

Ugo de Noyers

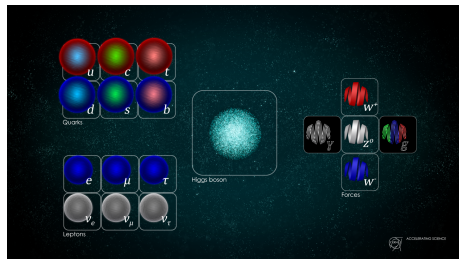
Unraveling Dark Matter and neutrino mysteries with a scotogenic approach

in collaboration with Maud Sarazin
and Björn Herrmann

June 4, 2024



Standard Model and its limitations



$$G_{\text{SM}} = SU(3)_C \times SU(2)_L \times U(1)_Y$$

Advantages

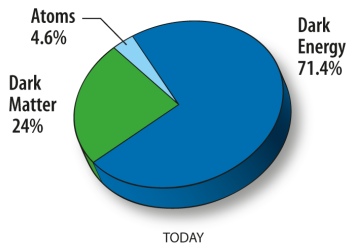
- Higgs boson prediction
- Observables well tested experimentally

Major problems

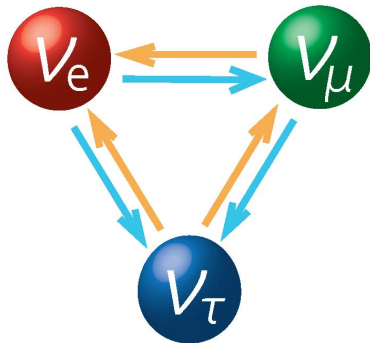
- Gravity not included
- Description of the visible matter only
- Neutrinos remain massless

Unexplained phenomena

Dark Matter nature unknown despite its abundance



Neutrino oscillations linked to them having a mass



Planck measurements^a :

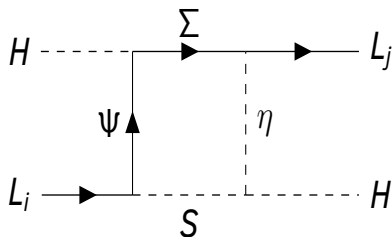
$$\Omega h^2 = 0.1200 \pm 0.0012$$

^a[arXiv:1807.06209v4](https://arxiv.org/abs/1807.06209v4)

$$P(\nu_i \longleftrightarrow \nu_j) \propto \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E} \right)$$

Scotogenic models: a possible extension of the SM

T12G topology



	Ψ_1	Ψ_2	Σ_1	Σ_2	S	η
$SU(3)_c$	1	1	1	1	1	1
$SU(2)_L$	2	2	3	3	1	2
$U(1)_Y$	1	-1	0	0	0	1
$[M]$	3/2	3/2	3/2	3/2	1	1

Same gauge symmetry
as SM

Addition of an extra
symmetry \mathbb{Z}_2 :

SM particles are even
under this symmetry

BSM particles are odd

Why this topology?

- T12A topology already been studied by Björn Herrmann and Maud Sarazin
- T12A model only allows for 2 neutrino masses to be generated
- Variant of T12A with extra fermion singlet can generate 3 neutrino masses
- Phenomenology of a fermion singlet F has been studied so now we wanted to study the one of a fermion triplet

Lagrangian of our model

$$\begin{aligned} -\mathcal{L}_{\text{scalar}} &\supset M_H^2 |H|^2 + \lambda_H |H|^4 + \frac{1}{2} M_S^2 S^2 + \frac{1}{2} \lambda_{4S} S^4 + M_\eta^2 |\eta|^2 + \lambda_{4\eta} |\eta|^4 \\ &+ \frac{1}{2} \lambda_S S^2 |H|^2 + \lambda_\eta |\eta|^2 |H|^2 + \frac{1}{2} \lambda_{S\eta} S^2 |\eta|^2 + \lambda'_\eta |\eta^\dagger H|^2 \\ &+ \frac{1}{2} \lambda''_\eta \left((\eta^\dagger H)^2 + \text{h.c.} \right) + \kappa \left(S \eta^\dagger H + \text{h.c.} \right) \\ -\mathcal{L}_{\text{fermion}} &\supset \frac{1}{2} M_{\Sigma_1} \bar{\Sigma}_1 \Sigma_1 + \frac{1}{2} M_{\Sigma_2} \bar{\Sigma}_2 \Sigma_2 \\ &+ M_\Psi \Psi_1 \Psi_2 + y_{1j} \Psi_1 \Sigma_j H + y_{2j} \bar{\Psi}_2 \Sigma_j H \\ &+ g_\Psi^k \Psi_2 L_k S + g_{\Sigma_j}^k \eta \Sigma_j L_k + g_R^k \tilde{\eta} \Psi_1 e_k^c + \text{h.c.} \end{aligned}$$

