

Impact of neutrino interaction uncertainties on oscillation measurements

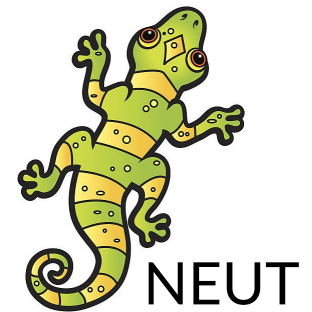
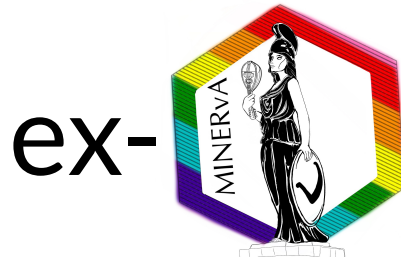
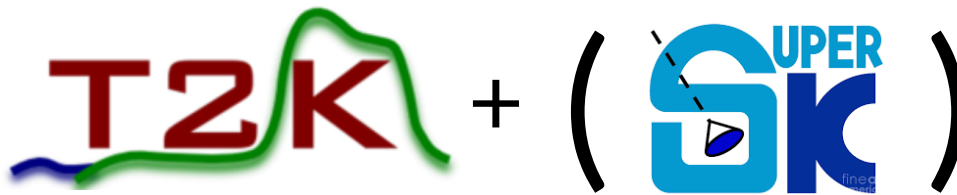
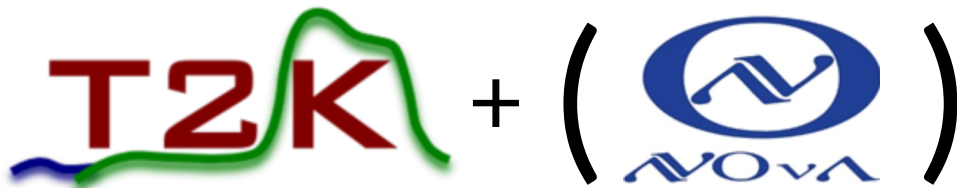
Clarence Wret
June 12 2024
NuSTEC school, CERN

Impact of neutrino interaction uncertainties on oscillation measurements

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Bias declarations:



Structure

- Recap of neutrino oscillations
 - What are we looking for and how?
 - How big are the effects?
- The role of the near detector
- Energy estimators
- What else can go wrong?

Neutrino oscillations 101

Neutrino oscillations 101

- Neutrino flavour and mass eigenstates are separated

$$|\nu_i\rangle = \sum_{\alpha}^n U_{\alpha i} |\nu_{\alpha}\rangle$$

Mass state α Flavour state
Mixing matrix

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix}$$

Neutrino oscillations 101

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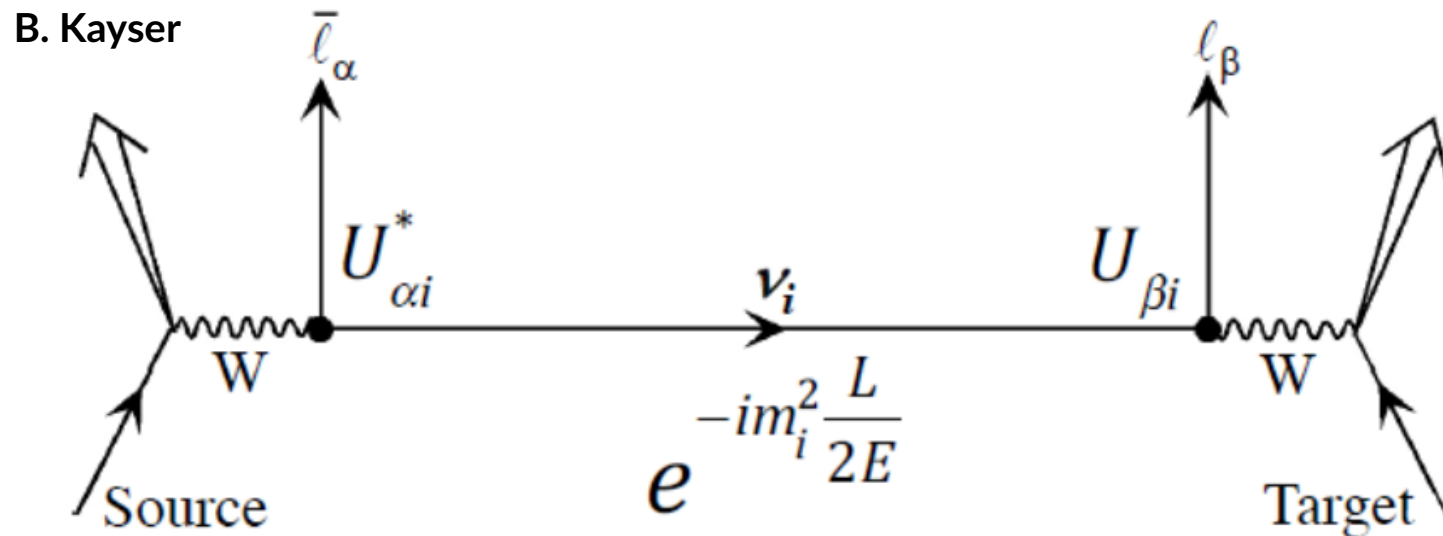
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Mass state ν_i $U_{\alpha i}$ Flavour state ν_{α}

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- Neutrinos propagate in **mass eigenstates**, but are born and detected in the **flavour eigenstate** via weak interaction



- Results in **oscillations** of the detected flavour eigenstates

Neutrino oscillations 101

- Express probability to detect a neutrino with flavour α and energy E , as flavour β after it's travelled distance L

$$P(\nu_\alpha \rightarrow \nu_\beta) = \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re} \left(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^* \right) \sin^2 \left(\Delta m_{ij}^2 \frac{L}{4E} \right)$$

$$\Delta m_{ij}^2 = m_i^2 - m_j^2 + (-)2 \sum_{i>j} \text{Im} \left(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^* \right) \sin \left(\Delta m_{ij}^2 \frac{L}{2E} \right)$$

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Mixing angles

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Mixing angles

Mass² difference between eigenstate i and j

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Neutrino oscillations 101

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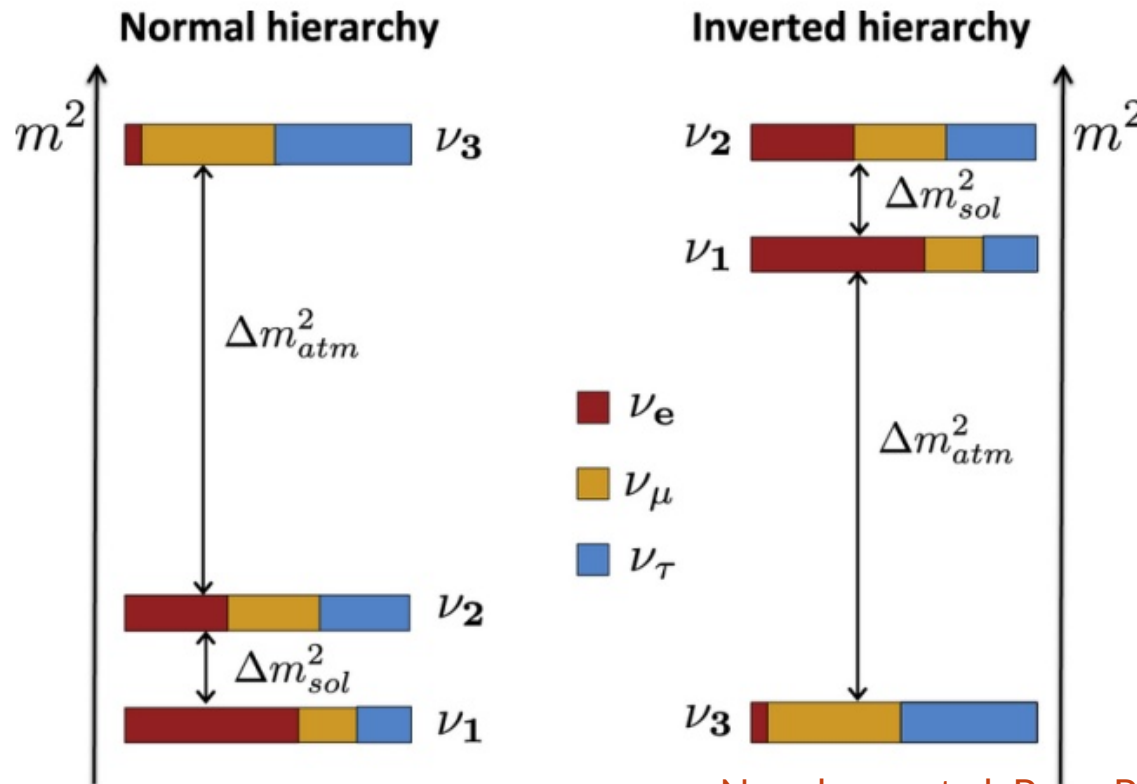
- Design of a neutrino oscillation experiment focusses on L/E
 - Determines sensitivity to mass squared splitting and mixing angles
 - Optimise L/E to match appearance/disappearance
 - Resolve neutrino energy adequately

Neutrino oscillations 101

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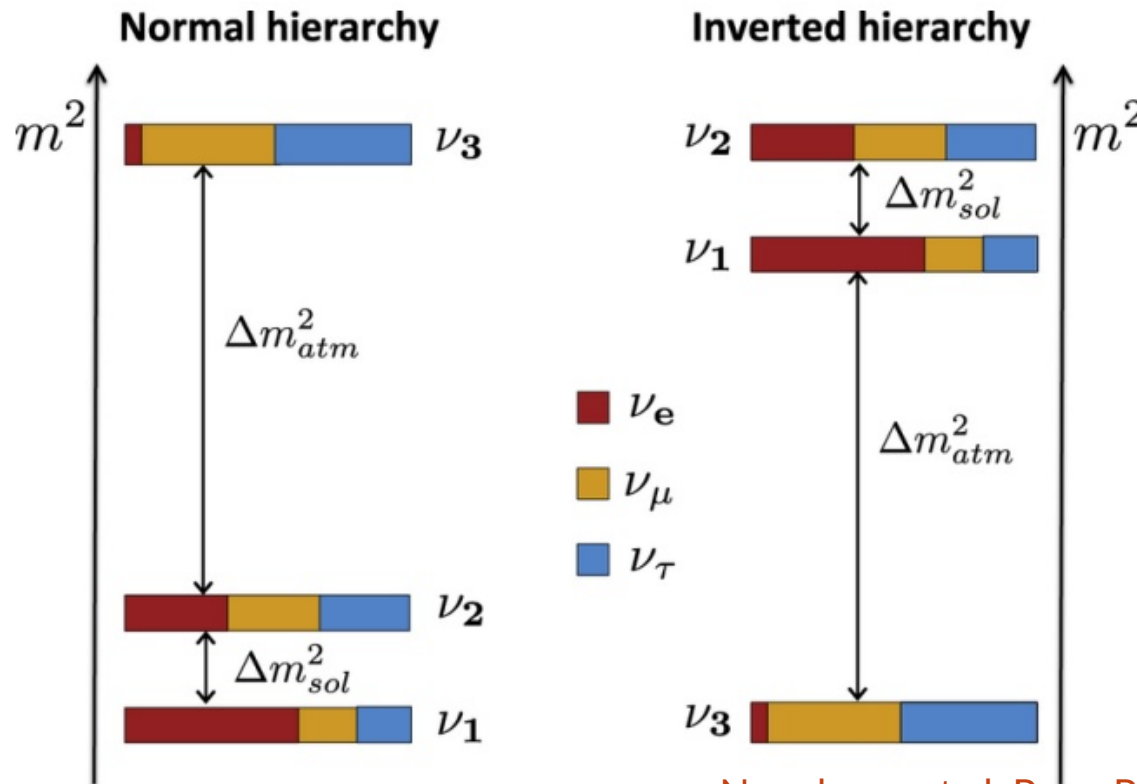
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Dominant effect from \sin^2 : to a unknown mass ordering: $\Delta m_{32}^2 > 0$?



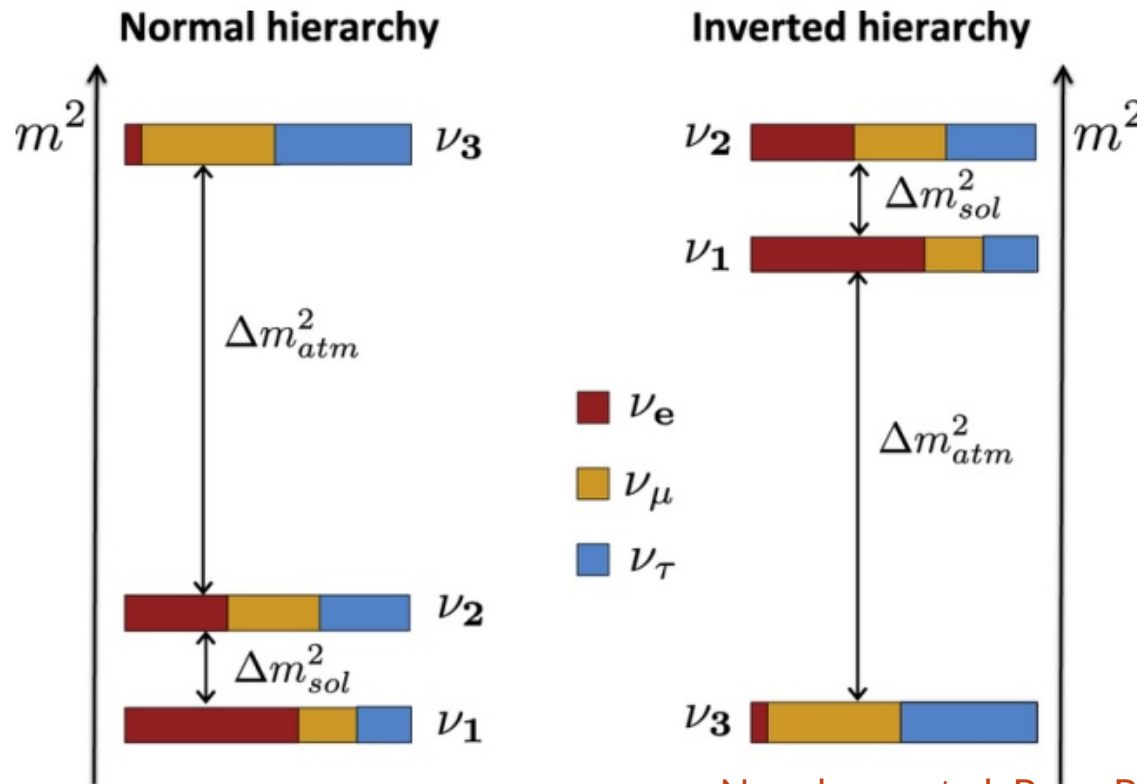
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\sin term resolves mass ordering, through second order

Know $\Delta m_{21}^2 > 0$ from SNO experiment

Neutrino oscillations 101

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Measure differences in $P(\nu_\mu \rightarrow \nu_e)$ and $P(\text{anti-}\nu_\mu \rightarrow \text{anti-}\nu_e)$
→ left with **single term**

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$$\Delta_{ij} \equiv \Delta m_{ij}^2 L / 4E$$

$$P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) = -16 J_{\alpha\beta} \sin \Delta_{12} \sin \Delta_{23} \sin \Delta_{31}$$

Sensitive to
CP violating phase

Sensitive to
mass ordering

$$J \equiv s_{12} c_{12} s_{23} c_{23} s_{13} c_{13}^2 \sin \delta$$

Neutrino oscillations 101

- But that was all in a **vacuum!**
- When **electron neutrinos** propagate through matter, they experience a different potential to the other flavours

$$\begin{aligned}
 P(\nu_\mu \rightarrow \nu_e) = & \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2(\Delta_{31} - aL)}{(\Delta_{31} - aL)^2} \Delta_{31}^2 \\
 & + \sin 2\theta_{23} \sin 2\theta_{13} \sin 2\theta_{12} \frac{\sin(\Delta_{31} - aL)}{(\Delta_{31} - aL)} \\
 & \times \Delta_{31} \frac{\sin(aL)}{aL} \Delta_{21} \cos(\Delta_{31} + \delta) \\
 & + \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(aL)}{(aL)^2} \Delta_{21}^2,
 \end{aligned}$$

(leading order calculation)

$$a \equiv G_F N_e / \sqrt{2}$$

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- For electron anti-neutrinos: $a \rightarrow -a$ and $\delta \rightarrow -\delta$
- Matter effect produces a difference between $P(\nu_\mu \rightarrow \nu_e)$ and $P(\text{anti-}\nu_\mu \rightarrow \text{anti-}\nu_e) \rightarrow$ **Same as CP violation signature**

Neutrino oscillations

- The most general form of mixing matrix is seldom used; instead separate into three mixing matrices

$$s_{ij} = \sin\theta_{ij}$$

$$c_{ij} = \cos\theta_{ij}$$

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

*Atmospheric or
"2,3" sector*

Reactor, or "1,3" sector

*Solar, or "1,2"
sector*

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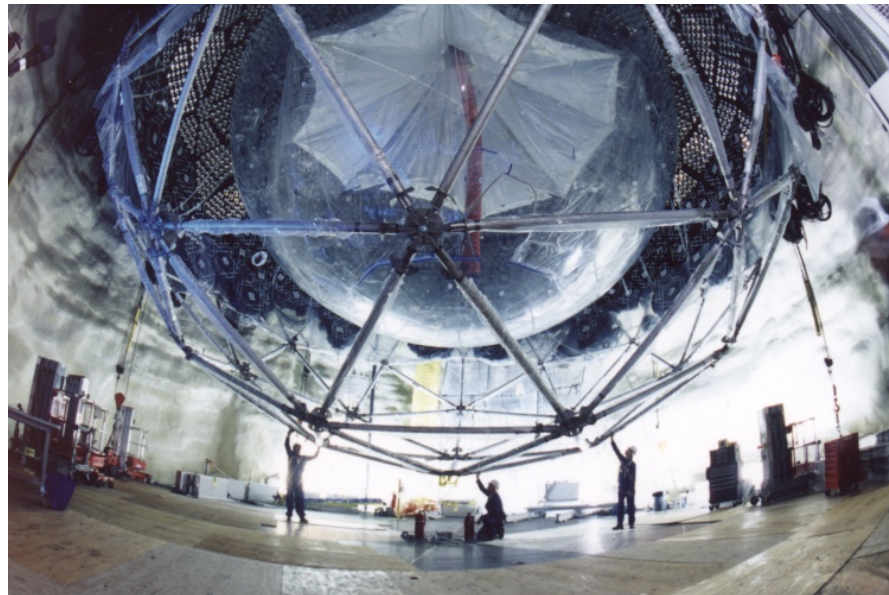
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Reactor, or "1,3" sector

Solar, or "1,2"
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Solar experiments (SNO, SK)
long baseline reactor
experiments (KamLAND,
JUNO)

L/E > 100km/MeV

From MIT

Neutrino oscillations

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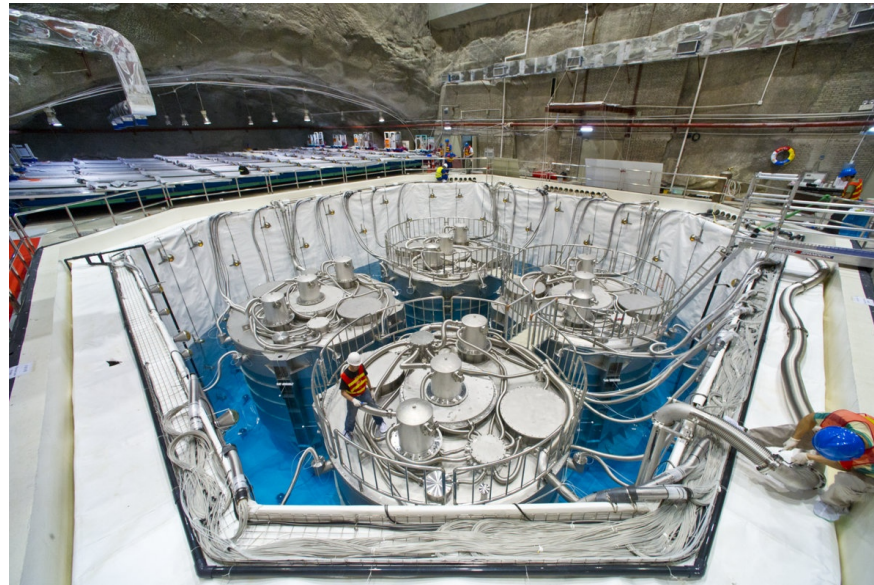
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$$s_{ij} = \sin\theta_{ij}$$

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Reactor experiments (Daya Bay, RENO, Double Chooz)

L/E ~ 1km/MeV



From LBL

Neutrino oscillations

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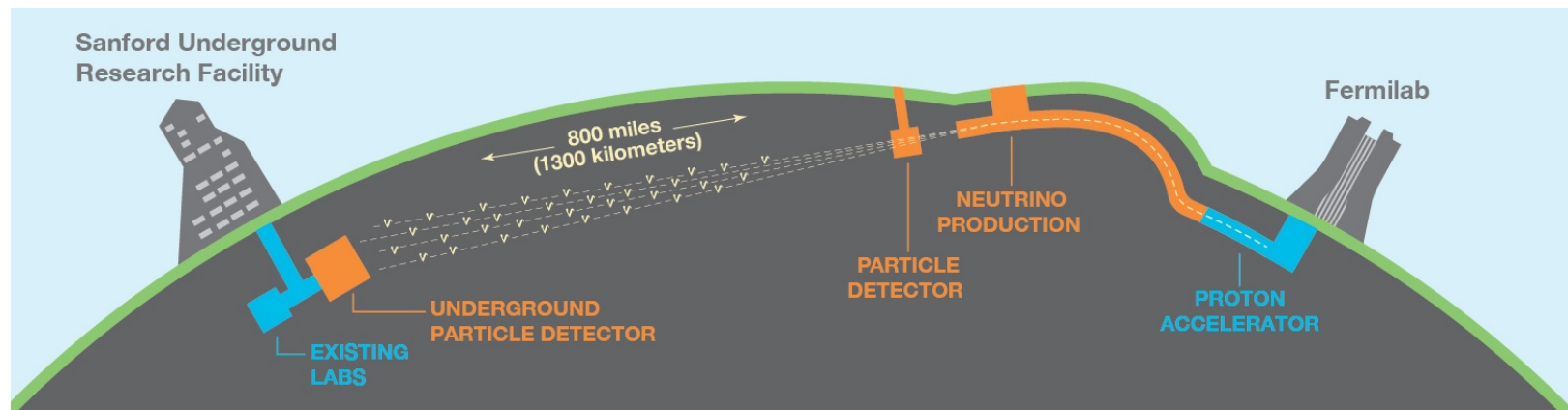
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Reactor, or "1,3" sector
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Long baseline experiments (K2K, T2K, NOvA, MINOS, DUNE, HK),
atmospheric experiments (SK, IceCube)

L/E ~ 400-500km/GeV



From DUNE

Neutrino oscillations

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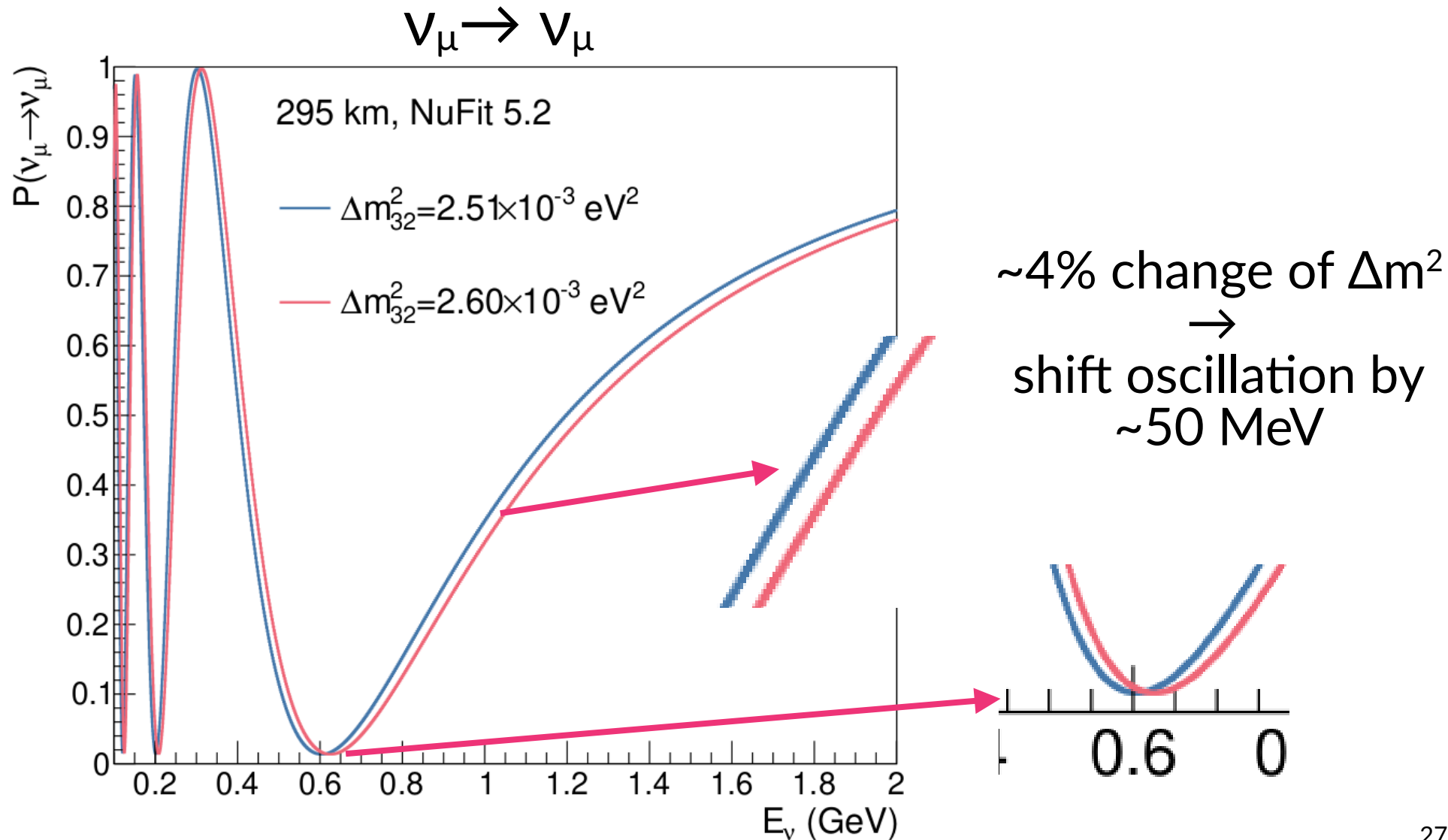
$L/E \sim 400\text{-}500\text{km/GeV}$

The focus of these lectures

From DUNE

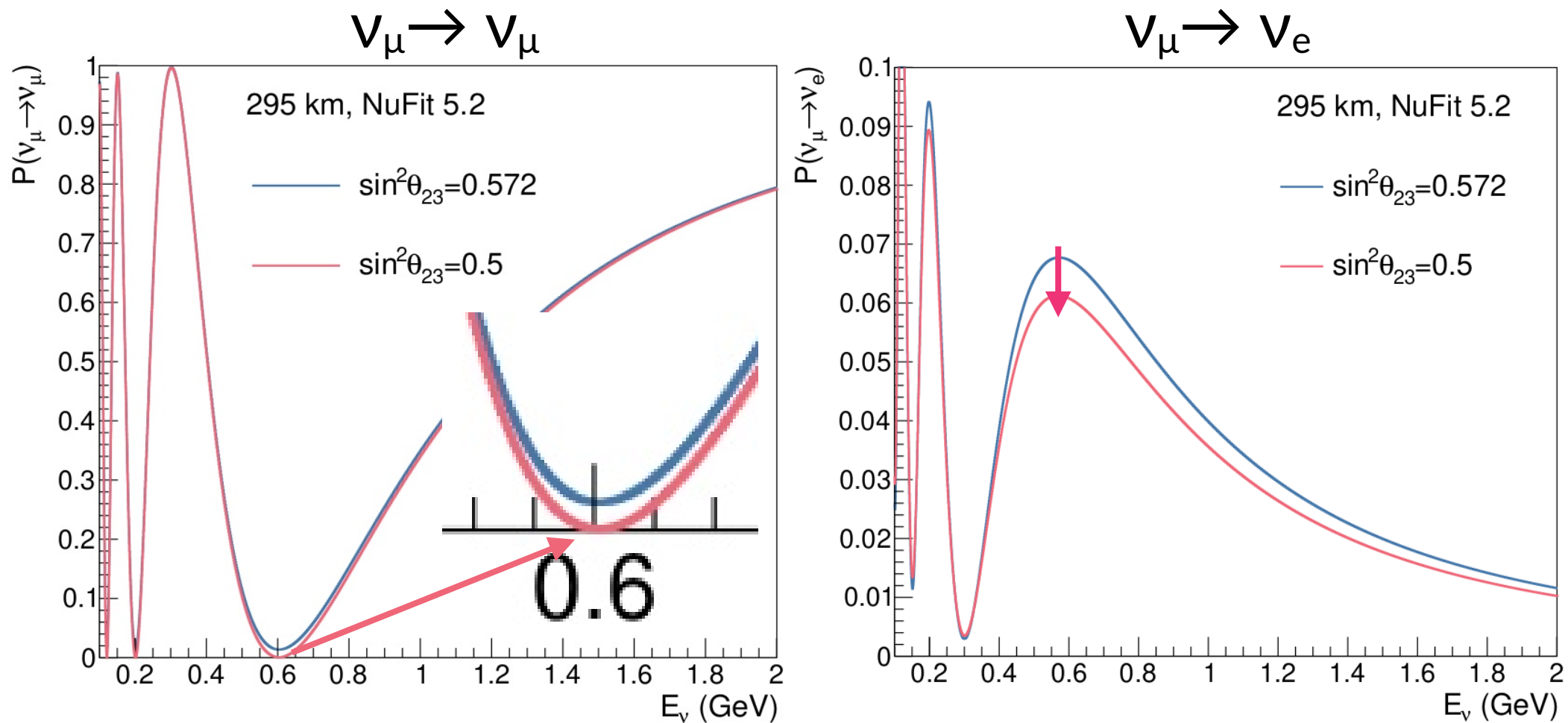
How big are the effects?

- Varying **mass-squared splitting** to see impact on muon neutrino oscillation probability
- **Induces a shift in energy around the main oscillation dip**



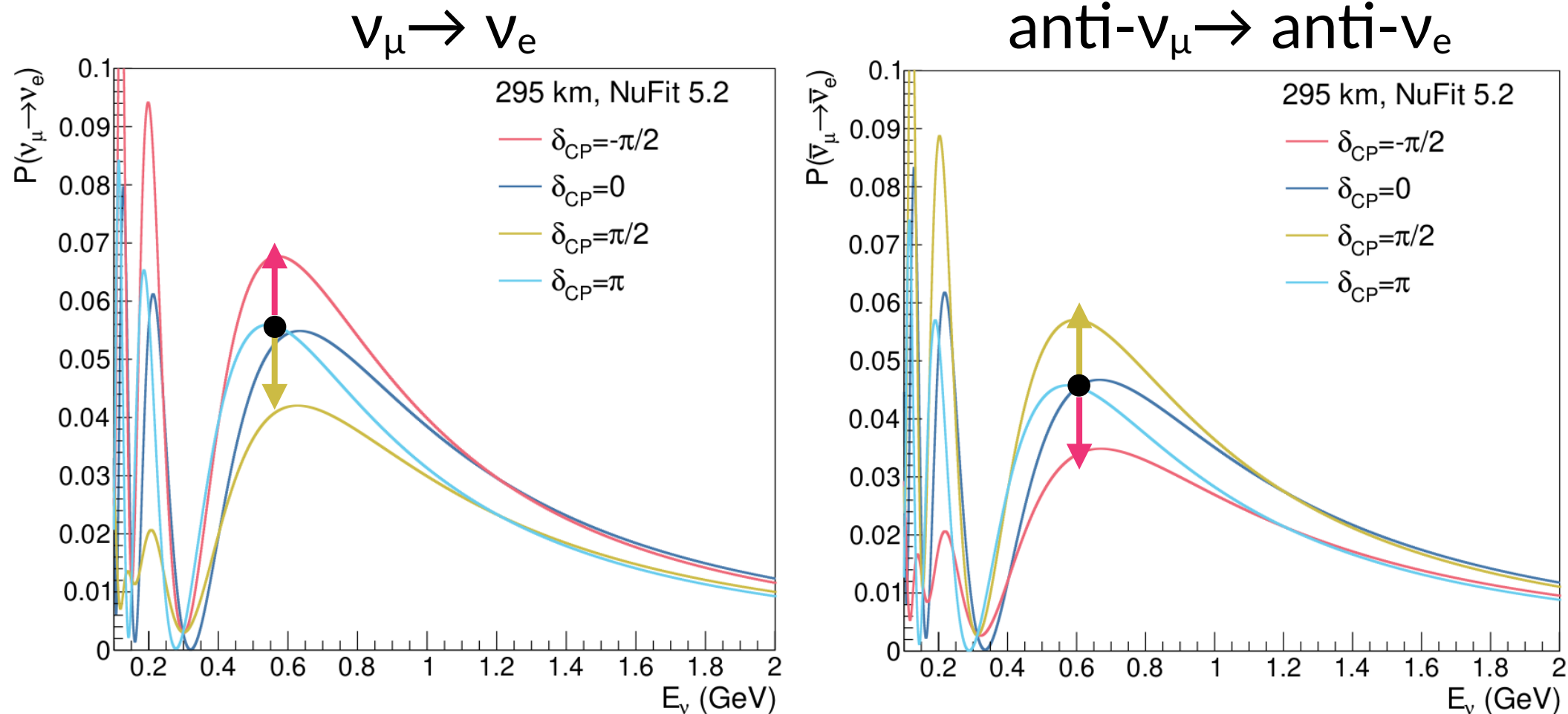
How big are the effects?

- Move from NuFit 5.2 to $\sin^2\theta_{23} = 0.5$ \rightarrow decrease probabilities for both flavours (increase $\nu_\mu \rightarrow \nu_\tau$ probability)
- **Overall decrease in normalisation, especially in dip region**



How big are the effects?

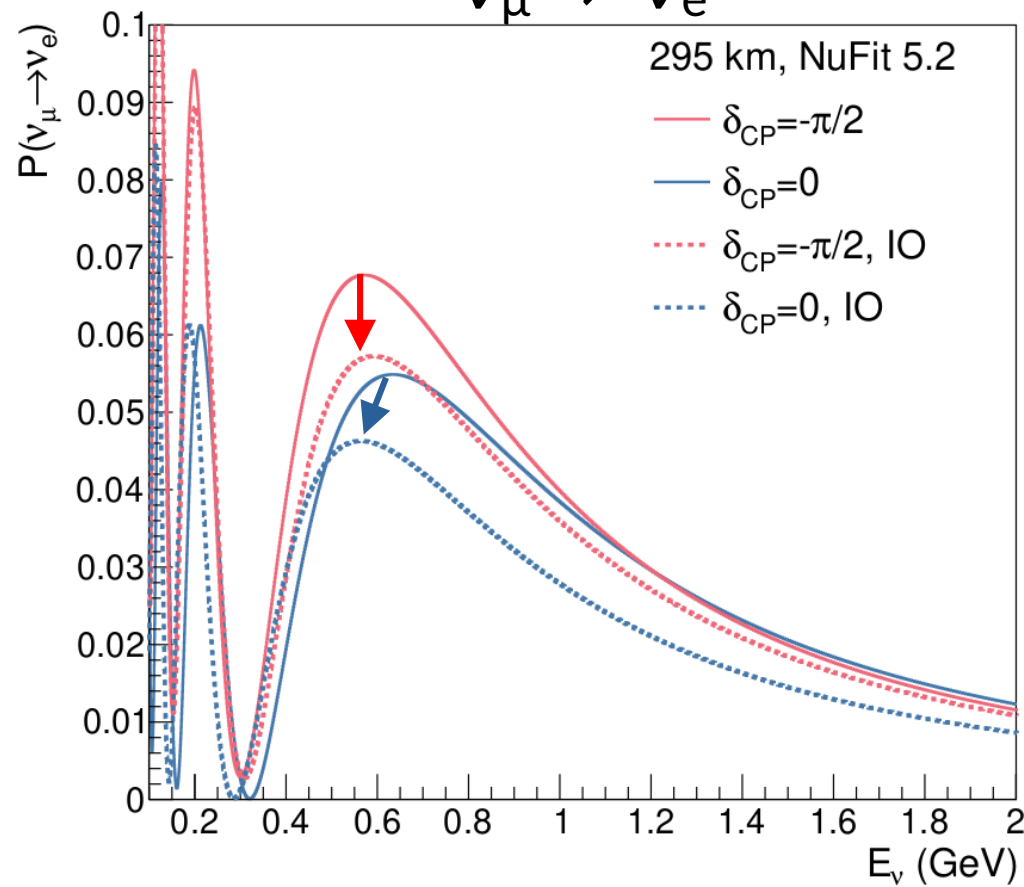
- Changing δ_{CP} cyclically from maximum to minimum effect, through the two CP-conserving points $\delta_{CP}=0, \pi$
- Opposite effect for electron neutrinos and anti-neutrinos



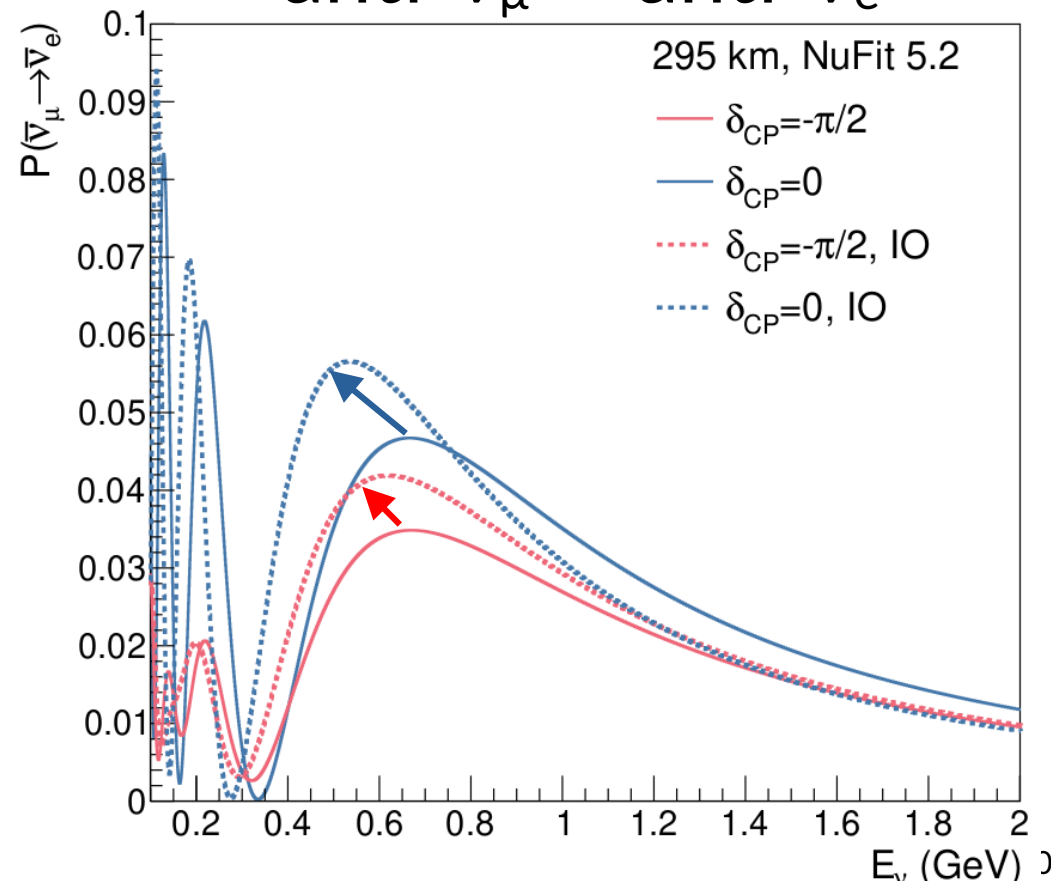
How big are the effects?

- Changing the mass ordering (NO, IO) and δ_{CP} from 0 to $-\pi/2$
- **Opposite effect** for electron neutrinos and anti-neutrinos
- **Degeneracy:** NO \rightarrow IO decreases electron neutrino; increases electron anti-neutrino. But, shape of spectrum changes
- $\delta_{CP}=0$, NO very similar to $\delta_{CP}=-\pi/2$, IO for neutrinos

$\nu_\mu \rightarrow \nu_e$

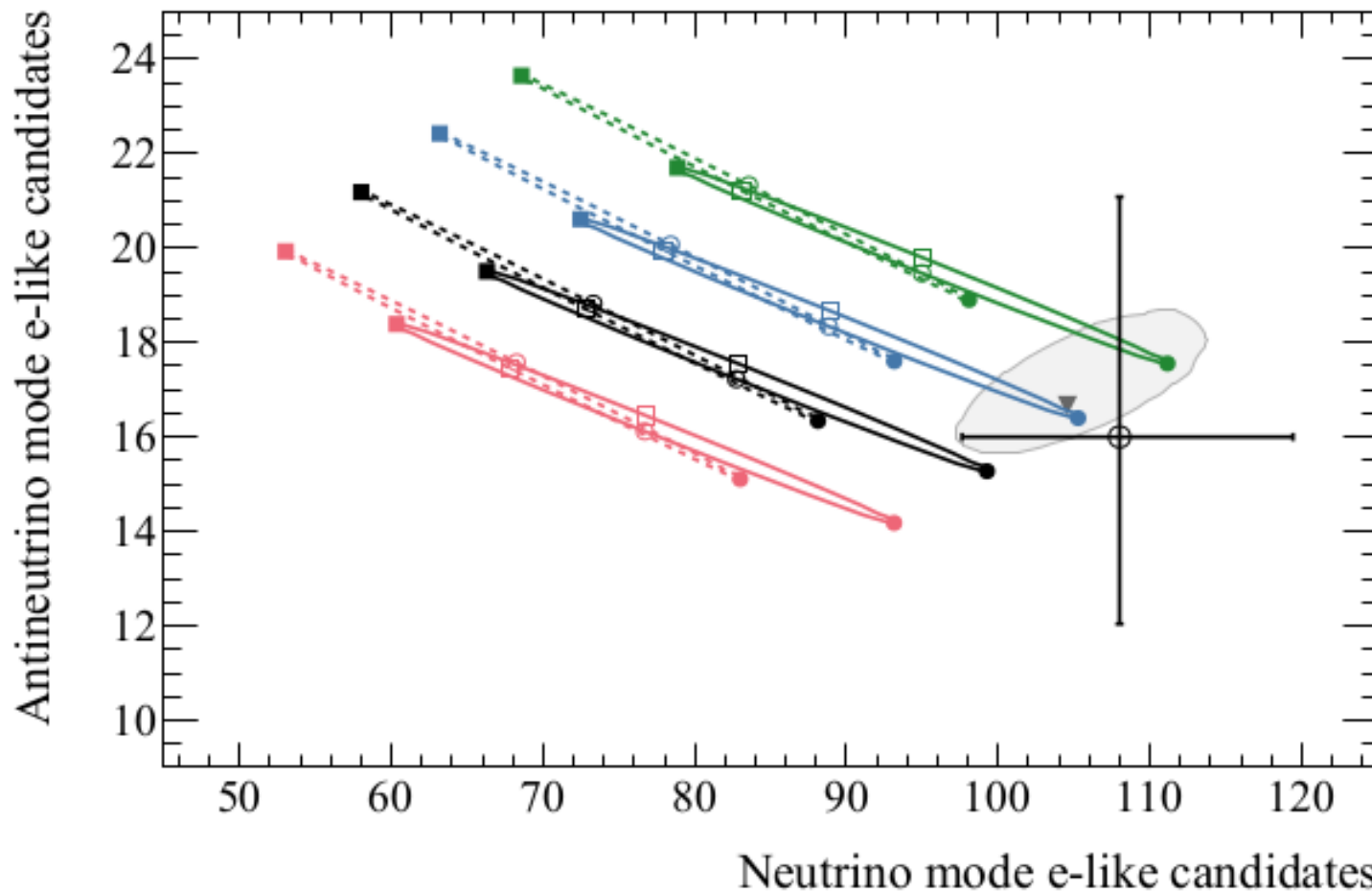
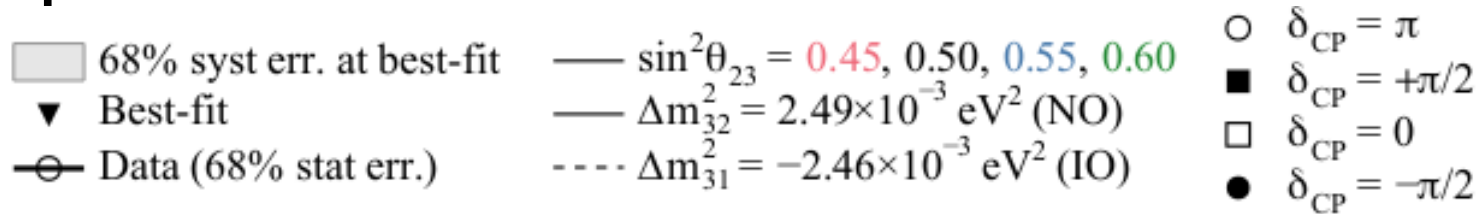


anti- $\nu_\mu \rightarrow$ anti- ν_e



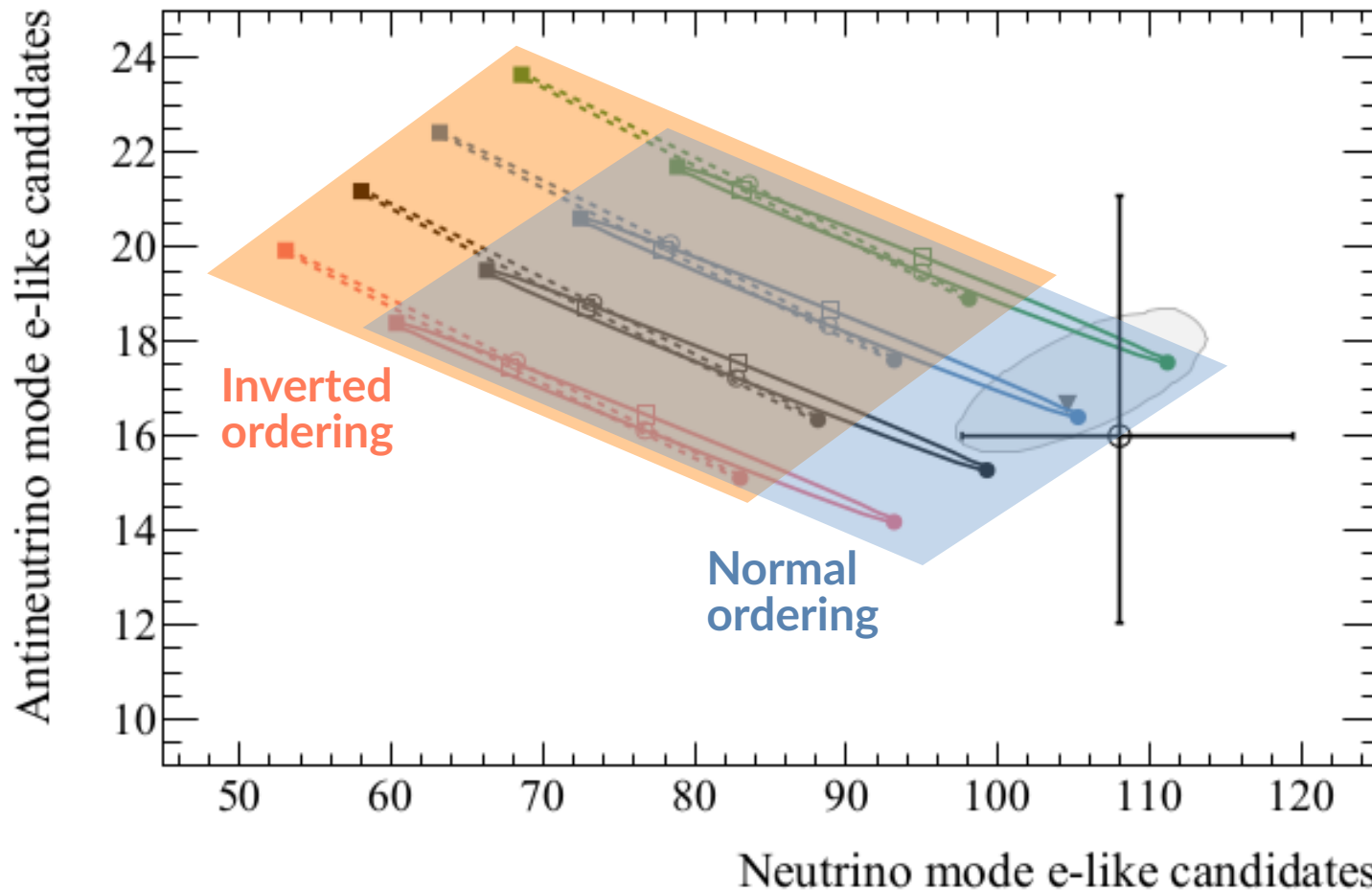
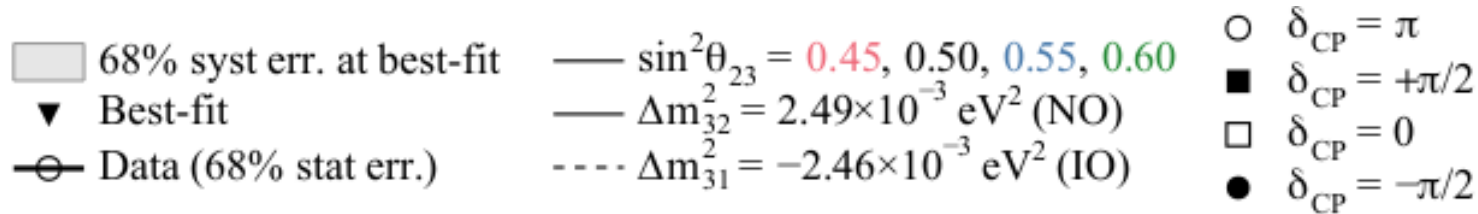
Exploring the degeneracies

