

Classifying the CP properties of the ggH coupling in H+2j production

WG2 / WG3 joint meeting

26.09.2023

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[2309.03146]

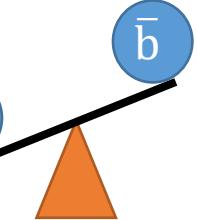


Overview & Motivation

The need for CP violation:

- CP violation is necessary to explain the baryon asymmetry of the universe
- SM does not contain enough CP violation for this
- Current experimental limits leave room for CP violation in the Higgs sector

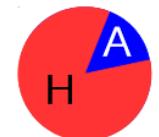
[CMS '21](#), [ATLAS '21](#), [ATLAS '22](#), [CMS '22](#)



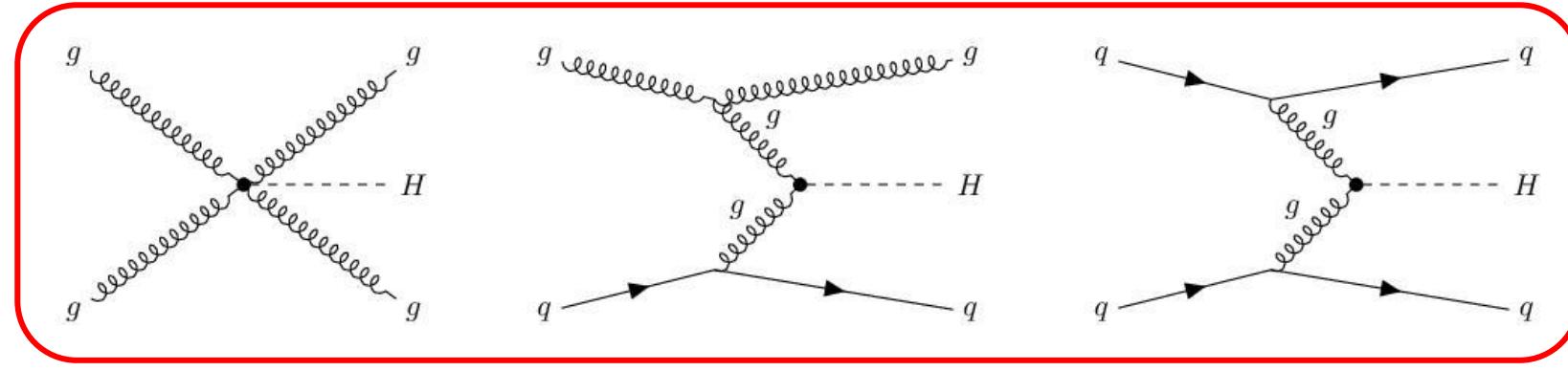
Why use ggF + 2j as a CP test?

- Gluon fusion is the main Higgs production mode at the LHC \Rightarrow lots of events
- 2 jets in the final state are needed to construct CP-odd observables
- More direct probe of the top-Yukawa coupling than simple rate information

[Klamke & Zeppenfeld '07](#), [Demartin et al. '14](#)

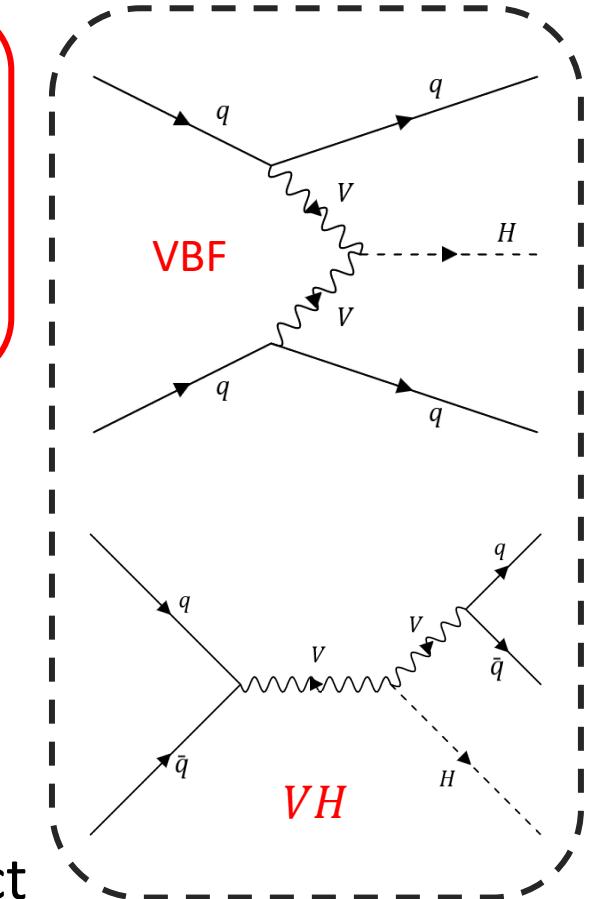


Signal & Background processes



$$\frac{\sigma_{\text{init}}}{\sigma_{\text{ggF2j}}} = \begin{array}{ll} 72\% & \\ \end{array} \quad \begin{array}{ll} 26.5\% & \\ \end{array} \quad \begin{array}{ll} 1.5\% & @13\text{TeV} \\ \end{array}$$

- **ggF2j** shows different topologies based on the initial state
- High CP sensitivity in $q\bar{q}$ initial states? [Hagiwara et al. '09](#)
- Considered background (BG) processes: VBF and VH
- $H \rightarrow \gamma\gamma$ decay: non-Higgs BG smoothly falling, easy to subtract
- Many analyses consider only signal region with VBF-like kinematics [CMS '22](#)



BSM framework

Free parameters:

- Higgs characterisation model: Higgs H assumed to be mixed CP state
- Effective Higgs-gluon coupling:

Artoisenet et al. '13

$$\mathcal{L}_{ggH} = -\frac{1}{4\nu} \left(-\frac{\alpha_s}{3\pi} \textcolor{blue}{c_g} G_{\mu\nu}^a G^{\mu\nu,a} + \frac{\alpha_s}{2\pi} \tilde{c}_g G_{\mu\nu}^a \tilde{G}^{\mu\nu,a} \right) H$$

- Effective CP-even ($\textcolor{blue}{c_g}$) and CP-odd (\tilde{c}_g) coupling modifiers
- SM obtained for $c_g = 1, \tilde{c}_g = 0$
- Higgs-gluon coupling corresponds to top-Yukawa in the heavy top limit and if there are no low-mass BSM particles in the ggF loop $\Rightarrow c_g = c_t, \tilde{c}_g = \tilde{c}_t$
- We impose a cut $p_T^H < 200\text{GeV}$ to remain in the heavy top limit

