

Precision Physics in ZH/ZZ/HH

Marco Vitti (Karlsruhe Institute of Technology, TTP and IAP)

1st COMETA General Meeting, Izmir, Feb 29 2024



- Recent developments in SM predictions
 - Emphasis on fixed-order
 - Massive loops in *gg*-initiated processes
-



See
[https://indico.cern.ch/
event/1360356/](https://indico.cern.ch/event/1360356/)
for dedicated discussions


Topical on-line WG1 meeting - 2 hours for 2 bosons






Wednesday Jan 17, 2024, 2:30 PM → 5:00 PM Europe/Rome

Giovanni Pelliccioli (Max-Planck-Institut für Physik), Ilaria Brivio (University & INFN Bologna),
Ramona Groeber (Università di Padova and INFN, Sezione di Padova)

Description Topical meeting of the COMETA WG1 (theoretical understanding) on diboson production (VV, VH, HH)

  Zoom recording

Videoconference  Topical on-line WG1 meeting - 2 hours for 2 bosons [Join](#)

2:30 PM → 2:40 PM	Introduction	🕒 10m
	Speakers: Giovanni Pelliccioli (Max-Planck-Institut für Physik), Ramona Groeber (Università di Padova and INFN, Sezione di Padova)	
	 gp_cometa_wg1_1s...	
2:40 PM → 3:10 PM	VV	🕒 30m
	Speaker: Marius Wiesemann (Max Planck Institute for Physics)	
	 Wiesemann_COME...	
3:10 PM → 3:40 PM	VLVL	🕒 30m
	Speaker: Rene Poncelet (IFJ PAN Krakow)	
	 2024.01.17-COMET...	
3:40 PM → 4:10 PM	VH	🕒 30m
	Speaker: Stephen Philip Jones (University of Durham (GB))	
	 talk_cometa_sj.pdf	
4:10 PM → 4:40 PM	HH	🕒 30m
	Speaker: Michael Spira (Paul Scherrer Institute (CH))	
	 spira.pdf	
4:40 PM → 5:00 PM	Discussion	🕒 20m

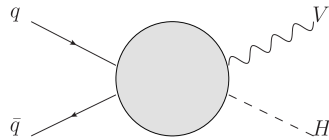
ZH

VH Production at the LHC

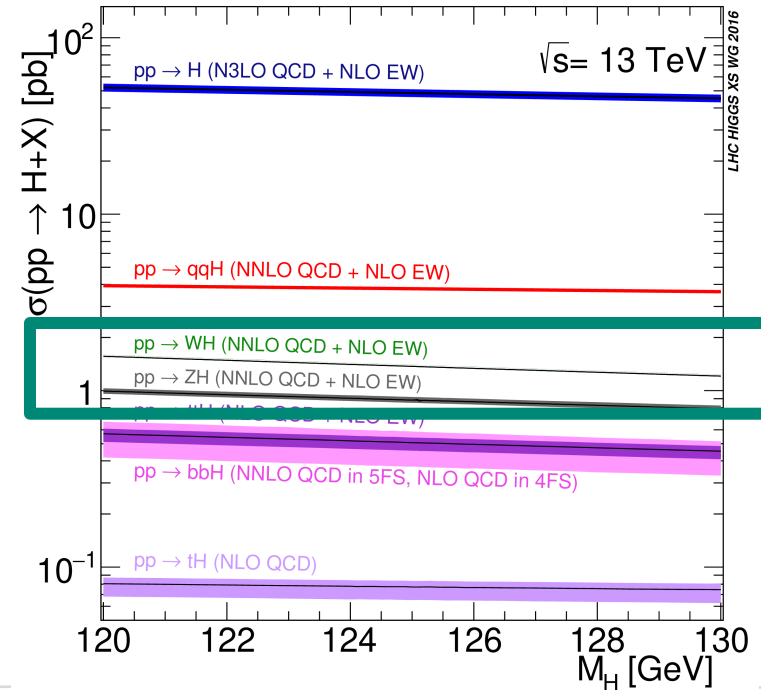
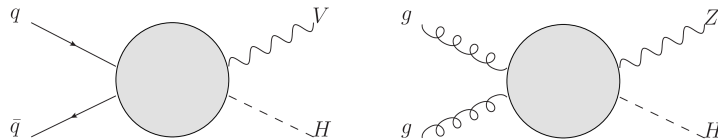
$pp \rightarrow VH$ is the most sensitive process to $H \rightarrow b\bar{b}$ [ATLAS-2007.02873, CMS-1808.08242]

- Work in progress on $H \rightarrow c\bar{c}$
[ATLAS-2201.11428, CMS-2205.0555]
- Probe of VVH coupling
- Partonic channels

$pp \rightarrow WH$



$pp \rightarrow ZH$



VH Production at the LHC

$pp \rightarrow VH$ is the most sensitive process to $H \rightarrow b\bar{b}$ [ATLAS-2007.02873, CMS-1808.08242]

- Work in progress on $H \rightarrow c\bar{c}$
[ATLAS-2201.11428, CMS-2205.0555]
- Probe of VVH coupling
- Larger scale uncertainties in ZH

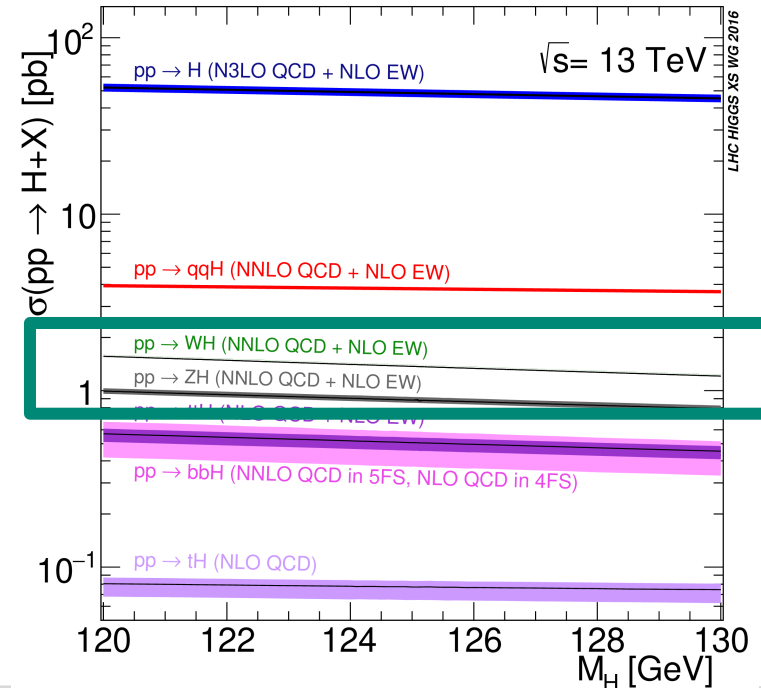
$pp \rightarrow WH$

\sqrt{s} [TeV]	$\sigma_{\text{NNLO QCD} \otimes \text{NLO EW}}$ [pb]	Δ_{scale} [%]	$\Delta_{\text{PDF} \oplus \alpha_s}$ [%]
13	1.358	+0.51 -0.51	1.35
14	1.498	+0.51 -0.51	1.35
27	3.397	+0.29 -0.72	1.37

$pp \rightarrow ZH$

\sqrt{s} [TeV]	$\sigma_{\text{NNLO QCD} \otimes \text{NLO EW}}$ [pb]	Δ_{scale} [%]	$\Delta_{\text{PDF} \oplus \alpha_s}$ [%]
13	0.880	+3.50 -2.68	1.65
14	0.981	+3.61 -2.94	1.90
27	2.463	+5.42 -4.00	2.24

[Cepeda et al. - 1902.00134]



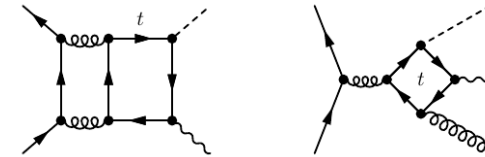
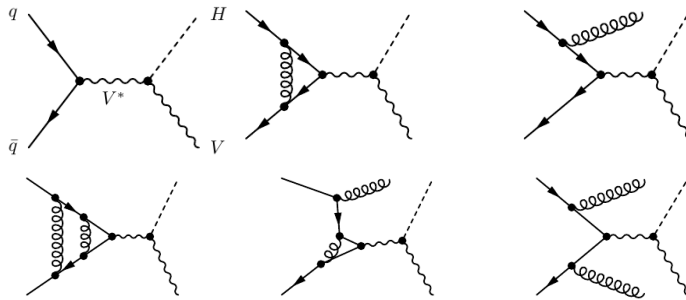
Theoretical Predictions for $pp \rightarrow ZH$

LO: quark-initiated tree-level contribution (purely EW)

QCD effects: mainly due to Drell-Yan (DY) production followed by $Z^* \rightarrow ZH$ decay

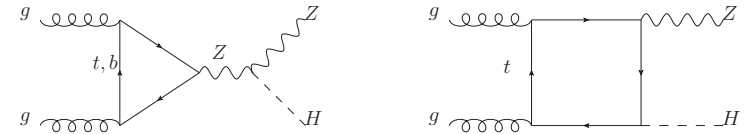
■ **Drell-Yan:** Known through N3LO (+30% wrt LO) ■ **Non Drell-Yan - quark-initiated** O(1%) wrt LO

[Brein, Harlander, Wiesemann, Zirke - 1111.0761]



[Han, Willenbrock ('91) ; Hamberg, van Neerven, Matsuura ('92) ;
Brein, Djouadi, Harlander – 0307206;
Baglio, Duhr, Mistlberger, Szafron - 2209.06138]

■ **Non Drell-Yan - gluon-initiated**



EW corrections: through NLO (-(5-10%) wrt LO)

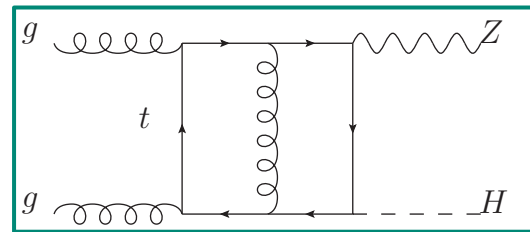
[Ciccolini, Dittmaier, Krämer - 0306234]

