

GRIFFIN: A C++ library for higher-order electroweak corrections

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L. Chen and A. Freitas,
SciPost Phys. Codeb. 2023, 18
[arXiv:2211.16272]

github.com/lisongc/GRIFFIN/releases



Tools for EW precision calculation pre-2022:

- **ZFITTER/DIZET, TOPAZØ, ...:**

Bardin et al. '99
Montagna et al. '98

- rad. corr. packages developed for LEP era
- SM prediction for EWPOs and (diff.) cross-sections ($e^+e^- \rightarrow f\bar{f}$)
- Full NLO corrections + partial higher orders
- QED ISR/FSR corrections through analyt. formulae
- Can be linked with MC codes (**KoralZ**, ...)
- Difficult to expand and maintain
(Fortran77, not fully gauge-invariant framework, ...)

Jadach et al. '80s–99

- Modern fitting tools (**Gfitter**, **HEPfit**, **GAPP**):

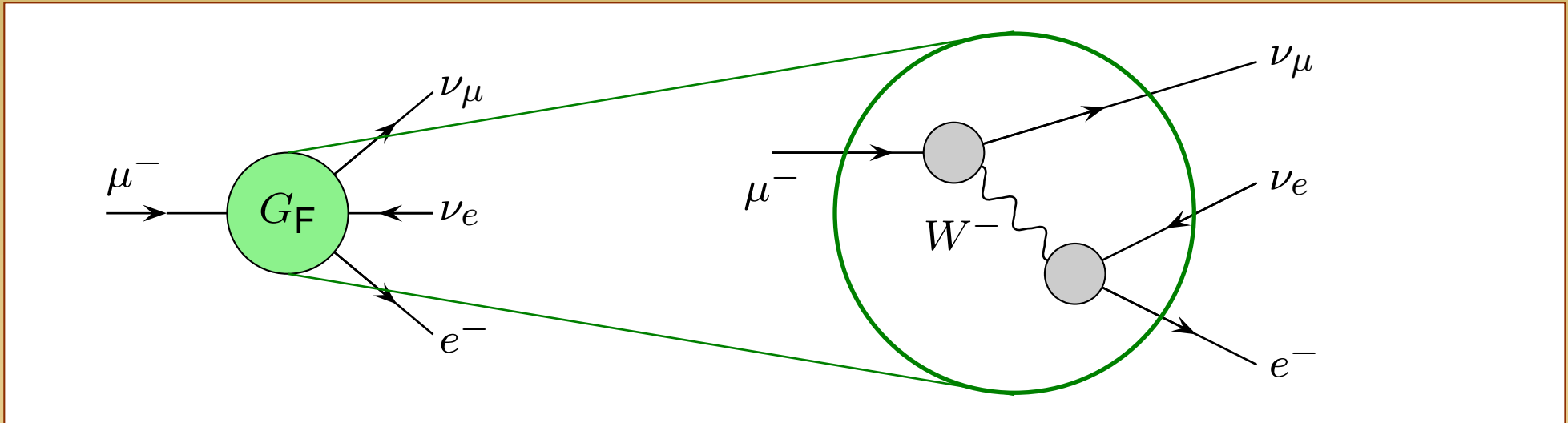
Baak et al. '14
de Blas et al. '19; Eler '00

- own implementations of rad. corr.
[only EWPOs (pseudo-obs.), not full observables]
- extensions to higher orders, different schemes and models require custom work

Goal of the GRIFFIN* project:

- New EW library that is modular / object-oriented (C++)
- Based on manifestly gauge-invariant setup
- Repository of existing calculations
- Can be extended to include ...
 - ... higher orders
 - ... different input parameter schemes
 - ... BSM physics (also SMEFT/HEFT)
 - ... new processes or new EWPOs
- Can be linked to MC generators and global fitting packages (QED more effectively handled with MC generators)

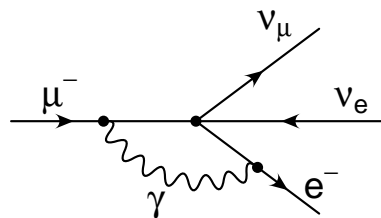
* Gauge-invariant Resonance In Four-Fermion Interactions