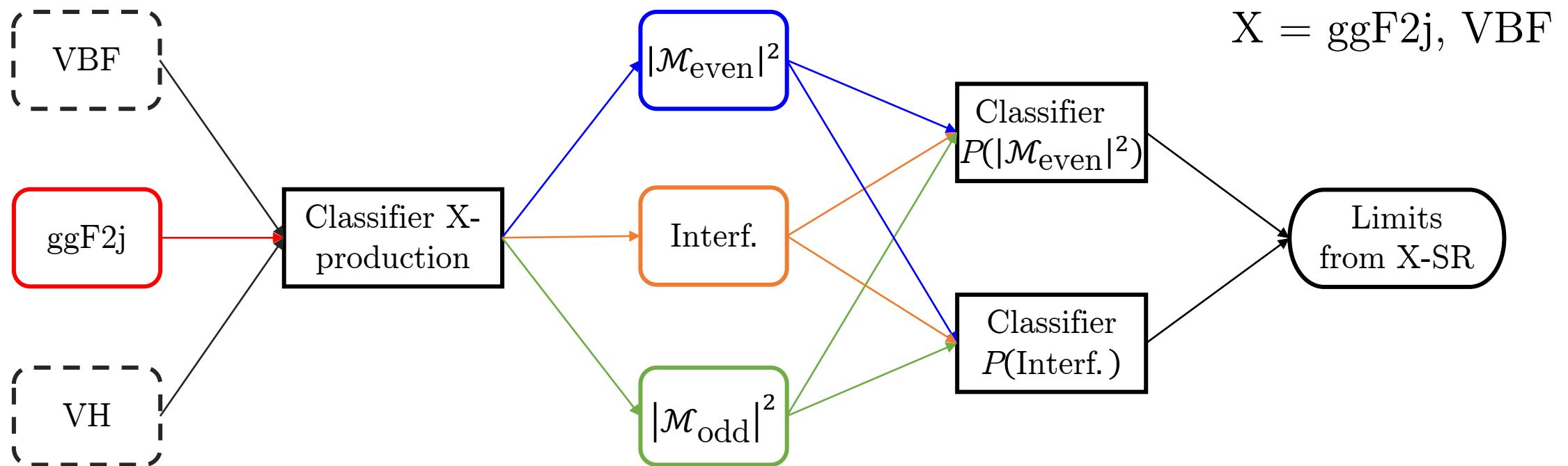
Analysis strategy

$$|\mathcal{M}_{\text{ggF2j}}|^2 = c_g^2 |\mathcal{M}_{\text{even}}|^2 + \underbrace{2c_g \tilde{c}_g \text{Re}[\mathcal{M}_{\text{even}} \mathcal{M}_{\text{odd}}^*]}_{\text{Interference}} + \tilde{c}_g^2 |\mathcal{M}_{\text{odd}}|^2$$



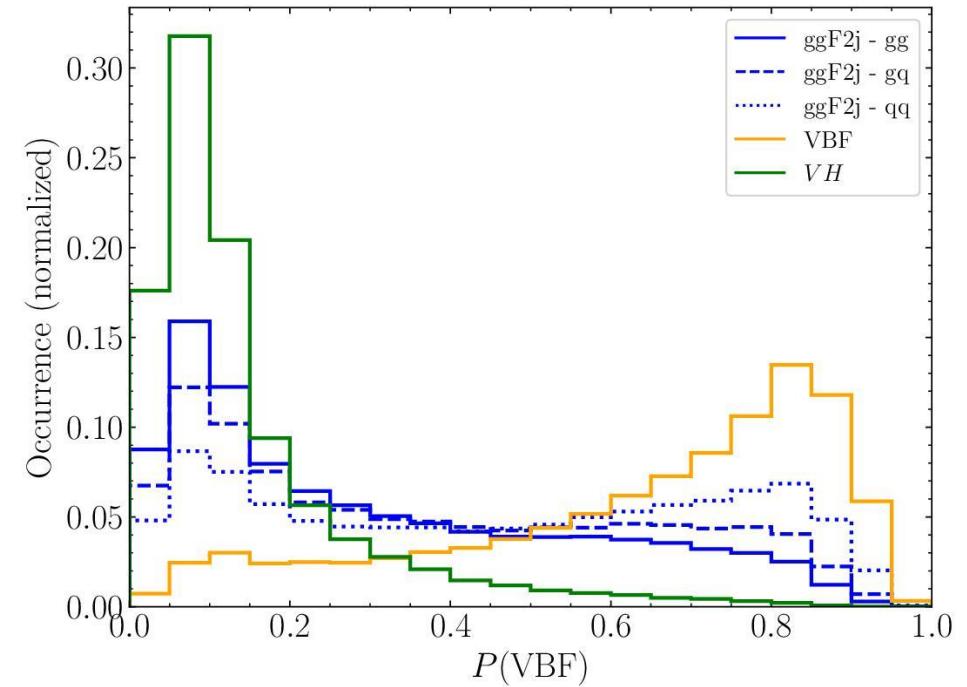
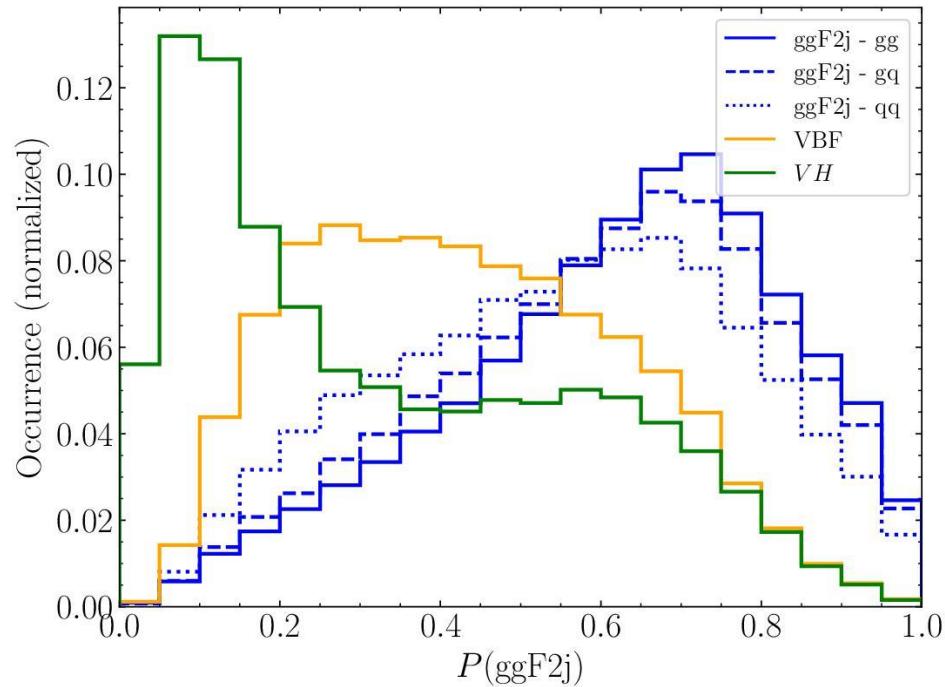
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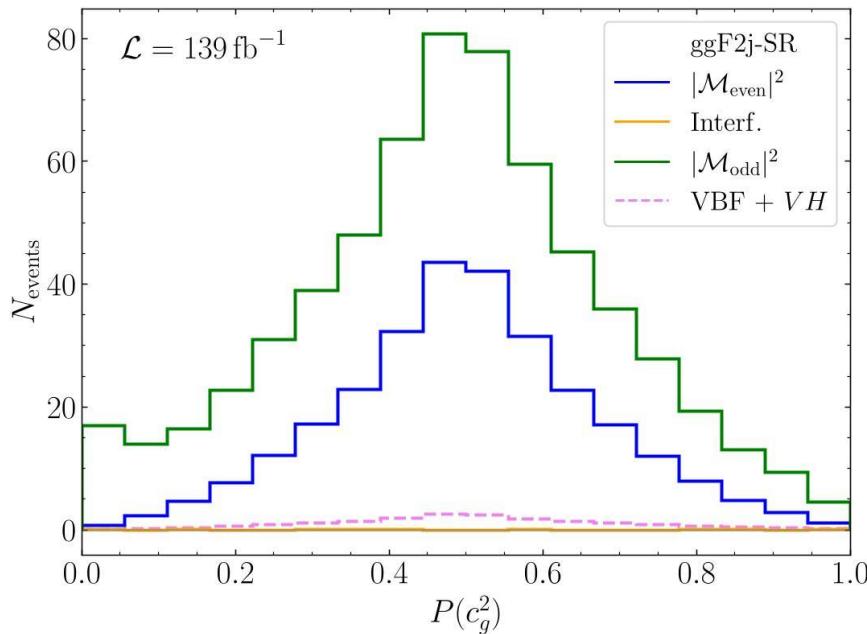
➤ Train a CP-even and a CP-odd classifier in a ggF2j-SR and a VBF-SR

