

Status of phenomenological studies of GPDs

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Forward Physics in ALICE 3 workshop, Heidelberg, Germany, October 18th, 2023

- Introduction
- Experimental campaign
- Recent progress
- New sources of GPD information
- Numerical tools
- Summary

Nucleon is not a point-like particle, it is made out of partons:

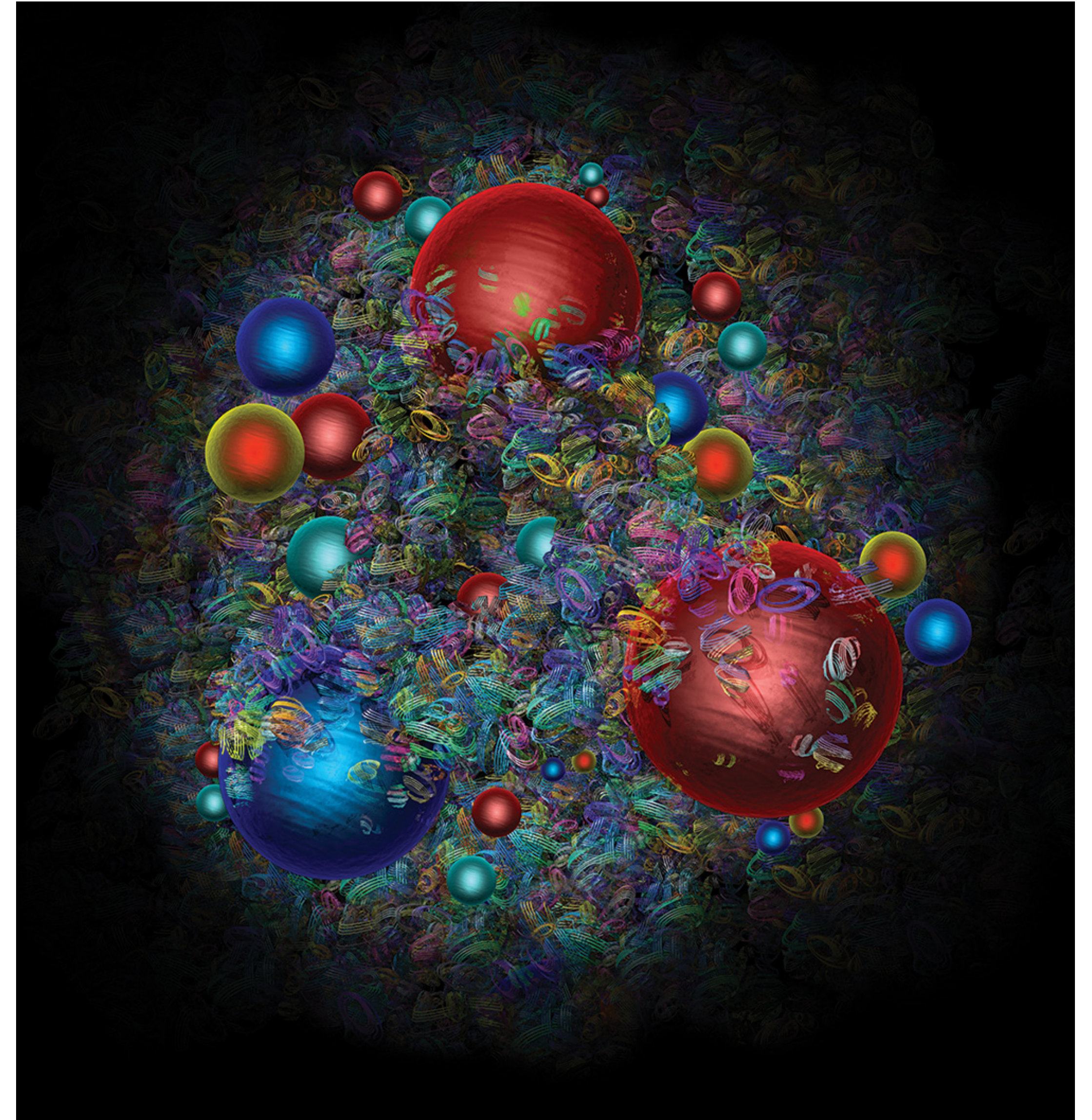
- quarks
- gluons

How can we recover basic properties of nucleon from those of its constituents?

- charge
- spin
- mass

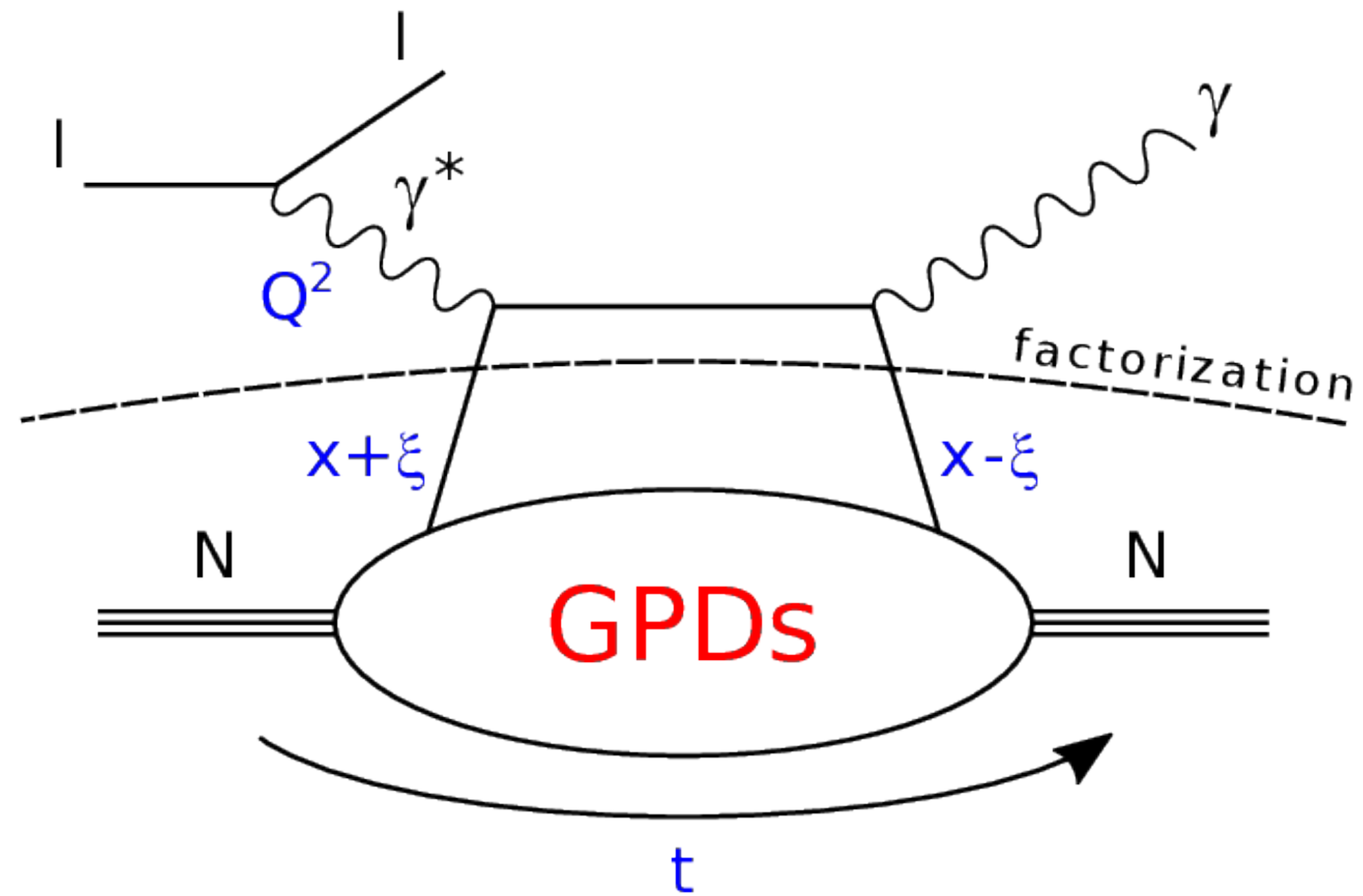
How partons are distributed inside nucleon?

- momentum (longitudinal and transverse)
- position
- polarisation
- "mechanical" properties



from CERN Courier / D. Dominguez

Deeply Virtual Compton Scattering (DVCS)



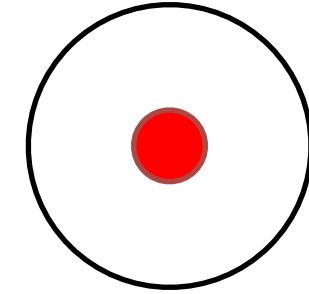
factorization for $|t|/Q^2 \ll 1$

Chiral-even GPDs:
(helicity of parton conserved)

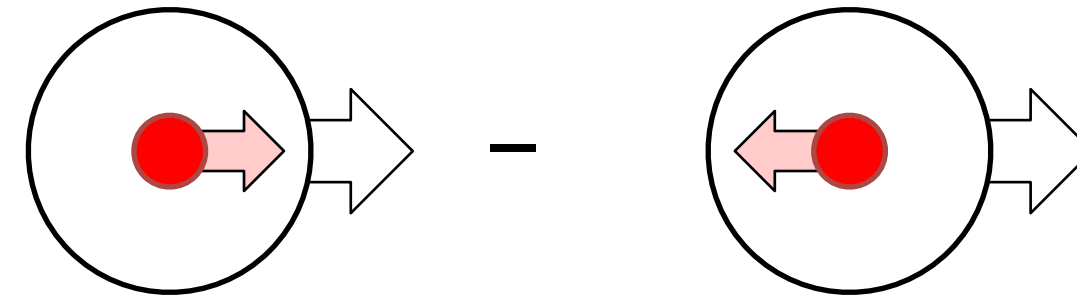
$H^{q,g}(x, \xi, t)$	$E^{q,g}(x, \xi, t)$	<i>for sum over parton helicities</i>
$\tilde{H}^{q,g}(x, \xi, t)$	$\tilde{E}^{q,g}(x, \xi, t)$	<i>for difference over parton helicities</i>
<i>nucleon helicity conserved</i>	<i>nucleon helicity changed</i>	

- Reduction to PDFs:

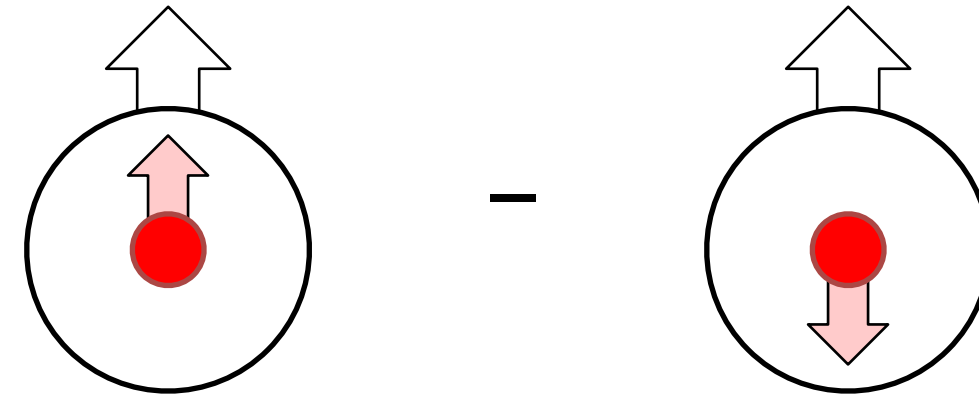
$$H^q(x, 0, 0) \equiv q(x)$$



$$\tilde{H}^q(x, 0, 0) \equiv \Delta q(x)$$



$$H_T^q(x, 0, 0) \equiv h_1(x)$$



no corresponding relations exist for other GPDs

- Reduction to Elastic Form Factors (EFFs):

$$\int_{-1}^1 dx H^q(x, \xi, t) \equiv F_1^q(t)$$

$$\int_{-1}^1 dx E^q(x, \xi, t) \equiv F_2^q(t)$$

$$\int_{-1}^1 dx \tilde{H}^q(x, \xi, t) \equiv g_A^q(t)$$

$$\int_{-1}^1 dx \tilde{E}^q(x, \xi, t) \equiv g_P^q(t)$$

Polynomiality - non-trivial consequence of Lorentz invariance:

$$\int_{-1}^1 dx x^n H^q(x, \xi, t) = h_0^{q,n}(t) + \xi^2 h_2^{q,n}(t) + \dots + \text{mod}(n, 2) \xi^{n+1} h_{n+1}^{q,n}(t)$$

$$\int_{-1}^1 dx x^n \tilde{H}^q(x, \xi, t) = \tilde{h}_0^{q,n}(t) + \xi^2 \tilde{h}_2^{q,n}(t) + \dots + \text{mod}(n+1, 2) \xi^n \tilde{h}_n^{q,n}(t)$$

