

Stability of multibrane models

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The banner features a yellow background with a dark silhouette of a city skyline on the left and a starburst of colorful arrows on the right.

- **5D models with branes**
 - multibrane models
- **Background configurations with Goldberger-Wise scalar**
- **Radion(s)**
 - equations of motion and boundary conditions
 - spectrum of radions
- **Conditions for stability of multibrane models**
- **Relation between radions and distances (between branes)**
- **Number of light radions**
- **Summary**

- 5D models with branes are very popular
 - **Randall Sundrum 1999**
25 years with more than 1 citation per day
- Approach to hierarchy problem
- Interesting phenomenology
 - field profiles in 5th dimension crucial for interactions
 - localization of different fields on different branes
- AdS/CFT interpretation
- Approach to phase transitions
- Dark matter
- ...

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- ...
- **More branes → more possibilities**

5D models with branes

Space-time is a warped product $\mathcal{M}^5 = \mathcal{M}^4 \times S^1/\mathbb{Z}_2$

with two orbifold-branes (fixed hyper-planes of \mathbb{Z}_2)

one (UV) located at $y = y_1$ and another (IR) located at $y = y_2$

and additional brane(s) in between, located at $y_{I_i}, i = 1, \dots$

All fields have well defined parities under \mathbb{Z}_2

$$S = \int_{\mathcal{M}^5} d^5x \sqrt{-g} \left[\frac{1}{2\kappa^2} R - \frac{1}{2} (\nabla\Phi)^2 - V(\Phi) - \sum_j \delta(y - y_j) U_j(\Phi) \right]$$

Φ is the Goldberger-Wise field

$j = 1, 2, I_1, \dots$ counts all branes

We are interested in background solutions of the form

(conformally flat coordinates)

$$ds^2 = e^{-2A(y)} (\eta_{\mu\nu} dx^\mu dx^\nu + dy^2) \quad \Phi = \phi(y)$$

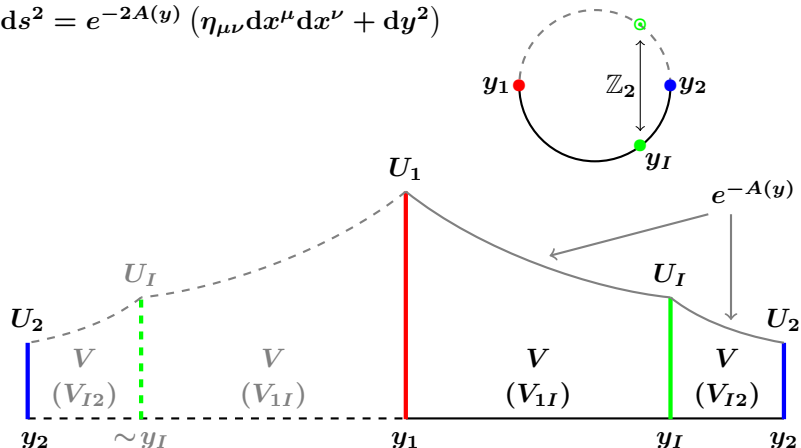
All results may be easily translated to often used coordinates

$$ds^2 = e^{-2\tilde{A}(z)} \eta_{\mu\nu} dx^\mu dx^\nu + dz^2$$

5D models with branes

$$\mathcal{M}^5 = \mathcal{M}^4 \times S^1/\mathbb{Z}_2$$

$$ds^2 = e^{-2A(y)} (\eta_{\mu\nu} dx^\mu dx^\nu + dy^2)$$



$$S = \int_{\mathcal{M}^5} d^5x \sqrt{-g} \left[\frac{1}{2\kappa^2} R - \frac{1}{2} (\nabla\Phi)^2 - V(\Phi) - \sum_j \delta(y - y_j) U_j(\Phi) \right]$$

