



ALP-Assisted Electroweak Phase Transition and Baryogenesis

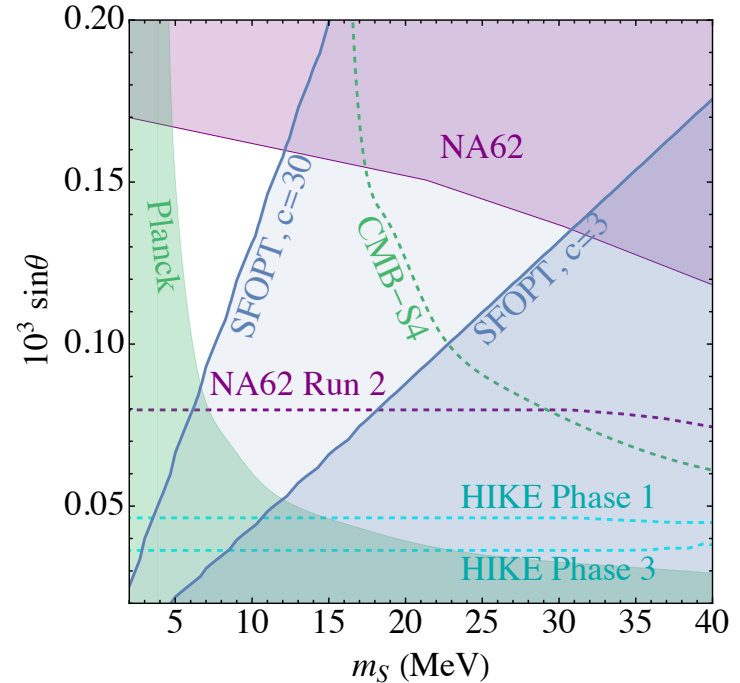
Isaac R. Wang

PIKIMO Fall 2023, 11/11/2023

Based on: 2309.00587, Collaborator: Prof. Keisuke Harigaya

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- Challenges of electroweak baryogenesis and phase transition (EWPT) $\delta=\pi/5$
 - Baryogenesis problem and principles of solution
 - Electroweak phase transition
 - Challenges
- Axion-like particle (ALP) couples with Higgs
- Experimental Probes
 - ΔN_{eff} constraint
 - Rare meson decay
 - Collider searches



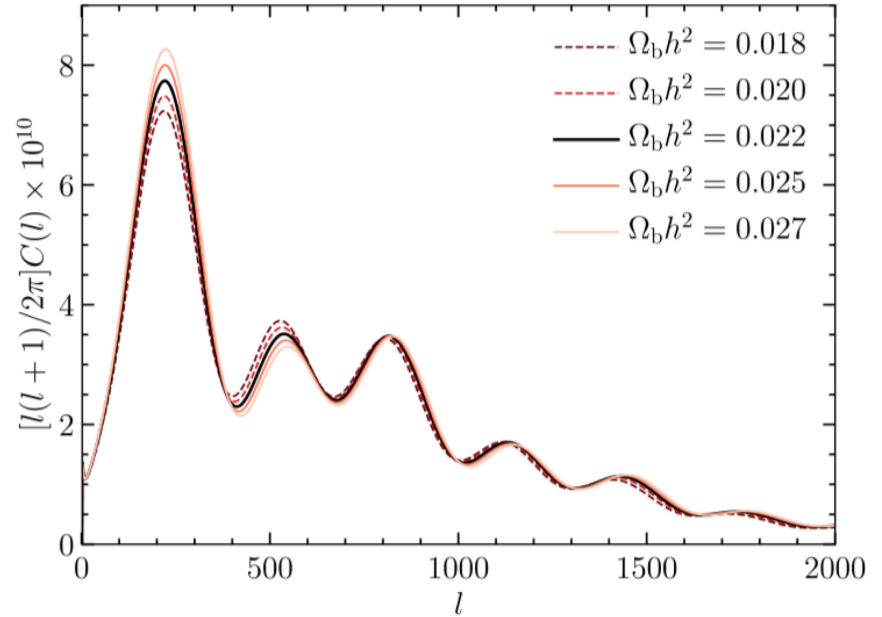
Part of the final results. White spaces are allowed.

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Matter-Antimatter Asymmetry: Baryogenesis Problem

- matter > anti-matter
- Define $n_B = n_{\text{baryon}} - n_{\text{anti-baryon}}$
- $\frac{n_B}{s} \simeq 9 \times 10^{-11}$
- What's the origin?



Planck 2018b

Sakharov Condition

- Baryon number violation

SM: sphaleron process. Violates $B + L$ but keeps $B - L$.

- C and CP violation

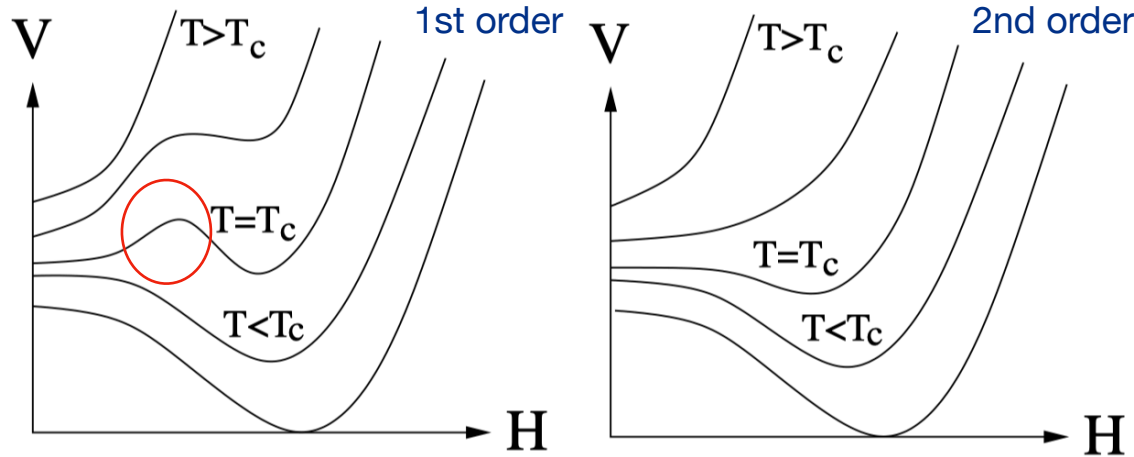
SM: CKM (too small), or UV scale new physics, model dependent.
(Not the main focus today.)

