### **Dissecting Neutrino Event Generators** A look at what's inside (and what could be wrong)





*The anatomy lesson of neutrino generators* **Rembrandt van Rijn (2024)** 

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# Focus of this lectures

- There is not time to cover all the interesting physics associated with neutrino interaction modelling in generators
- You'll get a slightly biased choice of topics!
- We'll stay mostly qualitative, I'll try to give a conceptual overview of the topics most relevant to ongoing experiments
- Lots of places to learn more:
- INSS lecture slides
- From eV to EeV (Formaggio and Zeller)
- NuSTEC White Paper
- <u>Xsecs for Oscillations (Katori and Martini)</u>
- <u>e-scat vs nu-scat (SuSAv2 group)</u>

- <u>K. McFarland's Lectures</u>
- <u>S. Boyd's Lectures</u>
- <u>T. Golan's thesis</u>
- <u>G. Megias' thesis</u>
- GiBUU based summaries (<u>1</u>,<u>2</u>)
- <u>Semi-inclusive interactions (Donnelly talk at ECT\* 2018)</u>

(Which I liberally borrowed material from when making these slides!)

# Overview

- Inputs to Generators: Neutrino-nucleon interactions
   QE, RES and DIS
- Inputs to Generators: Neutrino-nucleus interactions

   Nuclear effects
- Neutrino event generators

   Filling in gaps in our inputs
- Benchmarking generators with measurements
   o Inclusive successes and exclusive failures
- Why do we care?
   Neutrino interactions for neutrino oscillations
- Don't Panic! The future of neutrino interaction simulations

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• Don't Panic! The future of neutrino interaction simulations





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### Recap of neutrino scattering



Nucleon: mostly harmless

- Cross-Sections for point-like neutrino scattering with electrons or quarks are relatively easy to calculate
- In most experiments, **neutrinos interact with nucleons or a nucleus**

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- Cross-Sections for point-like neutrino scattering with electrons or quarks are relatively easy to calculate
- In most experiments, neutrinos interact with nucleons or a nucleus
- Next slides describe our baseline models for simulating neutrinonucleon interactions that go into event generators

# Neutrino nucleon scattering (See Aaron's lectures



Increasing Energy Transfer

### Quasi-elastic Scattering

#### See Aaron's lectures



 Let's work with the "easiest" neutrino-nucleon interaction: CCQE (= charged current quasi elastic)

$$M \sim \frac{g_{w}^{2}}{8} \frac{1}{M_{W}^{2}} [\bar{u}_{\mu} \gamma_{\mu} (1 - \gamma_{5}) u_{\nu}] [\bar{u}_{p} (\dots) u_{n}]$$

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### The hadronic current

Equation shamelessly lifted from G. Perdue's other 2012 INSS lecture

$$\begin{aligned} J_{H}^{\beta} &= \bar{u}_{p} \left[ f_{1V} \gamma^{\beta} + i \frac{\xi f_{2V}}{2M} \sigma^{\beta\delta} q_{\delta} + \frac{f_{3V}}{M} q^{\beta} + f_{A} \gamma^{\beta} \gamma_{5} + \frac{f_{p}}{M} q^{\beta} \gamma_{5} + \frac{f_{3A}}{M} \left( P_{p}^{\beta} + P_{n}^{\beta} \right) \gamma_{5} \right] u_{n} \\ M &= \left( M_{p} + M_{n} \right) / 2 \qquad q = p_{\nu} - p_{\mu} = P_{p} - P_{n} \qquad \xi = \mu_{p} - \mu_{n} \qquad \sigma^{\mu\nu} = \frac{i}{2} \left[ \gamma^{\mu}, \gamma^{\nu} \right] \\ \xi \text{ is the difference between proton and neutron anomalous magnetic moments} \end{aligned}$$

• A long and horrible expression for generalised scattering of an extended object

See Aaron's lectures

- The f factors are the "form factors"
  - Their Fourier transform represent a physical distribution
  - $\circ$  Dipole  $\rightarrow$  exponential
- Most can be constrained via electron scattering with one exception
   Their Fourier transform represent a physical distribution
  - $f_A$ , we guess the form of! One free parameter:  $M_A$

$$f_A\left(q^2\right) = \frac{f_A\left(0\right)}{\left(1 - \frac{q^2}{M_A^2}\right)^2}$$

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#### Aside: recent lattice QCD updates

- Some generators benefit from recent calculations from LQCD
- These suggest a dipole doesn't work
- See Aaron's slides for more details!



## Llewellyn-Smith CCQE

- See Aaron's lectures
- Putting this all together gets us to the cross section for neutrino-nucleon interactions we have in our generators:

$$\frac{d\sigma}{d|q^2|} {\nu n \to \ell^- p \choose \overline{\nu} p \to \ell^+ n} = \frac{M^2 G^2 \cos^2 \theta_c}{8\pi E_{\nu}^2} \left[ A(q^2) \mp B(q^2) \frac{(s-u)}{M^2} + \frac{C(q^2)(s-u)^2}{M^4} \right]$$
  
(s-u = 4ME<sub>\nu</sub> + q<sup>2</sup> - m<sup>2</sup>).

