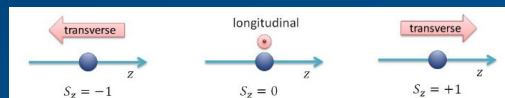


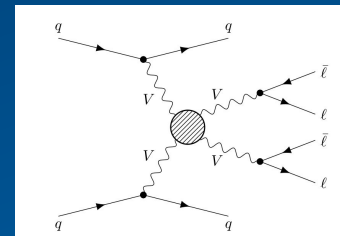
Mareen Hoppe<sup>1</sup>, Marek Schönherr<sup>2</sup>, Frank Siegert<sup>1</sup>,  
Max Stange<sup>1</sup>

<sup>1</sup>Institute of Nuclear and Particle Physics, Technische Universität Dresden

<sup>2</sup>Institute for Particle Physics Phenomenology, Durham University



Dr. Z. Zinonos: Tests of the Standard Model of Particles. <https://www.mpp.mpg.de/~zinonos/material/lecture10.pdf>



# Polarized cross sections for vector boson production with Sherpa



1st COMETA general meeting Izmir, February 29, 2024

based on the work  
published in arXiv:  
[2310.14803](https://arxiv.org/abs/2310.14803)

# Introduction

## Why vector boson polarization?

Vector boson polarization measurements can probe the electroweak sector:

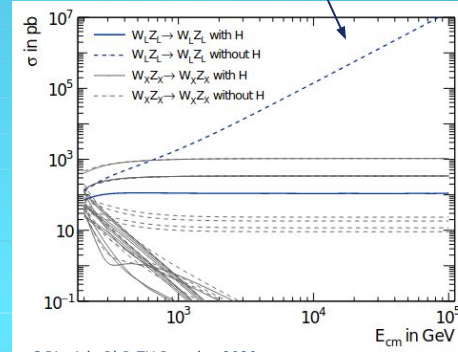
- ◆ concrete mechanism of EWSB  $\leftrightarrow$  longitudinal polarization
- ◆ EW gauge symmetry structure  $\leftrightarrow$  triple & quartic gauge couplings

## How to measure vector boson polarization?

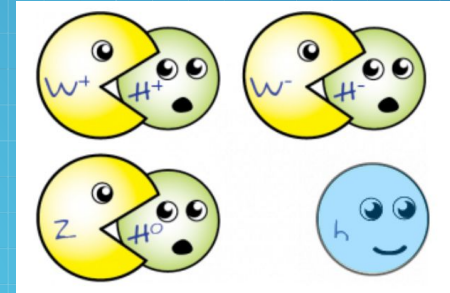
Fit with polarization templates from simulation

$$\frac{d\sigma}{dX} = f_L \frac{d\sigma_L}{dX} + f_R \frac{d\sigma_R}{dX} + f_0 \frac{d\sigma_0}{dX} \left( + f_{int.} \frac{d\sigma_{int.}}{dX} \right)$$

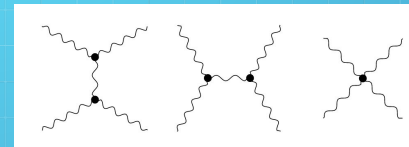
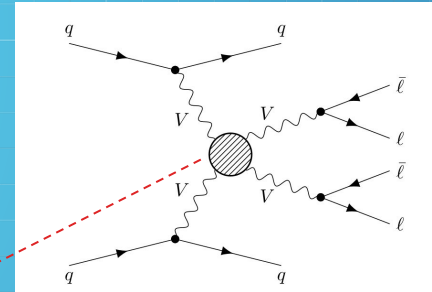
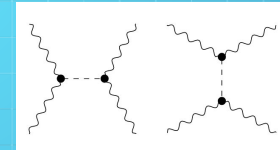
Unitarity violation



C.Bittrich, PhD TU Dresden 2020



J. Manjarrés Ramos, Talk at LHCP2022



# Generator landscape

- several private tools, see e.g. talk of [Rene Poncelet](#)
- public codes:

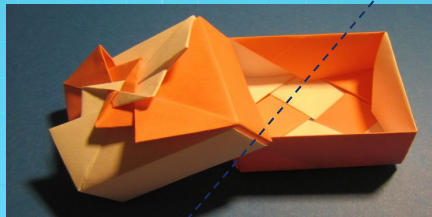
PHANTOM: a Monte Carlo event generator for six parton final states at high energy colliders.

Alessandro Ballestrero<sup>a</sup>, Aissa Bellouari<sup>c</sup>,  
Giuseppe Bevilacqua<sup>a,b</sup> Vladimir Kashkan<sup>a,b</sup> Ezio Maina<sup>a,b,1</sup>

<sup>a</sup>INFN, Sezione di Torino, Italy

<sup>b</sup>Dipartimento di Fisica Teorica, Università di Torino, Italy

<sup>c</sup>The Abdus Salam International Center for Theoretical Physics, Trieste, Italy



DPA

truncated propagator

NWA

- **PHANTOM**: 2->6 processes @ LO+PS [A. Ballestrero et al. [2008](#), [2017](#)]
- **mg5\_aMC@NLO**: arbitrary processes @ LO, PS matching, multi-jet merging [[D. Buarque Franzosi et al. 2020](#)]
- **POWHEG-BOX-RES**: diboson processes @NLO QCD+PS [[G. Pelliccioli, G. Zanderighi 2023](#)]

- **Sherpa**: arbitrary processes @nLO QCD, PS matching, multi-jet merging  
[[MH, M. Schönherr, F. Siegert 2023](#)]

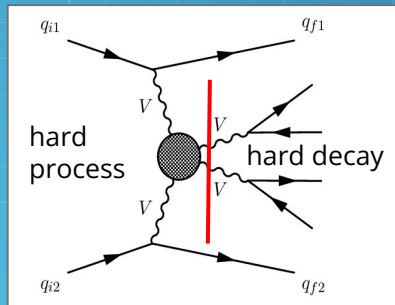
**THIS TALK**



# Calculation of polarized cross sections

- On-shell approximation: extended narrow-width approximation

$$\frac{1}{(q^2 - m_V^2)^2 + \Gamma_V^2 m_V^2} \rightarrow \frac{\pi \delta(q^2 - m_V^2)}{\Gamma_V m_V}$$



[Höche et al. 2014]

- spin correlation algorithm [Richardson 2001] to preserve spin correlations between VB production & decay
- off-shell effects via mass smearing of on-shell vector boson states

- Algorithm for polarized cross sections:

- simulation essentially unpolarized in NWA
- polarization fractions calculated on top
  - polarized amplitudes from spin correlation algorithm
  - keep whole amplitude tensor
  - multiplied by event XS
  - output as additional event weights

[MH, M. Schönherr, F. Siegert 2023]

$$|\mathcal{M}^{pol}|_{\lambda_I \dots \lambda_n; \lambda'_I \dots \lambda'_n}^2 \propto \mathcal{M}_{\lambda_I \dots \lambda_n}^P \mathcal{M}_{\lambda'_I \dots \lambda'_n}^{*P} \mathcal{M}_{\lambda_I \dots \lambda_n}^D \mathcal{M}_{\lambda'_I \dots \lambda'_n}^{*D}$$

