

Connections to BSM

Kevin J. Kelly, Texas A&M University
NuSTEC Summer School 2024
kjkelly@tamu.edu

Scenario 1: Little/No SM Background

Scenario 2: Theoretically Clean SM Background

(picking up with S2p2)

Scenario 3: Experimentally Measurable Bkg.

Back to our Simple "Dark Photon" $U(1)$

$$\mathcal{L} \supset -\frac{1}{4} F_{\mu\nu}' F^{\mu\nu} - \frac{\epsilon}{2} F_{\mu\nu}' F^{\mu\nu} + \frac{M_{A'}^2}{2} A_{\mu}' A^{\mu}$$

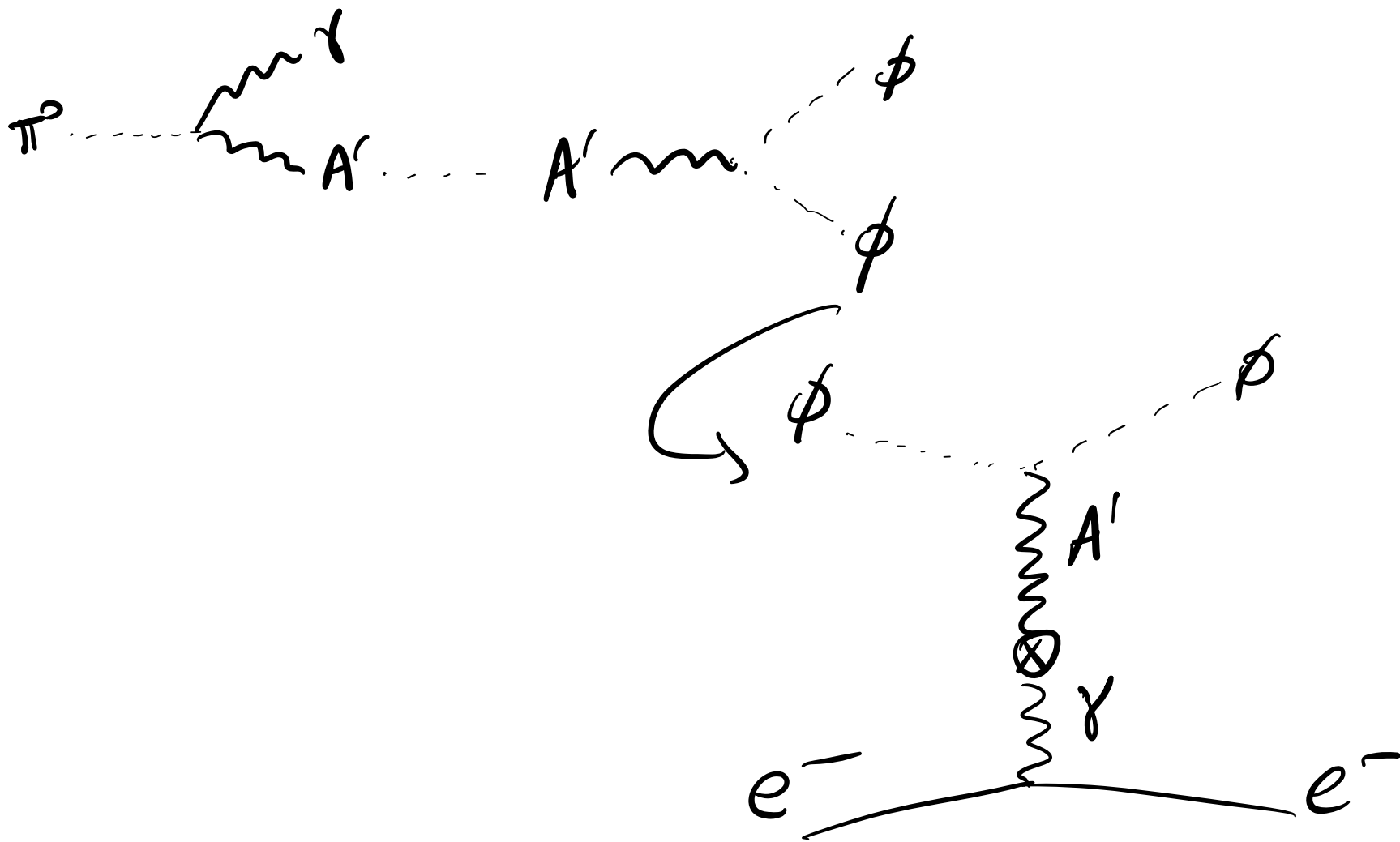
...plus one more ingredient

$$\mathcal{L} \supset |D_{\mu} \phi|^2 - M_{\phi}^2 |\phi|^2$$

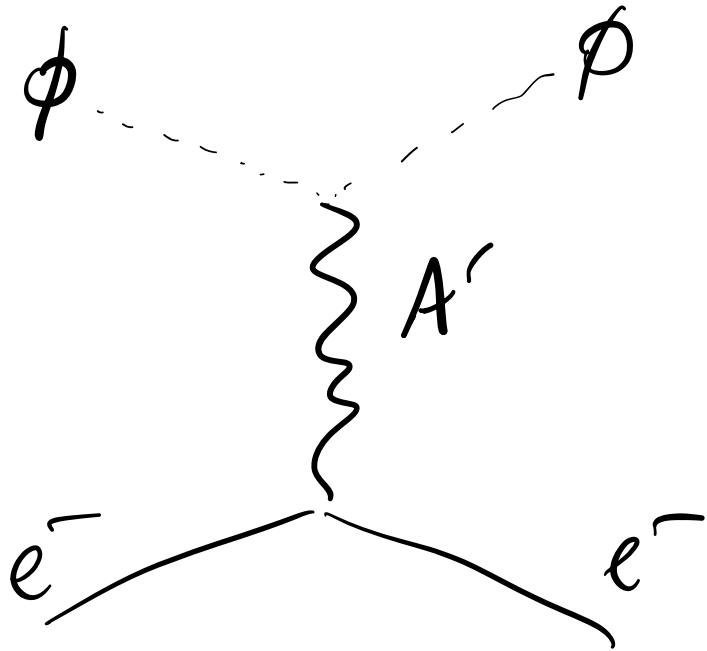
\hookrightarrow Scalar Particle
with Mass M_{ϕ}

$\phi - A'$ with Gauge Coupling
 $g_D \Rightarrow \alpha_D \equiv g_D^2 / 4\pi$

Impact in Neutrino Facilities?



Advantages of Light DM in Neutrino Beam

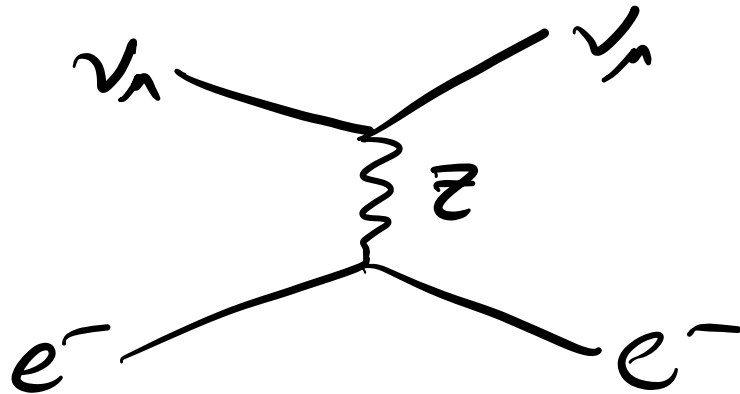


* ϕ Boosted: \sim GeV energies
for $g(\text{new})$ masses.

* $\mathcal{M} \sim \frac{1}{\underline{t} - m_{A'}^2}$
 $\hookrightarrow \underline{t} = -2m_e E_e^{\text{rec.}}$

* Many Target e^-

Problems of Light DM in Neutrino Beam!



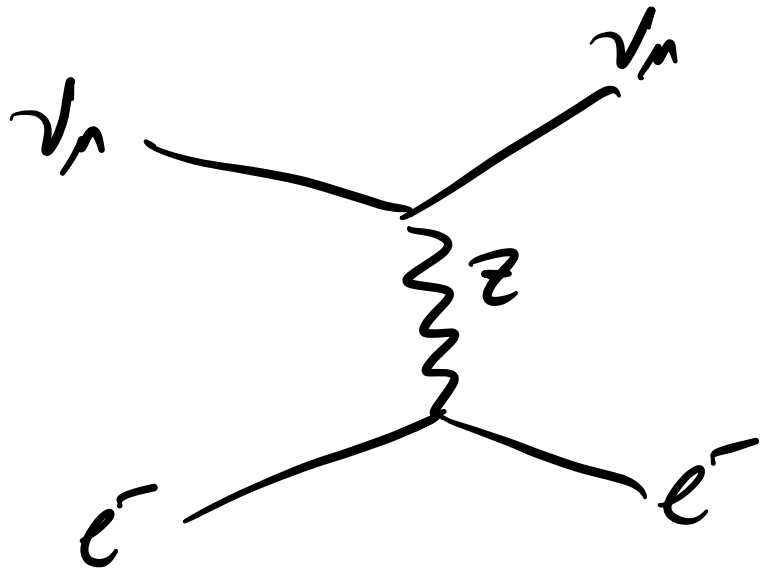
Why is this an existential problem?

* $\sigma(\nu_n e^- \rightarrow \nu_n e^-)$ very well understood

$\Rightarrow \Phi(E_\nu)$ constrained by this process

Aside: νe ES vs. ν_e CCQE

Consider the following two processes:

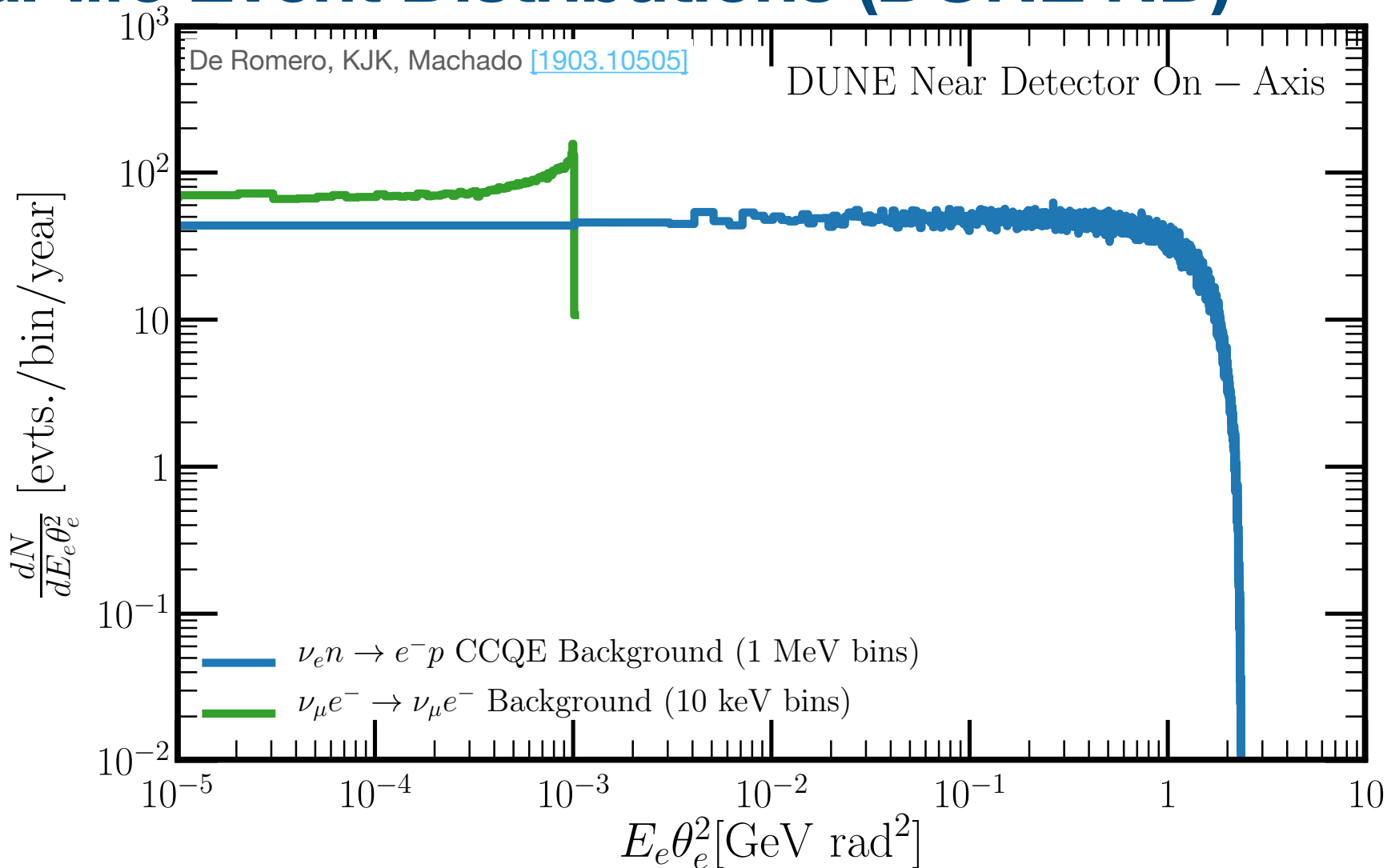


Below Threshold

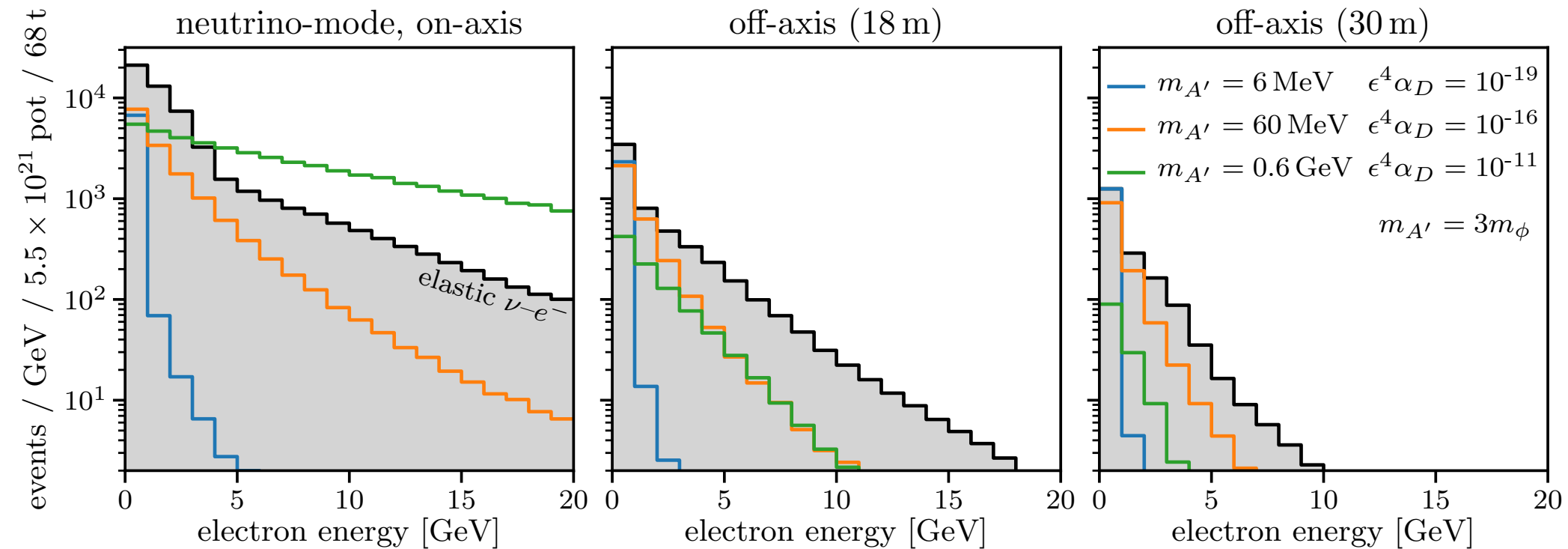
\Rightarrow Kinematics (S.R.)

$$\nu e \text{ ES: } E_e \times \theta_e^2 < 2m_e$$

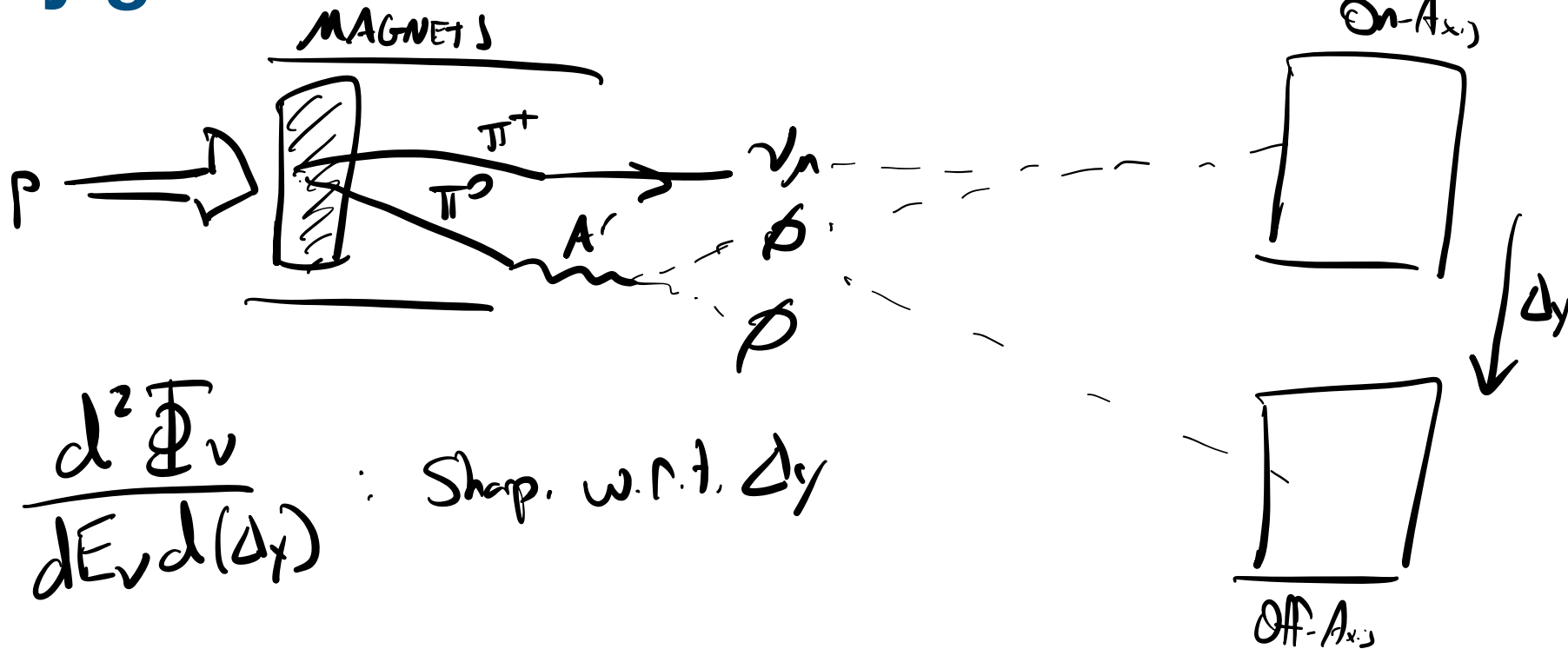
Real-life Event Distributions (DUNE ND)



DM Signal vs. Neutrino Background



Why go off-axis?



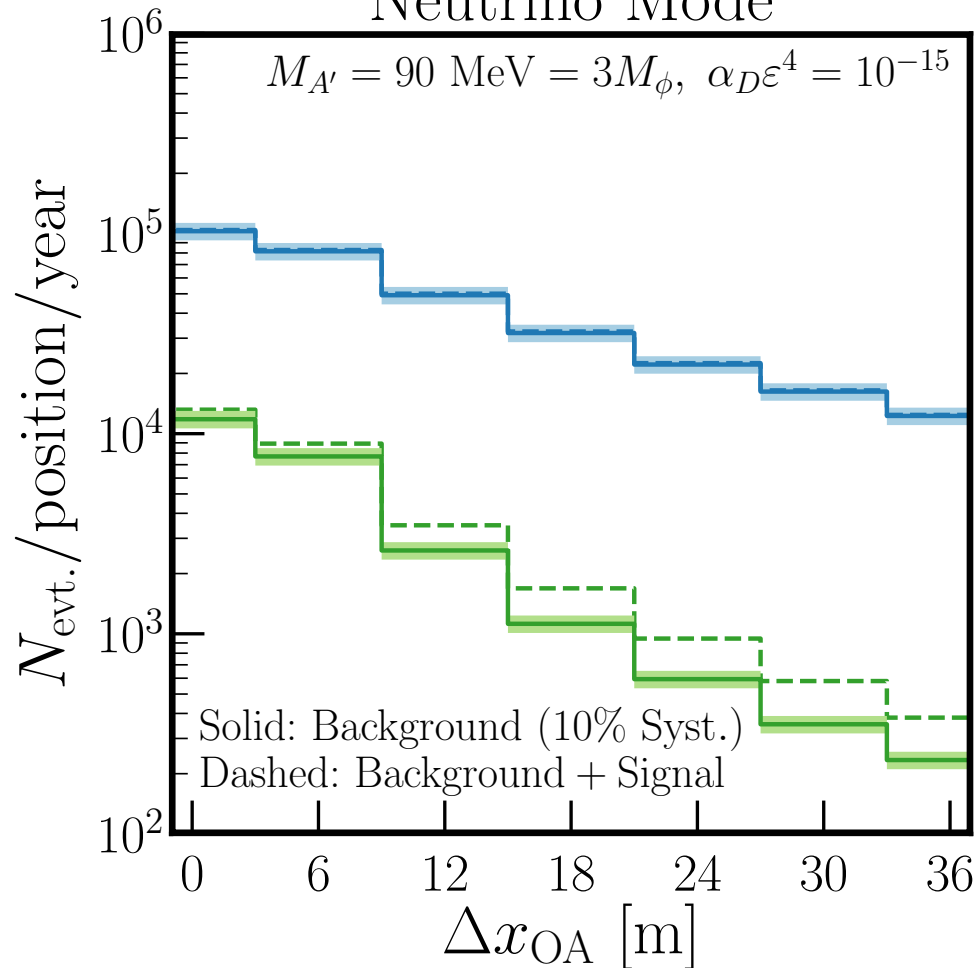
$$\frac{d^2 \Phi_v}{dE_v d(\Delta_y)} : \text{Sharp w.r.t. } \Delta_y$$