- The earlier features are often summarised in “bi-event plots”



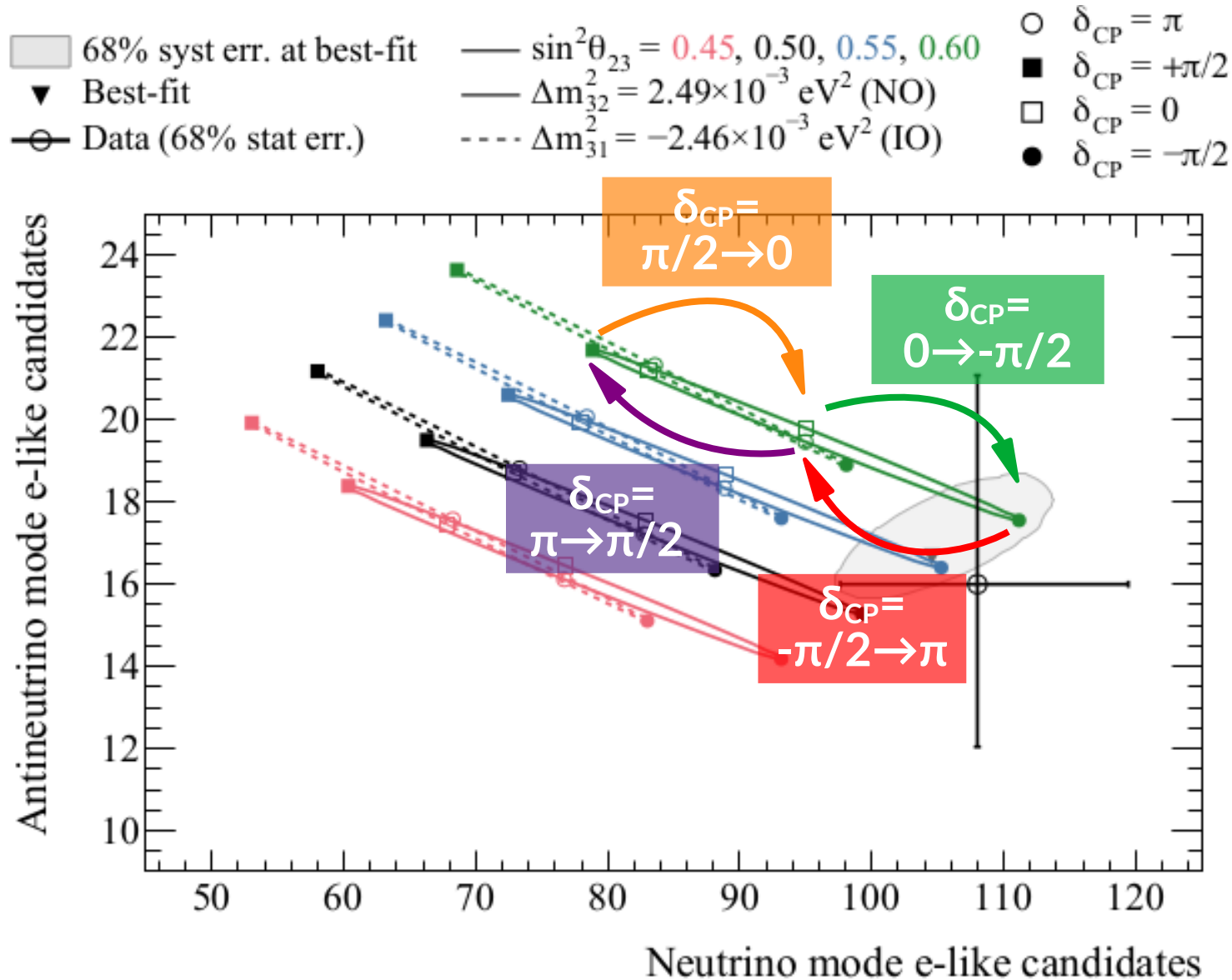
Exploring the degeneracies

- Separate by mass ordering scenarios



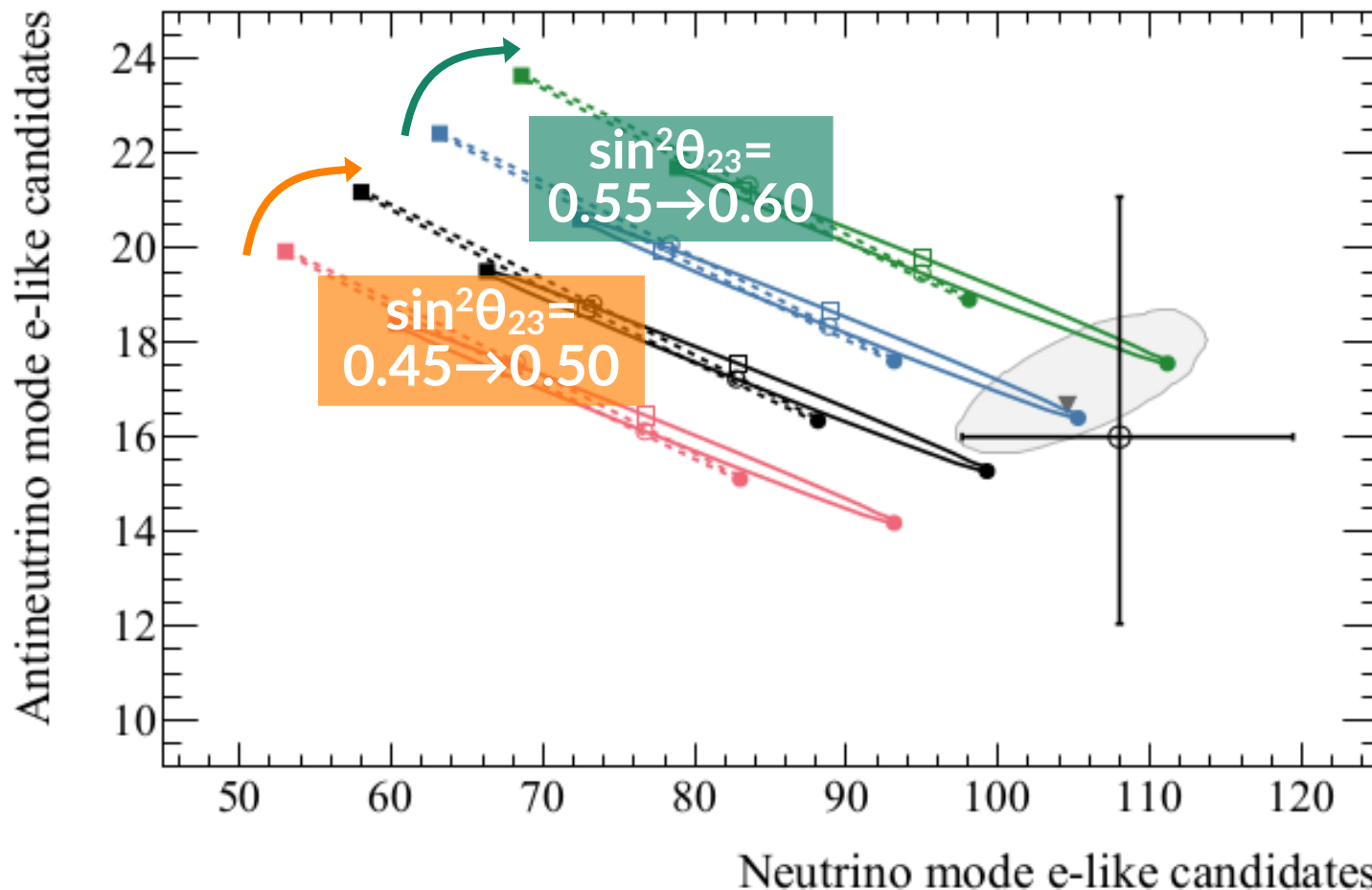
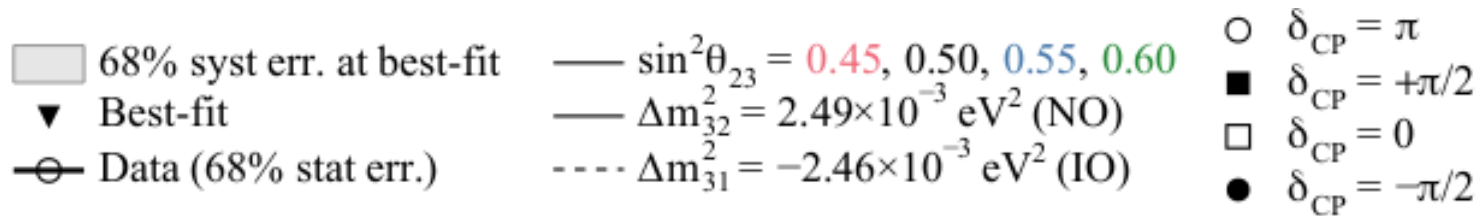
Exploring the degeneracies

- Separate by CP violating phase scenarios



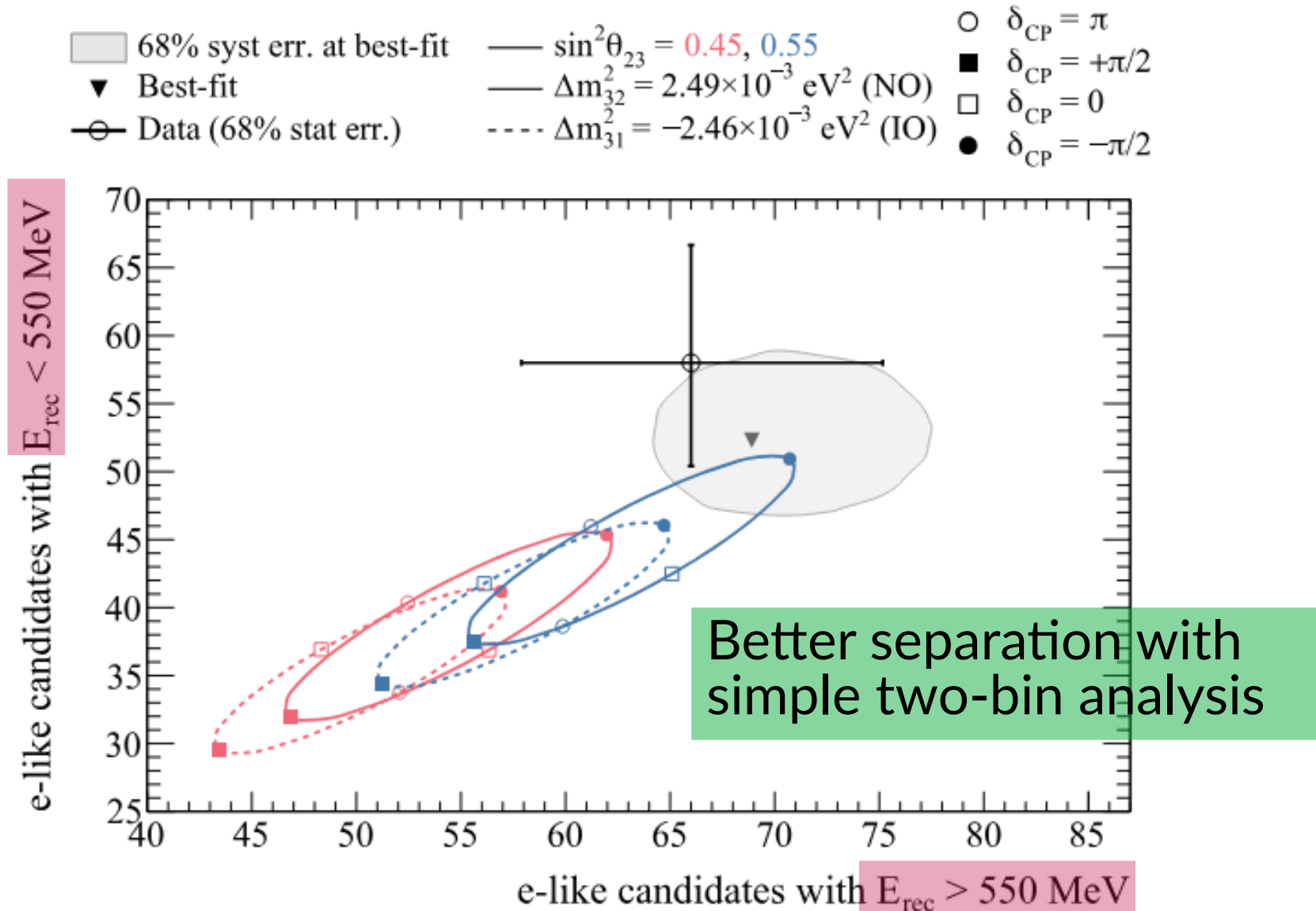
Exploring the degeneracies

- Separate by $\sin^2\theta_{23}$



Exploring the degeneracies

- But, these don't tell full story: **they ignore energy dependence** (simple counting experiment)

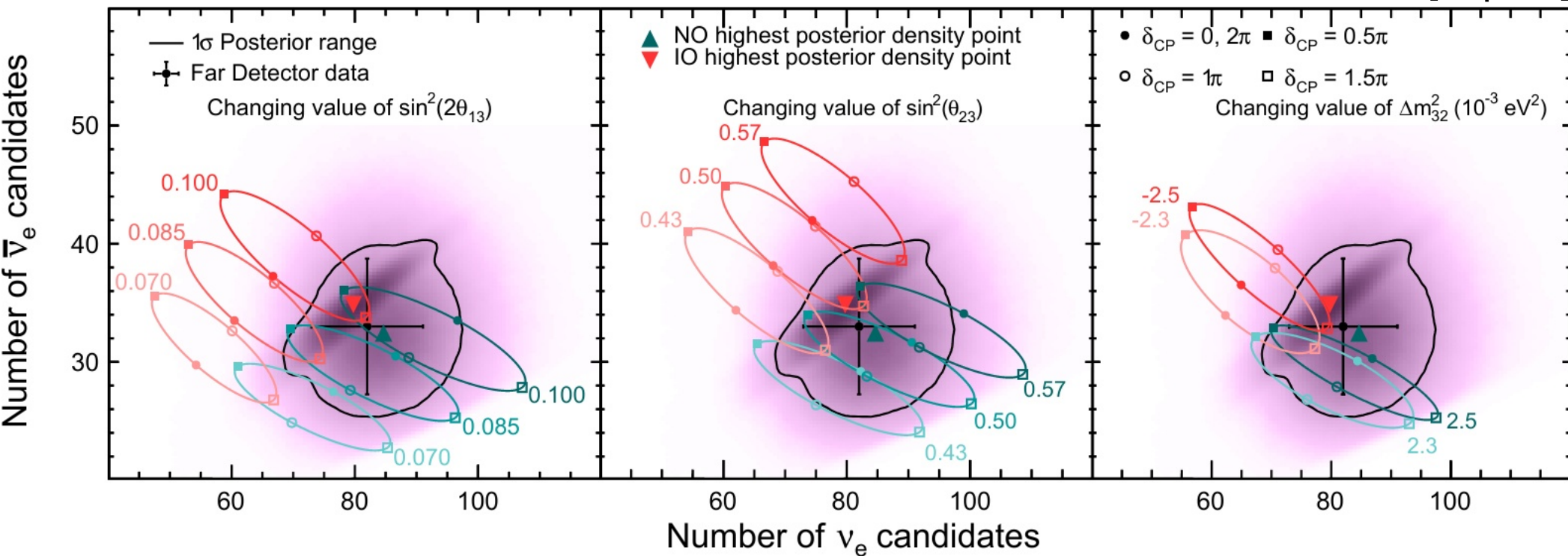


Exploring the degeneracies

- NOvA experiment has higher neutrino energy, and longer baseline compared to T2K
 - Stronger mass ordering sensitivity, weaker δ_{CP} sensitivity

Both Orderings

2311.07835 [hep-ex]

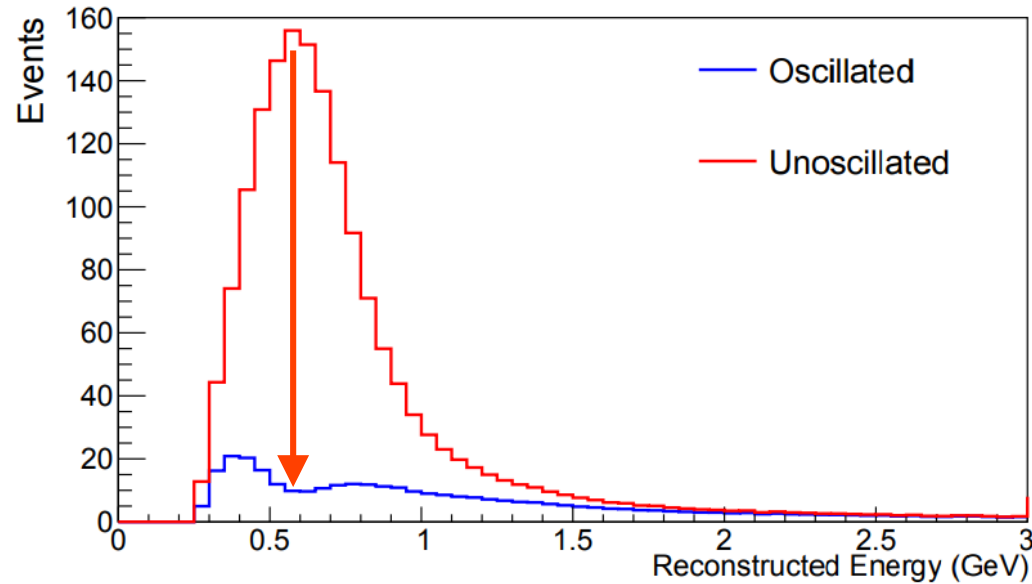


- Larger separation of δ_{CP} and mass ordering effects
- (the different sensitivity to δ_{CP} and MO makes joint T2K+NOvA fit very interesting, amongst other things)

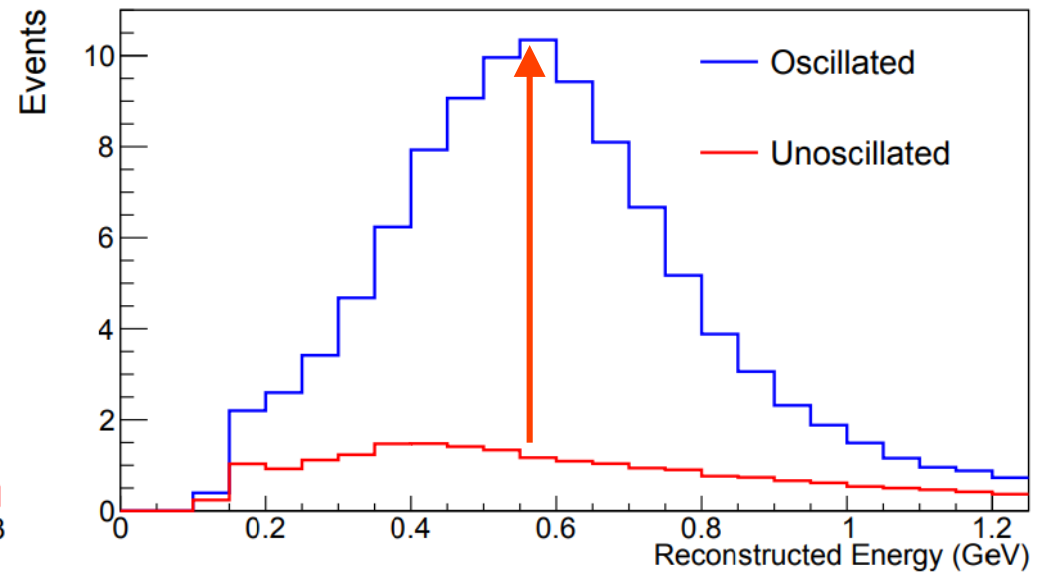
Introduction

- Oscillation parameters change the rate and shape of the appearing and disappearing neutrinos

T2K FHC 1R μ



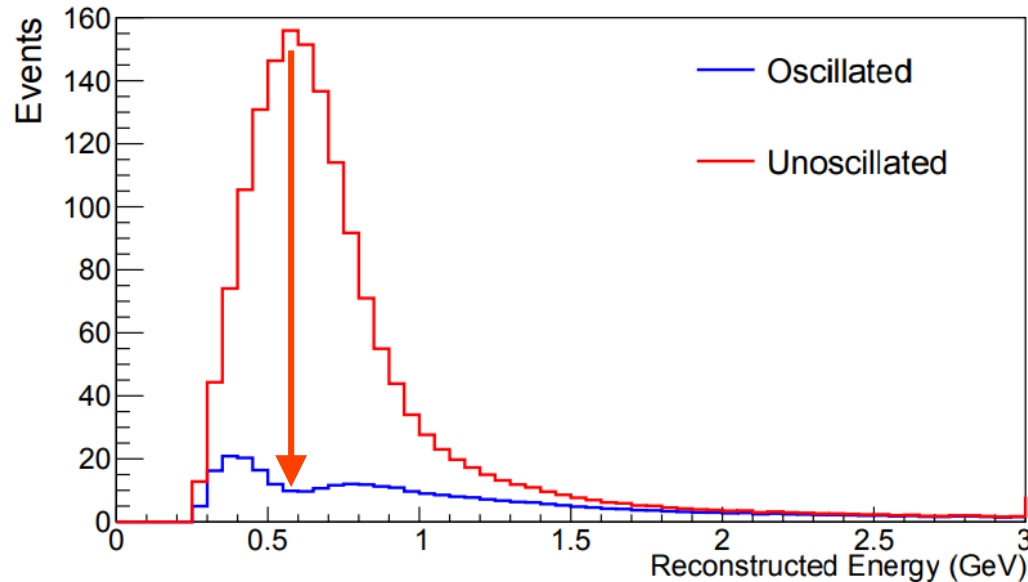
T2K FHC 1Re



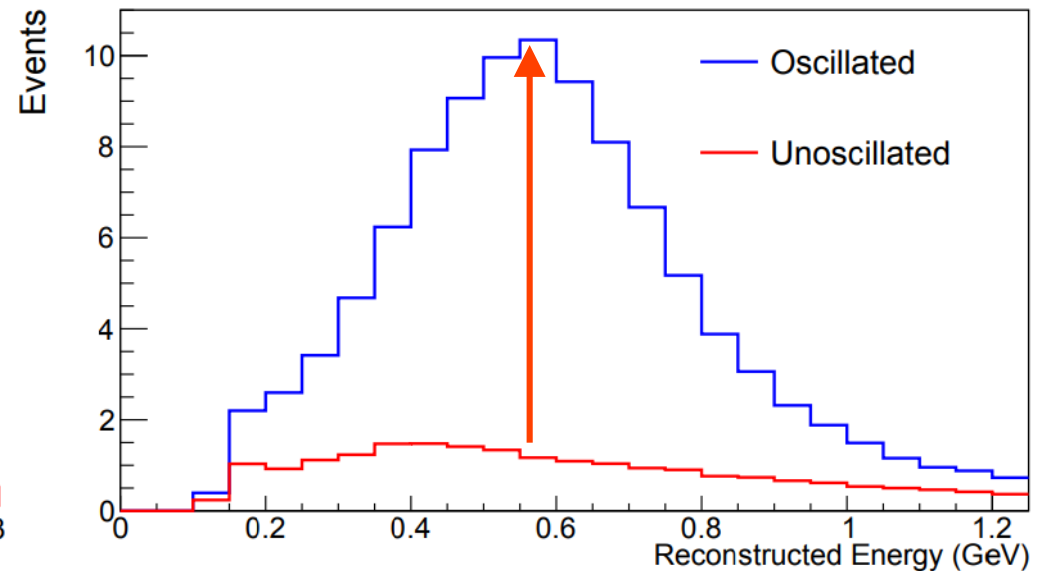
Introduction

- Oscillation parameters change the rate and shape of the appearing and disappearing neutrinos

T2K FHC 1R μ



T2K FHC 1Re

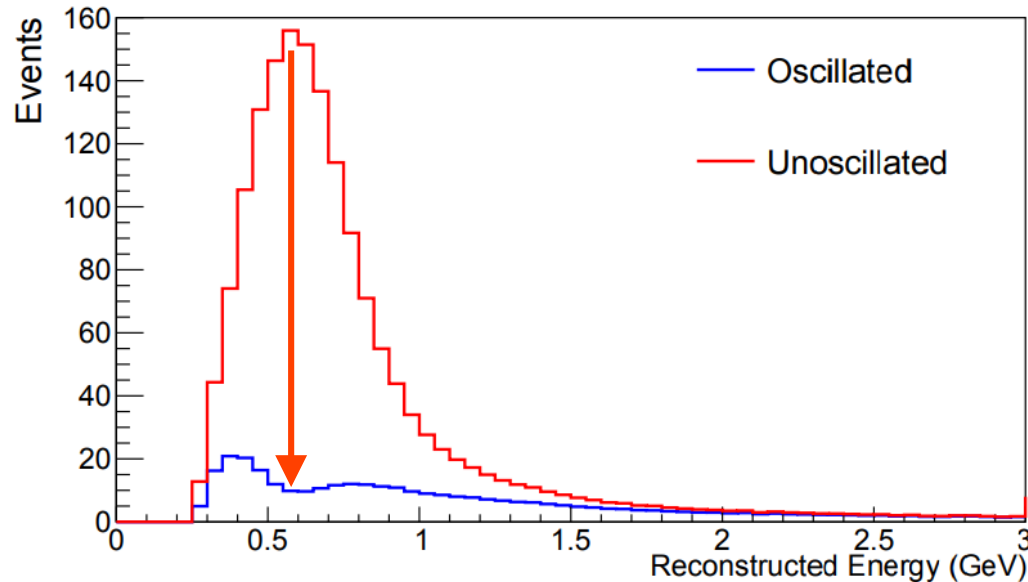


- Relies on the model prediction in the absence of oscillations
 - Constrain this model → constrain your oscillation parameters!

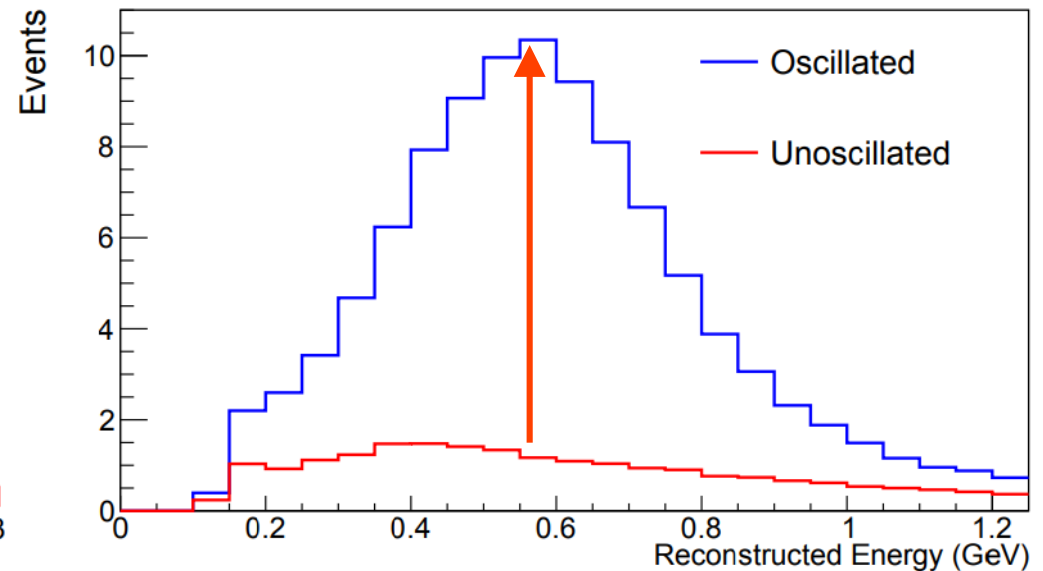
Introduction

- Oscillation parameters change the rate and shape of the appearing and disappearing neutrinos

T2K FHC 1R μ



T2K FHC 1Re



- Relies on the **model prediction** in the **absence of oscillations**
 - Constrain this model \rightarrow constrain your oscillation parameters!
- Finding cross-section effects which are degenerate with oscillation parameters is the **nightmare scenario**

Pause for air

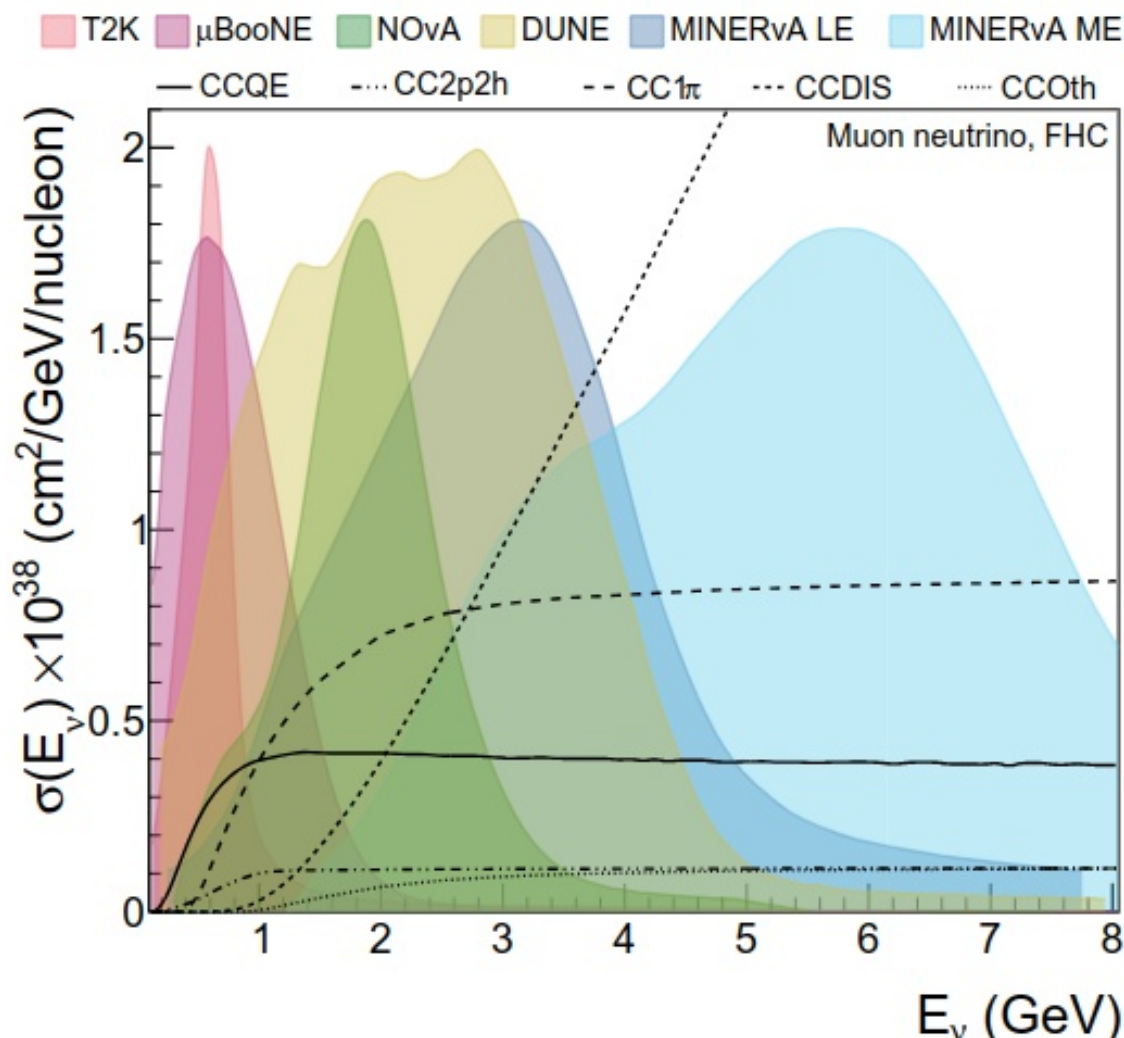
- **Muon** and **electron** (anti-)neutrinos respond differently to oscillation parameters
- **Electron (anti-)neutrinos** are the keys to unlocking δ_{CP} and **mass ordering** measurements
 - Both cause an **asymmetry** between electron neutrino and anti-neutrino oscillations; **it's not just the CP violating phase!**
- **The energy spectrum** of the electron neutrinos is important when disentangling the degeneracies
 - This is not obvious in the bi-event plots, although they are illustrative
- The **degeneracy improves** for NOvA and DUNE, which have longer baselines (larger matter effects)
 - However, they are less sensitive to δ_{CP}
 - Less events at far detector because much further away

If you enjoy playing with oscillation calculations, consider **Prob3++**, **NuFast**, and **many other calculators**

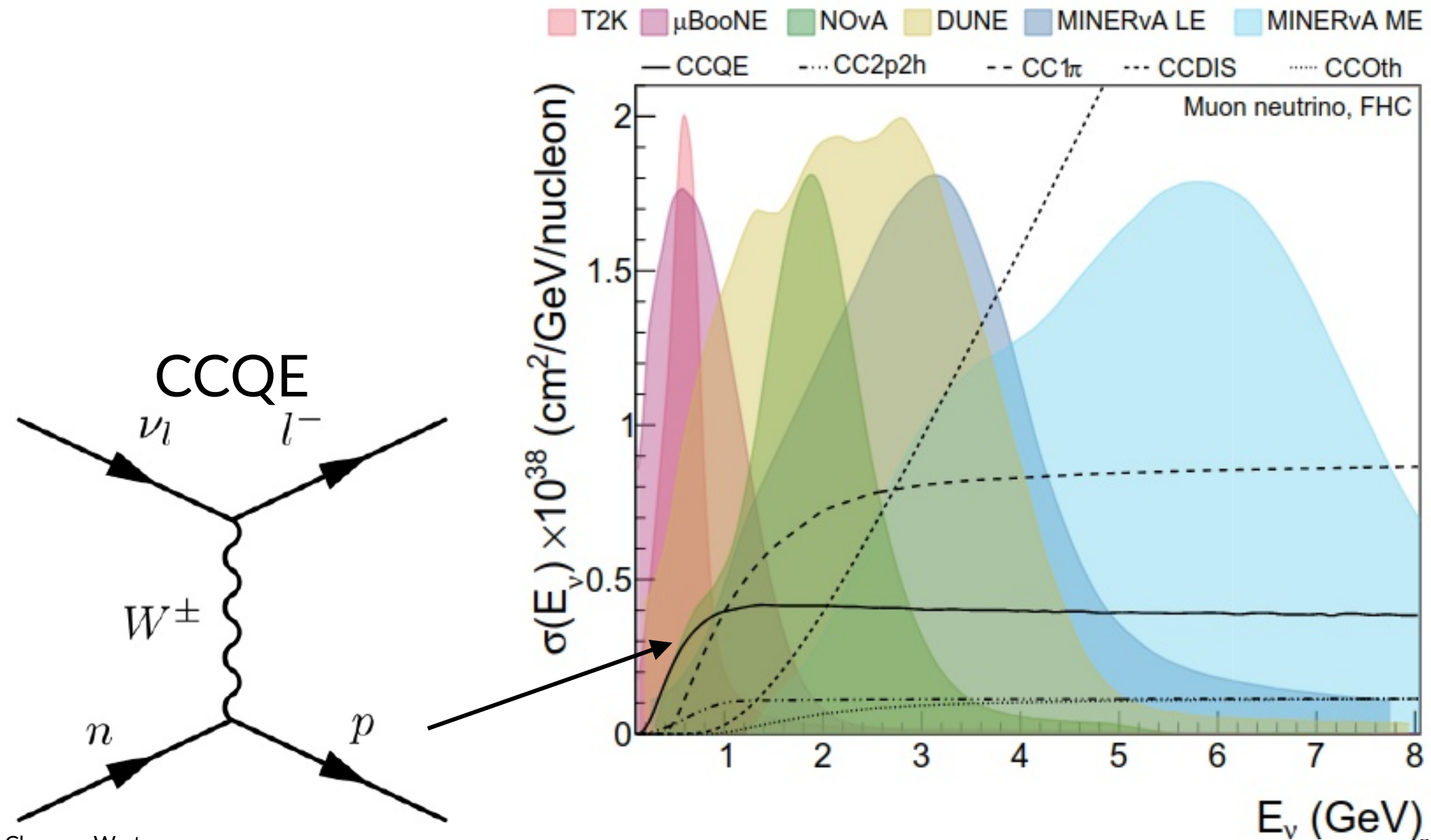
Experiments and how oscillations are measured

Neutrino fluxes from accelerators

- Accelerator neutrino oscillation experiments generally sit in the **0.5-5 GeV** region
 - Optimised for L/E ratio, matter effects, δ_{CP} sensitivity...
- The neutrino energy is a **key factor in dictating which interactions matter**
- Interaction mechanisms evolve differently in neutrino energy
- What matters for **T2K**, may not matter for **NOvA**, may not matter for **DUNE**
- Measurements from a cross-section experiment may not extrapolate well to oscillation experiment

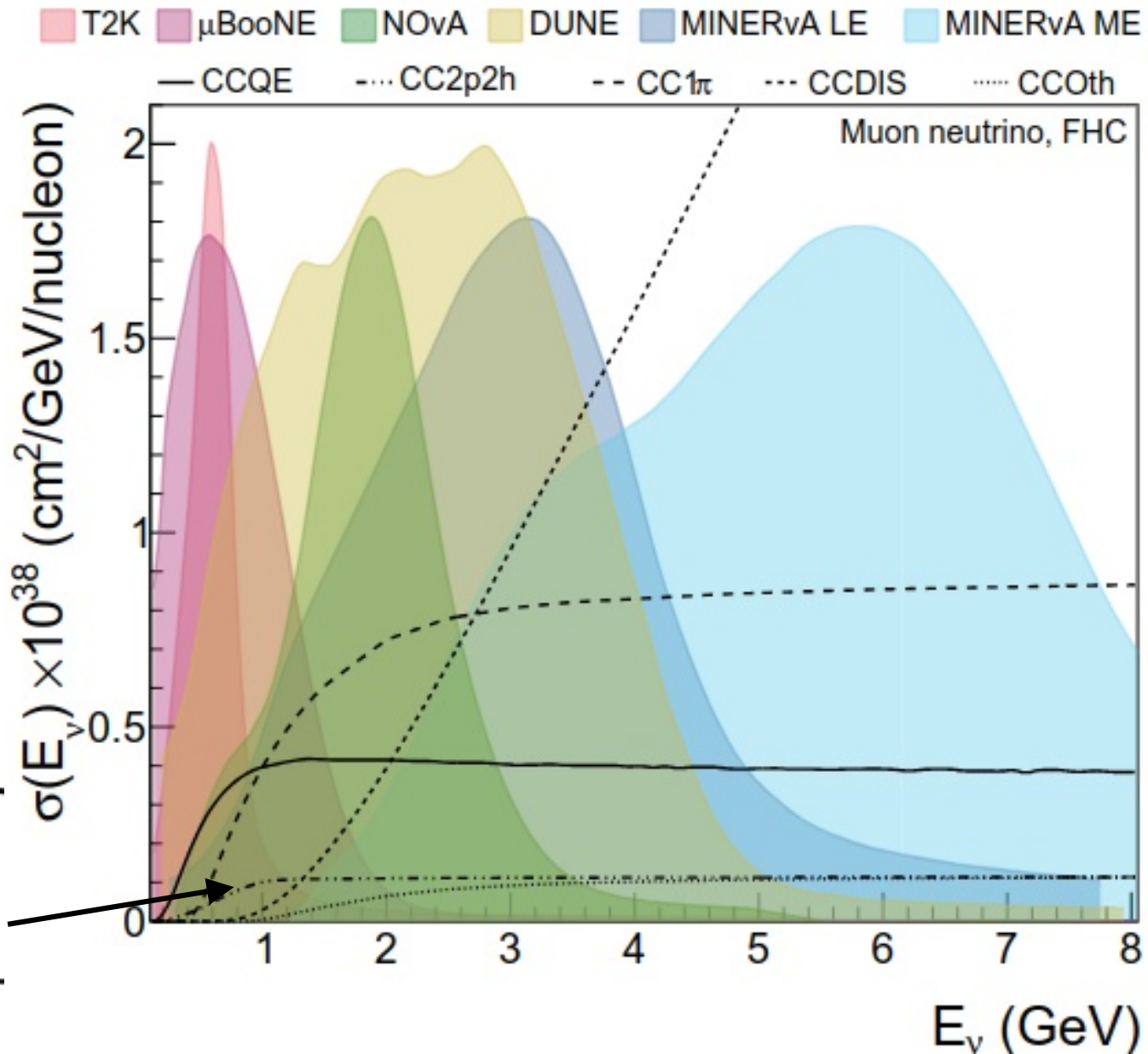
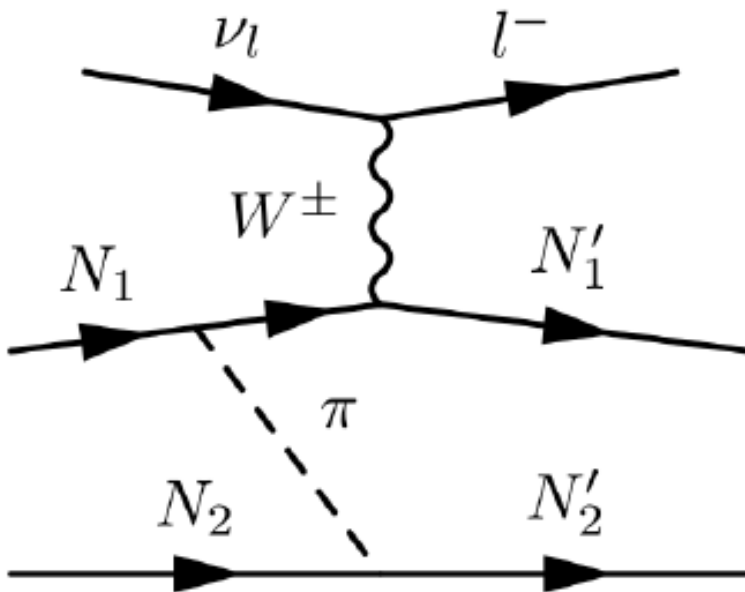


Neutrino fluxes from accelerators



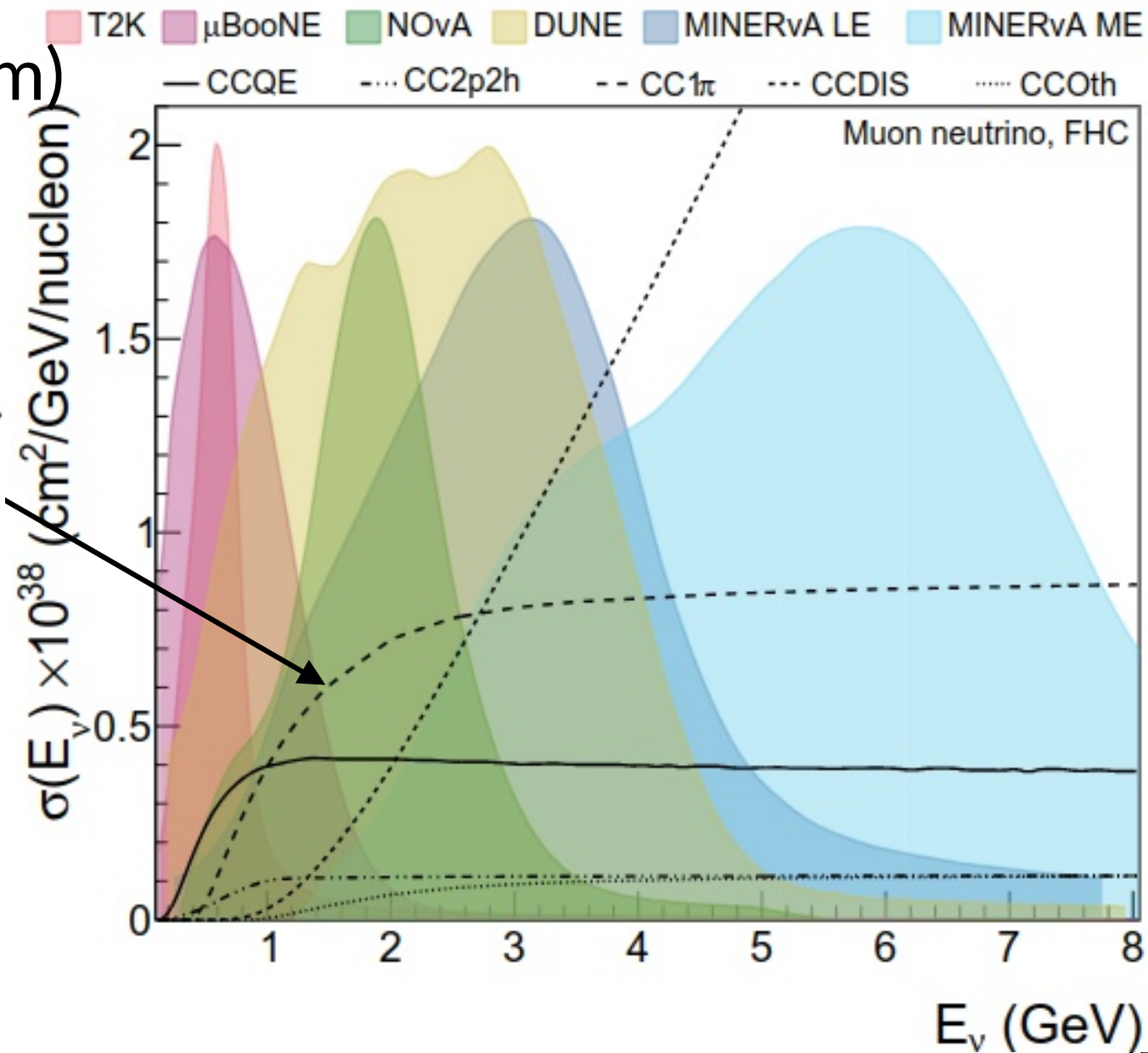
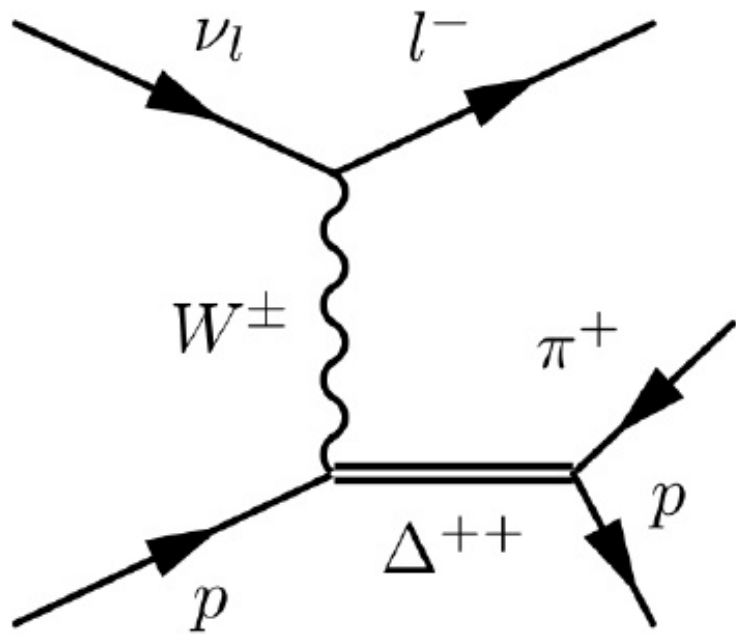
Neutrino fluxes from accelerators

2p2h (one diagram)