Nominal or variation name	XS [pb]
Nominal	0.00636840738801
PolWeight_COM.W+.W+.+	0.00060858451764
PolWeight_COM.W+.W+.-	0.00078445056672
PolWeight_COM.W+.W+.0	0.00037696877407
PolWeight_COM.W.-.W+.+	0.00077432059972
PolWeight_COM.W.-.W+.-	0.00148317472445
PolWeight_COM.W.-.W+.0	0.00063413615300
PolWeight_COM.W.0.W+.+	0.00037938319197
PolWeight_COM.W.0.W+.-	0.00063363366425
PolWeight_COM.W.0.W+.0	0.00065521961374
PolWeight_COM.W.0.W+.T	0.00101500799740
PolWeight_COM.W.0.W+.t	0.00101301695622
PolWeight_COM.W.T.W+.0	0.00101857271726
PolWeight_COM.W.T.W+.T	0.00364953844553
PolWeight_COM.W.t.W+.0	0.00101110492712
PolWeight_COM.W.t.W+.t	0.00365053040852



# Features

## Fully-realistic simulations for

- + arbitrary processes with intermediate VBs
- with
- + all polarizations in one run using weights
  - + direct calculation of interference templates
  - + multiple polarization bases in one run
- ✓ laboratory frame
  - ✓ center of mass frame of arbitrary combination of initial- / final state particles
  - ✓ parton-parton frame
  - ✓ easily extendable, if necessary



# Including QCD effects matched to PS

- based on Sherpa's MC@NLO [[S. Höche et al. 2012](#)]

$$\sigma_{\text{MC@NLO}}^{\text{NLO}} = \underbrace{\int d\tilde{\Phi}_n \bar{B}_n^A(\tilde{\Phi}_n) \left[ \bar{\Delta}^A(t_0) + \int_{t_0} d\tilde{\Phi}_1 \frac{D_n^A(\tilde{\Phi}_{n+1})}{B_n(\tilde{\Phi}_n)} \bar{\Delta}^A(t) \right]}_{=:\sigma_S} + \underbrace{\int d\tilde{\Phi}_{n+1} H_n^A(\tilde{\Phi}_{n+1})}_{=:\sigma_H}$$

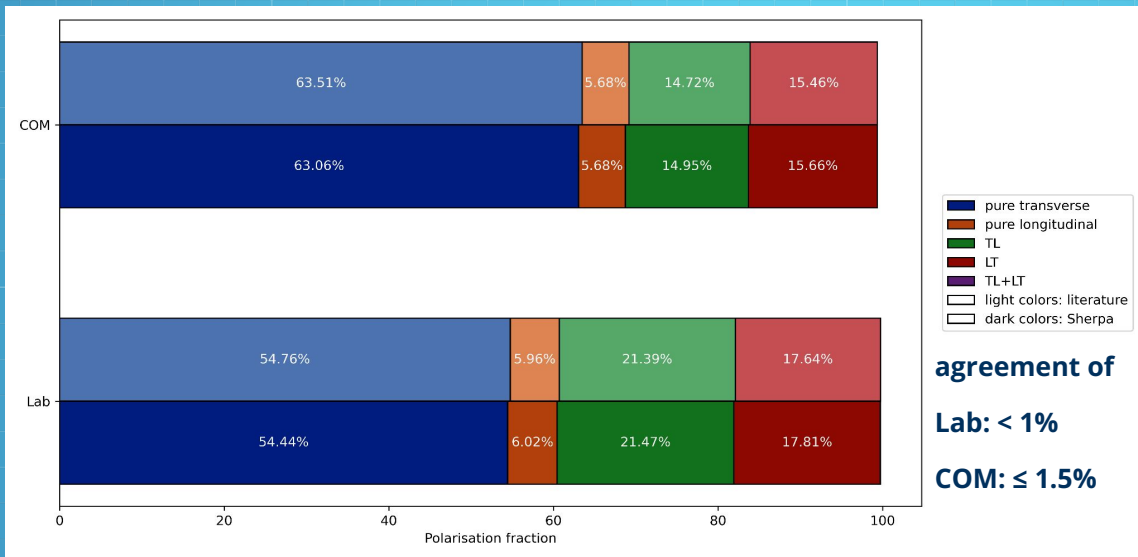
- unpolarized event XS fully NLO QCD accurate
- different amplitudes used for polarization fraction calculation in dependence of event type:
  - **H-events & resolved S-events:** real-emission amplitude
  - **Unresolved S-events:** Born amplitude
    - > virtual, ultra-soft & ultra-collinear emission effects neglected for polarization fractions



*dominant NLO effects captured through real emission MEs*

# nLO+PS vs. full fixed NLO - WZ-Diboson production

Comparison with NLO QCD fixed order calculation [A. Denner & G. Pelliccioli 2021]

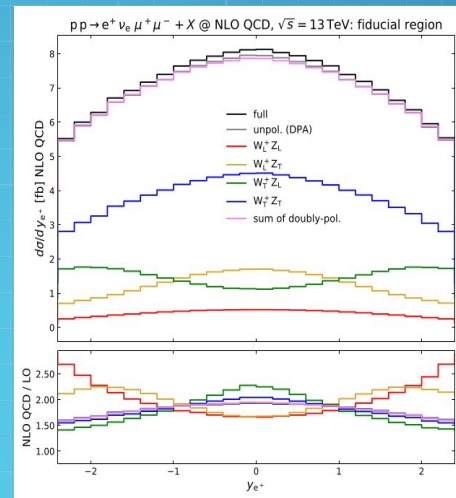


agreement of

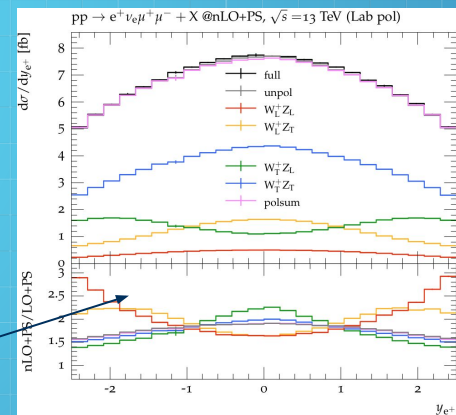
Lab: < 1%

COM: ≤ 1.5%

➤ can reproduce even non-trivial NLO effects!



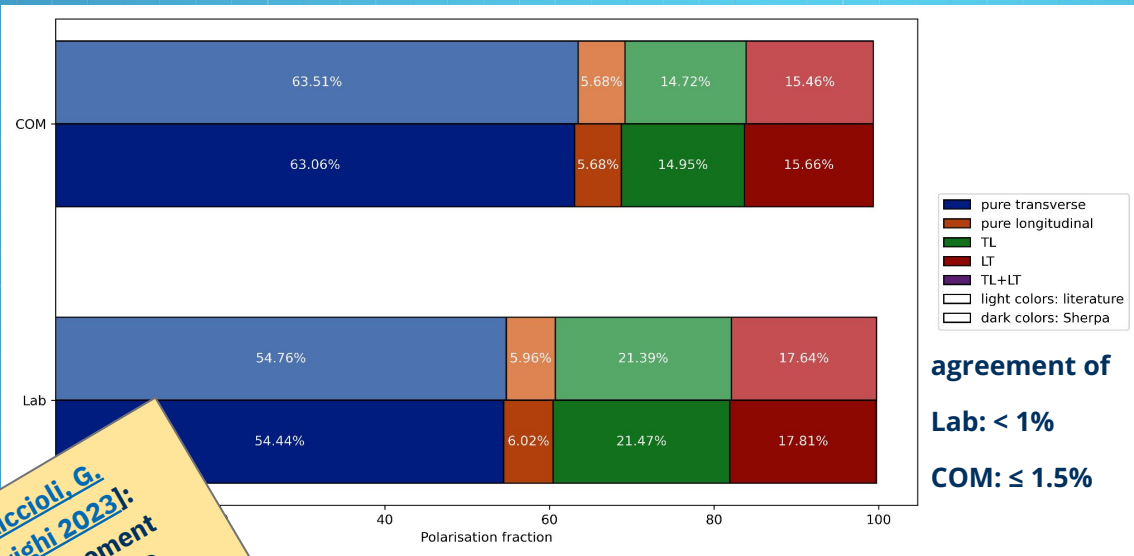
Literature



Sherpa

# nLO+PS vs. full fixed NLO - WZ-Diboson production

Comparison with NLO QCD fixed order calculation [A. Denner & G. Pelliccioli 2021]