$$\frac{d^2 \Phi_\phi}{dE_\phi d(\Delta_x)} : \text{Smooth w.r.t. } \Delta_x$$

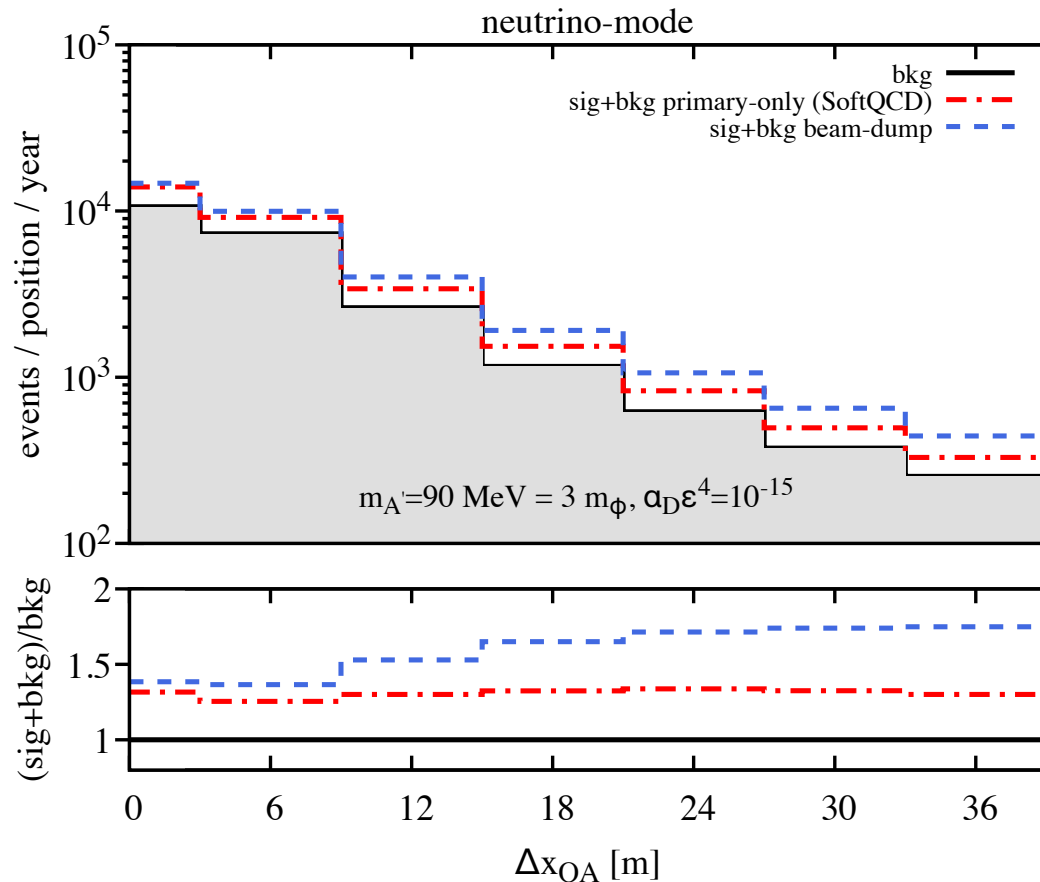
On- vs. off-axis

Neutrino Mode

$$M_{A'} = 90 \text{ MeV} = 3M_\phi, \quad \alpha_D \varepsilon^4 = 10^{-15}$$

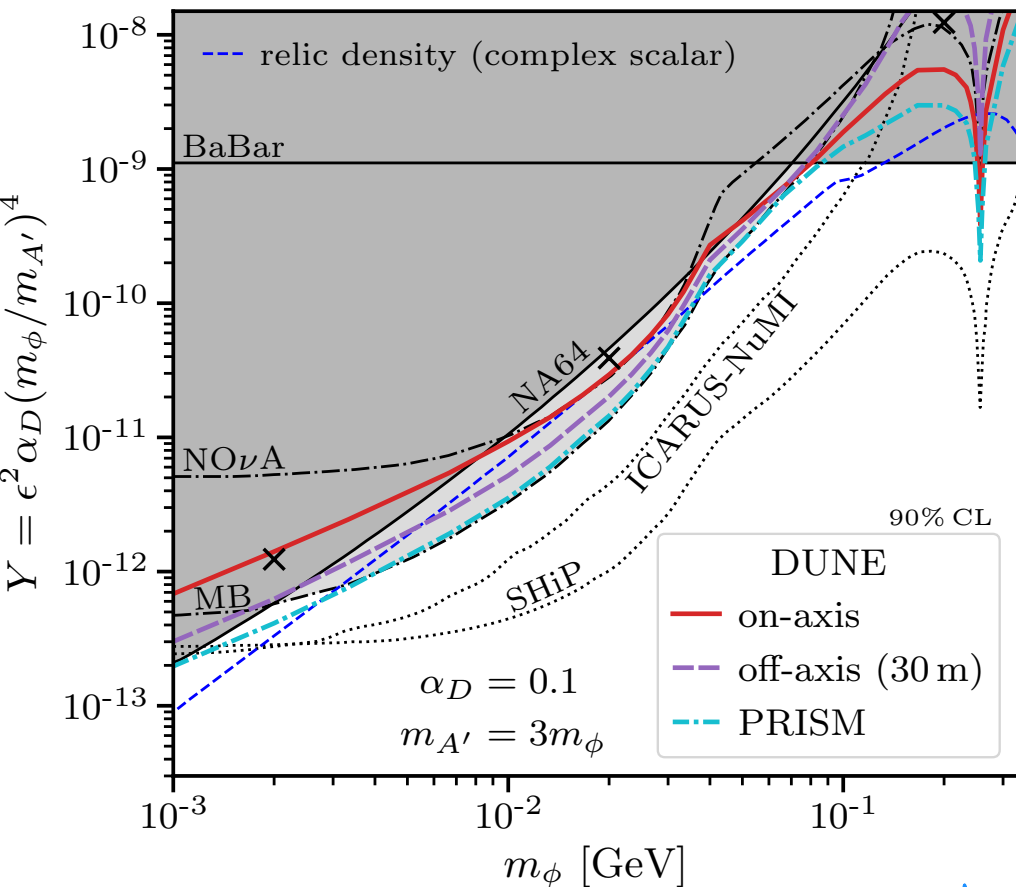


Breitbach et al [\[2102.03383\]](#)

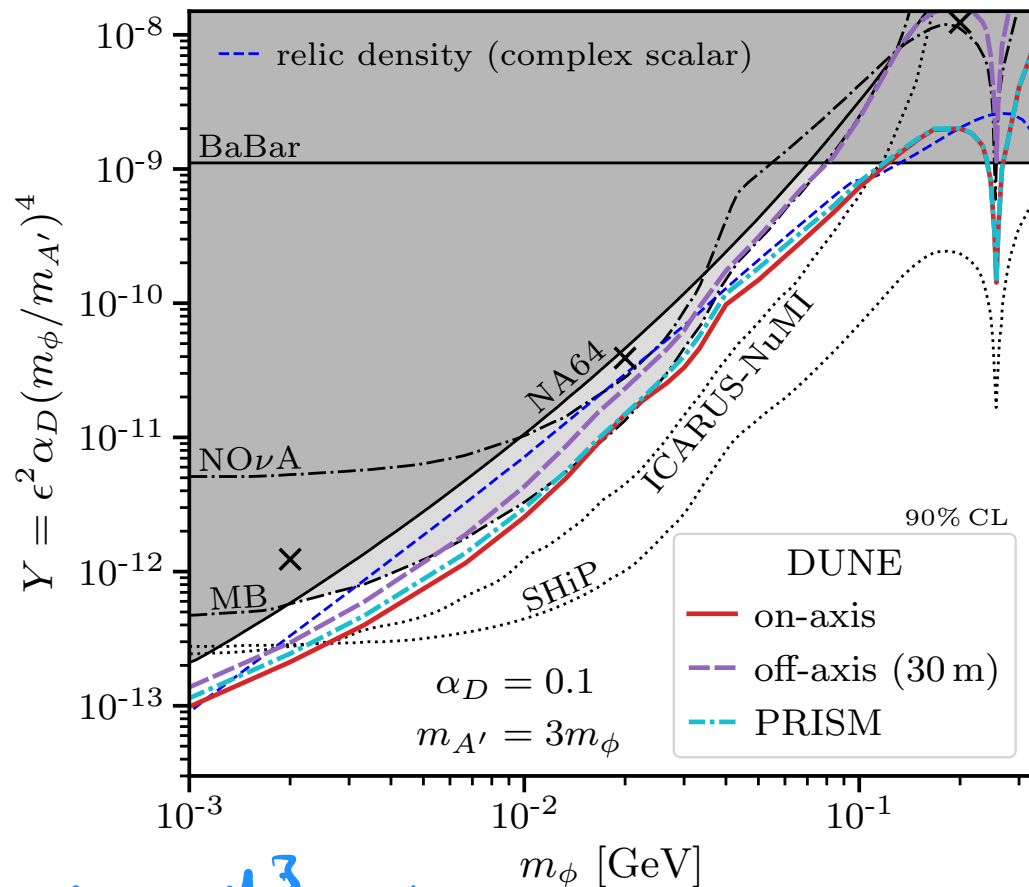


Sensitivity

total rates analysis



spectral analysis ($\Delta E = 250$ MeV)



LDMX, M^3 , etc...

Scenario 3: Challenging (but experimentally measurable) SM Background(s)

(following Coyle et al [[2210.03753](#)] for inspiration)

Sometimes, the desired physics "looks like" SM Neutrino Physics

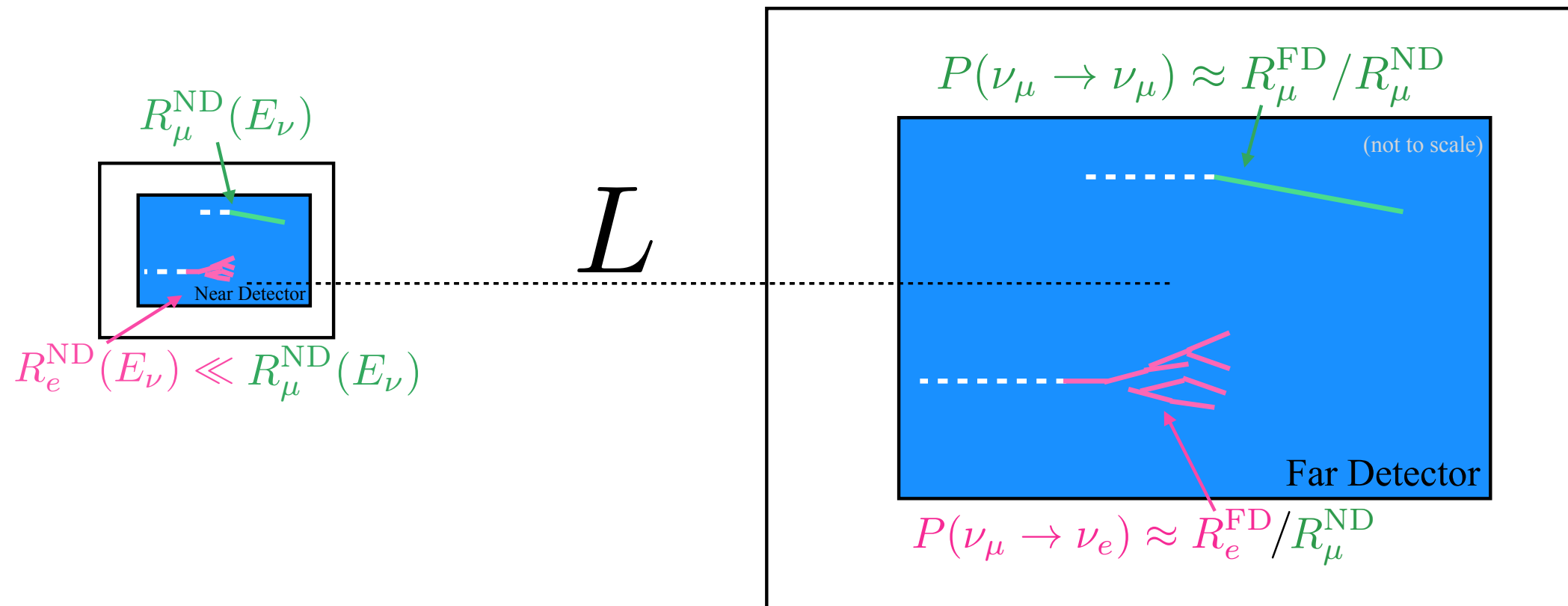
(sterile neutrinos 101)

\Rightarrow Fourth neutrino with mass $\sim m_4 \sim 1\text{eV}$

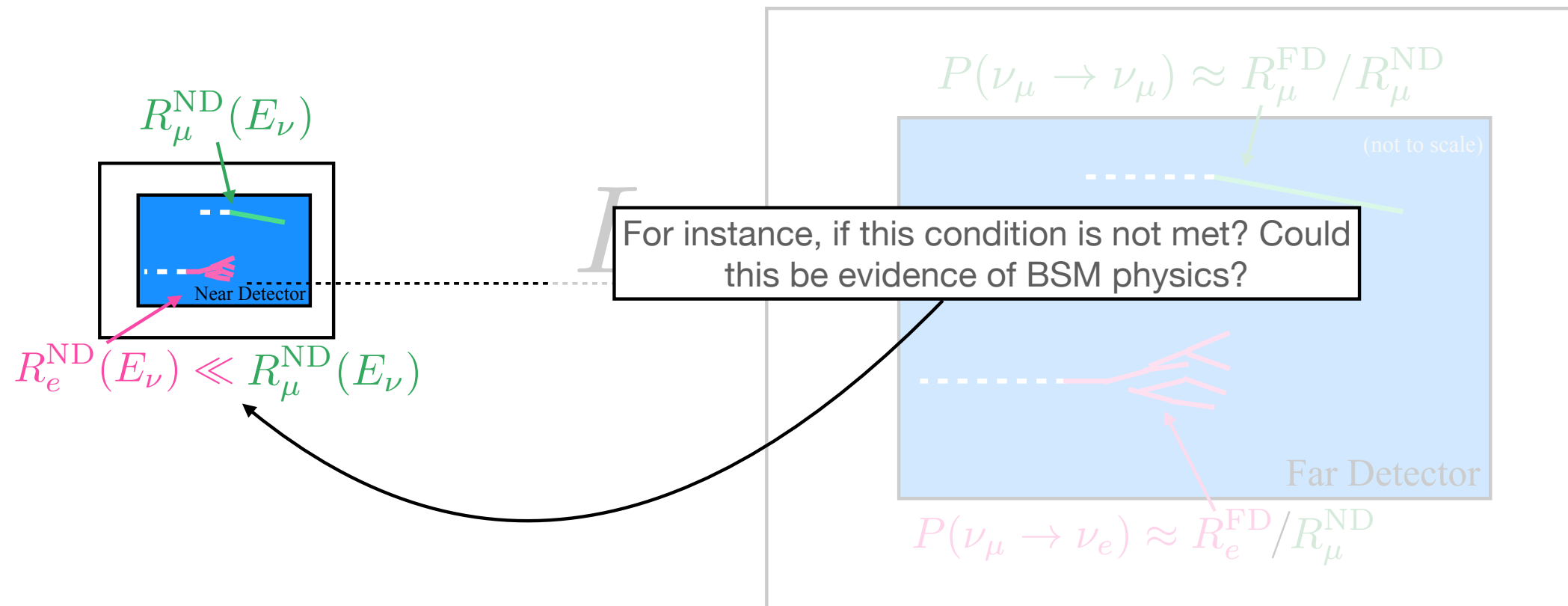
$$P_{\alpha\beta} \sim \sin^2(2\theta_{\alpha\beta}) \times \sin^2\left(\frac{\Delta m^2 L}{4E_\nu}\right)$$

\Rightarrow Energy dependence for our
Signal vs. SM expectation

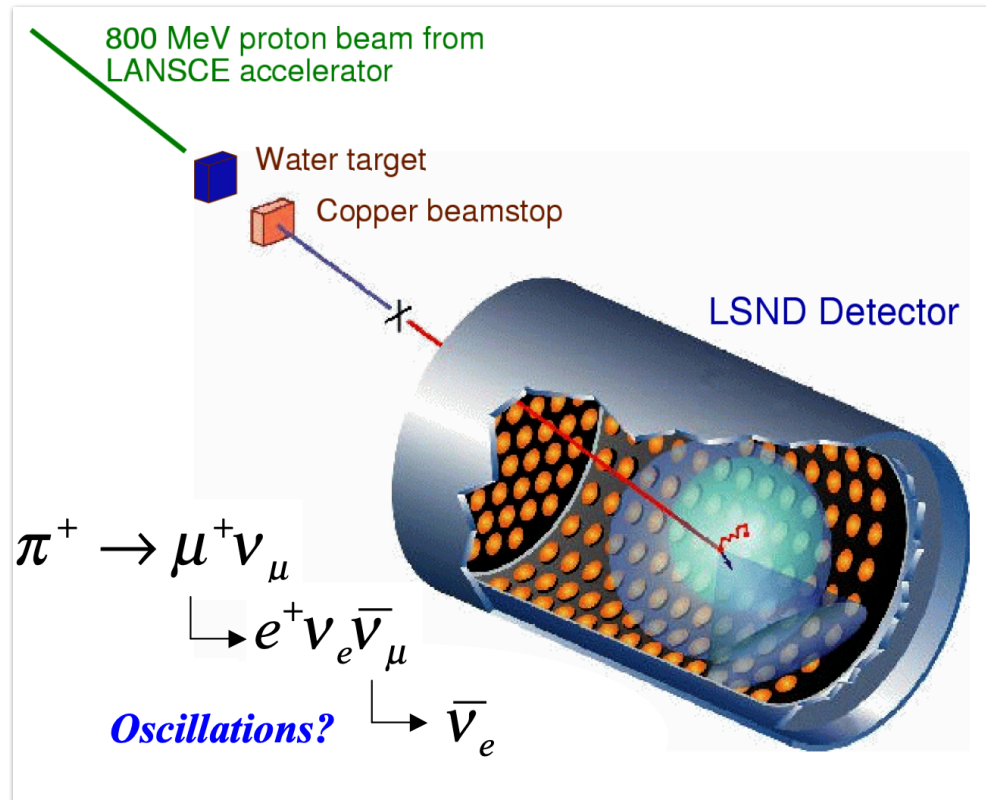
Deviations at Near Detector?



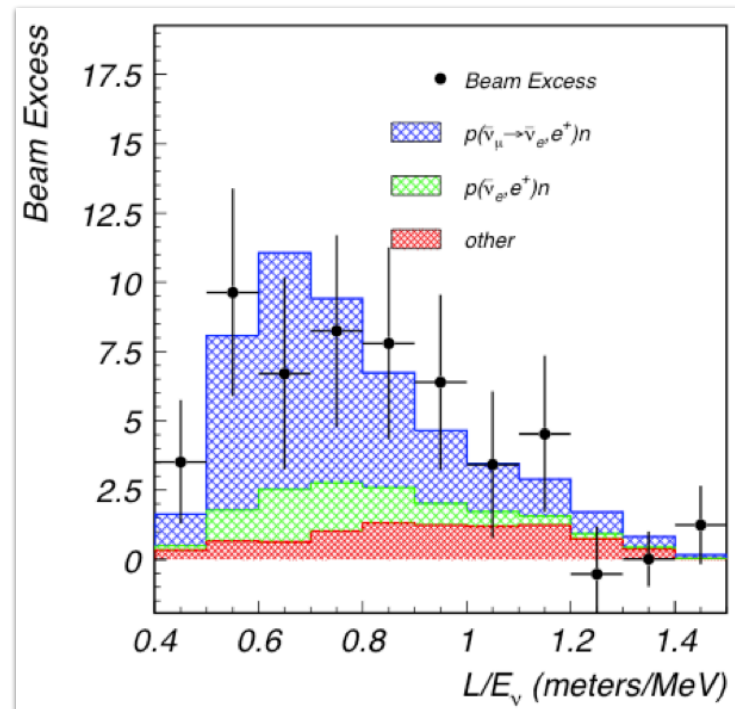
Deviations at Near Detector?



Liquid Scintillator Neutrino Detector (LSND)



$$\bar{\nu}_\mu \rightarrow \bar{\nu}_e?$$

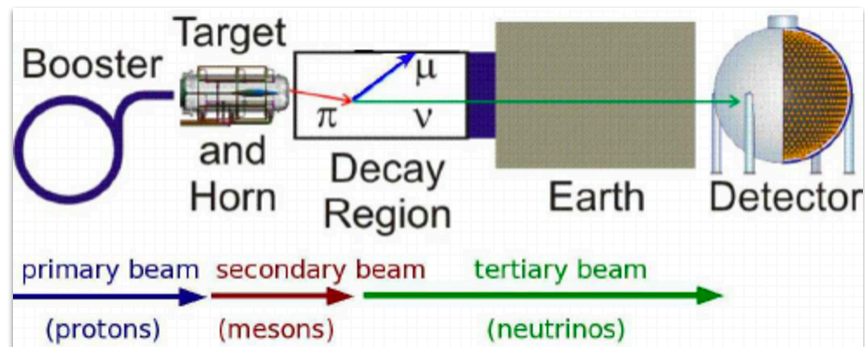


Neutrinos (mostly) from pion/muon decay-at-rest — O(30) MeV, roughly 50 meter baseline length.

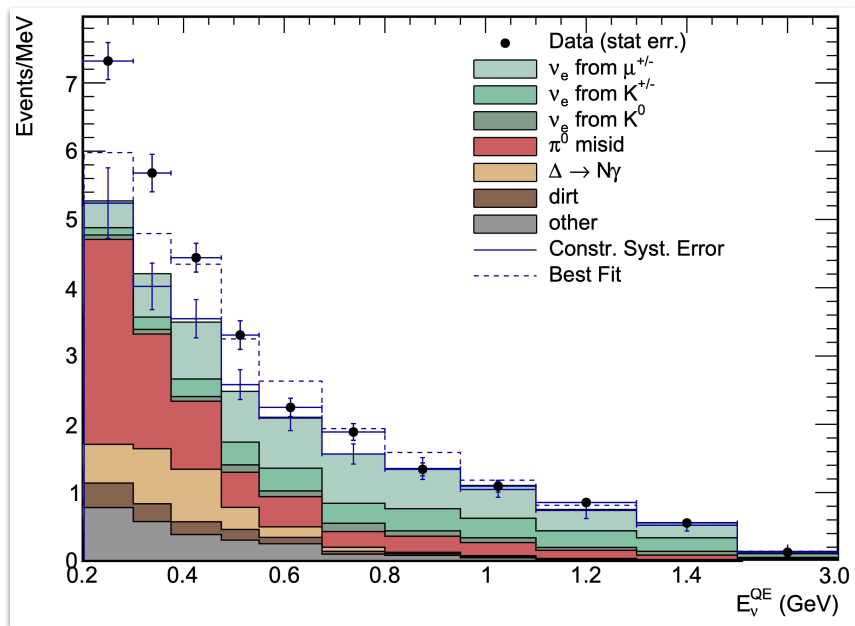
Observed excess — $87.9 \pm 22.4 \pm 6.0 \rightarrow P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \approx 2.6 \times 10^{-3}$

MiniBooNE

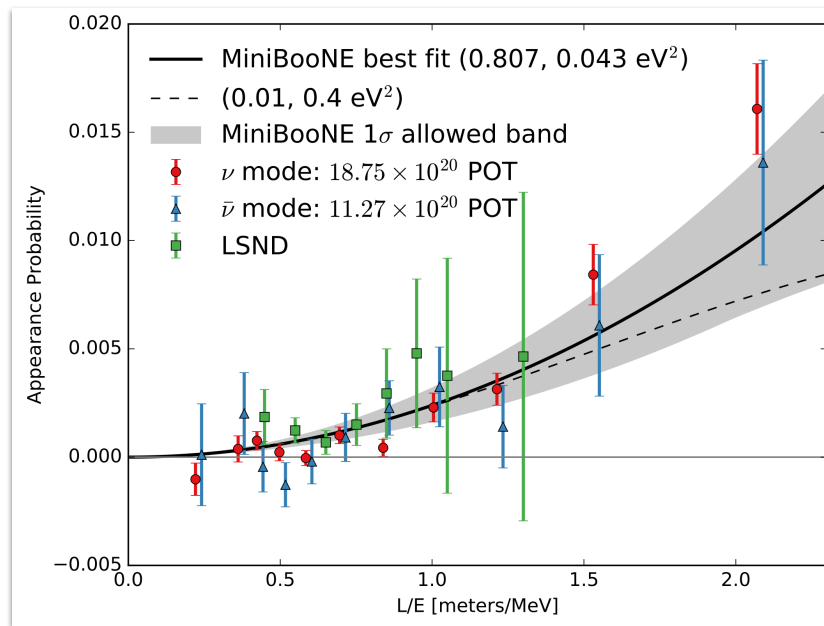
Designed to test the LSND anomaly — very different L, E, but similar L/E



