

Galactic Magnetic Field: some Puzzles

Michael Kachelrieß

NTNU, Trondheim

with C.Becker, G.Giacinti, D.Semikoz, ...

Outline:

- Introduction: CRs as probe for the GMF
- GMF models
 - ▶ input: RM, I_{syn} , and polarisation U, Q
 - ▶ status of models
- CR escape in GMF models
 - ▶ connection to diffusion picture
 - ▶ isotropic vs. anisotropic diffusion
- C-BASS and polarisation data
- TeV halos
- Extended halo model
- Summary

[Giacinti, MK, Semikoz '12ff]

[Becker]

CRs as probe for turbulent magnetic fields:

- Galactic **magnetic** field: **regular** + **turbulent** component
 turbulent: **fluctuations** on scales $l_{\max} \sim (10 - 150) \text{ pc}$ to $l_{\min} \ll l_{\max}$
- relevant scales:
 - ▶ fast modes with $kR_L \gg 1$: irrelevant
 - ▶ slow modes with $kR_L \ll 1$: act locally as uniform field B_0
 - ▶ CRs scatter on modes with $kR_L \sim 1$
- Larmor radius

$$R_L = \frac{cp_{\perp}}{ZeB} = \frac{\mathcal{R}}{B} \simeq 1.08 \text{ pc} \frac{\mathcal{R}}{\text{PV}} \frac{\mu\text{G}}{B}$$

⇒ CRs probe modes between $(10^{-7} - 1) \text{ pc}$

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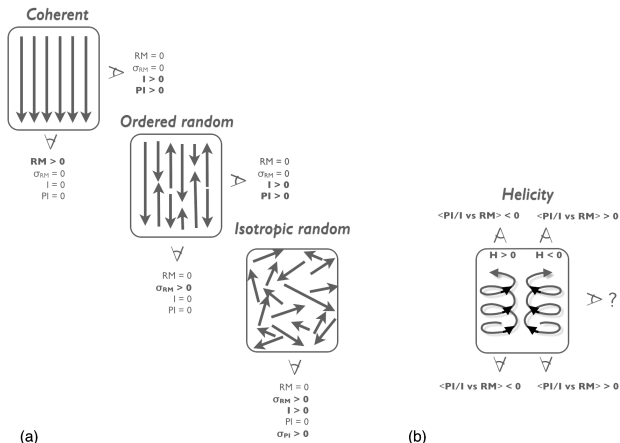
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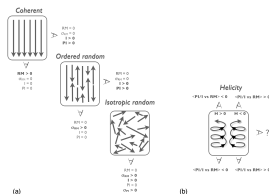
GMF observables

- unpolarised **synchrotron intensity** $I \sim \int_{\text{L.o.S.}} ds n_{\text{cre}}(\mathbf{x}, E) B_{\perp}^{\alpha}(\mathbf{x})$
- **polarised** synchrotron intensity P or Q, U
- **rotation measure** $\text{RM} \sim \int_{\text{L.o.S.}} ds n_e(\mathbf{x}) B_{\parallel}(\mathbf{x})$



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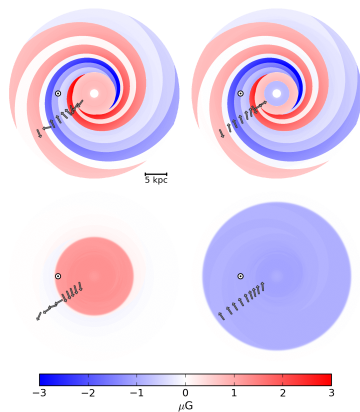


- standard approach:
 - ▶ only coherent field contributes to P and RM
 - ▶ CR electron density $n_{\text{cre}}(\mathbf{x}, E)$: fixed, independent of GMF

Status of GMF models

JF12 has become a “standard”:

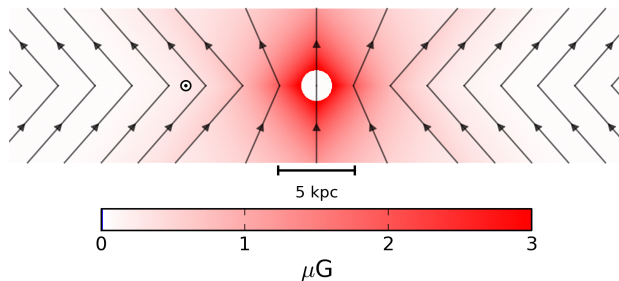
- fitted to RM and synchrotron data, 22 parameters for regular field
- (weak) spiral disk field:



Status of GMF models

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- fitted to RM and synchrotron data, 22 parameters for regular field
- halo + X-field



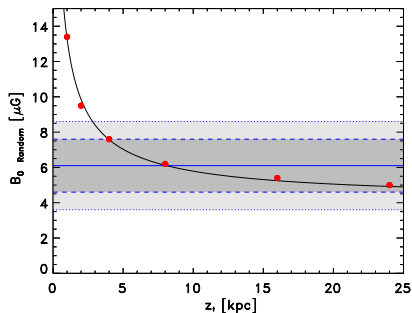
- propagation along X field eases CR escape

