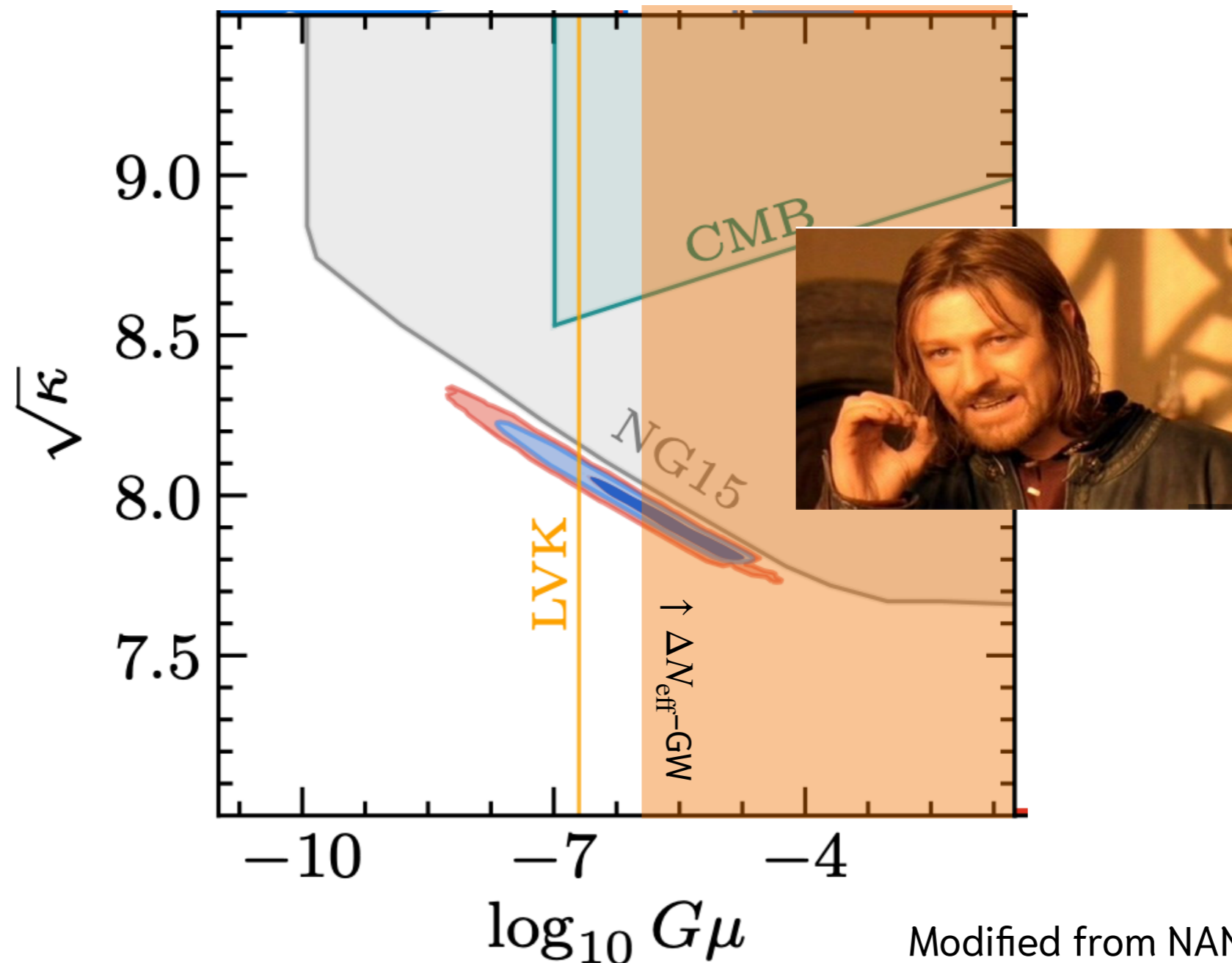
$$f_{\text{GW}}(T) \simeq \begin{cases} (2 \times 10^{-3} \text{ Hz}) \left(\frac{0.1 \times 50 \times 10^{-11}}{\alpha \times \Gamma G\mu} \right)^{1/2} \left(\frac{T}{\text{GeV}} \right) \left[\frac{g_*(T)}{g_*(T_0)} \right]^{1/4} & \text{(local strings),} \\ (4.7 \times 10^{-6} \text{ Hz}) \left(\frac{0.1}{\alpha} \right) \left(\frac{T}{\text{GeV}} \right) \left[\frac{g_*(T)}{g_*(T_0)} \right]^{1/4} & \text{(global strings),} \end{cases}$$

Temperature of the Universe

Local metastable strings can explain PTA data super well?

The best-fit region is excluded by LVK bound,
and on top of that the strings with $G\mu > 10^{-5}$ are in tension with ΔN_{eff} -GW bound

The Bayes factor for explaining the PTA data should be smaller than NG15 analysis.



Modified from NANOGrav [2306.16219]