

# A Numerical Implementation of the LASS Subtraction Scheme

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# Motivation

# Motivation

LASS : Local Analytic Sector Subtraction scheme [arXiv:2212.11190](#), [2010.14493](#), [1806.09570](#)

- Massless partons
- For final state radiation (i.e.  $e^+ e^-$  collisions)

Subtraction and slicing methods at NNLO QCD:

- Several methods for final/initial state both for massless and massive partons
- Several processes already calculated with several different methods
- Methods already started to be extended to  $N^3\text{LO}$

What is the point creating yet another?



# Motivation

Motto:

Zoltán Nagy: "We solve a **math** problem"

- Can have multiple solutions
- Can reach the same answer on a multitude of paths
- Physics and math is also about elegance
- Life will not stop at NNLO:
  - Can an  $N^n$ LO scheme **extended** to  $N^{n+1}$ LO?
  - Can a scheme for  $e^+ e^-$  extended to **pp collisions**?
  - At what **cost**?

What about **carbon-friendly** computing?



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## LASS, a practitioner's intro

# LASS, a practitioner's intro

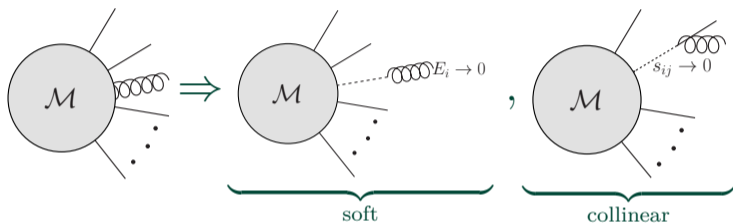
Key ingredients of the subtraction scheme:

- **Grouping** of the subtraction terms
- Phase space **mappings** used to get underlying Born and real configurations
- Phase space **partitioning**



# LASS, a practitioner's intro – grouping of terms

Kinematic singularity structure is easy at NLO (up to one unresolved emission):



$$d\sigma^R : |\mathcal{M}^R|^2 - K^{(1)}$$

$$K^{(1)} = \sum_i \mathcal{S}_i + \sum_{i,j < i} (c_{ij} - \mathcal{S}_i c_{ij} - \mathcal{S}_j c_{ij})$$

$\Rightarrow$  Removal of overlapping results in proliferation of terms!





# LASS, a practitioner's intro – grouping of terms

⇒ Removal of overlapping results in proliferation of terms!

$$K^{(1)} = \sum_i \mathcal{S}_i + \underbrace{\sum_{i,j < i} (c_{ij} - \mathcal{S}_i c_{ij} - \mathcal{S}_j c_{ij})}_{c_{ij}^{\text{hc}}} = \sum_i \mathcal{S}_i + \sum_{i,j < i} c_{ij}^{\text{hc}}$$

Simple replacement of AP kernels with their hard-collinear variants:

$$\langle P_{q\bar{q}} \rangle = T_R \left[ 1 - \frac{2x_q x_{\bar{q}}}{1 - \epsilon} \right], \quad \langle P_{qg} \rangle = C_F \left[ 2 \frac{x_q}{x_g} + (1 - \epsilon)x_g \right], \quad \langle P_{gg} \rangle = 2C_A \left[ \frac{x_1}{x_2} + \frac{x_2}{x_1} + x_1 x_2 \right],$$

