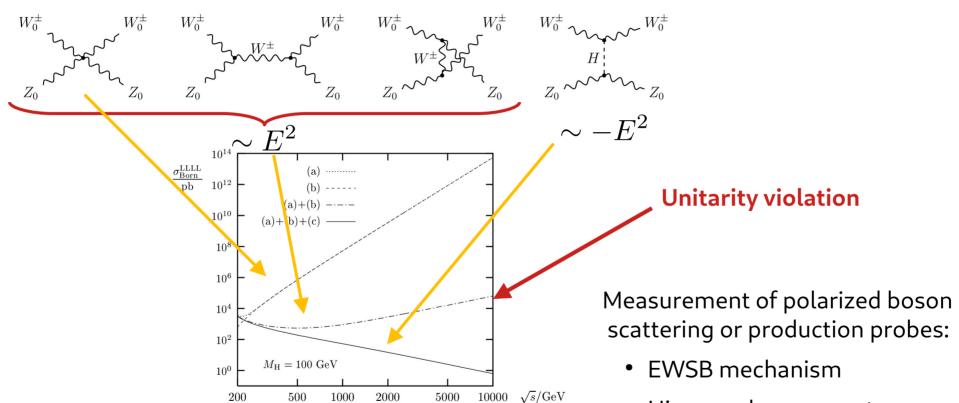
Precise polarisation predictions

Rene Poncelet



Longitudinal Vector-Boson-Scattering (VBS)

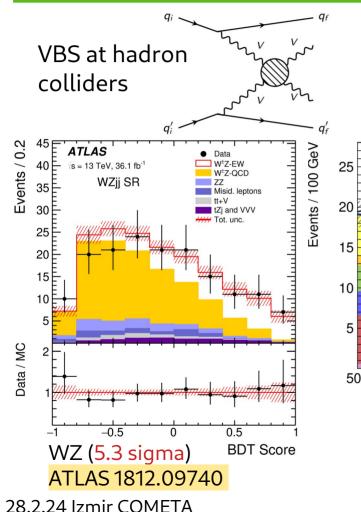


Radiative corrections to W+ W- → W+ W- in the electroweak standard model

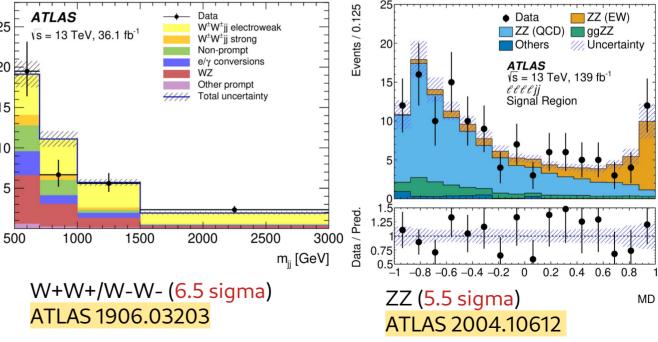
A. Denner, T. Hahn hep-ph/9711302

- Higgs and gauge sector
- New physics models

VBS at hadron colliders



Separate from background processes through VBS topology → a rare process, but observed.



Rene Poncelet – IFJ PAN

Polarised VBS at HL-LHC

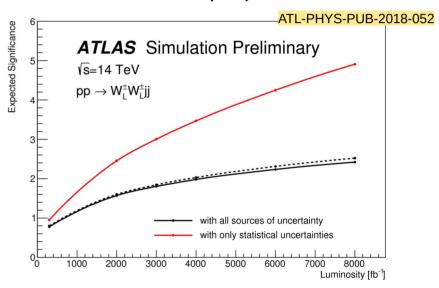
If we want to study unitarisation/EWSB we need to extract the longitudinal component

- only 5-10 % of the total rate
 - → very challenging

(remember: 130fb⁻¹ → ~5-7 sigma

- → naive improvement by factor 10 necessary for observation)
- Requires CMS/ATLAS combination and/or new techniques at HL-LHC
 - → improvement of systematic uncertainties!

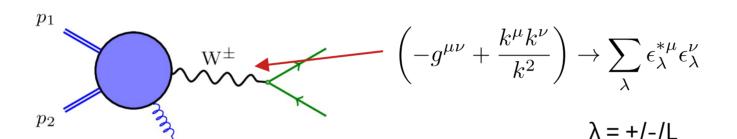
ATLAS HL-LHC projection



How to improve on the (theory) systematics?

- → Improved signal and background modelling
- → Effective separation of boson polarisations

Polarised boson production



Can we extract the longitudinal component?

Measurements of longitudinal polarisation fractions:

Measurement of the Polarization of W Bosons with Large Transverse Momenta in W+Jets Events at the LHC, CMS 1104.3829

Measurement of the polarisation of W bosons produced with large transverse momentum in pp collisions at \sqrt{s}=7 TeV with the ATLAS experiment, ATLAS 1203.2165

