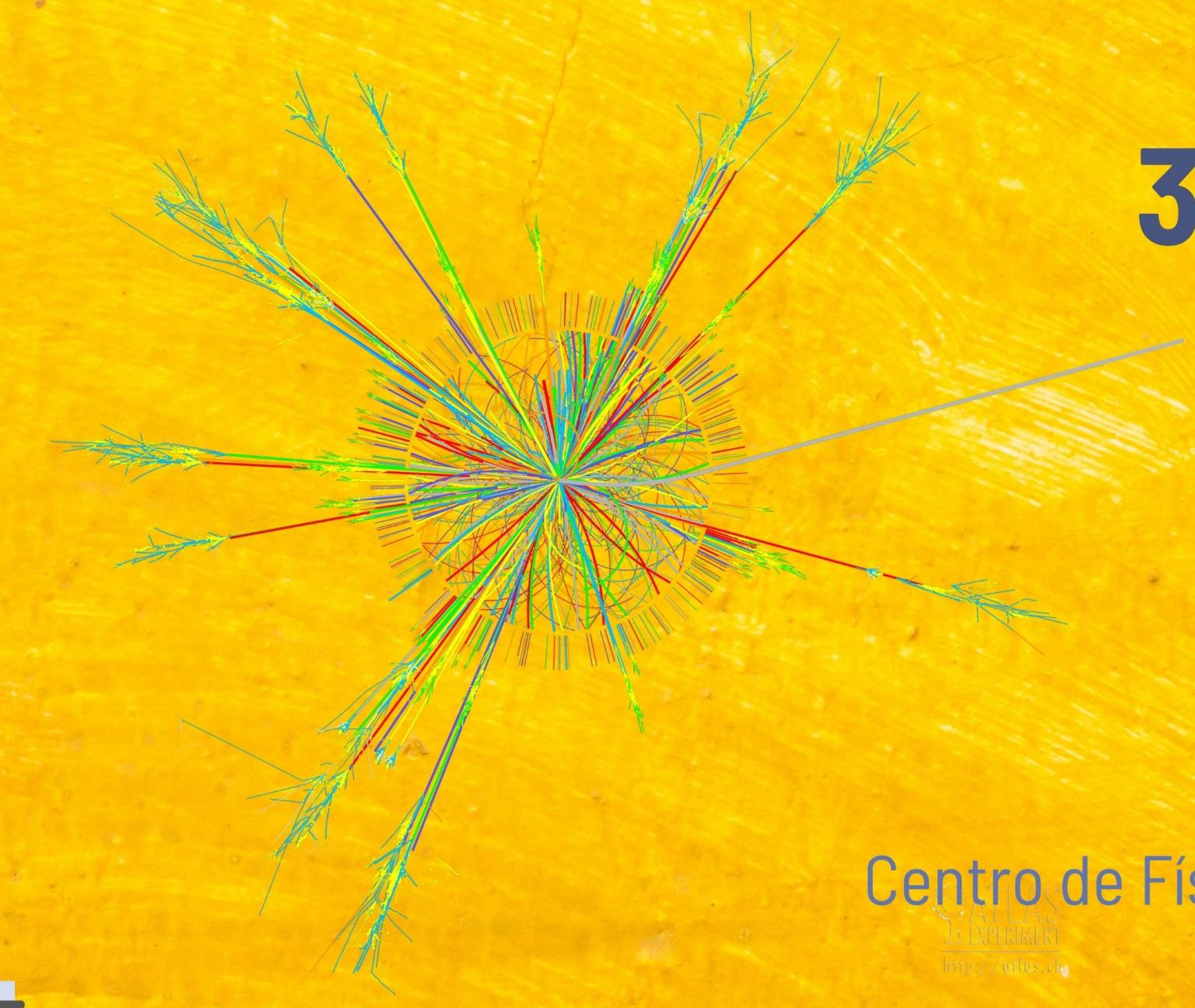


Gravitational Waves as Probes of New Physics



PLANCK2024

26th Conference “From the Planck Scale to
the Electroweak Scale”



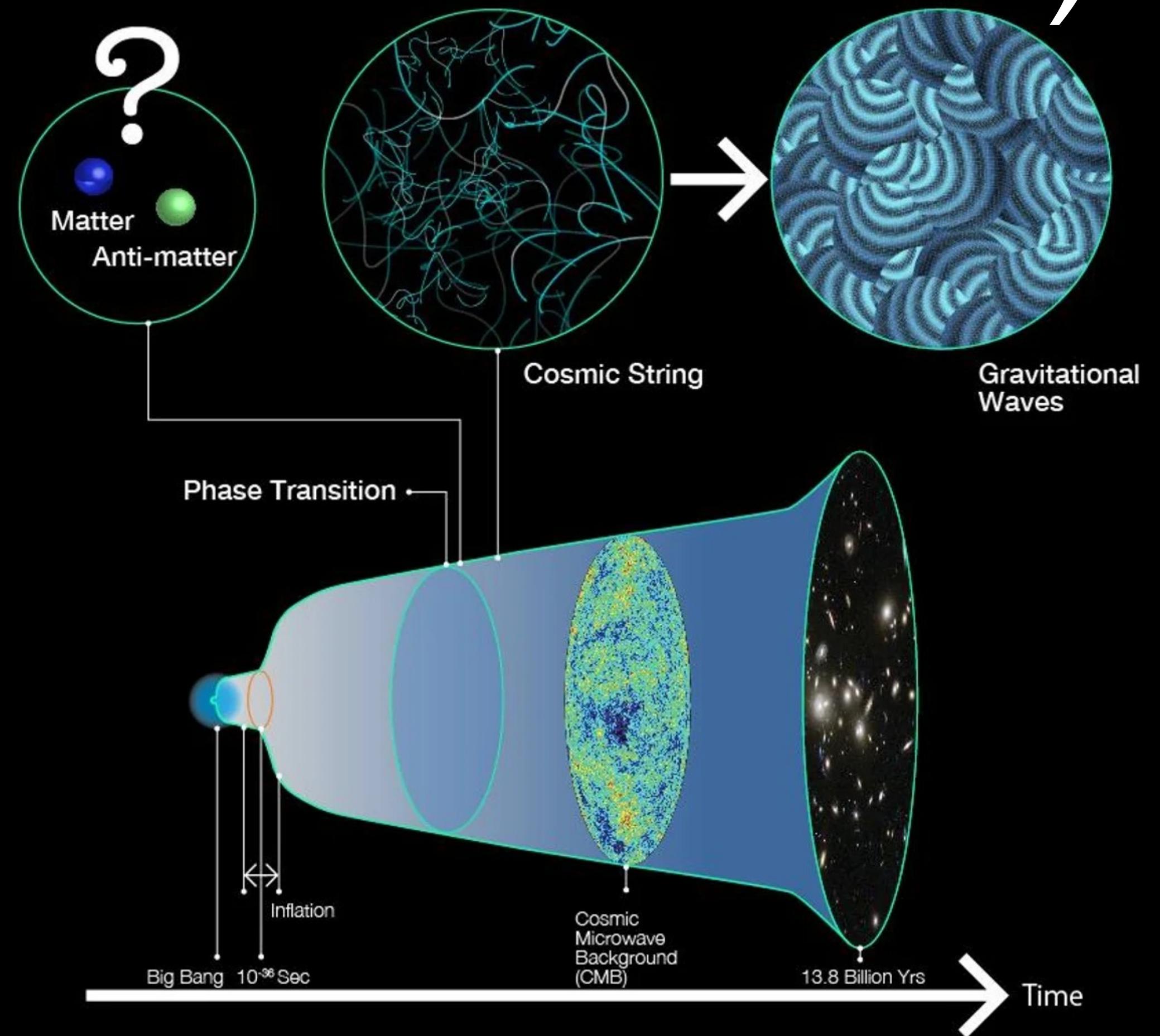
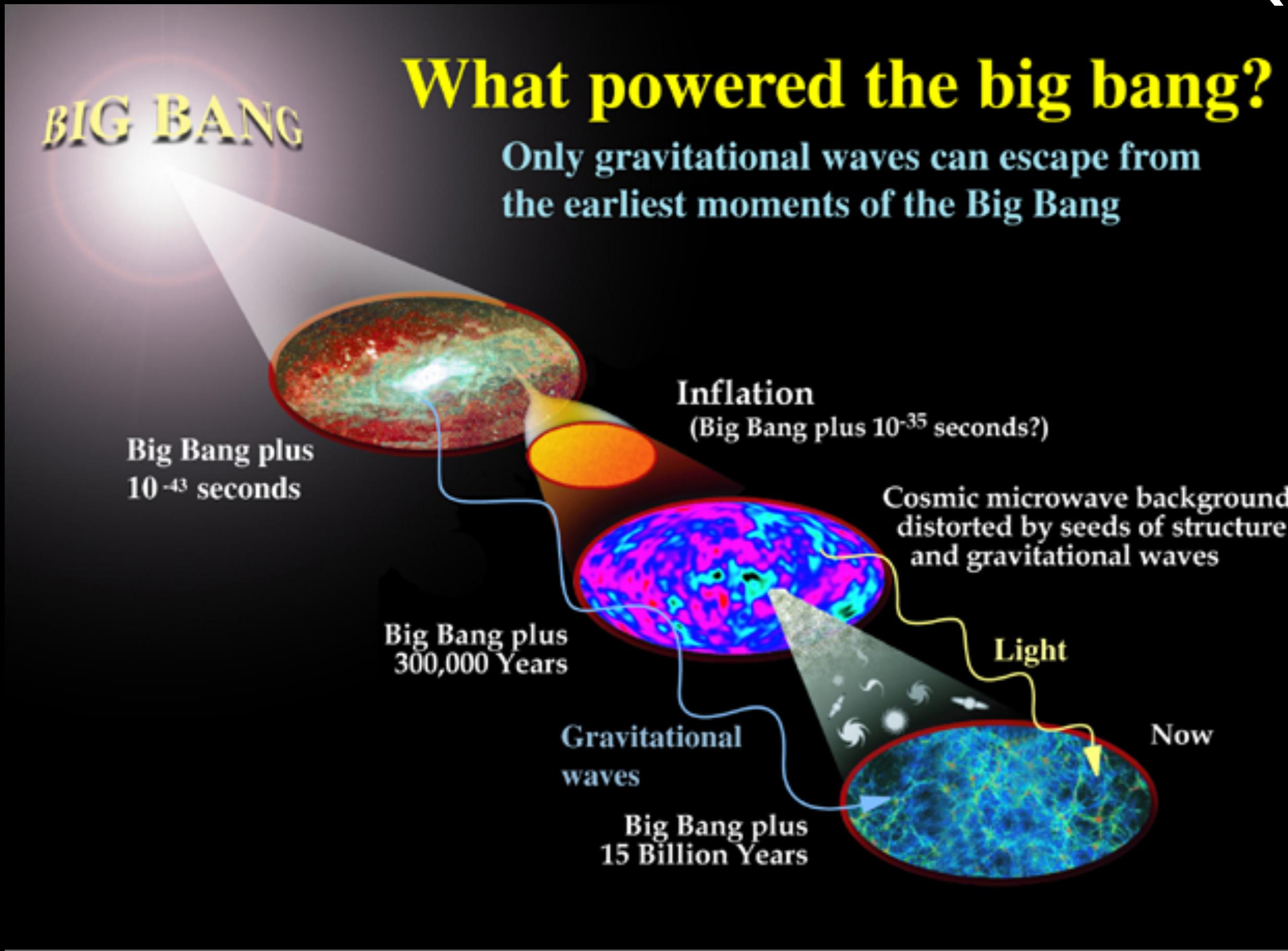
STEVE KING

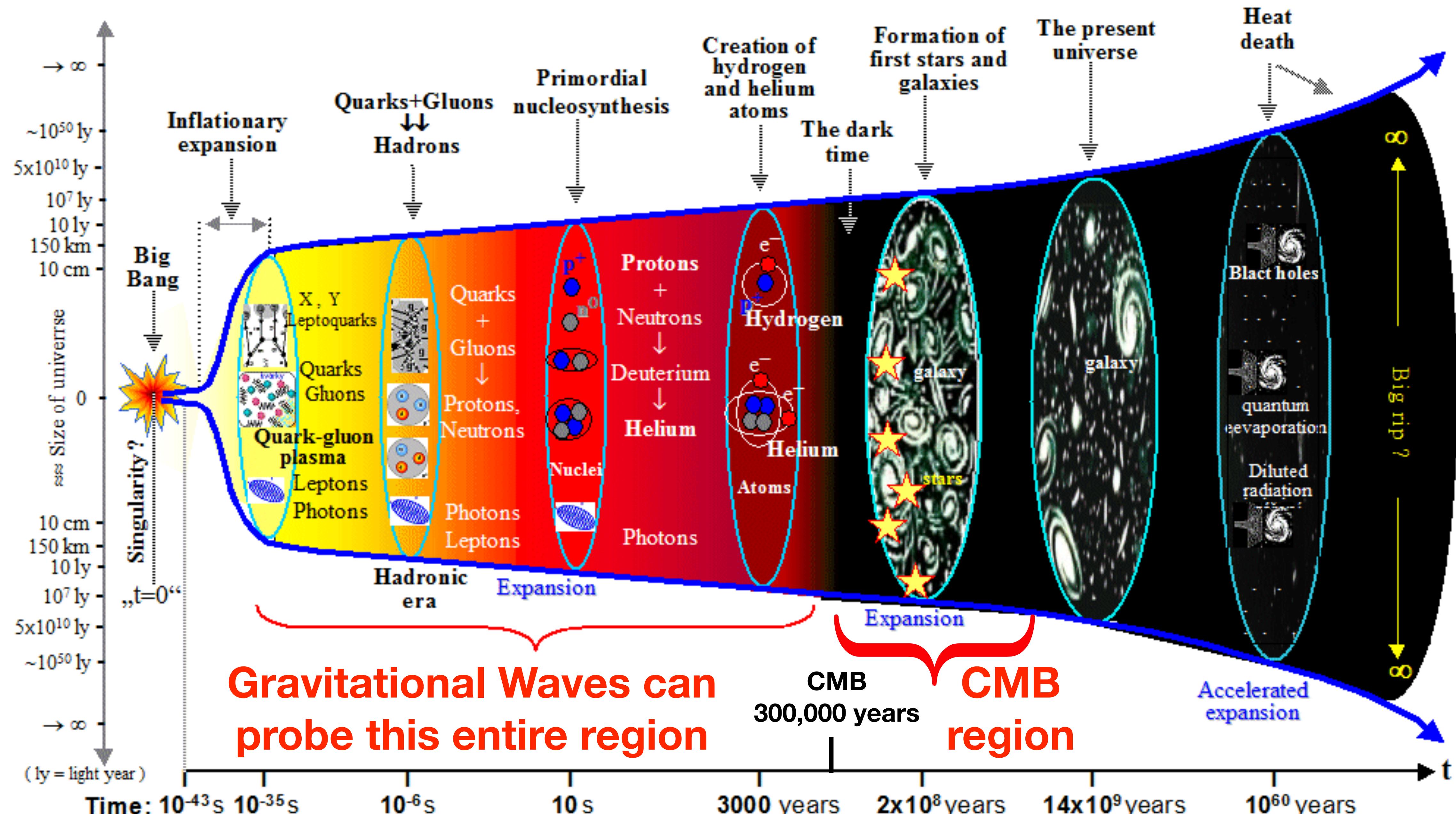
3-7 JUNE, 2024

Anfiteatro Abreu Faro,
Instituto Superior Técnico
Lisbon, Portugal

Organised by
Centro de Física Teórica de Partículas (CFTP)
Centro de Física Teórica de Partículas
cftp.ist.utl.pt

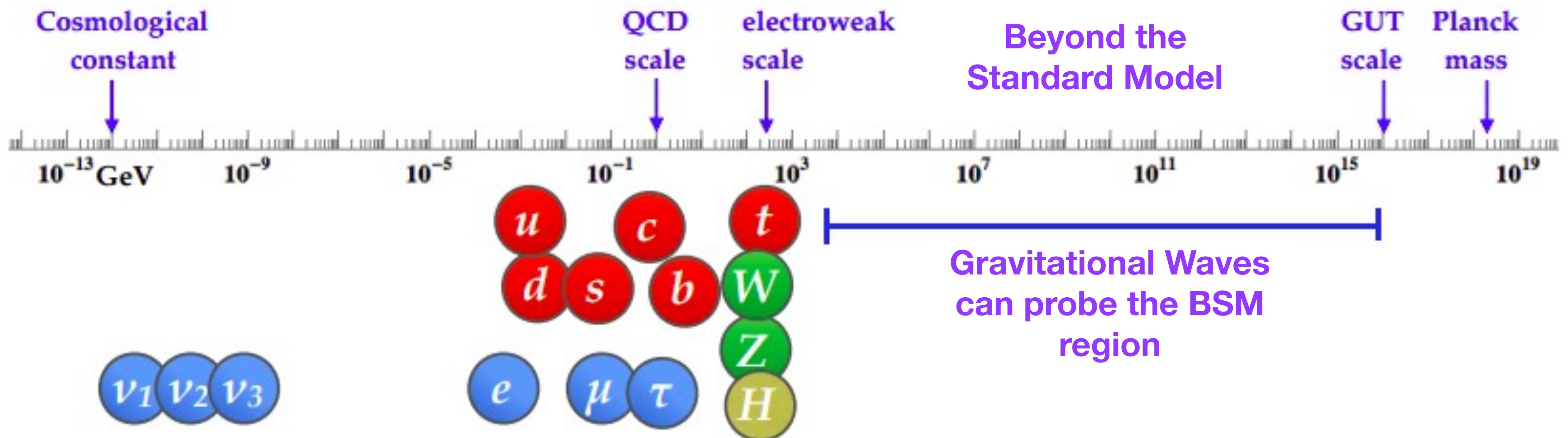
Gravitational Waves enable us to look back to the earliest moments of the Universe (back to Inflation)



