Accidentally Light Scalars from Large Representations

Giacomo Ferrante



- based on
- JHEP, vol.01, p. 075, 2024 w/ F. Brümmer, M. Frigerio & T. Hambye +
 - arXiv:2406.02531 w/ F. Brümmer & M. Frigerio



- 1. Nambu-Goldstone bosons:
 - SSB: $U(1) \rightarrow Ø$
 - Massless at all orders



Credit: CERN

- Nambu-Goldstone bosons: 1.
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- 2. Pseudo Nambu-Goldstone bosons:
 - SSB + explicit symmetry breaking terms



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- 3. Accidents



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 $\phi = (\phi_1, \phi_2, \phi_3, \phi_4, \phi_5)_1$

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$$V(\phi) = -\mu^2 S + \frac{1}{2} \left[\lambda S^2 + \kappa \left(S^2 - |S'|^2 \right) + \delta A^a \right]$$

No symmetry larger than $SU(2)_D \times U(1)_D$



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VEV:
$$\begin{cases} \langle \phi_1 \rangle = v \sin \alpha \\ \langle \phi_3 \rangle = v \cos \alpha \end{cases}$$

1-loop

$$\longrightarrow V_{\text{eff}}(\alpha) \simeq \frac{\alpha}{6}$$

No symmetry protection





Possible Applications

Dark matter



Accident \equiv Abelian Higgs of U(1)'

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Solution to Little Hierarchy Problem

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Solution to Little Hierarchy Problem

Small field Inflation requires: flat potential

K. Freese, J. A. Freeman, A. V. Olinto Phys.Rev.Lett. 65 (1990) 3233-3236

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 $V = V(\phi) + \frac{\lambda_{\chi}}{4} \left(|\chi|^2 - v_{\chi}^2 \right)^2 + \zeta T_{AB}^a T_{BC}^b \phi^{*A} \phi^C \chi^{*a} \chi^b$



$$V = V(\phi) + \frac{\lambda_{\chi}}{4} \left(|\chi|^2 - v_{\chi}^2 \right)^2 + \zeta T_{AB}^a T_{BC}^b \phi^{*A} \phi$$

Inflation:
$$\begin{cases} \langle \phi \rangle = v(\alpha) \text{ (Accidentally flat of } \alpha) \\ \langle \chi \rangle = 0 \end{cases}$$

"Accidental" Inflation

 $\phi^C \chi^{*a} \chi^b$

direction)



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Inflation: $\begin{cases} \langle \phi \rangle = v(\alpha) \text{ (Accidentally flat direction)} \\ \langle \chi \rangle = 0 \end{cases}$ $=m_{\chi^3}^2(\alpha)$ $V_{\text{inf}} = V_0 + \frac{c_1 v^4}{64\pi^2} \cos(6\alpha) - \frac{1}{2} \left[\mu_{\chi}^2 - \zeta v^2 \sin^2(\alpha) \right] |\chi^3|^2 + \dots$





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Protection from ALL Inflaton = Accident higher-order corrections







 n_s





 $\chi \sim 3$ $\mathrm{SU}(2) \times \mathbb{Z}_4^{(\chi)} : \chi \to i\chi$

 $\chi \sim \mathbf{3}$ SU(2) × $\mathbb{Z}_{4}^{(\chi)}$: $\chi \rightarrow i\chi$





$\chi \sim \mathbf{3}$ $SU(2) \times \mathbb{Z}_{4}^{(\chi)} : \chi \to i\chi$ \downarrow Add a soft breaking of \mathbb{Z}_{4} :

 $i m_{\chi}^2 \chi \chi + h \cdot c$.





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DWs annihilate and emit GWs!















Conclusions



Large representations => tree-level massless scalars



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- Possible applications:
 - (Abelian) Higgs model: little hierarchy problem
 - Dark Matter candidate



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- Possible applications:
 - (Abelian) Higgs model: little hierarchy problem
 - Dark Matter candidate
- "Accidental" Inflation
 - Flatness ensured by gauge symmetries
 - **GWs** from DWs \bullet
 - GWs from Tachyonic Preheating

Thank you for your attention!