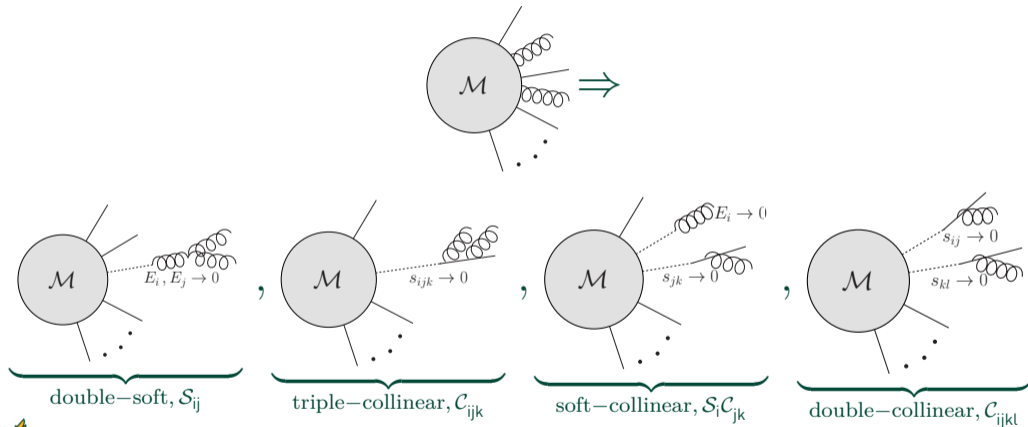
$$\langle P_{q\bar{q}} \rangle = \langle P_{q\bar{q}}^{\text{hc}} \rangle, \quad \langle P_{qg}^{\text{hc}} \rangle = C_F(1 - \epsilon)x_g, \quad \langle P_{gg}^{\text{hc}} \rangle = 2C_A x_1 x_2$$

$$c_{ij} \rightarrow c_{ij}^{\text{hc}} \quad : \quad P_{ij(r)}^{\mu\nu} \rightarrow P_{ij(r)}^{\mu\nu, \text{hc}} \quad , \quad \mathcal{S}_i c_{ij} \rightarrow \emptyset$$



# LASS, a practitioner's intro – grouping of terms

At NNLO we have **many more** kinematic singularities:



# LASS, a practitioner's intro – grouping of terms

With a much **richer overlap** between various terms:

$$\begin{aligned} K^{(2)} = & \sum_{i,j < i} \mathcal{S}_{ij} + \sum_j \sum_{k < j, i \neq j, k} \left( \mathcal{S}_i \mathcal{C}_{jk} - \mathcal{S}_i \mathcal{C}_{jk} \mathcal{S}_{ij} - \mathcal{S}_i \mathcal{C}_{jk} \mathcal{S}_{ik} \right) + \\ & + \sum_{i,j < i, k < j} \left( \mathcal{C}_{ijk} - \mathcal{C}_{ijk} \mathcal{S}_{ij} - \mathcal{C}_{ijk} \mathcal{S}_{ik} - \mathcal{C}_{ijk} \mathcal{S}_{jk} - \mathcal{C}_{ijk} \mathcal{S}_i \mathcal{C}_{jk} - \mathcal{C}_{ijk} \mathcal{S}_j \mathcal{C}_{ik} - \mathcal{C}_{ijk} \mathcal{S}_k \mathcal{C}_{ij} + \right. \\ & + \left. \mathcal{C}_{ijk} \mathcal{S}_{ij} \mathcal{S}_i \mathcal{C}_{jk} + \mathcal{C}_{ijk} \mathcal{S}_{ij} \mathcal{S}_j \mathcal{C}_{ik} + \mathcal{C}_{ijk} \mathcal{S}_{ik} \mathcal{S}_i \mathcal{C}_{jk} + \mathcal{C}_{ijk} \mathcal{S}_{ik} \mathcal{S}_k \mathcal{C}_{ij} + \mathcal{C}_{ijk} \mathcal{S}_{jk} \mathcal{S}_j \mathcal{C}_{ik} + \mathcal{C}_{ijk} \mathcal{S}_{jk} \mathcal{S}_k \mathcal{C}_{ij} \right) + \\ & + \sum_{i,j < i} \sum_{\substack{k,l < k \\ k < i, j \neq l}} \left( \mathcal{C}_{ijkl} - \mathcal{C}_{ijkl} \mathcal{S}_i \mathcal{C}_{kl} - \mathcal{C}_{ijkl} \mathcal{S}_j \mathcal{C}_{kl} - \mathcal{C}_{ijkl} \mathcal{S}_k \mathcal{C}_{ij} - \mathcal{C}_{ijkl} \mathcal{S}_l \mathcal{C}_{ij} + \right. \\ & + \left. \mathcal{C}_{ijkl} \mathcal{S}_{ik} + \mathcal{C}_{ijkl} \mathcal{S}_{il} + \mathcal{C}_{ijkl} \mathcal{S}_{jk} + \mathcal{C}_{ijkl} \mathcal{S}_{jl} \right) \end{aligned}$$