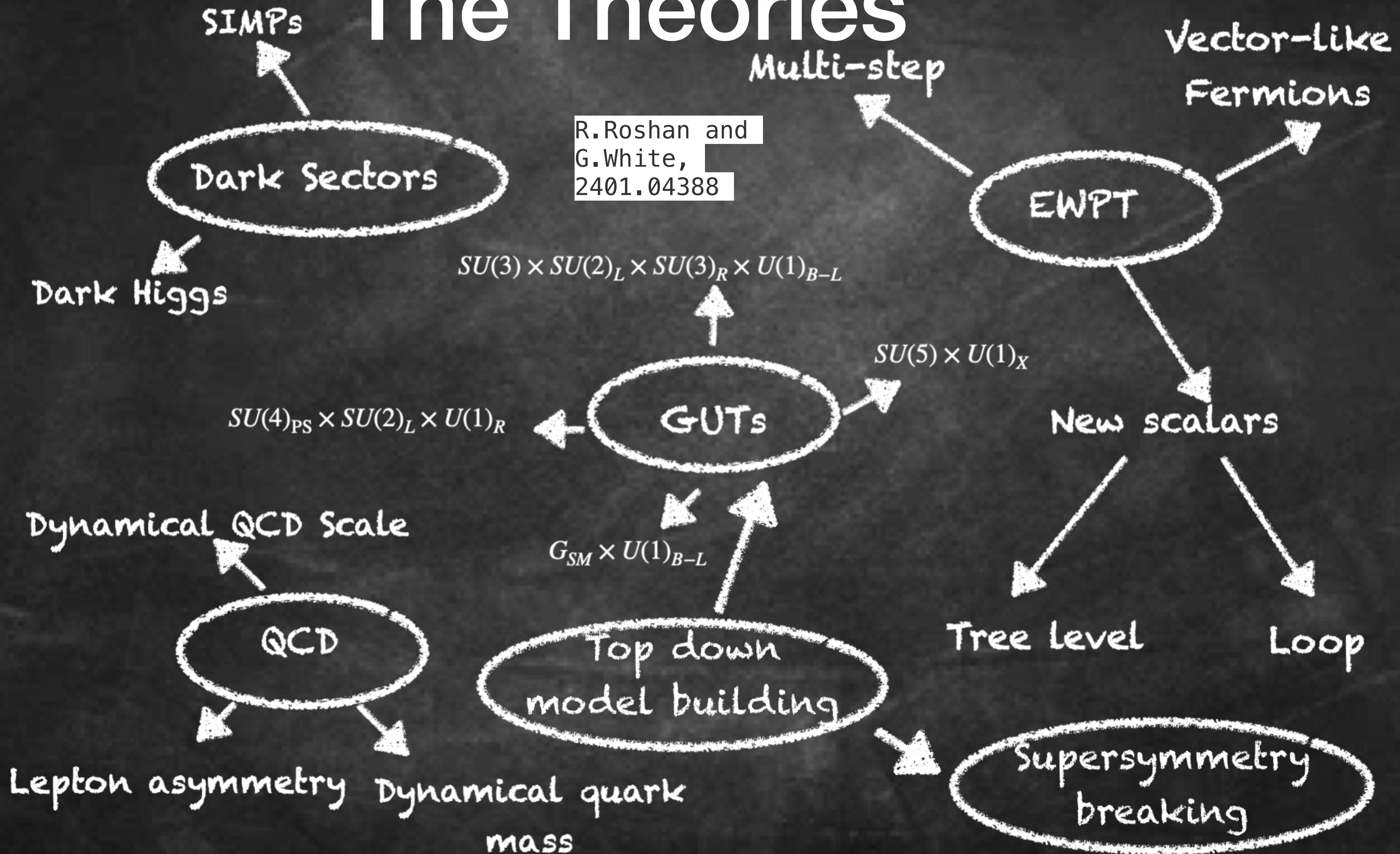


| Time: 10^{-43} s | Temperature: $10^{32} \text{ }^{\circ}\text{K}$ | Energy: 10^{19} GeV | Time: 10^{-35} s | Temperature: $10^{27} \text{ }^{\circ}\text{K}$ | Energy: 10^{15} GeV | Time: 10^{-6} s | Temperature: $10^{13} \text{ }^{\circ}\text{K}$ | Energy: 100 MeV | Time: 10 s | Temperature: $5 \times 10^9 \text{ }^{\circ}\text{K}$ | Energy: 200 keV | Time: 3000 years | Temperature: $3000 \text{ }^{\circ}\text{K}$ | Energy: $0,4 \text{ eV}$ | Time: $2 \times 10^8 \text{ years}$ | Temperature: $30 \text{ }^{\circ}\text{K}$ | Energy: $0,1 \text{ eV}$ | Time: $14 \times 10^9 \text{ years}$ | Temperature: $2,7 \text{ }^{\circ}\text{K}$ | Energy: 10^{-4} eV | Time: 10^{60} years | Temperature: $10^{-10} \text{ }^{\circ}\text{K}$ | Energy: ≈ 0 |
|----------------------------|---|-------------------------------|----------------------------|---|-------------------------------|---------------------------|---|---------------------------|----------------------|---|---------------------------|----------------------------|--|--------------------------|-------------------------------------|--|--------------------------|--------------------------------------|---|------------------------------|-------------------------------|--|---------------------|
| 10^{-43} s | $10^{32} \text{ }^{\circ}\text{K}$ | 10^{19} GeV | 10^{-35} s | $10^{27} \text{ }^{\circ}\text{K}$ | 10^{15} GeV | 10^{-6} s | $10^{13} \text{ }^{\circ}\text{K}$ | 100 MeV | 10 s | $5 \times 10^9 \text{ }^{\circ}\text{K}$ | 200 keV | 3000 years | $3000 \text{ }^{\circ}\text{K}$ | $0,4 \text{ eV}$ | $2 \times 10^8 \text{ years}$ | $30 \text{ }^{\circ}\text{K}$ | $0,1 \text{ eV}$ | $14 \times 10^9 \text{ years}$ | $2,7 \text{ }^{\circ}\text{K}$ | 10^{-4} eV | 10^{60} years | $10^{-10} \text{ }^{\circ}\text{K}$ | ≈ 0 |

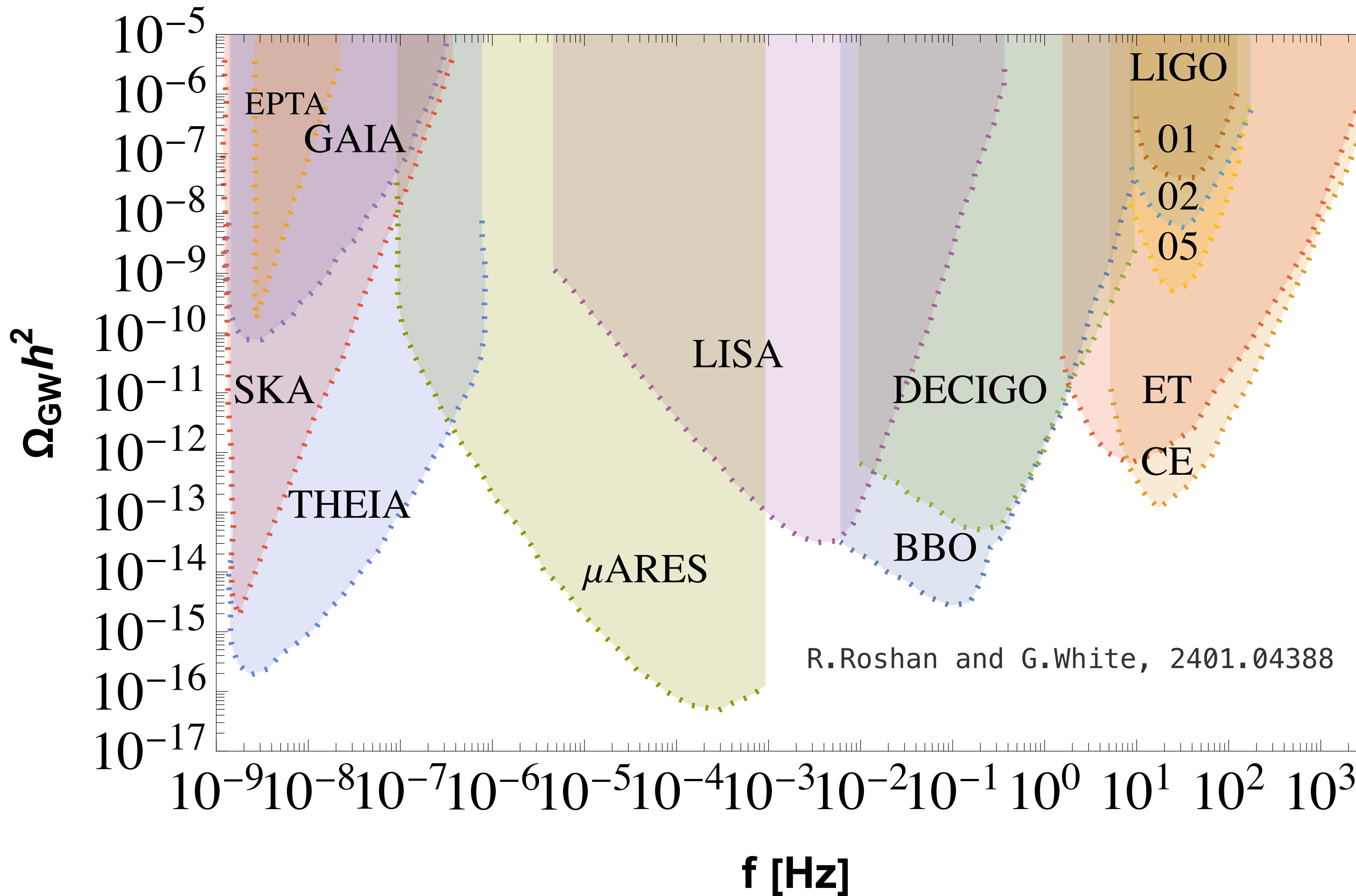
Gravitational Waves are sensitive to scales up to the Planck scale (well beyond the reach of colliders)



The Theories



The Detectors



Current observations are from LIGO 01-03

The stage is set for a bonanza of new results

Gravitational Waves are sensitive to :

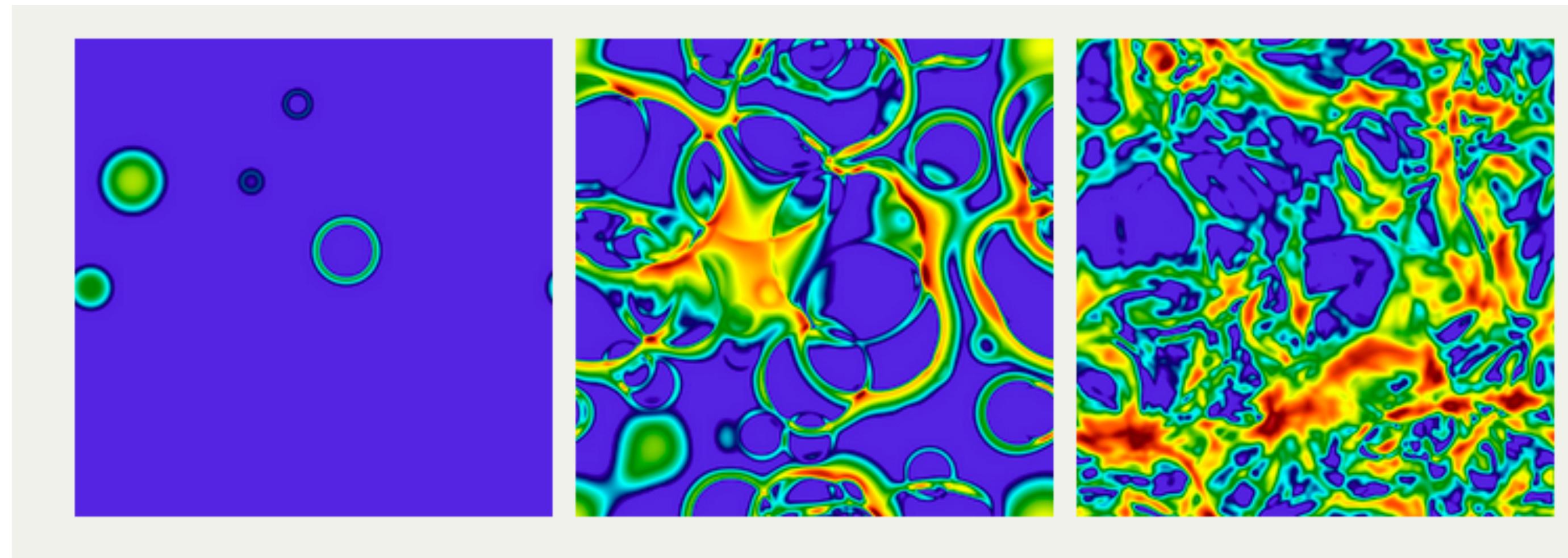
- First Order Phase transitions FOPT (e.g. QCD)
- Cosmic Strings CS (e.g. U(1) sym breaking)
- Domain Walls DW (e.g. Z_2 sym breaking)
- Inflation (e.g. with a kink or hybrid)
- Many other effects (e.g. PBHs,...)

