

Introdução à Física de Partículas

Introduction to particle Physics

(4/4)



FILIPE JOAQUIM

IST Dep. de Física e CFTP , Lisboa, Portugal

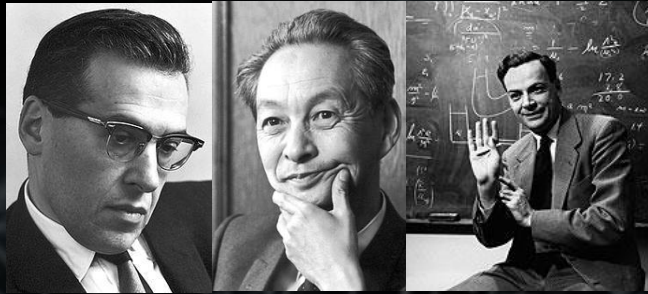


16ª Escola de Professores no CERN em Língua Portuguesa 2024

16th CERN Portuguese Language Teachers Programme 2024

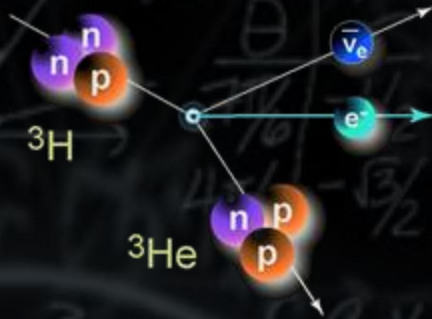
1-6 Setembro, CERN, Genebra

SITUAÇÃO NO FINAL DOS ANOS 40



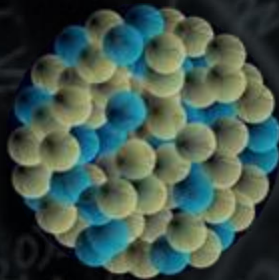
ELECTRODINÂMICA QUÂNTICA

Teoria quântica dos electrões, positrões, fotões e da interacção electromagnética.
(e^- , e^+ , γ)



FORÇA FRACA

Teoria do decaimento radioactivo descrita pela interacção de Fermi.



FORÇA FORTE

Força responsável pela coesão do núcleo descrita pelo potencial de Yukawa.

(n , p , π)

SIMETRIA

ΣΙΜΕΤΡΙΑ

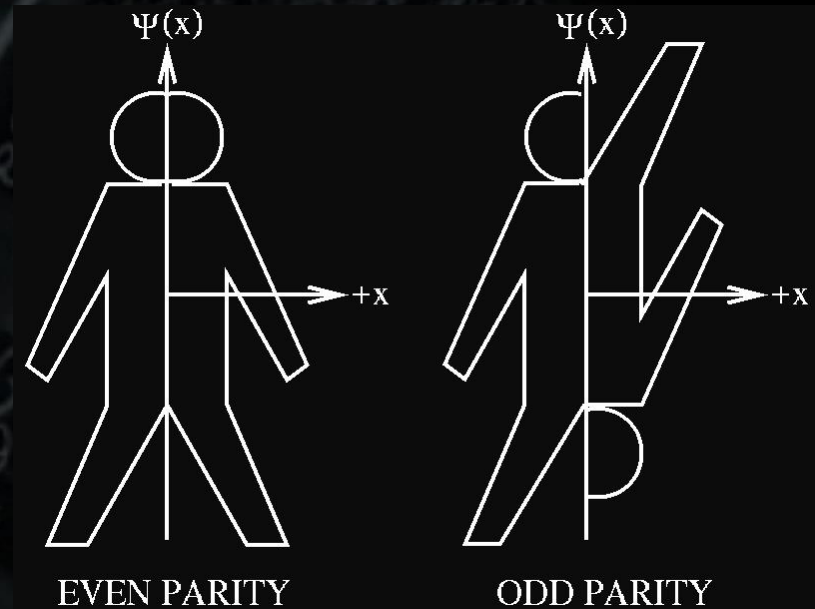
PARIDADE

Transformação de paridade P:

$$P : \vec{r} \rightarrow -\vec{r}$$

Uma dupla transformação de paridade corresponde a não fazer nada...

$$P^2 \psi(\vec{r}) = \psi(\vec{r}) \Rightarrow P \psi(\vec{r}) = \pm \psi(\vec{r})$$



Vector: $P(\vec{p}) = P\left(\frac{d\vec{r}}{dt}\right) = -\vec{p}$

Pseudo-vector: $\vec{L} = \vec{r} \times \vec{p}, \vec{S}$

A interacção electromagnética, a força forte e a gravidade são invariantes debaixo de paridade.

1948

Paridade intrínseca: As partículas elementares (e não só) têm uma paridade intrínseca.

$$P_T = P P'$$

PHYSICAL REVIEW

VOLUME 95, NUMBER 6

SEPTEMBER 15, 1954

Absorption of Negative Pions in Deuterium : Parity of the Pion*

W. CHINOWSKY AND J. STEINBERGER
Columbia University, New York, New York
(Received June 8, 1954)

The reaction $\pi^- + d \rightarrow 2n$ has been observed by detecting the two neutrons in coincidence with slow negative mesons incident on a liquid deuterium target. The observed angular correlation of the two neutrons confirms the identification of the process. The process is therefore not forbidden, and this fact may be used to establish the odd relative parity of the pion and the nucleon.

Dalitz (1954): Puzzle $\theta - \tau$ Duas partículas (mesões θ e τ) com a mesma massa decaíam para estados de paridade diferente.

$$\tau^+ \rightarrow \pi^+ \pi^- \pi^+ : P_T = (-1)^3 = -1 \quad \theta^+ \rightarrow \pi^+ \pi^0 : P_T = (-1)^2 = 1$$

1954

RECENT experimental data indicate closely identical masses¹ and lifetimes² of the θ^+ ($\equiv K_{\pi 2}^+$) and the τ^+ ($\equiv K_{\pi 3}^+$) mesons. On the other hand, analyses³ of the decay products of τ^+ strongly suggest on the grounds of angular momentum and parity conservation that the τ^+ and θ^+ are not the same particle. This poses a rather puzzling situation that has been extensively discussed.⁴

One way out of the difficulty is to assume that parity is not strictly conserved, so that θ^+ and τ^+ are two different decay modes of the same particle, which necessarily has a single mass value and a single lifetime. We wish to analyze this possibility in the present paper.

PHYSICAL REVIEW VOLUME 104, NUMBER 1 OCTOBER 1, 1954

Question of Parity Conservation in Weak Interactions*

T. D. LEE, *Columbia University, New York, New York*

AND

C. N. YANG, *Brookhaven National Laboratory, Upton, New York*

(Received June 21, 1954)

The question of parity conservation in β -decays and in hyperon and meson decays is examined. Possible experiments are suggested which might test parity conservation in these interactions.

RECENT experimental data indicate closely identical masses¹ and lifetimes² of the θ^+ ($\equiv K_{\pi 2}^+$) and the τ^+ ($\equiv K_{\pi 3}^+$) mesons. On the other hand, analyses³ of the decay products of τ^+ strongly suggest on the grounds of angular momentum and parity conservation that the τ^+ and θ^+ are not the same particle. This poses a rather puzzling situation that has been extensively discussed.⁴

One way out of the difficulty is to assume that parity is not strictly conserved, so that θ^+ and τ^+ are two different decay modes of the same particle, which necessarily has a single mass value and a single lifetime. We wish to analyze this possibility in the present paper against the background of the existing experimental evidence of parity conservation. It will become clear that existing experiments do indicate parity conservation in strong and electromagnetic interactions to a

PRESENT EXPERIMENTAL LIMIT ON PARITY NONCONSERVATION

If parity is not strictly conserved, all atomic and nuclear states become mixtures consisting mainly of the state they are usually assigned, together with small percentages of states possessing the opposite parity. The fractional weight of the latter will be called β^2 . It is a quantity that characterizes the degree of violation of parity conservation.

The existence of parity selection rules which work well in atomic and nuclear physics is a clear indication that the degree of mixing, β^2 , cannot be large. From such considerations one can impose the limit $\beta^2 \leq (\epsilon/\lambda)^2$, which for atomic spectroscopy is, in most cases, $\sim 10^{-4}$. In general a less accurate limit obtains for nuclear spectroscopy.

Parity nonconservation implies the existence of inter-



O prêmio Nobel da Física foi atribuído a C. N. Yang e T.D. Lee em 1957;



elementary particles".



1956

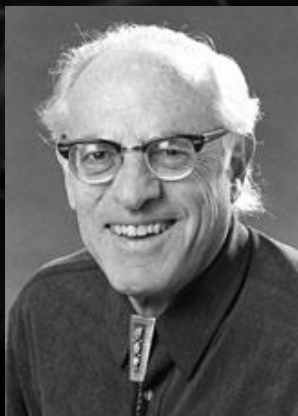
Em 1956 Cowan e Reines detectaram o neutrino do electrão usando como fonte os neutrinos provenientes de um reactor nuclear.