Problems of GMF models I

- thermal electron density $n_e(\mathbf{x})$ poorly constrained by DM's
- CR electron density $n_{\text{cre}}(\mathbf{x}, E)$ fixed
- $L_c \sim 150 \text{ pc}$ and $L \simeq 7 \text{ kpc} \Rightarrow N = L_{\text{max}}/L \sim 10$, $N_{\text{eff}} \sim \text{few}$

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- **B from RM is factor few smaller than from synchrotron:**



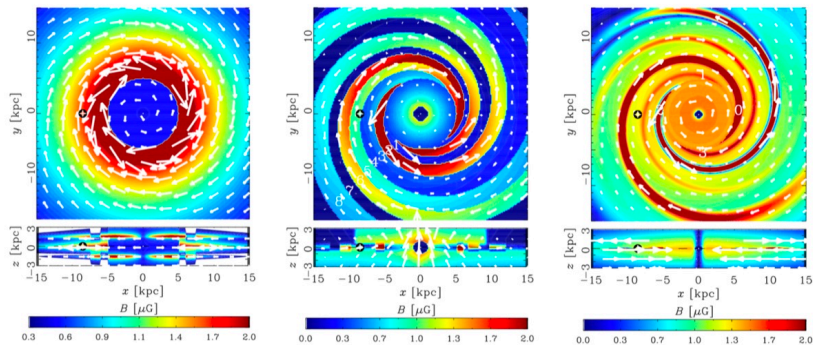
[Di Bernardo et al. '19]

$\Rightarrow b \gg B$

\Rightarrow isotropic diffusion

Problems of GMF models II

- **data are too sparse** to constrain (severely) models



- new suit of 8 UF24 models

⇒ talk by M. Unger

Approaches to CR propagation

1 UHECRs:

- ▶ use **model** for **Galactic Magnetic Field**
- ▶ **calculate trajectories** $\mathbf{x}(t)$ of **individual CRs** via $\mathbf{F}_L = q\mathbf{v} \times \mathbf{B}$.
- ▶ all **fluctuations** between l_{\max} and $\sim R_L/10$ have to be included
 \Rightarrow trajectory approach computationally very expensive for $E \searrow$

2 Galactic CR, low energies:

- ▶ CRs as relativistic fluid
- ▶ use effective diffusion picture
- ▶ connection to GMF only indirect:
 - ★ use quasi-linear theory to connect $D(E)$ and $P(k)$
 - ★ D is factor 50-100 too small *[Strong, Ptuskin, Moskalenko '07]*
 - ★ use D instead as fit parameter

Approaches to CR propagation

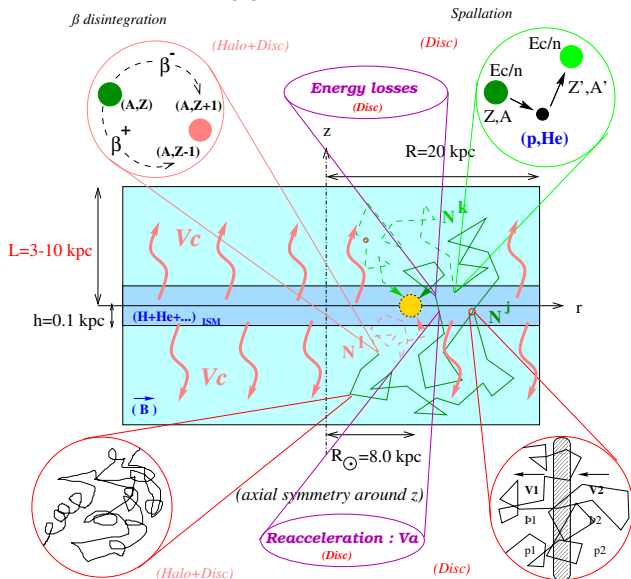
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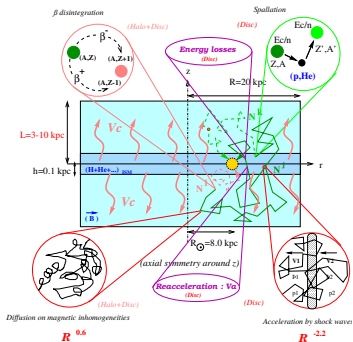
Standard diffusion approach:



Diffusion on magnetic inhomogeneities

Acceleration by shock waves

Standard diffusion approach:



- effective approach invites for simplifications:
- often $D_{ij}(E, \mathbf{x}) \rightarrow D(E)$, $\partial_t = 0$, etc.

How to connect diffusion and GMF?