MCMC algorithm

36 parameters to scan in the MCMC algorithm \implies consequent computational time to scan a 36-D hypercube

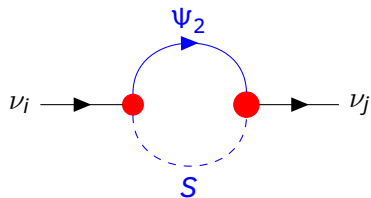
Use of Casas-Ibarra parametrization to take in account experimental constraints in neutrino sector

Likelihood computed as $\mathcal{L}_n = \prod_i \mathcal{L}_i^n$ with $\ln(\mathcal{L}_i^n) = -\frac{(\mathcal{O}_i^n - \mathcal{O}_i^{\text{exp}})^2}{2\sigma_i^2}$

33 experimental constraints were taken in account

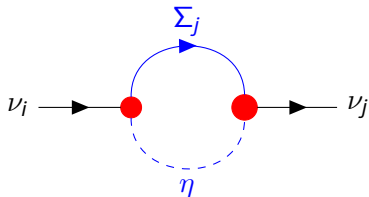
Observables computed thanks to **SPheno** and **micrOMEGAs**

Advantages: mass generation of 3 neutrinos



Couplings determined thanks
to Casas-Ibarra
parametrization

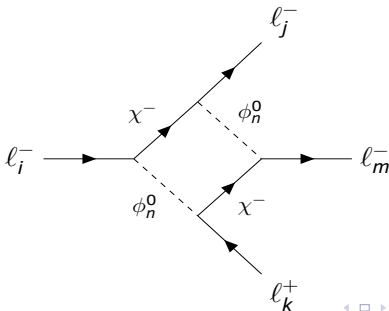
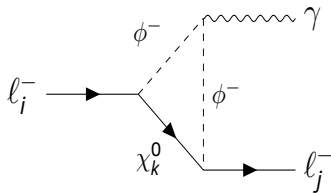
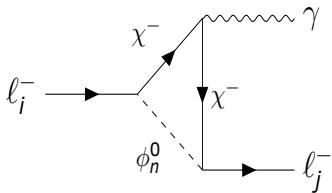
$$M_\nu = g^t M_L \mathcal{G}$$



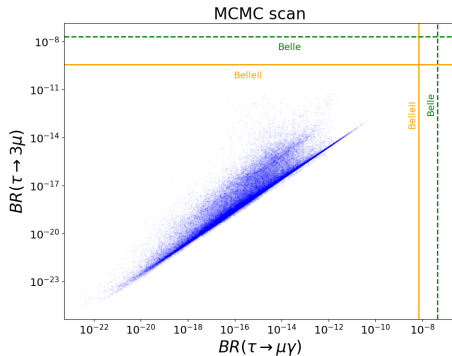
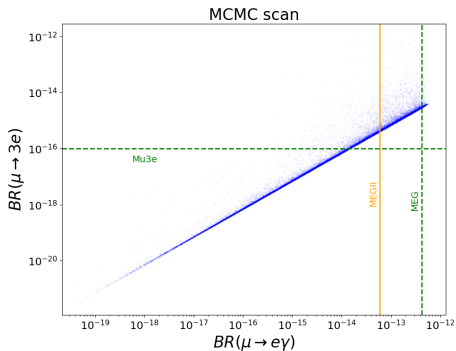
$$\mathcal{G} = \begin{pmatrix} g_\Psi^e & g_\Psi^\mu & g_\Psi^\tau \\ g_{\Sigma_1}^e & g_{\Sigma_1}^\mu & g_{\Sigma_1}^\tau \\ g_{\Sigma_2}^e & g_{\Sigma_2}^\mu & g_{\Sigma_2}^\tau \end{pmatrix}$$