Neutrino reactions at accelerator energies, Llewellyn Smith, 1972

$$\begin{split} A \simeq & \frac{t}{M^2} \left( |f_{1V}|^2 - |f_A|^2 \right) + \frac{t^2}{4M^2} \left( |f_{1V}|^2 + \xi^2 |f_{2V}|^2 + |f_A|^2 + 4\xi \operatorname{Re}\left(f_{1V}f_{2V}^*\right) \right) \\ & + \frac{t^3 \xi^2}{16M^6} |f_{2V}|^2 \\ B \simeq & \frac{1}{M^2} \left( \operatorname{Re}\left(f_{1V}f_A^*\right) + \xi \operatorname{Re}\left(f_{2V}f_A^*\right) \right) t \qquad C = & \frac{1}{4} \left( |f_{1V}|^2 + |f_A|^2 - \frac{\xi^2 |f_{2V}|^2}{4M^2} t \right) \end{split}$$

• It's a long expression, but only one unknown (in a dipole model):  $M_A$ 





#### Increasing Energy Transfer

### Resonant Pion Production

See Minoo's lectures

CCRES



CC Single Pion Production (SPP) final states

$$\begin{array}{ll}
\nu_{\mu} \, p \to \mu^{-} \, p \, \pi^{+}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, p \, \pi^{-} \\
\nu_{\mu} \, n \to \mu^{-} \, p \, \pi^{0}, & \overline{\nu}_{\mu} \, p \to \mu^{+} \, n \, \pi^{0} \\
\nu_{\mu} \, n \to \mu^{-} \, n \, \pi^{+}, & \overline{\nu}_{\mu} \, n \to \mu^{+} \, n \, \pi^{-}
\end{array}$$

D. Rein and L. Sehgal, Ann. Phys. 133, 79 (1981)

- Neutrinos can excite a nucleon into a resonance state, which decays to give a nucleon + meson final state
- The dominant resonance is  $\Delta(1232)$  but others can contribute, as can nonresonant pion production
- And the contributions from each should have interference terms ...
- Resonance models are complicated!
- Whilst CCQE scattering on the nucleon can described fully with 1 variable the multi-particle final state for SPP requires 4:



### **Resonant Pion Production**



**CCRES** 



Current Matrix Elements from a Relativistic Quark Model\*

R. P. Feynman, M. Kislinger, and F. Ravndal

Lauritsen Laboratory of Physics, California Institute of Technology, Pasadena, California 91109 (Received 17 December 1970)

The model's used in today's neutrino generators are often based on an approximate model from the 1970s

ficing theoretical adequacy for simplicity. We shall choose a relativistic theory which is naive and obviously wrong in its simplicity, but which is definite and in which we can calculate as many things as possible – not expecting the results to agree exactly with experiment, but to see how closely our "shadow of the truth" equation gives a partial reflection of reality. In our attempt to maintain simplicity, we shall evidently have to violate known principles of a complete relativistic field theory (for example, unitarity). We shall attempt to modify our calculated results in a general way to allow, in a vague way, for these errors.

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### **Resonant Pion Production**



CCRES

- New theory calculations tuned to precision electron scattering data are on the horizon
  - E.g.: <u>MK model at NuINT 2022</u>
- The axial component remains a challenge



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### Neutrino nucleon scattering



Increasing Energy Transfer

#### NuSTEC Summer School, CERN, June 2024

See Minoo's

lectures



**CCDIS** 



- Given enough energy, neutrinos can resolve the quarks within a nucleon. This is deep inelastic scattering.
- At high energies, the *inclusive* (i.e. integrating over possible hadronic final states) cross-section is fairly well understood (perturbative QCD):

$$\frac{d^2 \sigma^{\nu,\overline{\nu}}}{dx \, dy} = \frac{G_F^2 M E_{\nu}}{\pi \left(1 + Q^2 / M_{W,Z}^2\right)^2} \begin{bmatrix} \frac{y^2}{2} 2x F_1(x,Q^2) + \left(1 - y - \frac{Mxy}{2E}\right) F_2(x,Q^2) \\ \pm y \left(1 - \frac{y}{2}\right) x F_3(x,Q^2) \end{bmatrix}$$

#### Bjorken x and y

$$x = \frac{Q^2}{2M\nu} = \frac{Q^2}{2ME_\nu y}$$
$$y = E_{had}/E_\nu$$
$$Q^2 = -m_\mu^2 + 2E_\nu (E_\mu - p_\mu \cos \theta_\mu)$$



#### **CCDIS**



- At low energies (or actually low  $Q^2$ ) QCD becomes non-perturbative.
- Bodek-Yang: extrapolate down to low  $Q^2$ assuming some parametrised scaling. Fix the details with e-scattering, apply to  $\nu$ - scattering
- But this is an empirical treatment that comes with uncertainties



https://arxiv.org/abs/1012.0261

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#### **CCDIS**



- The hadronic side of DIS interactions requires more empirical treatments
- Often the PYTHIA generator is used, but this is really built for much higher energies than used in most neutrino experiments