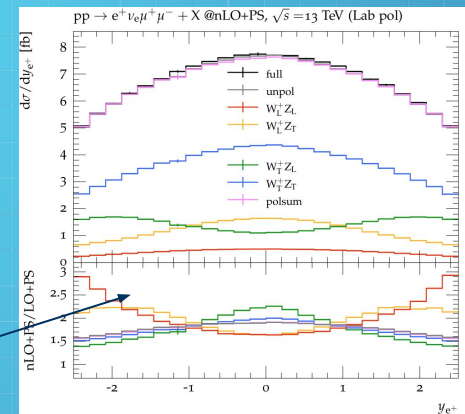
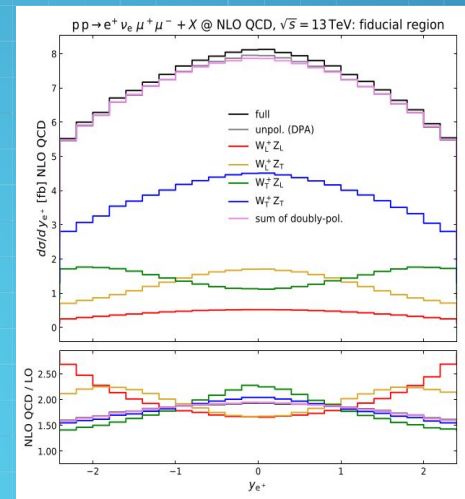


agreement of

Lab: < 1%

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Literature

Sherpa

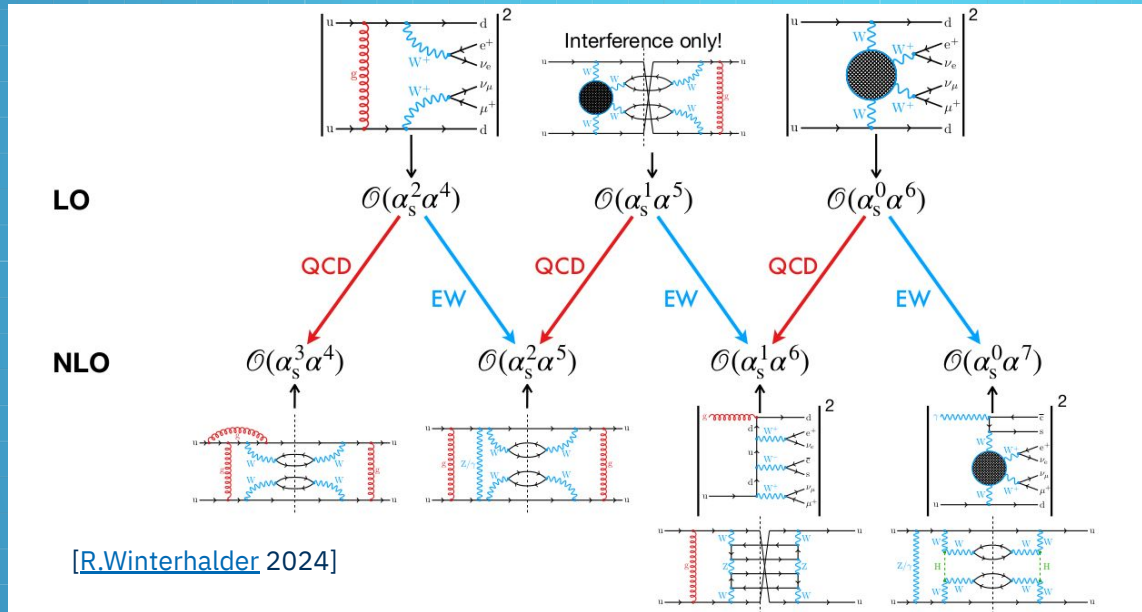
[G. Pelliccioli, G. Zanderighi 2023]: similar agreement compared to NLO QCD+PS with POWHEG+PYTHIA





# Higher order polarized cross sections for VBS processes

- Calculation of full NLO effects to VBS processes very challenging:



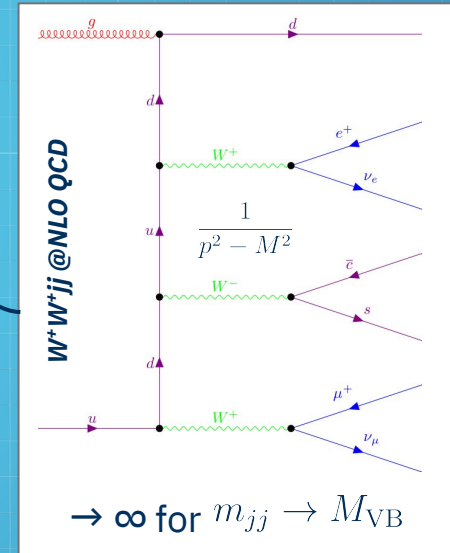
- currently only available at fixed order, unpolarized, e.g. [B. Biedermann et al. 2017, S. Dittmaier et al. 2023]
- polarized cross sections only @ LO [A. Ballestrero et al. 2020, 2017, 2019]

With the new framework  
 first study in  $W^+W^{+}jj$  :  
 *$\mathcal{LO}+1j$  merged calculation*

# First higher order polarized cross sections for $W^+W^+jj$

Preliminary results

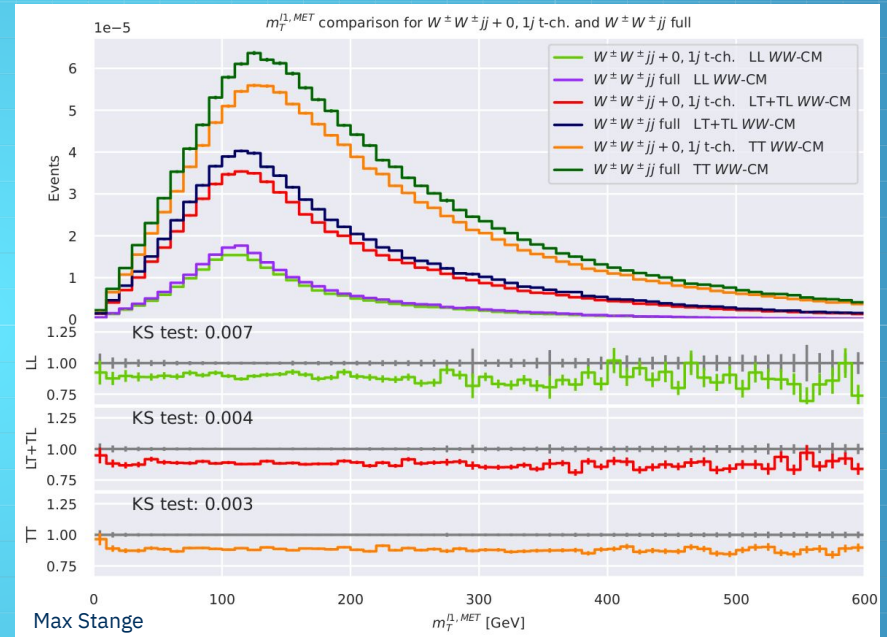
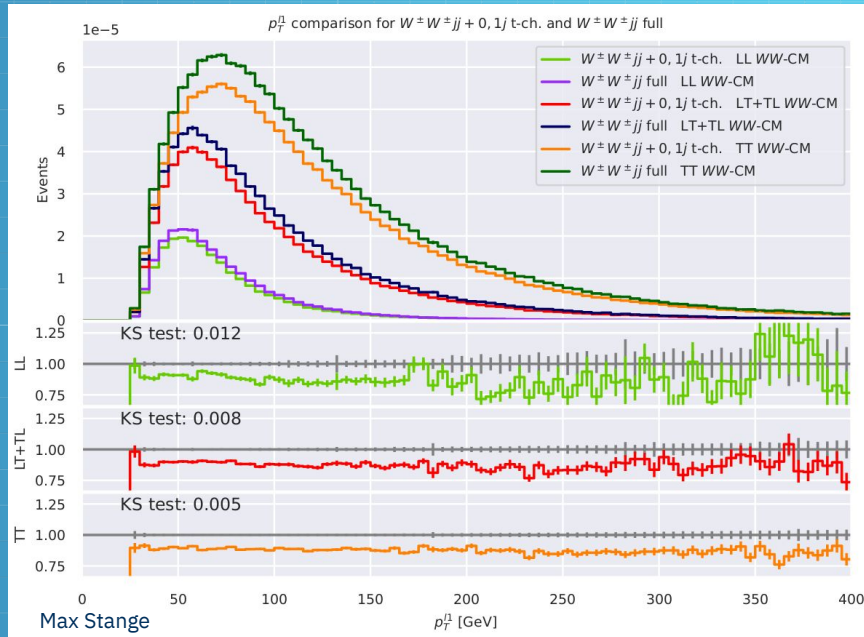
	unpol [fb]	LL [fb]	Fraction [%]	LT+TL [fb]	Fraction [%]	TT [fb]	Fraction [%]
$\mathcal{L}O+P8$	2.459	0.237	9.63	0.765	31.12	1.497	60.88
$\mathcal{L}O+1j$ VBS approx.	2.302	0.224	9.73	0.716	31.10	1.399	60.77