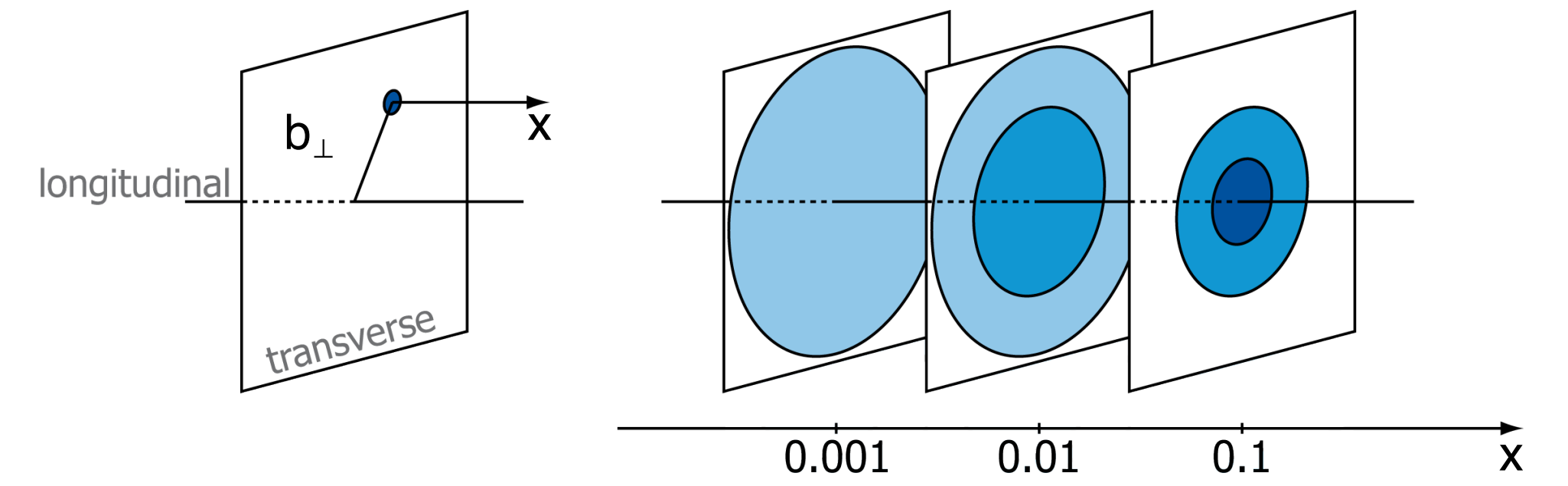
Positivity bounds - positivity of norm in Hilbert space, e.g.:

$$(1 - \xi^2) \left(H^q - \frac{\xi^2}{1 - \xi^2} E^q \right)^2 + \frac{t_0 - t}{4m^2} (E^q)^2 \leq q \left(\frac{x + \xi}{1 + \xi} \right) q \left(\frac{x - \xi}{1 - \xi} \right)$$

strong constraint on GPD parameterizations!

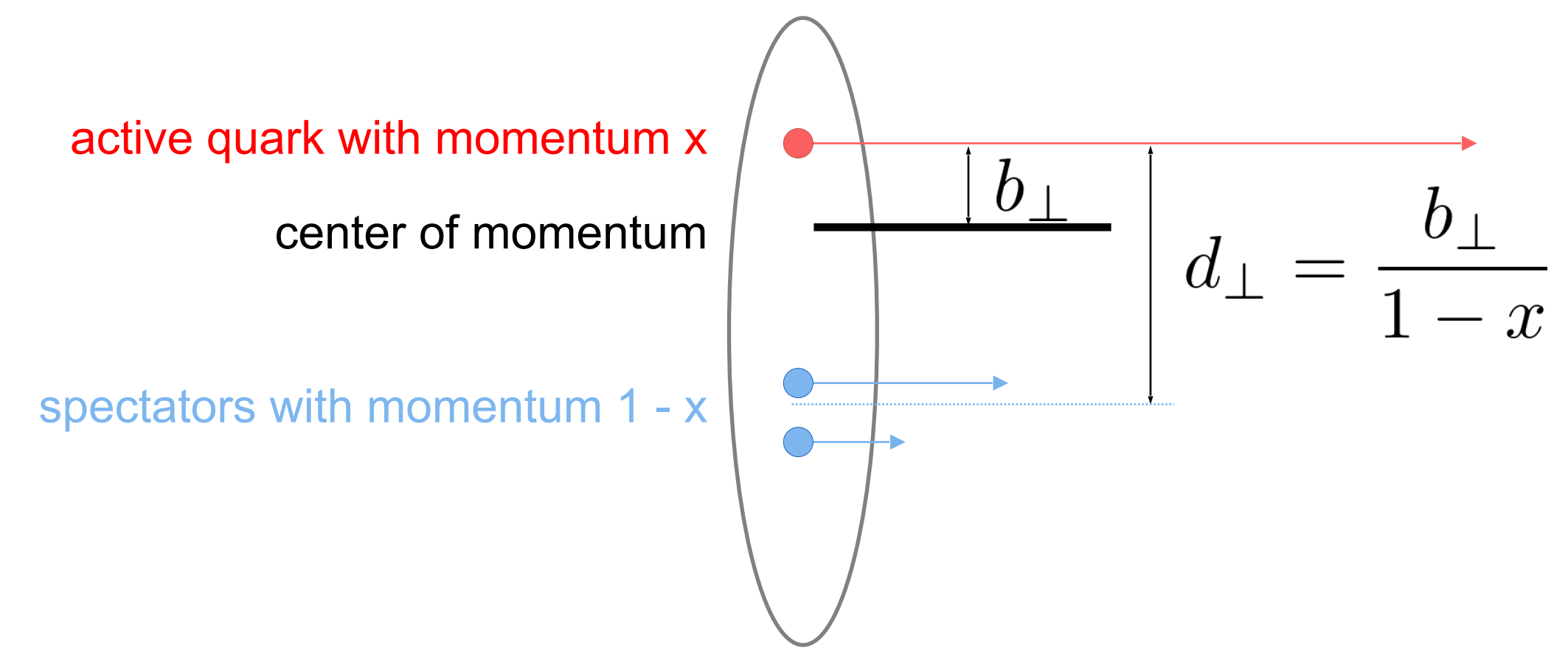
Nucleon tomography

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta}{4\pi^2} e^{-i\mathbf{b}_\perp \cdot \Delta} H^q(x, 0, t = -\Delta^2)$$



Study of long. polarization with GPD \tilde{H}
 Study of distortion in transv. polarized nucleon with GPD E

Impact parameter \mathbf{b}_\perp defined w.r.t. center of momentum, such as $\sum x \mathbf{b}_\perp = 0$



$$T^{\mu\nu} = \begin{bmatrix} \text{Energy density} & & & \\ T^{00} & T^{01} & T^{02} & T^{03} \\ \text{Energy flux} & \text{Momentum density} & & \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \\ \text{Momentum flux} & & & \\ & & & \text{Normal stress} \end{bmatrix}$$

Shear stress

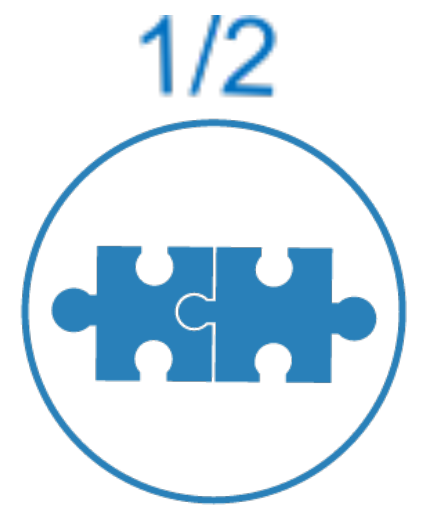
Energy momentum tensor in terms of form factors:

$$\langle p', s' | \hat{T}^{\mu\nu} | p, s \rangle = \bar{u}(p', s') \left[\frac{P^\mu P^\nu}{M} A(t) + \frac{\Delta^\mu \Delta^\nu - \eta^{\mu\nu} \Delta^2}{M} C(t) + M \eta^{\mu\nu} \bar{C}(t) + \frac{P^\mu i \sigma^{\nu\lambda} \Delta_\lambda}{4M} [A(t) + B(t) + D(t)] + \frac{P^\nu i \sigma^{\mu\lambda} \Delta_\lambda}{4M} [A(t) + B(t) - D(t)] \right] u(p, s)$$

Total angular momentum:

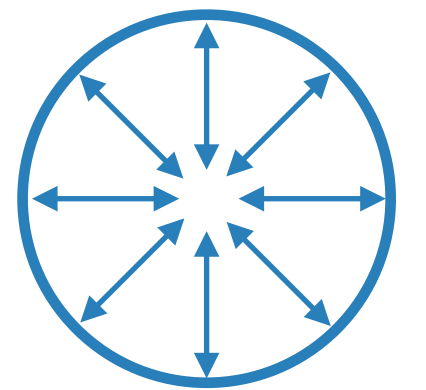
$$A^q(0) + B^q(0) = \int_{-1}^1 x [H^q(x, \xi, 0) + E^q(x, \xi, 0)] = 2J^q$$

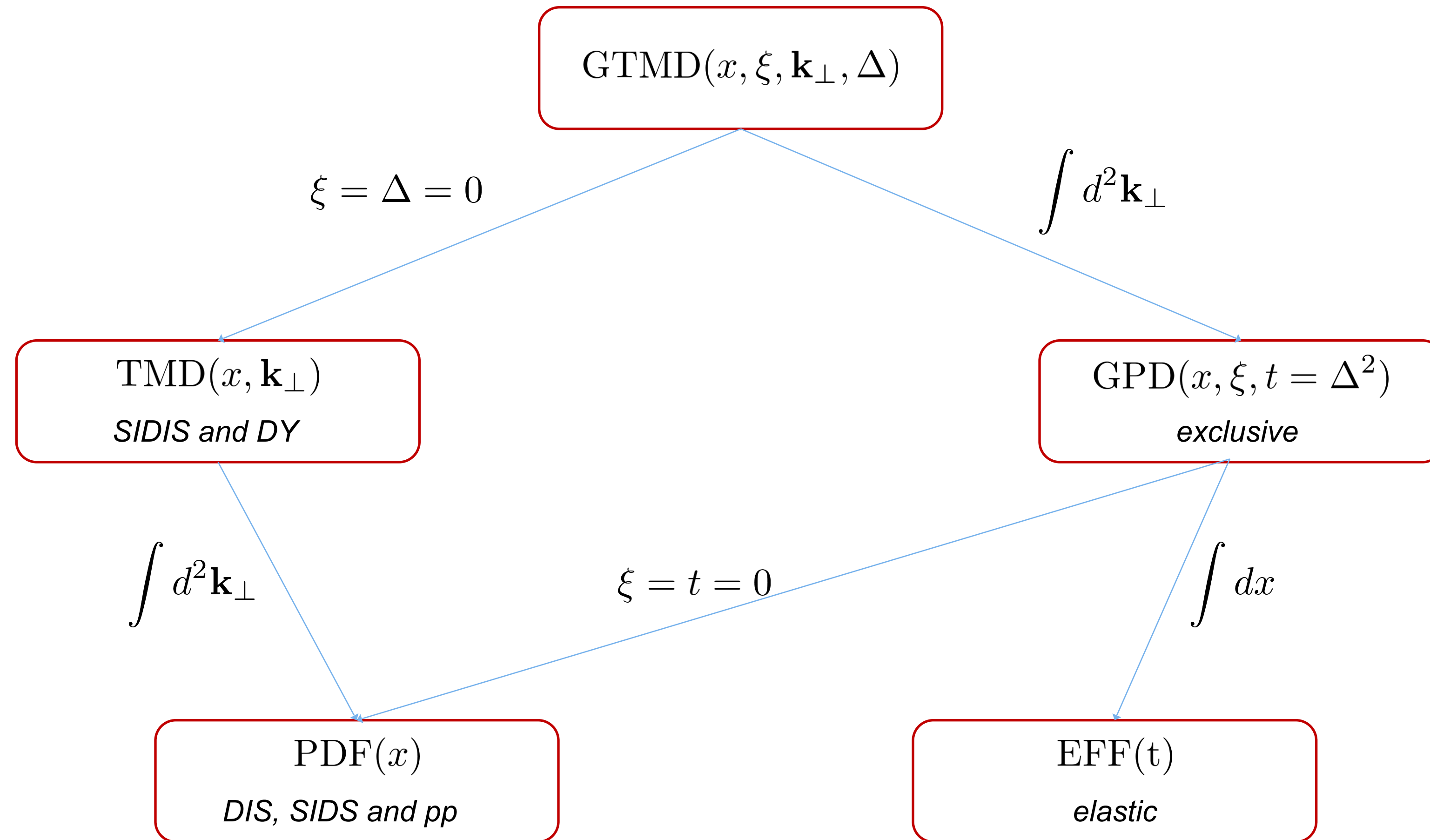
Ji's sum rule



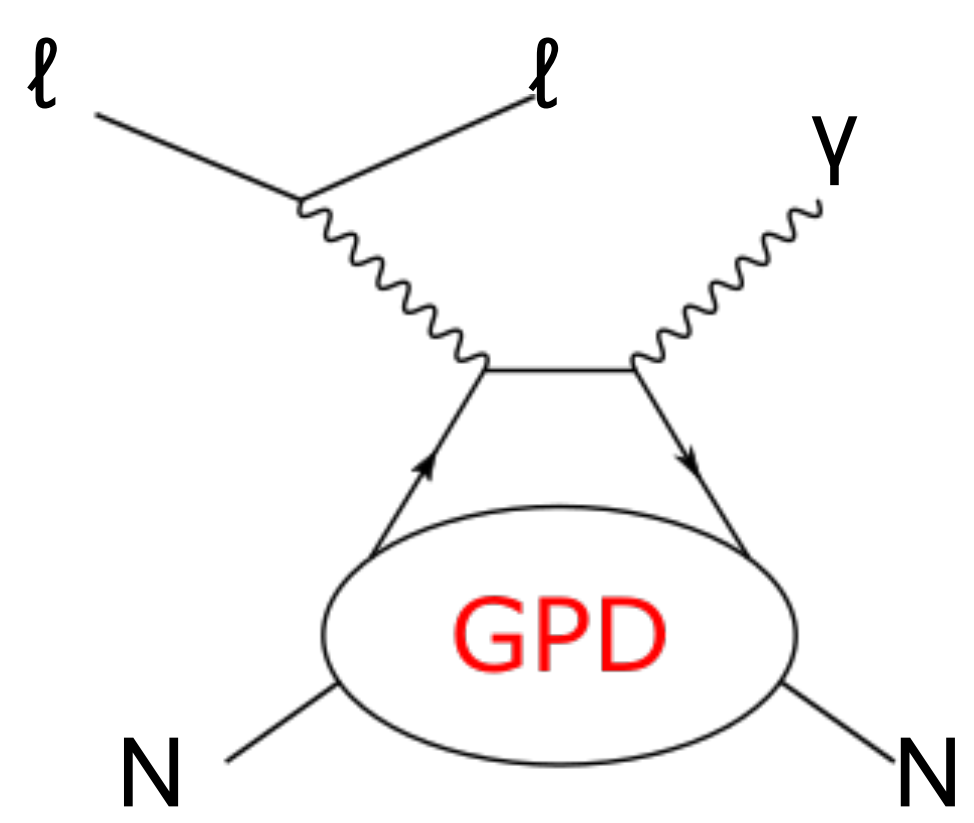
“Mechanical” forces acting on quarks, e.g. pressure in nucleon center