$gg \rightarrow ZH$ @ NLO QCD

Main problem in the virtual NLO calculation

Multi-scale (m_Z, m_H, m_t, s, t) two-loop box integrals

No full analytic results



Three (almost) independent calculations in agreement

- Small-mass expansion $m_Z, m_H \rightarrow 0$

[Wang, Xu, Xu, Yang - 2107.08206]

Elliptic integrals evaluated numerically

- Sector decomposition \oplus High-Energy expansion

[Chen et al. - 2011.12325; Davies et al. - 2011.12314; Chen, et al. - 2204.05225;]

$$m_Z^2, m_H^2 \ll m_t^2 \ll \hat{s}, \hat{t}$$

- pT expansion \oplus High-Energy expansion

[Alasfar et al. - 2103.06225; Bellafronte et al. - 2202.12157; Degraasi, Gröber, MV, Zhao - 2205.02769]

$$m_Z^2, m_H^2, p_T^2 \ll m_t^2, \hat{s} \quad m_Z^2, m_H^2 \ll m_t^2 \ll \hat{s}, \hat{t}$$

$gg \rightarrow ZH$ @ NLO QCD

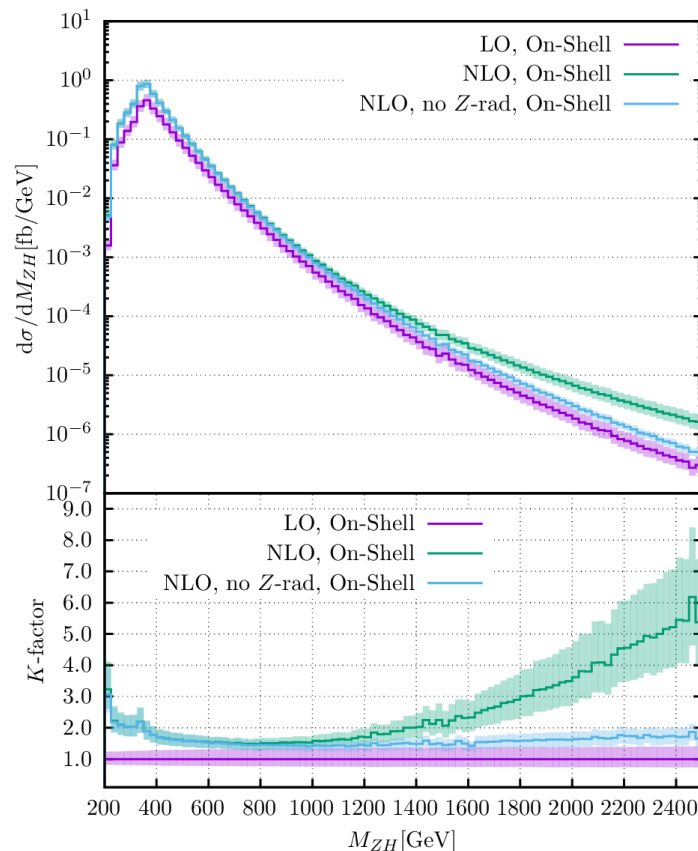
Inclusive cross section $\sqrt{s} = 13\text{TeV}$
 $\mu_r = \mu_f = M_{ZH}/2$

Top-mass scheme	LO [fb]	$\sigma_{LO}/\sigma_{LO}^{OS}$	NLO [fb]	$\sigma_{NLO}/\sigma_{NLO}^{OS}$	$K = \sigma_{NLO}/\sigma_{LO}$
On-Shell	64.01 ^{+27.2%} _{-20.3%}	—	118.6 ^{+16.7%} _{-14.1%}	—	1.85
$\overline{\text{MS}}, \mu_t = M_{ZH}/4$	59.40 ^{+27.1%} _{-20.2%}	0.928	113.3 ^{+17.4%} _{-14.5%}	0.955	1.91
$\overline{\text{MS}}, \mu_t = m_t^{\overline{\text{MS}}}(m_t^{\overline{\text{MS}}})$	57.95 ^{+26.9%} _{-20.1%}	0.905	111.7 ^{+17.7%} _{-14.6%}	0.942	1.93
$\overline{\text{MS}}, \mu_t = M_{ZH}/2$	54.22 ^{+26.8%} _{-20.0%}	0.847	107.9 ^{+18.4%} _{-15.0%}	0.910	1.99
$\overline{\text{MS}}, \mu_t = M_{ZH}$	49.23 ^{+26.6%} _{-19.9%}	0.769	103.3 ^{+19.6%} _{-15.6%}	0.871	2.10

- NLO corrections are the same size as LO ($K \sim 2$)
- Scale uncertainties reduced by 30% wrt LO

Invariant-mass distribution

- K -factor is not flat over M_{ZH} range
- Large NLO enhancement in the high-energy tail ($M_{ZH} > 1\text{TeV}$)



[Degrassi, Gröber, MV, Zhao - 2205.02769]

$gg \rightarrow ZH$ @ NLO QCD

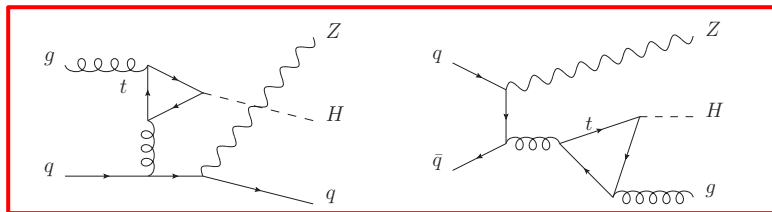
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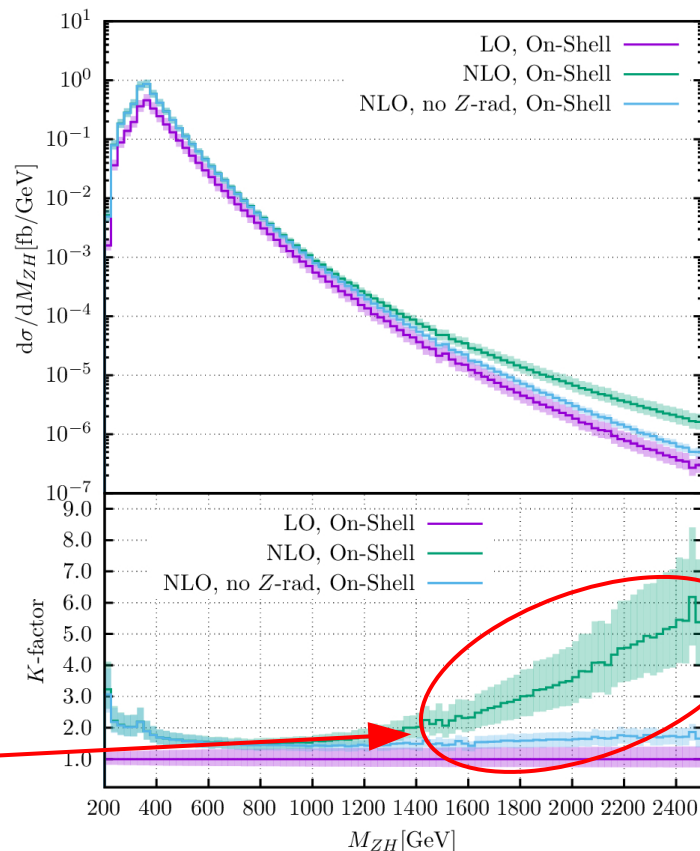
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Invariant-mass distribution



Dominant

PDF suppressed



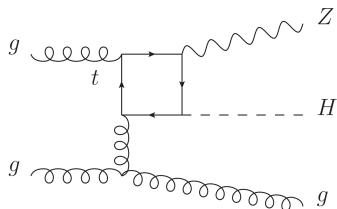
[Degrassi, Gröber, MV, Zhao - 2205.02769]

High-Energy Tails – pT Distributions

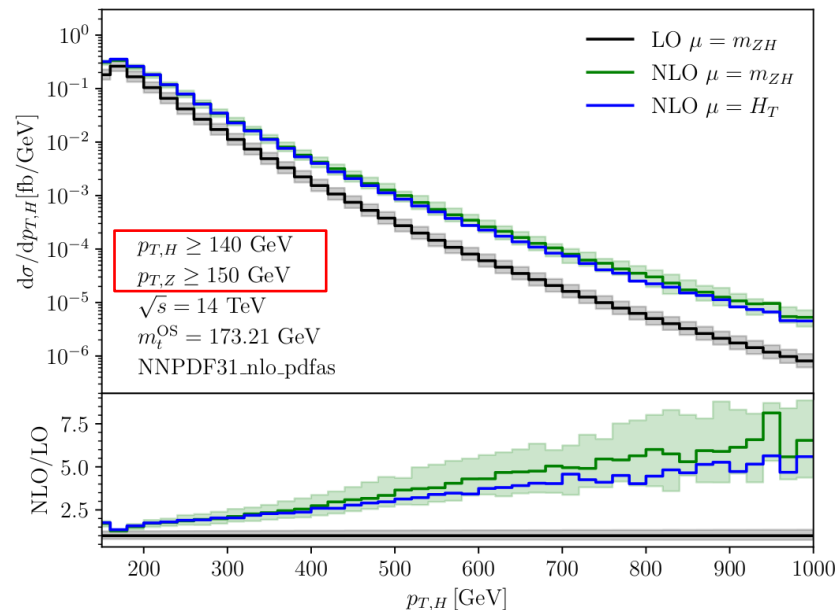
- Very large (~ 20) NLO corrections for $p_{T,H} > 400$ GeV
- Still K-factor of ~ 5 after pT cuts

- The pT cuts remove 2 \rightarrow 3 configurations with a hard jet and a soft Z

These are very likely in the high-energy region



Observed in
[Hespel, Maltoni, Vryonidou
-1503.01656]



Top Mass Scheme Uncertainty

- Envelope of deviations of $\overline{M_S}$ schemes wrt OS result
Same method already used for HH production

[Baglio et al. - 1811.05692, 2003.03227]

