# Quarkonium isolation

#### Lidia Carcedo\*

Universidad de Alcalá, Universidade de Santiago de Compostela \* lidia.carcedo.salgado@cern.ch

Quarkonia as Tools 2024 09/01/2024















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- 1. Introduction
- 2. Models
- 3. Measurements
- 4. Conclusions



### 1 - Introduction Why (I am interested in) quarkonium isolation?

- A part of my analysis consists of defining an isolation criterium for  $J/\psi$  and applying it to get a measurement of isolated  $J/\psi + \gamma$  in pp collisions at 13 TeV at LHCb
- Why isolated  $J/\psi$ ?
  - Access Transverse Momentum Dependent distributions (TMDs) won't cover the interest of this in today's talk

### Provide additional information on the production mechanism

What is covered by this talk:

Isolation criteria designed for photons Quarkonium fragmentation functions Modelization of quarkonia production Measurements of  $J/\psi$  in jets

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#### Review by Marius Höfer

- Lots of work has been done on photon isolation
  - Necessary in hadron collider experiments to distinguish photons produced in hard scattering processes from others
- Philosophy: a photon is isolated if most of the transverse energy around the photon is carried by the photon itself

### Types of isolation criteria

- (a) <u>Cone-based isolation</u>
  - i) Hard cone
  - ii) Smooth cone
  - iii) Modified versions (backup)

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- (b) <u>Clustering-based isolation</u>
  - i) Democratic isolation
  - ii) Softdrop isolation



**Review by Marius Höfer** 



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 $E_T^{\max}$ 

Ehad

### Cone-based isolation

#### Hard cone **i**)

 $\sum E_T^{had} \le E_T^{max} \quad \text{no dependence on how the} \\ \text{energy is distributed with } r$ 

includes fragmentation

can be implemented experimentally

### ii) Smooth cone

 $\sum E_T^{had} \le f(r \le R) \quad \begin{array}{l} \text{allowed energy around the} \\ \text{photon decreases with } r \end{array}$ 

removes fragmentation can't be implemented experimentally



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#### **Review by Marius Höfer**

### Clustering-based isolation



1) Cluster hadrons + photons democratically into a jet 2) Evaluate the z of jets 3) If one of them fulfils  $z > z_{min} \Longrightarrow$  isolated 'photon'



#### ii) Soft-drop isolation (smooth-like)

Cluster hadrons + photons democratically into a parent jet
 De-cluster into two subjets (with C/A algorithm)
 Evaluate the soft-drop condition

 a) Passed ⇒ keep parent jet
 b) Failed ⇒ keep softer subjet and repeat 1)

 If soft-drop condition failed at all steps and remaining constituent is a photon ⇒ isolated photon

#### Soft-drop condition



#### **Review by Marius Höfer**

#### Clustering-based isolation **Democratic isolation** (hard-like) 1) Cluster hadrons + photons democratically into a jet $z = \frac{z_{EM}}{z_{EM} + z_{HAD}} > z_{min}$ 2) Evaluate the z of jets 3) If one of them fulfils $z > z_{min} \Longrightarrow$ isolated 'photon' includes fragmentation ii) Soft-drop isolation (smooth-like) Soft-drop condition $\frac{\min(p_{T1}, p_{T2})}{p_{T1} + p_{T2}} \ge z_{cut} \left(\frac{R_{12}}{R_0}\right)^{\mu}$ 1) Cluster hadrons + photons democratically into a parent jet 2) De-cluster into two subjets (with C/A algorithm) 3) Evaluate the soft-drop condition a) Passed $\implies$ keep parent jet b) Failed $\implies$ keep softer subjet and repeat 1) 4) If soft-drop condition failed at all steps and remaining constituent is a photon $\implies$ isolated photon removes fragmentation

### 1-Introduction Non-Relativistic QCD (NRQCD)

The physics of quarkonia involves several energy scales:



### 1-Introduction Non-Relativistic QCD (NRQCD)

The physics of quarkonia involves several energy scales:

$$\begin{array}{lll} (m_Q v^2)^2 & \ll & (m_Q v)^2 & \ll & m_Q^2 \\ \hline \mbox{kinetic energy} & \mbox{momentum} & \mbox{mass} \\ \mbox{interaction} \\ \mbox{time} & \mbox{(spatial size)}^{-1} & \mbox{distance range for} \\ Q \overline{Q} \mbox{ creation} \end{array}$$

NRQCD keeps track of this scale hierarchy:

$$d\sigma(pp \to @+X) = \sum_{s,L,J} \left\langle \mathcal{O}^{@}[^{2S+1}L_{J}] \right\rangle \begin{array}{l} d\hat{\sigma}(pp \to Q\overline{Q}[^{2S+1}L_{J}] + X) \\ \text{long distance} \\ \text{effects} \end{array}$$