# LASS, a practitioner's intro – grouping of terms

Subterms can be grouped into **hard-collinear groups**:

$$\begin{aligned}
 K^{(2)} = & \sum_{i,j < i} S_{ij} + \sum_{j,k < j} \sum_{i \neq j,k} \underbrace{\left( S_i C_{jk} - S_i C_{jk} S_{ij} - S_i C_{jk} S_{ik} \right)}_{S_i C_{jk}^{hc}} + \\
 & + \sum_{\substack{i,j < i \\ k < j}} \underbrace{\left( C_{ijk} - C_{ijk} S_{ij} - C_{ijk} S_{ik} - C_{ijk} S_{jk} \right)}_{C_{ijk}^{hc}} - \underbrace{\left( C_{ijk} S_i C_{jk} + C_{ijk} S_{ij} S_i C_{jk} + C_{ijk} S_{ik} S_i C_{jk} - \dots \right)}_{S_i C_{jk}^{hc} C_{ijk}} + \\
 & + \sum_{i,j < i} \sum_{\substack{k,l < k \\ k < i \\ j \neq l}} \underbrace{\left( C_{ijkl} - C_{ijkl} S_i C_{kl} - \dots - C_{ijkl} S_l C_{ij} + C_{ijkl} S_{ik} + \dots + C_{ijkl} S_{jl} \right)}_{C_{ijkl}^{hc}}
 \end{aligned}$$



# LASS, a practitioner's intro – grouping of terms

Using the hard-collinear formalism the structure becomes very **compact**:

$$\begin{aligned} K^{(2)} = & \sum_{i,j < i} \mathcal{S}_{ij} + \\ & + \sum_{j,k < j} \sum_{i \neq j,k} \mathcal{S}_i c_{jk}^{\text{hc}} + \\ & + \sum_{\substack{i,j < i \\ k < j}} (\mathcal{C}_{ijk}^{\text{hc}} - \mathcal{S}_i c_{jk}^{\text{hc}} c_{ijk} - \mathcal{S}_j c_{ik}^{\text{hc}} c_{ijk} - \mathcal{S}_k c_{ij}^{\text{hc}} c_{ijk}) + \\ & + \sum_{i,j < i} \sum_{\substack{k,l < k \\ k < i \\ j \neq l}} c_{ijkl}^{\text{hc}} \end{aligned}$$



# LASS, a practitioner's intro – mappings

- The method heavily **relies on** single and iterated **Catani-Seymour** mappings
  - $K^{(1)}$  in the  $n + 2$  parton cont.:

$$\{k\}_{(n+2)} \xrightarrow{i,j,r \rightarrow \bar{j}, \bar{r}} \{\bar{k}\}_{(n+1)}^{(ijr)}, \quad \{k\}_{(n+2)} \xrightarrow{i,c,d \rightarrow \bar{c}, \bar{d}} \{\bar{k}\}_{(n+1)}^{(icd)}$$

- Also in  $K^{(2)}$  and in  $K^{(12)}$  in all possible ways:

$$\{k\}_{(n+2)} \xrightarrow{i,c,d \rightarrow \bar{c}, \bar{d}} \{\bar{k}\}_{(n+1)}^{(icd)} \xrightarrow{j,e,f \rightarrow \tilde{e}, \tilde{f}} \{\tilde{k}\}_{(n)}^{(icd,jef)},$$

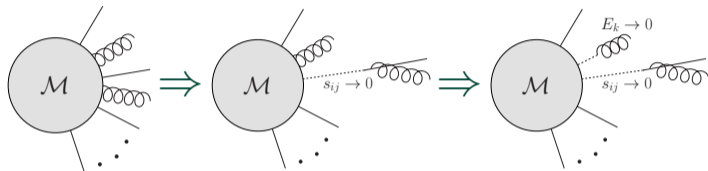
⋮

$$\{k\}_{(n+2)} \xrightarrow{k,r,j \rightarrow \bar{r}, \bar{j}} \{\bar{k}\}_{(n+1)}^{(krj)} \xrightarrow{i, \bar{r}, \bar{j} \rightarrow \tilde{r}, \tilde{j}} \{\tilde{k}\}_{(n)}^{(krj, i\bar{r}\bar{j})}$$