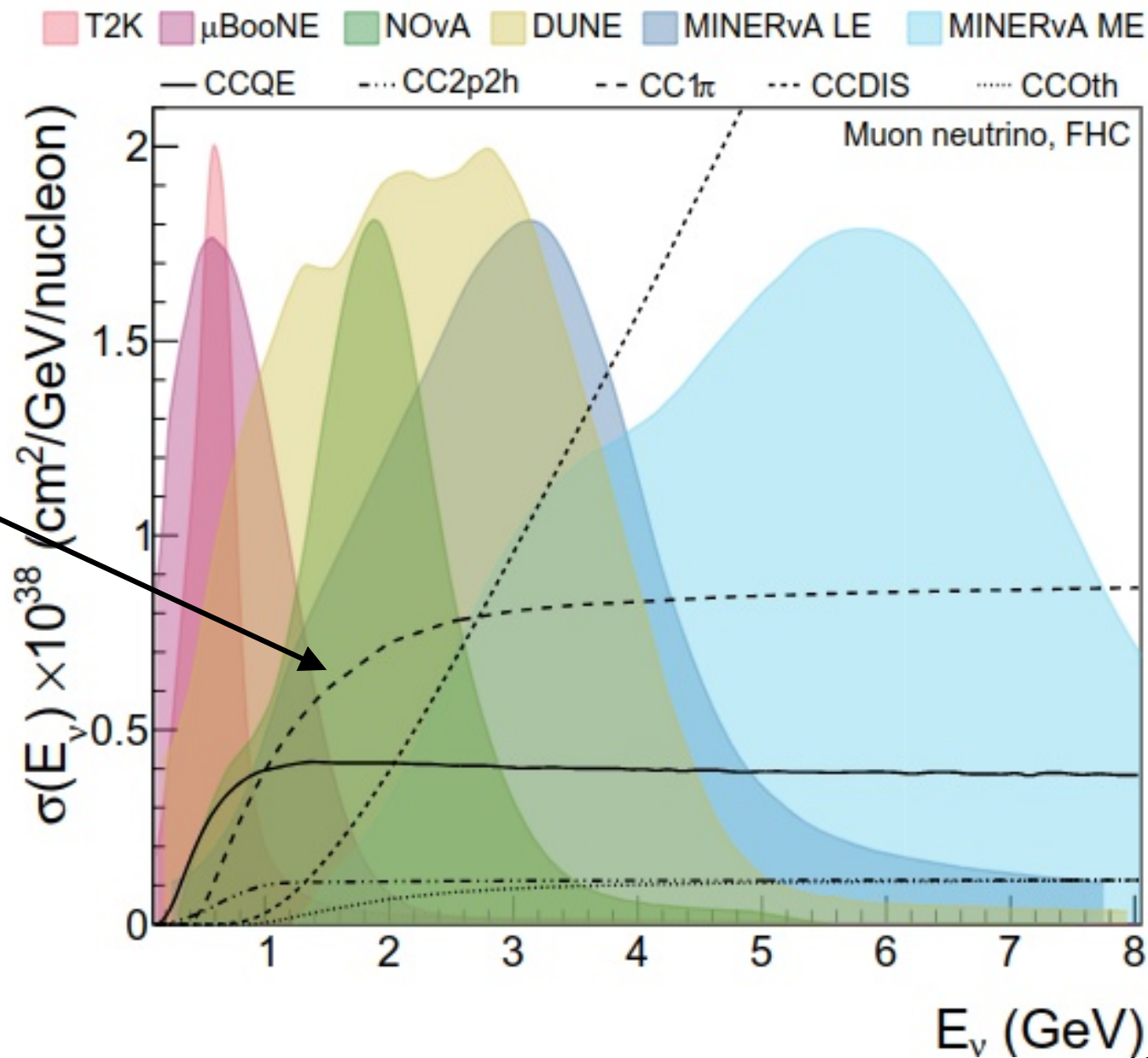
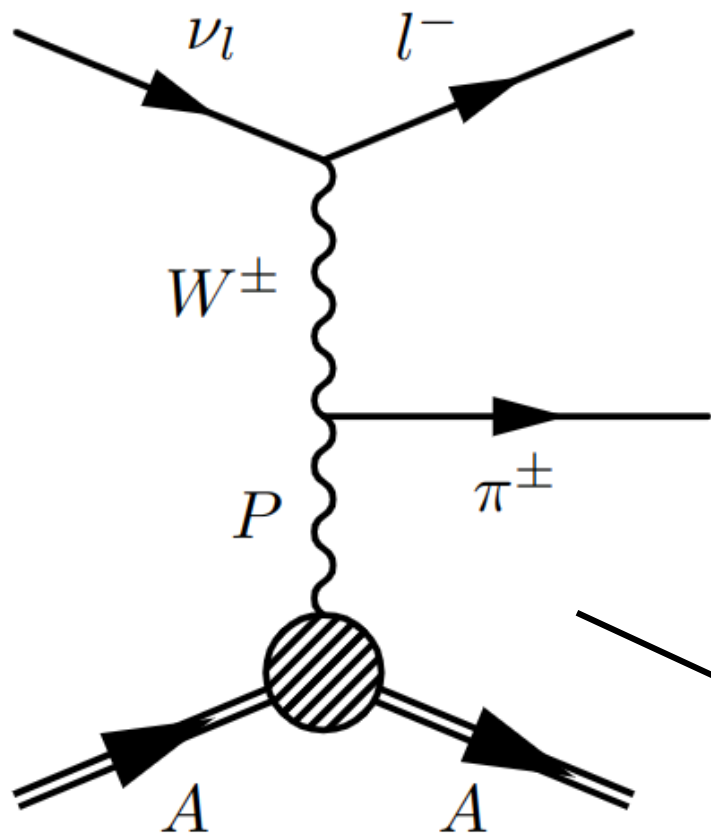
Neutrino fluxes from accelerators

CC1 π^+ (one diagram)

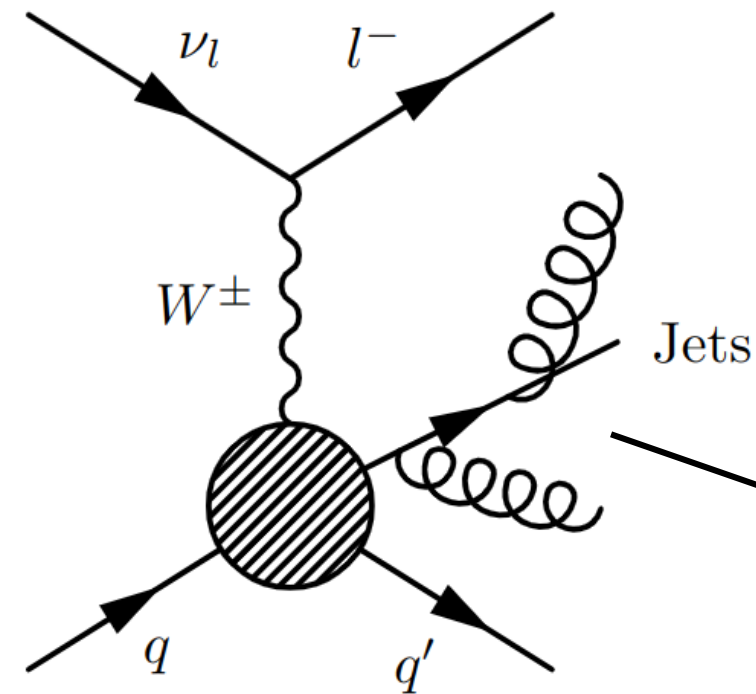


Neutrino fluxes from accelerators

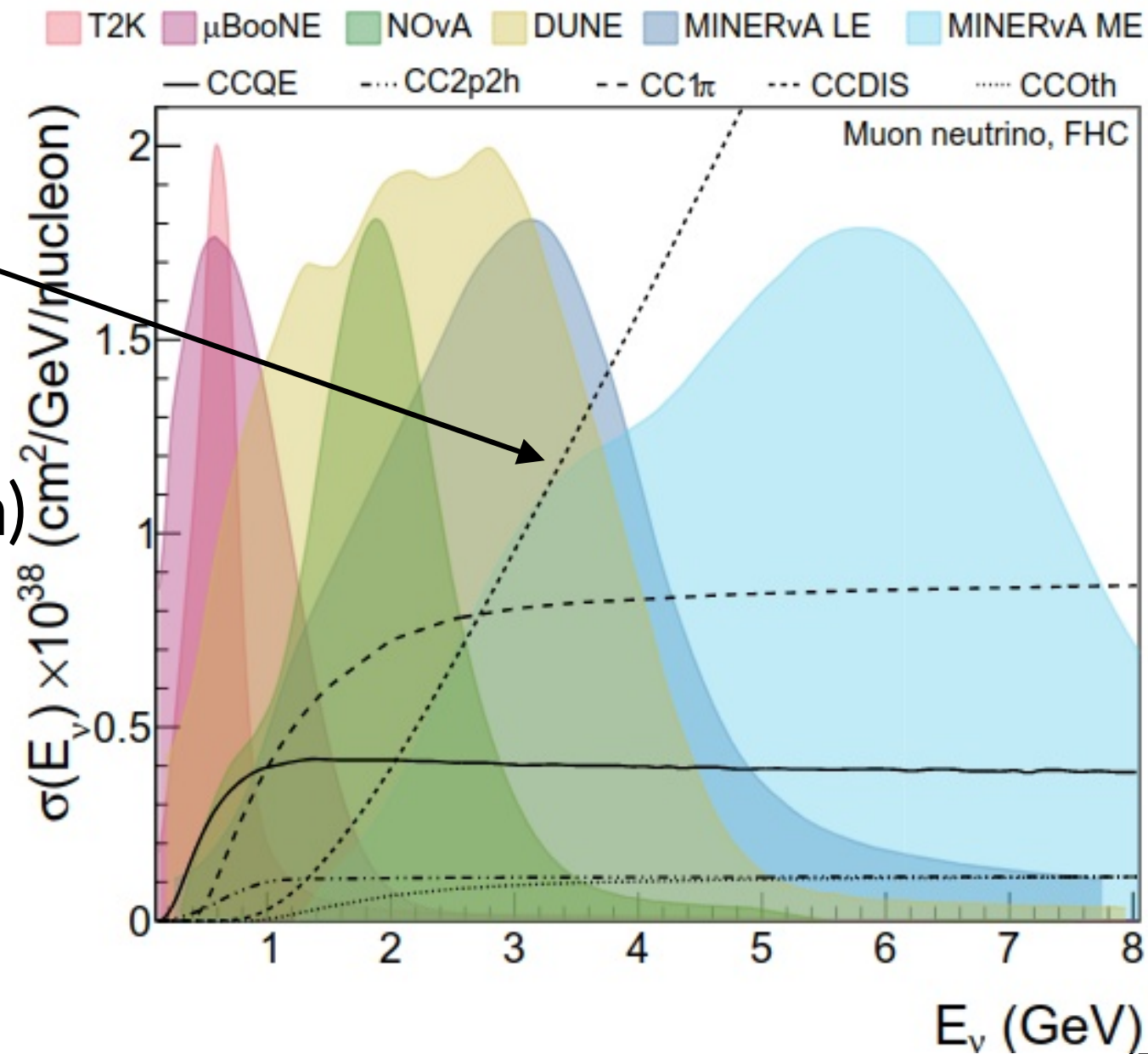
CC1 π^+ coherent



Neutrino fluxes from accelerators

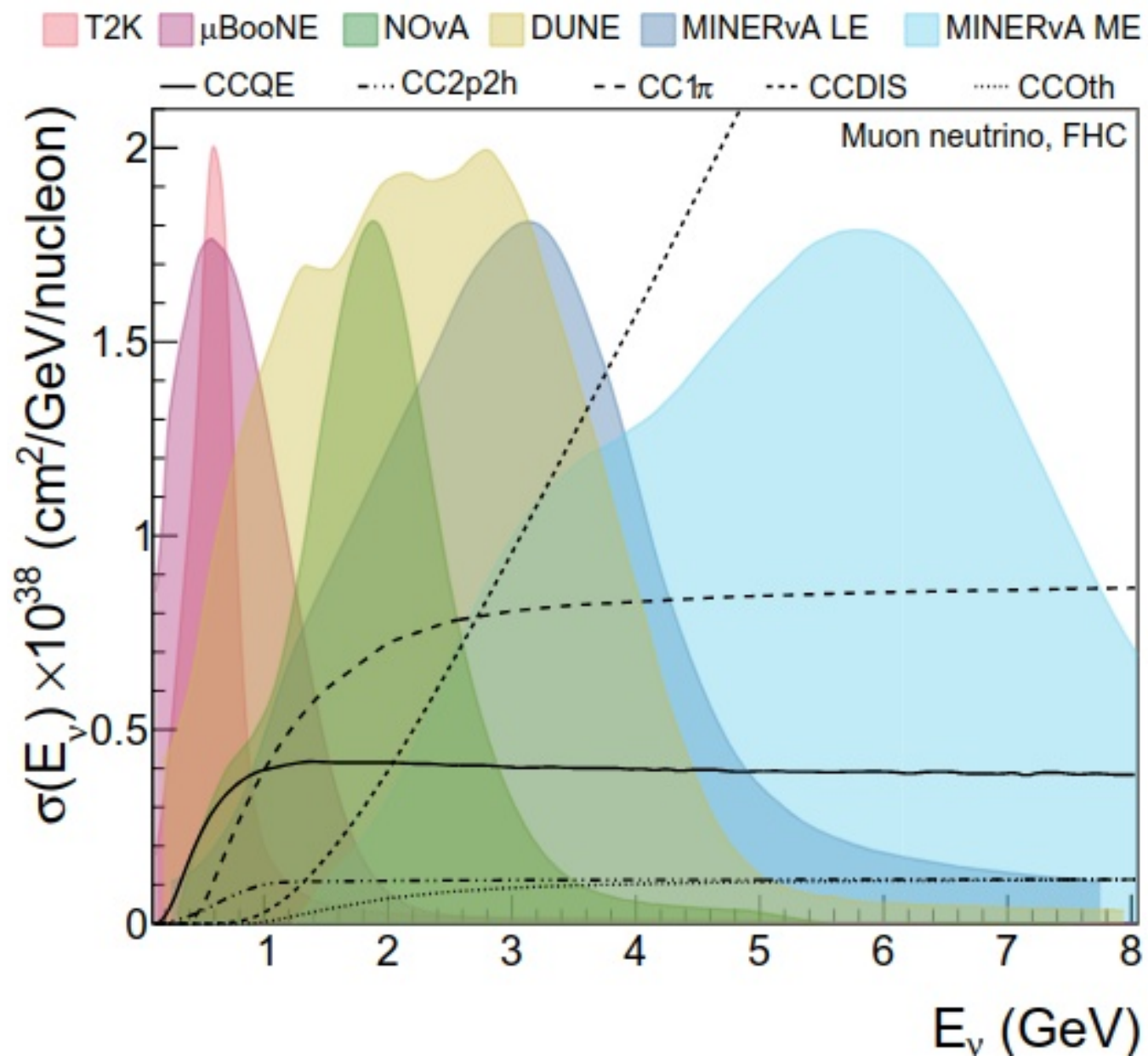


CCDIS (one diagram)



Neutrino fluxes from accelerators

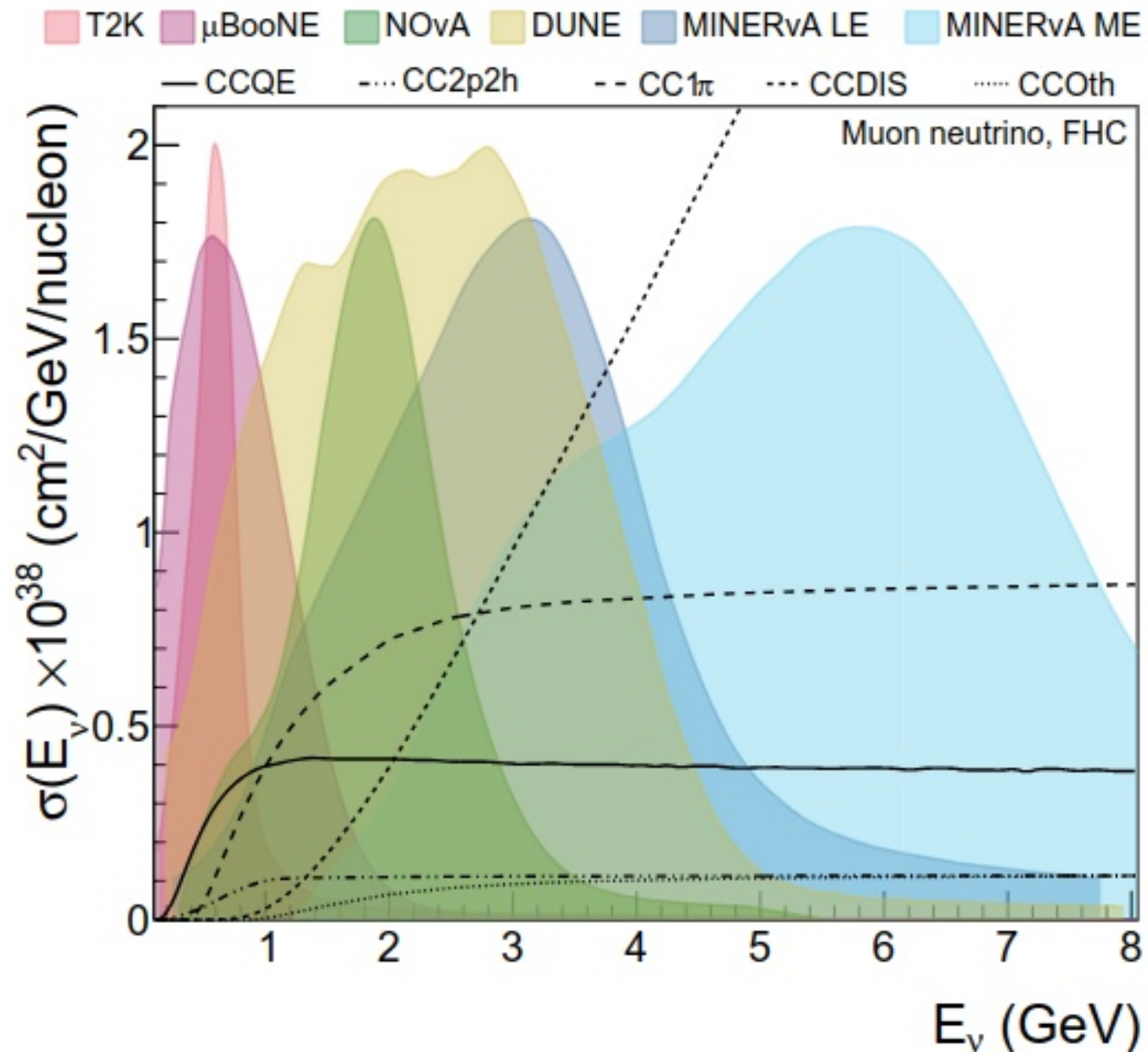
Which interactions do T2K need to worry about?



Neutrino fluxes from accelerators

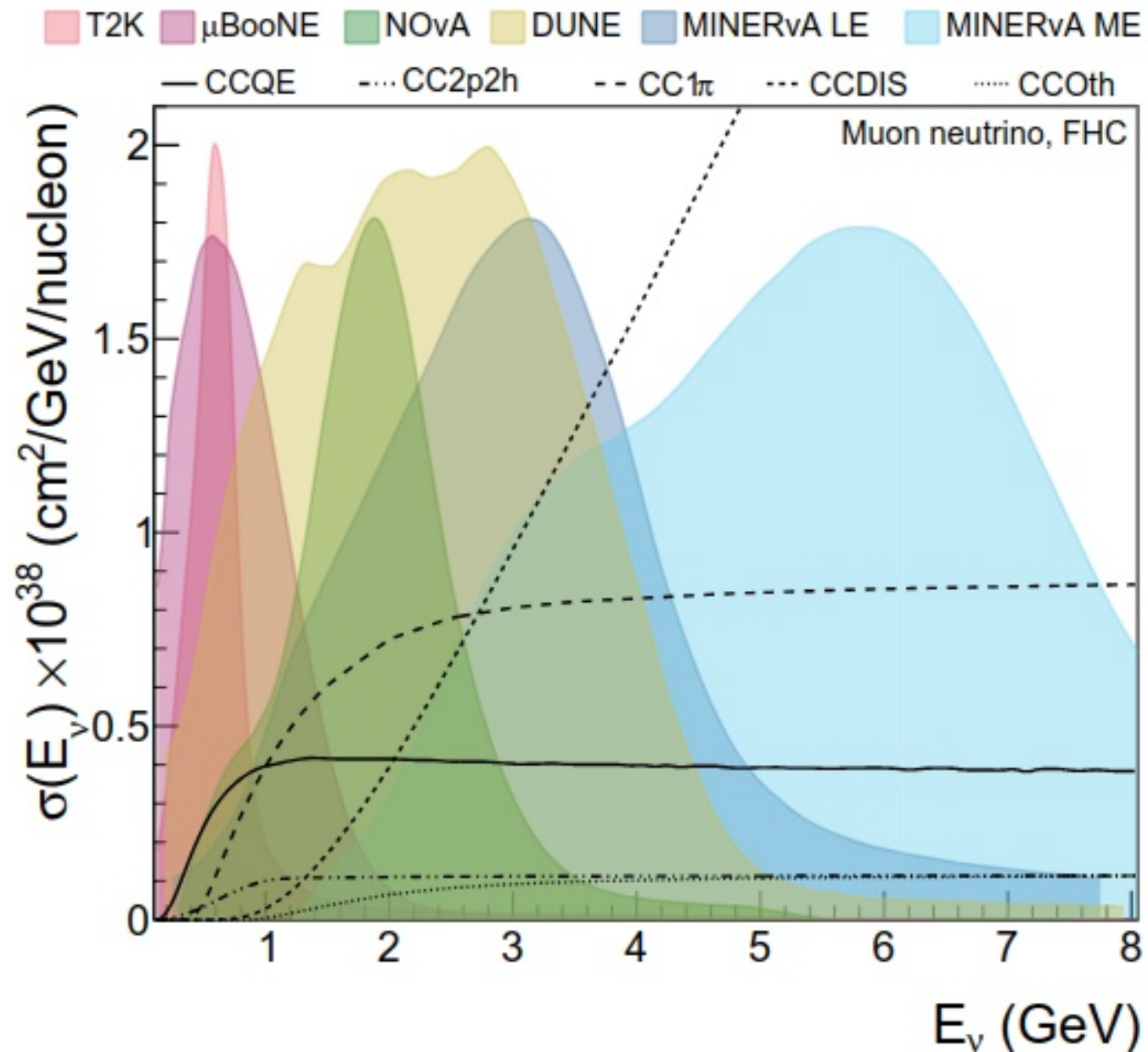
Which interactions do T2K need to worry about?

Are those shared with other experiments?

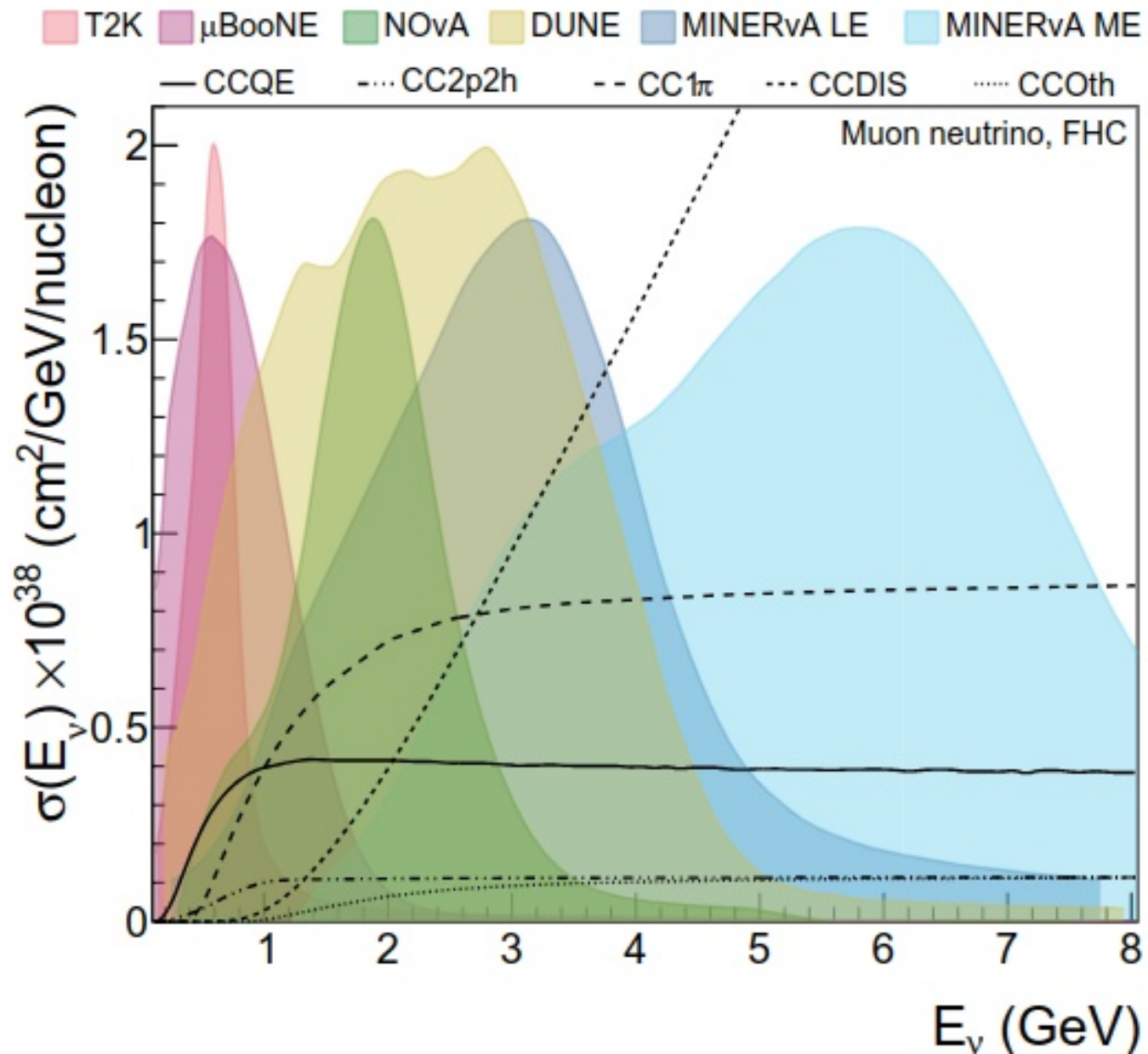


Neutrino fluxes from accelerators

What about NOvA?



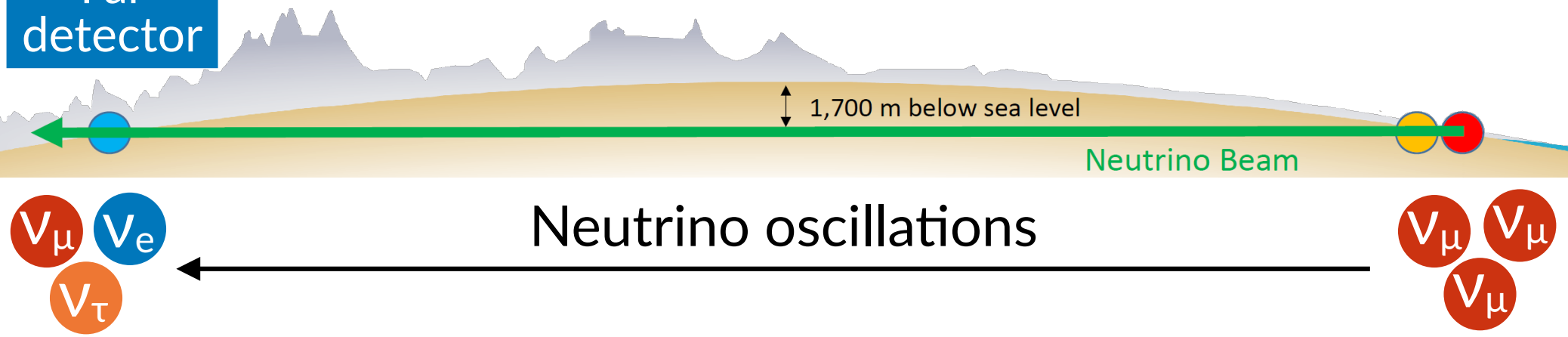
Neutrino fluxes from accelerators



And what about
DUNE?

Observations at a far detector

Far detector

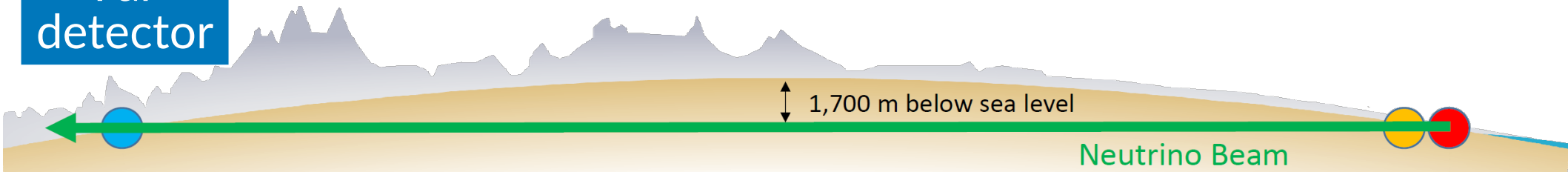


- Events observed at far detector depends on many factors

$$N_{\text{FD}}^{\alpha}(\vec{x}) = P(\nu_{\alpha} \rightarrow \nu_{\alpha}) \times \Phi^{\alpha}(E_{\nu}) \times \sigma^{\alpha}(\vec{x}) \times \epsilon_{\text{FD}}^{\alpha}(\vec{x})$$

Observations at a far detector

Far detector



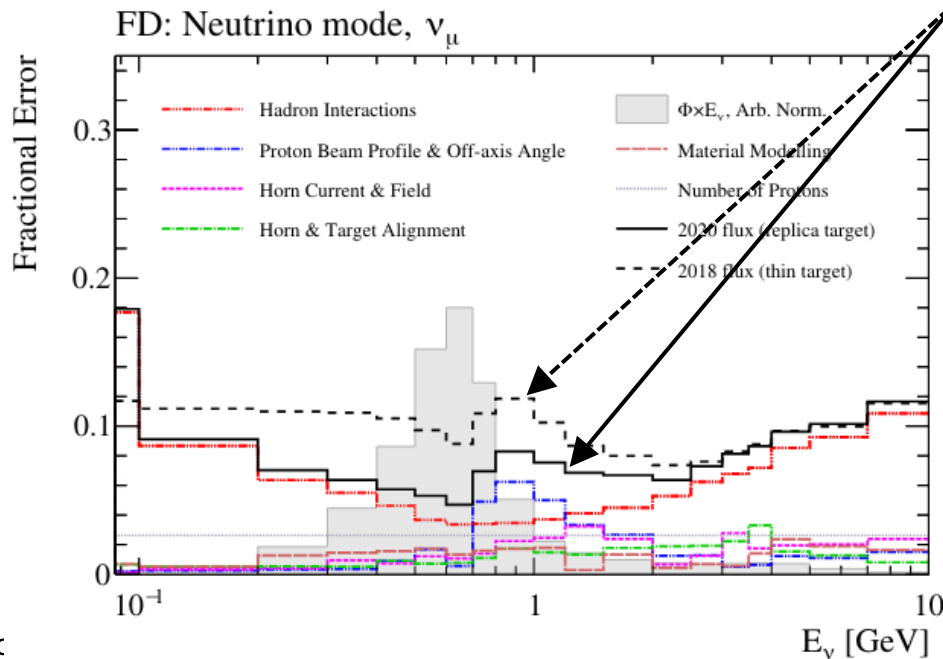
ν_μ ν_e
 ν_τ

Neutrino oscillations

ν_μ ν_μ
 ν_μ

- Events observed at far detector depends on many factors

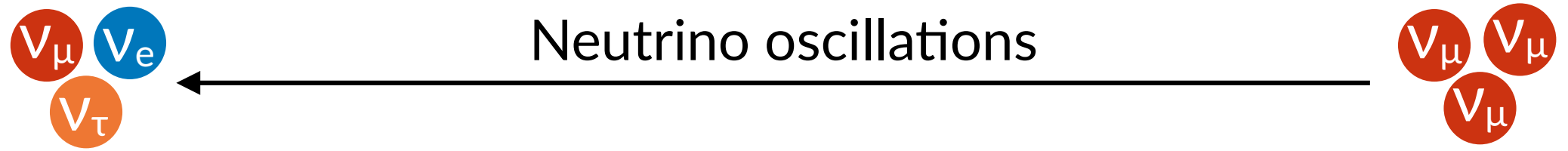
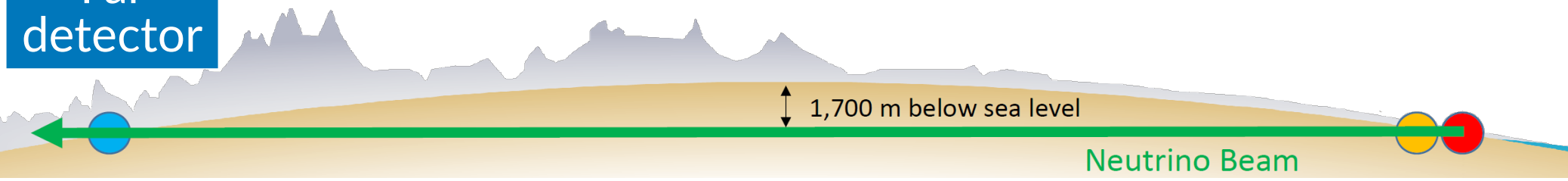
$$N_{\text{FD}}^\alpha(\vec{x}) = P(\nu_\alpha \rightarrow \nu_\alpha) \times \Phi^\alpha(E_\nu) \times \sigma^\alpha(\vec{x}) \times \epsilon_{\text{FD}}^\alpha(\vec{x})$$



5-10% absolute uncertainties on the neutrino flux

Observations at a far detector

Far detector



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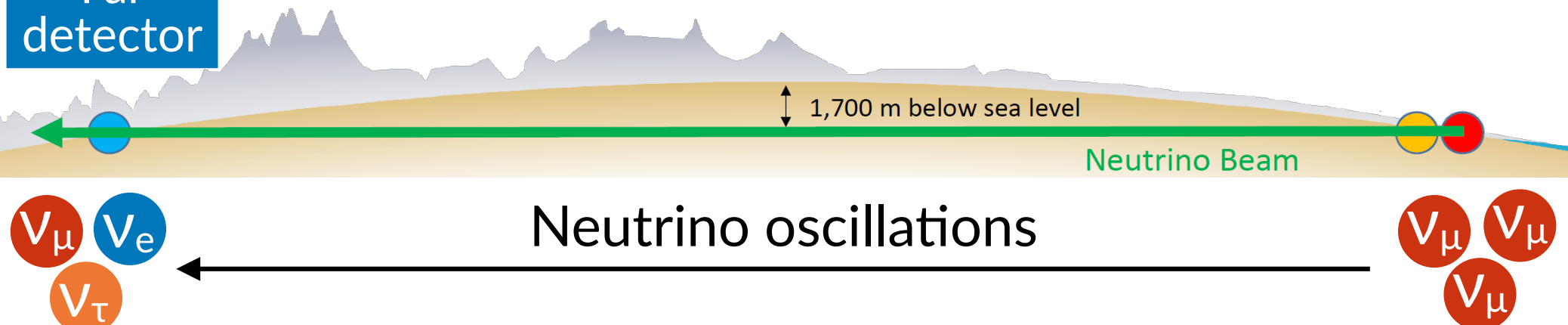
Sample	Interaction	
1R μ	ν	3.1 (11.7)
	$\bar{\nu}$	3.0 (10.8)
1Re	ν	3.2 (12.6)
	$\bar{\nu}$	3.1 (11.1)
1Re1de	ν	4.2 (12.1)

Complicated energy-dependent and selection-dependent **cross-sections**

~10% uncertainties

Observations at a far detector

Far detector



- Events observed at far detector depends on many factors

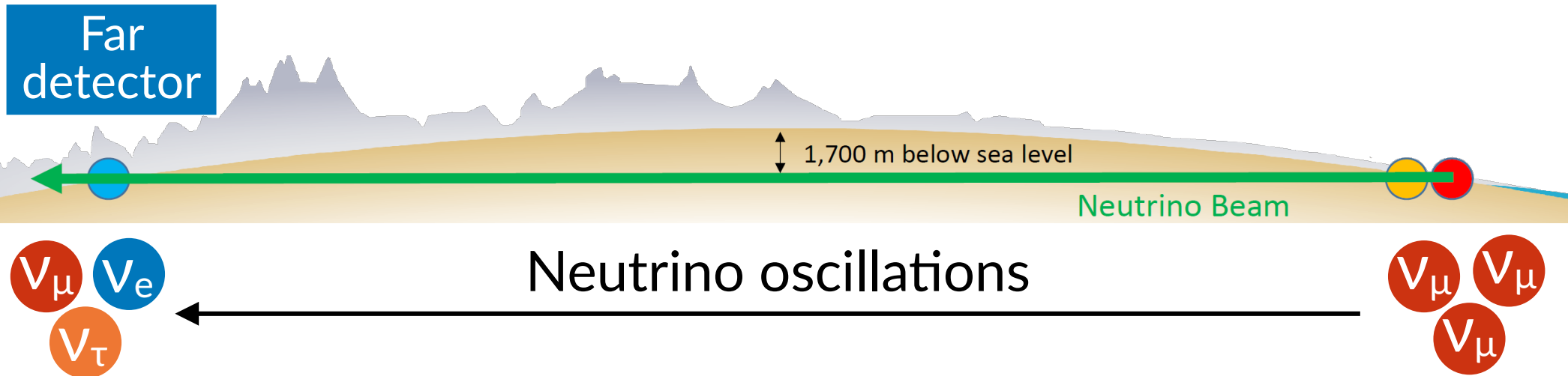
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Sample			
1R μ	ν	2.1	(2.7)
	$\bar{\nu}$	1.9	(2.3)
1Re	ν	3.1	(3.2)
	$\bar{\nu}$	3.9	(4.2)
1Re1de	ν	13.4	(13.4)

Particle acceptance may also depend on neutrino energy, and selection

3-15% uncertainty for T2K

Observations at a far detector

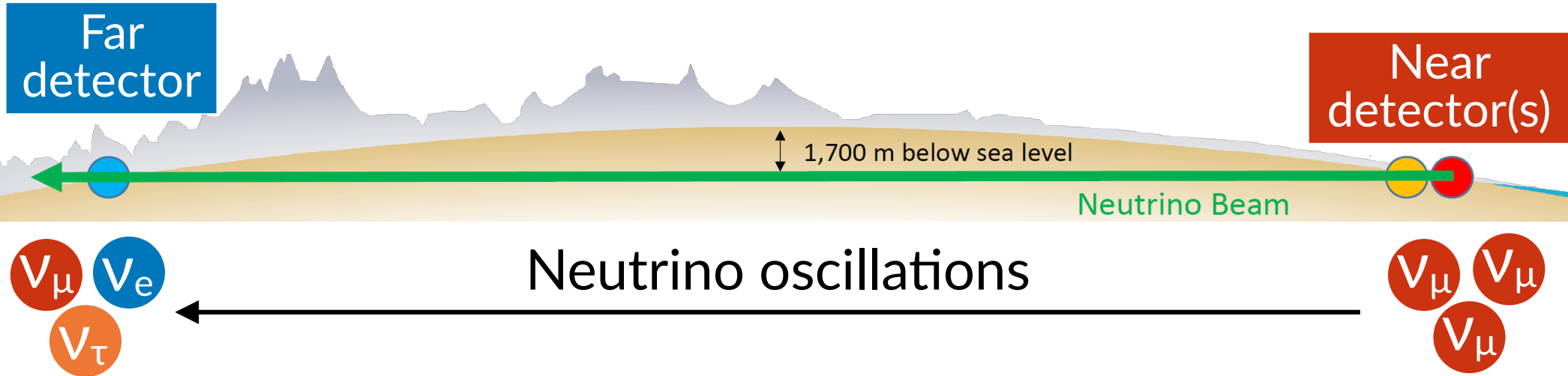


- Events observed at far detector depends on many factors

$$N_{\text{FD}}^\alpha(\vec{x}) = P(\nu_\alpha \rightarrow \nu_\alpha) \times \Phi^\alpha(E_\nu) \times \sigma^\alpha(\vec{x}) \times \epsilon_{\text{FD}}^\alpha(\vec{x})$$

- Difficult to accurately constraint **neutrino oscillations** with many **large uncertainties** getting in the way
 - Many effects may mimic the **oscillation signal**, especially if you only look at a single neutrino flavour

The near detector

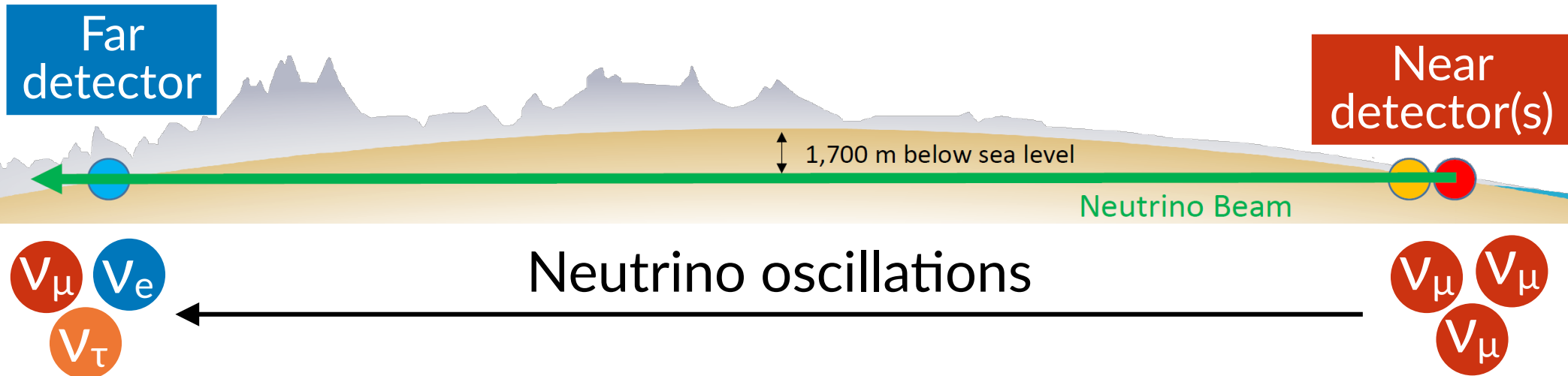


- But what if you have a **near detector**?

$$N_{\text{FD}}^\alpha(\vec{x}) = P(\nu_\alpha \rightarrow \nu_\alpha) \times \Phi^\alpha(E_\nu) \times \sigma^\alpha(\vec{x}) \times \epsilon_{\text{FD}}^\alpha(\vec{x})$$

$$N_{\text{ND}}^\alpha(\vec{x}) = \Phi^\alpha(E_\nu) \times \sigma^\alpha(\vec{x}) \times \epsilon_{\text{ND}}^\alpha(\vec{x})$$

The near detector



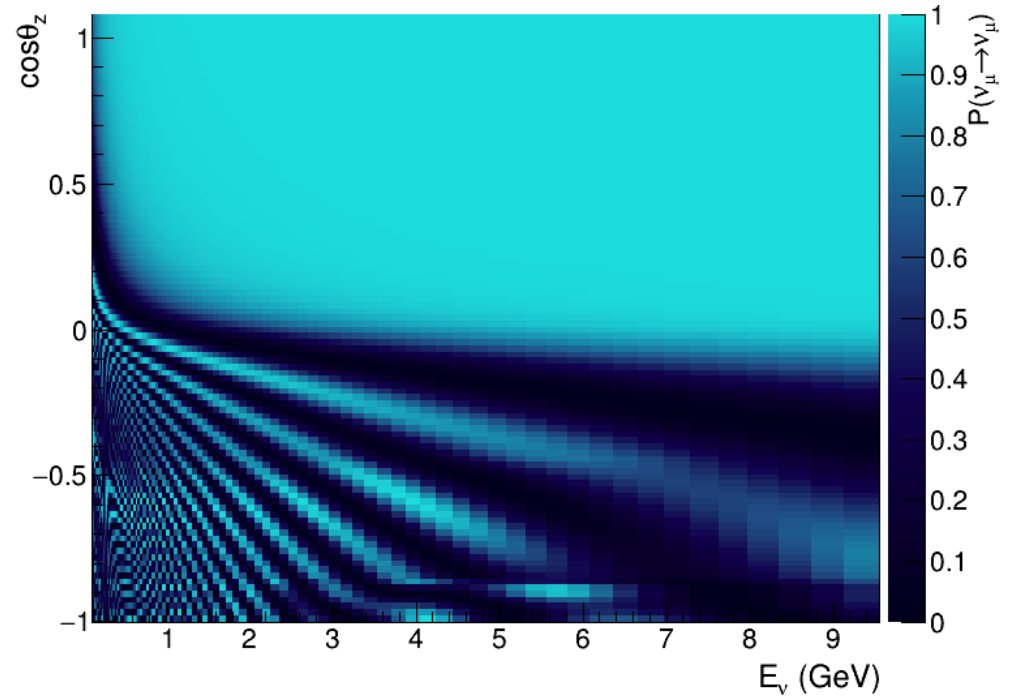
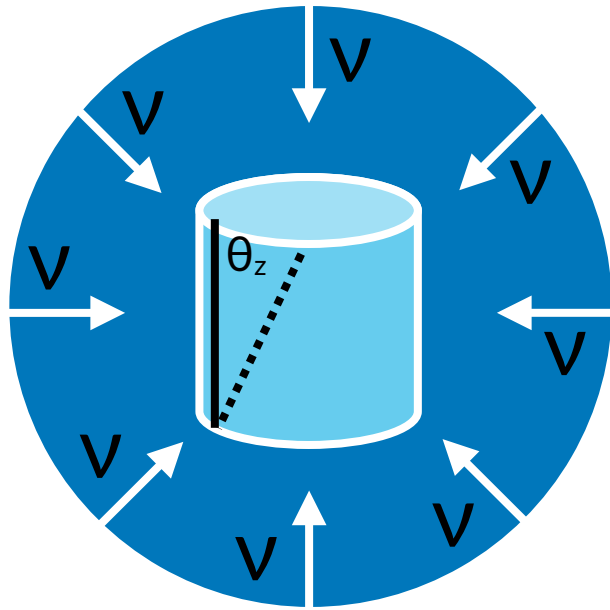
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$$N_{\text{ND}}^\alpha(\vec{x}) = \Phi^\alpha(E_\nu) \times \sigma^\alpha(\vec{x}) \times \epsilon_{\text{ND}}^\alpha(\vec{x})$$

- Events observed at the far detector have many **shared uncertainties** with the near detector
 - Constrain **flux and interaction model** using near detector data
- **Characterise neutrinos with high-statistics near-detector samples** before long baseline oscillations
- **Mitigates many of the issues**, e.g. size of cross sections, flux normalisation...

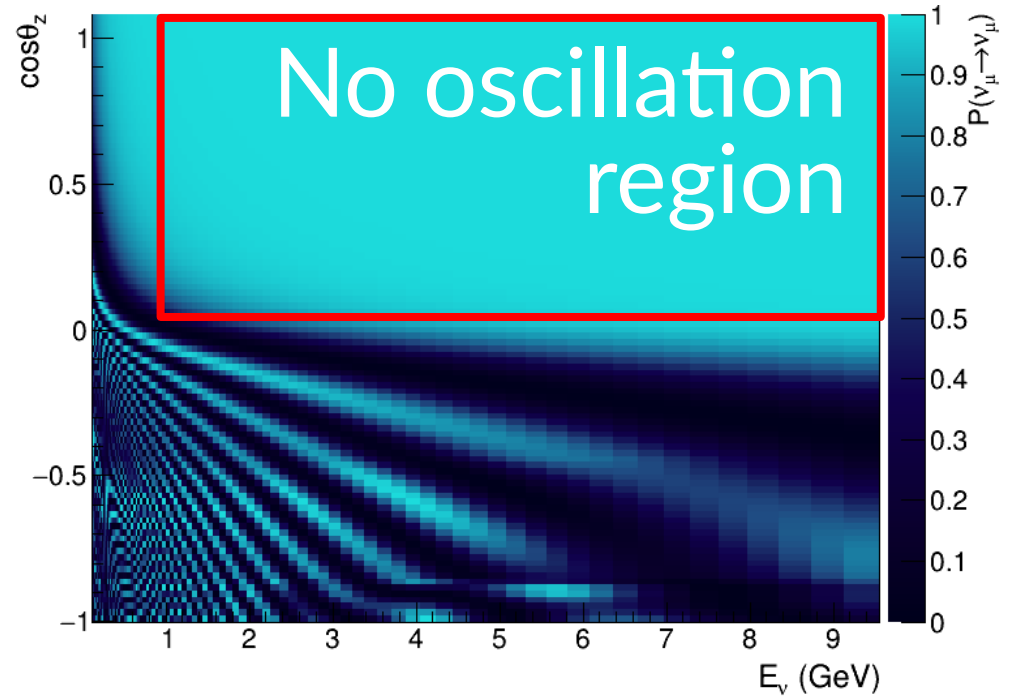
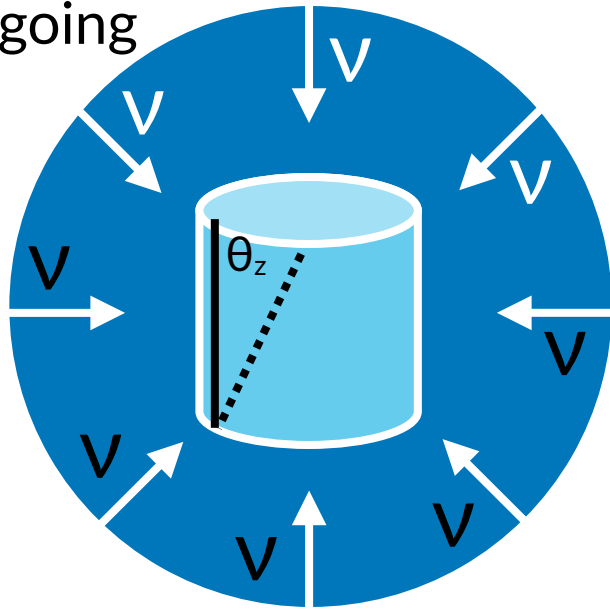
Aside: atmospheric near detector?



- For atmospheric neutrinos, there is no near detector

Aside: atmospheric near detector?

Down-going region



- For atmospheric neutrinos, there is no near detector
- Largely addressed by **down-going neutrinos**
 - Very small oscillation probability in region
 - **Effectively acts as a near-detector** constraint throughout a large neutrino energy range



**ALL SYSTEMATICS
CANCEL WITH A NEAR DETECTOR**



**YOU'VE THOUGHT ABOUT
ACCEPTANCE MATCHING, ENERGY
DEPENDENCE, INTRINSIC NUES, RIGHT?**



RIGHT?!

