

# Which vacuum pressure is acceptable in the booster?

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Acknowledgements:

G. Rumolo, F. Zimmermann

FCCIS 2023 WP2 workshop  
Rome, Italy  
13-15 November 2023

# **Which vacuum pressure is acceptable in the booster?**

## **Vacuum constraints from ion trapping and instability**

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Acknowledgements:

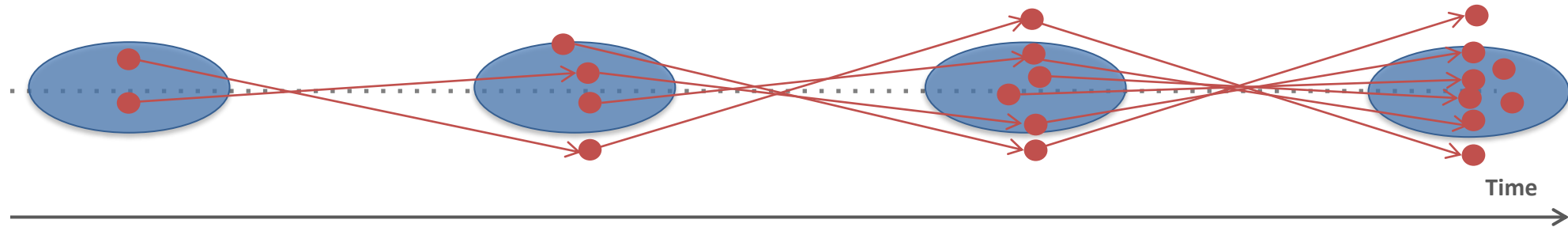
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# Ion trapping in electron machines

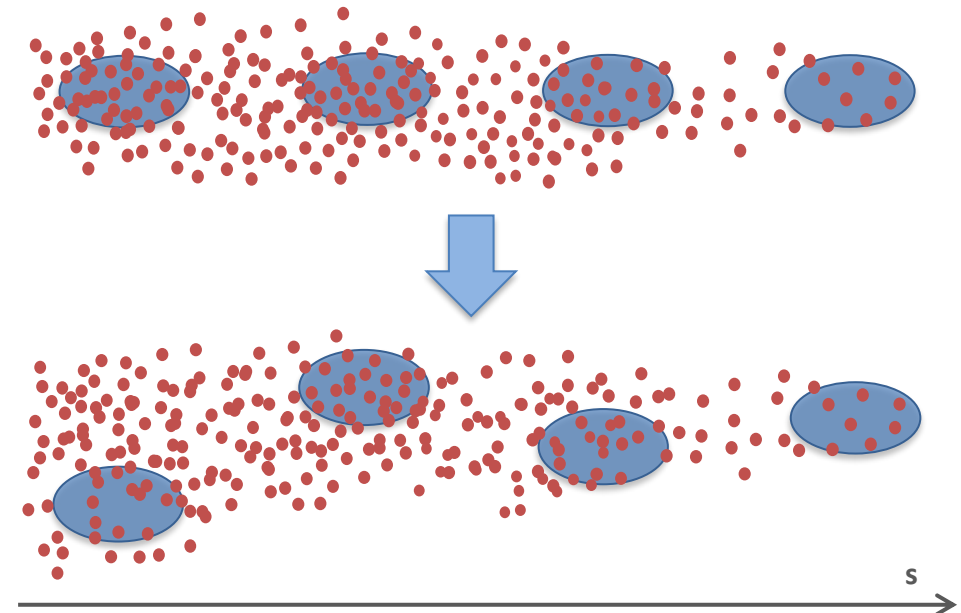
Beam-induced gas ionization gives rise to electrons and ions along the beam path

- The positive ions are attracted by the electron beam field and may be trapped in oscillation along the bunch train



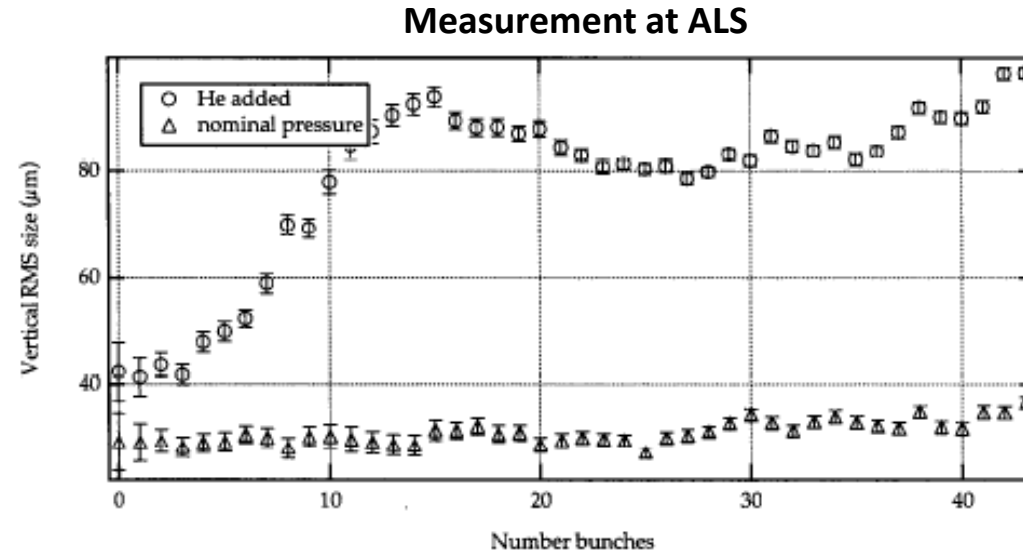
- Trapped ions accumulate over the passage of a bunch train and can lead to a fast beam-ion instability (FBII)

- Due to their large mass, the trapped ions barely move during the passage of a single bunch
- However, they transfer information on the offset of one bunch to the next, seeding a coupled-bunch instability
- Ion trapping also gives rise e.g. to emittance growth and tune shifts



# Machine observations

- Fast beam-ion instabilities have been observed since the 90's in several machines
  - ALS
  - PLS
  - BESSY II
  - ELETTRA
  - ALBA
  - SOLEIL
  - Cesr-TA