- Out of thermal equilibrium

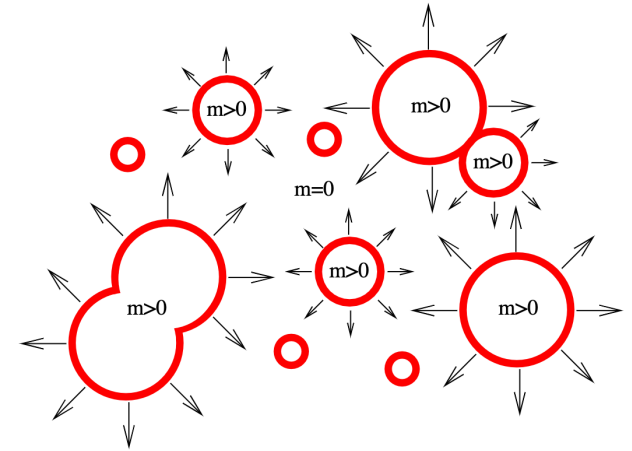
SM: a strong 1st-order electroweak phase transition (need BSM!)

[Sakharov, 1967]

EWPT: 1st order vs 2nd order



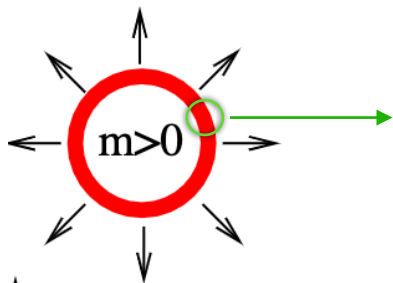
1st order: clear, well-defined barrier
2nd order: smooth crossover.



1st-order PT processes via *bubble nucleation*

Figure from:
[J Cline: hep-ph/0609145]

Electroweak baryogenesis from 1st-order EWPT



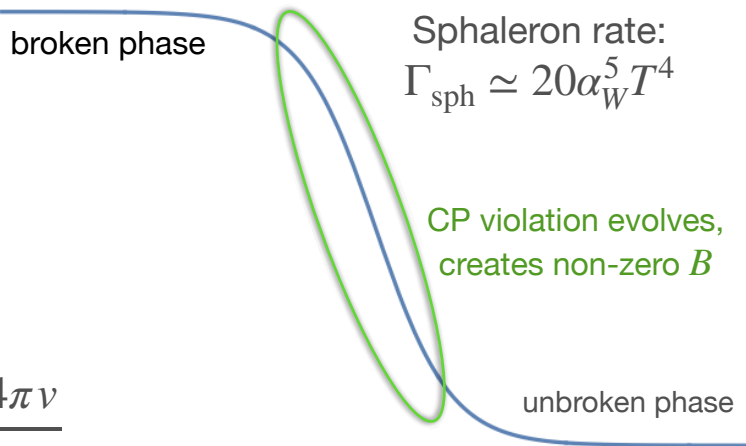
Sphaleron rate in the bubble:

$$\Gamma_{\text{sph}} \sim \exp(-E_{\text{sph}}/T), E_{\text{sph}} \propto \frac{4\pi v}{g}$$

suppressed by v

Need to be turned off inside the bubble!

Require $\frac{v_n}{T_n} \geq 1$, T_n : bubble nucleation



Sphaleron rate:
 $\Gamma_{\text{sph}} \simeq 20\alpha_W^5 T^4$

CP violation evolves,
 creates non-zero B

Wall profile

Typically (not supercooling):

$$\frac{v(T_c)}{T_c} \simeq \frac{v(T_n)}{T_n}$$

[F. R. Klinkhamer and N. S. Manton,
 Phys. Rev. D 30, 2212]

[V. A. Kuzmin, V. A. Rubakov and M. E. Shaposhnikov,
 Phys. Lett. B 155, 36 (1985)]

[M. E. Shaposhnikov, JETP Lett. 44, 465 (1986), Nucl.
 Phys. B 287, 757 (1987)]

[A. G. Cohen, D. B. Kaplan and A. E. Nelson, hep-ph/
 9302210]

[D'Onofrio et al: 1404.3565]

Phase Transition Strength

- Finite-T: thermal correction to the effective potential V .

- Boson contribution: $\frac{T^4}{2\pi^2} n_B \left(\frac{\pi^2}{12} \left(\frac{m}{T} \right)^2 - \frac{\pi}{6} \left(\frac{m}{T} \right)^3 + \dots \right)$,

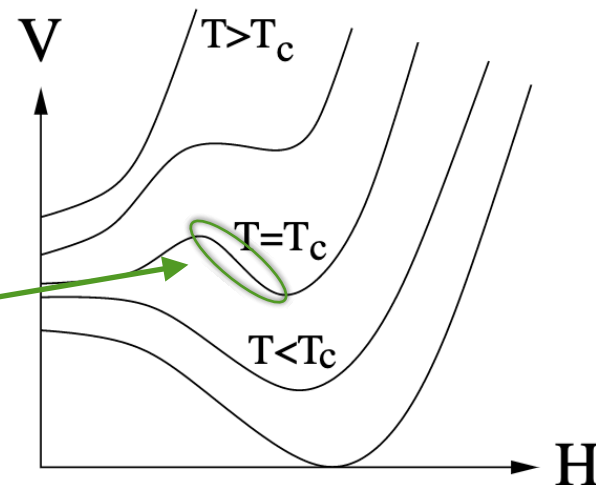
n_B : degree of freedom

- Fermion contribution: $\frac{T^4}{2\pi^2} n_F \left(\frac{\pi^2}{24} \left(\frac{m}{T} \right)^2 + \dots \right)$, n_F : degree of freedom.