Signal regions



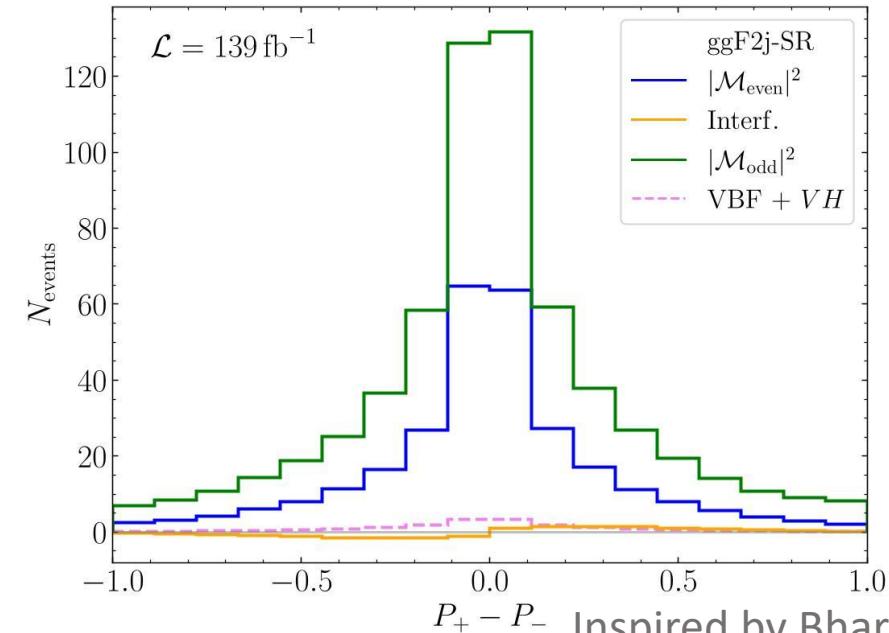
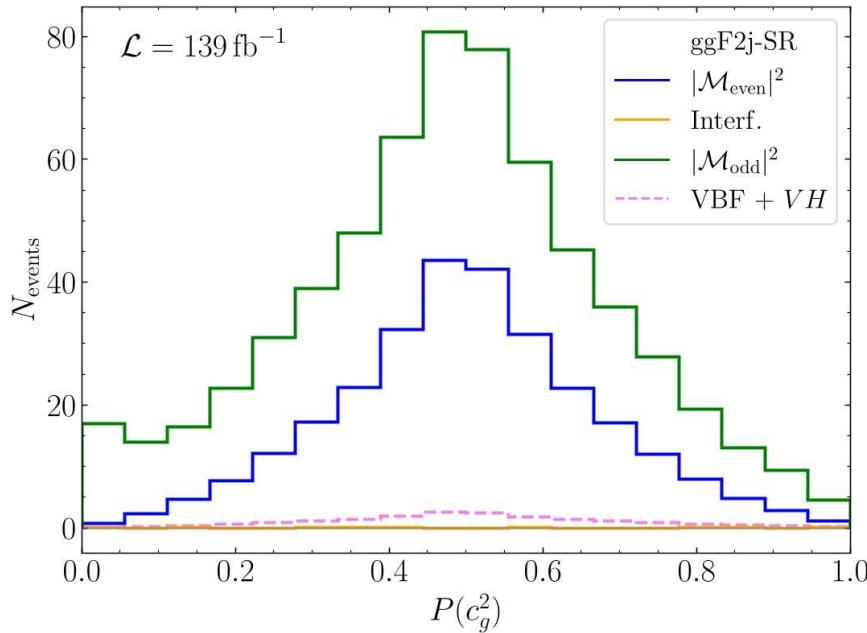
- Use $P(\text{signal}) > 0.5$ as a cut to define signal regions
- ggF2j with $q\bar{q}$ initial state are identified as VBF-like more often
- ggF2j interference events are also identified as VBF-like more often

ggF2j signal region



- $P(c_g^2)$ differentiates between $c_g^2 |\mathcal{M}_{\text{even}}|^2$ and $\tilde{c}_g^2 |\mathcal{M}_{\text{odd}}|^2$
- Kinematically very similar, but some separation in outer bins
- Interference term cancels out