Only sub-percent level effect of QCD corrections on integrated polarization fractions ...

# First higher order polarized cross sections for $W^+W^+jj$

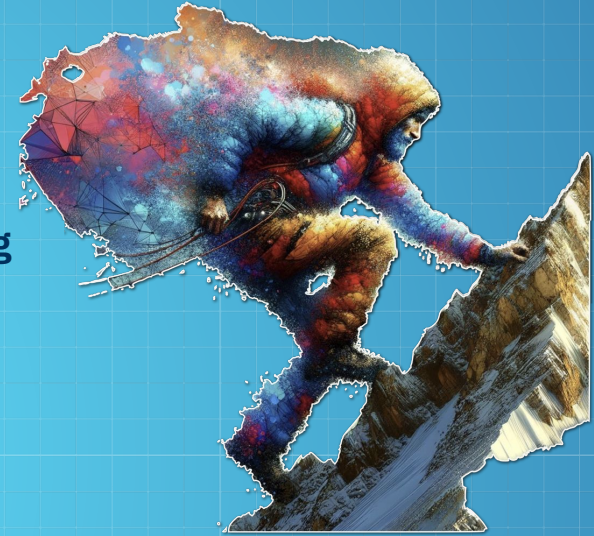
Preliminary results



... differential polarization fractions!

# Summary and Outlook

- ✓ first implementation of polarized cross sections in [Sherpa](#)
- ✓ higher order QCD corrections @nLO + PS & via multijet-merging
- first fully-realistic VB polarization samples including higher order QCD effects for VBS processes



<https://www.zeppelinschule-speyer.de/ausblick-auf-die-zeit-nach-pfingsten/>

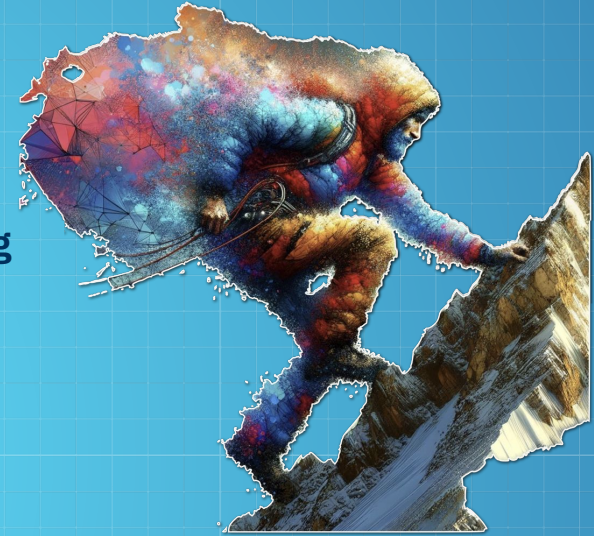
## What comes next ...

- ❑ Extension to NLO QCD and approximate NLO EW
- ❑ Applications in phenomenological analyses: BSM studies (UFO), full NLO effects for VBS processes



# Summary and Outlook

- ✓ first implementation of polarized cross sections in [Sherpa](#)
- ✓ higher order QCD corrections @nLO + PS & via multijet-merging
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<https://www.zeppelinschule-speyer.de/ausblick-auf-di-e-zeit-nach-pfingsten/>

## What comes next ...

- ❑ Extension to NLO QCD and approximate NLO EW
- ❑ Applications in phenomenological analyses: BSM studies (UFO), full NLO effects for VBS processes

*Thank you for your  
attention!  
Questions?*



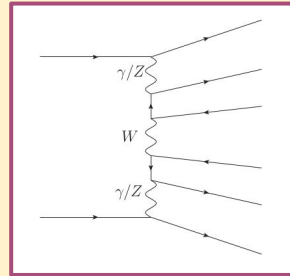
# Backup



# Polarized cross sections from fully exclusive events

## Polarization templates

- Interference between different polarizations
- Polarization fractions are frame dependent
- Separation of polarization not for all diagrams
- > approximations necessary



non-resonant diagrams

$$\frac{d\sigma}{dX} = f_L \frac{d\sigma_L}{dX} + f_R \frac{d\sigma_R}{dX} + f_0 \frac{d\sigma_0}{dX} \left( + f_{int.} \frac{d\sigma_{int.}}{dX} \right)$$

truncated propagator

## LO+PS Tools

- **PHANTOM**: 2->6 processes @ LO+PS [A. Ballestrero et al. 2008, 2017]
- **MG5\_aMC@NLO**: arbitrary processes @ LO, PS matching, multi-jet merging [D. Buarque Franzosi et al. 2020]



## Higher orders

- **POWHEG-BOX-RES**: diboson processes @NLO QCD+PS [G. Pelliccioli, G. Zanderighi 2023]
- **SHERPA**: arbitrary processes @nLO QCD, PS matching, multi-jet merging [MH, M. Schönherr, F. Siegert 2023]

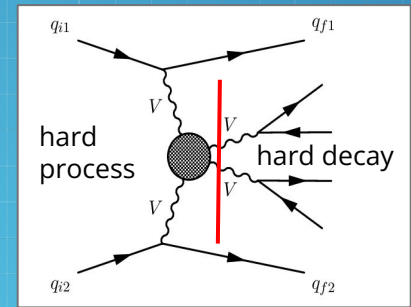


DPA

NWA

# On-shell Approximation in Sherpa [S. Höche et al. 2014]

- **extended narrow-width approximation:**
  - spin correlations preserved
  - off-shell effects via mass smearing of on-shell vector boson states



01

Simulation of unpolarized production process

$$|\mathcal{M}^P|^2 = \sum_{\lambda_1 \dots \lambda_n} \mathcal{M}_{\lambda_1 \dots \lambda_n}^P \mathcal{M}_{\lambda_1 \dots \lambda_n}^{*P}$$