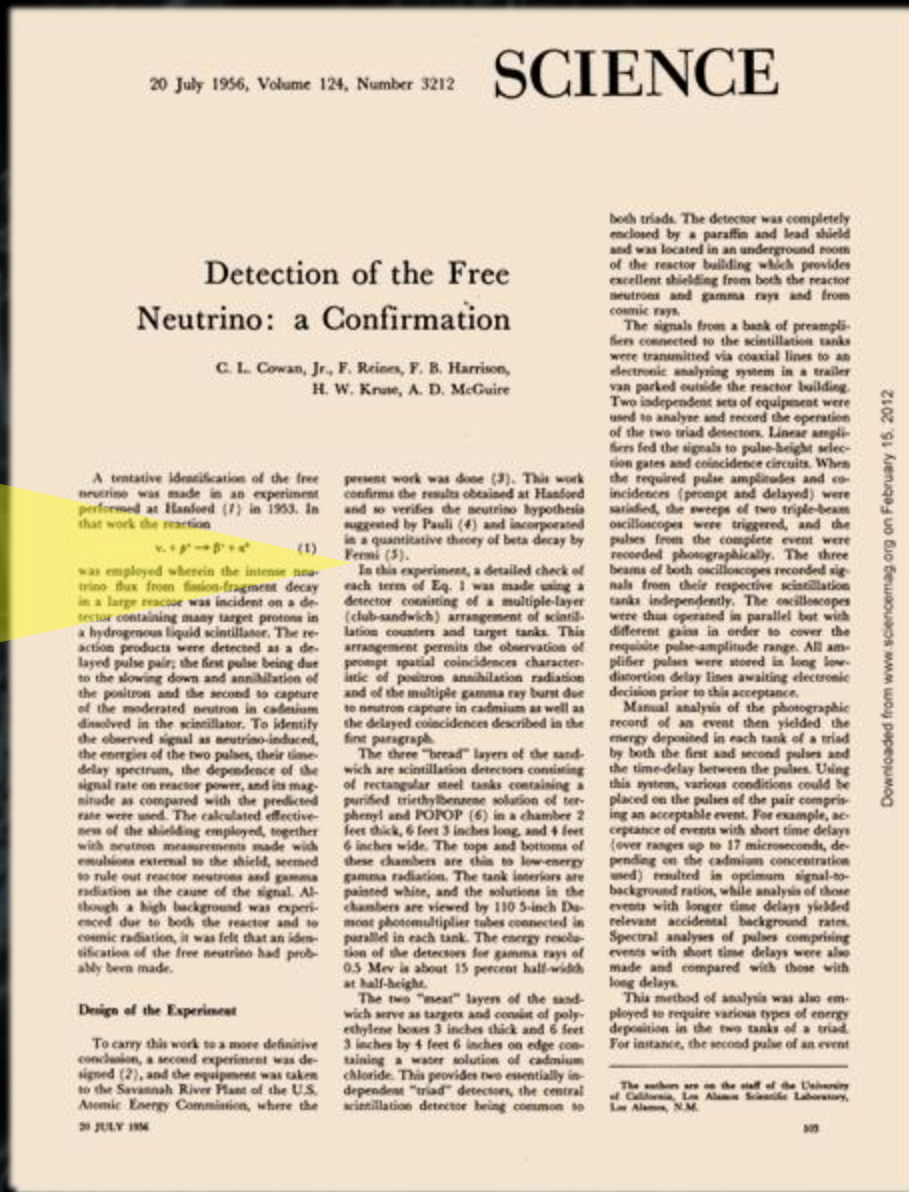
present work was done (3). This work confirms the results obtained at Hanford and so verifies the neutrino hypothesis suggested by Pauli (4) and incorporated in a quantitative theory of beta decay by Fermi (5).

O prémio Nobel da Física foi atribuído a F. Reines em 1995;

“for the detection of the neutrino”



26 anos depois: finalmente o neutrino



20 July 1956, Volume 124, Number 3212

SCIENCE

Detection of the Free Neutrino: a Confirmation

C. L. Cowan, Jr., F. Reines, F. B. Harrison, H. W. Kruse, A. D. McGuire

A tentative identification of the free neutrino was made in an experiment performed at Hanford (1) in 1953. In that work the reaction



was employed wherein the intense neutrino flux from fission-fragment decay in a large reactor was incident on a detector containing many target protons in a hydrogenous liquid scintillator. The reaction products were detected as a delayed pulse pair; the first pulse being due to the slowing down and annihilation of the positron and the second to capture of the moderated neutron in cadmium dissolved in the scintillator. To identify the observed signal as neutrino-induced, the energies of the two pulses, their time-delay spectrum, the dependence of the signal rate on reactor power, and its magnitude as compared with the predicted rate were used. The calculated effectiveness of the shielding employed, together with neutron measurements made with emulsions external to the shield, seemed to rule out reactor neutrons and gamma radiation as the cause of the signal. Although a high background was experienced due to both the reactor and to cosmic radiation, it was felt that an identification of the free neutrino had probably been made.

Design of the Experiment

To carry this work to a more definitive conclusion, a second experiment was designed (2), and the equipment was taken to the Savannah River Plant of the U.S. Atomic Energy Commission, where the

present work was done (3). This work confirms the results obtained at Hanford and so verifies the neutrino hypothesis suggested by Pauli (4) and incorporated in a quantitative theory of beta decay by Fermi (5).

In this experiment, a detailed check of each term of Eq. 1 was made using a detector consisting of a multiple-layer (club-sandwich) arrangement of scintillation counters and target tanks. This arrangement permits the observation of prompt spatial coincidences characteristic of positron annihilation radiation and of the multiple gamma ray burst due to neutron capture in cadmium as well as the delayed coincidences described in the first paragraph.

The three "bread" layers of the sandwich are scintillation detectors consisting of rectangular steel tanks containing a purified triethylbenzene solution of toluene and POPOP (6) in a chamber 2 feet thick, 6 feet 3 inches long, and 4 feet 6 inches wide. The tops and bottoms of these chambers are thin to low-energy gamma radiation. The tank interiors are painted white, and the solutions in the chambers are viewed by 110 5-inch Du-mont photomultiplier tubes connected in parallel in each tank. The energy resolution of the detectors for gamma rays of 0.5 Mev is about 15 percent half-width at half-height.

The two "meat" layers of the sandwich serve as targets and consist of polyethylene boxes 3 inches thick and 6 feet 3 inches by 4 feet 6 inches on edge containing a water solution of cadmium chloride. This provides two essentially independent "triad" detectors, the central scintillation detector being common to

both triads. The detector was completely enclosed by a paraffin and lead shield and was located in an underground room of the reactor building which provides excellent shielding from both the reactor neutrons and gamma rays and from cosmic rays.

The signals from a bank of preamplifiers connected to the scintillation tanks were transmitted via coaxial lines to an electronic analyzing system in a trailer van parked outside the reactor building. Two independent sets of equipment were used to analyze and record the operation of the two triad detectors. Linear amplifiers fed the signals to pulse-height selection gates and coincidence circuits. When the required pulse amplitudes and coincidences (prompt and delayed) were satisfied, the sweeps of two triple-beam oscilloscopes were triggered, and the pulses from the complete event were recorded photographically. The three beams of both oscilloscopes recorded signals from their respective scintillation tanks independently. The oscilloscopes were thus operated in parallel but with different gains in order to cover the requisite pulse-amplitude range. All amplifier pulses were stored in long low-distortion delay lines awaiting electronic decision prior to this acceptance.

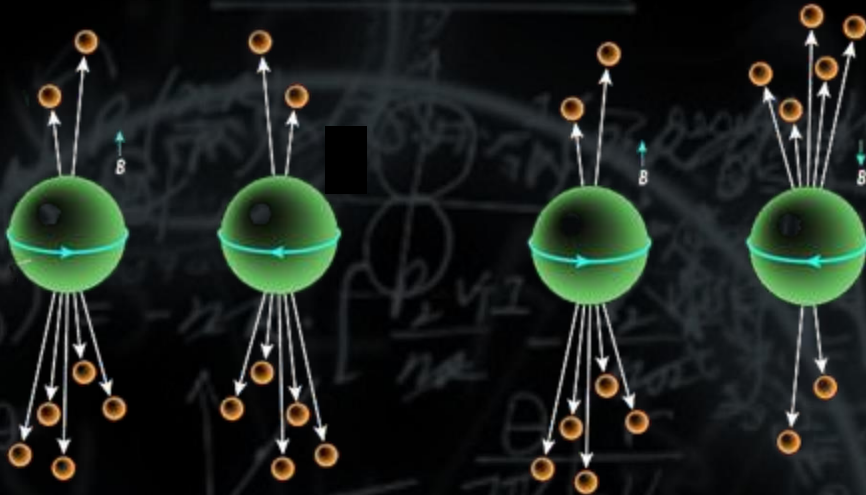
Manual analysis of the photographic record of an event then yielded the energy deposited in each tank of a triad by both the first and second pulses and the time-delay between the pulses. Using this system, various conditions could be placed on the pulses of the pair comprising an acceptable event. For example, acceptance of events with short time delays (over ranges up to 17 microseconds, depending on the cadmium concentration used) resulted in optimum signal-to-background ratios, while analysis of those events with longer time delays yielded relevant accidental background rates. Spectral analyses of pulses comprising events with short time delays were also made and compared with those with long delays.