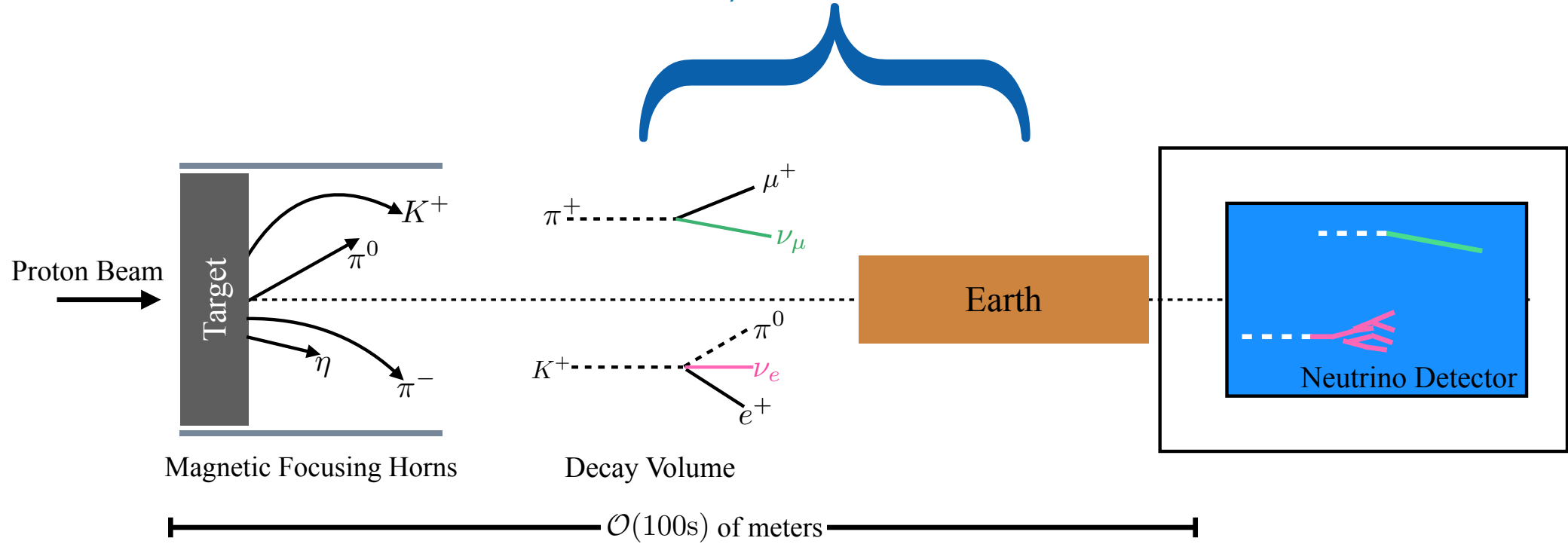
$$\nu_{\mu} \rightarrow \nu_e \text{ AND } \bar{\nu}_{\mu} \rightarrow \bar{\nu}_e?$$



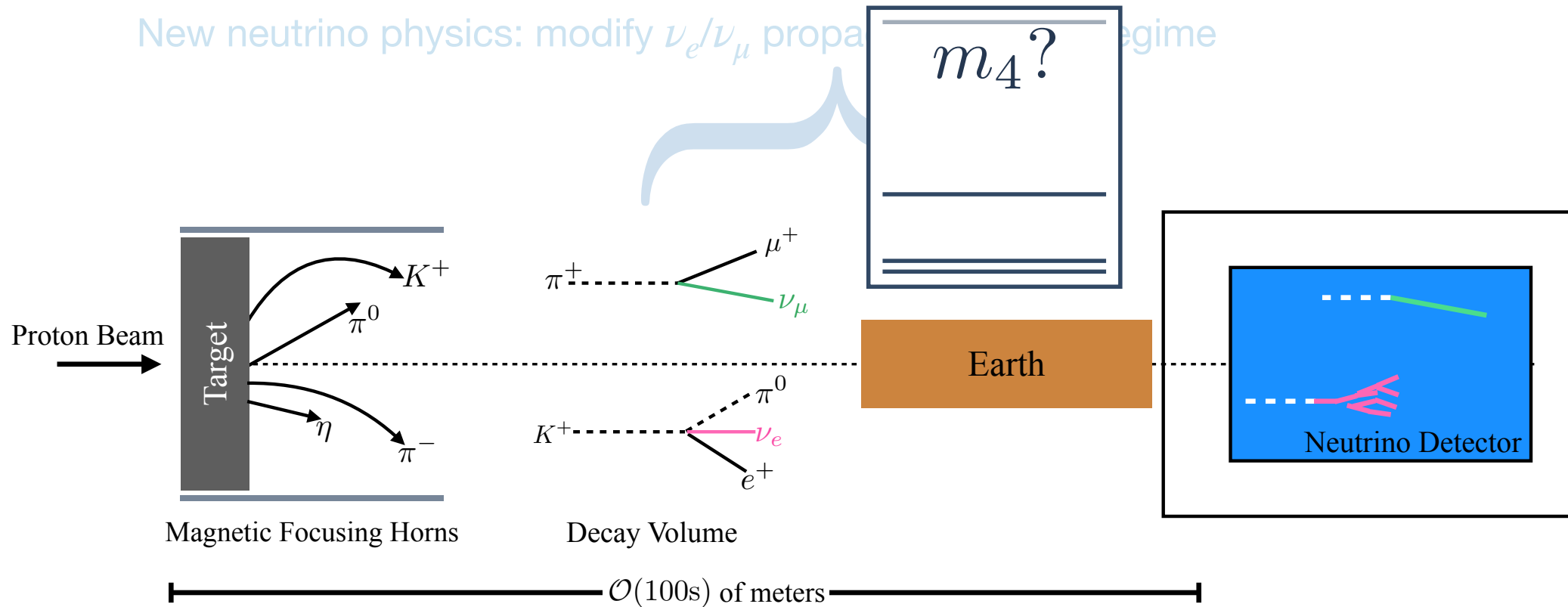
MiniBooNE Collab., [2006.16883]



New neutrino physics: modify ν_e/ν_μ propagation in this regime

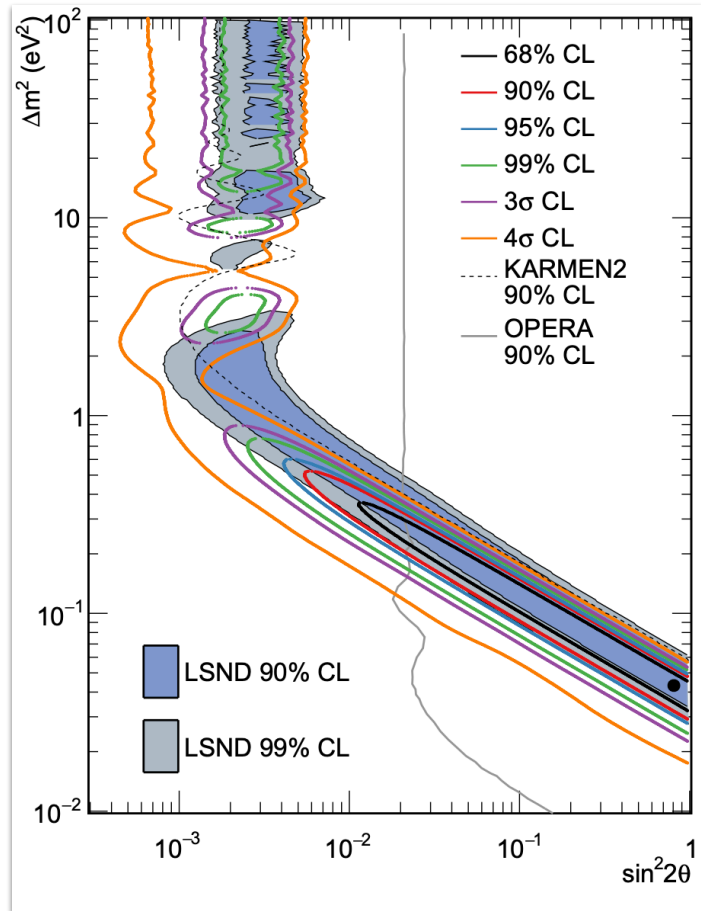


New neutrino physics: modify ν_e/ν_μ propagation regime



$$P_{\alpha\beta} \approx \sin^2(2\theta_{\alpha\beta}) \sin^2\left(\frac{\Delta m^2 L}{4E_\nu}\right) \quad \text{New mass scale } \Delta m^2 \gg \Delta m_{\text{SM}}^2$$

Fourth-neutrino interpretation



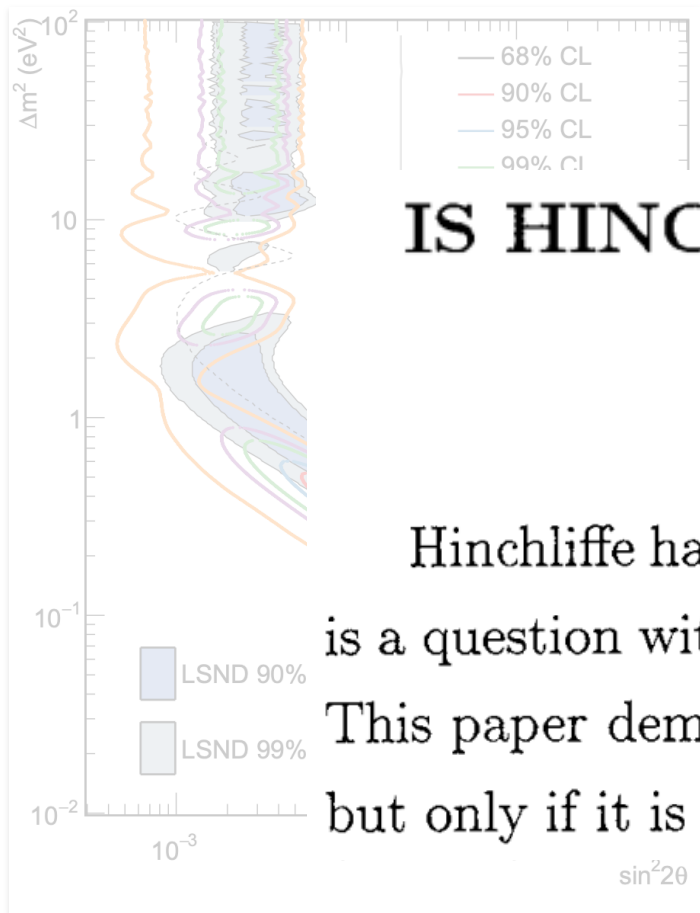
If coming from oscillations, the results from LSND and MiniBooNE require a new mass eigenstate around the eV scale.

Combined with the observed invisible width of the Z-boson (LEP), any additional light neutrino(s) must be sterile — gauge singlets.

Is this 3+1 scenario compatible with global data?

MiniBooNE Collab., [\[2006.16883\]](#)

Fourth-neutrino interpretation



IS HINCHLIFFE'S RULE TRUE? •

Boris Peon

Abstract

Hinchliffe has asserted that whenever the title of a paper is a question with a yes/no answer, the answer is always no. This paper demonstrates that Hinchliffe's assertion is false, but only if it is true.

ND and
und the eV

the Z-boson
sterile — gauge

gal data?

[\(in loving memory\)](#)

MiniBooNE Collab., [\[2006.16883\]](#)

Consequences of Invoking a light (sterile) Neutrino

$$P_{\alpha\beta} \sim \sin^2(2\theta_{\alpha\beta}) \times \sin^2\left(\frac{\Delta m^2 L}{4E\nu}\right)$$

$\hookrightarrow \alpha = \mu$
 $\hookrightarrow \beta = e$

$$\sin^2(2\theta_{\mu e}) = \frac{4 |U_{\mu 4}|^2 |U_{e 4}|^2}{2 \sin^2(\theta_{24}) \times \cos^2(\theta_{14})} \rightarrow \sin^2(\theta_{14})$$

$$P_{\alpha\alpha} \sim 1 - \sin^2(2\theta_{\alpha\alpha}) \sin^2\left(\frac{\Delta m^2 L}{4E\nu}\right)$$

$$\hookrightarrow \sin^2(2\theta_{\alpha\alpha}) = 4 |U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2)$$

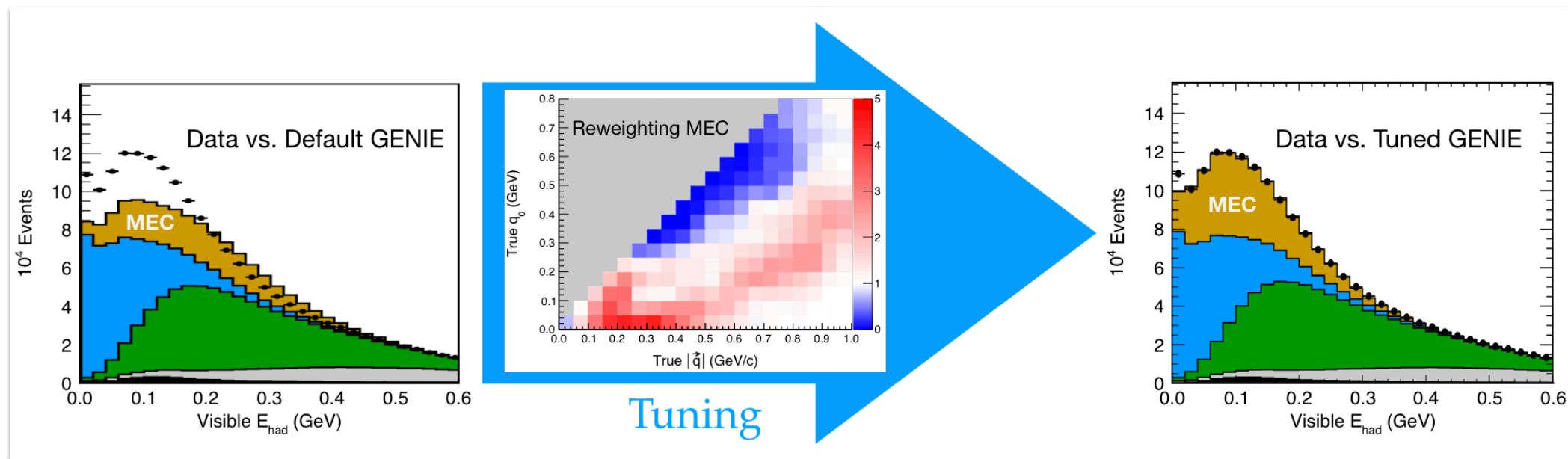
Impact at, e.g., NOvA ND

$$P_{\mu\mu} \approx 1 - \sin^2(2\theta_{\mu\mu}) \sin^2\left(\frac{\Delta m^2 L}{4E_\nu}\right)$$

$$\frac{\Delta m^2 L}{4E_\nu} \approx 4 \times \left(\frac{\Delta m^2}{\text{eV}^2}\right) \times \left(\frac{L}{\text{m}}\right) \times \left(\frac{\text{GeV}}{E_\nu}\right)$$

\Rightarrow Nontrivial E_ν dependence on
year ν_μ Event Rate

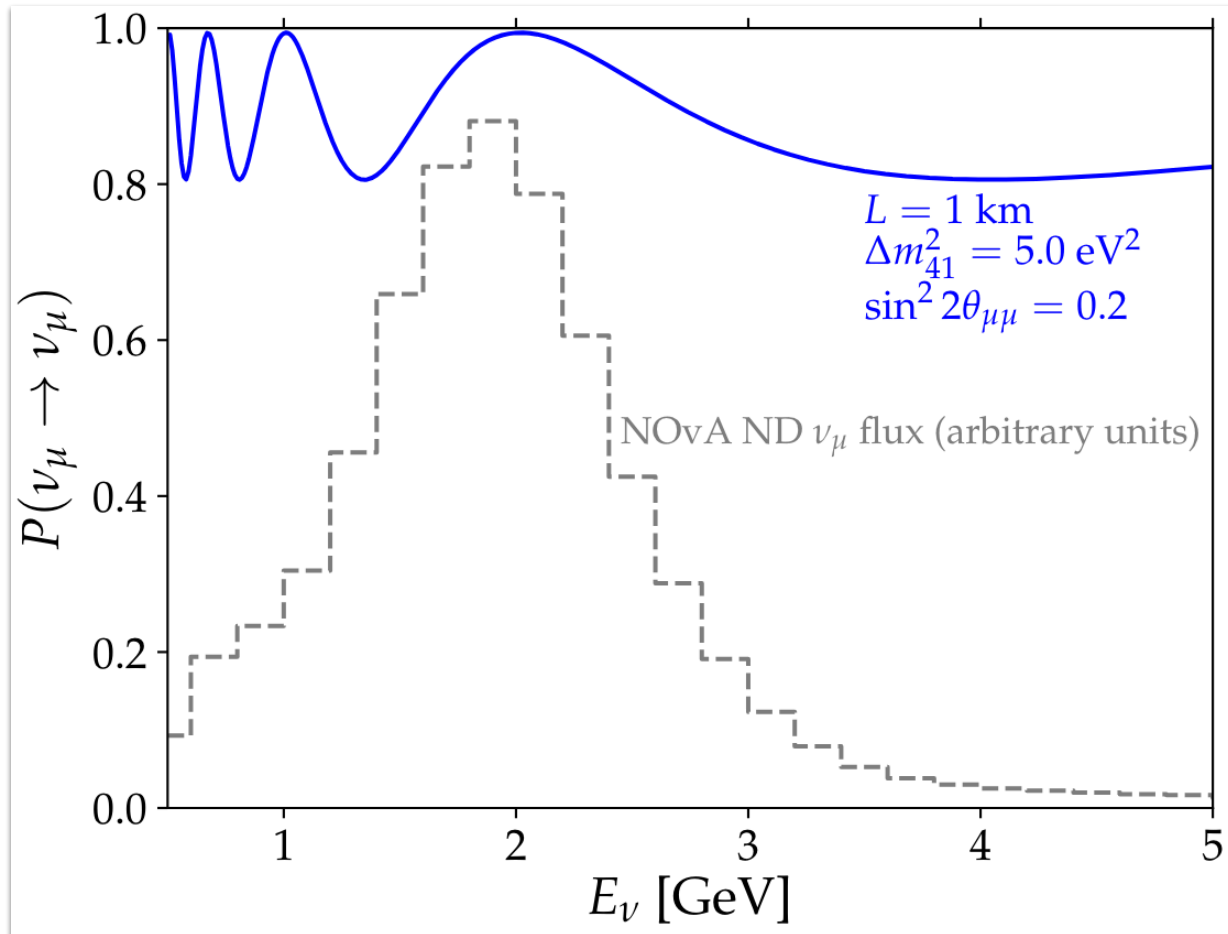
What do Experiments do with ND Data?



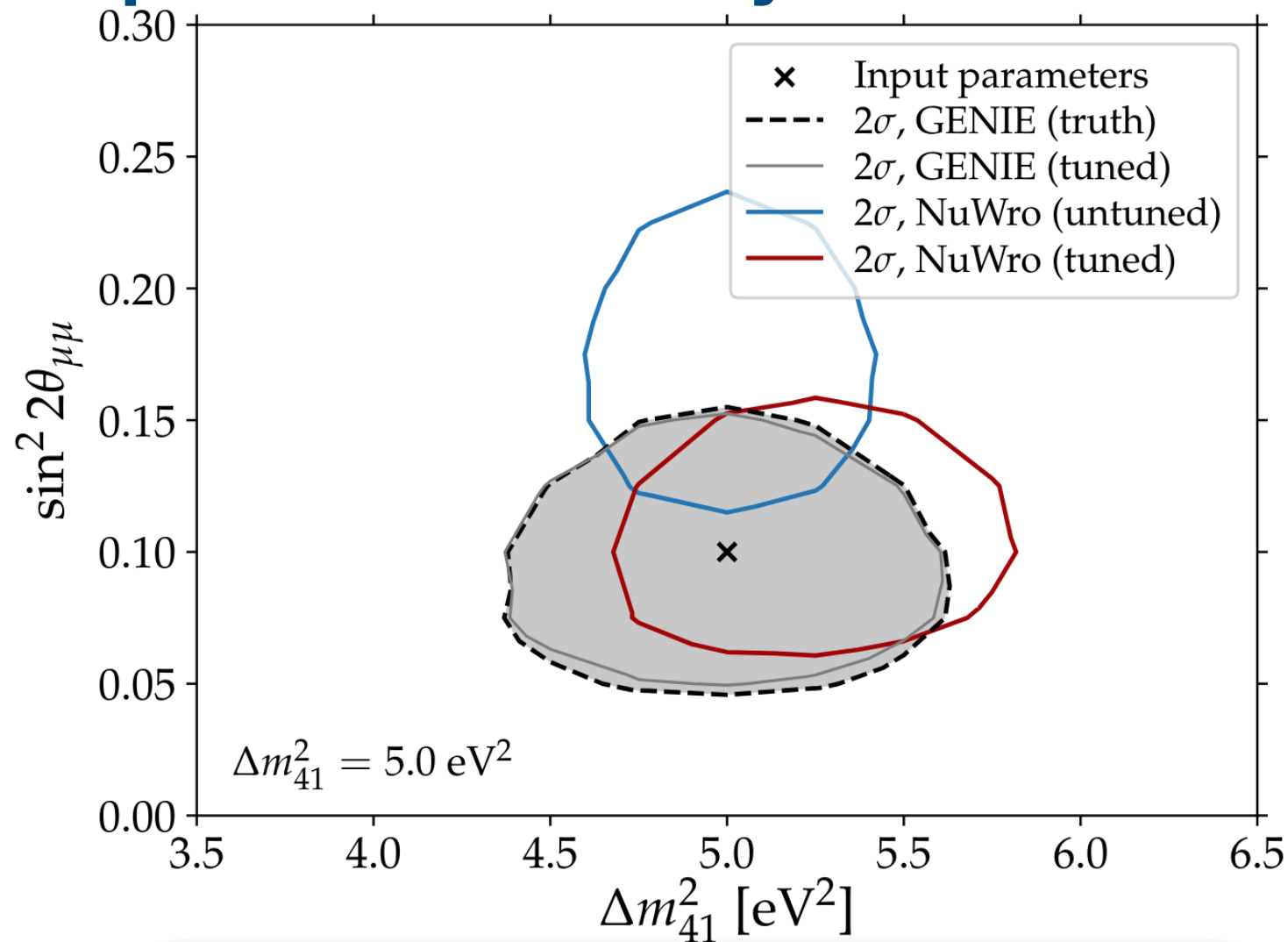
Coyle et al [\[2210.03753\]](#), adapted from NOvA [\[2006.08727\]](#)

What if there is underlying new physics in this data?