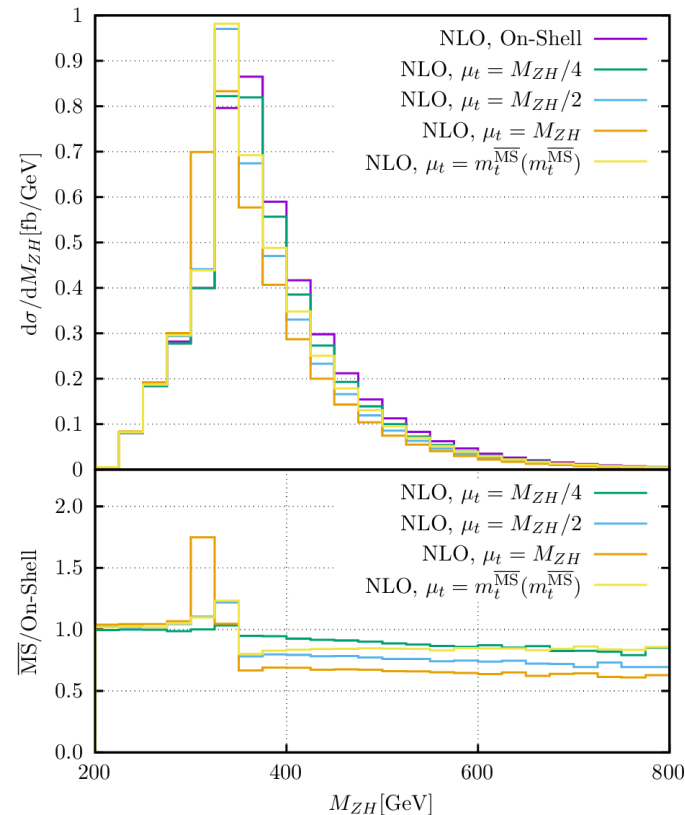
- Uncertainty sensitive to the binning of top-pair threshold peak

Bin Width [GeV]	LO	NLO
1	64.01 ^{+15.6%} -35.9%	118.6 ^{+17.2%} -27.0%
5	64.01 ^{+15.3%} -35.6%	118.6 ^{+14.7%} -24.9%
25	64.01 ^{+14.0%} -33.1%	118.6 ^{+10.9%} -20.8%
100	64.01 ^{+2.0%} -25.3%	118.6 ^{+0.6%} -13.7%
∞	64.01 ^{+0%} -23.1%	118.6 ^{+0%} -12.9%

Avoid overestimate
of uncertainty



- Top-mass uncertainty ~ scale uncertainty



[Degrassi, Gröber, MV, Zhao - 2205.02769]

ZZ

$pp \rightarrow ZZ$ at the LHC

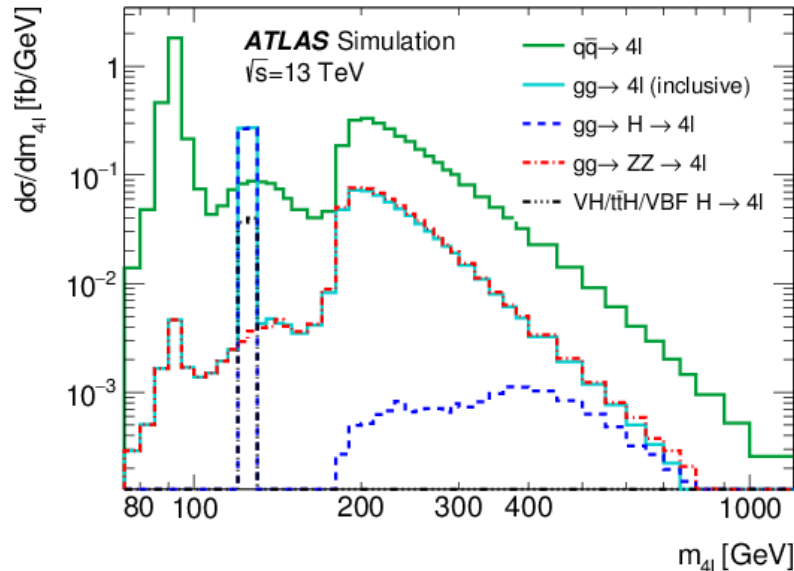
■ Probe of EW theory: VVV couplings, Higgs production, polarisations measurements...

■ Indirect access to Higgs width from off-shell measurements

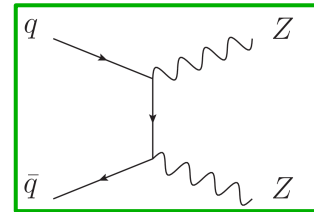
[Kauer, Passarino – 1206.4803]

[Caola, Melnikov – 1307.4935]

[Campbell, Ellis, Williams - 1311.3589]



[ATLAS – 1902.05892]



QCD: known through NNLO

[Cascioli et al. - 1405.2219; Heinrich et al. - 1710.06294;
Gehrmann et al. - 1404.4853; Caola et al. - 1408.6409;
Gehrmann et al. - 1503.04812; Grazzini et al. - 1507.06257;
Kallweit, Wiesemann - 1806.05941]

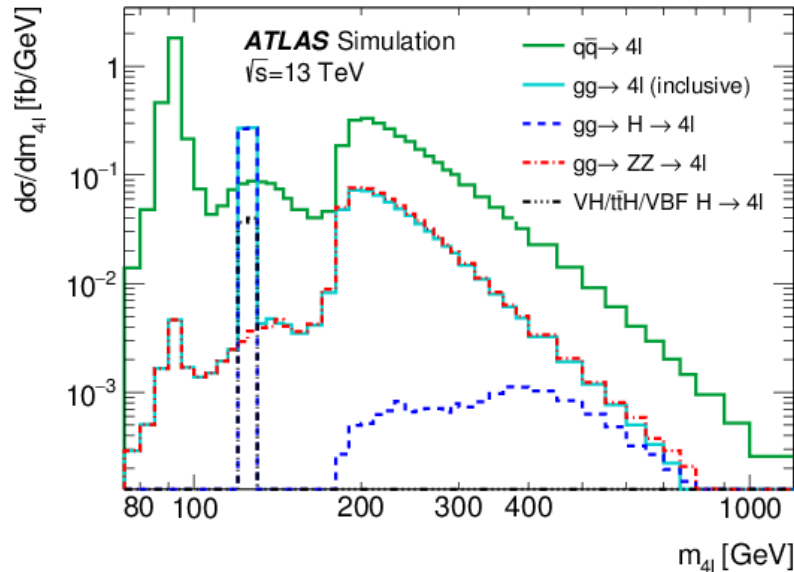
EW: known through NLO

[Bierweiler et al. – 1305.5402; Baglio, Ninh,
Weber – 1307.4331; Chiesa et al. - 2005.12146]

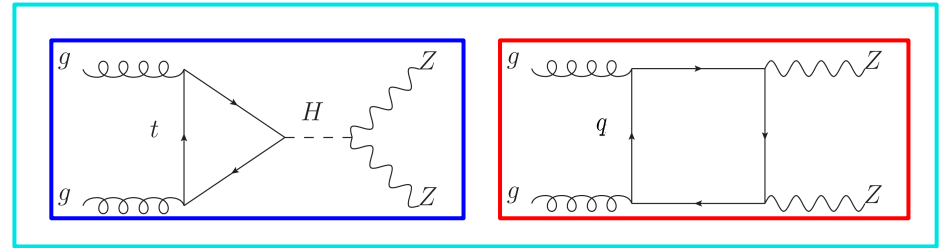
$pp \rightarrow ZZ$ at the LHC

■ Probe of EW theory: VVV couplings, Higgs production, polarisations measurements...

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[Campbell, Ellis, Williams - 1311.3589]



[ATLAS – 1902.05892]



Higgs-mediated: NLO QCD

[Spira et al. - 9504378 ; Aglietti et al. - 0611266 ;
Harlander, Kant - 0509189; Anastasiou et al. - 0611236]

Continuum ZZ production: NLO QCD

(light quark) [von Manteuffel, Tancredi – 1503.08835;
Caola et al. - 1509.06734]

Impact of $gg \rightarrow ZZ$

[Grazzini, Kallweit, Wiesemann, Yook - 1811.09593]

\sqrt{s}	8 TeV	13 TeV	8 TeV	13 TeV
	σ [fb]		$\sigma/\sigma_{\text{NLO}} - 1$	
LO	8.1881(8) ^{+2.4%} _{-3.2%}	13.933(1) ^{+5.5%} _{-6.4%}	-27.5%	-29.8%
NLO	11.2958(4) ^{+2.5%} _{-2.0%}	19.8454(7) ^{+2.5%} _{-2.1%}	0%	0%
$q\bar{q}$ NNLO	12.09(2) ^{+1.1%} _{-1.1%}	21.54(2) ^{+1.1%} _{-1.2%}	+7.0%	+8.6%
	σ [fb]		$\sigma/\sigma_{\text{ggLO}} - 1$	
gg LO	0.79355(6) ^{+28.2%} _{-20.9%}	2.0052(1) ^{+23.5%} _{-17.9%}	0%	0%
gg NLO _{gg}	1.4787(4) ^{+15.9%} _{-13.1%}	3.626(1) ^{+15.2%} _{-12.7%}	+86.3%	+80.8%
gg NLO	1.3892(4) ^{+15.4%} _{-13.6%}	3.425(1) ^{+13.9%} _{-12.0%}	+75.1%	+70.8%
	σ [fb]		$\sigma/\sigma_{\text{NLO}} - 1$	
NNLO	12.88(2) ^{+2.8%} _{-2.2%}	23.55(2) ^{+3.0%} _{-2.6%}	+14.0%	+18.7%
nNNLO	13.48(2) ^{+2.6%} _{-2.3%}	24.97(2) ^{+2.9%} _{-2.7%}	+19.3%	+25.8%

Top-quark loops \rightarrow permille level in inclusive cross section \rightarrow negligible

BUT important effects at high invariant masses

Interference @ NLO: massless vs massive quarks

Two-loop boxes are a problem (again)

- **Light**-quark (\sim massless)
known fully analytically
[Caola et al. - 1509.06734]