- Total: $V = DT^2 h^2 - \frac{1}{2} \mu^2 h^2 - ETH^3 + \frac{1}{4} \lambda h^4$, cubic from

bosons

- $\frac{v(T_c)}{T_c} = \frac{2E}{\lambda}$. **SM: 0.2! (Lattice confirmed.)**



[K. Jansen, hep-lat/9509018], [K. Kajantie et al, hep-lat/9510020]
 [K. Rummukainen, hep-lat/9608079], [K. Kajantie et al, hep-ph/9605288.]
 [M. Gurtler et al, hep-lat/9704013], [F. Csikor et al, hep-ph/9809291]
 [M. Laine and K. Rummukainen, hep-ph/9804255, hep-lat/9804019]
 [K. Rummukainen et al, hep-lat/9805013], [Z.Fodor, hep-lat/9909162]

Only for intuition. Needs full form to compute!

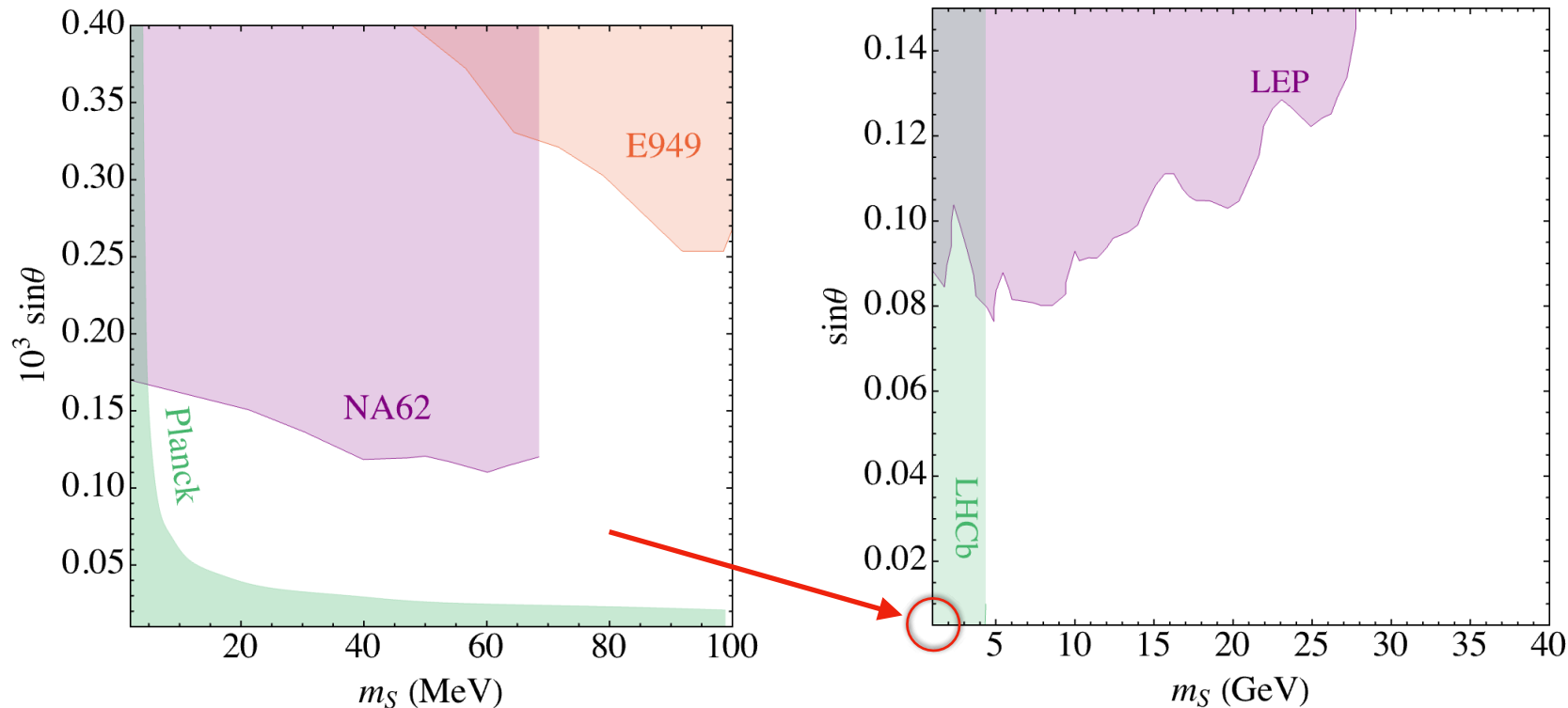
A simple solution: extra singlet scalar

- Introduce S : $(1,1,0)$ singlet
- General: $\mathcal{L} = \text{SM} + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{4}\lambda_S S^4 + \text{other interactions}$
- Only interact with Higgs: $SSH\bar{H}$, S^4 , S^3 , $SH\bar{H}$, depending on the symmetry
- Probe: mixing with Higgs, $\sin \theta$.

$$\frac{v(T_c)}{T_c} = \frac{2E}{\lambda}$$

What are the current constraints/status?