μ decay in Fermi Model

$$\Gamma_{\mu} = \frac{G_F m_{\mu}^5}{192\pi^3} F\left(\frac{m_e^2}{m_{\mu}^2}\right) (1 + \Delta q)$$

QED corrections (2-loop)

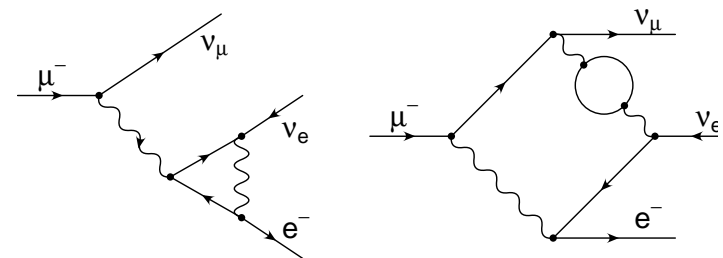


Ritbergen, Stuart '98
Pak, Czarnecki '08

μ decay in Standard Model

$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8s_W^2 M_W^2} (1 + \Delta r)$$

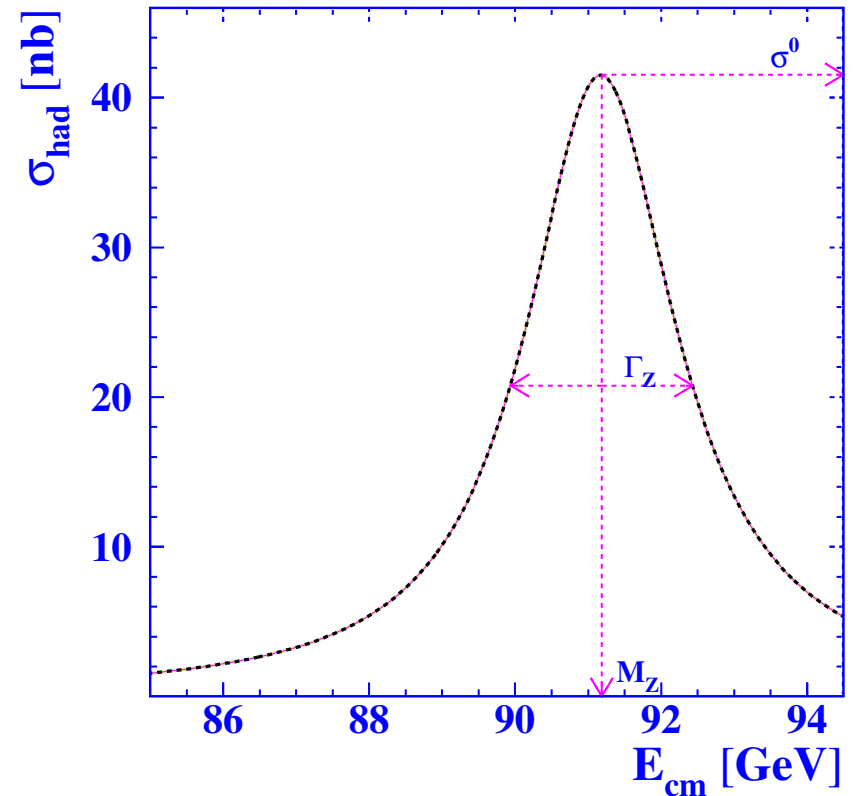
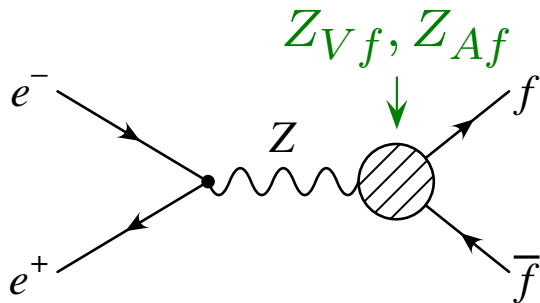
electroweak corrections



$e^+e^- \rightarrow f\bar{f}$ for $E_{\text{CM}} \sim M_Z$:

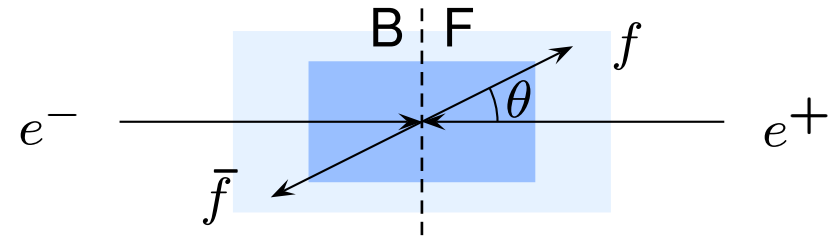
- Mass M_Z
- Width $\Gamma_Z = \sum_f \Gamma_{ff}$
- Branching ratio $R_f = \Gamma_{ff}/\Gamma_Z$
- $\sigma^0 \approx \frac{12\pi \Gamma_{ee} \Gamma_{ff}}{(s-M_Z^2)^2 + M_Z^2 \Gamma_Z^2} = \frac{12\pi}{M_Z^2} R_e R_f$

$$\Gamma_{ff} = C [(Z_{Vf})^2 + (Z_{Af})^2]$$



Forward-backward asymmetry:

$$A_{\text{FB}} \equiv \frac{\sigma_{\text{F}} - \sigma_{\text{B}}}{\sigma_{\text{F}} + \sigma_{\text{B}}} = \frac{3}{4} A_e A_f$$



$$A_f = \frac{2(1 - 4\sin^2 \theta_{\text{eff}}^f)}{1 + (1 - 4\sin^2 \theta_{\text{eff}}^f)^2}$$

$$\sin^2 \theta_{\text{eff}}^f = \frac{1}{4|Q_f|} \left[1 - \text{Re} \frac{Z_{Vf}}{Z_{Af}} \right]$$

Left-right asymmetry:

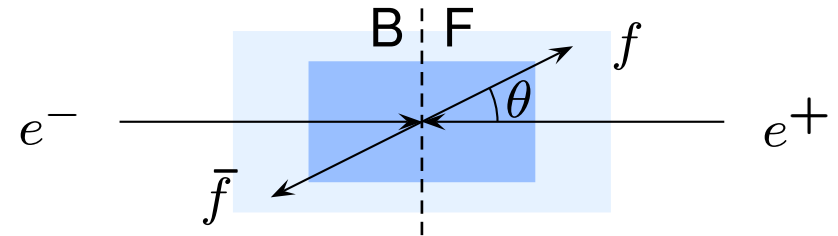
With polarized e^- beam:
$$A_{\text{LR}} \equiv \frac{\sigma_{\text{L}} - \sigma_{\text{R}}}{\sigma_{\text{L}} + \sigma_{\text{R}}} = A_e$$

Polarization asymmetry:

Average τ pol. in $e^+e^- \rightarrow \tau^+\tau^-$:
$$\langle \mathcal{P}_\tau \rangle = -A_\tau$$

Forward-backward asymmetry:

$$A_{FB} \equiv \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = \frac{3}{4} A_e A_f$$



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Polarization asymmetry:

Average τ pol. in $e^+e^- \rightarrow \tau^+\tau^-$:
$$\langle \mathcal{P}_\tau \rangle = -A_\tau$$

Decay widths in terms of $\sin^2 \theta_{\text{eff}}^f$:

$$\Gamma_{ff} = C \left[(Z_{Vf})^2 + (Z_{Af})^2 \right] = C (Z_{Af})^2 \left[(1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f)^2 + \left(\text{Im} \frac{Z_{Vf}}{Z_{Af}} \right)^2 \right]$$

Available results for Δr , $\sin^2 \theta_{\text{eff}}^f$, Γ_{ff} :

- Many seminal works on 1-loop and leading 2-loop corrections

Veltman, Passarino, Sirlin, Marciano, Bardin, Hollik, Riemann, Degrossi, Kniehl, ...

- Full 2-loop results

Freitas, Hollik, Walter, Weiglein '00

Awramik, Czakon '02

Onishchenko, Veretin '02

Awramik, Czakon, Freitas, Weiglein '04

Awramik, Czakon, Freitas '06

Hollik, Meier, Uccirati '05,07

Awramik, Czakon, Freitas, Kniehl '08

Freitas '14

Dubovyk, Freitas, Gluza, Riemann, Usovitsch '16,18

- Partial higher orders: $\mathcal{O}(\alpha_t \alpha_s^2)$, $\mathcal{O}(\alpha_t^2 \alpha_s)$, $\mathcal{O}(\alpha_t^3)$, $\mathcal{O}(\alpha_t \alpha_s^3)$,