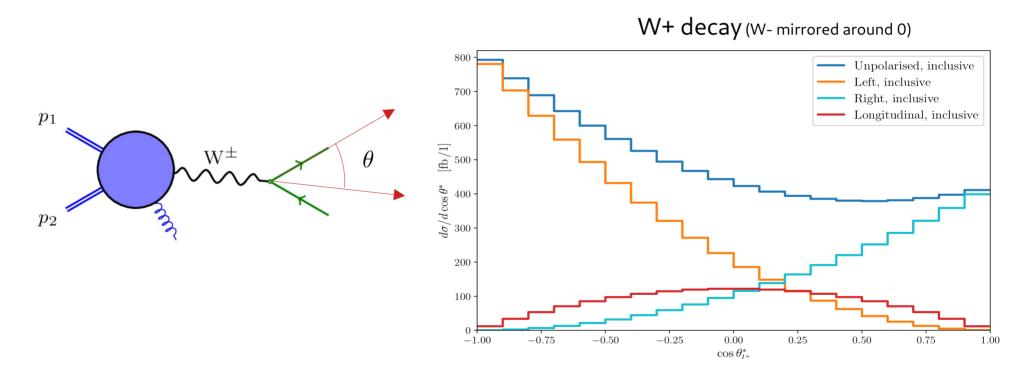
Measurement of WZ production cross sections and gauge boson polarisation in pp collisions at sqrt(s) = 13 TeV with the ATLAS detector, ATLAS 1902.05759

Measurement of the inclusive and differential WZ production cross sections, polarization angles, and triple gauge couplings in pp collisions at sqrt(s) = 13 TeV, CMS 2110.11231

Observation of gauge boson joint-polarisation states in WZ production from pp collisions at sqrt(s) = 13 TeV with the ATLAS detector ATLAS 2211.09435

How to measure polarized bosons?

- We can't measure boson polarization directly.
- Luckily decay products can be used as a "polarimeter":



Polarized cross sections

$$p_1$$
 p_2
 W^{\pm}

$$= \mathbf{P}_{\mu} \cdot \frac{-g_{\mu\nu} + \frac{\kappa \cdot \kappa}{k^2}}{k^2 - M_V^2 + iM_V\Gamma_V} \cdot \mathbf{D}_{\nu}$$

On-shell bosons: $\left(-g^{\mu\nu} + \frac{k^{\mu}k^{\nu}}{k^2}\right) \to \sum_{\lambda} \epsilon_{\lambda}^{*\mu} \epsilon_{\lambda}^{\nu}$ (DPA or NWA)

$$\left(-g^{\mu\nu} + \frac{k^{\mu}k^{\nu}}{k^2}\right) \to \sum_{\lambda} \epsilon_{\lambda}^{*\mu} \epsilon_{\lambda}^{\nu}$$

$$M = \mathbf{P}_{\mu} \cdot \frac{-g_{\mu\nu} + \frac{k^{\mu}k^{\nu}}{k^{2}}}{k^{2} - M_{V}^{2} + iM_{V}\Gamma_{V}} \cdot \mathbf{D}_{\nu} \qquad |M|^{2} = \sum_{\lambda} |M_{\lambda}|^{2} + \sum_{\lambda \neq \lambda'} M_{\lambda}^{*} M_{\lambda'}$$

→ polarised x-sections Interferences

Create samples of fixed polarisation:

$$\frac{\mathrm{d}\sigma}{\mathrm{d}X} = f_L \frac{\mathrm{d}\sigma_L}{\mathrm{d}X} + f_R \frac{\mathrm{d}\sigma_R}{\mathrm{d}X} + f_0 \frac{\mathrm{d}\sigma_0}{\mathrm{d}X} \left(+f_{int.} \frac{\mathrm{d}\sigma_{int.}}{\mathrm{d}X} \right)$$

Template fit f_L, f_R, f_0 to measured $\frac{\mathrm{d}\sigma^{exp.}}{\mathrm{d}X}$

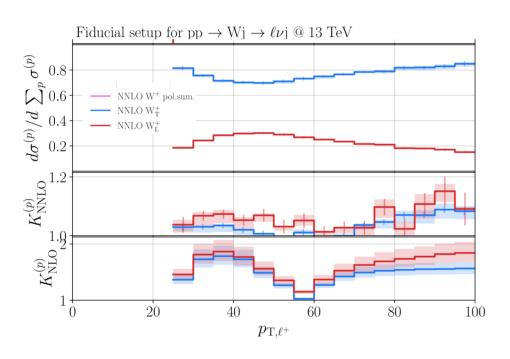
Polarized cross sections

$$\frac{\mathrm{d}\sigma}{\mathrm{d}X} = f_L \frac{\mathrm{d}\sigma_L}{\mathrm{d}X} + f_R \frac{\mathrm{d}\sigma_R}{\mathrm{d}X} + f_0 \frac{\mathrm{d}\sigma_0}{\mathrm{d}X} \left(+f_{int.} \frac{\mathrm{d}\sigma_{int.}}{\mathrm{d}X} \right)$$

- Interferences can be handled
- Does not rely on extrapolations to the full phase space
 X can be any observable → lab frame observables
- $\frac{\mathrm{d}\sigma_i}{\mathrm{d}X}$ can be systematically improved

Higher-order QCD/EW corrections + PS to minimize uncertainties from MHO (scale uncertainties)

Why do we need higher-order corrections?



