

# AXEL-2023

## Introduction to Particle Accelerators

### Longitudinal motion:

- ▣ *The basic synchrotron equations.*
- ▣ *What is Transition ?*
- ▣ *RF systems.*
- ▣ *Motion of low & high energy particles.*
- ▣ *Acceleration.*
- ▣ *What are Adiabatic changes?*

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# Motion in longitudinal plane

- # What happens when particle momentum increases?
  - ⇒ particles follow longer orbit (fixed B field)
  - ⇒ particles travel faster (initially)
- # How does the revolution frequency change with the momentum ?

$$\frac{df}{f} = \frac{dv}{v} - \frac{dr}{r}$$

Change in velocity

Change in orbit length

But

$$\frac{\Delta r}{r} = \alpha_p \frac{\Delta p}{p}$$

Momentum compaction factor

Therefore:

$$\frac{df}{f} = \frac{dv}{v} - \alpha_p \frac{dp}{p}$$

# The frequency - momentum relation

$$\frac{df}{f} = \frac{dv}{v} - \alpha_p \frac{dp}{p}$$

But

$$\frac{dv}{v} = \frac{d\beta}{\beta} \quad \left( \beta = \frac{v}{c} \right)$$

# The relativity theory says:

$$p = \frac{E_0 \beta \gamma}{c}$$

$$\frac{dv}{v} = \frac{d\beta}{\beta} = \frac{1}{\gamma^2} \frac{dp}{p}$$

$$\frac{dp}{p} = \gamma^2 \frac{d\beta}{\beta}$$

$$\frac{dp}{d\beta} = \frac{E_0 \gamma^3}{c}$$

$$\frac{df}{f} = \left( \frac{1}{\gamma^2} - \alpha_p \right) \frac{dp}{p}$$

varies with momentum  
( $\mathbf{E} = \mathbf{E}_0 \gamma$ )

fixed by the lattice

# Transition

# Lets look at the behaviour of a particle in a constant magnetic field.

# Low momentum ( $\beta \ll 1, \gamma \Rightarrow 1$ )  $\longrightarrow$   $\frac{1}{\gamma^2} > \alpha_p$

# The revolution frequency increases as momentum increases

# High momentum ( $\beta \approx 1, \gamma \gg 1$ )  $\longrightarrow$   $\frac{1}{\gamma^2} < \alpha_p$

# The revolution frequency decreases as momentum increases

# For one particular momentum or energy we have:

$$\frac{1}{\gamma^2} = \alpha_p$$

# This particular energy is called the Transition energy

# The frequency slip factor

# We found

$$\frac{df}{f} = \left( \frac{1}{\gamma^2} - \alpha_p \right) \frac{dp}{p} = \left( \frac{1}{\gamma^2} - \frac{1}{\gamma_{tr}^2} \right) \frac{dp}{p}$$

#  $\frac{1}{\gamma^2} > \alpha_p \longrightarrow$  Below transition  $\longrightarrow \eta = \text{positive}$

#  $\frac{1}{\gamma^2} = \alpha_p \longrightarrow$  Transition  $\longrightarrow \eta = \text{zero}$

#  $\frac{1}{\gamma^2} < \alpha_p \longrightarrow$  Above transition  $\longrightarrow \eta = \text{negative}$

$\eta$

# Transition is very important in proton machines.

■ A little later we will see why....

# In the PS machine :  $\gamma_{tr}$  is at  $\sim 6 \text{ GeV}/c$

# In the LHC machine :  $\gamma_{tr}$  is at  $\sim 55 \text{ GeV}/c$

# Transition does not exist in leptons machines, why?

# Radio Frequency System

## # Hadron machines:

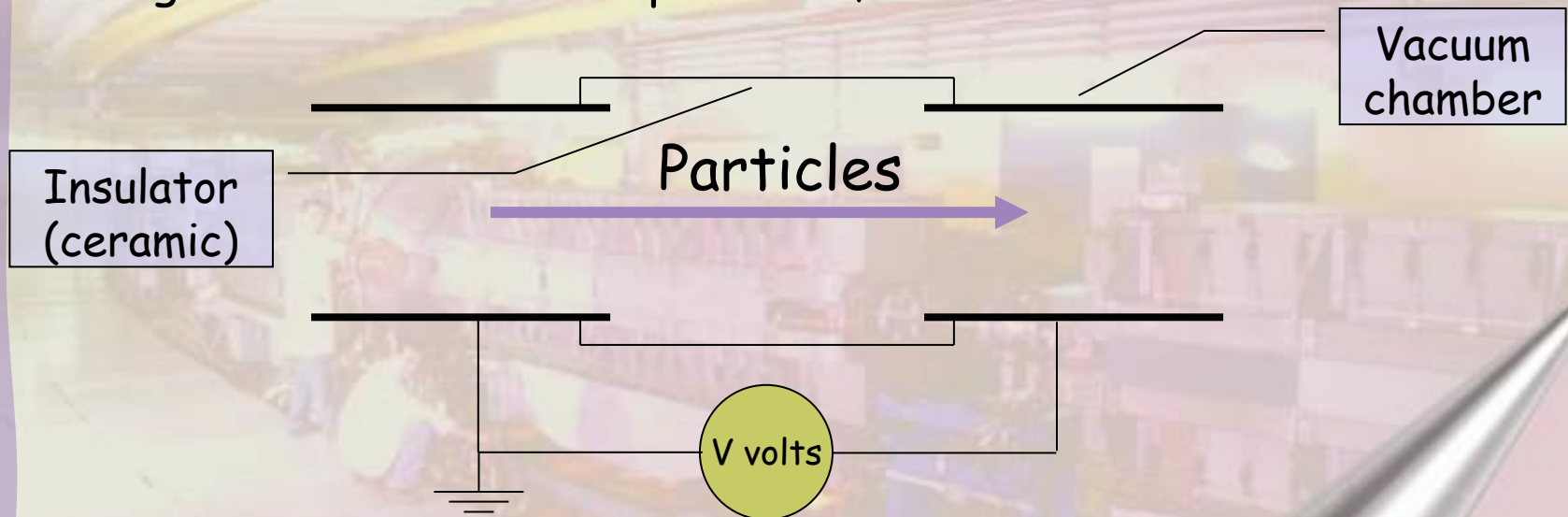
- ▣ Accelerate / Decelerate beams
- ▣ Beam shaping
- ▣ Beam measurements
- ▣ Increase luminosity in hadron colliders

## # Lepton machines:

- ▣ Accelerate beams
- ▣ Compensate for energy loss due to synchrotron radiation.

# RF Cavity

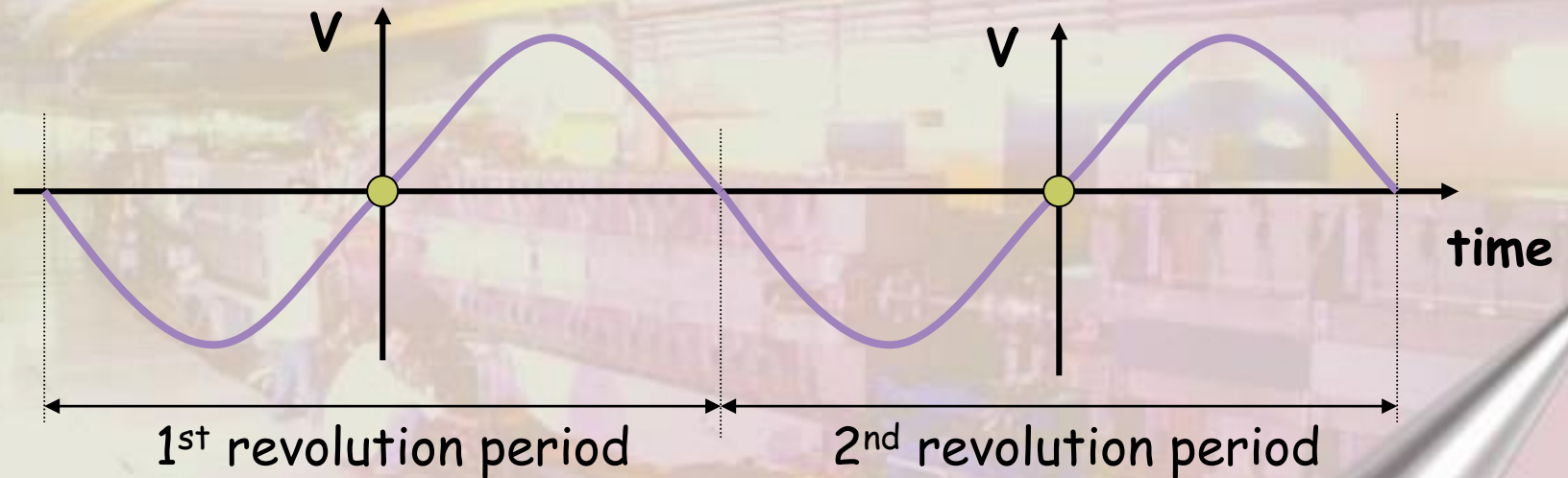
- # To accelerate charged particles we need a longitudinal electric field.
- # Magnetic fields deflect particles, but do not accelerate them.



- # If the voltage is DC then there is no acceleration !
  - The particle will accelerate towards the gap but decelerate after the gap.
- # Use an Oscillating Voltage with the right Frequency

# A Single particle in a longitudinal electric field

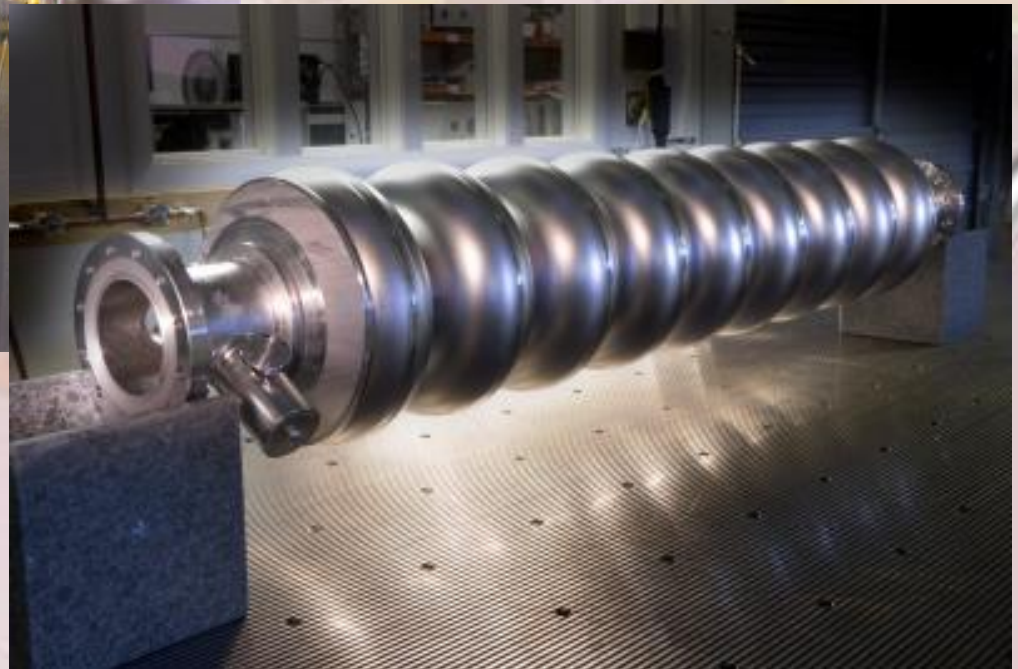
- # Lets see what a low energy particle does with this oscillating voltage in the cavity.



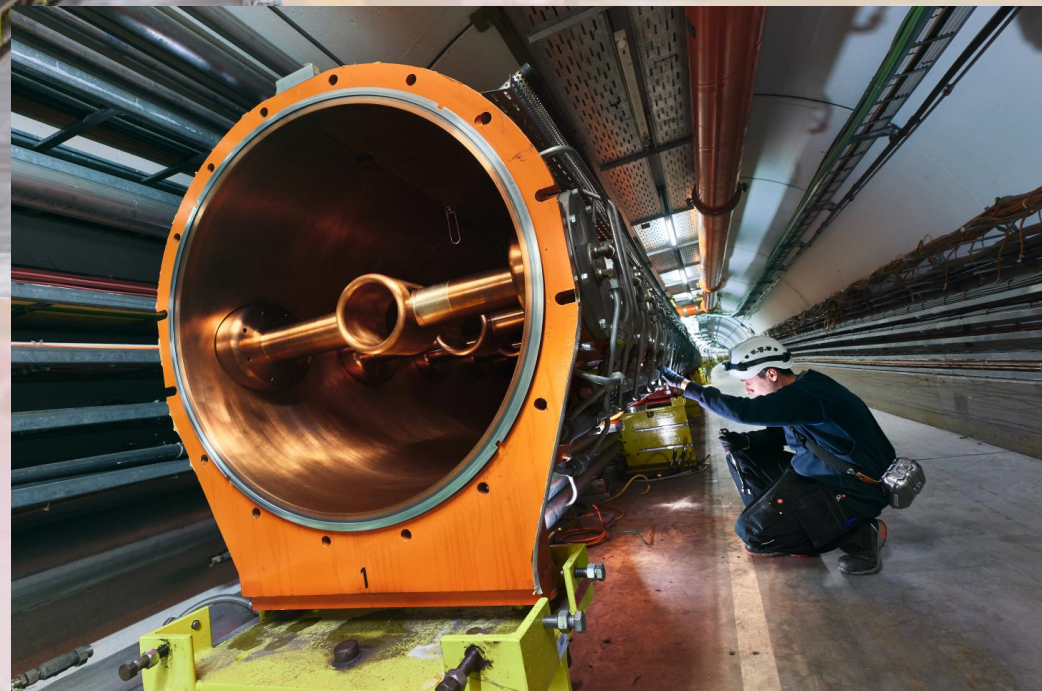
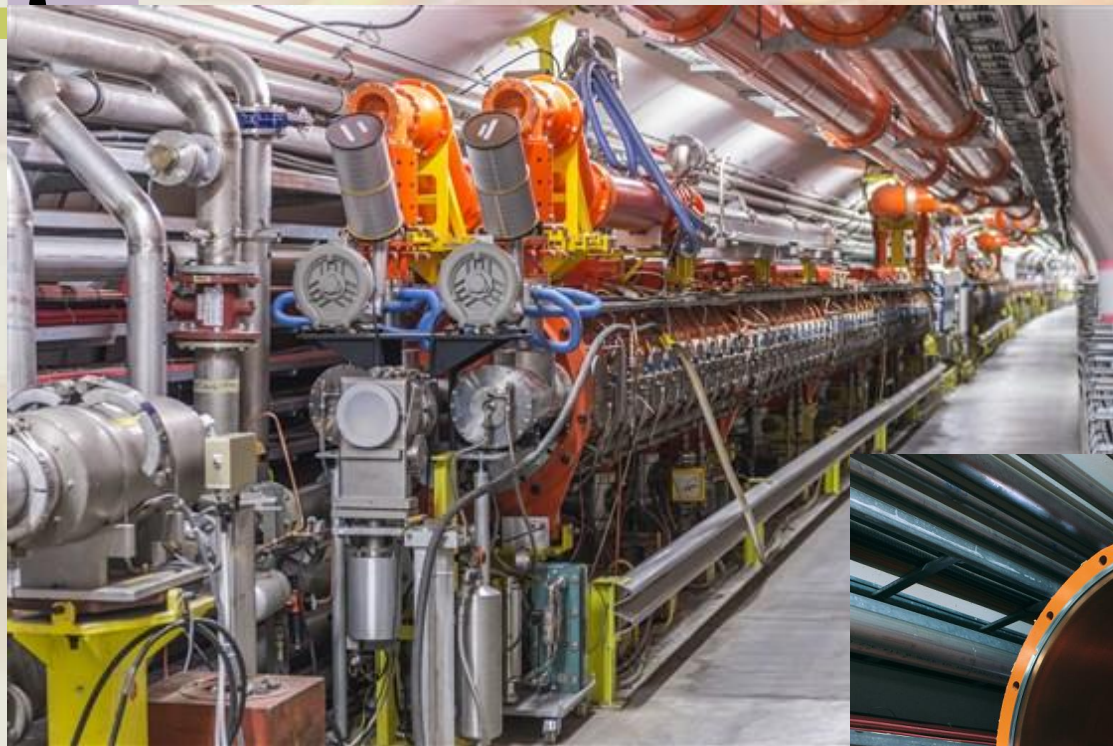
- # Set the oscillation frequency so that the period is exactly equal to one revolution period of the particle.



# LHC RF Cavities

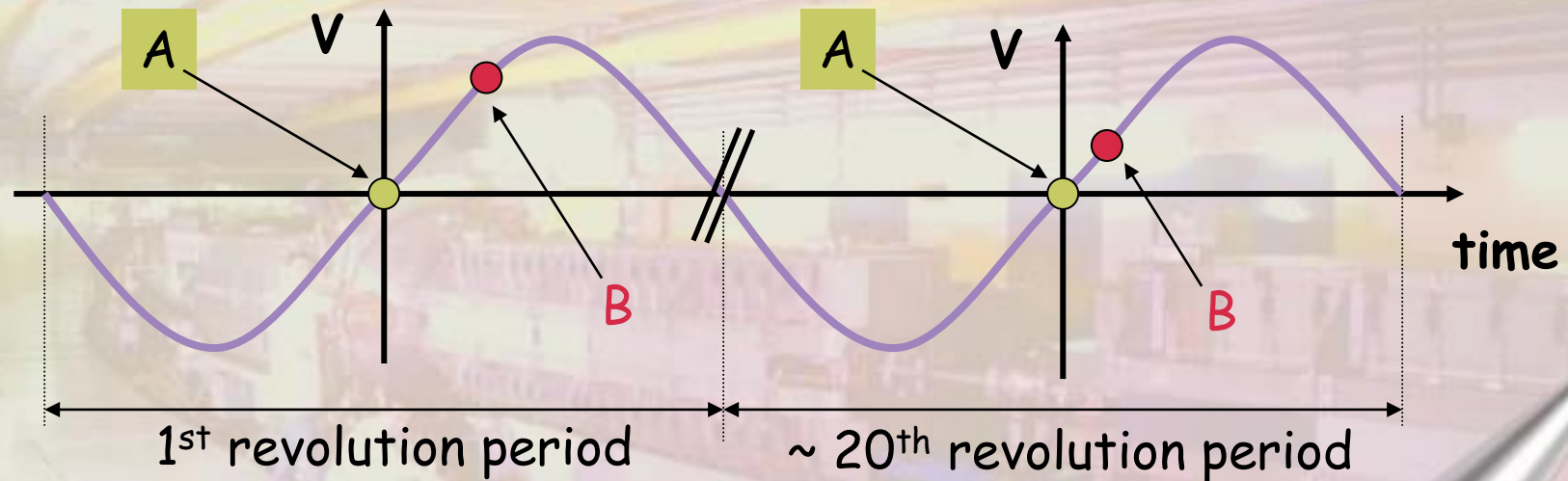


# SPS RF Cavities



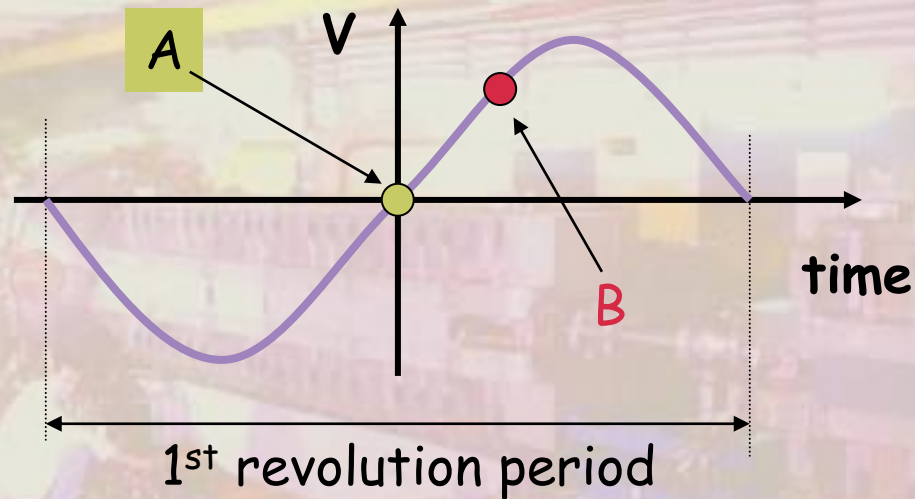
# Add a second particle to the first one

- # Lets see what a second low energy particle, which arrives later in the cavity, does with respect to our first particle.

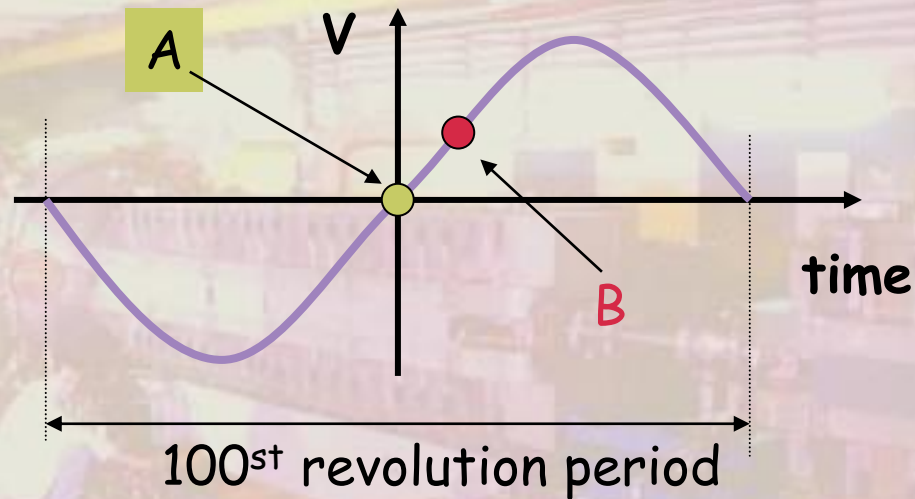


- # **B** arrives late in the cavity w.r.t. **A**
- # **B** sees a higher voltage than **A** and will therefore be accelerated
- # After many turns **B** approaches **A**
- # **B** is still late in the cavity w.r.t. **A**
- # **B** still sees a higher voltage and is still being accelerated

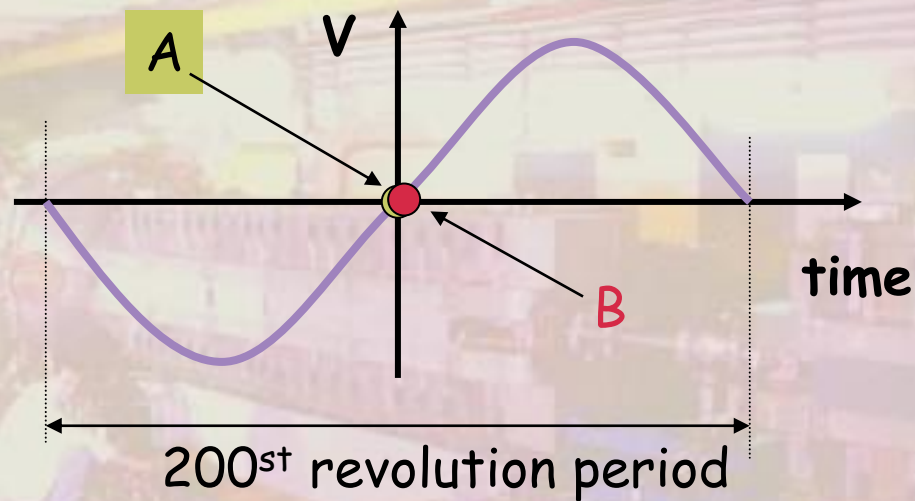
# Lets see what happens after many turns



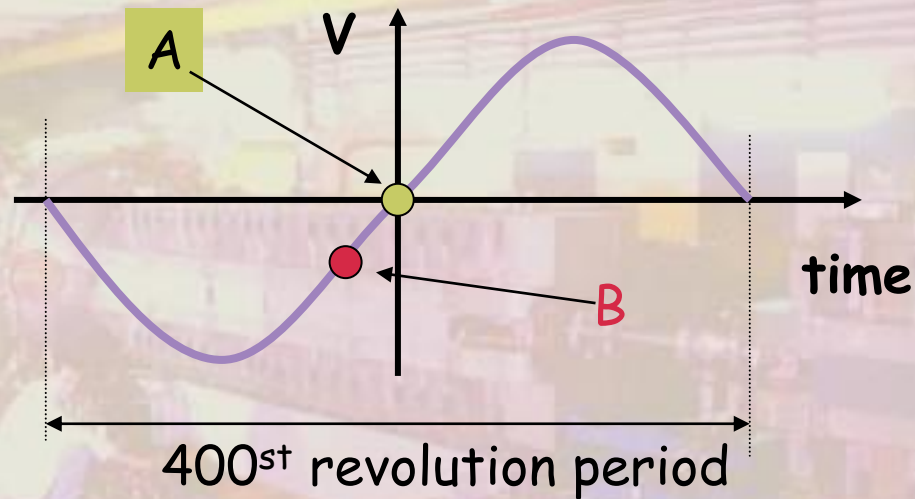
# Lets see what happens after many turns



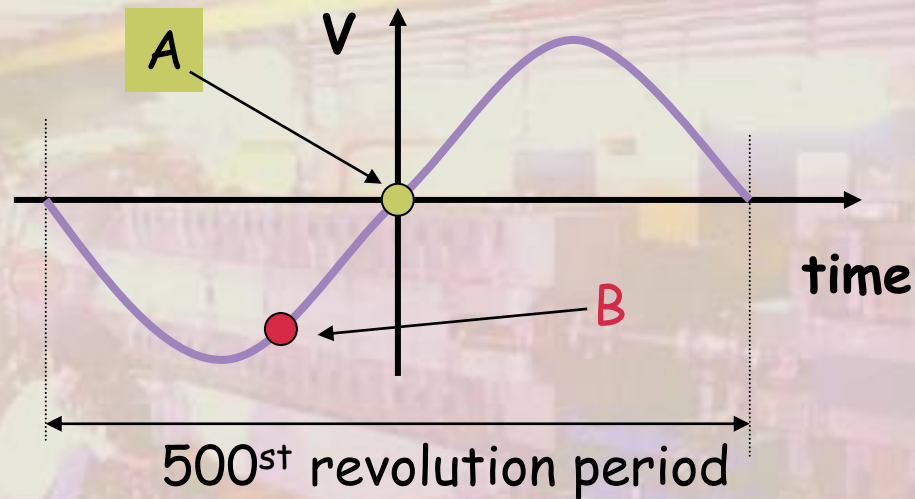
# Lets see what happens after many turns