Constraints: Lepton Flavor Violation

$$-\mathcal{L}_{\text{fermion}} \supset g_R^k e_k^c \tilde{\eta} \Psi_1 + (g_\Psi^k \Psi_2 L_k S + g_{\Sigma_j}^k \eta \Sigma_j L_k)$$



Constraints: Lepton Flavor Violation

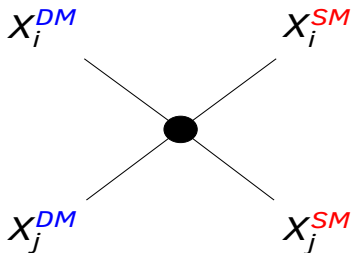


Advantages: stability of Dark Matter particle

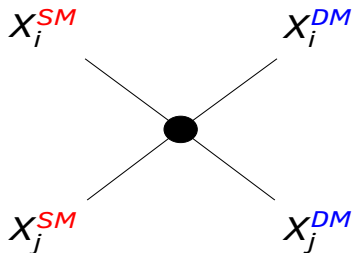
Boltzmann equation in FLRW model

$$\frac{dn}{dt} + 3Hn = -\langle\sigma v\rangle \left(n^2 - (n^{\text{eq}})^2\right)$$

$$\text{with } \Omega h^2 = \frac{nm_{\text{DM}}}{\rho_c}$$

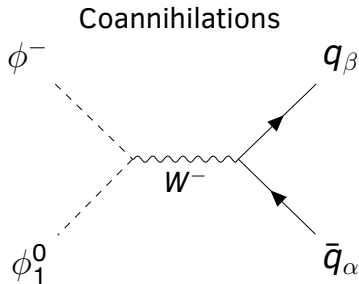
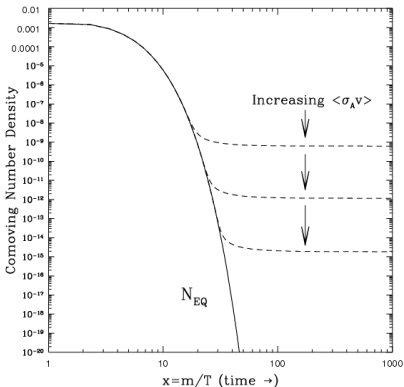


Freeze-out process



Freeze-in process

Satisfying DM relic density



$$(\phi_1^0, \phi_2^0, A^0) = U_\phi (S, \eta^0, A^0)$$

$$\frac{dn}{dt} = -3Hn - \sum_{i,j=1}^N \langle \sigma_{ij} v_{ij} \rangle (n_i n_j - n_i^{\text{eq}} n_j^{\text{eq}})$$

Dark matter candidate?

Diagonalization of mass matrices

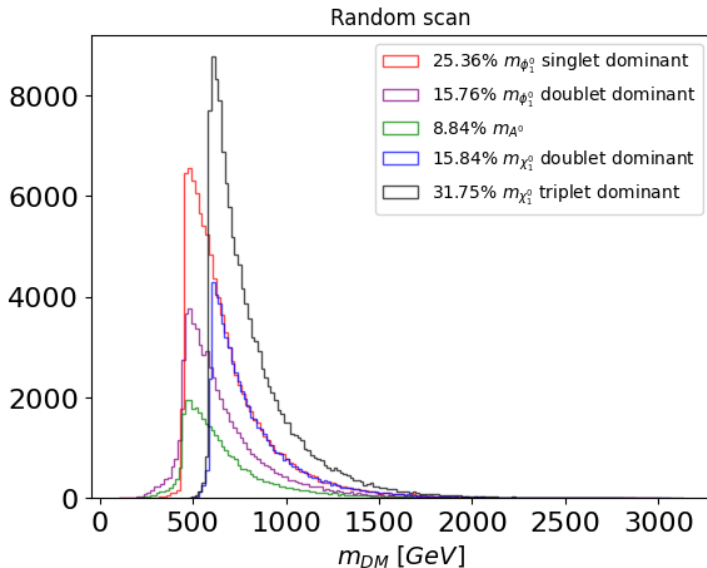
$$\mathcal{M}_{\chi^0} = \begin{pmatrix} M_{\Sigma_1} & 0 & \frac{v}{\sqrt{2}}y_{11} & \frac{v}{\sqrt{2}}y_{21} \\ 0 & M_{\Sigma_2} & \frac{v}{\sqrt{2}}y_{12} & \frac{v}{\sqrt{2}}y_{22} \\ \frac{v}{\sqrt{2}}y_{11} & \frac{v}{\sqrt{2}}y_{12} & 0 & M_{\Psi} \\ \frac{v}{\sqrt{2}}y_{21} & \frac{v}{\sqrt{2}}y_{22} & M_{\Psi} & 0 \end{pmatrix}$$