Far detector



• But wh

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• Charac
baselin

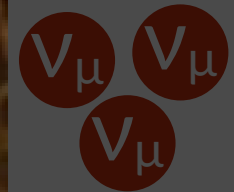
• Events
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- Cor

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Near tector(s)



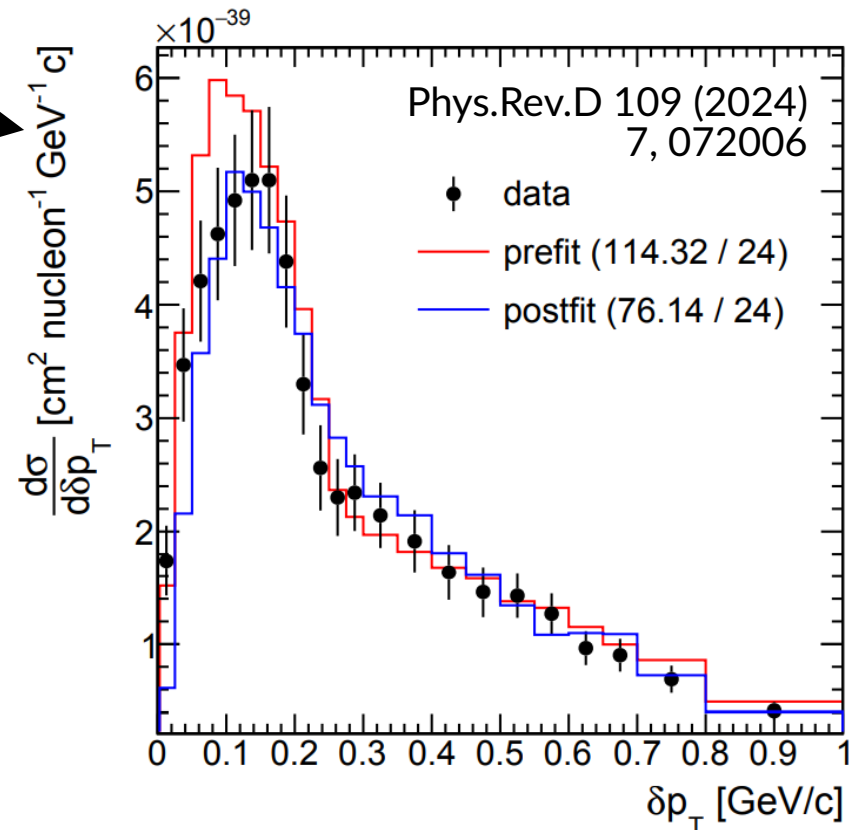
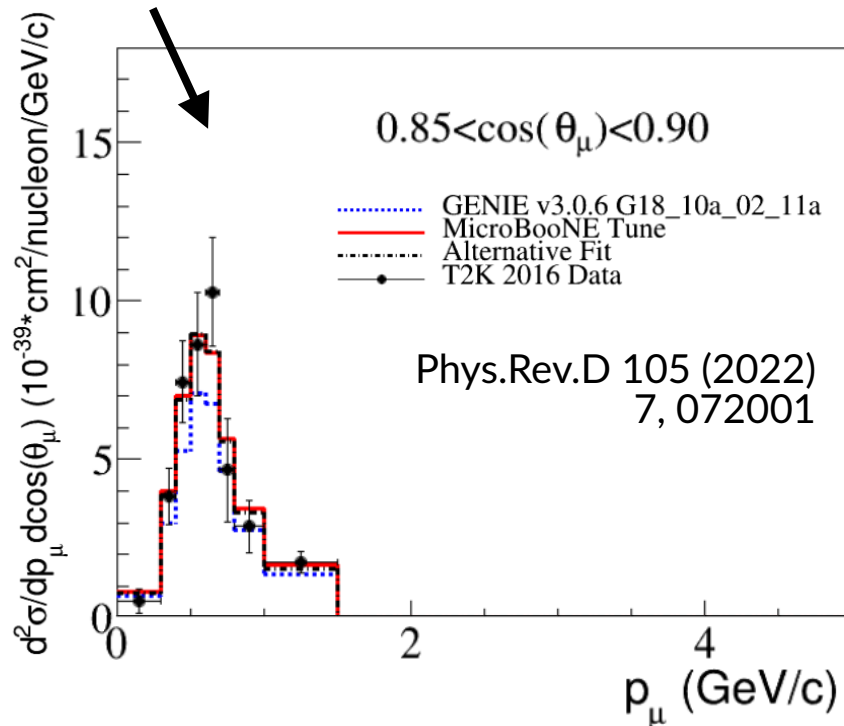
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Role of external data

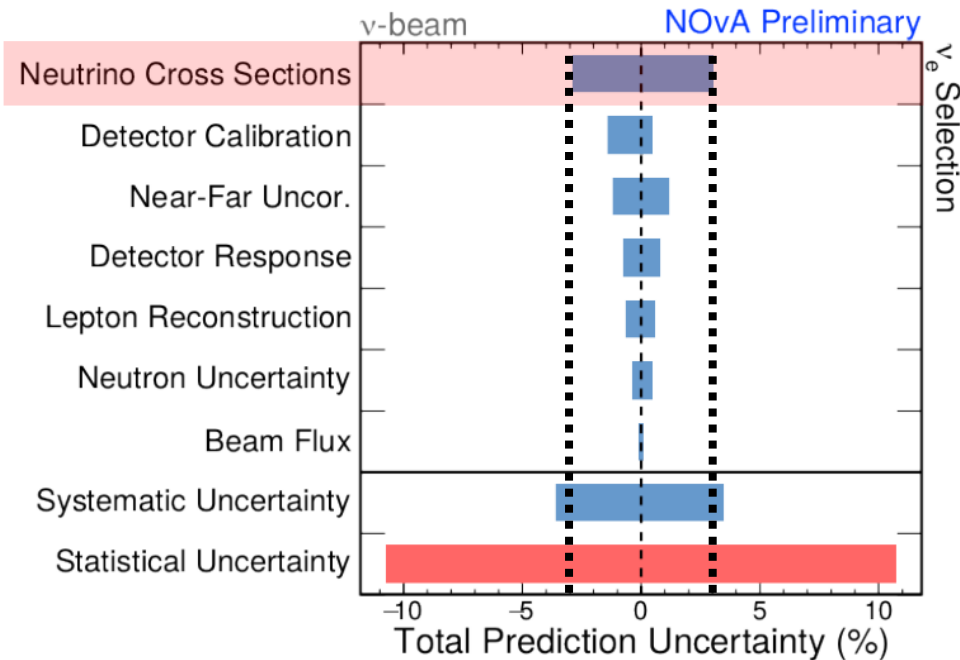
- You might not have a near detector; what do you do?
- Or in some cases, data from the **near detector might not suffice**
 - e.g. you have an unmagnetised detector, but want to estimate NC1 π^+ contribution to the background in ν_μ disappearance
- **External data is often used to estimate the cross section, and prevent a near-detector analysis from over-constraining the model**
 - T2K using MINERvA data
 - MicroBooNE using T2K data



Impact of systematics at the FD

- Neutrino cross-section uncertainties contribute ~3% to number of ν_e on NOvA

M. Elkins, T. Nosek, Neutrino 2020 poster

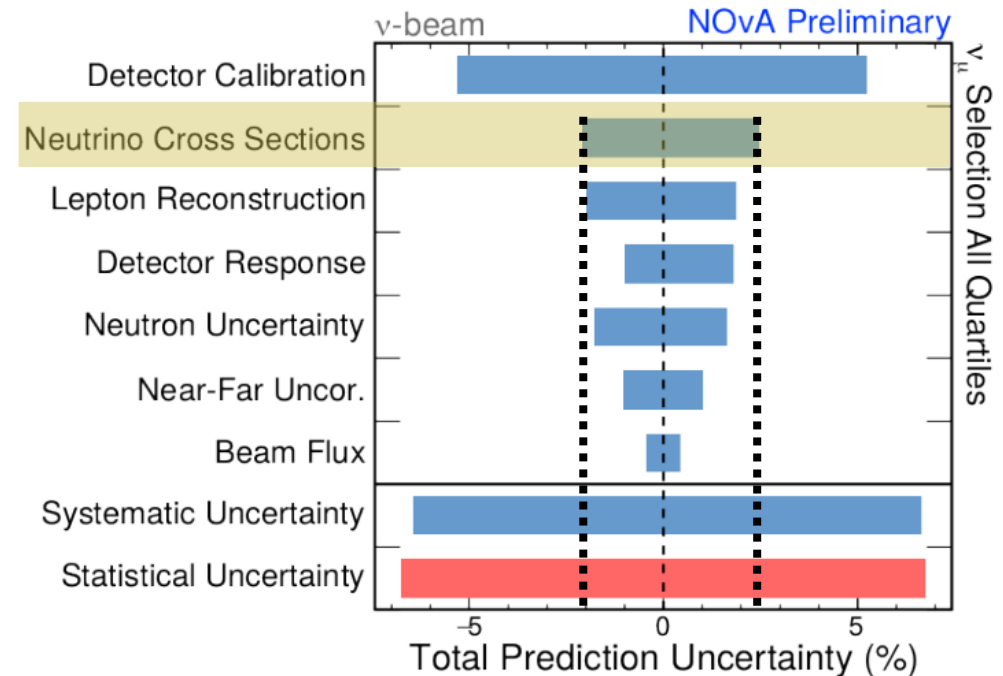
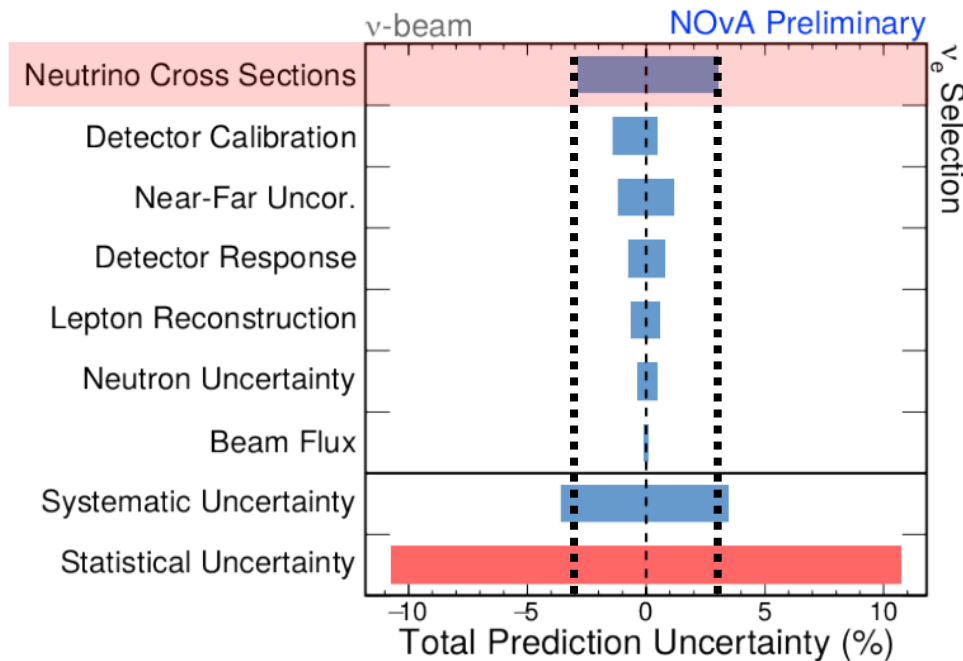


- Dominant systematic amongst all systematics
- But measurement significantly limited by statistics currently

Impact of systematics at the FD

- Neutrino cross-section uncertainties contribute $\sim 3\%$ to number of ν_e on NOvA

M. Elkins, T. Nosek, Neutrino 2020 poster



- Dominant systematic amongst all systematics
- But measurement significantly limited by statistics currently
- ν_μ roughly same systematic and statistical uncertainty!
 - Dominated by detector calibrations, followed by cross sections ($\sim 2\%$ level)

Impact of systematics at the FD

- On T2K, cross-section uncertainties contribute $\sim 3\%$ to ν_μ systematic uncertainty
 - In practice, slightly smaller because ND constrains convolution of flux * cross-section parameters

Sample		Uncertainty source (%)			Flux \otimes Interaction (%)	Total (%)
		Flux	Interaction	FD + SI + PN		
1R μ	ν	2.9 (5.0)	3.1 (11.7)	2.1 (2.7)	2.2 (12.7)	3.0 (13.0)
	$\bar{\nu}$	2.8 (4.7)	3.0 (10.8)	1.9 (2.3)	3.4 (11.8)	4.0 (12.0)

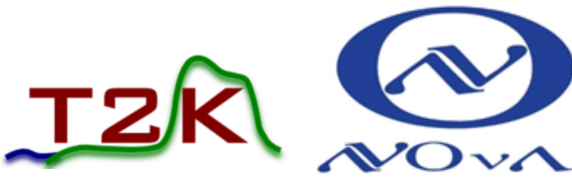
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1Re	ν	2.8 (4.8)	3.2 (12.6)	3.1 (3.2)	3.6 (13.5)	4.7 (13.8)
	$\bar{\nu}$	2.9 (4.7)	3.1 (11.1)	3.9 (4.2)	4.3 (12.1)	5.9 (12.7)
1Re1de	ν	2.8 (4.9)	4.2 (12.1)	13.4 (13.4)	5.0 (13.1)	14.3 (18.7)

- ν_e samples see 3-5% contribution to the 5-14% total
 - Detector systematics on-par with cross-section systematics
 - Small statistics means current measurements not limited by systematics
- But... we'll come back to this later with “fake-data studies”





Event counts at the FDs



Sample	T2K	NOVA
N_{μ}^{rec} FHC	318	211
N_{μ}^{rec} RHC	137	105
N_e^{rec} FHC	108	82
N_e^{rec} RHC	16	33

- ν_e measurements, especially in RHC, are heavily limited by statistics in current experiments
 - ~10-25%
- ν_{μ} measurements at the ~5% statistics level

Event counts at the FDs

Sample	 T2K	 NOVA	 Hyper-Kamiokande	 DUNE
N_{μ}^{rec} FHC	318	211	10000	7000
N_{μ}^{rec} RHC	137	105	14000	3500
N_e^{rec} FHC	108	82	3000	1500
N_e^{rec} RHC	16	33	3000	500

- HK and DUNE will have enough ν_e events to be limited by the $\sim 3\%$ (anti-) ν_e uncertainty
- ν_{μ} measurements on the 1% scale
- Current uncertainties at the 3-5% level uncertainties*

*Exception of T2K's single-pion-below-threshold sample (10-15%)

Event counts at the FDs

Sample



N_{μ}^{rec} FHC	7000
N_{μ}^{rec} RHC	5000
N_e^{rec} FHC	3000
N_e^{rec} RHC	2000

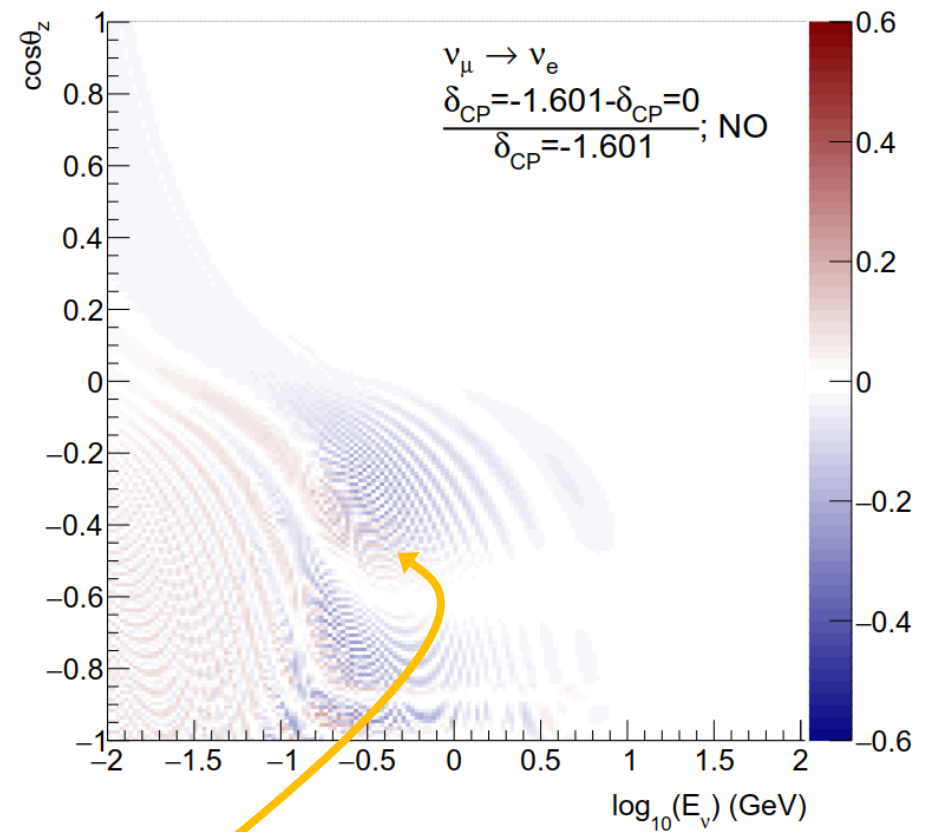
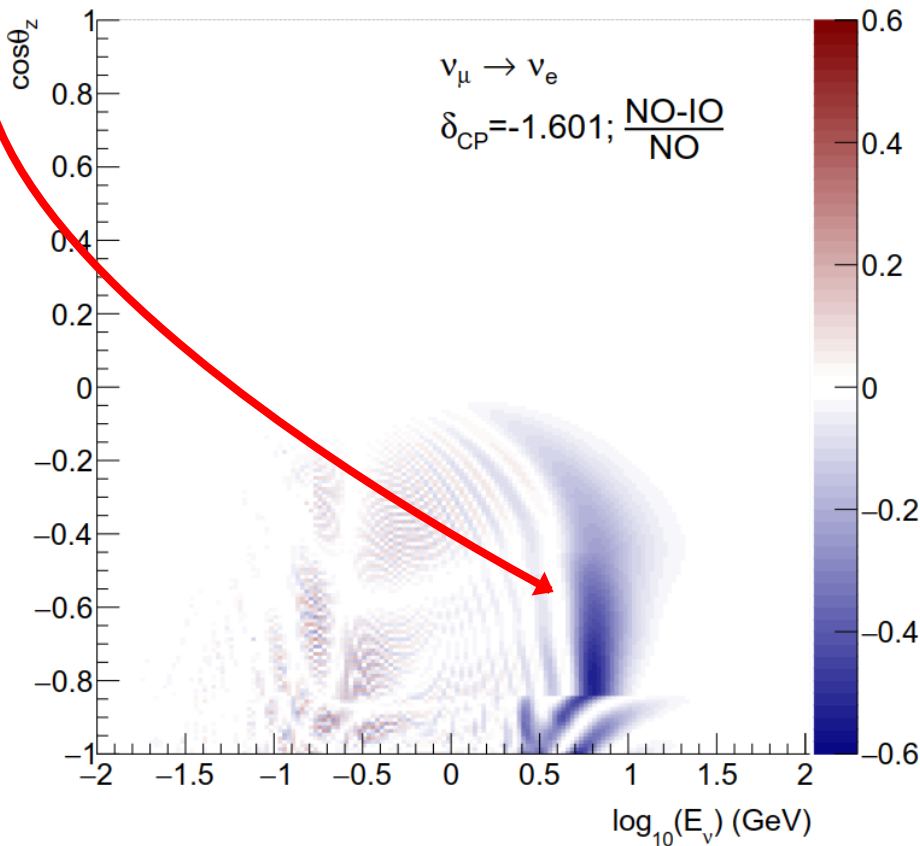
We've all got work to do!

- HK and RHC limited by ν_{μ} flux to be
- ν_{μ} measurement
- Current uncertainties at $\delta = 10\%$ level
- Current uncertainties* $\delta = 10\%$

*Exception of T2K's single-pion-below-threshold sample (10-15%)

Atmospheric neutrinos

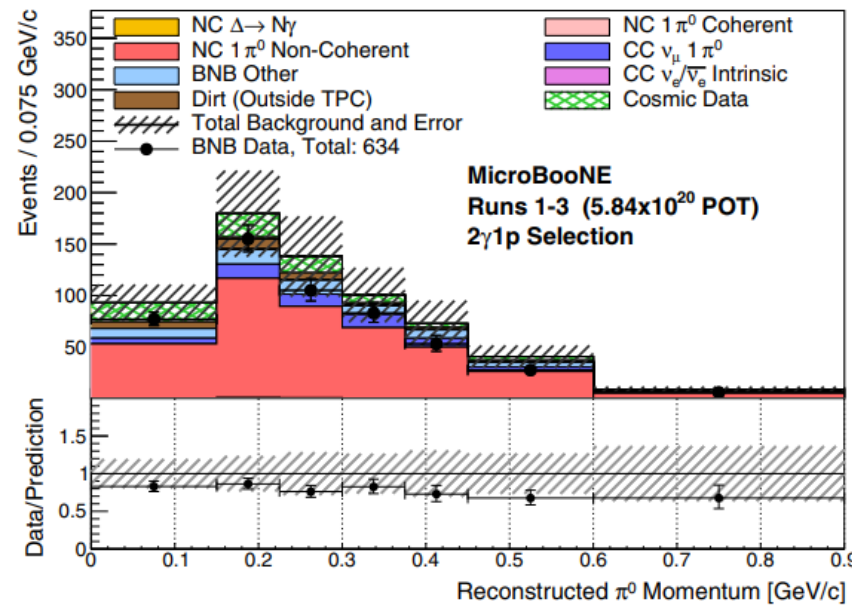
- Atmospheric neutrinos have sensitivity to mass ordering via 3-10 GeV resonance
 - Opposite effect for neutrino and anti-neutrinos: **need to separate**
 - Contribution from $\nu_\mu \rightarrow \nu_\tau$, where ν_τ enters multi-ring ν_e sample



- δ_{CP} sensitivity from ν_e below 1 GeV $\rightarrow \nu_e/\nu_\mu$ important
- **Neutrino flavour differences** also limiting atmospheric results

SBN

- For SBN programme and appearance searches, anything mimicking ν_e appearance is important
 - e.g. NC1 γ , NC1 π^0 DIS, NC1 π^0 resonant, NC1 π^0 coherent
 - Many constrained by dedicated measurements and sidebands



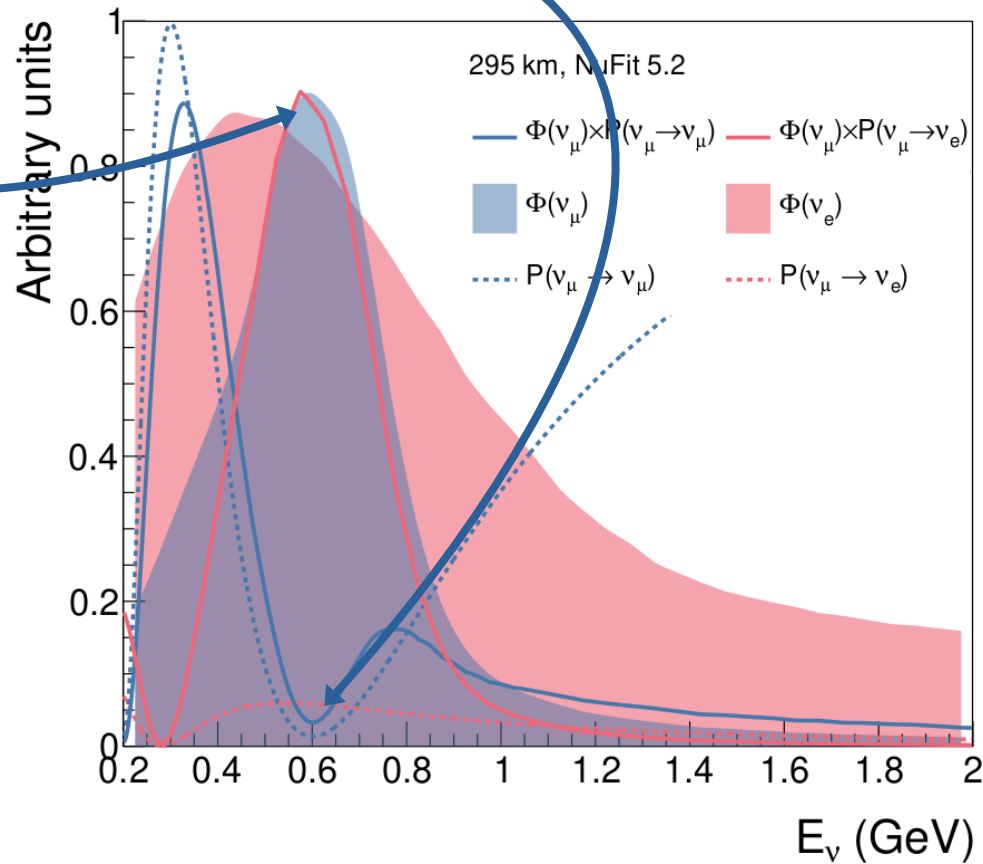
Phys. Rev. Lett.
128, 111801

- ν_e/ν_μ differences from nucleon and nuclear environment, especially considering ^{40}Ar
- Calorimetric energy reconstruction (see later)

Where does the
model dependence enter?

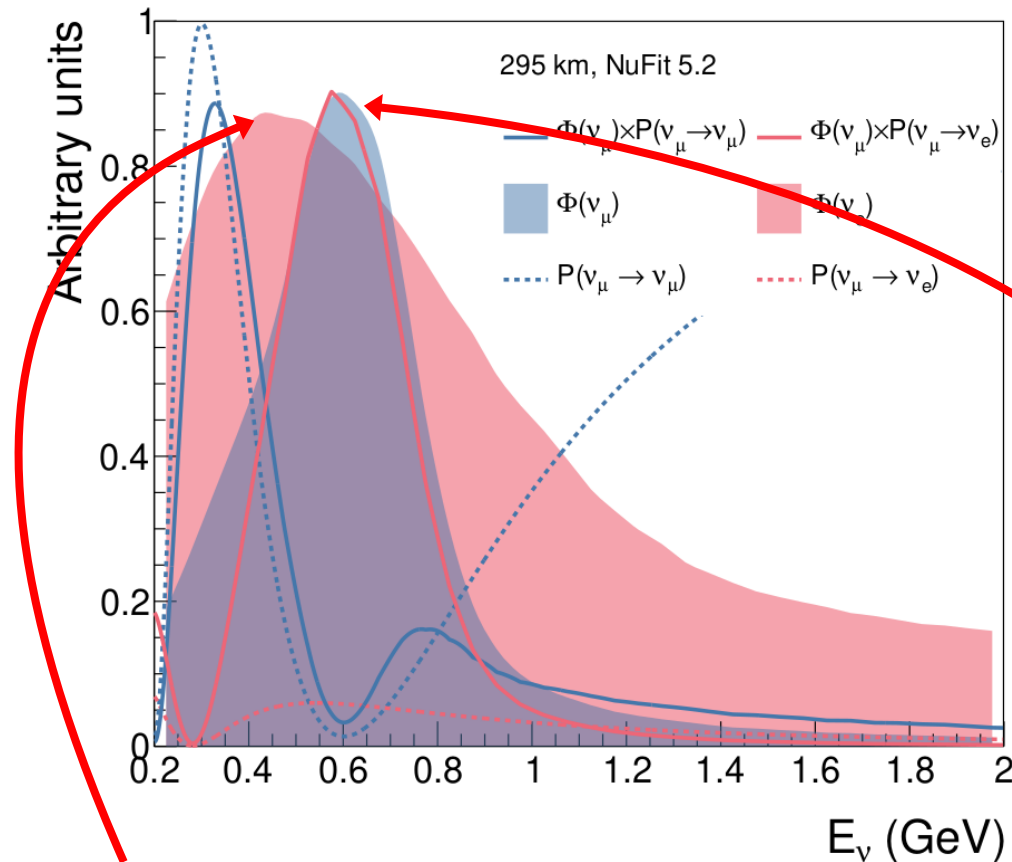
Issues with the near detector

- The ν_μ flux at the **FD** has a **minimum** where the ν_μ flux at the **ND** has a **maximum**



Issues with the near detector

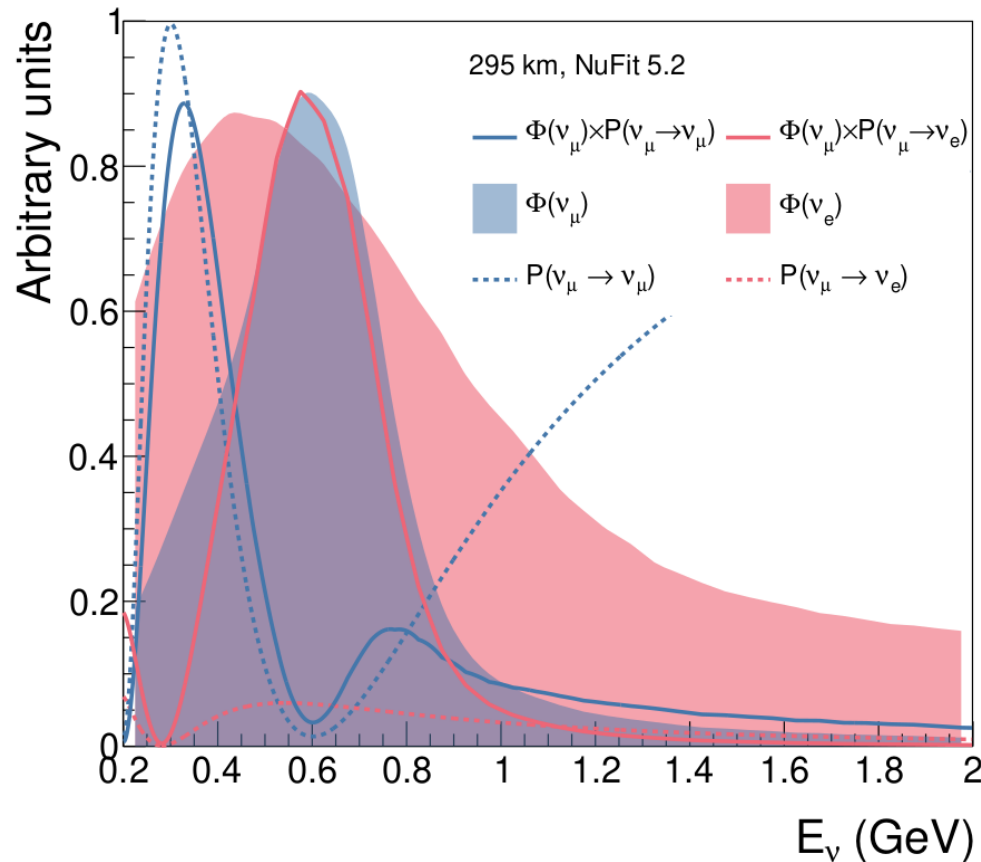
- The ν_μ flux at the **FD** has a **minimum** where the ν_μ flux at the **ND** has a **maximum**



- Similarly, the ν_e flux at the ND does not match the ν_e from $\nu_\mu \rightarrow \nu_e$ oscillations

Issues with the near detector

- The ν_μ flux at the FD has a minimum where the ν_μ flux at the ND has a maximum



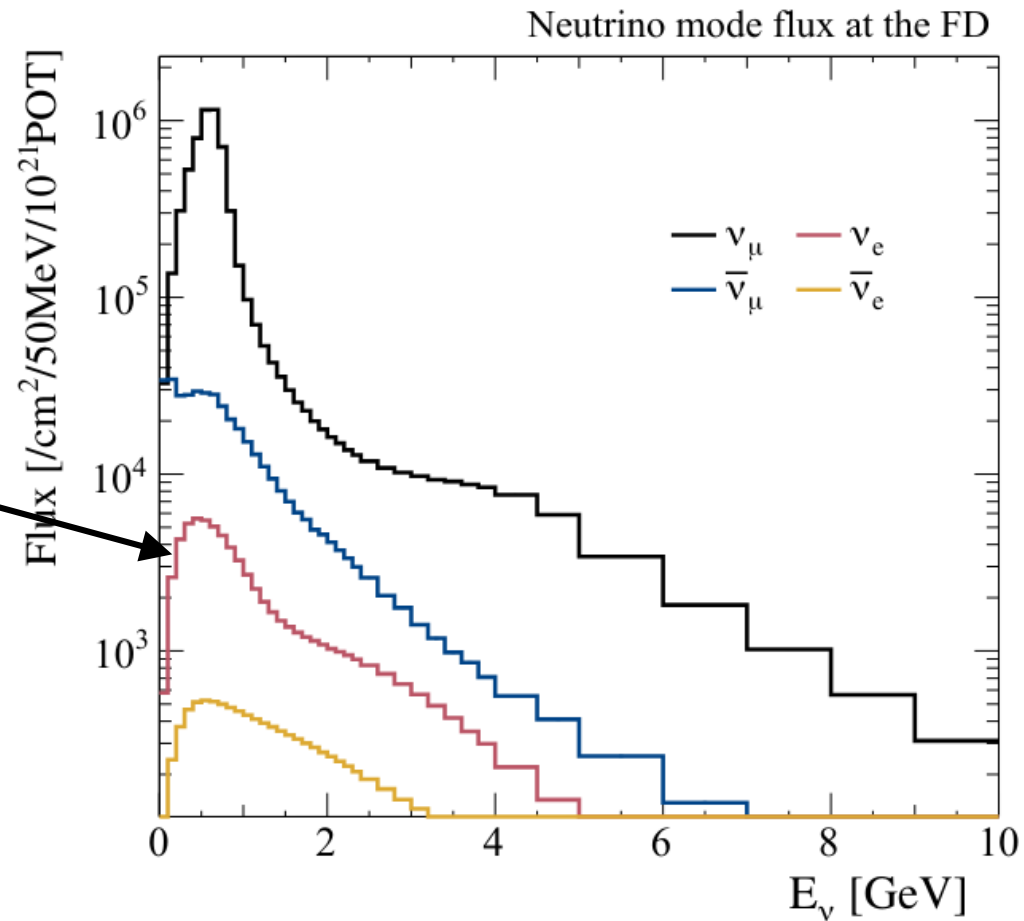
- Similarly, the ν_e flux at the ND does not match the ν_e from $\nu_\mu \rightarrow \nu_e$ oscillations
- Rely on model for extrapolating effects in neutrino energy, and ν_e at ND can't necessarily predict ν_e signal at FD

Electron/muon neutrino differences

- Appearing ν_e have different energy spectrum to ν_e at near detector

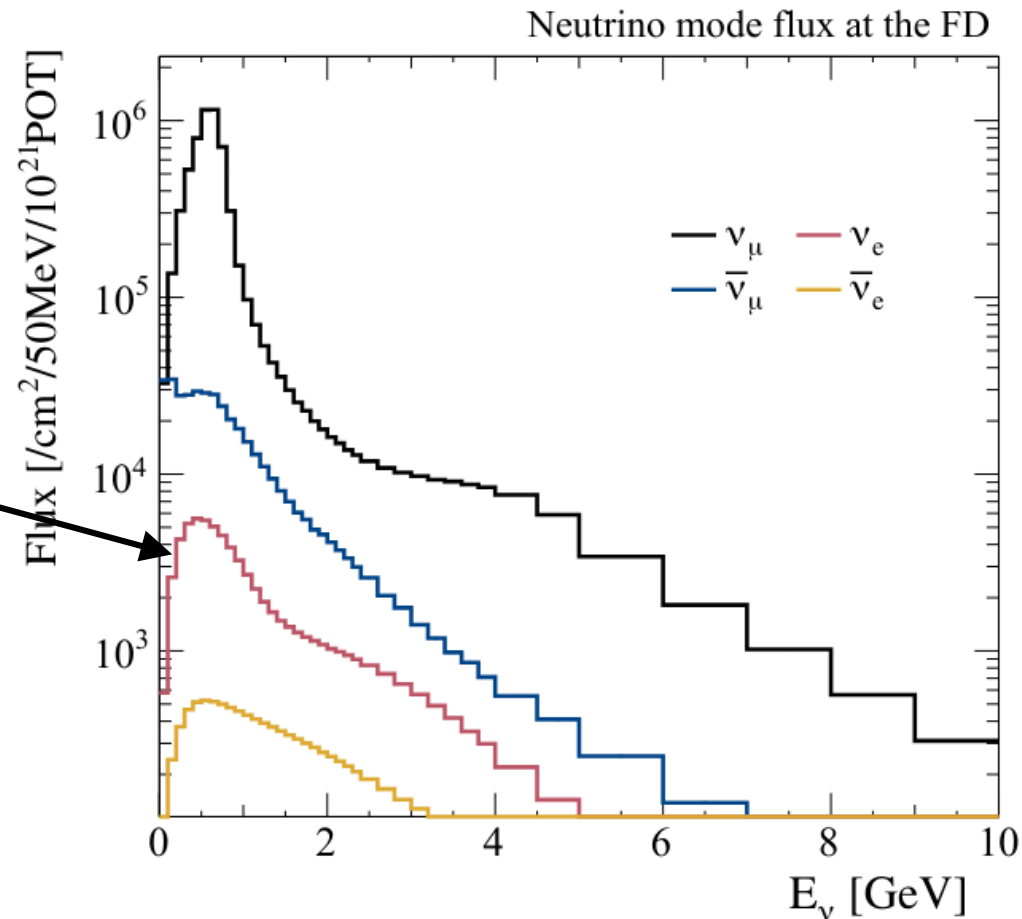
Electron/muon neutrino differences

- Appearing ν_e have different energy spectrum to ν_e at near detector
- ν_e at near detector used to understand **intrinsic** ν_e component in beam (irreducible background ν_e appearance)



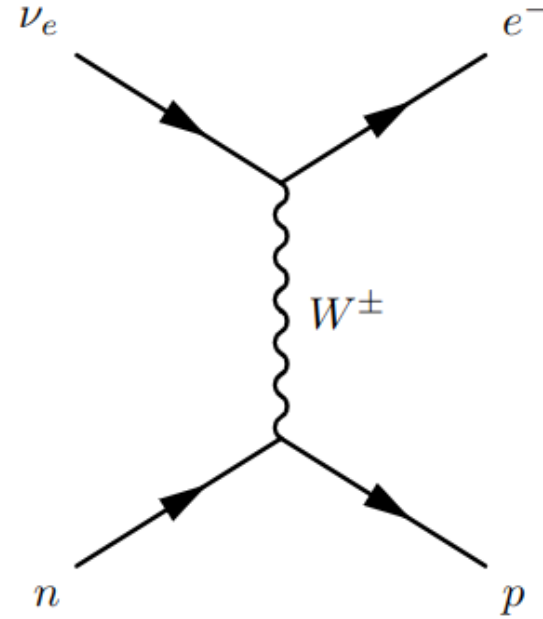
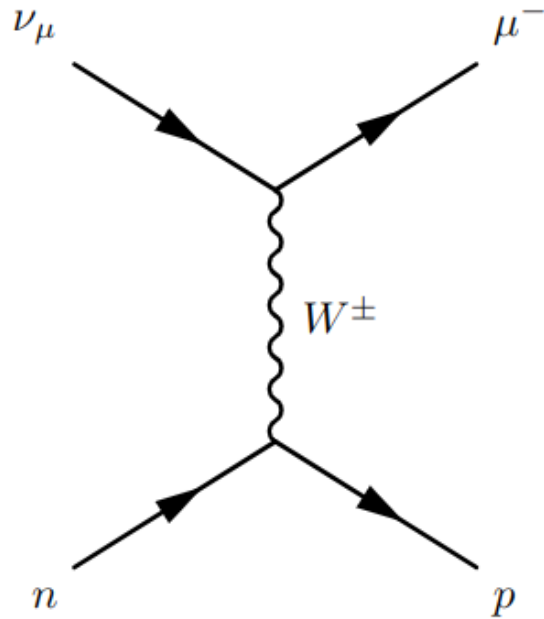
Electron/muon neutrino differences

- Appearing ν_e have different energy spectrum to ν_e at near detector
- ν_e at near detector used to understand **intrinsic** ν_e component in beam (irreducible background ν_e appearance)
- ν_μ at near detector to constrain appearing ν_e (~same flux): **explicit dependence on muon-to-electron mapping**



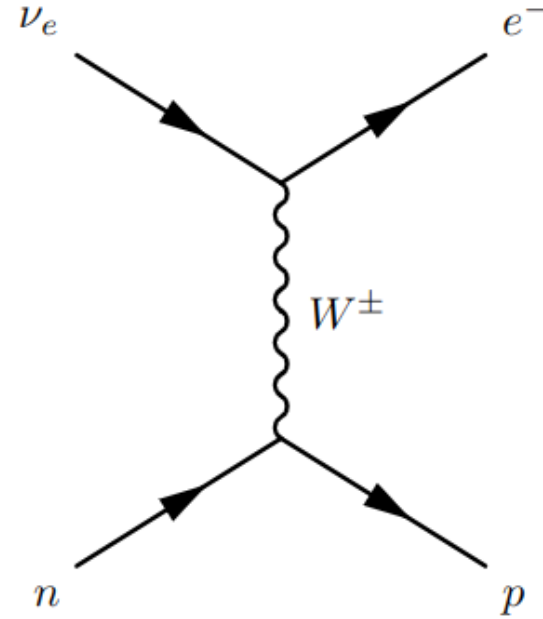
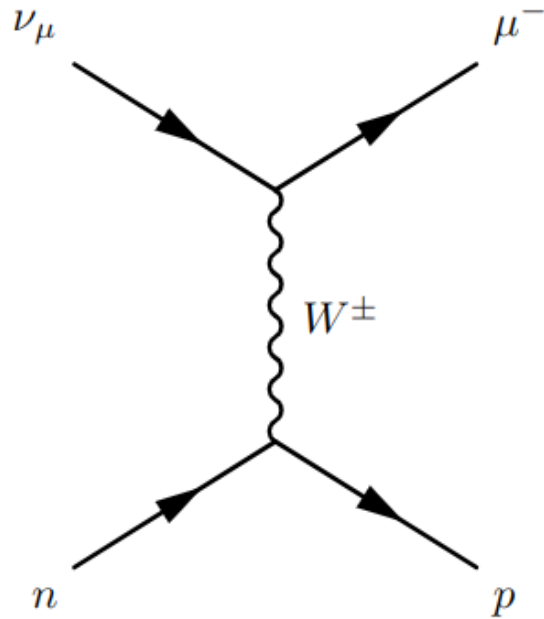
Electron/muon neutrino differences

- On the nucleon-level, pretty simple?



Electron/muon neutrino differences

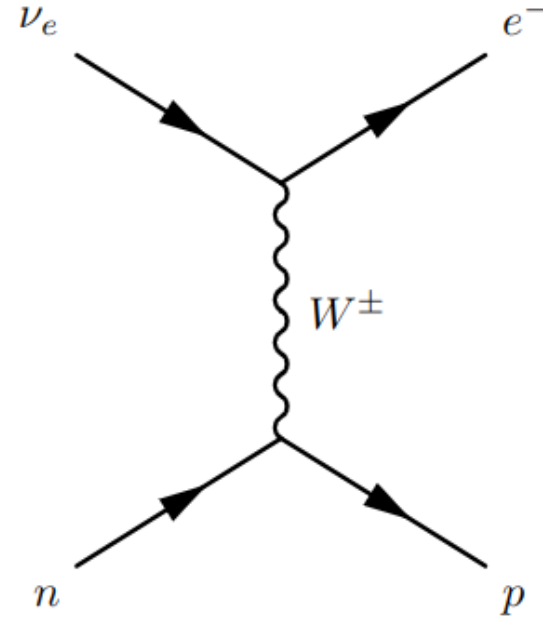
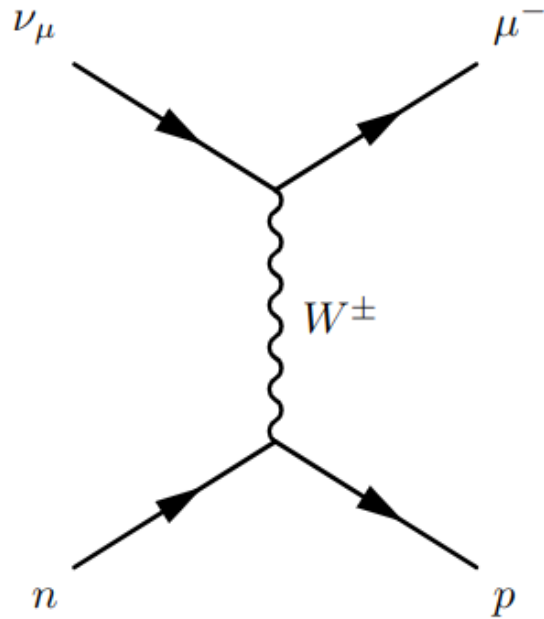
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- Account for **lighter lepton mass**

Electron/muon neutrino differences

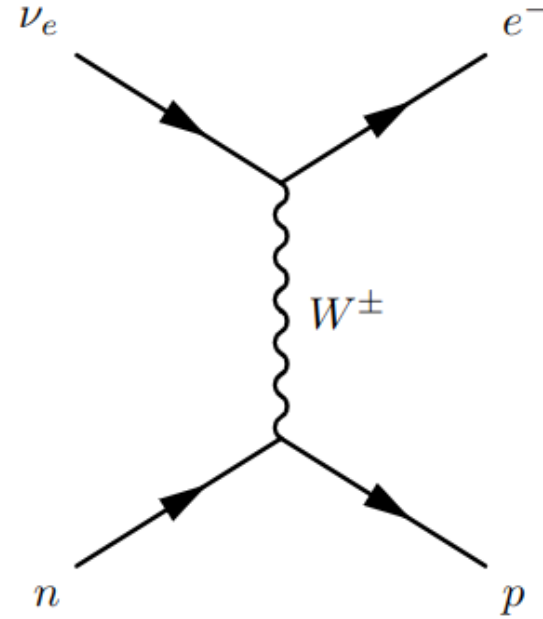
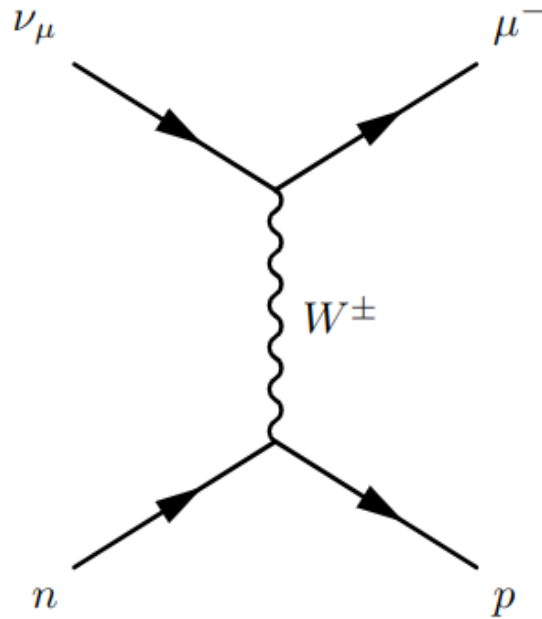
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- Account for lighter lepton mass
- Maybe some **Coulomb repulsion**, shifts energy by $\sim \text{MeV}$

Electron/muon neutrino differences

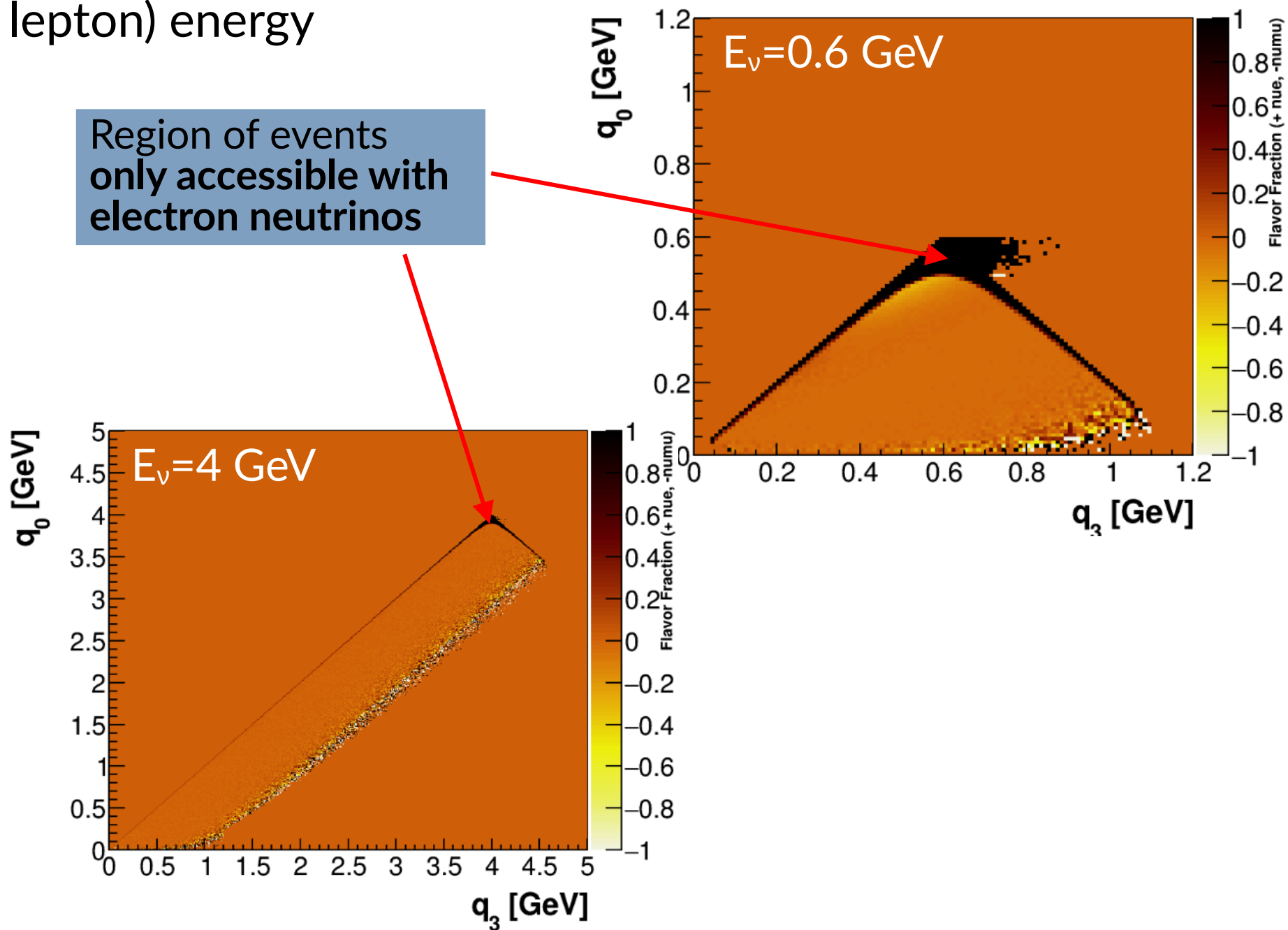
- On the nucleon-level, pretty simple?