- comparison of $D_{ij}(E)$:
 - ▶ **analytical** calculation: only approx. & limiting cases
 - ▶ **numerical** calculation straight-forward
- observable: grammage $\tau_{\text{esc}}(E) = L^2/(2D) \propto 1/X$

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Our approach:

- use model for Galactic magnetic field: Jansson-Farrar, Psirkhov et al.,...
- **calculate trajectories $\mathbf{x}(t)$ via $\mathbf{F}_L = q\mathbf{v} \times \mathbf{B}$.**

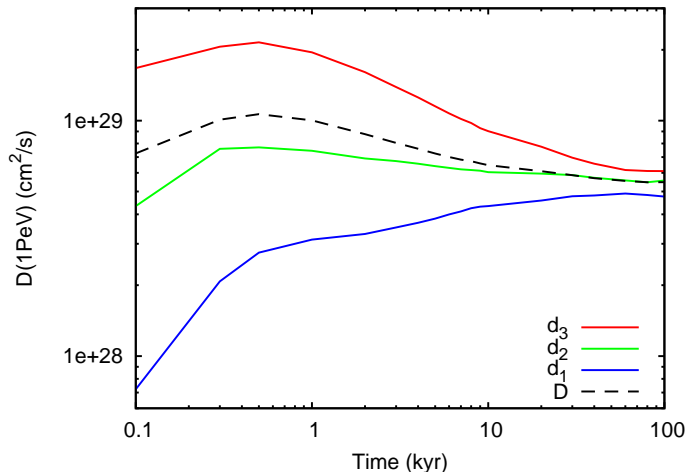
Our approach:

- use model for Galactic magnetic field: Jansson-Farrar, Psirkhov et al.,...
- calculate trajectories $x(t)$ via $F_L = qv \times B$.
- as preparation, let's **calculate diffusion tensor** in pure, isotropic turbulent magnetic field

Eigenvalues of $D_{ij} = \langle x_i x_j \rangle / (2t)$

$$E = 10^{15} \text{ eV}, B_{\text{rms}} = 4 \mu\text{G}$$

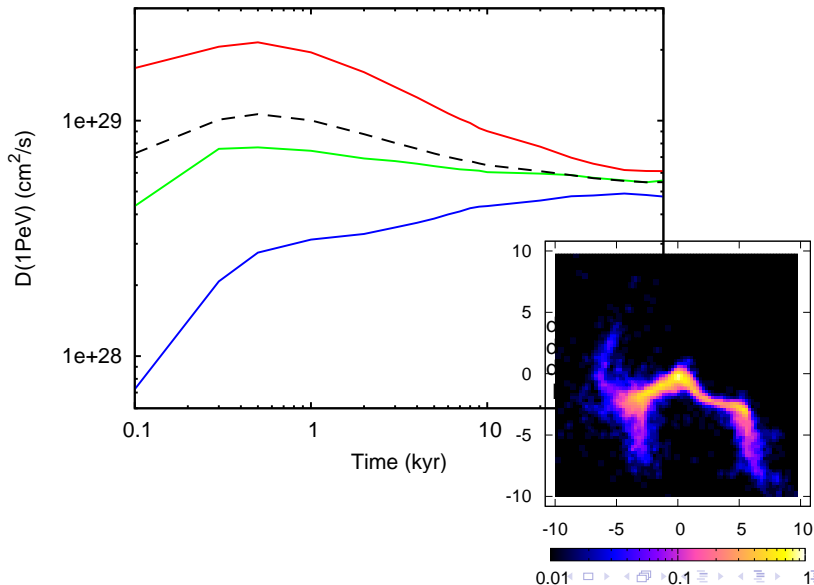
[Giacinti, MK, Semikoz ('12)]



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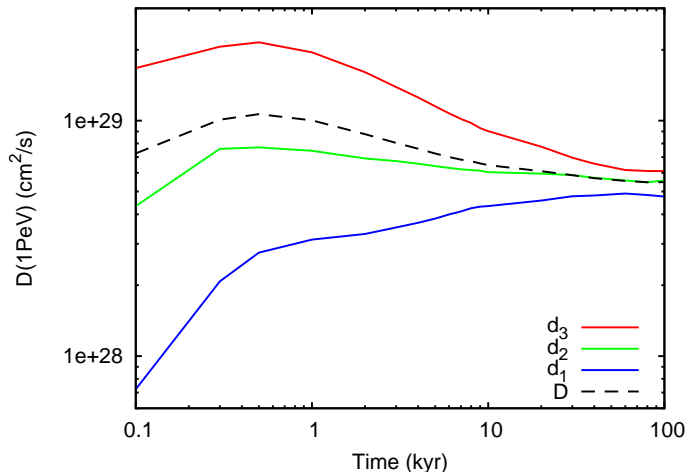
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[Giacinti, MK, Semikoz ('12)]



- asymptotic value is ~ 50 smaller than standard value

Is isotropic diffusion possible?