$$\Rightarrow (\chi_1^0, \chi_2^0, \chi_3^0, \chi_4^0) = U_{\chi} (\Sigma_1^0, \Sigma_2^0, \Psi_1^0, \Psi_2^0)$$

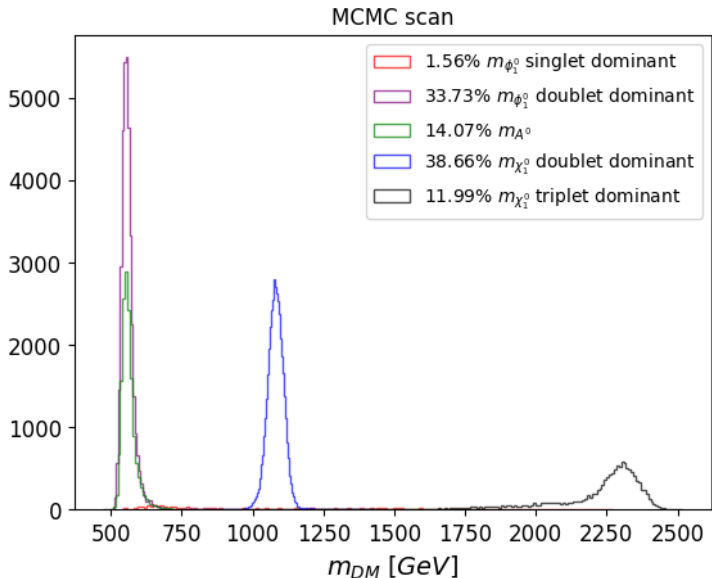
$$\mathcal{M}_{\phi^0}^2 = \begin{pmatrix} M_S^2 + \frac{1}{2}v^2\lambda_S & v\kappa & 0 \\ v\kappa & M_{\eta}^2 + \frac{1}{2}v^2\lambda_L & 0 \\ 0 & 0 & M_{\eta}^2 + \frac{1}{2}v^2\lambda_A \end{pmatrix}$$