$\mathcal{O}(N_f^3 \alpha^3)$, $\mathcal{O}(N_f^2 \alpha^2 \alpha_s)$

$$\alpha_t = \frac{y_t^2}{4\pi}$$

$N_f =$ closed fermion loop

Chetyrkin, Kühn, Steinhauser '95

Faisst, Kühn, Seidensticker, Veretin '03

Boughezal, Tausk, v. d. Bij '05

Schröder, Steinhauser '05

Chetyrkin et al. '06

Boughezal, Czakon '06

Chen, Freitas '20

EWPOs like $\sin^2 \theta_{\text{eff}}^f$, Γ_{ff} only contain **leading** EW corrections **on Z-pole**

For complete prediction need:

- Sub-leading corrections
- Off-peak matrix element
- QED/QCD ISR/FSR contributions

Expand amplitude for $e^+e^- \rightarrow f\bar{f}$ about **complex pole** $s_0 \equiv \overline{M_Z^2} + i\overline{M_Z}\overline{\Gamma_Z}$:

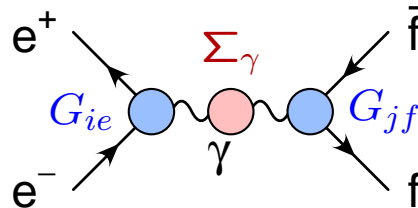
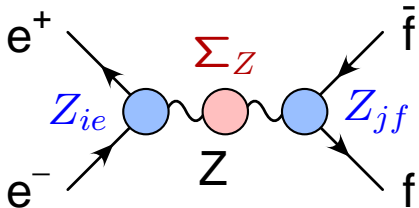
→ All terms are individually gauge-invariant

$$\mathcal{M}_{ij} = \frac{R_{ij}}{s - s_0} + S_{ij} + (s - s_0)S'_{ij} + \dots \quad (i, j = V, A)$$

$$R_{ij} = \left. \frac{Z_{ie}Z_{jf}}{1 + \Sigma'_Z} \right|_{s=s_0}$$

$$S_{ij} = \left[\frac{Z_{ie}Z'_{jf} + Z'_{ie}Z_{jf}}{1 + \Sigma'_Z} - \frac{Z_{ie}Z_{jf}\Sigma''_Z}{2(1 + \Sigma'_Z)^2} + \frac{G_{ie}G_{jf}}{s + \Sigma_\gamma} + B_{ij} \right]_{s=s_0}$$

$$S'_{ij} = \dots$$



Expand amplitude for $e^+e^- \rightarrow f\bar{f}$ about **complex pole** $s_0 \equiv \overline{M_Z^2} + i\overline{M_Z}\overline{\Gamma_Z}$:

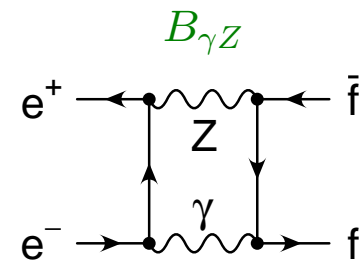
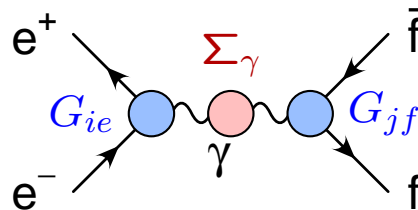
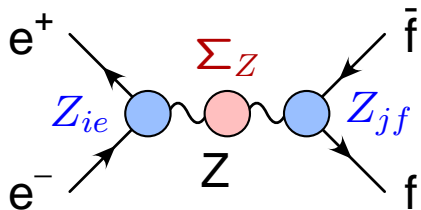
→ All terms are individually gauge-invariant

$$\mathcal{M}_{ij} = \frac{R_{ij}}{s - s_0} + S_{ij} + (s - s_0)S'_{ij} + \dots \quad (i, j = V, A)$$

$$R_{ij} = \left. \frac{Z_{ie}Z_{jf}}{1 + \Sigma'_Z} \right|_{s=s_0} + B_{\gamma Z, ij}^R + B_{\gamma Z, ij}^{RL} \ln\left(1 - \frac{s}{s_0}\right)$$