Important

Just using some K-factors is not enough

- 1) Differential polarization fraction have shapes (not just one number!)
- 2) Higher-order corrections dependent on polarization! Just using polarized K-factor would lead to distortion of spectrum.
- 3)NNLO QCD needed to reach percent-level scale-dependence → MHO

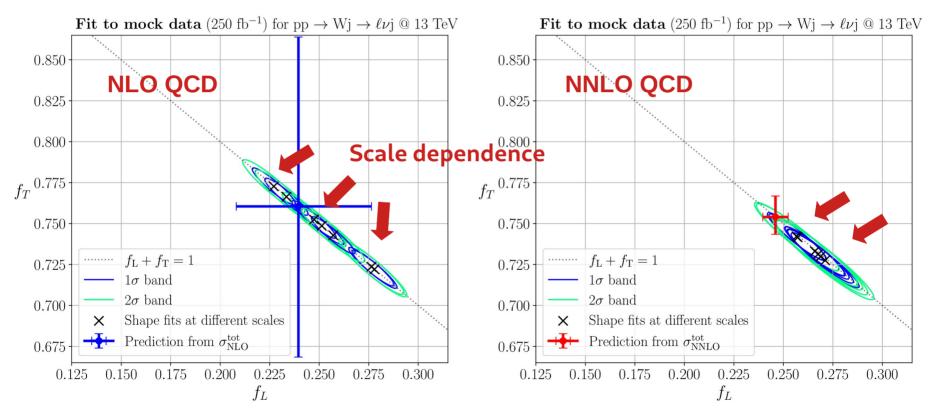
Polarised W+j production at the LHC: a study at NNLO QCD accuracy, Pellen, Poncelet, Popescu 2109.14336

W+jet: mock-data fit

Fit to mock-data (based on NNLO QCD and 250 fb⁻¹ stats):

Observable: $\cos(\ell, j_1)$

→ extreme case to see effect of scale dependence reduction



Status of polarization precision calculations

(Collection of papers in the backup)

					(Cottection
Process	LO	NLO	NLO EW	NNLO	+ PS
pp → WW	X	X	X	X	X
pp → ZZ	Χ	Χ	X		X
pp → WZ	X	X	X		X
pp → W/Z	X	X	X	(X)	X
pp → W+j	X	X	(X)	X	
pp → Z+j	X	(X)			
pp → VH	(X)				
pol. VBS	X	X			

Talks by Christopher, Christoph and Mareen on Thursday

Polarised nLO+PS: SHERPA

Polarised cross sections for vector boson production with SHERPA Hoppe, Schönherr, Siegert 2310.14803

- New bookkeeping of boson polarizations in SHERPA for LO MEs
- Approximate NLO corrections: nLO+PS
 - → Reals+matching are treated exact
 - → loop matrix elements unpolarised
- Comparison with multi-jet merged calculations

Comparison with literature

 nLO+PS approximation in fair agreement with full NLO
 → good for polarization fractions

$W^{+}Z$	$\sigma^{ m NLO}$ [fb]	Fraction [%]	K-factor	$\sigma_{ m SHERPA}^{ m nLO+PS}$ [fb]	Fraction [%]	K-factor
full	35.27(1)		1.81	33.80(4)		
unpol	34.63(1)	100	1.81	33.457(26)	100	1.79
Laboratory frame						
L-U	8.160(2)	23.563(9)	1.93	7.962(5)	23.796(25)	1.91
T-U	26.394(9)	76.217(34)	1.78	25.432(21)	76.01(9)	1.75
int	0.066(10) (diff)	0.191(29)	2.00	0.064(7)	0.191(22)	2.40(40)
U-L	9.550(4)	27.577(14)	1.73	9.275(16)	27.72(5)	1.72
U-T	25.052(8)	72.342(31)	1.83	24.156(18)	72.20(8)	1.81
int	0.028(10) (diff)	0.081(29)	-0.49	0.026(7)	0.079(22)	-0.471(34)

28.2.24 Izmir COMETA Rene Poncelet – IFJ PAN 12

Polarised NLO+PS: POWHEG

Polarised-boson pairs at the LHC with NLOPS accuracy Pelliccioli, Zanderighi 2311.05220

- NLO QCD + PS in POWHEG-BOX-RES framework
- Study of PS (Pythia8) + hadronisation effects on fractions and differential distributions WW/WZ/ZZ
 - → 1-5% effect on distributions, but generally small impact on fractions (~1% effects)

state	$\sigma[\mathrm{fb}]\mathrm{LHE}$	ratio [/unp., %] LHE	$\sigma[fb]$ PS+hadr	ratio [/unp., %] PS+hadr			
Inclusive setup							
full off-shell	$98.36(3)_{-3.9\%}^{+4.8\%}$	101.20	$95.27(3)_{-3.9\%}^{+4.9\%}$	101.28			
unpolarised	$97.20(3)_{-3.9\%}^{+4.8\%}$	100	$94.07(3)_{-3.9\%}^{+4.9\%}$	100			
LL	$4.499(2)_{-2.3\%}^{+2.8\%}$	$4.63^{+0.13}_{-0.13}$	$4.359(2)_{-2.2\%}^{+2.8\%}$	$4.63^{+0.13}_{-0.13}$			
LT	$13.151(4)_{-5.7\%}^{+7.0\%}$	$13.53^{+0.28}_{-0.27}$	$12.730(5)^{+7.0\%}_{-5.7\%}$	$13.53^{+0.28}_{-0.28}$			
\mathbf{TL}	$12.724(4)_{-5.9\%}^{+7.3\%}$	$13.09^{+0.32}_{-0.31}$	$12.314(5)^{+7.4\%}_{-5.9\%}$	$13.09^{+0.31}_{-0.32}$			
${f TT}$	$66.88(2)_{-3.3\%}^{+4.0\%}$	$68.81^{+0.47}_{-0.51}$	$64.74(2)_{-3.2\%}^{+4.1\%}$	$68.82^{+0.46}_{-0.51}$			
interference	-0.058	-0.06	-0.069	-0.06			

28.2.24 Izmir COMETA Rene Poncelet – IFJ PAN 13

NNLO QCD polarized WW production

NNLO QCD study of polarised W+W- production at the LHC, Poncelet, Popescu 2102.13583

Technical aspects:

- Implementation of NNLO QCD in c++ sector-improved residue subtraction framework [1408.2500,1907.12911]
- Massive b-quarks \rightarrow get rid of top production ($pp \rightarrow b\bar{b}W^+W^-$ enters at NNLO)
- NNPDF31 and a fixed renormalisation scale: $\mu_R = \mu_F = m_W$

Fiducial phase space

Measurement of fiducial and differential W+W- production crosssections at sqrt(s) = 13 TeV with the ATLAS detector ATLAS 1905.04242

- Leptons: $p_T(\ell) \geq 27 \text{ GeV}$ $|y(\ell)| < 2.5$ $m(\ell \bar{\ell}) > 55 \text{ GeV}$
- Missing transverse momentum: $p_{T, \text{miss}} = p_T(\nu_e + \bar{\nu}_\mu) \ge 20 \text{ GeV}$
- Jet-veto: $p_T(j) > 35 \text{ GeV} \quad |y(j)| < 4.5$

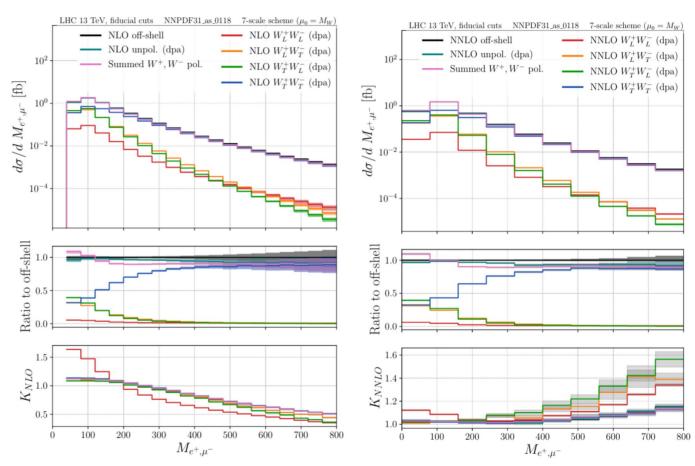
Doubly polarised cross sections

	NLO	NNLO	K_{NNLO}	LI	NNLO+LI
off-shell	$(220.060)^{+1.8\%}_{-2.3\%}$	$225.4(4)_{-0.6\%}^{+0.6\%}$	1.024	$13.8(2)^{+25.5\%}_{-18.7\%}$	$239.1(4)_{-1.2\%}^{+1.5\%}$
unpol. (nwa)	$221.85(8)_{-2.3\%}^{+1.8\%}$	$227.3(6)_{-0.6\%}^{+0.6\%}$	1.025	$13.68(3)^{+25.5\%}_{-18.7\%}$	$241.0(6)_{-1.1\%}^{+1.5\%}$
unpol. (dpa)	$214.55(7)_{-2.3\%}^{+1.8\%}$	$219.4(4)_{-0.6\%}^{+0.6\%}$	1.023	$13.28(3)^{+25.5\%}_{-18.7\%}$	$232.7(4)_{-1.1\%}^{+1.4\%}$
W_L^+ (dpa)	$57.48(3)_{-2.6\%}^{+1.9\%}$	$59.3(2)_{-0.7\%}^{+0.7\%}$	1.032	$2.478(6)_{-18.3\%}^{+25.5\%}$	$61.8(2)_{-0.8\%}^{+1.0\%}$
W_L^- (dpa)	$63.69(5)_{-2.6\%}^{+1.9\%}$	$65.4(3)^{+0.8\%}_{-0.8\%}$	1.026	$2.488(6)_{-18.3\%}^{+25.5\%}$	$67.9(3)_{-0.8\%}^{+0.9\%}$
W_T^+ (dpa)	$152.58(9)^{+1.7\%}_{-2.1\%}$	$155.7(6)_{-0.6\%}^{+0.7\%}$	1.020	$11.19(2)^{+25.5\%}_{-18.8\%}$	$166.9(6)_{-1.3\%}^{+1.6\%}$
W_T^- (dpa)	$156.41(7)^{+1.7\%}_{-2.1\%}$	$159.7(6)_{-0.6\%}^{+0.5\%}$	1.021	$11.19(2)^{+25.5\%}_{-18.8\%}$	$170.9(6)_{-1.3\%}^{+1.7\%}$
$W_L^+W_L^-$ (dpa)	$9.064(6)_{-3.0\%}^{+3.0\%}$	$9.88(3)_{-1.3\%}^{+1.3\%}$	1.090	$0.695(2)^{+25.5\%}_{-18.8\%}$	$10.57(3)^{+2.9\%}_{-2.4\%}$
$W_L^+W_T^-$ (dpa)	$48.34(3)_{-2.5\%}^{+1.9\%}$	$49.4(2)_{-0.7\%}^{+0.9\%}$	1.021	$1.790(5)^{+25.5\%}_{-18.3\%}$	$51.2(2)_{-0.8\%}^{+0.6\%}$
$W_T^+W_L^-$ (dpa)	$54.11(5)_{-2.5\%}^{+1.9\%}$	$55.5(4)^{+0.6\%}_{-0.7\%}$	1.025	$1.774(5)^{+25.5\%}_{-18.3\%}$	$57.2(4)_{-0.7\%}^{+0.7\%}$
$W_T^+W_T^-$ (dpa)	$106.26(4)_{-1.9\%}^{+1.6\%}$	$108.3(3)^{+0.5\%}_{-0.5\%}$	1.019	$9.58(2)_{-18.9\%}^{+25.5\%}$	$117.9(3)^{+2.1\%}_{-1.6\%}$

Small LL contribution, with large corrections (→ polarization frame)