# LASS, a practitioner's intro – phase space partitioning

- $K^{(1)}$  can also diverge in **doubly unresolved** limits:



- $K^{(2)}$  can also **diverge in singly unresolved** limits!
- $\Rightarrow$  These double-countings (**spurious singularities**) should also be removed:

$$|\mathcal{M}^{\text{RR}}|^2 - K^{(1)} - K^{(2)} - K^{(12)}$$

- Careful **synchronization** is needed in **mappings** and **parametrizations**!



# LASS, a practitioner's intro – phase space partitioning

- Careful **synchronization** is needed in **mappings** and **parametrizations**!
- ⇒ Introduction of **sector functions** can simplify the procedure:

$$1 = \sum_{i < j < k} z_{ijk} + \sum_{i < j} \sum_{\substack{k < l \\ i < k, j \neq l}} z_{ijkl}$$

- Certain sectors **only allow** certain limits/subterms (and overlaps):

$$z_{ijk} \left\{ \begin{array}{l} \mathcal{S}_i, \mathcal{S}_j, \mathcal{S}_k, \mathcal{C}_{ij}, \mathcal{C}_{ik}, \mathcal{C}_{jk}, \\ \mathcal{S}_{ij}, \mathcal{S}_{ik}, \mathcal{S}_{jk}, \\ \mathcal{S}_i \mathcal{C}_{jk}, \mathcal{S}_j \mathcal{C}_{ik}, \mathcal{S}_k \mathcal{C}_{ij}, \\ \mathcal{C}_{ijk} \end{array} \right. \quad z_{ijkl} \left\{ \begin{array}{l} \mathcal{S}_i, \mathcal{S}_j, \mathcal{S}_k, \mathcal{S}_l, \mathcal{C}_{ij}, \mathcal{C}_{kl}, \\ \mathcal{S}_{ik}, \mathcal{S}_{il}, \mathcal{S}_{jk}, \mathcal{S}_{jl}, \\ \mathcal{S}_i \mathcal{C}_{kl}, \mathcal{S}_j \mathcal{C}_{kl}, \mathcal{S}_k \mathcal{C}_{ij}, \mathcal{S}_l \mathcal{C}_{ij}, \\ \mathcal{C}_{ijkl} \end{array} \right.$$





# LASS, a practitioner's intro – phase space partitioning

- Subterms have **operator-like effect** on sector functions
- To ensure **analytic integrability** sum-rules should be fulfilled:

$$\sum_{k \neq i, j} [\mathcal{Z}_{ijk} + \sum_{l \neq i, j, k} \mathcal{Z}_{ikjl}] \mathcal{S}_{ij} = \mathcal{S}_{ij}, \quad [\mathcal{Z}_{ikl} + \sum_{j \neq i, k, l} \mathcal{Z}_{ijkl}] \mathcal{S}_i \mathcal{C}_{kl} = \mathcal{S}_i \mathcal{C}_{kl},$$
$$\mathcal{Z}_{ijk} \mathcal{C}_{ijk} = \mathcal{C}_{ijk}, \quad \mathcal{Z}_{ijkl} \mathcal{C}_{ijkl} = \mathcal{C}_{ijkl}$$

- Note: this is not the case in  $K^{(1)}$  and  $K^{(12)}$  !

$$\sum_{j < k} \mathcal{Z}_{ijk} \mathcal{S}_i = \mathcal{S}_i, \quad \sum_k \mathcal{Z}_{ijk} \mathcal{C}_{ij} = \mathcal{C}_{ij} \bar{\mathcal{Z}}_{jk} \left( \{\bar{k}\}^{(ijr)} \right),$$

...



