

Why do we need independent observables to improve our understanding of the QCD matter properties, how?

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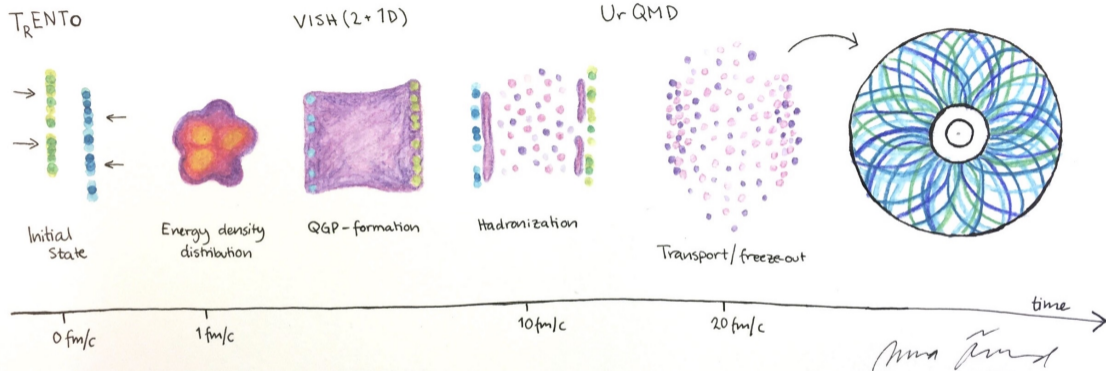
2. CERN, Switzerland

Thursday 12th October, 2023

Particle Physics Day, Jyväskylä, Finland



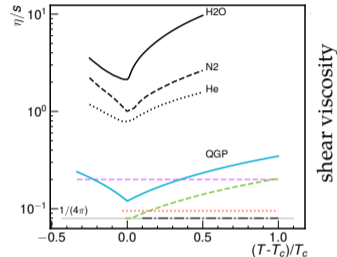
THE DIFFERENT STAGES OF HEAVY-ION COLLISIONS



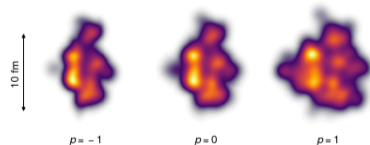
$$T^{\mu\nu} = eu^{\mu}u^{\nu} - (P + \Pi)\Delta_{\mu\nu} + \pi^{\mu\nu}, \quad \delta_{\mu}T^{\mu\nu} = 0$$

COLLECTION OF PARAMETERS

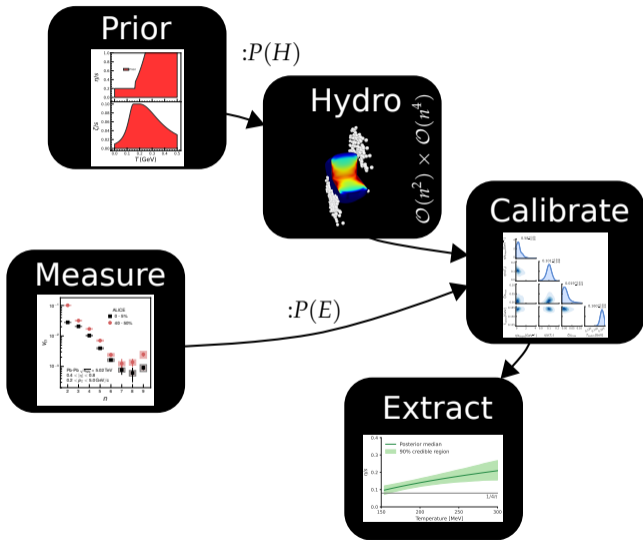
Parameter	Description
T_c	Temperature of const. $\eta/s(T)$, $T < T_c$
$\eta/s(T_c)$	Minimum $\eta/s(T)$
$(\eta/s)_{\text{slope}}$	Slope of $\eta/s(T)$ above T_c
$(\eta/s)_{\text{curve}}$	Curvature of $\eta/s(T)$ above T_c
$(\zeta/s)_{\text{peak}}$	Temperature of $\zeta/s(T)$ maximum
$(\zeta/s)_{\text{max}}$	Maximum $\zeta/s(T)$
$(\zeta/s)_{\text{width}}$	Width of $\zeta/s(T)$ peak
T_{switch}	Switching / particlization temperature
N(2.76 TeV)	Overall normalization (2.76 TeV)
N(5.02 TeV)	Overall normalization (5.02 TeV)
p	Entropy deposition parameter
w	Nucleon width
σ_k	Std. dev. of nucleon multiplicity fluctuations
d_{min}^3	Minimum volume per nucleon
τ_{fs}	Free-streaming time