This method of analysis was also employed to require various types of energy deposition in the two tanks of a triad. For instance, the second pulse of an event

The authors are on the staff of the University of California, Los Alamos Scientific Laboratory, Los Alamos, N.M.

1957

Em 1957 Wu obtém a prova experimental de que a paridade não é conservada pela interação fraca.



Experimental Test of Parity Conservation in Beta Decay*

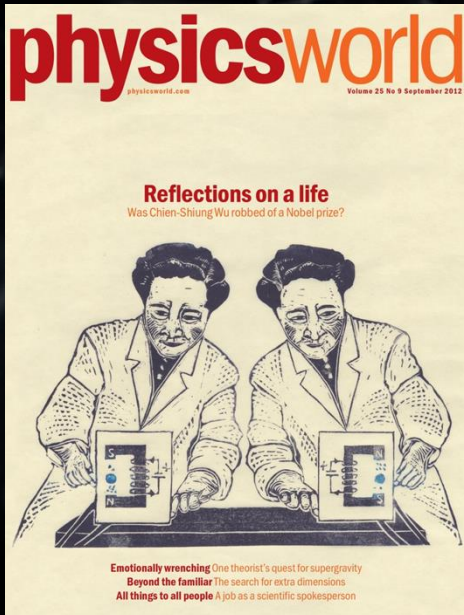
C. S. WU, *Columbia University, New York, New York*

AND

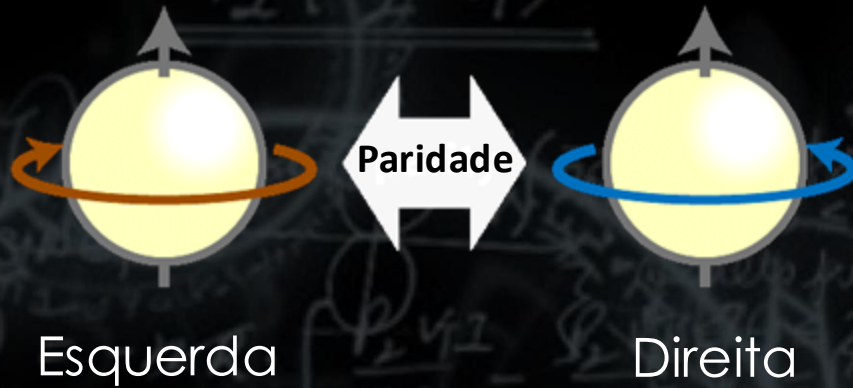
E. AMBLER, R. W. HAYWARD, D. D. HOPPES, AND R. P. HUDSON,
National Bureau of Standards, Washington, D. C.

(Received January 15, 1957)

In a recent paper¹ on the question of parity in weak interactions, Lee and Yang critically surveyed the experimental information concerning this question and reached the conclusion that there is no existing evidence either to support or to refute parity conservation in weak interactions. They proposed a number of experiments on beta decays and hyperon and meson decays which would provide the necessary evidence for parity conservation or nonconservation. In beta decay, one could measure the angular distribution of the electrons coming from beta decays of polarized nuclei. If an asymmetry in the distribution between θ and $180^\circ - \theta$ (where θ is the angle between the orientation of the parent nuclei and the momentum of the electrons) is observed, it provides unequivocal proof that parity is not conserved in beta decay. This asymmetry effect has been observed in the case of oriented Co^{60} .



PARIDADE E HELICIDADE



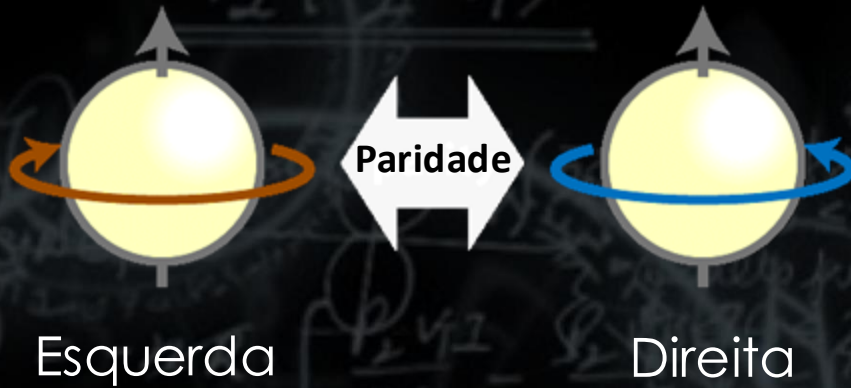
HELICIDADE: $h = \frac{\vec{s} \cdot \vec{p}}{|\vec{p}|}$

$h=1$: Direita

$h=-1$: Esquerda



PARIDADE E HELICIDADE



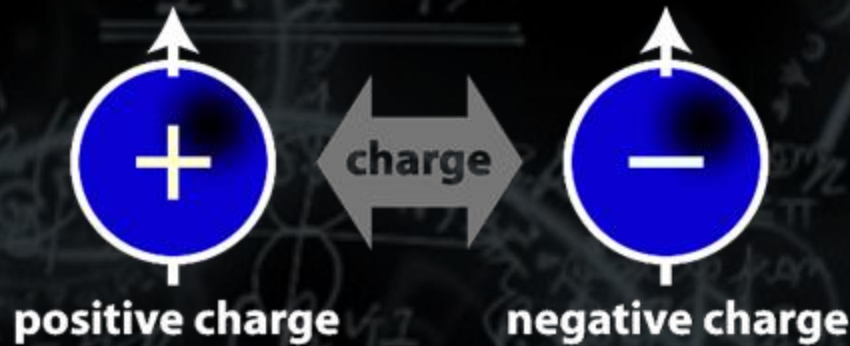
HELICIDADE: $h = \frac{\vec{S} \cdot \vec{p}}{|\vec{p}|}$