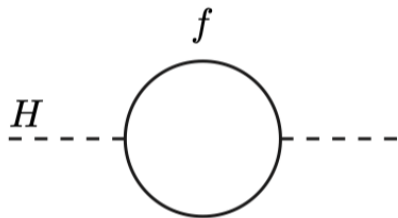
Experimental bound



Stringent experimental bound on extra singlet scalar!

Extra hierarchy problem

The SM electroweak hierarchy problem

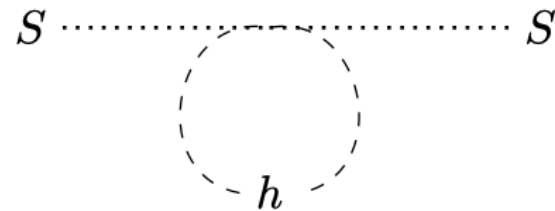


$$\delta m_H^2 = -\frac{y_f^2}{8\pi^2} \Lambda_{UV}^2$$

Quadratic sensitive to Λ_{UV} :
Huge quantum corrections!

The extra hierarchy problem from extra singlet

Typical models have $\frac{1}{4} \lambda_{hS} S^2 h^2$



$$\delta m_S^2 = \frac{\lambda_{hs}}{16\pi^2} \Lambda_{UV}^2$$

Typical singlet scalar model introduced
extra hierarchy problem!

Some traditional solutions: SUSY, compositeness....

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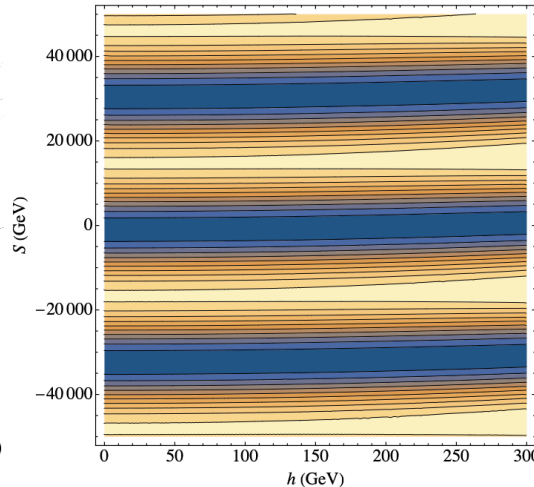
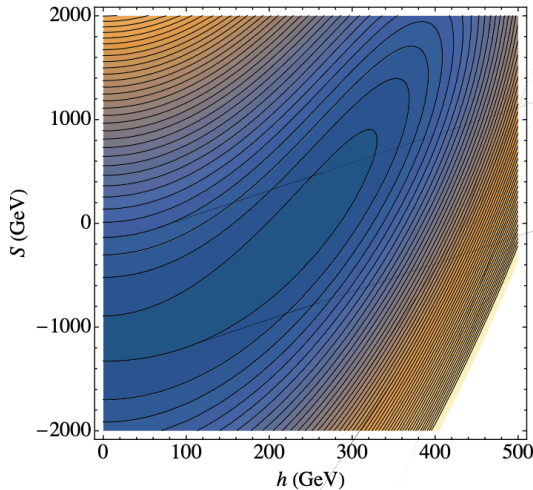
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Axion-like Particles (ALP) Couples with Higgs

ALP S : angular mode of P $P = (f + \Phi)^2 \exp(iS/f)$

pNGB. Weakly interacting. Naturally light.

$$V(H, S) = -(\mu_H^2 - Af \cos \delta) |H|^2 + \lambda |H|^4 + \mu_S^2 f^2 \left(1 - \cos\left(\frac{S}{f}\right) \right) - Af (|H|^2 - v^2) \cos\left(\frac{S}{f} - \delta\right),$$



EW scale vev: $\langle h \rangle = v$, $\langle S \rangle = 0$
 Periodic in S .

$$A = \frac{1}{2\sqrt{2}v} (m_h^2 - m_S^2) \sin(2\theta) \csc(\delta),$$

$$\lambda = \frac{1}{8v^2} (m_h^2 + m_S^2 + (m_h^2 - m_S^2) \cos(2\theta)),$$

$$\mu_H^2 = \frac{1}{4} (m_h^2 + m_S^2 + (m_h^2 - m_S^2) \cos(2\theta)),$$

$$\mu_S^2 = \frac{1}{2} (m_h^2 + m_S^2 - (m_h^2 - m_S^2) \cos(2\theta)).$$

[K. Harigaya and IRW, 2309.00587]

Generate Barrier at Tree Level

A barrier is generated at the tree level along the path $\partial V/\partial S = 0$

$$S = \frac{Af(h^2 - 2v^2) \sin \delta}{2f\mu_S^2 + A(h^2 - 2v^2) \cos \delta}$$
$$\simeq \frac{Afv^2 \sin \delta}{Av^2 \cos \delta - f\mu_S^2} + \frac{Af^2\mu_S^2 \sin \delta}{2(Av^2 \cos \delta - f\mu_S^2)^2} h^2 + \frac{A^2 f^2 \mu_S^2 \sin \delta \cos \delta}{4(Av^2 \cos \delta - f\mu_S^2)^3} h^4 + O(h^6).$$

Higgs Quartic term becomes: $\frac{1}{4} \left(\lambda - \frac{A^2 f^3 \mu_S^4 \sin^2 \delta}{2(\mu_S^2 f - Av^2 \cos \delta)^3} \right) h^4.$

Quartic becomes negative for viable parameter region where SF OPT is achieved!

Bounded from below condition protected by a positive $O(h^6)$ term!

Extra contribution to the barrier at the tree level.

[K. Harigaya and IRW, 2309.00587]

Thermal phase transition

High-T limit, large f limit.