In this talk we
are interested
in a few BSM
examples

Gravitational Waves from First Order Phase Transitions

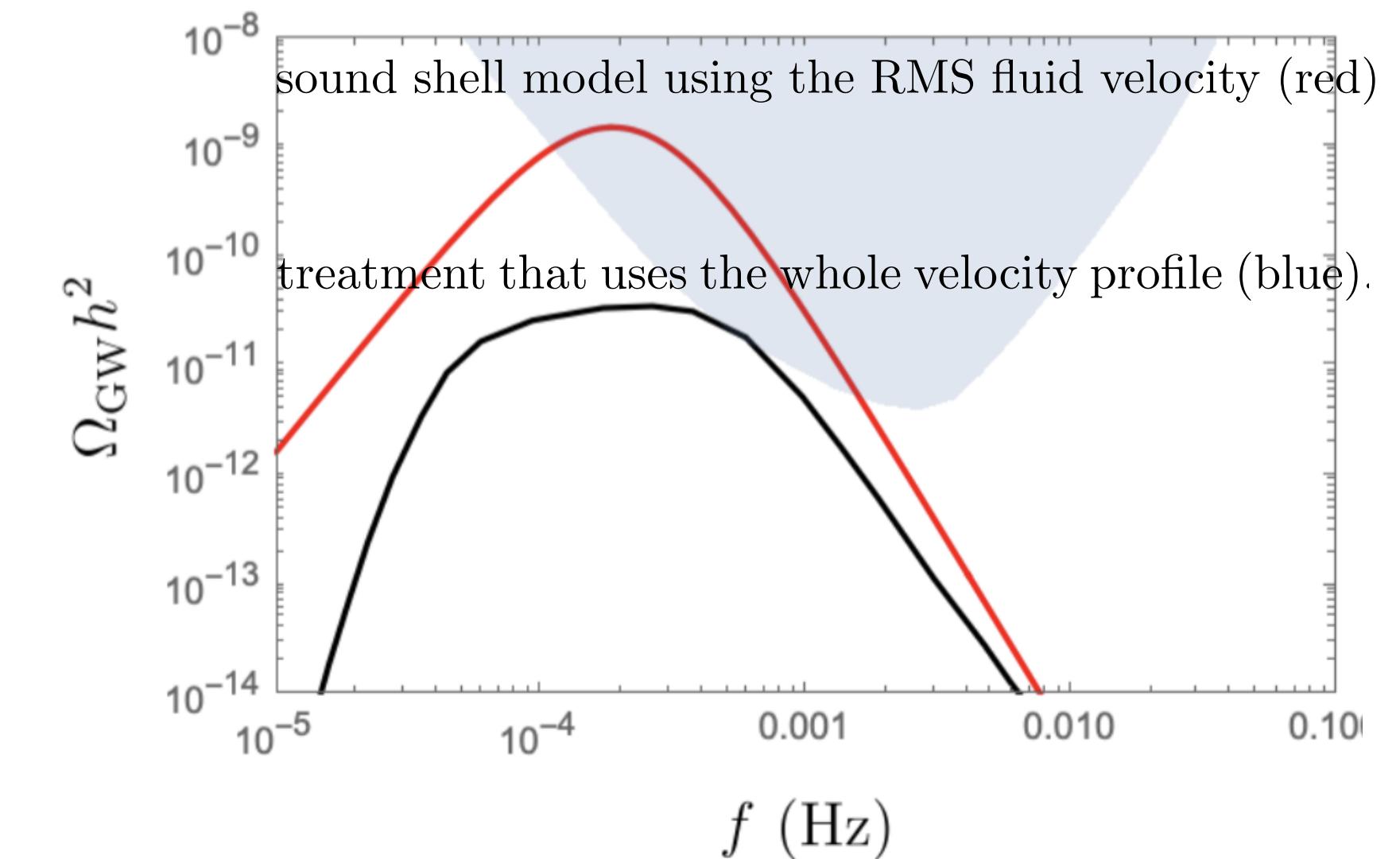
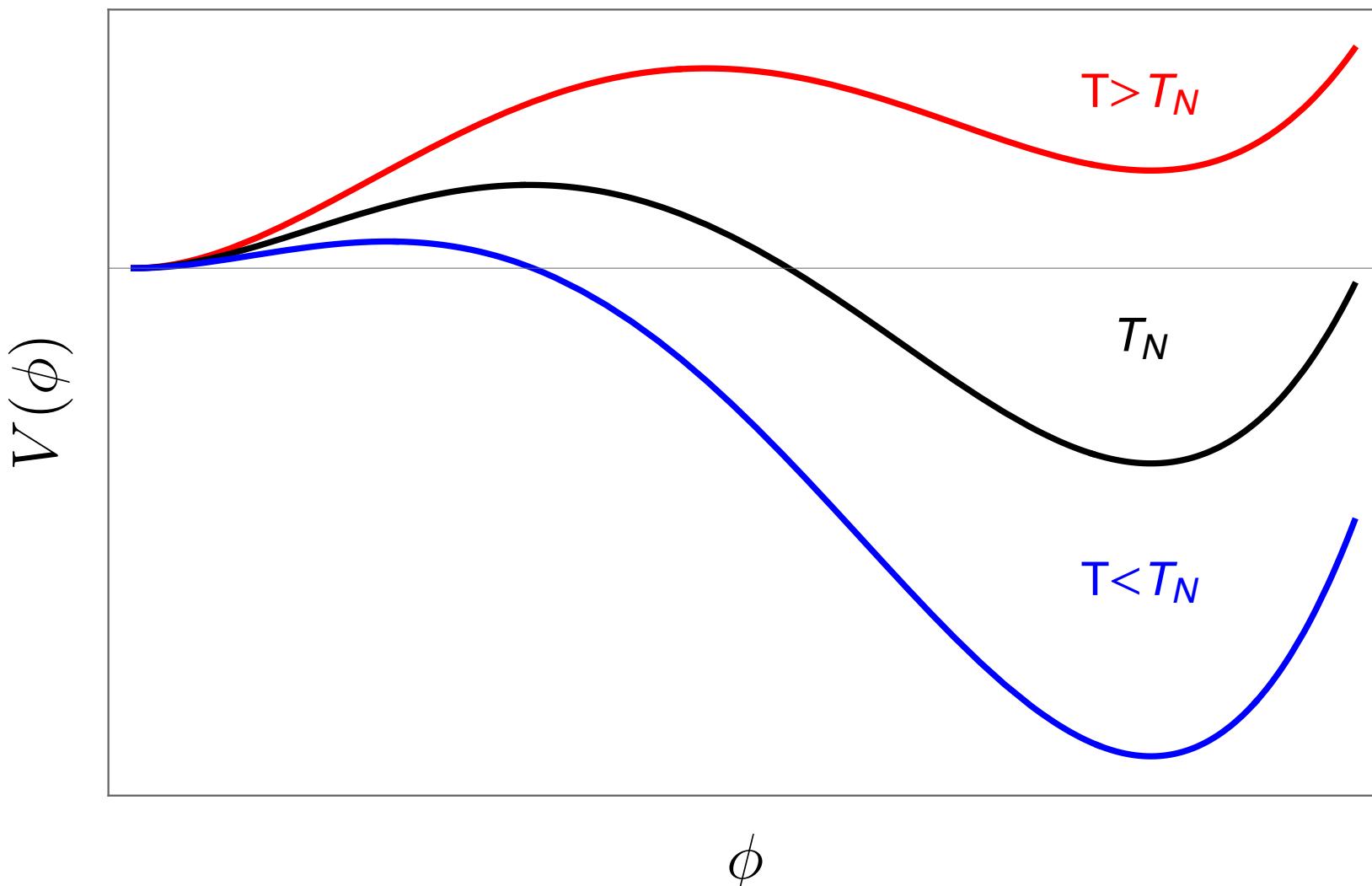
Phase Transitions:

- Bubbles nucleate and grow.
- Expand in plasma.
- Bubbles and fronts collide - - violent process.
- Sound Waves left behind in thermal plasma.
- Turbulence, damping.



$$\Omega_{\text{tot}}(f) = \Omega_{\text{coll}}(f) + \Omega_{\text{SW}}(f) + \Omega_{\text{turb}}(f)$$

See talk yesterday
by Pasquale Di Bari



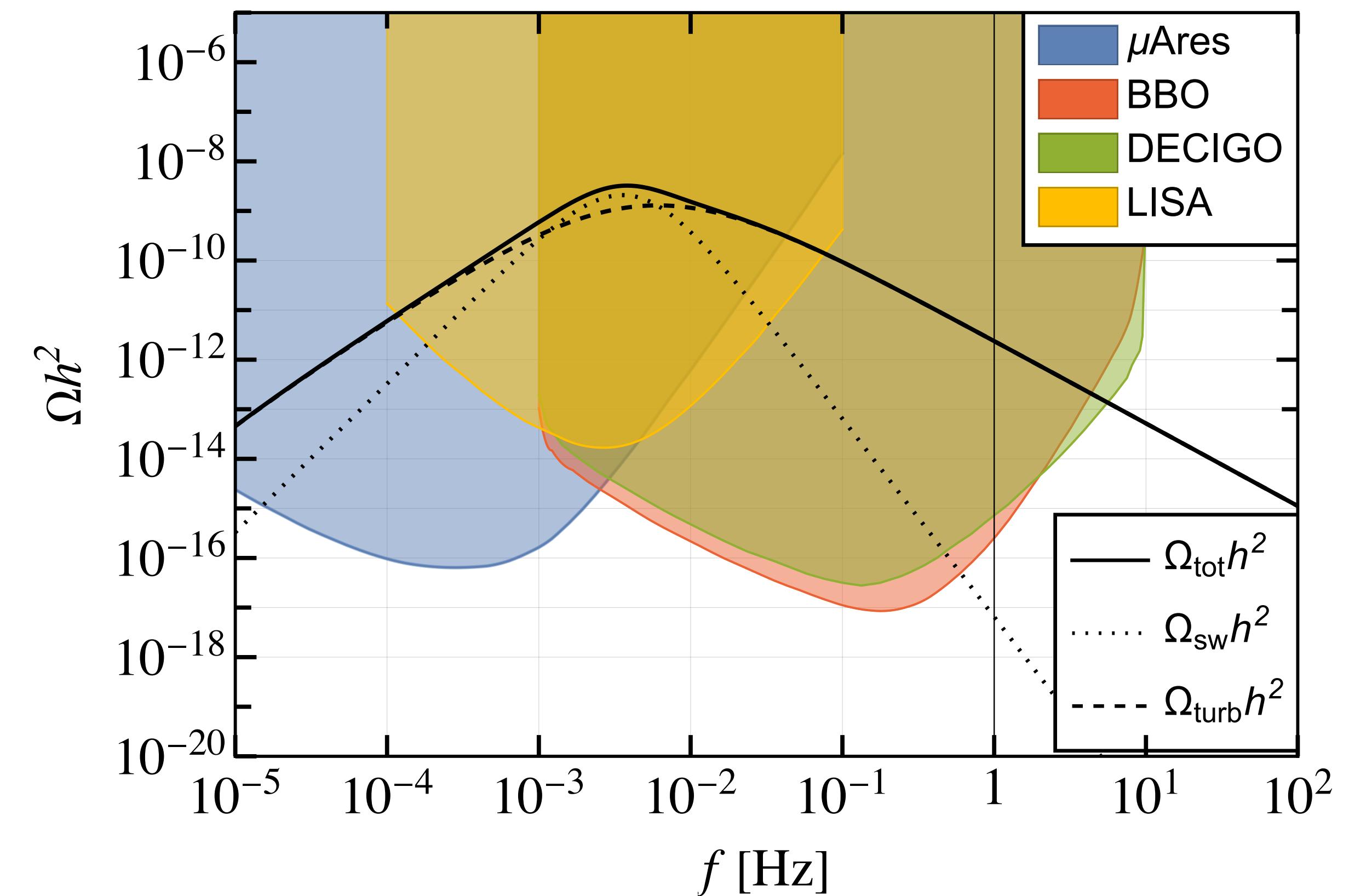
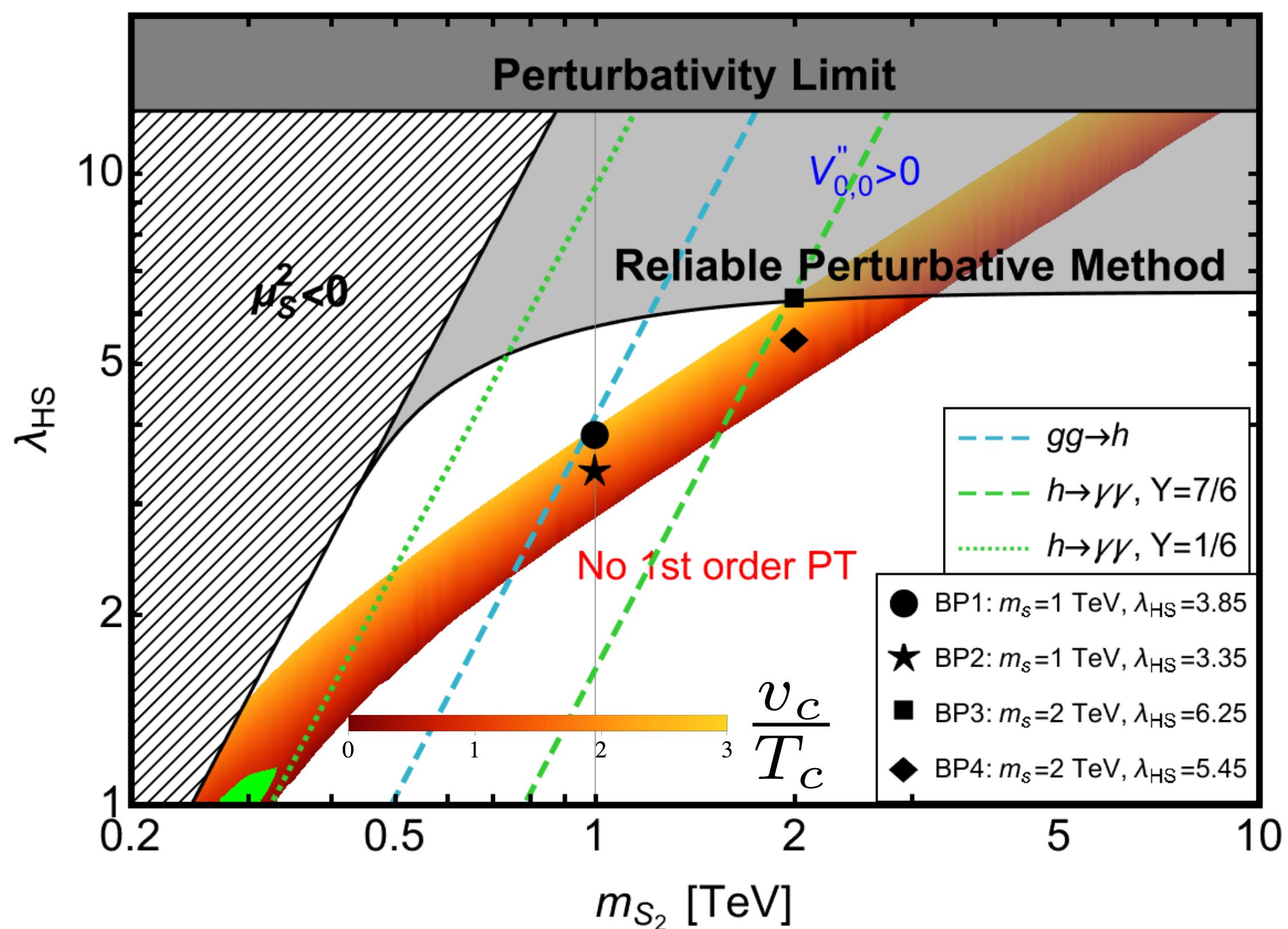
GW from leptoquark induced FOPT

$$V_0 = -\mu^2 |H|^2 + \lambda_H |H|^4 + \mu_S^2 |S_a|^2 + \lambda_S |S_a|^4 + 2\lambda_{HS} |H|^2 |S_a|^2$$

$$V_{\text{eff}}(h, T) = V_0 + \Delta V_0^{\text{1-loop}}(h) + \Delta V_T^{\text{1-loop}}(h, T)$$

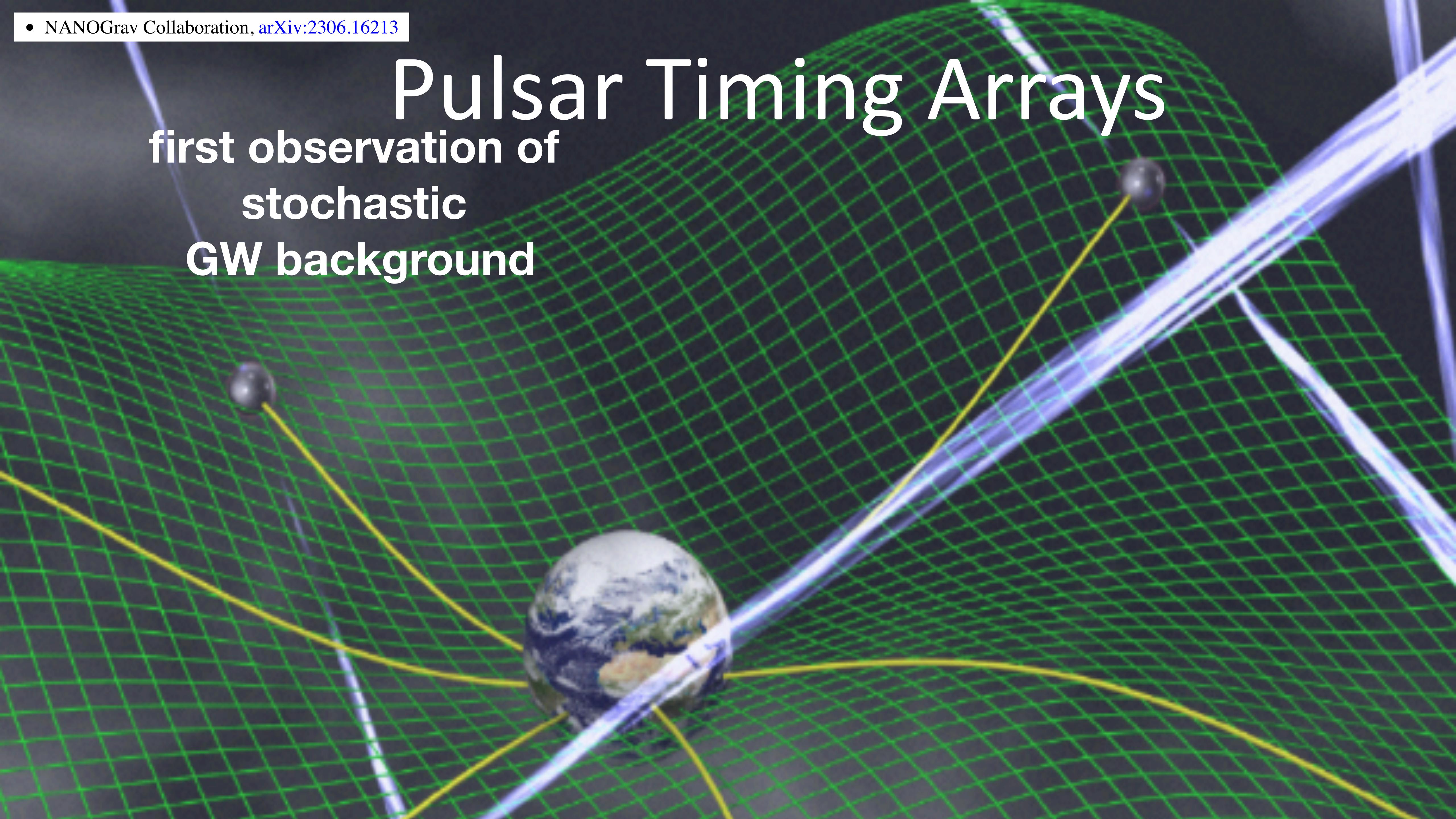
Leptoquark singlet S_1 , doublet S_2 or triplet S_3