$h=1$: Direita

$h=-1$: Esquerda

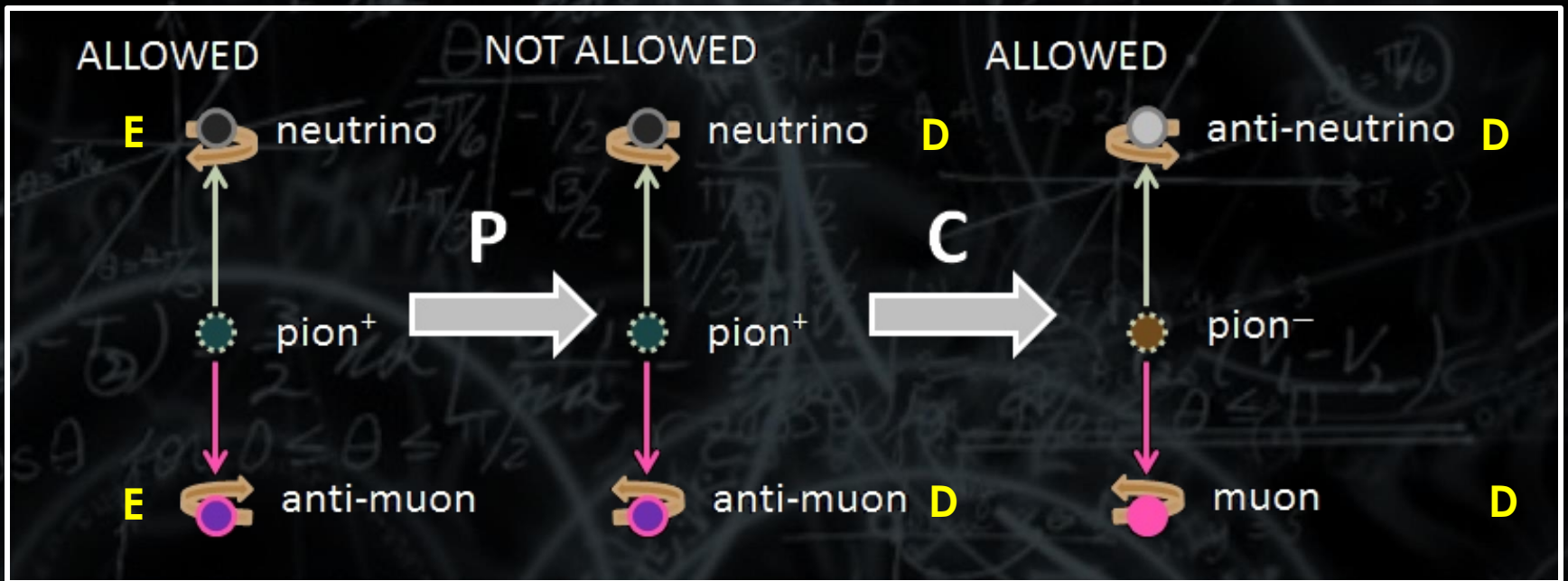


SIMETRIAS C (conjugação de carga) E CP



A conjugação de carga transforma uma partícula na sua antipartícula

$$C|p\rangle = |\bar{p}\rangle$$



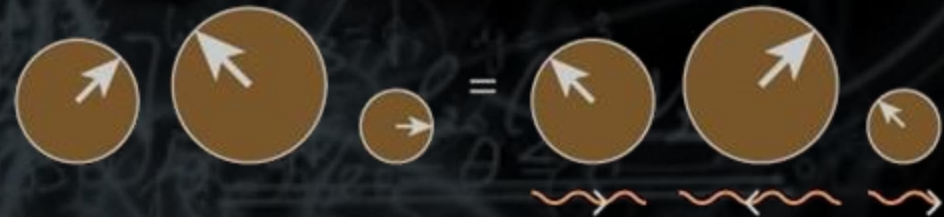
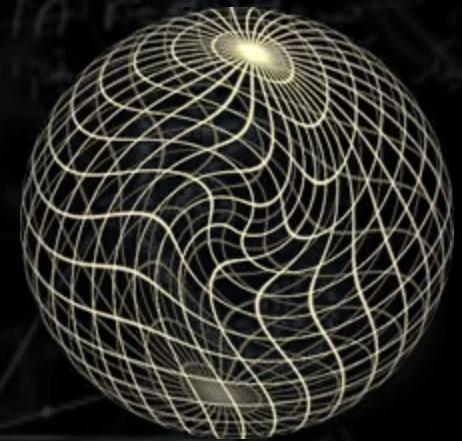
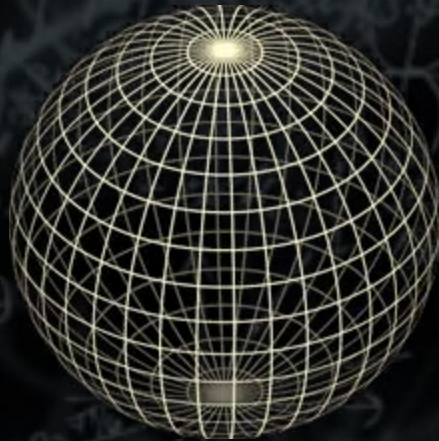
TALVEZ CP SEJA CONSERVADA...

TRANSFORMAÇÕES GLOBAIS E LOCAIS

Esfera original

Transformação global

Transformação local



Os bósons de gauge (e respectivas forças) surgem como consequência de impôr invariância debaixo de simetrias locais

O MODELO PADRÃO: A RECEITA

O GRUPO DE SIMETRIA DO MODELO PADRÃO É:

$$SU(2)_L \times U(1)_Y$$



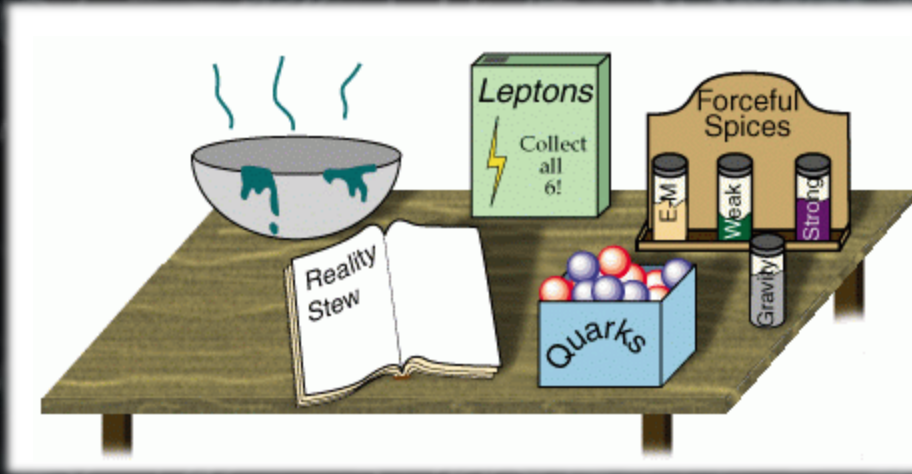
Glashow



Weinberg



Salam



- 1) Distribua as partículas elementares pelas “representações” do grupo de simetria.
- 2) Escreva todas as interações que são invariantes debaixo do grupo de simetria local.

RESULTADO: Teoria que descreve a interação dos quarks, leptões e bosões de gauge.

mas... Todas as partículas têm massa nula !

1964

O mecanismo ABEGHKK'tH

Anderson, Brout, Englert, Guralnik, Hagen, Higgs, Kibble and 't Hooft



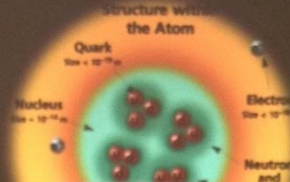
Como funciona então o mecanismo ABEGHKK'tH?

O CAMPO DE HIGGS...

Standard Model of FUNDAMENTAL PARTICLES AND INTERACTIONS

FERMIONS

Leptons $spin = 1/2$			Quarks $spin = 1/2$		
Flavor	Mass GeV/c ²	Electric charge	Flavor	Approx. Mass GeV/c ²	Electric charge
ν_e electron neutrino	< 0.17 eV	0	u up	0.002	2/3
ν_μ muon neutrino	0.000113	-1	d down	0.005	-1/3
ν_τ tau neutrino	< 0.0002	0	c charm	1.3	2/3
e^- electron	0.511	-1	s strange	0.1	-1/3
μ^- muon	0.106	-1	t top	175	2/3
τ^- tau	1.7771	-1	b bottom	4.3	-1/3



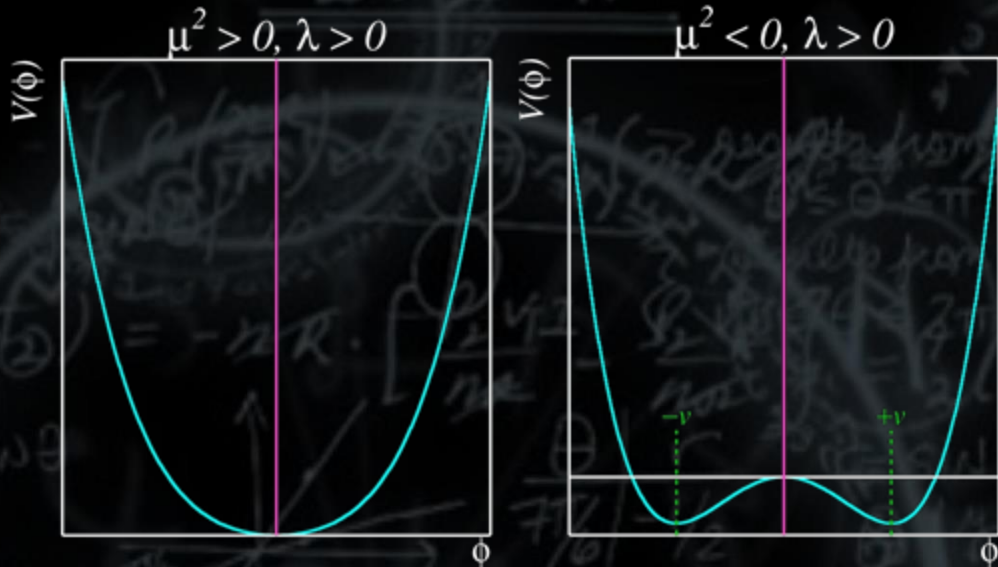
BOSONS

Unified Electroweak $spin = 1$			Strong (gluons) $spin = 1$		
Name	Mass GeV/c ²	Electric charge	Name	Mass GeV/c ²	Electric charge
γ photon	0	0	8 gluons	0	0
W^+	80.4	+1			
W^-	80.4	-1			
Z^0	91.187	0			

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi} \not{D} \psi + h.c. + \chi_i Y_{ij} \chi_j \phi + h.c. + |D_\mu \phi|^2 - V(\phi)$$

