

Listening for ultra-heavy DM with underwater acoustic detectors

Damon Cleaver, Christopher McCabe and Ciaran A.J. O'Hare NeXT Workshop, 15/07/24

What is Dark Matter?

[Planck, 2018]



$\Omega_{\rm CDM} h^2 \sim 0.120 \pm 0.001$

What (generic) DM properties do we need for DD?

Local DM Density

$\rho_{\chi} \sim 0.3 \text{ GeV/cm}^3$

Local DM Velocity

 $v_{\gamma} \sim 220 \text{ km/s}$

In truth, follows a (boosted) Maxwell Boltzmann distribution with peak v_{γ}

Ultra-heavy Dark Matter

- Ultra-heavy dark matter is necessarily composite (if thermally produced) due to s-wave unitarity
- Many different models for UHDM
 - PBH, Nuggets, Blobs, WIMPonium, Q-Balls etc...
- Is there a nice model-independent way to treat them?
- Answer: **Yes** (for some parts of parameter space)



- Consider parameters of models where:
 - The DM is Planck-mass or larger
 - DM Radius R_{γ} much larger than interaction length scale
 - Geometric cross section dominate
 - Parameterise the interaction in terms of R_{χ} , -> set by the theory -> make experimental statements about multiple models!

Macros

es i.e.
$$\sigma_{\chi} \approx \pi R_{\chi}^2$$



Interaction range



Macro Direct Detection

• Macro's often parameterised in grams (g)

DM Flux:
$$\phi_{\chi} \approx 6 \left(\frac{1 \text{ g}}{m_{\chi}}\right) \text{ km}^{-2} \text{ y}$$

- Need a *very* large detector (or very long integration time) to have significant number of events.
- No hope for conventional detectors (LUX-ZEPLIN, XENONnT etc.)

′r⁻¹

Macro Kinematics

$$E_{\chi}(L) =$$

With the DM final velocity after N collisions:

$$v_{\chi f} \simeq v_{\chi i}$$

$$\simeq v_{\chi i}$$

The energy of the macro after after traversing a path length L is given by:

$$\frac{1}{2}m_{\chi}v_{\chi f}(L)^2$$



Many collisions needed to slow down for $m_A \ll m_{\gamma}$

Macro Kinematics

Using the relation for the mean free path $\lambda = m_A / (\rho_{med} \sigma_{\gamma N})$ and the number of scatters as $N = L/\lambda$ yields:

 $v_{\chi f} \simeq v_{\chi i} e$

mass column density yields sea level velocity $v_{\gamma,SL}$:

$$v_{\chi,\text{SL}} \simeq v_{\chi i} \exp\left(-\frac{\sigma_{\chi N}}{m_{\chi}}X_0 \sec\theta\right)$$

$$\exp\left(-\frac{\sigma_{\chi N}}{m_{\chi}}\rho_{\rm med}L\right)$$

This is only for a constant density medium, for the atmospheric overburden, use

$$X(h) = \int_{h}^{\infty} dh' \rho_0 e^{-h'/H} = X_0 e^{-h/H}$$

Macro Kinematics

Total energy after traversing path length L in detector given by:

$$E_{\chi f}(L) = \frac{1}{2} m_{\chi} v_{\chi,SL}^2 \exp\left(-2\frac{\sigma_{\chi N}}{m_{\chi}}\rho_{\rm med}L\right)$$

Differentiating, can find the energy deposition rate into the medium:

$$\frac{dE}{dL} = -\frac{dE_{\chi}}{dL} = \rho_w \sigma_{\chi N} v_{\chi,\rm SL}^2 \exp\left(-2\frac{\sigma_{\chi N}}{m_{\chi}}\rho_{\rm med}L\right)$$

Current Constraints



- Annoying "gap" in constraints
- Mica underground too much overburden
- Radar not sensitive enough - not enough ionisation
- What phenomena could we use to constrain this region?

Acoustic Detection

- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium:
 the ocean
- DM deposits energy into the ocean creating pressure waves
- Detect pressure waves using a large hydrophone array





Neutrino Experiments

- Propositions for acoustic neutrino experiments with $\mathcal{O}(100~{\rm km^3})$ hydrophone arrays [Lahmann, 2016]
- Detect UHE neutrinos. Similar number density issues, but similarly high cross section
- Acoustic propagation distance in water much greater than light -> less dense instrumentation
- Energy deposition comes from hadronic showers

Proposed Acoustic Neutrino Experiment ~4km wide LUX-ZEPLIN (LZ) **1.5**m **1.5m**





[Learned, 1979]

$$\nabla^2 P(\mathbf{r}, t) - \frac{1}{c_s^2} \frac{\partial^2 P(\mathbf{r}, t)}{\partial t}$$
Acoustic pressure

General solution to this equation given by:

$$P(\vec{r},t) = \frac{\alpha}{4\pi C_p} \int_V \frac{dV'}{|\vec{r} - \vec{r'}|} \frac{\partial^2}{\partial t^2} q\left(\vec{r'},t'\right), \quad t' = t - \frac{|\vec{r} - \vec{r'}|}{c_s}$$

Pressure waves created from thermo-acoustic heating.