Bulk equations of motion (for one interbrane section):

$$\begin{aligned} A'' + (A')^2 - \frac{\kappa^2}{3}(\phi')^2 &= 0 \\ \phi'' - 3A'\phi' - e^{-2A} V' &= 0 \\ (A')^2 - \frac{\kappa^2}{12}(\phi')^2 + \frac{\kappa^2}{6}e^{-2A} V &= 0 \end{aligned}$$

Boundary conditions at j -th brane $[f]_j := f(y_j^+) - f(y_j^-)$

$$[A']_j = \frac{\kappa^2}{3} e^{-A} U_j \quad [\phi']_j = e^{-A} U'_j$$

which at orbifold branes (because of \mathbb{Z}_2 symmetry) reduce to:

$$\lim_{y \rightarrow y_i^\pm} A' = \pm \frac{\kappa^2}{6} e^{-A} U_i \quad \lim_{y \rightarrow y_i^\pm} \phi' = \pm \frac{1}{2} e^{-A} U'_i$$

Procedure do obtain background solution satisfying all EoM all BC

- Two BC and 3rd bulk EoM at the “first” brane
 $\Rightarrow A'(y_1^+), \phi(y_1^+)$ and $\phi'(y_1^+)$ ($A(y_1^+)$ overall normalization)

- integration of 2 bulk EoM in the first interbrane section
 $\Rightarrow A(y)$ and $\phi(y)$ in that section

- **Position of the first intermediate brane, y_{I_1} , is not a free parameter. $A(y_{I_1})$ and $\phi(y_{I_1})$ must satisfy the condition:**

$$e^A (4A'U_{I_1}(\phi) - \phi'U'_{I_1}(\phi)) + 4\kappa^2 U_{I_1}^2(\phi) - \frac{1}{2}(U'_{I_1}(\phi))^2 + [V(\phi)]_{I_1} = 0$$

$$\Rightarrow A(y_{I_1}), A'(y_{I_1}^-), \phi(y_{I_1}), \phi'(y_{I_1}^-)$$

- 2 BC at this brane $\Rightarrow A'(y_{I_1}^+), \phi'(y_{I_1}^+)$
- ... same procedure at each intermediate brane ...
- two BC to be fulfilled at the “last” brane but only one parameter to adjust – the position of the “last” brane y_2
- **One tuning of parameters necessary to have any solution**

Background solutions:

- $A(y)$ and $\phi(y)$ which satisfy bulk equations of motion and all boundary conditions
- One tuning of parameters necessary to have any solution
 - the effective 4D cosmological constant must vanish
- For a given background solution positions of all branes are fixed
- \Rightarrow all distances between branes are fixed

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It is not guaranteed that a given solution is stable

- It is necessary to consider perturbations of the solution
- All modes of such perturbations must have positive (at least non-negative) masses squared

Scalar perturbations of a given back-ground solution

$$ds^2 = e^{-2A(y)} \{ [(1 + 2f_1(x, y)) \eta_{\mu\nu} + \partial_\mu \partial_\nu e(x, y)] dx^\mu dx^\nu + (1 + 2f_2(x, y)) dy^2 \}$$

$$\Phi = \phi(y) + f_3(x, y)$$

Most general scalar perturbations (in longitudinal gauge)

$e(x, y)$ must be included in models with intermediate brane(s)

[Pilo, Rattazzi, Zaffaroni 2000; Lee, Nakai, Suzuki 2022]

We expand all 5D perturbations in 4D modes

$$f_k(x, y) = \sum_n f_n(x) F_k^n(y), \quad k = 1, 2, 3$$

$$e(x, y) = \sum_n f_n(x) E^n(y)$$

We consider EoM for 4D mode(s) with mass squared m^2 and profile functions $F_1(y)$, $F_2(y)$, $F_3(y)$, $E(y)$

Off-diagonal Einstein equations relate the profiles F_1 , F_2 , F_3 , E

$$2F_2 + 4F_1 - E'' + 3A'E' = 0$$

$$3F_1' + 3A'F_2 + \kappa^2\phi'F_3 = 0$$

Bulk EoM for $E(y)$ is automatically fulfilled (due to other EoM)