$$p(0) = \frac{1}{6\pi^2 M} \int_{-\infty}^0 dt \sqrt{-tt} C(t)$$

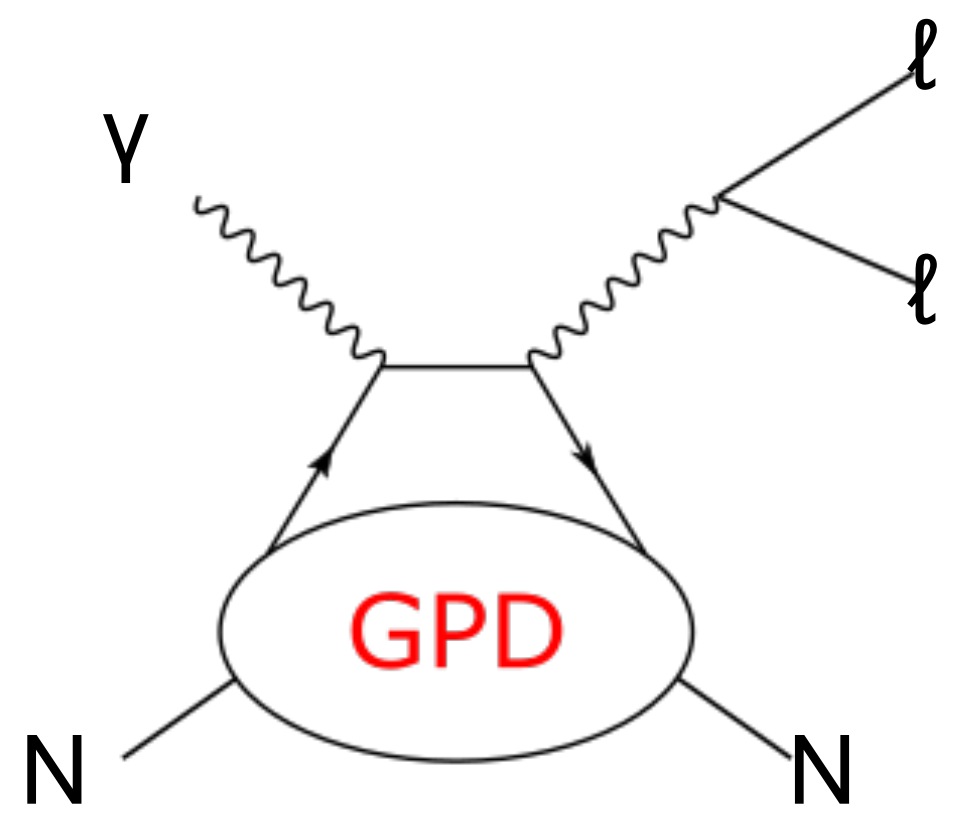




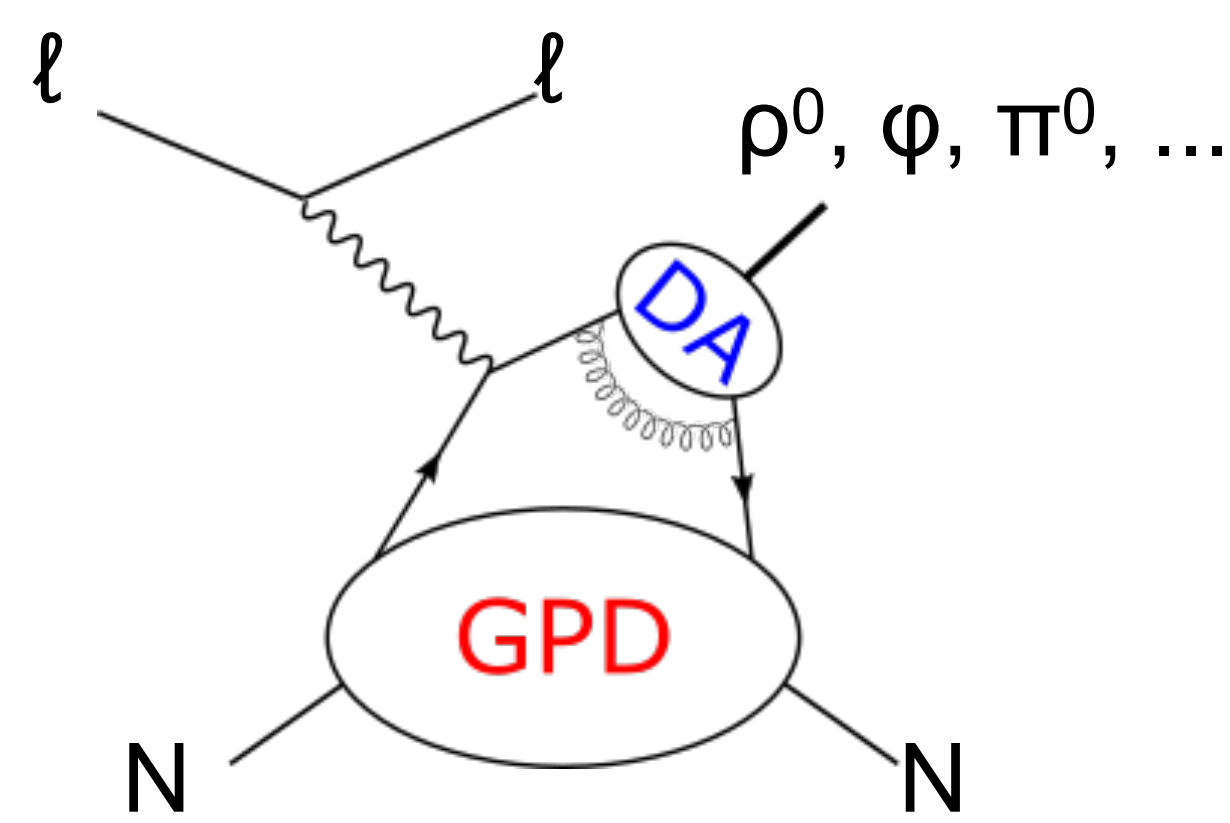
GPDs accessible in various production channels and observables
→ **experimental filters**



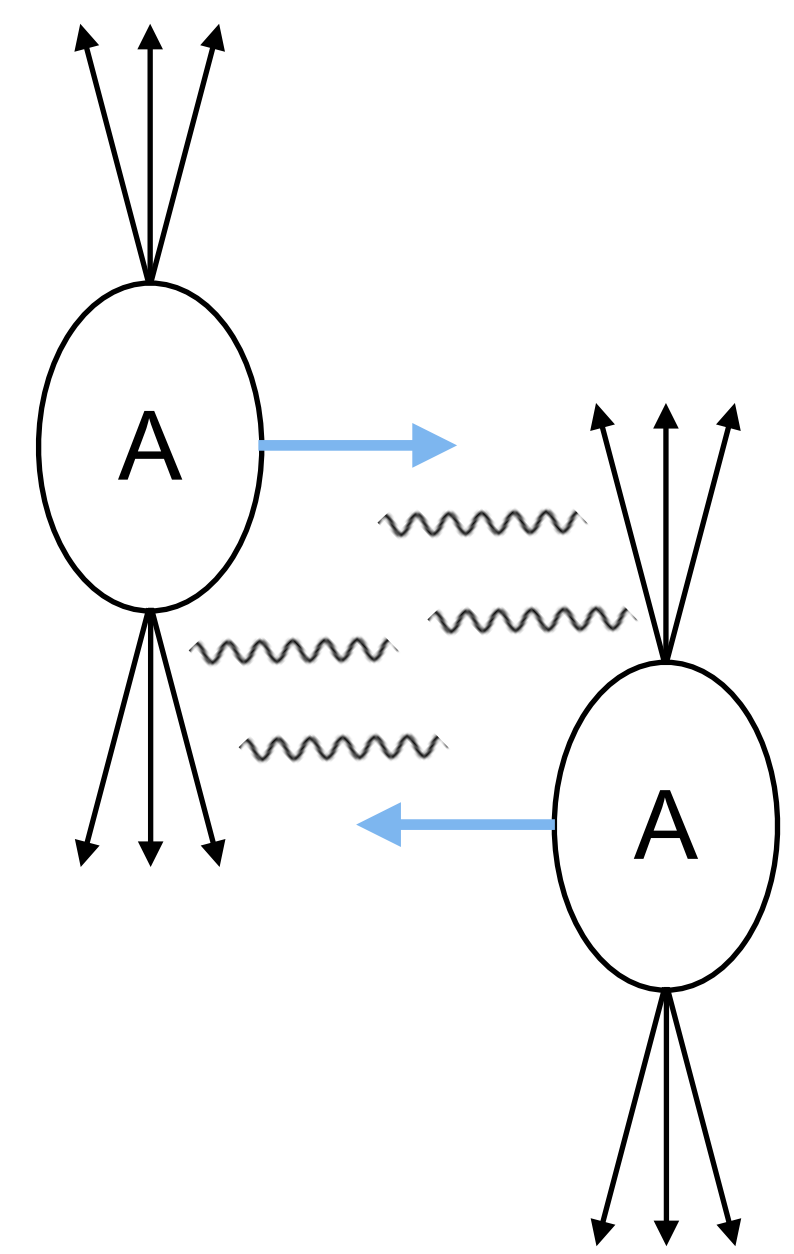
DVCS
Deeply Virtual Compton Scattering



TCS
Timelike Compton Scattering



HEMP
Hard Exclusive Meson Production



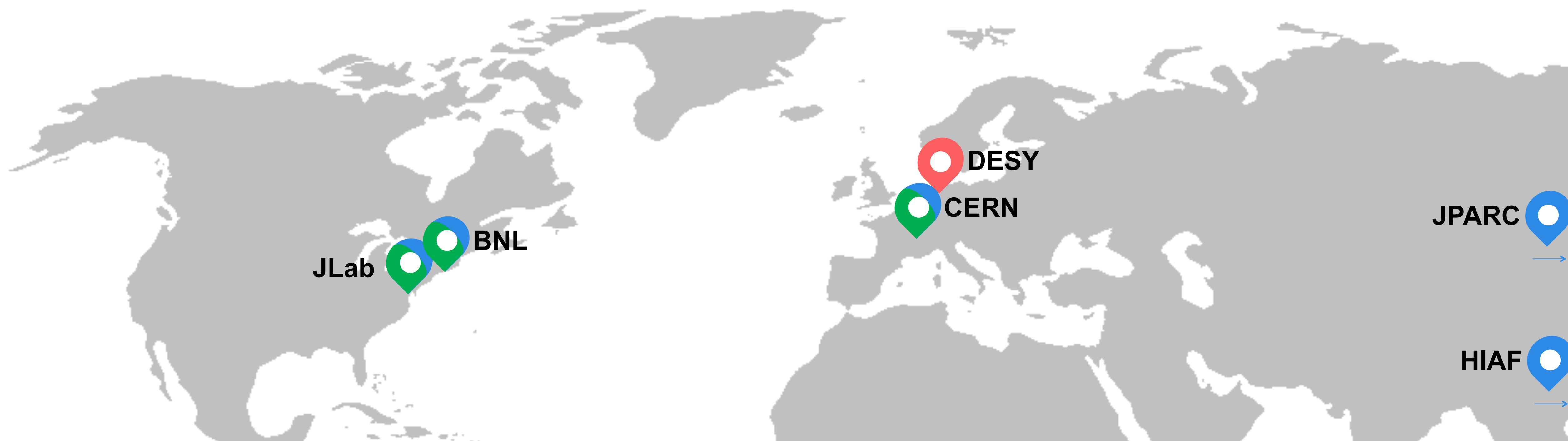
UPC
Ultra Peripheral Collisions

more production channels sensitive to GPDs exist!