$$S_{ij} = \left[\frac{Z_{ie}Z'_{jf} + Z'_{ie}Z_{jf}}{1 + \Sigma'_Z} - \frac{Z_{ie}Z_{jf}\Sigma''_Z}{2(1 + \Sigma'_Z)^2} + \frac{G_{ie}G_{jf}}{s + \Sigma_\gamma} + B_{ij} \right]_{s=s_0} + B_{\gamma Z, ij}^S + B_{\gamma Z, ij}^{SL} \ln\left(1 - \frac{s}{s_0}\right)$$

$$S'_{ij} = \dots$$



Expand amplitude for $e^+e^- \rightarrow f\bar{f}$ about **complex pole** $s_0 \equiv \overline{M}_Z^2 + i\overline{M}_Z\overline{\Gamma}_Z$:

$$\mathcal{M}_{ij} = \frac{R_{ij}}{s - s_0} + S_{ij} + (s - s_0)S'_{ij} + \dots \quad (i, j = V, A)$$

Current state of art: R @ NNLO + leading higher orders

S @ NLO

S' @ (N)LO

For future ee colliders: (at least) one order more!

Cross-section:
$$\sigma_Z = \frac{R}{(s - \overline{M}_Z^2)^2 + \overline{M}_Z^2\overline{\Gamma}_Z^2} + \sigma_{\text{non-res}}$$

In exp. studies:
$$\sigma \sim \frac{1}{(s - M_Z^2)^2 + s^2\Gamma_Z^2/M_Z^2}$$

$$\overline{M}_Z = M_Z / \sqrt{1 + \Gamma_Z^2/M_Z^2} \approx M_Z - 34 \text{ MeV}$$

$$\overline{\Gamma}_Z = \Gamma_Z / \sqrt{1 + \Gamma_Z^2/M_Z^2} \approx \Gamma_Z - 0.9 \text{ MeV}$$

Express R_{ij} in terms of $\sin^2 \theta_{\text{eff}}^f$ and F_A^f (accurate up to NNLO):

$$\sin^2 \theta_{\text{eff}}^f = \frac{1}{4|Q_f|} \left[1 - \text{Re} \frac{Z_{Vf}}{Z_{Af}} \right]_{s=\overline{M}_Z^2}$$

$$F_i^f = \left[\frac{|Z_{if}|^2}{1 + \text{Re} \Sigma'_Z} - \frac{1}{2} \overline{M}_Z \overline{\Gamma}_Z |Z_{if(0)}|^2 \text{Im} \Sigma''_{Z(1)} \right]_{s=\overline{M}_Z^2} + \mathcal{O}(\alpha^3)$$

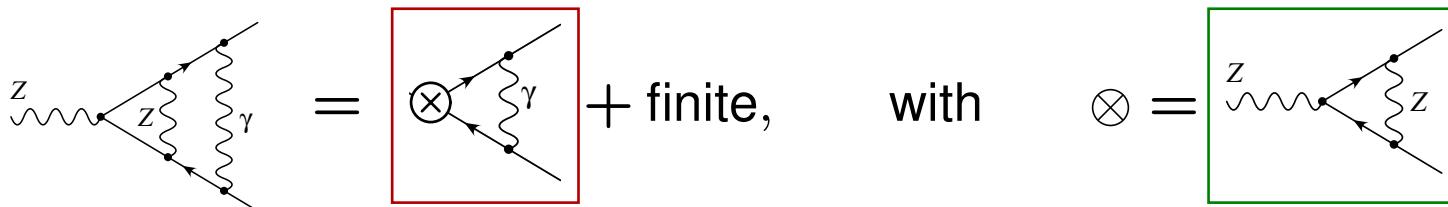
$$\begin{aligned} \Rightarrow R_{ij} &= \frac{Z_{ie} Z_{jf}}{1 + \Sigma'_Z} \Big|_{s=s_0} \\ &= 4I_e^3 I_f^3 \sqrt{F_A^e F_A^f} \left[Q_i^e Q_j^f + \dots \right], \quad \begin{aligned} Q_V^f &= 1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f, \\ Q_A^f &= 1 \end{aligned} \end{aligned}$$

terms related to
 $\text{Im} Z_{ie}, \text{Im} Z_{jf}, \text{Im} \Sigma_Z$
 and the γZ box

