

# $J/\psi$ electroproduction on light-nuclei at the EIC

Matteo Rinaldi  
INFN sezione di Perugia

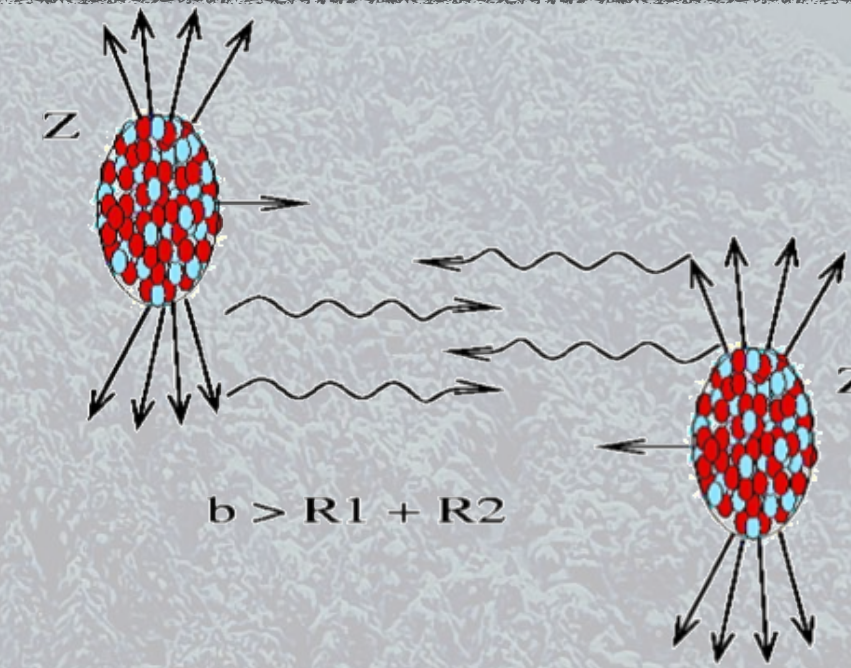


# J/ψ electroproduction on light-nuclei

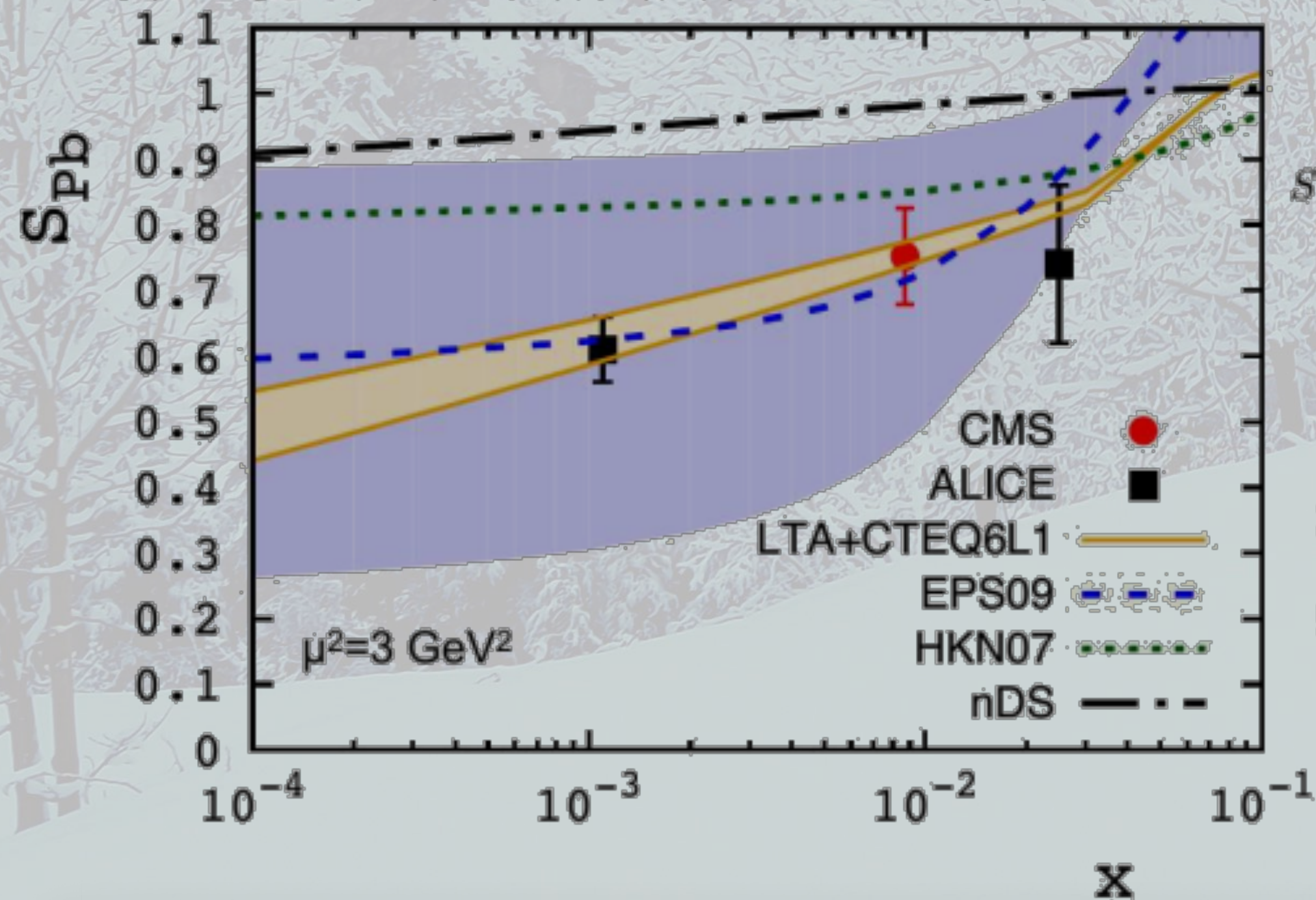
Gluon shadowing in UPC collisions @ LHC

Large (up to 40%) Leading twist (LT) shadowing in:  
 $\gamma + \text{Pb/Au} \rightarrow \rho(J/\psi) + \text{Pb/Au}$       Explained/predicted

(Frankfurt, Guzey, Strikman *Phys. Rep.* 512 (2012) 255)