ggF2j signal region



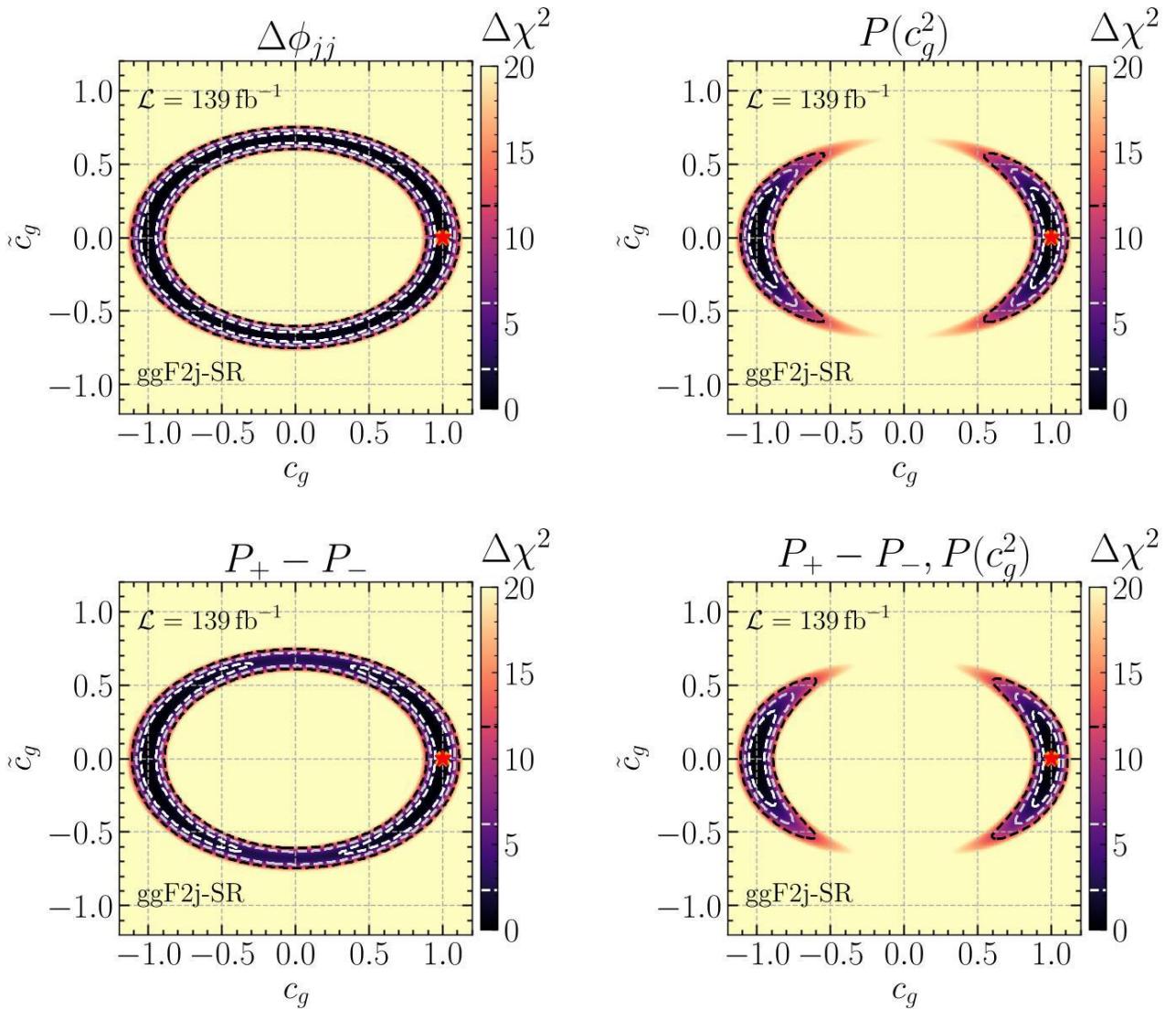
Inspired by [Bhardwaj et al. '21](#)

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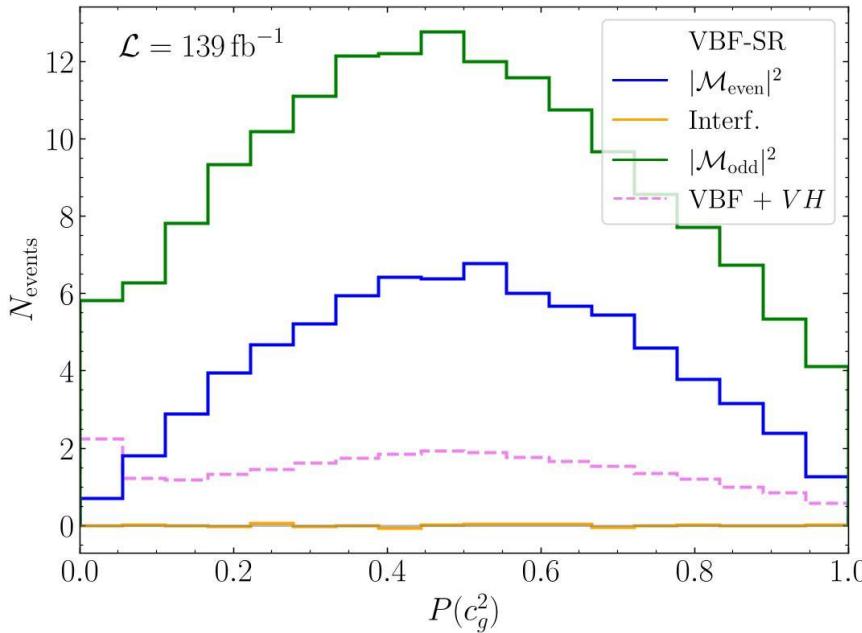
- $P_+ - P_-$ differentiates between **positive and negative interference**
- Interference barely visible due to low cross section & looks more VBF-like
- CP-even terms are symmetric

ggF2j signal region

- Ellipse from total rate
- $\Delta\phi_{jj}$ alone is not able to resolve the ellipse
- 2D-limits dominated by the $P(c_g^2)$ classifier (low interference contribution)
- $|\tilde{c}_g| \leq 0.32 @ 1\sigma$

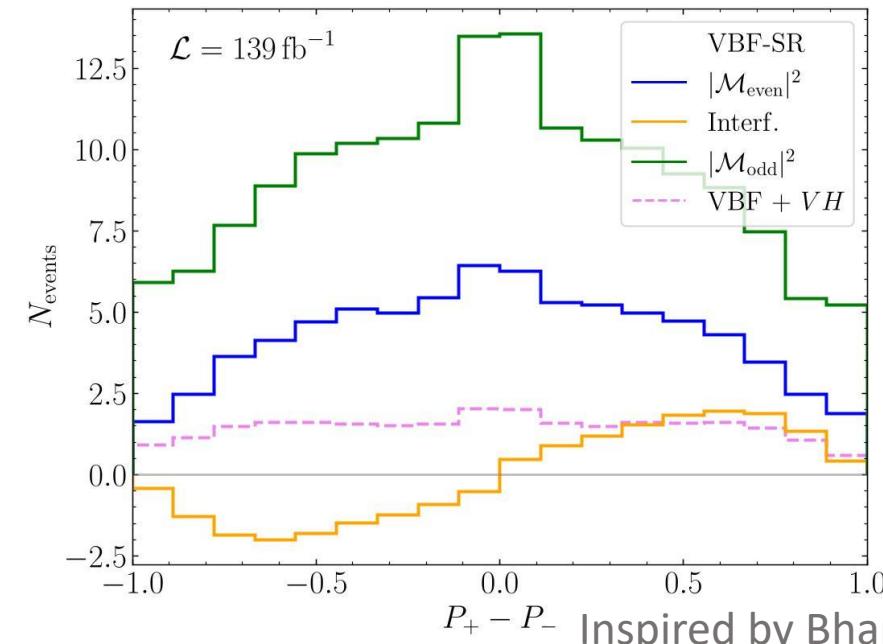
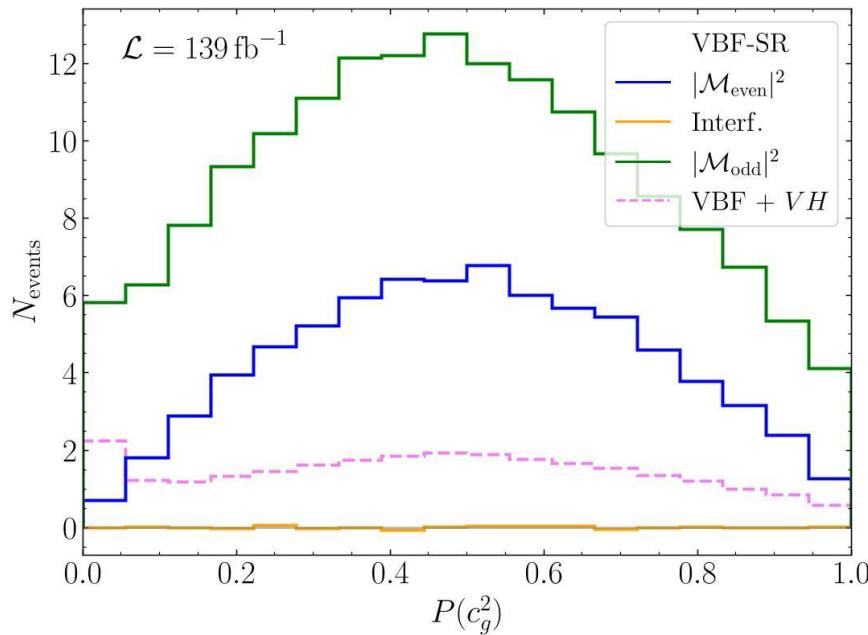