# Lets see what happens after many turns

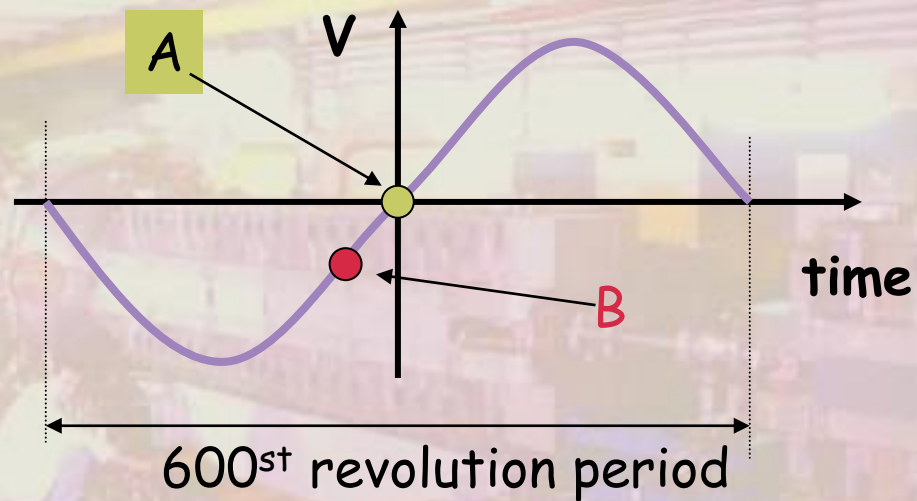


# Lets see what happens after many turns

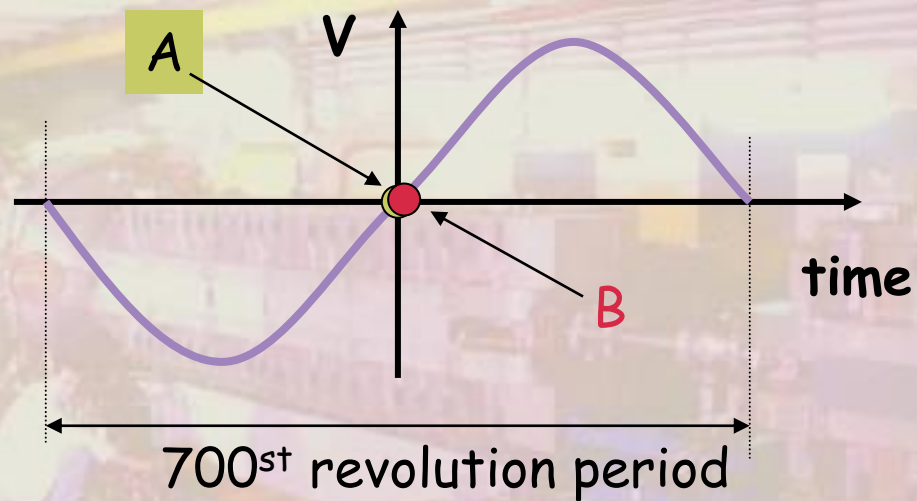




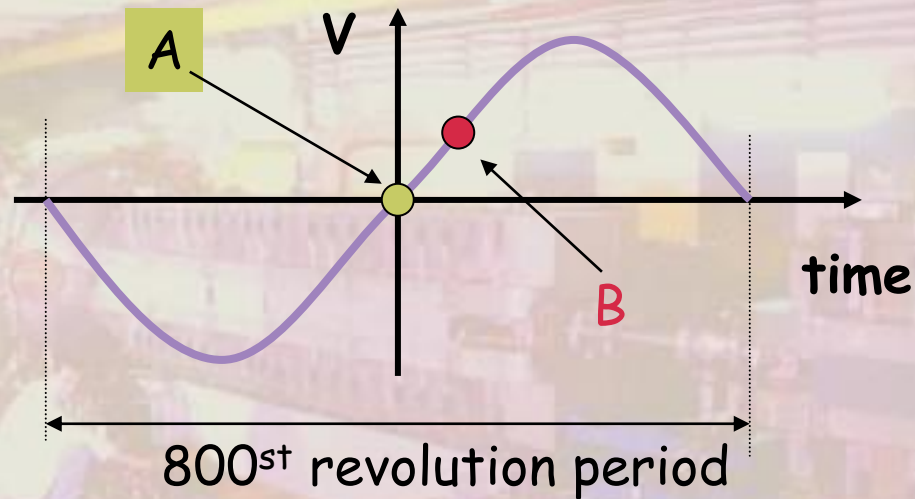
# Lets see what happens after many turns



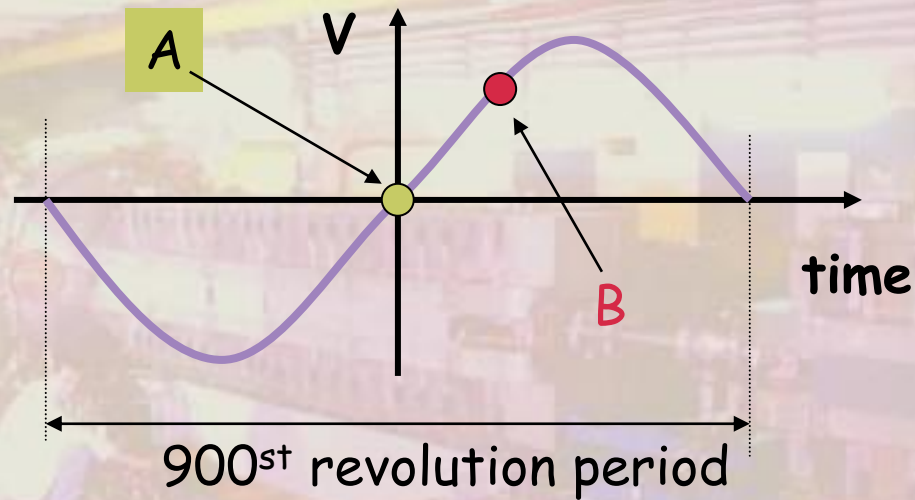
# Lets see what happens after many turns



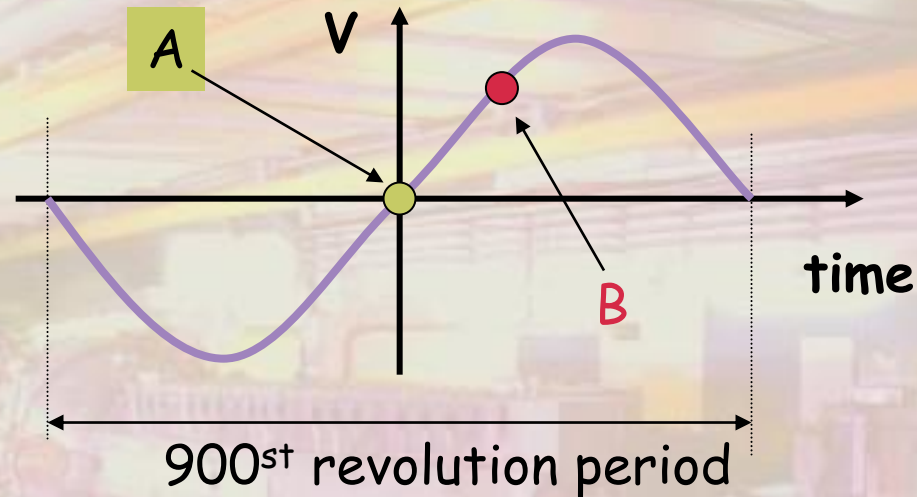
# Lets see what happens after many turns



# Lets see what happens after many turns



# Synchrotron Oscillations



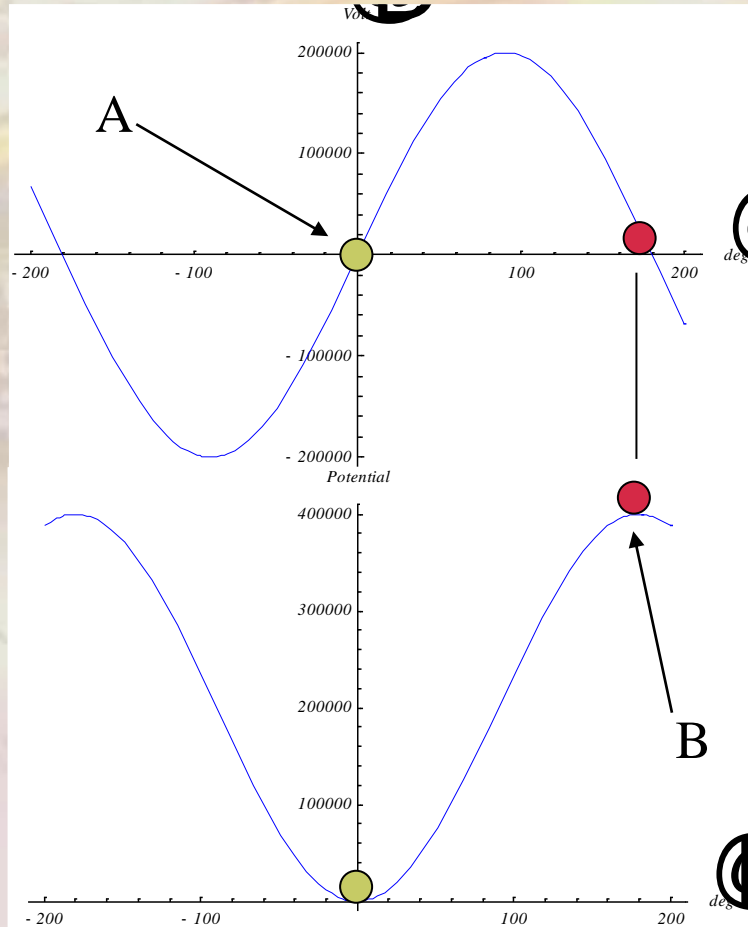
- # Particle B has made 1 full oscillation around particle A.
- # The amplitude depends on the initial phase.

Exactly like the pendulum

- # We call this oscillation:

Synchrotron Oscillation

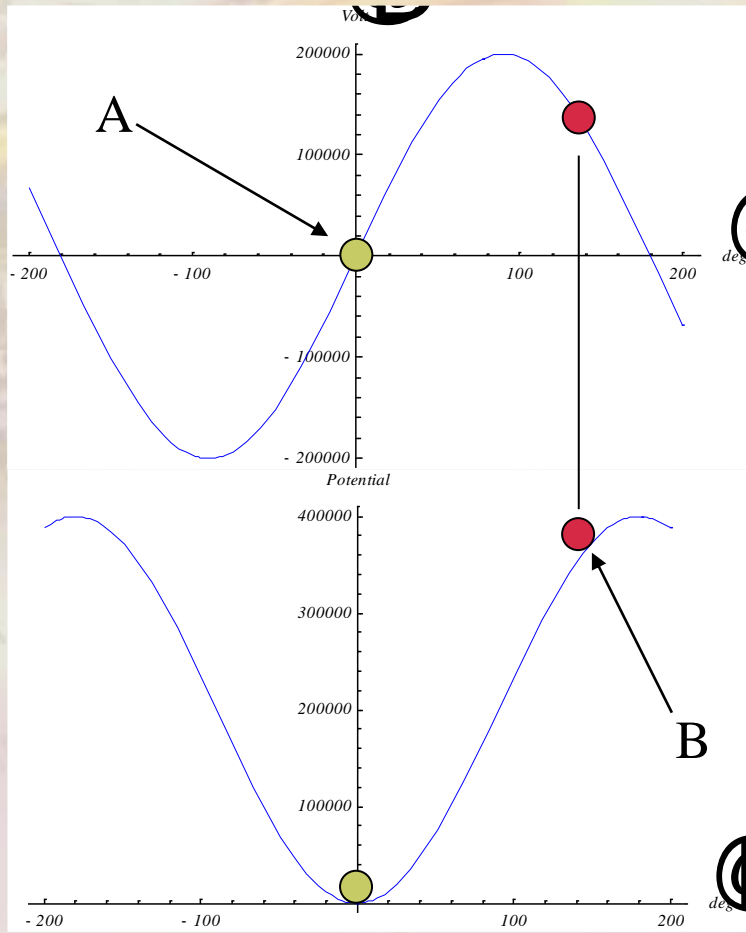
# The Potential Well (1)



Cavity voltage

Potential well

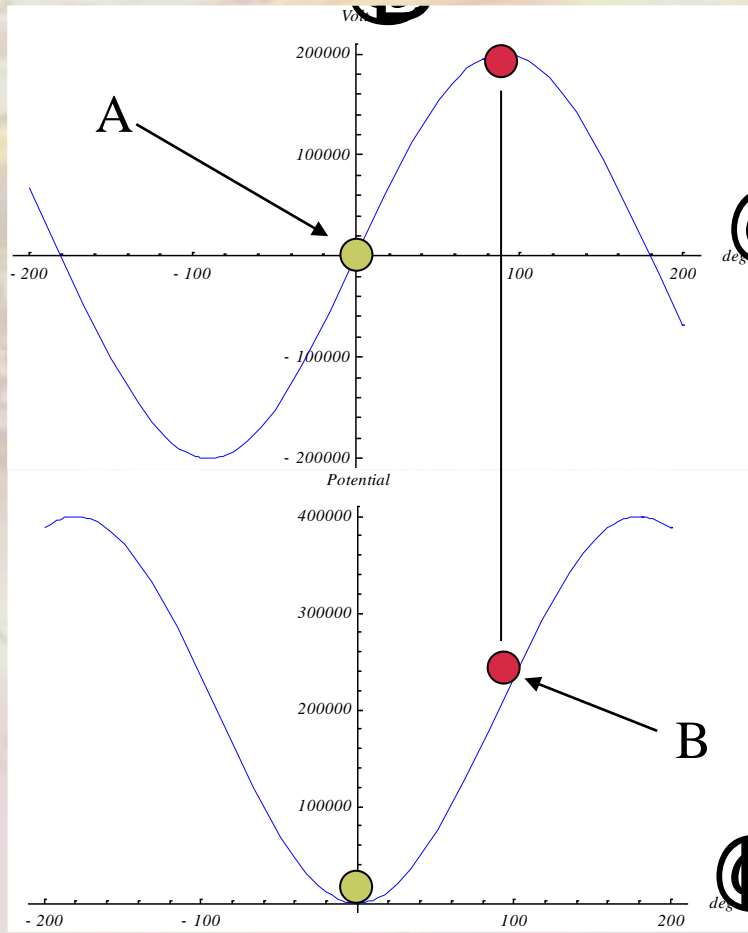
# The Potential Well (2)



Cavity voltage

Potential well

# The Potential Well (3)

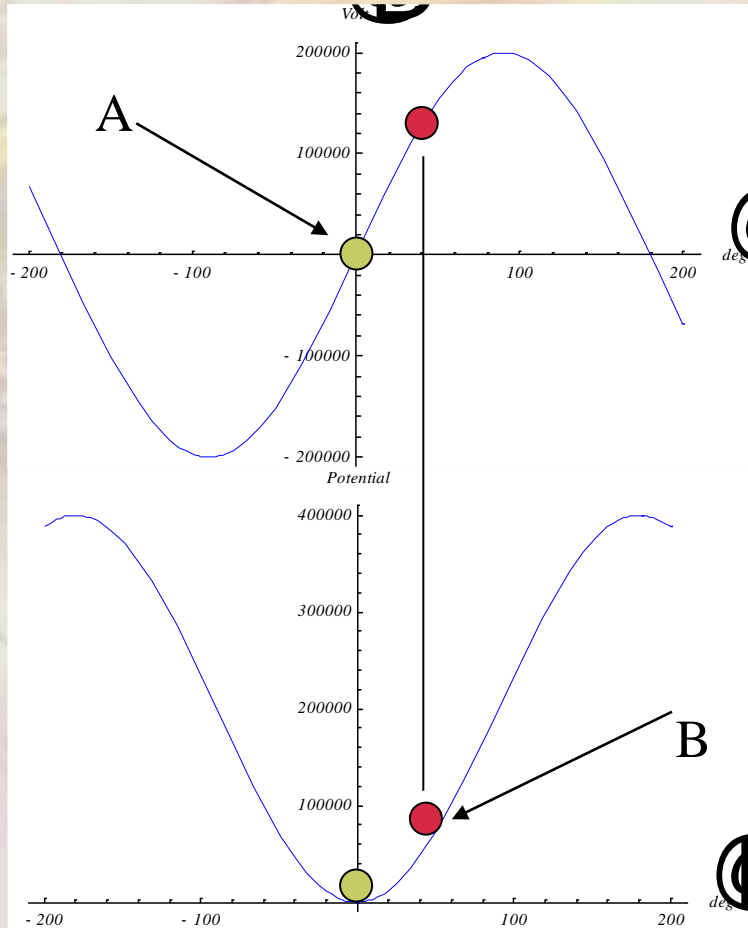


Cavity voltage

Potential well



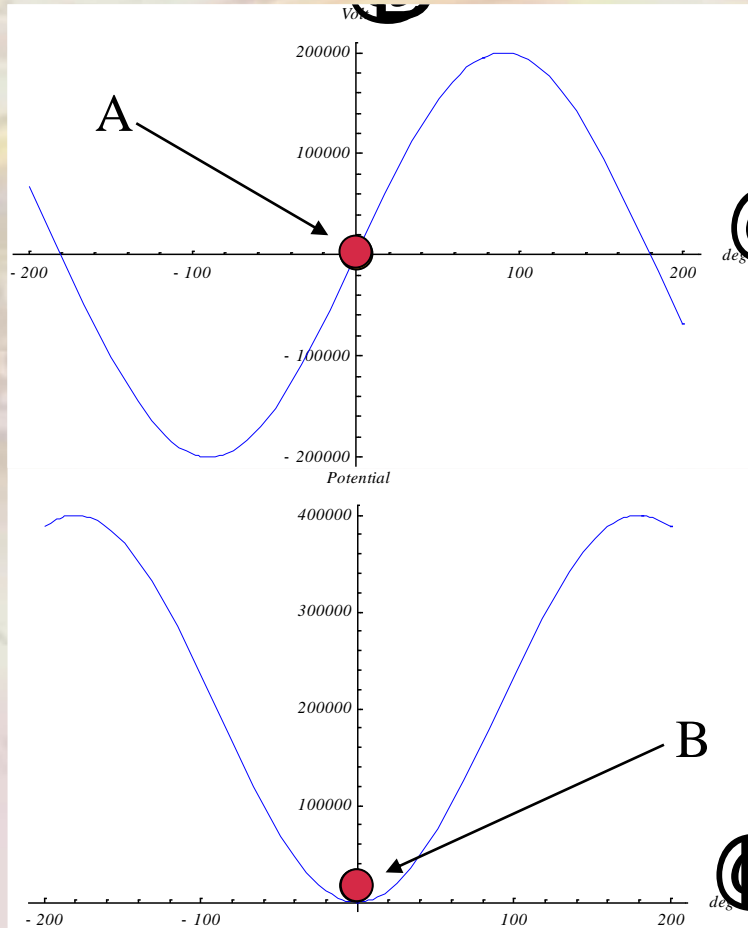
# The Potential Well (4)



Cavity voltage

Potential well

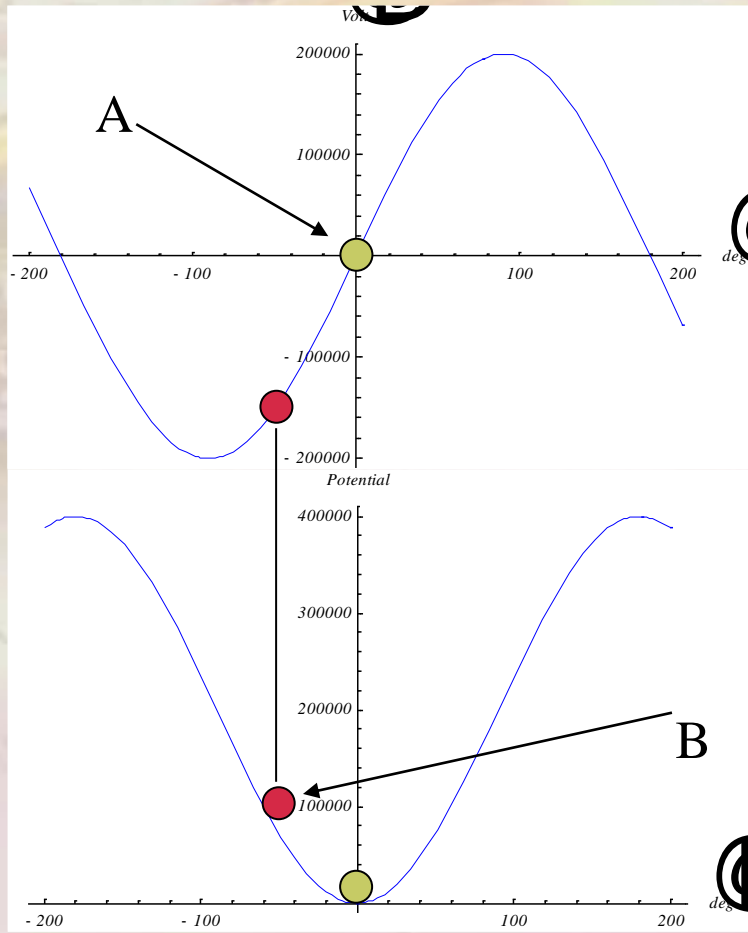
# The Potential Well (5)



Cavity voltage

Potential well

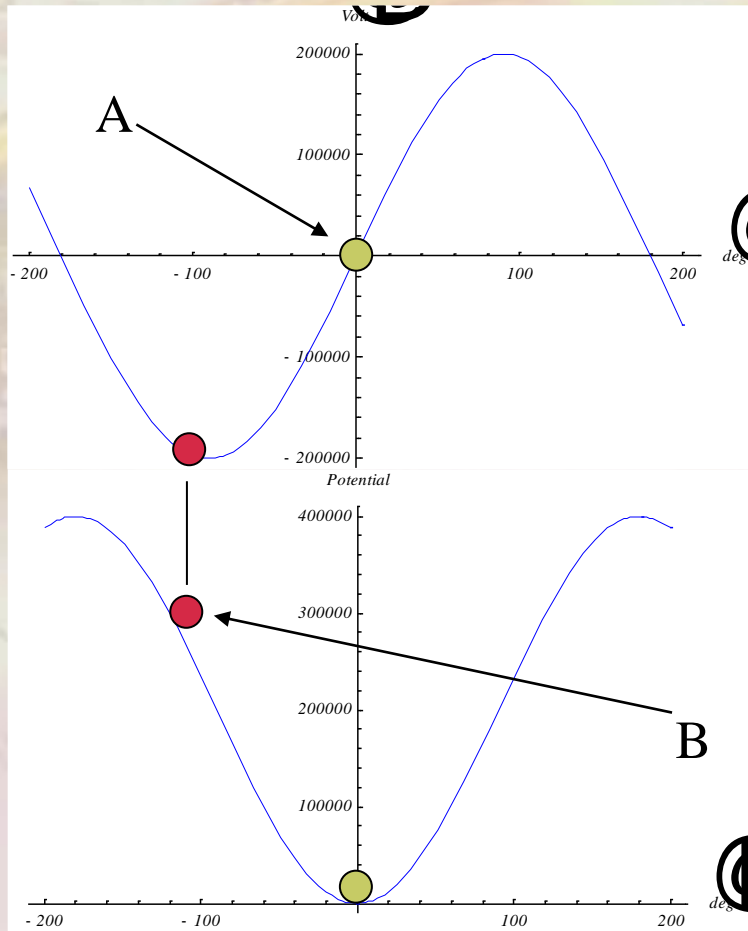
# The Potential Well (6)



Cavity voltage

Potential well

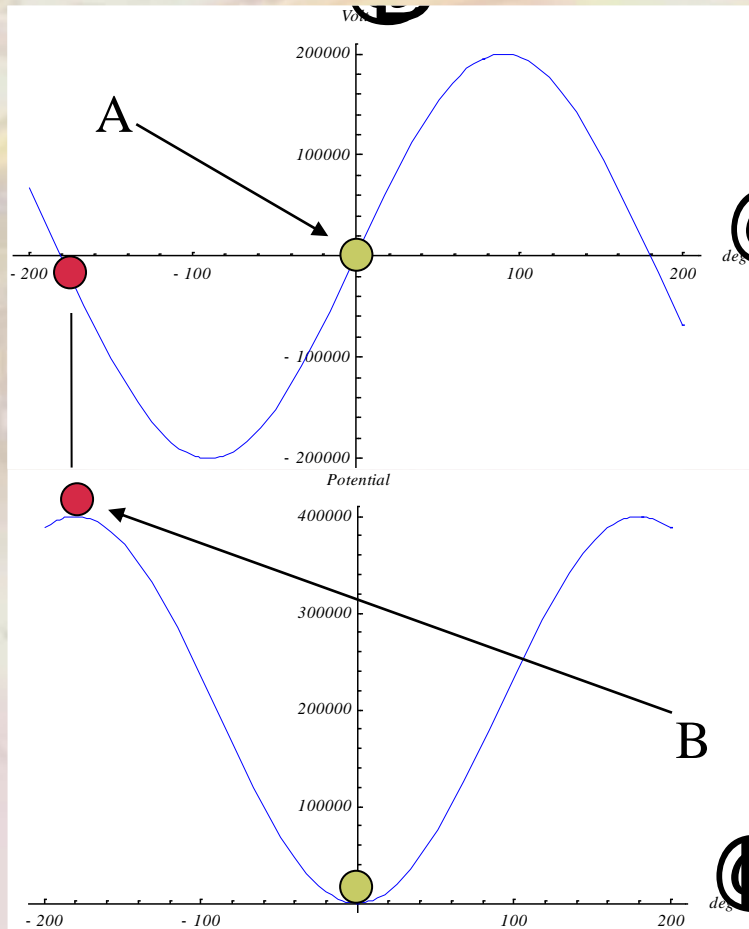
# The Potential Well (7)



Cavity voltage

Potential well

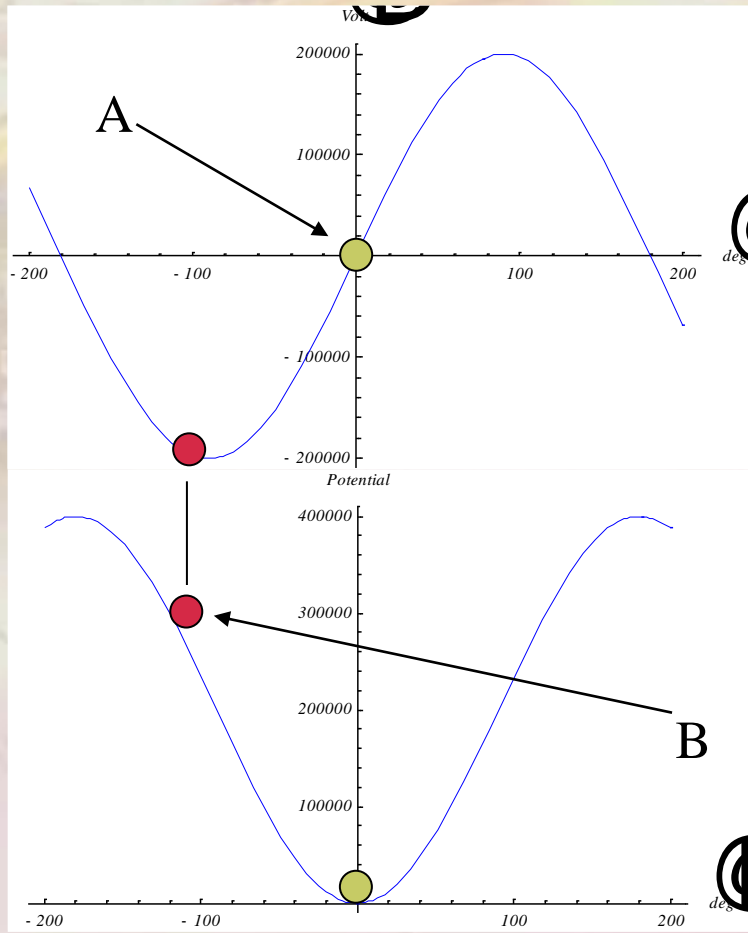
# The Potential Well (8)