Benchmark Point 1



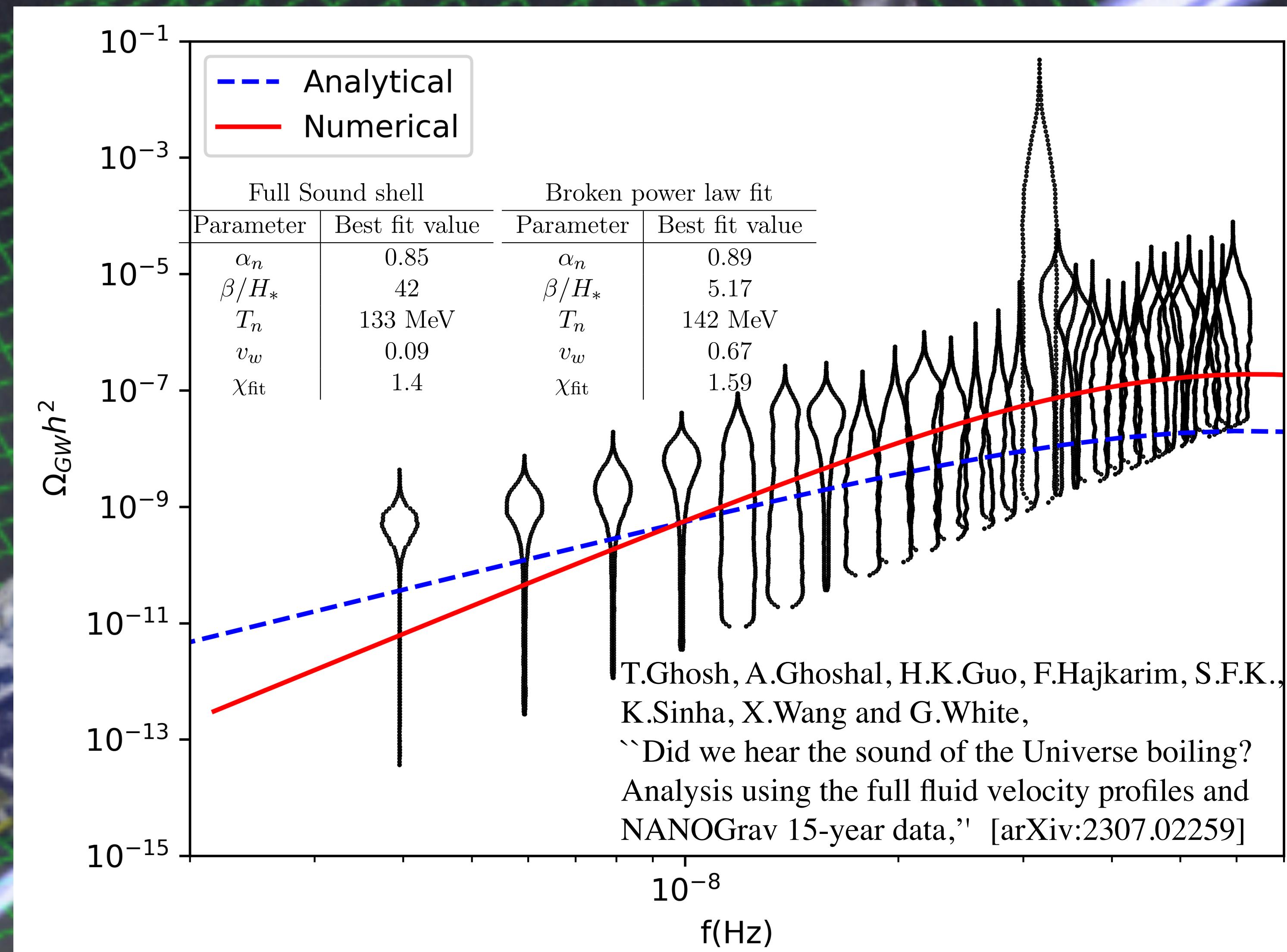
Pulsar Timing Arrays

**first observation of
stochastic
GW background**



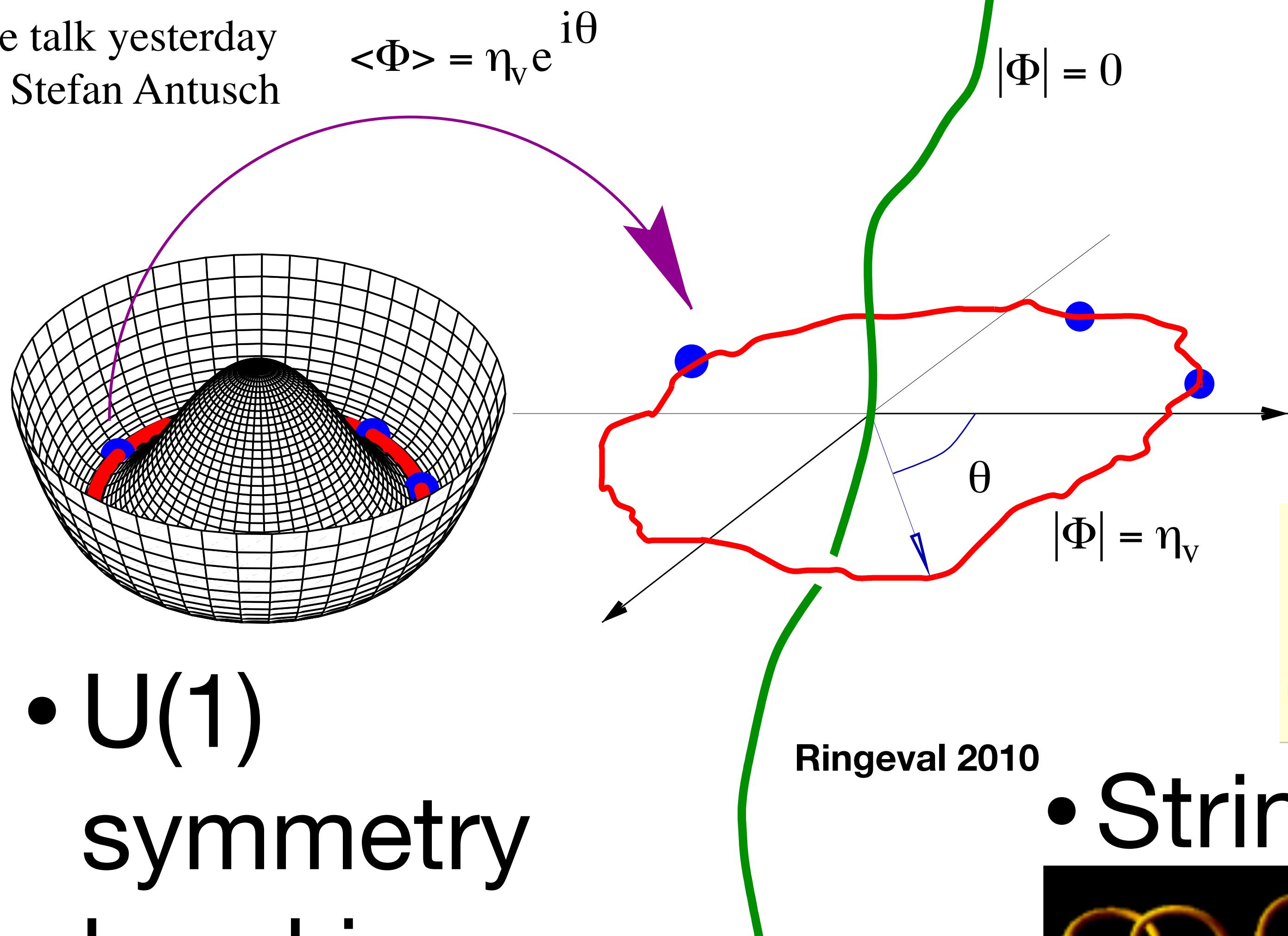
Pulsar Timing Arrays

first observation of stochastic GW background

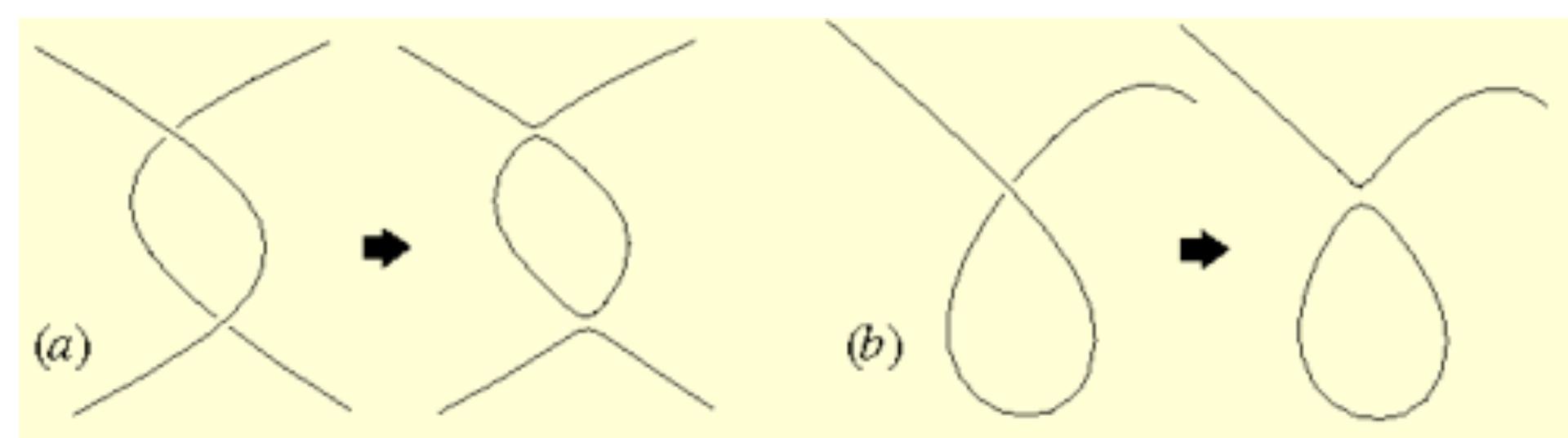
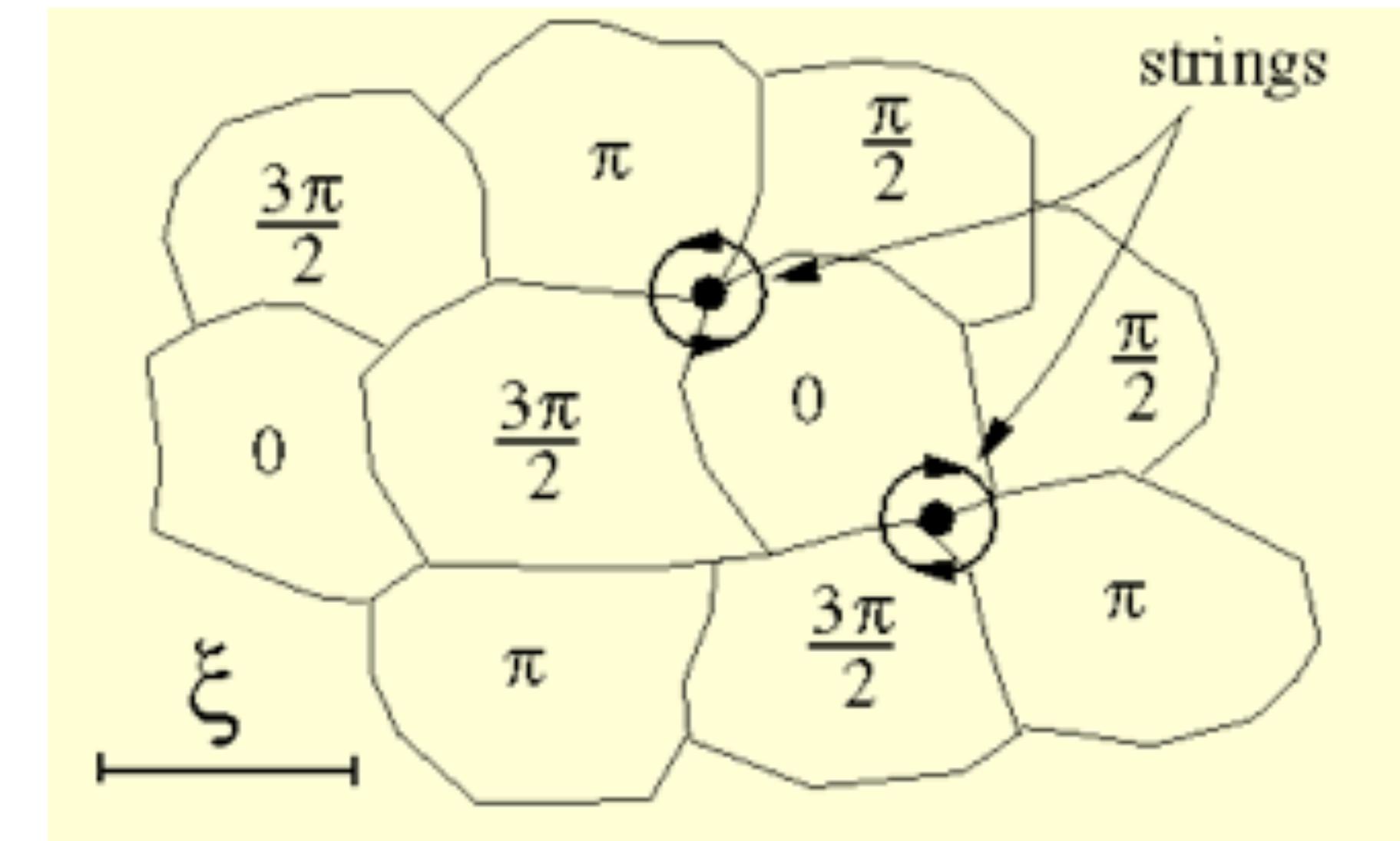