*J. Byrd, A. Chao, S. Heifets, et al.  
Phys. Rev. Lett. 79 (1997), 79-82*

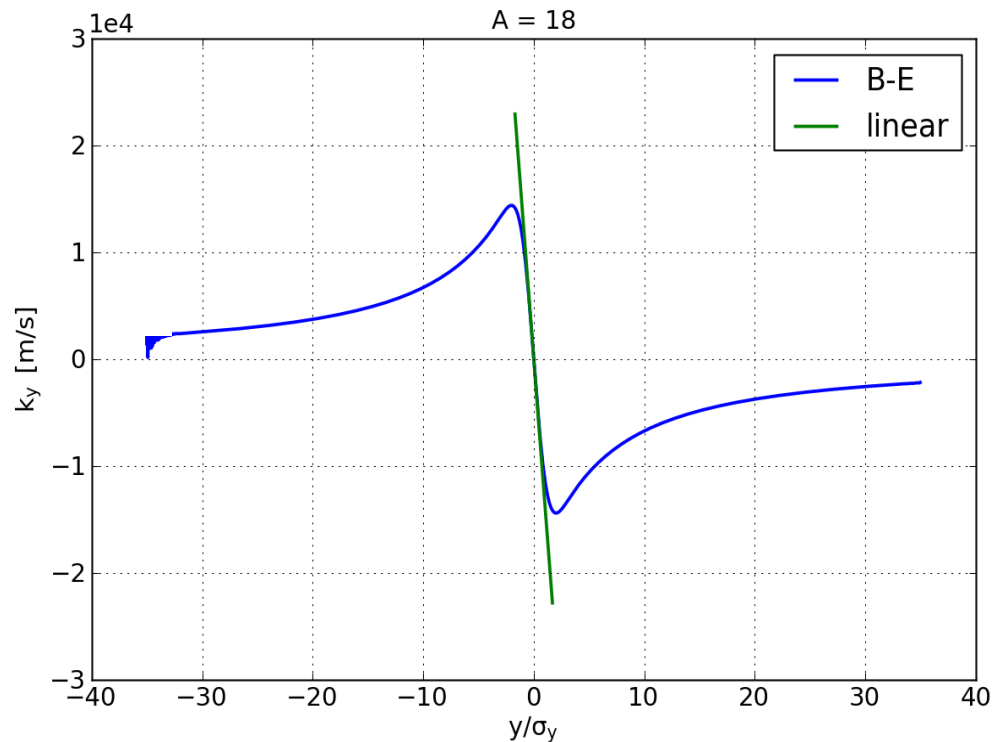
- In most cases the instability has been observed only in the presence of vacuum degradation
  - During commissioning
  - Due to a local pressure rise, e.g. due to impedance-induced heating
  - During experiments with additional injected gas

# Analytical model: Ion trapping

Ion trapping can be modelled using the linear approximation of the Bassetti-Erskine formula for a Gaussian beam field

- An ion with mass number  $A$ , receives a velocity kick by the beam

$$\frac{2N_b r_p c}{A} \frac{x,y}{(\sigma_x + \sigma_y) \sigma_{x,y}} = k_{x,y} * x, y$$



Raubenheimer et al., Phys. Rev. E 52, 5, 5487,  
Stupakov et al., Phys. Rev. E 52, 5, 5499

$N_b$  = bunch intensity  
 $r_p$  = classical proton radius  
 $\sigma_{x,y}$  = transverse beam size

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- During the bunch spacing  $T_b$ , the ions drift
  - The kick – drift – kick – drift... motion is stable if  $k_{x,y} T_b < 4$   
 → Trapping condition for the ion mass number

$$A > A_{\text{trap}} = \frac{N_b r_p T_b c}{2(\sigma_x + \sigma_y) \sigma_{x,y}}$$

- The instability rise time can also be estimated
  - Opposite dependence on bunch spacing, beam size and intensity w.r.t the trapping condition
  - Faster instability for large pressure, small mass (if trapped) and long trains

$$\tau_{\text{inst}}^2 \propto \frac{\gamma^2 A \omega_\beta}{n_b^4 N_b^3 P^2 T_b c} (\sigma_x + \sigma_y)^3 \sigma_{x,y}^3$$

$N_b$  = bunch intensity

$r_p$  = classical proton radius

$\sigma_{x,y}$  = transverse beam size

$n_b$  = nr of bunches

$P$  = pressure

$\omega_\beta$  = betatron frequency

*Raubenheimer et al., Phys. Rev. E 52, 5, 5487,  
 Stupakov et al., Phys. Rev. E 52, 5, 5499*

# Ion trapping in the booster

- Since ion trapping is more likely for small bunch spacing, it is of concern mainly for Z mode operation
  - We consider the booster parameters as presented at the FCC Week 2023 by A. Chance
    - Assuming the “standard” booster optics and LINAC as injector
  - For the filling scheme, we assume “case 4” from the MTR, with 10% of the collider bunch intensity

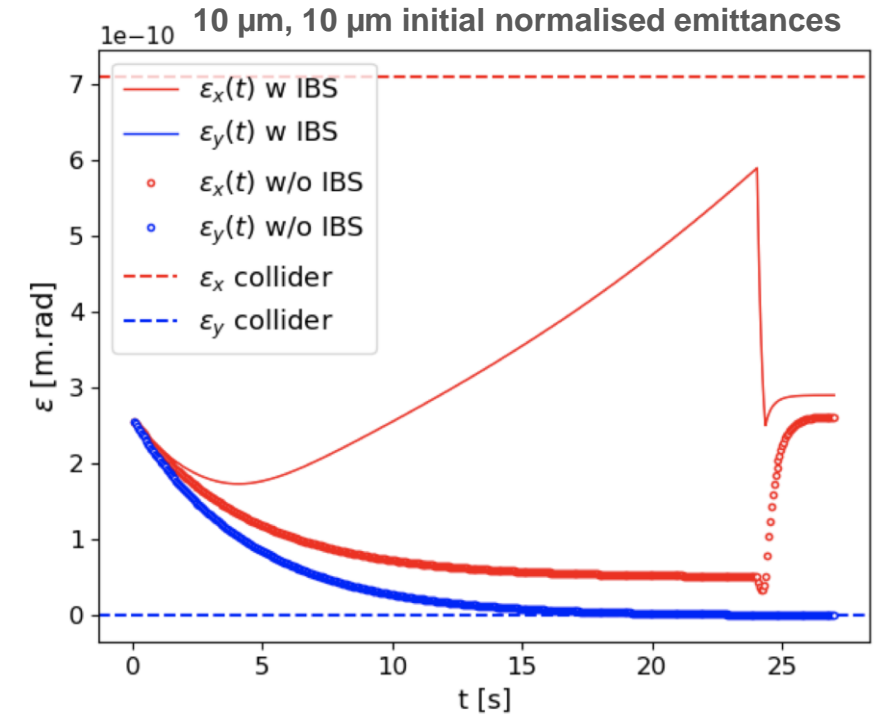
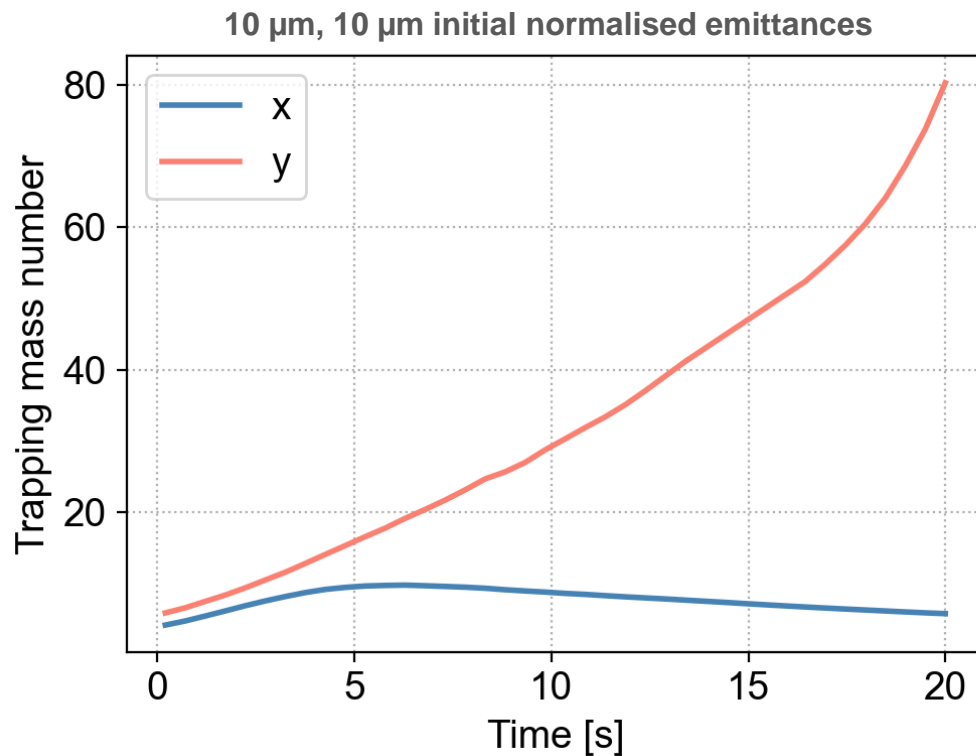
## A. Chance et al, FCC Week 2023