Polarised di-boson production

- Longitudinal contribution largest around production threshold.
- At high energy W
 effectively massless
 → transverse polarised



28.2.24 Izmir COMETA Rene Poncelet – IFJ PAN 16

Take home messages

- Higher-order corrections are vital to pin down polarization fractions and to minimize theory systematics
- NLO QCD/EW (+PS) are the state-of-the-art for polarized EW boson processes
 → WW is available at NNLO QCD
- Future/mid-term goals:
 - fixed-order: completion of di-boson processes @ NNLO QCD (+ NLO EW)
 → Comparisons between NNLO QCD and NLO+PS calculations
 - event-generators: NNLO QCD/EW-effects, SMEFT

Backup

Polarized VV @ (N)NLO QCD / NLO EW

Fiducial polarization observables in hadronic WZ production: A next-to-leading order QCD+EW study,

Baglio, Le Duc 1810.11034

Anomalous triple gauge boson couplings in ZZ production at the LHC and the role of Z boson polarizations,

Rahama, Singh 1810.11657

Polarization observables in WZ production at the 13 TeV LHC: Inclusive case,

Baglio, Le Duc 1910.13746

Unravelling the anomalous gauge boson couplings in ZW+- production at the LHC and the role of spin-1 polarizations,

Rahama, Singh 1911.03111

Polarized electroweak bosons in W+W- production at the LHC including NLO QCD effects,

Denner, Pelliccioli 2006.14867

NLO QCD predictions for doubly-polarized WZ production at the LHC,

Denner, Pelliccioli 2010.07149

NNLO QCD study of polarised W+W- production at the LHC,

Poncelet, Popescu 2102.13583

NLO EW and QCD corrections to polarized ZZ production in the four-charged-lepton channel at the LHC,

Denner, Pelliccioli 2107.06579

Breaking down the entire spectrum of spin correlations of a pair of particles involving fermions and gauge bosons,

Rahama, Singh 2109.09345

Doubly-polarized WZ hadronic cross sections at NLO QCD+EW accuracy,

Duc Ninh Le, Baglio 2203.01470

Doubly-polarized WZ hadronic production at NLO QCD+EW: Calculation method and further results

Duc Ninh Le, Baglio, Dao 2208.09232

NLO QCD corrections to polarised di-boson production in semi-leptonic final states

Denner, Haitz, Pelliccioli 2211.09040

Polarised cross sections for vector boson production with SHERPA

Hoppe, Schönherr, Siegert 2310.14803

Polarised-boson pairs at the LHC with NLOPS accuracy

Pelliccioli, Zanderighi 2311.05220

NLO EW corrections to polarised W+W- production and decay at the LHC

Denner, Haitz, Pelliccioli 2311.16031

NLO electroweak corrections to doubly-polarized W+W- production at the LHC

Thi Nhung Dao, Duc Ninh 2311.17027

Polarized ZZ pairs in gluon fusion and vector boson fusion at the LHC

Javurkova, Ruiz, Coelho, Sandesara 2401.17365

Other polarized cross section calculations

Polarised VBS (so far LO):

W boson polarization in vector boson scattering at the LHC,

Ballestrero, Maina, Pelliccioli 1710.09339

Polarized vector boson scattering in the fully leptonic WZ and ZZ channels at the LHC,

Ballestrero, Maina, Pelliccioli 1907.04722

Automated predictions from polarized matrix elements

Buarque Franzosi, Mattelaer, Ruiz, Shil 1912.01725

Different polarization definitions in same-sign WW scattering at the LHC,

Ballestrero, Maina, Pelliccioli 2007.07133

Single boson production

Left-Handed W Bosons at the LHC,

Z. Bern et. al. 1103.5445

Electroweak gauge boson polarisation at the LHC,

Stirling, Vryonidou 1204.6427

What Does the CMS Measurement of W-polarization Tell Us about the Underlying Theory of the Coupling of W-Bosons to Matter?,

Belyaev, Ross 1303.3297

Polarised W+j production at the LHC: a study at NNLO QCD accuracy,

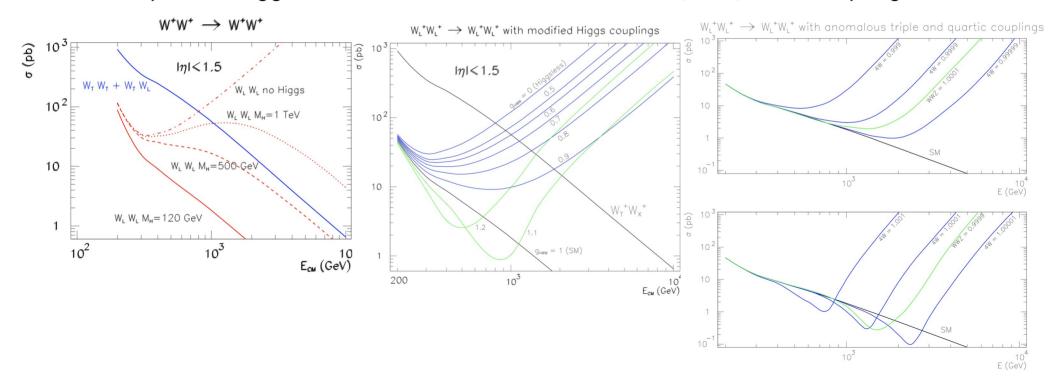
Pellen, Poncelet, Popescu 2109.14336

Longitudinal Vector-Boson-Scattering (VBS)

The Higgs boson and the physics of WW scattering before and after Higgs discovery M. Szleper 1412.8367