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### "I would not trust PYTHIA for anything with less than 6 pions"

S. Prestel (a PYTHIA author)

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- Low-W, CC0π and Δ region well covered by SBN data
- Timely data to guide ongoing theory efforts
- DUNE (ND here) has a significant high-W component
- Not well covered by theory
- No relevant data on <sup>40</sup>Ar

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#### C. Wilkinson, NuPhys 2018

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  - QE: almost calculable with some form factors
  - RES: much more difficult, lots of diagrams to consider
  - DIS: easy for *inclusive* high Q<sup>2</sup>, hard at low Q<sup>2</sup>, **hadronic side a total guess**

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Charge Exchange





### Beyond nucleon scattering



Nucleon: mostly harmless

- I leave most details to other talks (e.g. Noemi's and Kajetan's)
- Here we will just do a little bit of nuclear effect zoology to see what goes into our generators

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### Fermi Motion

- Nucleons are moving targets
- Their momenta are not so different than typical  $E_{\nu}$  for our experiments ...





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#### Nuclear removal energy

- Nucleons are bound inside the nucleus
- Some amount of energy is needed to free them
- Most models predict that removal energy and Fermi motion should be correlated





### Nuclear effects in a nutshell charge Exchange

#### Fermi Motion

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### Final Stat Interactions (FSI)

- Hadrons don't exit the nucleus cleanly
- They can re-interact inside the nucleus
- Distorts kinematics and changes the final state topology
- Full calculation also changes inclusive cross section



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#### Multi-nucleon Interactions

- Nucleons are interacting with each other inside the nucleus
- Some interactions are with nucleons bound together somehow
- o Multi-nucleon "2p2h" final states





#### N. Rocco INSS 2019

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- "long range" interactions between nucleons can act to shield target
- Difficult physics, usually parameterised and treated via "RPA" (random phase approximation)





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### Pauli Blocking

- Nucleons cannot be excited into nuclear states that are already filled
- Reduction of cross section at low energy transfer
   Borro





Borrowed from K. McFarland's INSS 2014 lectures

### Consequences of nuclear effects

#### Altered cross section

 Nuclear effects significantly alter the cross section with respect to the nucleon case

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#### Altered cross section

 Nuclear effects significantly alter the cross section with respect to the nucleon case

2p2h adds a contribution where there previously wasn't any "the dip region"





Fig. from N. Jachowicz

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#### Altered cross section



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#### Altered cross section

 Nuclear effects significantly alter the cross section with respect to the nucleon case

#### Increased dimensionality

Less constrained particle kinematics



CCQE on the nucleon is fully constrained with one kinematic variable



CCQE on the nucleus needs five!

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#### Altered cross section

 Nuclear effects significantly alter the cross section with respect to the nucleon case

#### Increased dimensionality

Less constrained particle kinematics

#### Altered hadronic final state

 Final state interactions hide/distort the interaction channel







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- Only QE for the moment, but with novel FSI
- e/nu-scattering equivalence built in from the ground-up
  - Full theory in its own right
  - Predicts nu/e/hadron scattering in the same framework
  - Very different philosophy to other generators
  - Few developers



- Most generators contain many possible model configurations
- "GENIE" or "NEUT" on a plot does not imply a particular one
- These configs sometimes have ... creative naming schemes
   E.g. GENIE's "G21\_11c\_02\_11b" model
  - This page decodes some of them



- To be used in **experimental analyses**, generators must be able to produce fully exclusive neutrino interactions. I.e.:
  - The full list of final state particles
  - The 4-momentum of each one
  - For all interaction channels



- To be used in experimental analyses, generators must be able to produce fully exclusive neutrino interactions. I.e.:
  - The full list of final state particles
  - The 4-momentum of each one
  - For all interaction channels
- The generators **take theory inputs where possible**, but ultimately ad-hoc approximations to "fill in the gaps" are needed

## Theory inputs

Four broad types of theory inputs to event generators:

- Nucleon-level calculation only
- Inclusive calculations
- Factorized calculations
- Exclusive calculations



## Neutrino-nucleon calculations

The most basic inputs are only neutrinonucleon calculations: no nuclear effects



$$\frac{d\sigma}{d|q^2|} \binom{\nu n \to \ell^- p}{\overline{\nu} p \to \ell^+ n} = \frac{M^2 G^2 \cos^2 \theta_c}{8\pi E_{\nu}^2} \left[ A(q^2) \mp B(q^2) \frac{(s-u)}{M^2} + \frac{C(q^2)(s-u)^2}{M^4} \right]$$
  
(s-u = 4ME<sub>\nu</sub> + q^2 - m^2).

Generators are forced to "dress" the interaction with nuclear effects themselves

#### Luke's python generator does exactly this!

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$$\frac{\mathrm{d}\sigma}{\mathrm{d}|q^2|} \binom{\nu \,\mathrm{n} \to \,\mathbb{Q}^- \,\mathrm{p}}{\overline{\nu} \,\mathrm{p} \to \,^* \,\mathbb{Q}^+ \,\mathrm{n}} = \frac{M^2 \,G^2 \,\cos^2\theta_c}{8\pi E_\nu^2} \left[ A(q^2) \,\overline{+} \,B(q^2) \,\frac{(s-u)}{M^2} + \frac{C(q^2)(s-u)^2}{M^4} \right]$$

$$(s-u = 4ME_\nu + q^2 - m^2) \,.$$

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#### Luke's python generator does exactly this!