VBF signal region



- $P(c_g^2)$ differentiates between $c_g^2 |\mathcal{M}_{\text{even}}|^2$ and $\tilde{c}_g^2 |\mathcal{M}_{\text{odd}}|^2$
- Less events than in ggF2j-SR, statistical fluctuations visible
- Wider peak around $P(c_g^2) = 0.5$

VBF signal region



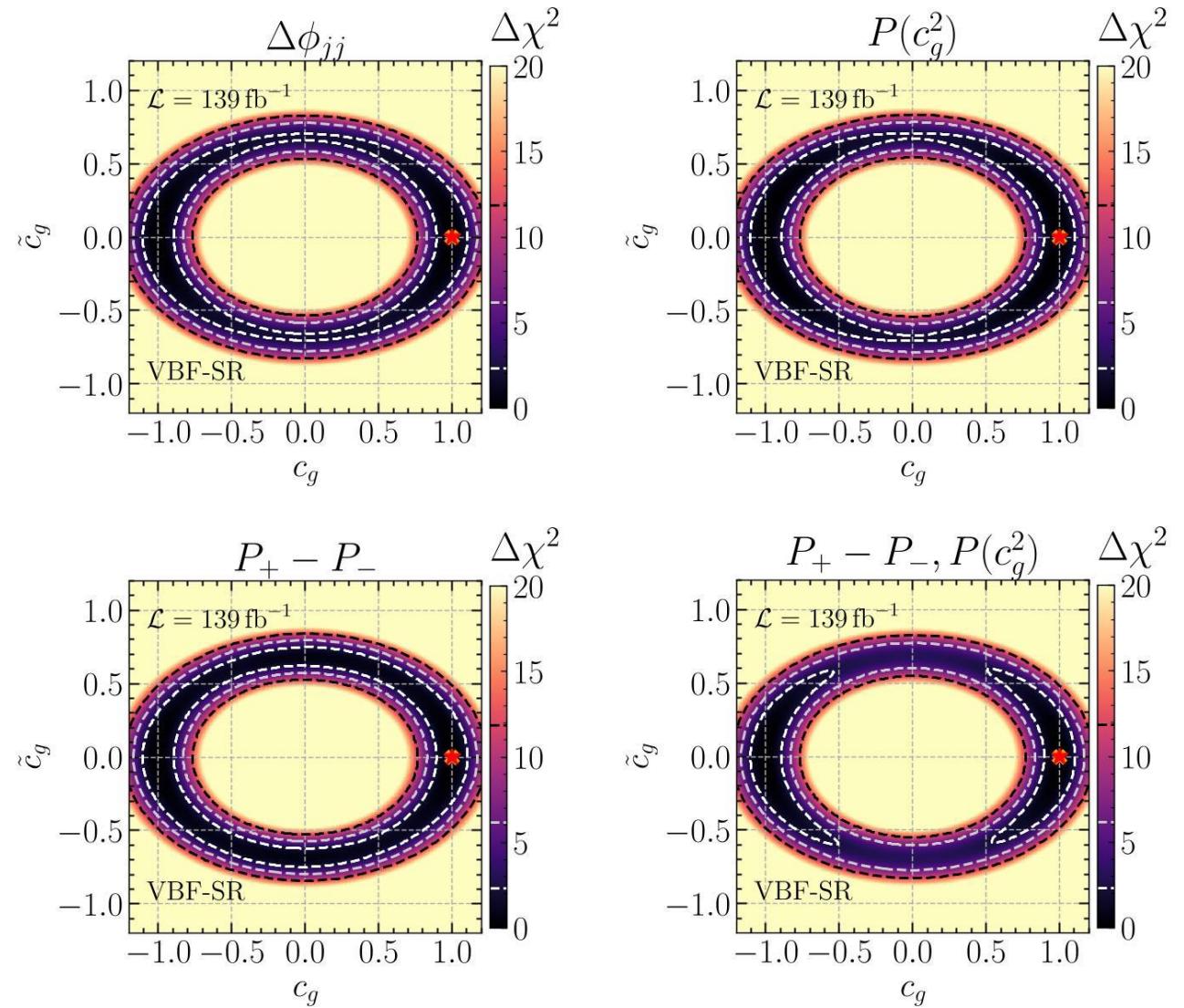
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- Less events than in ggF2j-SR, statistical fluctuations visible
- Wider peak around $P(c_g^2) = 0.5$

- $P_+ - P_-$ differentiates between **positive and negative interference**
- Interference much more pronounced due to overall lower events and their VBF-like kinematic