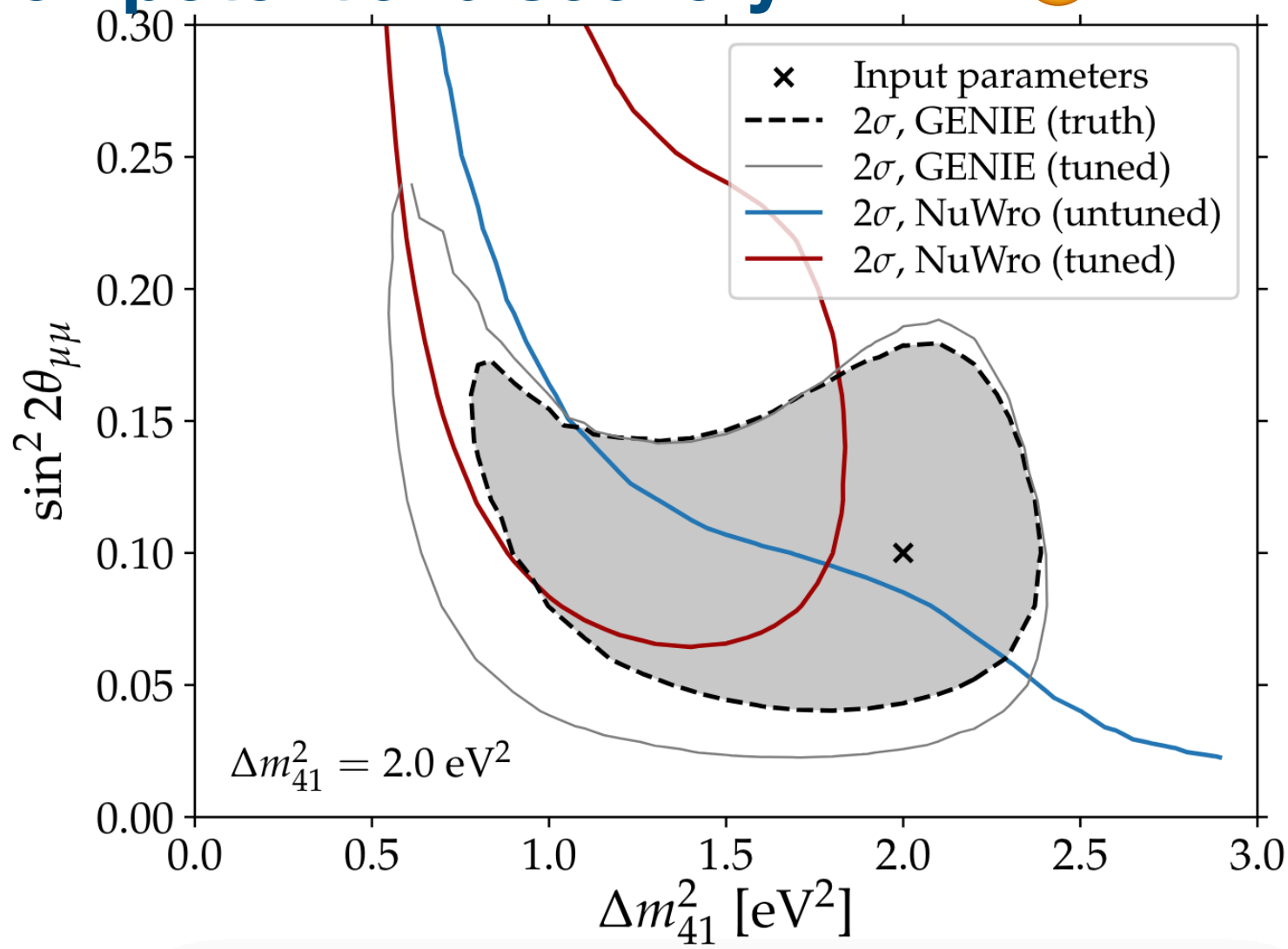
Coyle et al [\[2210.03753\]](#)



Impact on potential discovery? 🤔



Impact on potential discovery? v2... 🤪



A second ND "background-swamped" model

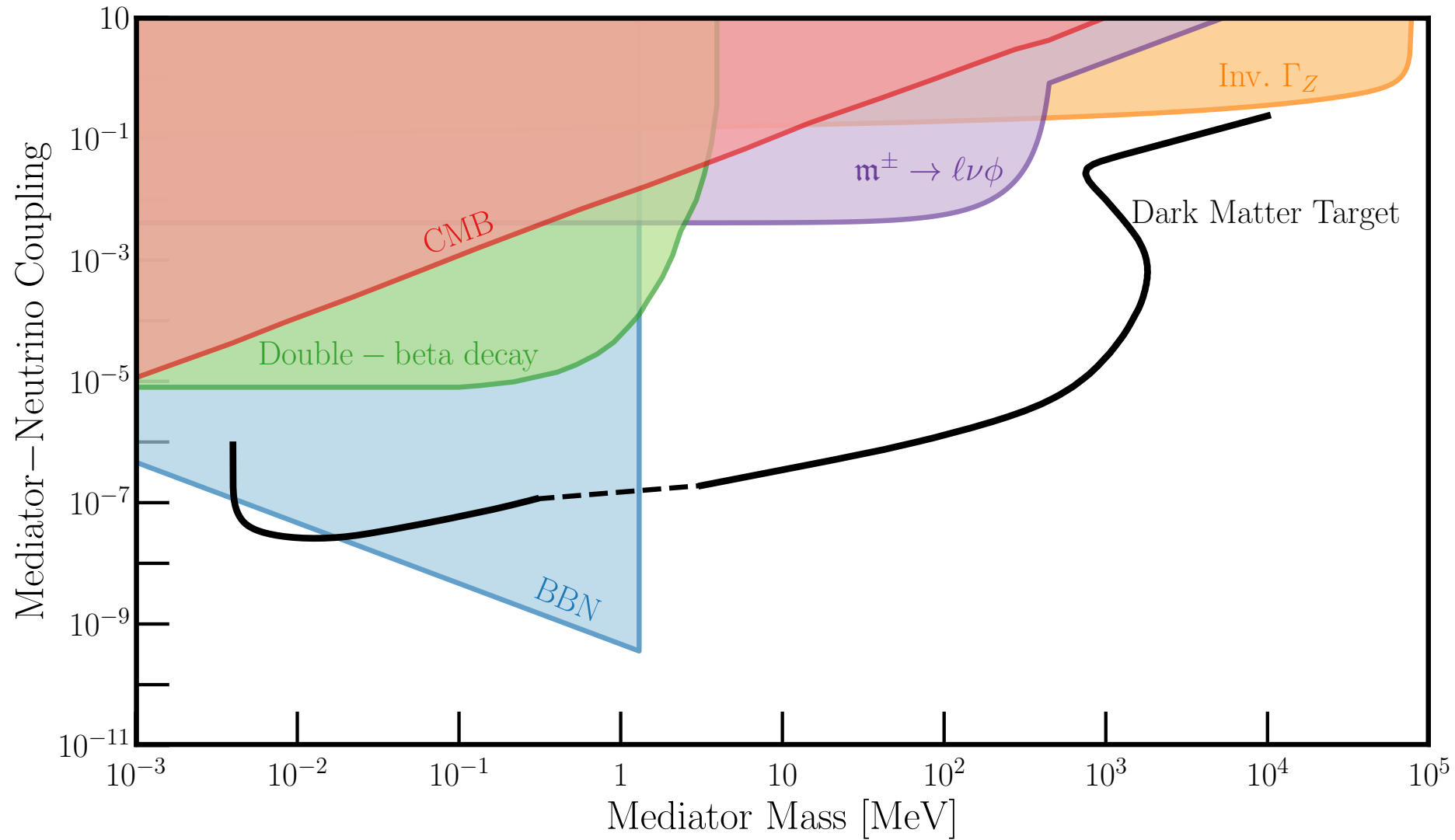
New particle ϕ scalar "neutrophilic"

$$\mathcal{L}_\alpha \frac{(L_\alpha H)(L_\beta H)\phi}{\Lambda_{\alpha\beta}^2}$$

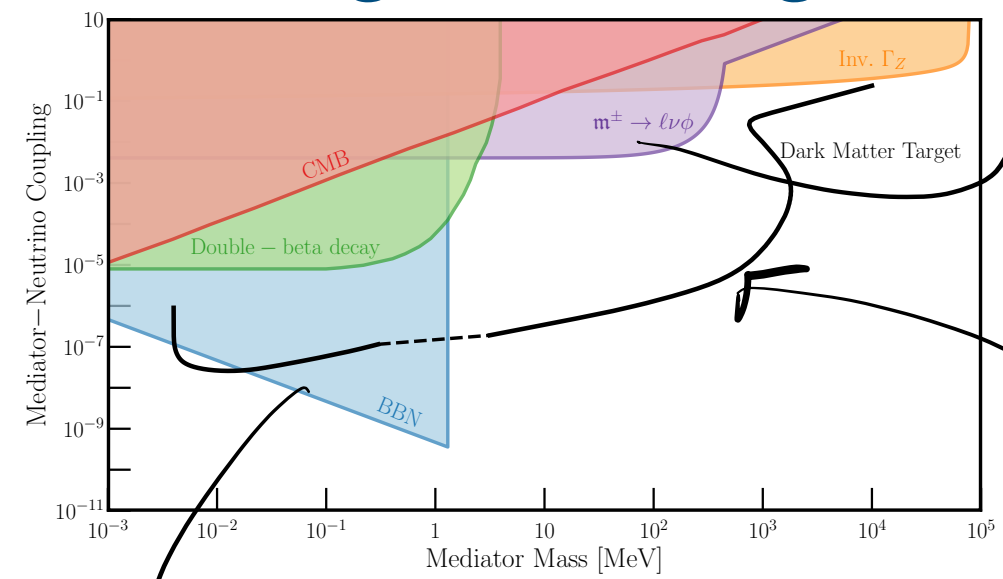
$$\Rightarrow \lambda_{\alpha\beta} \phi \nu_\alpha \nu_\beta$$

$$\phi \sim \text{MeV} - \text{GeV}$$

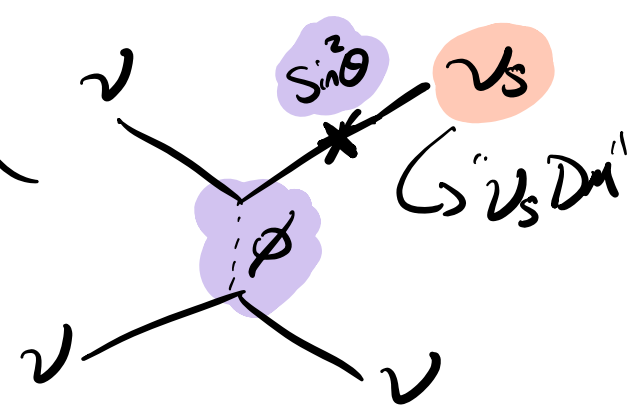
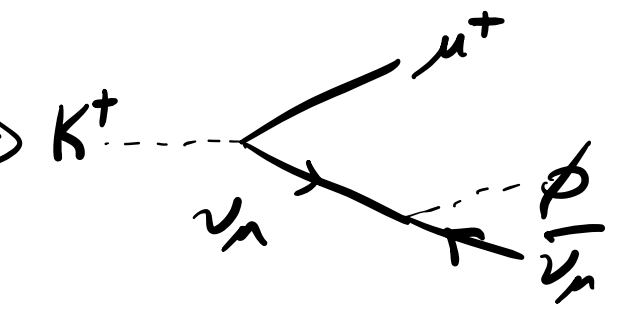
Broader motivation?



Breaking down regions...

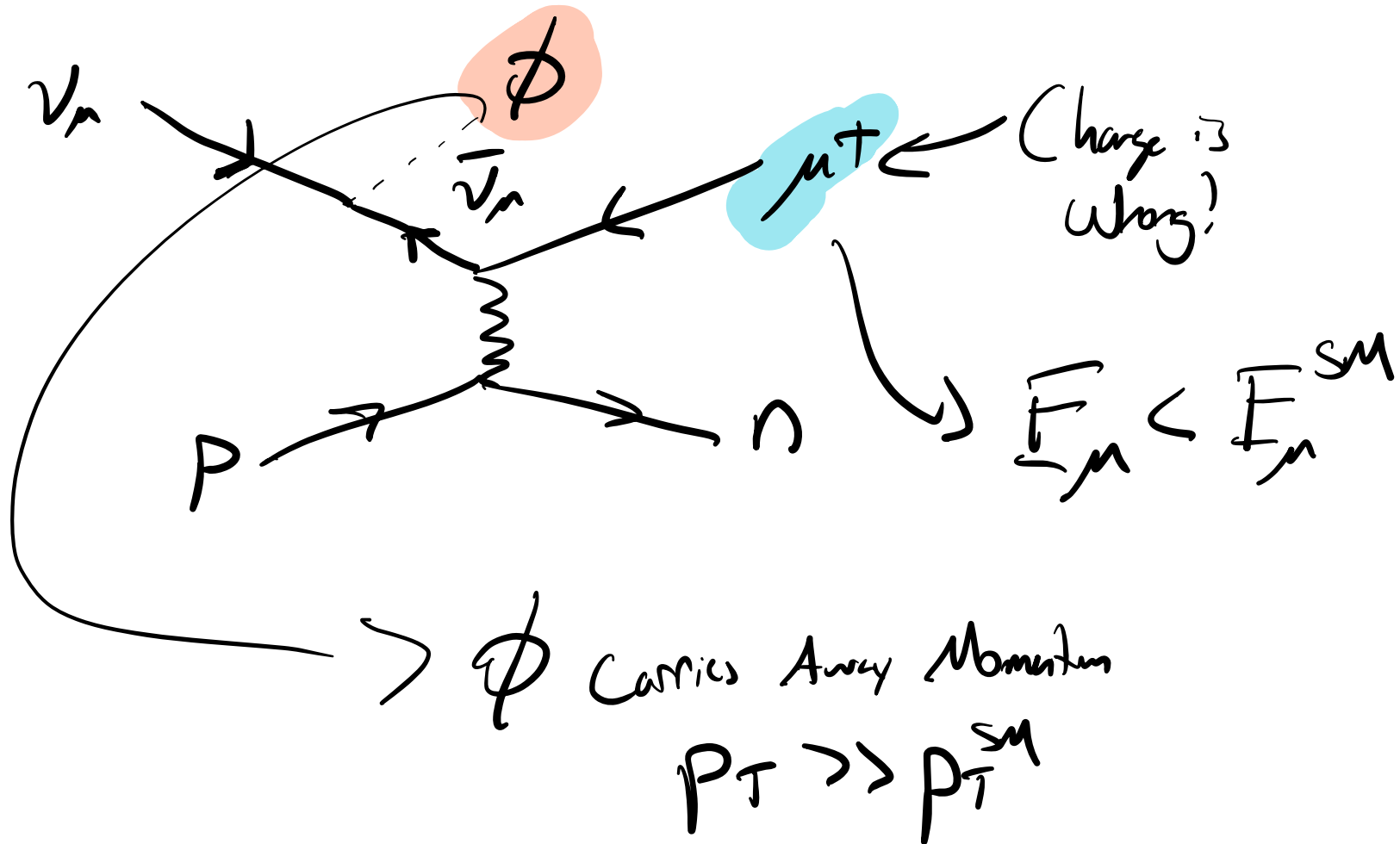


→ Thermalize in Early Universe

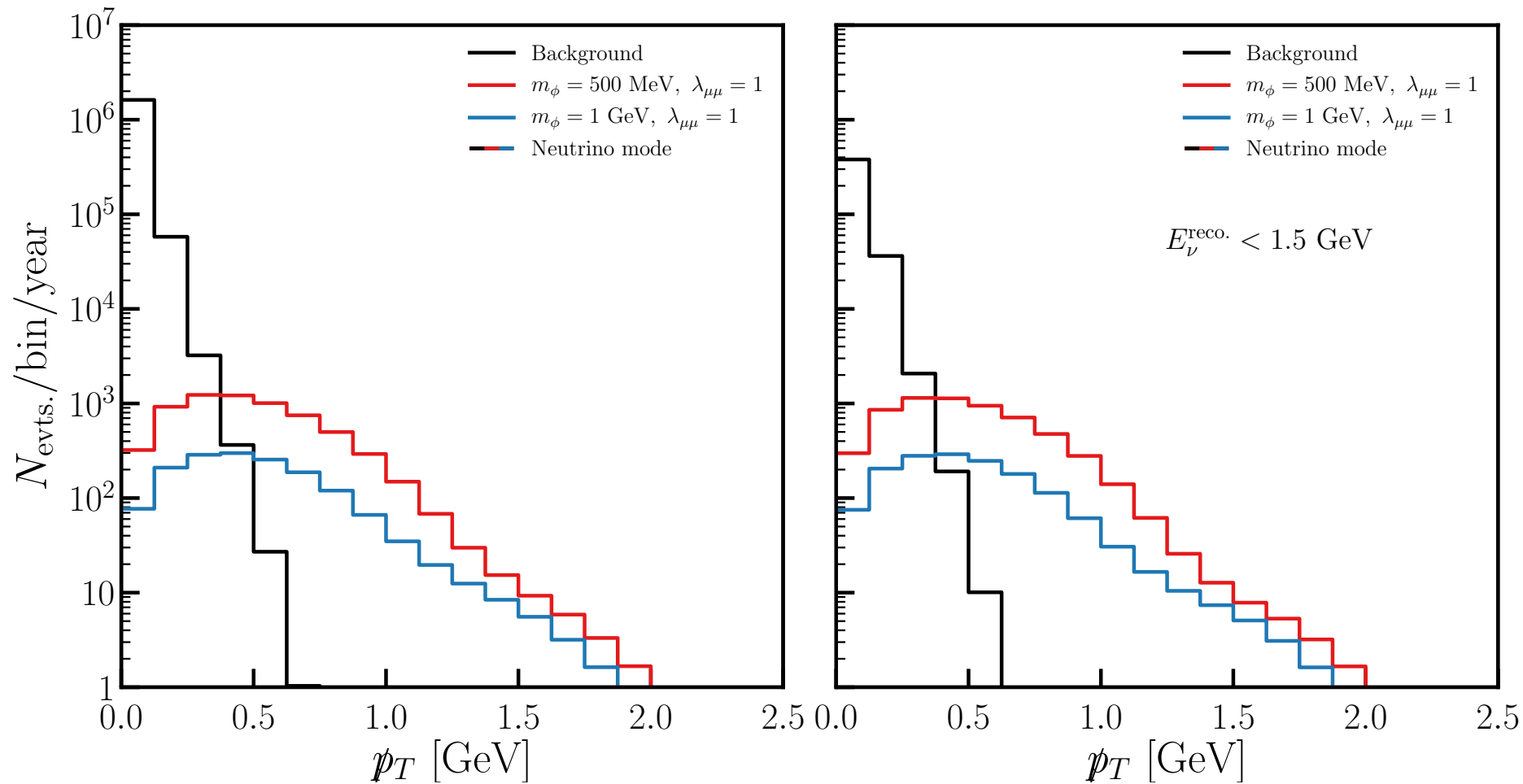


⇒ Hubble Tension!

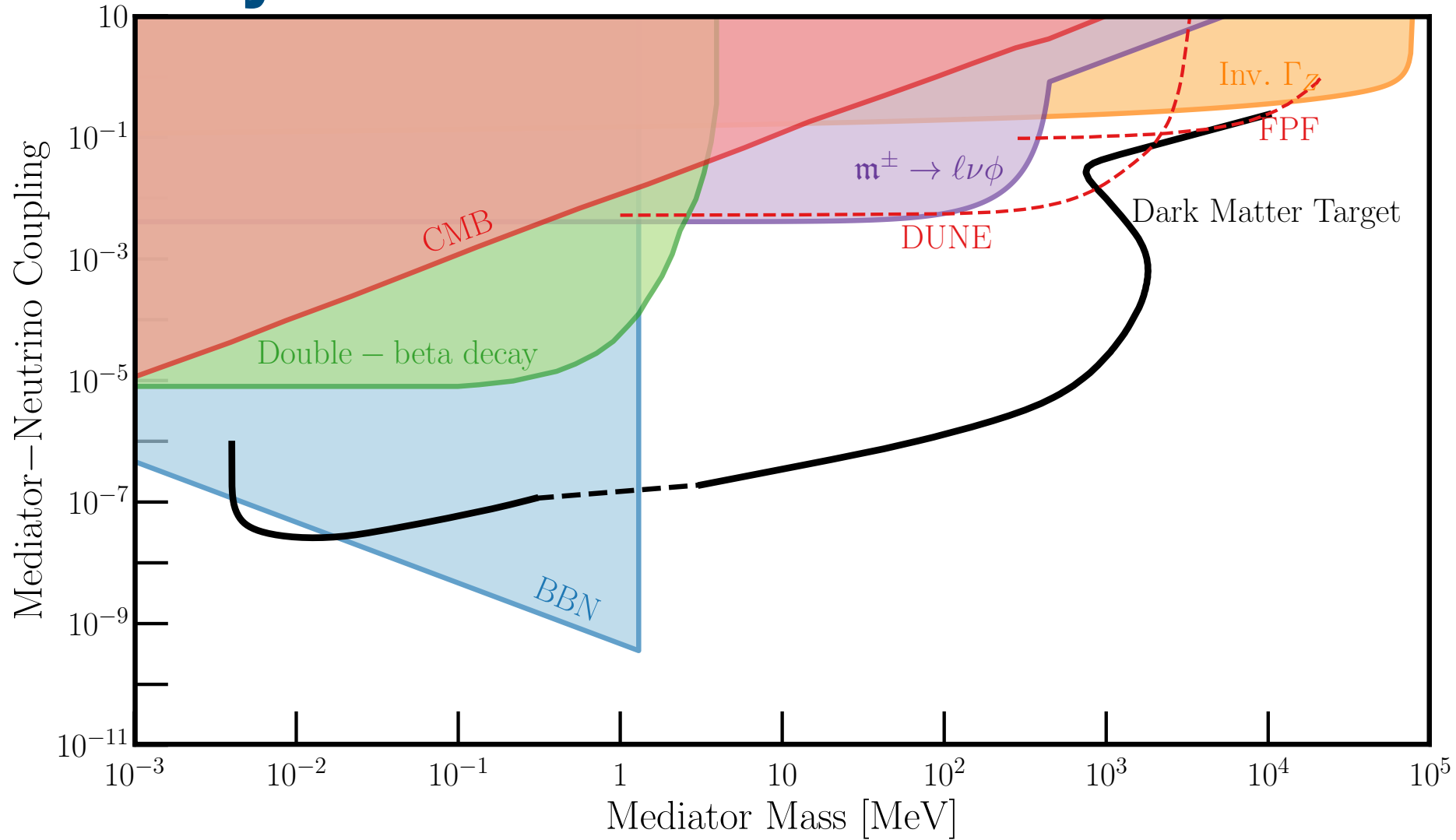
Neutrinophilic Scalars in Neutrino Facilities