$$\Rightarrow (\phi_1^0, \phi_2^0, A^0) = U_{\phi} (S, \eta^0, A^0)$$

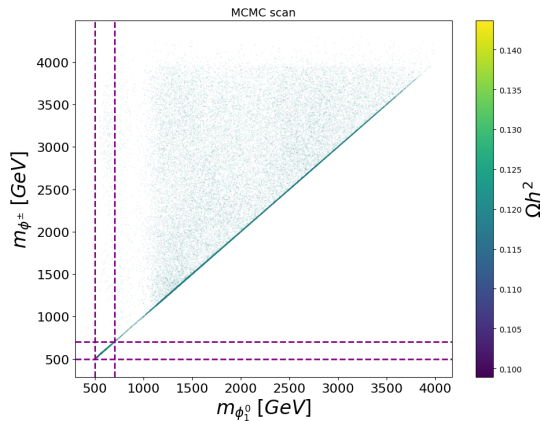
Dark matter candidate



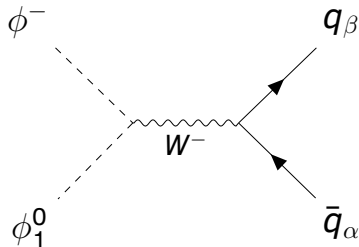
Dark matter candidate



Satisfying DM relic density: scalar sector



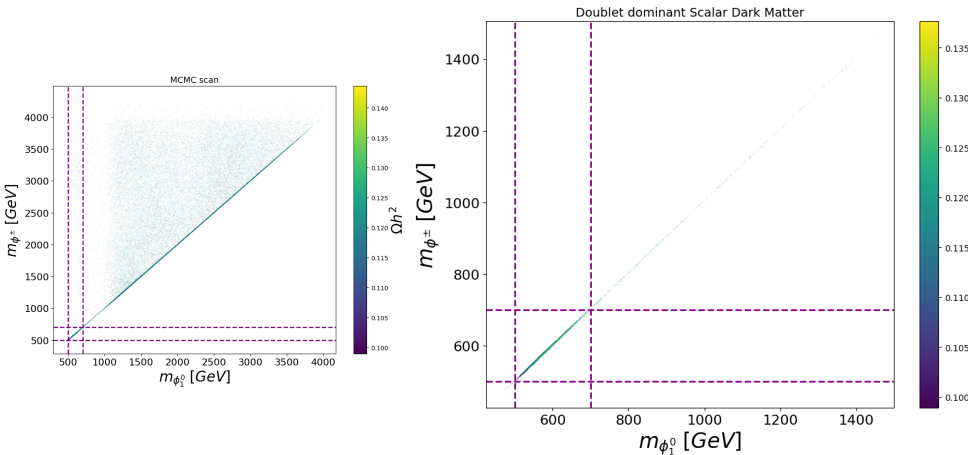
Coannihilations



$$(\phi_1^0, \phi_2^0, A^0) = U_\phi (S, \eta^0, A^0)$$

$$\eta = \begin{pmatrix} \eta^+ \\ \frac{1}{\sqrt{2}} (\eta^0 + iA^0) \end{pmatrix}$$

Satisfying DM relic density: scalar sector



Satisfying DM relic density: fermion sector

Coannihilations

$$(\chi_1^0, \chi_2^0, \chi_3^0, \chi_4^0) = U_{\chi^0} (\Sigma_1^0, \Sigma_2^0, \Psi_1^0, \Psi_2^0)$$

$$(\chi_1^+, \chi_2^+, \chi_3^+) = U_{\chi^+} (\Sigma_1^+, \Sigma_2^+, \Psi_2^+)$$

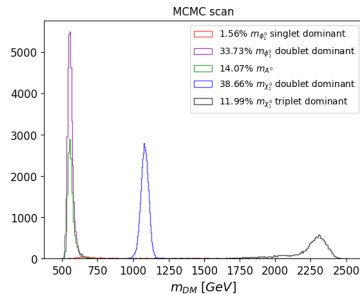
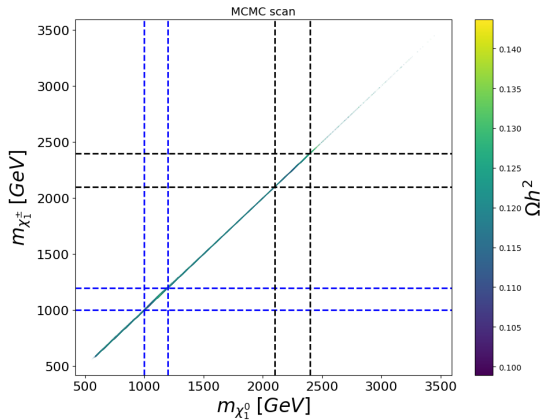
$$(\chi_1^-, \chi_2^-, \chi_3^-) = U_{\chi^-} (\Sigma_1^-, \Sigma_2^-, \Psi_1^-)$$