#### This is often still the level of input we work with

E.g.:

CCQE GENIEv2 or NEUT's "Smith-Moniz" Fermi gas model All RES interactions in GENIE or NEUT

Inclusive calculations come "preintegrated" over hadron kinematics All of the nuclear dynamics lives in here



E.g. Inclusive quasielastic charged-current neutrino-nucleus reactions, J. Nieves et al, 2004

Nuclear effects are "baked in" to the model used for the integration

Inclusive calculations come "preintegrated" over hadron kinematics

Only predicts lepton kinematics!



 $\frac{d^2\sigma_{\nu l}}{d\Omega(\hat{k}')dE_l'} = \frac{|\vec{k}'|}{|\vec{k}|} \frac{G^2}{4\pi^2} L_{\mu\sigma} W^{\mu\sigma}$ 

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Nuclear effects are "baked in" to the model used for the integration

This is what we have for most 2p2h and some CCQE models

E.g.:

SuSA or Valencia 2p2h SuSAv2 or CRPA in GENIE v3 All of the nuclear dynamics lives in here

Exclusive model: can describe all final state particle kinematics



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Exclusive model: can describe all final state particle kinematics



All of the nuclear dynamics still lives in here

$$\frac{d^5\sigma_{\nu\ell}}{d\Omega(\hat{k}')d\Omega(p_N)dE_{\ell'}} \sim L_{\mu\sigma}W^{\mu\sigma}$$

E.g. Semi-inclusive charged-current neutrino-nucleus reactions, O. Moreno et al, 2014

Some theory models can do this, e.g. Relativistic Mean Field (RMF)

But now this is much more challenging to calculate and to implement in event generators

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No generator does this ...

Exclusive model: can describe all final state particle kinematics



All of the nuclear dynamics still lives in here

 $\frac{d^5\sigma_{\nu\ell}}{d\Omega(\hat{k}')d\Omega(p_N)\mathrm{dE}_{\ell'}}\sim L_{\mu}$ 

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#### No generator does this ...



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## Factorized Calculations



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## Factorized Calculations



• Summarising what the theory inputs give us:

Theory input	What kinematics can I calculate?
Nucleon-level calculation	Lepton and nucleon before FSI
Inclusive calculation	Lepton only
Factorized calculation	Lepton and nucleon before FSI
Exclusive calculation	Lepton and nucleon

\* Possible to include in an ad-hoc way, but doesn't reliably allow for a calculation for alteration of outgoing nucleon kinematics See e.g.: Phys. Rev. D **91**, 033005

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• Summarising what the theory inputs give us:

Theory input	What kinematics can I calculate?	How accurate is the calculation?	
Nucleon-level calculation	Lepton and nucleon before FSI	on Do not trust!	
Inclusive calculation	Lepton only	As accurate as the underlying theory	
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• Summarising what the theory inputs give us:

Theory input	What kinematics can I calculate?	How accurate is the calculation?	FSI/RPA?	Example use in generators
Nucleon-level calculation	Lepton and nucleon before FSI	Do not trust!	Not included	Most nonQE/2p2h + older QE calcs.
Inclusive calculation	Lepton only	As accurate as the underlying theory	Can be included	Most 2p2h, SuSAv2 / CRPA QE in GENIEv3
Factorized calculation	Lepton and nucleon before FSI	Approximations can limit predications	Not without approximations*	SFQE in NEUT, NuWro, AChilLES. Default QE in GENIEv3
Exclusive calculation	Lepton and nucleon	As accurate as the underlying theory	Can be included	Not yet widely available

\* Possible to include in an ad-hoc way, but doesn't reliably allow for a calculation for alteration of outgoing nucleon kinematics See e.g.: Phys. Rev. D **91**, 033005

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# Filling in the gaps



- Generators take theory inputs where possible, but we found these are often limited:
  - Only capable of predicting a subset of observables
  - Only valid within some range of kinematic phase space
  - Only valid for certain processes
- Need to "fill in the gaps" to get to a useable event simulation

# Summary so far

- Weak Interactions with neutrinos
  - Point-like scattering is "easy" to calculate
  - o Interactions with nucleons is more challenging due to their finite extent
- Neutrino-nucleon interactions
  - QE: almost calculable with some form factors
  - RES: much more difficult, lots of diagrams to consider
  - DIS: easy for *inclusive* high Q<sup>2</sup>, hard at low Q<sup>2</sup>, hadronic side a total guess
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  - Nuclear effects: there are lots of them, they **significantly change the cross section**
  - Not all models can predict everything!