Abbas *et al.* [ALICE] *EPJ C* 73 (2013) 2617; CMS Collab. *PLB* 772 (2017) 489 → suppression factor  $S_{\text{Pb}}$



$$S(W_{\text{TP}}) = \left[ \frac{\sigma_{\gamma \text{Pb} \rightarrow \text{J}/\psi \text{Pb}}}{\sigma_{\gamma \text{Pb} \rightarrow \text{J}/\psi \text{Pb}}^{\text{IA}}} \right]^{1/2} = \kappa_{A/N} \frac{G_A(x, \mu^2)}{A G_N(x, \mu^2)}$$

LTA: Guzey, Zhavoronkov *JHEP* 1310 (2013) 207

EPS09: Eskola, Paukkunen, Salgado, *JHEP* 0904 (2009) 065

HKN07: Hirai, Kumano, Nagai, *PRC* 76 (2007) 065207

nDS: de Florian, Sassot, *PRD* 69 (2004) 074028

**Introduction.** Studies of nuclear shadowing have a long history [1–5]. In quantum mechanics and in the eikonal limit, it is manifested in the total hadron-nucleus cross section being smaller than the sum of individual hadron-nucleon cross sections. In essence, this is due to simultaneous interactions of the projectile with  $k \geq 2$  nucleons of the nuclear target, leading to a reduction (shadowing) of the total cross section. In this frame-

# Learning from light nuclei - I

- Problem:  
@ EIC/LHC it is challenging to measure coherent scattering at  $t \neq 0$  for  $A \approx 200$ ; Large coherence length: information on interactions with many nucleons, in average
- Solution:  
use the lightest nuclei, especially  $^3\text{He}$  and  $^4\text{He}$ , to study coherent effects for interactions with exactly 2 nucleons in the range of  $0 < -t < 0.5 \text{ GeV}^2$ .

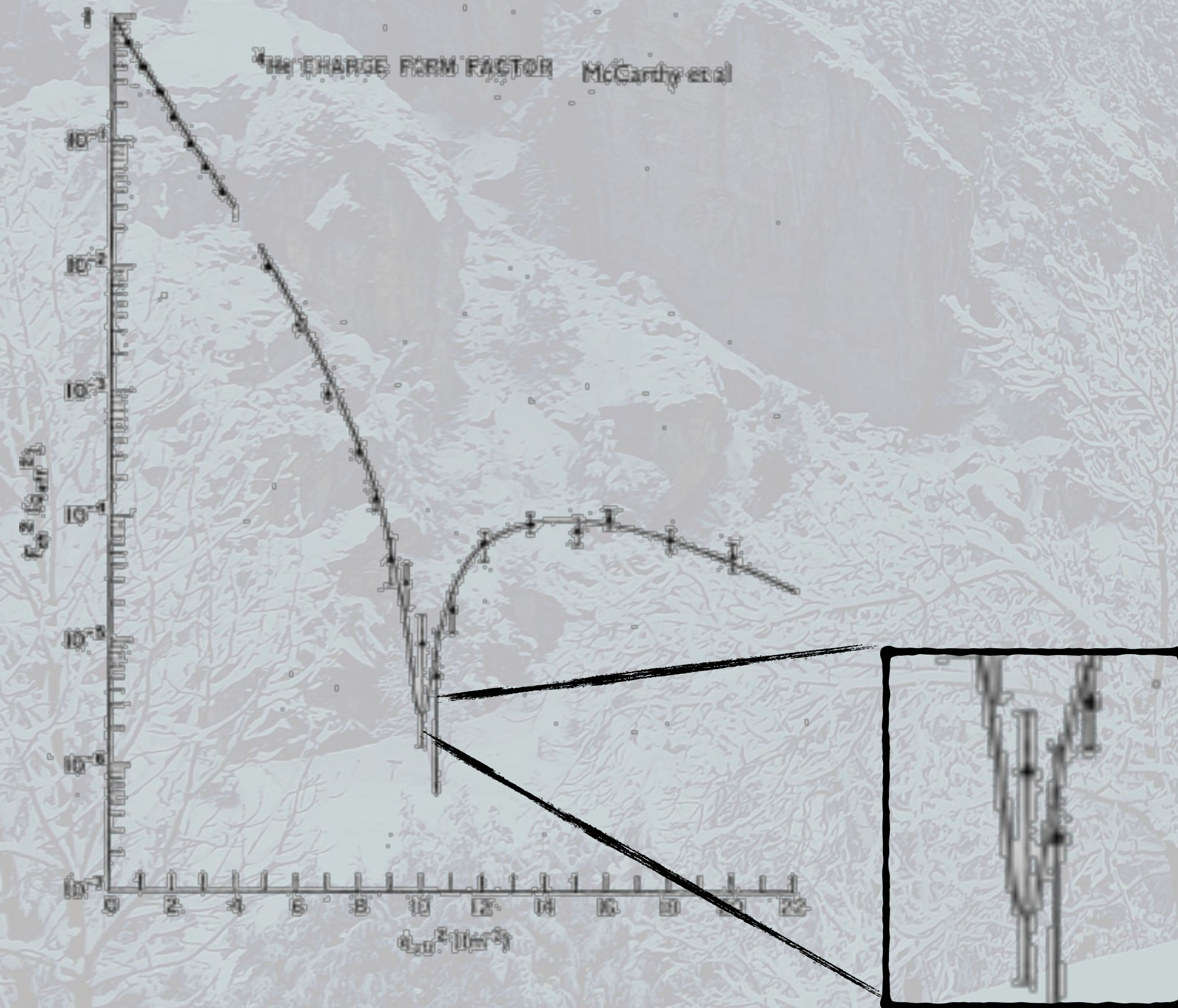
## Complementary measurements with light ion beams @ the EIC:

- Scattering off 2 and 3 nucleons can be separately probed
- no excited states -> easy to select coherent events

Here:

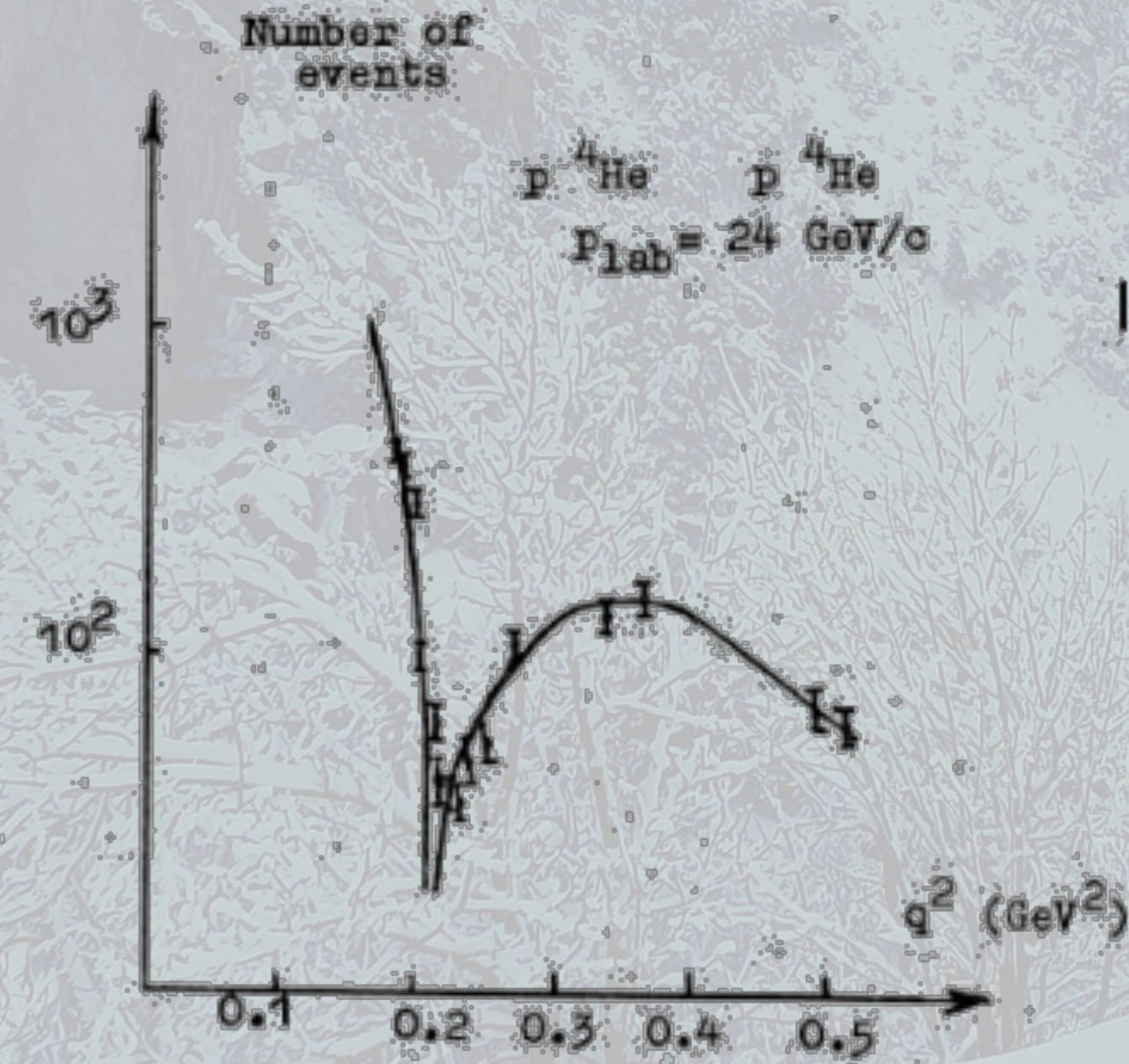
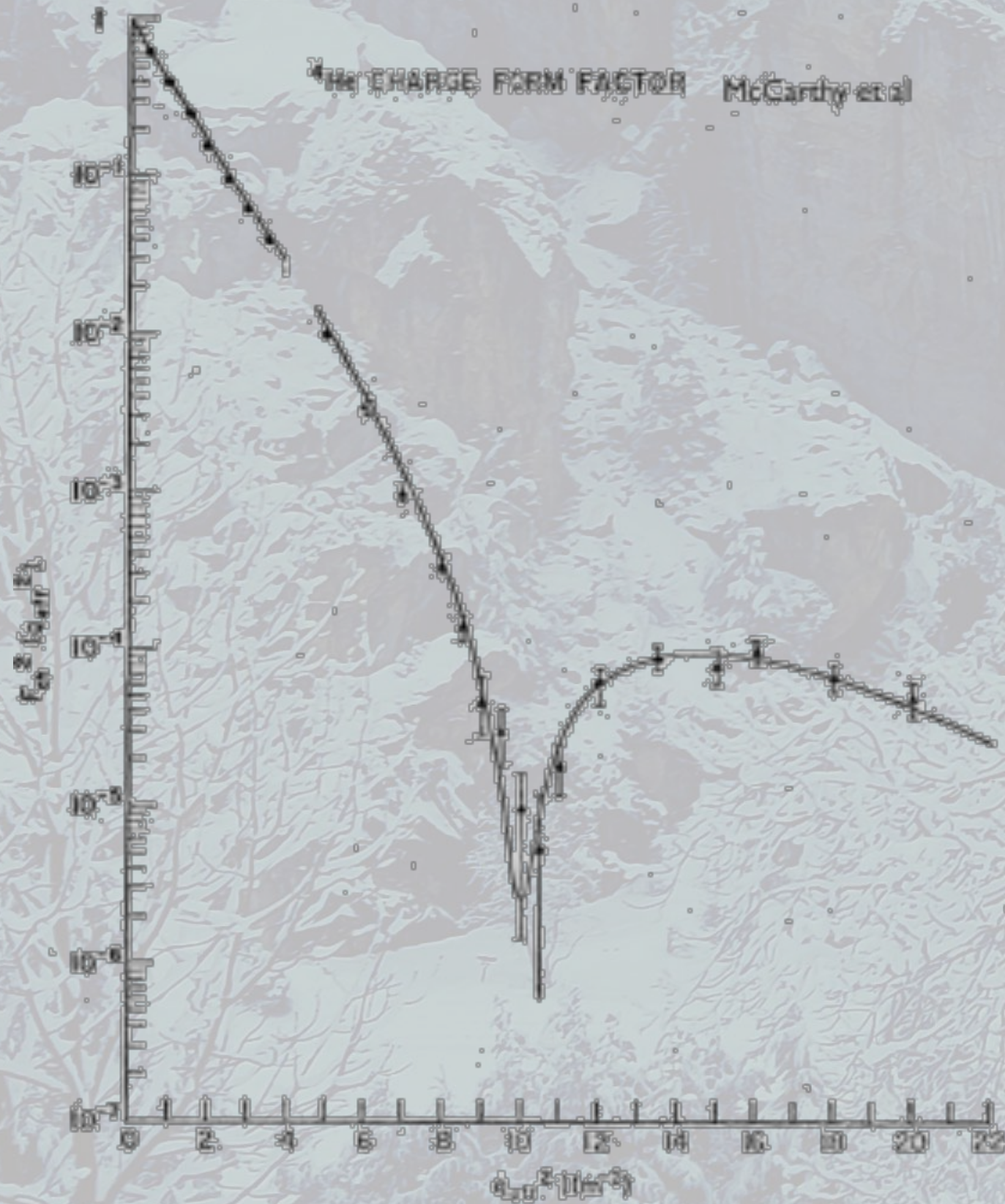
Results on  $J/\Psi$  diffractive electro-production off  $^3\text{He} - ^4\text{He}$   
V. Guzey, M. R., S. Scopetta, M. Strikman and M. Viviani, PRL 129 (2022) 24, 24503