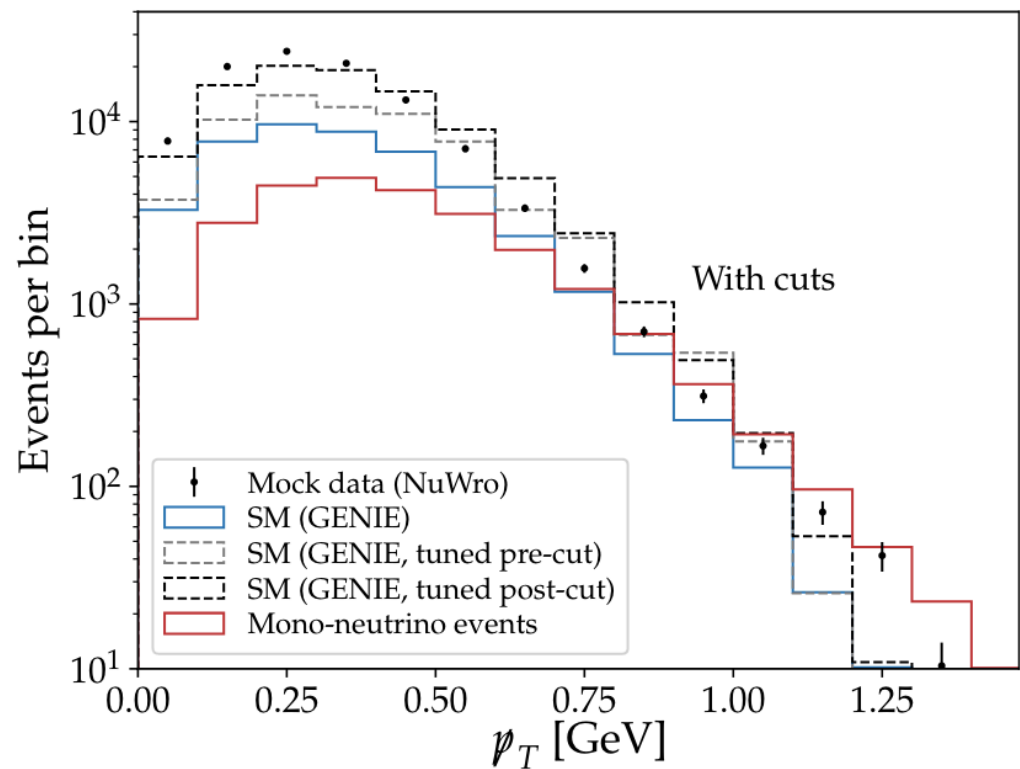
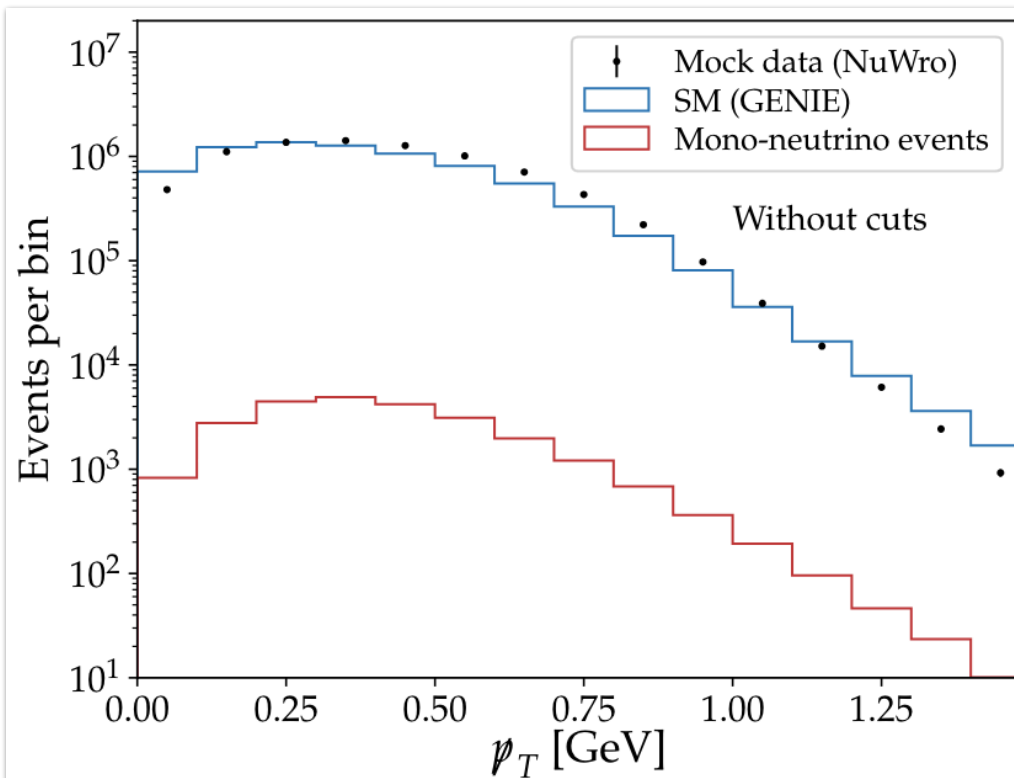
Signal vs. Background Kinematics



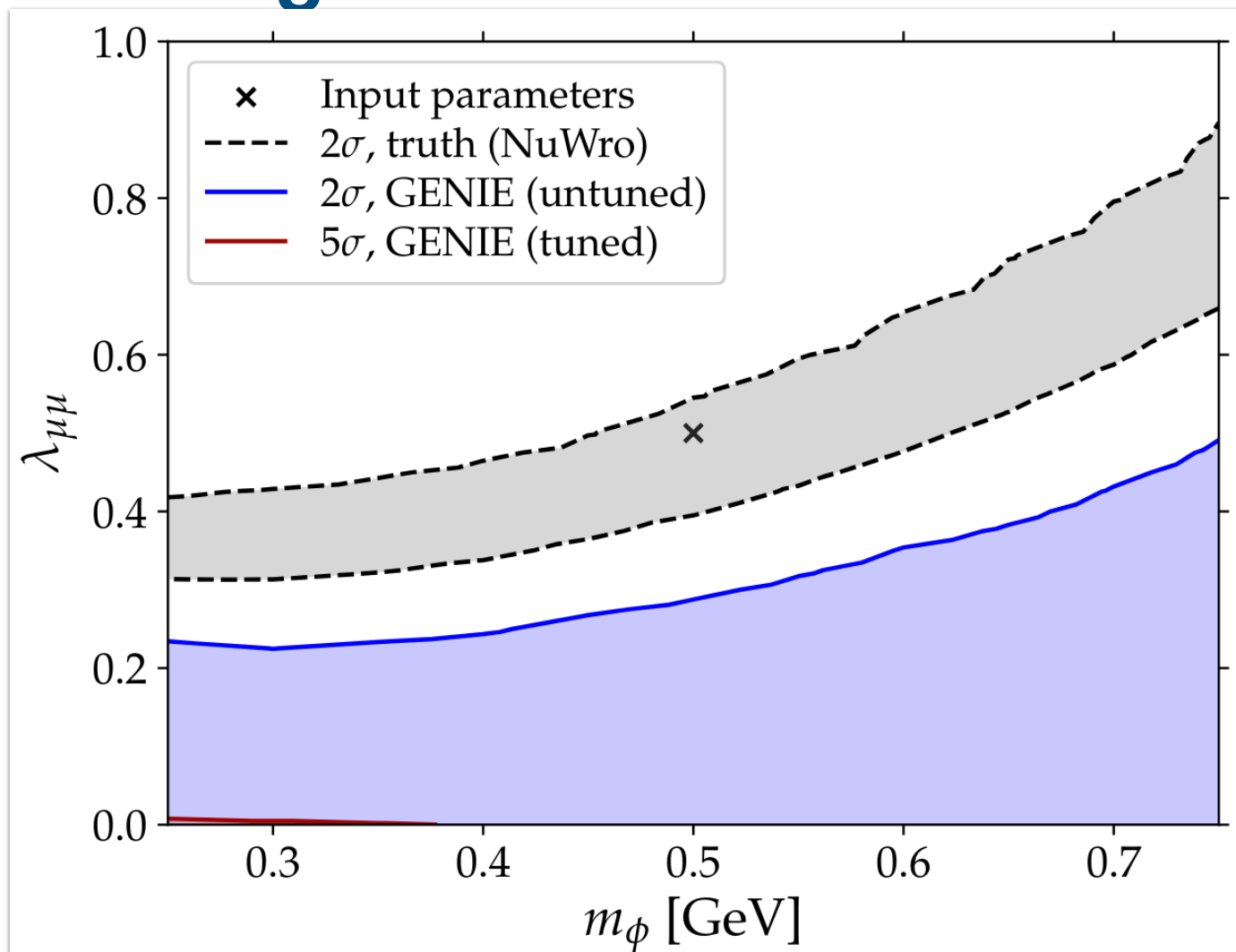
Sensitivity Estimates



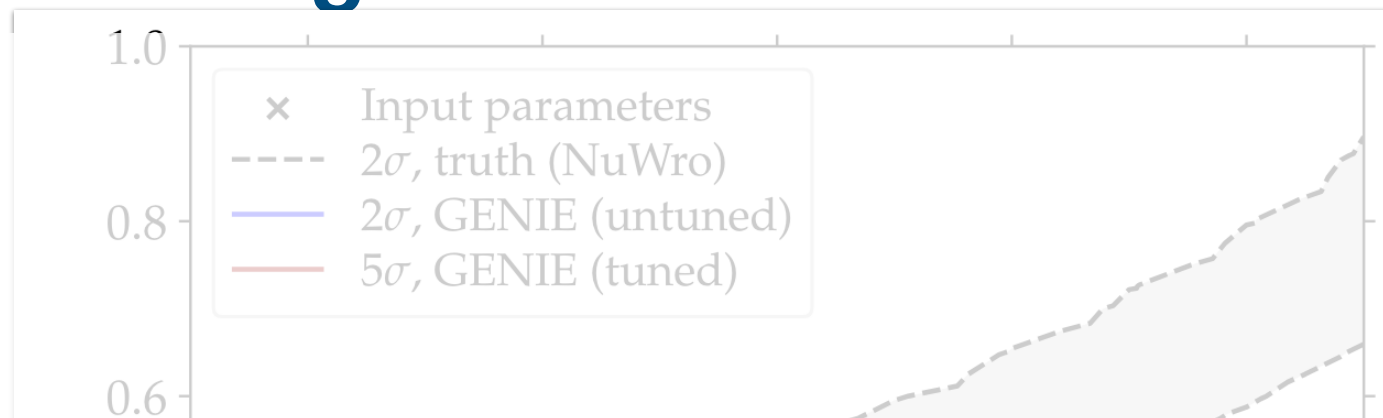
Re-enter Coyle/Li/Machado...



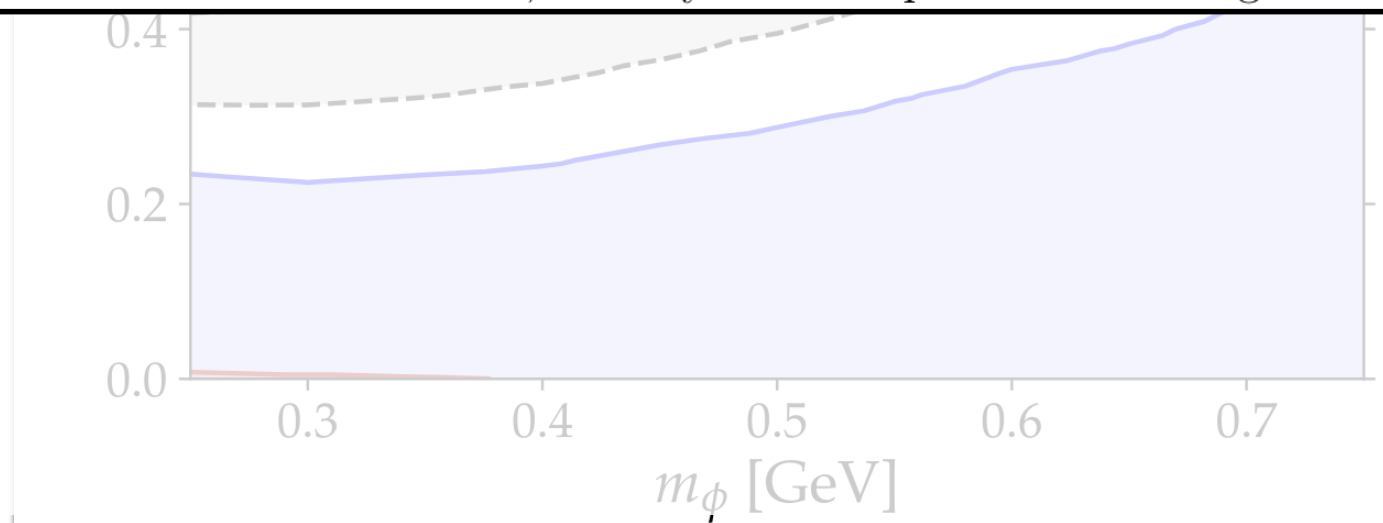
Impact of tuning on these searches



Impact of tuning on these searches



This shows that while near detectors have the potential to probe new physics, without proper modeling of neutrino-nucleus interactions, we may lose this potential. Tuning is not the solution.



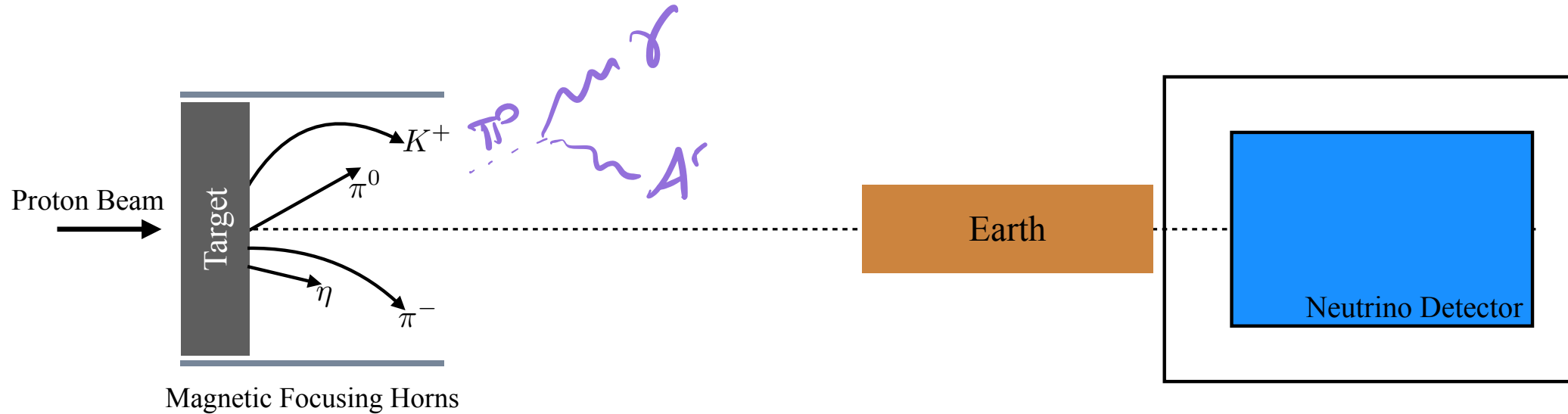
Backup

Novel BSM Production Mechanisms in Thick Targets — Electromagnetic Production

 <https://github.com/kjkellyphys/PETITE>

“Nonstandard” BSM Production?

What about new states produced in the target?



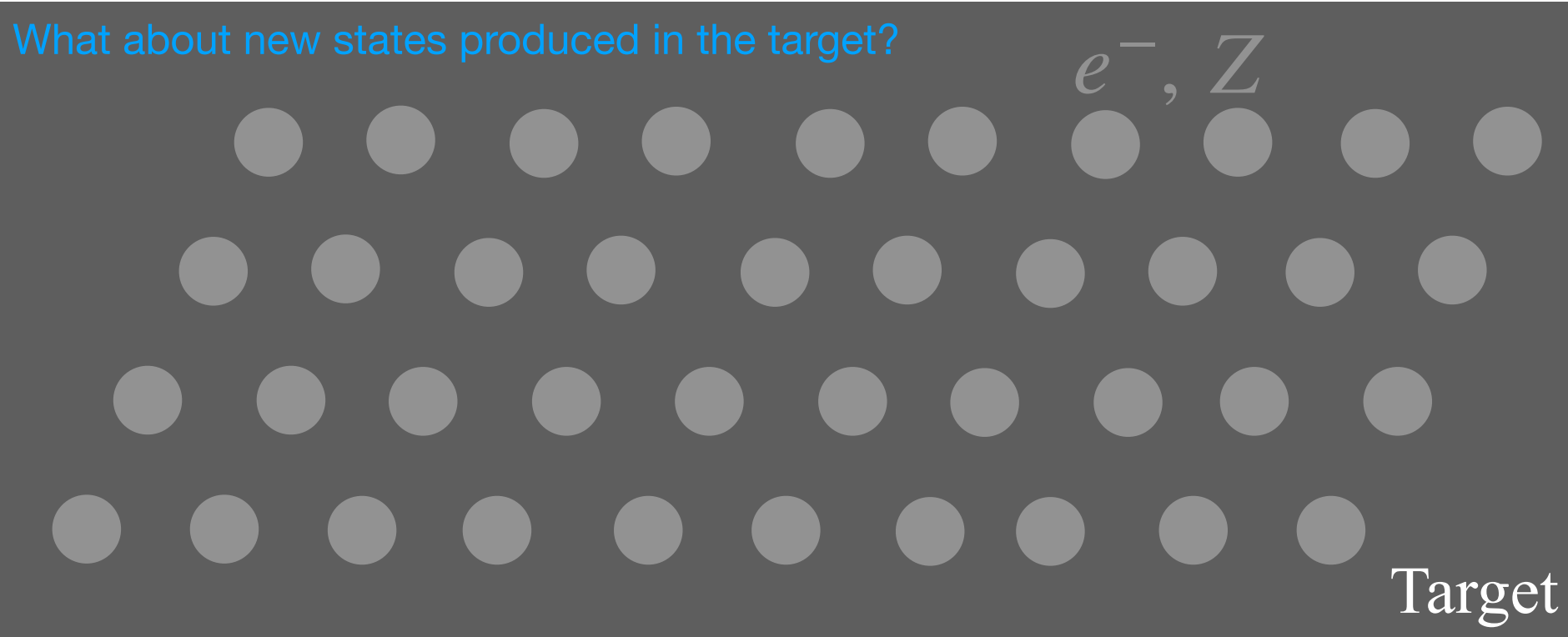
“Nonstandard” BSM Production?

What about new states produced in the target?

e^- , Z

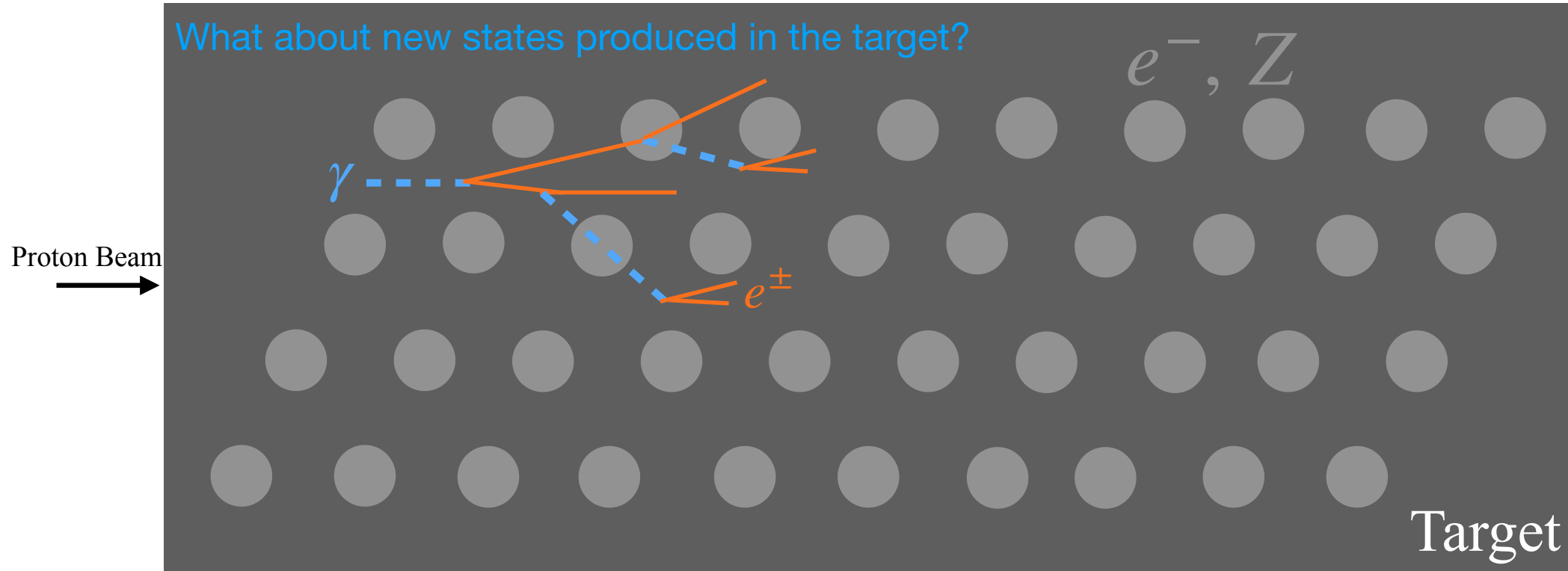
Proton Beam
→

Target

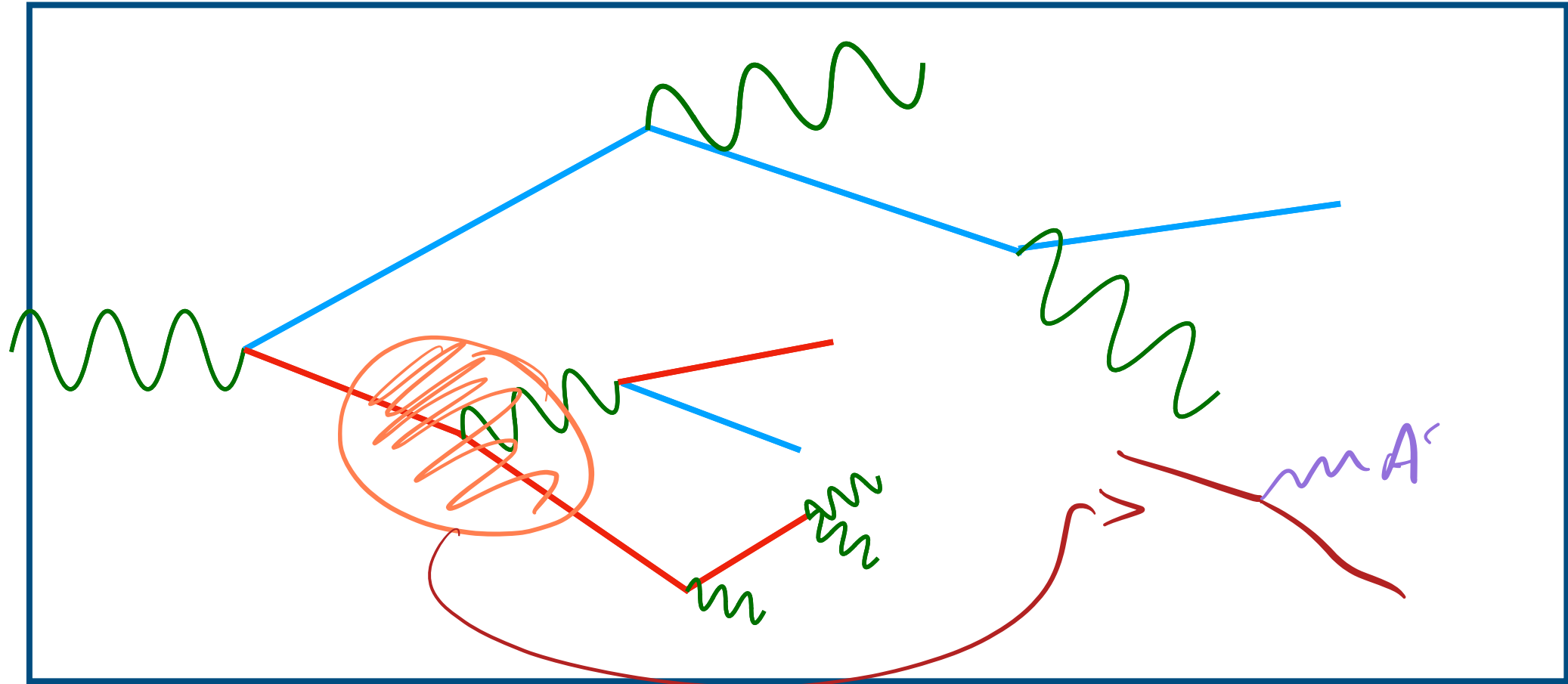


“Nonstandard” BSM Production?

Every Hadronic/Electromagnetic interaction in the target is a potential for BSM production!
many interactions = many opportunities for production



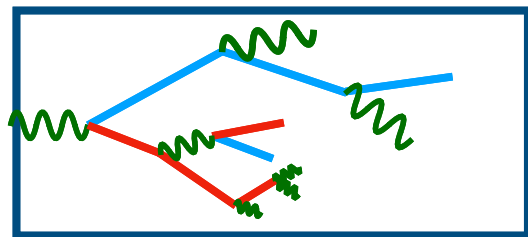
The big challenge: kinematics



The big challenge: kinematics

We are (often) interested in detectors in the \sim -forward region that have a small solid angle with respect to the incident beam.