More sources of
Coannihilations
processes
Direct impact on
triplet mass

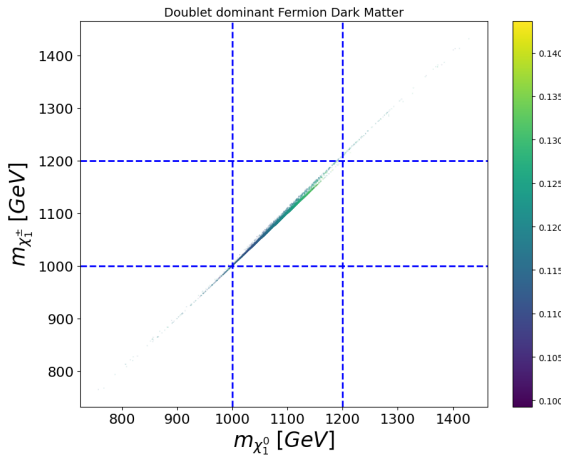
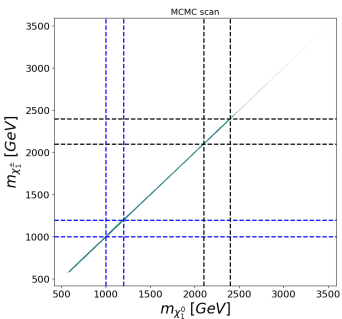
$$\Sigma_j = \begin{pmatrix} \frac{\Sigma_j^0}{\sqrt{2}} & \Sigma_j^+ \\ \Sigma_j^- & -\frac{\Sigma_j^0}{\sqrt{2}} \end{pmatrix}$$

$$\Psi_1 = \begin{pmatrix} \Psi_1^0 \\ \Psi_1^- \end{pmatrix}, \Psi_2 = \begin{pmatrix} -\Psi_2^+ \\ \Psi_2^0 \end{pmatrix}$$

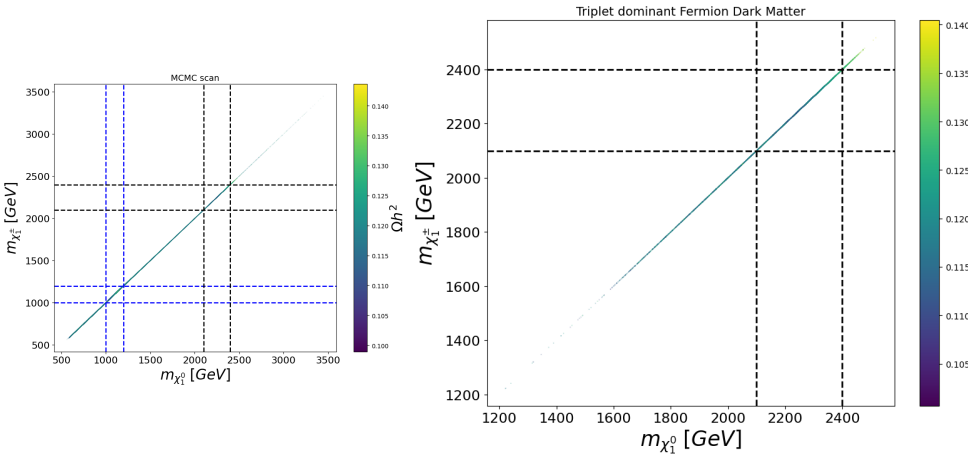
Satisfying DM relic density: fermion sector



Satisfying DM relic density: fermion sector

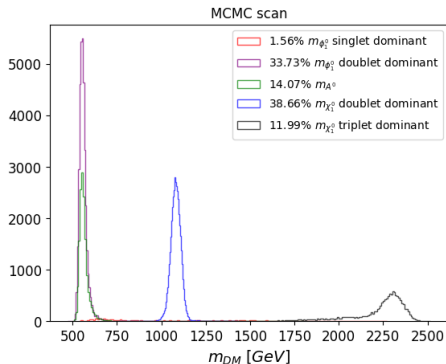


Satisfying DM relic density: fermion sector



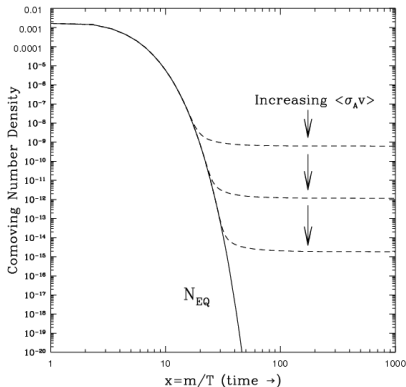
Summary

- Extension of SM that deals with DM and neutrino masses
- Regions in parameter space satisfying major constraints (DM relic density, direct detection, Higgs mass, neutrino mass differences, LFV)
- Mass of BSM particles are reachable at the LHC