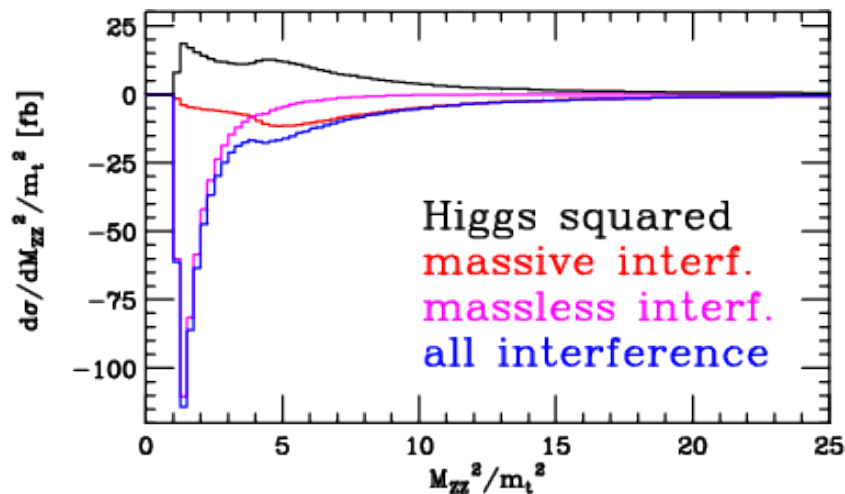
- **Heavy** quarks

- **Exact numerical** results available
[Agarwal, Jones, von Manteuffel - 2011.15113 ;
Brønnum-Hansen, Wang - 2101.12095]
- **Analytic approximations:**
 - $m_t \rightarrow \infty$ [Melnikov, Dowling - 1503.01274 ; Gröber,
Maier, Rauh - 1605.04610]
 - High-Energy exp [Davies et al. - 2002.05558]
 $m_Z^2 \ll m_t^2 \ll \hat{s}, \hat{t}$

$$2 \operatorname{Re} \left(\text{Diagram 1} * \text{Diagram 2} \right)$$

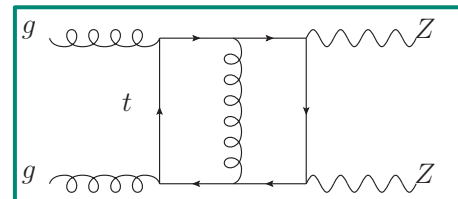
Diagram 1: A two-loop box diagram with a top quark loop and a Higgs boson exchange between two gluon lines.

Diagram 2: A two-loop box diagram with a top quark loop and a Z boson exchange between two gluon lines.



[Campbell et al. - 1605.01380]

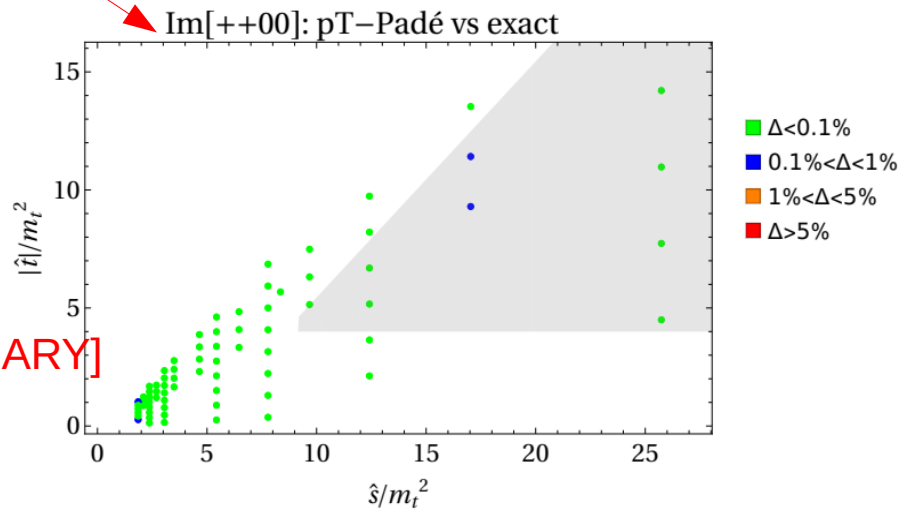
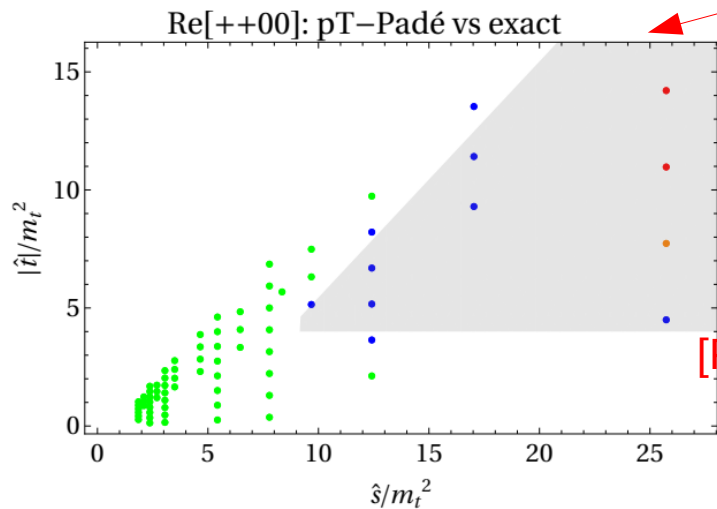
$gg \rightarrow ZZ$: Top Quark Loops @NLO QCD



Comparison with helicity amplitudes of

[Agarwal, Jones, von Manteuffel - 2011.15113]

$$\mathcal{M}_{\lambda_1, \lambda_2, \lambda_3, \lambda_4}^{\text{fin}} = \left(\frac{\alpha_s}{2\pi}\right) \mathcal{M}_{\lambda_1, \lambda_2, \lambda_3, \lambda_4}^{(1)} + \left(\frac{\alpha_s}{2\pi}\right)^2 \mathcal{M}_{\lambda_1, \lambda_2, \lambda_3, \lambda_4}^{(2)} + \mathcal{O}(\alpha_s^3)$$

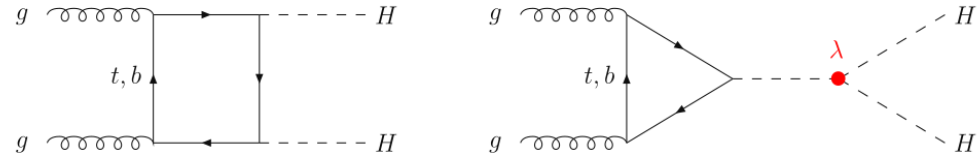


[PRELIMINARY]

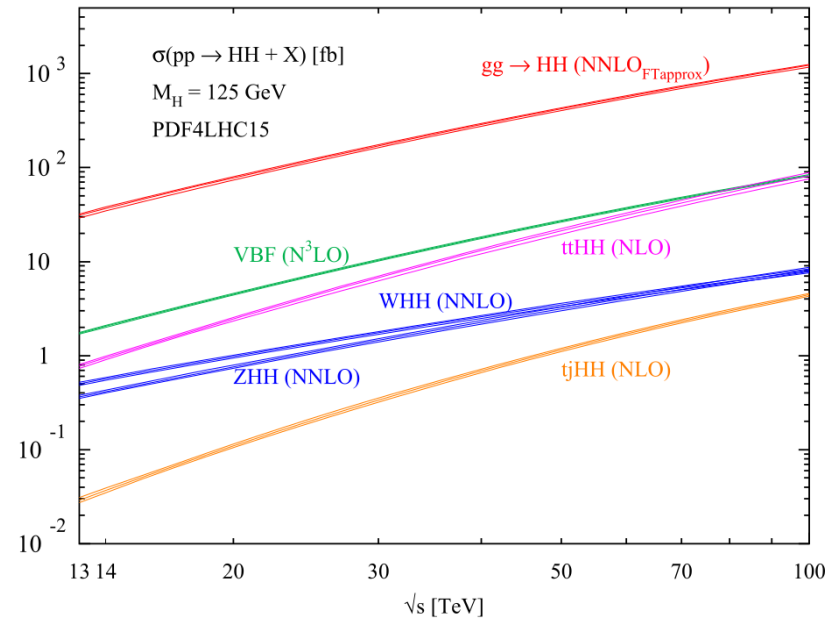
[Degrassi, Gröber, MV – in preparation]

HH

HH Production @LHC



- Best shot at measurement of Higgs trilinear self-coupling
- Gluon-initiated channel dominant



[Di Micco et al. - 1910.00012]

HH QCD corrections

$m_t \rightarrow \infty$ limit: N3LO

[De Florian, Mazzitelli 1305.5206 and 1309.6594; . Grigo, Melnikov and Steinhauser – 1408.2422; Chen et al. - 1909.06808 and 1912.13001;]

Finite $1/m_t$ effects

[Grigo, Hoff, Steinhauser - 1508.00909; Davies, Steinhauser – 1909.01361; Davies et al. 2110.03697]

Full top-mass dependence: NLO

■ Numerical evaluation

[Borowka et al. - 1604.06447, 1608.04798;
Baglio et al. - 1811.05692]