VBF signal region

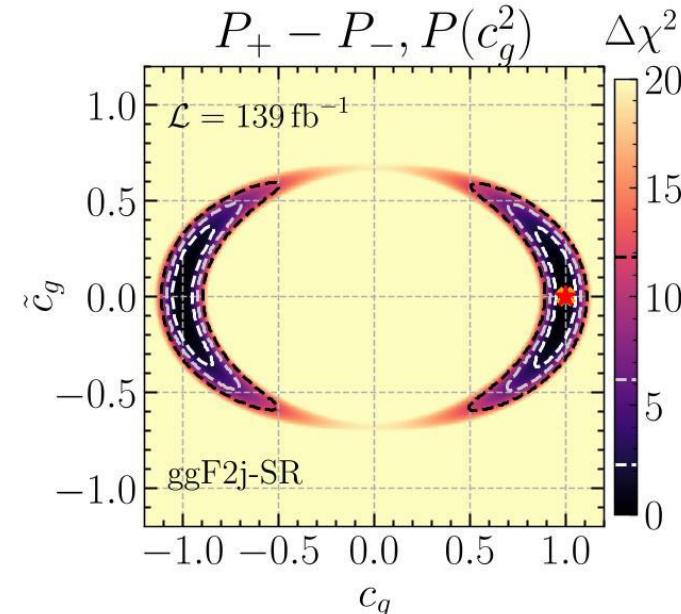
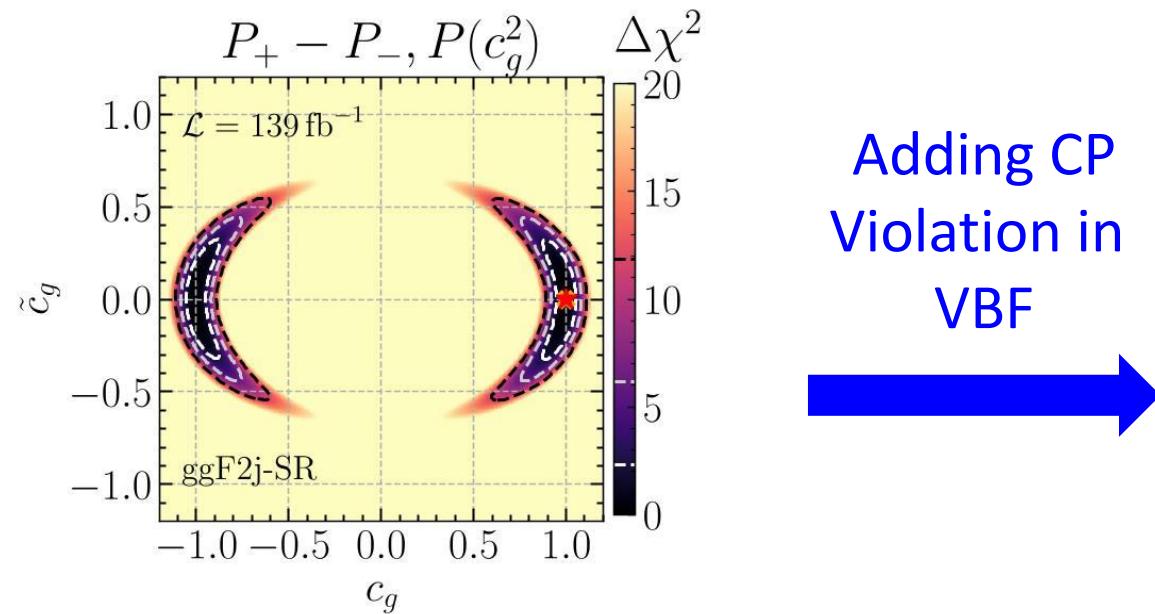
- Ellipse wider due to lower number of events
- Individual classifiers perform worse than $\Delta\phi_{jj}$ (less information) – large importance of $\Delta\phi_{jj}$
- 2D-limits are significantly stronger than 1D-limits
- $|\tilde{c}_g| \leq 0.58 @ 1\sigma$



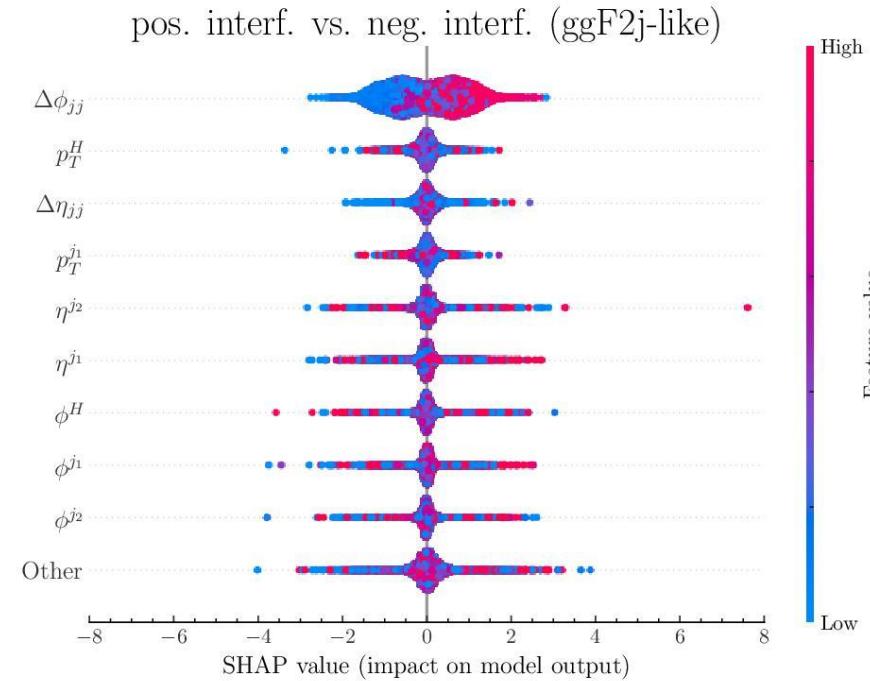
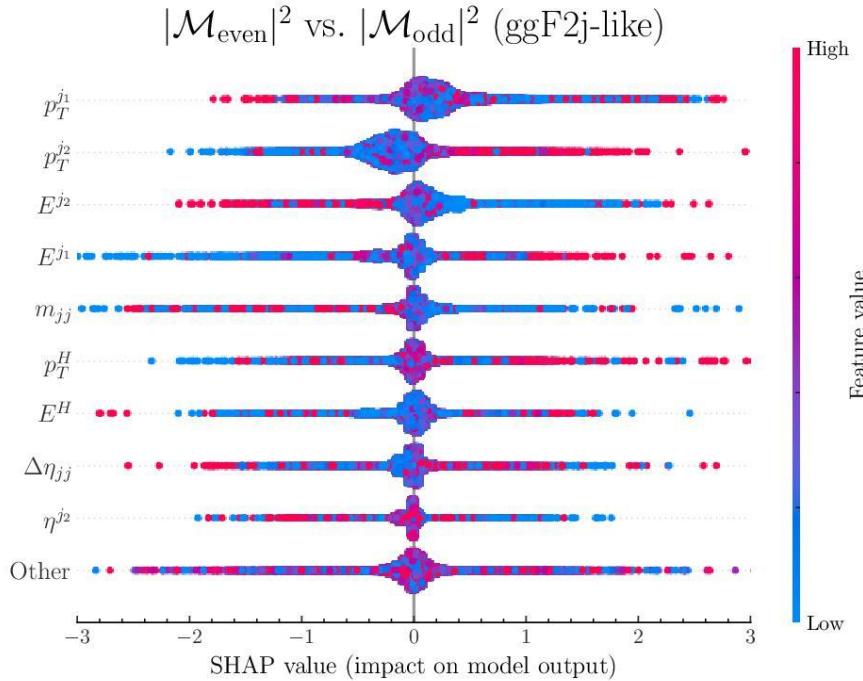
CP violation in HVV – ggF2j-SR