Injector		LINAC	LINAC	LINAC	Filling scheme	1	2	3	4
Injection energy	GeV	20	20	20	Bunch Intensity [ $\times 10^{10}$ ppb]	1.51	2.15	2.15	2.43
Extraction energy	GeV	45.6	45.6	45.6	Bunch spacing [ns]	15	20	25	25
Injection hor. emittance $\epsilon_{x,inj}$ (norm.)	$\mu\text{m}$	50	10	10	Train length [# bunches]	2.15	280	560	255
Injection vert. emittance $\epsilon_{y,inj}$ (norm.)	$\mu\text{m}$	50	10	1					
Injection energy spread $\delta_{p,inj}$	%	0.1-0.15	0.1-0.15	0.1-0.15	0.4				
Extraction hor emittance $\epsilon_{x,ext}$ (geom)	$\mu\text{m}$	<0.3	<0.3	<0.3	<0.3				
Extraction vert emittance $\epsilon_{y,ext}$ (geom)	pm	<1.42	<1.42	<1.42	<1.42				
Extraction energy spread $\delta_{p,ext}$	%	0.04	0.04	0.04	0.04				
Accumulation time	s	24	24	24	54				
Ramp time	s	0.32	0.32	0.32	0.37				
Flat top	s	2.6	1.9	1	1.9				
Total cycling time	s	26.92	26.22	25.32	56.27				

# Ion trapping in the booster

*A. Chance et al, FCC Week 2023*

- Based on the emittance evolution during the booster cycle, we can estimate the evolution of the trapping mass
  - The trapping condition is determined by the vertical plane
  - The trapping mass is lowest, i.e. most critical, right after injection and increases with the decrease in vertical emittance

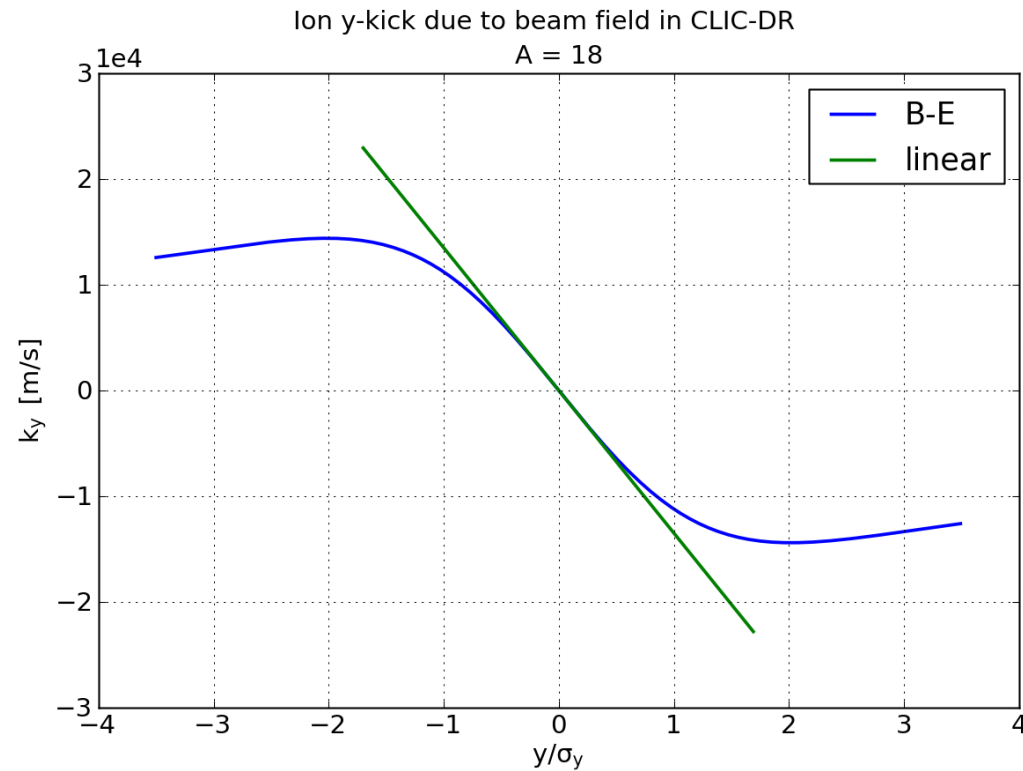


- In this case the initial trapping mass is around 5
  - Consequently,  $\text{H}_2$  ( $A = 2$ ) should not be trapped, while most heavier species would be trapped
  - However, in reality, the situation is more subtle



# Analytical model: beyond the linear approximation

The linear approximation is good only in a small region around the centre of the beam:  $x, y \leq 0.5 \sigma_{x,y}$