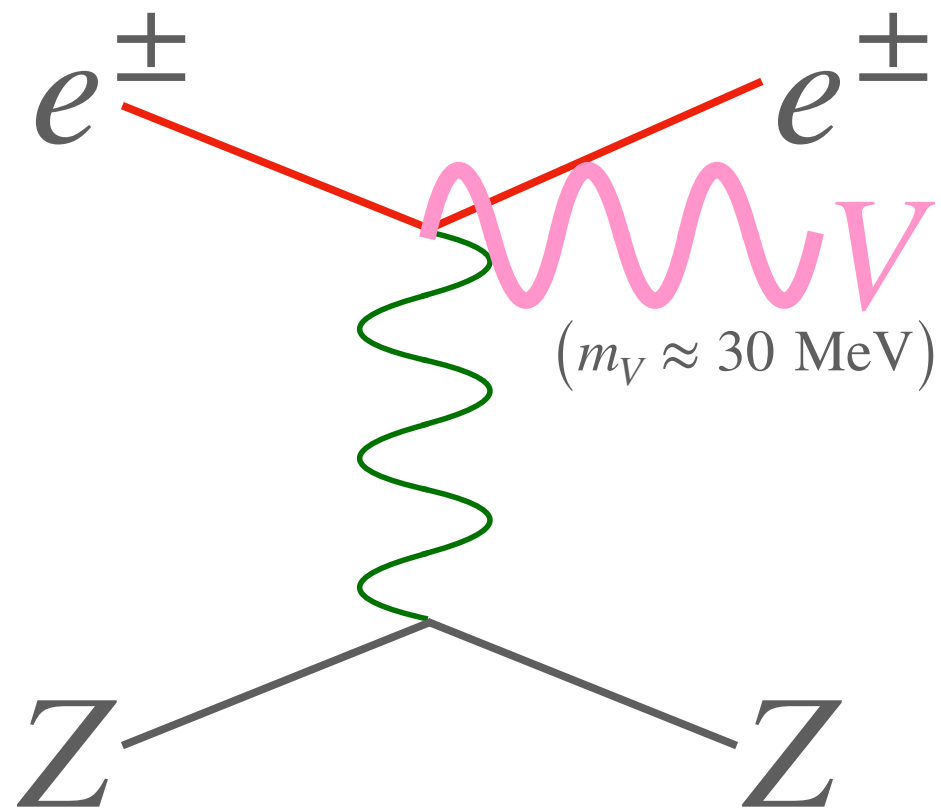
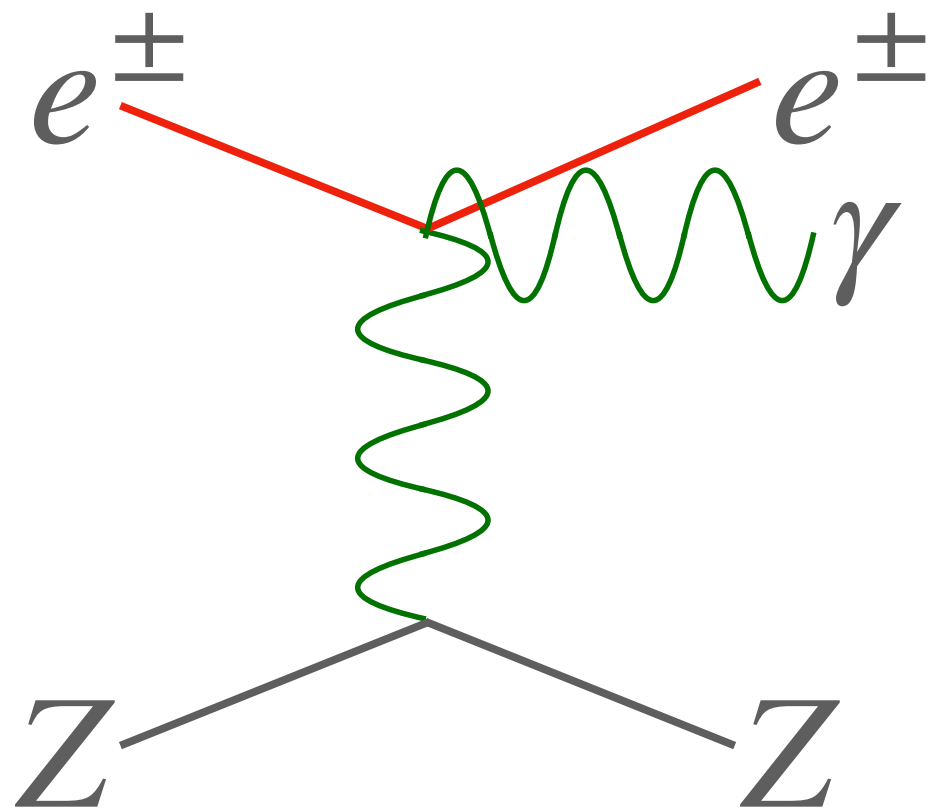
Any small effect in *directionality* of BSM production can have a profound effect.



We want to

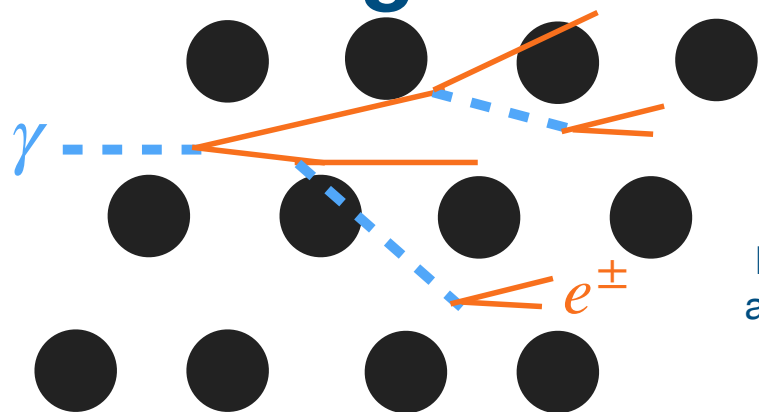
- (a) Generate sample SM showers, and
- (b) Resample those SM vertices to produce BSM states, tracking kinematics precisely.

SM vs. BSM Bremsstrahlung



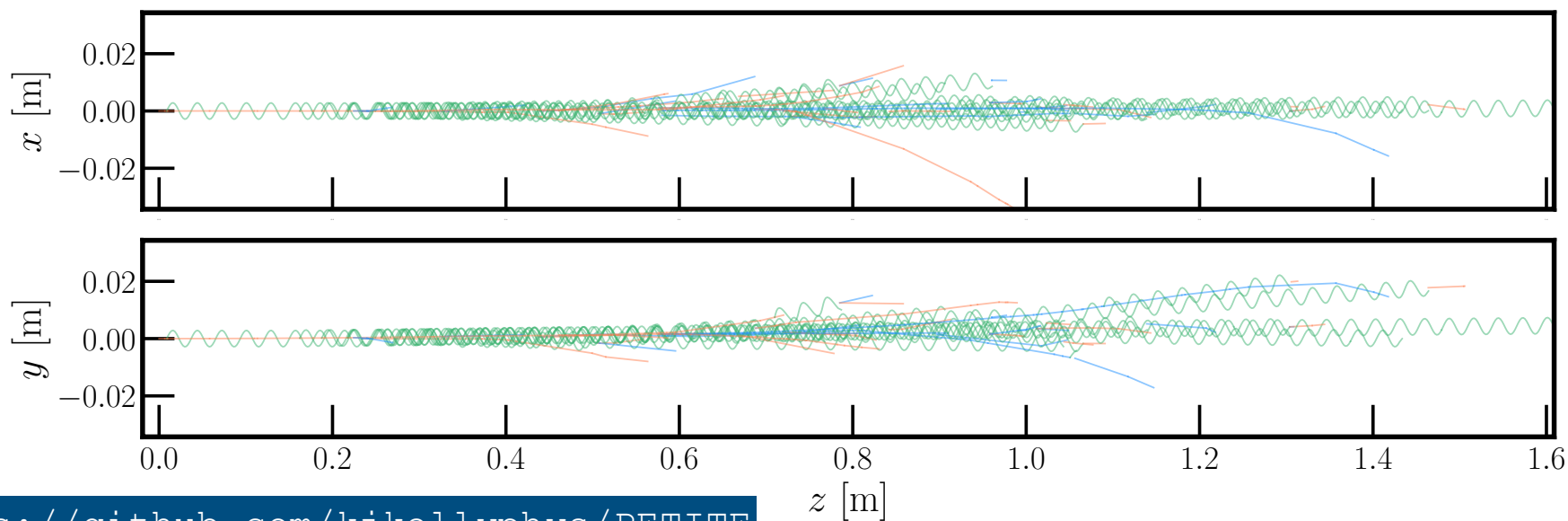
Comparing the two – big issue is kinematical distributions of outgoing particles

Introducing PETITE



PETITE allows for rapid simulation of EM cascades in thick targets that can be processed for determination of BSM flux predictions

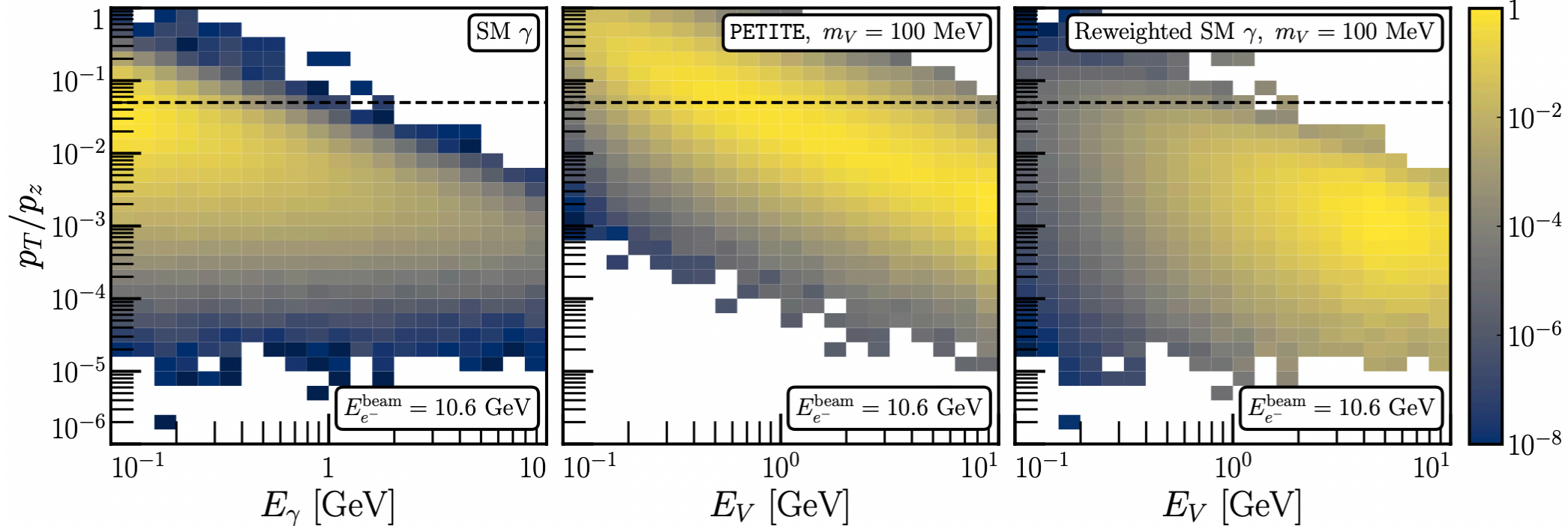
Includes SM effects for energy loss, multiple Coulomb scattering, as well as hard scattering processes. Compares extremely well against dedicated tools (e.g. GEANT-4) and analytic results ([Tsai/Whitis '66](#))



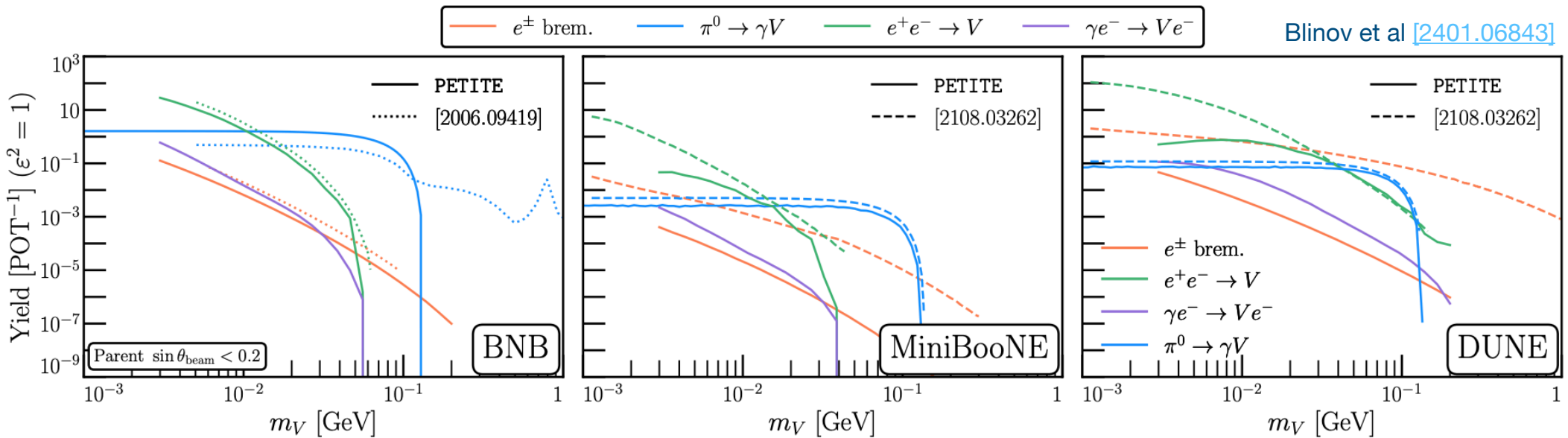
Care with Kinematics

Trying to turn daughter photons into daughter dark photons is tricky because of different kinematics. This has a significant impact, especially for very forward detectors.

Blinov, Fox, KJK, Machado, Plestid [\[2401.06843\]](#)

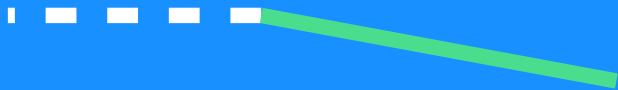
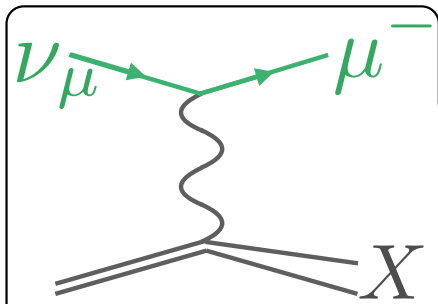


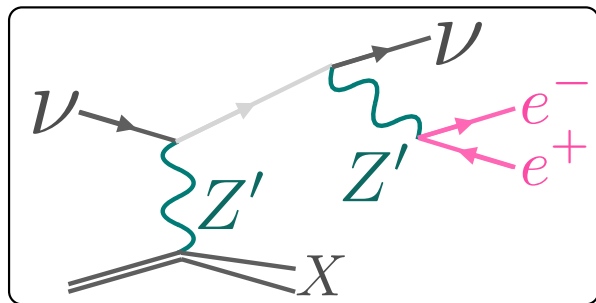
Yields from PETITE



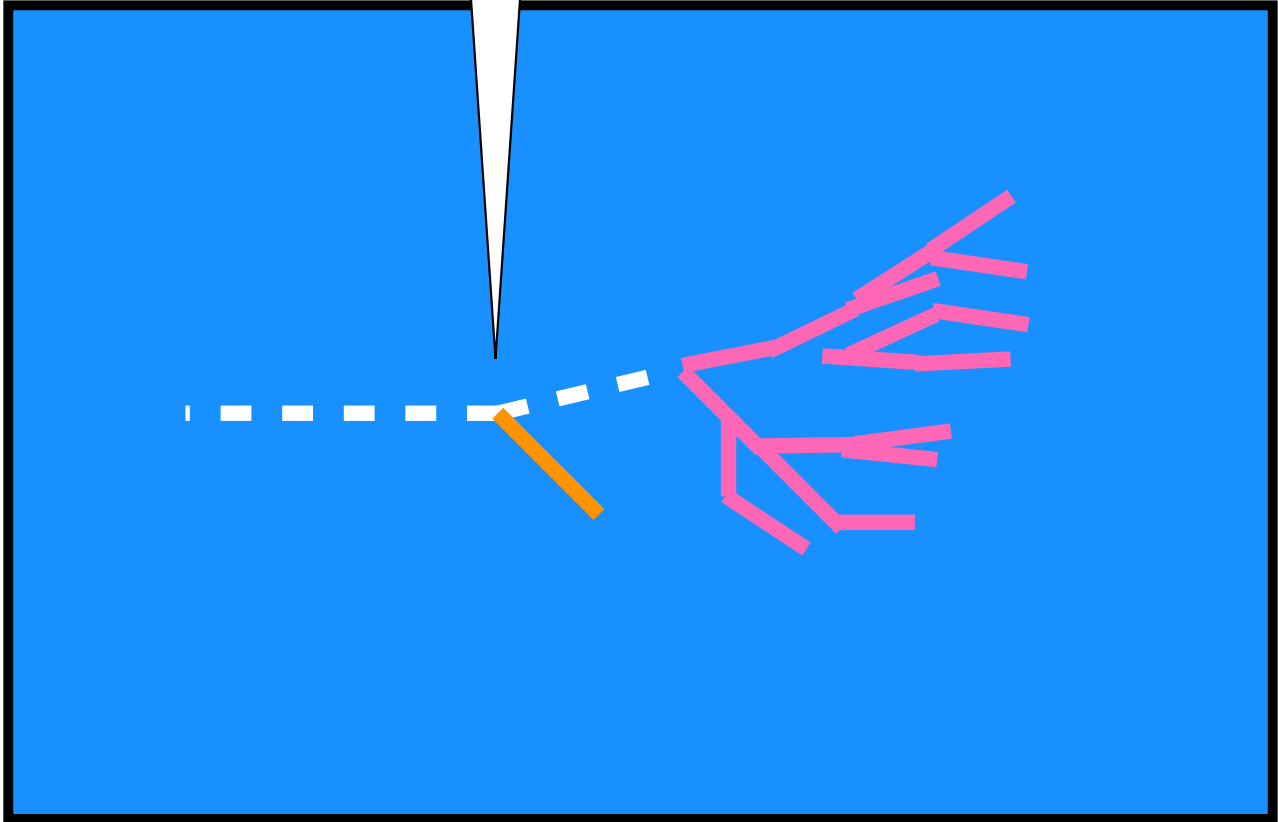
**More intricate BSM Scenarios/
Signatures —
“dark neutrinos”**

Back to our Detector





Back to our Detector



Unexpected neutrino-scattering can lead to novel signatures in the detector. Are we prepared to search for these?

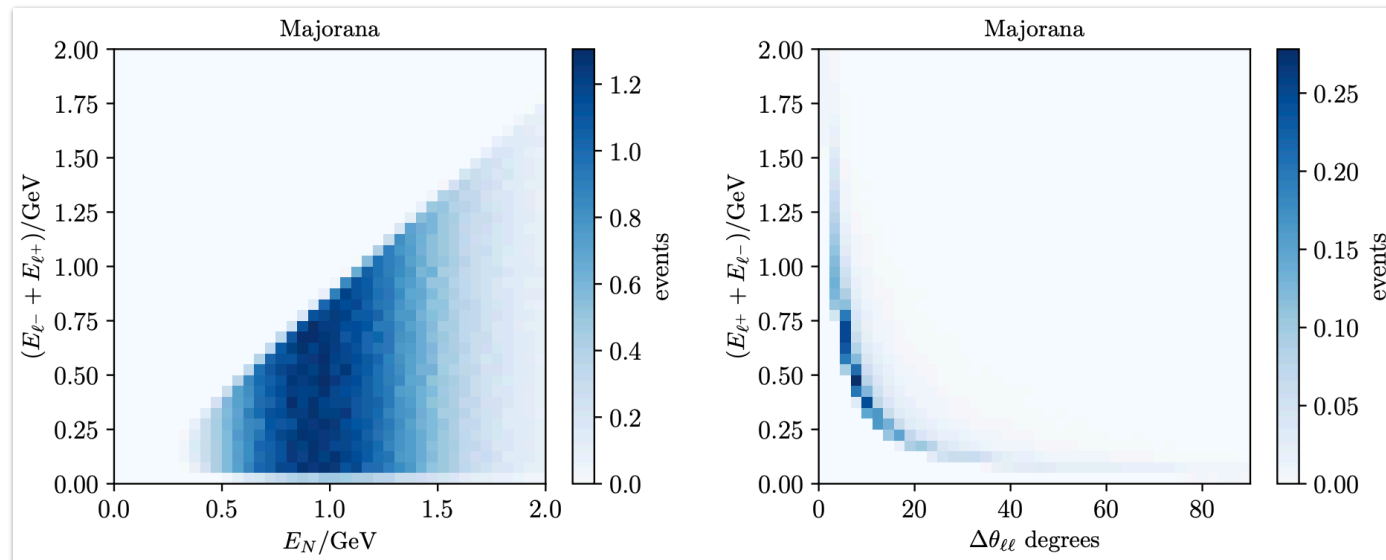
“Dark neutrinos” are a possible solution to the MiniBooNE low-energy excess (since to MiniBooNE, overlapping electron pairs look like a single electron)
 Bertuzzo et al [\[1807.09877\]](#)
 Ballett et al [\[1808.02915\]](#)

How do we simulate such BSM?