(n) = loop order

Factorization of massive EW corrections and QED/QCD ISR/FSR:

$$Z_{if}^{\text{tot}} = R_f^i \times Z_{if}, \quad R_f^V(s) \equiv \frac{\mathcal{M}_{V^* \rightarrow f\bar{f}}^{\text{QED/QCD}}}{\mathcal{M}_{V^* \rightarrow f\bar{f}}^{\text{Born}}}, \quad R_f^A(s) \equiv \frac{\mathcal{M}_{A^* \rightarrow f\bar{f}}^{\text{QED/QCD}}}{\mathcal{M}_{A^* \rightarrow f\bar{f}}^{\text{Born}}},$$



R_V^f, R_A^f : QED/QCD radiation factors;

FSR known inclusively to $\mathcal{O}(\alpha_s^4)$, $\mathcal{O}(\alpha^2)$, $\mathcal{O}(\alpha\alpha_s)$ Chetyrkin, Kühn, Kwiatkowski '96
Kataev '92; Baikov, Chetyrkin, Kühn, Rittinger '12

ISR via structure functions with LL resummation

Kureav, Fadin '85

Montagna, Nicosini, Piccinini '97

Ablinger, Blümlein, De Freitas, Schönwald '20

or compute exclusively using MC methods,

e.g. **KKMC**,

Arbuzov, Jadach, Wąs, Ward, Yost '20

SHERPA_YFS,

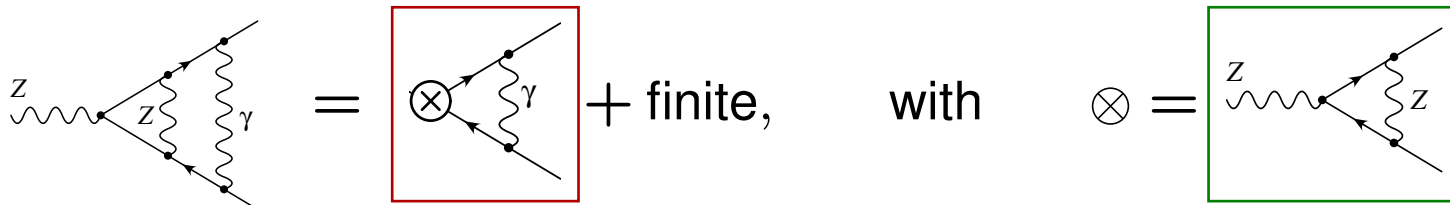
Krauss, Price, Schönherr '22

POWHEG_EW

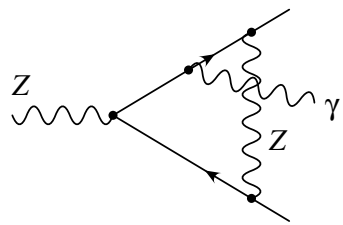
Barzè, Montagna, Nason, Nicosini, Piccinini '12,13

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$$Z_{if}^{\text{tot}} = R_f^i \times Z_{if}, \quad R_f^V(s) \equiv \frac{\mathcal{M}_{V^* \rightarrow f\bar{f}}^{\text{QED/QCD}}}{\mathcal{M}_{V^* \rightarrow f\bar{f}}^{\text{Born}}}, \quad R_f^A(s) \equiv \frac{\mathcal{M}_{A^* \rightarrow f\bar{f}}^{\text{QED/QCD}}}{\mathcal{M}_{A^* \rightarrow f\bar{f}}^{\text{Born}}},$$



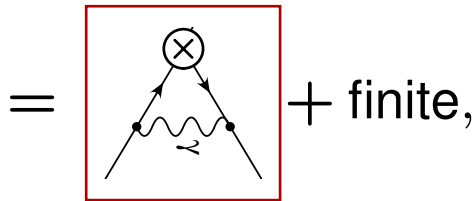
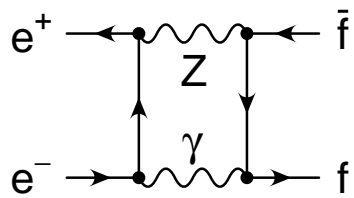
Additional non-factorizable contributions, e.g.