GPDs studied in various laboratories
→ need to cover a broad kinematic range

experiments

closed active planned

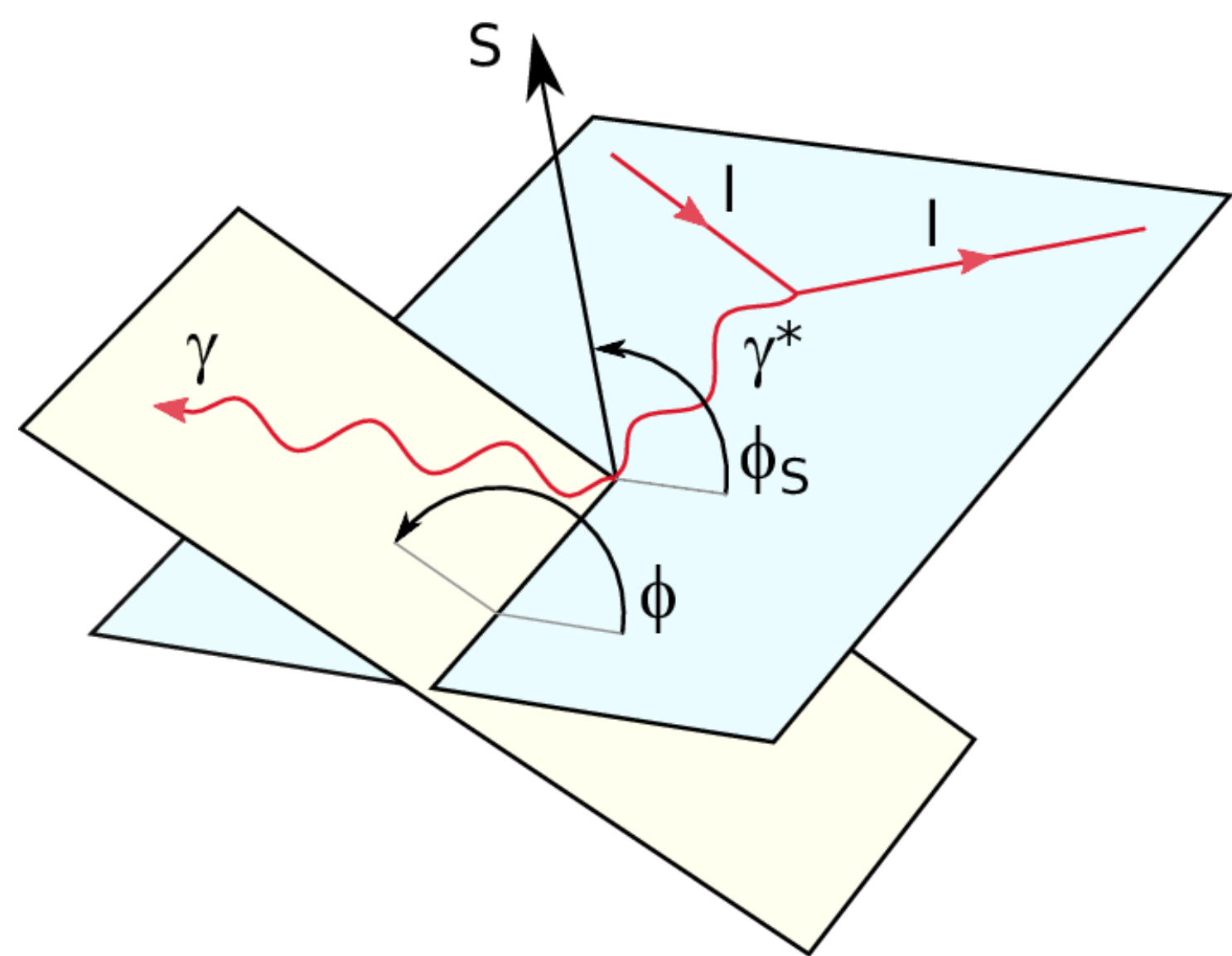


Kinematic cuts
used in our analyses:

$$Q^2 > 1.5 \text{ GeV}^2$$

$$-t/Q^2 < 0.2$$

Angles:

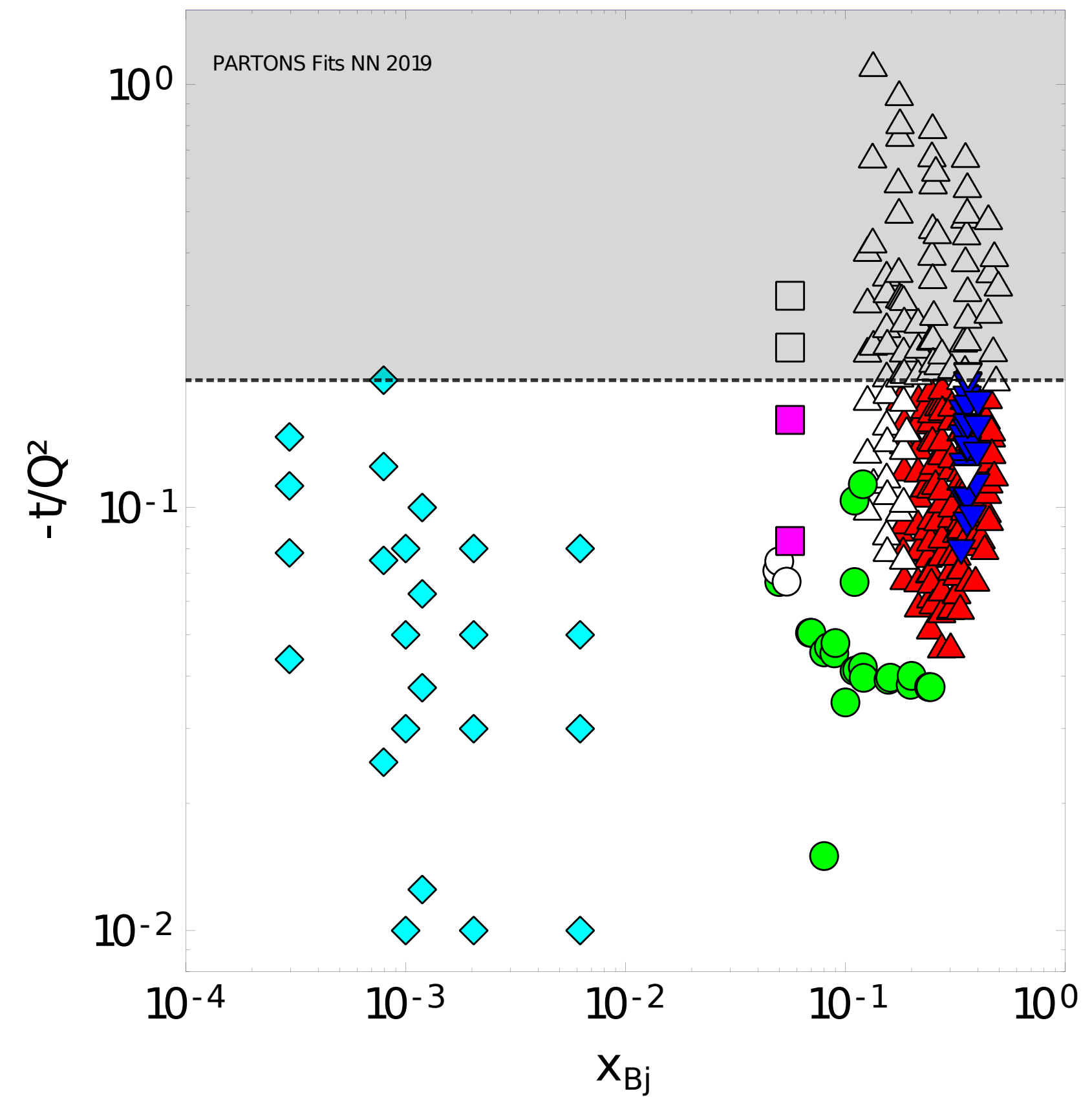
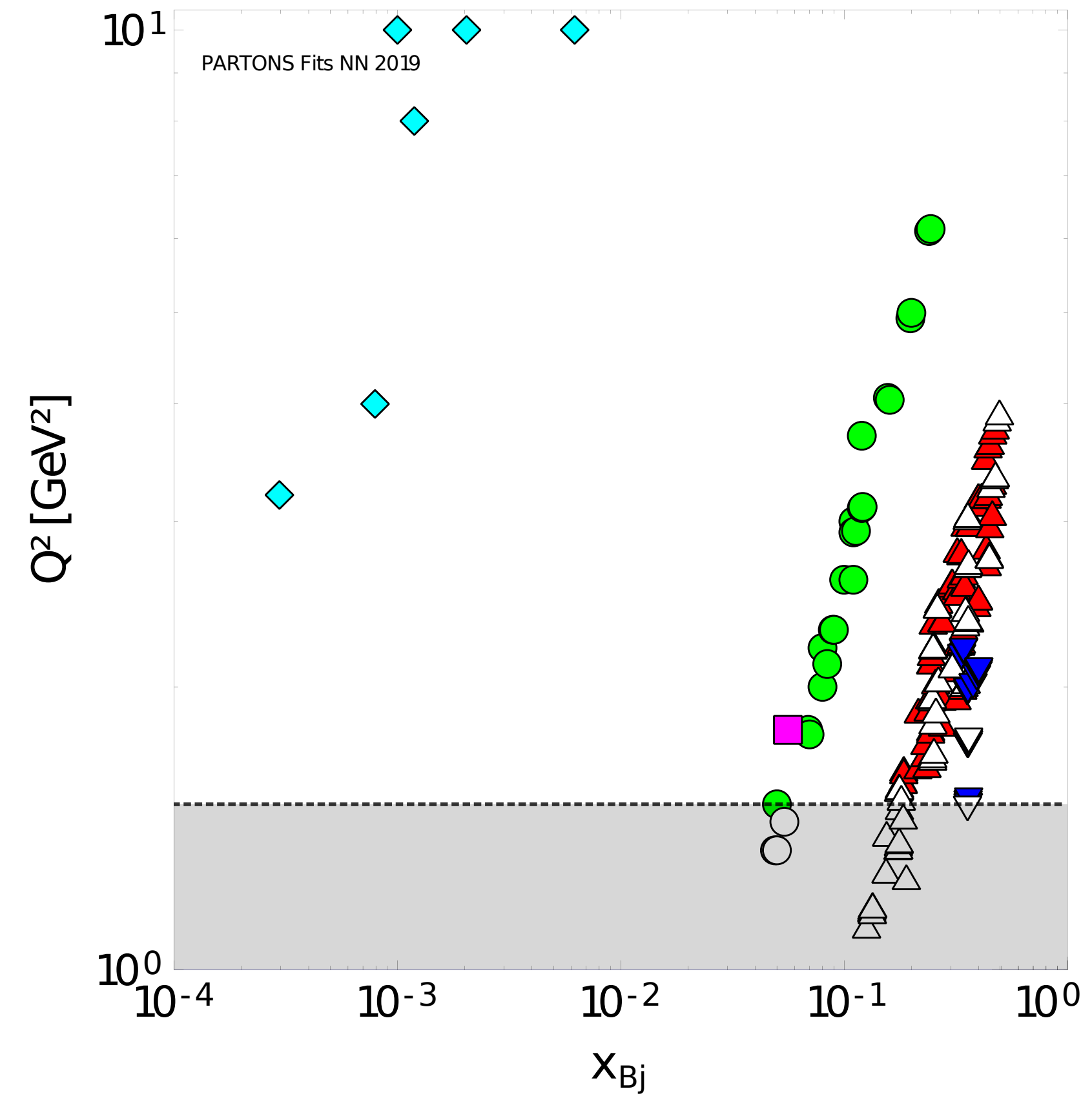


No.	Collab.	Year	Observable	Kinematic dependence	No. of points used / all
1	HERMES	2001	A_{LU}^+	ϕ	10 / 10
2		2006	$A_C^{\cos i\phi}$	t	4 / 4
3		2008	$A_C^{\cos i\phi}$	x_{Bj}	18 / 24
			$A_{UT,DVCS}^{\sin(\phi-\phi_S) \cos i\phi}$	$i = 0$	
			$A_{UT,I}^{\sin(\phi-\phi_S) \cos i\phi}$	$i = 0, 1$	
			$A_{UT,I}^{\cos(\phi-\phi_S) \sin i\phi}$	$i = 1$	
4		2009	$A_{LU,I}^{\sin i\phi}$	x_{Bj}	35 / 42
			$A_{LU,DVCS}^{\sin i\phi}$	$i = 1$	
			$A_C^{\cos i\phi}$	$i = 0, 1, 2, 3$	
5		2010	$A_{UL}^{+, \sin i\phi}$	x_{Bj}	18 / 24
			$A_{LL}^{+, \cos i\phi}$	$i = 0, 1, 2$	
6		2011	$A_{LT,DVCS}^{\cos(\phi-\phi_S) \cos i\phi}$	x_{Bj}	24 / 32
			$A_{LT,DVCS}^{\sin(\phi-\phi_S) \sin i\phi}$	$i = 1$	
			$A_{LT,I}^{\cos(\phi-\phi_S) \cos i\phi}$	$i = 0, 1, 2$	
			$A_{LT,I}^{\sin(\phi-\phi_S) \sin i\phi}$	$i = 1, 2$	
7		2012	$A_{LU,I}^{\sin i\phi}$	x_{Bj}	35 / 42
			$A_{LU,DVCS}^{\sin i\phi}$	$i = 1$	
			$A_C^{\cos i\phi}$	$i = 0, 1, 2, 3$	
8	CLAS	2001	$A_{LU}^{-, \sin i\phi}$	—	0 / 2
9		2006	$A_{UL}^{-, \sin i\phi}$	—	2 / 2
10		2008	A_{LU}^-	ϕ	283 / 737
11		2009	A_{LU}^-	ϕ	22 / 33
12		2015	$A_{LU}^-, A_{UL}^-, A_{LL}^-$	ϕ	311 / 497
13		2015	$d^4\sigma_{UU}^-$	ϕ	1333 / 1933
14	Hall A	2015	$\Delta d^4\sigma_{LU}^-$	ϕ	228 / 228
15		2017	$\Delta d^4\sigma_{LU}^-$	ϕ	276 / 358
16	COMPASS	2018	$d^3\sigma_{UU}^\pm$	t	2 / 4
17	ZEUS	2009	$d^3\sigma_{UU}^+$	t	4 / 4
18	H1	2005	$d^3\sigma_{UU}^+$	t	7 / 8
19		2009	$d^3\sigma_{UU}^\pm$	t	12 / 12
SUM:					2624 / 3996