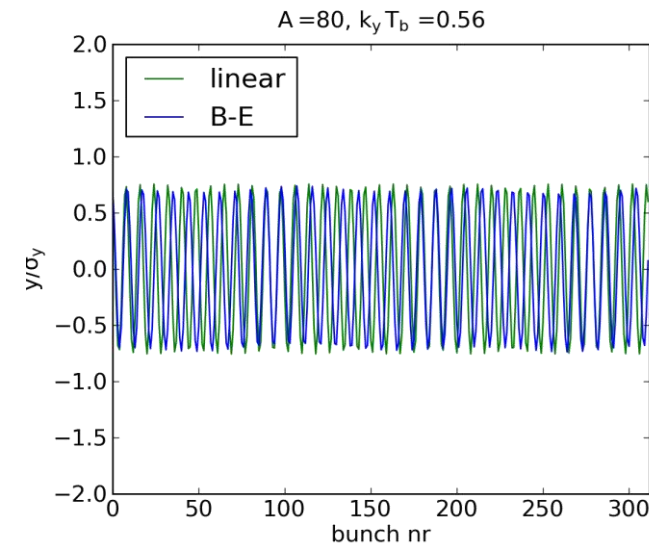
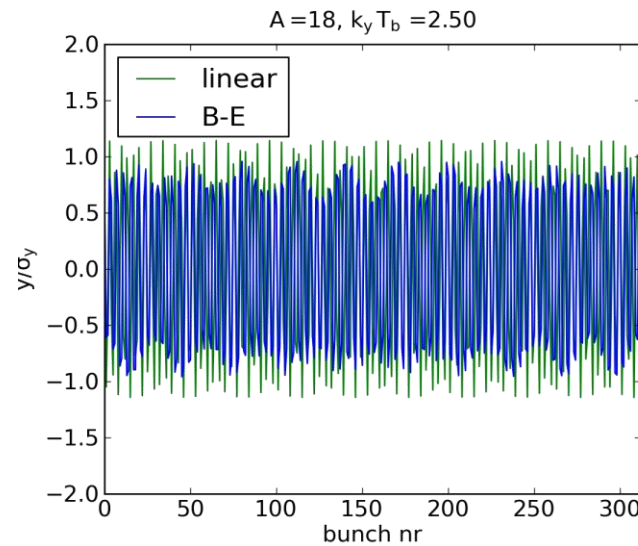
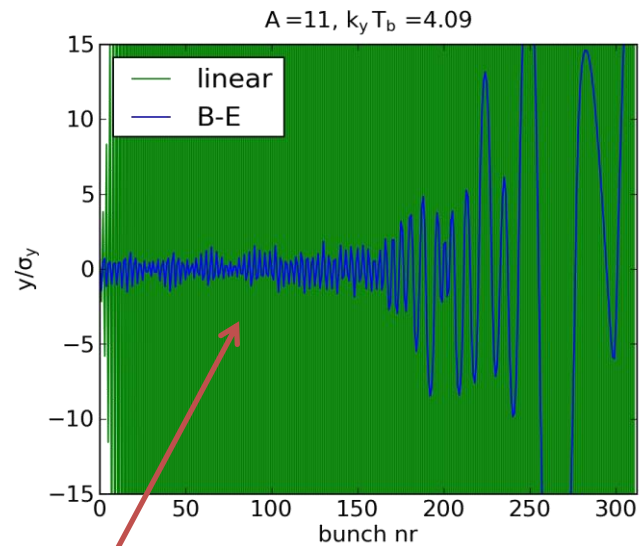
# Analytical model: beyond the linear approximation

The linear approximation is good only in a small region around the centre of the beam:  $x, y \leq 0.5 \sigma_{x,y}$

- In the non-linear regime, ion trajectories are altered and ions below the trapping mass can oscillate around a bunch train for some time (“weak trapping”)

## Trajectories for ion of mass A during the passage of a CLIC bunch train

$$x_0 = 0.7 \sigma_x, y_0 = 0.7 \sigma_y$$



Ion is trapped for several bunch passages

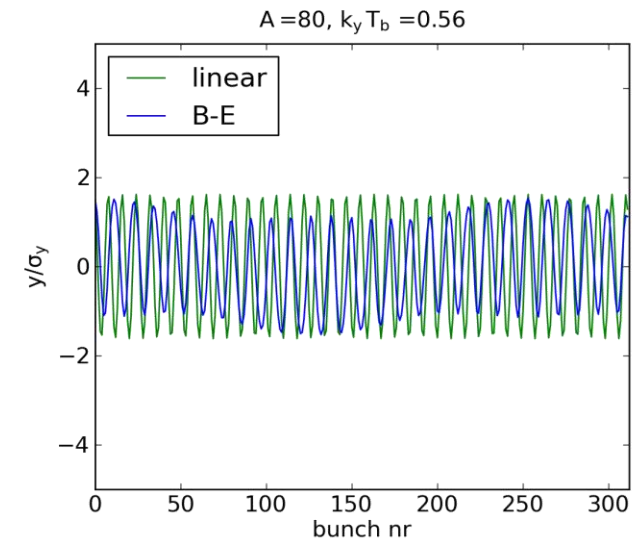
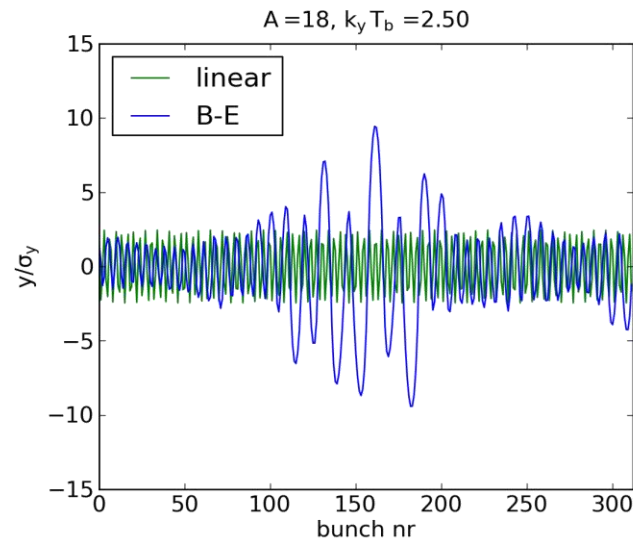
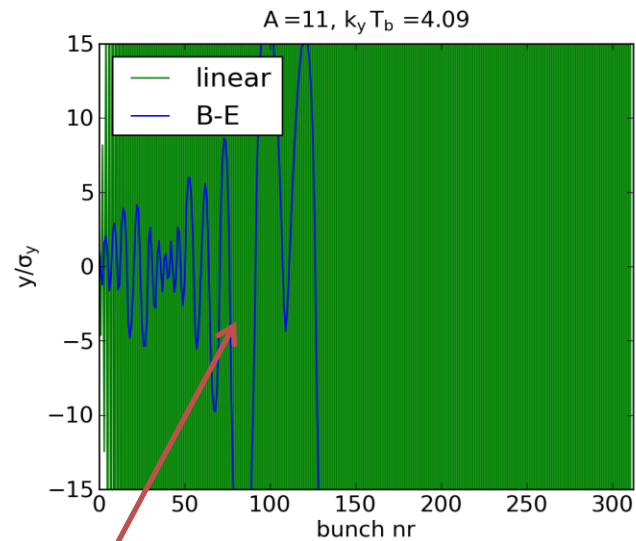
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- In the non-linear regime, ion trajectories are altered and ions below the trapping mass can oscillate around a bunch train for some time (“weak trapping”)