02

Recalculation of the production amplitude tensor

$$|\mathcal{M}^P|^2_{\lambda_1 \dots \lambda_n; \lambda'_1 \dots \lambda'_n} = \mathcal{M}_{\lambda_1 \dots \lambda_n}^P \mathcal{M}_{\lambda'_1 \dots \lambda'_n}^{*P}$$

03

Spin correlation algorithm generates decay chain, decay matrices [Richardson 2001]

$$\mathcal{D}_{\lambda_A \lambda'_A} = \frac{1}{\mathcal{N}_D} \mathcal{M}_{\lambda_A}^D \mathcal{M}_{\lambda'_A}^{D*}$$

04

Mass Smearing of the intermediate vector bosons according to Breit-Wigner distribution

$$\text{BW} \sim \frac{1}{(p^2 - M^2)^2 + \Gamma^2 M^2}$$

05

Reweighting of production cross section with decay branching ratio

$$|\mathcal{M}^{\text{NWA}}|^2 = |\mathcal{M}^P|^2 \times \text{BR}$$

# Spin-Correlation Algorithm

- here only for VB decaying into stable leptons

hard process final state particles & production (2→n) matrix element tensor  $|\mathcal{M}^{\mathcal{P}}|_{\lambda_1 \dots \lambda_n \lambda'_1 \dots \lambda'_n}^2$

choose one outgoing particle A  
randomly

$$\text{Spin density matrix } \rho_{\lambda_j \lambda'_j}(A) = \frac{1}{N_\rho} \mathcal{M}_{\kappa_1 \kappa_2; \lambda_1 \dots \lambda_j \dots \lambda_n}^{\mathcal{P}} \mathcal{M}_{\kappa_1 \kappa_2; \lambda'_1 \dots \lambda'_j \dots \lambda'_n}^{\mathcal{P}*} \prod_{i \neq j} \mathcal{D}_{\lambda_i \lambda'_i}^i$$

with  $\mathcal{D}_{\lambda_i \lambda'_i}^i = \frac{1}{n_{hel}} \delta_{\lambda_i \lambda'_i}$  if particle not chosen yet

choose decay channel of A according to branching ratios

Generate momenta of A's decay products according to  $\rho_{\lambda_A \lambda'_A} \mathcal{M}_{\lambda_A; \lambda_1 \dots \lambda_n}^{\mathcal{D}} \mathcal{M}_{\lambda'_A; \lambda_1 \dots \lambda_n}^{\mathcal{D}*}$

all decay products stable

Calculate A's decay matrix  $\mathcal{D}_{\lambda_A \lambda'_A} = \frac{1}{\mathcal{N}_D} \mathcal{M}_{\lambda_A; \lambda_1 \dots \lambda_n}^{\mathcal{D}} \mathcal{M}_{\lambda'_A; \lambda_1 \dots \lambda_n}^{\mathcal{D}*}$

# Key part of the polarization framework: Transformation of the matrix elements

- Two ways to change polarization definition in matrix elements:
  - a priori: change polarization definition directly in matrix element generator
  - **a posteriori: transformation of calculated production tensor, decay matrices**
- change of basis = basis transformation of polarization vectors

$$\begin{pmatrix} \tilde{\epsilon}_+^0 & \tilde{\epsilon}_-^0 & \tilde{\epsilon}_0^0 \\ \tilde{\epsilon}_+^1 & \tilde{\epsilon}_-^1 & \tilde{\epsilon}_0^1 \\ \tilde{\epsilon}_+^2 & \tilde{\epsilon}_-^2 & \tilde{\epsilon}_0^2 \\ \tilde{\epsilon}_+^3 & \tilde{\epsilon}_-^3 & \tilde{\epsilon}_0^3 \end{pmatrix} = \begin{pmatrix} \epsilon_+^0 & \epsilon_-^0 & \epsilon_0^0 \\ \epsilon_+^1 & \epsilon_-^1 & \epsilon_0^1 \\ \epsilon_+^2 & \epsilon_-^2 & \epsilon_0^2 \\ \epsilon_+^3 & \epsilon_-^3 & \epsilon_0^3 \end{pmatrix} \begin{pmatrix} a_{++} & a_{-+} & a_{0+} \\ a_{+-} & a_{--} & a_{0-} \\ a_{+0} & a_{-0} & a_{00} \end{pmatrix}$$

matrix of new polarization vectors

matrix of default polarization vectors

transformation coefficients

- Transformation of matrix elements

$$|\mathcal{M}|_{\lambda_1 \lambda'_1 \dots \lambda_n \lambda'_n}^2 = \sum_{\kappa_1 \kappa'_1 \dots \kappa_n \kappa'_n} a_{\lambda_1 \kappa_1}^{\text{part1}} a_{\lambda'_1 \kappa'_1}^{\text{part1}*} \dots a_{\lambda_n \kappa_n}^{\text{partn}} a_{\lambda'_n \kappa'_n}^{\text{partn}*} |\mathcal{M}|_{\kappa_1 \kappa'_1 \dots \kappa_n \kappa'_n}^2$$

matrix element in new basis

linear combination coefficients

matrix element in default basis



# Matching to parton shower: $\delta$ -MC@NLO

[S. Höche et. al, 2012]

[S. Frixione & B.R. Webber, 2002]

**Idea:** Double counting eliminated by subtraction of the shower contribution at fixed order from contributions of the hard process at higher orders

- Separation of real correction into IR-singular and IR-regular term:

$$R_n = D_n^A + H_n^A$$

- NLO QCD cross section:

$$\sigma^{\text{NLO}} = \int d\tilde{\Phi}_n \bar{B}^A(\tilde{\Phi}_n) + \int d\tilde{\Phi}_{n+1} H^A(\tilde{\Phi}_{n+1})$$

$$\bar{B}^A(\tilde{\Phi}_n) = \underbrace{B_n(\tilde{\Phi}_n)}_{\text{Born}} + \underbrace{\tilde{V}_n(\tilde{\Phi}_n)}_{\text{virtual connection}} + \underbrace{I_n^S(\tilde{\Phi}_n)}_{\text{integrated subtraction term}} + \int d\tilde{\Phi}_1 \left[ \underbrace{D_n^A(\tilde{\Phi}_{n+1})}_{\text{subtraction terms}} - D_n^S(\tilde{\Phi}_{n+1}) \right]$$

matched NLO QCD cross section up to first emission:

$$\sigma_{\text{MC@NLO}}^{\text{NLO}} = \underbrace{\int d\tilde{\Phi}_n \bar{B}_n^A(\tilde{\Phi}_n) \left[ \bar{\Delta}^A(t_0) + \int_{t_0} d\tilde{\Phi}_1 \frac{D_n^A(\tilde{\Phi}_{n+1})}{B_n(\tilde{\Phi}_n)} \bar{\Delta}^A(t) \right]}_{=:\sigma_S} + \underbrace{\int d\tilde{\Phi}_{n+1} H_n^A(\tilde{\Phi}_{n+1})}_{=:\sigma_H}$$