Cavity voltage

Potential well

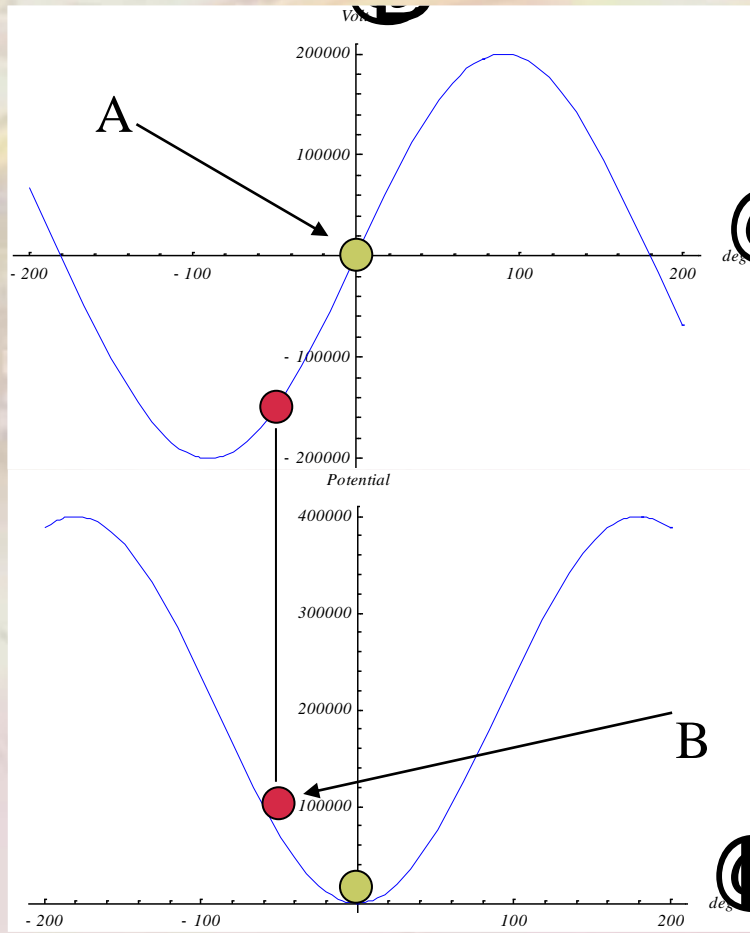
# The Potential Well (9)



Cavity voltage

Potential well

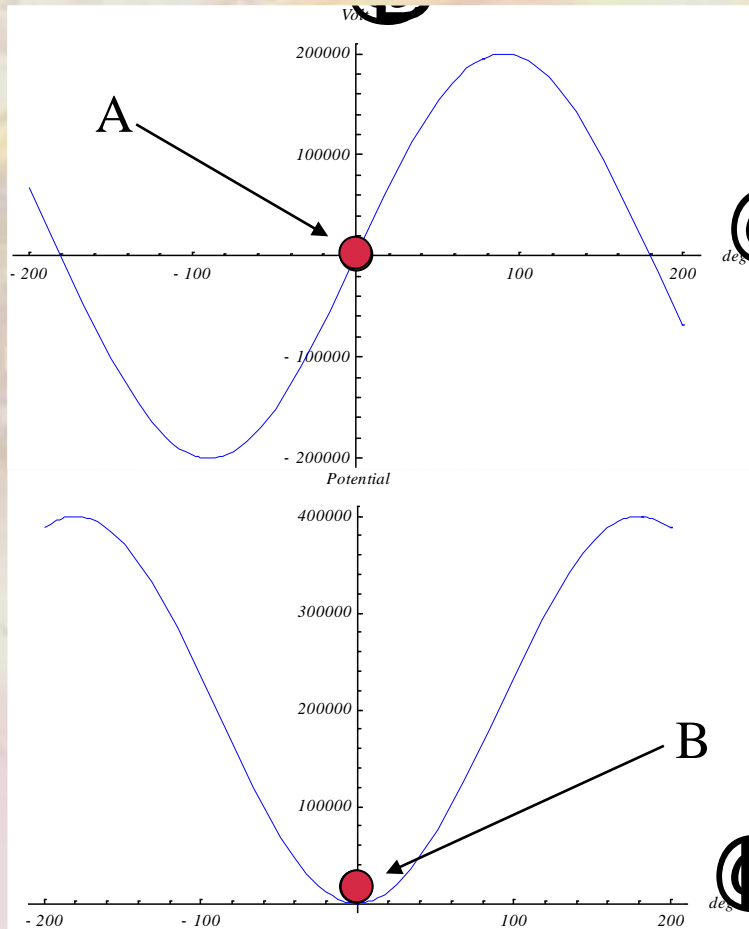
# The Potential Well (10)



Cavity voltage

Potential well

# The Potential Well (11)

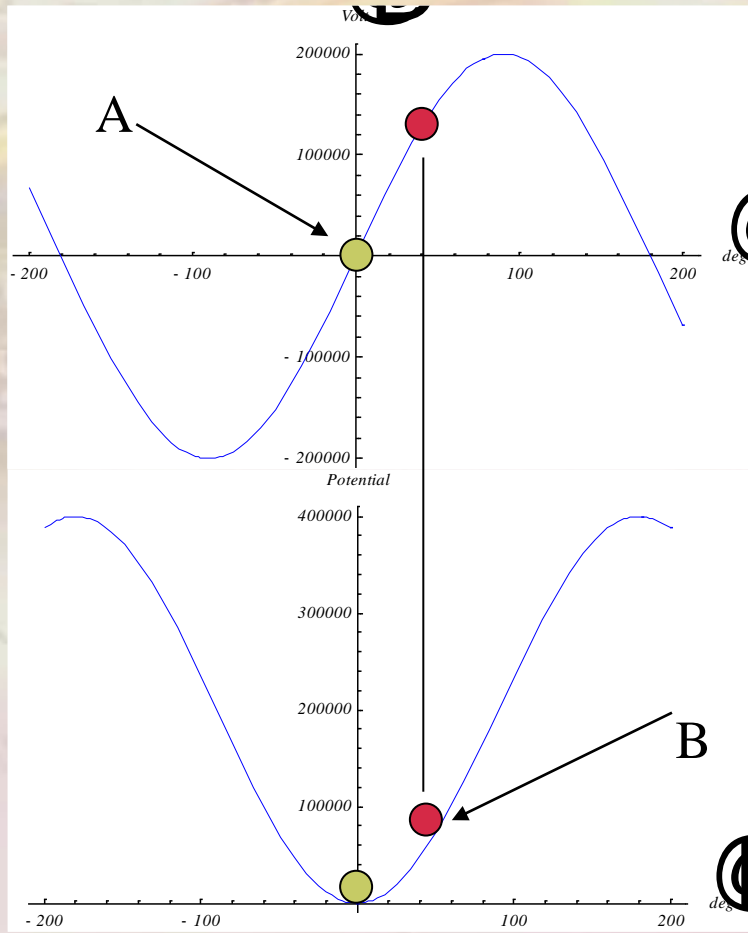


Cavity voltage

Potential well



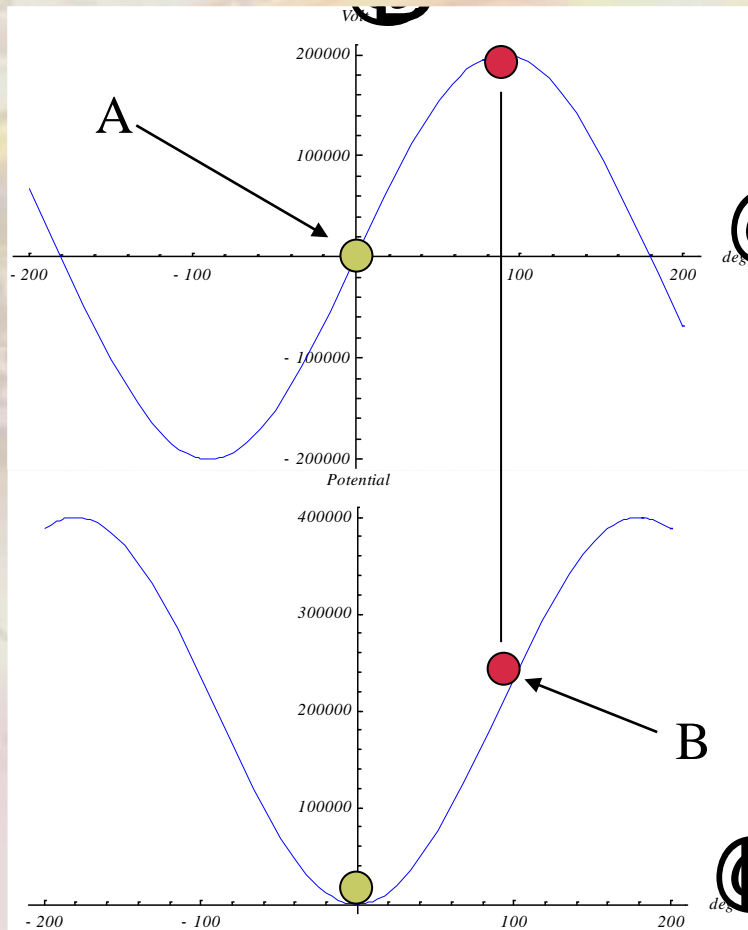
# The Potential Well (12)



Cavity voltage

Potential well

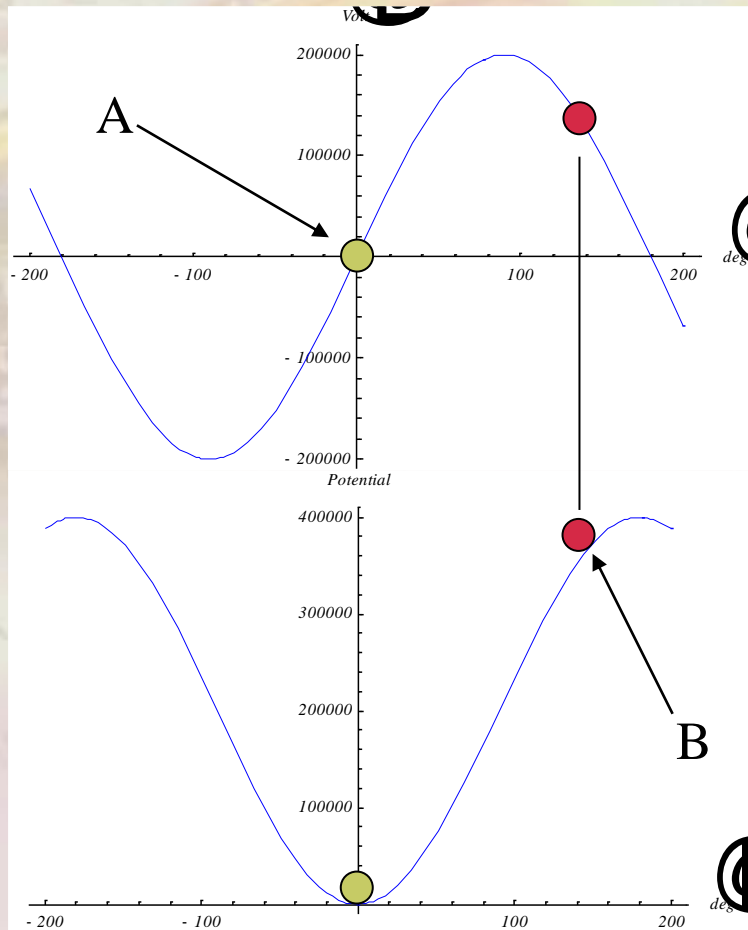
# The Potential Well (13)



Cavity voltage

Potential well

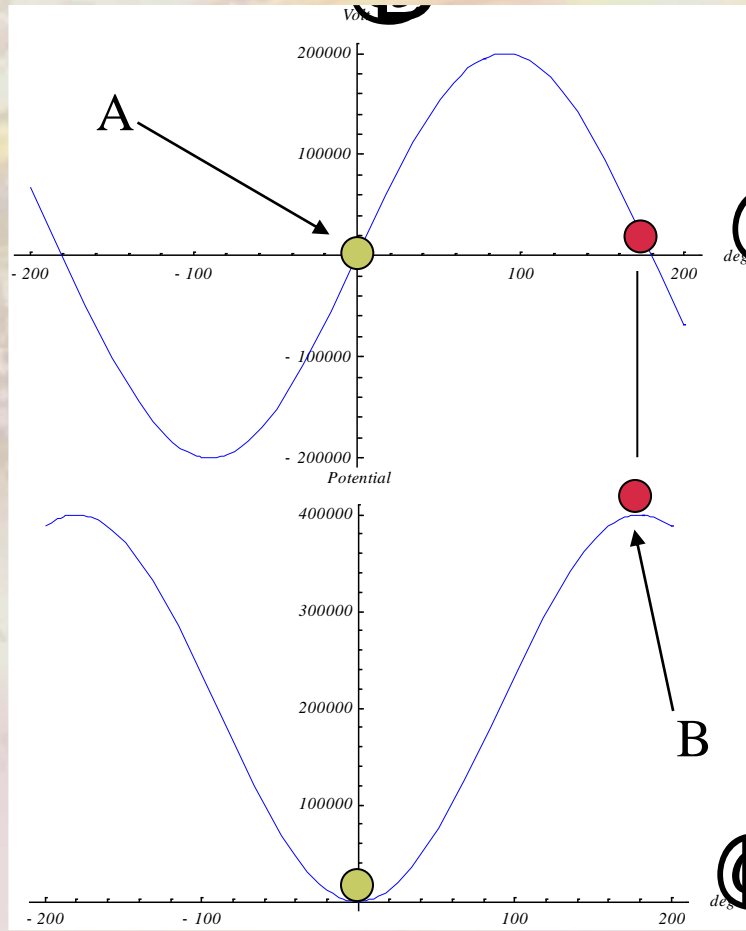
# The Potential Well (14)



Cavity voltage

Potential well

# The Potential Well (15)

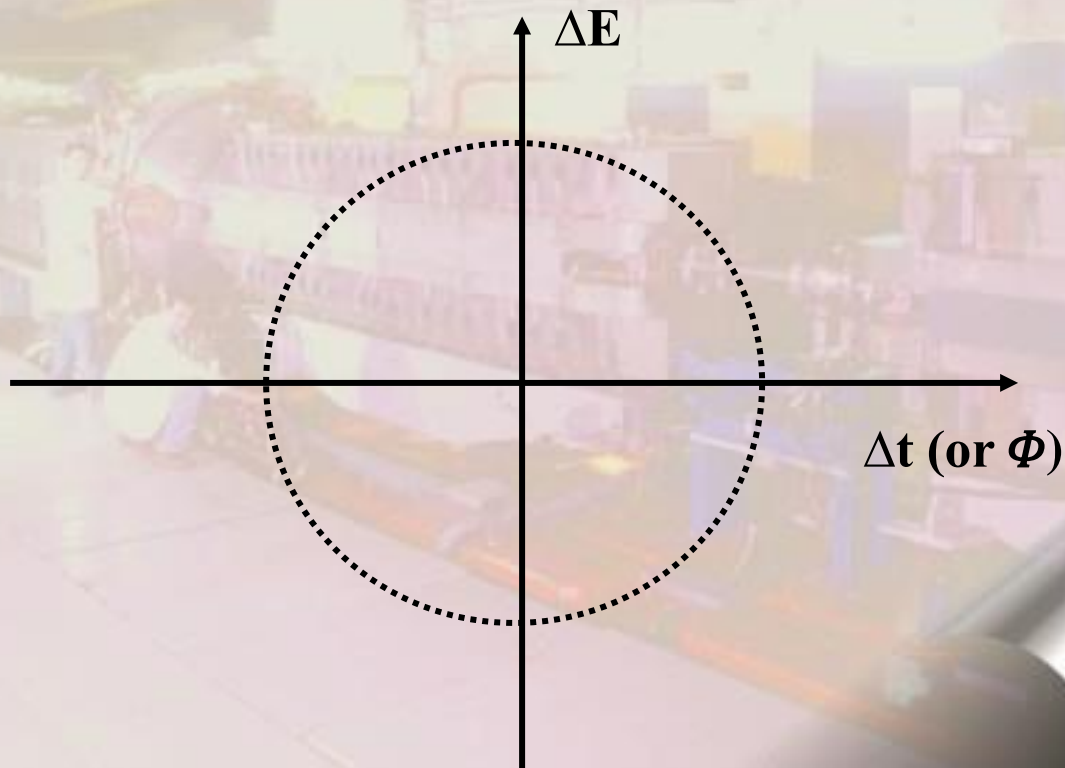


Cavity voltage

Potential well

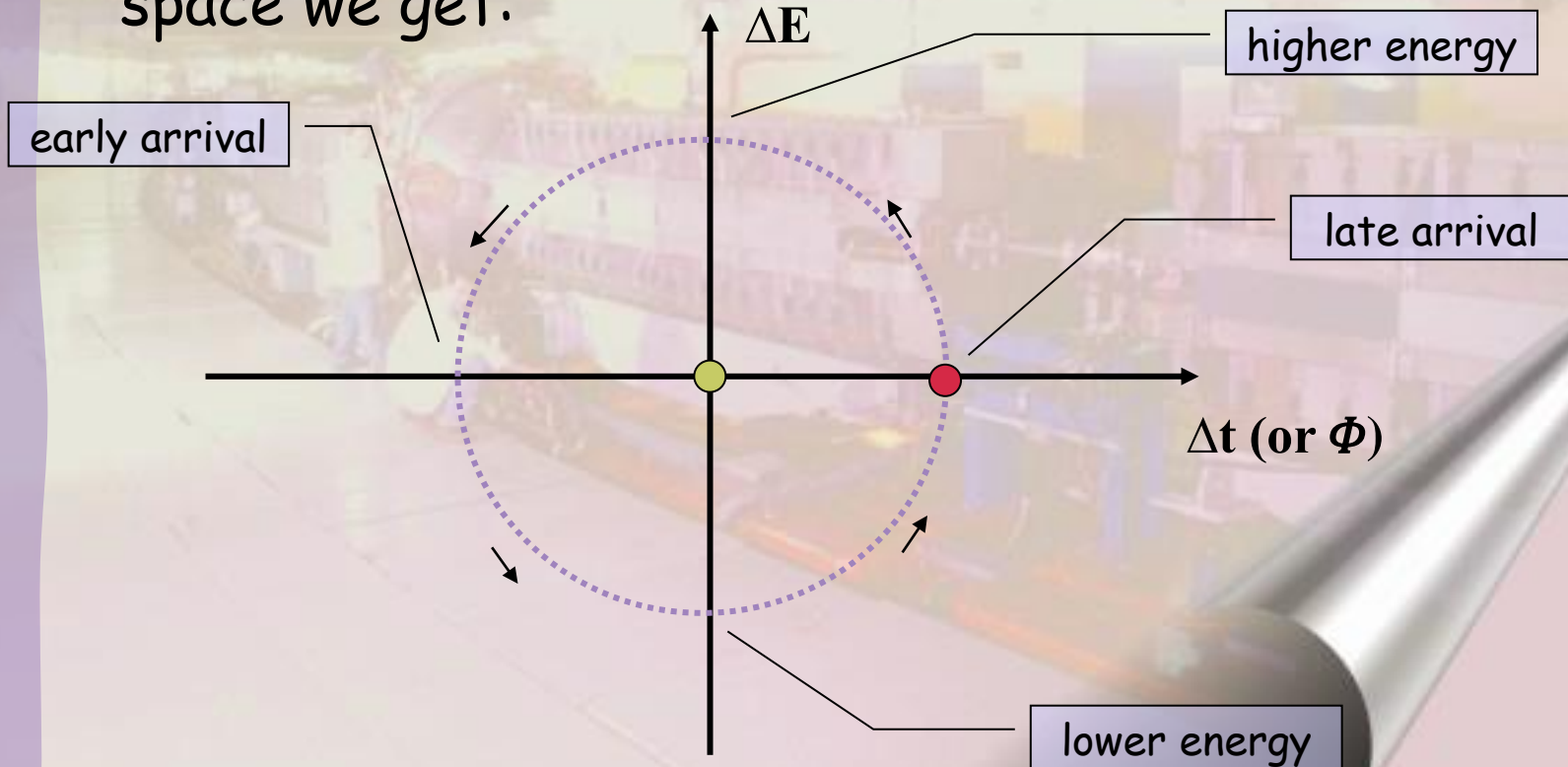
# Longitudinal Phase Space

- # In order to be able to visualize the motion in the longitudinal plane we define the longitudinal phase space (like we did for the transverse phase space)



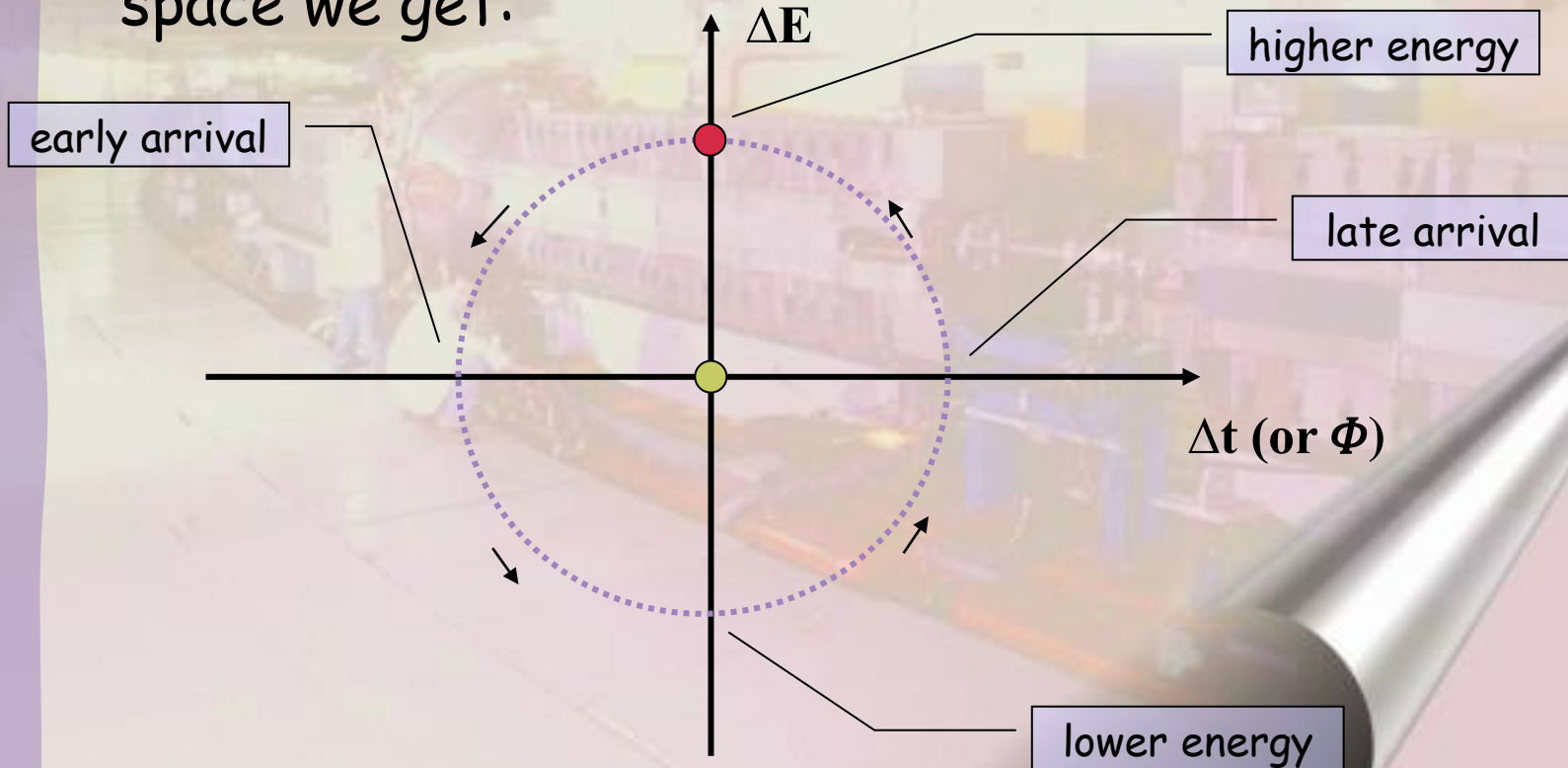
# Phase Space motion (1)

- # Particle B oscillates around particle A
  - This is synchrotron oscillation
- # When we plot this motion in our longitudinal phase space we get:



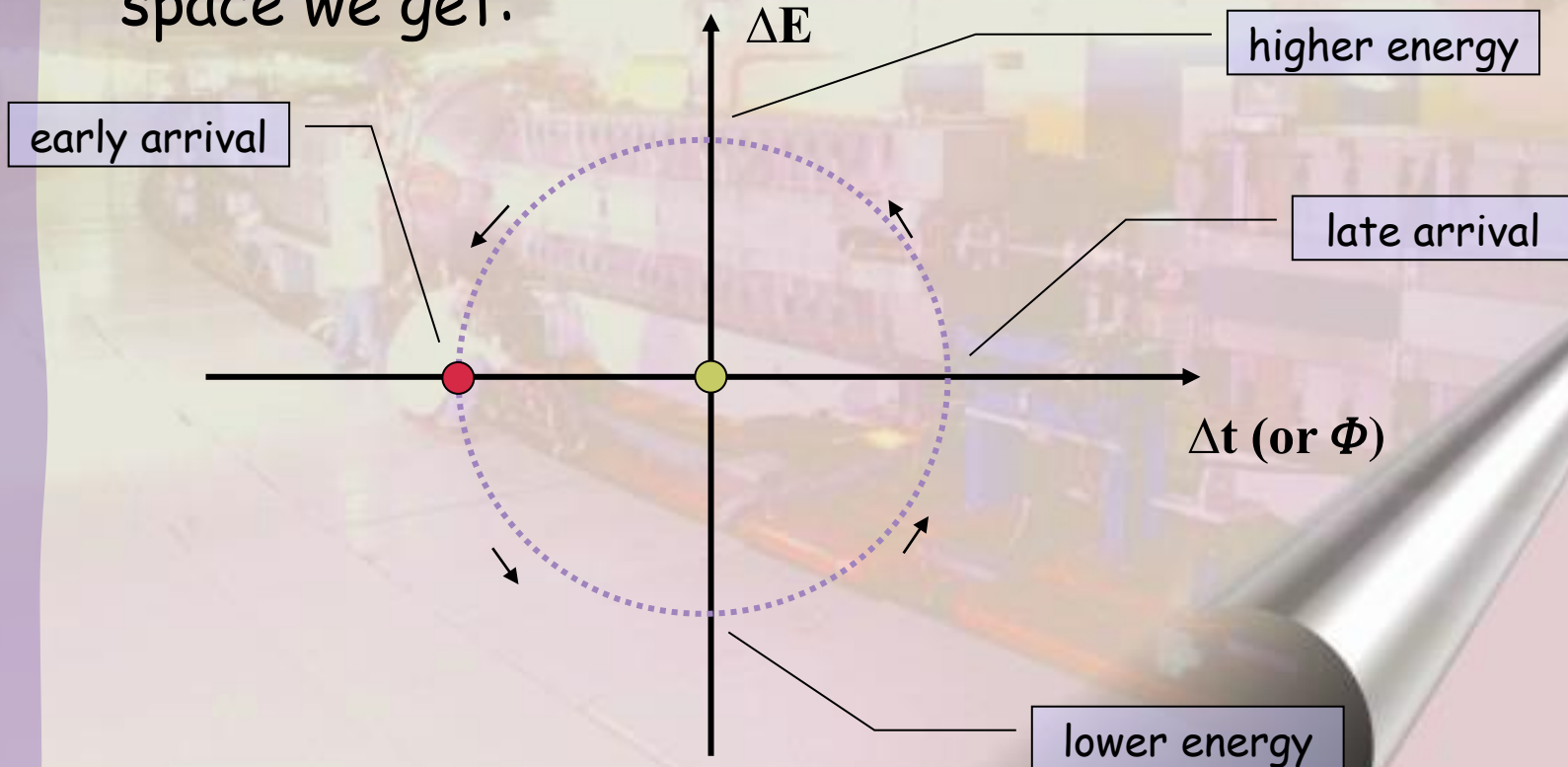
# Phase Space motion (2)