A way to check subtractions



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# A way to check subtractions

- **Numerical checks are essential** for validating both **individual** subterms and **whole** contributions:
  - Individual subterms:
    - Proper usage of mappings
    - Parameters are calculated correctly
    - Normalizations are correct
  - Full ensemble:
    - Parametrizations are **synchronized** (where need be)
    - **Spurious cancellations** are happening:

$$K^{(1)} \text{ limits : } K^{(2)} \leftrightarrow K^{(12)}, \quad K^{(2)} \text{ limits : } K^{(1)} \leftrightarrow K^{(12)}$$



# A way to check subtractions

Idea: check everything **numerically**, but with **arbitrary precision**!

- Fortran90 and MPFUN20 by David Bailey to have **arbitrary precision floating numbers**
- **50-60 digits** suffice for most checks (only for checks, **MC in double precision**)
- Starting with an  $n + 2$  or  $n + 1$  parton PS point we create a **sequence** of PS points bringing partons to specific limits
- Ratio is monitored between SME and subterm/full contribution:

$$n + 2 \text{ parton line : } \lim_{5||6||7} \frac{|\mathcal{M}^{\text{RR}}|^2}{\mathcal{C}_{765}}, \quad \lim_{5,6 \rightarrow 0} \frac{\mathcal{S}_{65}\mathcal{C}_{765}}{\mathcal{C}_{765}}, \quad \frac{|\mathcal{M}^{\text{RR}}|^2}{K^{(1)} + K^{(2)} + K^{(12)}}$$



# A way to check subtractions – example

- LASS currently is for **final state** radiation only ( $e^+ e^-$  collisions)
- **Three-jet production** is the obvious choice
- Most complicated subprocess:  $e^+ e^- \rightarrow q\bar{q} ggg$ 
  - **187** subtraction terms
  - **49** singular regions



## A way to check subtractions – example

Ratio sequence for a single soft subtraction term and the RR SME:

$$\lim_{5 \rightarrow 0} \frac{|\mathcal{M}^{\text{RR}}|^2}{\mathcal{S}_5}$$

Ratio sequence for the **complete**  $K^{(1)} + K^{(2)} + K^{(12)}$  and the RR SME in a collinear limit:

$$\lim_{7||3} \frac{|\mathcal{M}^{\text{RR}}|^2}{K^{(1)} + K^{(2)} + K^{(12)}}$$



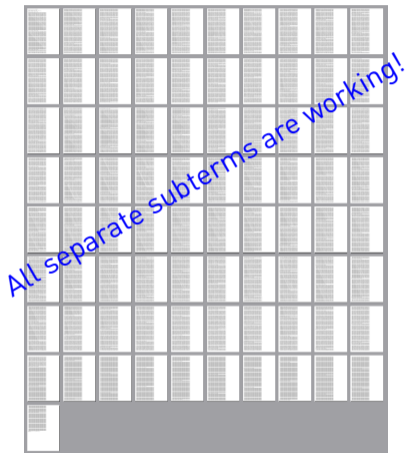
# A way to check subtractions – example

```
Checking subtractions for ep em > d db g g g
Checking individual subterm (1/187)
1.3615262894238763826759820725791757465583666344146e1
1.0694493566690845438842431546992987055507713973236e0
1.0007162821758278182424090928223132563018751174295e0
1.0000071987341108237419042328666792441638763318880e0
1.0000000720247434870412374459056904352354450789094e0
1.0000000007202849869842647183822196539254192619824e0
1.0000000000072028874369372901453904752500972080358e0
1.000000000000000720289119379656962975720606869247628e0
1.000000000000000007202891569483995737357639735843981e0
1.000000000000000000072028916070527533297166023327910e0
1.000000000000000000000720289161080962923876849506640e0
1.00000000000000000000000720289161185316831171844311e0
1.000000000000000000000000072028916112228855904271608e0
1.00000000000000000000000000720289161122664246635284e0
1.00000000000000000000000000007202891611227018153945e0
1.00000000000000000000000000000072028916112270557227e0
1.00000000000000000000000000000000720289161122705947e0
1.0000000000000000000000000000000007202891611227059e0
1.00000000000000000000000000000000072028916112270e0
```

```
Checking singular region (12/49)
9.6594481743241900295619875419238279005187022510109e-1
1.0058615538802959144867884196860137973711485348134e0
9.9231059807099869356172685061099655988514093224880e-1
9.9583088426373405730413153367459343558862822235912e-1
9.9849562883100781816326951516801463419402897054026e-1
9.9950556967204941845739013900159932542190523615067e-1
9.9984177766130266861510092596371921504512639386555e-1
9.9994977880695929580926317624625089086796062180328e-1
9.9998409997765460232995100855871423405820977458026e-1
9.9999497010288688320443743161728335546792151263641e-1
9.9999840922001909000444820286260247019193046061875e-1
9.9999949693251516037014963280570874098054795389092e-1
9.9999984091422459144755648094664046742776493134045e-1
9.9999994969247378539788299063072534264967525981580e-1
9.9999998409134468611615545698564930910050993985655e-1
9.9999999496923960123804514553411354665657800543057e-1
9.9999999840913369088147655043925236337199932696573e-1
9.9999999949692388235079173196393701217379754206853e-1
9.9999999984091336131084641210578013734477707761172e-1
9.9999999994969238745734905001952273575863643999079e-1
```