Gravitational Waves from Cosmic Strings

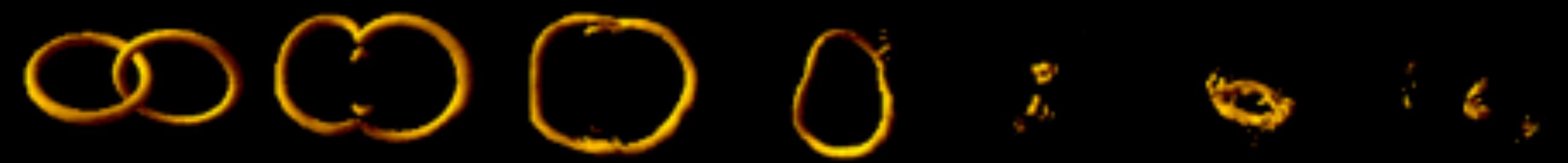
See talk yesterday
by Stefan Antusch



- U(1) symmetry breaking

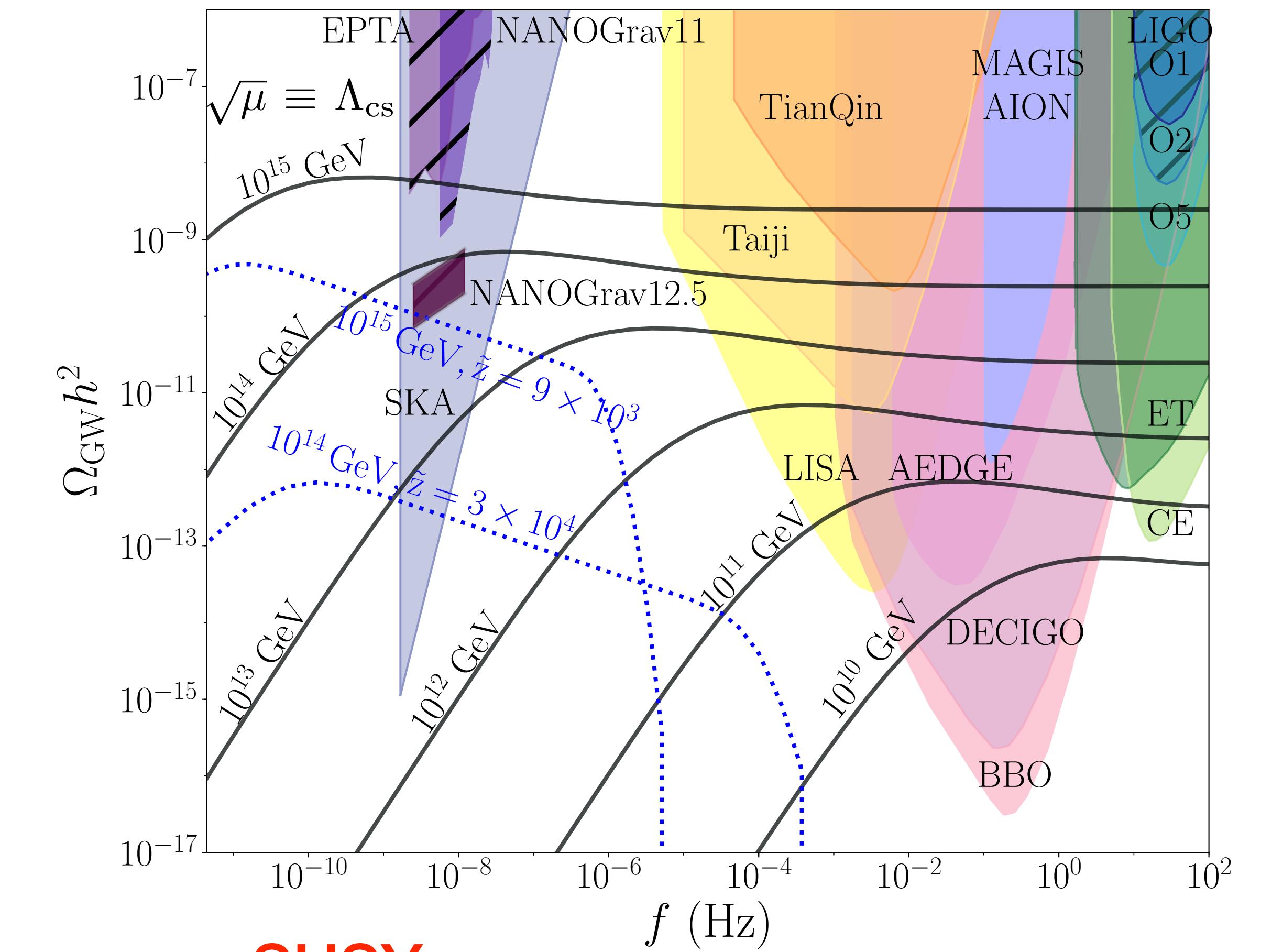
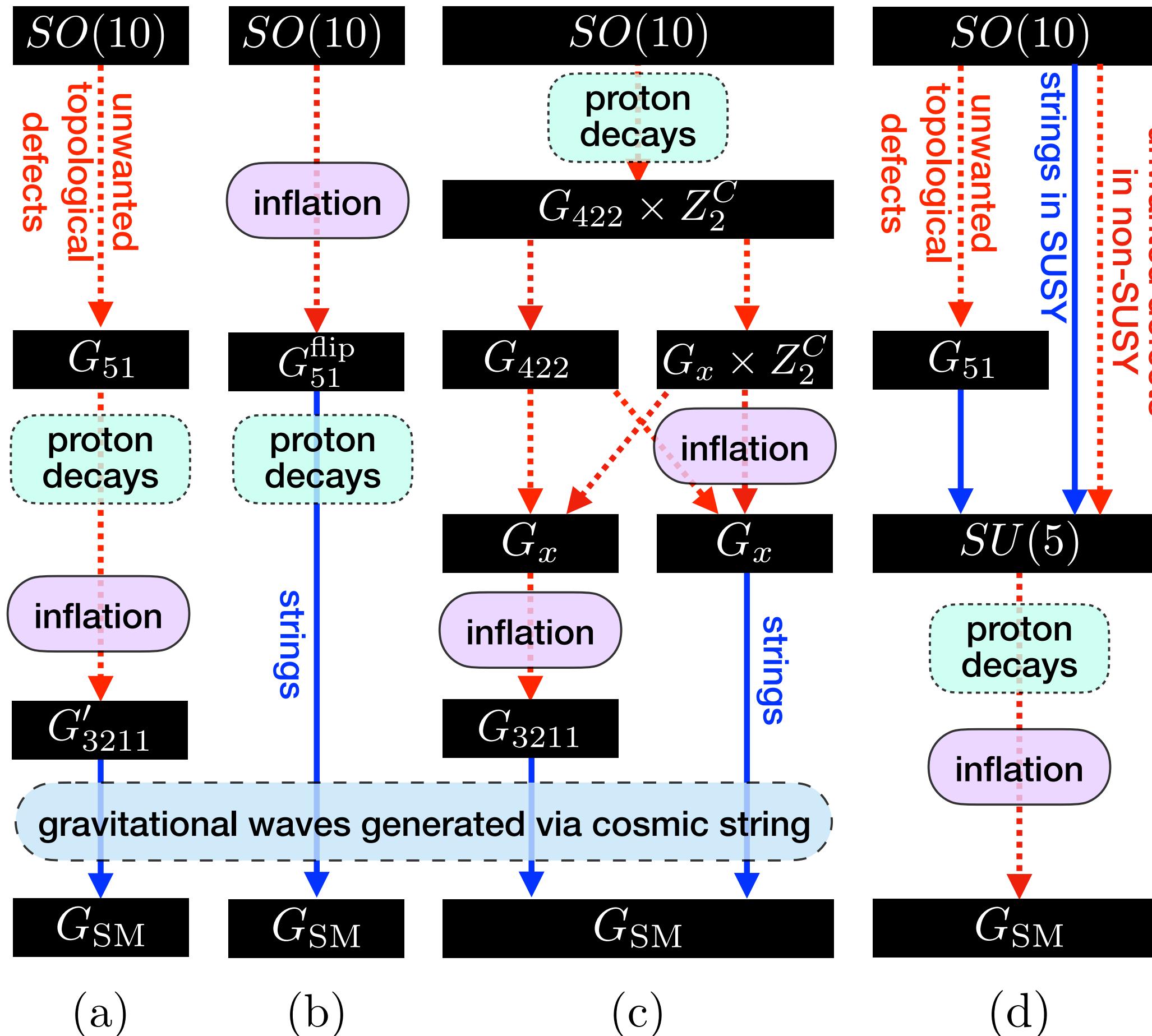


- String loops radiate GWs

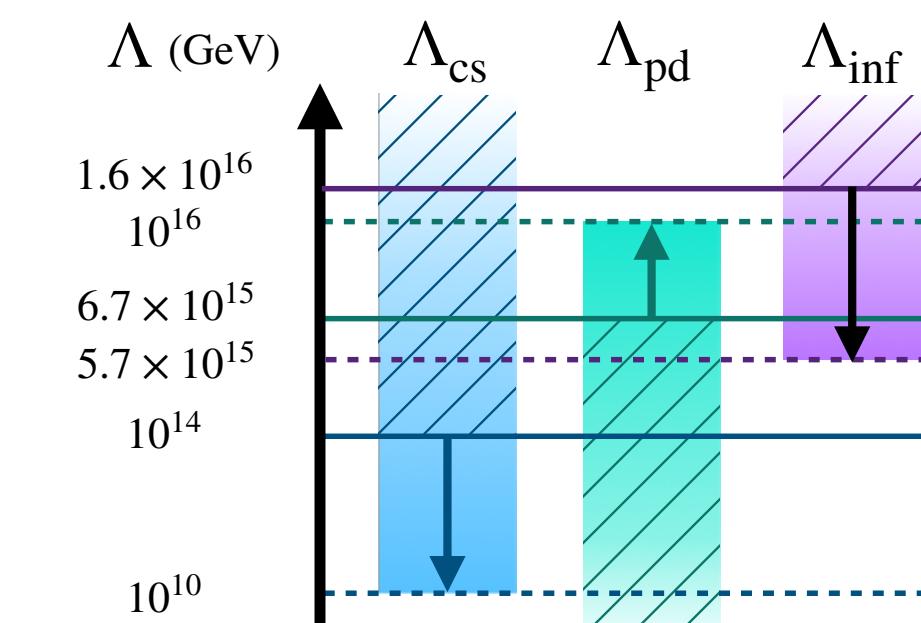


GWs via CSs from gauged $U(1)_{B-L}$ in $SO(10)$ GUTs

S.F.K., S.Pascoli, J.Turner and Y.L.Zhou, 2005.13549;
2106.15634; w/ Marsili 2209.00021; 2308.05799



non-SUSY



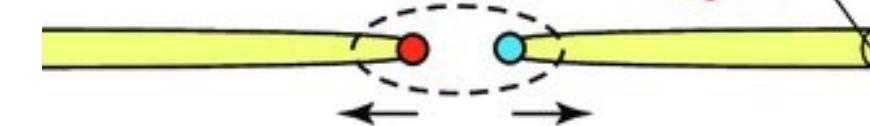
| Observables | Proton decays | |
|-------------|--|--|
| | $p \rightarrow \pi^0 e^+$ observed \Rightarrow non-SUSY contribution indicated | |
| GWs | Observed | <ul style="list-style-type: none"> types (a) and (c) favoured types (b) and (d) excluded |
| | Marginal | <ul style="list-style-type: none"> types (a) and (c) favoured type (d) excluded type (b) allowed if $p \rightarrow K^+ \bar{\nu}$ not observed and $\Lambda_{\text{pd}} \sim \Lambda_{\text{cs}}$ |

Flipped SU(5): unification, proton decay, fermion masses and gravitational waves

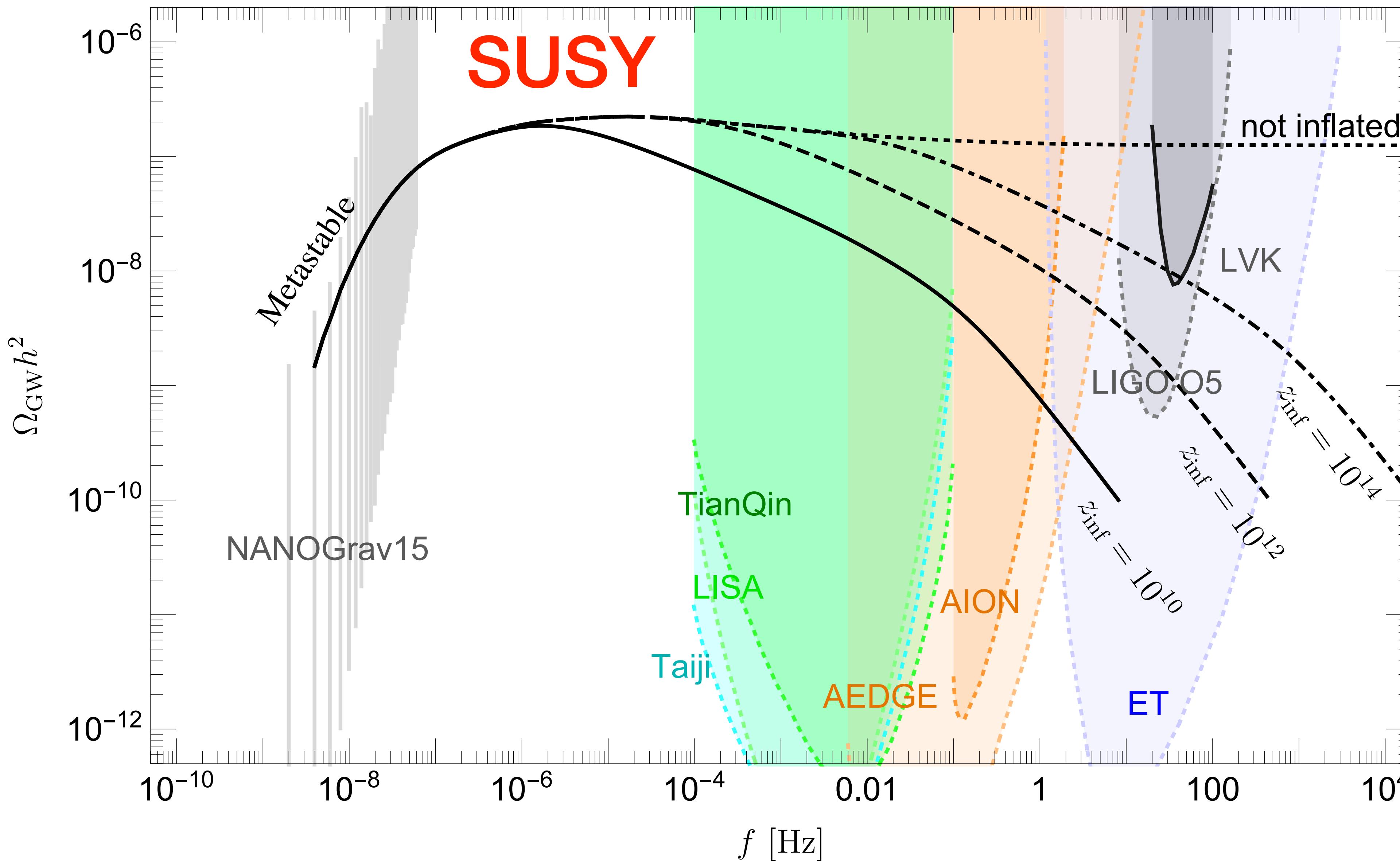
S.F.K., G.K.Leontaris and Y.L.Zhou, 2311.11857

See talk yesterday
by Stefan Antusch

Metastable
cosmic strings



**SUSY GUTs
with
metastable
cosmic strings
and inflation
dilution can fit
NANOGrav 15
year data**

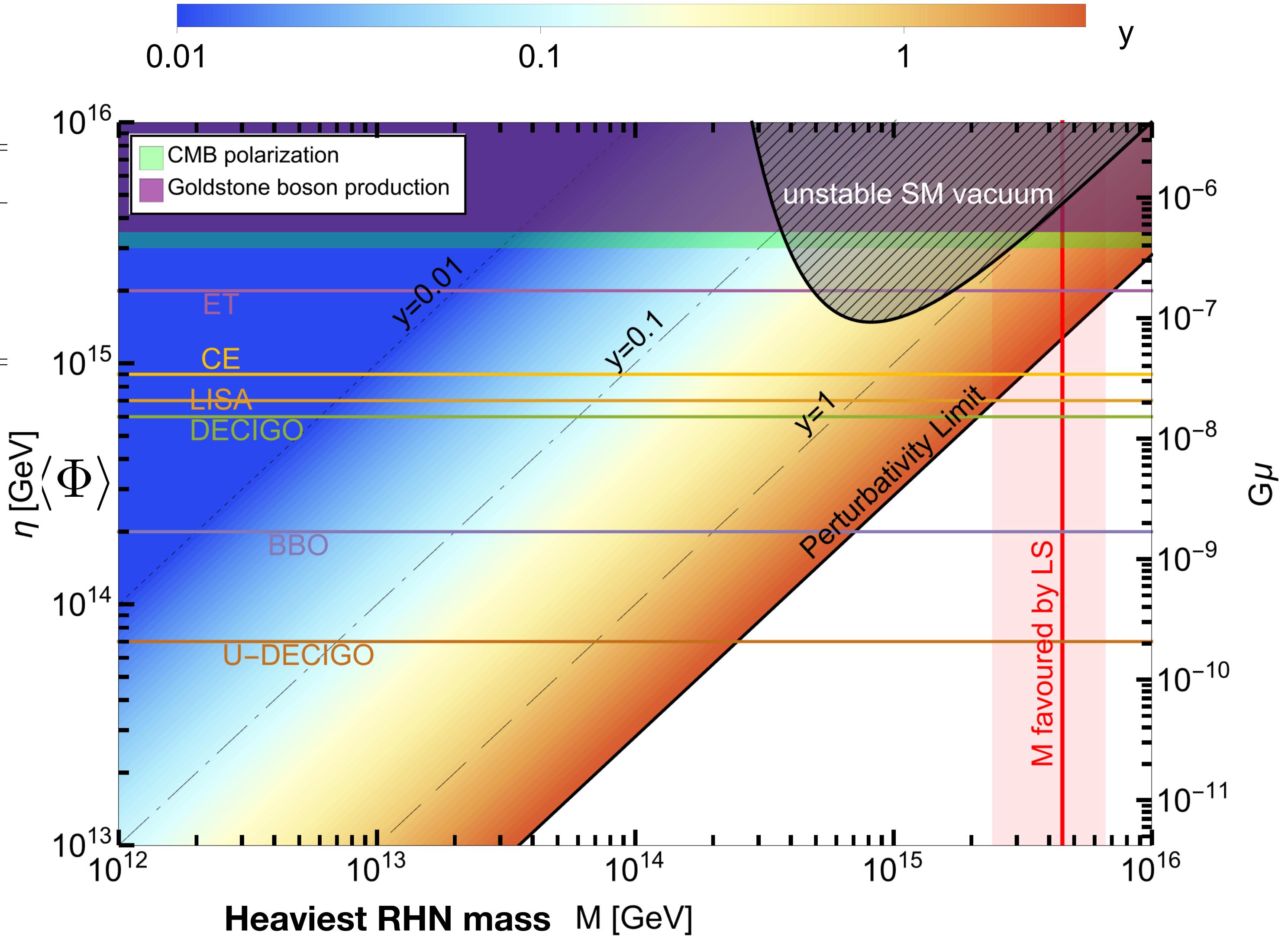
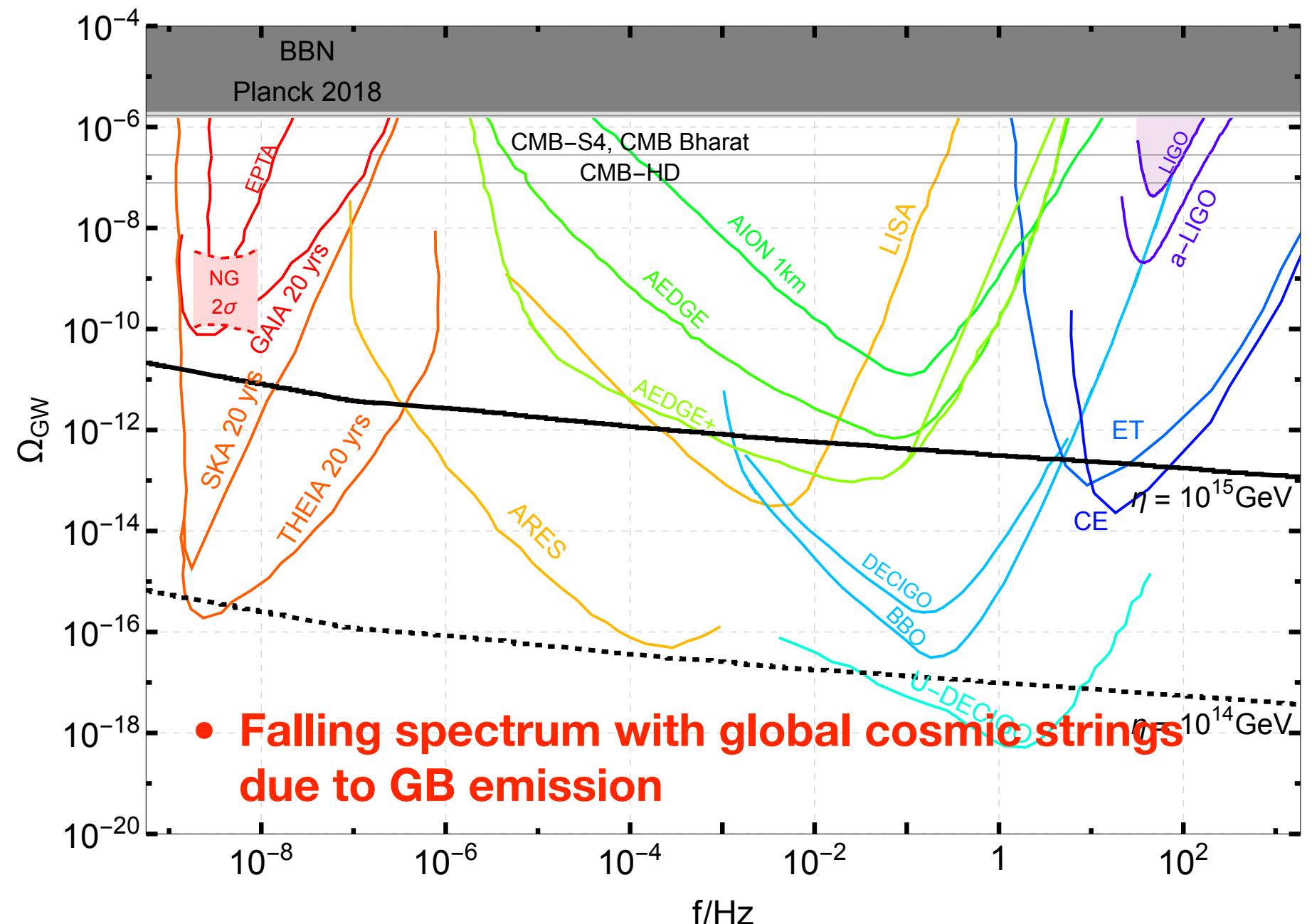