applied to expectation value of an IR-safe observable  $O$

particles  $a_i$  with flavours  $f_i$  in the Born process with momenta  $p_i$

$$\langle O \rangle^{(\text{NLOMC})} = \sum_{\{\vec{f}\}} \int d\Phi_B(\{\vec{p}\}) \bar{B}^{(A)}(\{\vec{a}\}) \left[ \underbrace{\bar{\Delta}^{(A)}(t_0; \{\vec{a}\})}_{\text{unresolved}} O(\{\vec{p}\}) \right. \\ \left. + \sum_{\{\vec{r}, \vec{k}\}} \sum_{F_i=q,g} \int d\Phi_{R/B}^{ij,k} \Theta(t(\Phi_{R/B}^{ij,k}) - t_0) O(r_{\vec{r}, \vec{k}}(\{\vec{p}\})) \right. \\ \left. \times \underbrace{\frac{1}{S_{ij}} \frac{S(r_{\vec{r}, \vec{k}}(F_i; \{\vec{f}\}))}{S(\{\vec{f}\})} \frac{D_{ij,k}^{(A)}(r_{\vec{r}, \vec{k}}(F_i; \Phi_{R/B}^{ij,k}; \{\vec{a}\}))}{B(\{\vec{a}\})}}_{\text{resolved, singular}} \bar{\Delta}^{(A)}(t; \{\vec{a}\}) \right] \\ \left. + \sum_{\{\vec{F}\}} \int d\Phi_R(\{\vec{p}\}) \left[ \underbrace{R(\{\vec{A}\}) - \sum_{ij,k} D_{ij,k}^{(A)}(\{\vec{A}\})}_{\text{resolved, non-singular}} \right] O(\{\vec{p}\}) \right]$$

$\delta$ -event

$H$ -event

particles  $a_i$  with flavours  $F_i$  in the real process with momenta  $p_i$

Sudakov factor

# Polarized cross sections matched to PS - MC@NLO

[S. Höche et al. 2012]

**Idea MC@NLO:** Double counting eliminated by subtraction of the shower contribution at fixed order from contributions of the hard process at higher orders

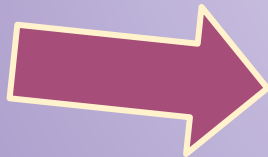
$$\sigma_{\text{MC@NLO}}^{\text{NLO}} = \underbrace{\int d\tilde{\Phi}_n \bar{B}_n^A(\tilde{\Phi}_n) \left[ \bar{\Delta}^A(t_0) + \int_{t_0} d\tilde{\Phi}_1 \frac{D_n^A(\tilde{\Phi}_{n+1})}{B_n(\tilde{\Phi}_n)} \bar{\Delta}^A(t) \right]}_{=:\sigma_{\text{S}}} + \underbrace{\int d\tilde{\Phi}_{n+1} H_n^A(\tilde{\Phi}_{n+1})}_{=:\sigma_{\text{H}}}$$

## Resolved S-events

- universal-soft-collinear radiation pattern in PS approximation above PS IR cutoff

- construction from complete *real emission* amplitude

virtual, ultra-soft & ultra-collinear emission effects neglected for polarisation fractions



## Recalculation of the amplitude tensor in dependency of event type

### Unresolved S-events

- emissions below PS IR cutoff
- virtual corrections

- construction from *Born* amplitude

### H-events

- hard, well-separated emissions beyond PS starting scale
- process-specific corrections to universal soft-collinear emission pattern

- construction from complete *real emission* amplitude

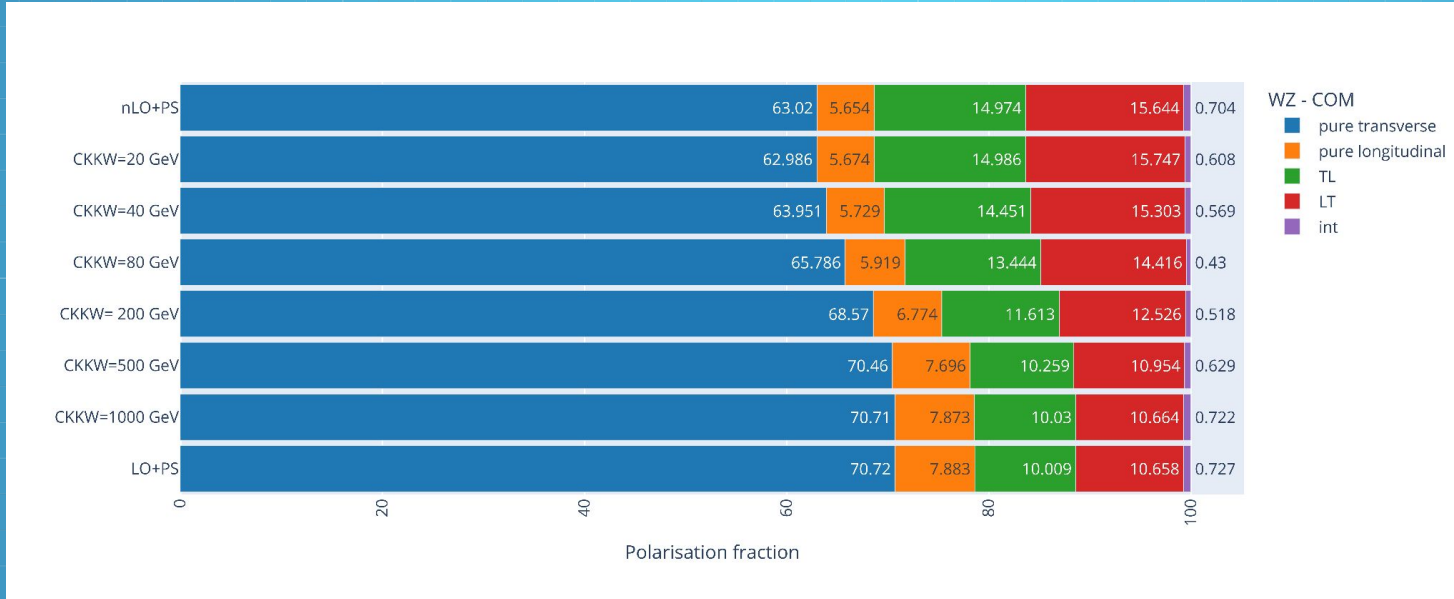
# nLO+PS (Sherpa) vs. full NLO+PS (POWHEG+PYTHIA)

Literature: [\[G. Pelliccioli, G. Zanderighi 2023\]](#)

## ZZ-Production

	$\sigma_{\text{NLO+PS}}$ [fb]	Fraction [%]	$\sigma_{\text{nLO+PS}}$ [fb]	Fraction [%]	$\frac{\text{nLOPS}}{\text{NLOPS}} - 1$ [%]	
					XS	Frac
<i>unpol</i>	14.02(1)	100	14.017(17)	100	-0.0	
<i>LL</i>	0.819(1)	5.84	0.8404(12)	5.996(11)	+2.6	+2.7
<i>LT + TL</i>	3.565(3)	25.43	3.6177(39)	25.81(4)	+1.5	+1.5
<i>TT</i>	9.47(1)	67.52	9.370(14)	66.85(13)	-1.1	-1.0
<i>int</i>	0.171	1.28	0.1886(24)	1.345(17)	+10.3	+5.1