Trento p-value, <http://qcd.phy.duke.edu/trento/>



BAYESIAN PARAMETER ESTIMATION



Bayes' theorem:

$$P(H|E) = \frac{P(E|H) \cdot P(H)}{P(E)}$$

$$P(E) = \sum_{i=1}^n P(E|H_i)P(H_i)$$

- Find optimal set of model parameters that best reproduce the experimental data.
- Utilize constraints, such as flow observables, to help narrow down the $\eta/s(T)$ and such.

Testing a single set of parameters requires $\mathcal{O}(10^4)$ hydro events, and evaluating eight different parameters five times each requires $5^8 \times 10^4 \approx 10^9$ hydro events.

That's roughly 10^5 CPU years!

OUR ARSENAL OF OBSERVABLES - STOCHASTIC APPROACH

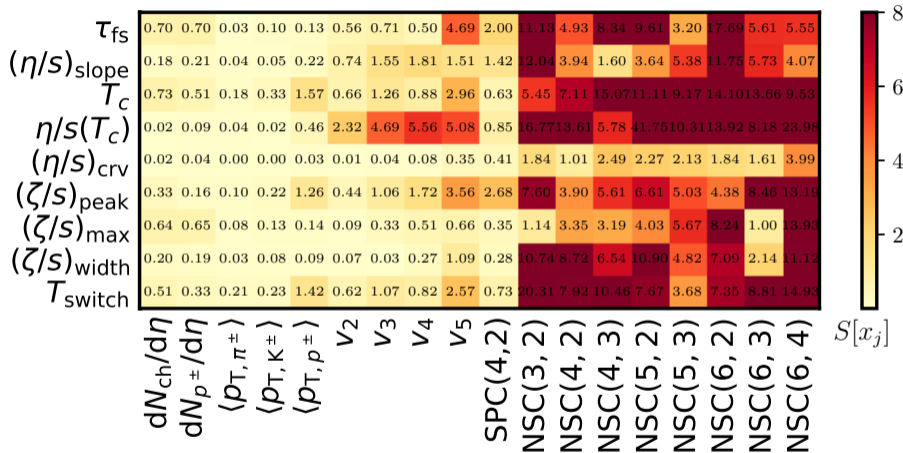
- Together, various flow observables cover the sensitivity for all components of transport properties.

Name	Symbol	Measure	Sensitivity-stochastic approach
Flow coefficients	v_n	System expansion and anisotropy of the flow	Average $\langle \eta/s \rangle$ and $\zeta/s(T)$ peak
(Normalized) Symmetric cumulants	(N)SC(k, l, m)	Correlation between magnitudes of flow harmonics	$\eta/s(T)$ temperature dependence
Non-linear flow mode coefficients	$\chi_{n,mk}$	Quantification of the non-linear response	$\eta/s(T)$ at the freeze-out
Symmetry plane correlations	$\rho_{n,mk}$	Correlations between the directions of flow harmonics	$\eta/s(T)$

Thanks to excellent ALICE papers over years:

- Phys.Rev.Lett. 117 (2016) 182301, Phys.Lett. B773 (2017) 68, Phys.Rev. C 97 (2018) 024906, JHEP05 (2020) 085, Phys.Lett. B818 (2021) 136354, Phys.Rev.Lett. 127 (2021) 092302 - [flow](#)
- Phys.Rev.Lett. 106 (2011) 032301, Phys.Rev.C 88 (2013) 044910, Phys.Lett. B772 (2017) 567-577, Phys.Rev.C 101, 044907 (2020) - [N_{ch} and \$\langle p_T \rangle\$](#)

OBSERVABLE SENSITIVITIES



Sensitivity

$S[x_j] = \Delta/\delta.$, where

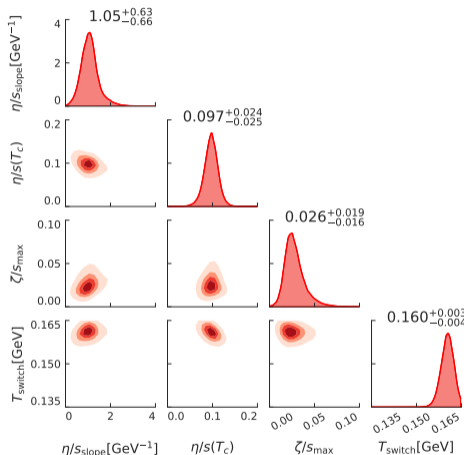
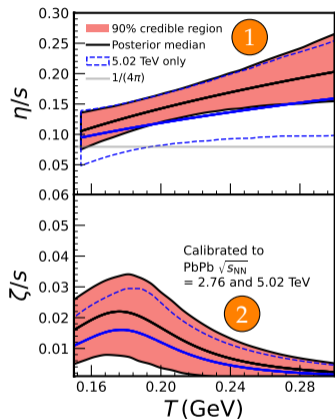
$$\Delta = \frac{|\hat{O}(\vec{x}') - \hat{O}(\vec{x})|}{|\hat{O}(\vec{x})|}$$

- NSCs most sensitive to multiple different parameters
- v_n s show sensitivity to specific shear viscosity

$$(\eta/s)(T) = (\eta/s)(T_c) + (\eta/s)_{slope}(T - T_c) \left(\frac{T}{T_c}\right)^{(\eta/s)_{curve}}$$

$$(\zeta/s)(T) = \frac{(\zeta/s)_{max}}{1 + \left(\frac{T - (\zeta/s)_{T_{peak}}}{(\zeta/s)_{width}}\right)^2}$$

RESULTS: JYVÄSKYLÄ (2022) – COMBINED COLLISION ENERGY ANALYSIS (2.76 + 5.02 TeV)



- 1 Significantly improved $\eta/s(T)$ uncertainty
- 2 Non-zero $\zeta/s(T)$
- 3 Overall better convergence for parameter components

Together with two collision energies and added observables, the uncertainty has reduced!

How can we further reduce the uncertainties?