- for **isotropic** diffusion:

$$D = \frac{cL_0}{3} [(R_L/L_0)^{2-\alpha} + (R_L/L_0)^2]$$

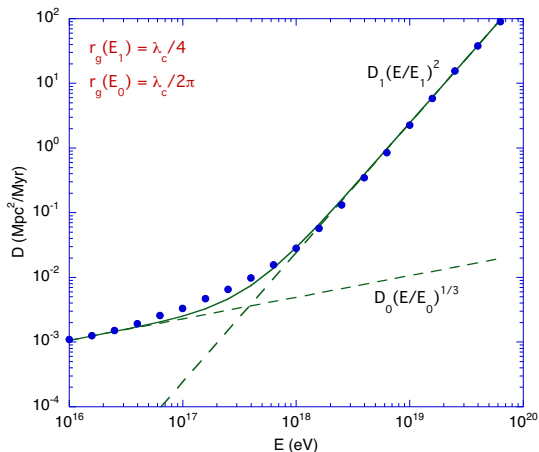
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- for isotropic diffusion:

for $\alpha = 5/3$

$$D = \frac{cL_0}{3} \left[(R_L/L_0)^{1/3} + (R_L/L_0)^2 \right]$$

with $L_0 \simeq L_c/(2\pi)$

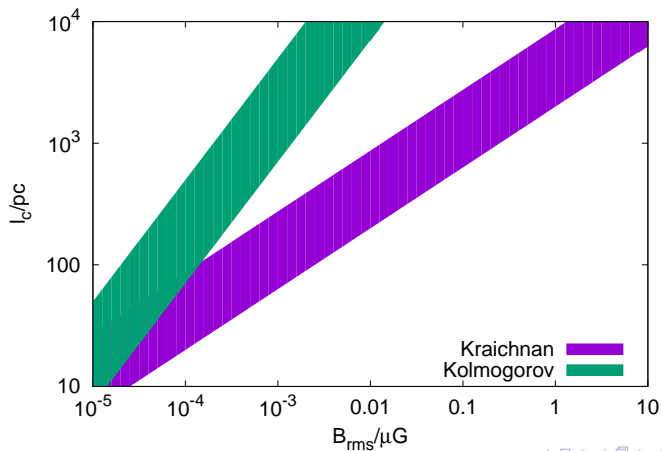


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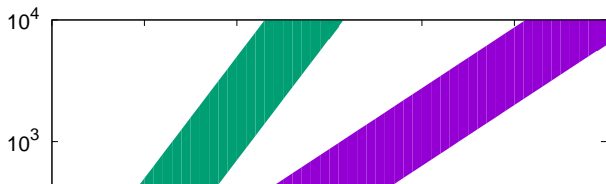
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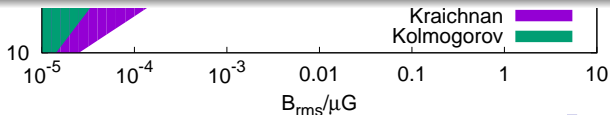
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isotropic diffusion is excluded:

- ▶ which effects do we miss?
- ▶ regular field \Rightarrow anisotropic diffusion

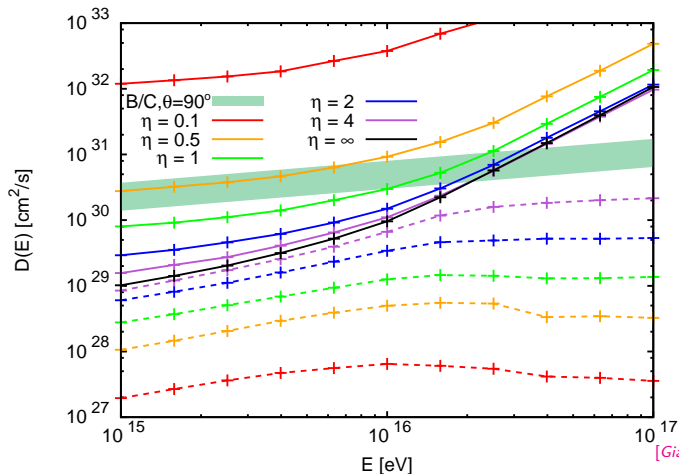


Anisotropic diffusion – 2 options:

- anisotropic turbulence
- dominance of regular field, $B_{\text{rms}} = \eta B_0 \ll B_0 \Rightarrow D_{\parallel} \gg D_{\perp}$

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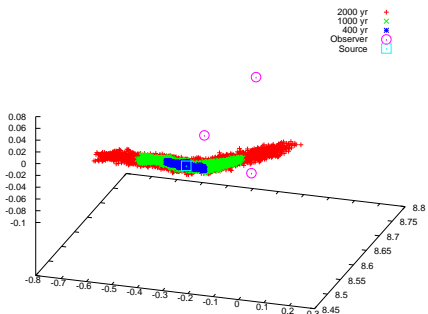
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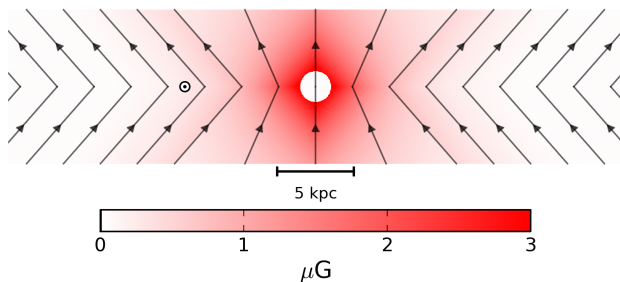