HS INTRODUCTION TO ELEMENTARY PARTICLES

O MECANISMO DE HIGGS



$$V(\phi) = \mu^2 |\phi|^2 + \lambda |\phi|^4$$

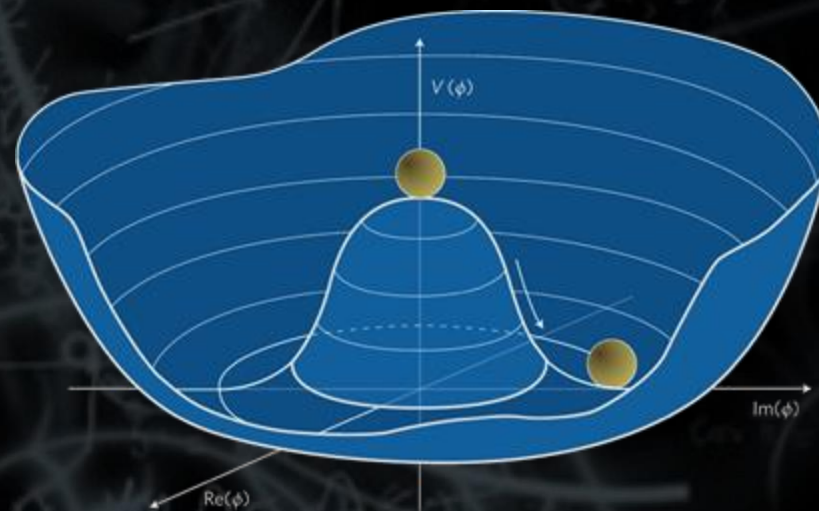
$$\text{"NO VÁCUO"}: v = \sqrt{-\frac{\mu^2}{2\lambda}}$$

Essencial para o mecanismo de Higgs funcionar.

A simetria é quebrada espontaneamente!!

Os bósons de gauge (W e Z) e os fermiões adquirem massa!

E... O FOTÃO PERMANECE SEM MASSA!!!



O MECANISMO DE HIGGS



O MECANISMO DE HIGGS

Algumas previsões da teoria:

$$M_Z \cos \theta_W = M_W, \quad \sin^2 \theta_W = 1 - M_W^2 / M_Z^2$$

$$M_W^2 \sin^2 \theta_W = \frac{e^2}{4\pi\sqrt{2} G_F}$$

Os bósons W e Z foram descobertos no CERN em 1983.

$$M_W = 80.385 \pm 0.015 \text{ GeV}$$

$$M_Z = 91.1876 \pm 0.086 \text{ GeV}$$



O detector gargamelle

O prêmio Nobel da Física foi atribuído a Rubbia e Van De Meer em 1984;

"for their decisive contributions to the large project, which led to the discovery of the field particles W and Z, communicators of weak interaction"



O NOBEL PARA O MODELO PADRÃO



O prémio Nobel da Física foi atribuído a Glashow, Weinberg e Salam em 1979;

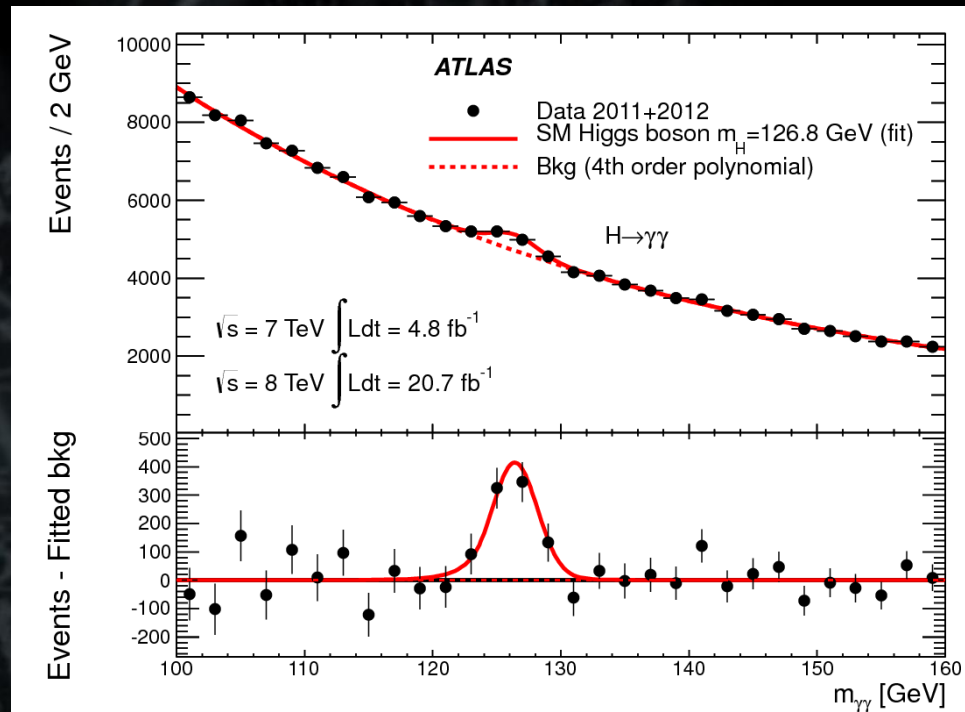
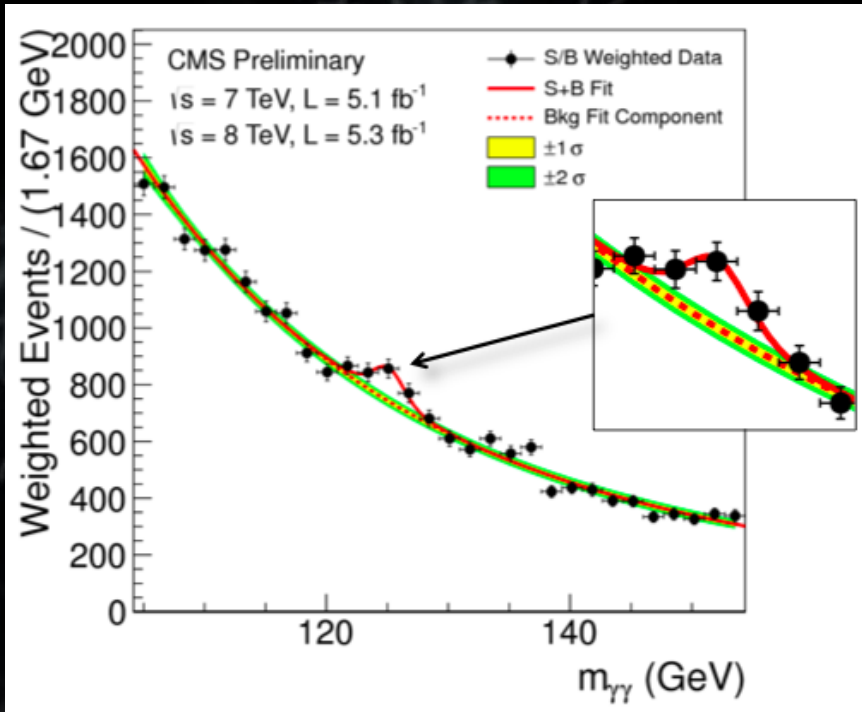


"for their contributions to the theory of the unified weak and electromagnetic interaction between elementary particles, including, the prediction of the weak neutral current".

Até ao dia 4 de Julho de 2012 não se sabia nada sobre o que estava por detrás da quebra de simetria electrofraca.

Até que...

BORN ON THE 4TH OF JULY



“The discovery of a particle consistent with the Higgs boson opens the way to more detailed studies, ... , and is likely to shed light on other mysteries of our Universe.”

Rolf Heuer, CERN D.G., Press Release July 4, 2012

“We are reaching into the fabric of the Universe at the level never done before... We are in the edge of a new exploration.”

Joe Incandela, CMS spokesperson, Press Conference, July 4, 2012

O ACONTECIMENTO CIENTÍFICO MAIS MEDIÁTICO

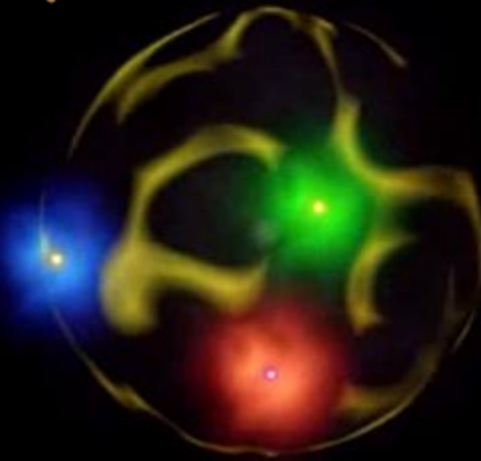


Le Matin
CERN
Le boson
de Higgs
fêté comme
une rock-star

“ENTÃO A MASSA VEM TODA DO HIGGS?”