IMPROVING THE UNCERTAINTIES

Add more data



IMPROVING THE UNCERTAINTIES

Add more data



IMPROVING THE UNCERTAINTIES

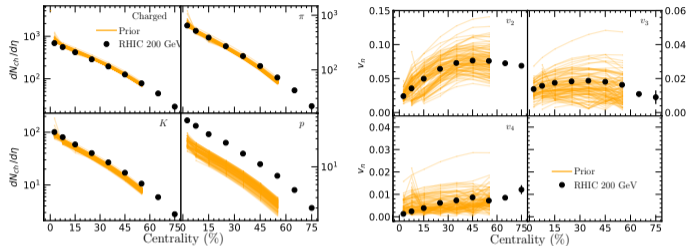
Add more data



Add independent data



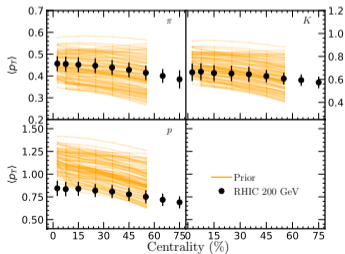
INCLUDING RHIC AuAu 200 GeV



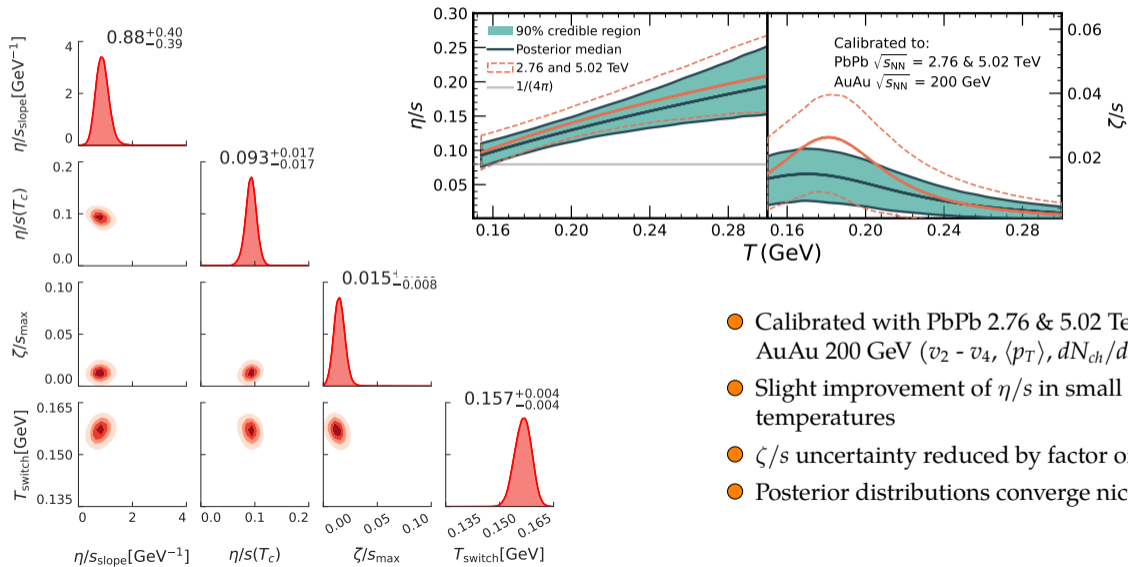
Binnings

Separate centrality binnings for all parametrisations

- Included observables for RHIC $\sqrt{s_{NN}} = 200$ GeV data: $v_2 - v_4$, $\langle p_T \rangle$ and N_{ch} for charged particles and PID
- Priors underestimate proton $dN_{ch}/d\eta$ and overestimate proton p_T
- v_n s are covered well



PRELIMINARY RESULTS



NEW OBSERVABLES

Are there any new observables that could be included?

SYMMETRY PLANE CORRELATIONS

- Previous scalar product method is biased, due to $\langle v_n^2 v_m^2 \rangle \neq \langle v_n^2 \rangle \langle v_m^2 \rangle$

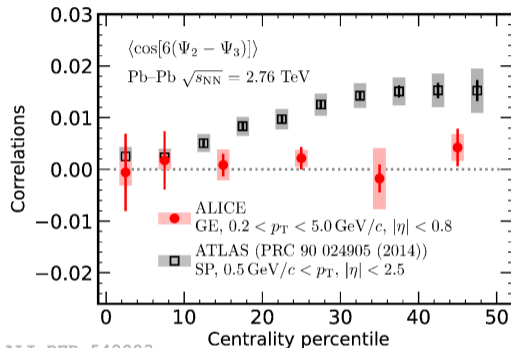
ALICE Collaboration, *Phys. Rev. C* 97 (2018) 024906

- Newly developed Gaussian Estimator (GE)

A. Bilandzic, M. Lesch, and S. F. Taghavi, *Phys. Rev. C* 102 (2020) 024910

$$\langle \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle_{GE}$$

$$\approx \sqrt{\frac{\pi}{4}} \frac{\langle v_{n_1}^{a_1} \dots v_{n_k}^{a_k} \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle}{\sqrt{\langle v_{n_1}^{2a_1} \dots v_{n_k}^{2a_k} \rangle}}$$