- # Particle B oscillates around particle A
  - This is synchrotron oscillation
- # When we plot this motion in our longitudinal phase space we get:



# Phase Space motion (3)

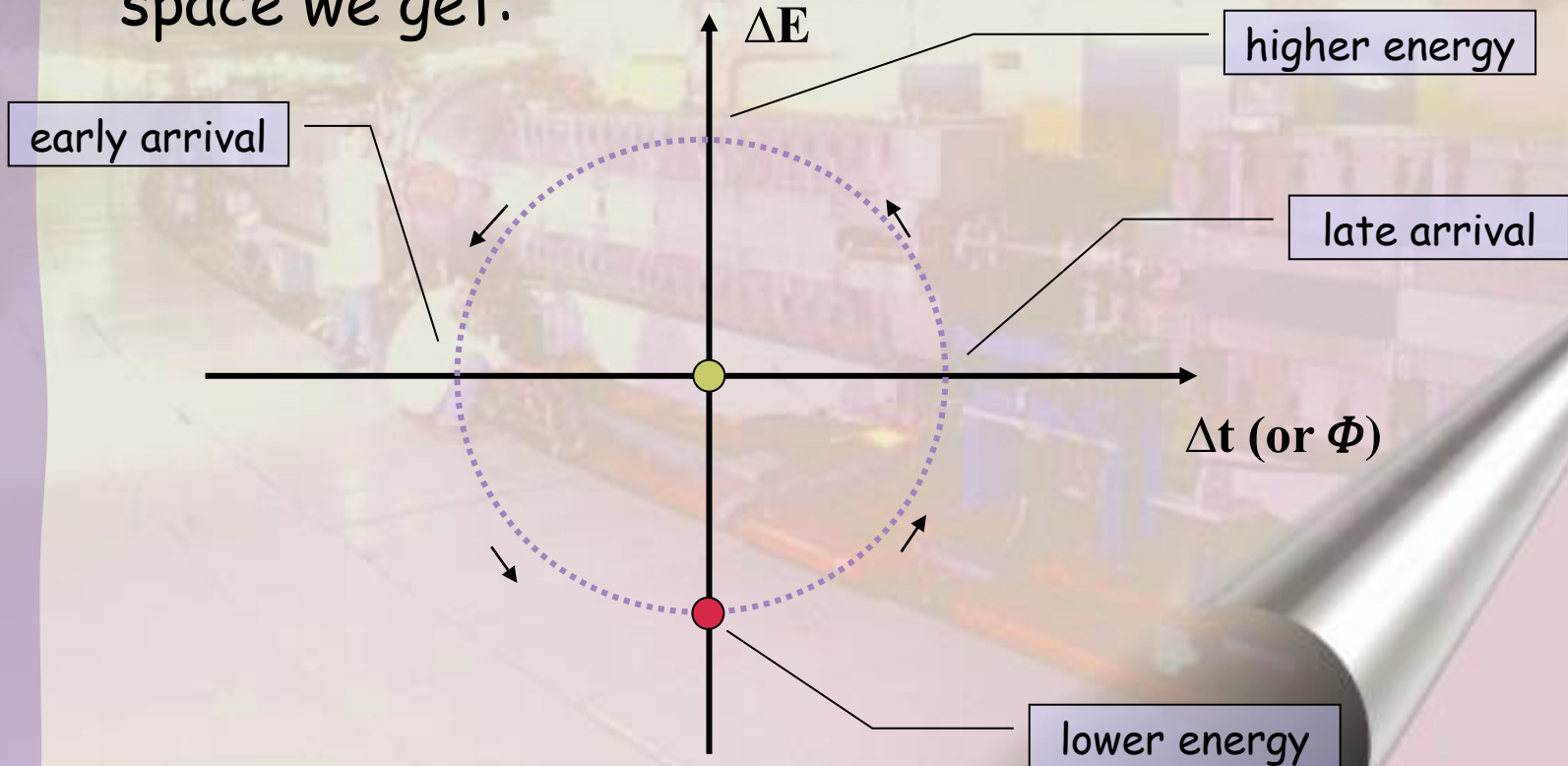
- # Particle B oscillates around particle A
  - This is synchrotron oscillation
- # When we plot this motion in our longitudinal phase space we get:





# Phase Space motion (4)

- # Particle B oscillates around particle A
  - This is synchrotron oscillation
- # When we plot this motion in our longitudinal phase space we get:



# Quick intermediate summary...

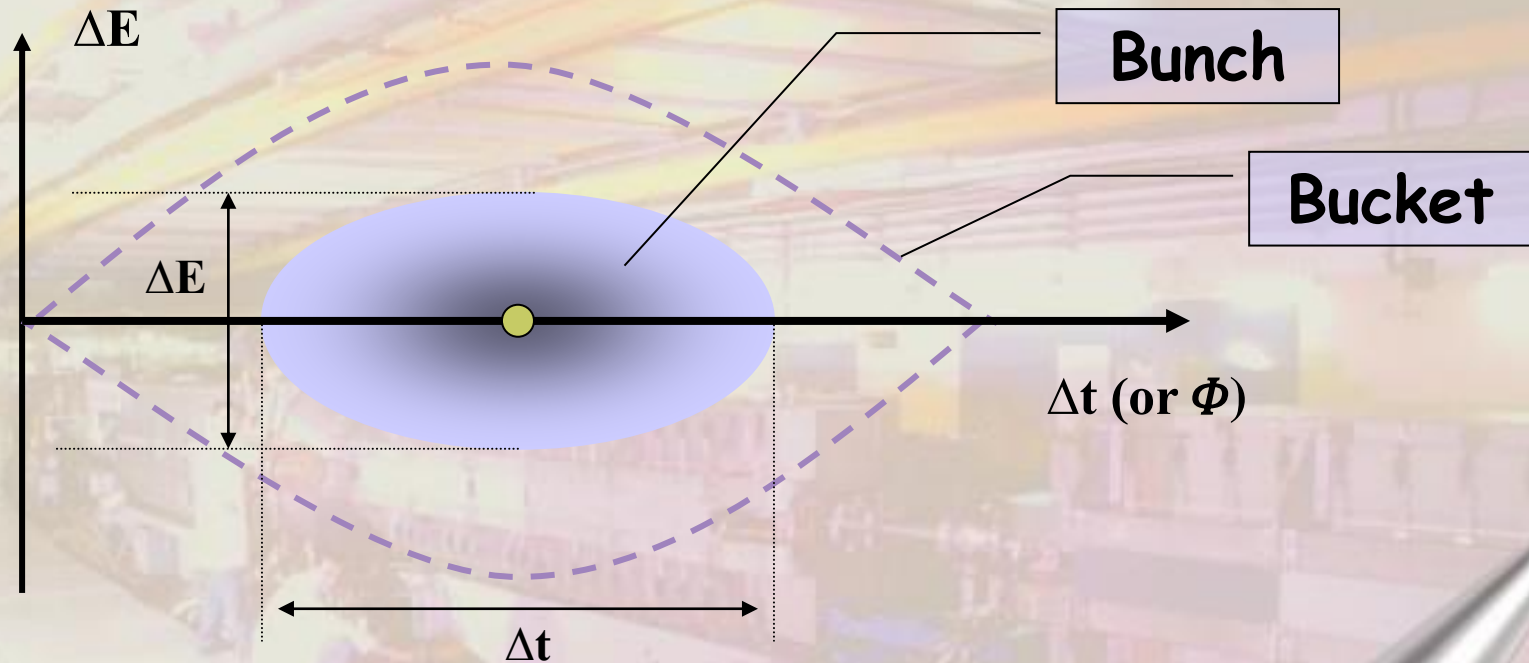
## # We have seen that:

- The RF system forms a potential well in which the particles oscillate (synchrotron oscillation).
- We can describe this motion in the longitudinal phase space (energy versus time or phase).
- This works for particles below transition.

## # However,

- Due to the shape of the potential well, the oscillation is a non-linear motion.
- The phase space trajectories are therefore no circles nor ellipses.
- What when our particles are **above transition** ?

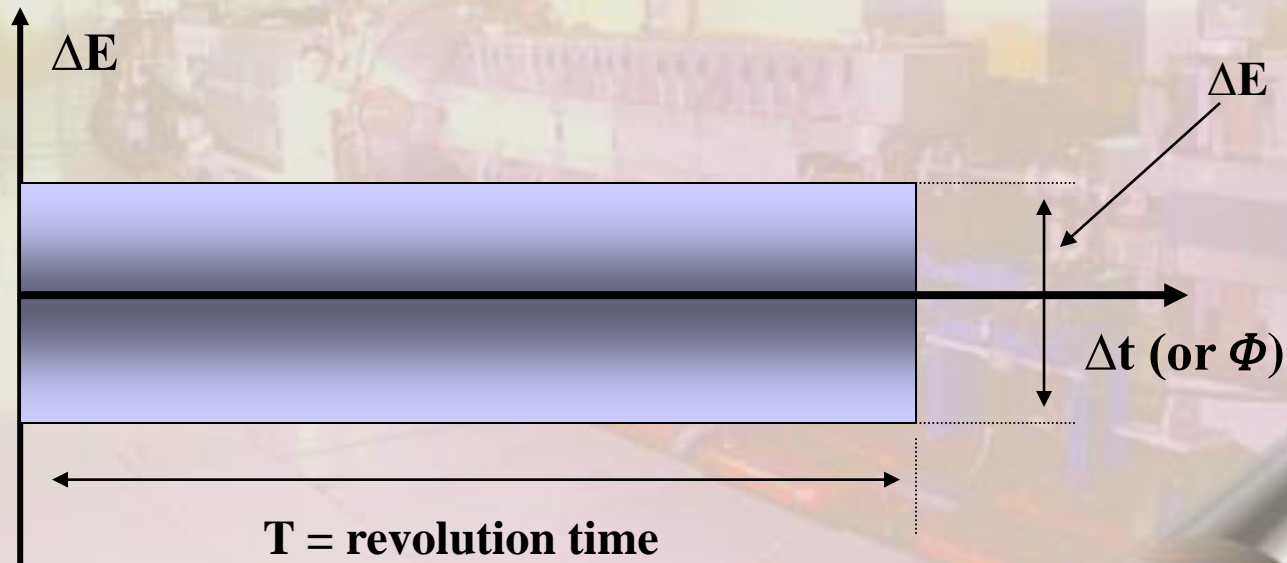
# Stationary bunch & bucket



- # Bucket area = longitudinal Acceptance [eVs]
- # Bunch area = longitudinal beam emittance =  $\pi \cdot \Delta E \cdot \Delta t / 4$  [eVs]

# Unbunched (coasting) beam

- # The emittance of an unbunched beam is just  $\Delta E T$  eVs
  - $\Delta E$  is the energy spread [eV]
  - $T$  is the revolution time [s]



# What happens beyond transition ?

- # Until now we have seen how things look like below transition

$\eta = \text{positive}$

Higher energy  $\Rightarrow$  faster orbit  $\Rightarrow$  higher  $F_{\text{rev}}$   $\Rightarrow$  next time particle will be **earlier**.

Lower energy  $\Rightarrow$  slower orbit  $\Rightarrow$  lower  $F_{\text{rev}}$   $\Rightarrow$  next time particle will be **later**.

- # What will happen above transition ?

$\eta = \text{negative}$

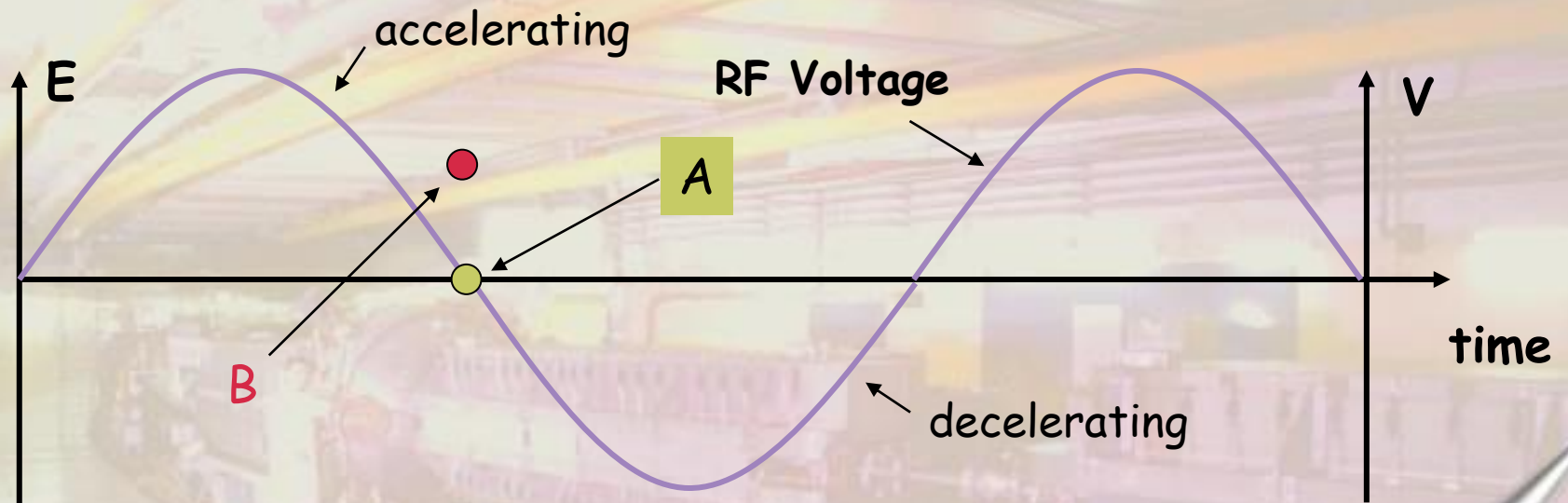
Higher energy  $\Rightarrow$  longer orbit  $\Rightarrow$  lower  $F_{\text{rev}}$   $\Rightarrow$  next time particle will be **later**.

Lower energy  $\Rightarrow$  shorter orbit  $\Rightarrow$  higher  $F_{\text{rev}}$   $\Rightarrow$  next time particle will be **earlier**.

# What are the implication for the RF ?

- # For particles below transition we worked on the rising edge of the sine wave.
- # For Particles above transition we will work on the falling edge of the sine wave.
- # We will see why.....

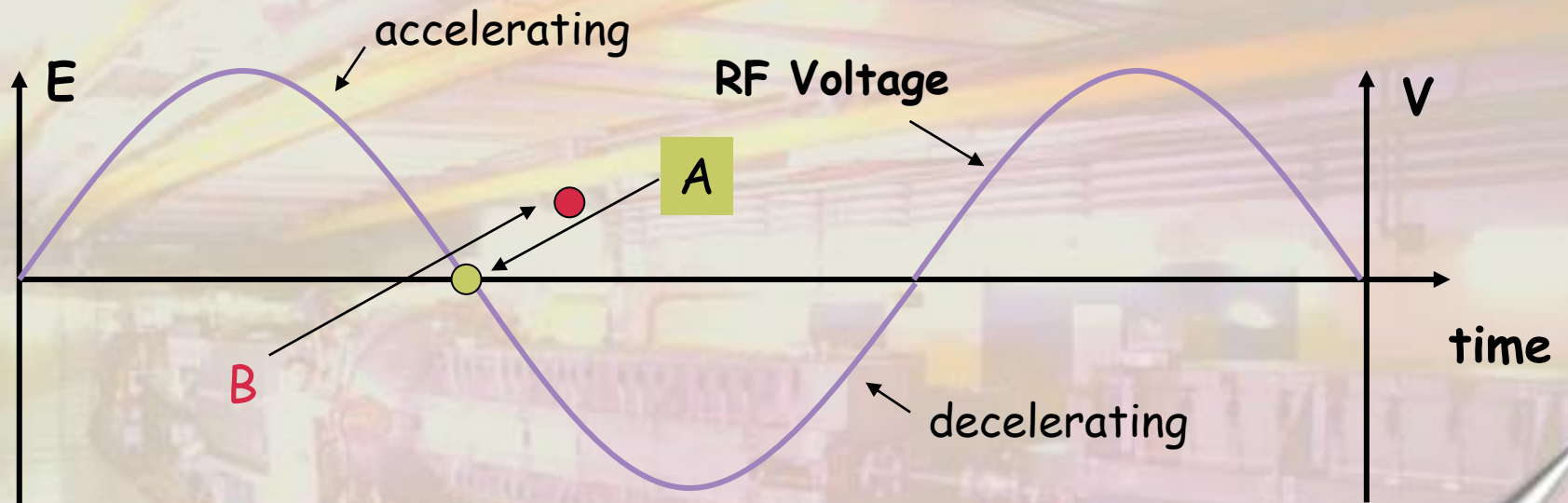
# Longitudinal motion beyond transition (1)



# Imagine two particles A and B, that arrive at the same time in the accelerating cavity (when  $V_{rf} = 0V$ )

- For A the energy is such that  $F_{rev A} = F_{rf}$ .
- The energy of B is higher  $\rightarrow F_{rev B} < F_{rev A}$

# Longitudinal motion beyond transition (2)

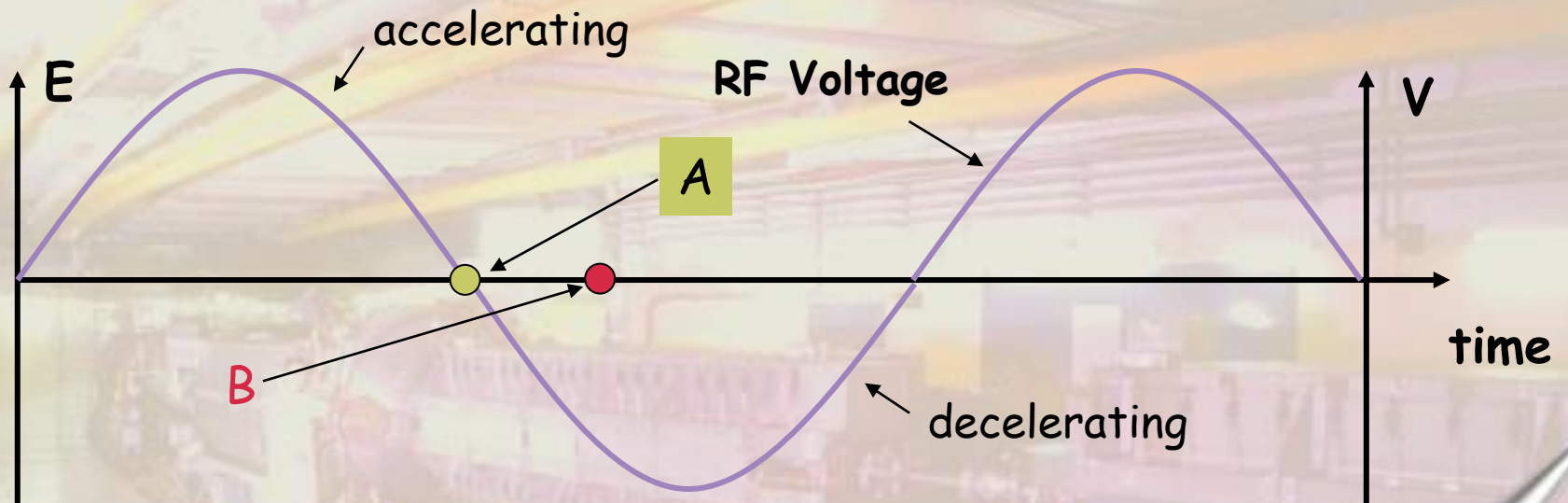


# Particle B arrives after A and experiences a decelerating voltage.

■ The energy of B is still higher, but less  $\rightarrow F_{\text{rev B}} < F_{\text{rev A}}$



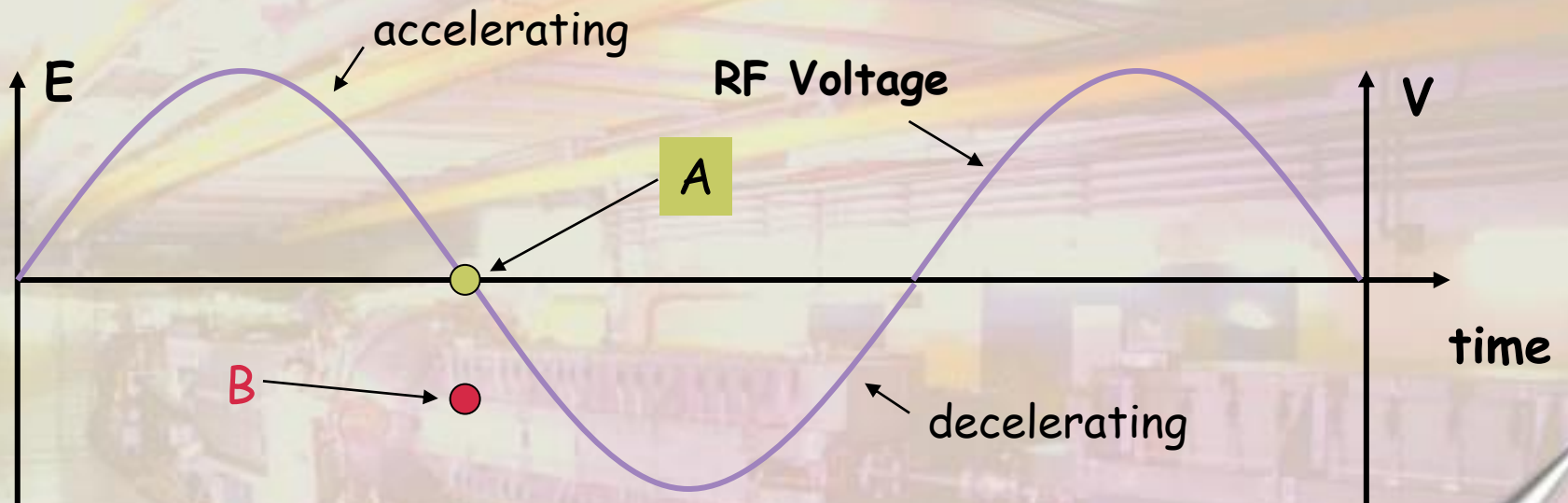
# Longitudinal motion beyond transition (3)



# B has now the same energy as A, but arrives still later and experiences therefore a decelerating voltage.

■  $F_{\text{rev B}} = F_{\text{rev A}}$

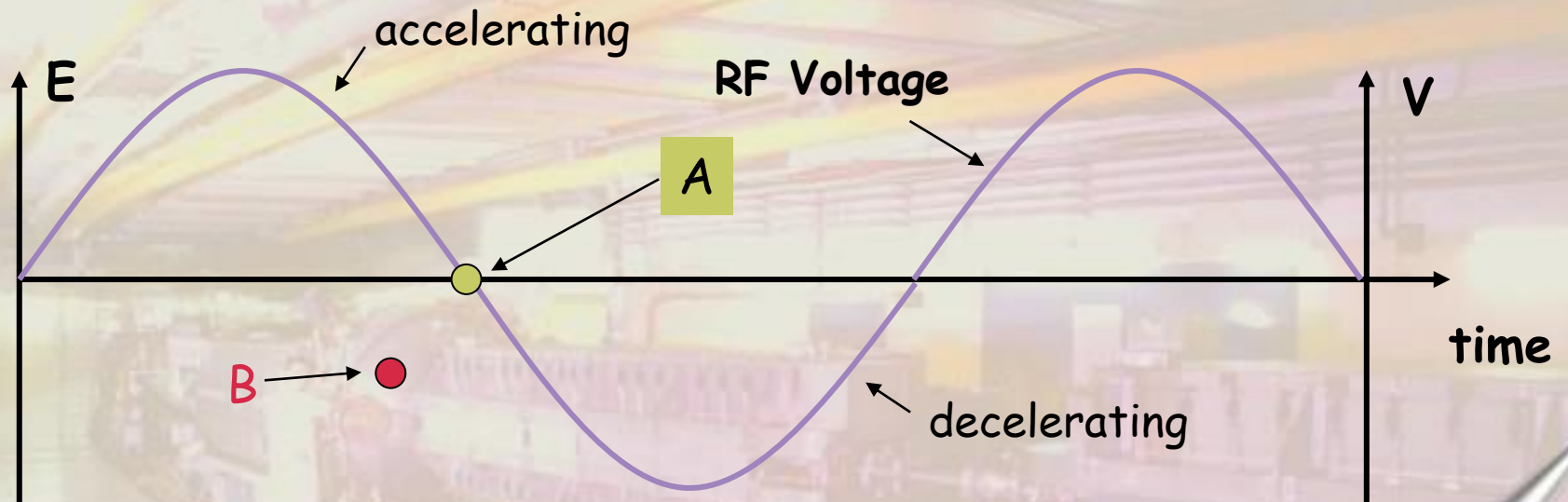
# Longitudinal motion beyond transition (4)



# Particle B has now a lower energy as A, but arrives at the same time

$$\blacksquare F_{\text{rev B}} > F_{\text{rev A}}$$

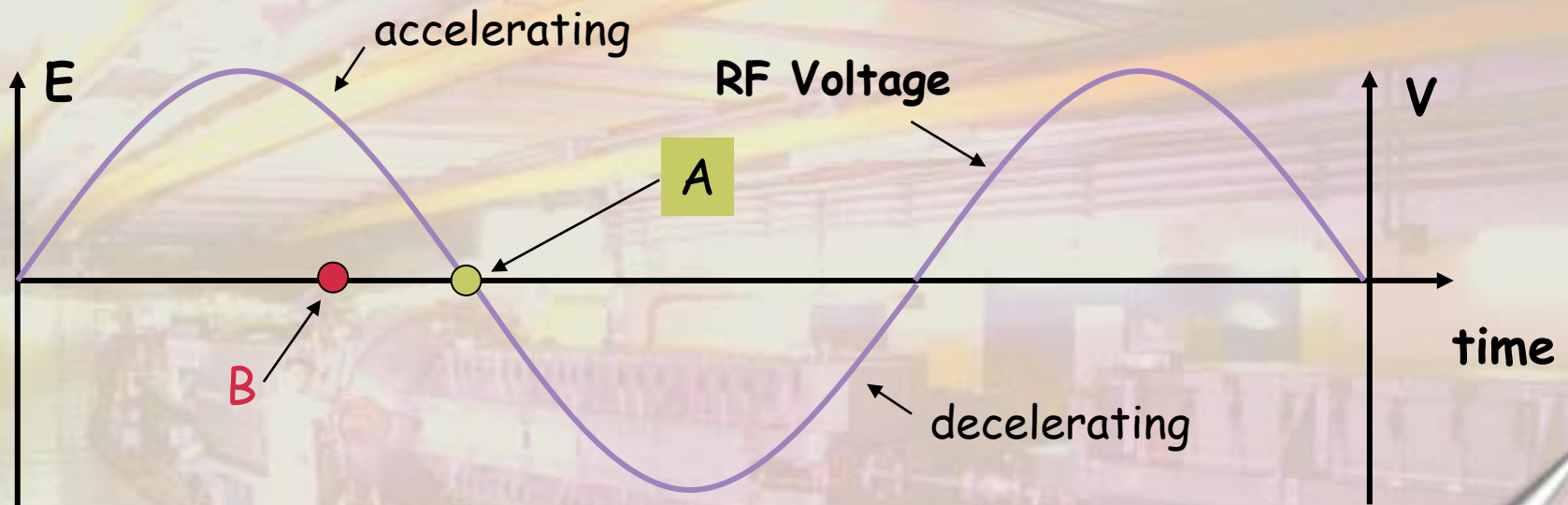
# Longitudinal motion beyond transition (5)