# An old idea (Levin and Strikman 1975)



- 4He charge FF, dominated by one-body dynamics (IA) presents the first diffraction minimum at:  
 $-t \approx 0.4 \text{ GeV}^2$

# An old idea (Levin and Strikman 1975)



⚙️ <sup>4</sup>He charge FF, dominated by one-body dynamics (IA) presents the first diffraction minimum at:  
 $-t \approx 0.4 \text{ GeV}^2$

⚙️ around this value of  $t$ , the cross section in  $p + ^4\text{He} \rightarrow p + ^4\text{He}$  is dominated by effects beyond IA:  
multinucleon interactions, gluon shadowing for hard processes

**no contribution of impulse approximation term**

# The X-section for $J/\psi$ exclusive production @ EIC

LT parton shadowing for  $J/\psi$  coherent production off He (gluon GPDs in He)

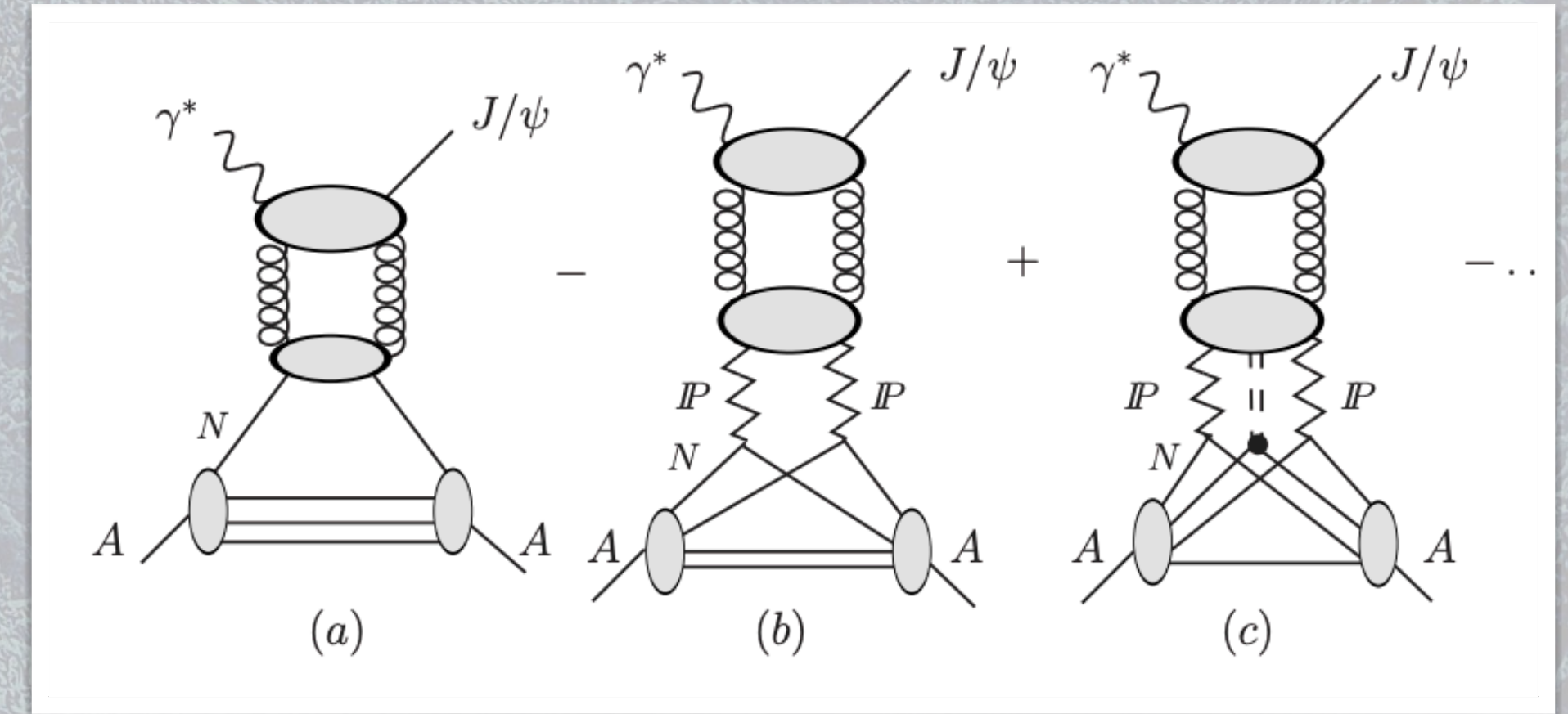
(Frankfurt, Guzey, Strikman Phys. Rep. 512 (2012) 255)

$$\frac{d\sigma_{\gamma^* A \rightarrow VA}}{dt} = \frac{d\sigma_{\gamma^* N \rightarrow VN}}{dt}(t=0) \left| F_1(t) e^{(B_0/2)t} + \sum_{k=2}^4 F_k(t) \right|^2$$

$$F_k(q) = \left( \frac{i}{8\pi^2} \right)^{k-1} C_n^k A_k \int \prod_{l=1}^k d^2 q_l f(q_l) \Phi_k(q, q_l) \delta \left( \sum_l q_l - q \right) \quad k = 2, 3, 4$$

$$F_1(q) = 4\Phi_1(q) \quad f(q_l) = \text{scattering amplitude for } J/\psi N \rightarrow J/\psi N$$

$$A_{k>1} = \frac{\langle \sigma^k \rangle}{\langle \sigma \rangle} \frac{(1 - i\eta)^k}{1 - i\eta_0}; \text{ the same used in UPC studies!}$$



Parameters:

- $B_0$
- $\eta (\eta_0) = \text{Re}(f) / \text{Im}(f)$  for  $\gamma p \rightarrow J/\psi p$  ( $J/\psi p \rightarrow J/\psi p$ )

- moments  $\langle \sigma^i \rangle$  chosen for the specific final state and the specific kinematics

(Guzey et al. PRC 93 (2016) 055206).