Thanks for your attention

Details on Freeze-out



Features

- DM in thermal equilibrium with thermal bath deep within the radiation-dominated epoch
- as $\Gamma \lesssim H$ DM decouples first chemically and then kinematically from the thermal bath
- Correct relic density can be reached by different ways

Casas-Ibarra parametrization

- Diagonalization of the neutrino mass matrix by the PMNS matrix

$$M_\nu = V_{\text{PMNS}}^\dagger D_\nu V_{\text{PMNS}}^*$$

Pontecorvo-Maki-Nakagawa-Sakata (PMNS) matrix expresses the mismatch between the rotations of the LH charged leptons and the neutrinos

$$V_{\text{PMNS}} = V_{eL}^\dagger V_{\nu L}$$

- Decomposition of neutrino mass matrix in different terms

$$M_\nu = \mathcal{G}^t M_L \mathcal{G}$$

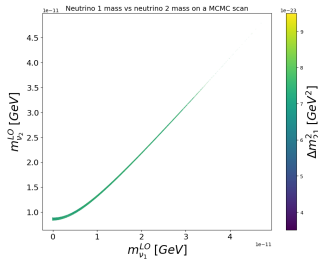
Combining those two statements leads to

$$\mathcal{G} = U_L D_L^{-1/2} R D_\nu^{1/2} V_{\text{PMNS}}^*$$

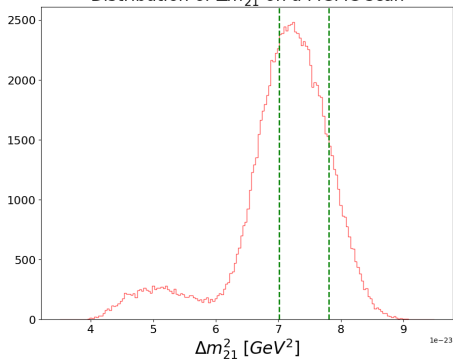
$$\text{with } R^t R = R R^t = \mathbb{I}_3$$

Satisfying neutrino mass differences

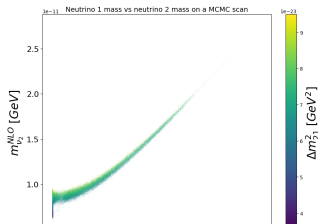
At LO level



Distribution of Δm_{21}^2 on a MCMC scan



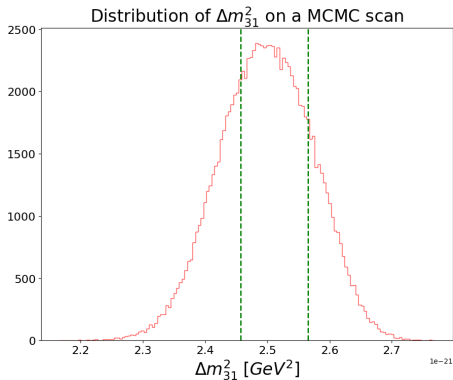
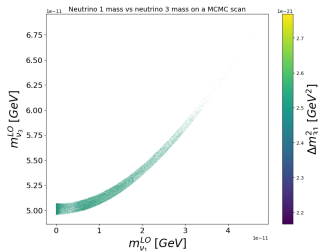
At NLO level



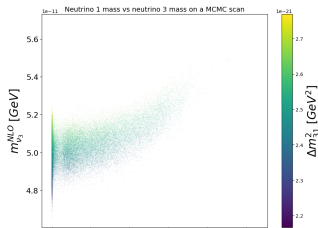
Second bump might be linked to loop corrections

Satisfying neutrino mass differences

At LO level



At NLO level



Clean peak