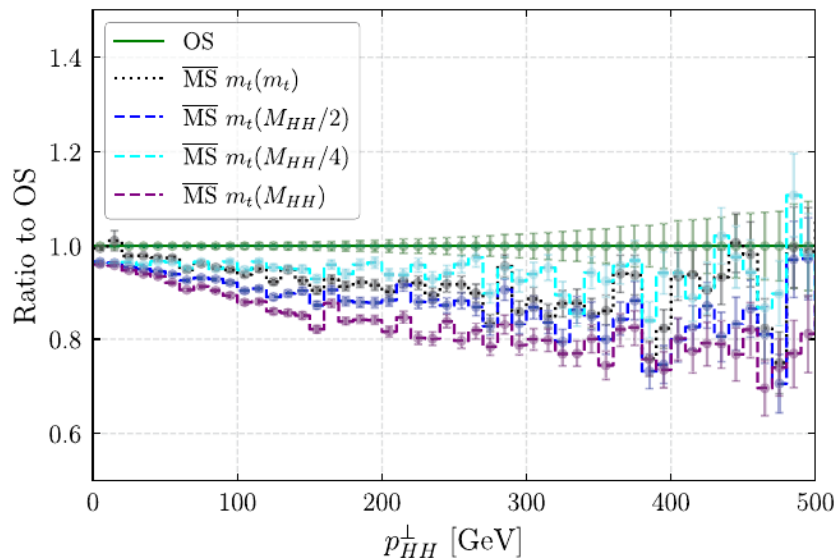
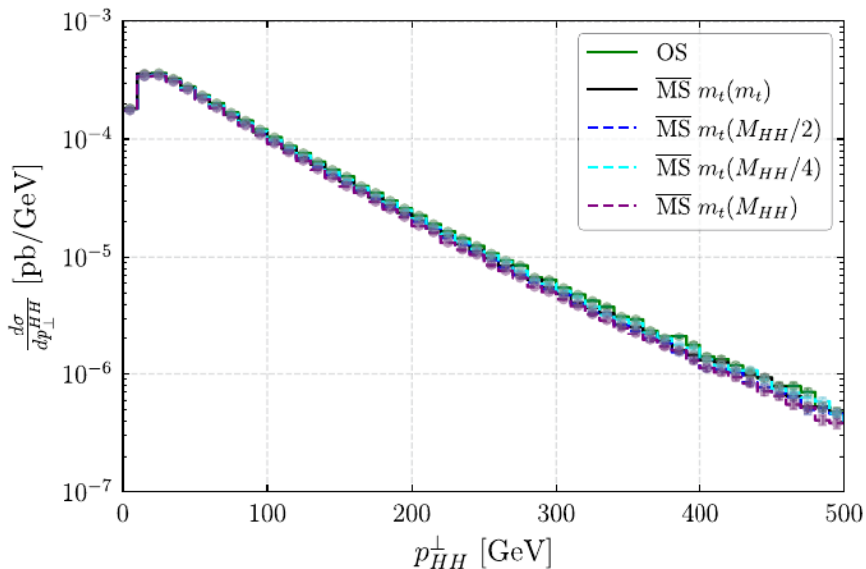
■ Analytic approximations

[Davies et al. - 1811.05489; Bonciani et al. - 1806.11564; Wang et al. - 2010.15649]
(High-energy exp) (pT expansion) Small-mass expansion

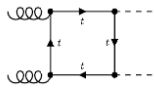
Full phase space covered in [Bellafrente et al. – 2202.12157; Davies et a. - 2302.01356]

HH @NLO QCD (+ PS)

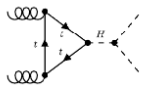
- Full flexibility over choice of trilinear and top mass
- Top mass scheme uncertainty studied differentially (see also [\[Baglio et al. - 2008.11626\]](#))



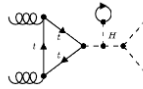
HH NLO EW



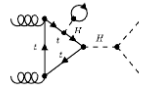
(a-1)



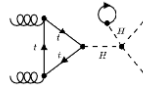
(a-2)



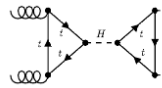
(b-1)



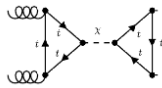
(b-2)



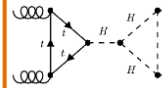
(b-3)



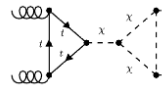
(c-1)



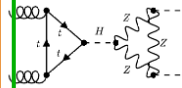
(c-2)



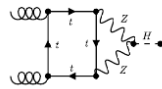
(c-3)



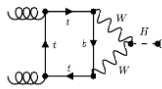
(c-4)



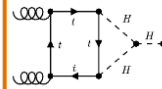
(c-5)



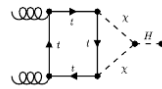
(d-1)



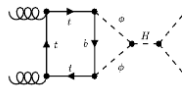
(d-2)



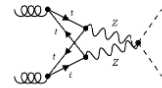
(d-3)



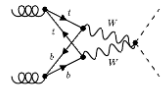
(d-4)



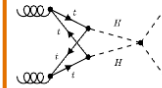
(d-5)



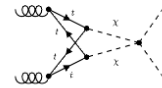
(e-1)



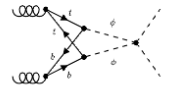
(e-2)



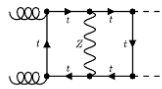
(e-3)



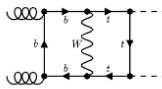
(e-4)



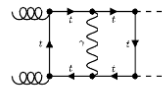
(e-5)



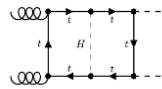
(f-1)



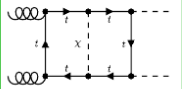
(f-2)



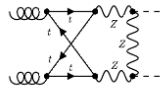
(f-3)



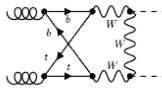
(f-4)



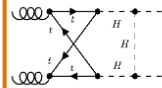
(f-5)



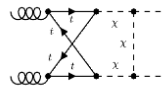
(g-1)



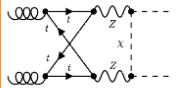
(g-2)



(g-3)



(g-4)



(g-5)

Trilinear and quadrilinear contributions
at two loops (Sector Decomposition)

[Borowka et al. - 1811.12366]

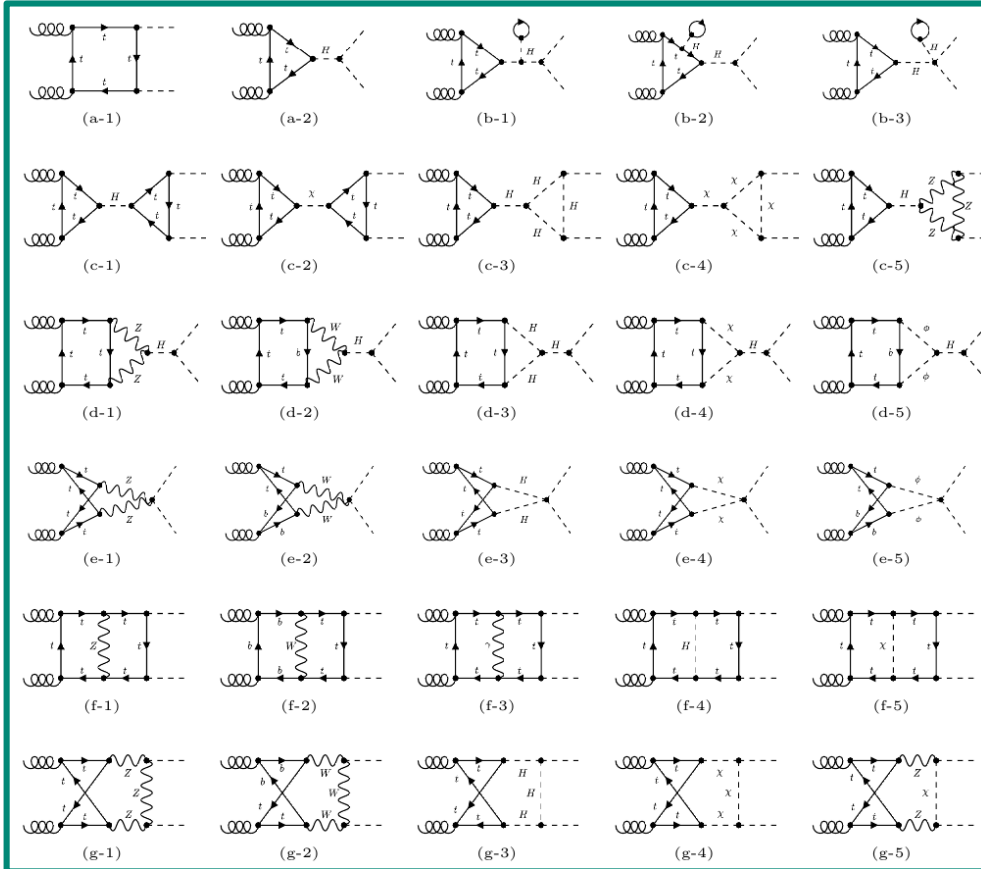
Top-Yukawa-induced corrections
(Heavy-Top EFT)

[Mühlleitner, Schlenk, Spira – 2207.02524]

Two loop boxes with virtual H
(High-Energy expansion)

[Davies, et al. - 2207.02587]

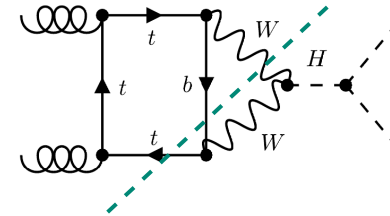
HH NLO EW



Complete NLO EW in Large Top Expansion

$$m_t^2 \gg \xi_W m_W^2, \xi_Z m_Z^2 \gg \hat{s}, \hat{t}, m_W^2, m_Z^2, m_H^2$$

- Typically valid for $\sqrt{\hat{s}} \lesssim 2m_t \sim 350\text{GeV}$
Compromised by diagrams with t-W-b cut



$$\sqrt{\hat{s}} \lesssim m_t + m_W + m_b \sim 250\text{GeV}$$

HH NLO EW

Full numerical evaluation

Inclusive cross-section

-4%, unaffected by choice of parameters

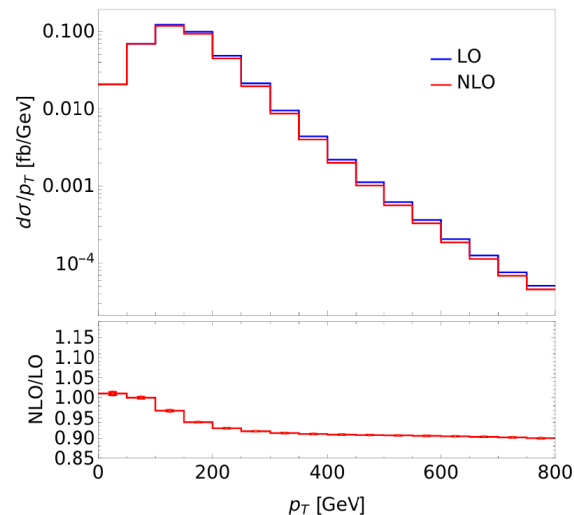
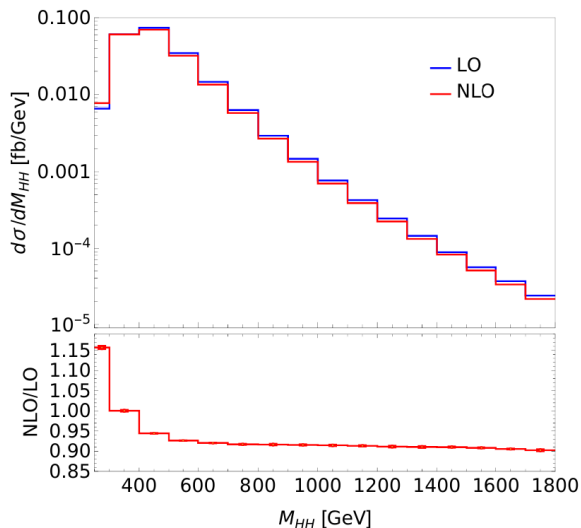
Invariant-mass