Sensitivity to the Higgs mass

Modified HVV, VVV, VVVV couplings



EWSB

The reason is the EWSB in the SM:

$$\mathcal{L}_{EW} = -\frac{1}{4} (W_{\mu\nu}^i)^2 - \frac{1}{4} (B_{\mu\nu}^i)^2 + (D_{\mu}\phi)^2 - V(\phi^{\dagger}\phi)$$

Higgs potential and minimum:

$$V(\phi^{\dagger}\phi) = -\mu^2(\phi^{\dagger}\phi)^2 + \lambda(\phi^{\dagger}\phi)^4$$
 $\phi = U(\pi^i) \begin{pmatrix} 0 \\ \frac{v+H}{\sqrt{2}} \end{pmatrix}$ VEV: $\phi^{\dagger}\phi = \frac{\mu^2}{2\lambda} \equiv \frac{v^2}{2}$

• Goldstone bosons can be absorbed via gauge transformation (unitary gauge). This gives rise to massive gauge bosons:

$$\phi = U^{-1}(\pi^i)\phi, \qquad W_{\mu} = U^{-1}W_{\mu}U - \frac{i}{g_W}U^{-1}\partial_{\mu}U$$

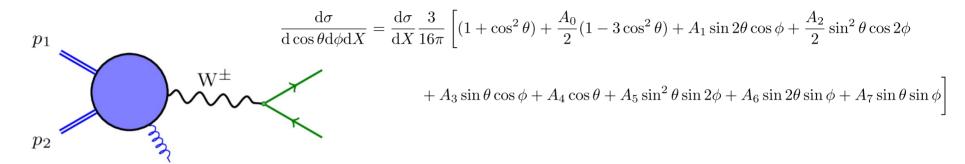
$$|D_{\mu}\phi|^{2} \ni \frac{v^{2}}{8} \left[2g_{W}^{2}W_{\mu}^{+}W^{-\mu} + (g_{W}W_{\mu}^{3} - g_{W}^{\prime}B_{\mu})^{2} \right] \longrightarrow M_{W} = \frac{1}{2}vg_{W} , \quad M_{Z} = \frac{M_{W}}{\cos\theta_{W}}$$

Restores renormalizability and unitarity

22

Angular coefficients

Angular decomposition of 2-body W decay:



After azimuthal integration:

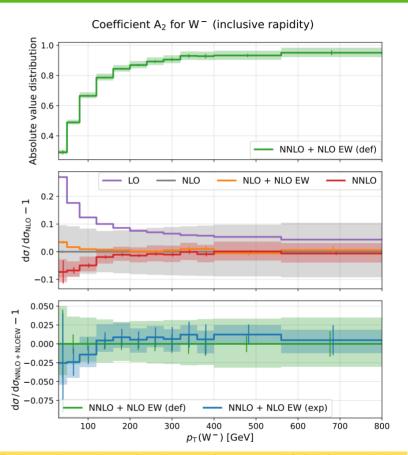
$$\frac{1}{\sigma} \frac{d\sigma}{\cos \theta} = \frac{3}{4} \sin \theta f_0 + \frac{3}{8} (1 - \cos \theta)^2 f_L + \frac{3}{8} (1 + \cos \theta)^2 f_R$$

Idea: Suitable projections (or fits) extract fractions of left, right and longitudinal components.

Angular coefficients as function of V kinematics

Keeping azimuthal dependence & boson kinematics:

$$\begin{split} \frac{\mathrm{d}\sigma}{\mathrm{d}p_{\mathrm{T,W}}\,\mathrm{d}y_{\mathrm{W}}\,\mathrm{d}m_{\ell\nu}\,\mathrm{d}\Omega} = & \frac{3}{16\pi} \frac{\mathrm{d}\sigma^{U+L}}{\mathrm{d}p_{\mathrm{T,W}}\,\mathrm{d}y_{\mathrm{W}}\,\mathrm{d}m_{\ell\nu}} \bigg((1+\cos^2\theta) + \mathrm{A}_0 \frac{1}{2} (1-3\cos^2\theta) \\ & + \mathrm{A}_1\sin2\theta\cos\phi + \mathrm{A}_2 \frac{1}{2}\sin^2\theta\cos2\phi + \mathrm{A}_3\sin\theta\cos\phi + \mathrm{A}_4\cos\theta \\ & + \mathrm{A}_5\sin^2\theta\sin2\phi + \mathrm{A}_6\sin2\theta\sin\phi + \mathrm{A}_7\sin\theta\sin\phi \bigg), \end{split}$$



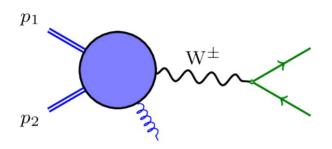
Angular coefficients in W+j production at the LHC with high precision Pellen, Poncelet, Popescu, Vitos, 2204.12394

Angular coefficients, practical considerations

This simple idea suffers from:

- Fiducial phase space requirements on the leptons:
 - → Interferences do not cancel
 - \rightarrow Correspondence between fractions (f_0, f_L, f_R) and angular distributions broken.
- Higher order corrections to decay (QED radiation or QCD in hadronic decays)
 - \rightarrow Decomposition in $\{A_i\}$ does not hold any more
- Angles in boson rest frame
 - → Z rest frame accessible, but W more difficult to reconstruct

Polarised W+jet cross sections



Why looking at polarised W+jet with leptonic decays?