\Rightarrow **only one independent scalar bulk perturbation**

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⇒ **only one independent scalar bulk perturbation**

One tower of scalar 4D KK modes

- infinitely many radions (scalar KK modes)
 - similarly to infinitely many gravitons (tensor KK modes)
- the number of radions does not depend on the number of branes

It is convenient to define $Q := e^{-2A} (F_1 + \frac{1}{2} A' E')$

Bulk EoM

$$-(p Q')' + q Q = m^2 p Q$$

$$p := \frac{3}{2\kappa^2} \frac{e^A}{(\phi')^2}$$
$$q := e^A$$

Radions

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Boundary conditions at j -th brane (non-trivial E' plays role)

$$\left[\frac{Q'}{\phi'} \right]_j = \frac{U'_j}{U_j} [Q]_j$$

$$B_j \left(\left\langle \frac{Q'}{\phi'} \right\rangle_j - \left[\frac{Q'}{\phi'} \right]_j \frac{\langle \phi' \rangle_j}{[\phi']_j} \right) + m^2 \left[\frac{Q}{\phi'} \right]_j = 0$$

$$B_j := e^{-A} U_j'' - [A']_j - \left[\frac{\phi''}{\phi'} \right]_j \quad \langle f \rangle_j = \frac{1}{2} (f(y_j^-) + f(y_j^+))$$

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For orbifold branes boundary conditions simplify to

$$0 = 0 \quad \frac{1}{2} B_i Q'(y_i^\pm) \pm m^2 Q(y_i^\pm) = 0 \quad +(-) \text{ for } i = 1(2)$$

Background configuration is stable (unstable)



the smallest radion mass squared is positive (negative)

Using solution of bulk EoM with $m^2 = 0$ satisfying BC at the orbifold brane at y_i :

$$Q_0(y) = e^{-2A(y)} + 2A'(y)e^{A(y)} \int_{y_i}^y e^{-3A(y')} dy'$$
$$Q'_0(y) = \frac{2}{3}\kappa^2(\phi'(y))^2 e^{A(y)} \int_{y_i}^y e^{-3A(y')} dy'$$

one can prove that the necessary and sufficient condition for the existence of at least one massless radion is

$$\prod_j B_j \phi'(y_j^-) \phi'(y_j^+) = 0$$

Analogous result for 2-brane models e.g. [Lesgourgues, Sorbo 2004]

Conditions for stability of multibrane models

It is also possible to show that the smallest radion m^2 is given by

$$m^2 = \min_Q \frac{\int p(Q')^2 + \int qQ^2 + \sum_I q \frac{[Q]_I^2}{[A']_I}}{\int pQ^2 + \frac{2}{B_1} pQ^2 \Big|_{y_1^+} + \frac{2}{B_2} pQ^2 \Big|_{y_2^-} + \sum_I \frac{3e^A}{2\kappa^2} \frac{1}{B_I} \left[\frac{Q}{\phi'} \right]_I^2}$$

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At least one radion is tachyonic if any of the following conditions is satisfied

- $B_j < 0$ at any brane
- $[A']_I \propto U_I < 0$ negative tension of any intermediate brane
- $\phi'(y) = 0$ at any bulk point

The last point is subtle because $p = \frac{3e^A}{2\kappa^2(\phi')^2}$ diverges for $\phi' = 0$.