Protão $p = uud$: $2 m_u + m_d = 11 \text{ MeV}$

$$m_p = 938 \text{ MeV}$$



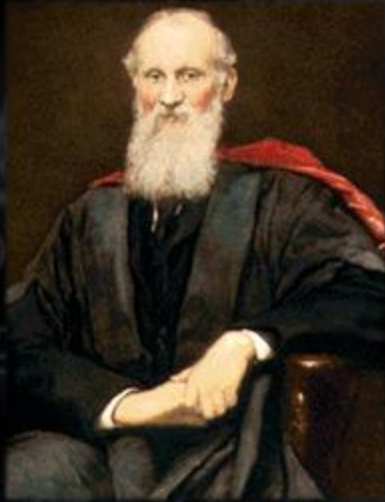
Só 1% da massa do protão é devida à massa em repouso dos quarks, ou seja...

Apenas uma infima parte da massa é devida ao mecanismo de Higgs...

E AGORA?

“Nineteenth-Century Clouds over the Dynamical Theory of Heat and Light”

Lord Kelvin, 27 de Abril 1900



As núvens negras de Kelvin

Incapacidade de detectar o Éter e a
“Catástrofe ultra-violeta”

A Física estaria limitada à medição de quantidades conhecidas com grande precisão...

Kelvin não podia estar mais enganado...

Stephen Hawking (1998)

“Com a descoberta iminente do bosão de Higgs não há nada fundamentalmente novo a ser descoberto. Tudo o que há a fazer é medir com mais precisão.”

REPETIÇÃO DA HISTÓRIA?

**STEPHEN
HAWKING**

AUTOR DE
BREVE
HISTÓRIA
DO TEMPO

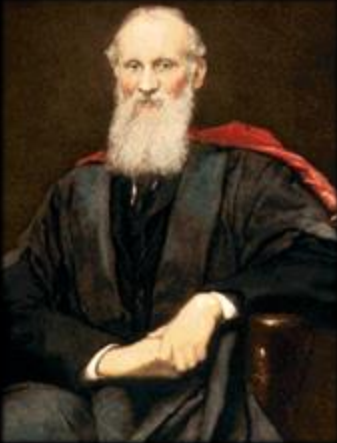


**O FIM
DA FÍSICA**

CONCEITOS
grátis

O QUE DIRIA KELVIN AGORA?

“Twentieth first-Century Clouds over the electroweak theory”



“A beleza e a clareza da teoria electrofraca está obscurecida por algumas núvens”

As núvens do Pedro:

- Matéria escura e energia escura
- Porque existe mais matéria que anti-matéria no Universo?
- Porquê 3 famílias?;
- Problema da Hierarquia;
- Porque é que as massas das partículas elementares são o que são;
- Porque é que os neutrinos são muito mais leves do que os leptões carregados e os quarks;
- Será que as 3 (ou 4) forças se unificam a alguma escala?;
- Será que as partículas elementares são mesmo elementares?;

PORQUE FÍSICA PARA ALÉM DO MP?

EVIDÊNCIAS EXPERIMENTAIS PARA A EXISTÊNCIA DE NOVA FÍSICA



FERMIÕES

Leptões spin = 1/2			Quarks spin = 1/2		
Sabor	Massa GeV/c ²	Carga Eléctrica	Sabor	Massa Aprox. GeV/c ²	Carga Eléctrica
ν_L neutrino* mais leve	$(0-2)\times 10^{-9}$	0	u up	0.002	2/3
e electrão	0.000511	-1	d down	0.005	-1/3
ν_M neutrino* intermedio	$(0.009-2)\times 10^{-9}$	0	c charm	1.3	2/3
μ muão	0.106	-1	s strange	0.1	-1/3
ν_H neutrino* pesado	$(0.05-2)\times 10^{-9}$	0	t top	173	2/3
τ tau	1.777	-1	b bottom	4.2	-1/3

Assimetria matéria-antimatéria

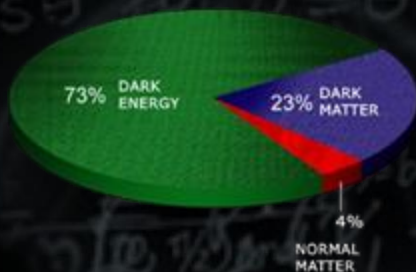
O MP falha em explicar porque razão existe um excesso de matéria no Universo ou “porque estamos de facto aqui!”

Massas de neutrinos

No MP os neutrinos não têm massa, mas de facto hoje sabemos que estas partículas são massivas (com uma massa muito menor que os restantes fermiões).

Problema da matéria escura

23% do budget de energia do Universo surge sob a forma de matéria escura. O MP não tem um candidato para a matéria escura.



1964: MISTÉRIO DOS NEUTRINOS DESAPARECIDOS



Bachall



N_{ν_e} N'_{ν_e}



Davis

$$\frac{N'_{\nu_e}}{N_{\nu_e}} = \frac{1}{3}$$

- O número de neutrinos que saem do Sol está mal calculado, OU
- A experiência está errada.

OSCILAÇÕES DE NEUTRINOS



$\nu_e \rightarrow \nu_\mu$

$$P = \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2}{4E} L\right)$$

$$\Delta m^2 = m_2^2 - m_1^2$$





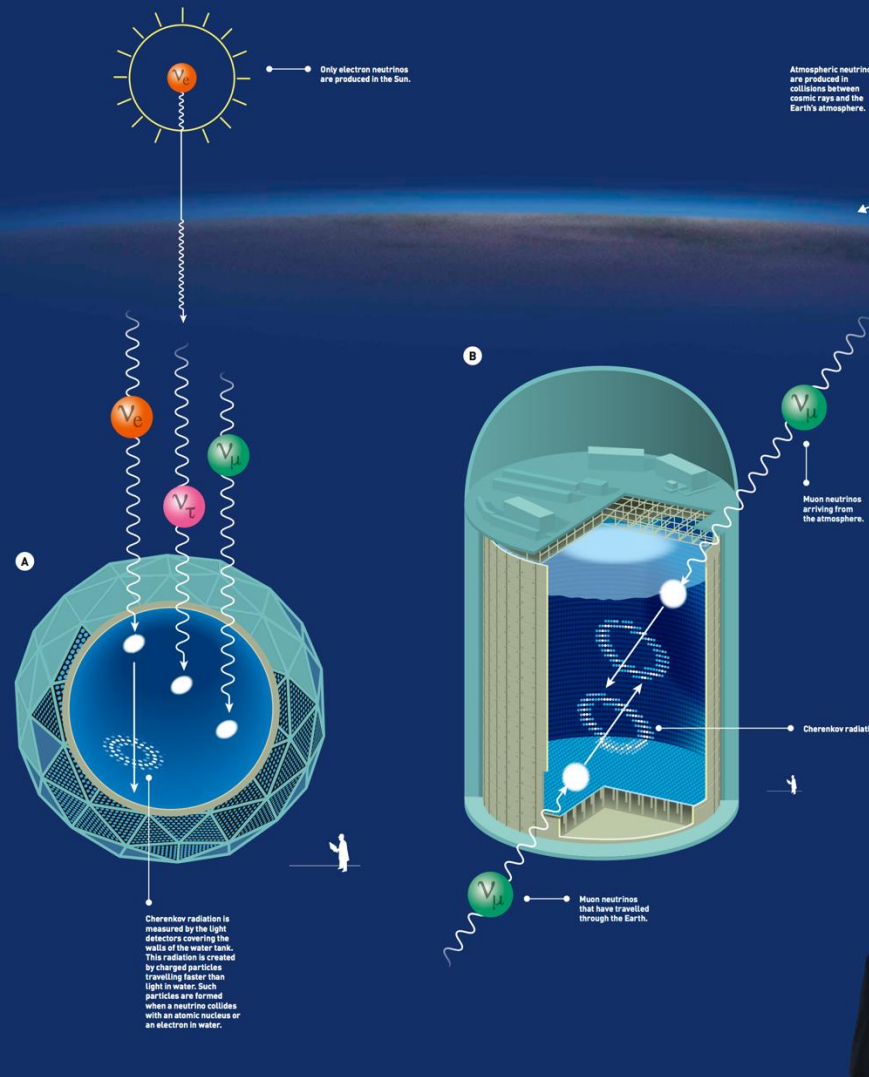
Chameleons of space



Takaaki Kajita in Japan and Arthur B. McDonald in Canada were key scientists in two large research groups that discovered that neutrinos change identities, which requires that neutrinos have mass. The discovery has changed our understanding of the innermost workings of matter and may prove crucial to our view of the universe.