• The sum in  $d\sigma(pp \rightarrow Q + X)$  for a given physical quarkonium state is over the expansion of its Fock states:

$$\begin{split} | \mathcal{Q}[^{2S+1}L_J] \rangle &= \mathcal{O}(1) | \mathcal{Q}\overline{\mathcal{Q}}[^{2S+1}L_J^{(1)}] \rangle + \mathcal{O}(v) | \mathcal{Q}\overline{\mathcal{Q}}[^{2S+1}(L\pm 1)_{J'}^{(8)}]g \rangle \\ &+ \mathcal{O}(v^2) | \mathcal{Q}\overline{\mathcal{Q}}[^{2S+1}(L\pm 1)_{J'}^{(8)}]gg \rangle + \dots \end{split}$$

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### 1 - Introduction Fragmentation Functions (FFs)

• Collins and Soper (1981):  $d\sigma[H+X] = \sum d\hat{\sigma}[parton + X] \otimes D_{i \to H}(z)$ 

i = parton

**FFs** (non-perturbative functions describing the formation of hadrons from partons)

### **1-Introduction Fragmentation Functions (FFs)**

• Collins and Soper (1981):  $d\sigma[H+X] = \sum d\hat{\sigma}[parton + X] \otimes D_{i \to H}(z)$ 

i = parton **FFs** (non-perturbative functions describing the formation of hadrons from partons)

What distinguishes quarkonium FFs is the existence of the subset of scales, which allows (applying NRQCD methodology) to factorise them



Thus, the FFs for quarkonium are calculable up to a set of LDMEs

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### 1 - Introduction Fragmentation Functions (FFs)

**Gluon FFs into various states at**  $\mu = 2m_c$ 

z = fraction of the fragmenting gluon's energy carried by the  $c\overline{c}$ 



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#### Potential discrimination power! \*





•  $J/\psi$  produced...

- In the hard interaction
  - If in an octet-state, it radiates soft gluons



- $J/\psi$  produced...
  - In the hard interaction
    - If in an octet-state, it radiates soft gluons
  - From the fragmentation of a gluon produced in the hard interaction
    - ► Dominates at high- $p_T$



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- PYTHIA8 implements a LO NRQCD-based prediction which includes both coloursinglet (cs) and colour-octet (co) mechanisms
  - If the  $Q\overline{Q}$  is produced in cs state, it is treated as a cs particle
  - ▶ If the  $Q\overline{Q}$  is produced in co state, it is treated as a single coloured particle that showers with 2 × splitting function of  $q \rightarrow qg$

 $c\bar{c}[{}^{3}S_{1}^{(8)}] \xrightarrow{\text{N recoils}} c\bar{c}(8) \xrightarrow{\text{N recoils}} c\bar{c}(8) \xrightarrow{\text{N recoils}} c\bar{c}(8) \xrightarrow{\text{N recoils}} J/\psi \longrightarrow \mu^{+}\mu^{-}$ 

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- Let's take a closer look into this modelization...
  - Simulation produced with Helac-Onia (hard process) + PYTHIA8 (shower)
  - ► Hard process:  $g + g \rightarrow c\overline{c}[^{2S+1}L_I^{(c)}] + g$  in pp collision at  $\sqrt{s} = 13$  TeV
  - Particles considered for the activity around the  $J/\psi$ : final charged particles with  $p > 2 \text{ GeV}, p_T > 200 \text{ MeV}$
  - More details on the simulation parameters in the backup

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In an event where the  $J/\psi$  was produced via singlet, the considered final state charged particles (p > 2 GeV,  $p_T > 200$ MeV) come from the underlying event

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cone of radius 0.5  $MC \text{ priv, } \Delta R = 0.5$  $\int 0^{-1} \int 0^{-1}$ 

Multiplicity inside

This provides already an indication of the limited activity surrounding the  $J/\psi$ , as simulated by PYTHIA8  $\implies$  In most of the events, the cone is empty

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- ► That was on PYTHIA8's side =>  $J/\psi$  from hard interaction
- Now, what if the  $J/\psi$  comes from  $c\overline{c}$  produced within a jet?