- Account for lighter lepton mass
- Maybe some Coulomb repulsion, shifts energy by \sim MeV
- **Radiative corrections**, emitting collinear or virtual photon

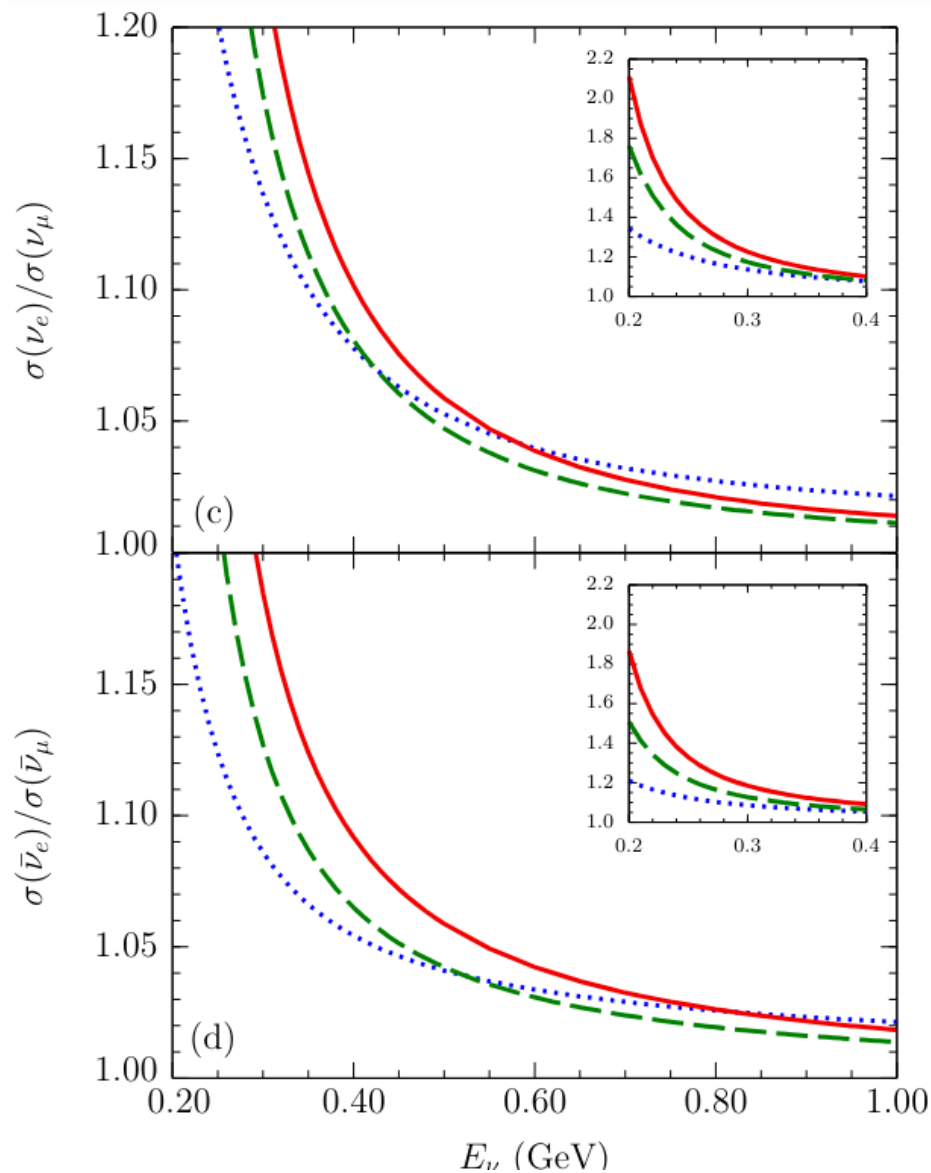
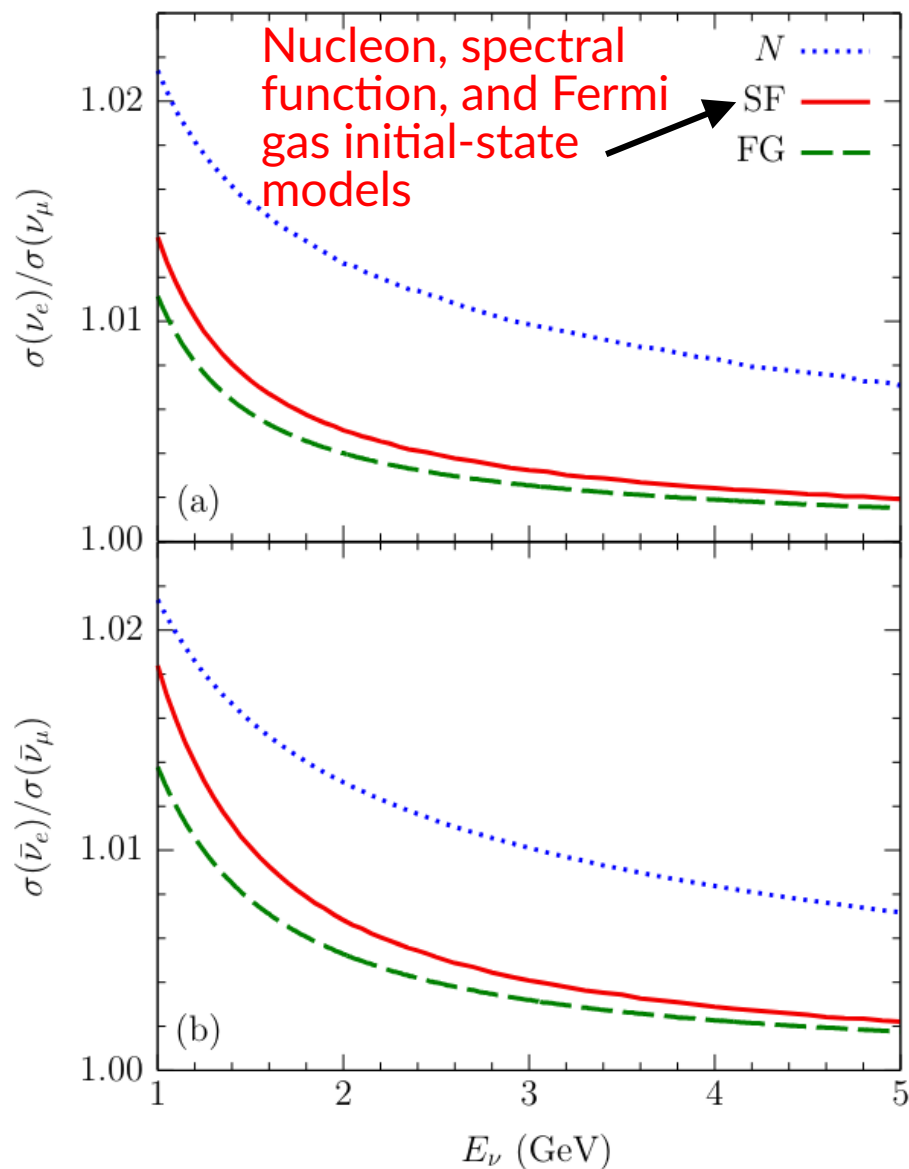
Electron/muon neutrino differences

- Due to lepton mass, access different (q_0, q_3) for a fixed neutrino (or lepton) energy



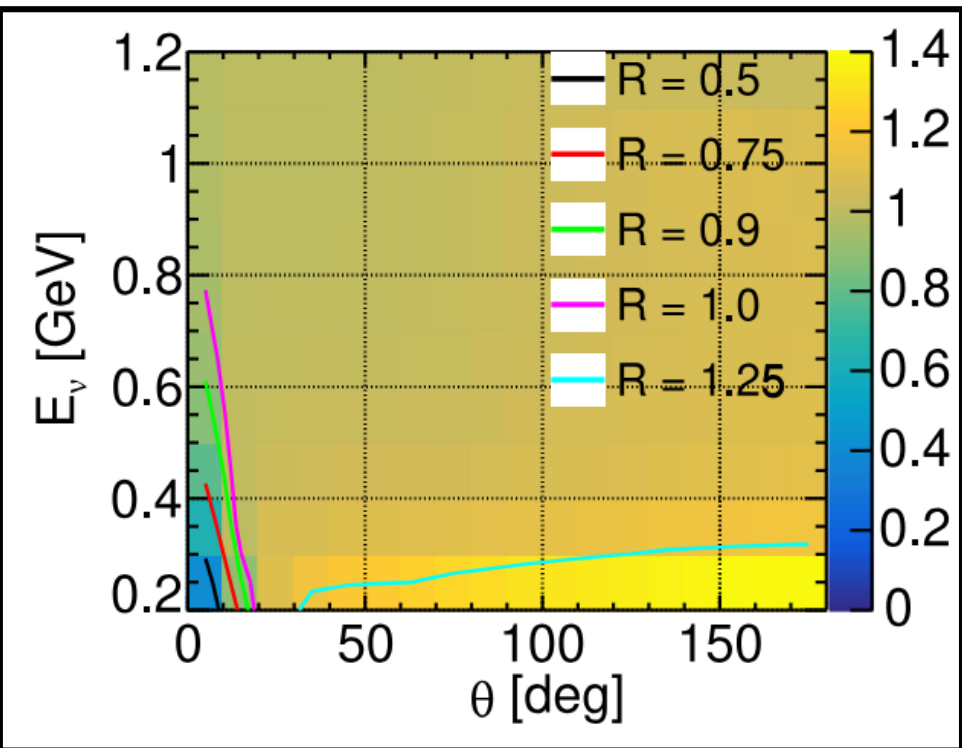
Electron/muon neutrino differences

- Due to lepton mass, access different (q_0, q_3) for a fixed neutrino (or lepton) energy
- Especially bad when the cross-section rapidly rising, and lepton mass is non-negligible energy (T2K)



Electron/muon neutrino differences

- Extension of study, focussing on nuclear models and kinematics accessible in experiments, found similar effects, $\sim 3-4\%$ level



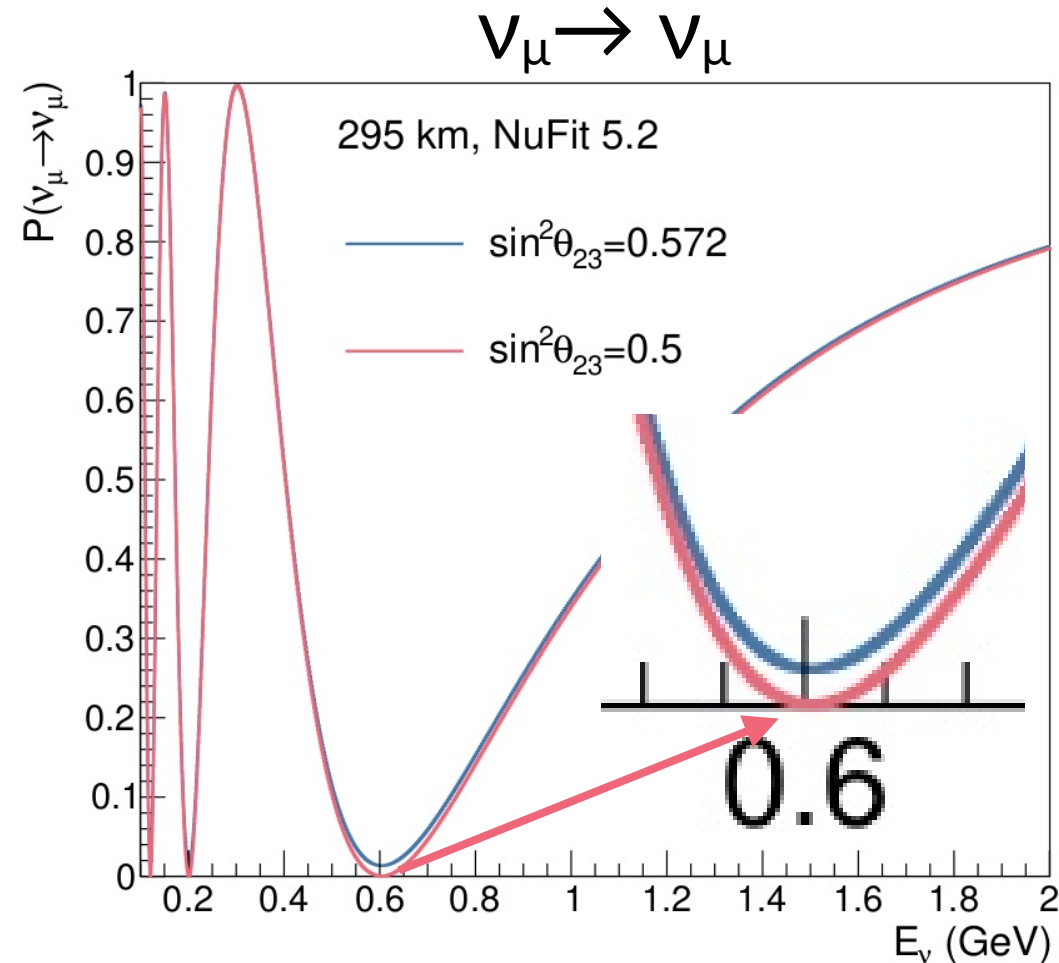
ν_e/ν_μ uncertainty [%]

SUSA	3.2	2.7	2.7	2.4	0.1	0.2	0.0	0.1	
HF-CRPA PW	3.4	2.8	2.9	2.4	0.2	0.0	0.1		0.1
HF-CRPA C	3.2	2.7	2.7	2.3	0.1	0.1		0.1	0.0
HF-CRPA	3.4	2.8	2.9	2.4	0.2		0.1	0.0	0.1
HF	3.2	2.6	2.7	2.2		0.2	0.1	0.2	0.1
SF w/o PB	0.5	0.1	0.1		2.5	2.7	2.6	2.6	2.6
SF	0.4	0.0		0.3	2.6	2.9	2.7	2.8	2.7
SF M_A^{QE} 1.03	0.4		0.0	0.2	2.6	2.8	2.7	2.8	2.7
LFG		0.4	0.4	0.8	3.0	3.2	3.1	3.2	3.1

Different nuclear models

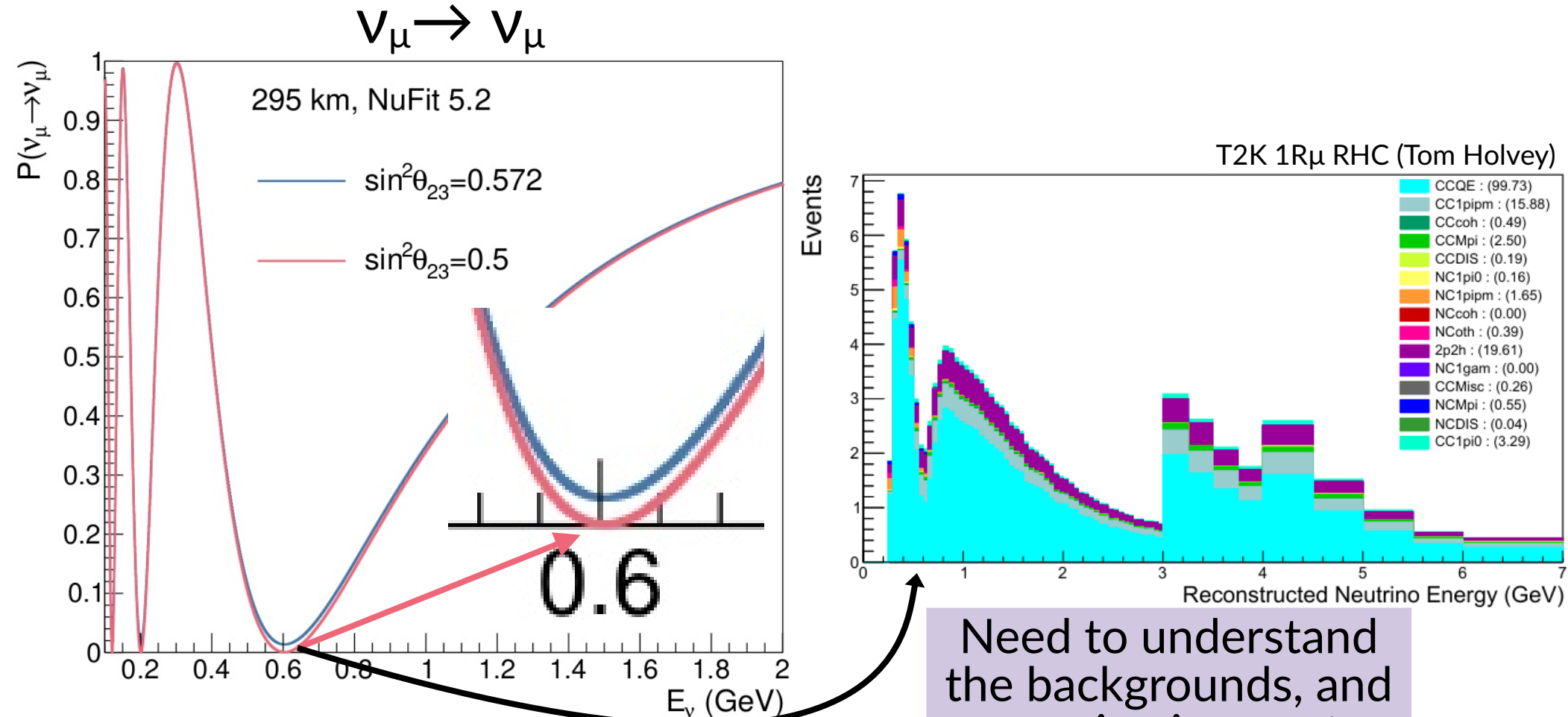
Issues with the near detector

- For accurate measurements of the dip (e.g. $\sin^2\theta_{23}$), the modelling of the few events in the dip becomes important



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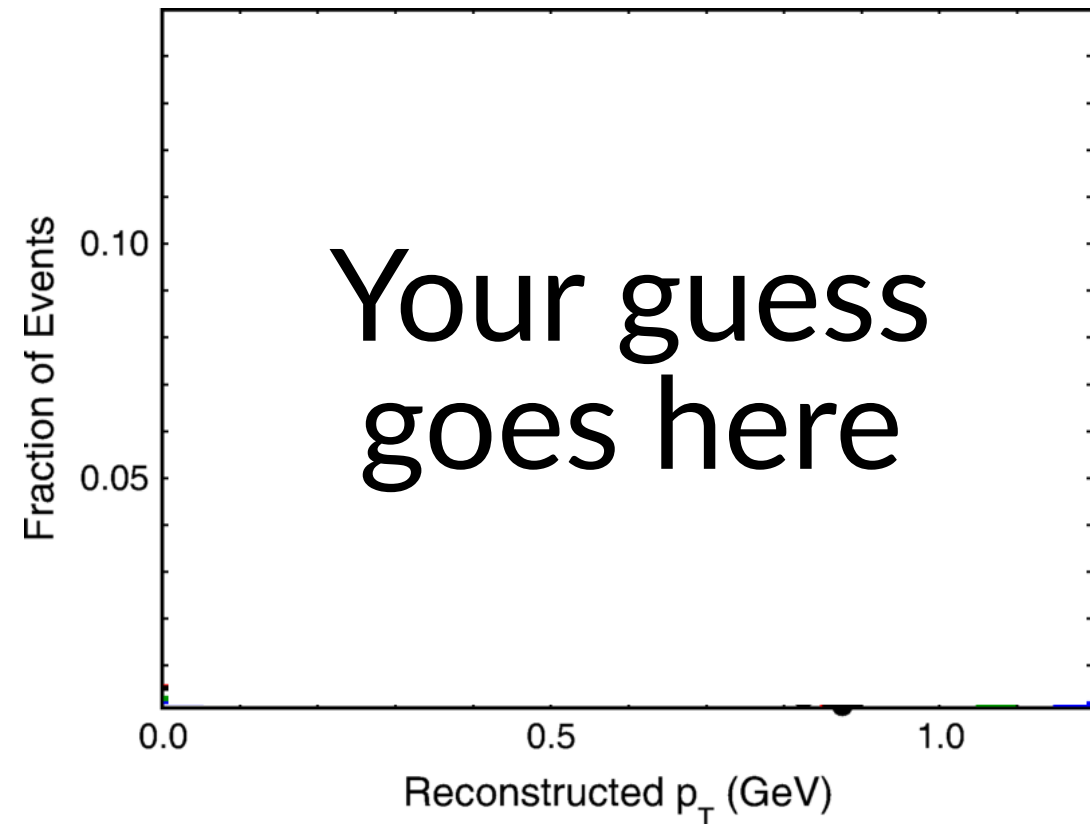
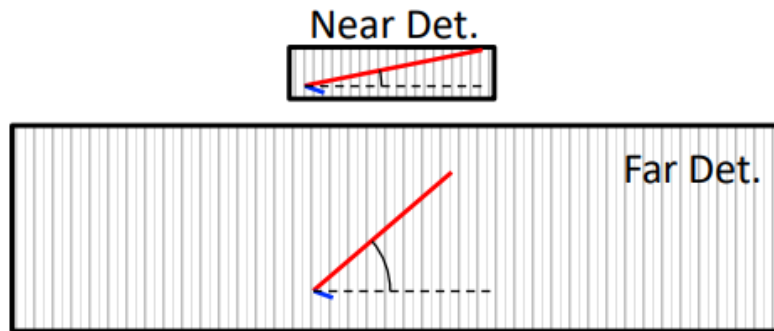


In reality, this minimum never reaches zero

Need to understand the backgrounds, and smearing in reco to true energy

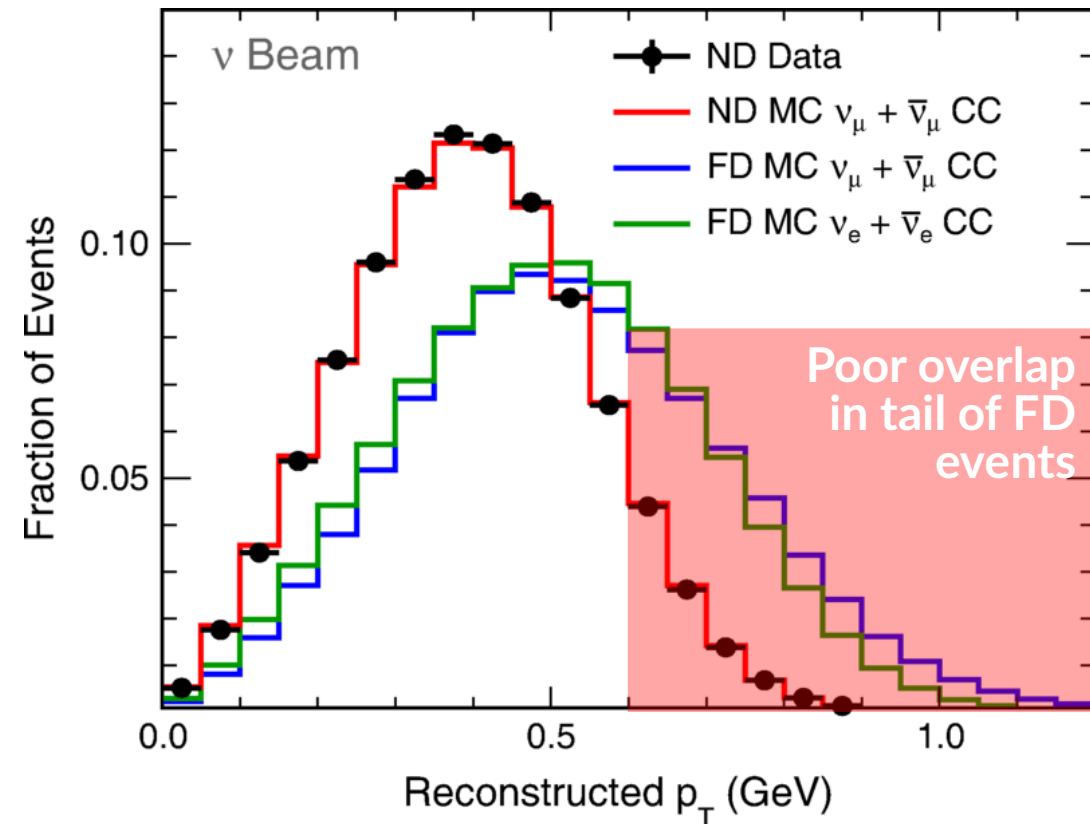
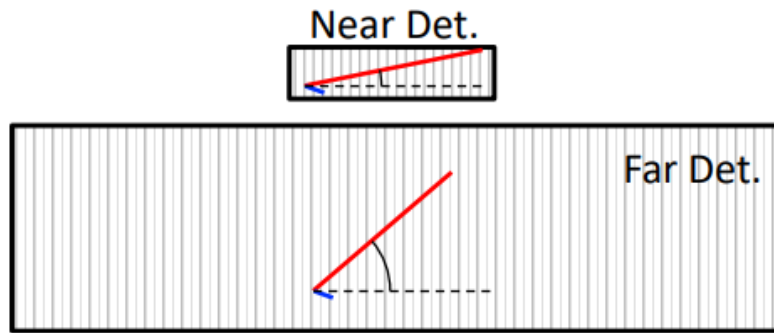
Acceptance mismatch

- Acceptance differences from **differently sized detectors**



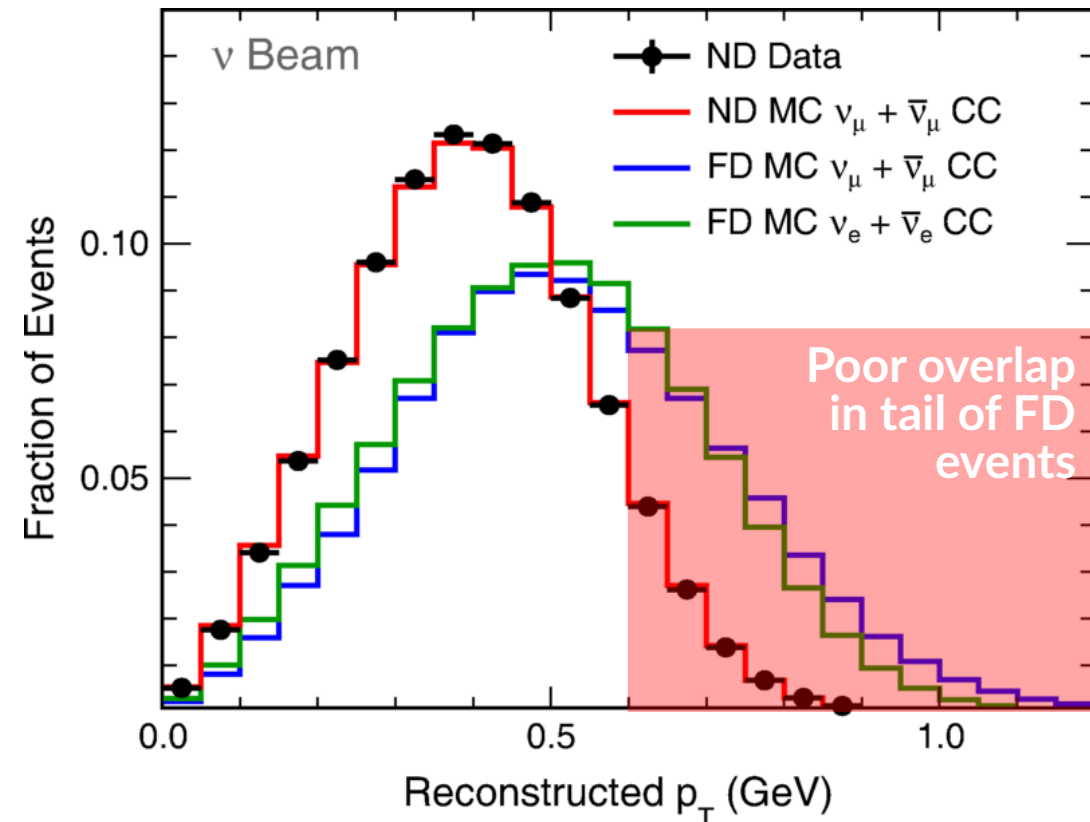
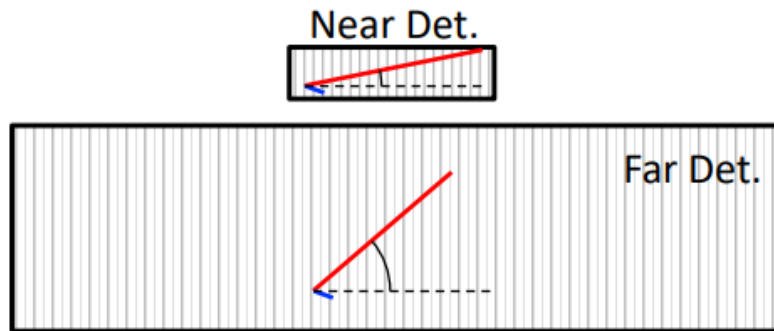
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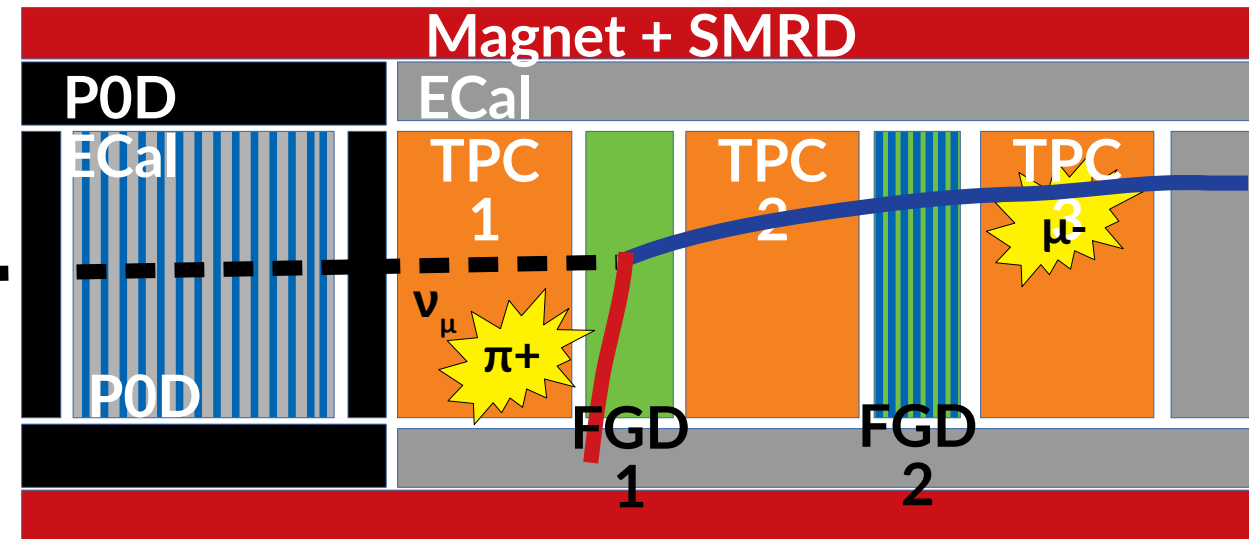
- Acceptance differences from **differently sized detectors**
 - Functionally identical does not mean identical acceptance



- **Different target material and detector design** means additional model dependence in $CH \rightarrow H_2O$
- **Different detector technologies and geometry** may mean different particle acceptance

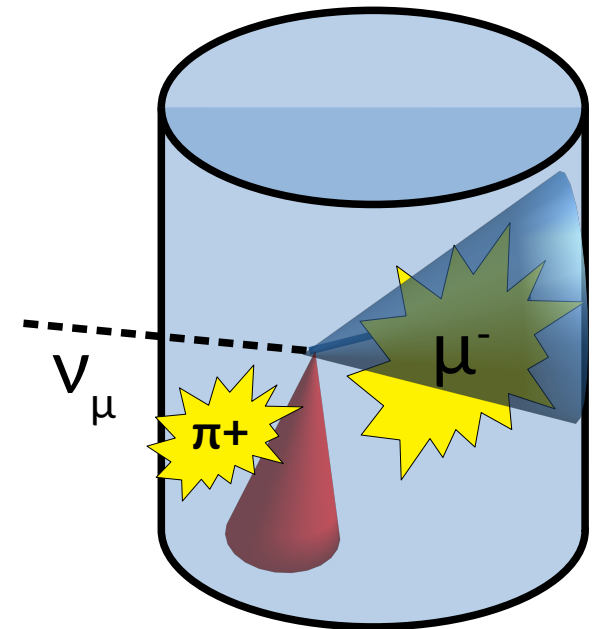
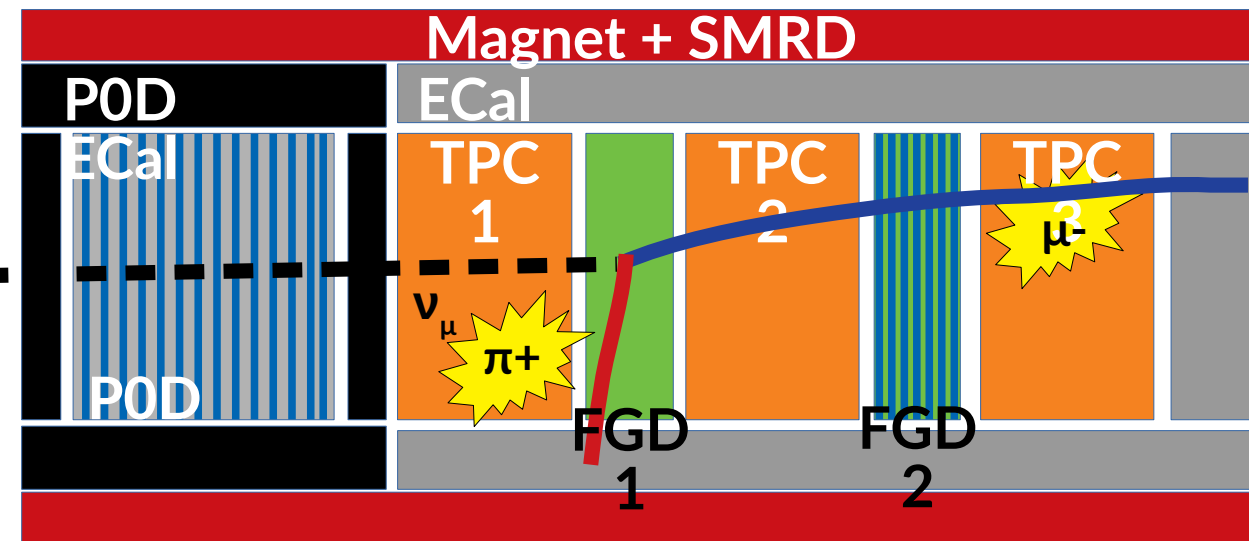
Acceptance mismatch

- Issue is present in T2K too, potentially even larger
 - Near detector very forward-oriented
 - High-angle tracks challenging to reconstruct



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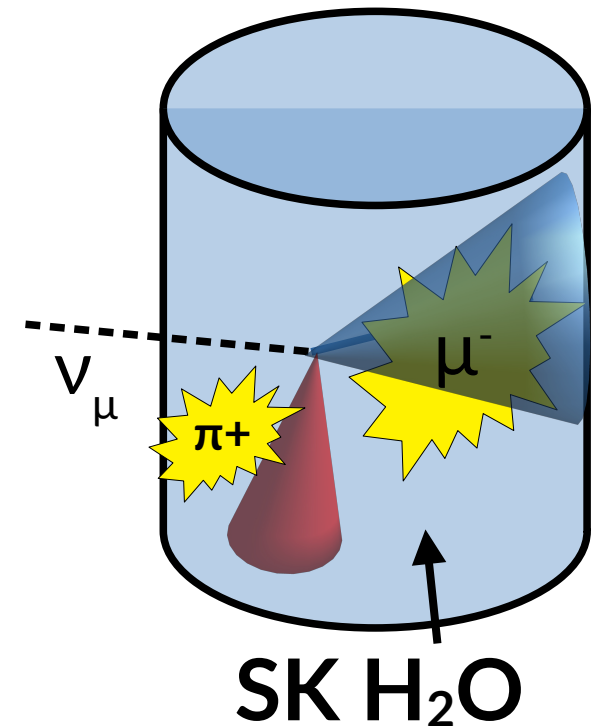
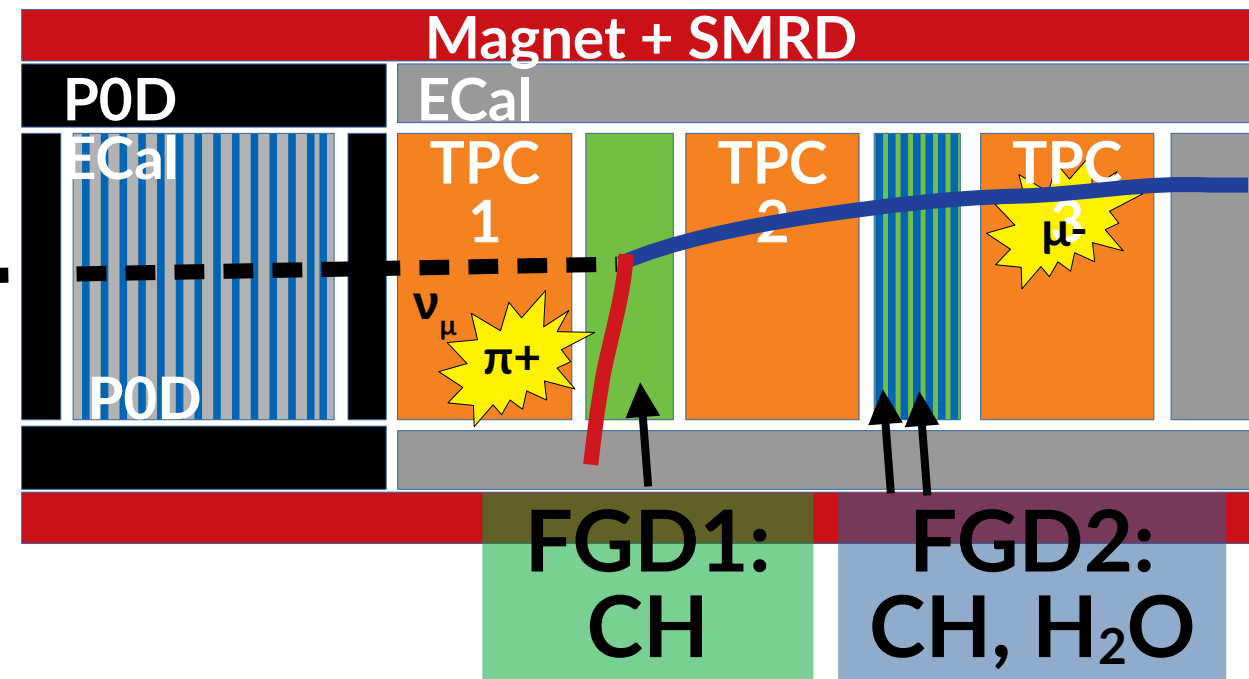
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- SK is instead very symmetric and isotropic
 - Good acceptance forward, backward, upward and downward

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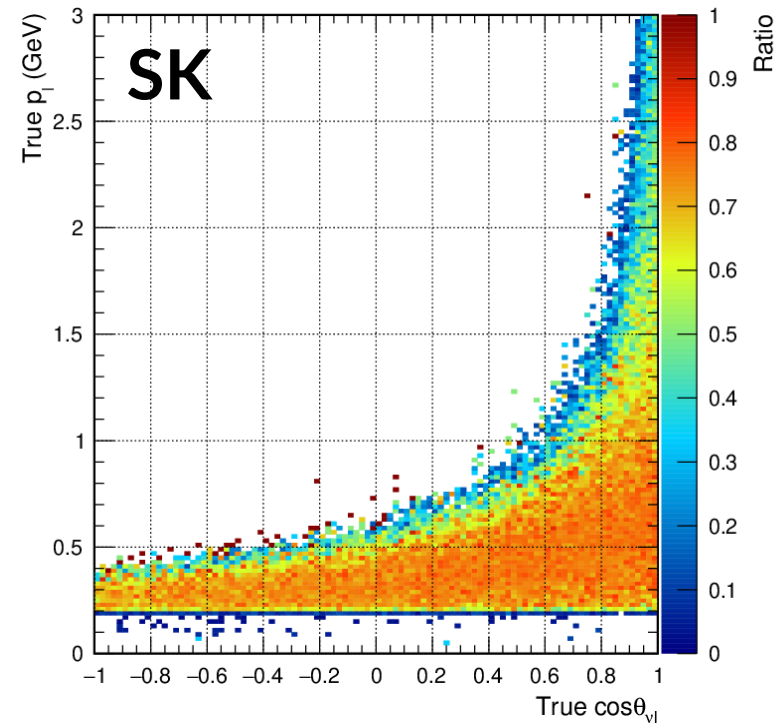
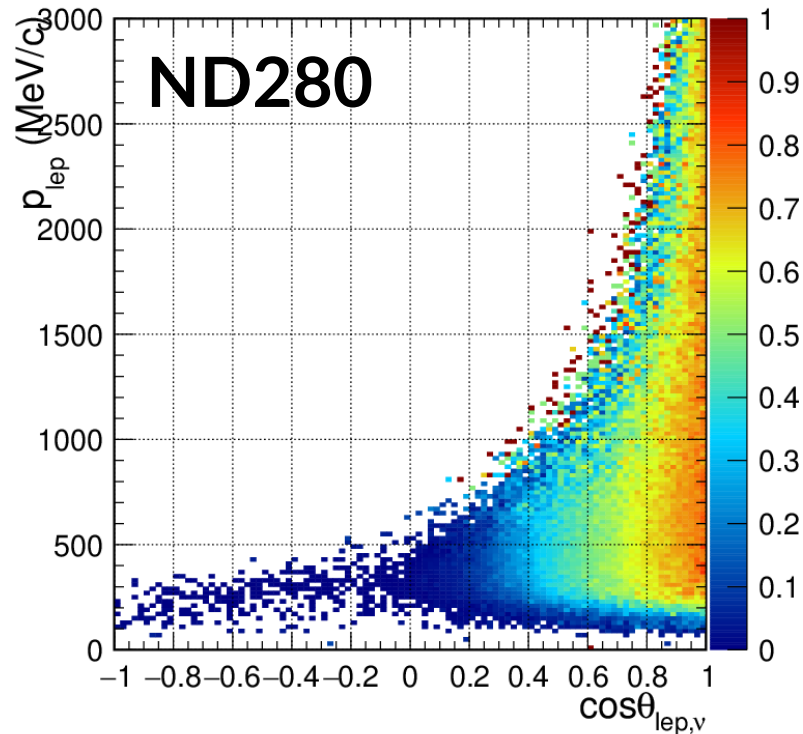
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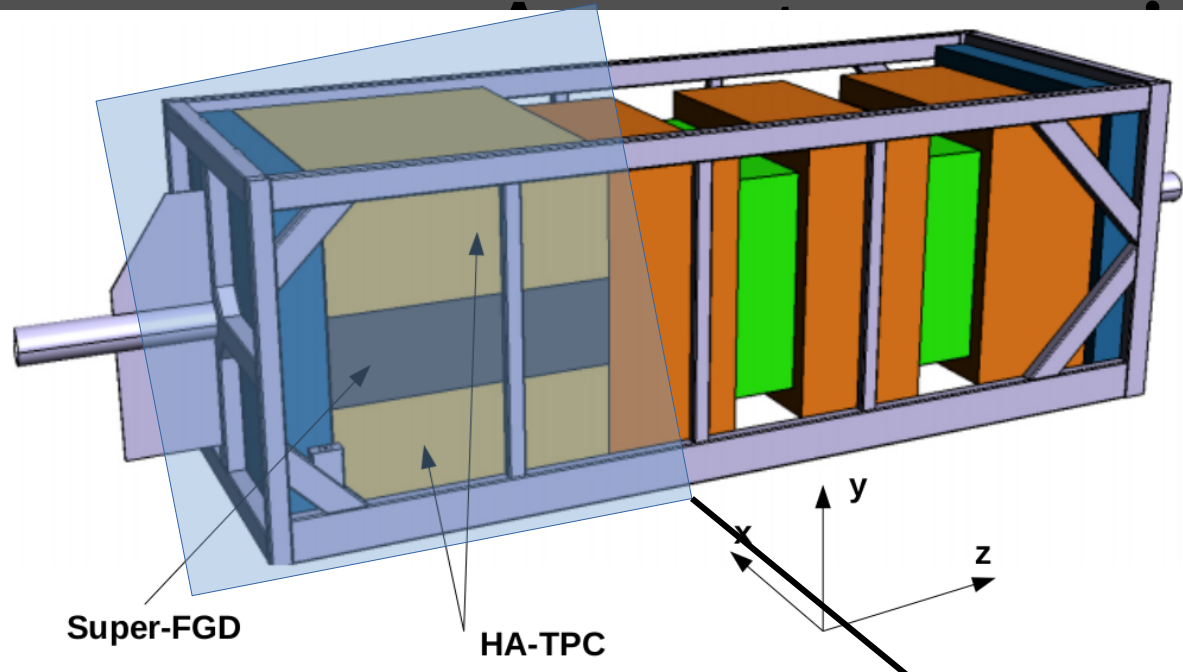
- SK is instead very symmetric and isotropic
 - Good acceptance forward, backward, upward and downward, and a water target

Acceptance mismatch

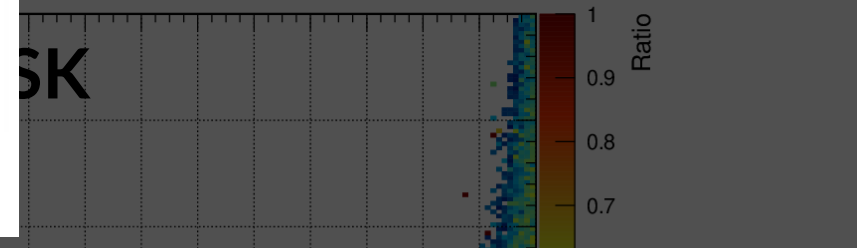
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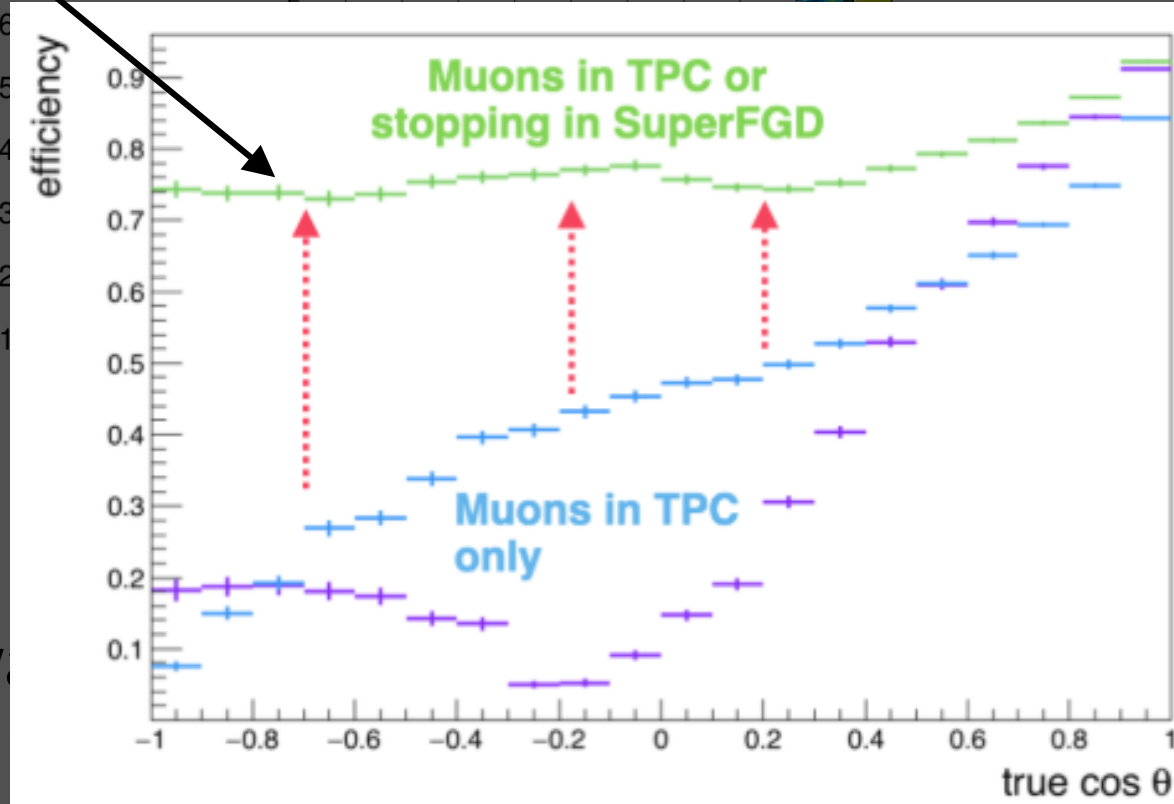
smatch
 essentially even larger
 ted
 reconstruct



Poor acceptance at high and backward angles in ND280 largely addressed with ND280 upgrade

SK is instead very sym

– Good acceptance forward downward

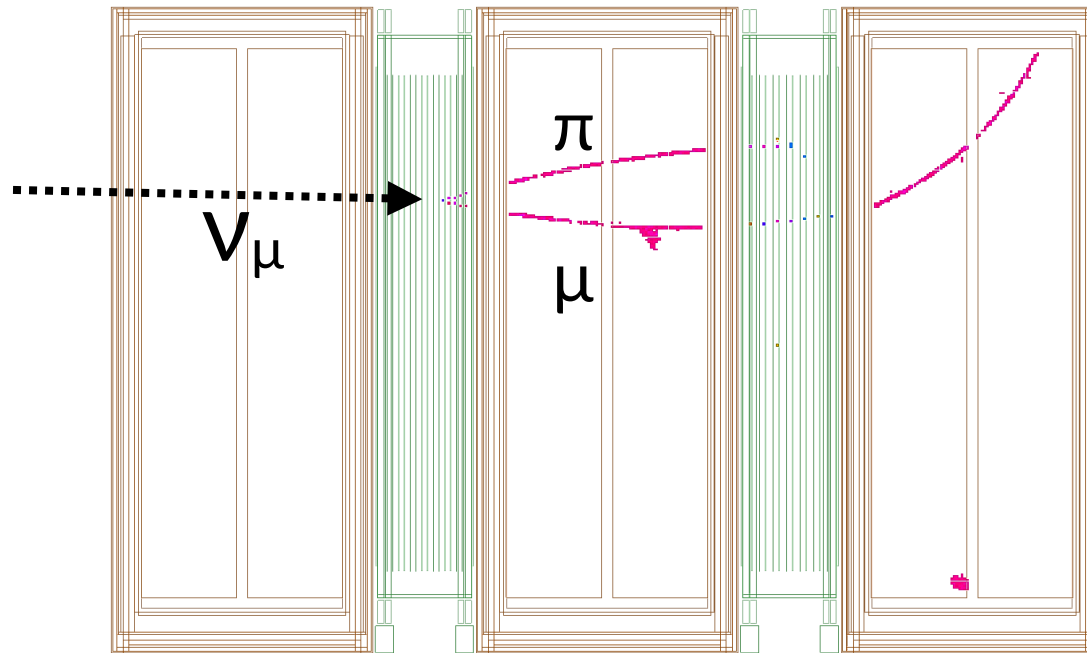


Acceptance mismatch

- Use forward-going events to model backward-going events
 - If this correlation is poorly modelled, issues!
- Similar argument goes for counting particles
 - If particles were emitted backwards in ND280, poorly reconstructed background
- DUNE's near and far detectors will have similar issues to NOvA
- Intermediate Water Cherenkov Detector (IWCD) addresses this on HK
 - Basically a small Super-K near detector