At high- T : $\langle h \rangle = 0$, restored symmetry! $V_{\text{high}T} = -\frac{1}{2}(\mu_H^2 - Af \cos \delta)h^2 + D_{\text{SM}}T^2h^2 - E_{\text{SM}}Th^3 + \frac{1}{4}\lambda h^4$

At low- T : $\langle h \rangle = v(T)$

$$-\frac{1}{2}Af \cos \left(\frac{S}{f} - \delta \right) \left(h^2 - 2v^2 + \frac{1}{3}T^2 \right)$$

High- f limit and high- T expansion: $\langle S \rangle = \frac{A}{2\mu_S^2}(h^2 - v^2)$

$$-f^2\mu_S^2 \cos \left(\frac{S}{f} \right).$$

1-step PT: $(0, \langle S \rangle(0, T)) \rightarrow (v(T), \langle S \rangle(v(T), T))$

$$f_c \equiv \frac{A}{\mu_S^2} v^2 \sin \delta,$$

Again. For a quick look, not for careful computation!

Parametrize: $f \equiv cf_c$

[K. Harigaya and IRW, 2309.00587]

Full Effective Potential Computation

$$V_{\text{eff}} = V_0 + V_{\text{CW}} + V_{\text{FT}}$$

$$V_{\text{CW}} = \frac{1}{64\pi^2} \left(\sum_B n_B \left(\log \left(\frac{m_B^2(h, S)}{Q^2} \right) - c_B \right) - \sum_F n_F \left(\log \left(\frac{m_F^2(h, S)}{Q^2} \right) - c_F \right) \right)$$

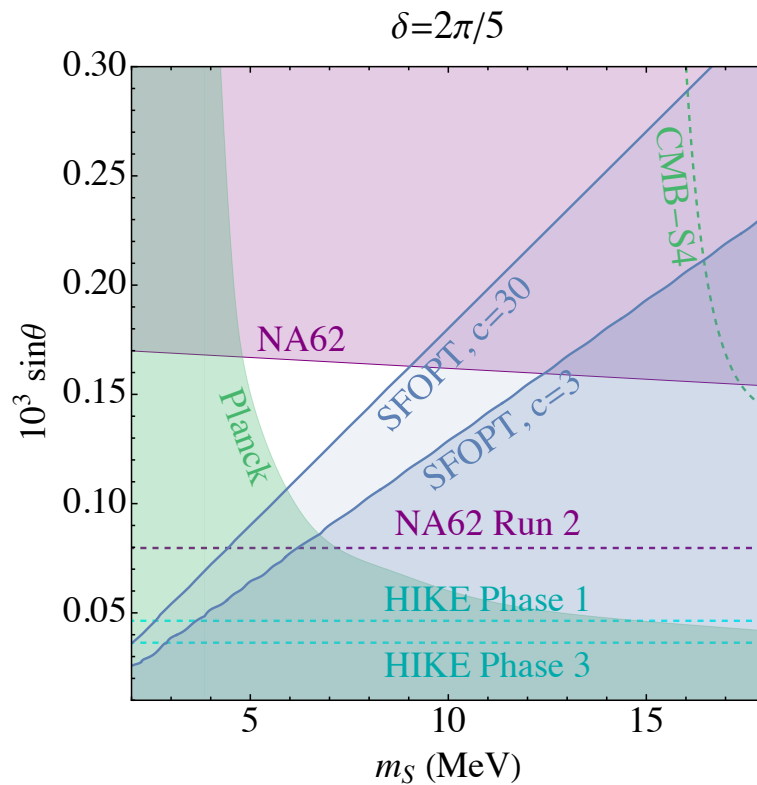
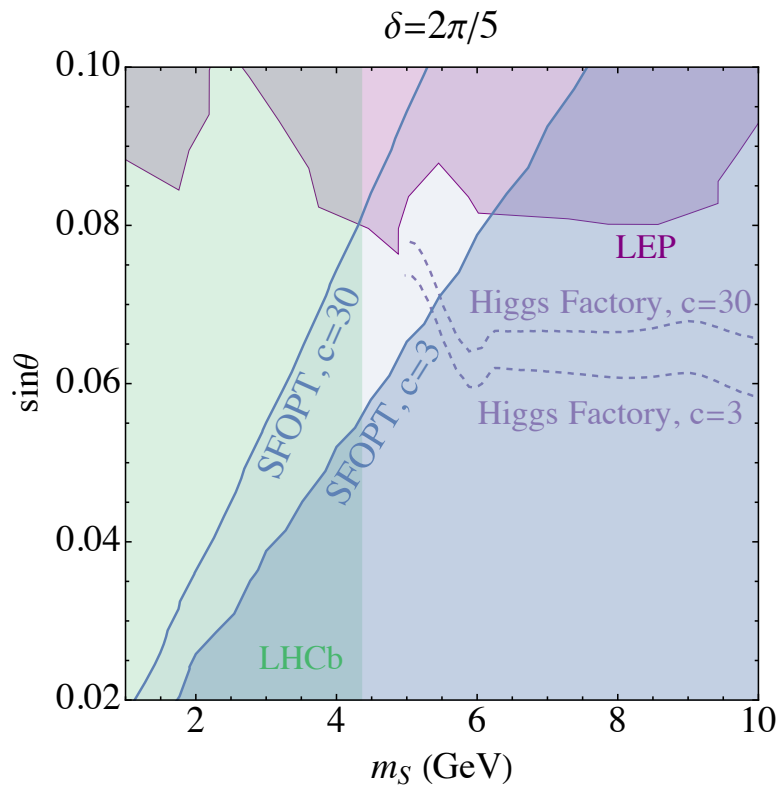
$$V_{\text{FT}} = \frac{T^4}{2\pi^2} \left(\sum_B n_B J_B \left(\frac{m_B^2(h, S)}{T^2} \right) + \sum_F n_F J_F \left(\frac{m_F^2(h, S)}{T^2} \right) \right),$$

$$J_{B,F}(x^2) = \pm \int_0^\infty dy y^2 \log \left(1 \mp \exp \left(-\sqrt{y^2 + x^2} \right) \right).$$