Cosmic string gravitational waves from **global** U(1)_{B-L} symmetry breaking as a probe of the type I seesaw scale

Majoron breaks U(1)_{B-L}

$$Y_{\alpha i} \overline{L}_\alpha \tilde{H} N_i + \frac{1}{2} y_i \Phi \overline{N}_i^c N_i + h.c.$$

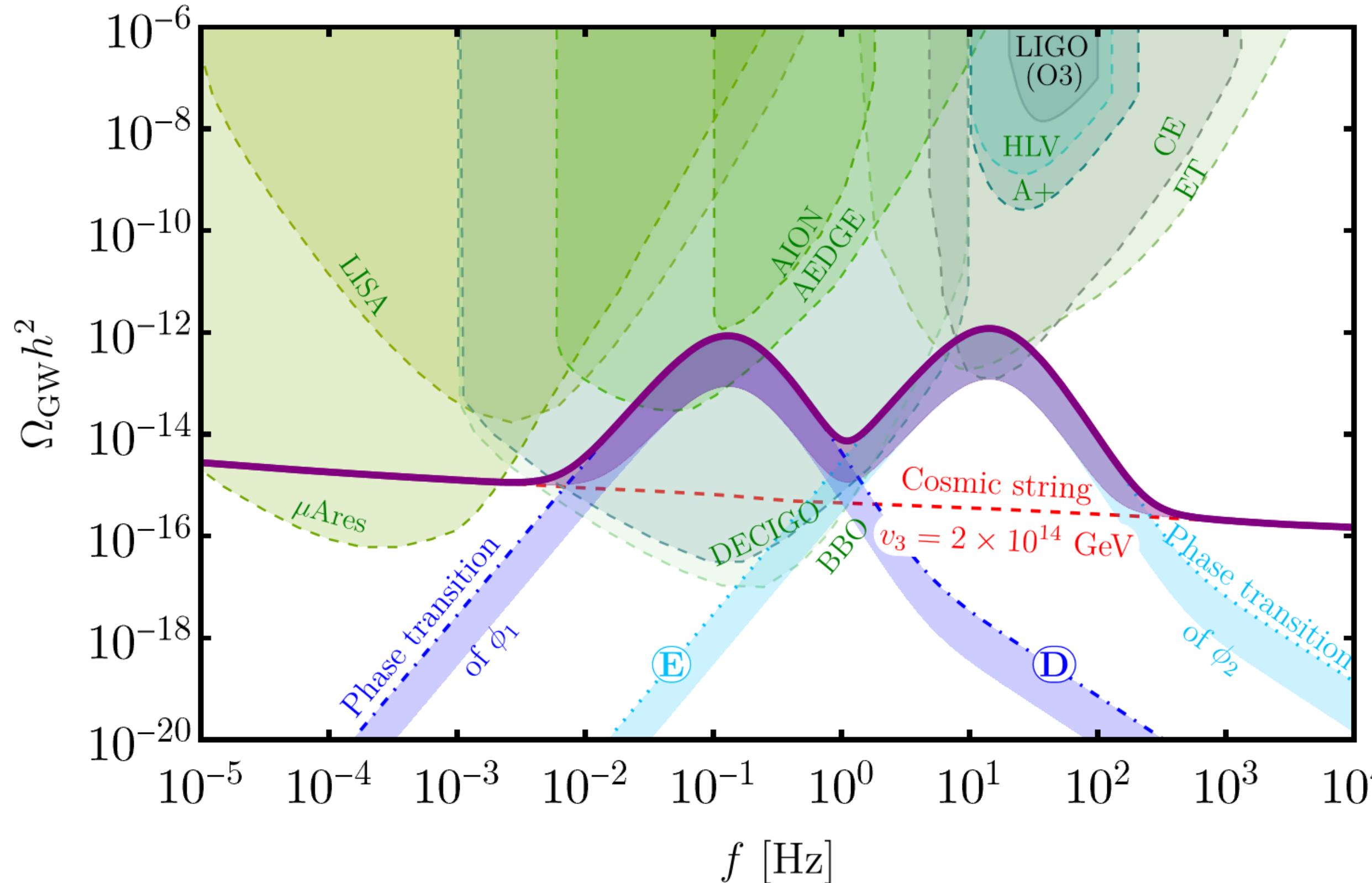
| | Q | u_R | d_R | L | e_R | H | N | Φ |
|--------------|---------------|---------------|----------------|----------------|-------|----------------|-----|--------|
| $SU(2)_L$ | 2 | 1 | 1 | 2 | 1 | 2 | 1 | 1 |
| $U(1)_Y$ | $\frac{1}{6}$ | $\frac{2}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{2}$ | -1 | $-\frac{1}{2}$ | 0 | 0 |
| $U(1)_{B-L}$ | $\frac{1}{3}$ | $\frac{1}{3}$ | $\frac{1}{3}$ | -1 | -1 | 0 | -1 | 2 |



Gravitational waves from phase transitions and cosmic strings in neutrino mass models with multiple Majorons

$$U(1)_{L_1} \times U(1)_{L_2} \times U(1)_{L_3}$$

$$\left(\overline{L}_a h_{aI} H N_I + \frac{y_1}{2} \phi_1 \overline{N_1^c} N_1 + \frac{y_2}{2} \phi_2 \overline{N_2^c} N_2 + \frac{y_3}{2} \phi_3 \overline{N_3^c} N_3 + \text{h.c.} \right) + V_0(\phi_1, \phi_2, \phi_3)$$



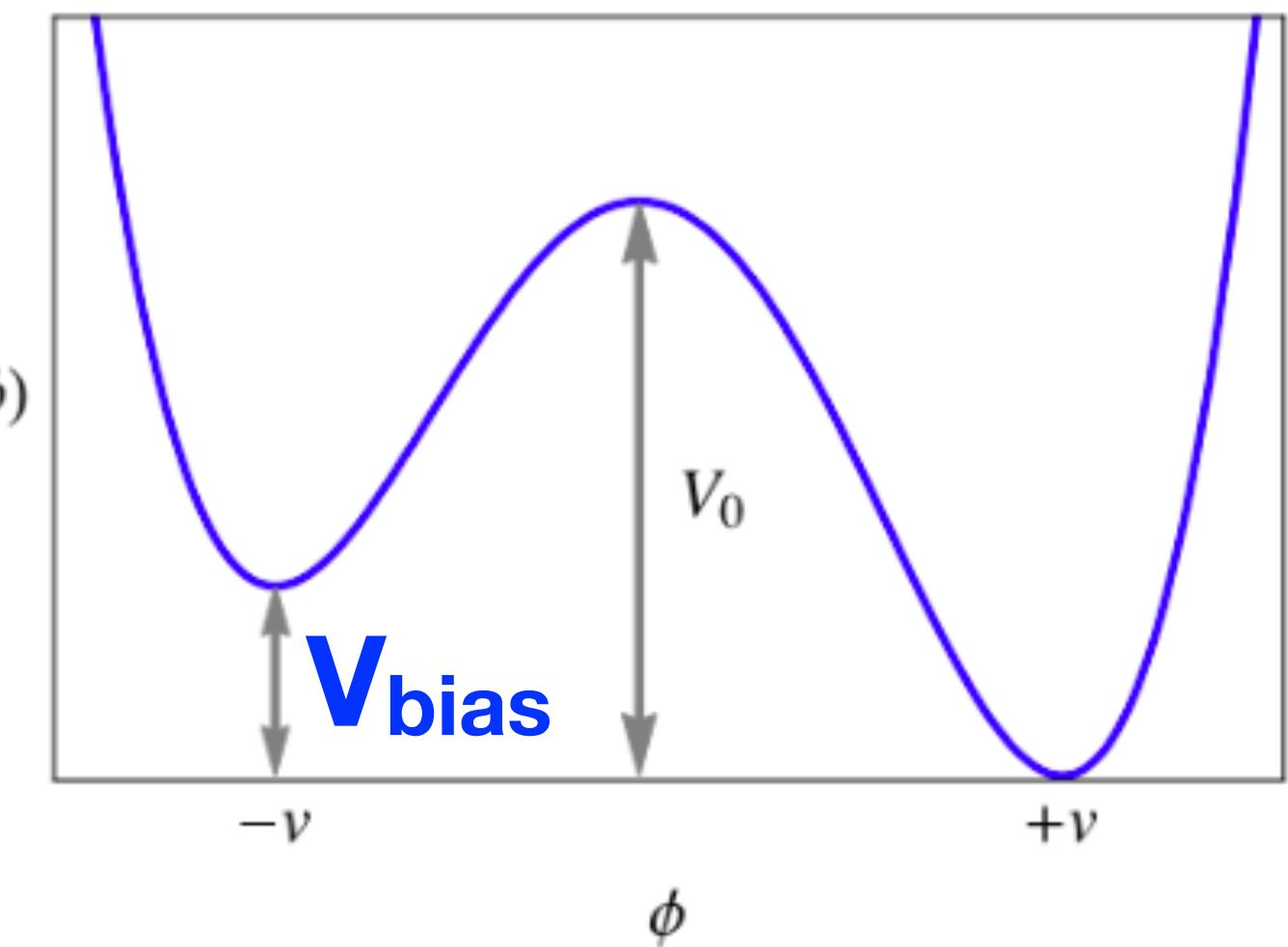
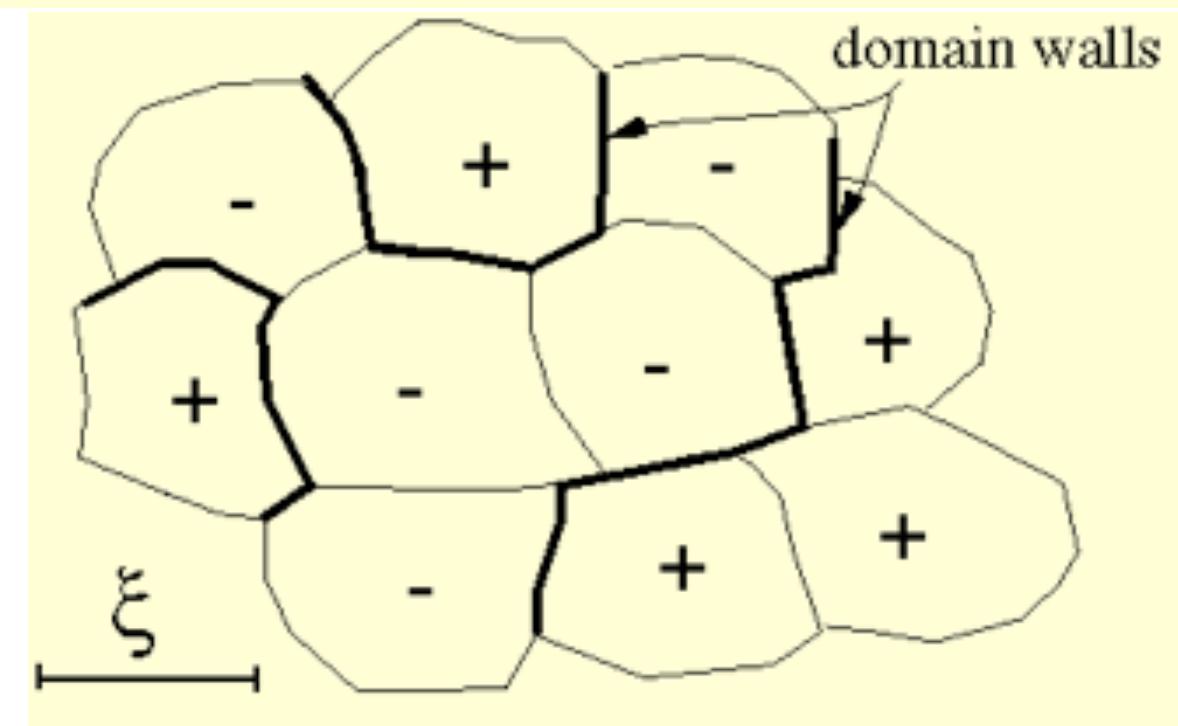
$$\sum_{I=1,2,3} [-\mu_I^2 \phi_I^* \phi_I + \lambda_I (\phi_I^* \phi_I)^2] + \sum_{I,J,I \neq J}^{1,2,3} \frac{\zeta_{IJ}}{2} (\phi_I^* \phi_I)(\phi_J^* \phi_J)$$

| | λ_I | v_I [GeV] |
|-----|-------------|--------------------|
| (D) | 0.00027 | 1188.2 |
| (E) | 0.00029 | 2.32×10^5 |

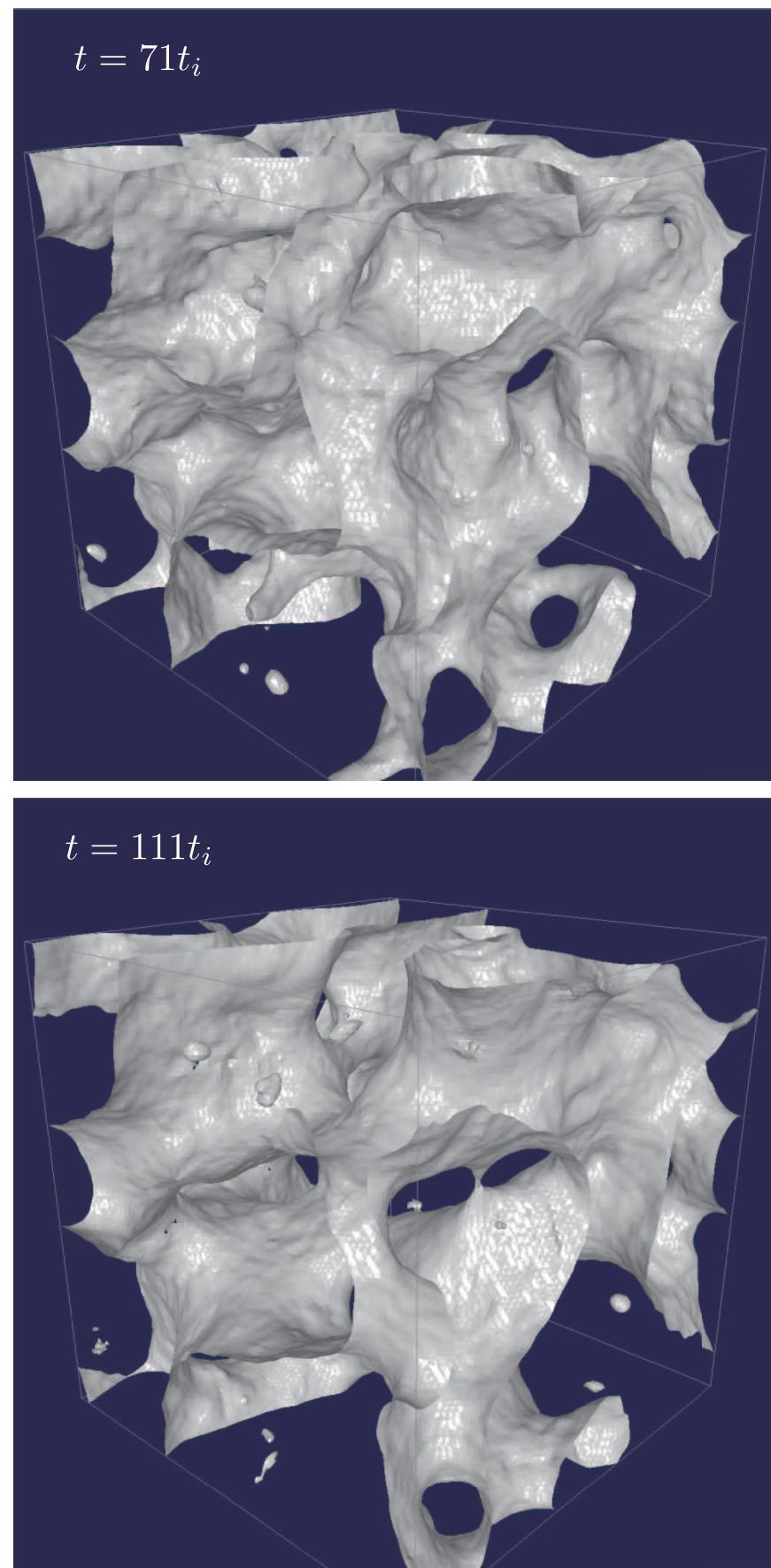
Phase transitions for Φ_1, Φ_2 enhanced by Φ_3
 See talk yesterday by Pasquale Di Bari