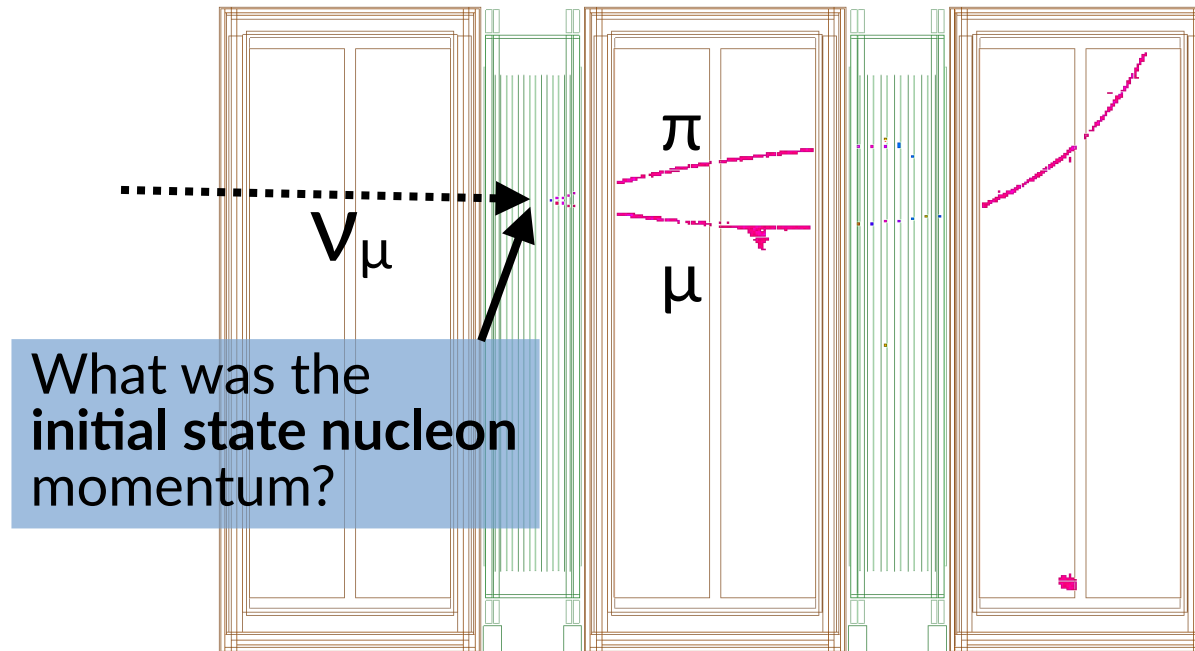
Energy reconstruction

- Energy reconstruction method is function of **selection** and **detector technology**
- Need to understanding mapping between **observed events** and the **not-observed neutrino energy**



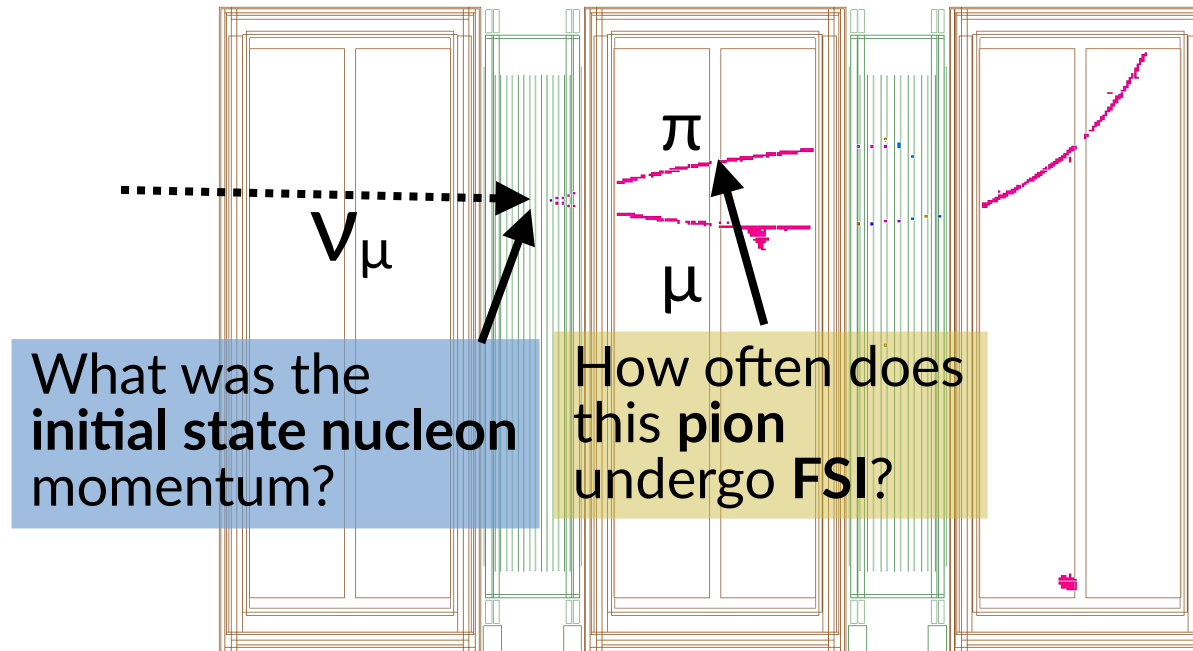
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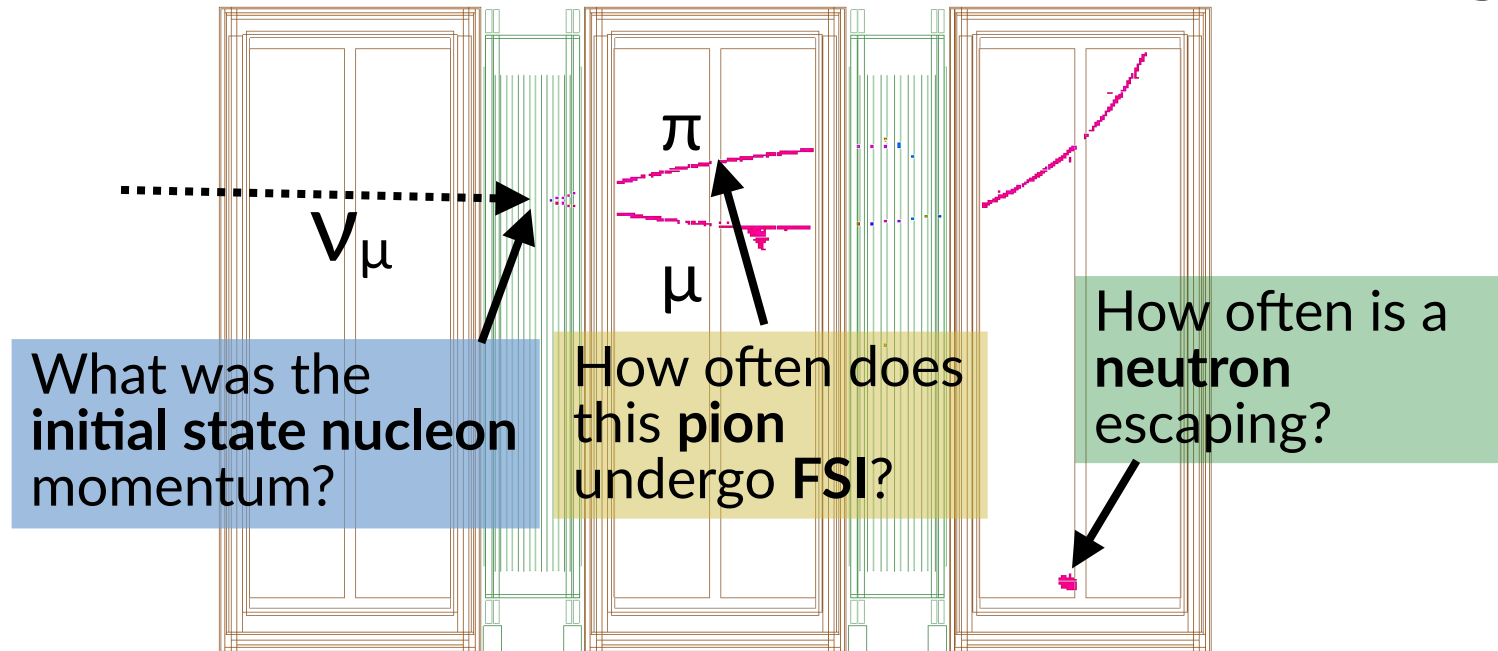
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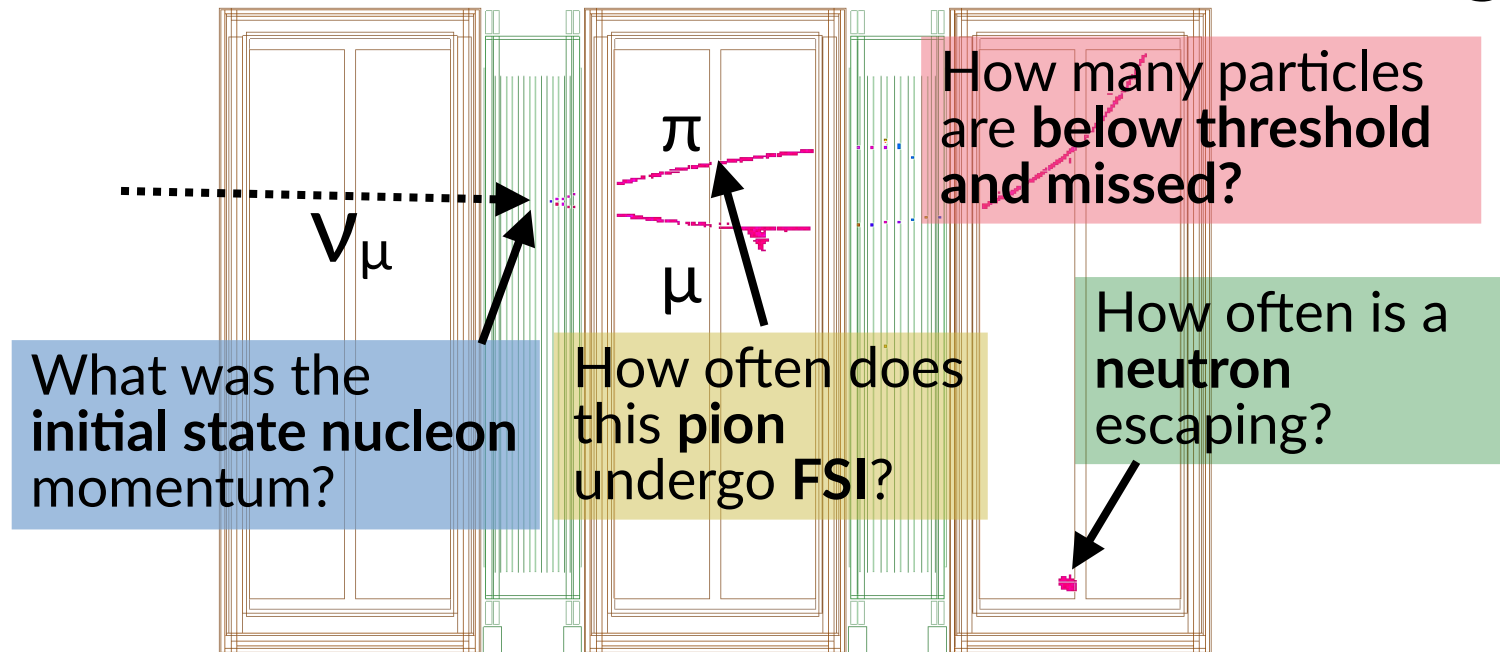
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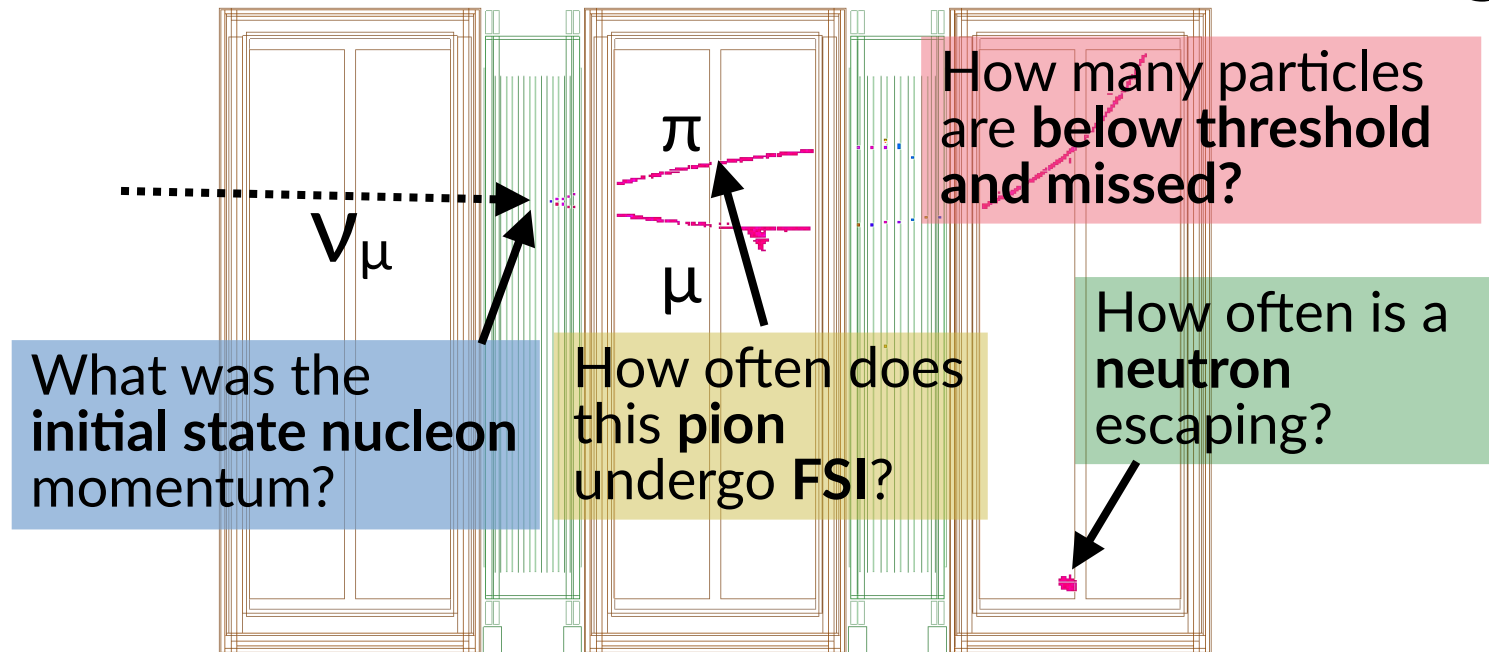
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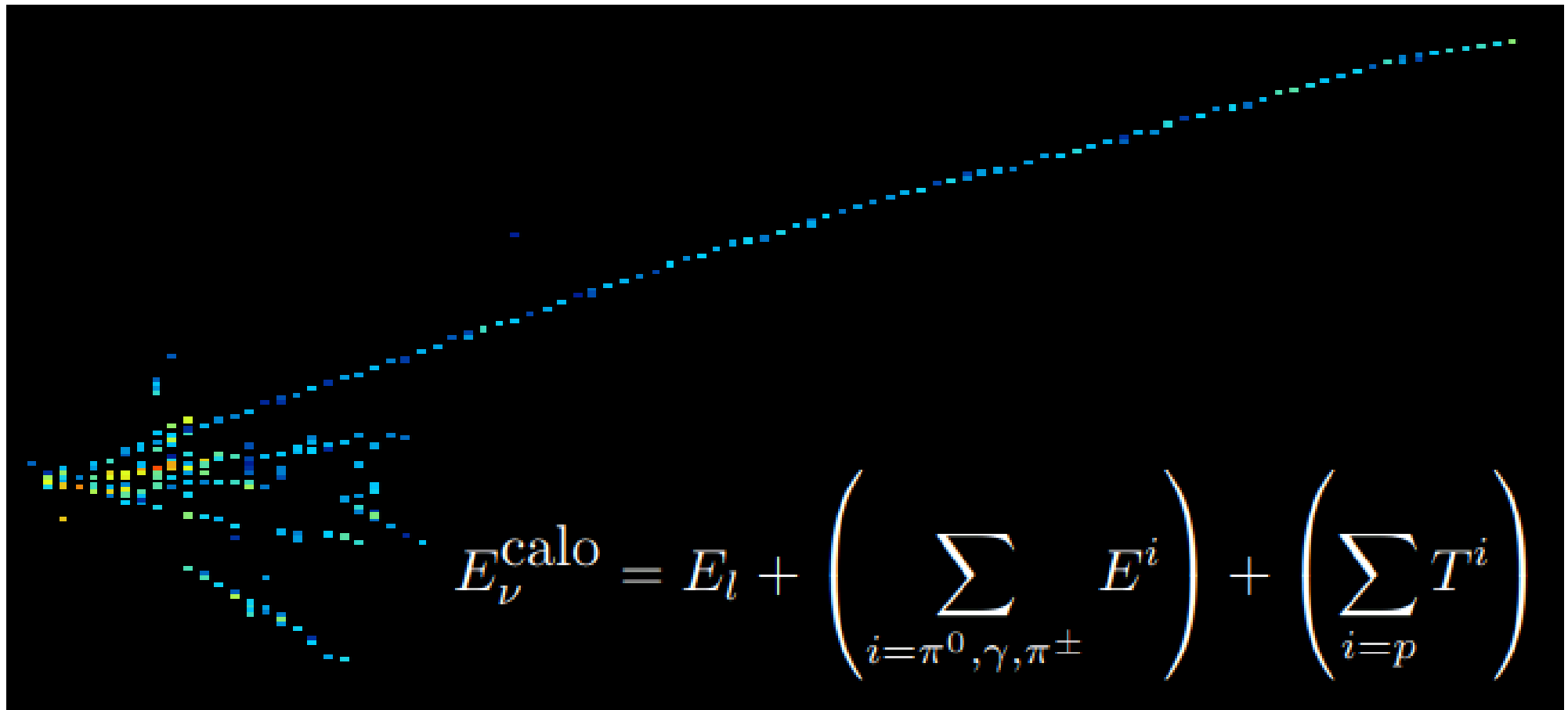
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- All estimators are biased
 - Try to **reduce** the amount of bias
 - Understand the **uncertainty on the bias**

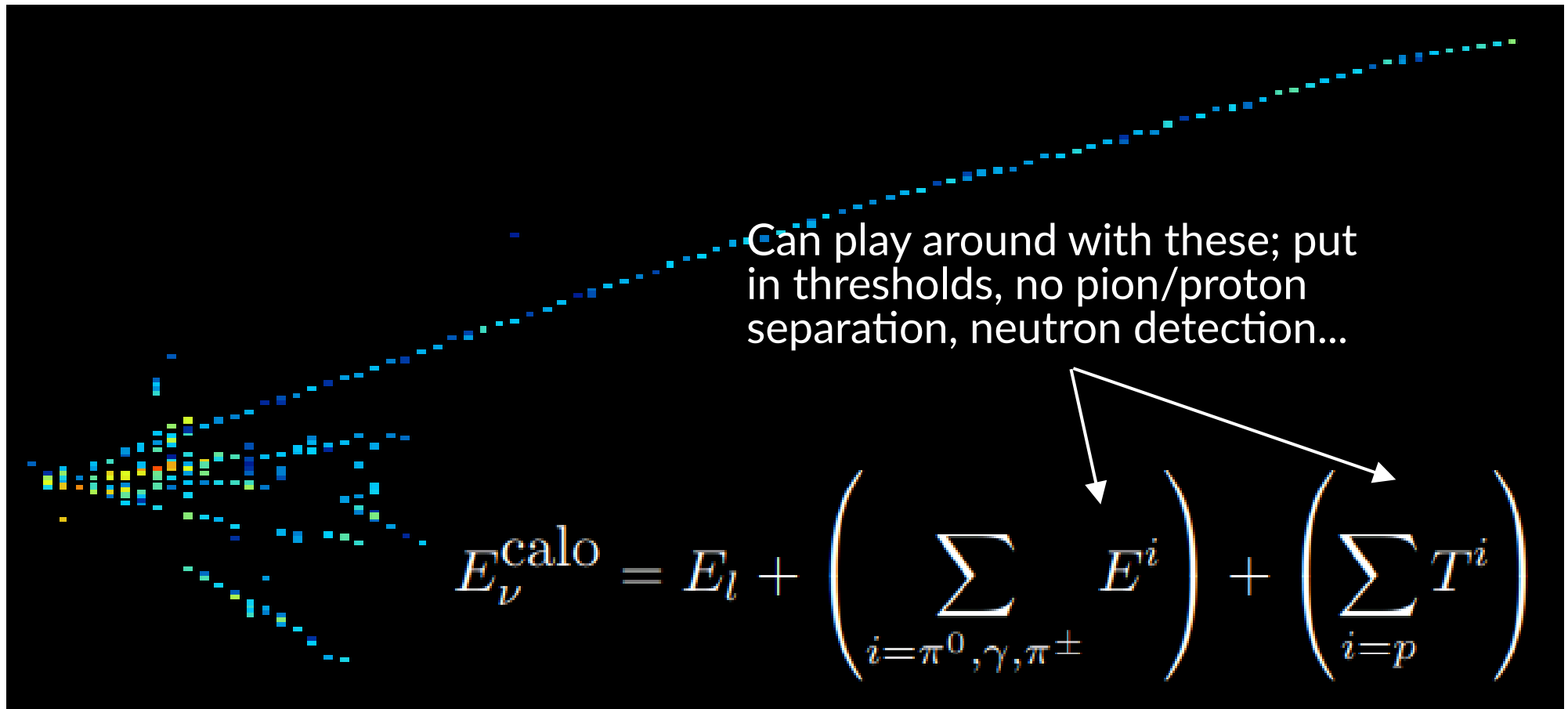
Calorimetric energy reconstruction

- NOvA, DUNE and SBN have **sampling calorimeters** and often **events with multiple tracks**
 - **CC-inclusive** selection
 - Energy estimator which **sums up energy deposits**



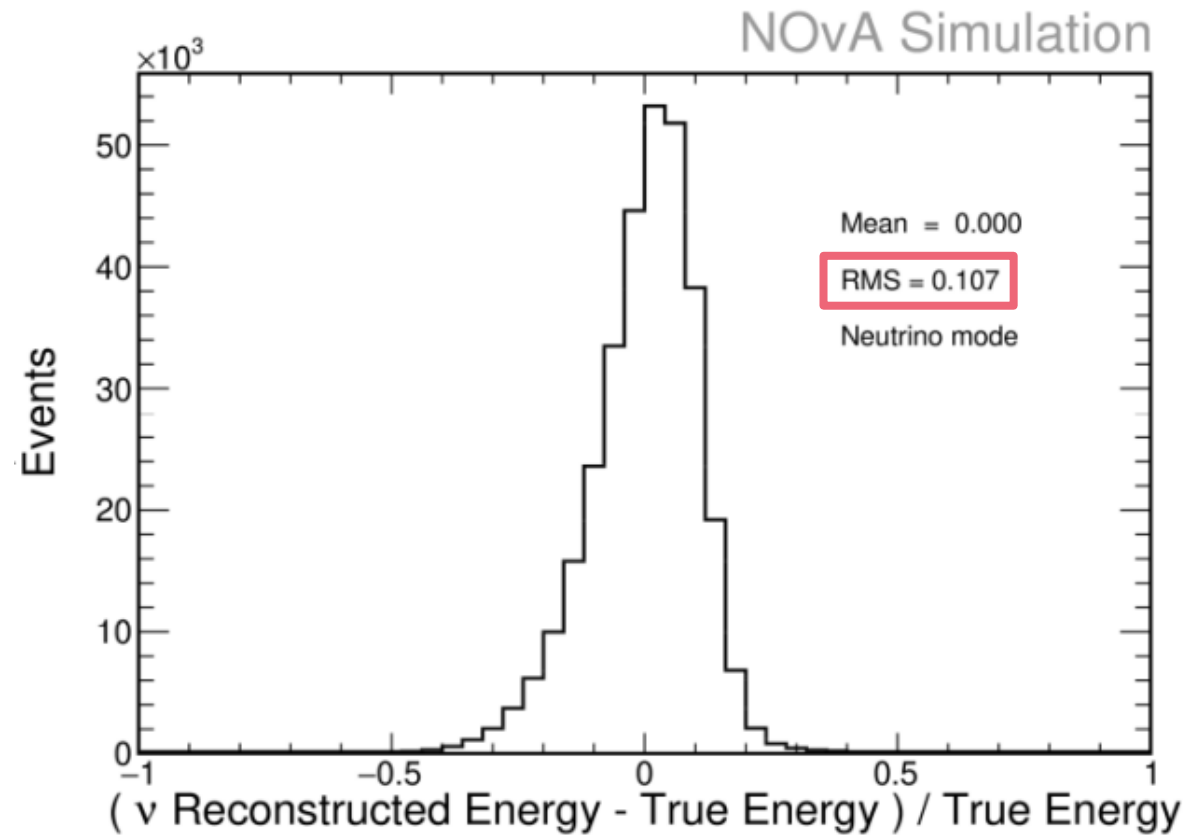
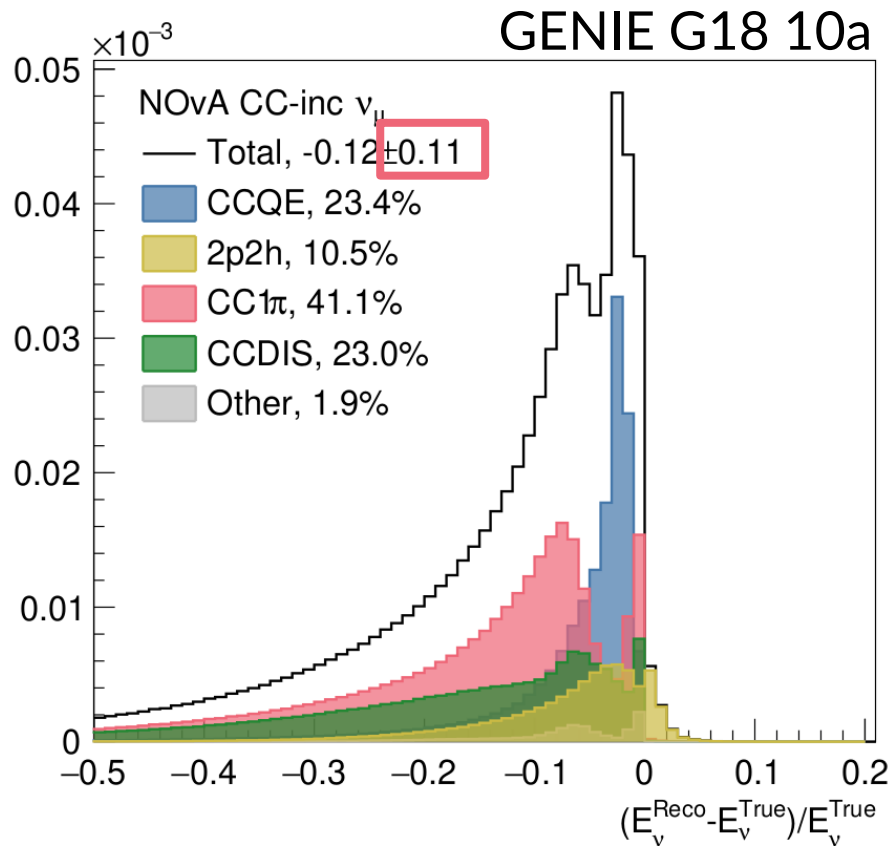
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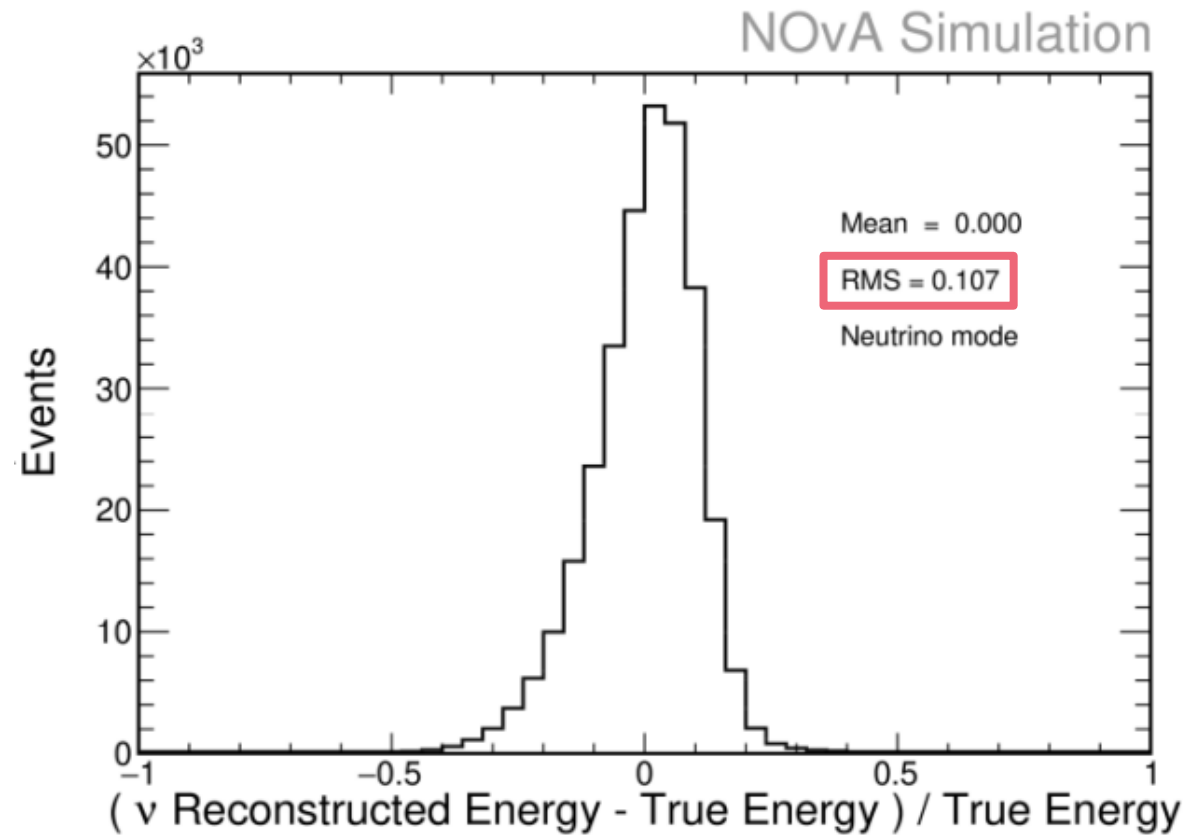
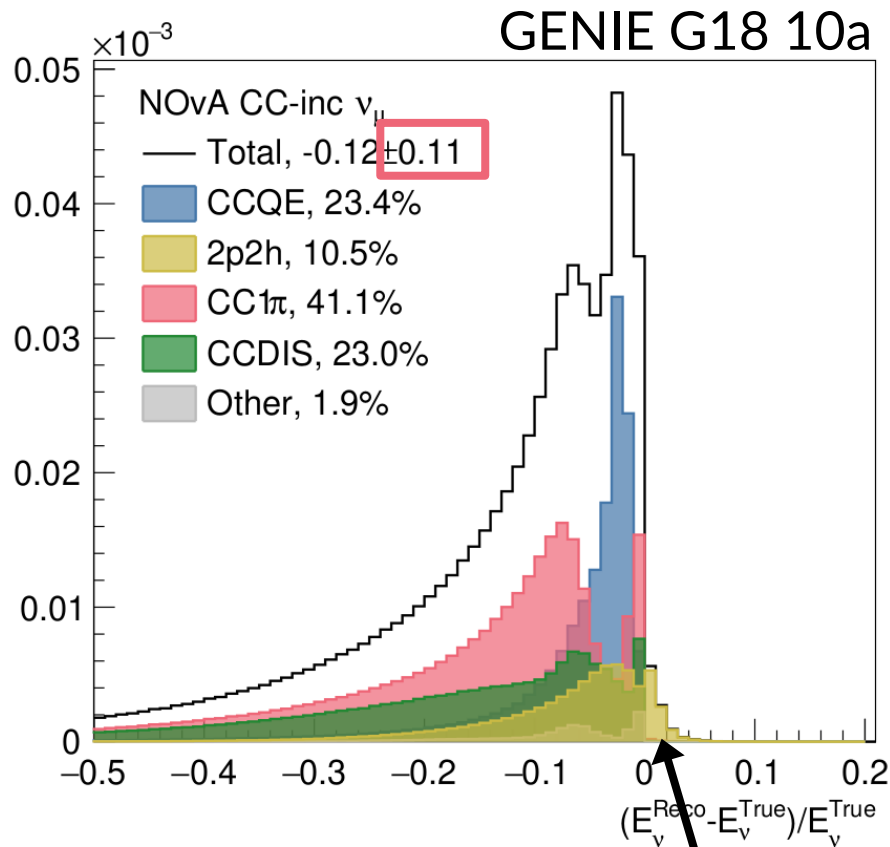
Calorimetric energy reconstruction

- Simple simulation result agrees well with NOvA official figure: ~11% RMS



Calorimetric energy reconstruction

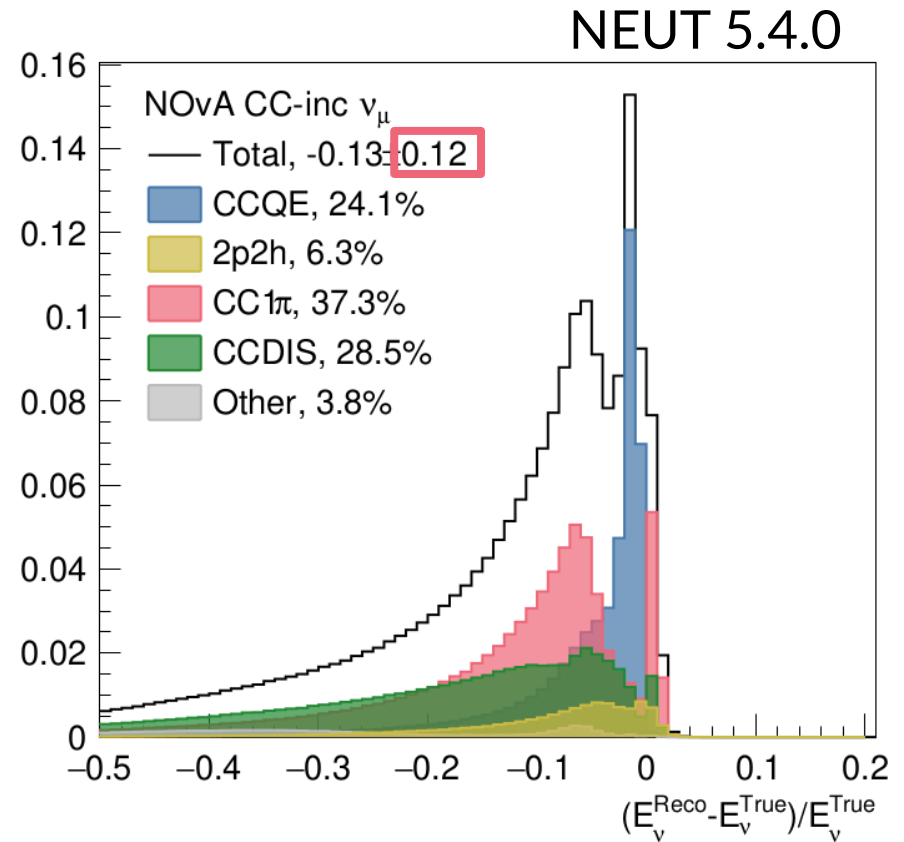
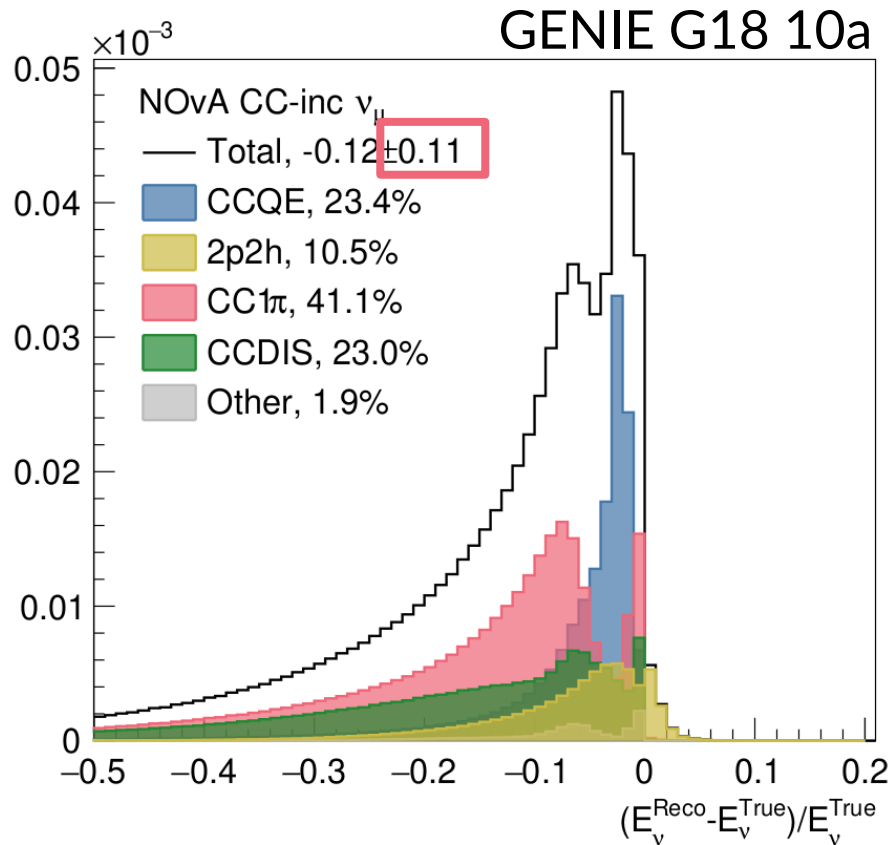
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- Interaction modes bias differently, e.g. DIS has multiple neutrons and and pion that may undergo FSI

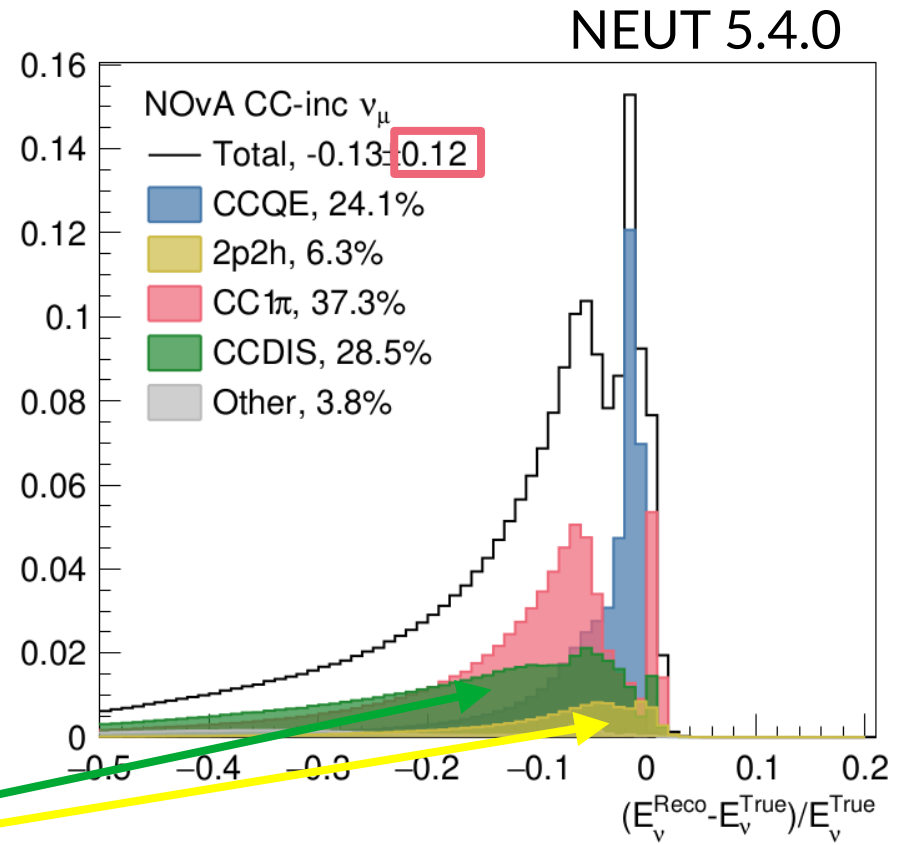
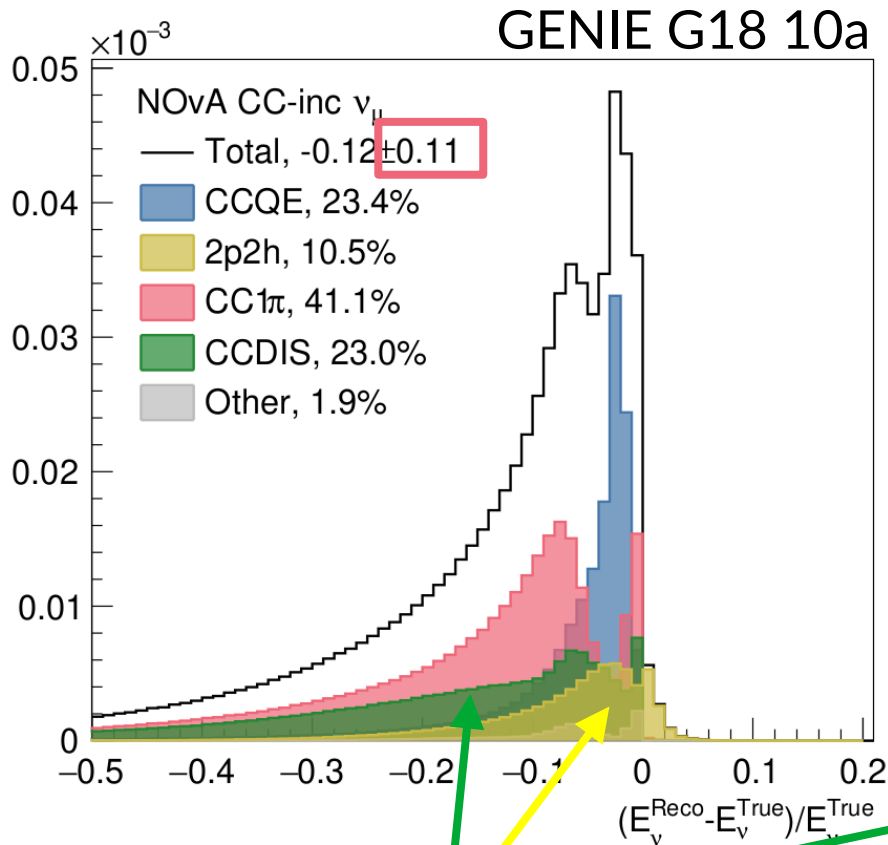
Calorimetric energy reconstruction

- Use a different generator (NEUT), approximately the same result as GENIE G18 10a



Calorimetric energy reconstruction

- Use a different generator (NEUT), approximately the same result as GENIE G18 10a



- Or... is it the same result?

- Bias in the tail clearly different; source of uncertainty

Calorimetric energy reconstruction

- Generally **more precise energy estimate** than kinematic method (shown next)
- Susceptible to **missing neutrons** and other particles
- **Final-state interactions** directly bias the estimator
 - Absorption, charge exchange, nucleon knock-out, energy lost from rescattering
- Relies on **correct PID of every track**, otherwise risk bias by rest mass (e.g. mistake proton for pion)
- Will always have bias from **initial state motion**
 - Smaller impact at higher energies, e.g. NOvA and DUNE
- CC-inclusive selection means **complex contributions from multiple interaction modes**
 - Especially for DUNE and NOvA (many interaction modes)

Paraphrasing from Stephen:
 “When we look at external data, the lepton kinematics are often OK. But the hadrons are a mess!”

Absorption, charge exchange, nucleon knock out, energy loss from rescattering
 General rule of thumb
 Generators vs data: a horror story



Lepton kinematics (except maybe at low energy transfers)

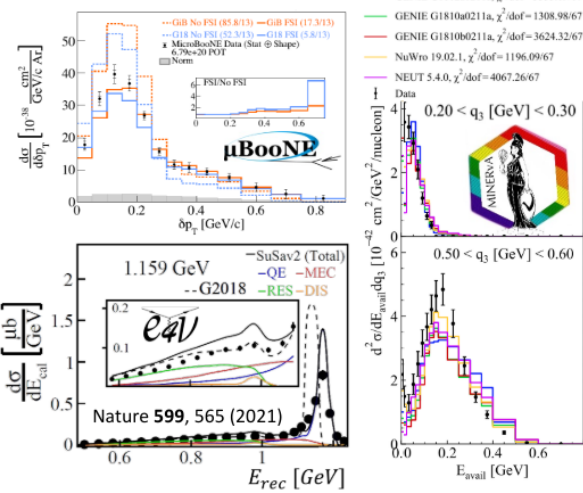
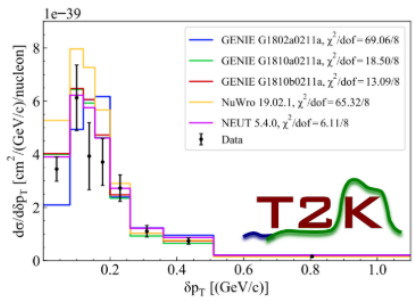


Lepton-hadron correlations

- No generator can come close to describing global data measuring lepton-hadron correlations
- All models are “wrong”, but they are each wrong in different ways

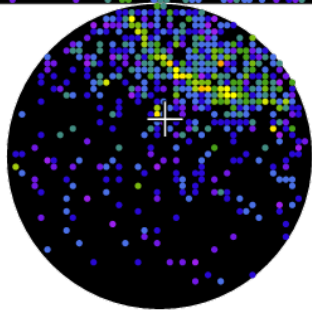
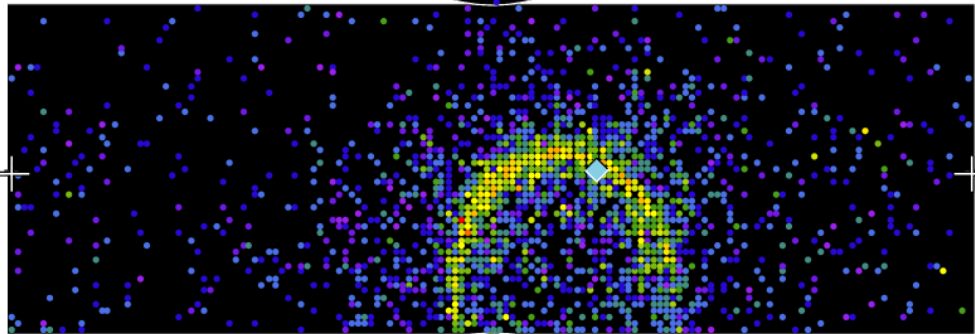
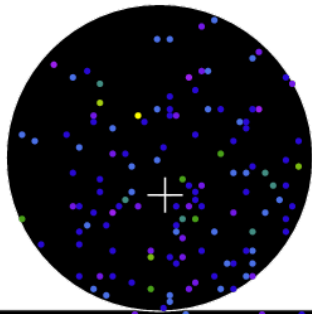
See many more informative generator comparisons in the TENSIONS 2019 report (arXiv:2112.09194)

$$\chi^2_{allModels} \gg N_{bins}$$



Kinematic energy reconstruction

- Energy reconstruction method is function of **selection** and **detector technology**
- T2K and HK are dominated by **CC0 π final state**, and Cherenkov threshold for proton is >1 GeV in H₂O

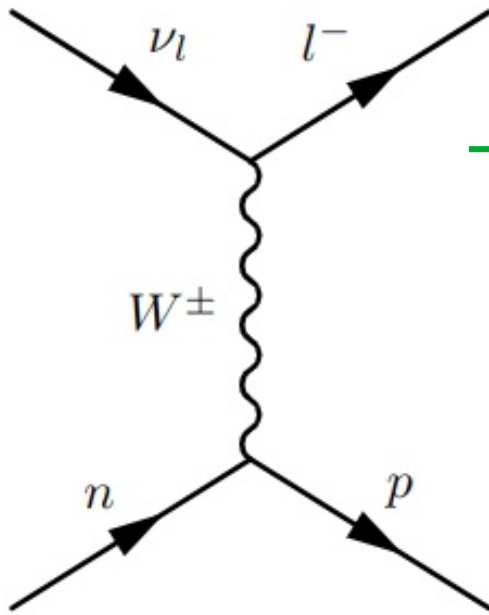


- **Single-track** events
- Kinematic reconstruction using **only lepton** information
- Assumes **4 legged CCQE** interaction, and **initial state nucleon at rest**

$$E_{\nu}^{\text{CCQE}} = \frac{2m_N E_l - m_l^2 + m_{N'}^2 - m_N^2}{2(m_N - E_l + p_l \cos \theta_{\nu,l})}$$

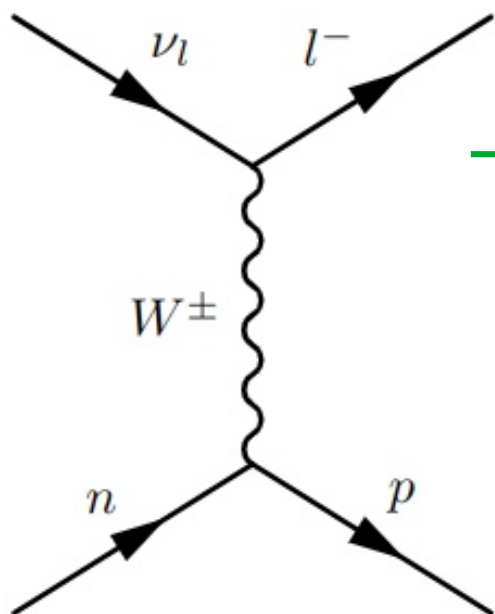
Kinematic energy reconstruction

Any interest in deriving this?