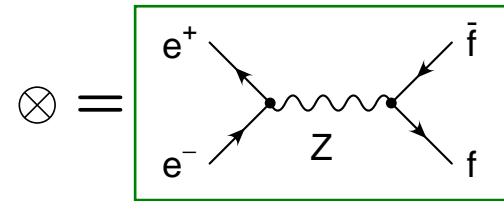
- Incorporated in F_A^f , $\sin^2 \theta_{\text{eff}}^f$ form factors
- Known at $\mathcal{O}(\alpha\alpha_s)$ Czarnecki, Kühn '96
Harlander, Seidensticker, Steinhauser '98
- Currently not known at $\mathcal{O}(\alpha^2)$ and beyond

QED soft IR singular pieces for IFI also factorize:

$$B_{ij(1)} = B_{ij(1)}^{\text{tot}} - \mathcal{M}_{ij(0)} 2Q_e Q_f [R_{e(1)}(t) - R_{e(1)}(u)]$$



with



Pole expansion works well in window of few GeV about Z pole, but not beyond

$$\mathcal{M}_{ij}^{\text{exp},s_0} = \frac{R_{ij}}{s - s_0} + S_{ij} + (s - s_0)S'_{ij} + \dots$$

Outside this window, use $\mathcal{M}_{ij}^{\text{noexp}}$ without expansion in s and Dyson summation:

$$\mathcal{M}_{ij} = \mathcal{M}_{ij}^{\text{exp},s_0} + \mathcal{M}_{ij}^{\text{noexp}} - \mathcal{M}_{ij}^{\text{exp},\bar{M}_Z}$$



To avoid double counting and cancel unphys. pole at $s = \bar{M}_Z^2$ in $\mathcal{M}_{ij}^{\text{noexp}}$:

$$M_{ij}^{\text{exp},\bar{M}_Z^2} = \mathcal{T}_\alpha \left\{ \mathcal{M}_{ij}^{\text{exp},s_0} \Big|_{s_0 = \bar{M}_Z^2 - i\bar{M}_Z \alpha \bar{\Gamma}_Z(1)} \right\}$$

$\mathcal{T}_x =$ Taylor operator in x

- See also [Dittmaier, Huber '09](#)
- Could also use complex-mass scheme to compute $\mathcal{M}_{ij}^{\text{noexp}}$

[Denner, Dittmaier, Roth, Wieders '05](#); [Denner, Dittmaier '06](#)

- Implementation in GRIFFIN v1.0:

$$\mathcal{M}_{ij}^{\text{exp},s_0} = \frac{R_{ij}}{s - s_0} + S_{ij} + (s - s_0)S'_{ij} + \dots$$

\uparrow \uparrow
@NNLO @NLO

$$\mathcal{M}_{ij} = \mathcal{M}_{ij}^{\text{exp},s_0} + \mathcal{M}_{ij}^{\text{noexp}} - \mathcal{M}_{ij}^{\text{exp},\bar{M}_Z}$$

\nwarrow \nearrow
@NLO

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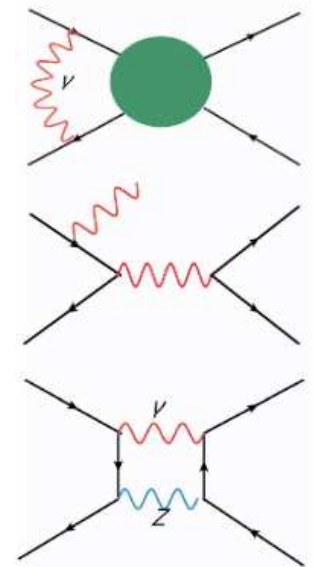
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@NNLO @NLO

$$\mathcal{M}_{ij} = \mathcal{M}_{ij}^{\text{exp},s_0} + \mathcal{M}_{ij}^{\text{noexp}} - \mathcal{M}_{ij}^{\text{exp},\bar{M}_Z}$$

\nwarrow \nearrow
@NLO

- QED contributions have been subtracted



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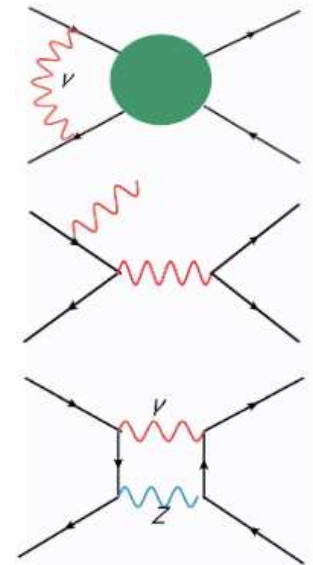
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\uparrow
 \uparrow
@NNLO
@NLO