The model has been tested in  $J/\psi$  photoproduction in Pb-Pb UPCs at the LHC (v. Guzey and M. Zhalov, JHEP 10, 207 (2013))

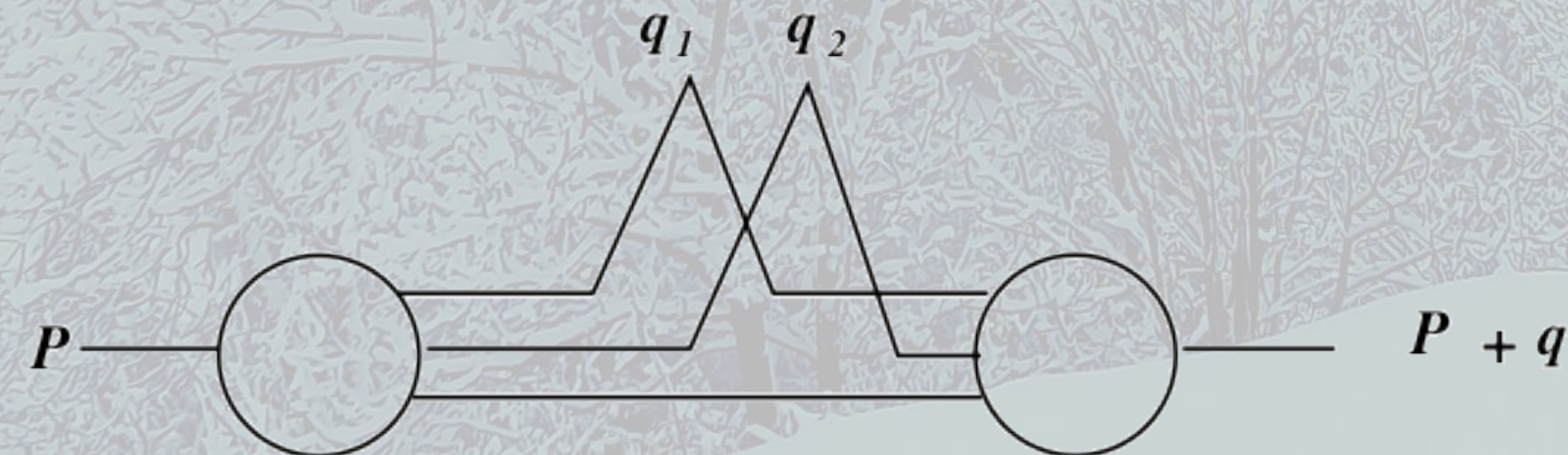
- $\Phi_k$  "k-body form factor", is the nuclear input

# Nuclear Physics input

$$\Phi_k(\vec{q}_1, \dots, \vec{q}_k) = \int \prod_{i=1}^N \left\{ \frac{d\vec{p}_i}{(2\pi)^3} \right\} \psi_{P'}^*(\vec{p}_1 + \vec{q}_1, \dots, \vec{p}_k + \vec{q}_k, \dots, \vec{p}_N) \psi_P(\vec{p}_1, \dots, \vec{p}_k, \dots, \vec{p}_N) \delta\left(\sum_{i=1}^N \vec{p}_i\right)$$