## Trajectories for ion of mass A during the passage of a CLIC bunch train

$$x_0 = 1.5 \sigma_x, y_0 = 1.5 \sigma_y$$

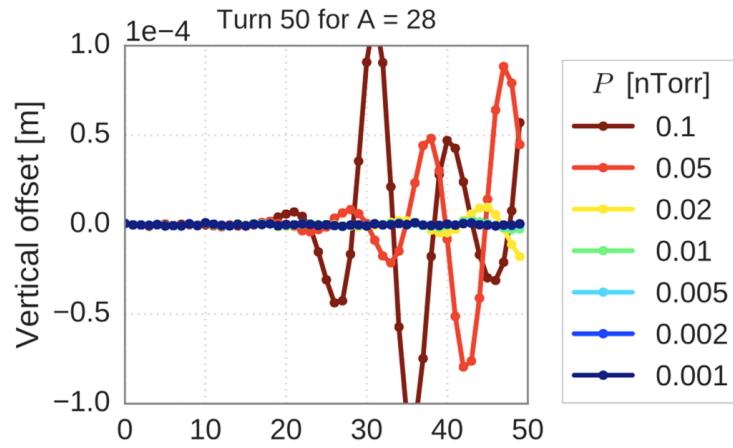


**Ion is trapped for several bunch passages**

# Weak trapping in simulations

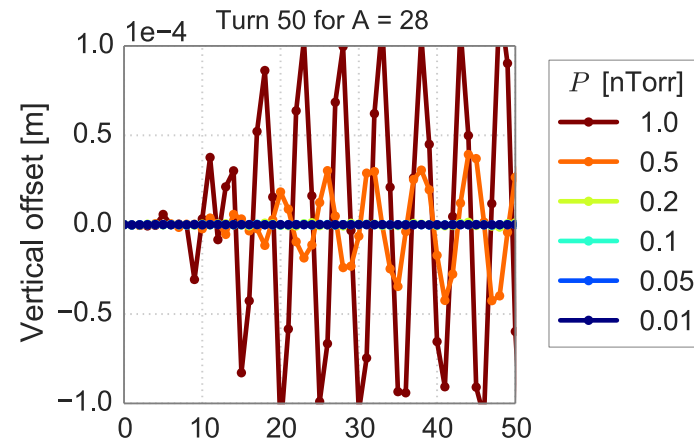
- This effect has been observed also in a past simulation study for the FCC-ee from 2017 (with obsolete parameters)
  - Weak trapping can also cause instability, but with different characteristics than for fully trapped ions

$A_{\text{trap}} \approx 10$ ,  $N_2$  and  $CO_2$  are fully trapped

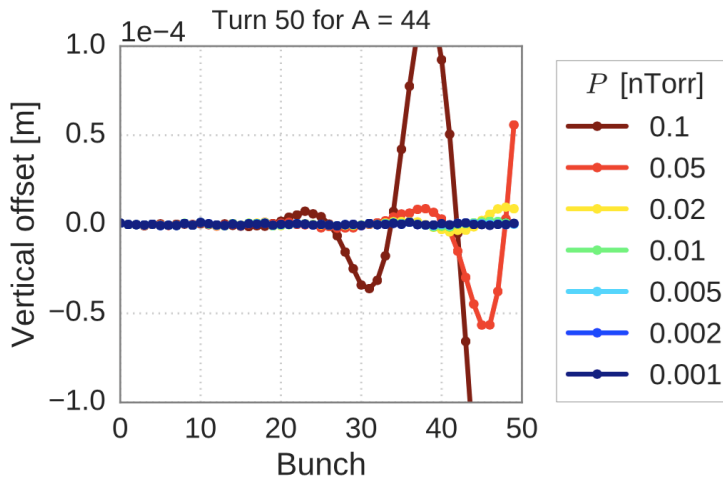


The instability starts at the end of trains, moving towards the front with increasing pressure

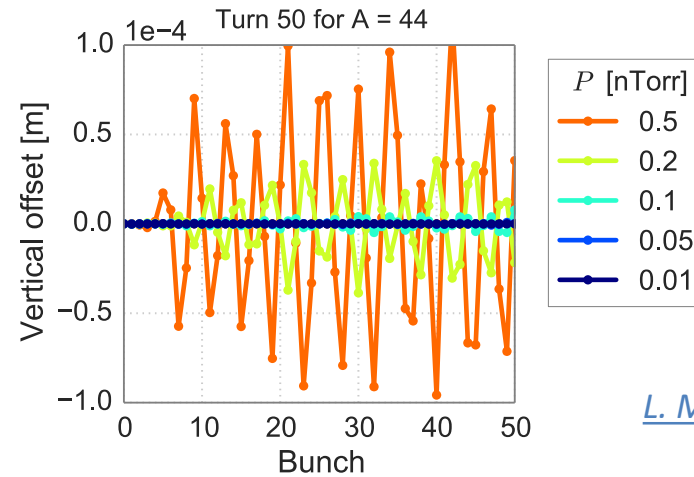
$A_{\text{trap}} \approx 80$ ,  $N_2$  and  $CO_2$  are weakly trapped



The instability develops simultaneously over most of the train



The instability is stronger for **lower A**



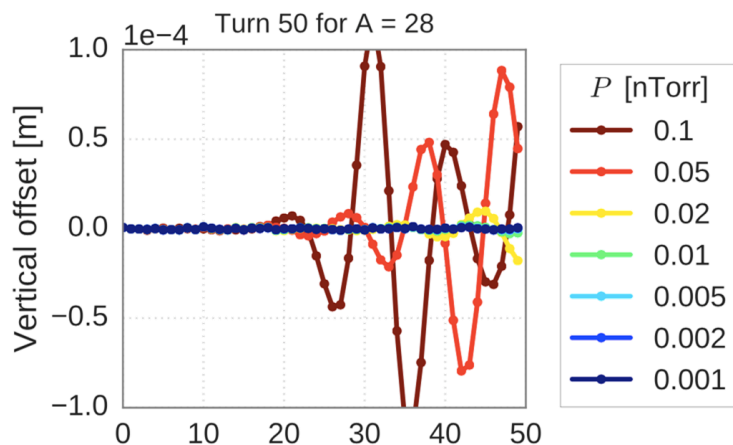
The instability is stronger for **higher A**