The discovery of neutrino identity changes has resolved a neutrino puzzle that physicists had wrestled with for decades. Compared to theoretical calculations of the number of neutrinos, up to two-thirds of them were missing in measurements performed on Earth. The two research groups discovered that the neutrinos had changed identities, which led to the conclusion that neutrinos must have some mass, however small. This discovery was historic for particle physics, as its Standard Model requires neutrinos to be massless. Thus new physics is now needed. The Earth is constantly bombarded by neutrinos. Many are created in reactions

between cosmic radiation and the Earth's atmosphere. Others are produced in nuclear reactions inside the Sun. Thousands of billions of neutrinos stream through our bodies every second. The combined weight of neutrinos is estimated to be roughly equal to that of all visible stars in the universe. Hardly anything can stop the neutrinos; they are amongst nature's most elusive elementary particles. Experiments are continuing to uncover the all but hidden world of neutrinos. New discoveries about their deepest secrets are expected to change our current understanding of the history, structure and future of the universe.

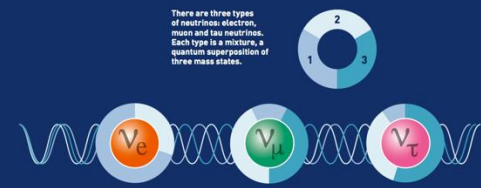


A Sudbury Neutrino Observatory
The detector measured neutrinos from the Sun. Its tank, filled with heavy water, was placed two kilometres under the surface of the Earth. Signals from all three types of neutrinos were registered in the tank. The sum of the neutrinos corresponded to what was expected, but there were not enough electron neutrinos – they must have changed identity.

B Super-Kamiokande
The detector measured atmospheric neutrinos. Its tank, filled with water, was placed one kilometre under the surface of the Earth. The muon neutrinos that arrived straight at Super-Kamiokande from the atmosphere were more numerous than those that arrived at the detector after passing through the Earth. The muon neutrinos that travelled further thus had time to change identity and become another type of neutrino.

Arthur B. McDonald
Canadian citizen. Born 1943 in Sydney, Canada. Professor Emeritus at Queen's University, Kingston, Canada.

Takaaki Kajita
Japanese citizen. Born 1959 in Higashimatsuyama, Japan. Director of Institute for Cosmic Ray Research and Professor at University of Tokyo, Kashiwa, Japan.



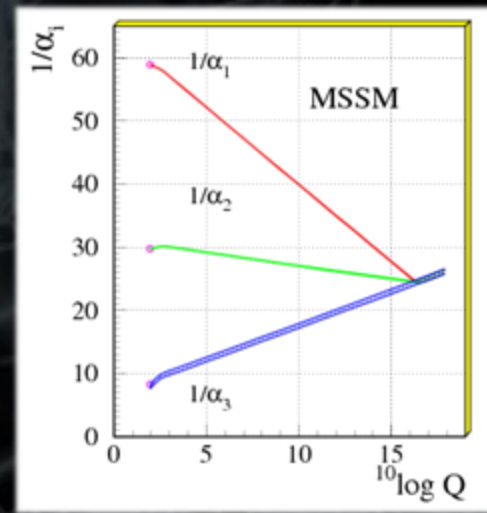
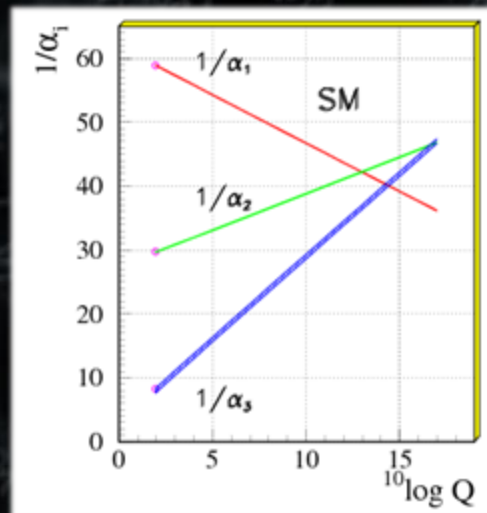
Neutrino oscillations
Neutrinos change identities as they travel through space. Quantum physics is required to explain this magic, where the neutrinos are represented by superposed waves that correspond to neutrino states with different masses. When the neutrinos travel, these waves go out of phase and are superposed in different ways. The superposition in any given location yields the probability of which type of neutrino is most likely to be found there. These probabilities vary from one location to another – oscillate – and the neutrinos appear in their various identities. This is only possible if neutrinos have mass.

UNIFICAÇÃO DAS INTERAÇÕES FUNDAMENTAIS

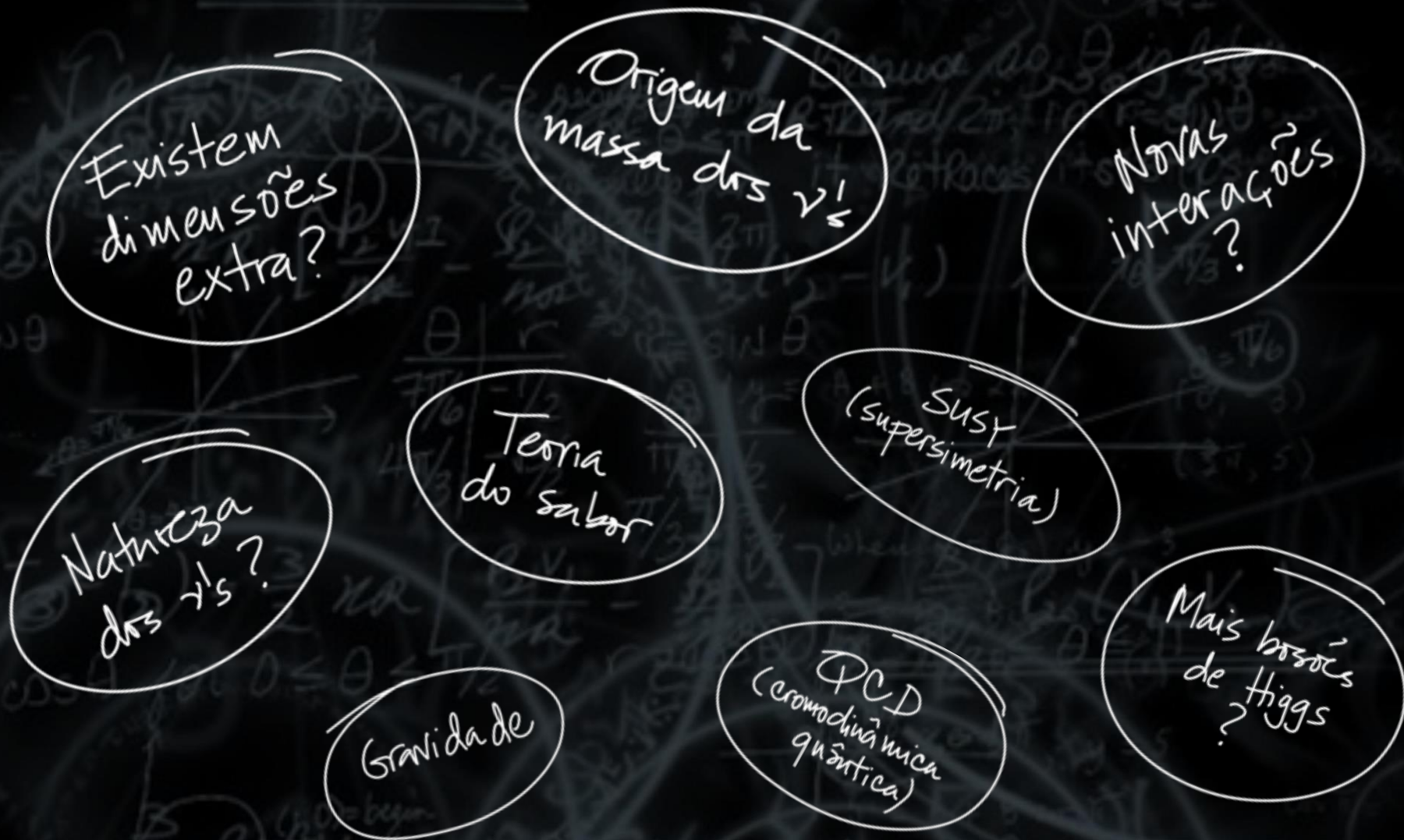
UNIFICAÇÃO DAS INTERAÇÕES FUNDAMENTAIS

Propriedade	Interação Gravítica	Interação Fraca (Electrofraca)	Interação Electromagnética	Interação Forte
Actua em:	Massa – Energia	Sabor	Carga Eléctrica	Carga de cor
Partículas afectadas:	Todas	Quarks, Leptões	Electricamente carregadas	Quarks, Gluões
Partículas mediadoras:	Gravitão (ainda por observar)	W^+ W^- Z^0	γ	Gluões
Intensidade a $\left\{ \begin{array}{l} 10^{-18} \text{ m} \\ 3 \times 10^{-17} \text{ m} \end{array} \right.$	10^{-41}	0.8	1	25
	10^{-41}	10^{-4}	1	60

SE O MP FOR VÁLIDO ATÉ À ESCALA DE PLANCK AS 3 (4) FORÇAS FUNDAMENTAIS UNIFICAM-SE OU NÃO?