# Particle B has now a lower energy as A, but B arrives before A and experiences an accelerating voltageage.

$$\blacksquare F_{\text{rev B}} > F_{\text{rev A}}$$

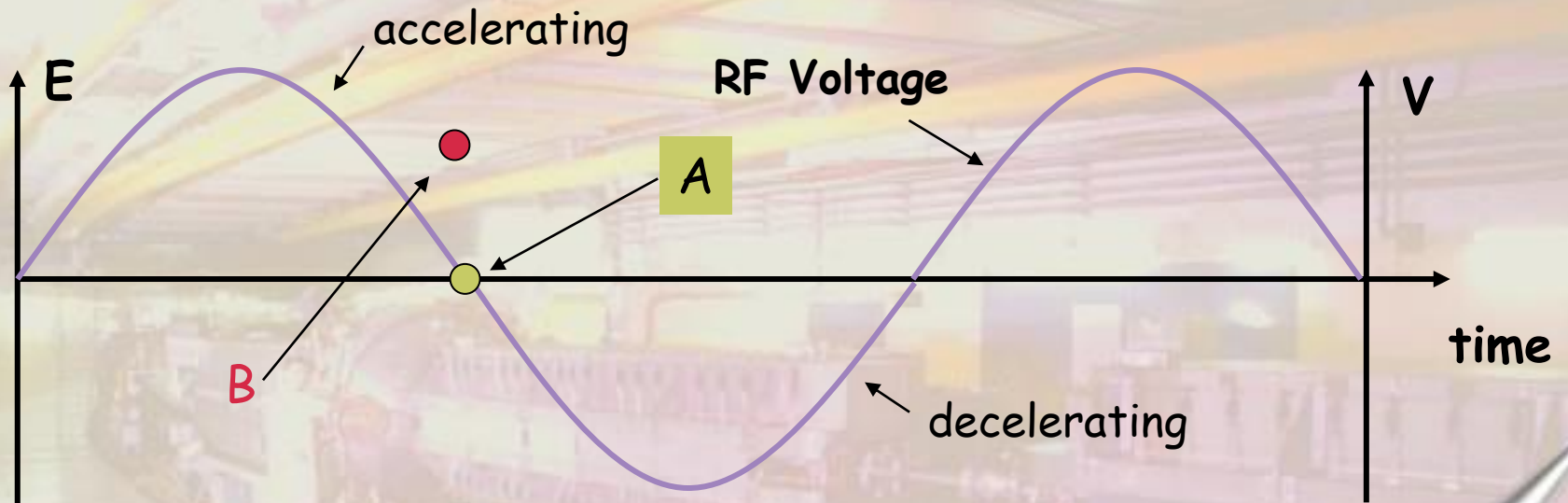
# Longitudinal motion beyond transition (6)



# Particle B has now the same energy as A, but B still arrives before A and experiences an accelerating voltage.

$$\blacksquare F_{\text{rev B}} > F_{\text{rev A}}$$

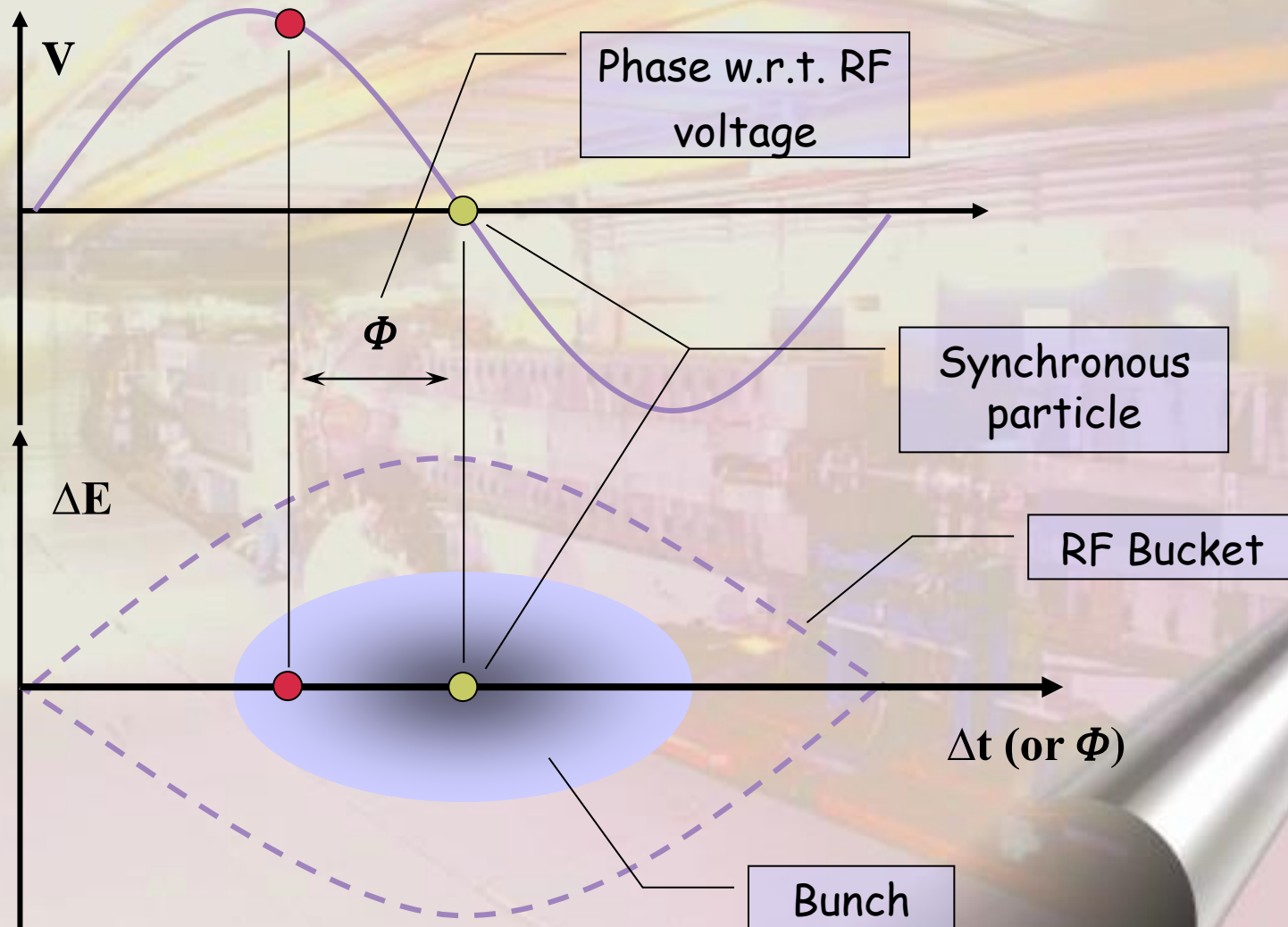
# Longitudinal motion beyond transition (7)



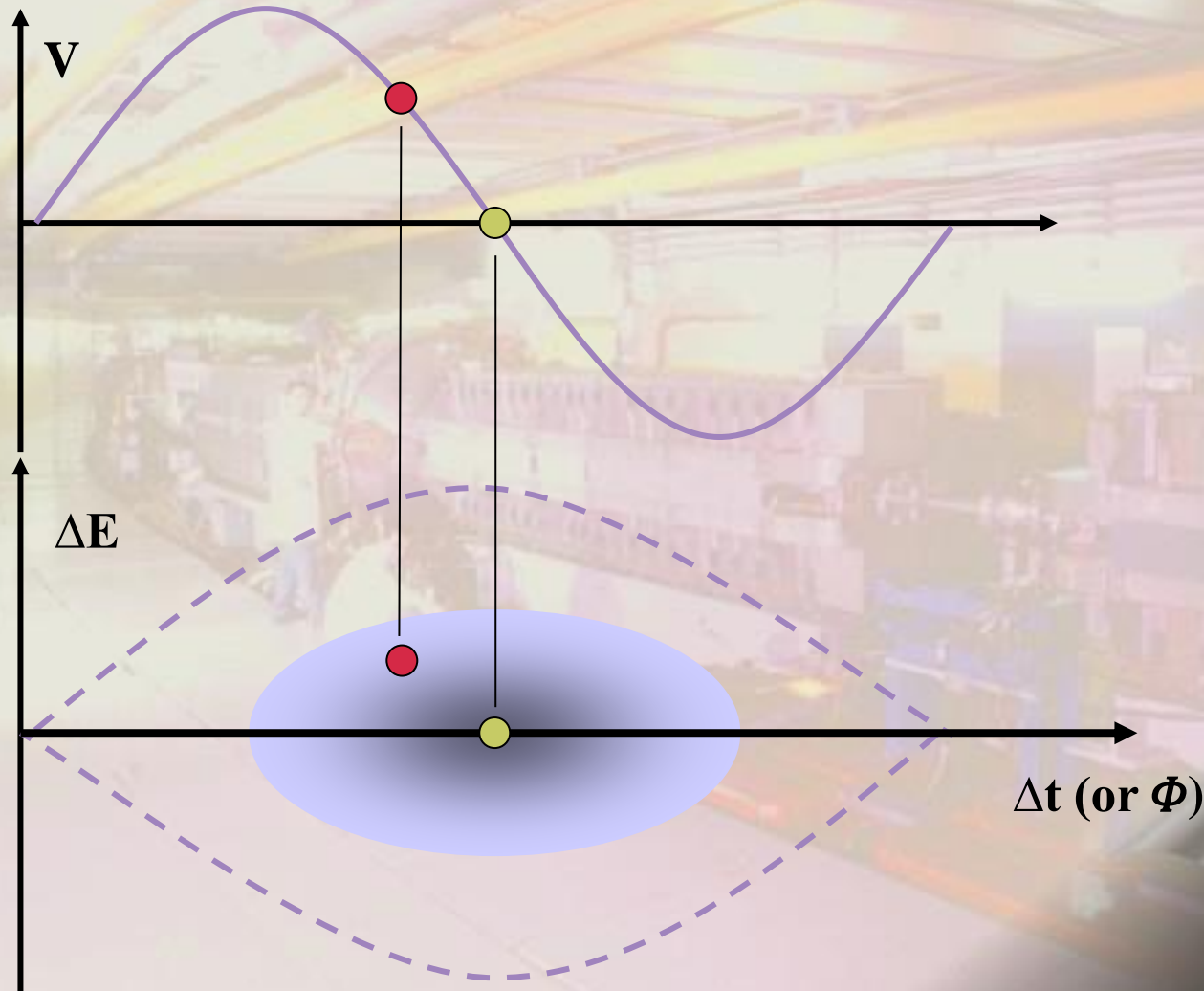
# Particle B has now a higher energy as A and arrives at the same time again....

$$\square F_{\text{rev } B} < F_{\text{rev } A}$$

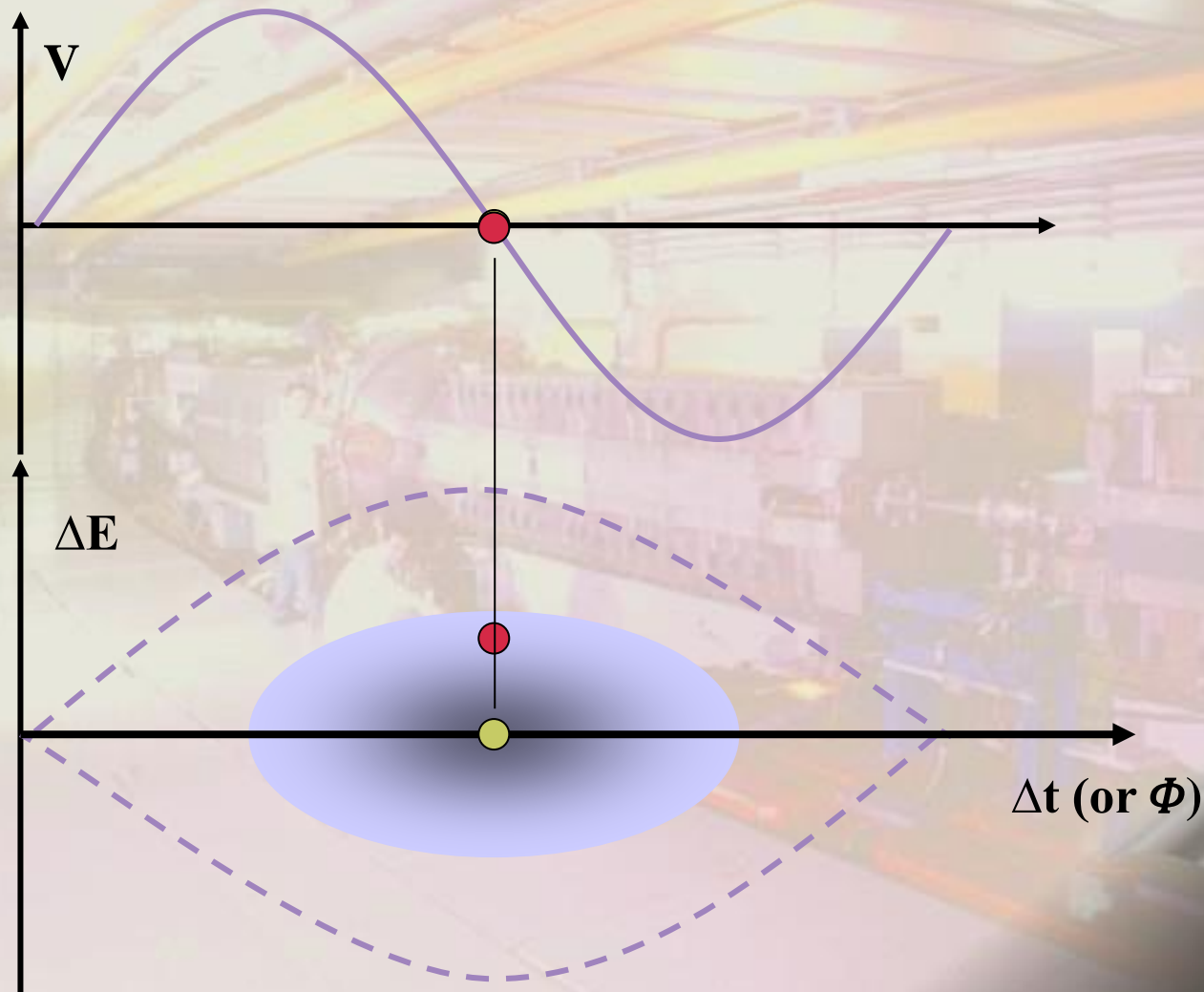
# The motion in the bucket (1)



# The motion in the bucket (2)

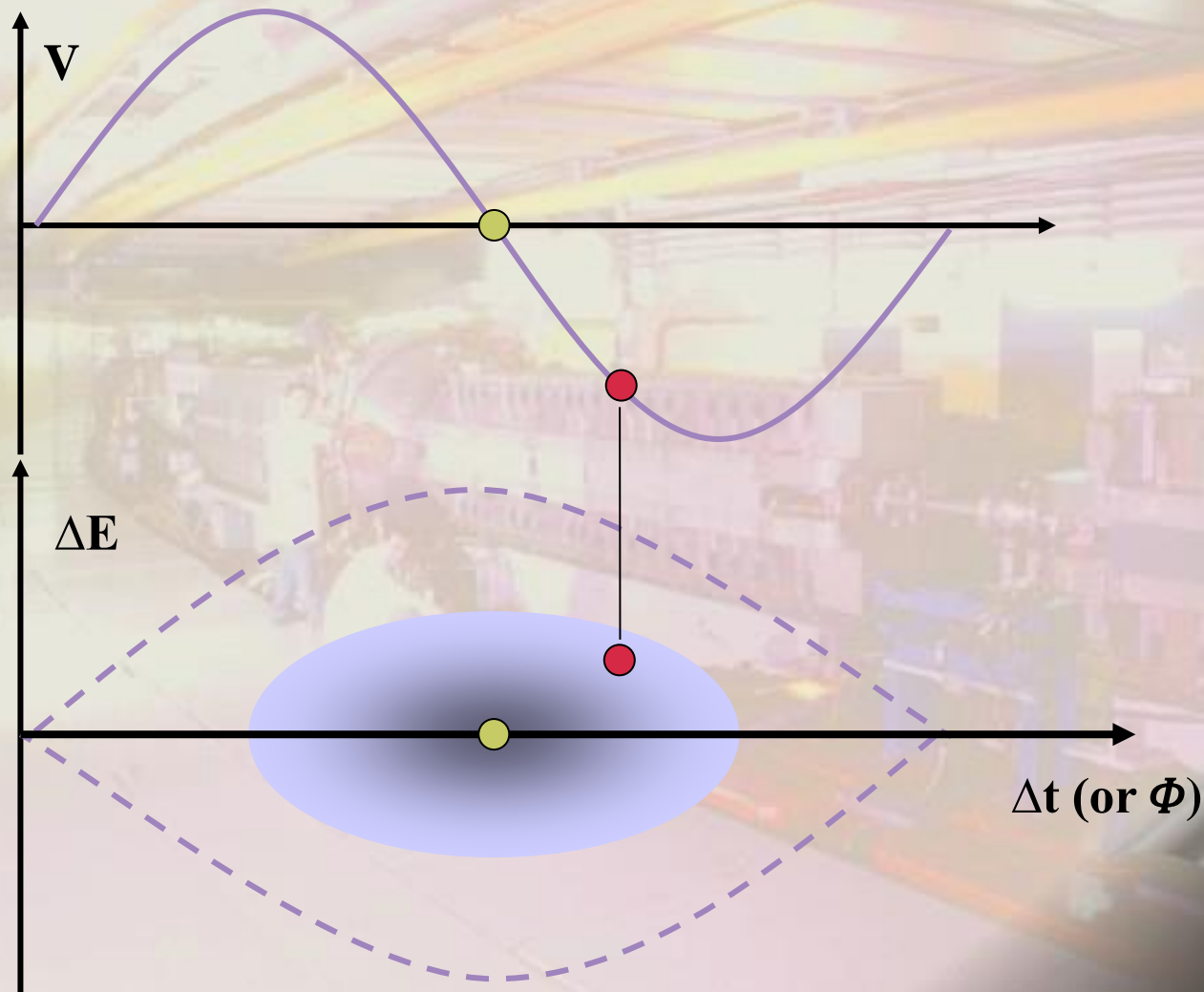


# The motion in the bucket (3)

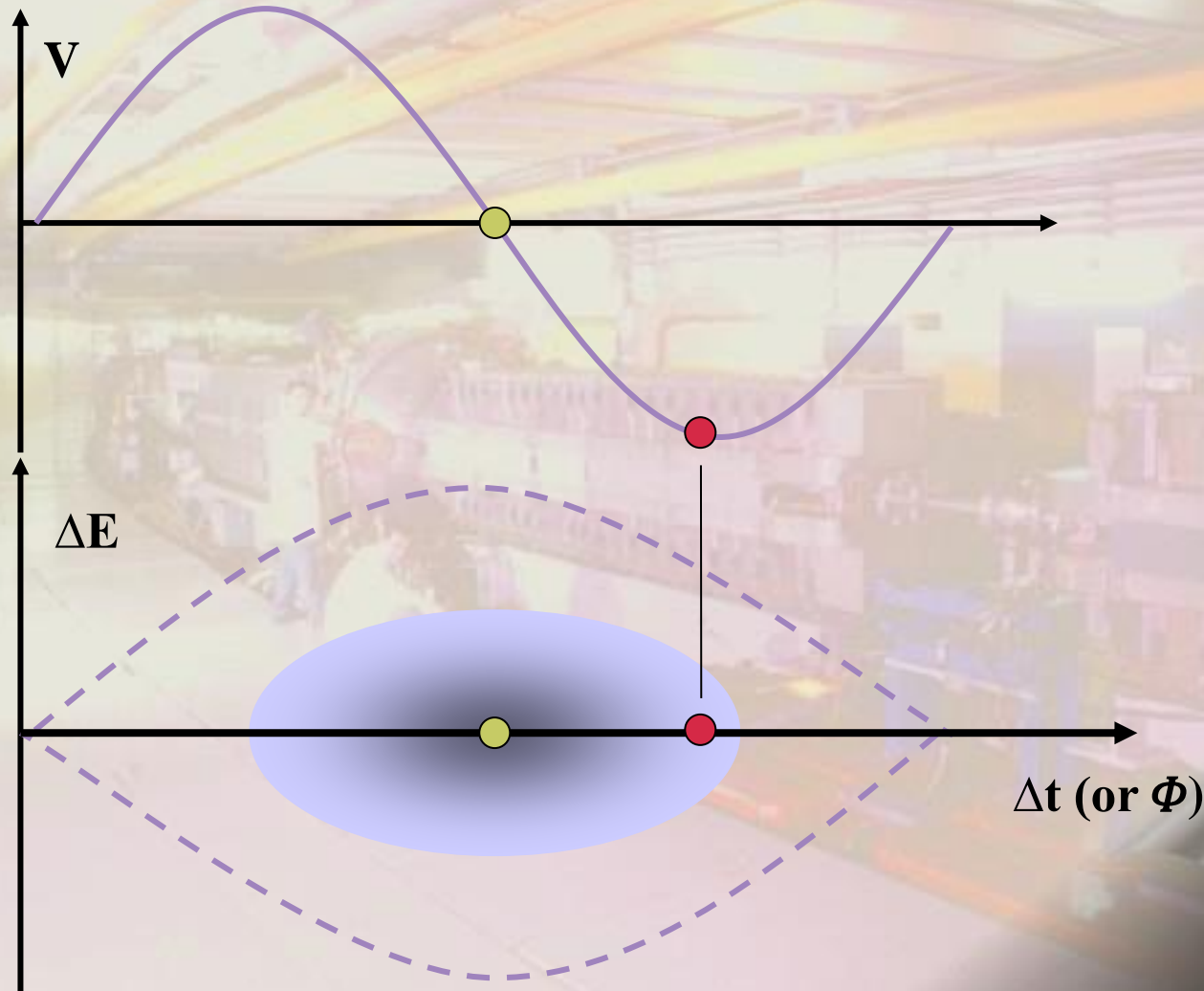




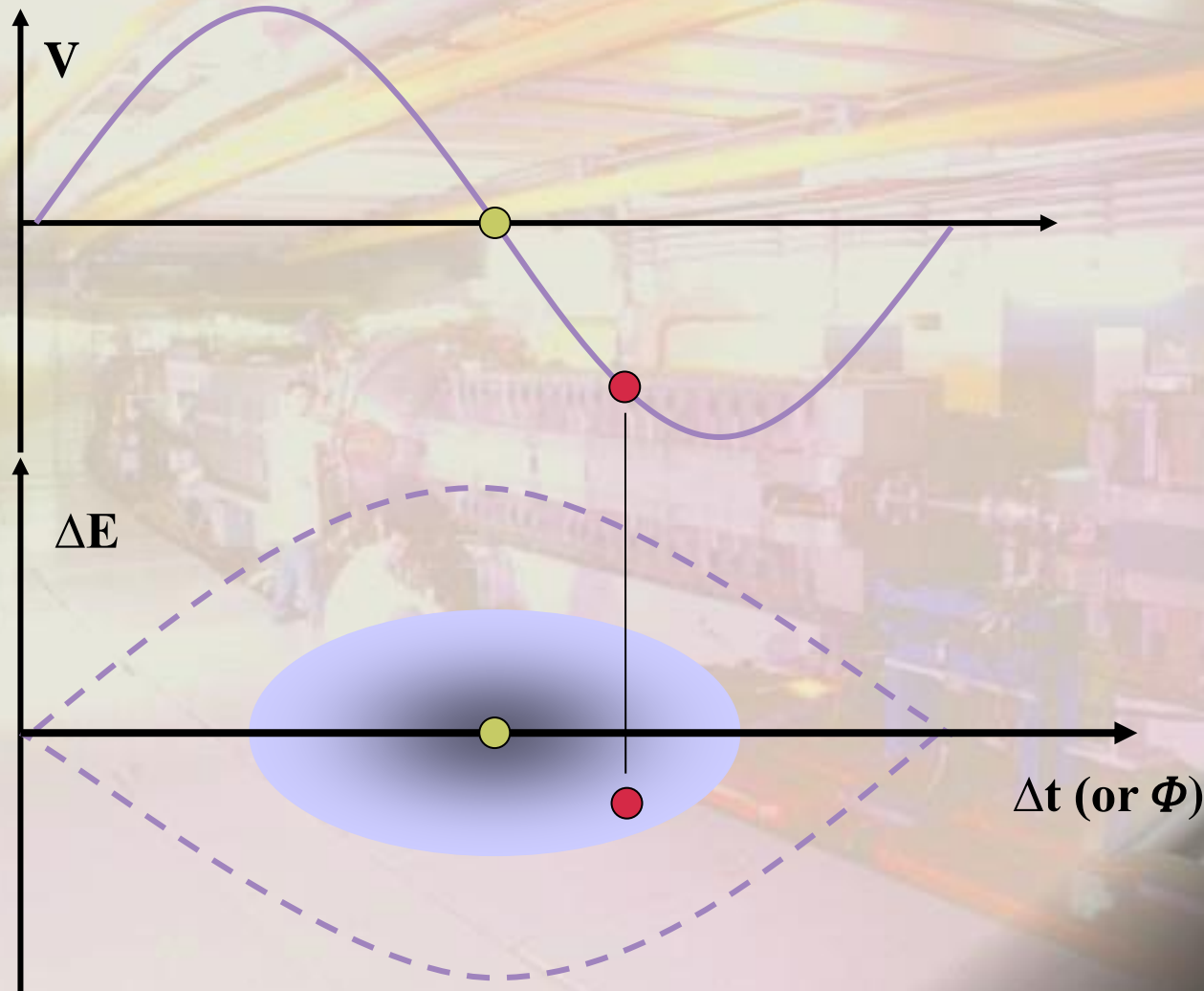
# The motion in the bucket (4)