Backup Slides





Vacuum Manifold



 $\Delta V_{\rm CW}(\alpha) = \frac{1}{64\pi^2} \operatorname{Str}\left(\mathscr{M}(\alpha)^4 \log \frac{\mathscr{M}(\alpha)^2}{\Lambda^2}\right)$

Fermions: $\psi \sim \mathbf{3}_{+1/2}, \quad \xi \sim \mathbf{3}_{-1/2}$

 $\mathscr{L} \supset y \left(\psi^T \phi \psi + \chi^T \phi^* \chi \right) + M \psi^T \chi + h.c.$

Effective Potential



Possible Applications Abelian Higgs



$V_{\rm eff}(\alpha) \simeq c_1 \cos(6\alpha) + c_2 \cos(12\alpha)$ \downarrow Tuning c_1 against c_2 : breaking of U(1)' at a scale $v' \ll v$

We can identify the accident with the Abelian Higgs

Possible Applications Dark Matter





The SU(3) ten-plet

$$V = -\mu^2 S + \frac{1}{2}(\lambda S^2 + \delta A^a A^a)$$

 $SU(3) \times U(1) \longrightarrow U(1)_3 \times U(1)_8$ & 6 accidents • ESP: • Generic point: $SU(3) \times U(1) \longrightarrow Ø$ & 2 accidents

Scalar one-loop corrections -----> The ESP is stabilised

Invariant ONLY under $SU(3) \times U(1)$

"Accidental" Inflation "Real" Model

 $\phi \sim 5, \quad \chi \sim 3$

 $G = \operatorname{SO}(3) \times \mathbb{Z}_2^{(\phi)}$

 $V = -\frac{1}{2}\mu_{\phi}^2\phi^2 - \frac{1}{2}\mu_{\chi}^2\chi^2 + \frac{\lambda_{\phi}}{\Lambda}(\phi^2)^2 + \frac{\lambda_{\chi}}{\Lambda}(\chi^2)^2 + \frac{\varepsilon}{\Lambda}\phi^2\chi^2 + \frac{\zeta}{\Lambda}T_{AC}^aT_{CB}^b\phi_A\phi_B\chi^a\chi^b$

No Topological Defect Production

Parameter Space



 n_r

 n_r

Accidental Inflation GWs from Tachyonic Preheating

End of inflation: $m_{\chi^3}^2 < 0$

Tachyonic Preheating

G. N. Felder et al. *Phys. Rev. Lett.* 87, 011601 (2001) G. N. Felder et al. *Phys. Rev. D 64, 123517 (2001)*

Large "bubbly" inhomogeneities: $R_* \sim$

$$\frac{1}{k_*}$$

$$\begin{aligned} & \iint_{\substack{Phys. \; Rev. \; D, \; vol. \; 76, \; p. \; 123517, \; 2007.}} \\ & \\ \text{GWs:} & \begin{cases} & \nu_* \simeq 4 \times 10^{10} \; \text{Hz} \frac{k_*}{\rho_{\inf}^{1/4}} \\ & & h^2 \Omega_* \simeq 10^{-6} \left(\frac{H_{\inf}}{k_*}\right)^2 \end{cases} \end{aligned}$$






$(G\mu)^{\text{CMB}} \lesssim 10^{-7} \Longrightarrow V_0 \lesssim \text{few} \times 10^{14} \text{GeV}$

Cosmic Strings





 $V = -\frac{1}{2}\mu_{\phi}^{2}\phi^{2} + \frac{\lambda_{\phi}}{4}(\phi^{2})^{2} - \mu_{\chi}^{2}\chi^{*}\chi + \lambda_{\chi}(\chi^{*}\chi)^{2} + \delta\chi^{*2}\chi^{*}\chi^{2}$

 $f_p \simeq 1.6 \times 10^{-7}$

 $h^2 \Omega_{\rm GW} \left(f_p \right) \simeq 1.6 \times 10^{-5}$

$$\alpha_{\rm ann} \equiv \frac{\rho_{\rm DW}(t_{\rm ann})}{\rho_r(t_{\rm ann})} \simeq \frac{4}{3} C_d \mathcal{A}^2 \frac{\sigma^2}{M_P^2 \Delta V}$$

Domain Walls

$$\chi^2 + \frac{1}{2} \left(\kappa \chi^2 \chi^2 + \text{h.c.} \right) + \frac{\varepsilon}{2} \phi^2 \left(\chi^* \chi \right) + \frac{\zeta}{2} T^a_{AC} T^b_{CB} \phi_A \phi_B \chi$$

$$\operatorname{Hz}\left(\frac{g_*(T_{\mathrm{ann}})}{100}\right)^{1/6} \frac{T_{\mathrm{ann}}}{\operatorname{GeV}}$$
$$\left(\frac{100}{g_*(T_{\mathrm{ann}})}\right)^{1/3} \frac{3}{32\pi} \tilde{\epsilon} \,\alpha_{\mathrm{ann}}^2 \,\mathcal{S}(f/f_p)$$

$$T_{\rm ann} = \left[\frac{45}{2\pi^2} \frac{g_*(T_{\rm ann})^{-1}}{C_d^2 \mathcal{A}^2} \frac{M_{\rm P}^2 \Delta V^2}{\sigma^2}\right]^{1/4}$$