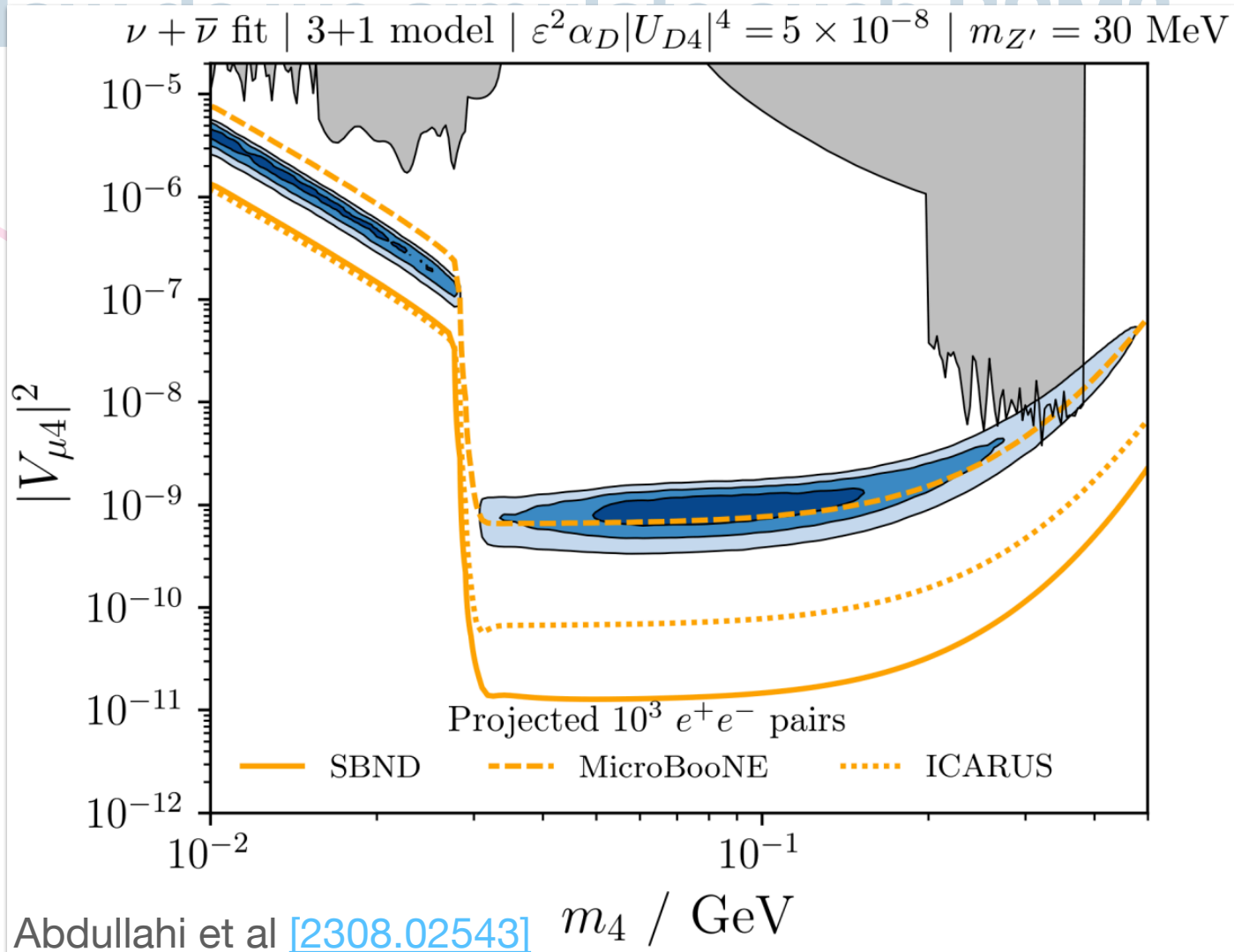
[DarkNews Generator](#)



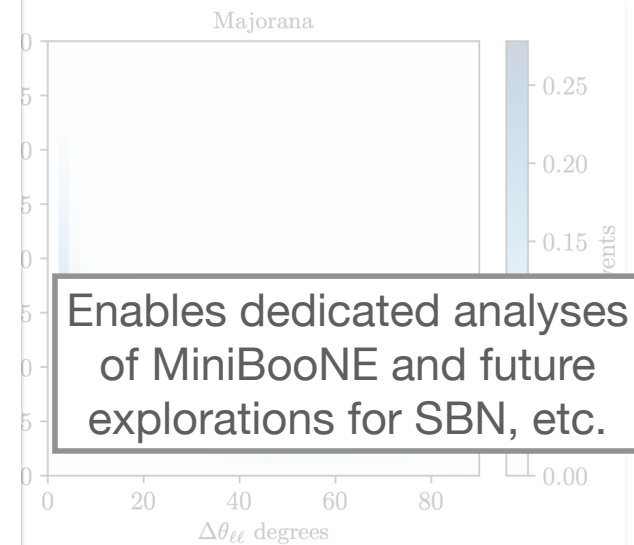
Abdullahi et al [[2207.04137](#)]



DarkNews — purpose-built tool for upscattering-type signatures.



Abdullahi et al [2308.02543]



Enables dedicated analyses of MiniBooNE and future explorations for SBN, etc.

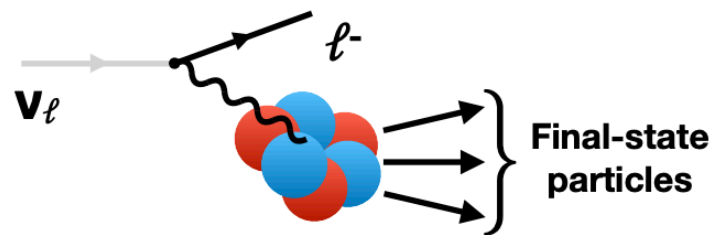
pe signatures.

Another approach?

Isaacson et al [\[2205.06378\]](#)



Factorizability + Modularity!



$$|\mathcal{M}|^2 = L_{\mu\nu} \frac{1}{P^2} W^{\mu\nu} \rightarrow \langle \Psi_0 | J_\mu^\dagger(q) | \Psi_f \rangle \langle \Psi_f | J_\nu(q) | \Psi_0 \rangle$$

Leptonic tensor:
 Calculable analytically
 in SM or BSM
 scenario.

Hadronic tensor:
 Complicated multi-scale
 object, encoding all the
 hadronic/nuclear physics

$|\Psi_0\rangle$: Initial state (say, ^{40}Ar or H_2O)
 $|\Psi_f\rangle$: Final state (nuclear remnant +
 outgoing pions, kaons, etc...)

Version: 1.0.0
 Authors: Joshua Isaacson, William Jay, Alessandro Lovato,
 Pedro A. Machado, Luke Pickering, Noemi Rocco,
 Noah Steinberg

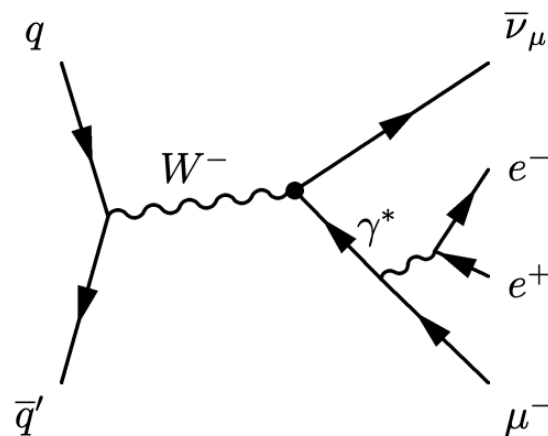
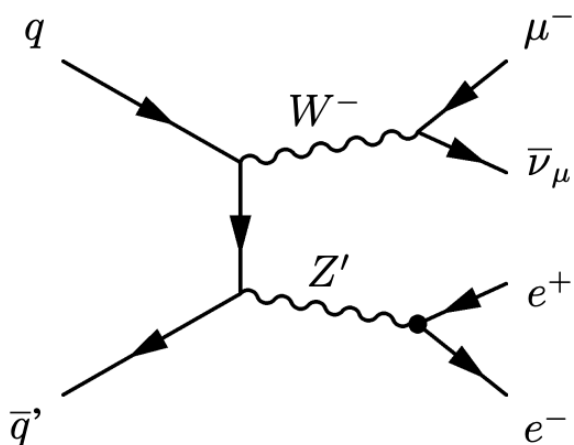
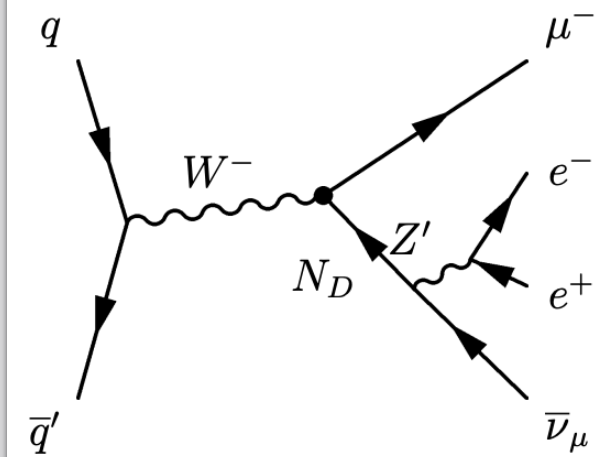
Undergraduate Student Contributions:
 Diego Lopez Gutierrez, Sherry Wang, Russell Farnsworth

Interfaces well with NuHepMC event record format, Gardiner et al [\[2310.13211\]](#)

Another approach?

Isaacson et al [\[2205.06378\]](#)