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  - Nuclear effects: there are lots of them, they significantly change the cross section
  - Not all models can predict everything!
- Neutrino event generators
  - Many generators on the market, each with different use cases
  - Take theory where possible, but **need to "fill the gaps"** for a complete calculation
  - This limits generators predictive power

# Example: 2p2h

- Theory give us: • (The "inclusive" cross section)
- Exclusive cross section: (what we actually need)
- $d^2\sigma$  $dq_0 dq_3$  $d^8\sigma$



How do we get there!? •

### A generator's view of an interaction



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### A generator's view of an interaction



• Generators need to "dress" our primary interaction with extra physics

### Fermi motion



Arbitrary units

### Final state interactions



• We now have a nucleon inside the nucleus, but it still needs to gets out: **Final State Interactions** 



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## Final state interactions

• Intranuclear cascade models: classical billiard ball scattering within the nucleus







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• Intranuclear cascade models: classical billiard ball scattering within the nucleus





1. Step the particle through the nucleus a distance equal to its mean free path between interactions







- 1. Step the particle through the nucleus a distance equal to its mean free path between interactions
- 2. Check whether it's outside the nucleus, if it is add this particle to the final state and stop FSI for it





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Reactive (TunedFSI) Reactive (OldFSI)

1000 1200 1400

 $\pi^+$  Initial Momentum (MeV/c)

1600

Juasi-elastic

Single CX



. Step the particle through the nucleus a distance equal to its mean free path between interactions

- 2. Check whether it's outside the nucleus, if it is add this particle to the final state and stop FSI for it
- 3. Use MC methods to determine if it interacts or not, if it does choose a process according to its cross section



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a (mb)

450E

400

350

300 250

200 150 100

200

400

600

800

 $\pi^{+} \frac{12}{6}C$ 

• Intranuclear cascade models: classical billiard ball scattering within the nucleus





Absorption



- 1. Step the particle through the nucleus a distance equal to its mean free path between interactions
- 2. Check whether it's outside the nucleus, if it is add this particle to the final state and stop FSI for it
- 3. Use MC methods to determine if it interacts or not, if it does choose a process according to its cross section
- 4. Generate the interaction

• Intranuclear cascade models: classical billiard ball scattering within the nucleus





Charge Exchange



- 1. Step the particle through the nucleus a distance equal to its mean free path between interactions
- 2. Check whether it's outside the nucleus, if it is add this particle to the final state and stop FSI for it
- 3. Use MC methods to determine if it interacts or not, if it does choose a process according to its cross section
- 4. Generate the interaction

• Intranuclear cascade models: classical billiard ball scattering within the nucleus





Pion Production



- 1. Step the particle through the nucleus a distance equal to its mean free path between interactions
- 2. Check whether it's outside the nucleus, if it is add this particle to the final state and stop FSI for it
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- 4. Generate the interaction

• Intranuclear cascade models: classical billiard ball scattering within the nucleus





Elastic scatter



- 1. Step the particle through the nucleus a distance equal to its mean free path between interactions
- 2. Check whether it's outside the nucleus, if it is add this particle to the final state and stop FSI for it
- 3. Use MC methods to determine if it interacts or not, if it does choose a process according to its cross section
- 4. Generate the interaction

• Intranuclear cascade models: classical billiard ball scattering within the nucleus







- 1. Step the particle through the nucleus a distance equal to its mean free path between interactions
- 2. Check whether it's outside the nucleus, if it is add this particle to the final state and stop FSI for it
- 3. Use MC methods to determine if it interacts or not, if it does choose a process according to its cross section
- 4. Generate the interaction
- 5. Return to 1.







- Note that FSI is totally factorised from the rest of the interaction
- Unlike theory-treatments of FSI, cascades don't change the cross section as a function of lepton kinematics

# Example: 2p2h

- Theory give us: • (The "inclusive" cross section)
- Exclusive cross section: (what we actually need)
- $d^2\sigma$  $dq_0 dq_3$  $d^8\sigma$



How do we get there!? •

# Example: 2p2h

- Theory give us: (The "inclusive" cross section)
  - Exclusive cross section: (what we actually need)

 $\frac{d}{dq_0 dq_3 d\boldsymbol{p}_1 d\boldsymbol{p}_2}$ 

 $d^8\sigma$ 

 $d^2\sigma$ 

 $dq_0 dq_3$ 

• How do we get there!?

#### Generate remaining particle kinematics using a best-guess approach

- Sample struck nucleons 4-momenta independently from some spectral function and combine into a 2-nucleon "cluster"
  - Assumption: no correlations between nucleon's momentum/energy
- Give 4-momentum transfer  $(q_0, q_3)$  to the cluster
- "Decay" the cluster to two nucleons
  - Assumption: momentum transfer shared evenly between the nucleons
- Put both nucleons through an FSI cascade
  - Assumption: the FSI model is reasonable
  - Assumption: FSI doesn't change the inclusive cross section

# Example: 2p2h

1.6

uoto1.4 1.2 1

lqng0.8

 $\substack{b_2 \\ b_2 \\ b_3 \\ b_4 \\ c_4 \\ c_5 \\ c_6 \\ c_$ 

0.2

Generator attempts at semi-exclusive cross section



Recent theory calculation of semi-exclusive cross section

#### Generate remaining particle kinematics using a best-guess approach

- Sample struck nucleons 4-momenta independently from some spectral function and combine into a 2-nucleon "cluster"
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  - Assumption: the FSI model is reasonable
  - Assumption: FSI doesn't change the inclusive cross section

### A generator's view of $\nu N$ scattering



- This approach allows generators to produce complete simulations from incomplete theory inputs
- But, does it work?