ALI-DER-542003

- Gaussian estimator (GE) ALICE analysis at 2.76 TeV, *Eur.Phys.J.C* 83 (2023) 7 [M. Lesch, A. Bilandzic, D.J. Kim, S.F. Taghavi]
- Scalar product (SP) method (ATLAS: PRC 90, 024905 (2014), ALICE, PLB 773: 68 (2017) [J.E. Parkkila, D.J. Kim])

ASYMMETRIC CUMULANTS

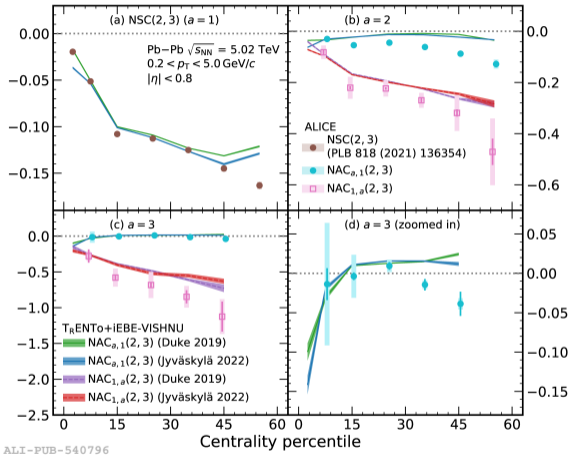
Normalised Asymmetric Cumulants

$$\text{NAC}_{a,1}(m, n) = \frac{\langle (v_m^2)^a (v_n^2)^c \rangle}{\langle v_m^2 \rangle^a \langle v_n^2 \rangle^c}, \quad a = 1 - 4$$

- First measurements in Pb-Pb $\sqrt{s_{NN}} = 5.02$ TeV

ALICE Collaboration, arXiv:2303.13414 (recommended for publication in PRC)

- Model predictions disagree with the data for $\text{NAC}_{2,1}(2,3)$ and $\text{NAC}_{3,1}(2,3)$
- Will help constrain the initial condition modeling
- Duke 2019 parametrization Nature Phys. 15, 1113-1117 (2019) [J.E. Bernhard *et al.*]
- Jyväskylä 2022 parametrization Phys. Lett. B 835 (2022) 137485 [J.E. Parkkila *et al.*]



OUTLOOK

Experiments

- RHIC data (AuAu collisions) - Energy and system size dependence (ongoing)
- Use new observables in BA
 - Higher order ($n > 5$) Symmetric cumulants (Anna)
 - Improved Symmetric Plane Correlation (SPC) : independent from flow magnitude correlations (Maxim)
 - Asymmetric Cumulants (AC) (Cindy)

Theory

- Improving the initial conditions with
 - EKRT
 - IP+Glasma
- Testing hydro limit with small systems?
- Study the substructure

SUMMARY

Summary

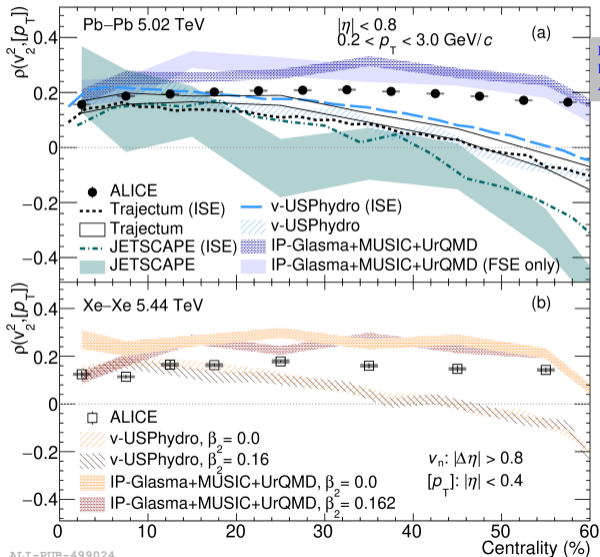
- RHIC data have been added with separate centralities.
- Need to do the centrality calibration to PbPb as well.
- More observables will be added in future.