Gravitational Waves from Domain Walls

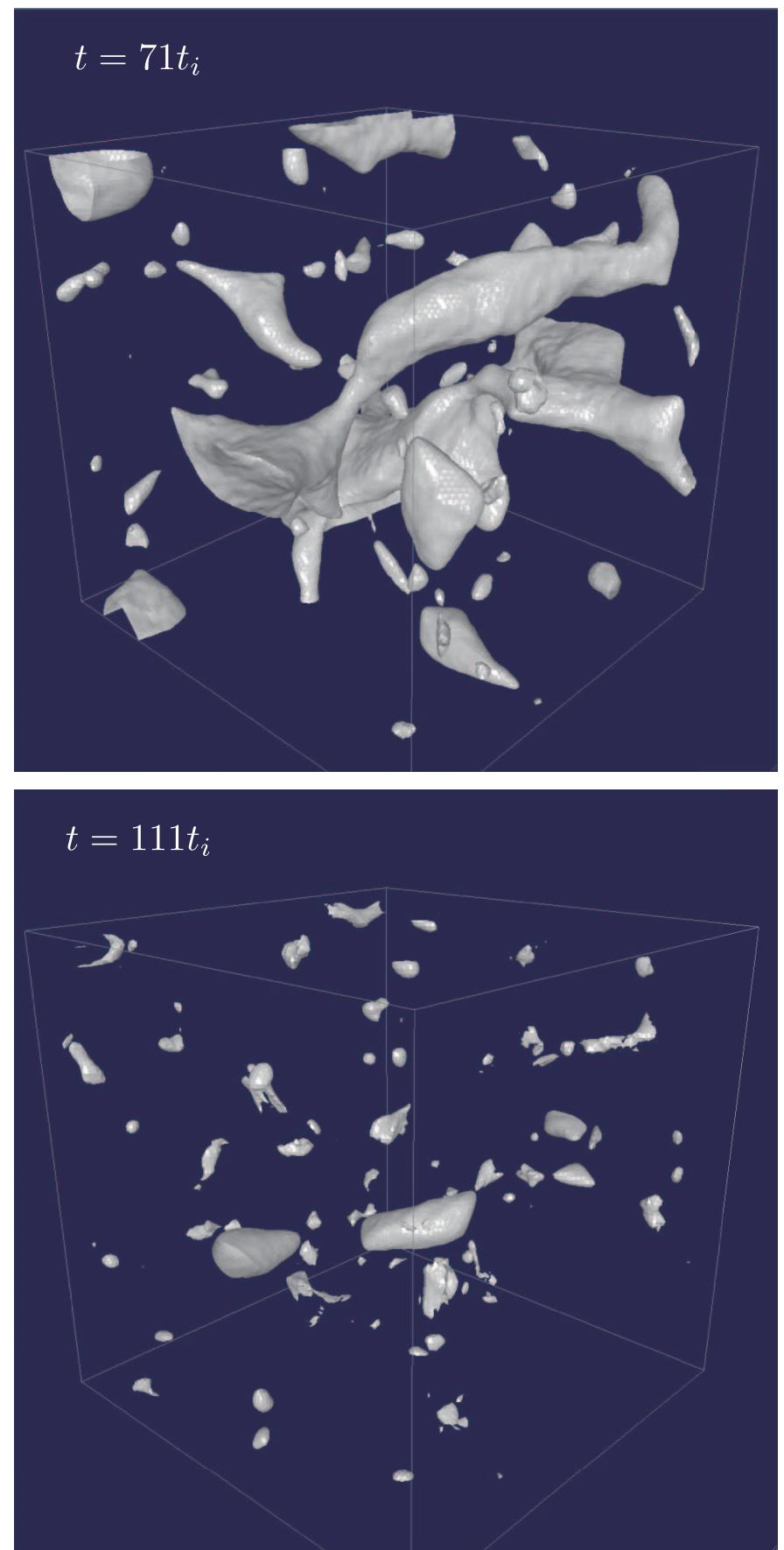
- Z_2 symmetry breaking



Introduce V_{bias} to allow domain walls to decay (otherwise dominate energy density of Universe)



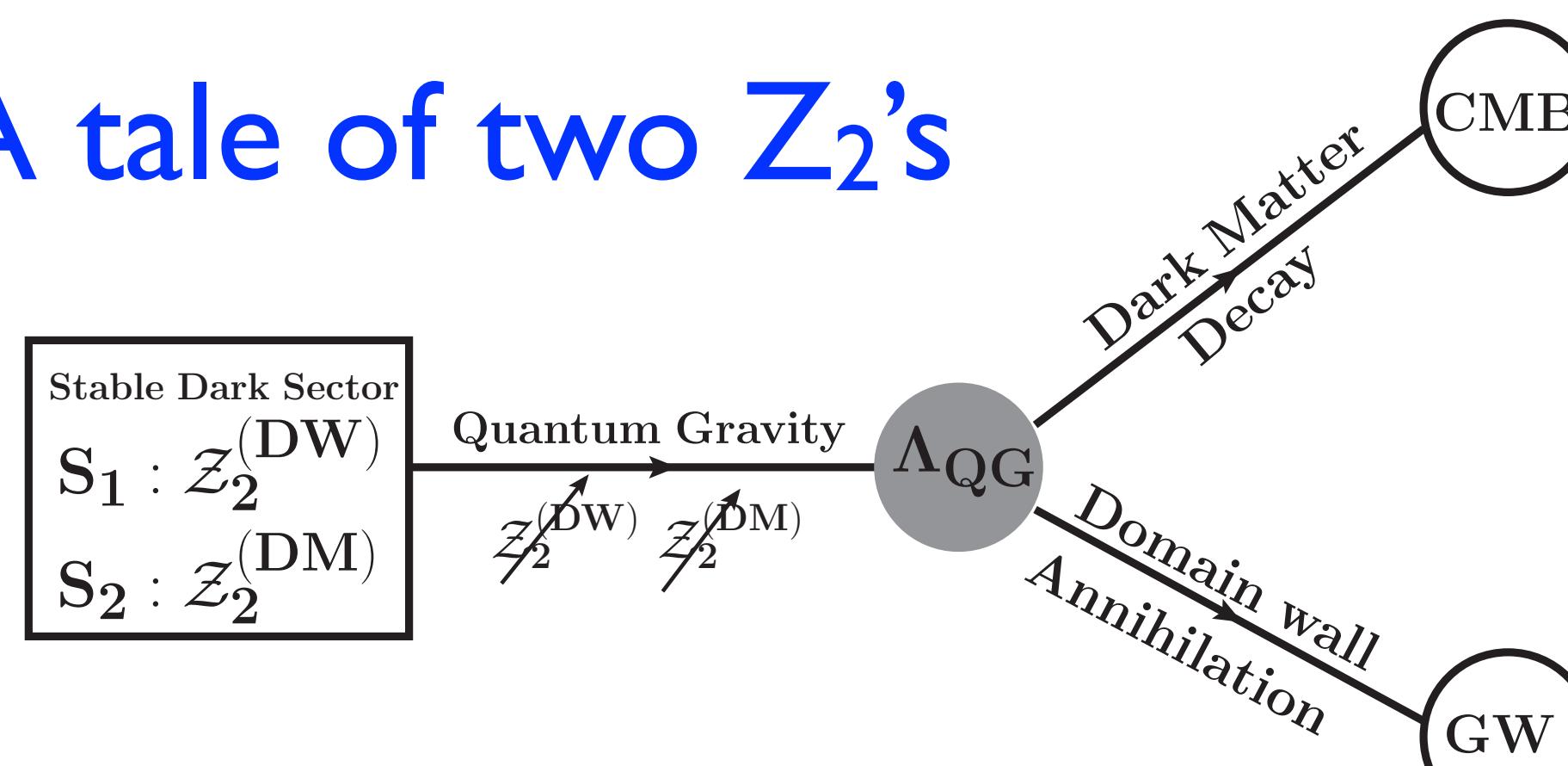
Stable DW on the left with $V_{\text{bias}} = 0$



Unstable DW decay via GWs due to V_{bias}

Quantum gravity effects on dark matter and gravitational waves

A tale of two Z_2 's



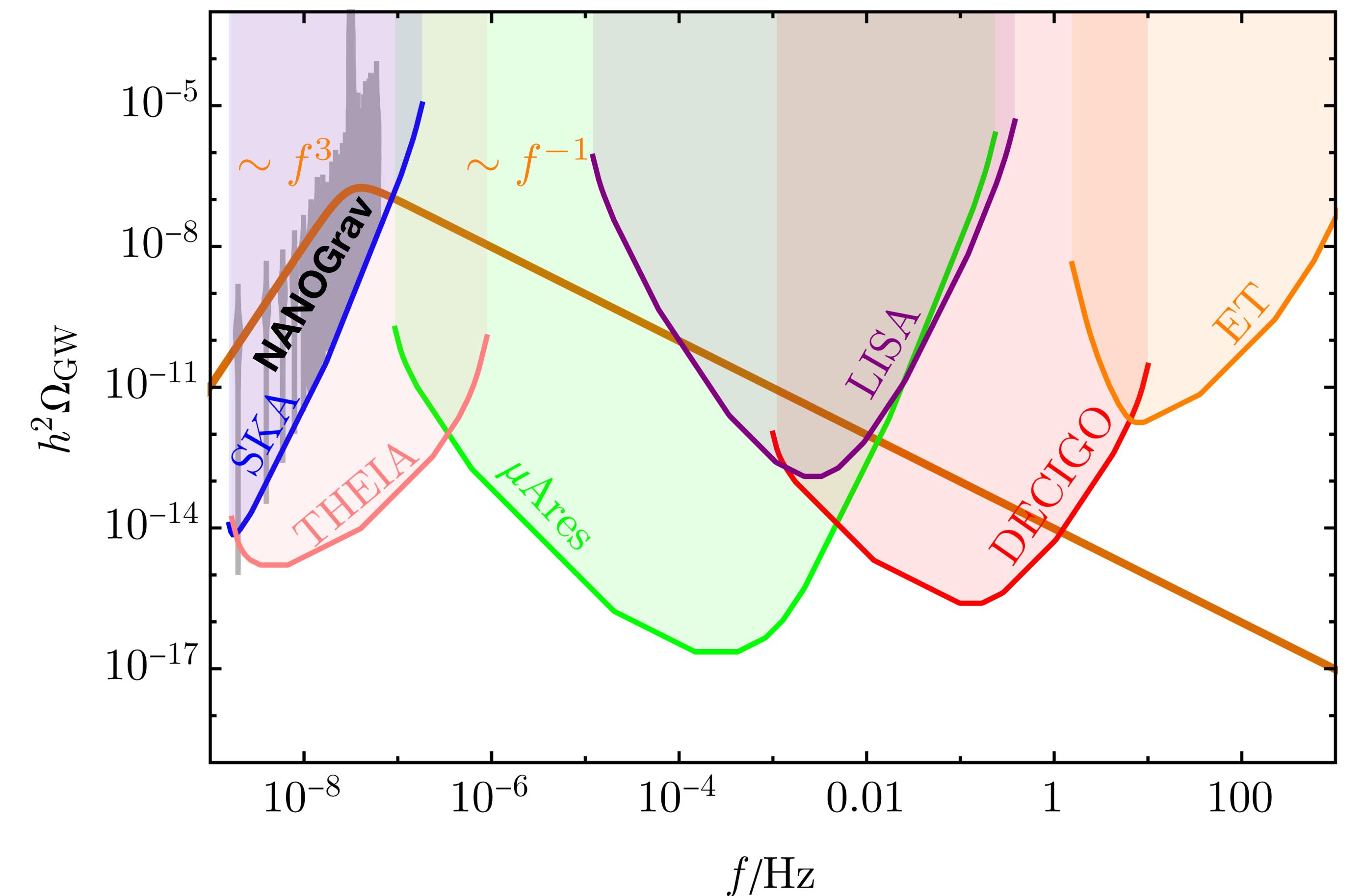
Both broken by QG effects

$$\mathcal{L}_{Z_2} = \frac{1}{\Lambda_{\text{QG}}} \mathcal{O}_5 \rightarrow V_{\text{bias}}$$

Due to instanton effects

$$\Lambda_{\text{QG}} \sim M_{\text{Pl}} e^{\mathcal{S}} \gg M_{\text{Pl}}$$

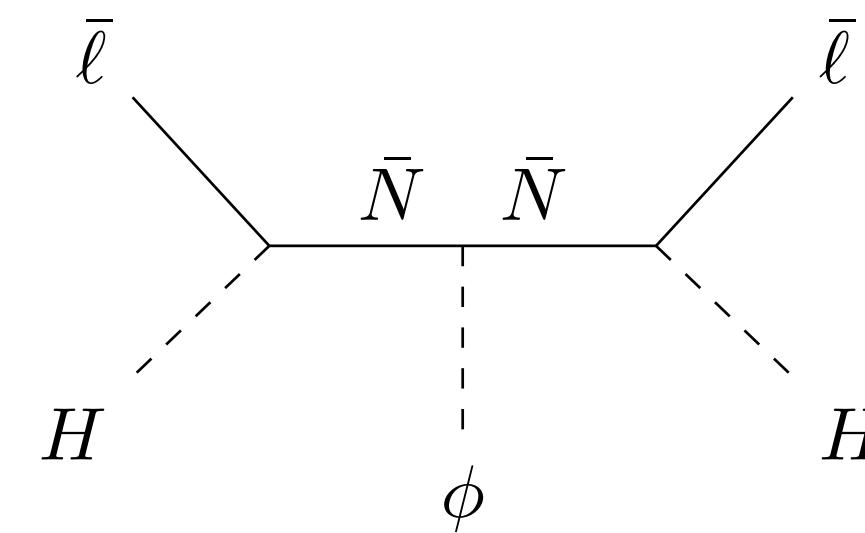
Peak occurs when volume pressure $\sim V_{\text{bias}}$



Toward distinguishing Dirac from Majorana neutrino mass with gravitational waves

Majorana seesaw

$$-\mathcal{L}_M \supset \mathcal{Y} \bar{\ell} H \bar{N} + \bar{N} \bar{N}^T \phi$$

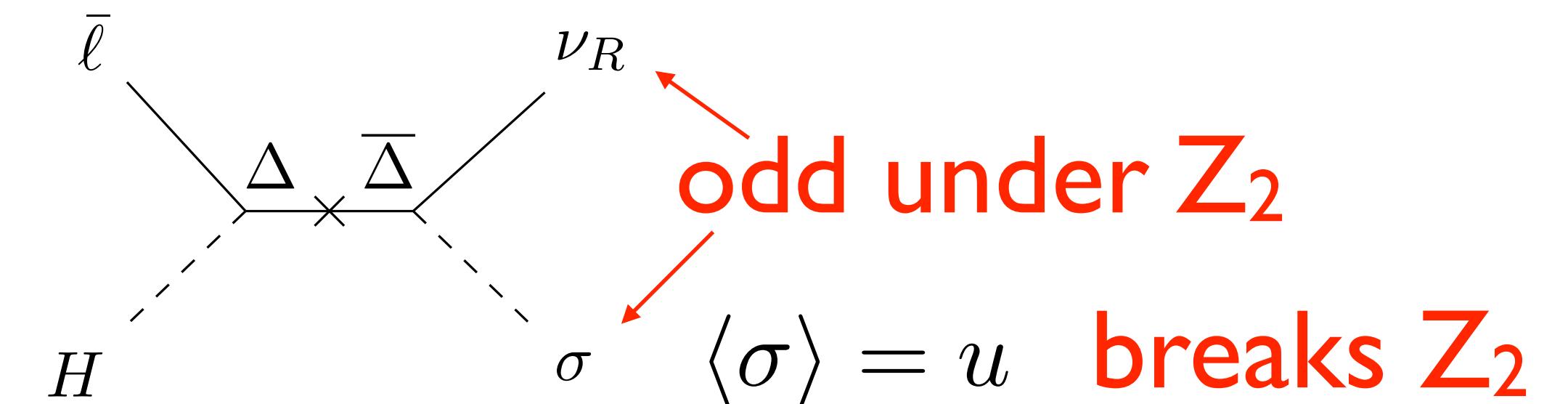


$$\mathcal{M}_M = \frac{1}{\sqrt{2}} v^2 \mathcal{Y} \mathcal{M}_N^{-1} \mathcal{Y}^T$$