# A way to check subtractions – example





# A first glimpse at numerical integrations



# A first glimpse at numerical integrations

- Checking subterms and complete subtractions are important but **thorough numerical checks** are also needed
  - A sequence does not **cover whole phase space**
  - A sequence is too **artificial**
  - What happens close to **two-jet limit**?
  - ...
- One possible check is the **saturation** of cross section:

$$\min_{i,j} y_{ij} = \min_{i,j} \frac{S_{ij}}{S} \geq y_{\min}$$

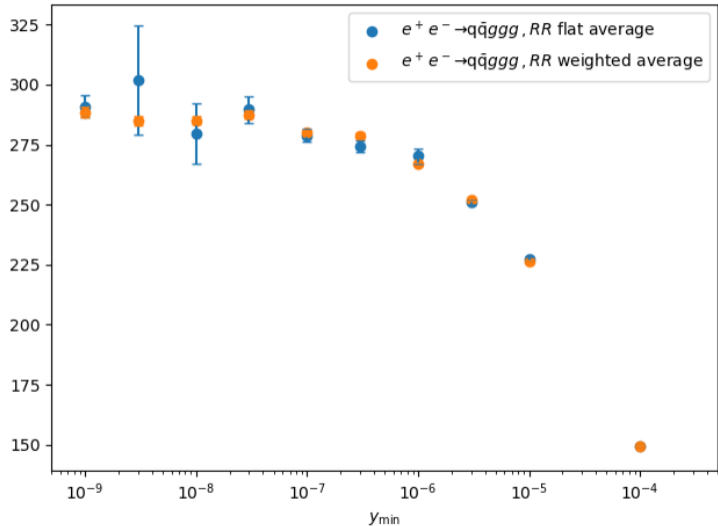
- Cut at edge of phase space, than **decrease** the cut
- In case all singularities are regularized contribution should **stabilize** (saturate)



# A first glimpse at numerical integrations

$n+2$  parton contribution  
(double-real):

Saturation plot for  
 $e^+ e^- \rightarrow q\bar{q}ggg$

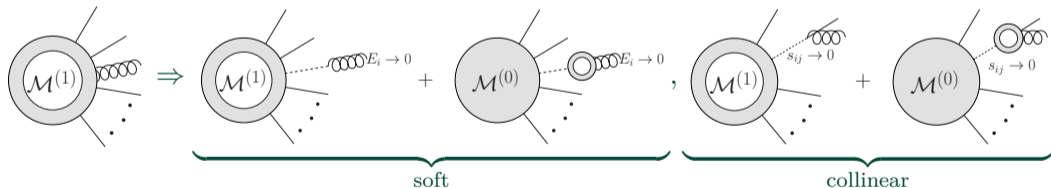


# A first glimpse at numerical integrations

Contributions at NNLO:

$$d\sigma = d\sigma^{\text{RR}} + d\sigma^{\text{RV}} + d\sigma^{\text{VV}}$$

- $d\sigma^{\text{RR}}$  is the **most difficult**, double-unresolved radiation
- $d\sigma^{\text{RV}}$  only has single-unresolved radiation but with **one-loop amplitudes!**