Kinematic cuts
used in our analyses:

$$Q^2 > 1.5 \text{ GeV}^2$$
$$-t/Q^2 < 0.2$$

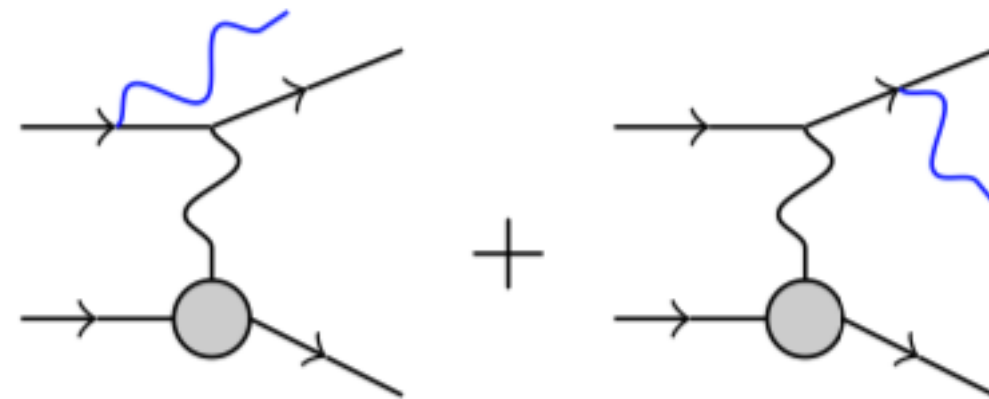
- ▼ HALL A
- ▲ CLAS
- HERMES
- COMPASS
- ◆ H1 and ZEUS



Cross-section for single photon production ($l + N \rightarrow l + N + \gamma$):

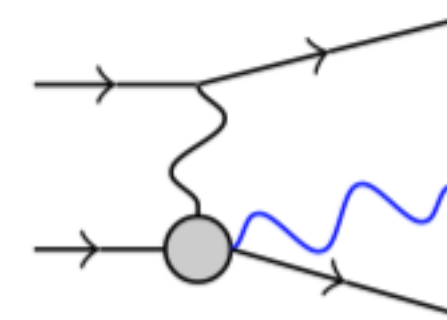
$$\sigma \propto |\mathcal{A}|^2 = |\mathcal{A}_{BH} + \mathcal{A}_{DVCS}|^2 = |\mathcal{A}_{BH}|^2 + |\mathcal{A}_{DVCS}|^2 + \mathcal{F}$$

Bethe-Heitler process



*calculable within QED
parametrised by elastic FFs*

DVCS



*calculable within QCD
parametrised by CFFs*

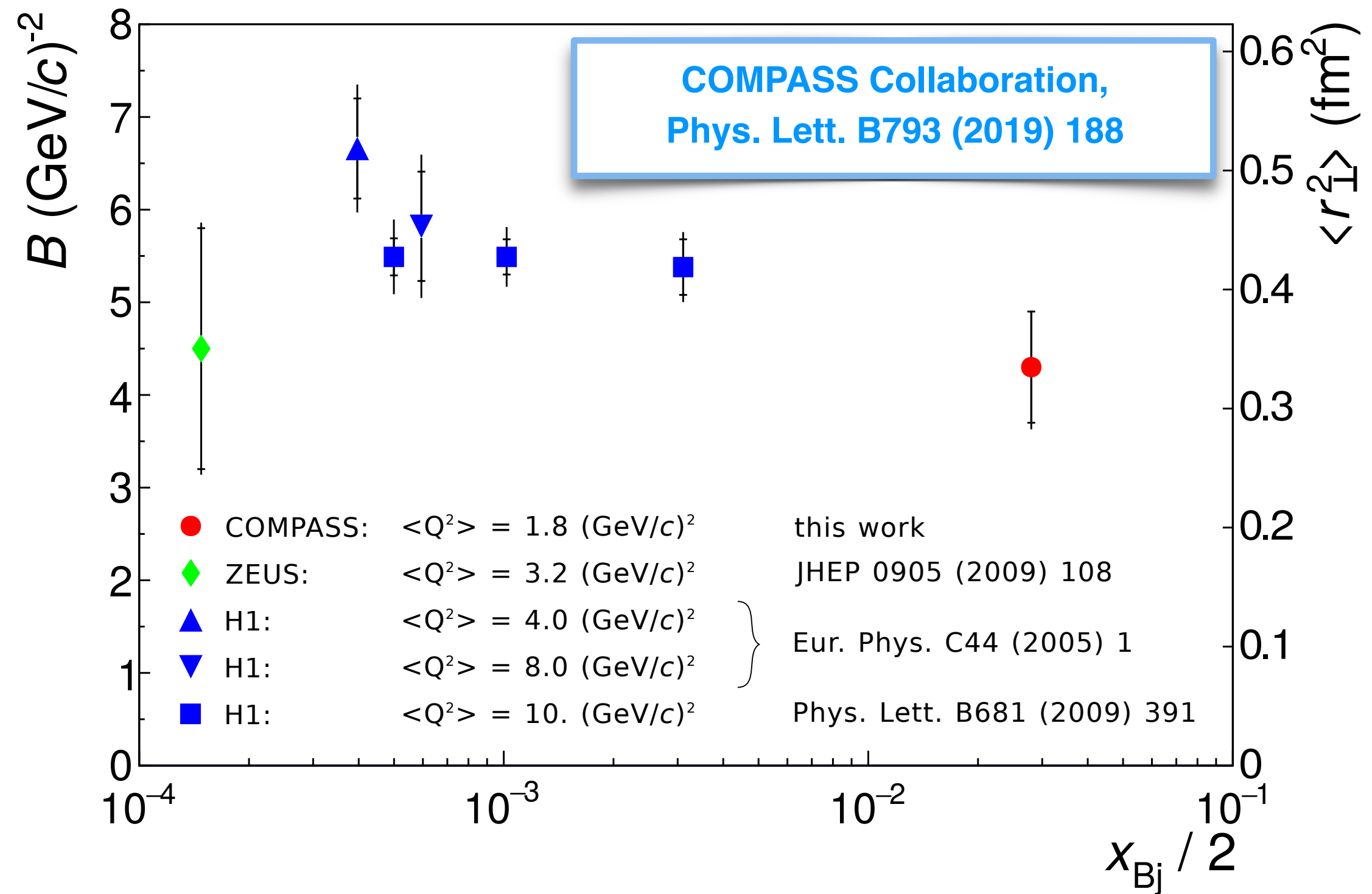
see e.g. NPB 878 (2014) 214
for more details

$$\text{Im}\mathcal{H}(\xi, t) \stackrel{\text{LO}}{=} \pi \sum_q e_q^2 H^{q(+)}(\xi, \xi, t)$$

$$\text{Re}\mathcal{H}(\xi, t) = \text{PV} \int_0^1 \frac{d\xi'}{\pi} \text{Im}\mathcal{H}(\xi', t) \left(\frac{1}{\xi - \xi'} - \frac{1}{\xi + \xi'} \right) + C_H(t)$$

What we can learn from DVCS amplitudes?

- nucleon tomography at low- x_B
- "mechanical" properties



Under following assumptions:

- single-exponential dependence
- dominance of CFF $\text{Im} \mathcal{M}$
- negligible "skewness effect"
 $H(x, x, t) \sim H(x, 0, t)$

From t-dependance of cross-section:

$$\frac{d\sigma^{\gamma^*p \rightarrow \gamma p}}{dt} \propto e^{-Bt}$$

related to transverse extension of quarks:

$$\langle r_{\perp}^2(x_{Bj}) \rangle \approx 2 \langle B(x_{Bj}) \rangle$$

What we can learn from DVCS amplitudes?

- nucleon tomography at low-xB
- "mechanical" properties

This is example of **model dependent analysis**, based on specific Ansatz of GPD

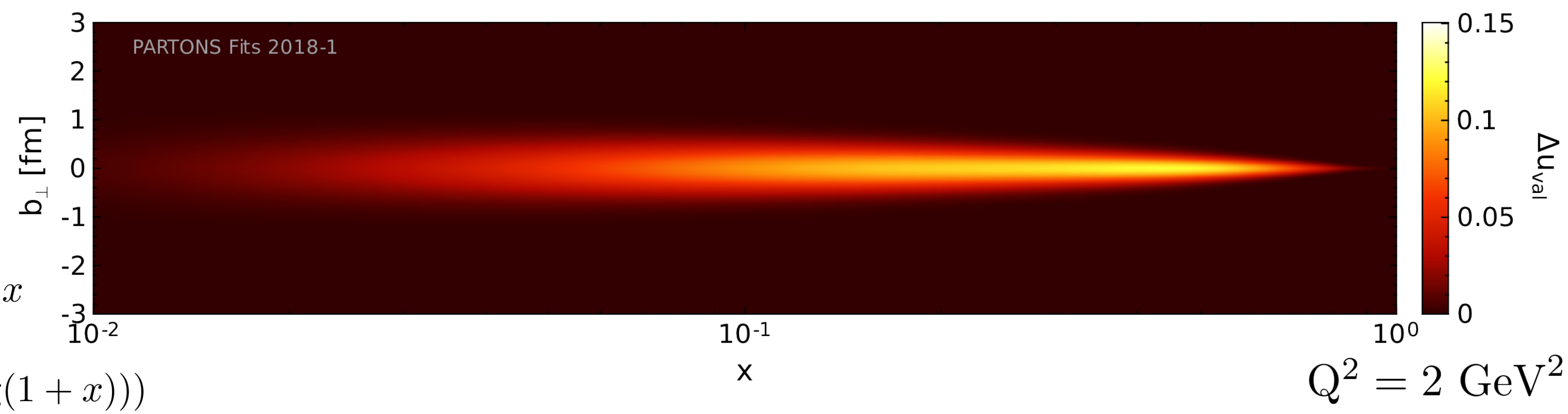
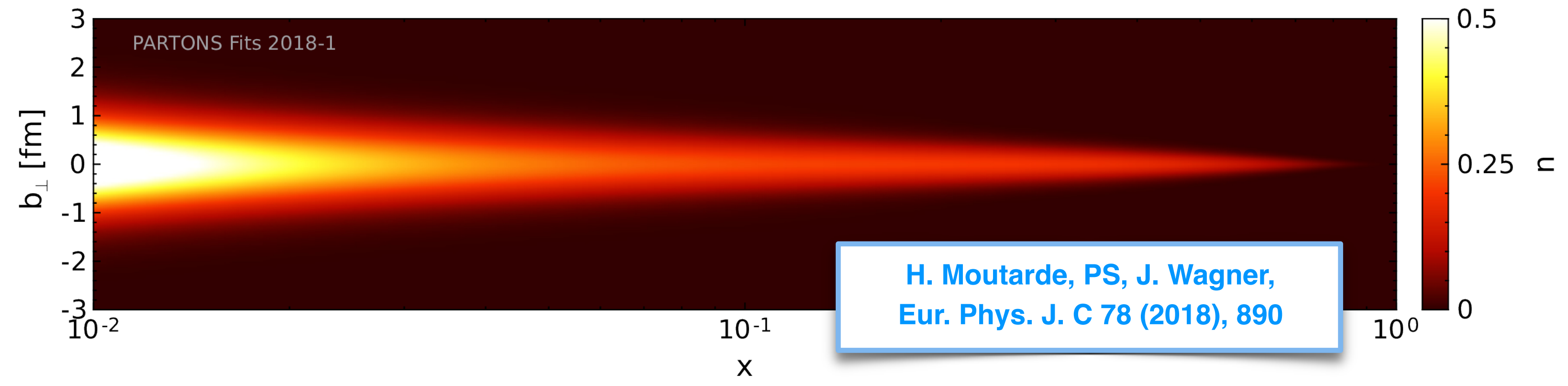
$$G^q(x, x, t) = G^q(x, 0, t) g_G^q(x, x, t)$$

$$G^q(x, 0, t) = \text{pdf}_G^q(x) \exp(f_G^q(x)t)$$

$$f_G^q(x) = A_G^q \log(1/x) + B_G^q (1-x)^2 + C_G^q (1-x)x$$

$$g_G^q(x, x, t) = \frac{a_G^q}{(1-x^2)^2} (1 + t(1-x)(b_G^q + c_G^q \log(1+x)))$$

 free parameters (for given GPD "G" and quark flavour "q") fixed by DVCS and elastic form factor data



What we can learn from DVCS amplitudes?