### A generator's view of $\nu N$ scattering



- · Generators do what they can with what they work with
- But the "gap filling" implies significant approximations which limit their predictive power



• When relying on generators, it's crucial to consider what these approximations are and to assign associated systematic uncertainties

### General rule of thumb





#### Lepton kinematics (except maybe at low energy transfers)

#### Lepton-hadron correlations

Stephen Dolan



Stephen Dolan

# Overview

- Inputs to Generators: Neutrino-nucleon interactions
   QE, RES and DIS
- Inputs to Generators: Neutrino-nucleus interactions

   Nuclear effects
- Neutrino event generators
   Filling in gaps in our inputs
- Benchmarking generators with measurements
   o Inclusive successes and exclusive failures
- Why do we care?
   Neutrino interactions for neutrino oscillations
- Don't Panic! The future of neutrino interaction simulations









#### An experiment's view of an interaction



• See Deborah's talk for details on how these measurements are made!

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### Which observables?



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### Our current models vs data



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### Our current models vs data



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### Forward Angles

- The very forward region is especially sensitive to interactions with low energy transfer ( $\omega$ )
- Things don't look so good here ...





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We describe intermediate muon kinematics in CC0 $\pi$  measurements quite well with most models

```
Expected?
```



We describe intermediate muon kinematics in CC0 $\pi$  measurements quite well with most models

Expected?

Yes!

- Generator approximations are reasonable
- The details of the hadron kinematics don't matter so much
- The impact of FSI is small

We describe forward going muon kinematics in  $CC0\pi$  measurements badly in many models



#### Expected?

Models with RPA do better here

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### We describe forward going muon kinematics in $CC0\pi$ measurements badly in many models



Expected?

Yes!

 Generator approximations for many models are not valid at low momentum transfer

Models with RPA do better here

Makes sense!

• Provides some modelling of physics beyond the generator's approximations

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### Which observables?

Lepton and proton?



Correlations between the muon and proton kinematics allow us to disentangle nuclear effects from neutrino energy

## Generators vs data: a horror story

 No generator can come close to describing global data measuring lepton-hadron correlations

See many more informative generator comparisons in the TENSIONS 2019 report (arXiv:2112.09194)



Stephen Dolan

NuSTEC Summer School, CERN, June 2024

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D 98. 032003 6  $\frac{1}{\sigma} \frac{d\sigma}{d\delta p_{T}}$ TZK NuWro 11q 5 -+- Result 4 LFG RFG 3 SF 2 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1 0 0.2 0 0.4 0.6 0.8  $\delta p_{T} (GeV)$ Phys. Rev. Lett. 121, 022504 ର୍ବୁ ବୃଷ୍ଣ NuWro 11a - Result -10 LFG 3.5 RFG 0.6 3 SF 0.4 2.5 n 2 2 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1.5 1 0.5 0 0.2 0.8 0 0.4 0.6 1.2  $\delta p_{\tau} (GeV)$ 

We describe lepton-nucleon correlations badly

Exclusive or factorized models do better here



### We describe lepton-nucleon correlations badly

Expected?

Yes!

• Most of our models rely on *ad-hoc* model combinations to predict nucleon kinematics

#### Exclusive or factorized models do better here

#### Makes sense!

 Fewer approximations in predictions of nucleon kinematics



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# Pion production measurements

#### Similar story:

- Models generally able to predict lepton kinematics reasonably well
  - Even in the forward region!





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# Pion production measurements

#### Similar story:

- Models generally able to predict lepton kinematics reasonably well

   Even in the forward region!
- But pion kinematics are poorly described across experiments




### So, how did we do?





#### Lepton kinematics (except maybe at low energy transfers)

#### Lepton-hadron correlations

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NuSTEC Summer School, CERN, June 2024

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- Why do we care?
   Neutrino interactions for neutrino oscillations
- Don't Panic! The future of neutrino interaction simulations



- 1. The energy dependence of neutrino cross sections
  - So we know how to extrapolate from our near to far detectors

See Clarence's lectures

- 1. The energy dependence of neutrino cross sections
  - So we know how to extrapolate from our near to far detectors
- 2. The smearing of our neutrino energy reconstruction
  - So we can infer the shape of the oscillated spectrum

See Clarence's lectures

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- 3. Differences in the cross section for  $v_e/v_{\mu}$  (and  $v/\bar{v}$ )
  - So we can use  $v_e$  appearance to probe CP-violation

See Clarence's lectures

- 1. The energy dependence of neutrino cross sections
  - So we know how to extrapolate from our near to far detectors
  - Models differ by 5-10% in the region of interest for DUNE and Hyper-K
  - Significant, given the expected statistics for DUNE and Hyper-K



- 2. The smearing of our neutrino energy reconstruction
  - So we can infer the shape of the oscillated spectrum



- 3. Differences in the cross section for  $v_e/v_\mu$  (and  $v/\bar{v}$ )
  - So we can use  $v_e$  appearance to probe CP-violation



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• Don't Panic! The future of neutrino interaction simulations

## Path to Precision Measurements



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# Undetectable, you say?