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\Rightarrow anisotropic CR propagation

$\Rightarrow D_{\parallel} \sin(2\vartheta)$ reduces grammage

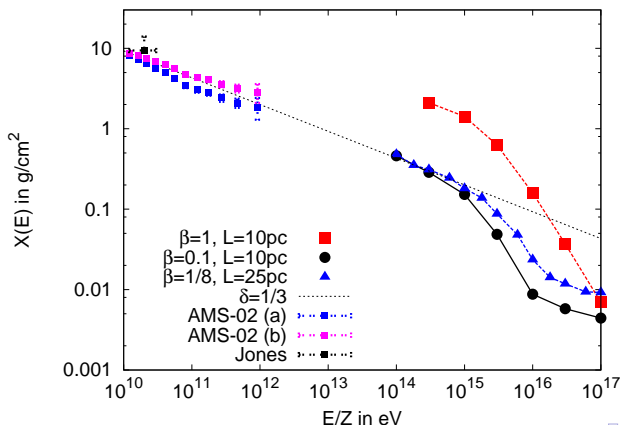


How dominant is the regular field?

- LOFAR: $l_{\text{coh}} \lesssim 10 \text{ pc}$ in disc
- use JF12 model and rescale turbulent field
- determine magnitude of random $B_{\text{rms}}(\mathbf{x})$ from grammage $X(E)$

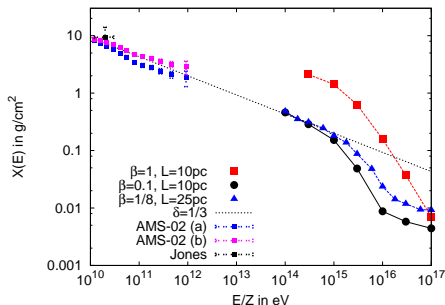
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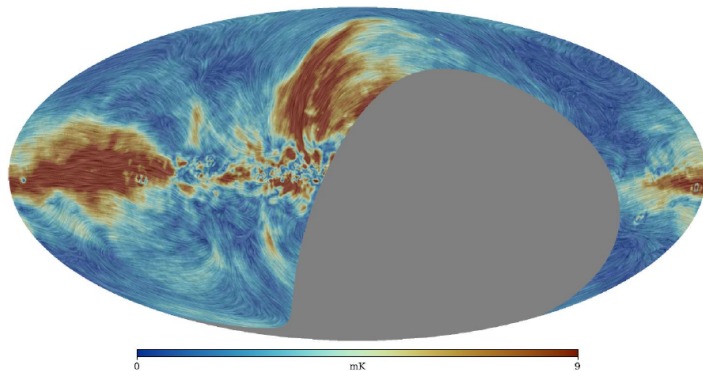


⇒ prefers **weak turbulent fields**

⇒ contradiction to **synchrotron** intensity

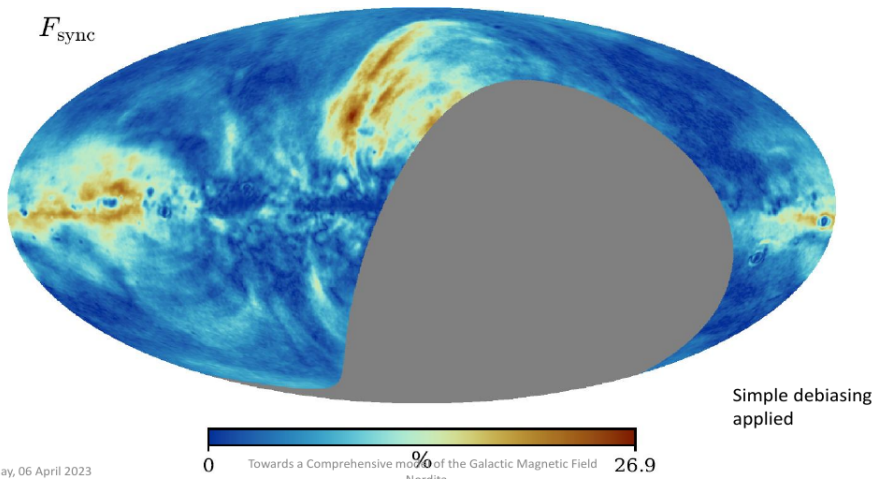
C-BASS polarisation experiment at 5 GHz:

[P. Leahy '23]