BSM physics might introduce CP violation in multiple Higgs couplings
⇒ inject CP violation in VBF production, evaluate effect on c_g, \tilde{c}_g limits

- Limits slightly weaken from $|\tilde{c}_g| \leq 0.32 @ 1\sigma$ to $|\tilde{c}_g| \leq 0.37 @ 1\sigma$
- ggF2j-SR provides conservative results for limited knowledge of the HVV coupling
- Opposite effect in VBF-SR ⇒ fakes much stronger limits on \tilde{c}_g (see backup)



Feature importance (SHAPLEY) – ggF2j-SR



- No clear „best“ variables for differentiating $c_g^2 |\mathcal{M}_{\text{even}}|^2$ vs $\tilde{c}_g^2 |\mathcal{M}_{\text{odd}}|^2$
- $\Delta\phi_{jj}$ is most important for interference classification (CP-odd)
- $\Delta\phi_{jj}$ much less significance than in the VBF-SR (see backup)

Results & Conclusion

Expected limits on $|\alpha^{Htt}| = \left| \tan^{-1} \frac{\tilde{c}_t}{c_t} \right| :$

Process / SR	Our work		Literature		References
	68% CL	95% CL	68% CL	95% CL	
ggF2j (ggF2j-SR)	15°	25°	--	--	
ggF2j (VBF-SR)	25°	--	26°	--	<u>CMS '22</u>
<i>tth</i>	--	--	35°	43°	<u>ATLAS '20</u>
Global fit	21°	28°	--	--	Update from Bahl et al. '22

- A multivariate analysis in the ggF2j-SR is expected to give strongest limits on α^{Htt}
- The ggF2j-SR is robust against CP violation in HVV mimicking CP violation in ggH
- Our findings consolidate the need for a dedicated ggF2j signal region

Backup

Backup: Other notations

Our work: $\mathcal{L}_{ggH} = -\frac{1}{4v} \left(-\frac{\alpha_s}{3\pi} \textcolor{blue}{c}_g G_{\mu\nu}^a G^{\mu\nu,a} + \frac{\alpha_s}{2\pi} \tilde{c}_g G_{\mu\nu}^a \tilde{G}^{\mu\nu,a} \right) H$

Separate top loop: $\mathcal{L}_{ggH} = -\frac{\alpha_s \pi}{v} (\textcolor{blue}{c}_{gg} G_{\mu\nu}^a G^{\mu\nu,a} + \tilde{c}_{gg} G_{\mu\nu}^a \tilde{G}^{\mu\nu,a}) H$

$$\Rightarrow \textcolor{blue}{c}_g = 1 + 12\pi^2 \textcolor{blue}{c}_{gg}, \quad \tilde{c}_g = -8\pi^2 \tilde{c}_{gg}$$

Used in CMS '22

SMEFT: $\textcolor{blue}{c}_g \sim 1 + \sum_i c_i / \Lambda^2 + \dots, \quad \tilde{c}_g \sim \sum_i \tilde{c}_i / \Lambda^2 + \dots$

c_i, \tilde{c}_i : Wilson coefficients

Backup: Cutflow

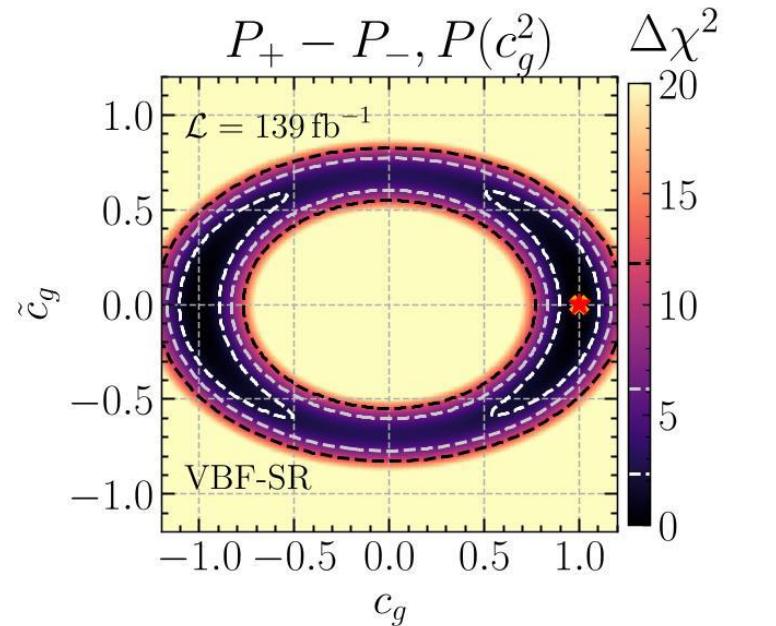
	Fraction of accepted events				
Applied cut	ggF2j $ \mathcal{M}_{\text{even}} ^2$	ggF2j Interf.	ggF2j $ \mathcal{M}_{\text{odd}} ^2$	VBF	VH
Initial events	100%	100%	100%	100%	100%
$N_j \geq 2, N_\gamma \geq 2$	48.1%	50.8%	48.1%	62.6%	49.8%
$100\text{GeV} \leq m_{\gamma\gamma}$ $m_{\gamma\gamma} \leq 140\text{GeV}$	47.8%	50.5%	47.9%	62.0%	49.4%
$p_T^{\gamma_1}/m_{\gamma\gamma} \geq 0.35$ $p_T^{\gamma_2}/m_{\gamma\gamma} \geq 0.25$	39.4%	40.9%	39.8%	50.0%	40.5%
$p_T^{j_1} \geq 30\text{GeV}$ $p_T^{j_2} \geq 20\text{GeV}$	38.6%	40.2%	38.6%	49.7%	39.9%
$ \eta_j \leq 2.5$ $ \eta_\gamma \leq 2.5$	22.9%	21.5%	22.7%	39.8%	31.2%
$p_T^H \leq 200\text{GeV} *$	18.6%	18.4%	18.3%	34.4%	26.8%

* It is possible to relax this cut using FT_{approx} → see [Maltoni et al. '14](#)