"I have done something very bad today by proposing a particle that cannot be detected; it is something no theorist should ever do." *Wolfgang Pauli, 1930* 

# Well, have I got vs for you!

"I have done something very bad today by proposing a particle that cannot be detected; it is something no theorist should ever do." *Wolfgang Pauli, 1930* 

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Phys. Rev. D 104, 092007

Using these criteria, a sample of 4,105,696 interactions was selected. The simulation predicts an average selection efficiency of 64% in the  $p_t$ - $p_{||}$  phase space, where

"I have done something very bad today by proposing a particle that cannot be detected; it is something no theorist should ever do." *Wolfgang Pauli, 1930* 

## Well, have I got vs for you!





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# Tailored electron scattering $e_{4\nu}$



• Our models are becoming more able to make neutrino and electron scattering predictions in the same framework

• New data from CLAS (e-scatting): specifically to help better understand neutrino scattering



## New models, new constraints



- New models, successful in describing electron scattering data, are now being implemented in neutrino interaction simulations
- Such models that describe  $e^-$  and v interactions in the same framework can be directly constrained by precision  $e^-$  data
- New theoretical efforts are allowing models to be more predictive

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### Improved near detectors

#### **DUNE Near Detector**

#### Upgraded T2K/Hyper-K Near Detector



System for moving the LArTPC and tracker up to 30m transverse to the beam



#### Stephen Dolan





Stephen Dolan



Stephen Dolan



Stephen Dolan



Stephen Dolan



Stephen Dolan



Stephen Dolan

# Take-home messages

I think I can safely say that nobody understands neutrino-nucleus interactions

(But an understanding of them is crucial for oscillation experiments)

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(anytime soon, at least for SIS/DIS)

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# Take-home messages

All neutrino event generator models are wrong, but some are more wrong\* than others

G. Orwelll

\*or wrong in different ways

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# Take-home messages

The golden age (of neutrino interaction models and measurements) is before us, not behind us. W. Shakespeare



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### Backups

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## The precision era of $\nu$ oscillations?

#### Latest results

- Indication of CP violation!
- Currently largely limited by statistics ... but not for long!





#### **Current systematic uncertainties**

Source ( <u>TZ</u> R)	$N(v_e)$	
$\sigma_{\!\scriptscriptstyle {\cal V}N}$ and FSI	3.8%	
Total Syst.	5.2%	
NEUTRINO 2022		
Source ( <u>)</u>	$N(v_e)$	
$\sigma_{ u N}$ and FSI	7.7%	
Total Syst.	9.2%	
Dhuc Dov D 08 032012	•	

- Tables show **largest** and **total** syst. uncertainty on samples most sensitive to CP-violation
- Current results have  $\sim 100 v_e$  events, expect **1000-2000** for DUNE/HK

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## The precision era of $\nu$ oscillations?



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### Neutrino nucleon scattering



See Aaron's

lectures



#### Increasing Energy Transfer

## Neutrino-nucleon cross sections

- Discussed neutrino-nucleon interactions
- But it's been a long time since we've measured this process!
- Almost all modern experiments
   use nuclear targets





### **Resonant Pion Production**

CCRES



Current Matrix Elements from a Relativistic Quark Model\*

R. P. Feynman, M. Kislinger, and F. Ravndal

Lauritsen Laboratory of Physics, California Institute of Technology, Pasadena, California 91109 (Received 17 December 1970)

The model's used in today's neutrino experiments are based on an approximate model from the 1970s

gence of the axial-vector current matrix elements. Starting only from these two constants, the slope of the Regge trajectories, and the masses of the particles, 75 matrix elements are calculated, of which more than  $\frac{3}{4}$  agree with the experimental values within 40%. The prob-

ficing theoretical adequacy for simplicity. We shall choose a relativistic theory which is naive and obviously wrong in its simplicity, but which is definite and in which we can calculate as many things as possible – not expecting the results to agree exactly with experiment. but to see how closely our "shadow of the truth" equation gives a partial reflection of reality. In our attempt to maintain simplicity, we shall evidently have to violate known principles of a complete relativistic field theory (for example, unitarity). We shall attempt to modify our calculated results in a general way to allow, in a vague way, for these errors.

The model includes its own form factors, including an axial part with an analogous  $M_A$  (and an additional uncertainty in the form factor numerator)  $f_A(q^2) = \frac{f_A(0)}{\left(1 - \frac{q^2}{M_{\star}^2}\right)^2}$ 

Theoretical developments are underway but it's safe to say CCRES is less well understood than CCQE!