It is necessary to use the Mukhanov-Sasaki variable $v(y)$ satisfying in this case: $Q = \lambda^{-1} e^{-2A} (\phi')^2 (A')^{-1} (e^{3A/2} A' (\phi')^{-1} v)'$ for which bulk EoM is regular also at points at which $\phi' = 0$

Necessary and sufficient conditions for stability of (multibrane) configuration

- each intermediate brane must have positive tension

$$U_{I_i}(\phi(y_{I_i})) > 0$$

- at each brane

$$e^{-A(y_j)} U_j''(\phi(y_j)) - [A']_j - \left[\frac{\phi''}{\phi'} \right]_j > 0$$

- everywhere

$$\phi'(y) \neq 0$$

Background configurations DO NOT depend on quantities U_j''

For any given background solution decreasing of any U_j'' may result in arbitrarily light radion and finally destabilize the system

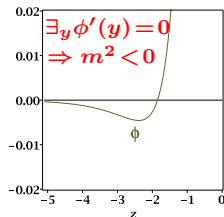
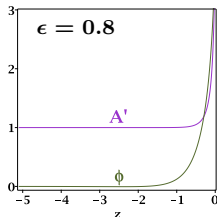
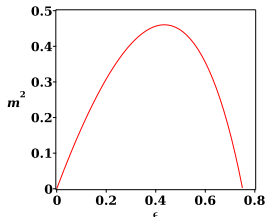
Upper bounds on radion(s) masses are obtained in the stiff potential approximation i.e. $U_j'' \rightarrow \infty$

However, even $U_j'' \rightarrow \infty$ DOES NOT guarantee stability

Conditions for stability of multibrane models

Example: simple 2-brane model with quadratic potentials

$$V(\Phi) = -\frac{6k^2}{\kappa^2} - \frac{1}{2}\epsilon k^2 \Phi^2 \quad U_i(\Phi) = \lambda_i + \mu_i(\Phi - v_i)^2$$

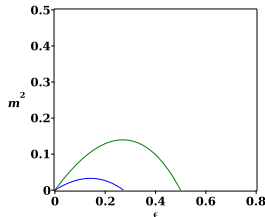


$$\frac{e^{-A}U_2''}{|[A']_2|} \approx 10^5 \gg 1$$

$$B_j = e^{-A}U_j'' - [A']_j - \left[\frac{\phi''}{\phi'}\right]_j$$

\sim stiff potential approximation

$$\frac{e^{-A}U_2''}{|[A']_2|} \approx 20$$
$$\frac{e^{-A}U_2''}{|[A']_2|} \approx 10$$



Relation between radions and distances

- Distance between points A and B with coordinates (x, y_A) and (x, y_B) is given by: $d_{AB} = \int_{y_A}^{y_B} \sqrt{g_{55}} dy$
- In our metric

$$\sqrt{g_{55}} = \sqrt{e^{-2A(y)} (1 + 2f_2(x, y))} \approx e^{-A(y)} (1 + f_2(x, y))$$

$$d_{AB} \approx \int_{y_A}^{y_B} e^{-A(y)} + \sum_k f_k(x) \int_{y_A}^{y_B} e^{-A(y)} F_2^k(y)$$

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Excitations of **each radion** field change distances between **fixed points on two different** branes.

Changes depend on t and \vec{x}

Unless for given k , y_A and y_B $\int_{y_A}^{y_B} e^{-A(y)} F_2^k(y) = 0$

Typically lighter radions (less zeros of $F_2(y)$) have stronger impact on distances than heavier radions (more zeros of $F_2(y)$)

We should express $F_2(y)$ in terms of the radion field $Q(y)$

Relation between radions and distances

Excitations of one radion with the profile $Q_{m^2}(y)$ change distances as follows:

$$\Delta d_{12} = -2f_{m^2}(t, \vec{x}) \int_{y_1}^{y_2} e^A Q_{m^2} dy$$

$$\Delta d_{1I(I_2)} = f_{m^2}(t, \vec{x}) \left(-2 \int_{y_1(y_I)}^{y_I(y_2)} e^A Q_{m^2} dy + (-) e^{A(y_I)} \frac{[Q_{m^2}]_I}{[A']_I} \right)$$

In general, the distance between fixed points on two different branes is changed by a radion even if (the integral of) its profile between those branes vanishes