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rsday, 06 April 2023

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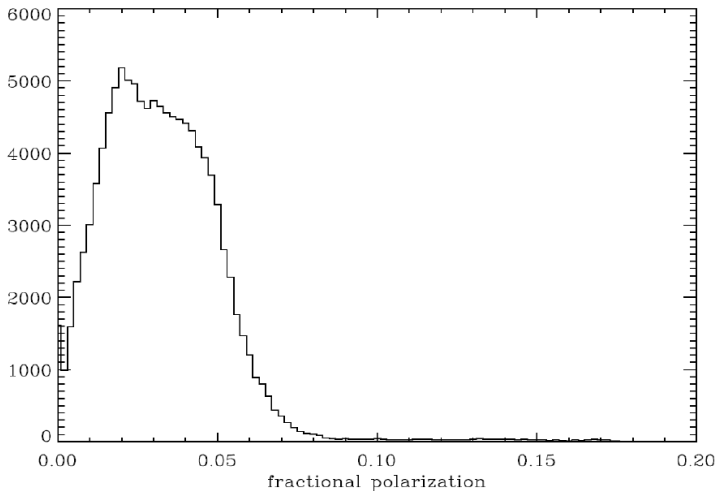
High-latitude sky mask

- $b > 30^\circ$
- Avoids obvious structures:
 - Loop I
 - Loop III
 - Virgo A



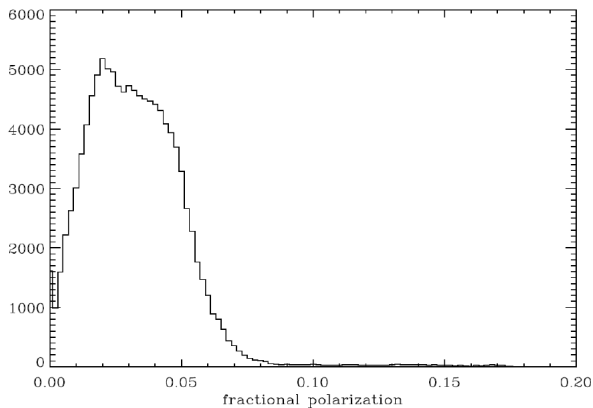
C-BASS: Polarisation degree per pixel

[P. Leahy '23]



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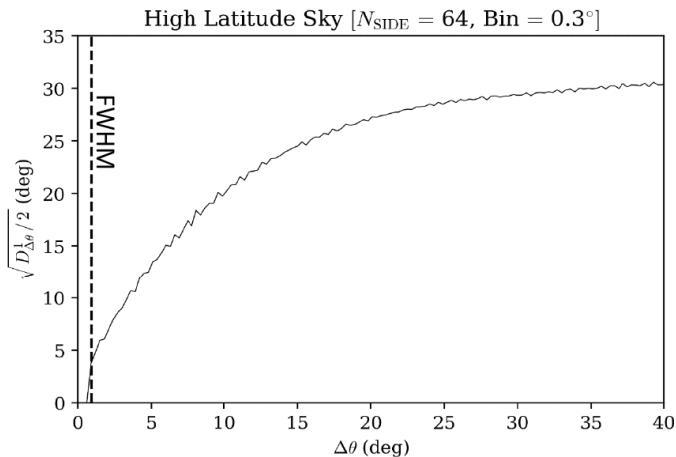
[P. Leahy '23]



- average polarisation $\langle P \rangle \simeq 3.3\%$, almost everywhere $< 10\%$.
- Local Spur, fan regions: $P \sim 30\%$

C-BASS: Structure function

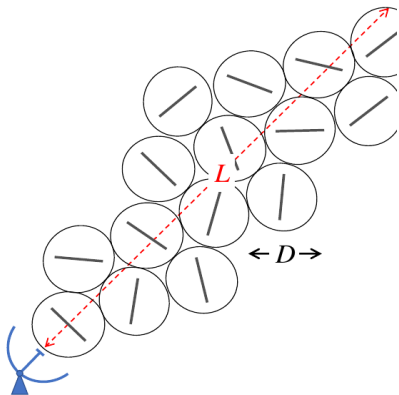
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- typical angular scale $\vartheta \simeq 15^\circ$

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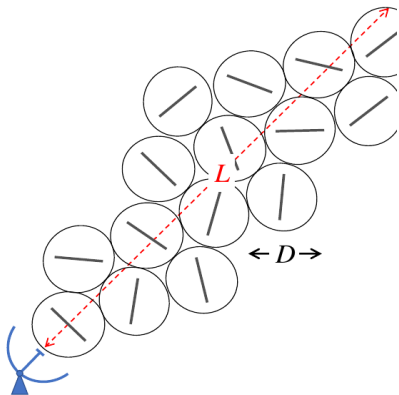
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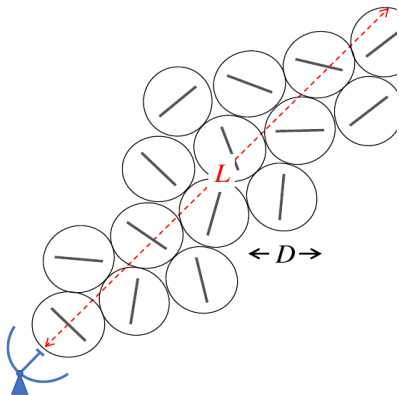
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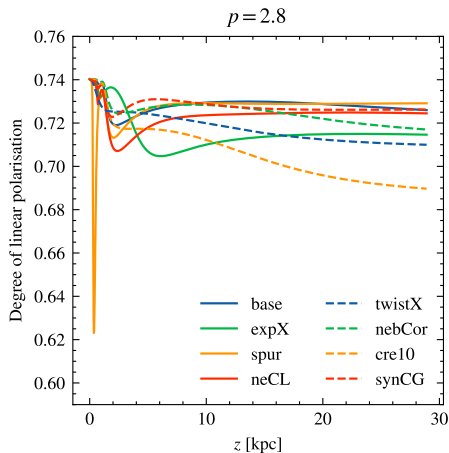
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- typical angular scale $\vartheta \simeq 15^\circ$
- $\sin \vartheta/2 \simeq D/L \simeq 1/N$
- $\vartheta \simeq 15^\circ \Rightarrow N \simeq 8$
- but $N \simeq 8$ gives $\langle P \rangle \simeq 0.7/\sqrt{N} \simeq 27\%$