- Large positive corrections at threshold
- -10% at high M_{HH}

Negative corrections (-10%) at high p_T

μ	$M_{HH}/2$	$\sqrt{p_T^2 + m_H^2}$	m_H
LO	19.96(6)	21.11(7)	25.09(8)
NLO	19.12(6)	20.21(6)	23.94(8)
\mathcal{K} -factor	0.958(1)	0.957(1)	0.954(1)

[Bi, Huang, Huang, Ma, Yu - 2311.16963]



Conclusions

- ZH and ZZ: NNLO accuracy well established
 - *gg*-initiated channel important ingredient for N3LO
 - Top mass effects and relative scheme uncertainties are important

 - HH: NLO QCD with full top dependence
NLO EW recently achieved

 - Getting closer to theorists' goal of 1% accuracy
-

Thank you for your attention

Backup

ZH

Drell-Yan @ N3LO for ZH

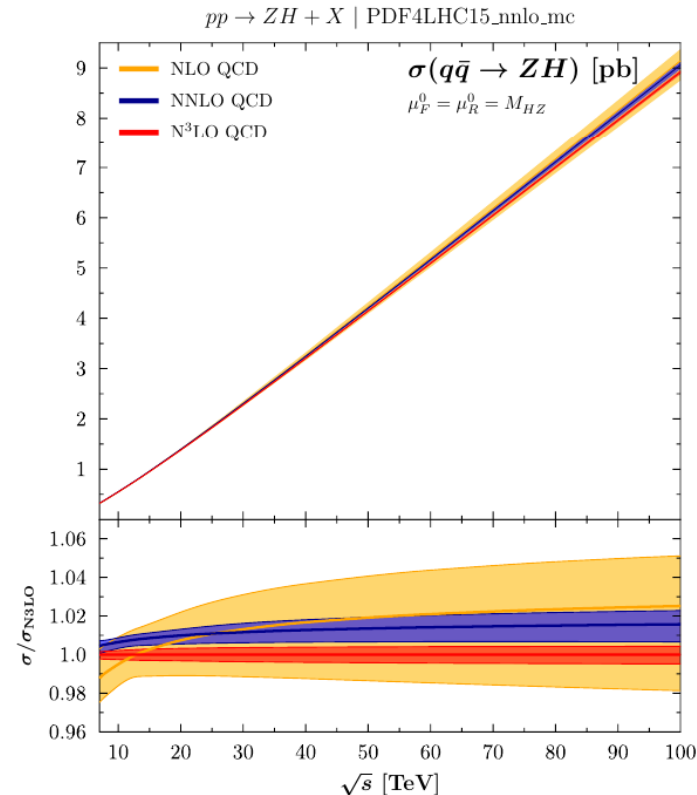
[Baglio, Duhr, Mistlberger, Szafron - 2209.06138]

$$\sqrt{s} = 13\text{TeV}$$

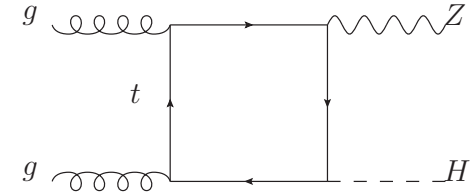
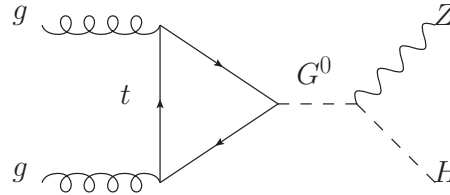
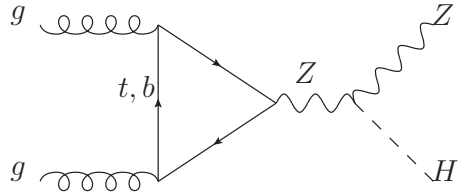
$$\mu_0 = M_{VH}$$

Process	σ^{LO} [pb]	σ^{NLO} [pb]	K^{NLO}	σ^{NNLO} [pb]	K^{NNLO}	$\sigma^{\text{N}^3\text{LO}}$ [pb]	$K^{\text{N}^3\text{LO}}$
W^+H	$0.758^{+2.43\%}_{-3.13\%}$	$0.883^{+1.38\%}_{-1.20\%}$	1.16	$0.891^{+0.28\%}_{-0.34\%}$	1.18	$0.884^{+0.27\%}_{-0.30\%}$	1.17
W^-H	$0.484^{+2.50\%}_{-3.26\%}$	$0.560^{+1.34\%}_{-1.23\%}$	1.16	$0.564^{+0.27\%}_{-0.34\%}$	1.17	$0.559^{+0.30\%}_{-0.33\%}$	1.16
ZH	$0.678^{+2.40\%}_{-3.11\%}$	$0.786^{+1.33\%}_{-1.16\%}$	1.16	$0.792^{+0.25\%}_{-0.32\%}$	1.17	$0.786^{+0.26\%}_{-0.29\%}$	1.16

- N3LO corrections larger than NNLO corrections
- No reduction in scale uncertainties at N3LO
- Fixed ($\mu_0 = m_V + m_H$) and dynamical ($\mu_0 = M_{VH}$) scale choice yield comparable results
- NNLO and N3LO scale bands do not overlap



$gg \rightarrow ZH$ @ LO

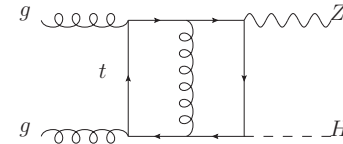
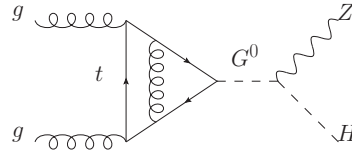
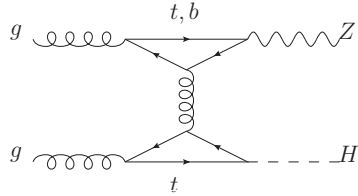


- Top-quark loops give dominant contribution [Kniehl ('90) - Dicus, Kao ('88)]
- $O(\alpha_s^2)$ correction to $\sigma(pp \rightarrow ZH)$
- NNLO suppression wrt to $q\bar{q} \rightarrow ZH$ compensated by larger gluon luminosity
- Contributes to $\sim 6\%$ of $\sigma(pp \rightarrow ZH)$ for $\sqrt{s} = 14$ TeV
- Only LO included in MC \rightarrow scale variation leads to **25%** relative uncertainties

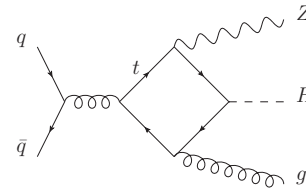
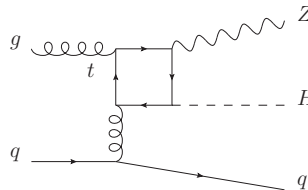
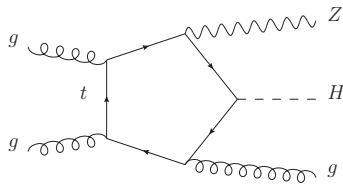
\sqrt{s} [TeV]	$\sigma_{\text{NNLO QCD} \otimes \text{NLO EW}}$ [pb]	Δ_{scale} [%]	$\Delta_{\text{PDF} \oplus \alpha_s}$ [%]
13	0.123	+24.9 -18.8	4.37
14	0.145	+24.3 -19.6	7.47
27	0.526	+25.3 -18.5	5.85

$gg \rightarrow ZH$ @ NLO in QCD - Ingredients

Virtual corrections ($2 \rightarrow 2$, two loops) - interference with LO



Real emission ($2 \rightarrow 3$, one loop) - squared amplitudes



Two-loop Massive Boxes for $gg \rightarrow ZH$

Numerical Evaluation [Chen, Heinrich, Jones, Kerner, Klappert, Schlenk - 2011.12325]

- Exact results
- Demanding in terms of computing resources and time
- Issues with flexibility

Analytic Approximations: exploit hierarchies of masses/kinematic invariants

- Reduce the number of scales in Feynman integrals
- Proliferation of integrals
- Restricted to specific phase-space regions

■ Limit $m_t \rightarrow \infty$

[Altenkamp, Dittmaier, Harlander, Rzehak, Zirke - 1211.50]

■ Large mass expansion

[Hasselhuhn, Luthe, Steinhauser - 1611.05881]

■ High-energy expansion: $m_Z^2, m_H^2 \ll m_t^2 \ll \hat{s}, \hat{t}$

[Davies, Mishima, Steinhauser - 2011.12314]

■ Small-mass expansion: $m_Z, m_H \rightarrow 0$

[Wang, Xu, Xu, Yang - 2107.08206]

■ pT expansion: $m_Z^2, m_H^2, p_T^2 \ll m_t^2, \hat{s}$

[Alasfar, Degrossi, Giardino Groeber, MV - 2103.06225]

pT Expansion - Calculation Overview

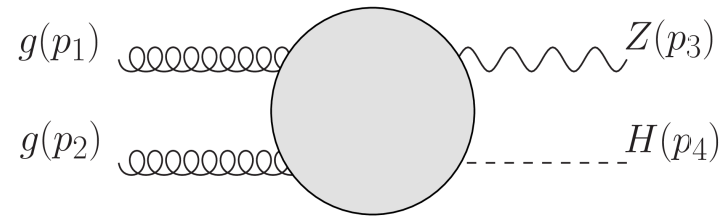
1. Generation of Feynman diagrams - $O(100 \text{ diags})$ (FeynArts [Hahn - 0012260])
2. Lorentz decomposition of the amplitude: **projectors** and **scalar form factors** (FeynCalc [Mertig et al. ('91) ; Shtabovenko et al. - 1601.01167]): contractions, Dirac traces...