Exception: the total distance between the “end-of-the-world” orbifold branes

$$e^{A(y_I)} \frac{[Q_{m^2}]_I}{[A']_I} = e^{2A(y_I)} \frac{3[Q_{m^2}]_I}{\kappa^2 U_I}$$

Brane with smaller tension U_I is more easy to bend
(distance becomes x -dependent)

Relation between radions and distances

Non-zero radion field Q_{m^2}

- $F_2 \neq 0$
 g_{55} is changed \Rightarrow distances **between** branes are affected
- $F_1 \neq 0$ $E \neq 0$
 $g_{\mu\nu}$ is changed \Rightarrow distances **along** branes are also affected

Distance between points C and D with coordinates (t, x_C^1, x^2, x^3, y) and (t, x_D^1, x^2, x^3, y)

$$\begin{aligned}d_{CD} &\approx e^{-A(y)}(x_D^1 - x_C^1) \\ &+ e^{+A(y)} \left[\frac{Q_{m^2}}{A'} \right]_j \left[\frac{1}{A'} \right]_j^{-1} \int_{x_C^1}^{x_D^1} f_{m^2}(t, \vec{x}) dx^1 \\ &+ \frac{1}{2} e^{-A(y)} E_{m^2}(y) \int_{x_C^1}^{x_D^1} \partial_{xx} f_{m^2}(t, \vec{x}) dx^1\end{aligned}$$

Excitations of (each) radion field modify distances **along** branes

How many radions may be light?

- number of massless radions in 2-brane models:

- **0**

- if $B_1 \cdot B_2 \cdot \phi'(y_1^+) \cdot \phi'(y_2^-) \neq 0$

- **1**

- if $B_1 \cdot B_2 \cdot \phi'(y_1^+) \cdot \phi'(y_2^-) = 0$ but at least one of $B_i \neq 0$

- **2**

- if $B_1 = B_2 = 0$

How many radions may be light?

- number of massless radions in 2-brane models:

- 0

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- 1

- if $B_1 \cdot B_2 \cdot \phi'(y_1^+) \cdot \phi'(y_2^-) = 0$ but at least one of $B_i \neq 0$

- 2

- if $B_1 = B_2 = 0$

- number of massless radions in 3-brane models

- 0

- if $B_1 \cdot B_I \cdot B_2 \cdot \phi'(y_1^+) \cdot \phi'(y_I^-) \cdot \phi'(y_I^+) \cdot \phi'(y_2^-) \neq 0$

- 1

- if $B_1 \cdot B_I \cdot B_2 \cdot \phi'(y_1^+) \cdot \phi'(y_I^-) \cdot \phi'(y_I^+) \cdot \phi'(y_2^-) = 0$
but at least two of $B_j \neq 0$

- 2

- if exactly two $B_j = 0$

- one localized in one interbrane space if $B_I \cdot \phi'(y_I^\pm) = 0$

- 4

- if $B_1 = B_I = B_2 = 0$

- two localized in interbrane spaces

- **5D multibrane models**

- one fine tuning necessary for solutions with flat 4D sections
- all distances between branes are fixed by a given background solution – warp factor $e^{-2A(y)}$ and GW scalar $\phi(y)$

- **Radions: 4D modes of scalar perturbations**

- one more 5D scalar perturbation must be taken into account (as compared to 2-brane models)
- still one (infinite) tower of KK modes – 4D radions

- **Bulk EoM and especially BC for radions**

are generalized in the presence of intermediate branes

- **Conditions for stability of a given background configuration**

$$U_{I_i} > 0 \quad e^{-A} U_j'' - [A']_j - \left[\frac{\phi''}{\phi'} \right]_j > 0 \quad \phi' \neq 0$$

- background configurations do not depend on U_j''
- radion(s) may be light or even tachyonic also for $U_j'' \rightarrow \infty$

- **Each radion field excitation deforms branes**

changing distances in all directions: **between and along** branes

- in most cases in a non-trivial way

- **There may be up to 2^{N-1} massless (light) radions in a model with N branes**