Thank you for your attention!

Acknowledgments:

- CSC for providing the ~24 million CPU hours
- Harri Niemi, Kari Eskola, Jonah E. Bernhard, J. Scott Moreland and Steffen A. Bass for their useful comments

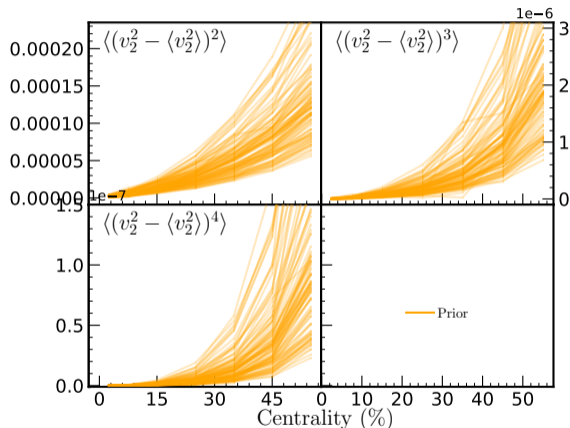
$v_n, [p_T]$ CORRELATION - SHORTAGE OF T_{RENT0} MODEL



P. Bozek, R. Samanta, Phys. Rev. C 102, 034905
 B. Schenke, C. Shen, D. Teaney, Phys. Rev. C 102, 034905
 ALICE, arXiv:2111.06106

$$\rho(v_2^2, [p_T]) = \frac{\langle \delta v_2^2 \delta [p_T] \rangle}{\sqrt{\langle (\delta v_2^2)^2 \rangle \langle (\delta [p_T])^2 \rangle}} \quad (1)$$

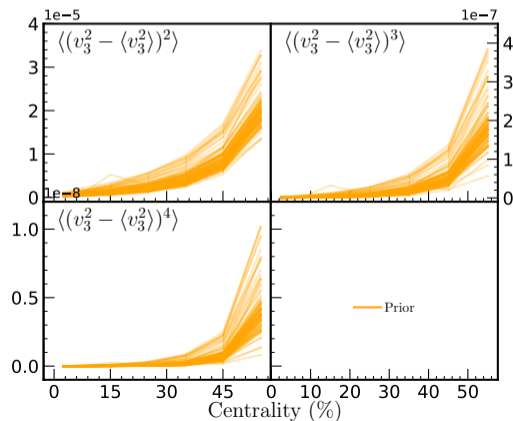
- Correlation between $[p_T]$ and v_2 :
 - can be used to differentiate initial state models
 - More peripheral \rightarrow best described by models with IP-Glasma
 - strong centrality dependence on the models with Trento
- Ongoing progress:
 - Calculate sensitivity
 - Adapt it to the Bayesian Analysis

MOMENTS OF δv_n 

- Characterizes the fluctuation of different order
- $|\eta| < 0.8$ and $0.2 \text{ GeV} < p_T < 5.0 \text{ GeV}$