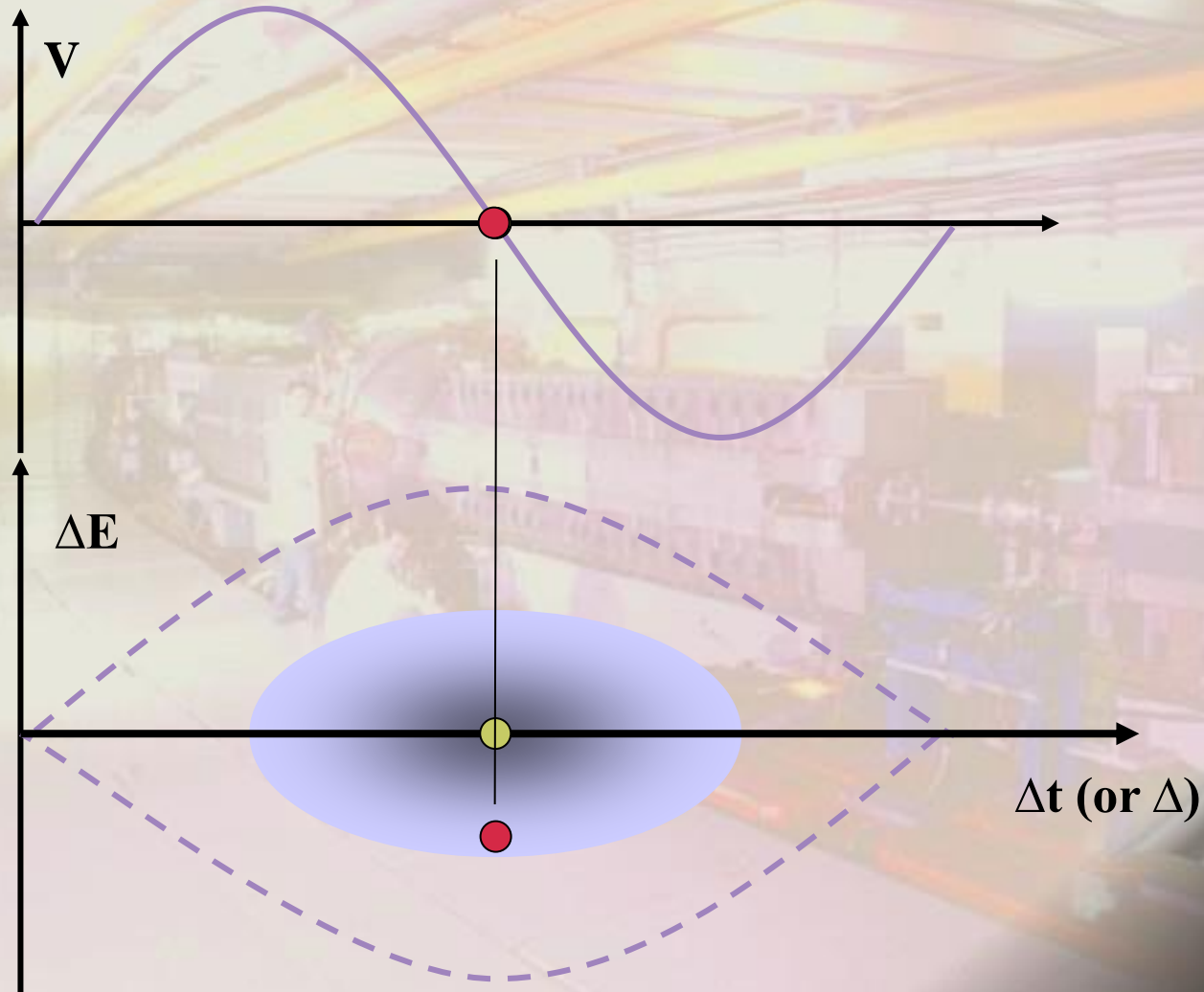
# The motion in the bucket (5)



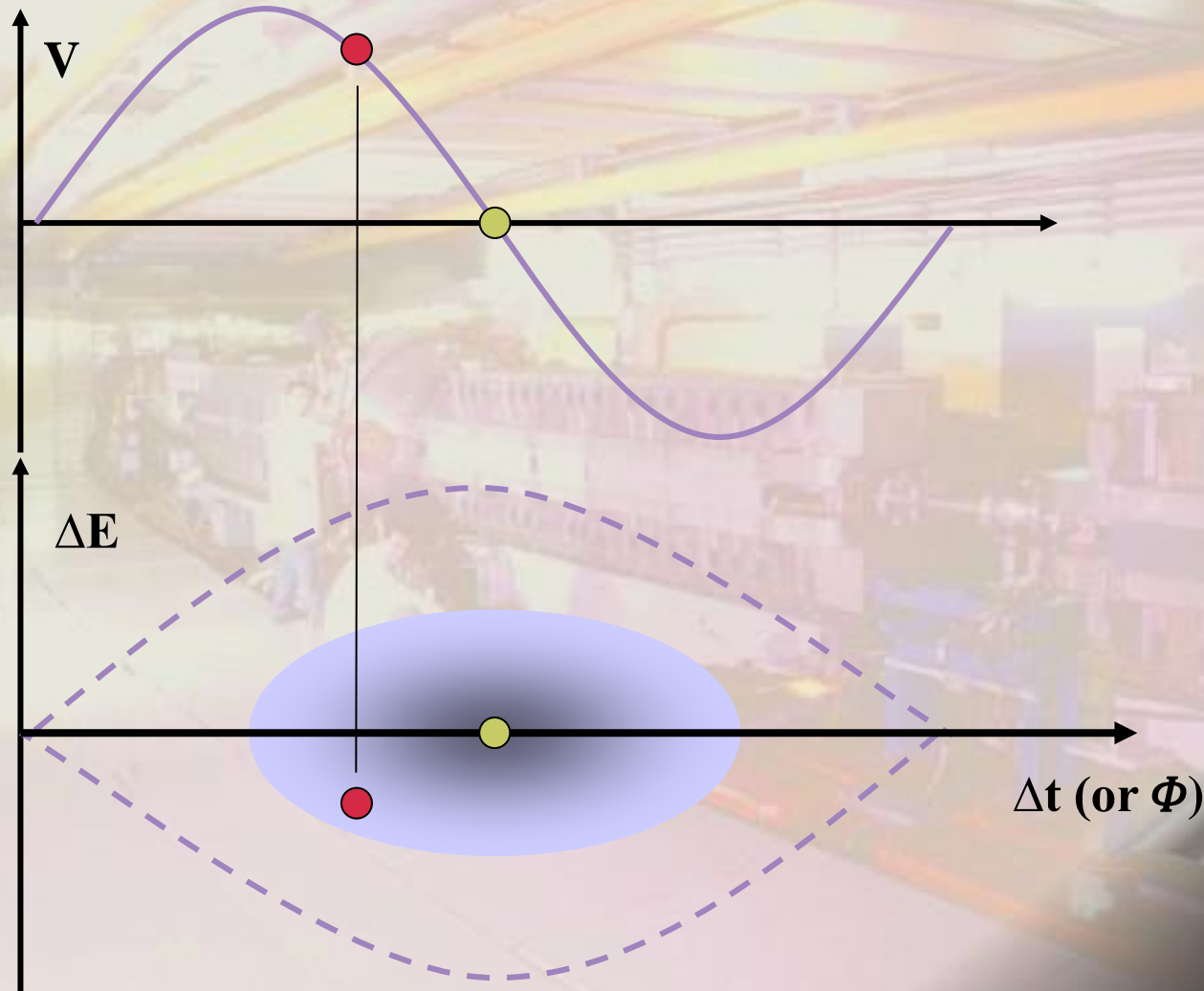
# The motion in the bucket (6)



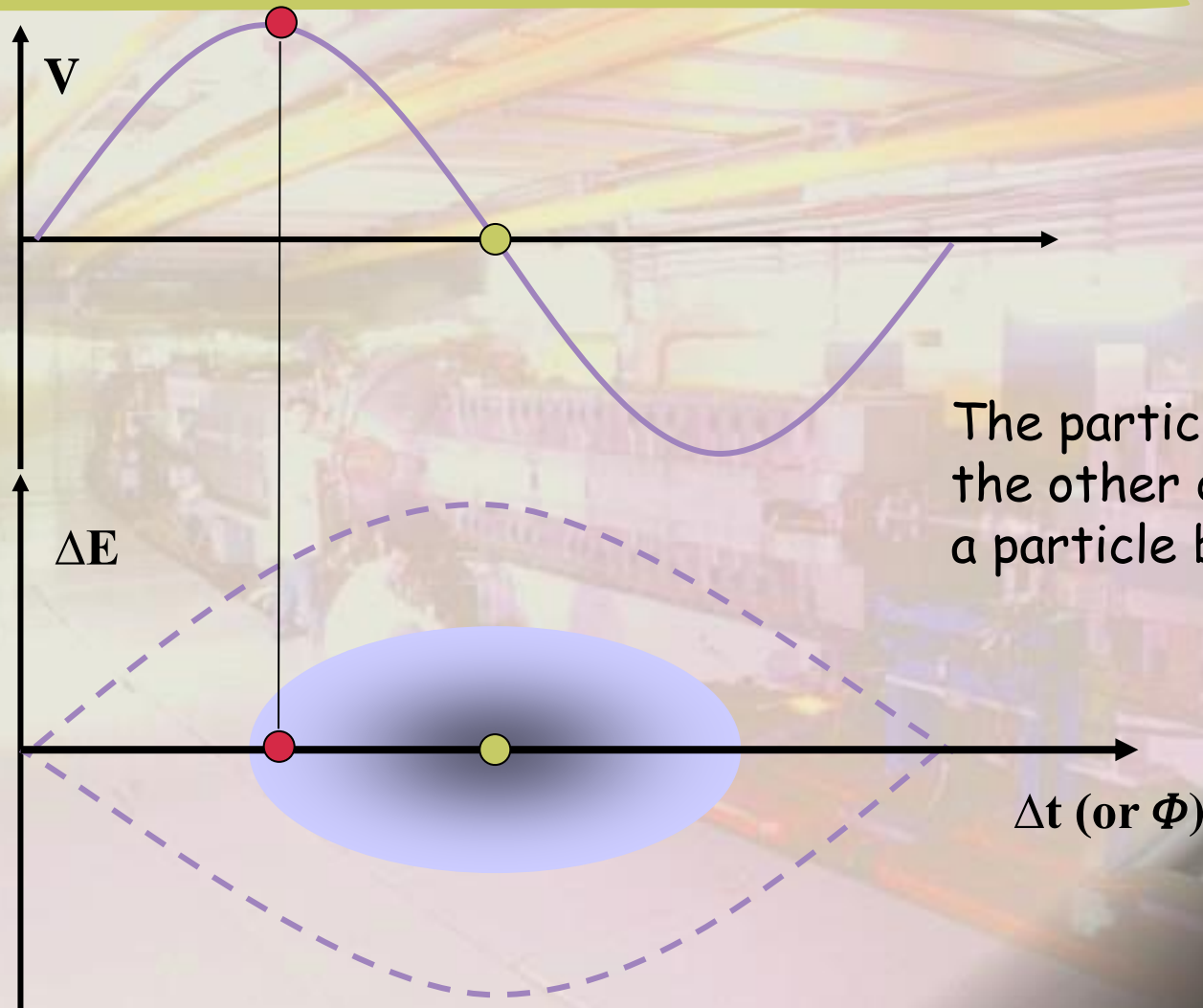
# The motion in the bucket (7)



# The motion in the bucket (8)

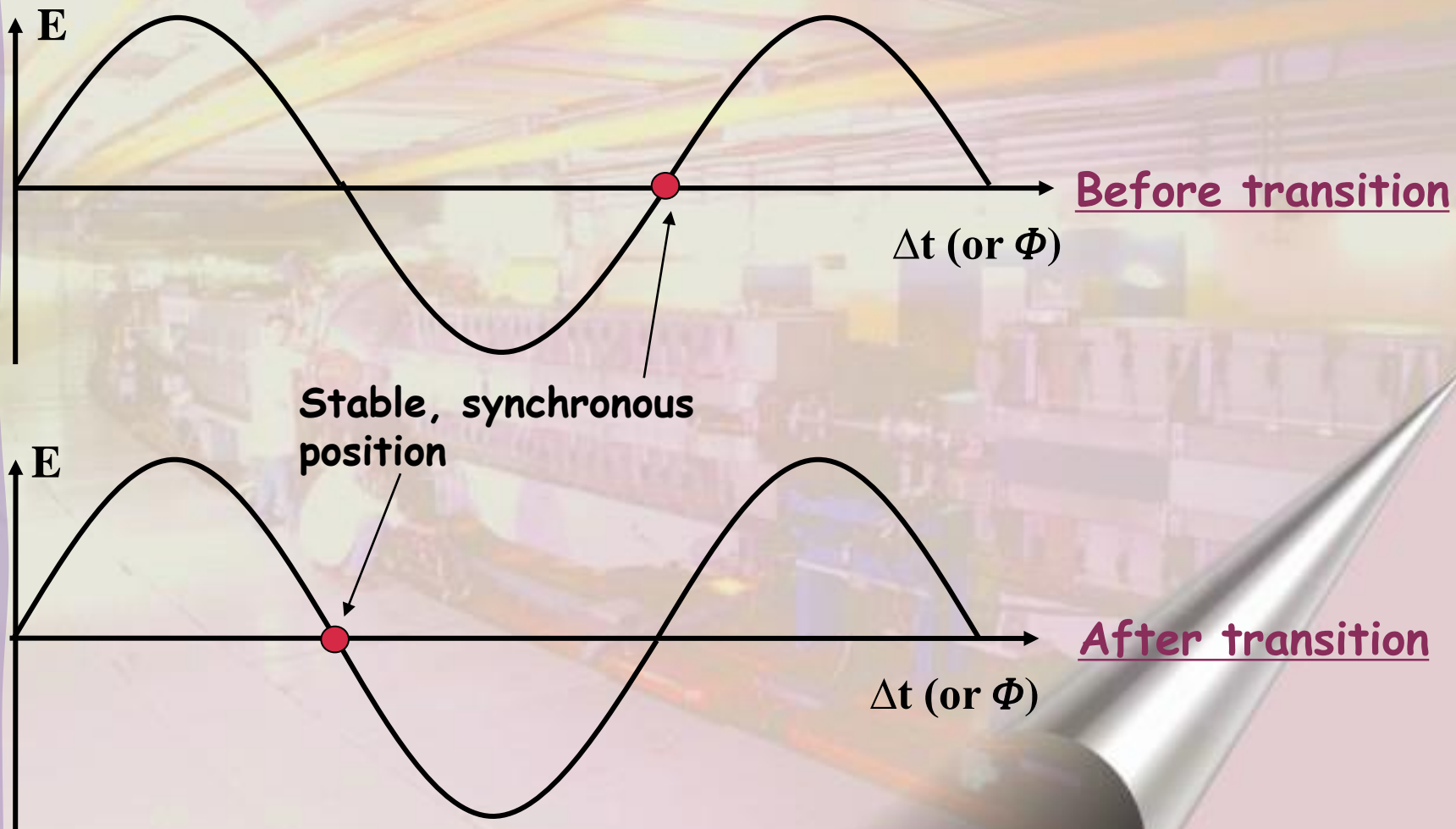


# The motion in the bucket (9)



The particle now turns in the other direction w.r.t. a particle below transition

# Before and After Transition



# Transition crossing in the PS

- # Transition in the PS occurs around 6 GeV/c
  - Injection happens at 2.12 GeV/c
  - Ejection can be done at 3.5 GeV/c up to 26 GeV/c
- # Therefore the particles in the PS must nearly always cross transition.
- # The beam must stay bunched
- # Therefore the phase of the RF must “jump” by  $\pi$  at transition



# Harmonic number (1)

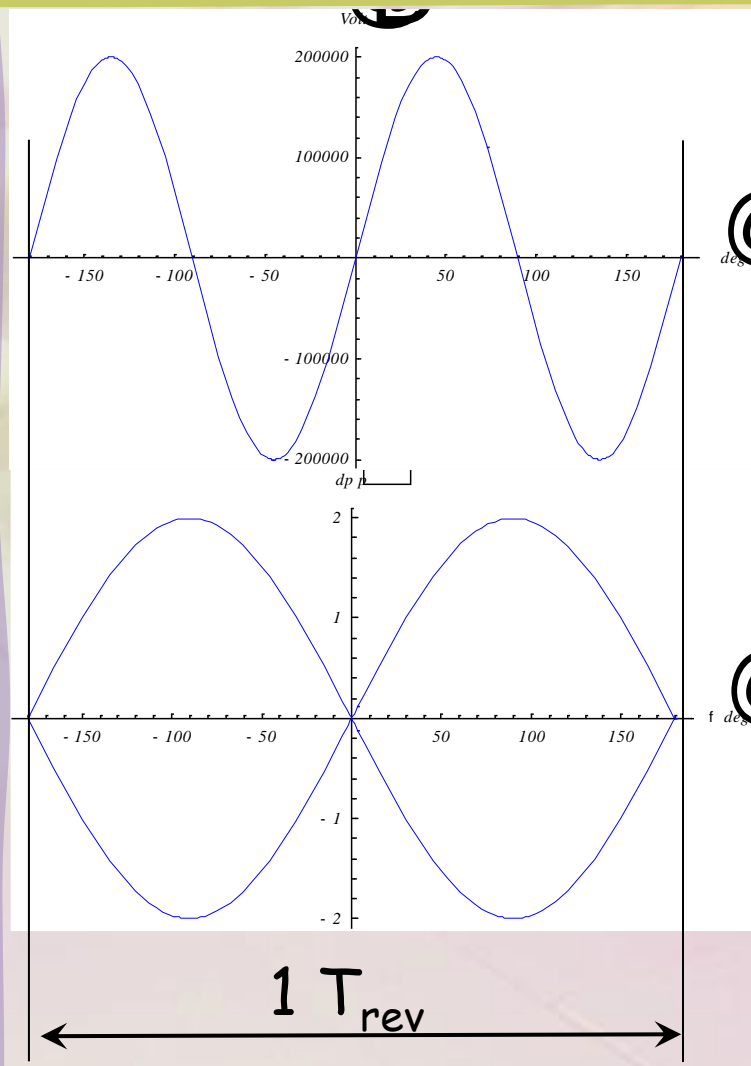
- # Until now we have applied an oscillating voltage with a frequency equal to the revolution frequency.

$$F_{\text{rf}} = F_{\text{rev}}$$

- # What will happen when  $F_{\text{rf}}$  is a multiple of  $f_{\text{rev}}$  ???

$$F_{\text{rf}} = h \times F_{\text{rev}}$$

# Harmonic number (2)



$$F_{rf} = h \times F_{rev}$$

Frequency of cavity voltage

Harmonic number

Variable for  $\beta < 1$

# Then we will have h buckets

# Frequency of the synchrotron oscillation (1)

- # On each turn the phase,  $\Phi$ , of a particle w.r.t. the RF waveform changes due to the synchrotron oscillations.

$$\frac{d\phi}{dt} = 2\pi h \Delta f_{rev}$$

Harmonic number

Change in revolution frequency

- # We know that  $\frac{df_{rev}}{f_{rev}} = -\eta \frac{dE}{E}$

- # Combining this with the above  $\therefore \frac{d\phi}{dt} = \frac{-2\pi h \eta}{E} \cdot dE \cdot f_{rev}$

- # This can be written as

$$\frac{d^2\phi}{dt^2} = \frac{-2\pi h \eta}{E} \cdot f_{rev} \cdot \frac{dE}{dt}$$

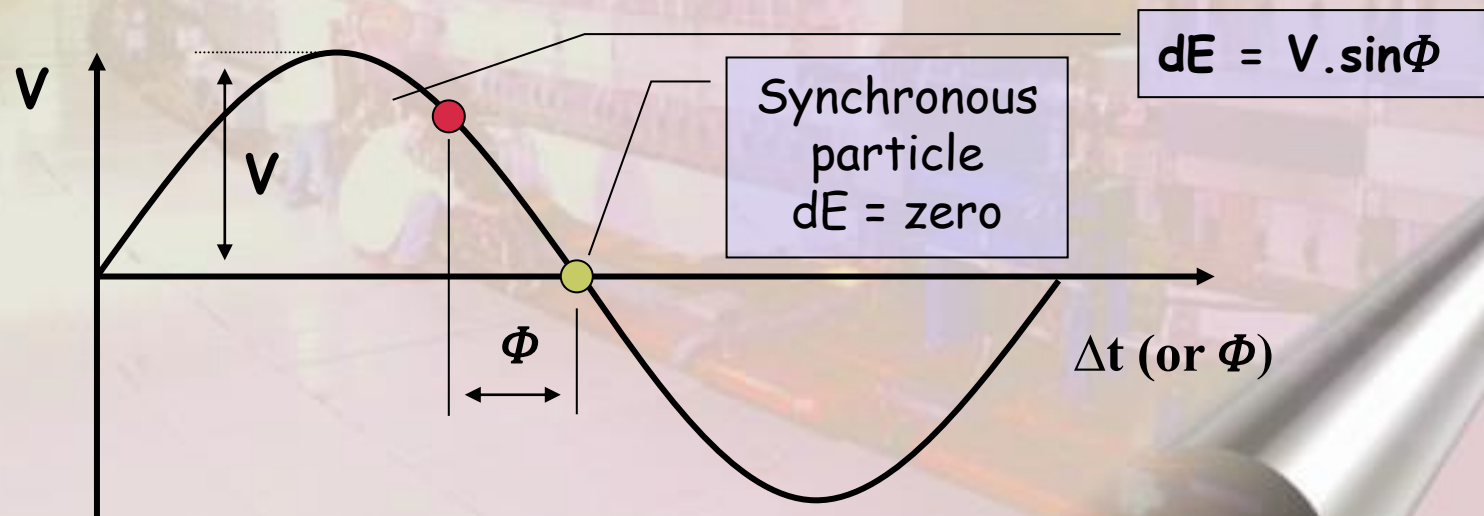
Change of energy as a function of time

# Frequency of the synchrotron oscillation (2)

# So, we have:

$$\frac{d^2\phi}{dt^2} = \frac{-2\pi h\eta}{E} \cdot f_{rev} \cdot \frac{dE}{dt}$$

# Where  $dE$  is just the energy gain or loss due to the RF system during each turn

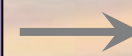


# Frequency of the synchrotron oscillation (3)

$$\frac{d^2\phi}{dt^2} = \frac{-2\pi h\eta}{E} \cdot f_{rev} \frac{dE}{dt}$$

and

$$dE = V \sin \phi$$



$$\frac{dE}{dt} = f_{rev} V \sin \phi$$

$$\frac{d^2\phi}{dt^2} = \frac{-2\pi h\eta}{E} \cdot f_{rev}^2 \cdot V \cdot \sin \phi$$

# If  $\phi$  is small then  $\sin\phi = \phi$

$$\frac{d^2\phi}{dt^2} + \left( \frac{2\pi h\eta}{E} \cdot f_{rev}^2 \cdot V \right) \phi = 0$$

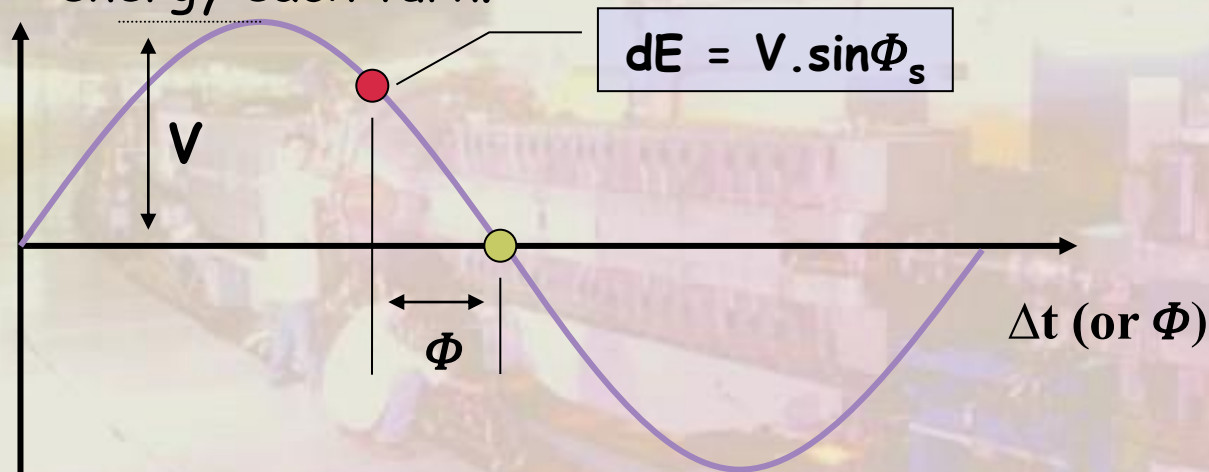
# This is a SHM where the synchrotron oscillation frequency is given by:

Synchrotron  
tune  $Q_s$

$$\left( \sqrt{\frac{2\pi h\eta V}{E}} \right) \cdot f_{rev}$$

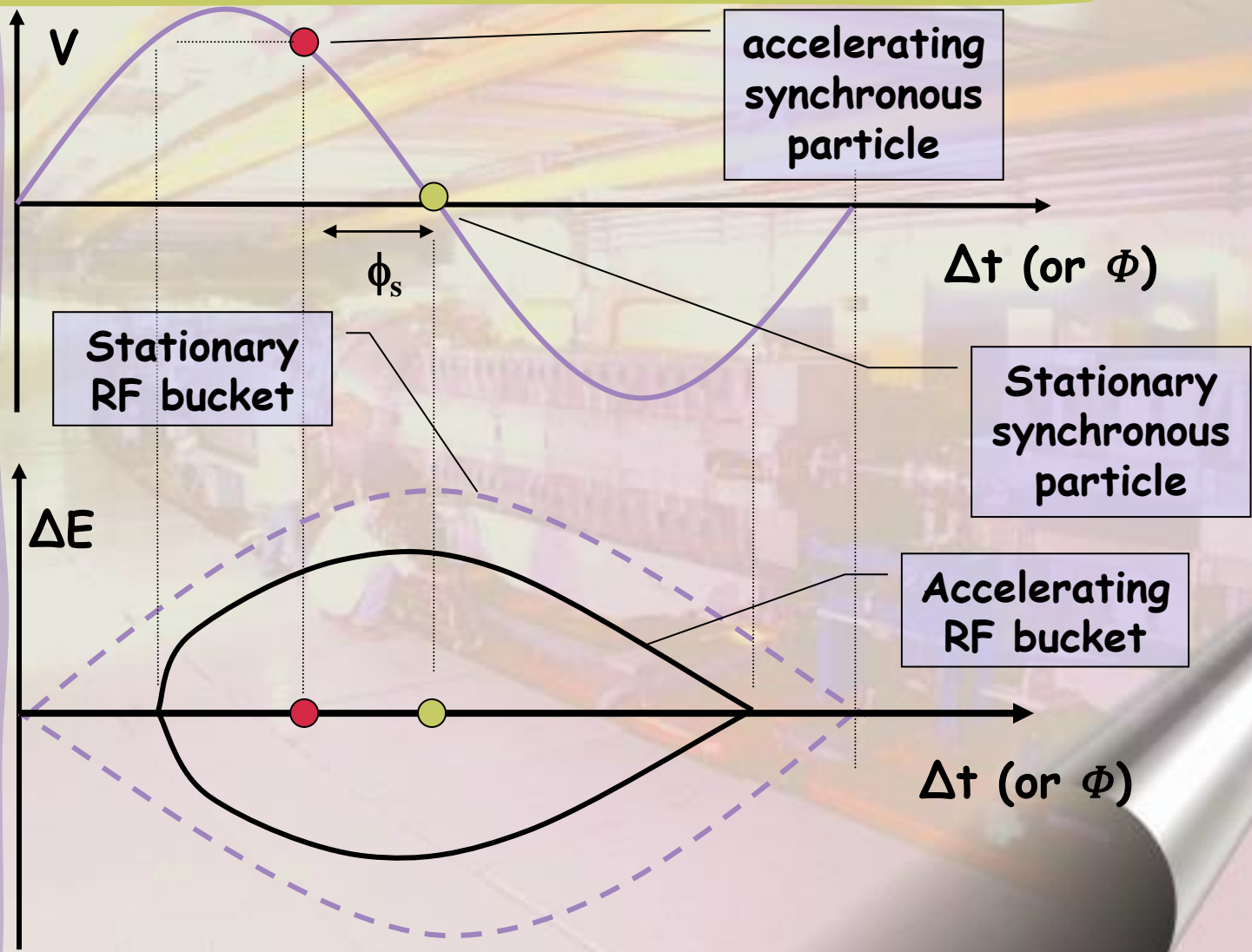
# Acceleration

- # Increase the magnetic field slightly on each turn.
- # The particles will follow a shorter orbit. ( $F_{\text{rev}} < F_{\text{synch}}$ )
- # Beyond transition, early arrival in the cavity causes a gain in energy each turn.