•

## **DIS-RES Transition Region**

- There is no cut off where we better describe interactions in a DIS framework compared to In a RES framework
- In general we use models that extrapolate between regions which are definitely DIS (e.g. W>5 GeV) and that are definitively RES (e.g. W<2 GeV)</li>
- Different simulations use 4000 different ad-hoc methods PYTHIA KNO 3500 Transitio - Total of dealing with this 3000 Quasi-elastic Resonance 2500 But this is a region that will DIS 2000 F be important for DUNE! 1500F 1000 500 Invariant mass W (GeV/c<sup>2</sup>)

### Example: nucleon axial mass "puzzle"

- Some heavier nuclear target experiments also try to measure  $M_A$
- The MiniBooNE experiment (carbonbased target) prefers a much higher  $M_A$  to the bubble chambers




• What MiniBooNE really measured wasn't CCQE, they just looked for interactions with no mesons in the final state



#### Stephen Dolan N

- What MiniBooNE really measured wasn't CCQE, they just looked for interactions with no mesons in the final state
- This should include contributions from 2p2h (and FSI with pion absorption)!





#### Stephen Dolan

- What MiniBooNE really measured wasn't CCQE, they just looked for interactions with no mesons in the final state
- This should include contributions from 2p2h (and FSI with pion absorption)!



**Stephen Dolan** 

#### **Dramatic Conclusion**

 For the first time, we have multiple observables pointing to a two body current





NuInt 2015 Summary by Kevin McFarland  It's time to say goodbye to M<sub>A</sub><sup>effective</sup>



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#### Inclusive calculations

Inclusive calculations come "preintegrated" over hadron kinematics

Only predicts lepton kinematics!



 $\frac{d^2\sigma_{\nu l}}{d\Omega(\hat{k}')dE_l'} = \frac{|\hat{k}'|}{|\hat{k}|}\frac{G^2}{4\pi^2}L$ 

E.g. Inclusive quasielastic charged-current neutrino-nucleus reactions, J. Nieves et al, 2004



All of the nuclear dynamics lives in here

In some calculations, the nuclear effects considered includes the impact of Final State Interactions (FSI) with a QM treatment  $\mathsf{RMF}\text{-}\mathsf{FSI}\text{:}\mathsf{Scattered}$  nucleon w.f. is solution of Dirac eq. in presence of the same potentials used to describe the bound nucleon w.f.

Like this, FSI changes the matrix element

Affects cross section as a function of lepton and hadron kinematics!

#### Inclusive calculations

Inclusive calculations come "preintegrated" over hadron kinematics

Only predicts lepton kinematics!



 $\frac{d^2\sigma_{\nu l}}{d\Omega(\hat{k}')dE_l'} = \frac{|\hat{k}'|}{|\hat{k}|}\frac{G^2}{4\pi^2}L_{\mu}$ 

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In some calculations, the nuclear effects considered **includes the impact of Final State Interactions** (FSI) with a QM treatment

 $\mathsf{RMF}\text{-}\mathsf{FSI}\text{:}$  Scattered nucleon w.f. is solution of Dirac eq. in presence of the same potentials used to describe the bound nucleon w.f.

#### FSI like this is included in QE models, but not 2p2h

SuSA or Valencia 2p2h – no consideration of FSI SuSAv2 or CRPA in GENIE v3 – impact of FSI on inclusive cross section is considered

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- Randomly select  $E_{\nu}$  based on the product of an input flux and total  $\sigma(E_{\nu})$  model

— CC Inclusive



Example for if we generate only CC events

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— CC Inclusive



Example for if we generate only CC events

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• Randomly select interaction channel based on their cross sections for the chosen  $E_{\nu}$ 



• Randomly select interaction channel based on their cross sections for the chosen  $E_{\nu}$ 



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• Randomly select interaction channel based on their cross sections for the chosen  $E_{\nu}$ 



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• Select **outgoing particle kinematics** according to differential cross section for the chosen interaction channel at the chosen  $E_{\nu}$ 



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• Select **outgoing particle kinematics** according to differential cross section for the chosen interaction channel at the chosen  $E_{\nu}$ 



Stephen Dolan

- Select **outgoing particle kinematics** according to differential cross section for the chosen interaction channel at the chosen  $E_{\nu}$ 
  - **Nucleon-level input**: for CCQE we sample in 1 dimension (more for nonQE):  $d\sigma/dQ^2$
  - Inclusive input: for any interaction channel we sample in 2, e.g.:  $d\sigma/dq_0dq_3$
  - o **Factorized/exclusive input**: for CCQE we sample in 5 dimensions (more for nonQE)



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- Select **outgoing particle kinematics** according to differential cross section for the chosen interaction channel at the chosen  $E_{\nu}$ 
  - **Nucleon-level input**: gives us lepton kinematics in the struck nucleon rest-frame
  - o Inclusive input: gives us lepton kinematics in the lab frame
  - **Factorized/exclusive input**: lepton + pre-FSI nucleon kinematics in the lab frame



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# A bright future for Argon

Short Baseline Program: Fermilab liquid Argon detectors in "Booster" beam (~0.8 GeV)



- MicroBooNE: already producing interesting results
- ICARUS: taking physics data
- SBND: enormous event rates coming soon (1M  $\nu$ /y)

Beyond SBN:

• DUNE "2x2" prototype: measurements at DUNE energies

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