Resummation: all boson masses replaced by thermal mass: $m^2 \rightarrow m^2 + \Pi$, $\Pi \sim g^2 T^2$

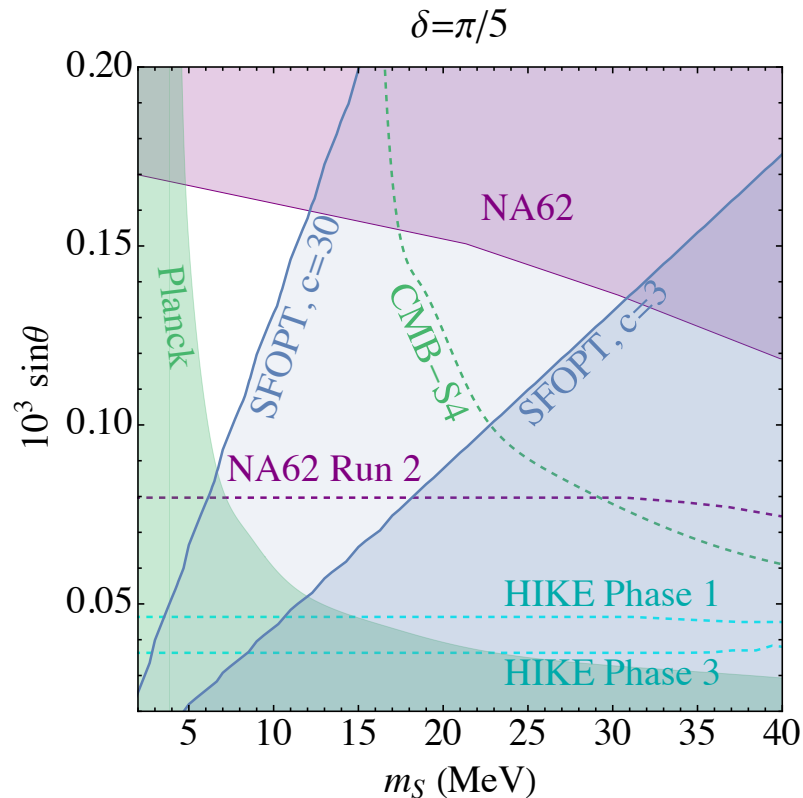
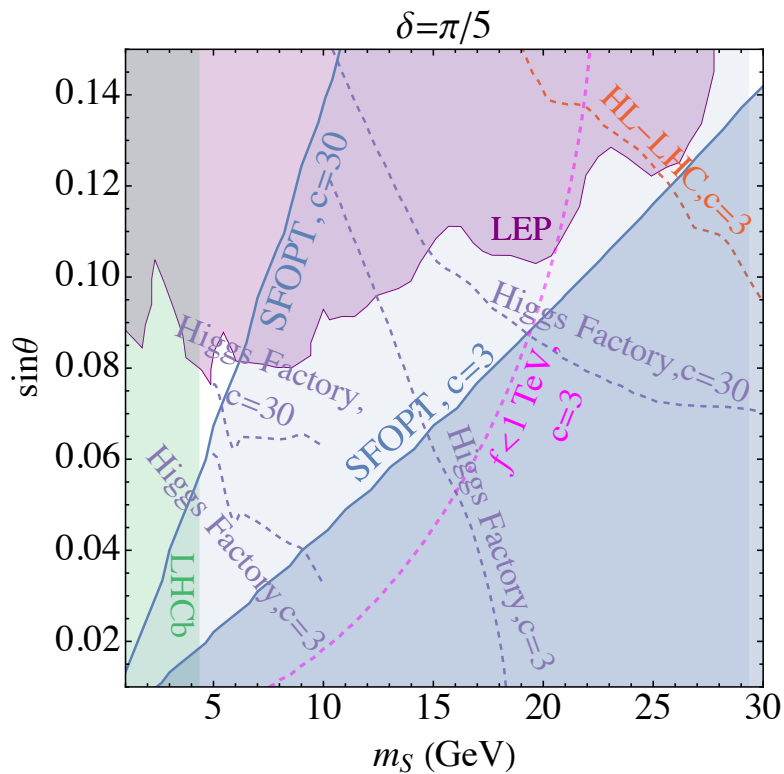
[K. Harigaya and IRW, 2309.00587]