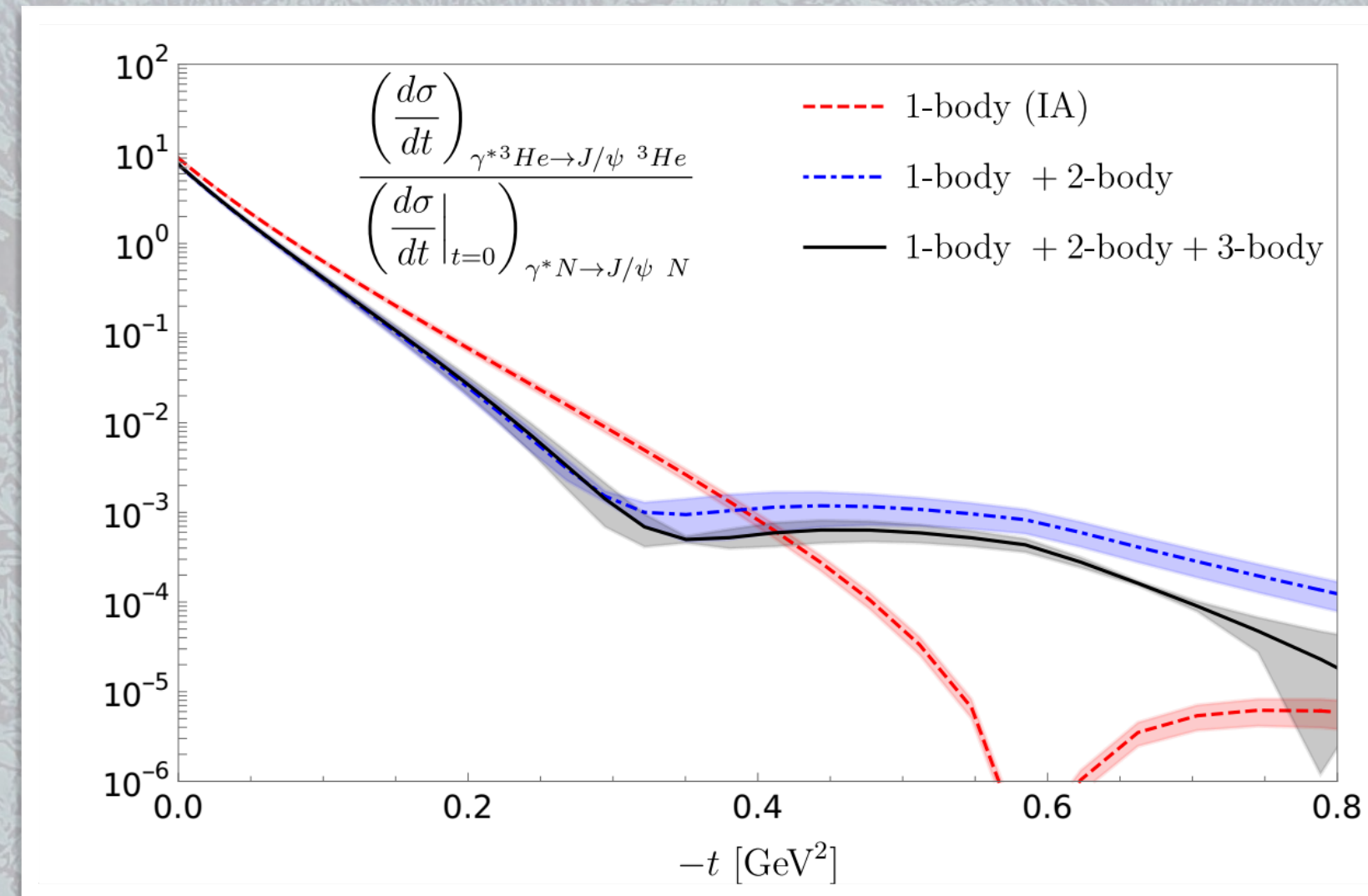
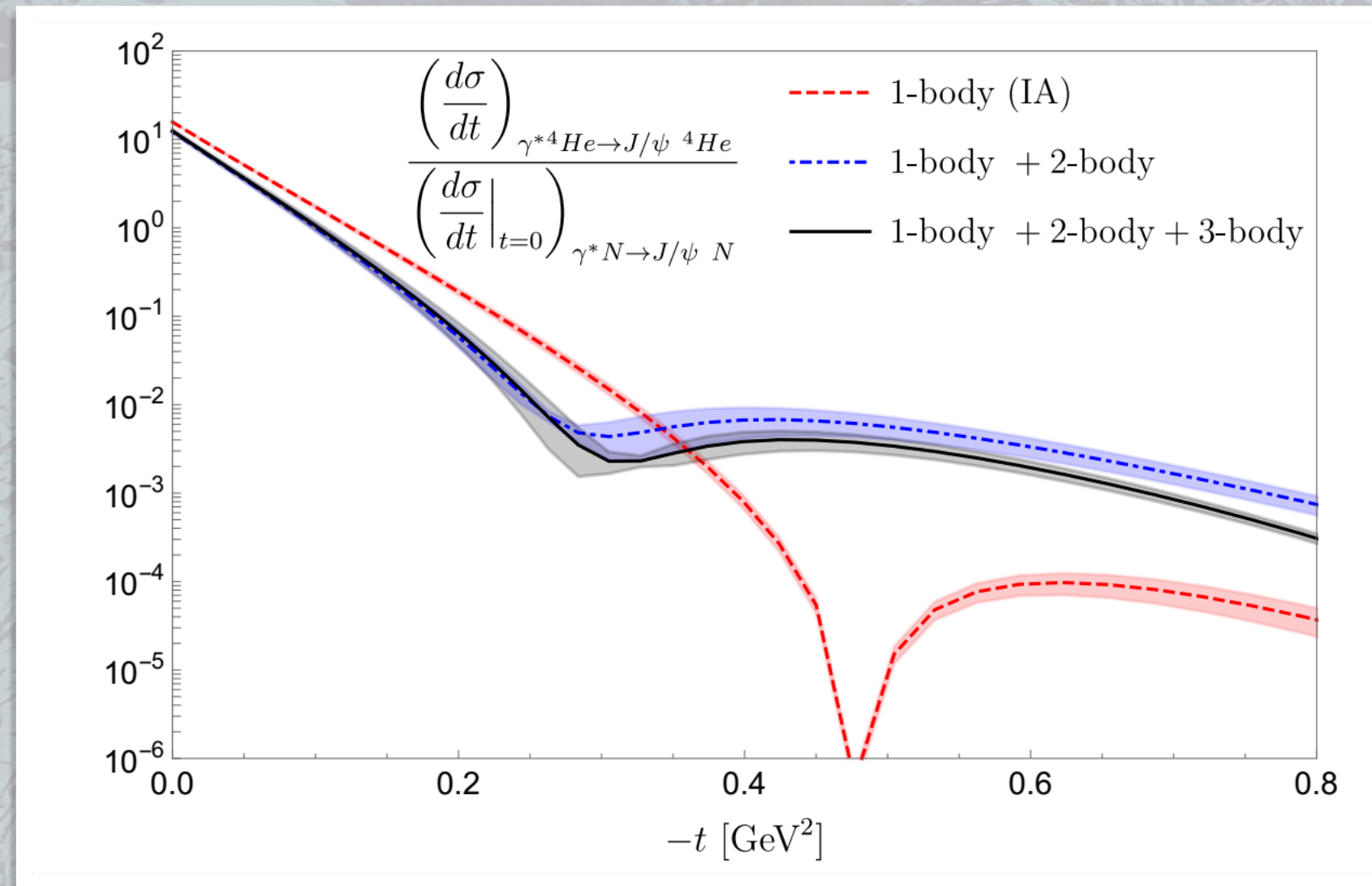
- #  $\Phi_1$  (IA, very important here),  $\Phi_2$  and  $\Phi_3$  evaluated using the realistic w. f. obtained by the Pisa group using:  
a) Av18 for  ${}^3\text{He}$  b) the N4LO chiral potential (D. R. Entem, R. Machleidt, Y. Nosyk, Phys. Rev. C 96, 024004 (2017)) for  ${}^4\text{He}$

- # Example of  $\Phi_2$  :



- # we remark that  $\Phi_2(k_\perp, -k_\perp)$  is the same quantity appearing in the DPS2 mechanism

# Results for J/Ψ exclusive production @EIC: $x_B \approx 10^{-3}$



Error bars account:

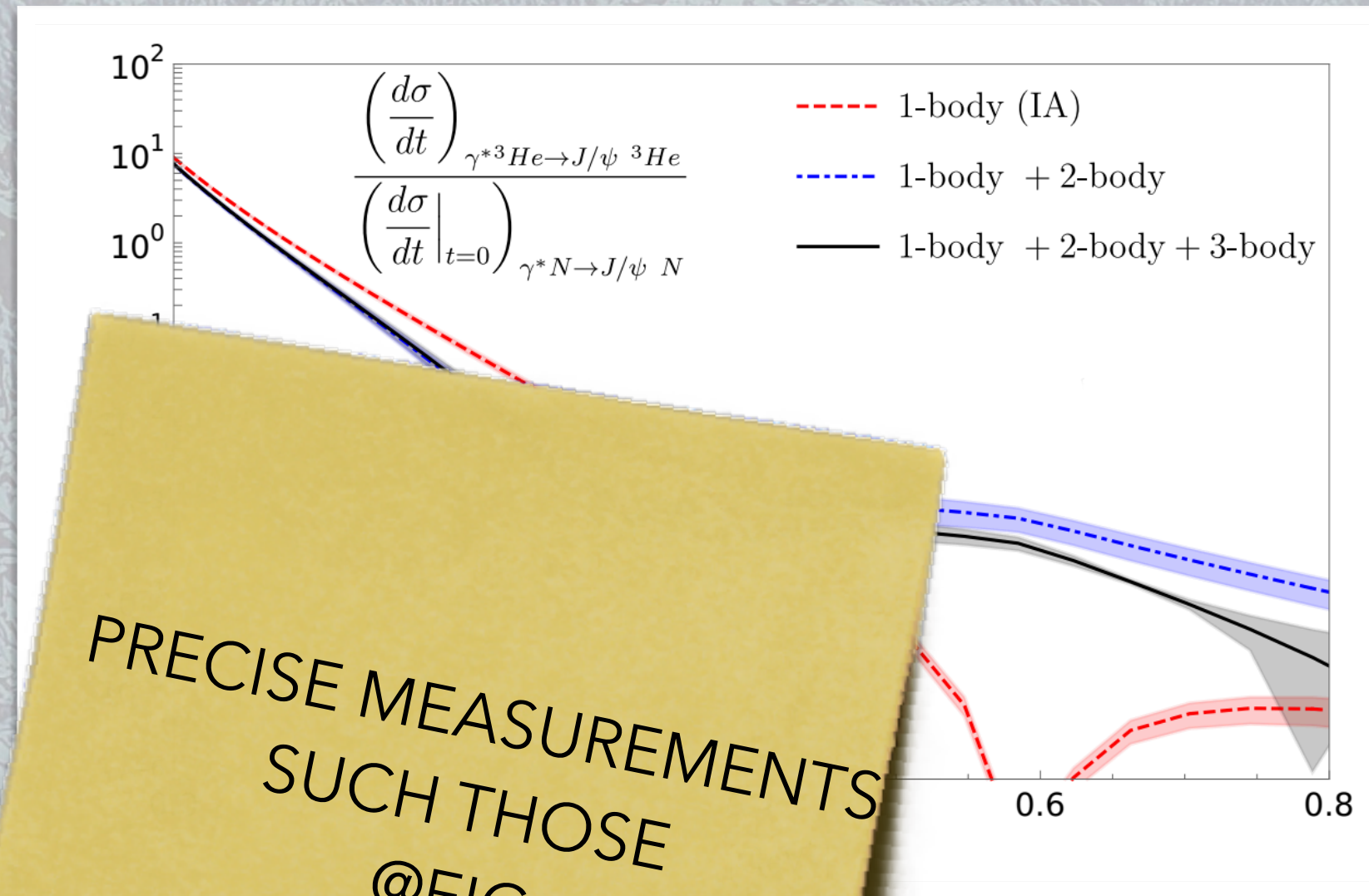
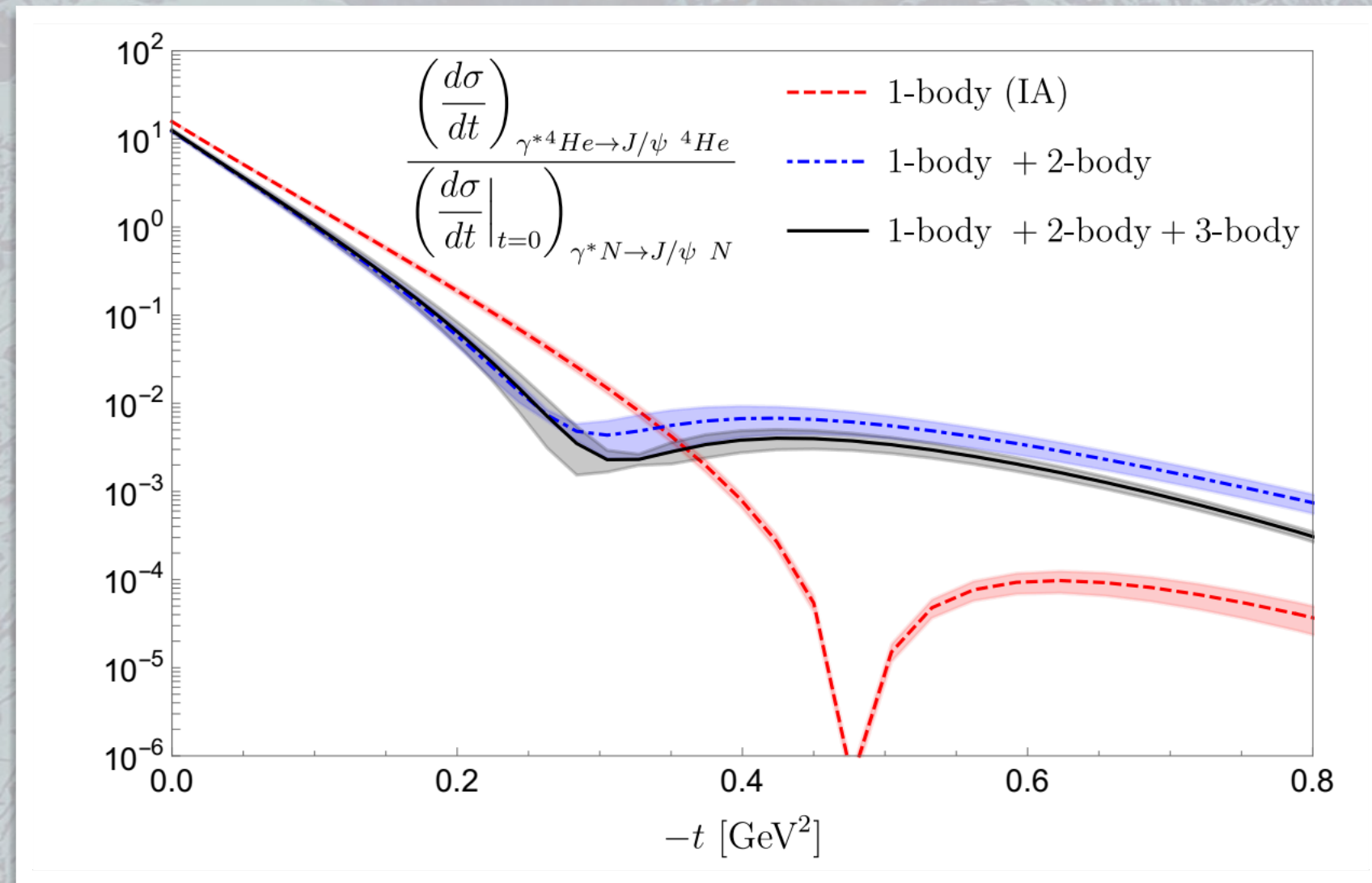
-10% of variation for  $B_0$

-15 of variation in  $\langle \sigma^2 \rangle$

- ✓ 1-body + 2-body re-scatterings dominate the cross-sections shift of the minimum due to 2-body dynamics
- ✓ 1-body dynamics under theoretical control: very good chances to disentangle
- ✓ 2-body dynamics (LT gluon shadowing)
- ✓ unique opportunity to access the real part of the scattering amplitudes in a wide range of  $t$
- ✓ The position of the minimum is extremely sensitive to dynamics and the structure!