$$\mathcal{A}_{\mu\nu\rho} = \sum_{i=1}^6 \mathcal{P}_{\mu\nu\rho}^{(i)} F^{(i)} \qquad F^{(i)} = \sum_{i=1}^n C^{(i)} I^{(i)}(\hat{s}, \hat{t}, m_Z^2, m_H^2, m_t^2)$$

3. Expansion of the form factors in the limit of small pT
4. Decomposition of scalar integrals using integration-by-parts (IBP) identities (LiteRed [Lee - 1310.1145])
5. Evaluation of master integrals

Steps implemented in **Mathematica** code on a **desktop machine**

pT Expansion - Details



- We assume the limit of a **forward kinematics**

$$(p_1 + p_3)^2 \rightarrow 0 \Leftrightarrow \hat{t} \rightarrow 0 \Rightarrow p_T \rightarrow 0$$

- Then Taylor-expand the form factors in the ratios

$$\frac{m_H^2}{\hat{s}}, \frac{m_Z^2}{\hat{s}}, \frac{p_T^2}{\hat{s}} \ll 1$$

$$\frac{p_T^2}{4m_t^2} \ll 1$$

Expansion at
integrand level

- Now scalar loop integrals depend on fewer scales

$$I(\hat{s}, \hat{t}, m_Z^2, m_H^2, m_t^2) \rightarrow I'(\hat{s}, \hat{t}, m_t^2)$$

- The new scalar integrals are decomposed in MIs using IBP relations

- The MIs depend on the ratio $\hat{s}/m_t^2 \Rightarrow$ **only one scale**

$$I(\hat{s}, \hat{t}, m_Z^2, m_H^2, m_t^2) \rightarrow I'(\hat{s}, \hat{t}, m_t^2) \rightarrow \text{MI}(\hat{s}/m_t^2)$$

- 52 MIs already known in the literature

SAME MIs FOR $gg \rightarrow HH$, $gg \rightarrow ZH$, $gg \rightarrow ZZ$

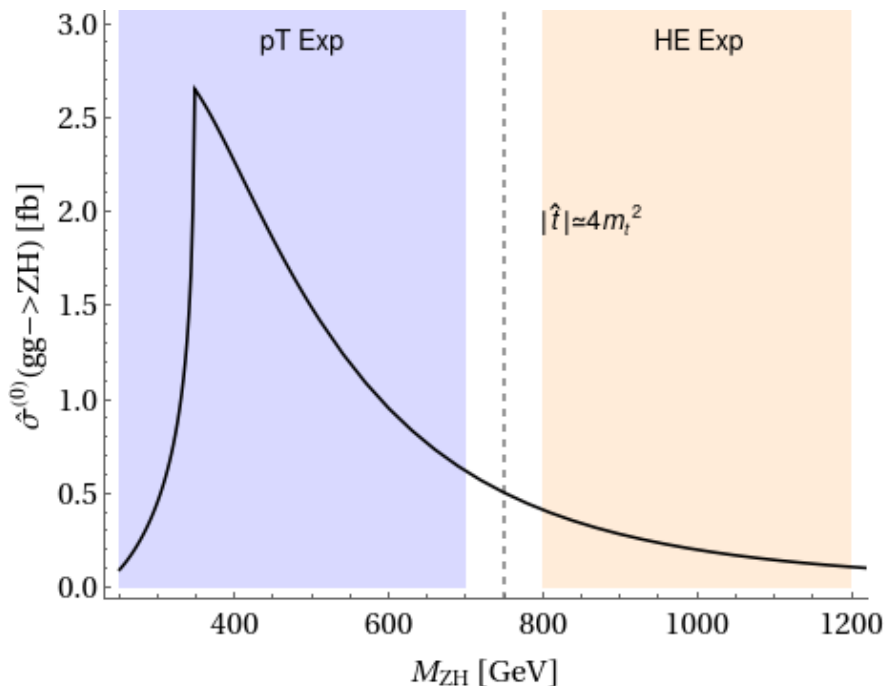
Comparing Validity Ranges

■ p_T exp: valid for

$$p_T^2 \lesssim 4m_t^2$$

or

$$|\hat{t}| \lesssim 4m_t^2$$



■ High-Energy exp:

$$|\hat{t}| \gtrsim 4m_t^2$$

[Davies, Mishima, Steinhauser - 2011.12314]

[Alasfar, Degrassi, Giardino
Groeber, MV - 2103.06225]

The two expansions can be combined

(Needed refinement using Padé approximants)

[Bellafronte, Degrassi, Giardino, Groeber, MV -2103.06225]

■ Accuracy at % level or below

⇒ OK for phenomenology

■ Evaluation time for a phase-space point
below 0.1 s ⇒ suitable for Monte Carlo

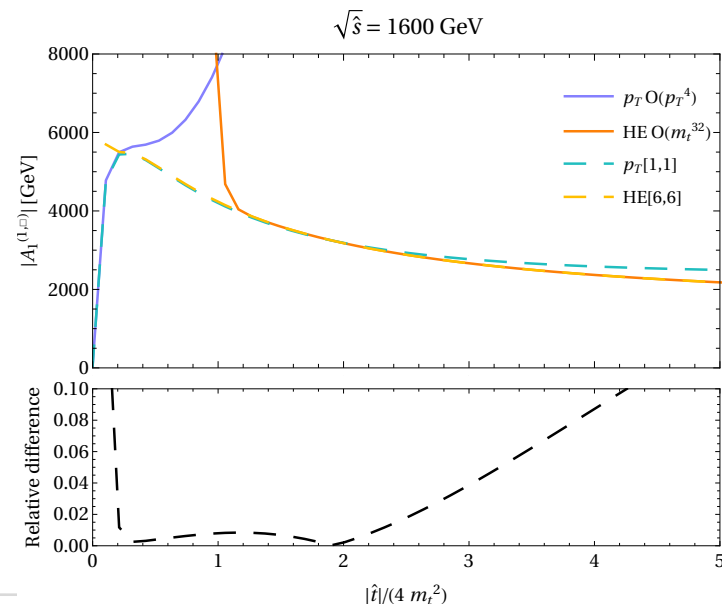
Merging pT and HE Expansions at NLO

Improve the convergence of a series expansion by matching the coefficients of the **Pade approximant** [m/n] [e.g. Fleisher, Tarasov ('94)]

$$f(x) \stackrel{x \rightarrow 0}{\simeq} c_0 + c_1 x + \dots + c_q x^q \quad f(x) \simeq [m/n](x) = \frac{a_0 + a_1 x + \dots + a_m x^m}{1 + b_1 x + \dots + b_n x^n} \quad (q = m + n)$$

[Bellafronte, Degrassi, Giardino, Gröber, MV -2103.06225]

- For each FF we merged the following results
 - pT exp improved by [1/1] Padé
 - HE exp improved by [6/6] Padé
- Padé results are stable and comparable in the region $|\hat{t}| \sim 4 m_t^2 \rightarrow$ can switch without loss of accuracy (% level or below)
- Evaluation time for a phase-space point below 0.1 s \Rightarrow suitable for Monte Carlo

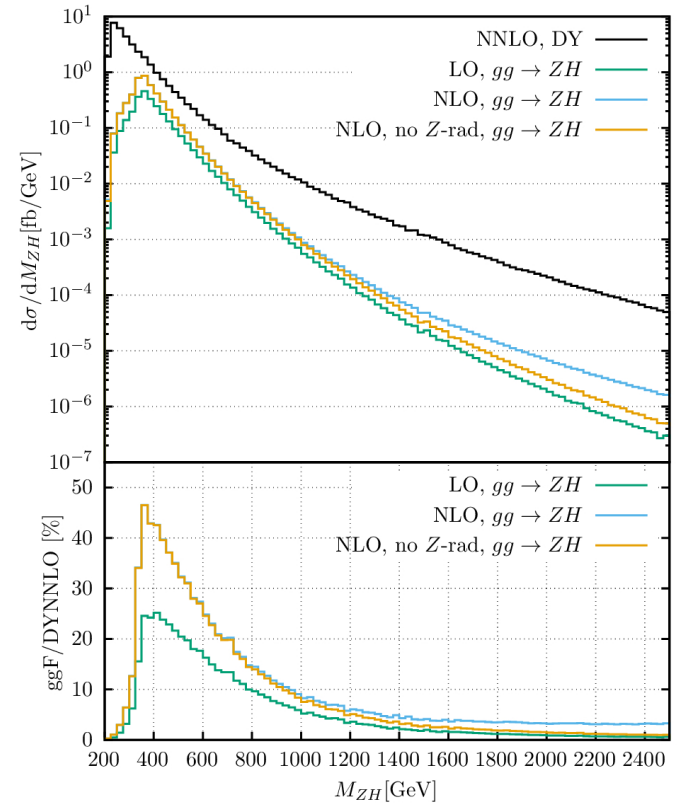


$gg \rightarrow ZH @$ vs Drell-Yan contribution

■ $gg \rightarrow ZH$ is almost 50% of DY near $M_{ZH} \sim 2 m_t$

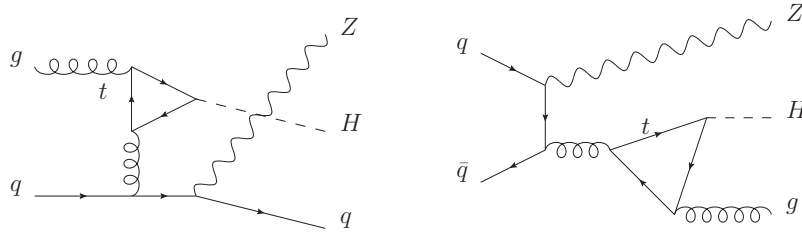
DY obtained using **vh@nnlo**

[Harlander et al - 1802.04817]