### 2 - Modelization of quarkonia production Probing Quarkonium production Mechanisms with jet substructure

Consider that the QQ pair is not produced directly in the hard scattering but is a fragmenting product of a high-p<sub>T</sub> jet

- Thus, a generic cross section for producing a  $J/\psi$  within a jet (energy  $E_{jet}$ , size R) can be factorised into
  - ▶ Hard and soft functions describing the production of the  $c\overline{c}[n]$  (+ other jets)

Fragmenting Jet Function (FJF) describing the non-perturbative fragmentation of the  $c\overline{c}[n]$  into the  $J/\psi$ 



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### 2 - Modelization of quarkonia production Probing Quarkonium production Mechanisms with jet substructure



E<sub>jet</sub> dependent distributions show a high discriminating power among contributions to quarkonium production

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# 3-Measurements (of $J/\psi$ in jets) Summary

Ref.	Collaboration, system, energy	Jet	J/ψ
<u>Phys. Rev. Lett. 118 (2017) 192001</u>	LHCb pp 13 TeV	R = 0.5, anti-kt pt > 20 GeV 2.5 < eta < 4.0	J/ψ : 2.5 < eta < 4.0
<u>Phys. Lett. B 825 (2022) 136842</u>	CMS pp (and PbPb) 5.02 TeV	R = 0.3, anti-kt 20 < pt < 40 GeV  eta  < 2	J/ψ : pt > 6.5 GeV
<u>Phys. Lett. B 804 (2020) 135409</u>	CMS pp 8 TeV	R = 0.5, anti-kt pt > 25 GeV  eta  < 1	J/ψ : E > 15 GeV;  y <1
<u>D. Bjergaard, PhD thesis, Duke U., 2017.</u> * Unpublished	ATLAS pp 8 TeV	R = 0.4, anti-kt pt > 45 GeV  eta  < 2.5	J/ψ : pt > 45 GeV;  y <2
<u>PoS (HardProbes2020), vol 387, 072</u> * Unpublished	STAR pp 500 GeV	R = 0.2 / 0.4 / 0.6, anti-kt pt > 10 GeV  eta  < 1 - R	J/ψ : pt > 5 GeV;  eta  < 1

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# 3 - Measurements Study of $J/\psi$ production in jets



Jet:

- Anti-kt, R = 0.5,
- ▶  $p_T(jet) > 20 \text{ GeV}$
- ►  $2.5 < \eta(jet) < 4.0$
- $J/\psi$  + charged +neutral (GhostProb<0.1, p < 1 TeV)
- Results:
  - Prompt–J/\u03c6 distributions show a much softer trend than Pythia8 predictions, which underestimate the activity around the J/\u03c6
- Limitations:
  - Cross-section not provided
  - 'J/ψ in jet' notion

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# **3-Measurements** Other measurements of promt $J/\psi$ in jets



- No significant dependence
  - More statistics
- Less isolated than **PYTHIA8** predictions
  - Point at z = 1 represent events where  $J/\psi = jet$



Similar to LHCb results • Peak at  $z \approx 0.5$ 

Peak at higher z than LHCb results

... more in the backup

0.8

0.6

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Jet z

### **3 - Measurements** Study of $J/\psi$ meson production inside jets in pp collisions at $\sqrt{s} = 8$ TeV

- Comparison of data with FJF predictions
- Observable: self-normalized E<sub>jet</sub>-dependence in bins of z

### Jet:

- Anti-kt
- ► R = 0.5
- ▶  $p_T > 25 \text{ GeV}$
- J/ψ in jet means that both muons are constituents of the jet

### Results

- Only FJF predictions using the **BCKL** (Phys. Rev. Lett. 113, 022001) LDME set, based on  $p_T(J/\psi) > 10$  GeV data, matches data in all bins
- ln  $(94.2 \pm 0.1)$ % of events with a jet, the  $J/\psi$  is inside the jet

#### 19.1 fb<sup>-1</sup> (8 TeV) 0.15 CMS 0.40 < z < 0.45 $\Xi(E_{jet};\ 0.425)$ in 8 GeV bins Phys. Rev. Lett. 113, 022001 Mod. Phys. Lett. A, Vol. 28, Data 0.14 Phys. Rev. Lett. 108, 242004 (2012)0.13 0.12 0.15 CMS 0.50 < z < 0.55 $\Xi(E_{jet}; 0.525)$ in 8 GeV bins BCKL FJF TOTAL Data 0.14 BK FJF TOTAL CHAO FJF TOTAL 0.13 0.12 0.11 0.15⊢ CMS 0.60 < z < 0.65 $\Xi(E_{jet}; 0.625)$ in 8 GeV bins BCKL FJF TOTAL Data BK FJF TOTAL 0.14 CHAO FJF TOTAL 0.13 0.12 0.11 60 70 80 90 100 120 110 E<sub>iet</sub> (GeV)

Phys. Lett. B 804 (2020) 135409

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## 4 - Conclusions and outlook

- The study of the hadronic activity around the  $J/\psi$  in our Helac-Onia+PYTHIA8 simulation reveals that the octet-shower is very similar to the underlying event and gives largely isolated  $J/\psi$
- Measurements of  $J/\psi$  in jets show that this quarkonium state is largely produced within jets
- Predictions based on quarkonium produced at later stages of the evolution within a jet show good agreement with data
- Outlook
  - Ongoing LHCb analysis for the fragmentation functions of  $J/\psi$ ,  $\psi(2S)$ ,  $\chi_{c1}(3872)$  and the three  $\Upsilon$ s in pp collisions at  $\sqrt{s} = 13$ TeV Eliane Epple, Naomi Cooke, Philip Ilten, J. Matthew Durham, Cesar da Silva
  - My ongoing analysis... ... for which input is welcome!