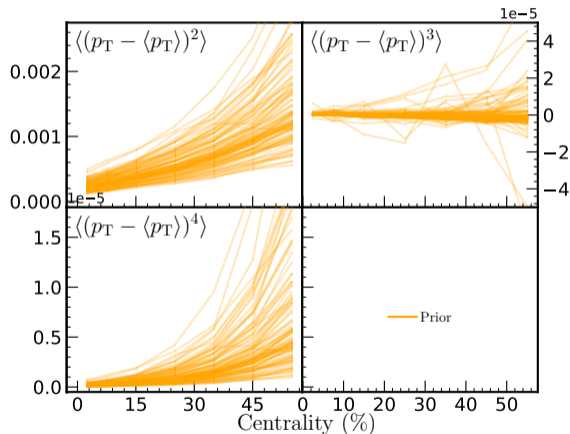
Moments of a distribution

Variance: $(X - \mu)^2$, **skewness:** $(X - \mu)^3$
and kurtosis: $(X - \mu)^4$



PRIOR DISTRIBUTION OF $\rho(v_n^2, [p_T])$

Moments of $[p_T]$ distribution

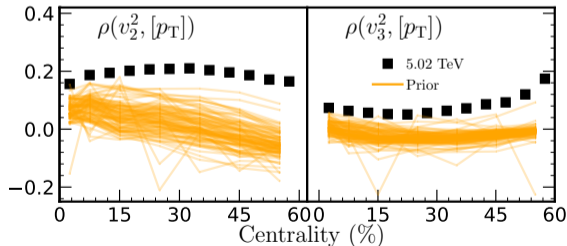


■ $|\eta| < 0.8$ and $0.2 \text{ GeV} < p_T < 5.0 \text{ GeV}$

$v_n^2, [p_T]$ correlation

$$\rho(v_2^2, [p_T]) = \frac{\langle (v_2^2 - \langle v_2^2 \rangle) ([p_T] - \langle [p_T] \rangle) \rangle}{\sqrt{\langle (v_2^2 - \langle v_2^2 \rangle)^2 \rangle \langle ([p_T] - \langle [p_T] \rangle)^2 \rangle}}$$

- Clear centrality dependence
- Gains negative values



REQUIREMENT OF $\langle p_T \rangle$ EVENT BY EVENT

Full equation

$$\rho(v_2^2, [p_T]) = \frac{\langle (v_2^2 - \langle v_2^2 \rangle) ([p_T] - \langle [p_T] \rangle) \rangle}{\sqrt{\langle (v_2^2 - \langle v_2^2 \rangle)^2 \rangle \langle ([p_T] - \langle [p_T] \rangle)^2 \rangle}}$$

- With one set of hydro runs?
 - Store all particles from hydro
 - Uses too much space
- Use mean values from previous runs!

SPC SENSITIVITY COMPARISON

- Previous biased scalar product method

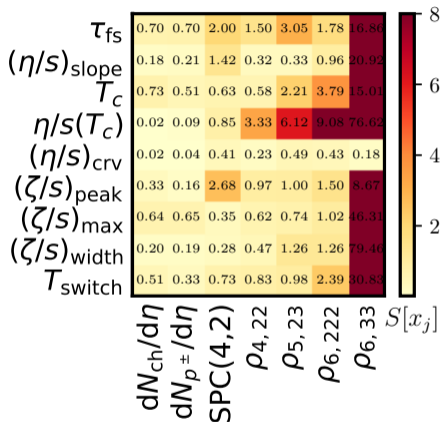
$$\langle \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle_{SP} = \frac{\langle v_{n_1}^{a_1} \dots v_{n_k}^{a_k} \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle}{\sqrt{\langle v_{n_1}^{2a_1} \rangle \dots \langle v_{n_k}^{2a_k} \rangle}}$$

- New Gaussian distribution method

$$\begin{aligned} \langle \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle_{GE} &= \int d\Theta N_{\Theta}(\Theta) \cos \Theta \\ &\approx \sqrt{\frac{\pi}{4}} \frac{\langle v_{n_1}^{a_1} \dots v_{n_k}^{a_k} \cos(a_1 n_1 \Psi_1 + \dots + a_k n_k \Psi_k) \rangle}{\sqrt{\langle v_{n_1}^{2a_1} \rangle \dots \langle v_{n_k}^{2a_k} \rangle}} \end{aligned}$$

- Way to probe the non-linear response, e.g.

$$v_2^2 v_4 e^{i4(\Psi_4 - \Psi_2)} = \omega_2 \omega_4 c_2^2 c_4 e^{i4(\Phi_4 - \Phi_2)} + \omega_{422} \omega_2^2 c_2^2$$



- The new SPC shows higher dependence on the specific bulk viscosity rather than the specific shear viscosity.