PORQUE FÍSICA PARA ALÉM DO MP?



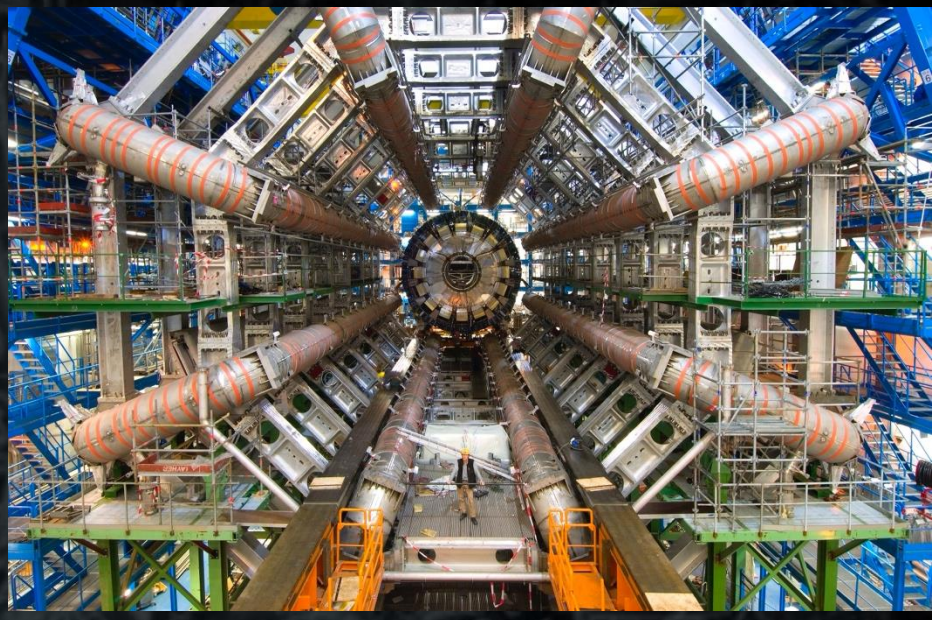
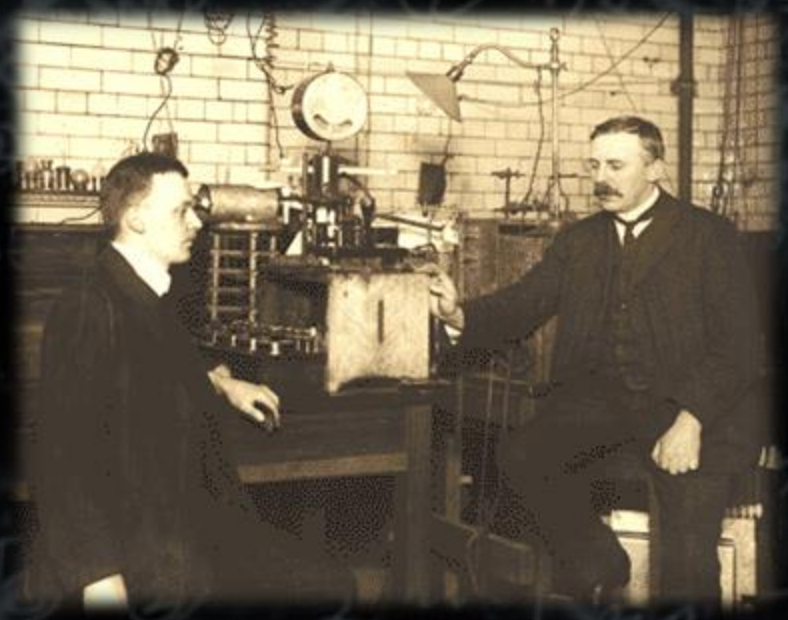


“We are reaching into the fabric of the Universe at the level never done before... We are in the edge of a new exploration.”

Joe Incandela, CMS spokesperson, Press Conference, July 4, 2012

$$\underline{P_2 (V_2 - V_1)}$$

Because $\rho_2 > \rho_1$ before

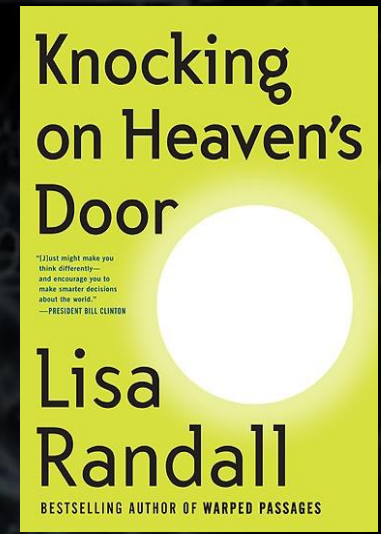
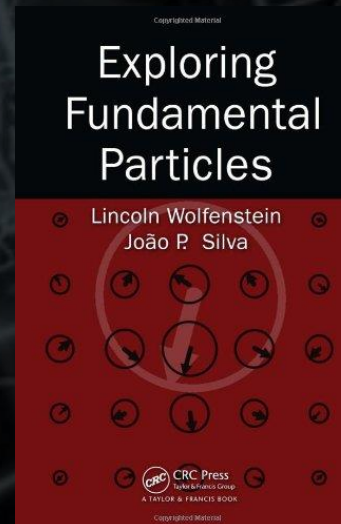
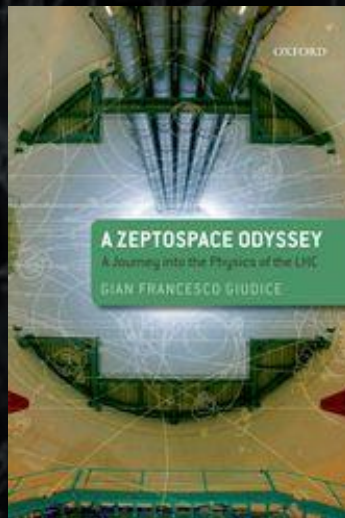
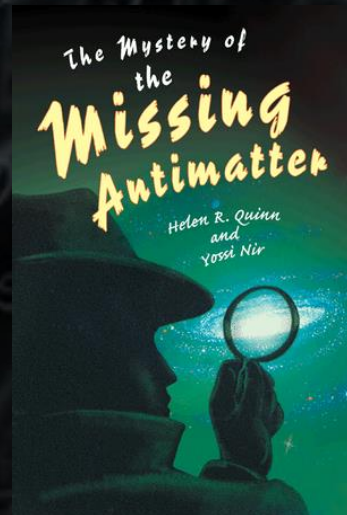
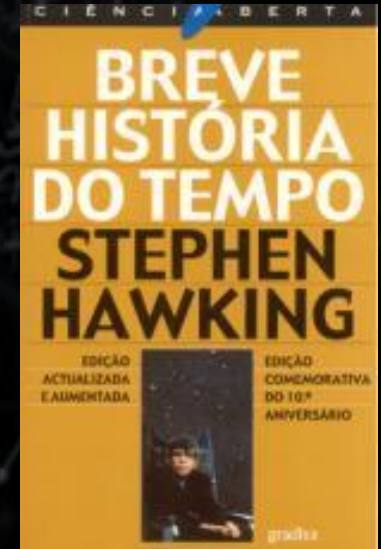
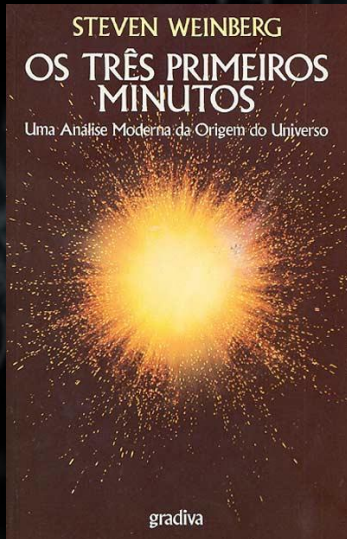


$$\cos \theta \quad \text{for } 0 \leq \theta \leq \pi/2$$

$(1,0) = \text{begin}$
 $\frac{1}{2} \pi \leq \theta \leq \pi$
end

scope, for $0 \leq \theta \leq \pi$
 $\theta = \pi$

ALGUNS LIVROS (Os mais actuais em inglês)



<https://www.facebook.com/PhysicsTecnico>

https://www.instagram.com/physics_at_tecnico/



filipe.joaquim@tecnico.ulisboa.pt