Results



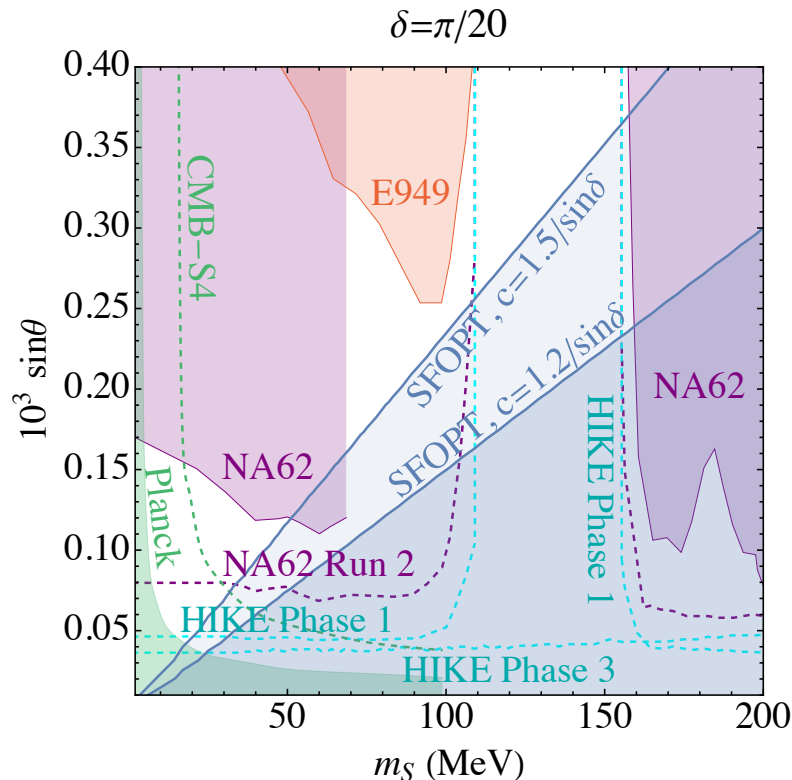
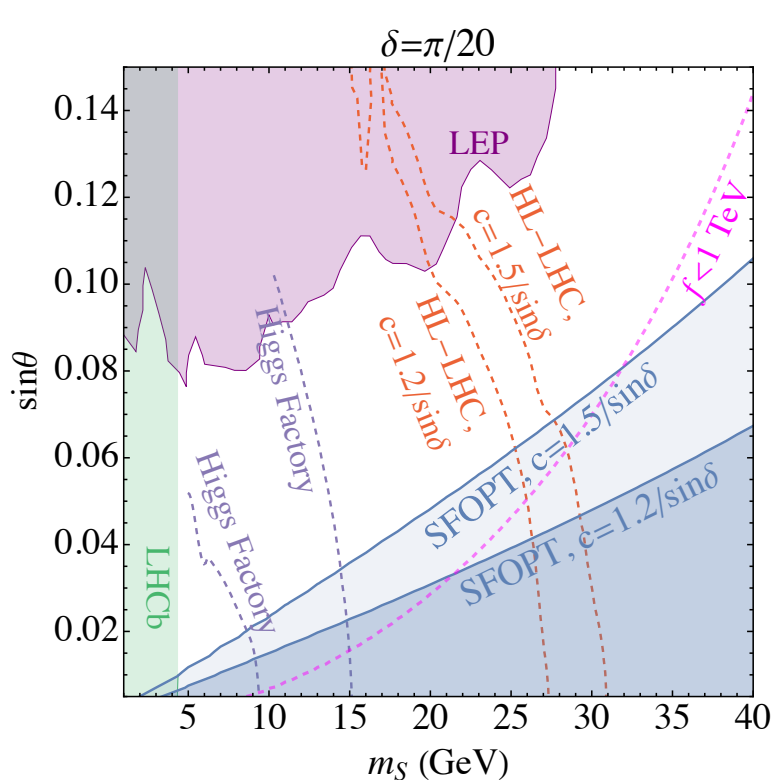
[K. Harigaya and IRW, 2309.00587]

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Collider Signals

S can mix with h , mixing angle $\sin \theta$.

Generating vertex:

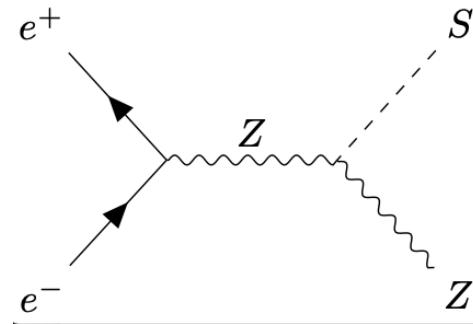
$$hZZ \rightarrow \sin \theta SZZ$$

$$h^3 \rightarrow \sin \theta hhS, \sin^2 \theta hSS$$

General probe:

- scalar production: $SM \rightarrow S + SM$
- Higgs exotic decay: $h \rightarrow SS$

Collider search can probe the extra singlet scalar at **GeV** scale



[K. Harigaya and IRW, 2309.00587]

Rare Meson Decay and ΔN_{eff}

- Extra decay channel for B meson:
- $B^0 \rightarrow K^0 S, B^+ \rightarrow K^+ S$, searched by LHCb at $200 \text{ MeV} < m_S < 4 \text{ GeV}$
- Extra decay channel for Kaon:
- $K^+ \rightarrow \pi^+ S, K^0 \rightarrow \pi^0 S$, searched by NA62, KLEVER.... for MeV scale.
- MeV scale m_S : large energy density when neutrino decouples.
- S decays into γ : negative ΔN_{eff}

Review: [PBC Group, 1901.09966], [E. Goudzovski et al, 2201.07805]
[LHCb Collaboration, 1508.04094, 1612.07818, 1703.08501]
[NA62 Collaboration, 2010.07644, 2103.15389]
[KLEVER Project Collaboration, 1901.03099]
[M. Ibe et al, 2112.11096], [Planck Collaboration, 1807.06209]
[CMB-S4 Collaboration, 1610.02743],[K. Harigaya and IRW, 2309.00587]

$$\partial_\mu J_\mu^B = \partial_\mu J_\mu^L = \frac{3g^2}{32\pi^2} W\tilde{W}$$

A general effective CP-violating operator: $\mathcal{L} = \frac{\alpha_2}{8\pi} \frac{S}{M} W\tilde{W}$, $\frac{\partial_\mu S}{M} q^\dagger \bar{\sigma}^\mu q$, or $\frac{\partial_\mu S}{M} \ell^\dagger \bar{\sigma}^\mu \ell$, M : UV scale.

CP-violation comes from $W\tilde{W}$ or $f^\dagger \bar{\sigma} f$! This operator itself is CP-violating.

Rewrite: $\mathcal{L}_{CP} \propto \frac{1}{M} (\partial_0 S) n_B$ in thick-wall regime!