*L. Mether et al, FCC Week 2018*

# Weak trapping in simulations

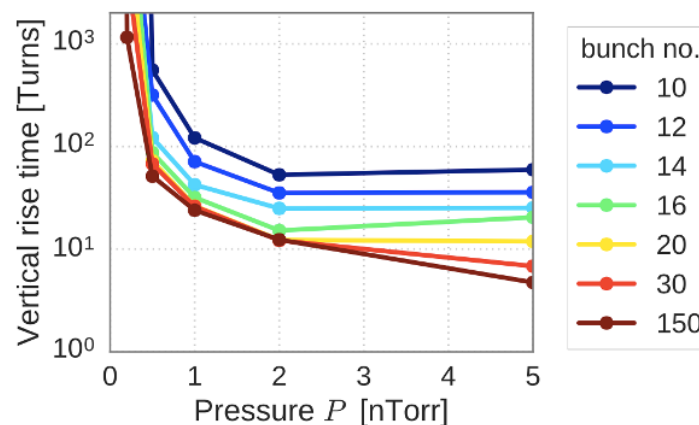
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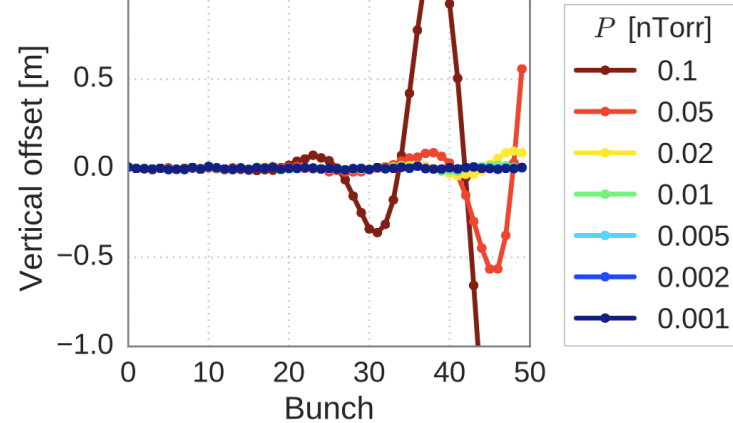
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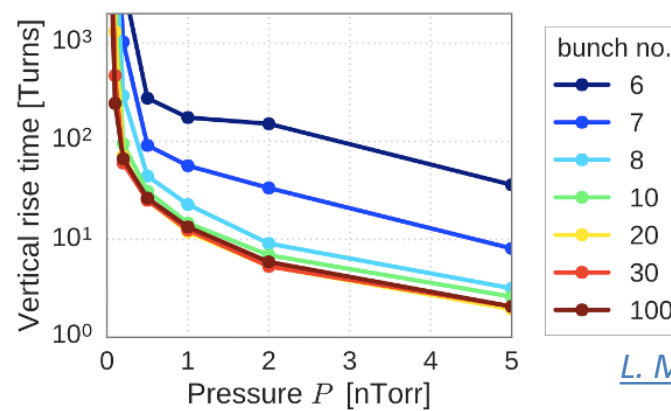


Bunch-by-bunch rise times (based on emittance) saturate after a small part of the train

$A_{\text{trap}} \approx 44$ ,  $N_2$  and  $CO_2$  are weakly trapped



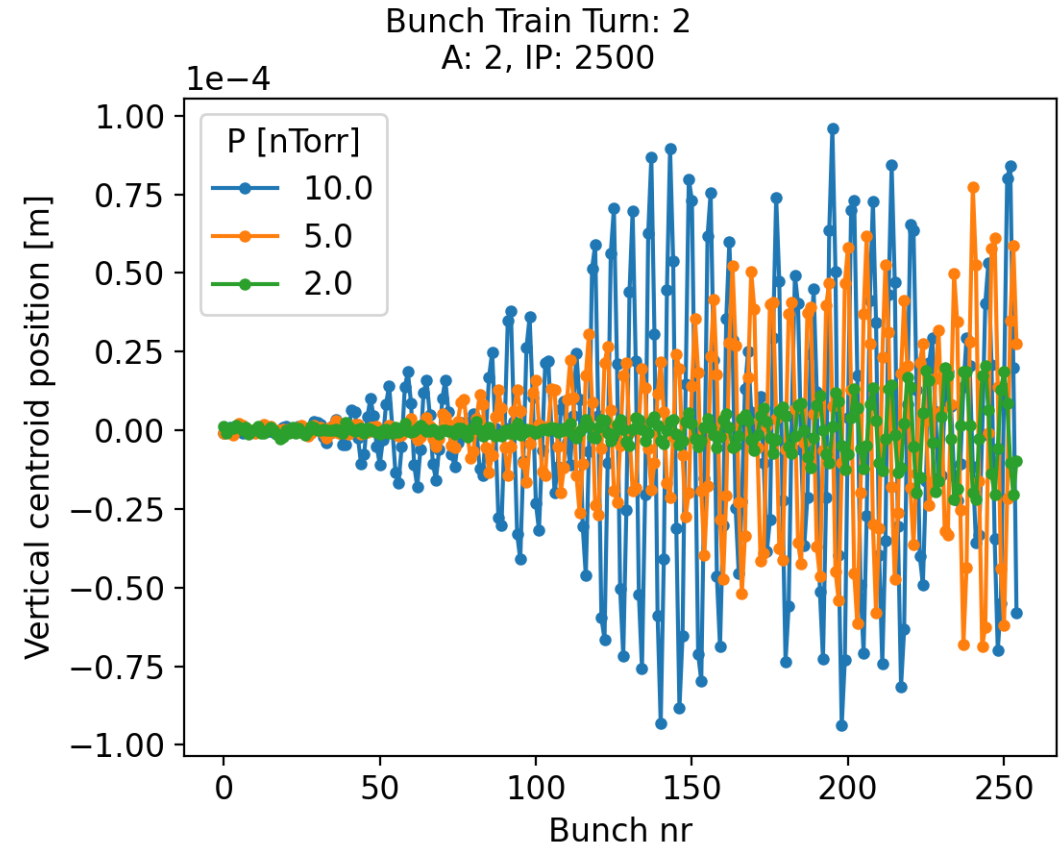
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*L. Mether et al, FCC Week 2018*