Model the energy deposition as instantaneous, while heat dissipation as slow

$$q(\vec{r},t) = q(\vec{r})\Theta(t-t_0) \implies \partial_t q(\vec{r},t) = q(\vec{r})\delta(t-t_0)$$

Decompose surfaces into constant propagation time surfaces

$$P(\vec{r},t) = \frac{\alpha}{4\pi C_p} \int_V \frac{dV'}{|\vec{r} - \vec{r'}|} \frac{\partial}{\partial t} q(\vec{r'}) \delta(t' - t_0),$$

$$= \frac{c_s^2}{4\pi} \frac{\alpha}{C_p} \frac{\partial}{\partial R} \int_{S_{\vec{r}}^{R_{t_0}}} \frac{dS'}{R_{t_0}} q(\vec{r'}),$$

15



Model energy deposition rate as:

$$q(x, y, z) = \frac{dE_w}{dz} \frac{1}{(\lambda\sqrt{2\pi})^2} \exp\left(-\frac{1}{2}\frac{x^2 + \lambda^2}{\lambda^2}\right)$$

Substituting macro result for dE/dz

$$q(x, y, z) = \frac{\rho_w \sigma_{\chi N} v_{\chi i}^2}{(\lambda \sqrt{2\pi})^2} \exp\left(-\frac{1}{2} \frac{x^2 + y^2}{\lambda^2}\right)$$

 $\frac{y^2}{}$







In far field, spherical surfaces approximated as y-z planes

$$P(x_0\vec{e_x},t)$$

Align detector element along xaxis at position X_0





Solving in this case we find the following time-domain signal:

$$P(x_0 \vec{e_x}, t) = -\frac{\gamma_G}{4\pi x_0} \frac{\Delta E_{\chi}}{\sqrt{2\pi}} \frac{c_s t}{\lambda^3} \exp\left(-\frac{1}{2} \frac{(c_s t)^2}{\lambda^2}\right)$$

Which has frequency content (Fourier transform)

$$\tilde{P}(\omega) = i\omega \frac{\gamma_G}{4\pi c_s^2} \frac{\Delta E_{\chi}}{r_0} \exp\left(-\frac{\lambda^2 \omega^2}{2c_s^2}\right)$$

Noise and Transmission Loss

The pressure wave will have some **frequency dependent absorption** due to the chemical content of the sea water

Convenient to use a **decibel formalism**, as this is most often used in acoustics. Sound pressure level:

SPL = 20

Time Domain

$$P_{\rm ref} = 1 \mu P a$$

$$\log_{10}\left(\frac{P}{P_{\rm ref}}\right)$$

Frequency Domain

$$P_{\rm ref} = 1\mu Pa/Hz$$

Transmission Loss

The Transmission loss over a propagation distance r, is given by a frequency dependent parameter α

$TL = \alpha r$

[Ainslie, McColm 1998]

Dependent on the temperature T, salinity S, depth D and pH. Frequency here in kHz. After transmission losses, DM signal peaks at ~15kHz

$$\begin{split} \alpha &= 0.106 \frac{f_1 f^2}{f^2 + f_1^2} e^{(\text{pH}-8)/0.56} \\ &+ 0.52 \left(1 + \frac{T}{43}\right) \left(\frac{S}{35}\right) \frac{f_2 f^2}{f^2 + f_2^2} e^{-D/6} \\ &+ 0.00049 f^2 e^{-(T/27 + D/17)}, \end{split}$$

Sea State Noise

(surface agitation due to wind etc). Parameterised by Knudsen curves:

$NL(f, n_s) =$

There will also be transient noise sources from ocean wildlife e.g. dolphins and sperm wales. Assume differentiable from DM signal using algorithms being developed for neutrino detection.

The dominant background noise source 10-100 kHz band is sea state noise

$$n_s - 17 \log_{10}(f)$$

Signal and Noise Characteristics



- The signal is **broadband**
- After transmission losses, signal power greatest in 10-30kHz band
- Dominant ambient noise in this band is sea-state noise (surface agitation from wind)



Discerning Neutrinos from DM



- DM is galactic in origin. Neutrinos are extragalactic.
- Direction of dark matter in the direction of the Cygnus constellation.
- So UHDM flux modulates
- Varies throughout the year due to following the sidereal day over solar day



Constraint Recipe

- Assume DM cross section and mass
- Impose SNR requirement, hydrophone sensitivity and flux requirement
- Assume array of particular geometry (volume, cross sectional area, hydrophone density)

• Enhance signal by $\sqrt{N_{\rm hydro}}$ (Poisson statistics)

- If DM signal amplitude fulfils criteria, **CONSTRAIN**!



Preliminary Sensitivities

- Assuming proposed acoustic neutrino experiment parameters, could constrain the gap!
- Complementary to Humans, Mica, Ohya and Cosmological Bounds

Punchline:

Future acoustic neutrino experiments could have the power to constrain macro DM candidates

Thank you for listening! Any Questions?