From minimizing free energy, n_B gets a minimum $n_B^0 \propto \frac{1}{M} (\partial_0 S) \langle S \rangle \simeq \frac{A}{2\mu_S^2} h^2$

Thus $\dot{n}_B \propto \frac{\Gamma_{\text{sph}}}{T^3} (n_B - n_B^0) \simeq \frac{\Gamma_{\text{sph}}}{T^3} n_B^0 = \frac{\Gamma_{\text{sph}}}{T^3} \boxed{\dot{S}}$ Large field-value shift! For old literature using $\mathcal{L}_{CP} \propto \frac{\sin(\delta)}{M^2} h^2 W\tilde{W}$:

[A. Cohen and B. Kaplan, Phys.Lett.B 199 (1987) 251-258,

Nucl.Phys.B 308 (1988) 913-928]

[M. Dine et al, Phys.Lett.B 257 (1991) 351-356]

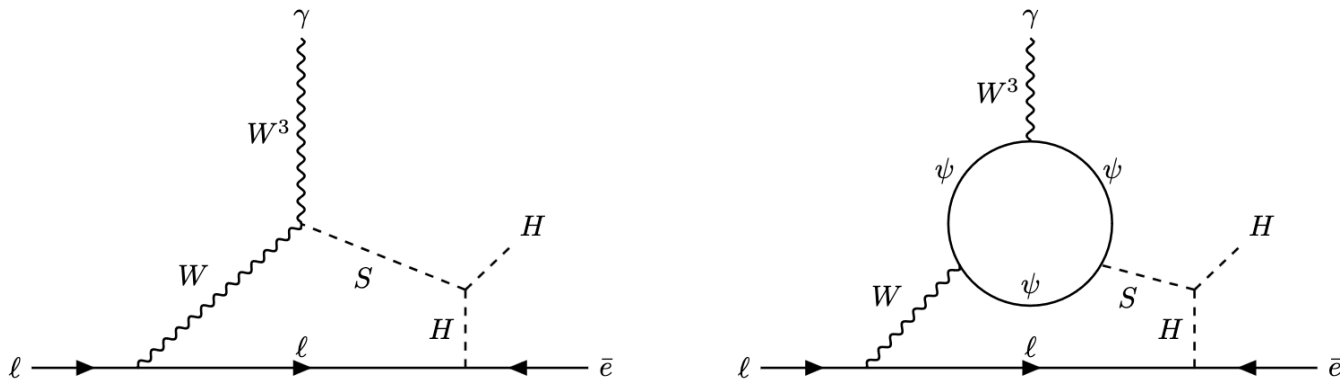
[M. Dine, hep-ph/9206220]

[K. Harigaya and IRW, 2309.00587]

$$\frac{n_B}{s} \simeq \frac{3\Gamma_{\text{ws}} T^2 \Delta S}{Ms}$$

BAU properly produced for $M \gtrsim f$

The electric dipole moment



$$iF_{\mu\nu}\ell\sigma^{\mu\nu}\bar{e}H \times \frac{ey_e}{(16\pi^2)^2} \frac{A \sin \delta}{Mv^2}.$$

when M consistent with n_B :

EDM around $10^{-32} - 10^{-30}$ for GeV scale m_S , can be probed by future experiments.

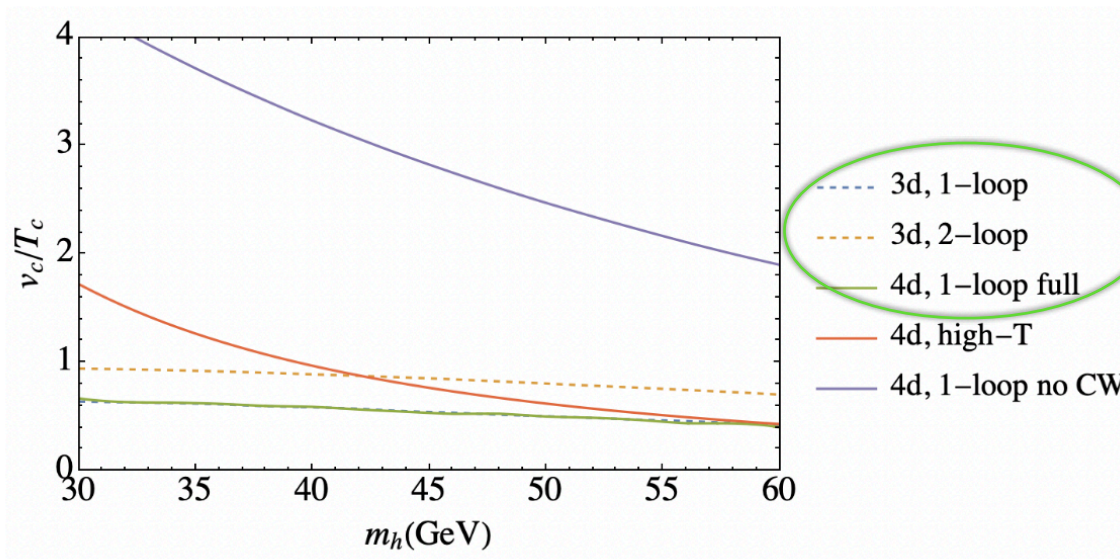
[K. Harigaya and IRW, 2309.00587]

Conclusion

- Baryon number asymmetry remains unsolved. SM EWBG fails due to lack of a SFOPT at the EW scale.
- We propose an ALP to enhance the EWPT strength.
- Successful baryogenesis can be achieved via the local baryogenesis mechanism.
- Various experiments can be applied to probe the parameter space.

Backups

SM EWPT strength with lighter Higgs



Most trustable.
Main suppression:
top Yukawa correction

[P. Arnold, O. Espinosa, hep-ph/9212235]
[K. Harigaya and IRW, 2309.00587]