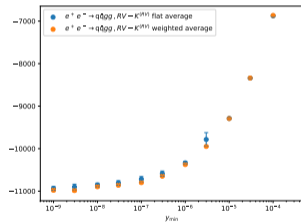
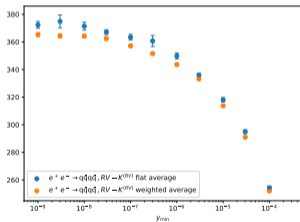
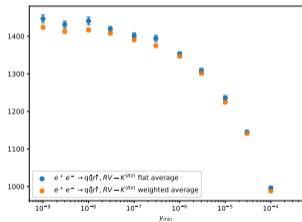
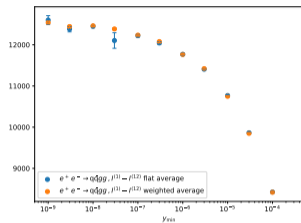
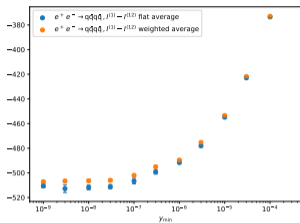
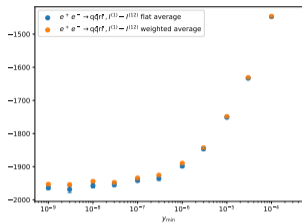


# A first glimpse at numerical integrations

- The  $n + 1$  parton contribution is **NLO in nature**
- Subterms are **more complicated** (and have more)
  - Kernels from **one-loop factorization**
  - Integrated subterms ( $I^{(1)}$ , from  $n + 2$  parton cont.) also diverge



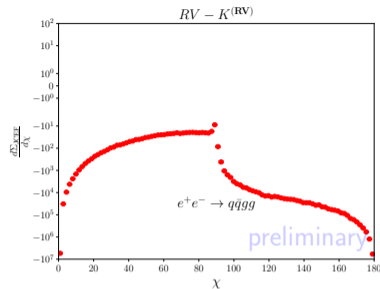
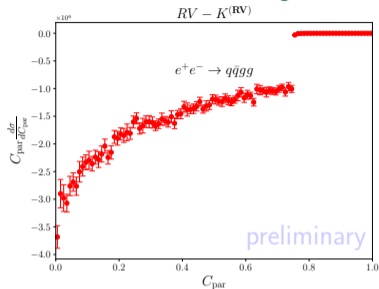
# A first glimpse at numerical integrations



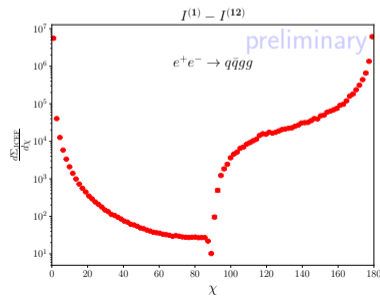
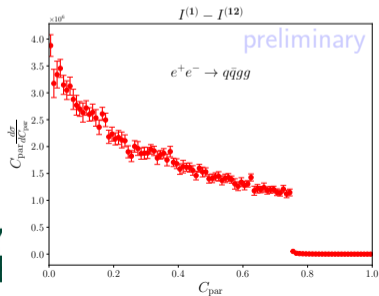
Saturation works for all contributions!

# A first glimpse at numerical integrations

One-loop part:



Integrated subterms:



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# Summary



# Summary

- **First numerical** implementation of LASS (at NNLO) in the making
- Numerically **proved**: subterms regularize kinematic singularities
- **Framework** is created: generate subtractions automatically to arbitrary processes
- Applied to three-jet production in  $e^+ e^-$
- Cross section contributions **saturate** in the  $n + 2$  and  $n + 1$  parton contributions
- **Histograms** can already be **produced** for the  $n + 1$  parton contribution



Thank you for your attention!