- The EW part is simple:
 - no non-resonant backgrounds
 - neutrino momentum approx. accessible (missing ET)
- Large cross section → precise measurements

Goals:

- Use W+j data to extract the longitudinal polarisation fraction (done before by exp.)
 → understand impact of NNLO QCD corrections (reduced scale dependence)
- Study inclusive (in terms of W decay products) and fiducial phase spaces
 → How does the sensitivity to longitudinal Ws depend on this?
 Which observables have small interference/off-shell effects?
- Are there any differences between W+ and W-?
 From PDFs and the fact that we cut on the charged lepton?

Setup W+jet: LHC @ 13 TeV

Polarised W+j production at the LHC: a study at NNLO QCD accuracy, Pellen, Poncelet, Popescu 2109.14336

Inclusive phase space:

• At least one jet with $|y(j)| \le 2.4$ and $p_T(j) \ge 30 \text{ GeV}$

Fiducial phase space:

Measurement of the differential cross sections for the associated production of a W boson and jets in proton-proton collisions at \sqrt{s}=13 TeV, CMS 1707.05979

- Lepton cuts: $p_T(\ell) \geq 25 \; \mathrm{GeV}$, $|\eta(\ell)| \leq 2.5$ and $\Delta R(\ell,j) > 0.4$
- Transverse mass of the W: $M_T(W) = \sqrt{m_W^2 + p_T^2(W)} \ge 50 \; \mathrm{GeV}$

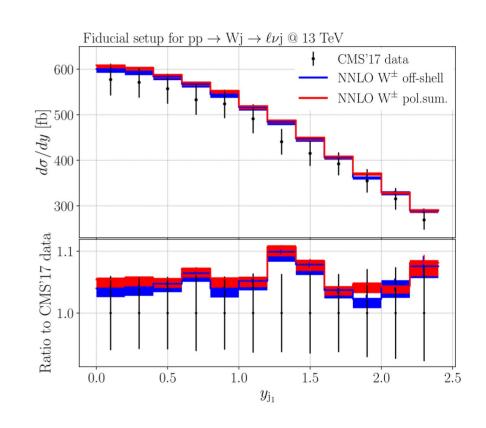
Technical aspects:

- NNPDF31 and dynamical scale choice: $\mu_R = \mu_F = \frac{1}{2} \left(m_T(W) + \sum p_T(j) \right)$
- Implementation in STRIPPER framework (NNLO QCD subtractions) [1408.2500]
 - Narrow-Width-Approximation and OSP/Pole-Approximation
 - Matrix elements from: AvH[1503.08612], OpenLoops2 [1907.13071](cross checks with Recola [1605.01090]) and VVamp [1503.04812]

Extraction of polarisation fractions

Identified 4 observables (ranges) with

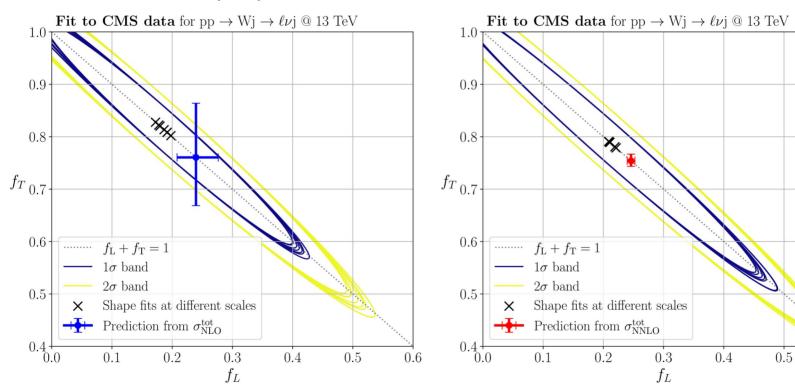
- → Small interference effects (<2%)
- → Small off-shell effects (<2%)
- → Shape differences between L and T
 - $\Delta \phi(\ell, j_1) \ge 0.3$
 - $25 \text{ GeV} \leq p_T(\ell) < 70 \text{ GeV}$
 - $\cos(\theta_{\ell}^*) \ge -0.75$
 - $|y(j_1)| \le 2$



28

W+jet: fit to CMS data

Fit to actual data, here $|y(j_1)|$ \rightarrow dominated by experimental uncertainties (no correlations available)