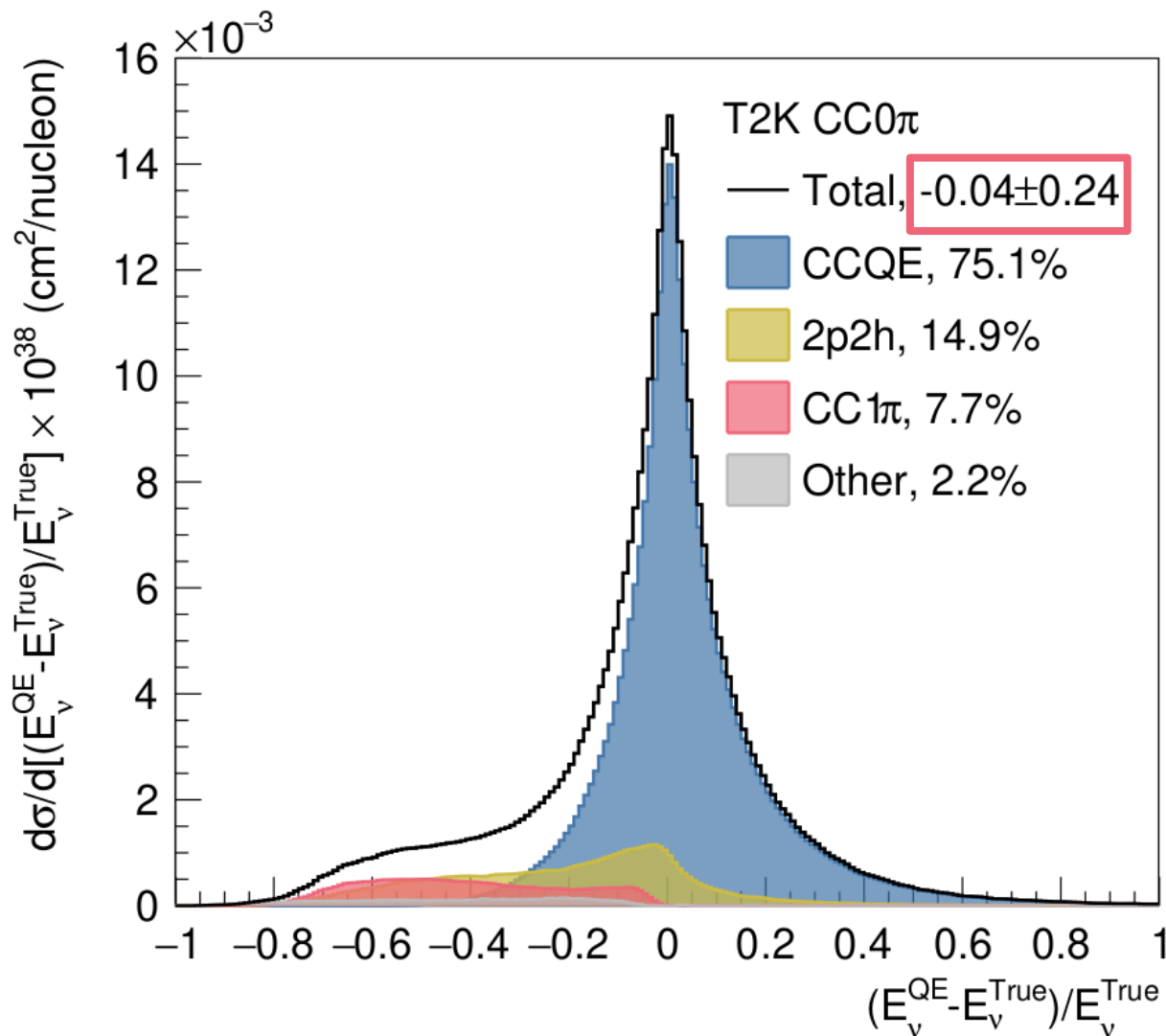
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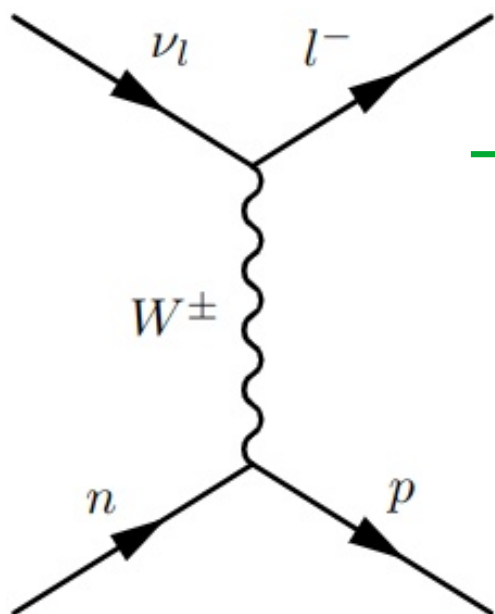


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- **CCQE** contribution largely unbiased
- **20-25% RMS**
- **CC1 π +FSI and 2p2h contribution less than 25% of total signal**

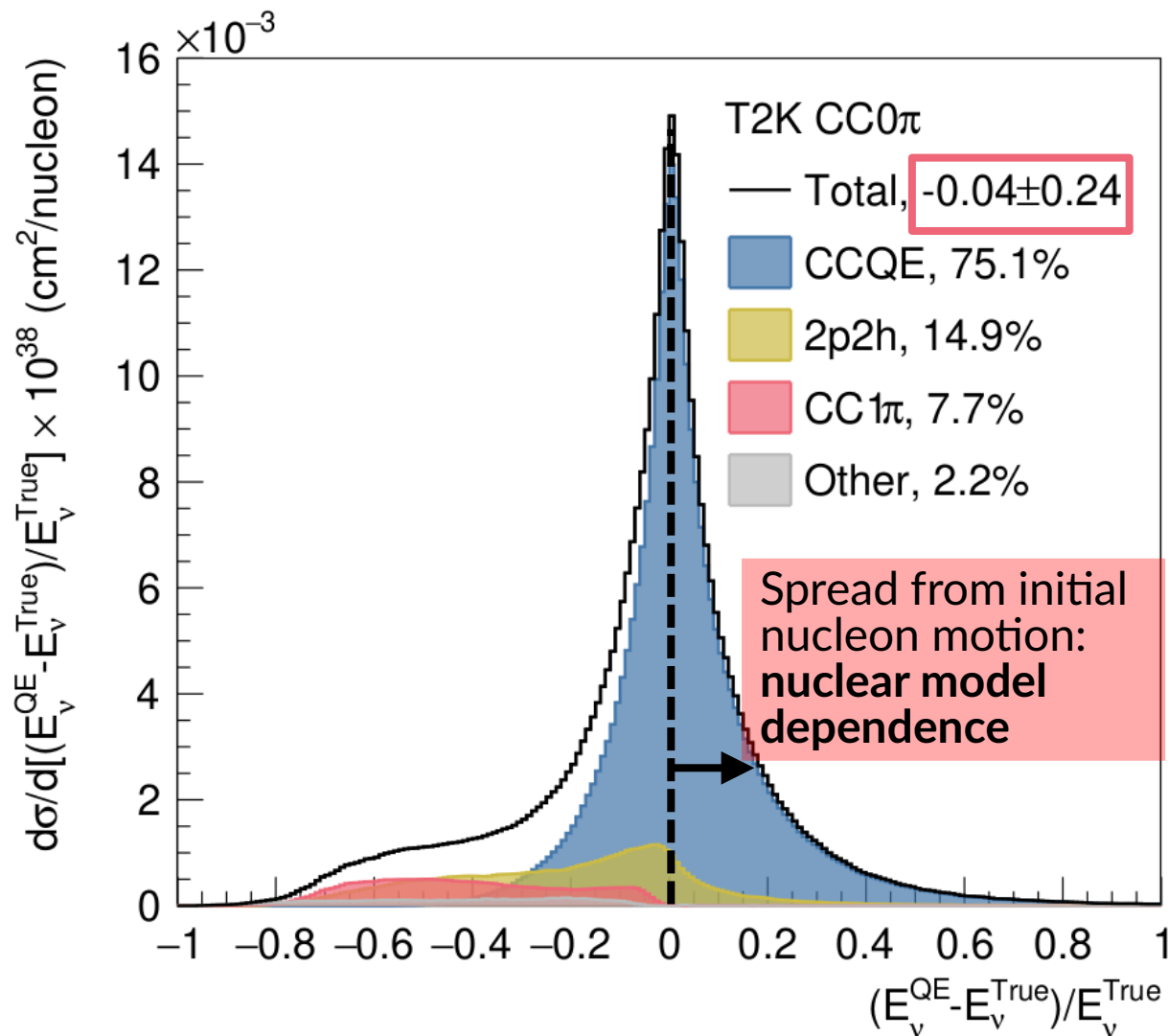


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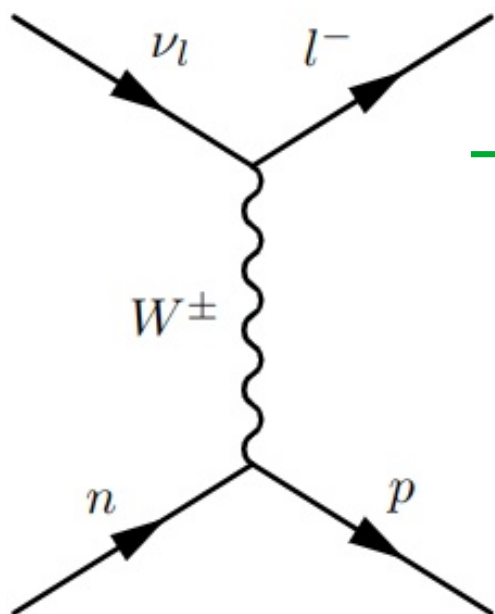


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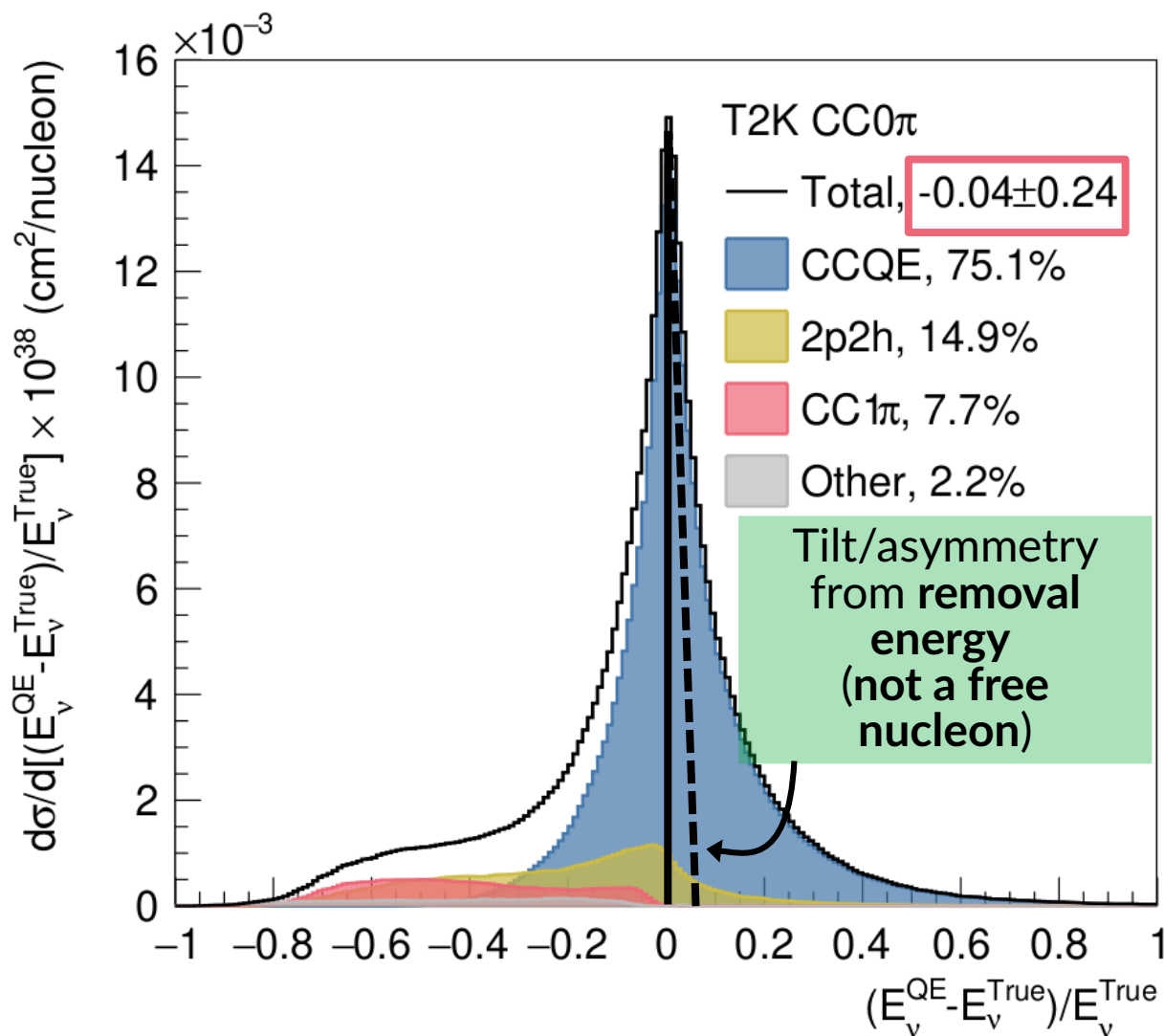


Kinematic energy reconstruction

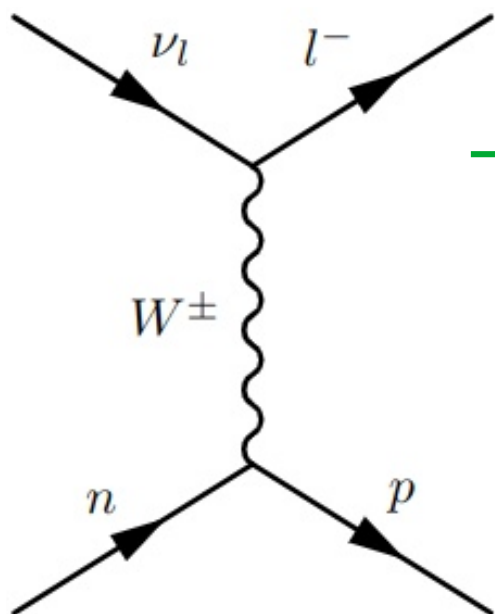


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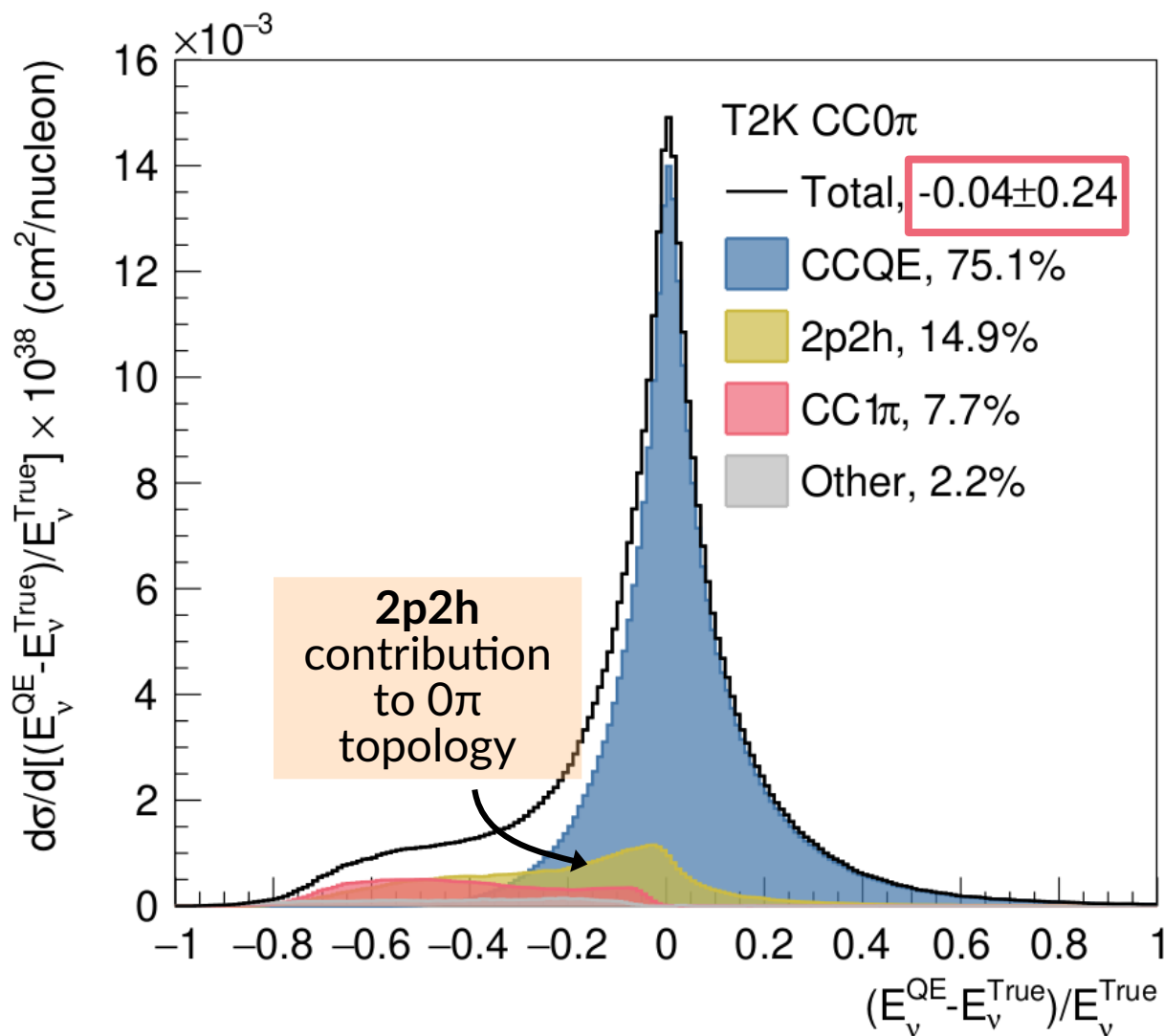


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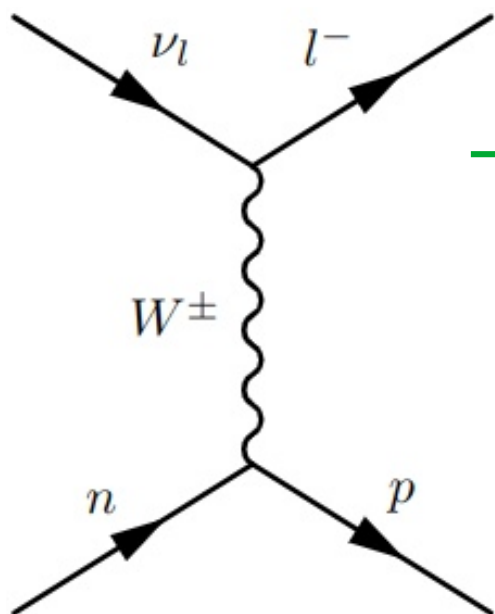


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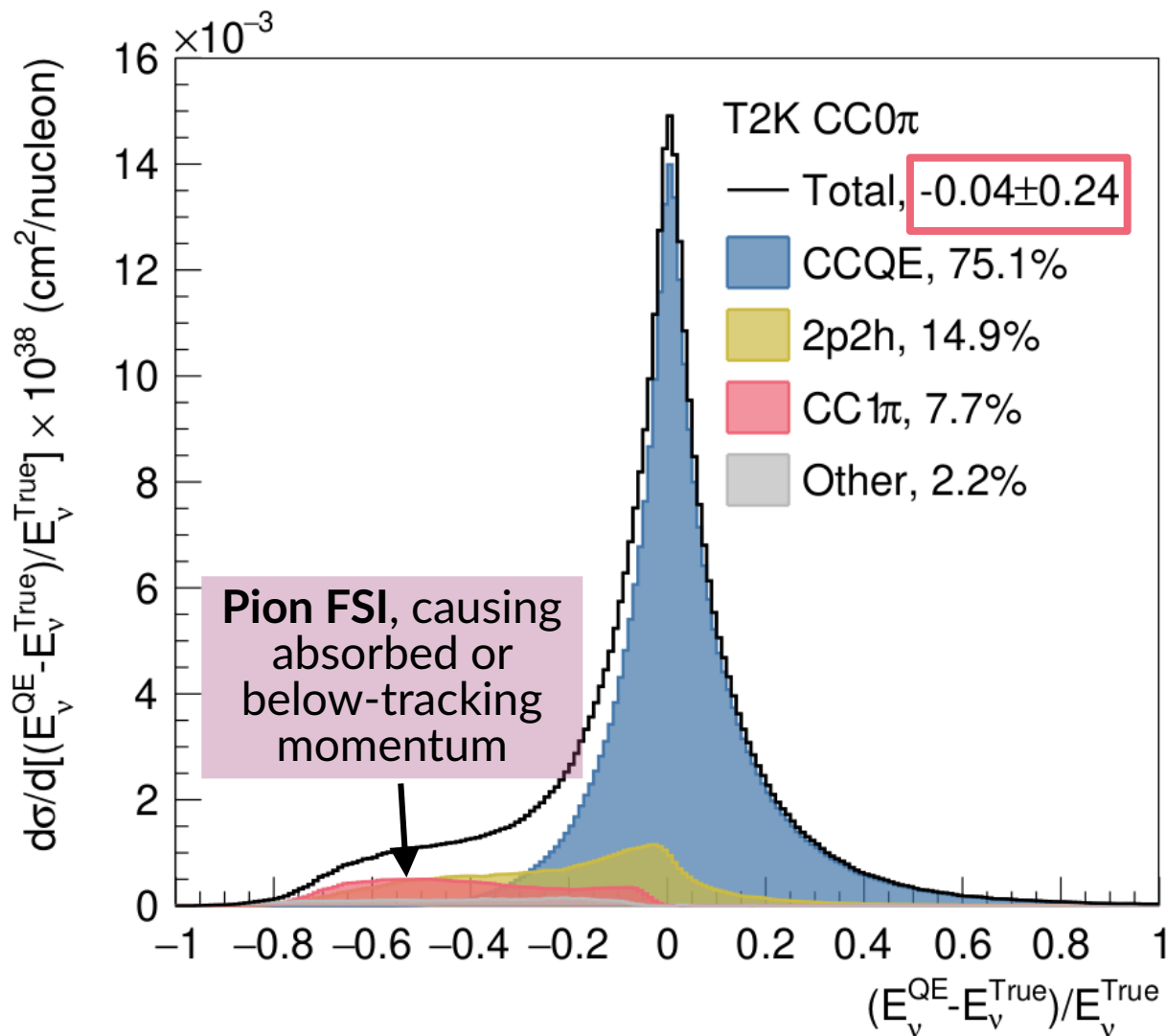


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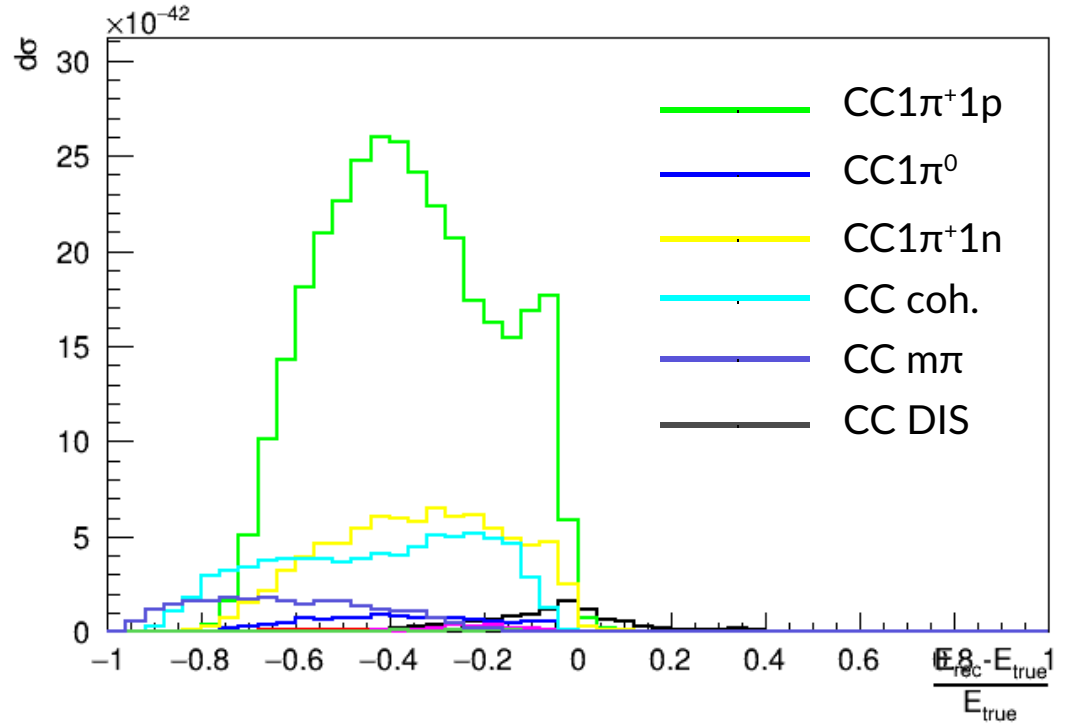
$$E_{\nu}^{\text{CCQE}} = \frac{2m_N E_l - m_l^2 + m_{N'}^2 - m_N^2}{2(m_N - E_l + p_l \cos \theta_{\nu,l})}$$

- **CCQE** contribution largely unbiased
- **20-25% RMS**
- **CC1π+FSI** and **2p2h** contribution less than **25%** of total signal



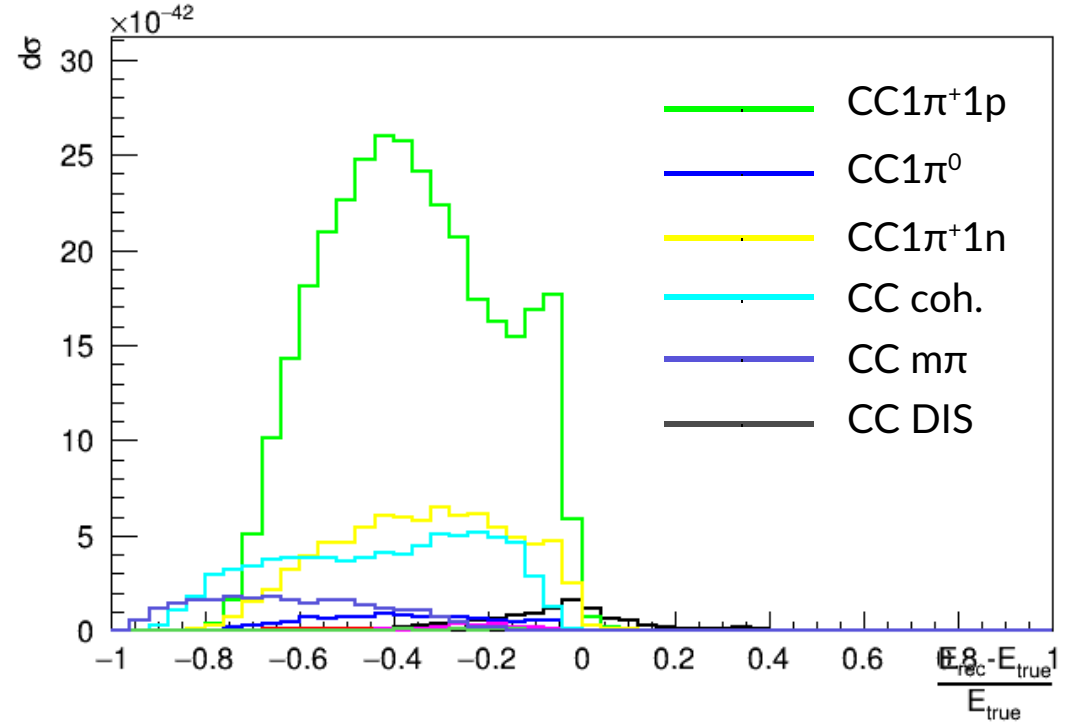
Kinematic energy reconstruction

- When applied to T2K's CC1 π sample, we get a large bias
 - This is for CC1 π events with a pion below 200 MeV/c momentum
- How can we improve?

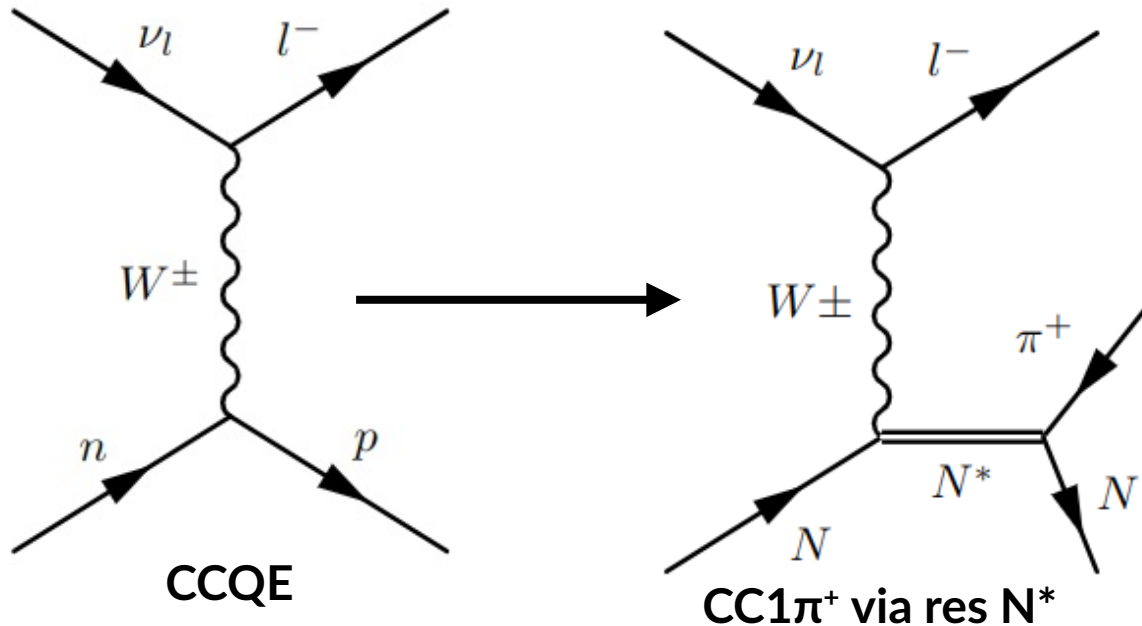


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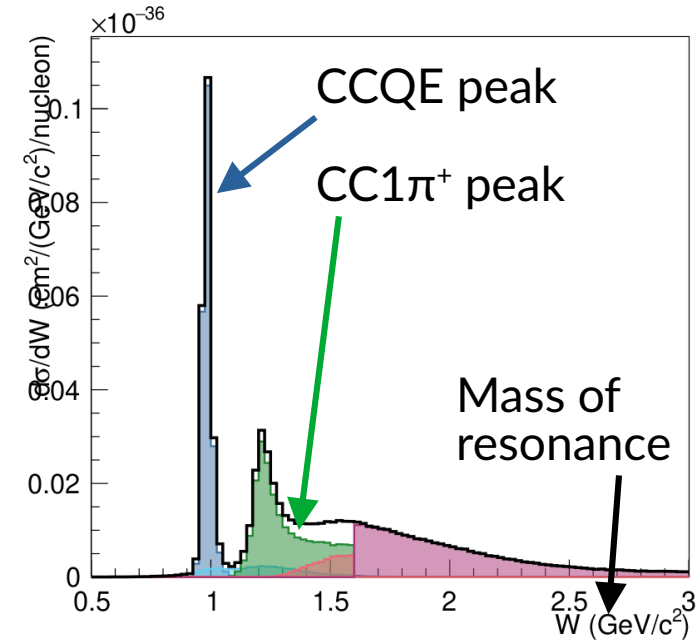
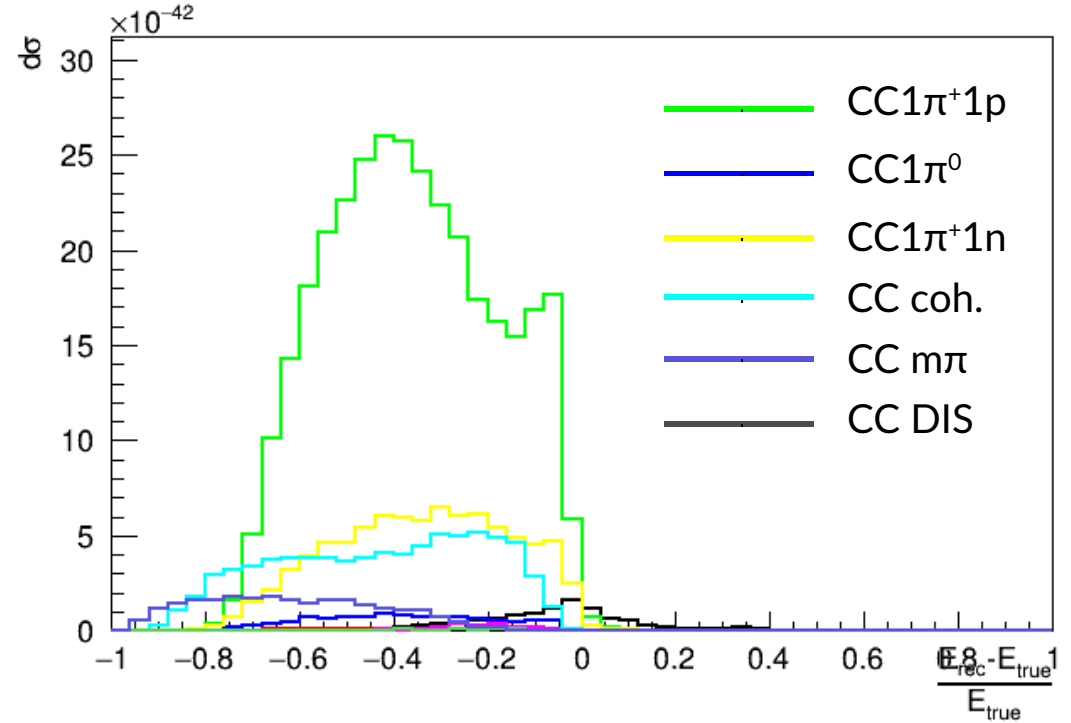
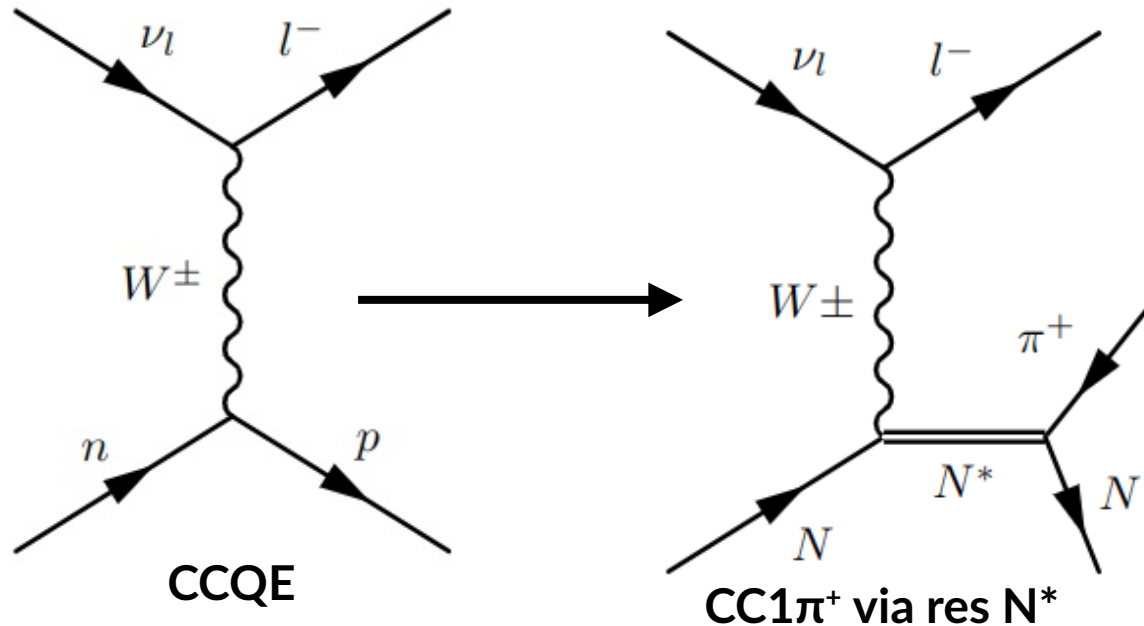
Clue:



Kinematic energy reconstruction

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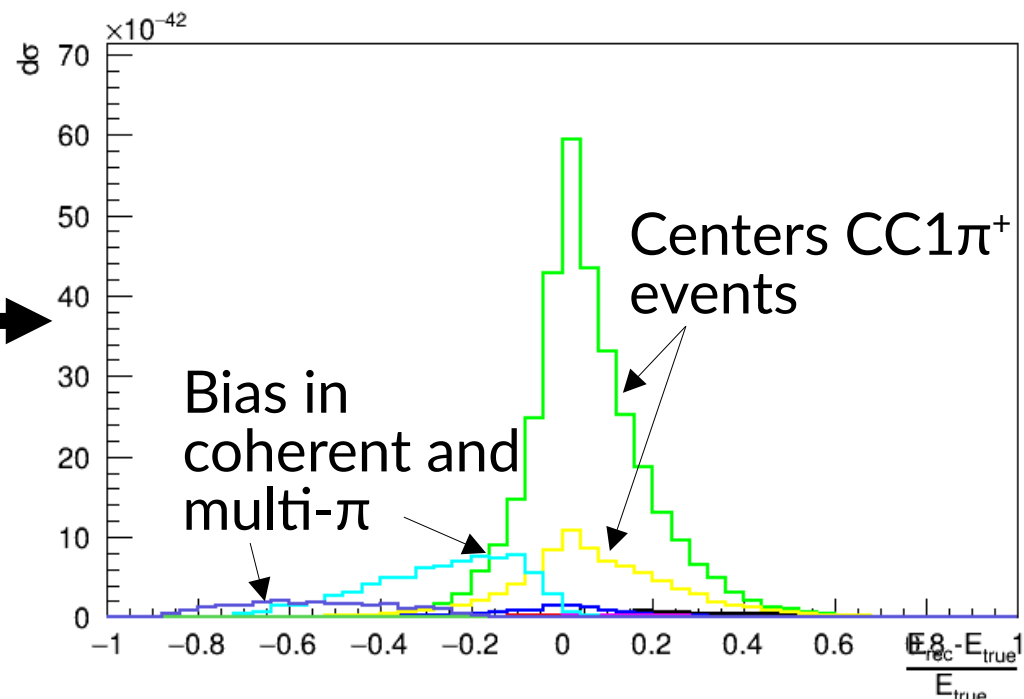
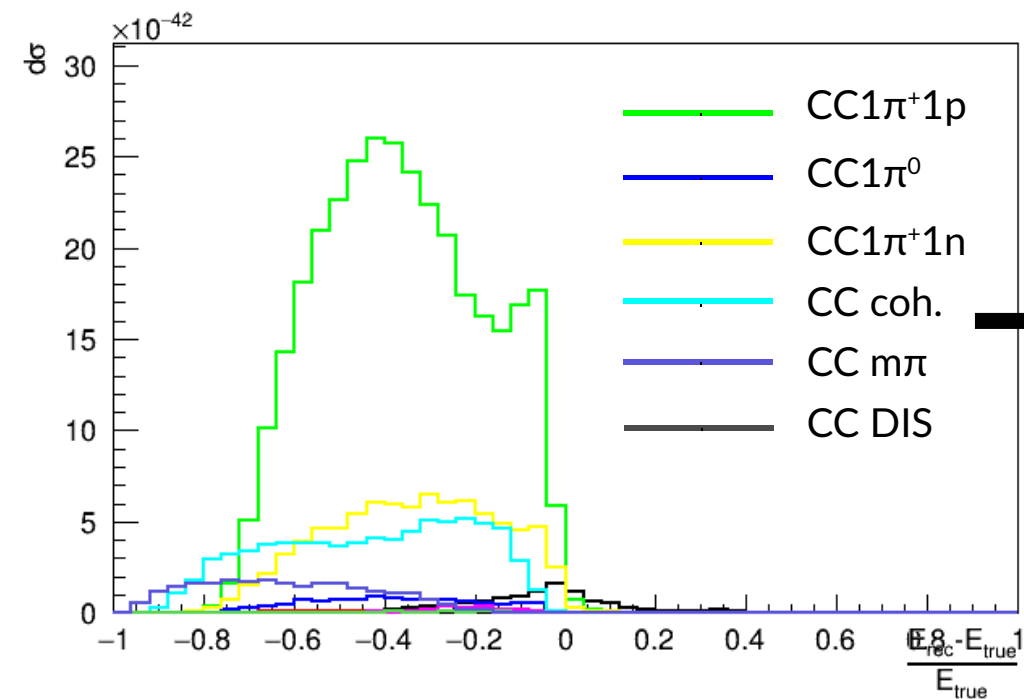
Clue:



Kinematic energy reconstruction

Replace $m_{N'}$ ($\sim 0.938 \text{ GeV}/c^2$)
by m_{Δ} ($\sim 1.232 \text{ GeV}/c^2$)

$$E_{\nu}^{\text{CCQE}} = \frac{2m_N E_l - m_l^2 + m_{N'}^2 - m_N^2}{2(m_N - E_l + p_l \cos \theta_{\nu,l})}$$

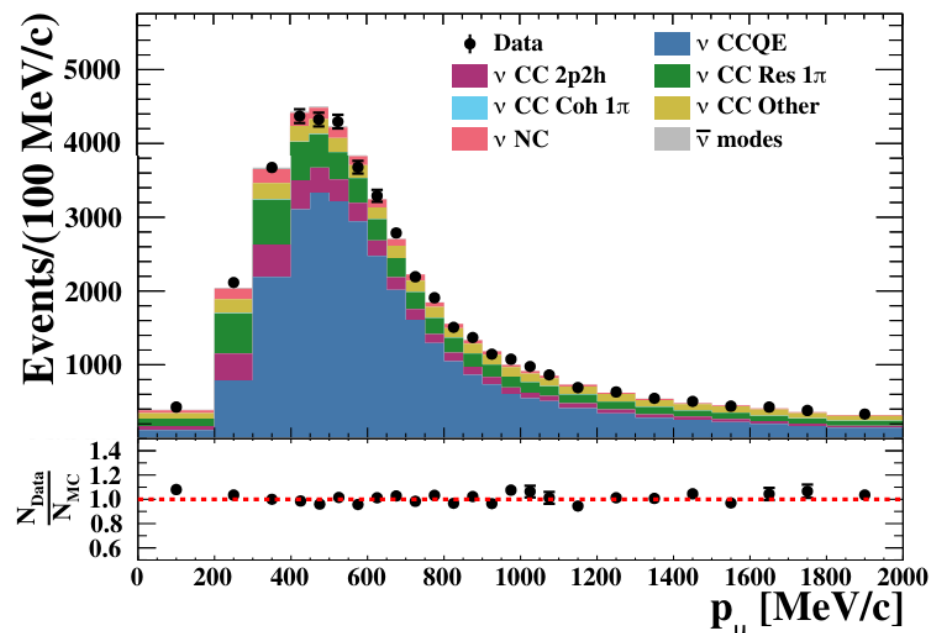
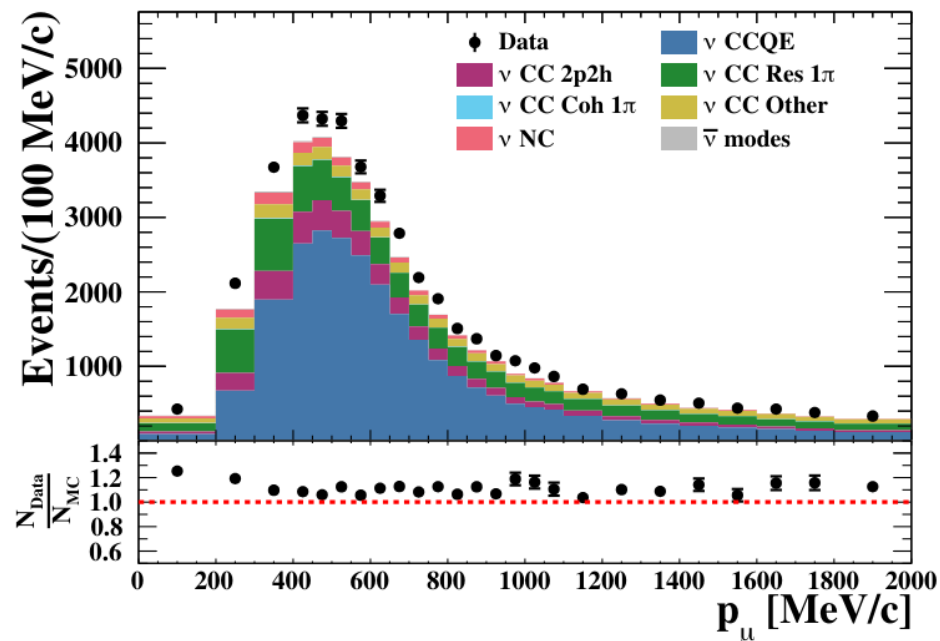


Kinematic energy reconstruction

- Important to get the **CCQE, 2p2h and CC1 π contributions correct**
 - They bias the estimator differently: mistaking non-CCQE for CCQE imposes a bias
- **Direct dependence on nuclear initial-state model**
 - Relatively large contribution at $E_\nu=0.6$ GeV
- Only dependent on **FSI in the absorption**
 - Proton may lose energy to nucleus; does not matter in estimator
 - Secondary dependence on FSI through **missing particles**: think it's four-limbed interaction when it was not
- **Small contribution** from higher W resonances, SIS and DIS contributions (if T2K energies!)

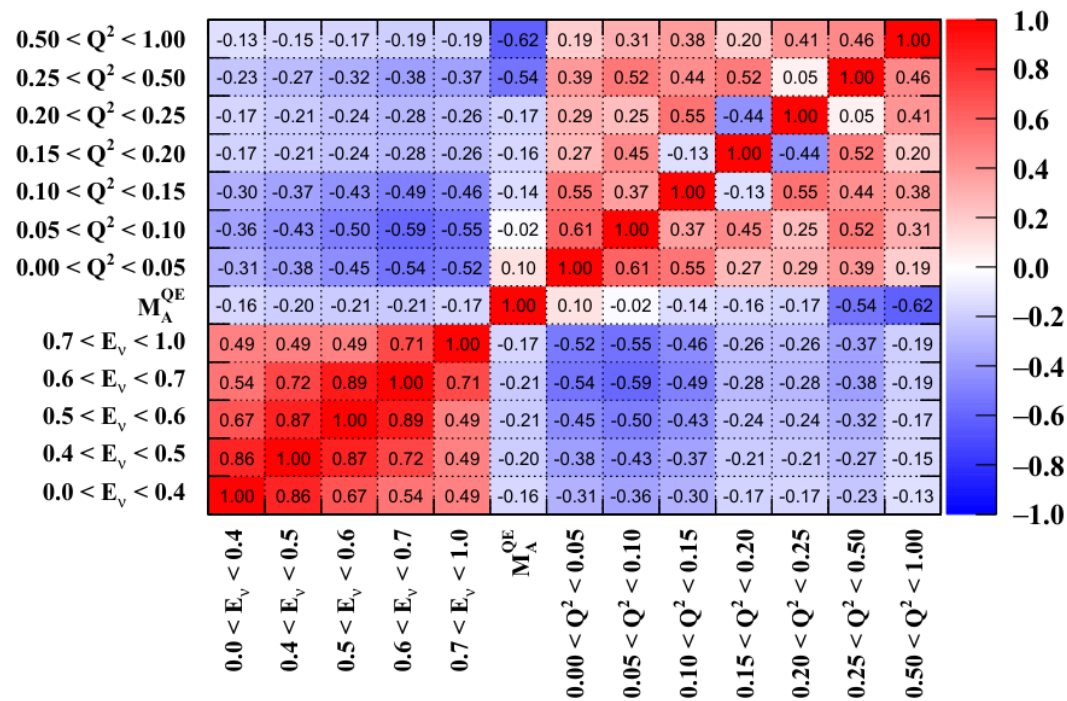
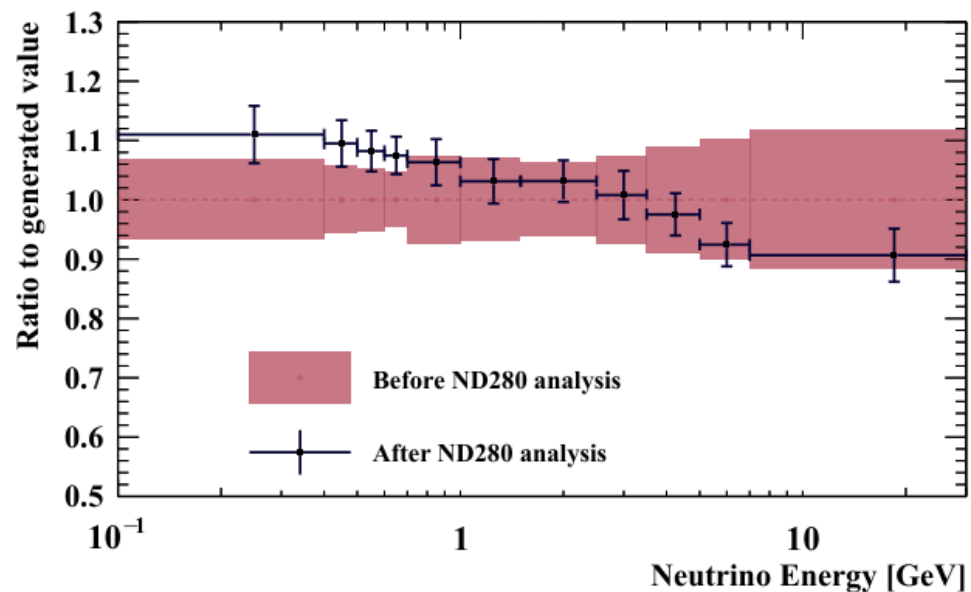
Using the near detector in analysis

- T2K builds prediction for data at the ND using model parameters
 - e.g. Nieves 2p2h normalisations, CCQE mean-field parameters, single pion production, final-state interactions...