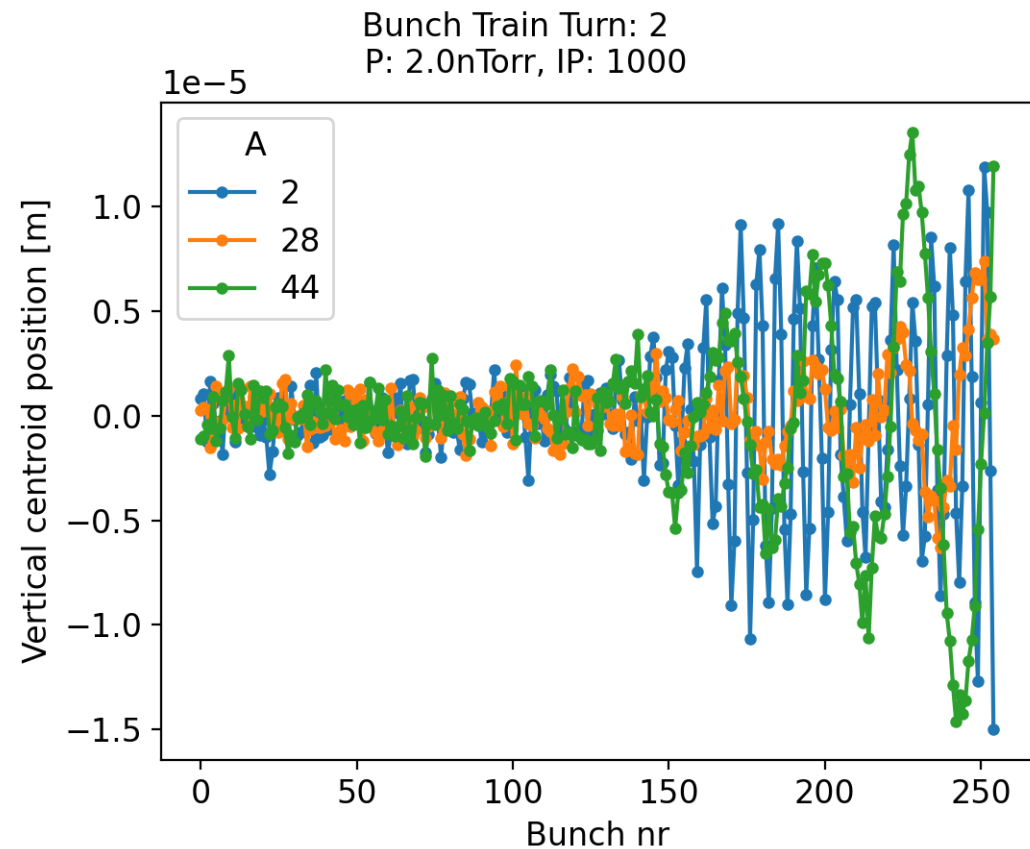
# First instability simulations for booster

- Ideally, we should simulate the FBII with a realistic composition of vacuum species for different total pressures over at least a damping time
  - FBII simulations are computationally very heavy  
→ such a study takes months to perform
- First simulations cover only a few cases over a small number of turns
  - Focus on H<sub>2</sub>, despite weaker trapping, since heavier species make up a much smaller fraction of the total pressure
  - With initial emittances of 10 μm in both planes, an instability develops within 2 turns even for a pressure of 2 nTorr  $\approx 2.7e-9$  mbar



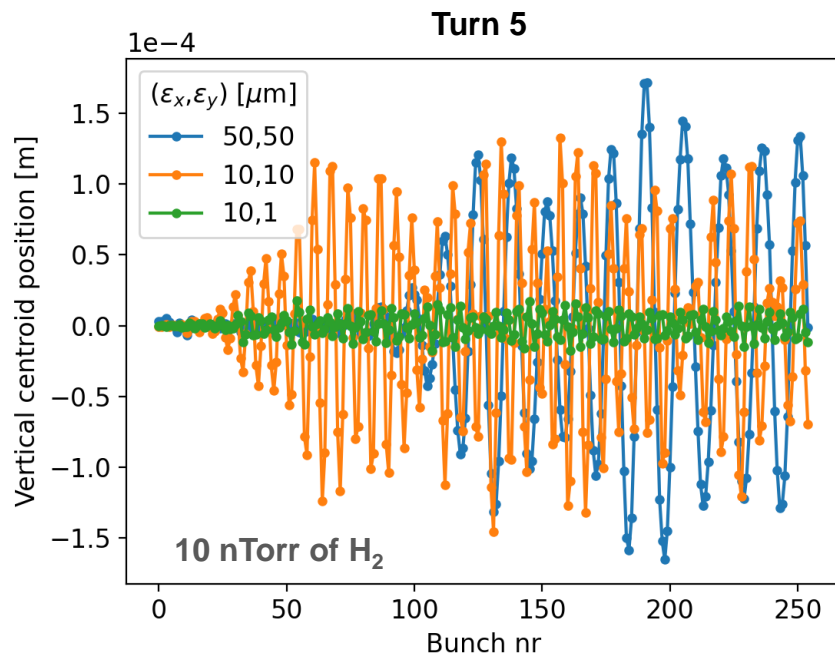
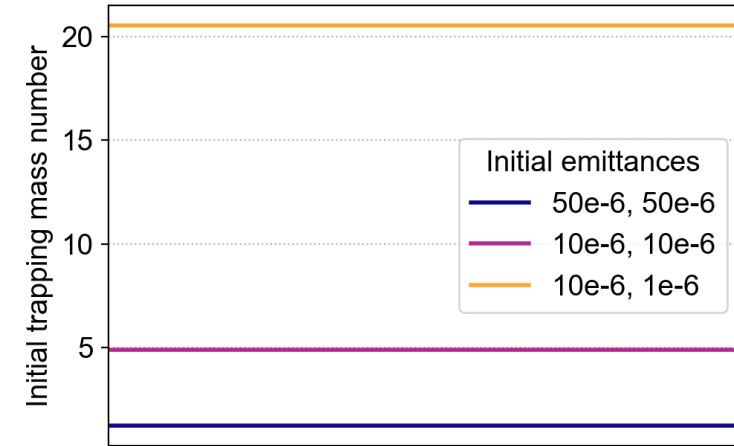
# Trapping of heavier species

- Heavier species, such as CO, N<sub>2</sub> and CO<sub>2</sub>, are strongly trapped with a round beam of 10 μm
- Also for these species, an instability develops within 2 turns for a pressure of 2 nTorr ≈ 2.7e-9 mbar
  - However, this would correspond to larger total pressures, since the fraction of heavier species is typically much lower

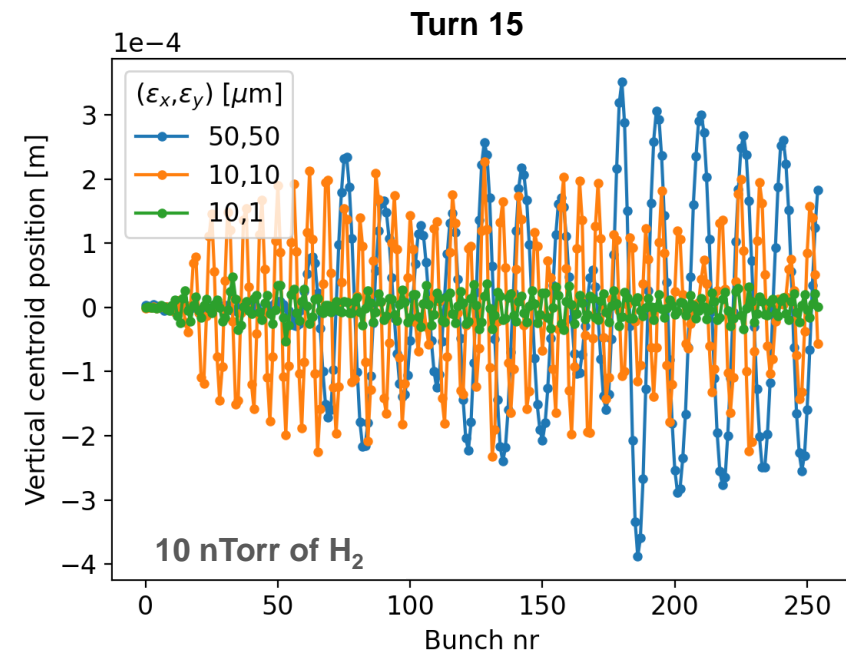


# Impact of initial emittances

- Since the trapping mass number increases with decreasing emittance, a lower (vertical) emittance at injection should alleviate the instability, at least for  $H_2$
- Simulations confirm that the option with  $1\ \mu\text{m}$  vertical emittance at injection is significantly less impacted by  $H_2$  trapping, although smaller amplitude oscillations still occur