# Results for J/ψ exclusive production @EIC: $x_B \approx 10^{-3}$



Error bars account:

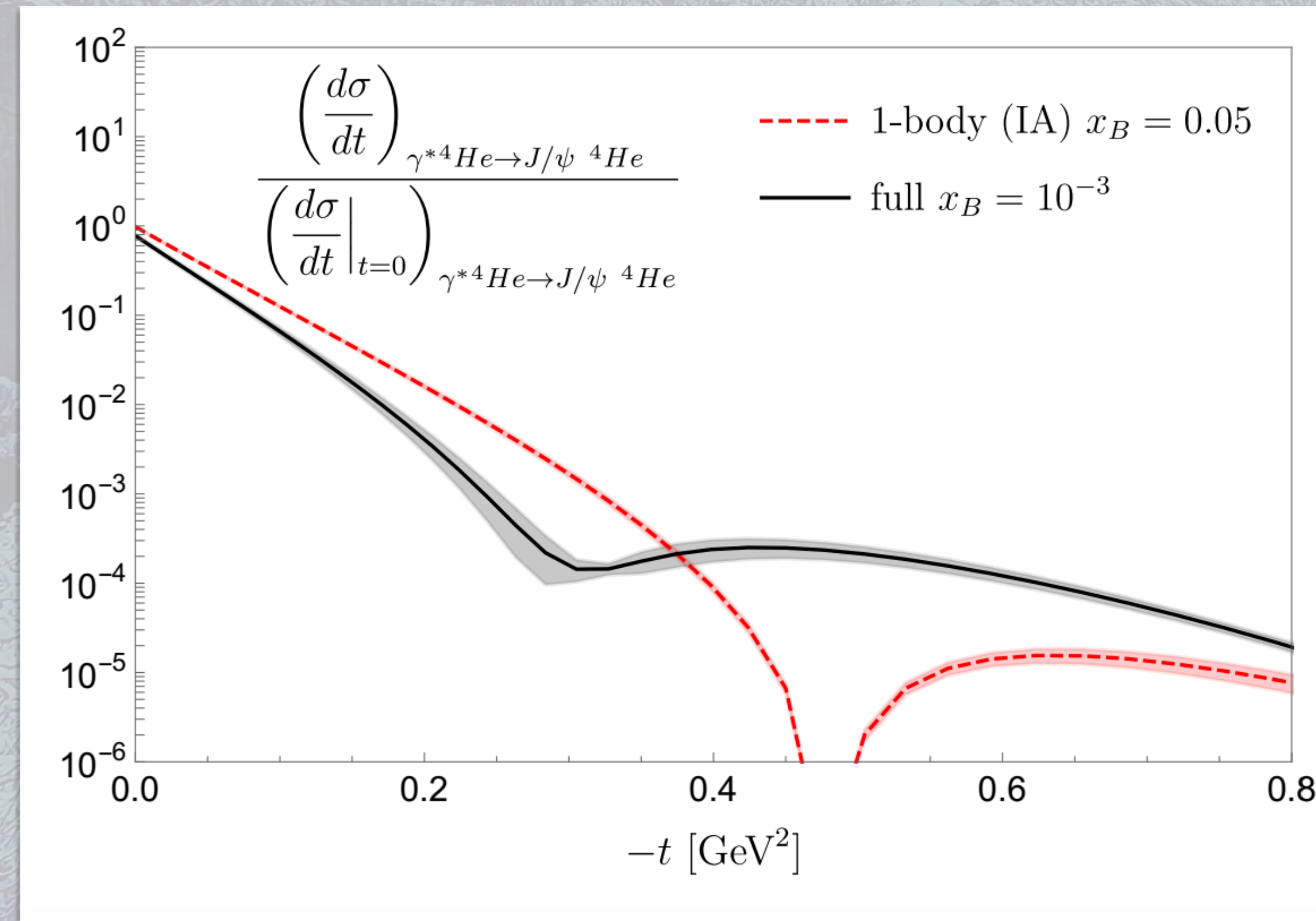
-10% of variation for  $B_0$

-15 of variation in  $\langle \sigma^2 \rangle$

PRECISE MEASUREMENTS  
SUCH THOSE  
@EIC  
ARE NECESSARY!

- ✓ 1-body + 2-body re-scatterings dominate the cross-section
- ✓ 1-body dynamics under theoretical control: very good chances to
- ✓ 2-body dynamics (LT gluon shadowing)
- ✓ unique opportunity to access the real part of the scattering amplitudes in a wide range of  $t$
- ✓ The position of the minimum is extremely sensitive to dynamics and the structure!

# Results for $J/\psi$ exclusive production @EIC: $x_B \approx 0.05$ ( $x_B$ -evolution)



Error bars account:

-10% of variation for  $B_0$

-15 of variation in  $\langle \sigma^2 \rangle$

- ✓  $x_B$ -evolution of the  $t$ -dependence predicted in Frankfurt, Strikman, Weiss PRD 83 (2011) 054012 , considering HERA data: possible check of the model at the EIC
- ✓ 1-body dominates the cross-section at  $x_B \approx 0.05$ ; no shadowing at  $t = 0$
- ✓ possible interpretation: at low  $x_B$  , significant broadening in the impact parameter space of the nuclear gluon distribution (see Guzey, Strikman, Zhalov PRC 95 (2017) 2,025204 for heavy nuclei)

# Backup

Good-Walker formalism of eigenstates of the scattering operator which allows one to characterize the interaction with  $k$  nucleons by the  $k$ th moment  $\langle \sigma^k \rangle$

$$\langle \sigma^k \rangle = \int d\sigma P(\sigma) \sigma^k$$

parametrizes the hadronic structure of the virtual photon and gives the probability for the photon to interact with a nucleon with the cross section  $\sigma$

**M. L. Good and W. D. Walker, Phys. Rev. 120, 1857- 1860 (1960)**  
**B. Blaettel, et al, Phys. Rev. D 47, 2761-2772 (1993)**

virtual interaction with individual nucleons, giving rise to the Glauber model [6, 7]. However, it was demonstrated by Mandelstam [8] and Gribov [9] that the contribution of eikonal diagrams in quantum field theory models tends to zero at high energies because, qualitatively, there is not enough time between interactions with two nucleons for the projectile to transform back into itself. As a result, shadowing in the high energy limit is determined by the totality of diffractive interactions of the projectile in different configurations [9].