Polarisation degree in UF24 models

- regular field B_0 from UF24 model



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- **add turbulent field b** such that locally RMS $b = \beta B_0$

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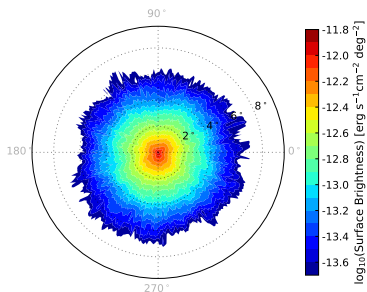
- regular field B_0 from UF24 model
- add turbulent field b such that locally RMS $b = \beta B_0$
- calculate P for various field realisations:

β	base	neCL	expX	spur	cre10	synCG	twistX	nebCor
1	0.36	0.37	0.36	0.39	0.33	0.39	0.36	0.35
1.5	0.22	0.23	0.22	0.24	0.2	0.24	0.22	0.21

- agrees with $\langle P \rangle$ in Local Spur & fan regions $|b| > 30^\circ$
- small $\langle P \rangle$ would require $b \gg B_0$

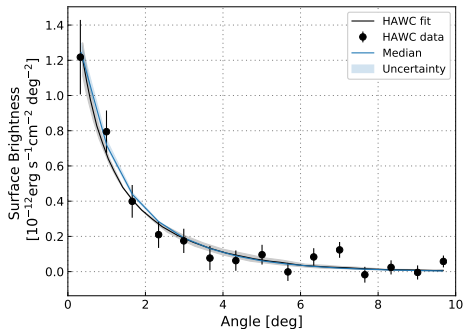
TeV halos (around PWNe?)

- HAWC: **slow diffusion around Geminga: $D \sim D_0/100$**



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- **consistent** with expectation for **isotropic diffusion**, $B = 3\mu\text{G}$ and $L_c = 1\text{ pc}$



[López-Coto, Giacinti '17]

TeV halos (around PWNe?)

- HAWC: slow diffusion around Geminga: $D \sim D_0/100$
- consistent with expectation for isotropic diffusion, $B = 3\mu\text{G}$ and $L_c = 1\text{ pc}$
- **three options:**
 - ▶ regular field “expelled” around SNR
 - ▶ self-generated turbulence close to SNR/PWNe
 - ▶ typical situation in disk

Three-component model for the GMF:

add extended halo/corona field:

- **disk** field: small L_c & **turbulent** dominated such $D \simeq D_{\text{iso}} \sim D_0/100$
- **halo** field: large L_c , dominated by **regular** field, $\sin(2\vartheta)D_{\parallel} \simeq D_0$
- **extended halo/corona**: large L_c , **turbulent** field up to 200 kpc

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simplifications:

- power law $dN/dE \sim E^{-p}$ for e^{\pm}
- profiles $B(z) \sim \exp[-(z - z_0)/z_t]$ and B_{min}
- determine $n_e(z)$ from stationary 1d advection-diffusion equation

Stationary 1d advection-diffusion equation

- solution for $v = 0$:

$$n(z) = n_0 - j_0 \int_0^z \frac{dz'}{D(z')}$$

- use $D(z) = D_0 \exp(z/z_0)$:

$$n(z) = n_0 + \frac{j_0 z_0}{D_0} [\exp(-z/z_0) - 1]$$

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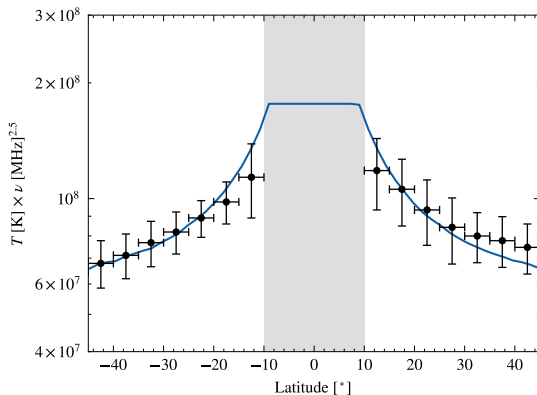
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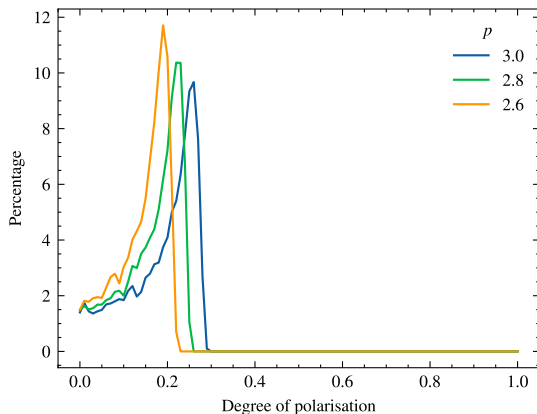
- **two cases:**

- ▶ $j_0 z_0 / D_0 n_0 > 1$: **free-escape boundary**
- ▶ $j_0 z_0 / D_0 n_0 < 1$: **non-zero $n(z)$ for $z \rightarrow \infty$**

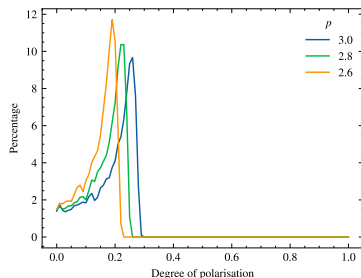
Three-component model for the GMF: fit to synchrotron



Three-component model for the GMF: polarisation



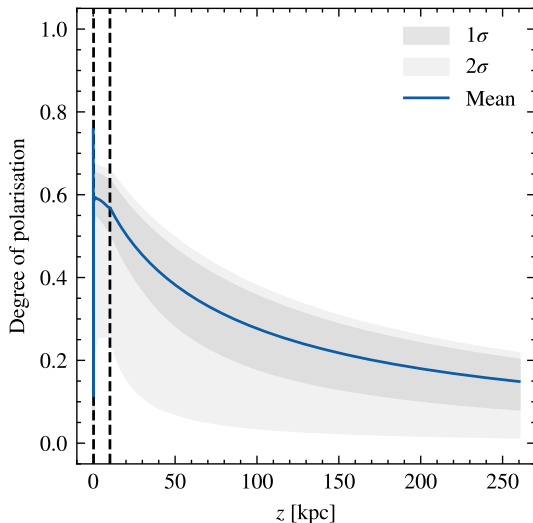
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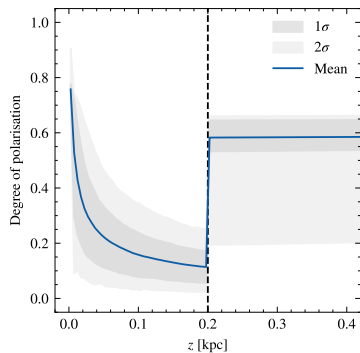
distribution of P :

- ▶ very small P requires partial **cancellation** of disk and halo contribution

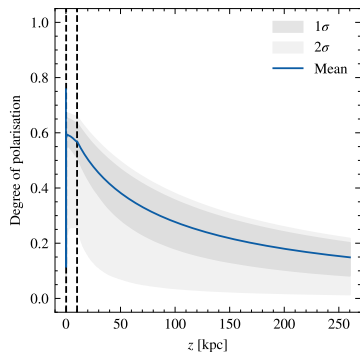
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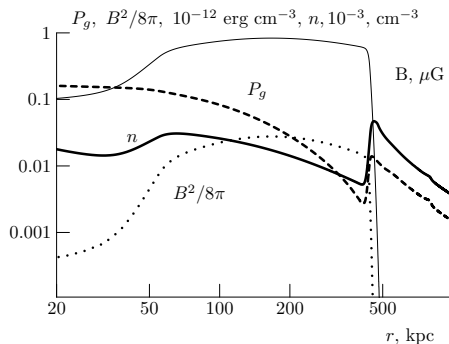
Three-component model for the GMF: polarisation



- representative values: $B_{\min} = 0.1\mu\text{G}$ and $n_{\min} = 0.1n_0$
- compatible with CR escape

Other suggestions for an extended halo

- ..., Taylor et al. '14
- Zirakashvili, Ptuskin, Rogovaya '23:



Consequences:

- additional time-delays (and deflections) for UHECRs
- diffuse photon and neutrino fluxes

Conclusions

- 1 CR escape requires fast diffusion
 - ▶ regular field should dominate in halo
 - ▶ total contribution to I subdominant
- 2 Polarisation degree in C-Bass very low
 - ▶ large contribution to intensity from turbulent dominated region
 - ▶ region should be large $N = L/L_c \gg 1$
- 3 disk turbulent dominated
 - ▶ $N = L/L_c \simeq \text{few}$
- 4 shape of CR sources in photons:
 - ▶ disk: spherical
 - ▶ halo: elongated

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