- nucleon tomography at low-xB
- "mechanical" properties

$$\mathcal{C}^q(t) \stackrel{\text{LO}}{=} 2 e_q^2 \int_{-1}^1 dz \frac{D^q(z, t)}{1-z} \equiv 4D^q(t) \quad z = \frac{x}{\xi}$$

$$D^q(z, t) = (1-z^2) \sum_{i=0}^{\infty} d_i^q C_{2i+1}^{3/2}(z)$$

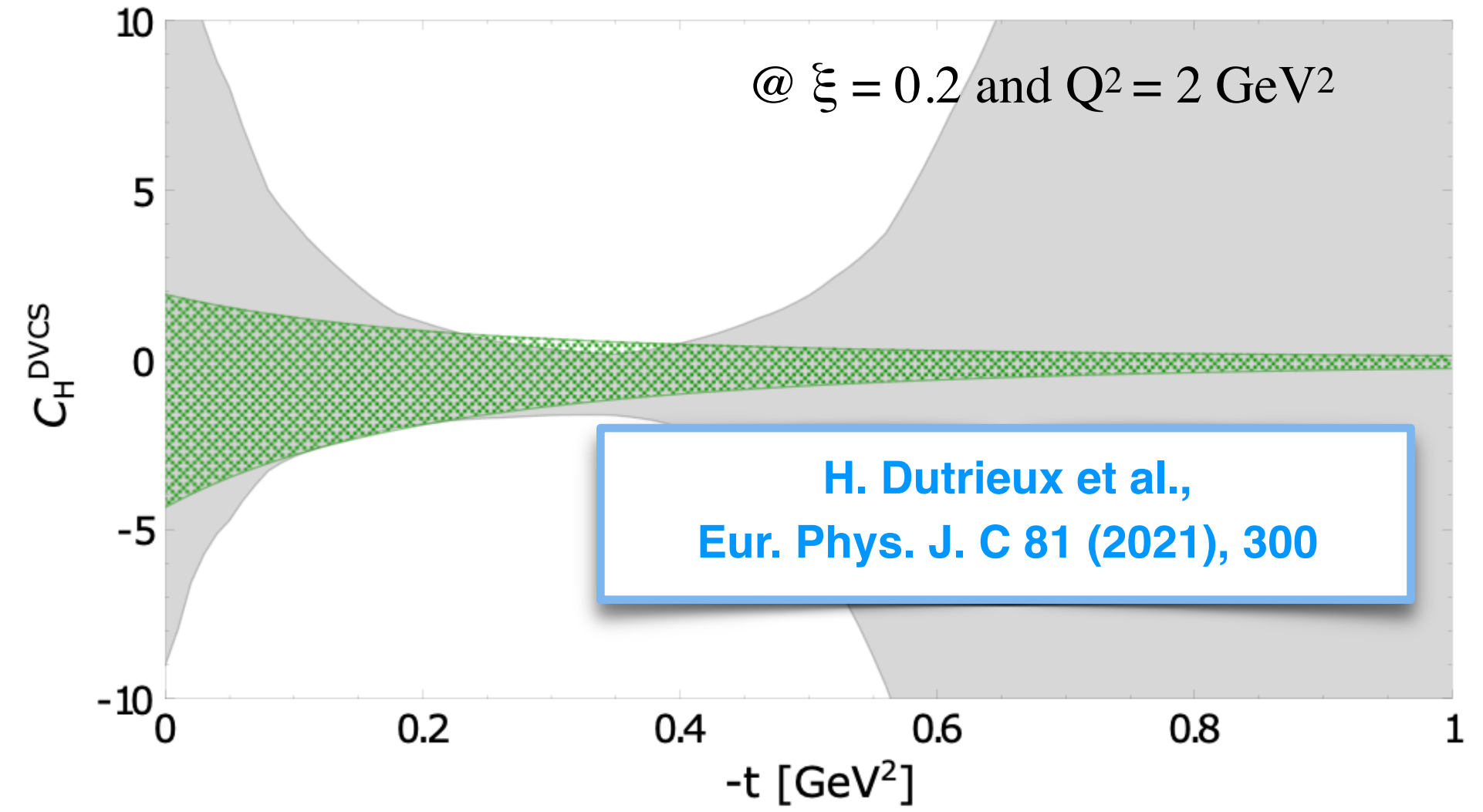
$$D^q(t) = \sum_{\substack{i=1 \\ \text{odd}}}^{\infty} d_i^q(t) \quad d_1^q(t) \equiv 5C_q(t)$$

ANN analysis

Model dependent extraction

$$d_1^{uds}(t, \mu_F^2) = d_1^{uds}(\mu_F^2) \left(1 - \frac{t}{\Lambda^2}\right)^{-\alpha} \quad \alpha = 3$$

$$\Lambda = 0.8 \text{ GeV}$$



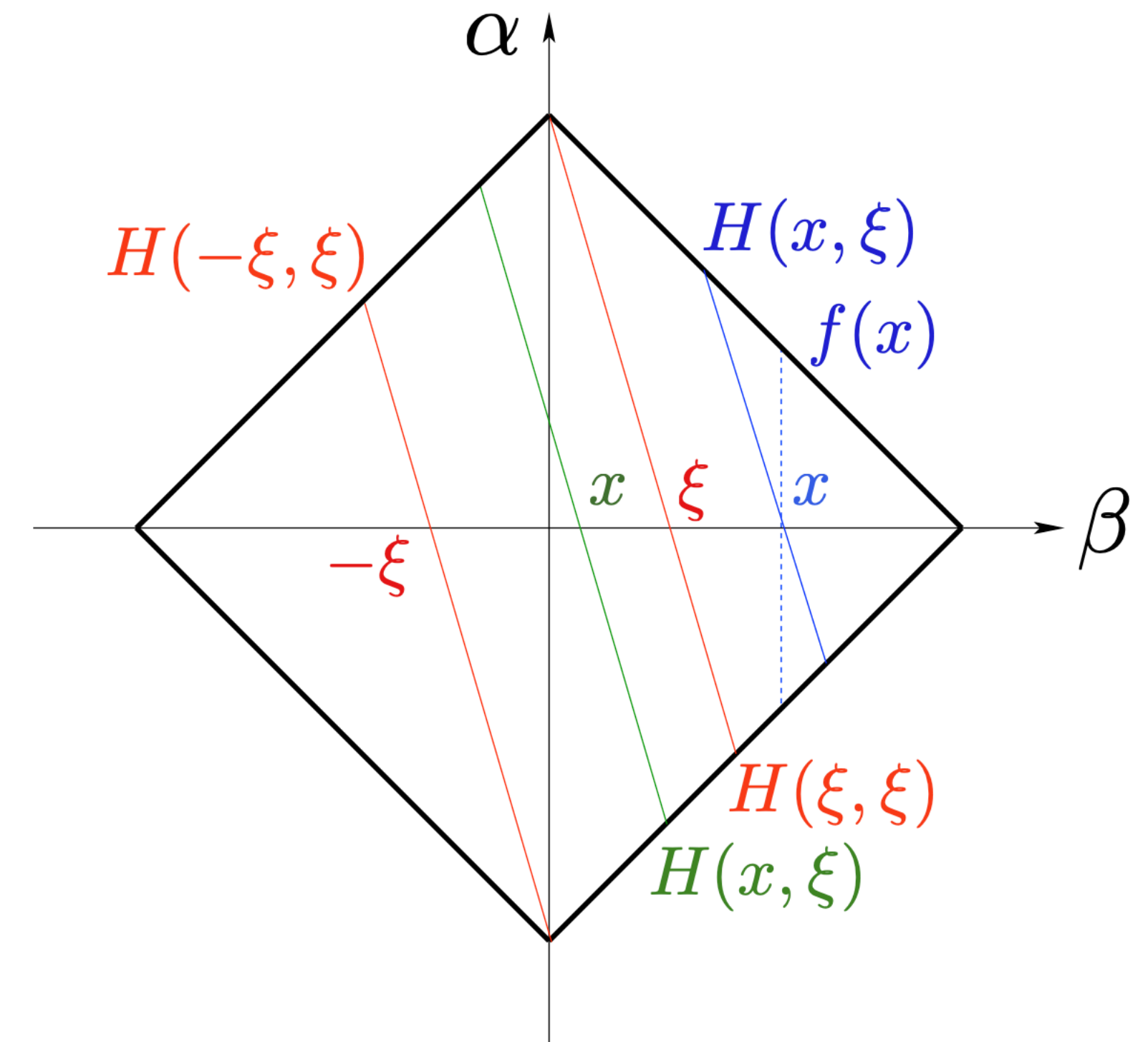
Double distribution:

$$H(x, \xi, t) = \int d\Omega F(\beta, \alpha, t)$$

where:

$$d\Omega = d\beta d\alpha \delta(x - \beta - \alpha\xi)$$

$$|\alpha| + |\beta| \leq 1$$



from PRD83, 076006, 2011

Double distribution:

$$(1 - x^2)F_C(\beta, \alpha) + (x^2 - \xi^2)F_S(\beta, \alpha) + \xi F_D(\beta, \alpha)$$

Classical term:

$$F_C(\beta, \alpha) = f(\beta)h_C(\beta, \alpha)\frac{1}{1 - \beta^2}$$

$$f(\beta) = \text{sgn}(\beta)q(|\beta|)$$

$$h_C(\beta, \alpha) = \frac{\text{ANN}_C(|\beta|, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha \text{ANN}_C(|\beta|, \alpha)}$$

Shadow term:

$$F_S(\beta, \alpha) = f(\beta)h_S(\beta, \alpha)$$

$$f(\beta) = \text{sgn}(\beta)q(|\beta|)$$

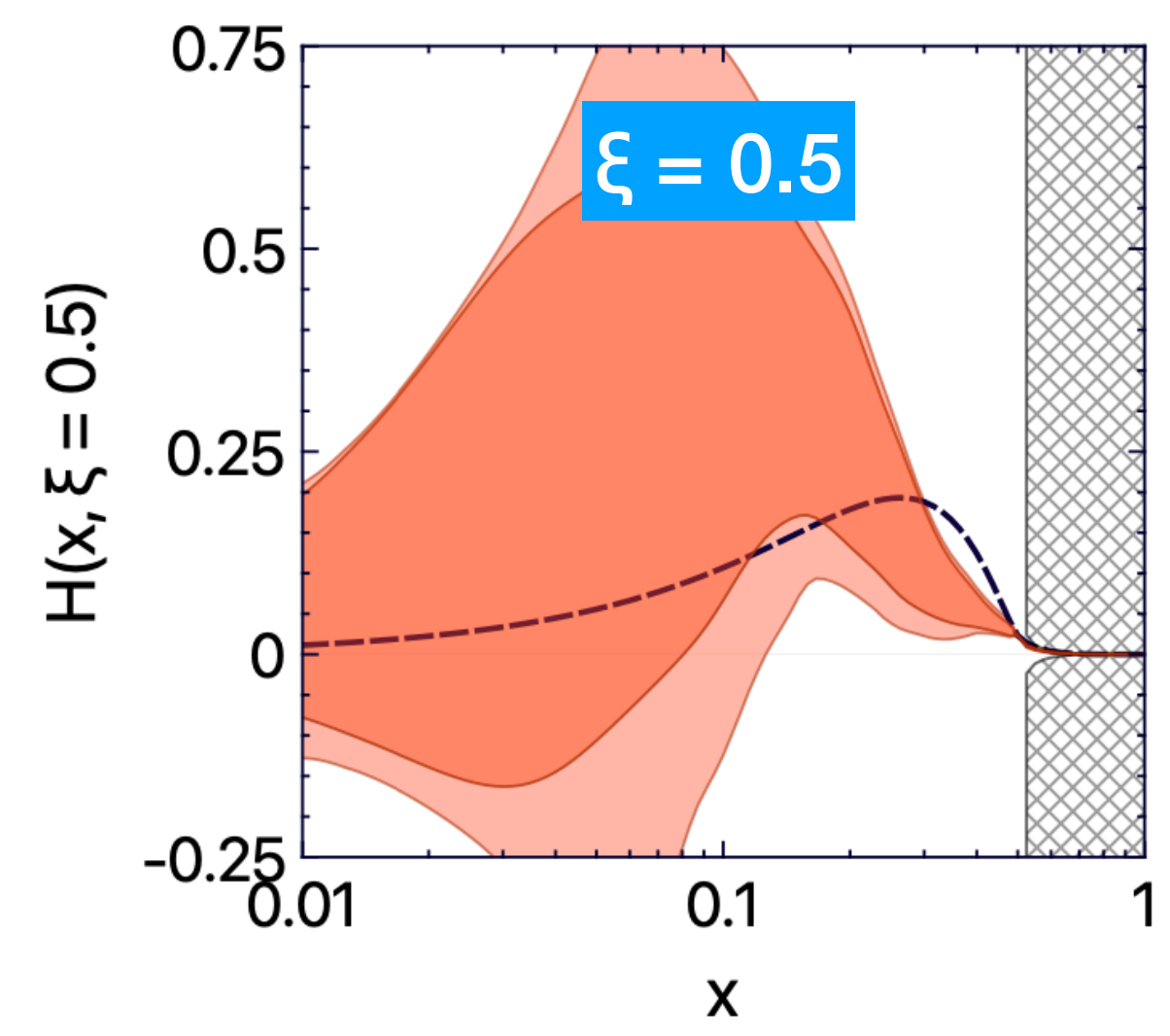
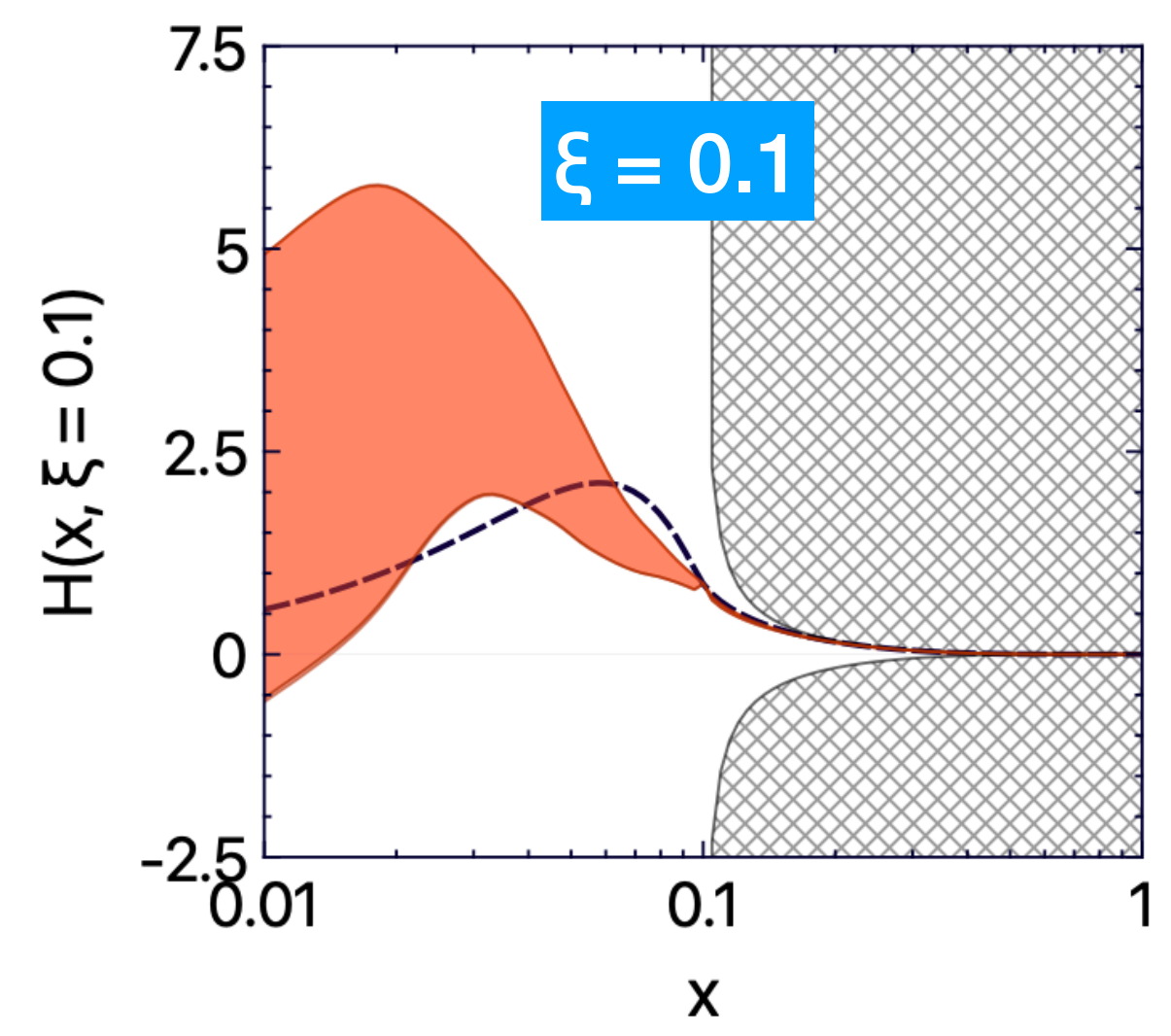
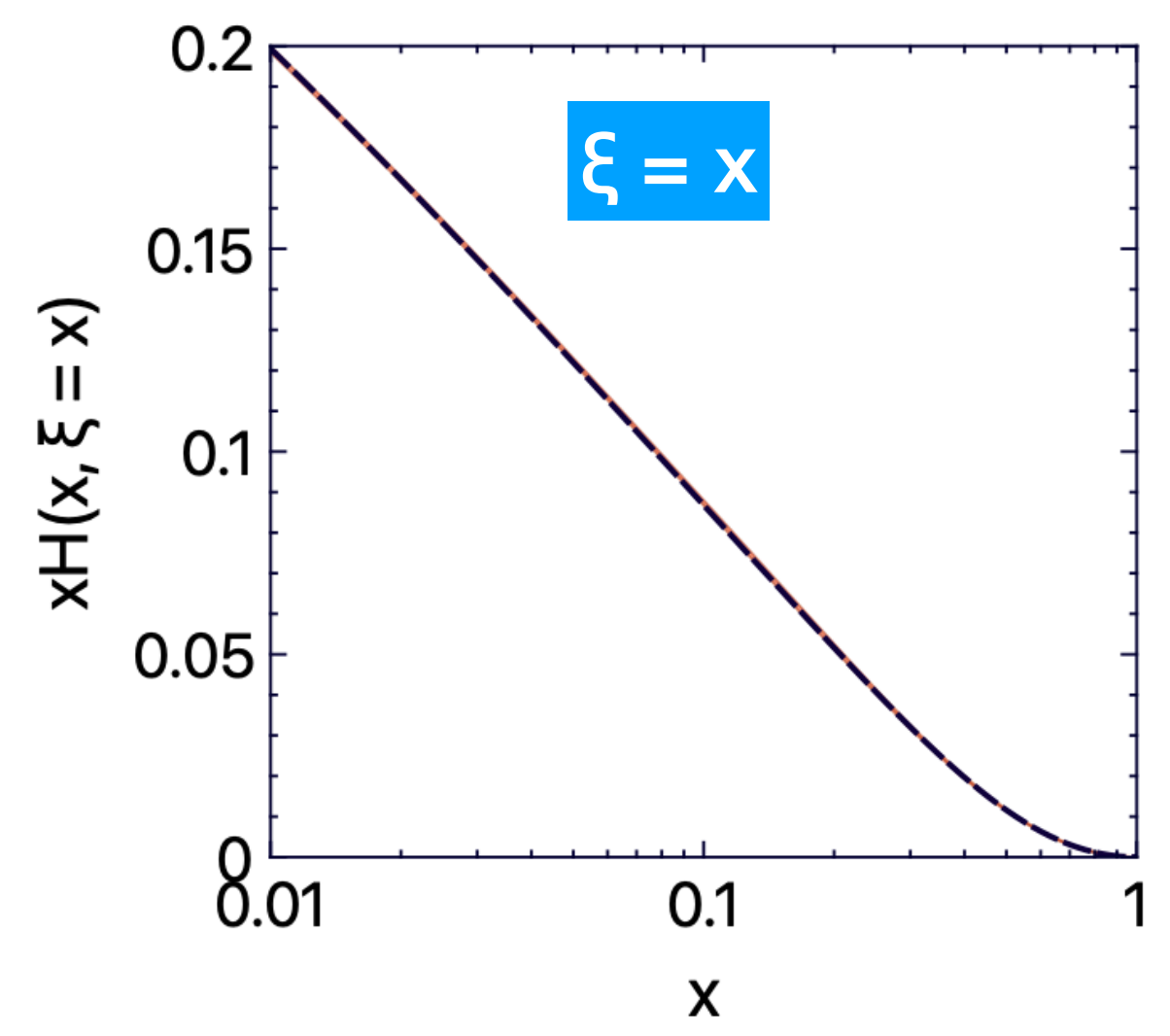
$$h_S(\beta, \alpha)/N_S = \frac{\text{ANN}_S(|\beta|, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha \text{ANN}_S(|\beta|, \alpha)} \cdot \frac{\text{ANN}_{S'}(|\beta|, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha \text{ANN}_{S'}(|\beta|, \alpha)}$$

$$\text{ANN}_{S'}(|\beta|, \alpha) \equiv \text{ANN}_C(|\beta|, \alpha)$$

D-term:

$$F_D(\beta, \alpha) = \delta(\beta)D(\alpha)$$

$$D(\alpha) = (1 - \alpha^2) \sum_{\substack{i=1 \\ \text{odd}}} d_i C_i^{3/2}(\alpha)$$



Conditions:

- Input: 200 $x = \xi$ points generated with GK model
- Positivity **forced**

GK
 ANN model 68% CL F_C
 Excluded by positivity
 ANN model 68% CL $F_C + F_S$

- The process allows to directly probe GPDs outside $x=\xi$ line, but is much more challenging experimentally

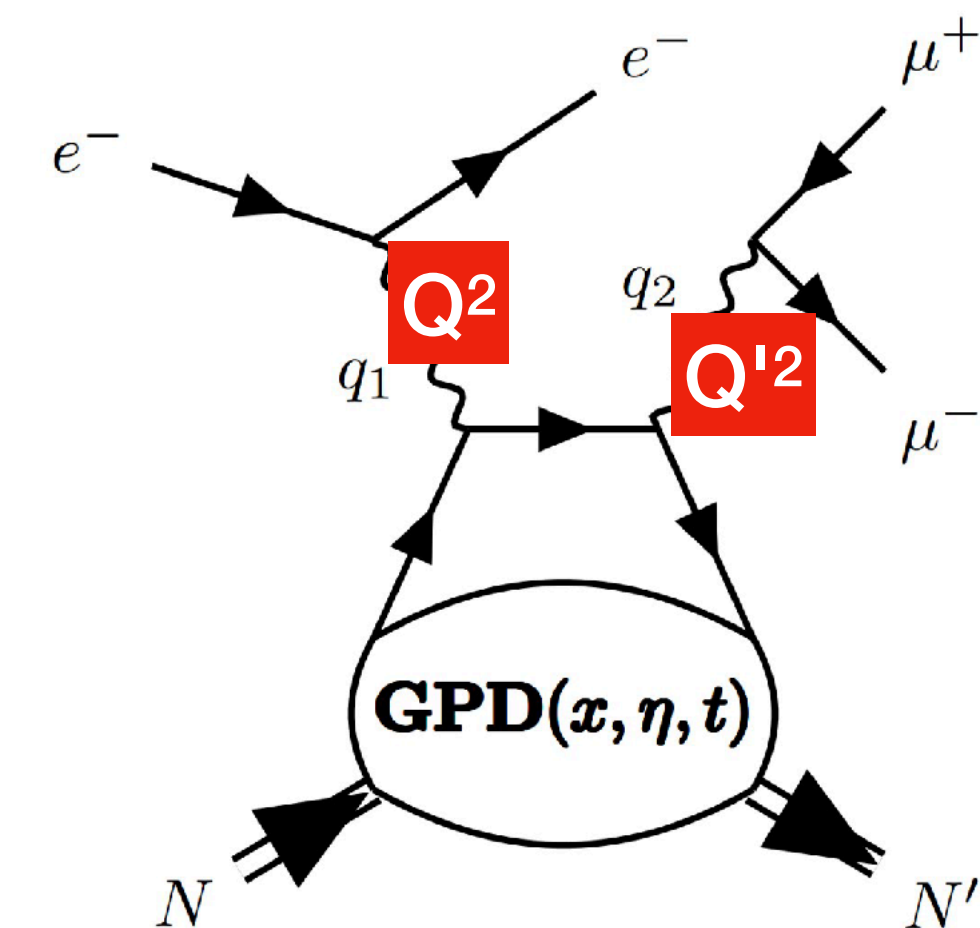
$$(\mathcal{H}, \mathcal{E})(\rho, \xi, t) = \sum_{f=\{u,d,s\}} \int_{-1}^1 dx C_f^{(-)}(x, \rho)(H_f, E_f)(x, \xi, t)$$

$$C_f^{(\pm)}(x, \rho) \stackrel{LO}{=} \left(\frac{e_f}{e}\right)^2 \left(\frac{1}{\rho - x - i0} \pm \frac{1}{\rho + x - i0} \right)$$

- We revisit DDVCS phenomenology in view of new experiments, including reevaluation of DDVCS and BH cross-sections with Kleiss-Stirling spinor techniques
- Obtained results are available in PARTONS and EpIC MC generator

$$\xi = \frac{Q^2 + Q'^2}{2Q^2/x_B - Q^2 - Q'^2}$$

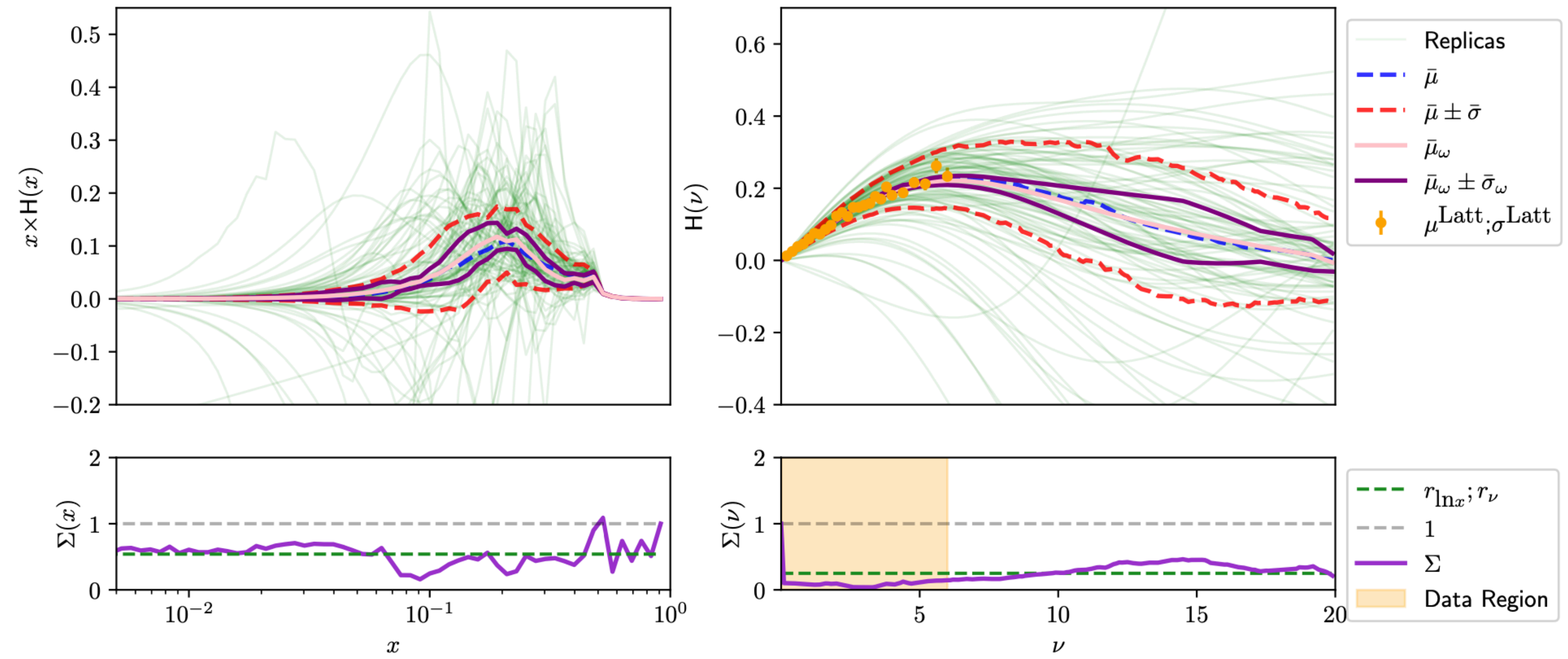
$$\rho = \xi \frac{Q^2 - Q'^2}{Q^2 + Q'^2}$$