# WZ - LO + 1jet merged with PS compared with (n)LO+PS - COM



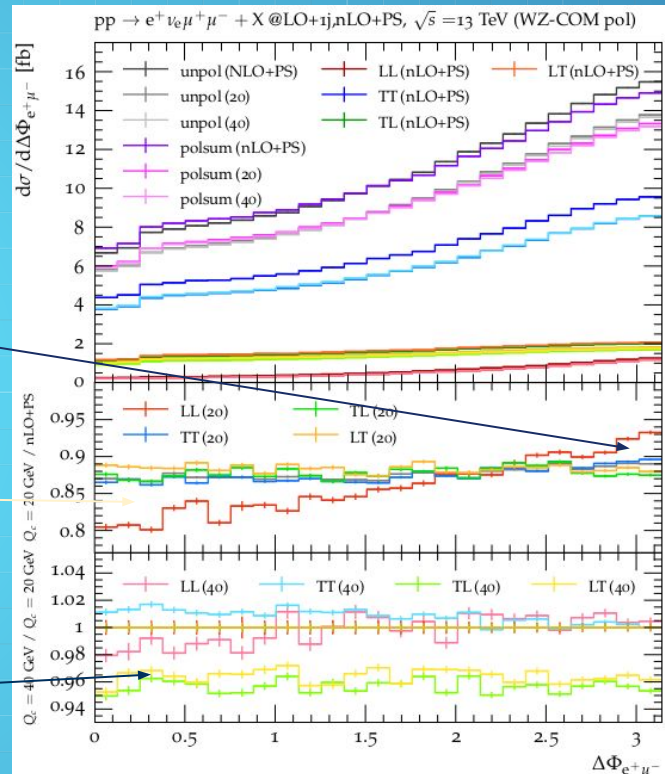
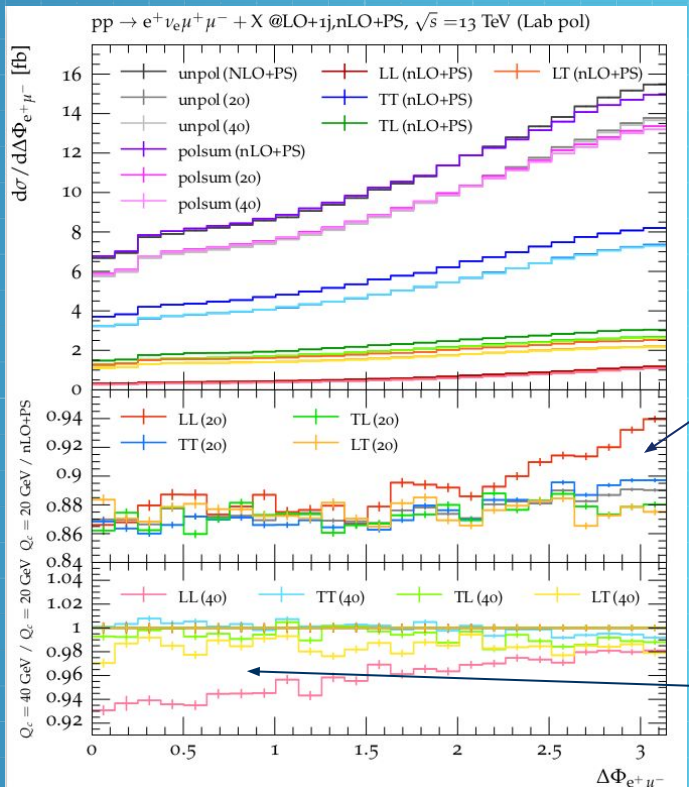
**small merging scale  
(20-40 GeV):**

- good approximation to nLO+PS result  
- no large dependency on concrete scale value

**larger merging scale:**

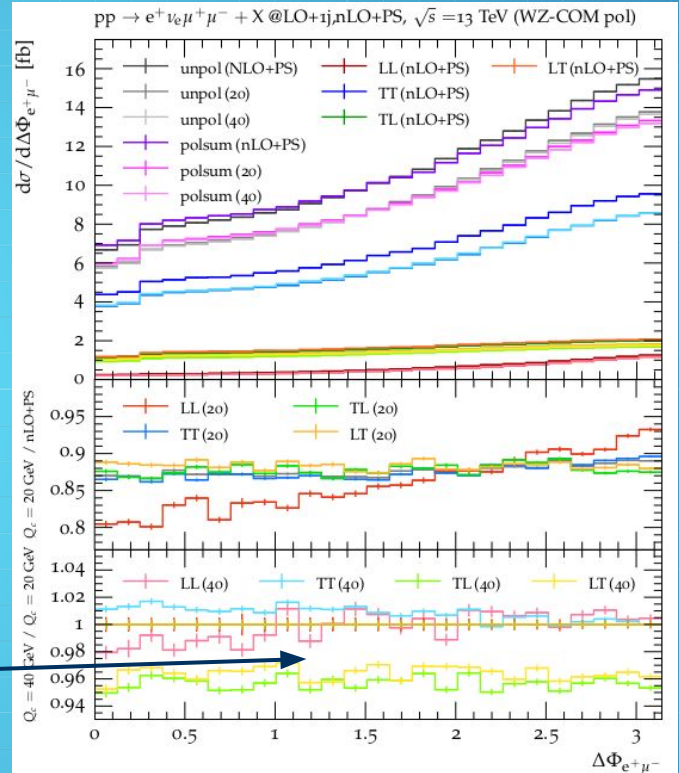
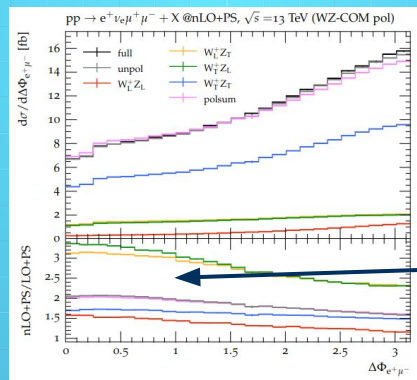
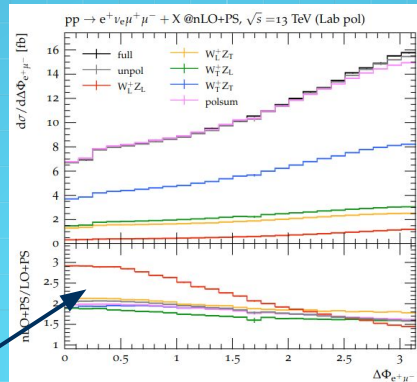
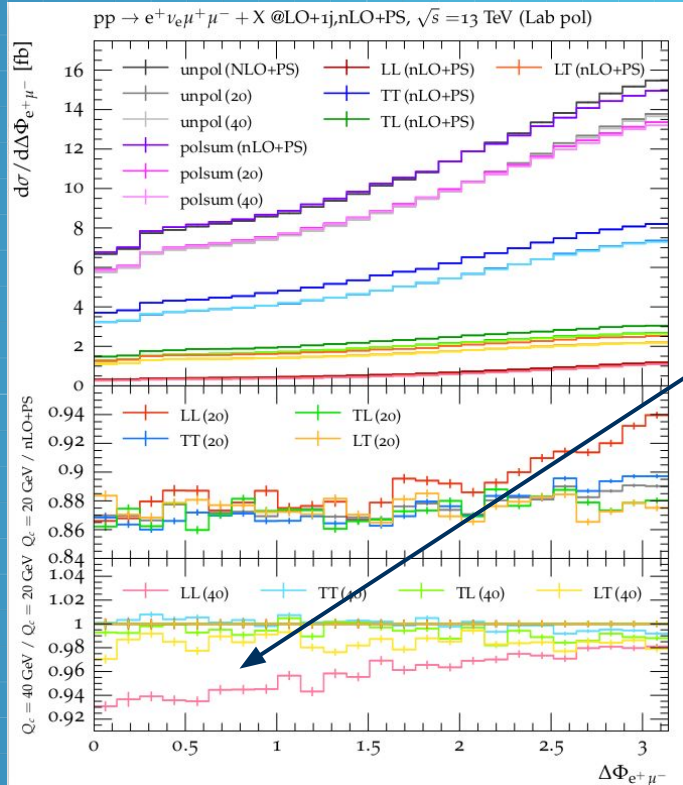
- transition to LO behavior