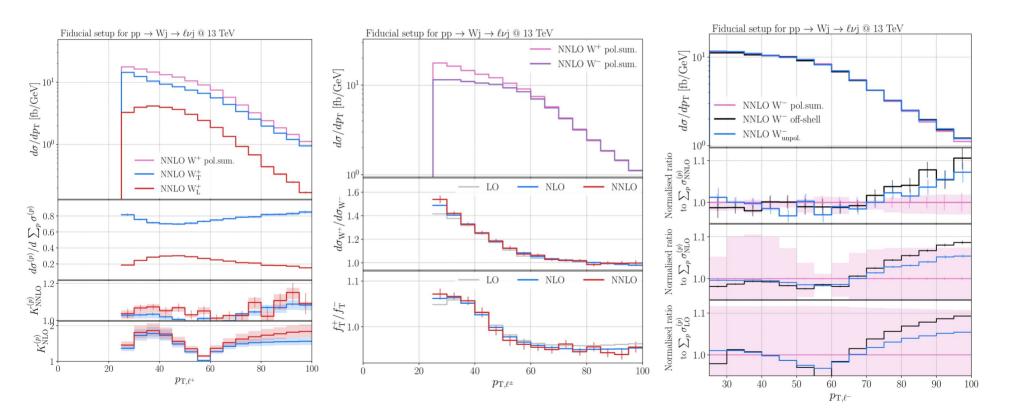
0.6

Example: lepton transverse momentum



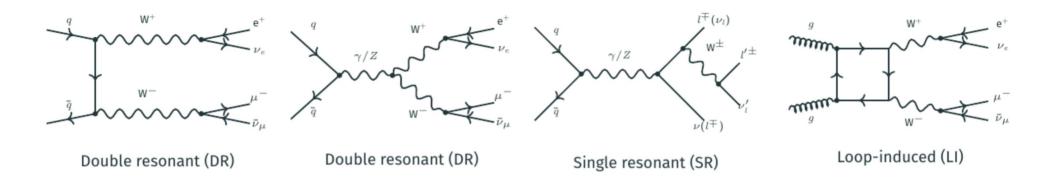
Charge differences

Off-shell/Interference effects

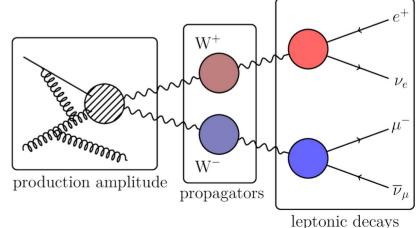


Rene Poncelet – IFJ PAN 30

W-boson pair production

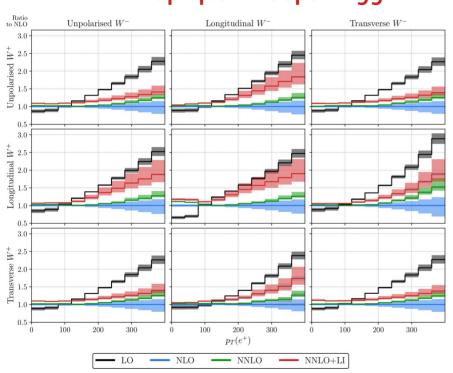


- Single resonant backgrounds: Definition of polarizations states in DPA [1710.09339] and NWA
- LI enters at NNLO → large corrections

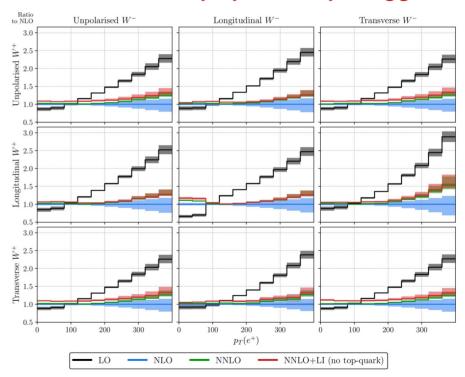


Loop induced gg → WW contributions

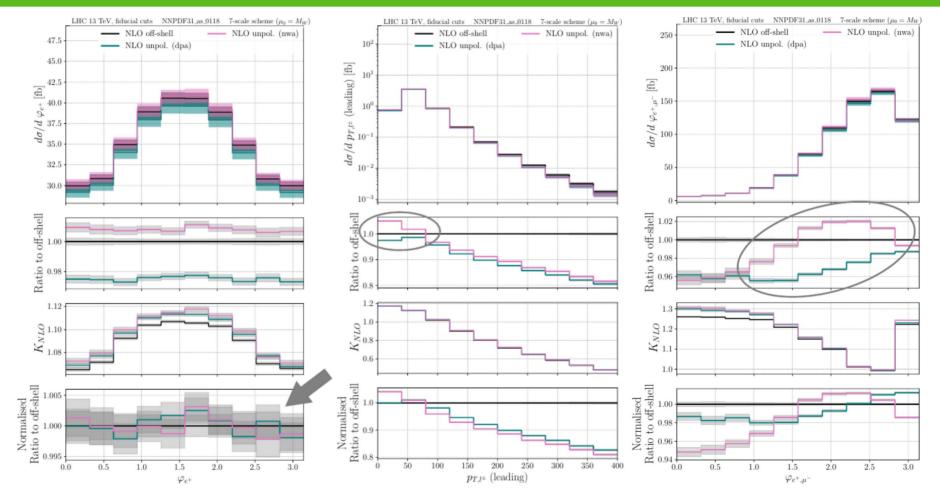
With top-quark loops in gg LI



Without top-quark loops in gg LI



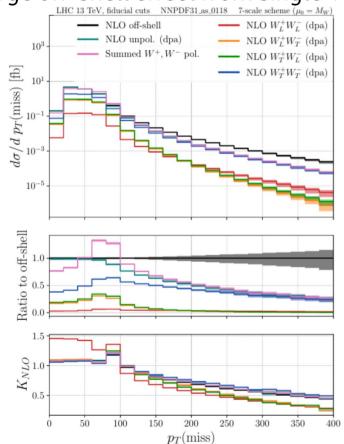
NWA vs. DPA

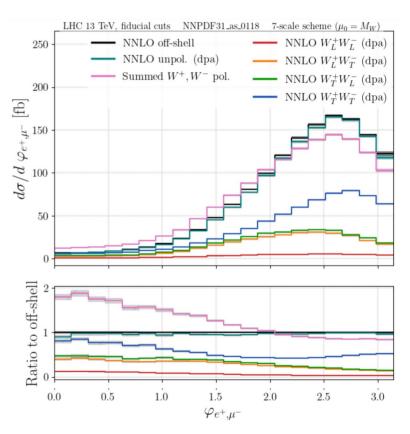


Rene Poncelet – IFJ PAN

Interference and off-shell effects

Large off-shell effect from single-resonant contributions





Large interference effects through phase space constraints

28.2.24 Izmir COMETA Rene Poncelet – IFJ PAN