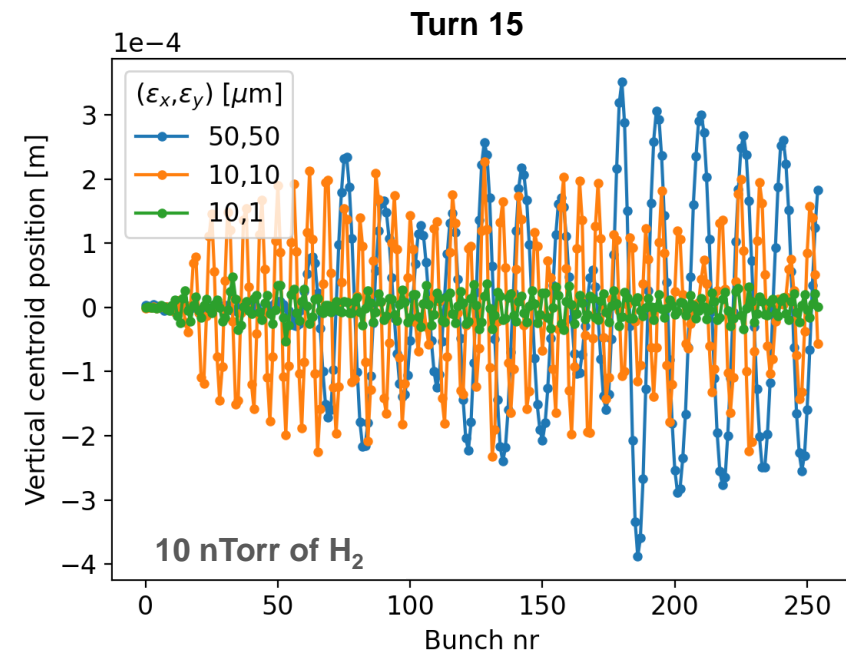
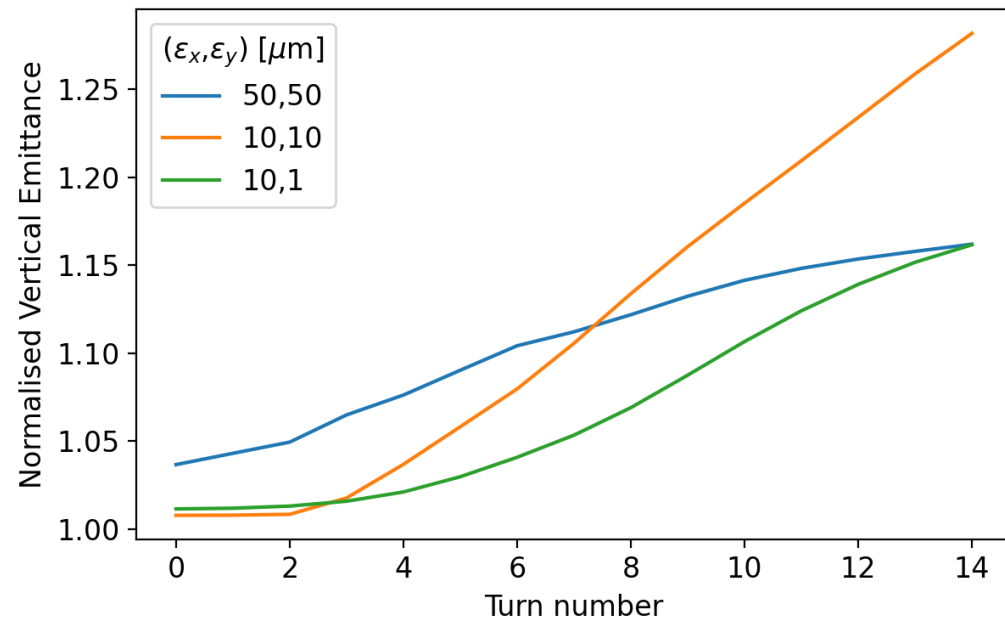
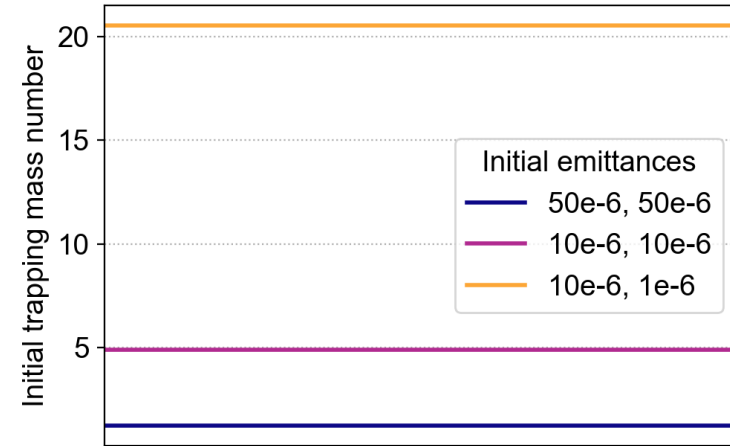
Note the change in scale





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  - Emittance growth occurs in all cases



# Conclusion and outlook

- The fast beam-ion instability is likely to cause problems with the currently foreseen vacuum system
  - Ion trapping is strongest just after injection
    - Heavier ion species, e.g. CO, N<sub>2</sub>, CO<sub>2</sub> are trapped for all considered emittances
    - Hydrogen is weakly trapped for emittance of ~30 μm or smaller, but can still cause instability
  - Opting for a smaller vertical emittance at injection reduces the impact, but it is probably not sufficient
- Mitigation measures that don't rely on low vacuum pressure can be considered
  - The instability can generally be mitigated with a transverse feedback, but ion trapping will still occur and may still lead to tune shift and emittance growth
  - Clearing electrodes could be another solution, but design and impact e.g. on impedance need to be considered
- A more extensive study (with the latest parameters) is needed for more accurate constraints
  - The need for mitigation measures should be assessed and their effectiveness evaluated
  - Mitigation strategy should be included in the overall design process, including in cost estimates