Gauged U(1)_{B-L} broken
→ Cosmic strings

Dirac seesaw

$$-\mathcal{L}_D \supset \mathcal{Y}_L \bar{\ell} H \Delta_R + \mathcal{Y}_R \bar{\Delta}_L \sigma \nu_R + \mathcal{M}_\Delta \bar{\Delta} \Delta$$

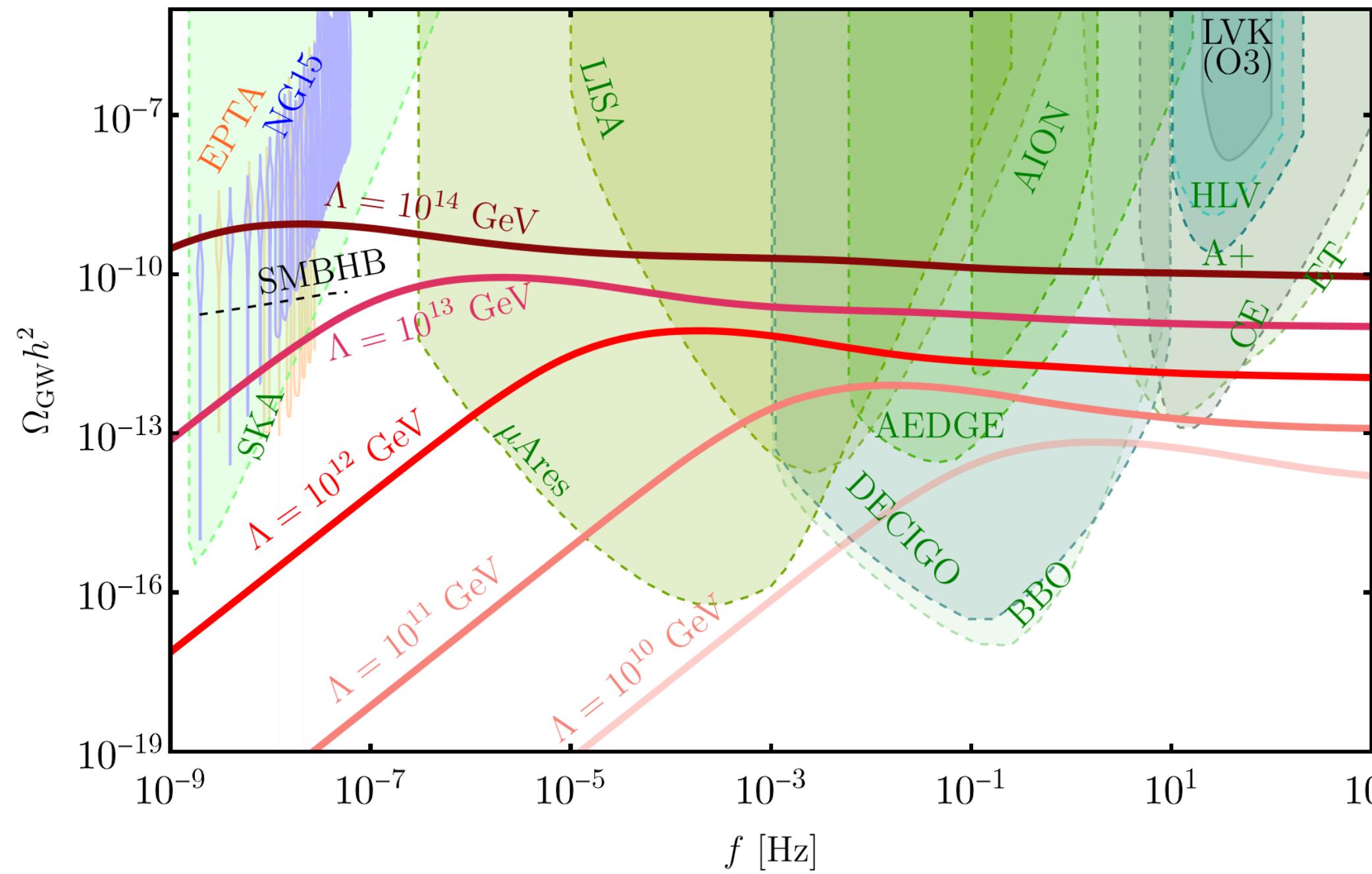


$$\mathcal{M}_D = \frac{1}{\sqrt{2}} v u \mathcal{Y}_L \mathcal{M}_\Delta^{-1} \mathcal{Y}_R$$

Gauged U(1)_{B-L} preserved
 Z_2 broken → Domain Walls

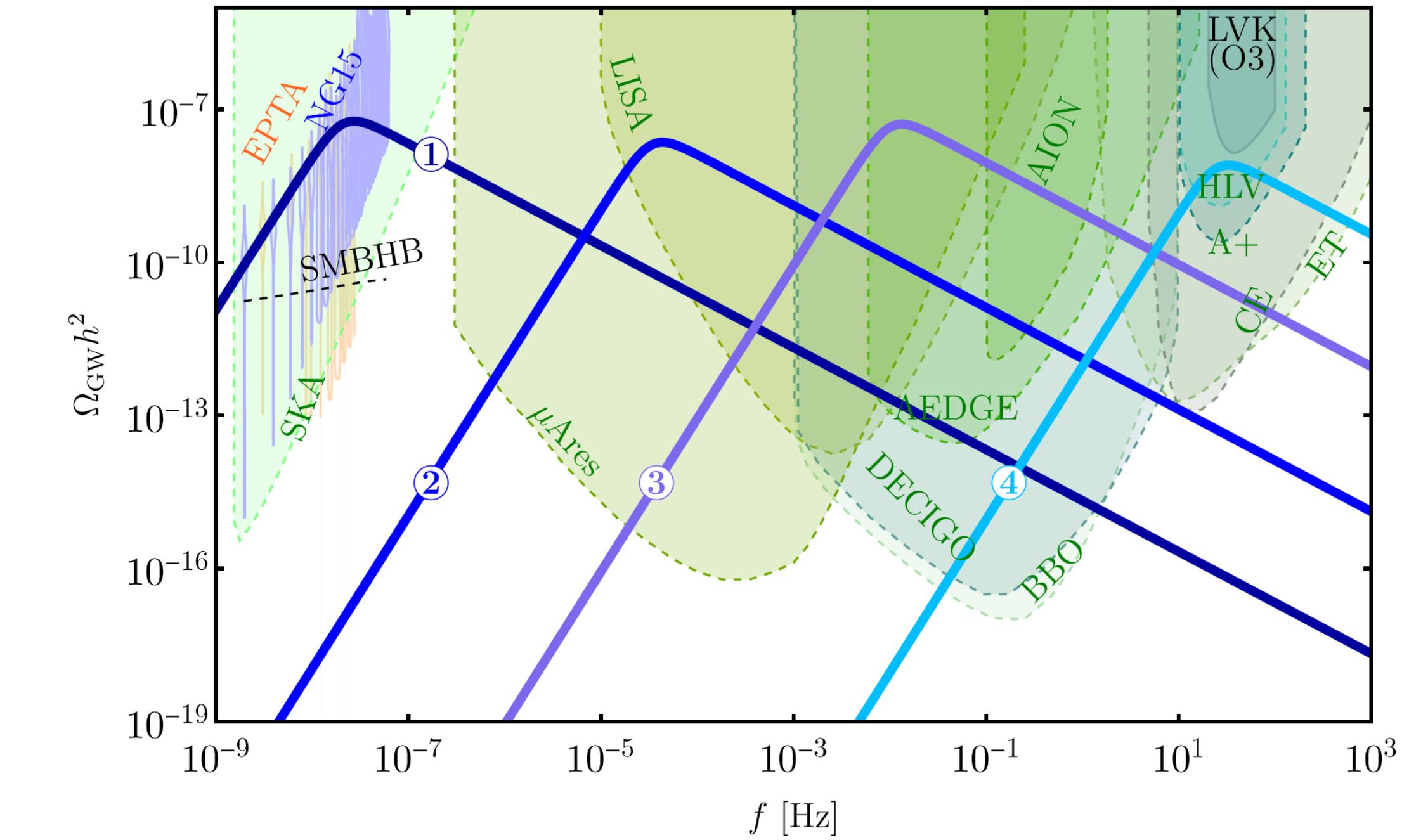
Toward distinguishing Dirac from Majorana neutrino mass with gravitational waves

Majorana seesaw



Majorana vs **Dirac** can be distinguished
from shape of GW spectrum
- - Dirac is better fit to NANOGrav

Dirac seesaw



| Benchmark Point | u [GeV] | V_{bias} [GeV^4] | $y_{\text{max}}(M_\Delta < M_{\text{Pl}})$ |
|-----------------|--------------------|--------------------------------------|--|
| ① | 10^5 | 10^{-5} | 4.93 |
| ② | 5.2×10^7 | 7.14×10^{10} | 0.216 |
| ③ | 1.2×10^9 | 10^{19} | 0.045 |
| ④ | 2×10^{11} | 2.5×10^{32} | 0.0035 |

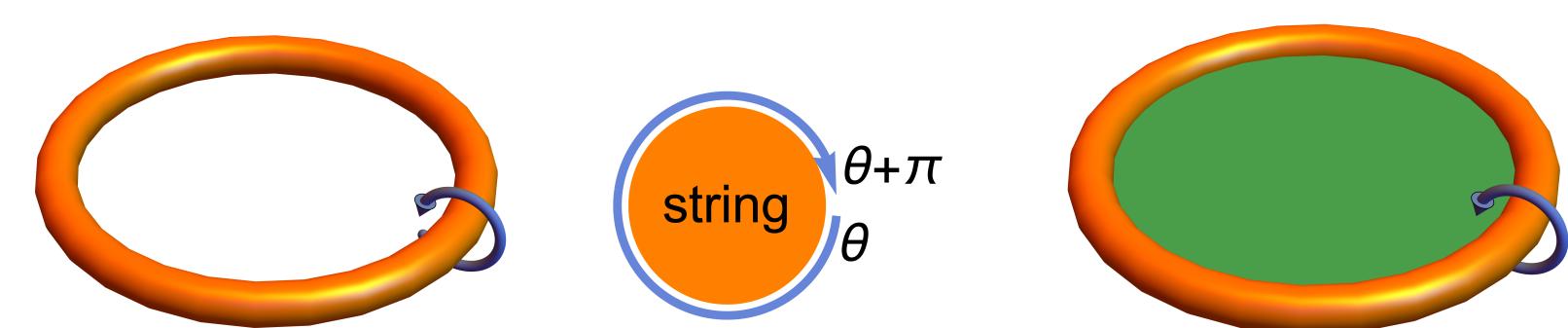
$$V(\sigma) = \frac{\lambda}{4}(\sigma^2 - u^2)^2 \quad \Delta V(\sigma) = \epsilon u \sigma \left(\frac{\sigma^2}{3} - u^2 \right)$$

Type-I two-Higgs-doublet model and gravitational waves from domain walls bounded by strings

$$Y_u \overline{Q} \tilde{\Phi}_2 u_R + Y_d \overline{Q} \Phi_2 d_R + Y_e \overline{L} \Phi_2 e_R + Y_N \overline{L} \tilde{\Phi}_2 N_R + y_N \phi N_R N_R$$

| | $u_{R\beta}$ | $d_{R\beta}$ | Q_α | L_α | $e_{R\beta}$ | $N_{R\beta}$ | Φ_2 | Φ_1 | ϕ |
|----------------|---------------|----------------|---------------|----------------|--------------|--------------|---------------|---------------|--------|
| $SU(2)_L$ | 1 | 1 | 2 | 2 | 1 | 1 | 2 | 2 | 1 |
| $U(1)_Y$ | $\frac{2}{3}$ | $-\frac{1}{3}$ | $\frac{1}{6}$ | $-\frac{1}{2}$ | -1 | 0 | $\frac{1}{2}$ | $\frac{1}{2}$ | 0 |
| $U(1)_R$ | 1 | -1 | 0 | 0 | -1 | 1 | 1 | -1 | -2 |
| residual Z_2 | - | - | + | + | - | - | - | - | + |

$$U(1)_R \xrightarrow{\langle \phi \rangle} Z_2 \xrightarrow{\langle \Phi_i \rangle} \text{nothing}$$

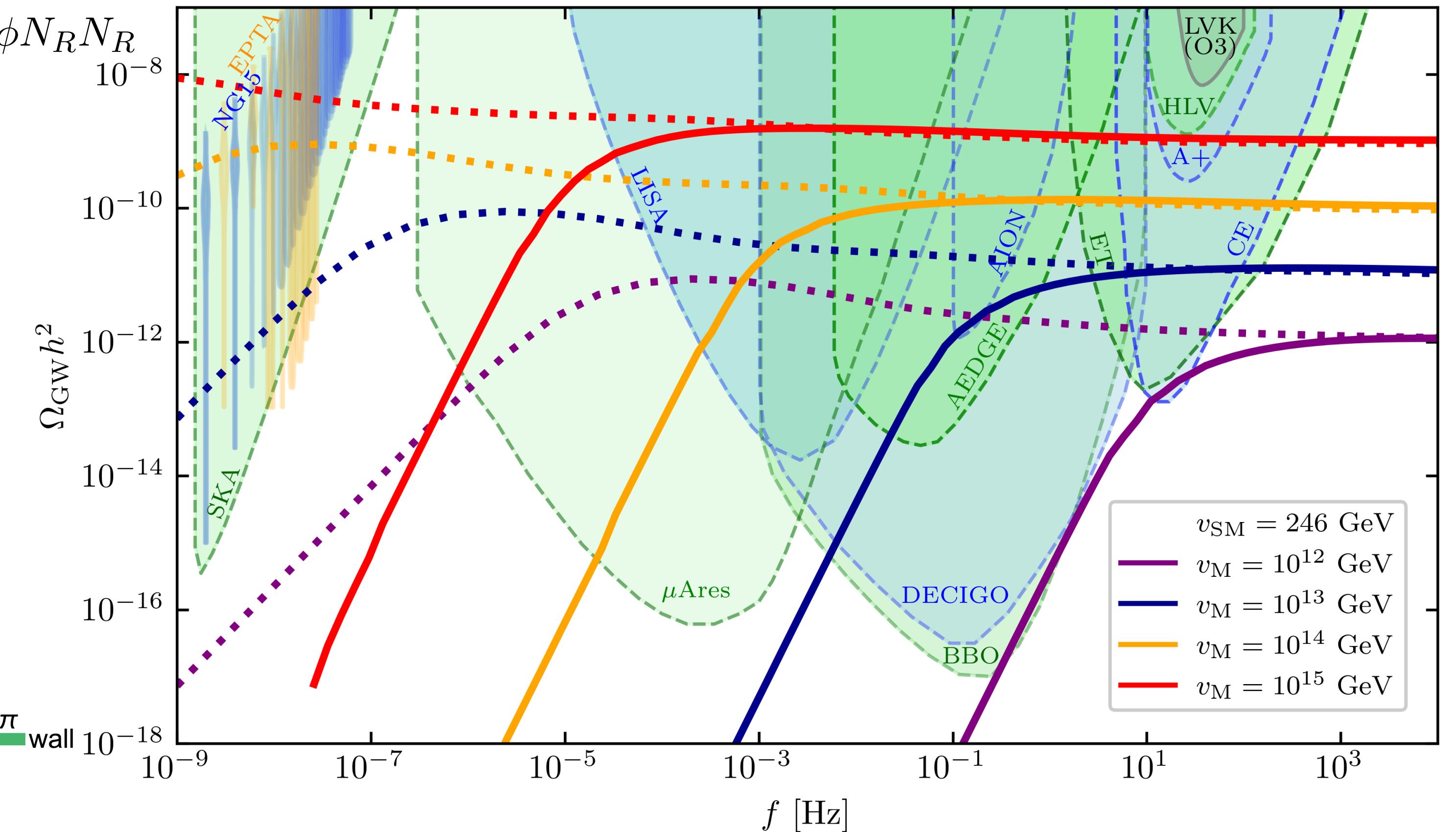


Vilenkin and Everett 1982

(a) Before Z_2 symmetry breaking.

(b) After Z_2 symmetry breaking.

DW decay without V_{bias} !



Surface tension in the walls causes the combined relic to decay earlier than strings

Conclusion

- GWs can probe new physics BSM at HE, only a few examples here: FOPT, CS, DW (+ combos)
- FOPT at QCD scale can describe NANOGrav
- CS $U(1)_{B-L}$ gauged w/GUTs; global w/Majorons
- DW Z_2 w/QG bias; Majorana vs Dirac
- DW bounded by CS in 2HDM (type I)