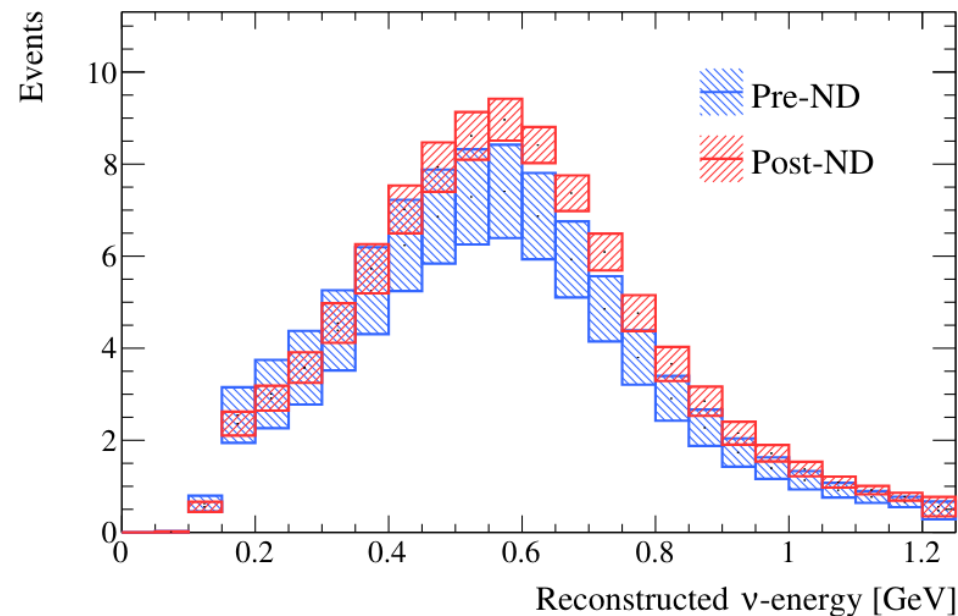
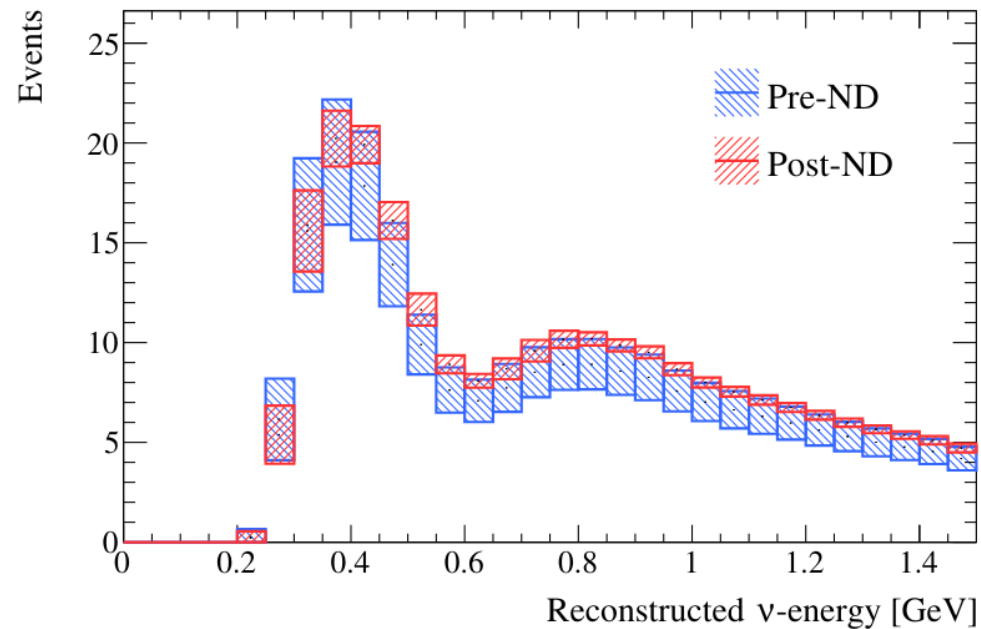
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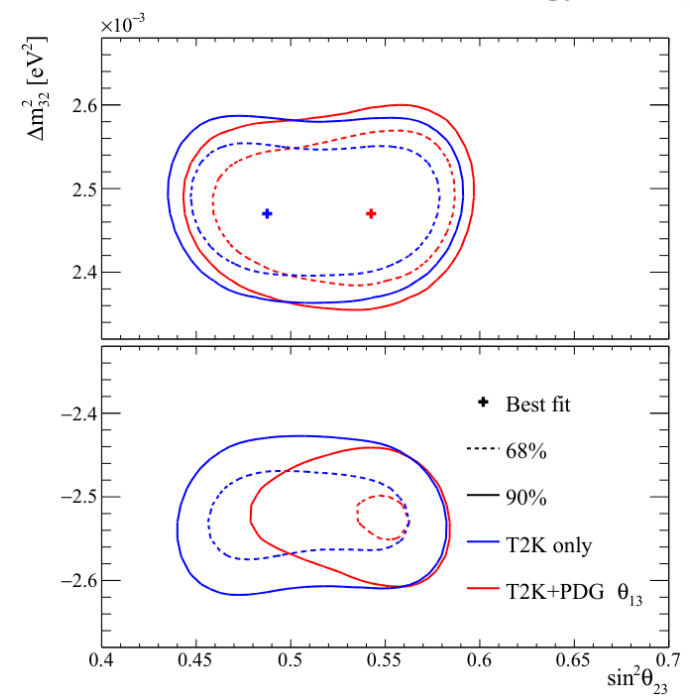
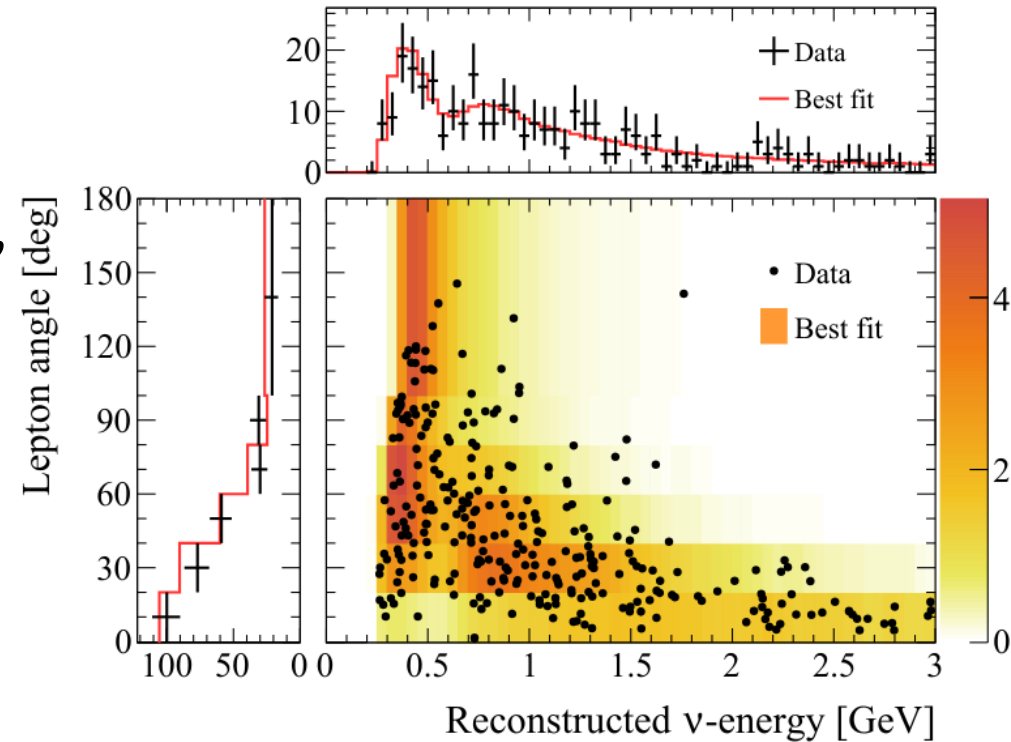
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 - Using the adjusted model
- Build the predictions at the FD against data, after the ND fit to data
 - Using the adjusted model



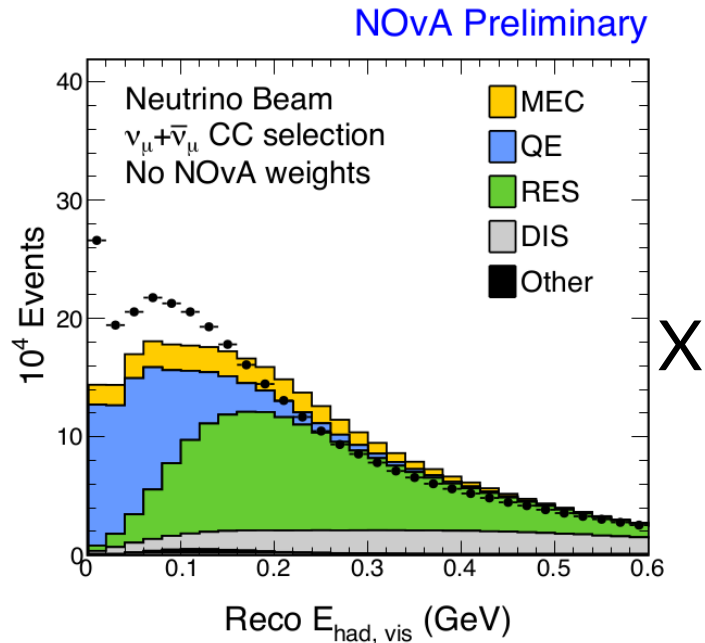
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- Get a set of parameter values fitted to data, and their correlations
 - Using the adjusted model
- Build the predictions at the FD against data, after the ND fit to data
 - Using the adjusted model
- Fit the oscillation parameters!



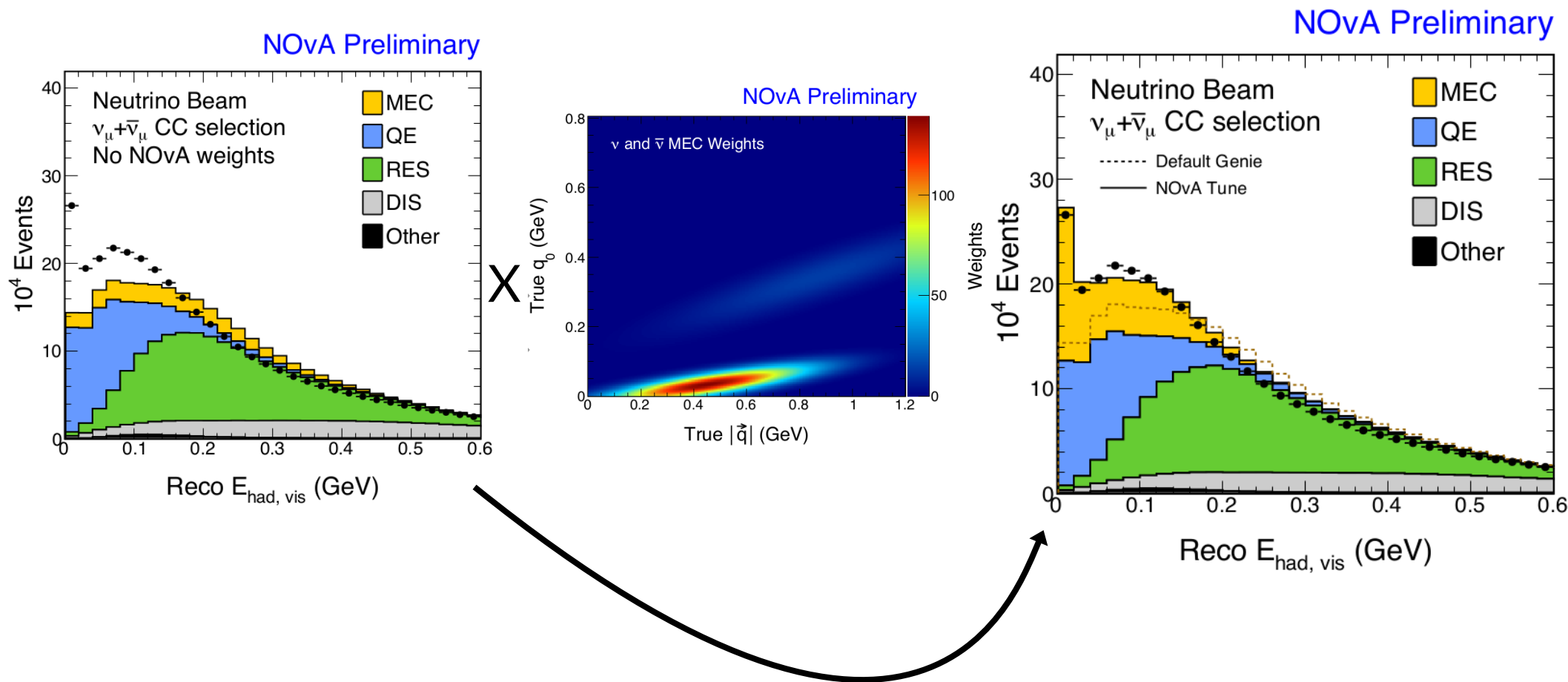
Using the near detector in analysis

- NOvA instead first tune **2p2h** model to **data** in **reconstructed hadronic energy**



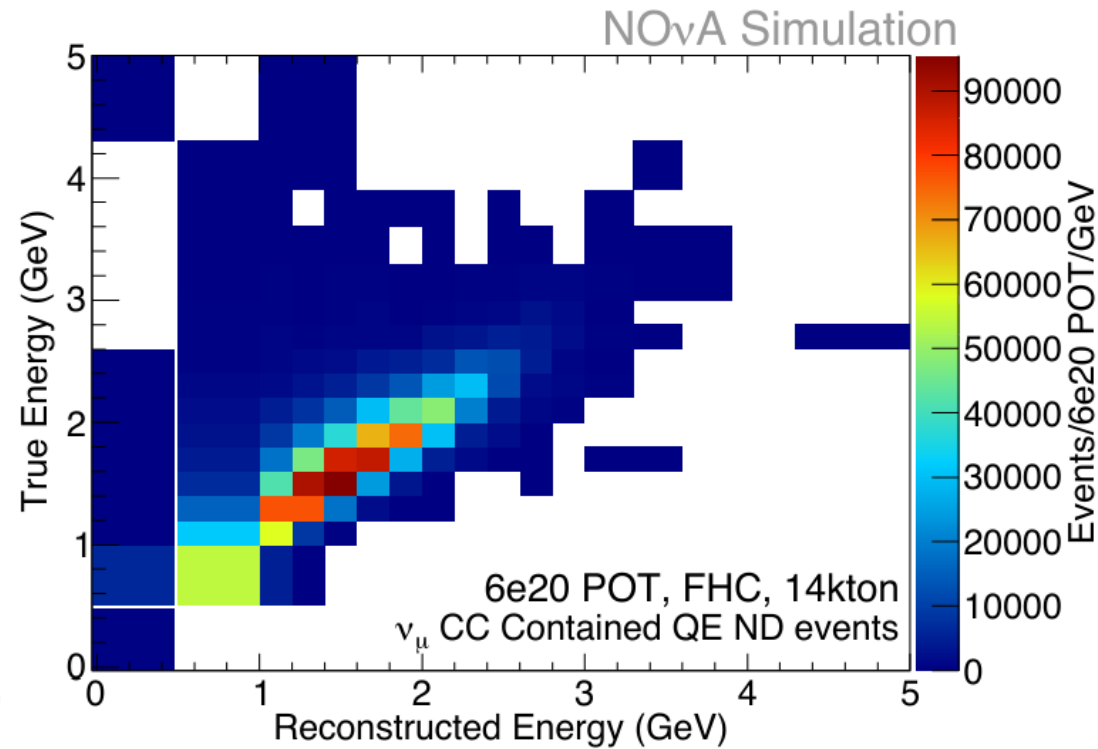
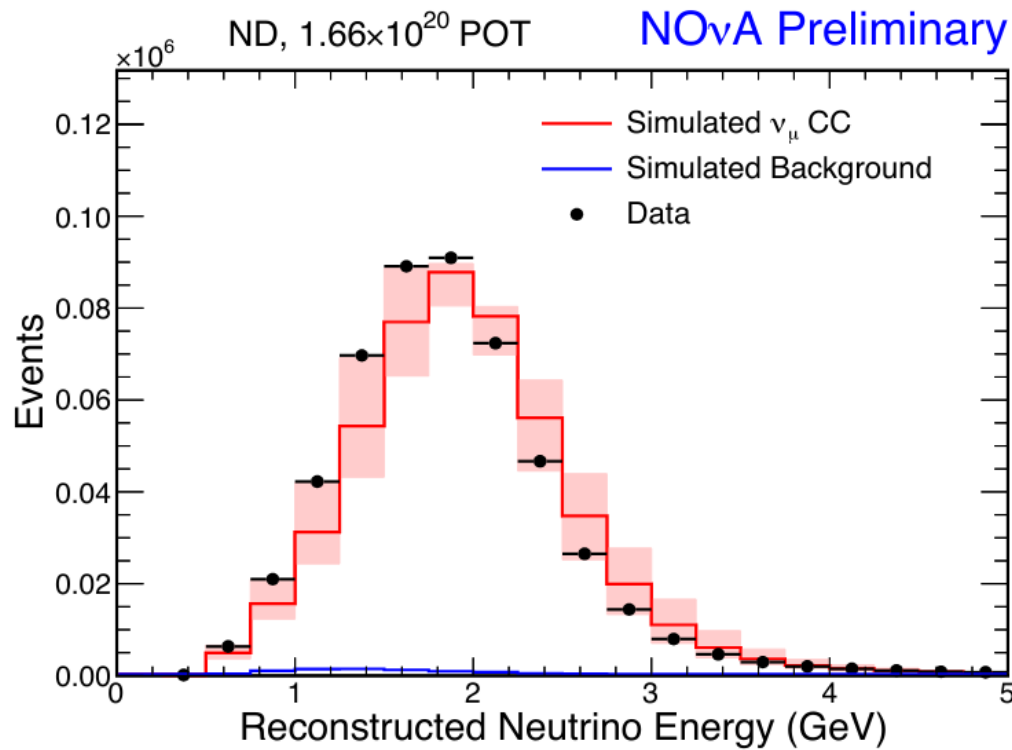
Using the near detector in analysis

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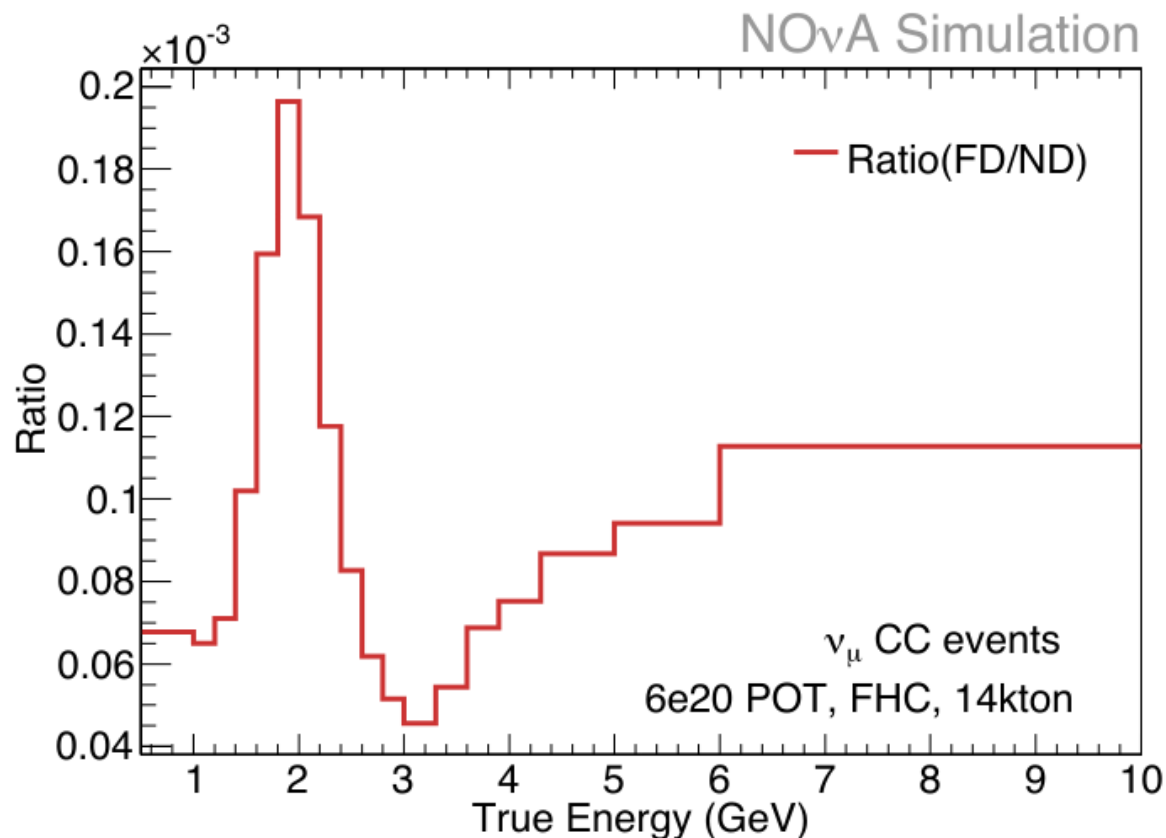
Using the near detector in analysis

- NOvA instead first tune 2p2h model to data in reconstructed hadronic energy
- Unfold reco neutrino energy to true neutrino energy via ND smearing matrix



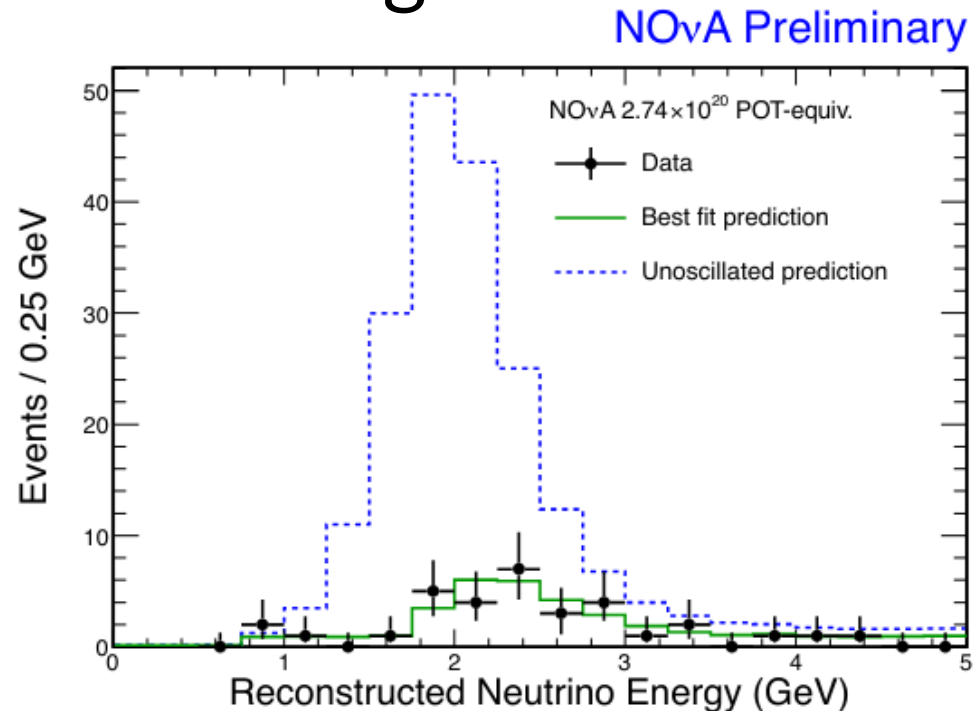
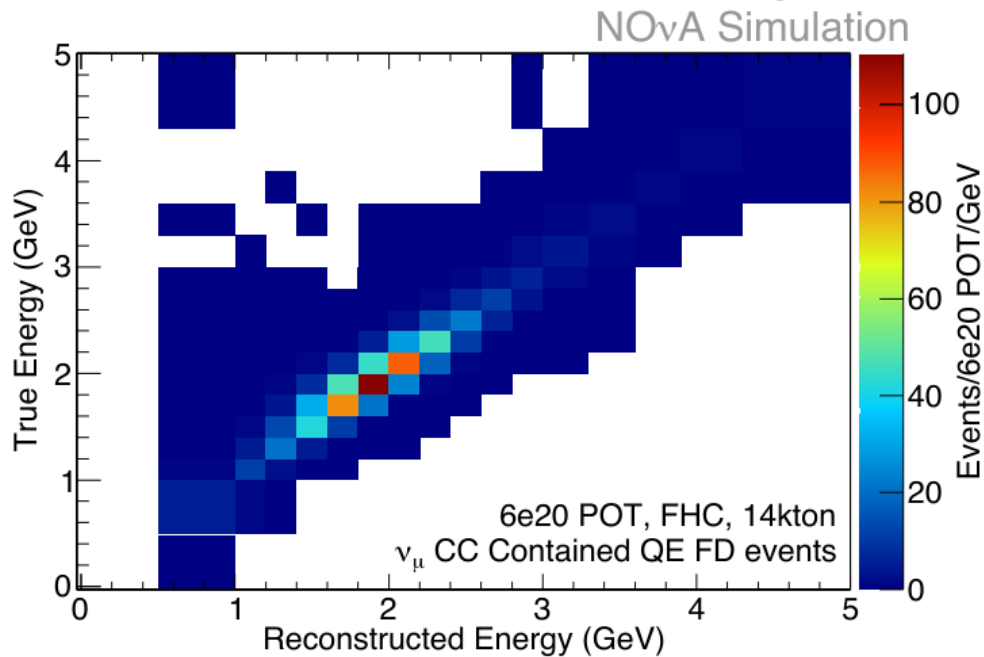
Using the near detector in analysis

- NOvA instead first tune **2p2h** model to **data** in **reconstructed hadronic energy**
- Unfold **reco neutrino energy** to **true neutrino energy** via ND smearing matrix
- Apply “near-to-far” scaling



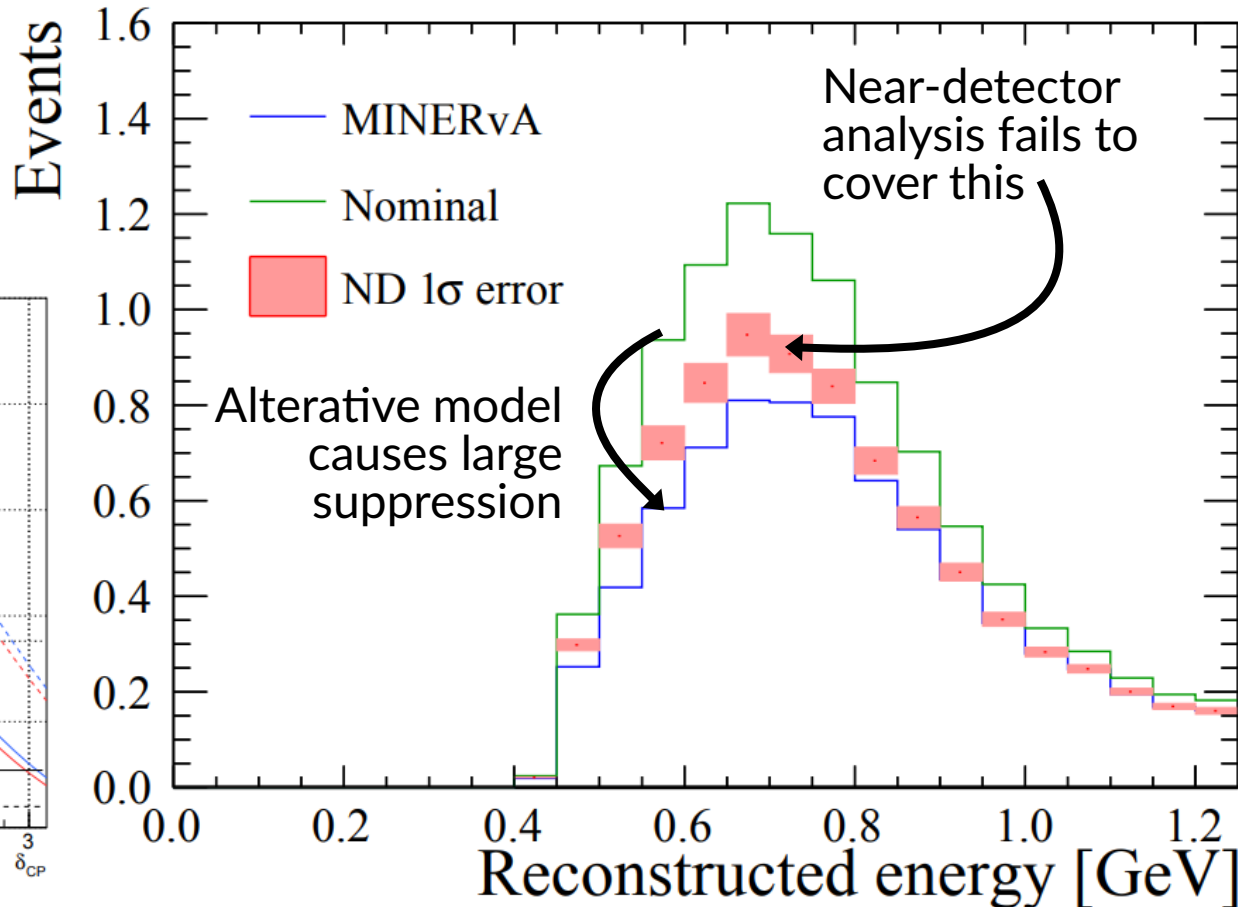
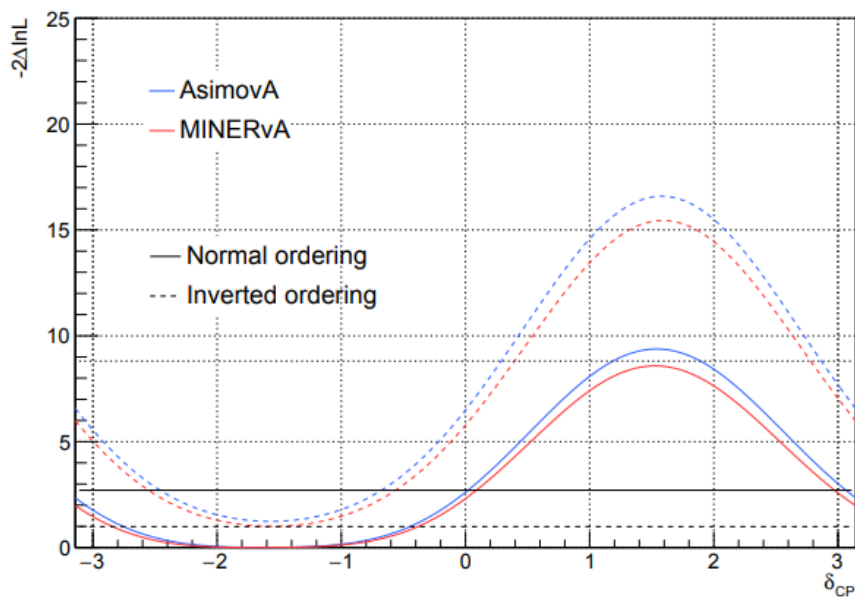
Using the near detector in analysis

- NOvA instead first tune 2p2h model to data in reconstructed hadronic energy
- Unfold reco neutrino energy to true neutrino energy via ND smearing matrix
- Apply “near-to-far” scaling
- Fold back into reconstructed neutrino energy from true neutrino energy, via FD smearing matrix



Fake-data studies

- Use an alternative model to make a prediction for near and far detectors
- Fit to the alternative model at the near detector
 - Set of parameters that best describe the alternative model
- Propagate result to far detector, perform oscillation analysis

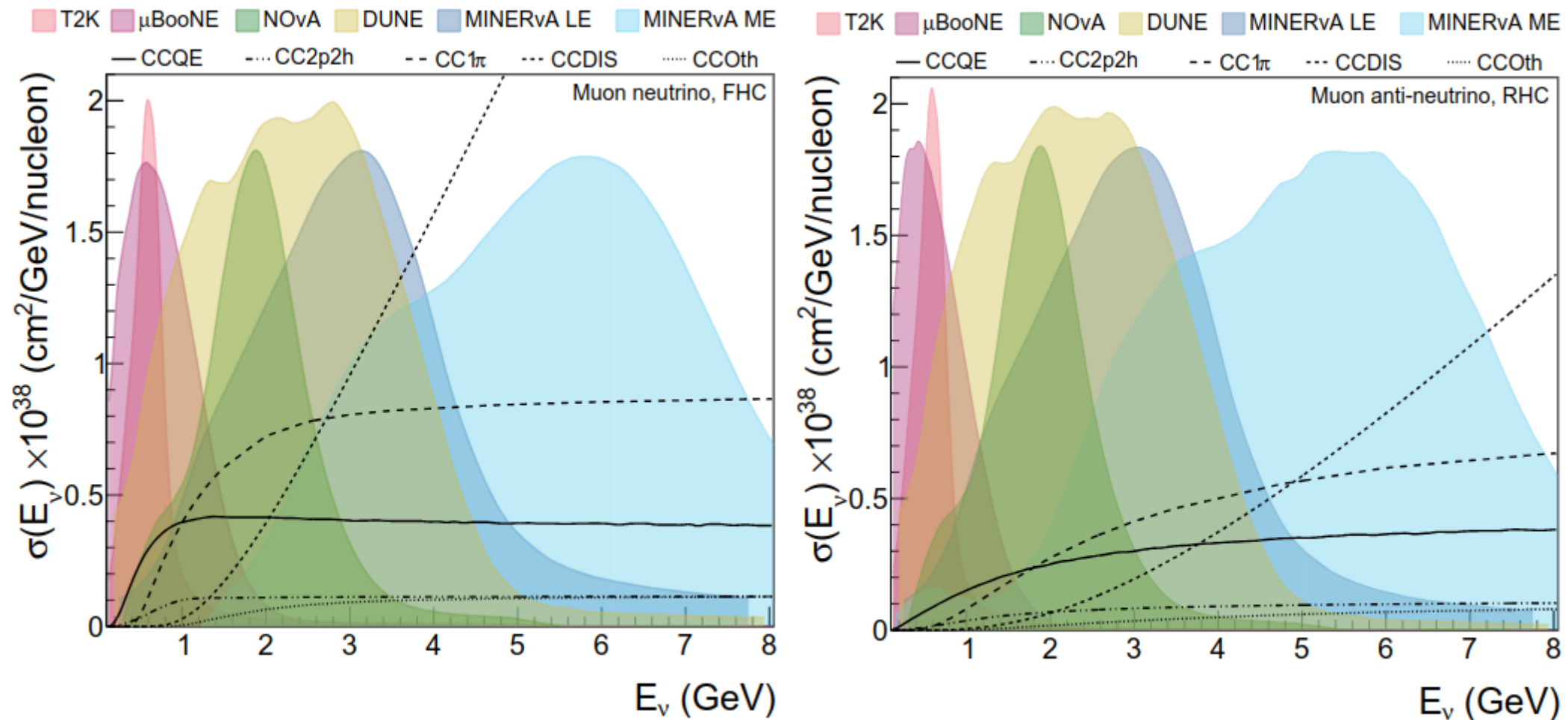


Concluding points

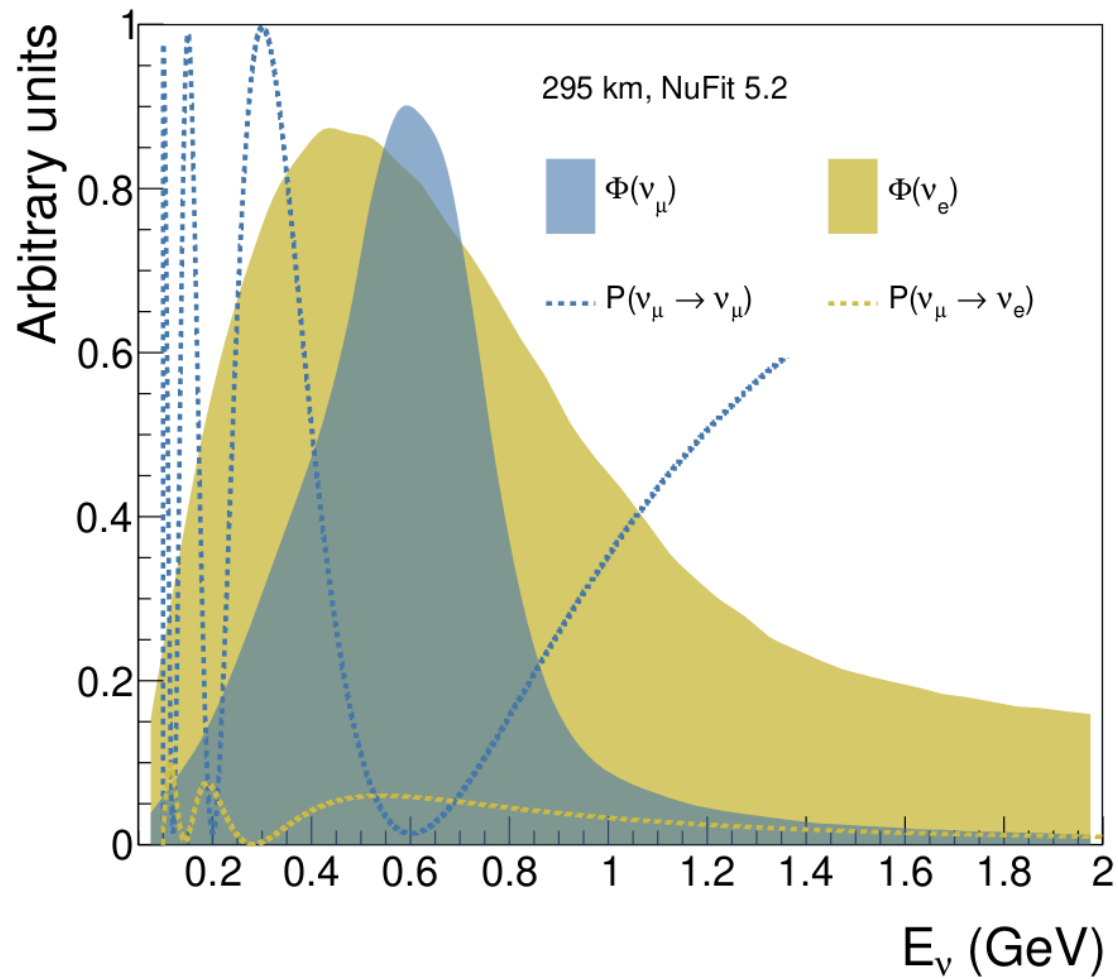
- We're in for a **statistical treat** with HK and DUNE, and the final results from NOvA and T2K!
- In the next 10 years, model uncertainties on neutrino interaction cross sections need to be reduced from **3-5%**, to **1-2%** level
 - Electron neutrino interaction cross sections
 - Carbon to Oxygen scaling (T2K)
 - Neutrino energy reconstruction, **hadrons**
 - Cross section evolution in neutrino energy
- Rich interaction physics contribute across the range
 - Initial state model, nuclear processes (e.g. FSI, in-medium corrections), transition between resonances and DIS...
 - What needs to be prioritised varies greatly between experiments
- Ongoing and upcoming experimental and theoretical programmes aim at addressing these
 - So speak to your new colleagues, and join in on the fun!

Backups

Neutrino fluxes

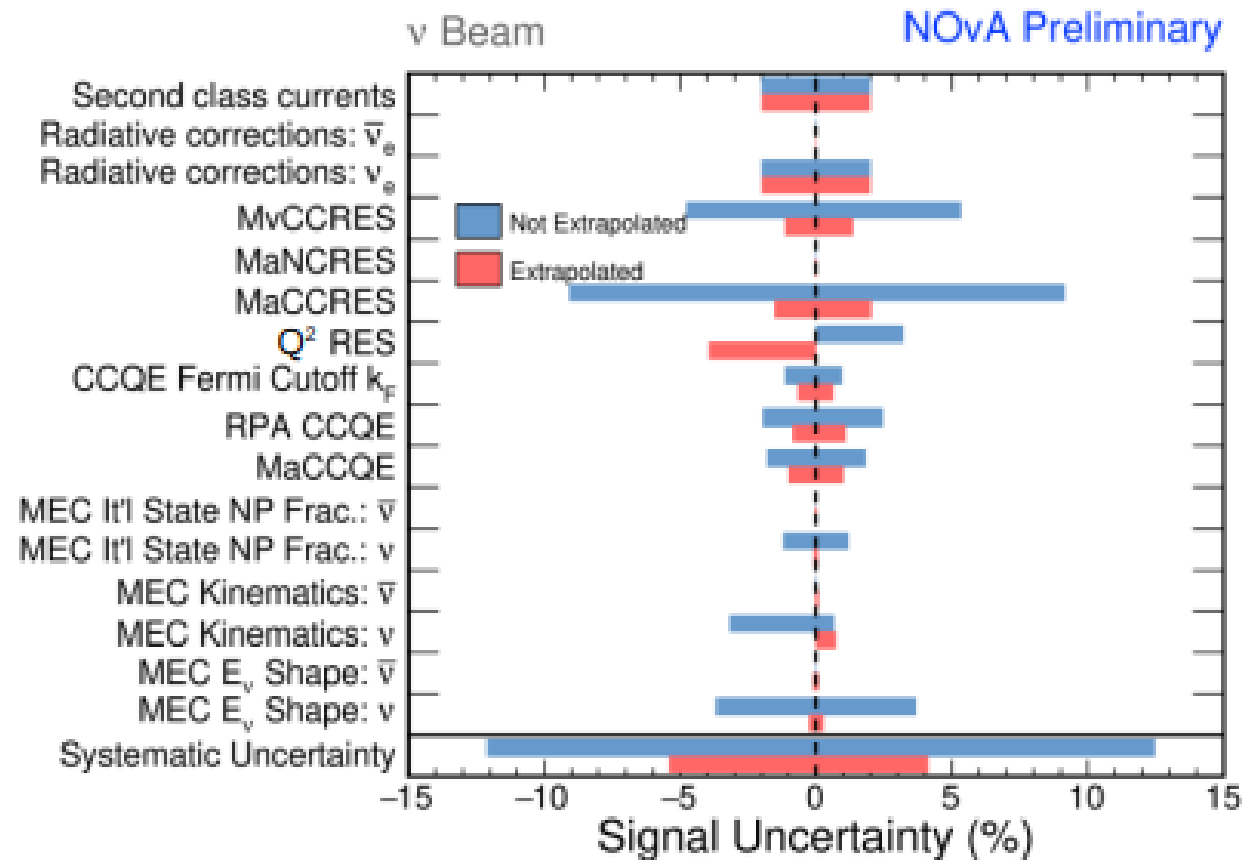


Neutrino fluxes



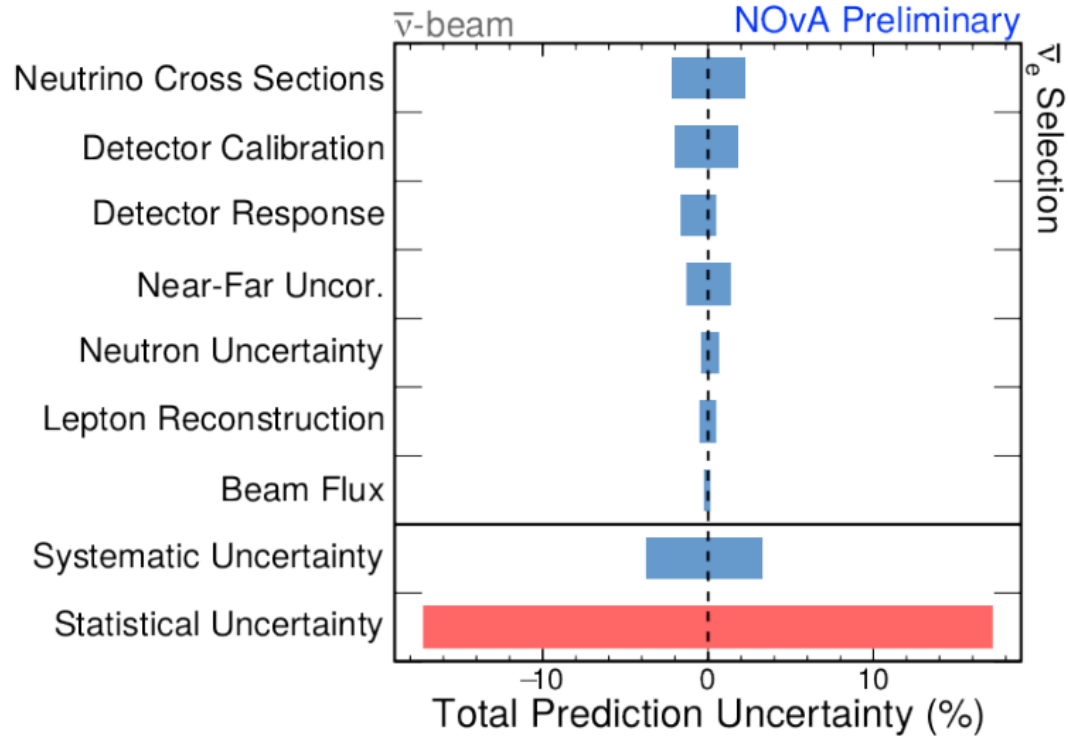
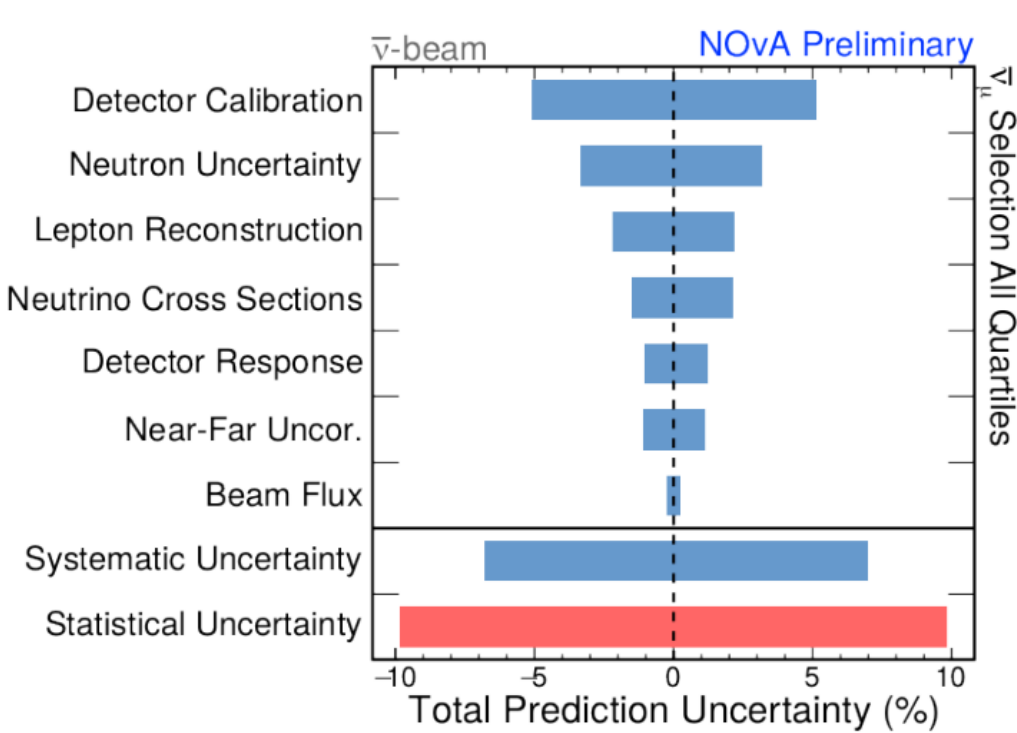
NOvA

- Jeremy Wolcott, NuInt17



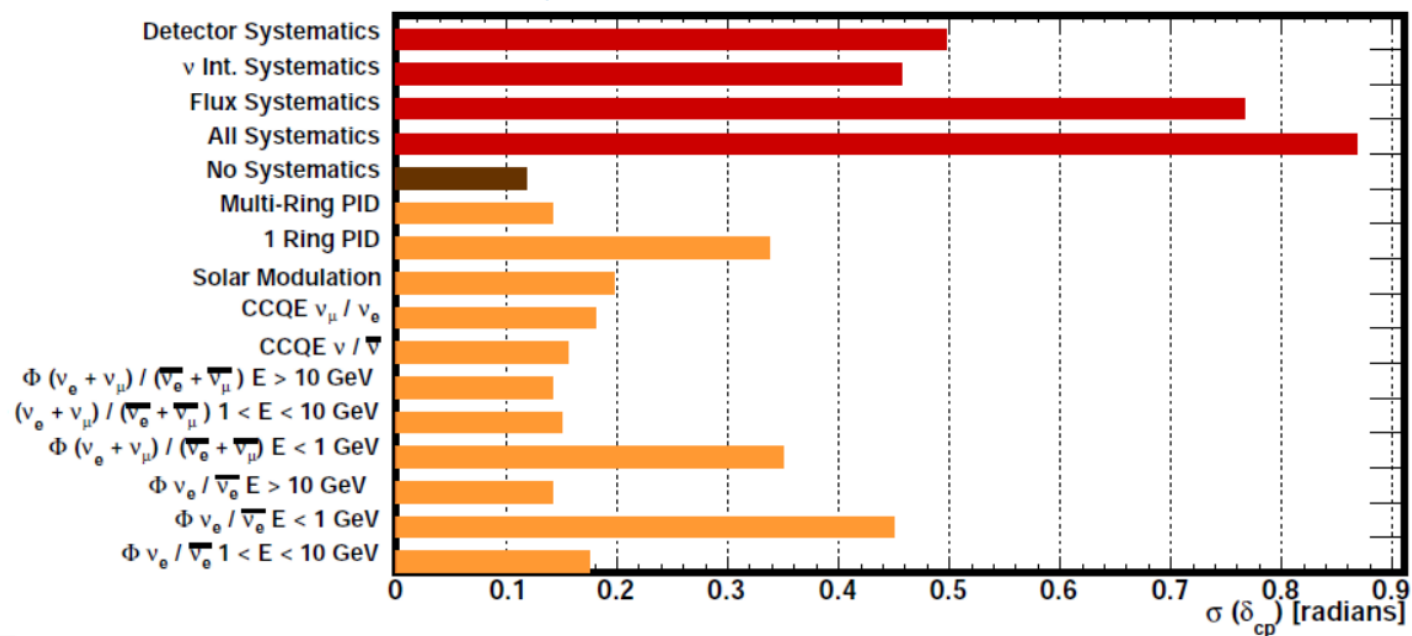
NOvA

M. Elkins, T. Nosek, Neutrino 2020 poster



Atmospheric

Hyper-K's Sensitivity to δ_{cp} with Atmospheric neutrinos



Systematic Effect on Hierarchy Sensitivity at Super-K