- # We change the phase of the cavity such that the new synchronous particle is at  $\Phi_s$  and therefore always sees an accelerating voltage
- #  $V_s = V \sin \Phi_s = V \Gamma = \text{energy gain/turn} = dE$

# Acceleration & RF bucket shape (1)



# Acceleration & RF bucket shape (2)

- # The modification of the RF bucket reduces the acceptance
- # The faster we accelerate (increasing  $\sin \phi_s$ ) the smaller the acceptance
- # Faster acceleration also modifies the synchrotron tune.
- # For a stationary bucket ( $\phi_s = 0$ ) we had:

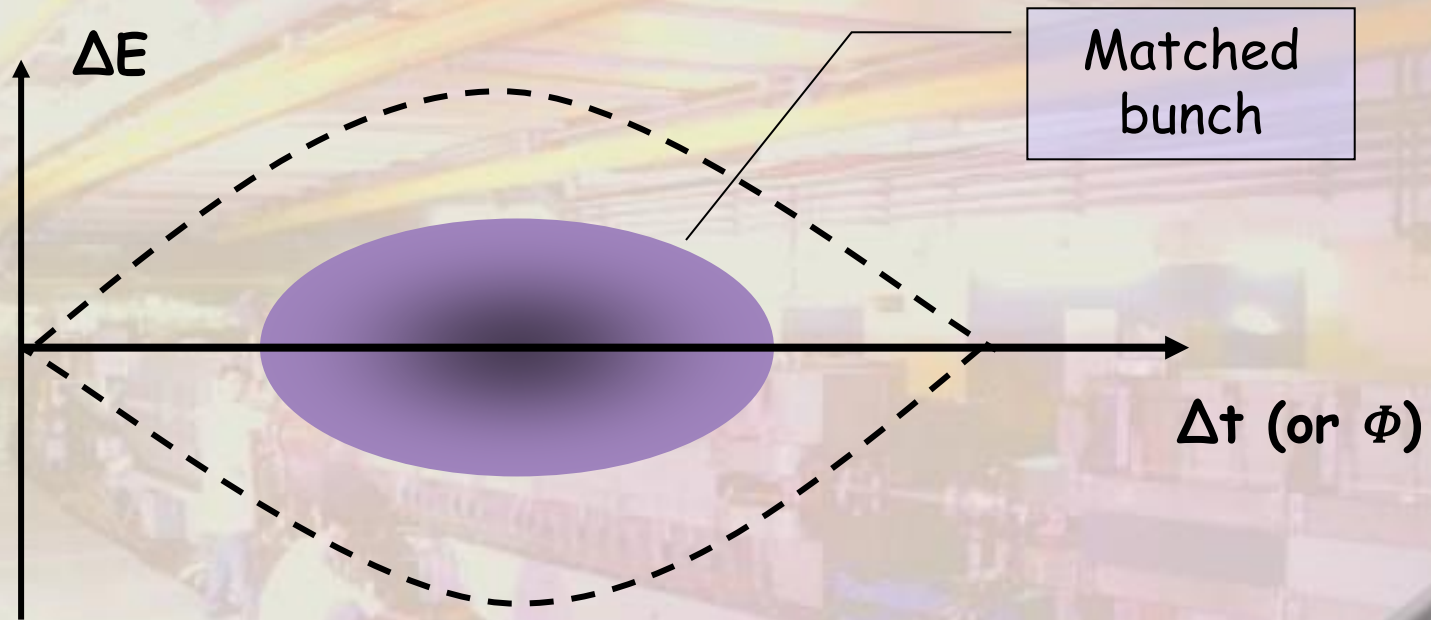
$$\left( \sqrt{\frac{2\pi h \eta}{E}} \right) \cdot f_{rev}$$

- # For a moving bucket ( $\phi_s \neq 0$ ) this becomes:

$$\left( \sqrt{\frac{2\pi h \eta}{E}} \right) \cdot f_{rev} \cos \phi_s$$

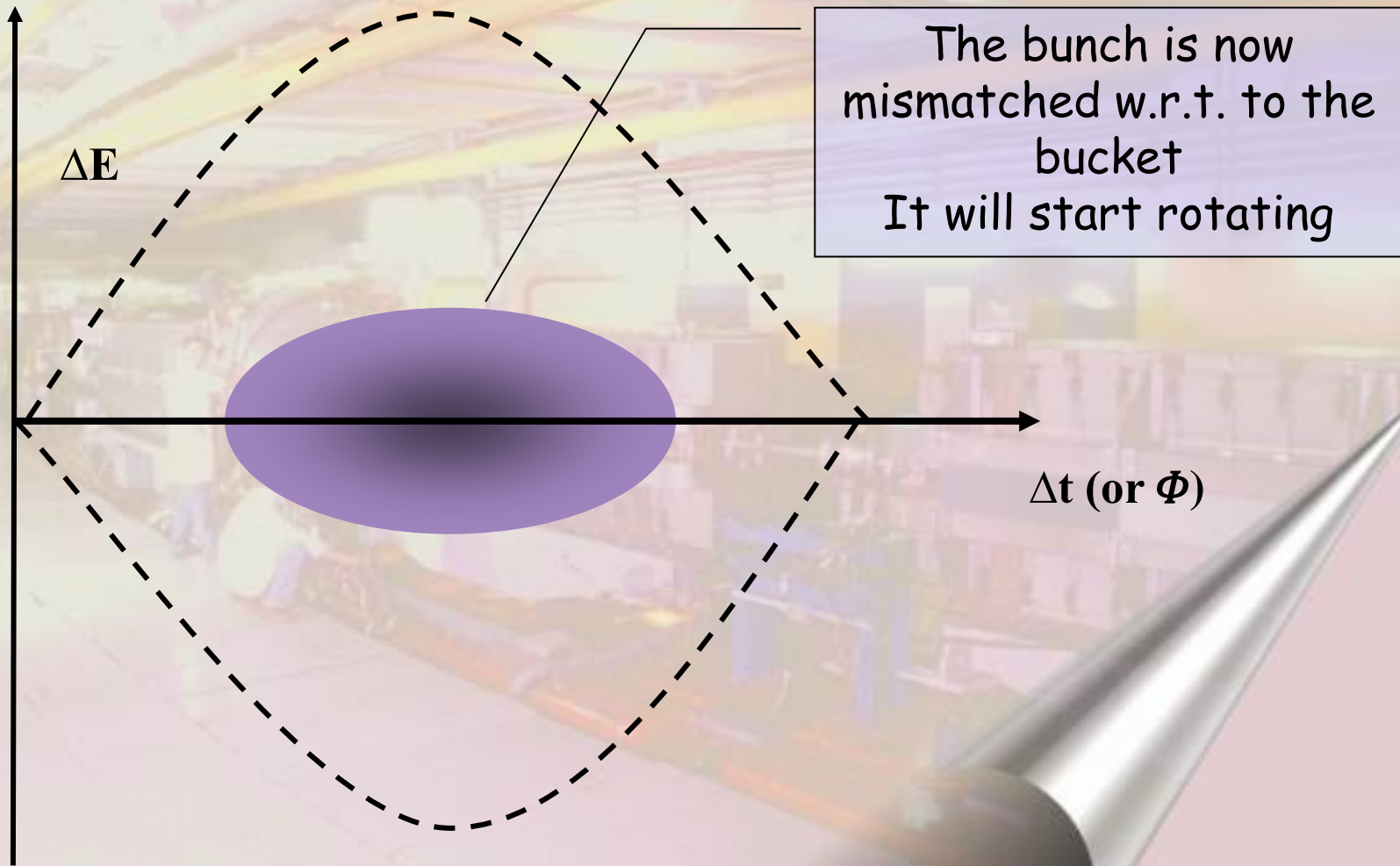


# Non-adiabatic change (1)

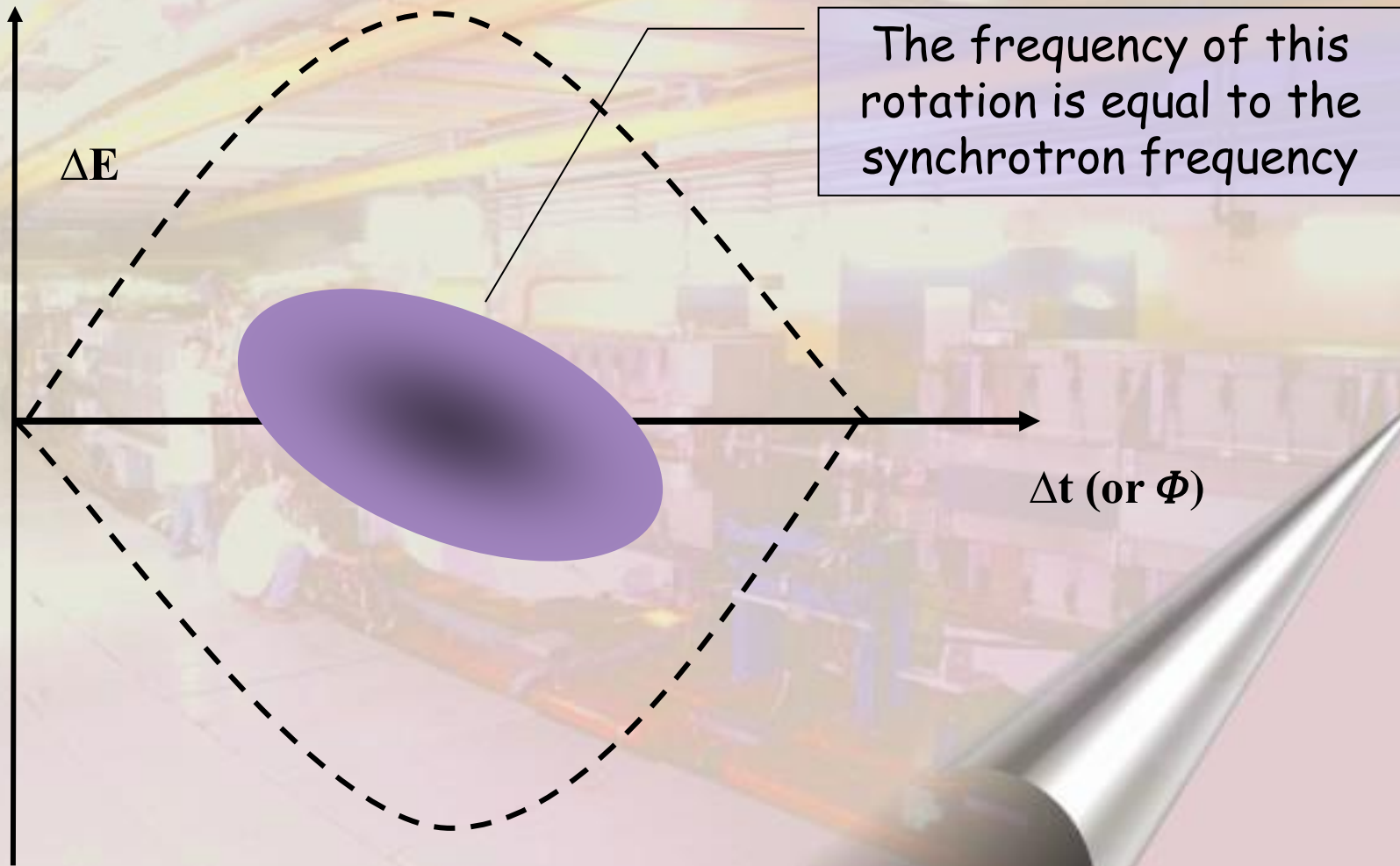


- # What will happen when we increase the voltage rapidly ?

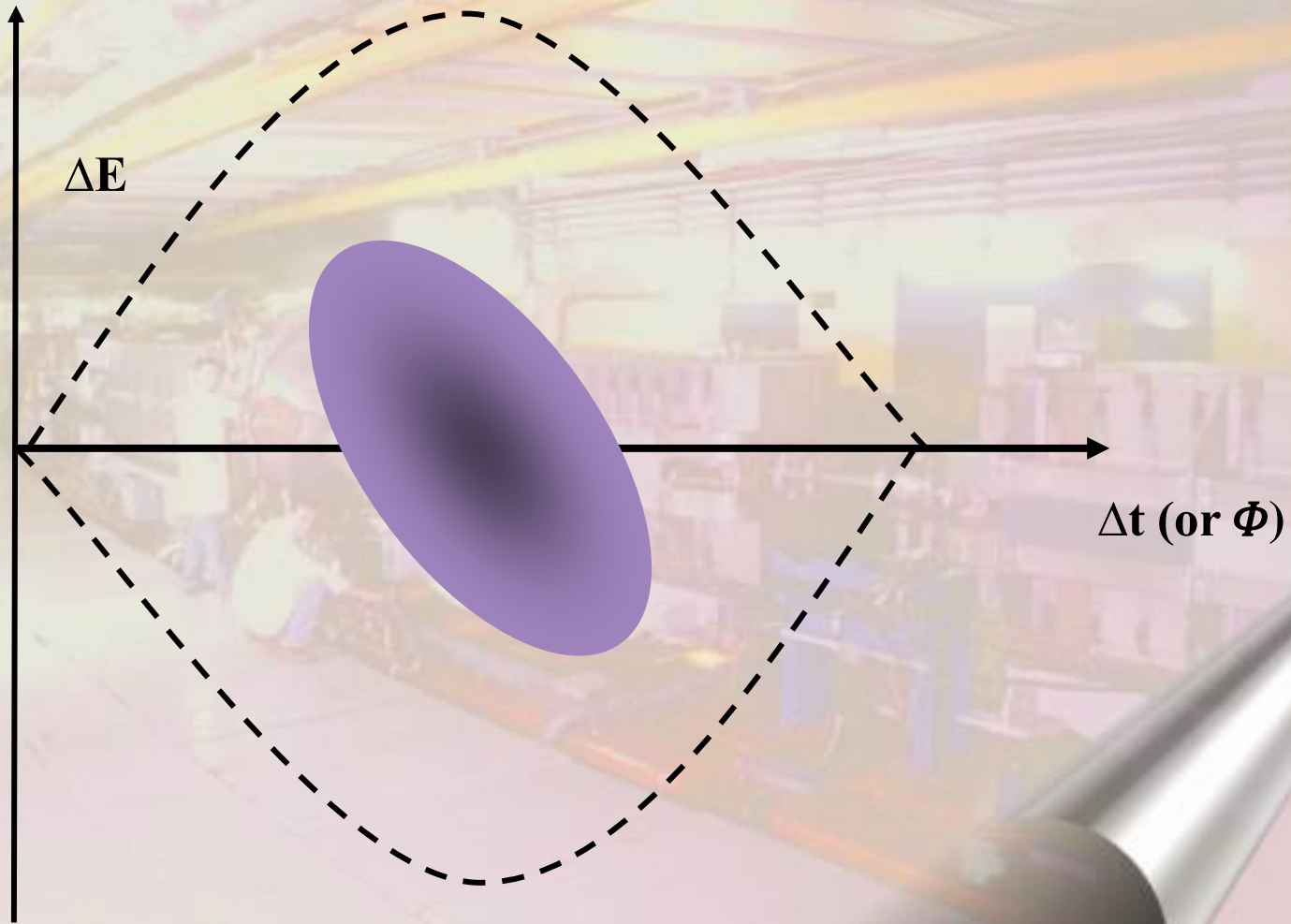
# Non-adiabatic change (2)



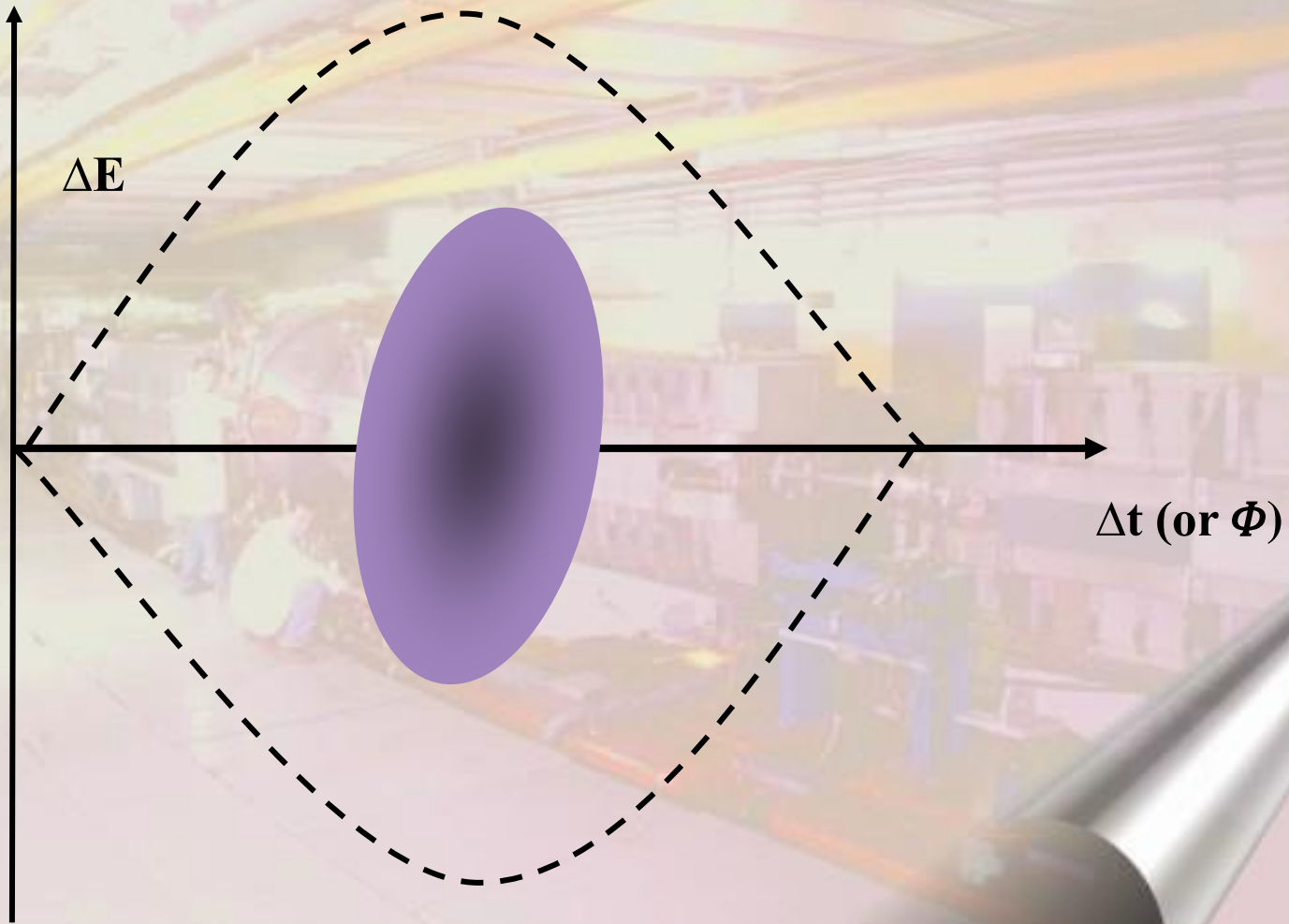
# Non-adiabatic change (3)



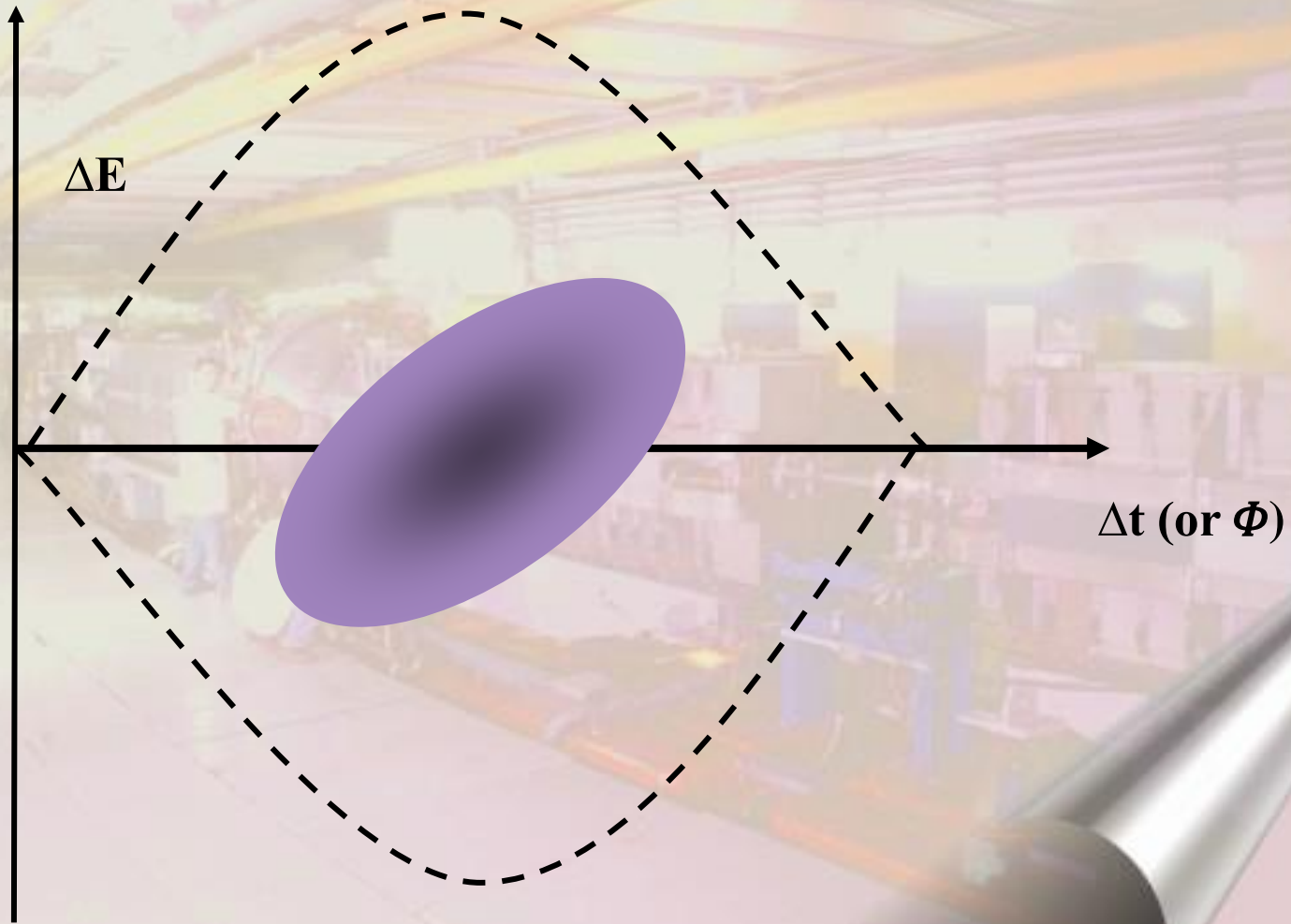
# Non-adiabatic change (4)



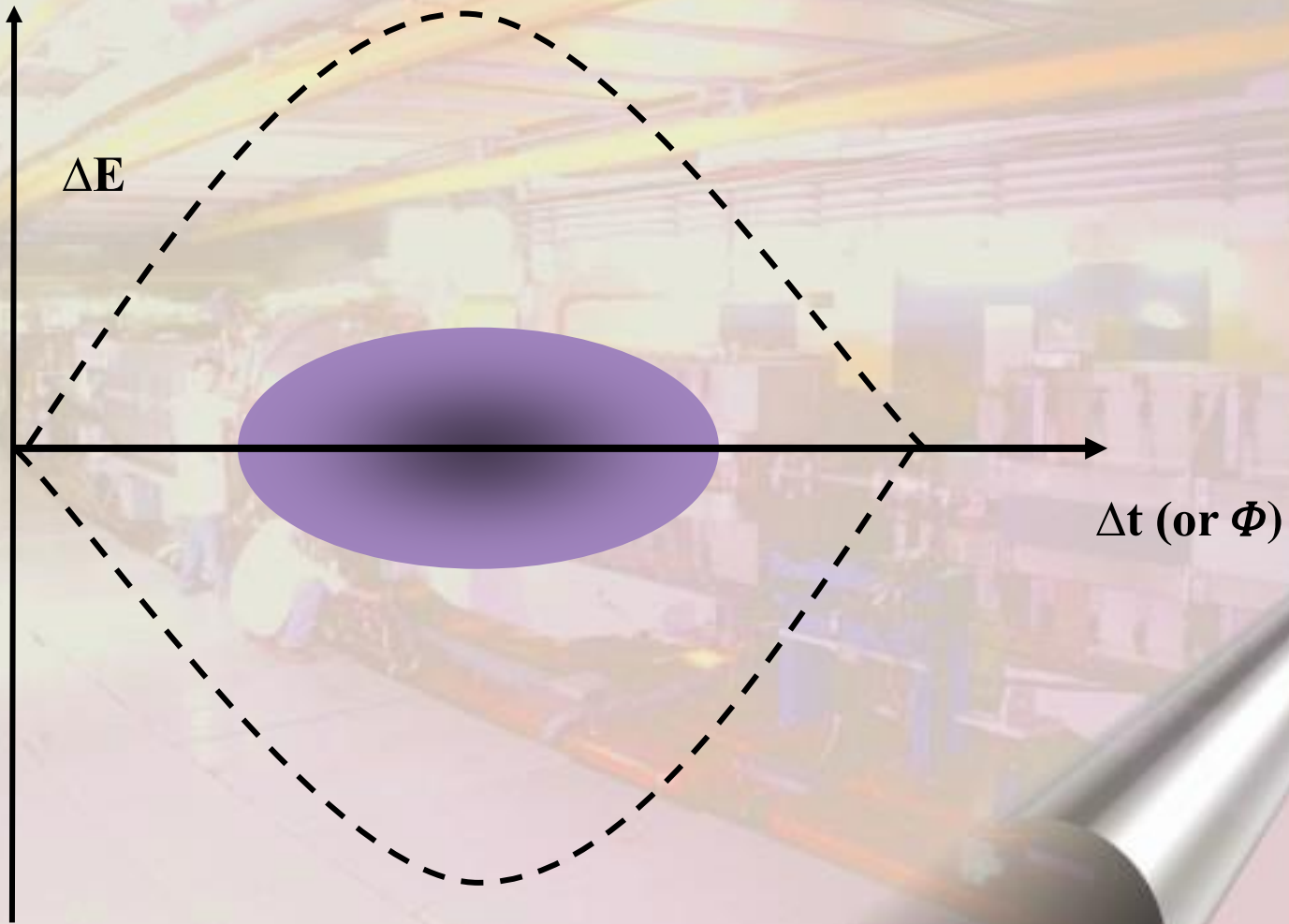
# Non-adiabatic change (5)