# $WZ - LO + 1jet$ merged with PS compared with $(n)LO+PS$

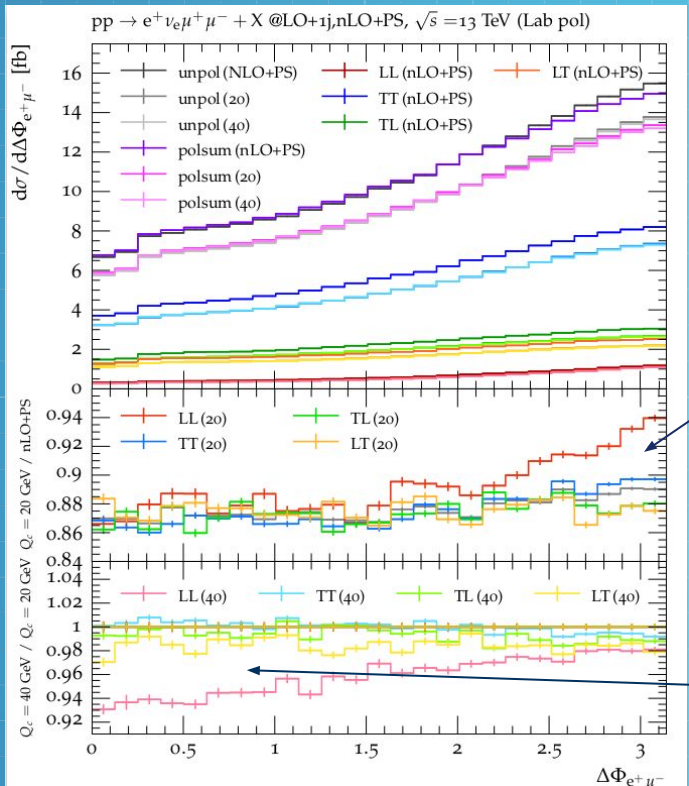




# WZ - LO + 1jet merged with PS compared with (n)LO+PS



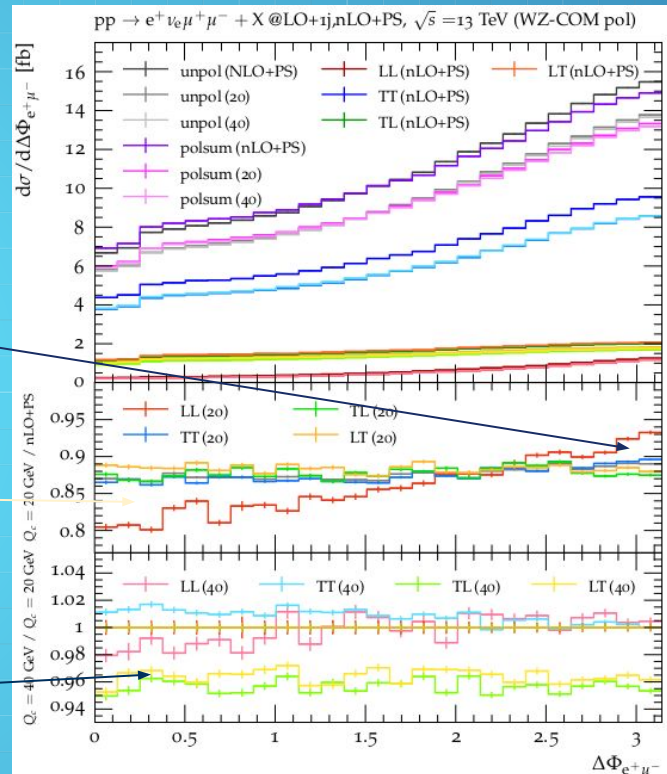
# WZ - LO + 1jet merged with PS compared with (n)LO+PS



impact of very soft emissions below 20 GeV different on LL component

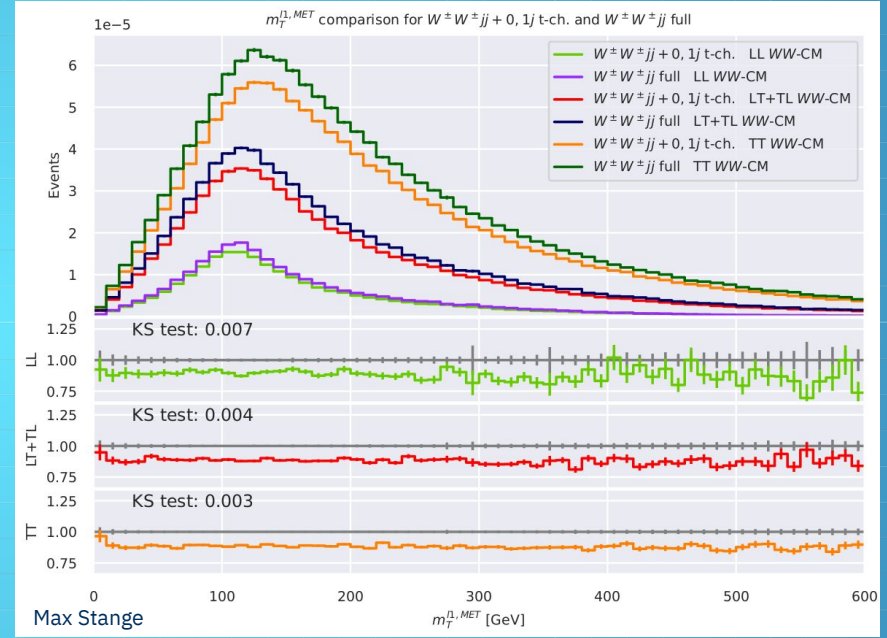
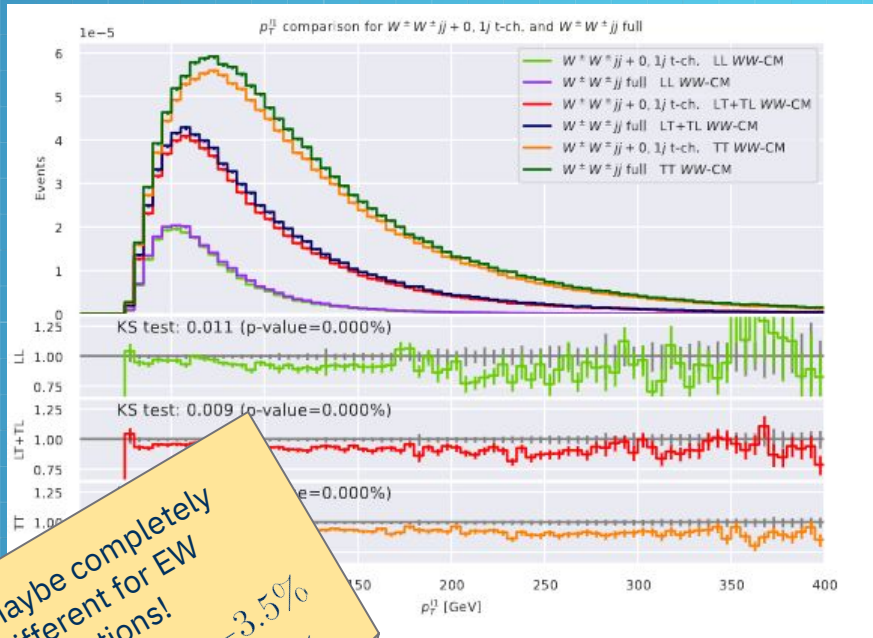
frame dependent effect

effect from emissions between 20 & 40 GeV extremely frame dependent



# First higher order polarized cross sections for $W^+W^+jj$

Preliminary results



Maybe completely different for EW corrections!

$\mathcal{O}(\alpha^6 \alpha_s) = -3.5\%$

$\mathcal{O}(\alpha^7) = -13.2\%$

B. Biedermann et al. 2017

... differential polarization fractions!