[Degrassi, Gröber, MV, Zhao - 2205.02769]

High-Energy Tails II – Z Radiation



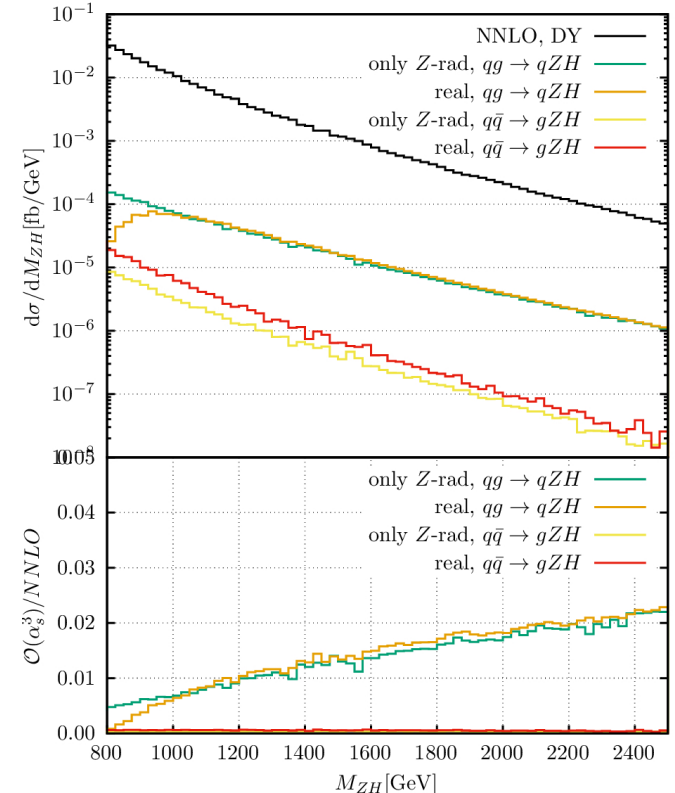
In the high-energy tail ($M_{ZH} > 1$ TeV)

■ $qg \rightarrow ZHq$ channel

- Z-radiated diagrams dominate
- Non-negligible contribution (up to 2% wrt DY)

■ $q\bar{q} \rightarrow ZHg$ channel

- Z-radiated diagrams dominate
- Negligible (PDF suppression)



$gg \rightarrow ZH$ @ NLO QCD – Inclusive Cross Section

[Wang, Xu, Xu, Yang - 2107.08206]

$\sqrt{s} = 13\text{TeV}$

$\mu_r = \mu_f$	σ_{LO}^{gg}	σ_{NLO}^{gg}	$\sigma_{pp \rightarrow ZH}^{\text{no } gg}$	$\sigma_{pp \rightarrow ZH}$	$\sigma_{\text{NLO}}^{gg, m_t \rightarrow \infty}$	$\sigma_{pp \rightarrow ZH}^{m_t \rightarrow \infty}$
$M_{ZH}/3$	73.56(7)	129.4(3)	784.0(7)	913.4(7)	133.6(6)	917.6(9)
M_{ZH}	51.03(5)	101.7(2)	781.1(7)	882.9(7)	106.0(4)	887.2(8)
$3M_{ZH}$	36.62(4)	80.4(2)	780.7(8)	861.1(8)	84.0(3)	864.8(9)

[Chen et al. - 2204.05225]

$\mu_r = \mu_f = M_{ZH}$

\sqrt{s}	LO [fb]	NLO [fb]
13 TeV	52.42 ^{+25.5%} _{-19.3%}	103.8(3) ^{+16.4%} _{-13.9%}
13.6 TeV	58.06 ^{+25.1%} _{-19.0%}	114.7(3) ^{+16.2%} _{-13.7%}
14 TeV	61.96 ^{+24.9%} _{-18.9%}	122.2(3) ^{+16.1%} _{-13.6%}

■ General agreement between the three results when same inputs are adopted

[Degrandi, Gröber, MV, Zhao - 2205.02769]

Top-mass scheme	LO [fb]	$\sigma_{\text{LO}}/\sigma_{\text{LO}}^{\text{OS}}$	NLO [fb]	$\sigma_{\text{NLO}}/\sigma_{\text{NLO}}^{\text{OS}}$	$K = \sigma_{\text{NLO}}/\sigma_{\text{LO}}$
On-Shell	64.01 ^{+27.2%} _{-20.3%}	—	118.6 ^{+16.7%} _{-14.1%}	—	1.85
$\overline{\text{MS}}, \mu_t = M_{ZH}/4$	59.40 ^{+27.1%} _{-20.2%}	0.928	113.3 ^{+17.4%} _{-14.5%}	0.955	1.91
$\overline{\text{MS}}, \mu_t = m_t^{\overline{\text{MS}}}(m_t^{\overline{\text{MS}}})$	57.95 ^{+26.9%} _{-20.1%}	0.905	111.7 ^{+17.7%} _{-14.6%}	0.942	1.93
$\overline{\text{MS}}, \mu_t = M_{ZH}/2$	54.22 ^{+26.8%} _{-20.0%}	0.847	107.9 ^{+18.4%} _{-15.0%}	0.910	1.99
$\overline{\text{MS}}, \mu_t = M_{ZH}$	49.23 ^{+26.6%} _{-19.9%}	0.769	103.3 ^{+19.6%} _{-15.6%}	0.871	2.10

$\sqrt{s} = 13\text{TeV}$
 $\mu_r = \mu_f = M_{ZH}/2$

$gg \rightarrow ZH$ @ NLO QCD – Inclusive Cross Section

[Wang, Xu, Xu, Yang - 2107.08206]

$\sqrt{s} = 13\text{TeV}$

$\mu_r = \mu_f$	σ_{LO}^{gg}	σ_{NLO}^{gg}	$\sigma_{pp \rightarrow ZH}^{\text{no } gg}$	$\sigma_{pp \rightarrow ZH}$	$\sigma_{\text{NLO}}^{gg, m_t \rightarrow \infty}$	$\sigma_{pp \rightarrow ZH}^{m_t \rightarrow \infty}$
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[Chen et al. - 2204.05225]

$\mu_r = \mu_f = M_{ZH}$

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14 TeV	61.96 ^{+24.9%} _{-18.9%}	122.2(3) ^{+16.1%} _{-13.6%}

■ NLO corrections are the same size as LO

[Degrassi, Gröber, MV, Zhao - 2205.02769]

$\sqrt{s} = 13\text{TeV}$
 $\mu_r = \mu_f = M_{ZH}/2$

Top-mass scheme	LO [fb]	$\sigma_{\text{LO}}/\sigma_{\text{LO}}^{\text{OS}}$	NLO [fb]	$\sigma_{\text{NLO}}/\sigma_{\text{NLO}}^{\text{OS}}$	$K = \sigma_{\text{NLO}}/\sigma_{\text{LO}}$
On-Shell	64.01 ^{+27.2%} _{-20.3%}	—	118.6 ^{+16.7%} _{-14.1%}	—	1.85
$\overline{\text{MS}}, \mu_t = M_{ZH}/4$	59.40 ^{+27.1%} _{-20.2%}	0.928	113.3 ^{+17.1%} _{-14.5%}	0.955	1.91
$\overline{\text{MS}}, \mu_t = m_t^{\overline{\text{MS}}}(m_t^{\overline{\text{MS}}})$	57.95 ^{+26.9%} _{-20.1%}	0.905	111.7 ^{+17.7%} _{-14.6%}	0.942	1.93
$\overline{\text{MS}}, \mu_t = M_{ZH}/2$	54.22 ^{+26.8%} _{-20.0%}	0.847	107.9 ^{+18.4%} _{-15.0%}	0.910	1.99
$\overline{\text{MS}}, \mu_t = M_{ZH}$	49.23 ^{+26.6%} _{-19.9%}	0.769	103.3 ^{+19.6%} _{-15.6%}	0.871	2.10

$gg \rightarrow ZH$ @ NLO QCD – Inclusive Cross Section

[Wang, Xu, Xu, Yang - 2107.08206]

$\sqrt{s} = 13\text{TeV}$

$\mu_r = \mu_f$	σ_{LO}^{gg}	σ_{NLO}^{gg}	$\sigma_{pp \rightarrow ZH}^{\text{no } gg}$	$\sigma_{pp \rightarrow ZH}$	$\sigma_{\text{NLO}}^{gg, m_t \rightarrow \infty}$	$\sigma_{pp \rightarrow ZH}^{m_t \rightarrow \infty}$
$M_{ZH}/3$	73.56(7)	129.4(3)	784.0(7)	913.4(7)	133.6(6)	917.6(9)
M_{ZH}	51.03(5)	101.7(2)	781.1(7)	882.9(7)	106.0(4)	887.2(8)
$3M_{ZH}$	36.62(4)	80.4(2)	780.7(8)	861.1(8)	84.0(3)	864.8(9)

[Chen et al. - 2204.05225]

$\mu_r = \mu_f = M_{ZH}$

\sqrt{s}	LO [fb]	NLO [fb]
13 TeV	52.42 ^{+25.5%} _{-19.3%}	103.8(3) ^{+16.4%} _{-13.9%}
13.6 TeV	58.06 ^{+25.1%} _{-19.0%}	114.7(3) ^{+16.2%} _{-13.7%}
14 TeV	61.96 ^{+24.9%} _{-18.9%}	122.2(3) ^{+16.1%} _{-13.6%}

■ Scale uncertainties reduced by ~30% wrt LO

[Degrassi, Gröber, MV, Zhao - 2205.02769]

Top-mass scheme	LO [fb]	$\sigma_{\text{LO}}/\sigma_{\text{LO}}^{\text{OS}}$	NLO [fb]	$\sigma_{\text{NLO}}/\sigma_{\text{NLO}}^{\text{OS}}$	$K = \sigma_{\text{NLO}}/\sigma_{\text{LO}}$
On-Shell	64.01 ^{+27.2%} _{-20.3%}	—	118.6 ^{+16.7%} _{-14.1%}	—	1.85
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$\sqrt{s} = 13\text{TeV}$

$\mu_r = \mu_f = M_{ZH}/2$