M. J. Riberdy, H. Dutrieux, C. Mezrag, PS,
 hep-ph/2306.01647

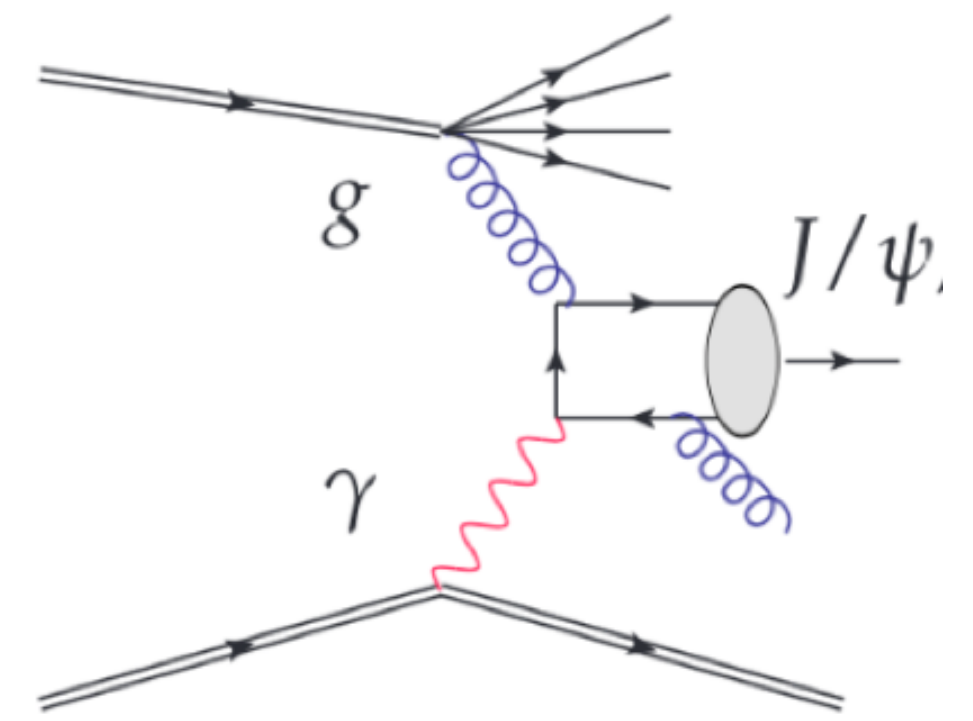
- Exploratory study to include lattice-QCD results!

Reduction of GPD model uncertainties due to inclusion of pseudo-latticeQCD results

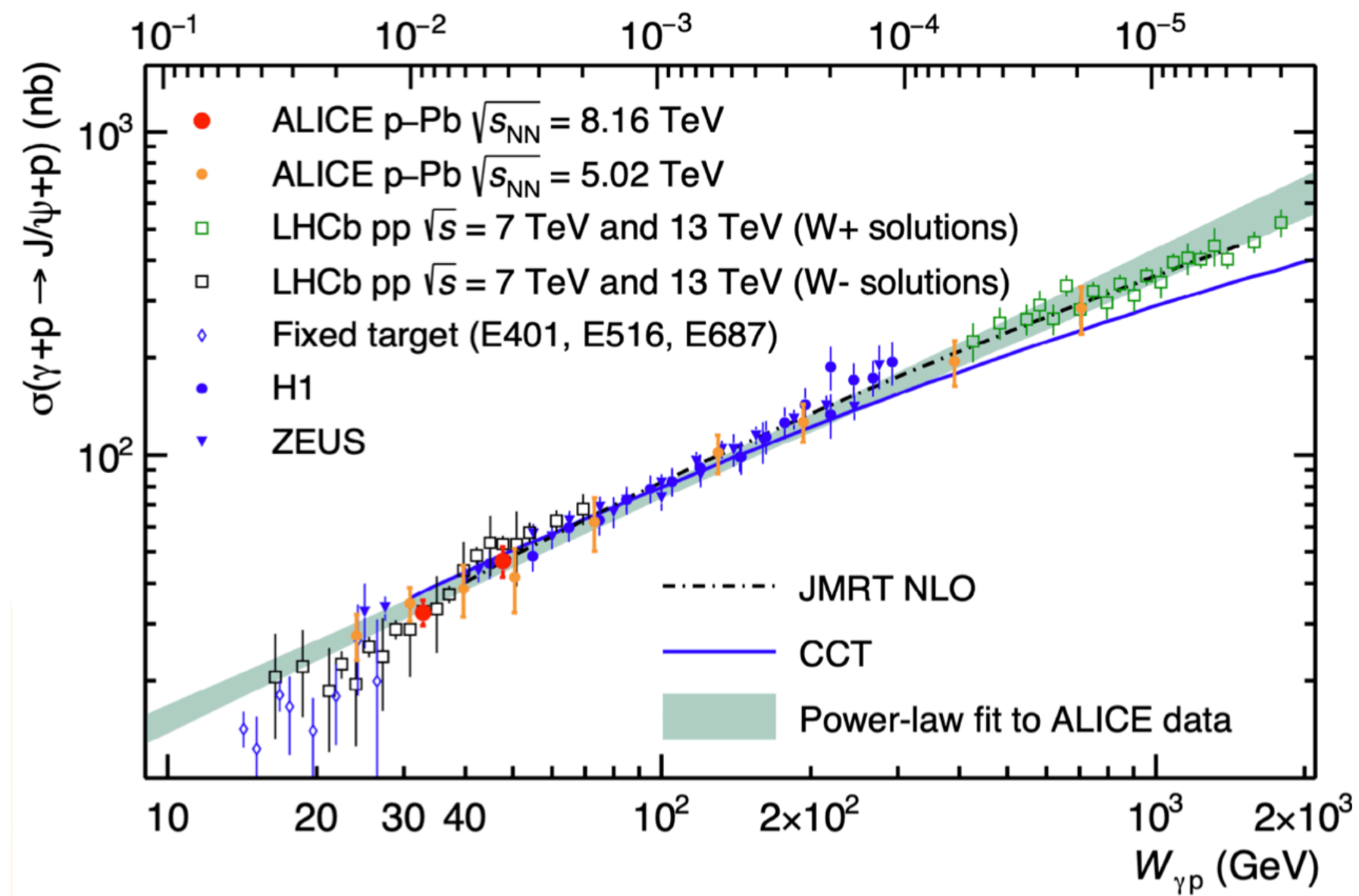
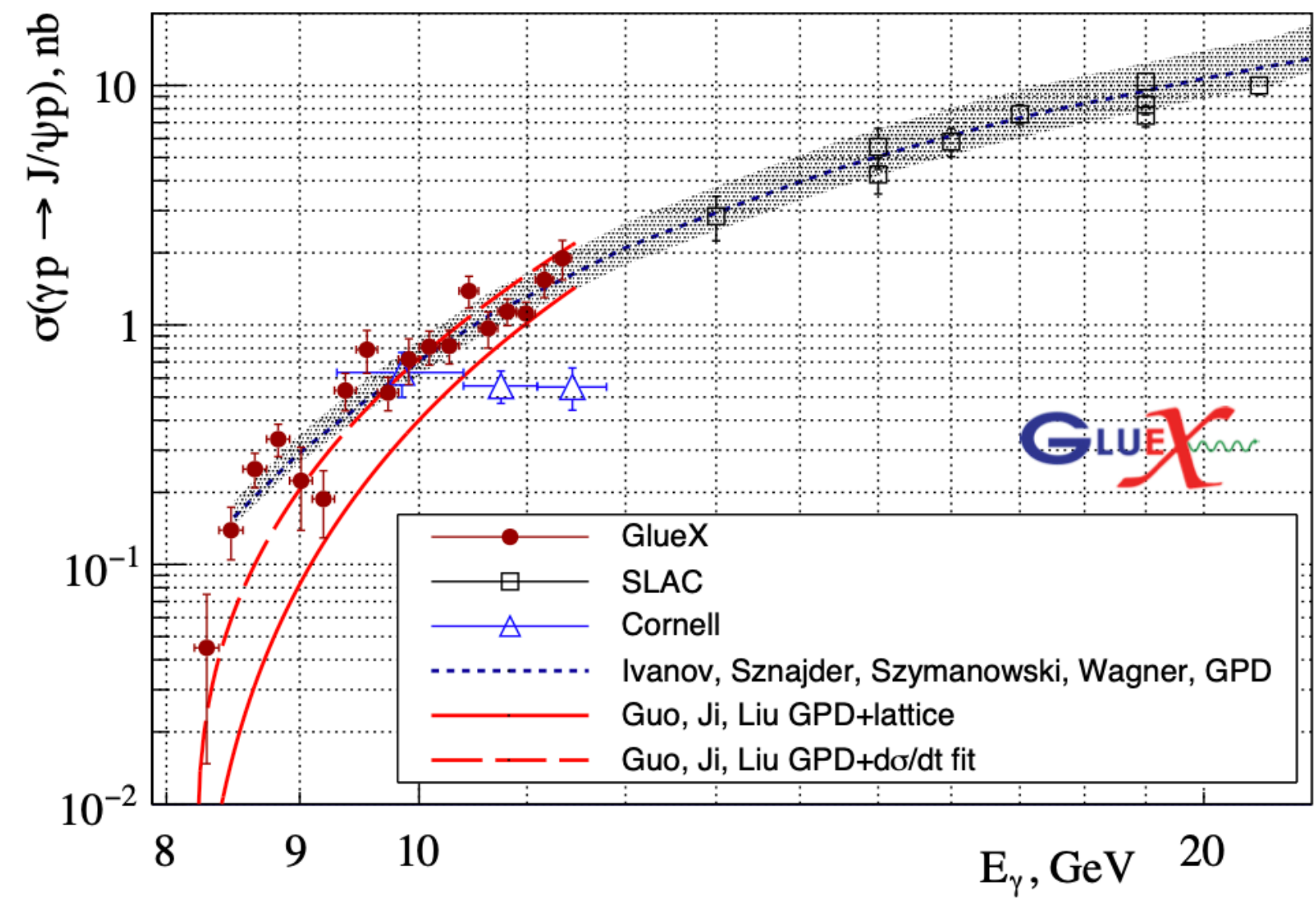


- The energy frontier for electromagnetic probes -> 10 times higher CM energy than available in HERA
- Allows to probe x_Bj down to a few 10^{-6} at moderate Q^2
- Electromagnetic probes have $a_{EM} \sim 1/137$, so are less affected by multiple interactions than hadronic interactions
- Also important source of PDF information

H. Dutrieux, M. Winn, V. Bertone,
Phys. Rev. D 107 (2023) 11, 114019



Exclusive photoproduction of J/Psi



- PARTONS - open-source framework to study GPDs
→ <http://partons.cea.fr>
- Come with number of available physics developments implemented
- Written in C++, also available via virtual machines (VirtualBox) and containers (Docker)
- Addition of new developments as easy as possible
- Developed to support effort of GPD community,
can be used by both theorists and experimentalists
- v3 version of PARTONS is now available!



- Novel MC generator called EpIC released
→ <https://pawelsznajder.github.io/epic>
- EpIC is based on PARTONS
- EpIC is characterised by:
 - flexible architecture that utilises a modular programming paradigm
 - a variety of modelling options, including radiative corrections
 - multichannel capability (now: DVCS, TCS, $DV\pi^0P$, diphoton; coming soon: DDVCS, J/ψ)
- This is the new tool to be use in the precision era commenced by the new generation of experiments



- Substantial progress in:
 - understanding of fundamental problems, like deconvolution of CFFs, and analysis methods
→ important for extraction of GPDs
 - modelling of GPD, fulfilling all theory-driven constraints (including positivity)
→ subject not touched enough in the current literature
→ developed in mind for easy inclusion of latticeQCD data
 - addressing the long-standing problem of model dependency of GPDs
→ nontrivial and timely analysis
 - description of exclusive processes
→ new sources of GPD information
 - delivering open-source tools for the community
→ to support both experimentalists and theoreticians

This progress is important for the precision era of GPD extraction allowed by the new generation of experiments