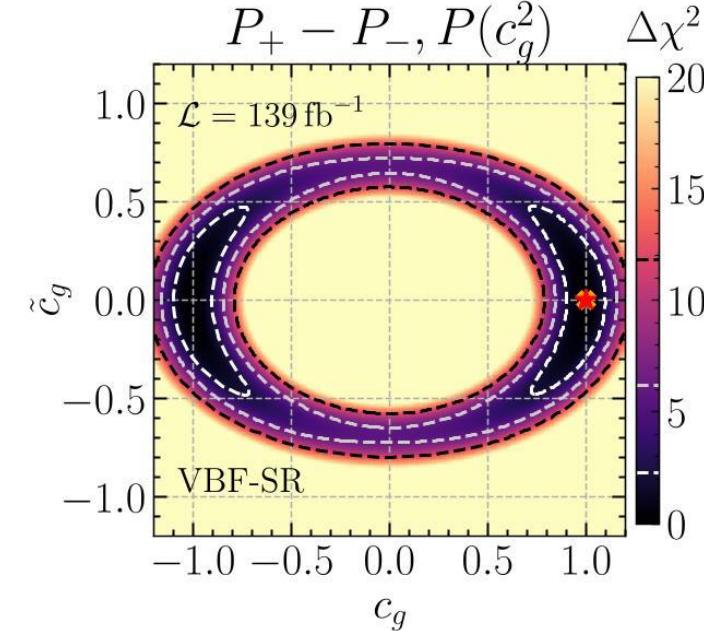
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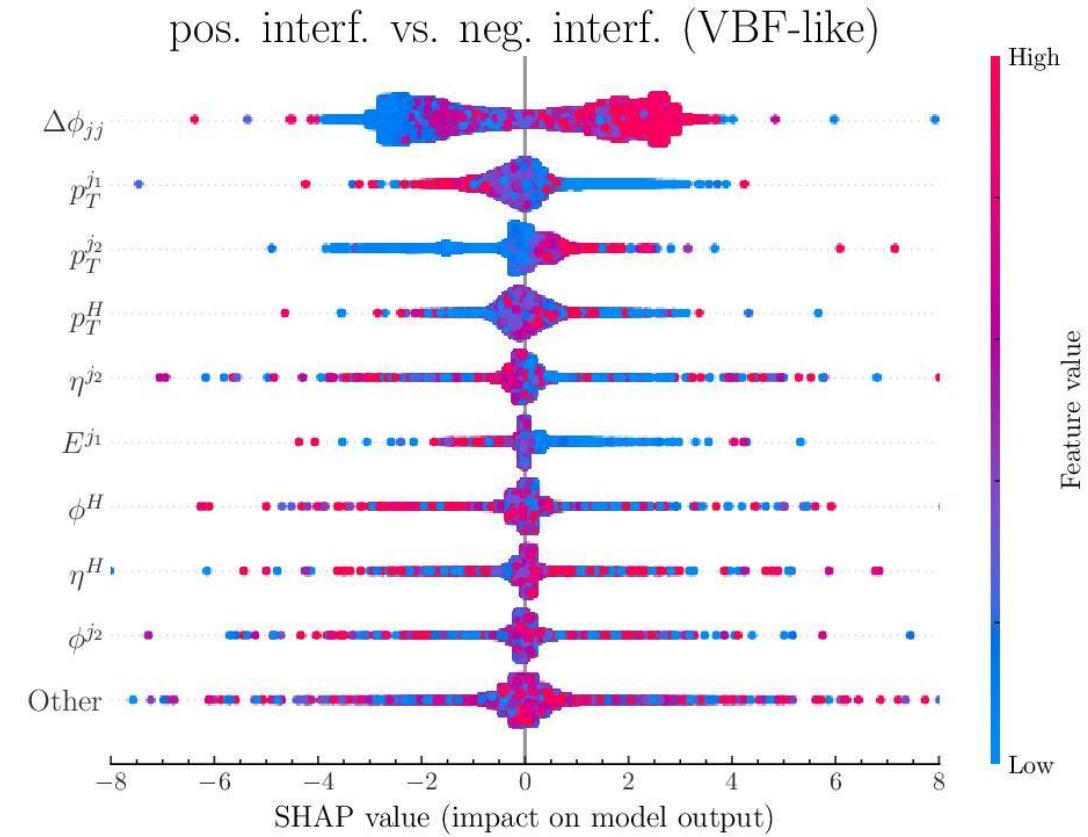
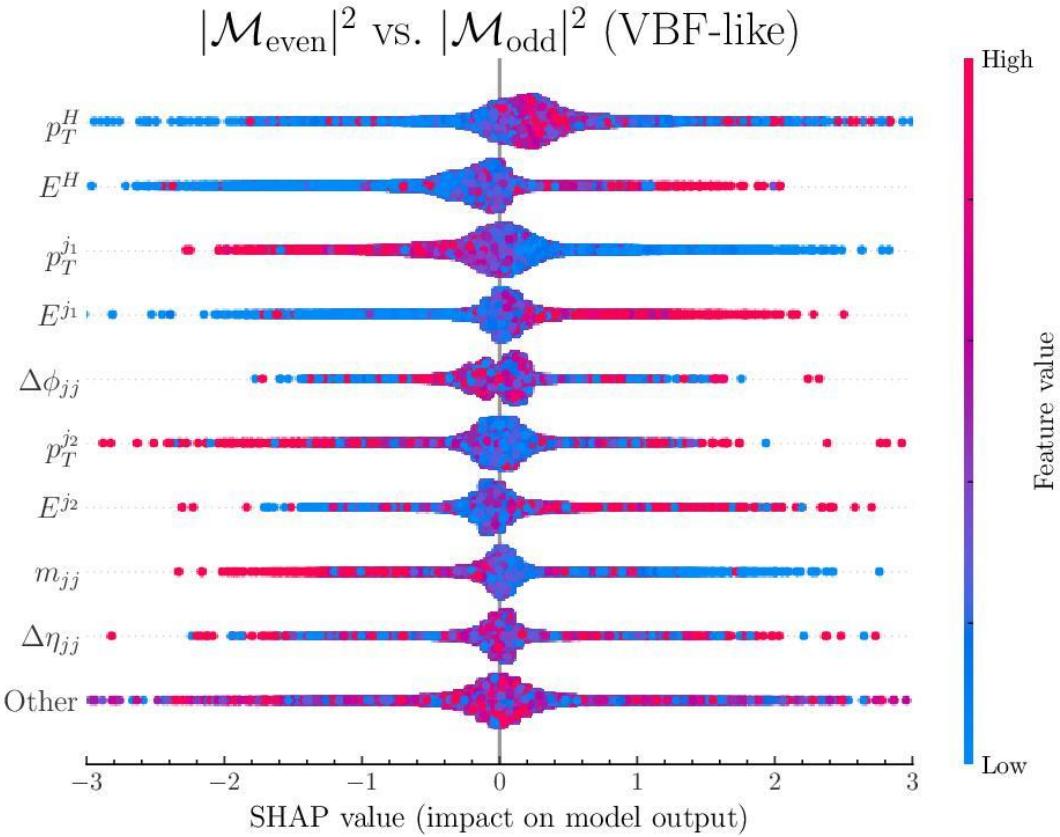
- Limits tighten from $|\tilde{c}_g| \leq 0.58$ @ 1σ to $|\tilde{c}_g| \leq 0.48$ @ 1σ
- VBF-SR might “fake” stronger limits not originating from CP violation in ggH



Adding CP
Violation in
VBF



Backup: Feature importance – VBF-SR



Backup: Comparison to global fit

- Interpret c_g, \tilde{c}_g as top-Yukawa coupling modifiers c_t, \tilde{c}_t (heavy top limit / strong limits on colored BSM particles)
- Global fit based on experimental results for Higgs signal rates
- Performed with HiggsTools [Bahl et al. '22](#)
- Different form of ellipsis due to recent $t\bar{t}H(bb)$ measurements [ATLAS '19](#), [ATLAS '22](#)
- Similar constraints from our analysis