Complementary model searches with the LHC — Herwig et al [\[2310.13042\]](#) + Modularity!



```

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```

Version: 1.0.0
 Authors: Joshua Isaacson, William Jay, Alessandro Lovato,
 Pedro A. Machado, Luke Pickering, Noemi Rocco,
 Noah Steinberg

Undergraduate Student Contributions:
 Diego Lopez Gutierrez, Sherry Wang, Russell Farnsworth

Leptonic tensor:
 Calculable analytically
 in SM or BSM
 scenario.

Hadronic tensor:
 Complicated multi-scale
 object, encoding all the
 hadronic/nuclear physics

$|\Psi_0\rangle$: Initial state (say, ⁴⁰Ar or H₂O)
 $|\Psi_f\rangle$: Final state (nuclear remnant +
 outgoing pions, kaons, etc...)

Interfaces well with NuHepMC event record format, Gardiner et al [\[2310.13211\]](#)

MicroBooNE Recast:

Higgs-Portal Scalar \longrightarrow

Heavy Neutral Lepton

[\[2106.06548\]](#) KJK & P.A.N. Machado

MicroBooNE Search for Higgs-Portal Scalars

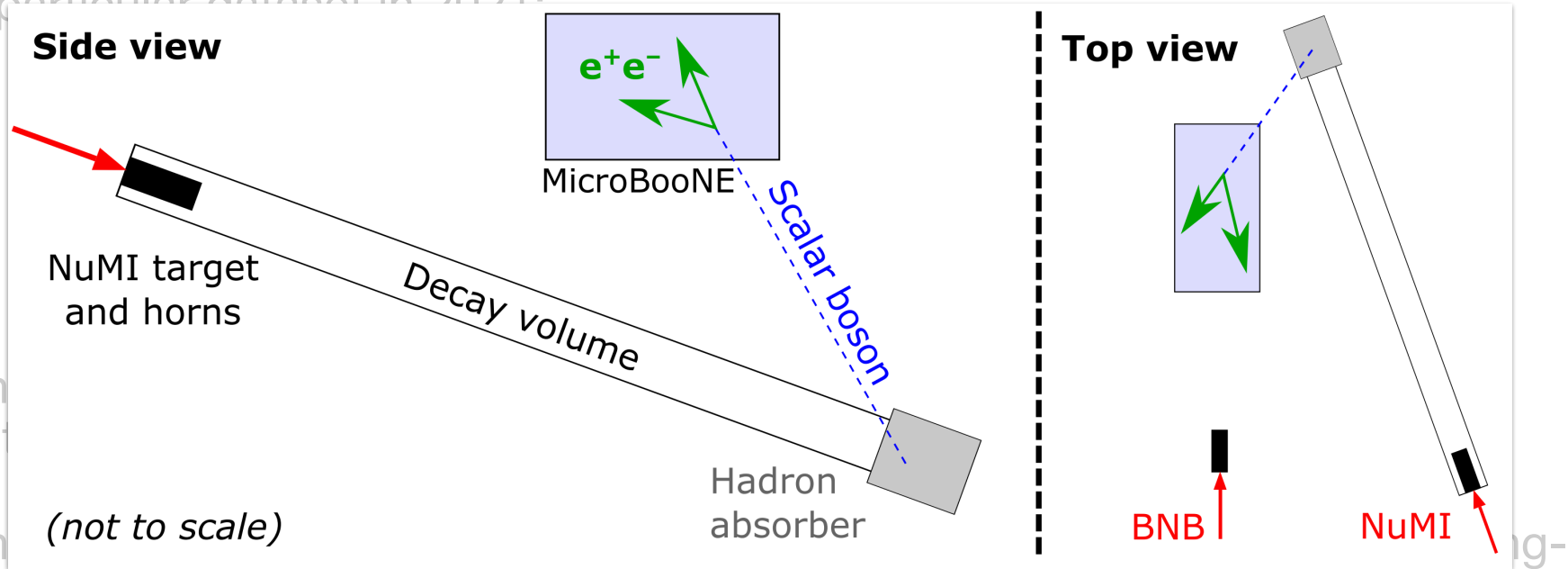
- Inspired by Batell et al [\[1909.11670\]](#), MicroBooNE sought a BSM signature in a particular dataset in 2021:

$$K^+ \rightarrow \pi^+ S, \quad S \rightarrow e^+ e^-$$

- These kaons are produced in the NuMI beam line or absorber, and decay within the absorber.
- The absorber is 100 m from MicroBooNE — the S must be moderately long-lived to reach MicroBooNE and decay inside.

MicroBooNE Search for Higgs-Portal Scalars

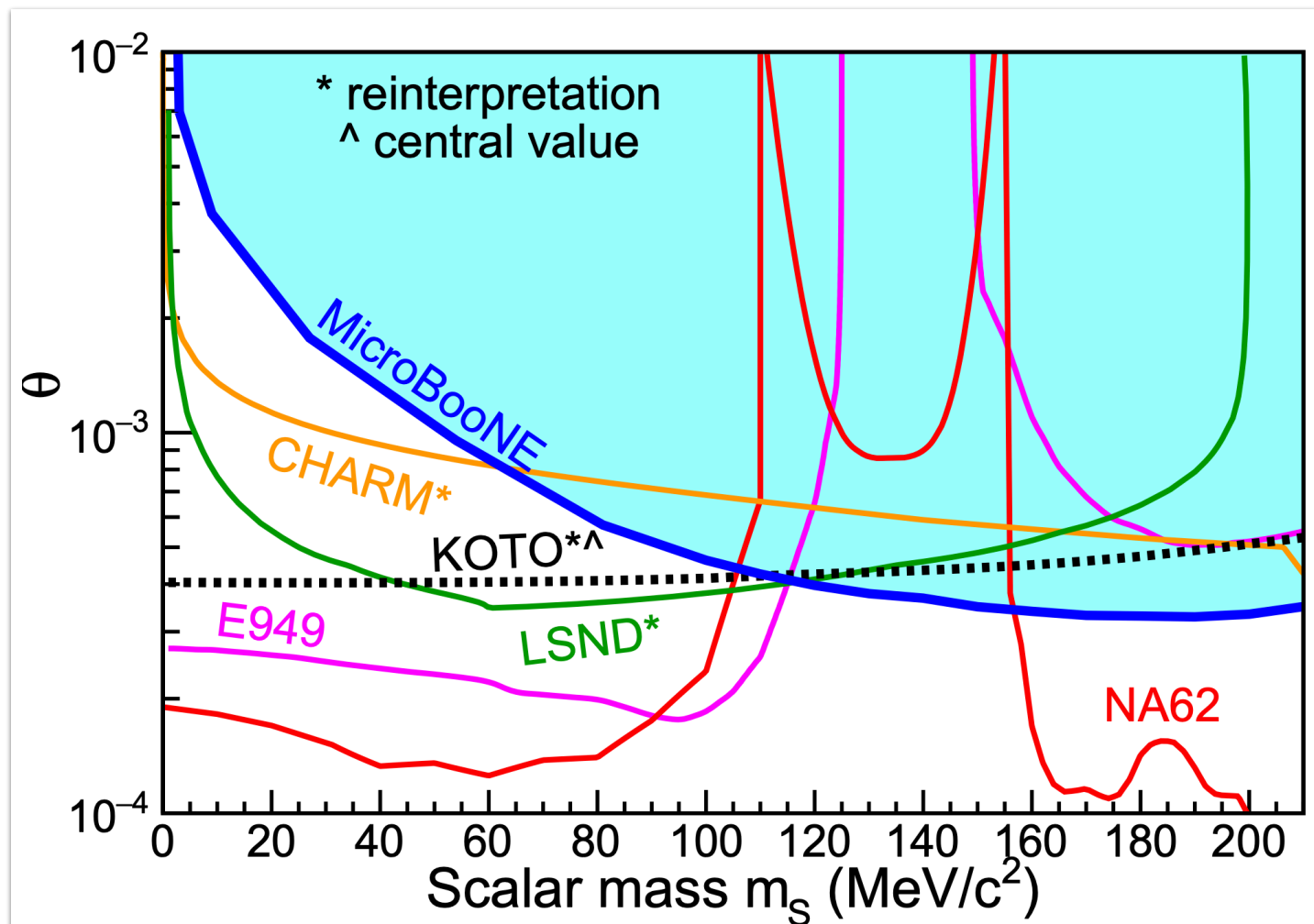
- Inspired by Batell et al [[1909.11670](#)], MicroBooNE sought a BSM signature in a particular dataset in 2021.



- The scalar boson lived to reach MicroBooNE and decay inside.

MicroBooNE Constraint

MicroBooNE [\[2106.00568\]](#)



Signal Rate from KDAR BSM

signal rate

flux of new particles

detector area

$$R_X = \Phi_X A_{\text{det.}} P(X \rightarrow e^+ e^-) \varepsilon(m_X)$$

probability of decay happening in detector

signal efficiency

$$\Phi_X = \frac{N_{\text{KDAR}} \text{Br}(K^+ \rightarrow X)}{4\pi D^2}$$

$$P \approx \frac{L_{\text{det.}}}{\gamma} \Gamma(X \rightarrow e^+ e^-)$$

Flux Example

$$\Phi_X = \frac{N_{K\text{DAR}} \text{Br}(K^+ \rightarrow X)}{4\pi D^2}$$

Scalar Model

$$\text{Br}(K^\pm \rightarrow \pi^\pm \varphi) = 2 \times 10^{-3} \sin^2 \vartheta \rho_\varphi \left(\frac{M_\varphi^2}{m_{K^\pm}^2}, \frac{m_{\pi^\pm}^2}{m_{K^\pm}^2} \right)$$

HNL Model

$$\text{Br}(K \rightarrow \mu N) \simeq \text{Br}(K \rightarrow \mu \nu) |U_{\mu 4}|^2 \rho_N \left(\frac{m_\mu^2}{m_K^2}, \frac{m_N^2}{m_K^2} \right)$$

How to recast HPS to HNL?

- Without digging into the weeds, MicroBooNE recorded data and performed an analysis looking for $S \rightarrow e^+e^-$ signal events.
- This included a boosted decision tree (BDT) trained on signal and background Monte Carlo.
- After cutting on the BDT score, two* candidate events pass, on a background expectation of 1.9 ± 0.8 events.

TABLE II. Estimated signal selection efficiency (eff.) for a scalar boson decay inside the TPC, and event yield [unweighted (unwt.) and beam-on exposure-weighted (exp. wt.), with the expected signal for θ_{KCV}].

Category	Eff. (%)	Event count	
		Unwt.	Exp. Wt.
Beam-off dataset		10	1.1 ± 0.4
Neutrino simulation		16	0.8 ± 0.7
Signal (120 MeV/ c^2)	14.0 ± 0.8	7268	4.9 ± 1.5
Signal (160 MeV/ c^2)	14.9 ± 0.9	7654	12.2 ± 3.6

Goal: as a function of mass, determine the HNL model parameters that predict the *same* signal rate that MicroBooNE has excluded for the Higgs-portal scalar model.

Same Rate?

$$R_X = \Phi_X A_{\text{det.}} P(X \rightarrow e^+e^-) \varepsilon(m_X)$$

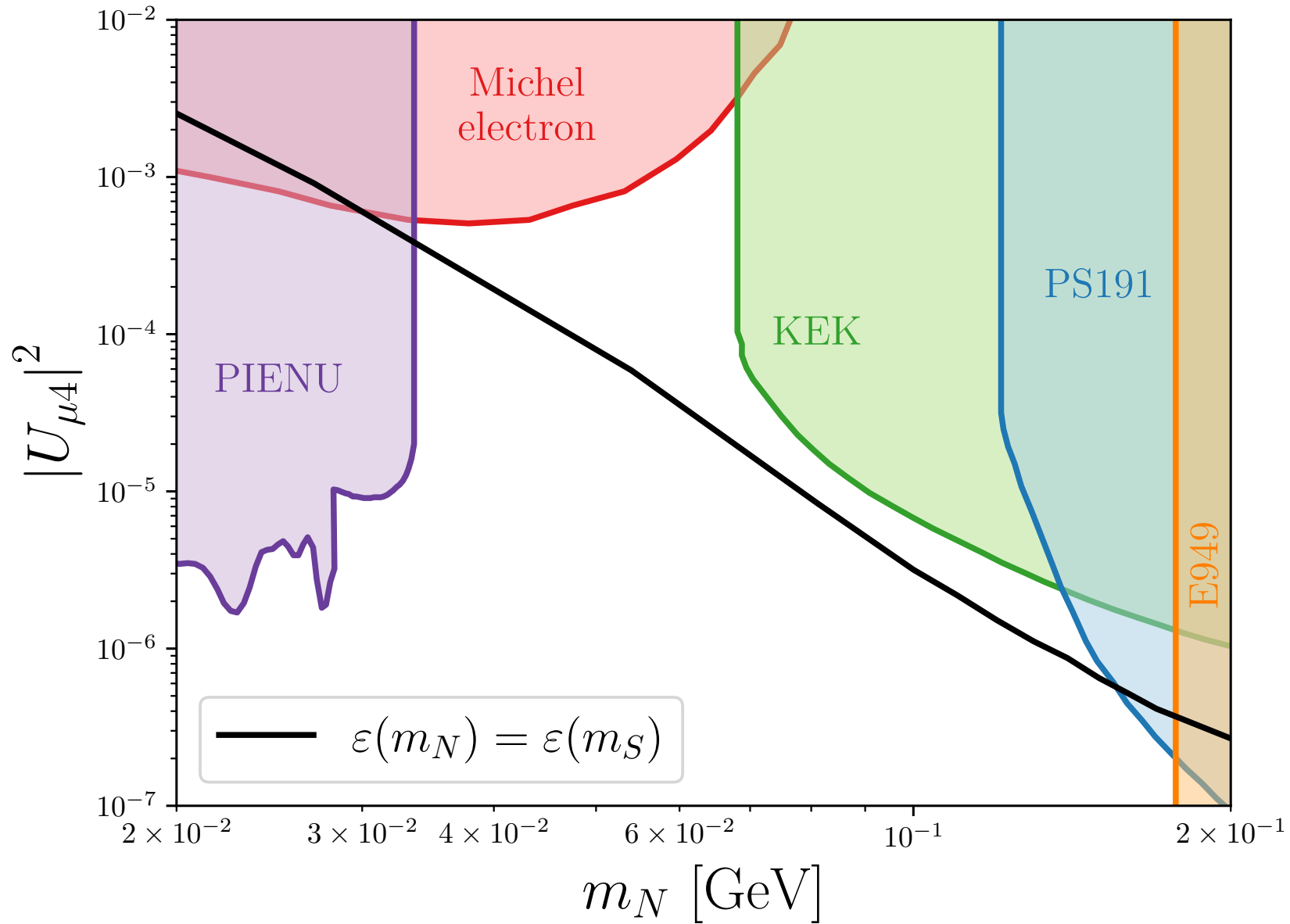
$$\frac{R_N}{R_S} \approx \frac{\text{Br}(K \rightarrow \mu N) \quad m_N E_S \Gamma(N \rightarrow \nu e^+ e^-) \quad \varepsilon(m_N)}{\text{Br}(K \rightarrow \pi S) \quad m_S E_N \Gamma(S \rightarrow e^+ e^-) \quad \varepsilon(m_S)}$$

Given or Calculable

Calculable given $|U_{\mu N}|^2$ (and proportional to that)

Pretend it's equal to $\varepsilon(m_S)$ for now

HPSRecast.ipynb



OK, what about efficiency?

$$R_X = \Phi_X A_{\text{det.}} P(X \rightarrow e^+e^-) \varepsilon(m_X)$$

$$\frac{R_N}{R_S} \approx \frac{\text{Br}(K \rightarrow \mu N) m_N E_S \Gamma(N \rightarrow \nu e^+e^-) \varepsilon(m_N)}{\text{Br}(K \rightarrow \pi S) m_S E_N \Gamma(S \rightarrow e^+e^-) \varepsilon(m_S)}$$

What goes into signal efficiency?

Training Information from MicroBooNE

MicroBooNE [\[2106.00568\]](#)

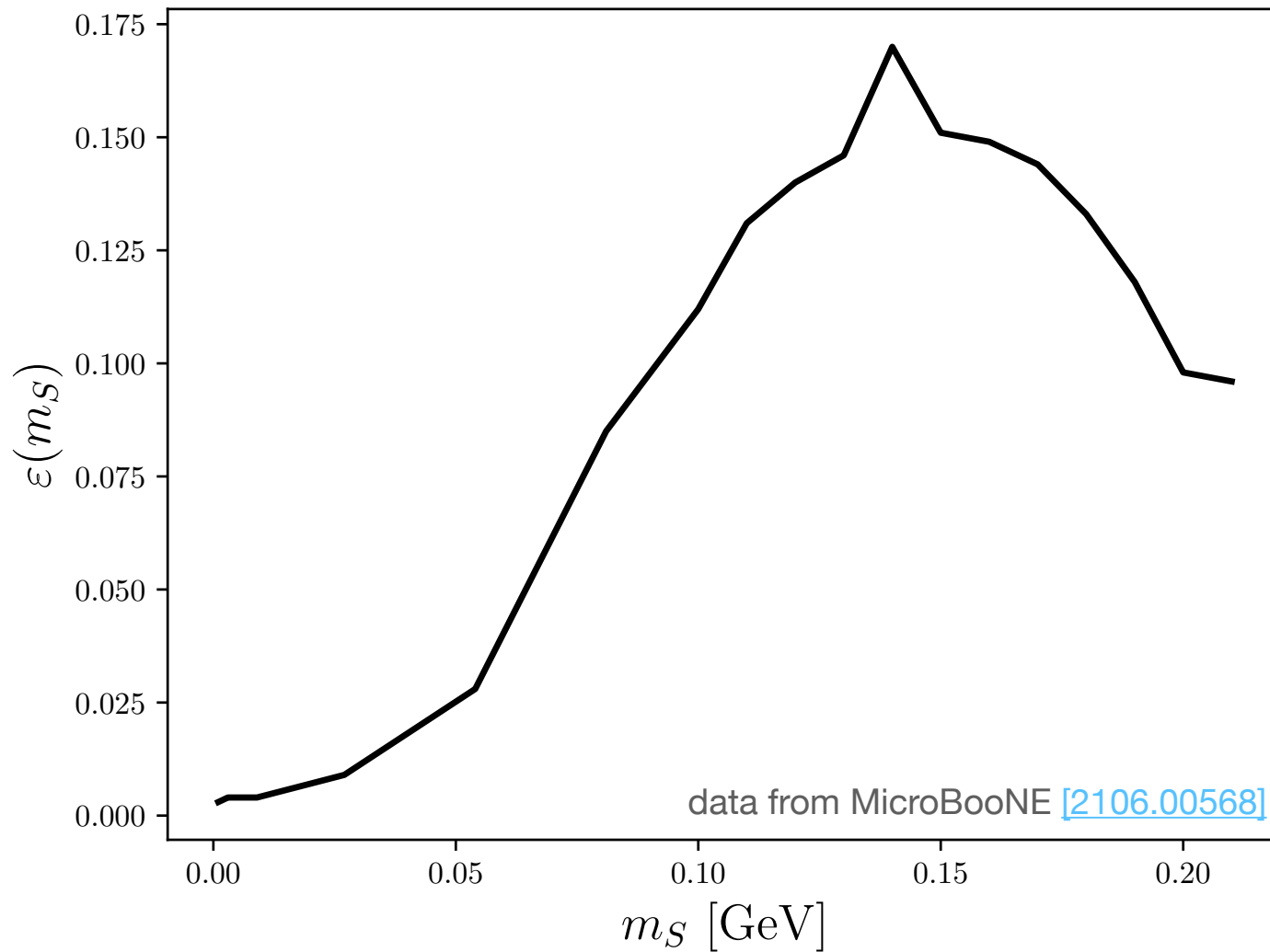
We apply two different BDTs to the preselected candidates: one trained against cosmic backgrounds and one trained against neutrino interactions simulated inside the cryostat. Each BDT is trained separately over the run 1 events and run 3 events, i.e., there are four BDTs in total. We split the run periods because the use of the CRT in run 3 and the differences between forward and reverse horn current operations can change the topologies and properties of the background distributions that the BDTs are trained against. We use `xgboost` [27] to train and apply the BDTs. We train the BDTs on ten input variables each. Nine of the ten input variables are the same for the cosmic-focused and neutrino-focused BDTs. These are (1) the opening angle between the two reconstructed objects; (2) the opening angle in the plane transverse to the

hadron absorber direction from the detector center; (3,4) the two angles between the two objects and the hadron absorber direction; (5) the Pandora track or shower score of the larger of the two objects (when ordered by number of hits); (6) the number of hits of the larger object; (7) the total number of hits contained in other objects in the slice, not including the two objects that form the decay candidate; (8) the maximum y coordinate, relative to the decay vertex position, of shower start positions or track start or end positions, for any other objects in the slice; and (9) the minimum z coordinate, relative to the decay vertex position, of shower start positions or track start or end positions, for any other objects in the slice. The last two variables are treated as “missing” within `xgboost` if the slice contains only two objects. The tenth input variable of the cosmic-focused BDT is the length of the larger object. The tenth input variable of the neutrino-focused BDT is the number of tracks in the slice. For all

Ansatz:

(1) dominates the signal efficiency as a function of BSM particle mass

HPS Efficiency



Generate e^+e^- events in the rest-frame of the decaying HPS/HNL

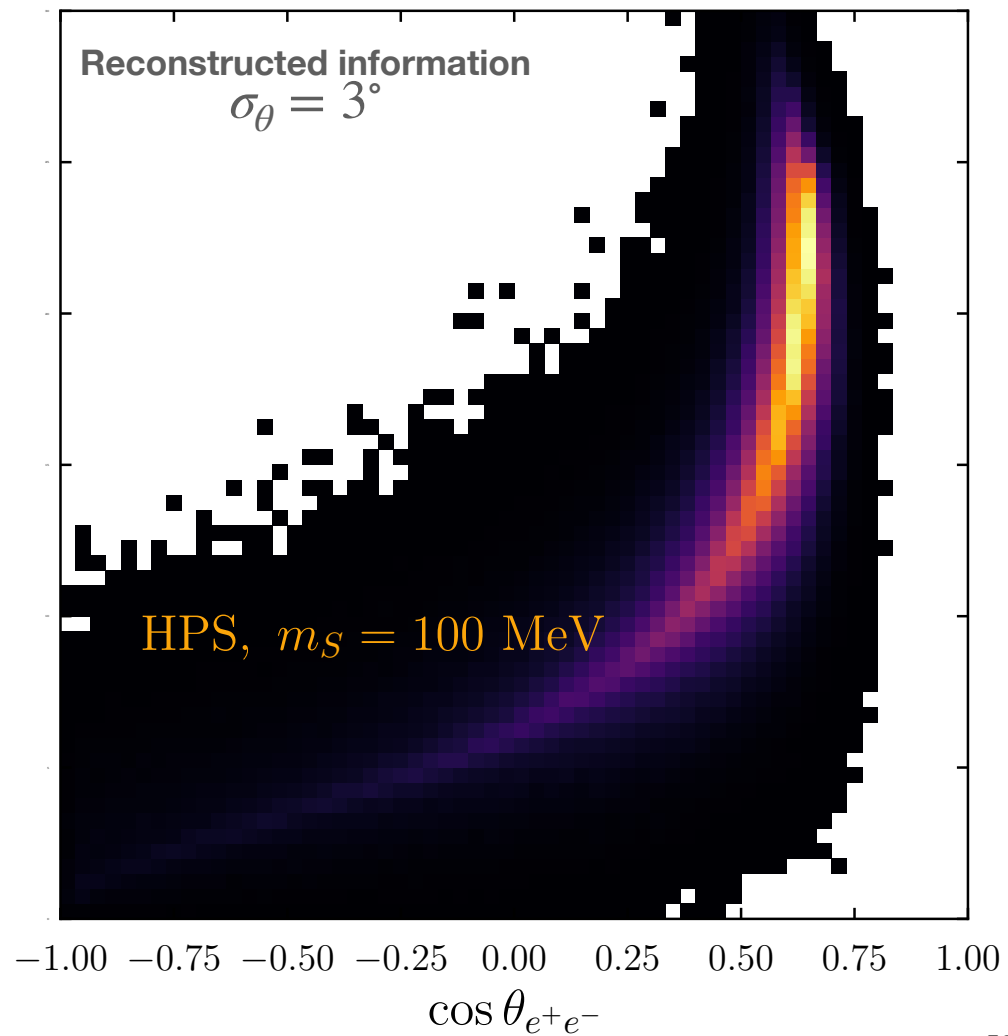
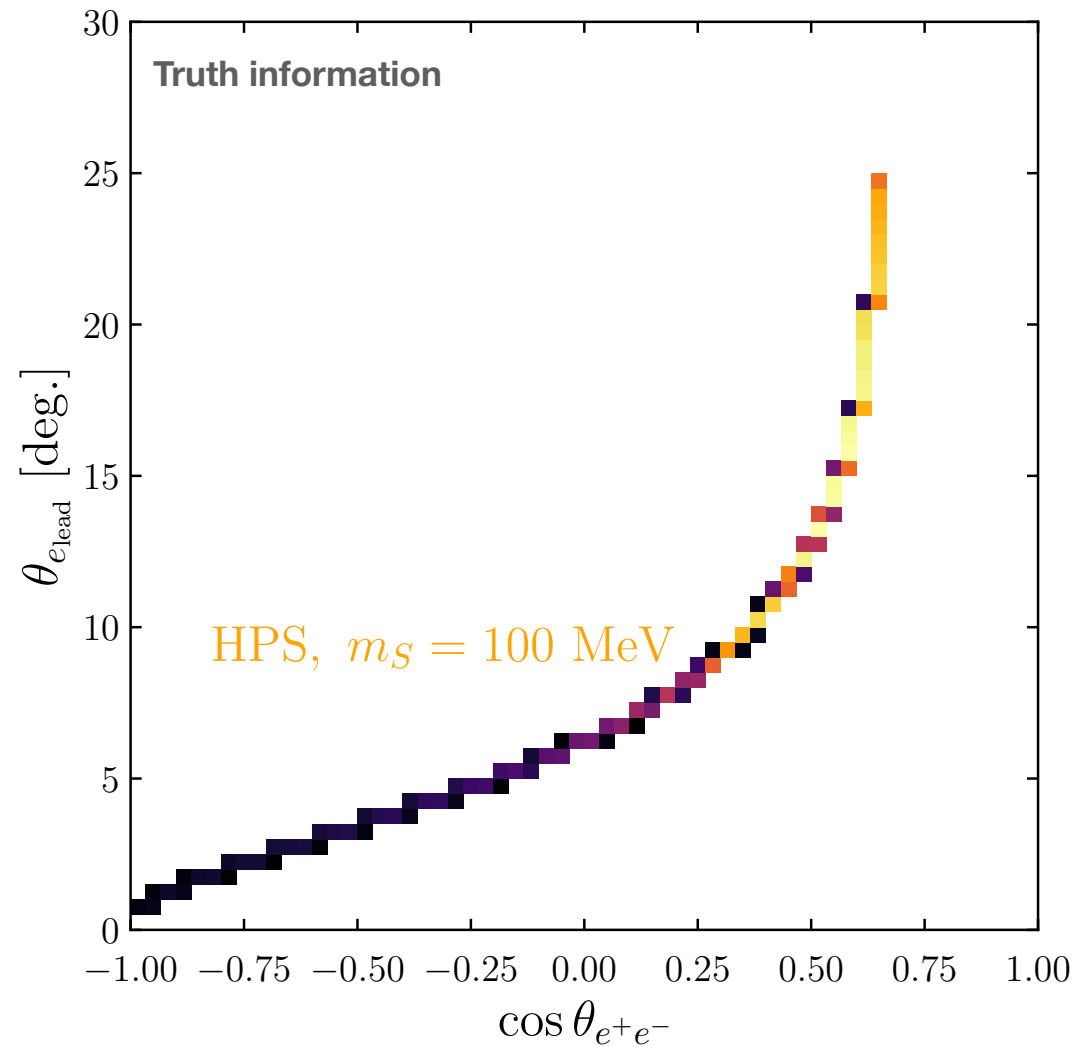
(depends on `RestFrame.py` for HNL three-body kinematics and `vegas` for phase-space sampling)

`RestFrame.py`

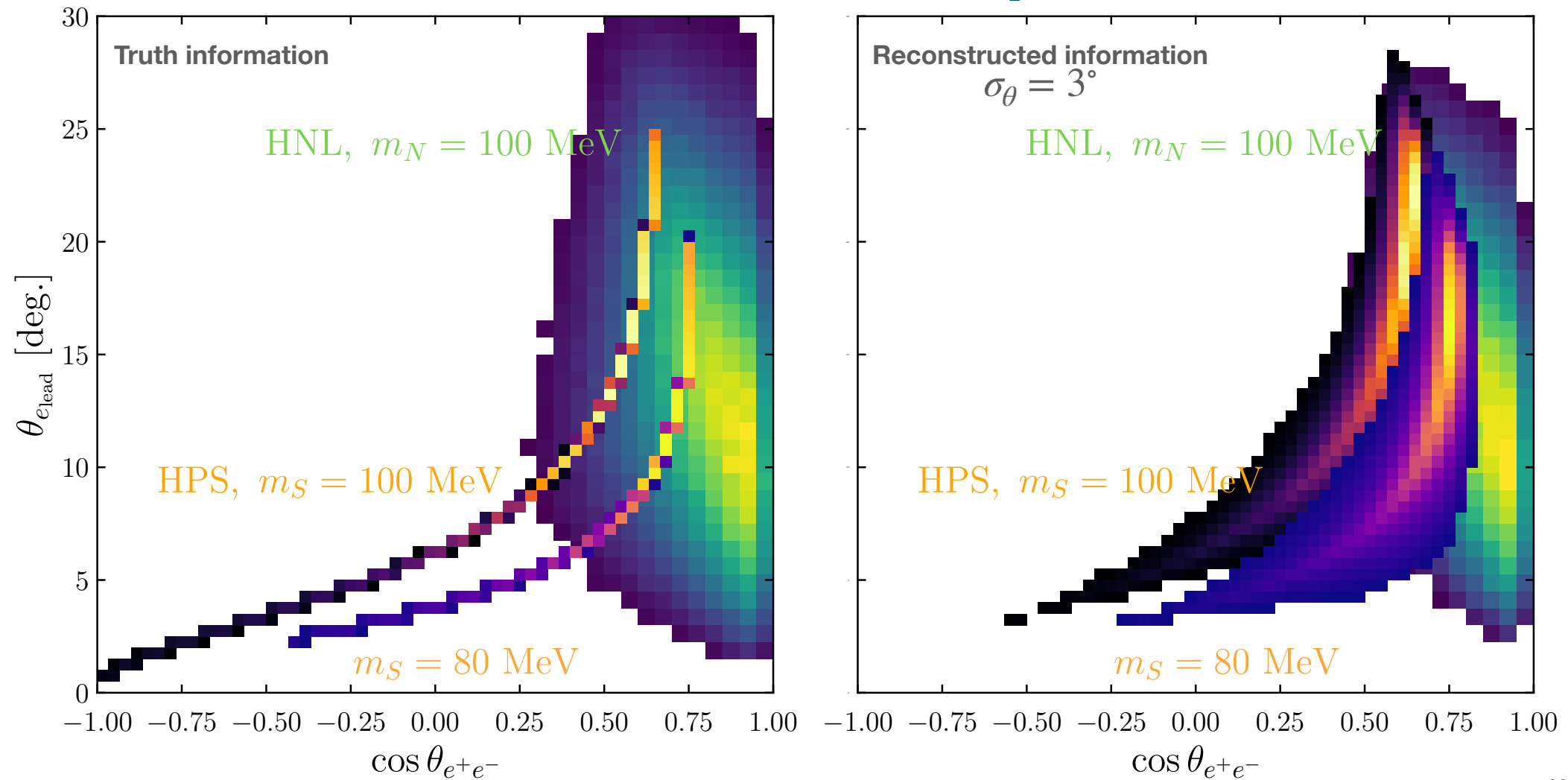
`LabFrame.py`

Transforms event to the laboratory frame, smears events, and performs different reconstruction/analyses on the events.

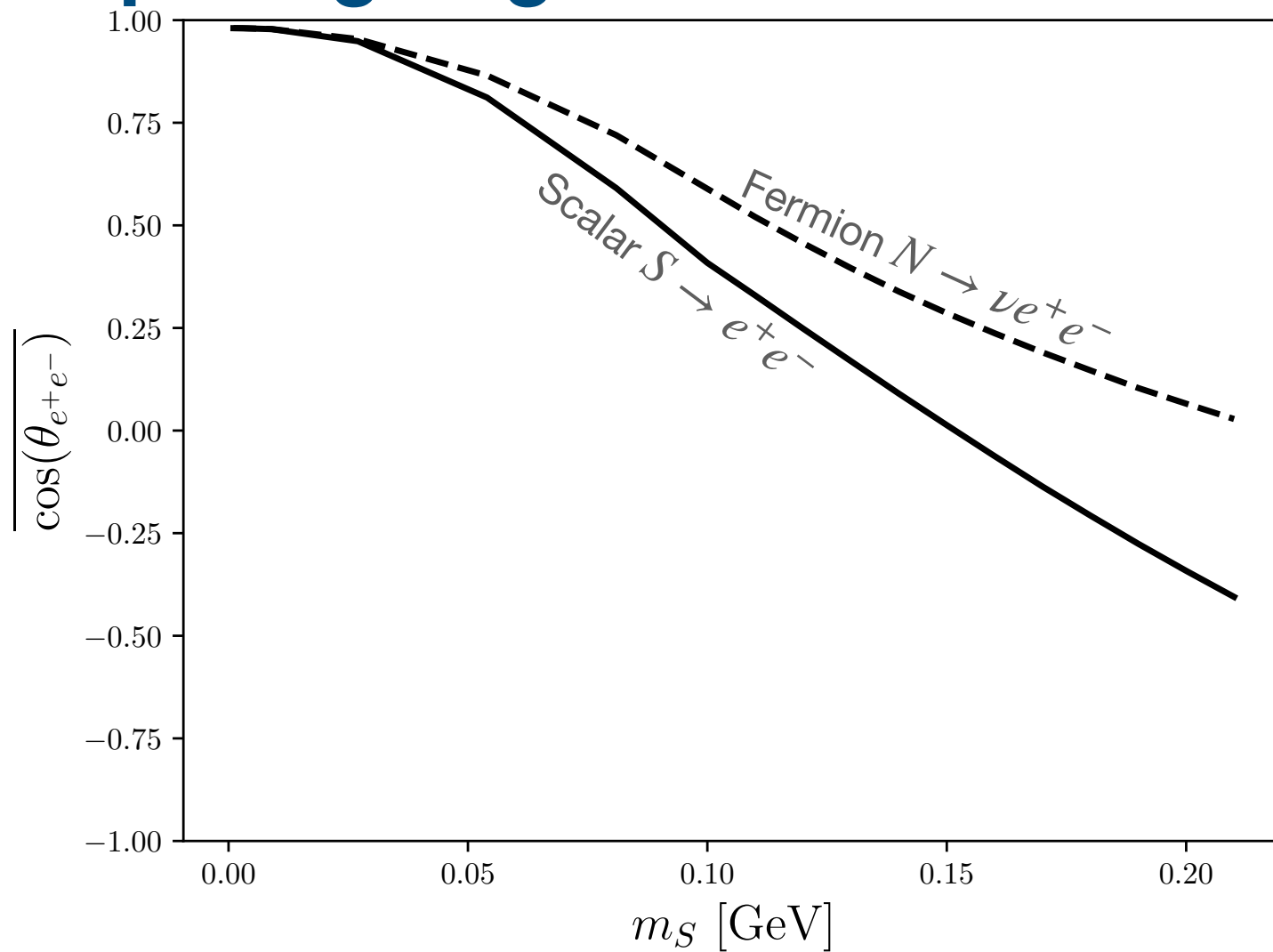
Event Distributions



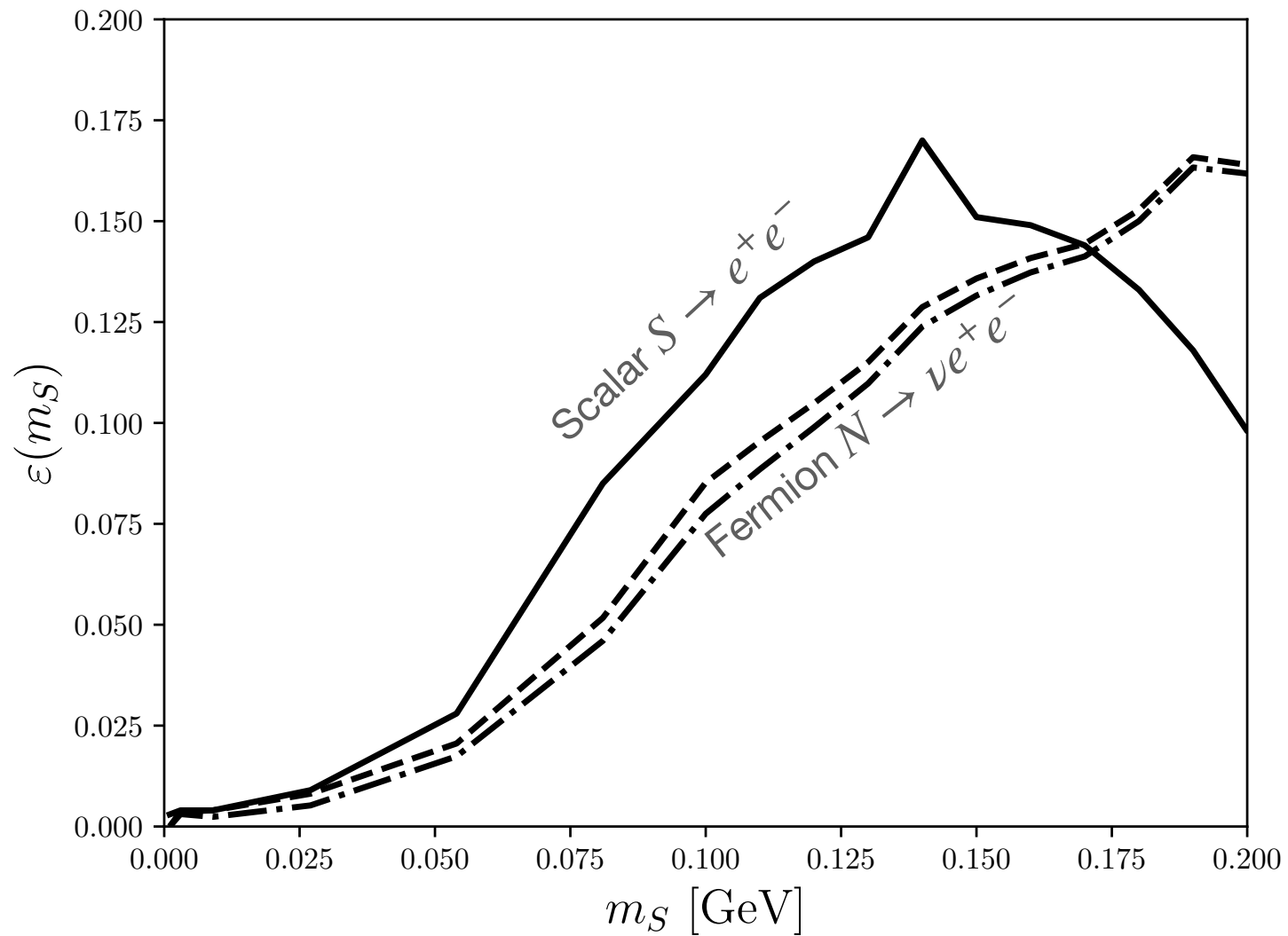
HNL vs HPS in this Kinematic Space



Average Opening Angles



Updated Efficiency for HNLs



New Constraint on HNLs