# Thank you for your attention!

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### **Isolation criteria**

#### Cone-based isolation **iii)** Modified versions Hollow cone **Discretised smooth cone** Hybrid cone designed to mimic detector outer cone mimics detector designed to mimic detector inner soft cone granularity granularity includes fragmentation includes fragmentation? removes fragmentation parameter dependence $\varepsilon_d E_T^{\gamma}$ \_\_\_\_\_ $\vec{E}_{T}^{had}$ $E_{\tau}^{\rm had}$ $-E_T^{\max}$ Enad $E_T^{\max}$

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### 1 - Introduction Fragmentation Functions (FFs)

**Gluon FFs into various states at**  $\mu = 2m_c$ 



### Fragmentation Functions (FFs)



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### Probing Quarkonium production Mechanisms with jet substructure

- Consider that the QQ pair is not produced directly in the hard scattering but is a fragmenting product of a high-p<sub>T</sub> jet
- The di-jet cross section for a jet (energy  $E_{jet}$ , size R) and a  $J/\psi$  with energy fraction z in a pp collision is schematically

$$\frac{d^2\sigma}{dE_{jet}dz} = \sum_{A,B,i,j} f_{A/p} f_{B/p} \cdot d\sigma(c\overline{c}[n]X, J_j^{recoil}) \otimes \mathcal{F}_S \otimes \mathcal{G}_i^{J/\psi}(E_{jet}, z \mid R, \mu)$$

Cross section for the production of a  $c\overline{c}[n]$  inside a jet + a recoil jet Fragmenting Jet Function (**FJF**) for the jet containing the  $J/\psi$ 

- FJF contains two scales:
  - Energy of the jet
  - Hadron mass

### Our simulation - main characteristics



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# NRQCD confronts LHCb Data on Quarkonium production within jets

#### Gluon Fragmentation Improved PYTHIA (GFIP):

Hard gluons produced in the short distance process with with virtuality of order  $E_{jet}$  are allowed to shower until a gluon with virtuality ~  $2m_c$  hadronizes into the  $J/\psi$ 



FIG. 1. PYTHIA predictions for c quark and gluon z distributions (where z is the fraction of the energy of the parton initiating the jet) after showering to the scale  $2m_c$ .

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# NRQCD confronts LHCb Data on Quarkonium production within jets



### **Gluon Fragmentation Improved PYTHIA (GFIP)**:

Hard gluons produced in the short distance process with virtuality of order  $E_{jet}$  are allowed to shower until a gluon with virtuality ~  $2m_c$  and convolution with LO-NRQCD FFs to obtain  $J/\psi$  distributions

### FJF

Combine FJFs with hard events and evolve FFs from  $2m_c$  to  $E_{jet}$ 

Better agreement than default PYTHIA

High- $p_T$  LDME sets better agreement

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Results for  $J/\psi$  from b-hadrons

### Study of $J/\psi$ production in jets

### Fiducial region

Table 3: Fiducial region employed to measure  $p_{\rm T}(J/\psi)/p_{\rm T}(j)$ .

object	requirement(s)
jets	$p_{\rm T}(j) > 20 { m GeV}, \ \ 2.5 < \eta(j) < 4$
$J\!/\!\psi$	$2 < \eta(J\!/\!\psi) < 4.5$
muons	$p_{\rm T}(\mu) > 0.5 \text{GeV}, \ p(\mu) > 5 \text{GeV}, \ 2 < \eta(\mu) < 4.5$

### Unfolding matrices



### Fragmentation of jets containing a prompt $J/\psi$ meson in PbPb and pp collisions at $\sqrt{s} = 5.02$ TeV

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 $J/\psi$  differential cross section in pp and the  $T_{AA}$  scaled yield in PbPb collisions

 $T_{AA}$  = average effective nucleon-nucleon luminosity delivered by a single heavy ion collision for a given centrality selection



### **3-Measurements** J/ $\psi$ production in jets in p+p collisions at $\sqrt{s} = 500$ GeV by STAR









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