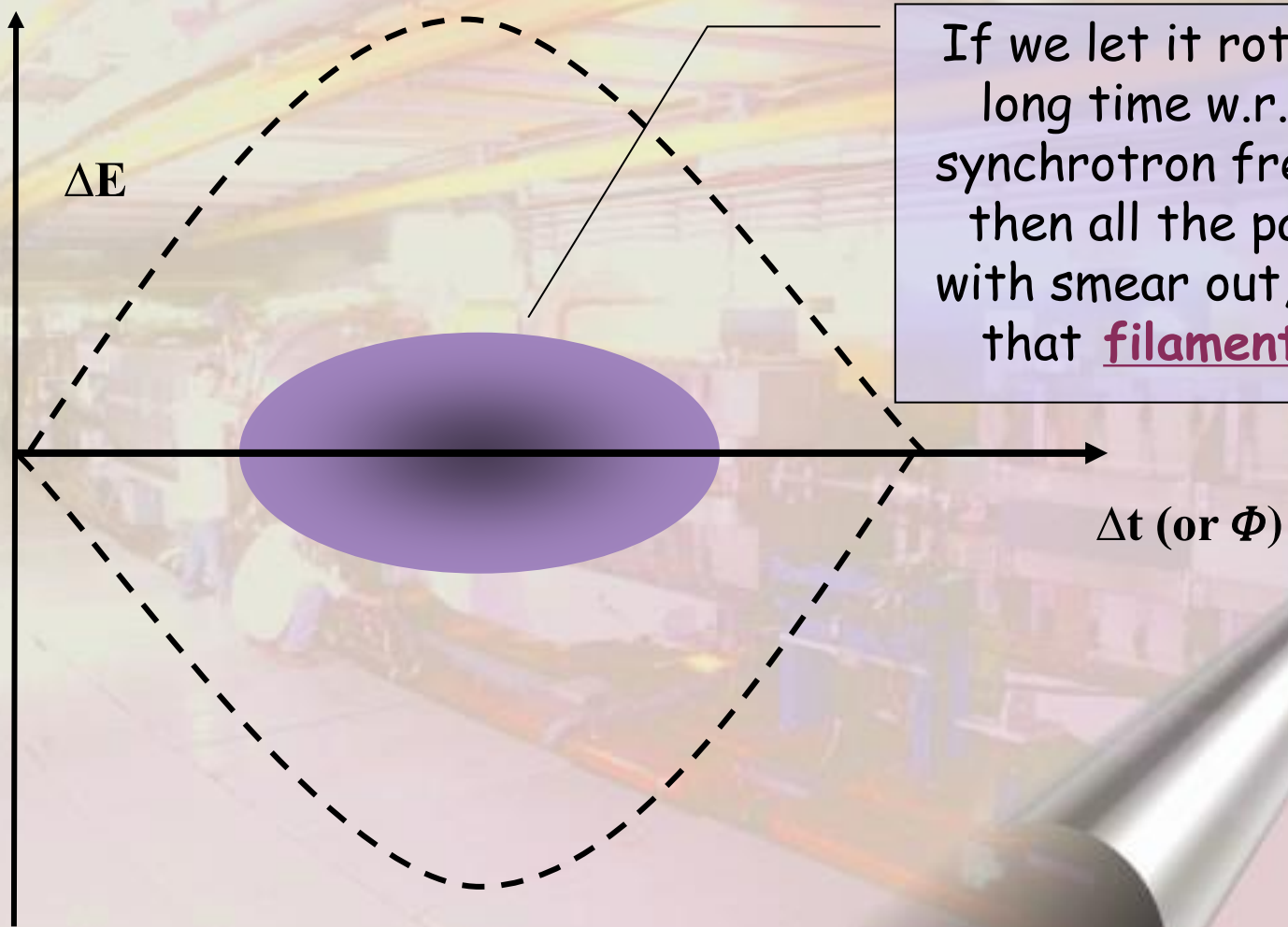
# Non-adiabatic change (6)



# Non-adiabatic change (7)



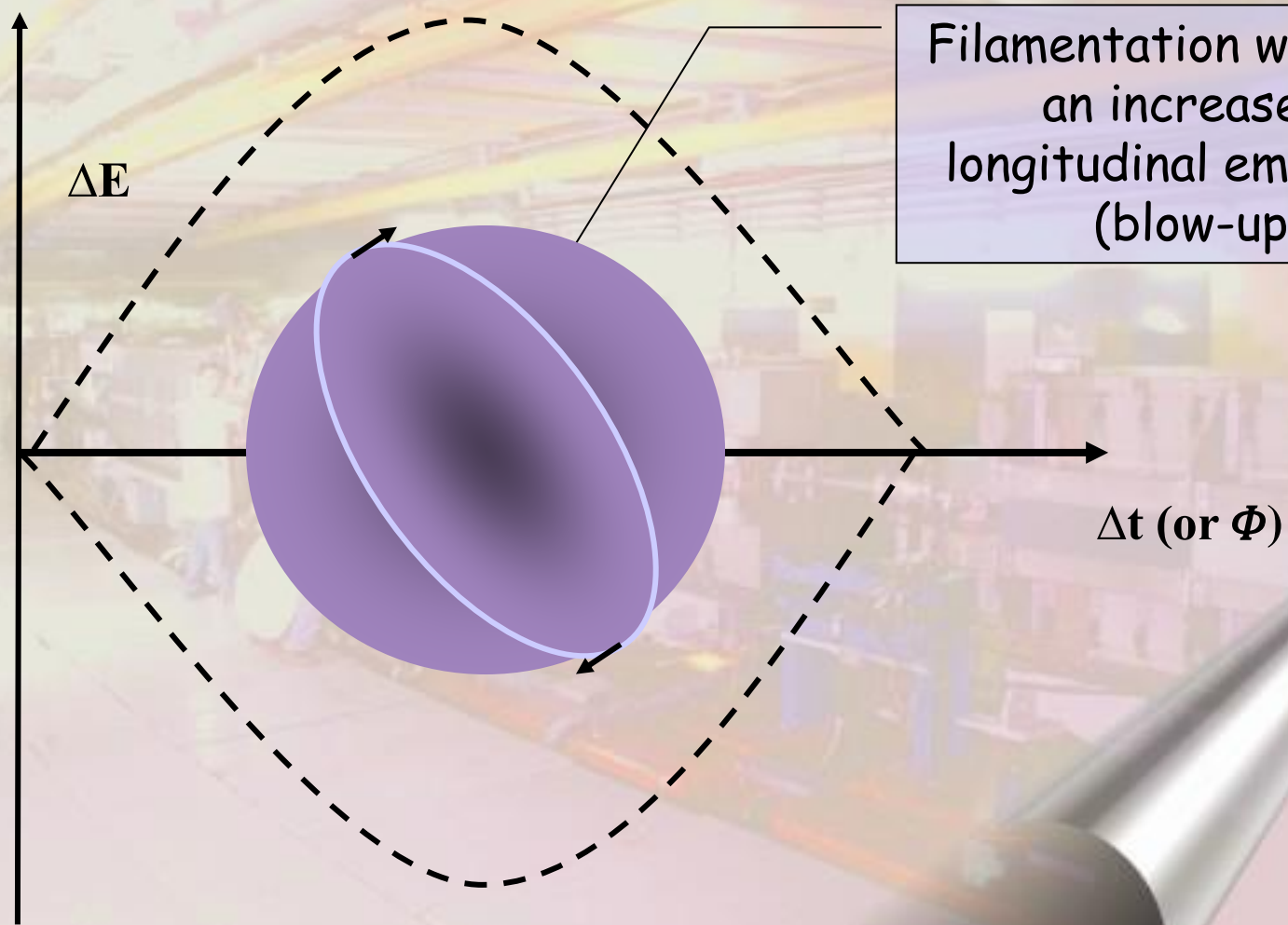
# Non-adiabatic change (8)



If we let it rotate for long time w.r.t. the synchrotron frequency then all the particle smear out, we call that filamentation

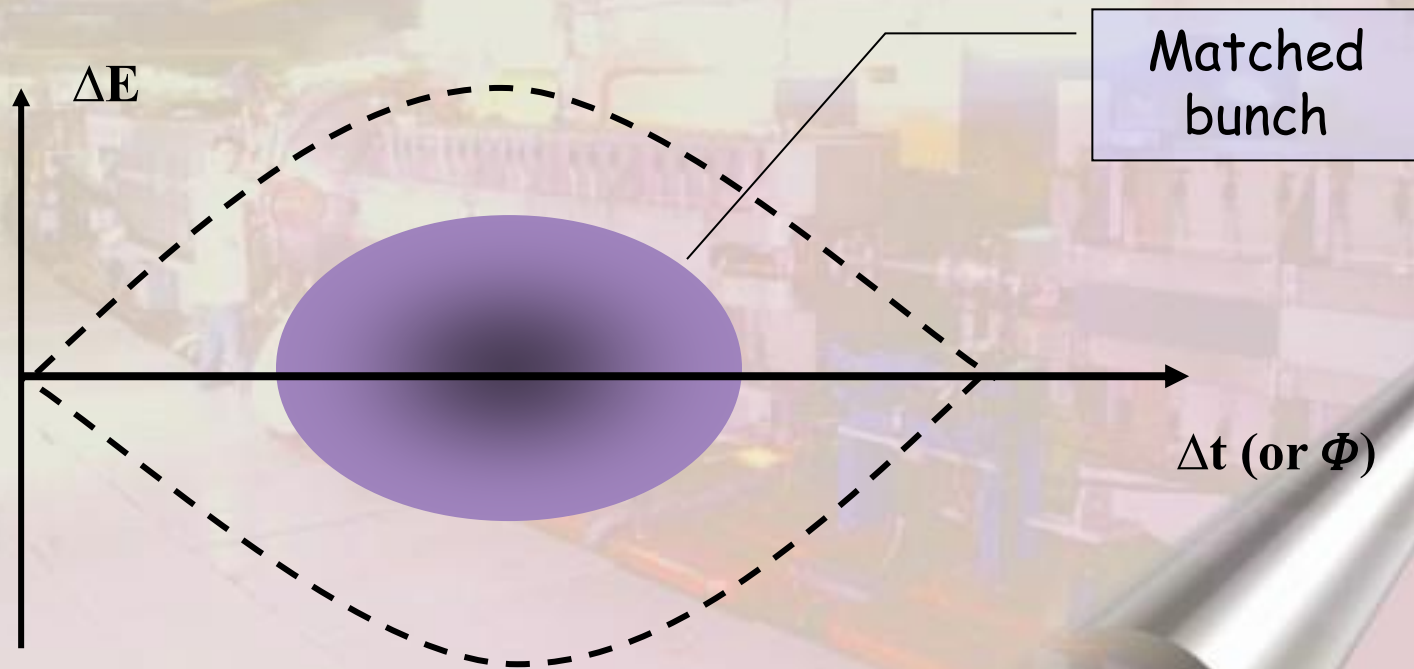


# Non-adiabatic change (9)

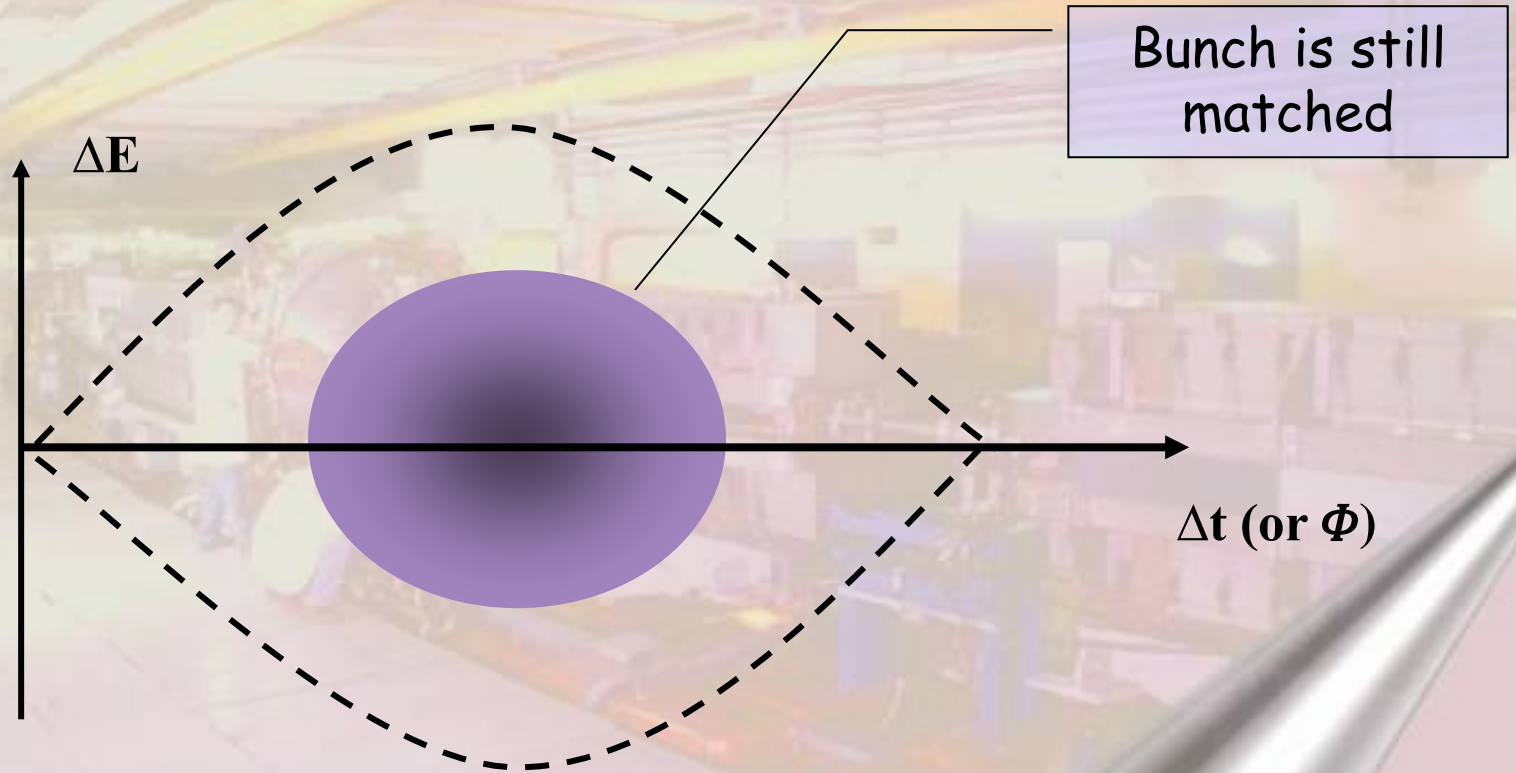


# Adiabatic change (1)

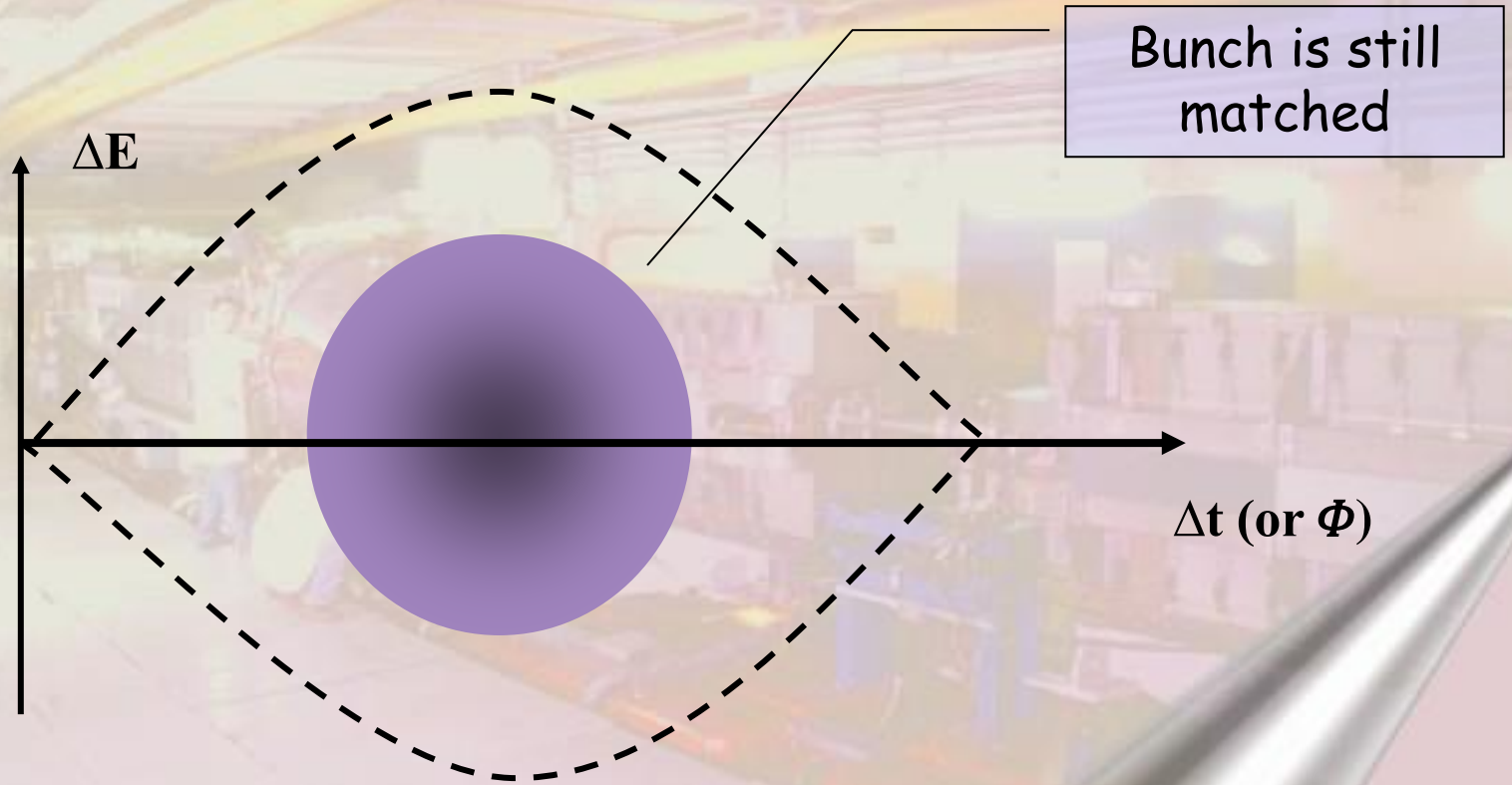
- # To avoid this filamentation we have to change slowly w.r.t. the synchrotron frequency.
- # This is called 'Adiabatic' change.



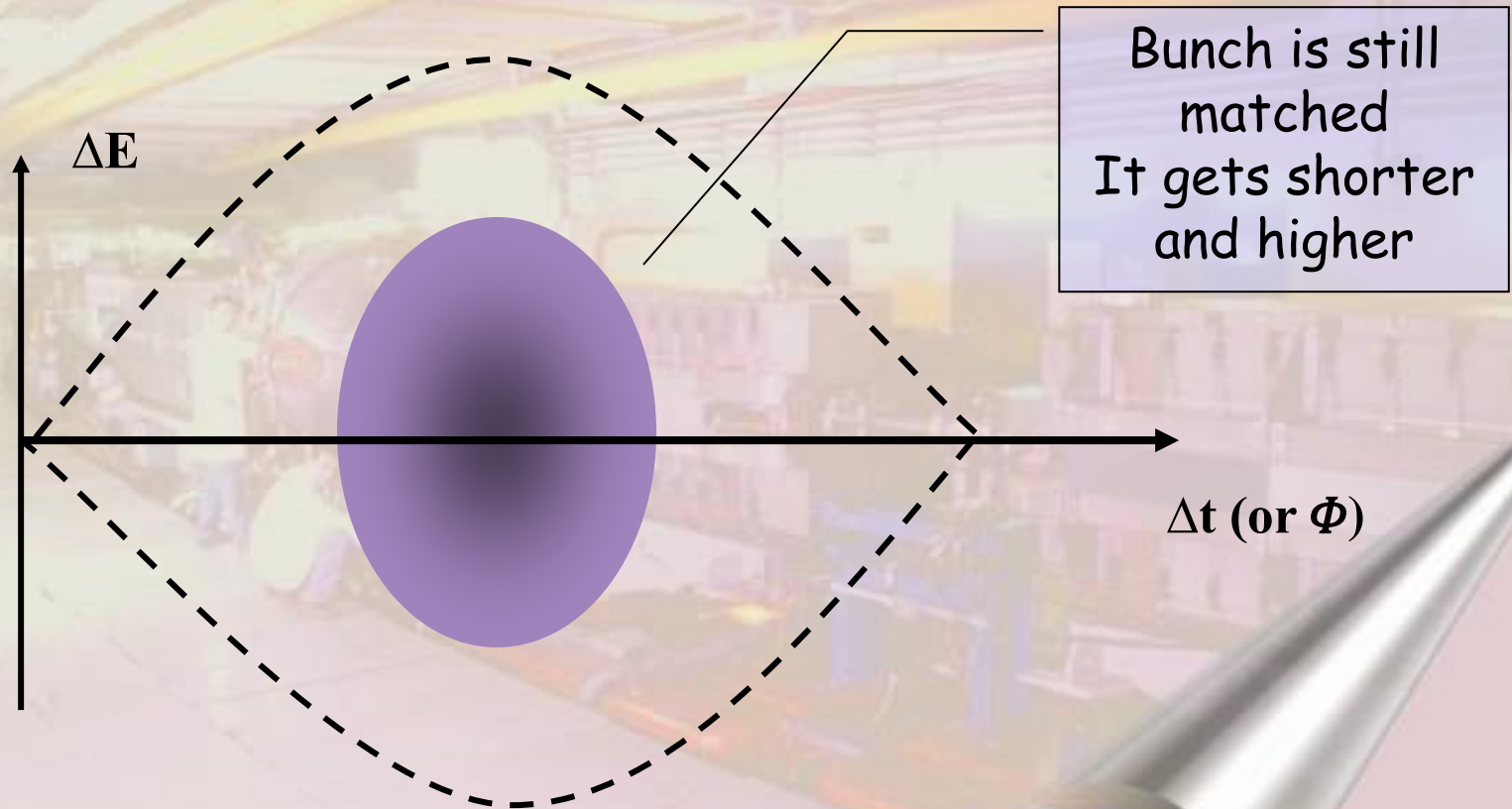
# Adiabatic change (2)



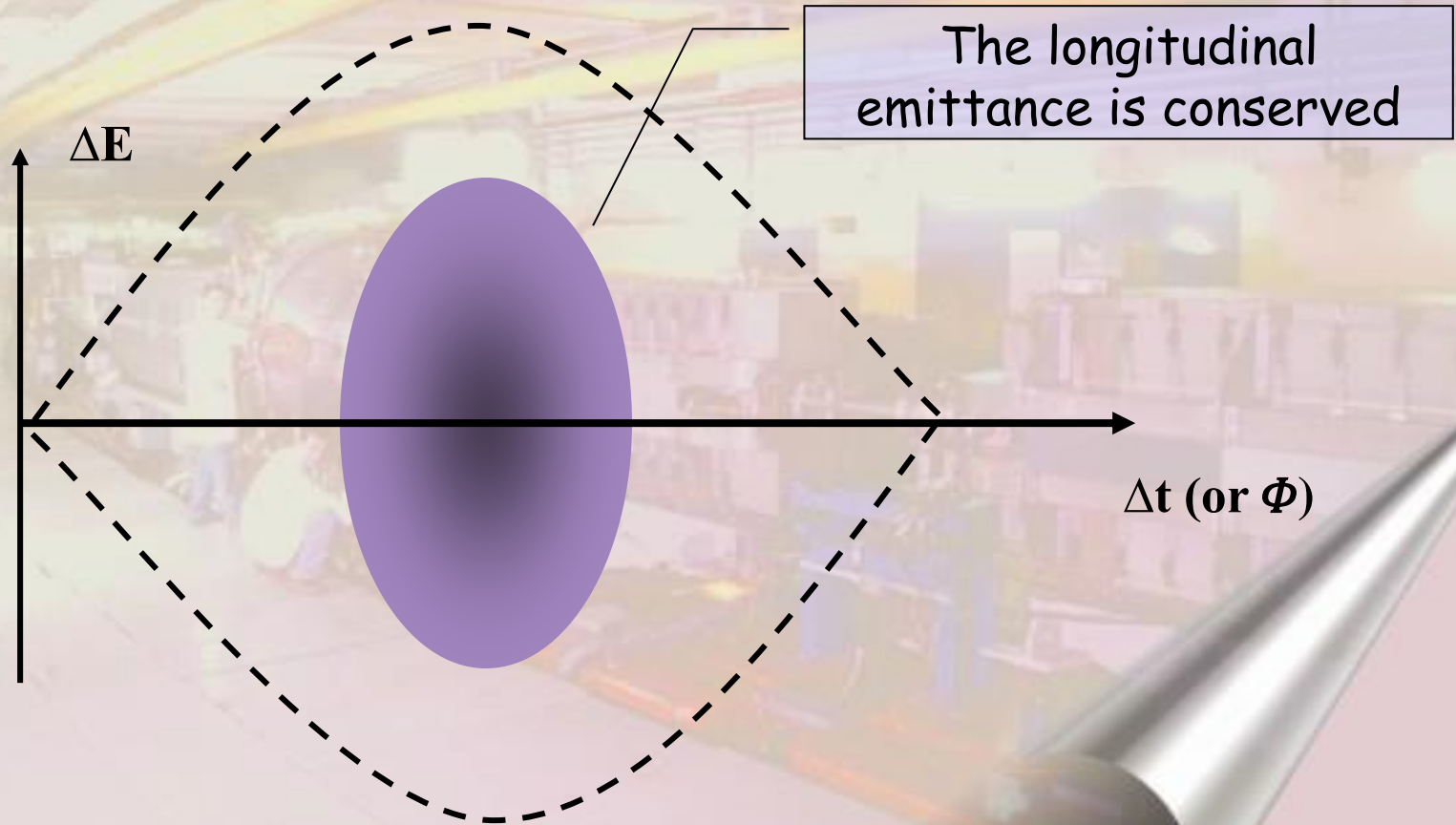
# Adiabatic change (3)



# Adiabatic change (4)



# Adiabatic change (5)



# Questions....,Remarks...?

*Longitudinal  
Phase space*

*Transition*

*Acceleration*

*Adiabatic &  
non-adiabatic  
changes*