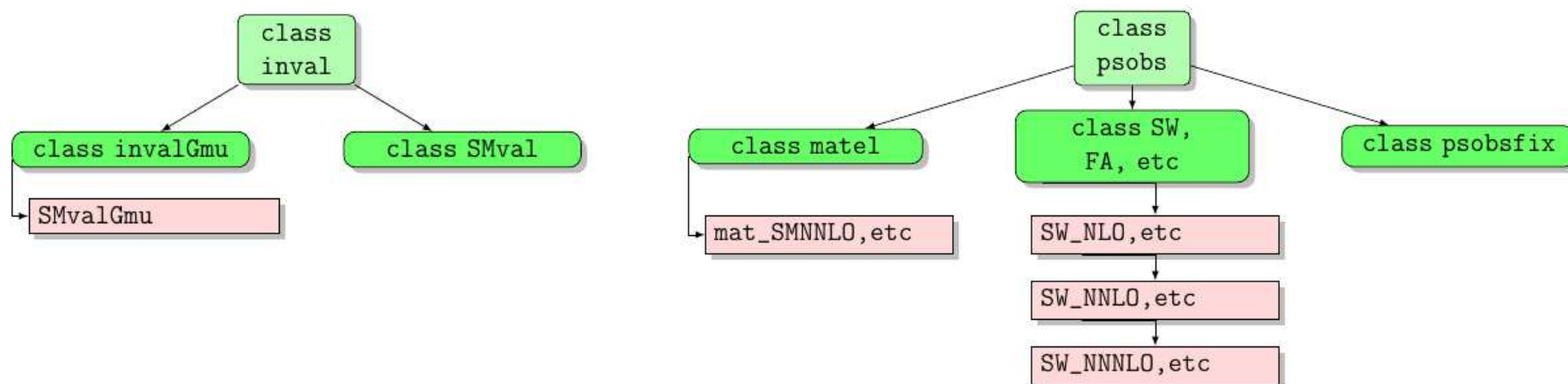
$$\mathcal{M}_{ij} = \mathcal{M}_{ij}^{\text{exp},s_0} + \mathcal{M}_{ij}^{\text{noexp}} - \mathcal{M}_{ij}^{\text{exp},\bar{M}_Z}$$

\nwarrow
 \nearrow
@NLO

- QED contributions have been subtracted
- SM predictions for EWPOs (Δr , $\sin^2 \theta_{\text{eff}}^f$, F_A^f) at NNLO+



class inval		class psobs	
input parameters (in the SM)		output observables	
Boson masses and widths	$M_{W,Z,H}$ $\Gamma_{W,Z}$	pseudo-observables defined at Z-peak	$F_{V,A}, \sin^2 \theta_{eff}^f$ $\Gamma_{Z \rightarrow f\bar{f}}, \Delta r,$ etc.
Fermion masses	$m_{e,\mu,\tau}^{OS}$ $m_{d,u,s,c}^{\overline{MS}}(M_Z)$ m_t^{OS}	amplitude coefficients under pole scheme	$R, S, \text{ and } S'$
Couplings	$\alpha(0)$ $\Delta\alpha \equiv 1 - \alpha(0)/\alpha(M_Z^2)$ $\alpha_s^{\overline{MS}}(M_Z^2), G_\mu$	(polarized) matrix element square near Z-peak	$\text{Re } M_{ij} M_{kl}^*$



```
#include <iostream>
using namespace std;

#include "EWPOZ2.h"
#include "xscnnlo.h"
#include "SMval.h"

int main()
{
    SMval myinput; // convert masses from PDG values to complex pole scheme
    myinput.set(a1, 1/137.03599976);
    myinput.set(MZ, 91.1876);
    myinput.set(MW, 80.377);
    myinput.set(GamZ, 2.4952);
    myinput.set(GamW, 2.085);
    myinput.set(MH, 125.1);
    myinput.set(MT, 172.5);
    myinput.set(MB, 2.87);
    myinput.set(Delal, 0.059);
    myinput.set(als, 0.1179);

    cout << endl << "Complex-pole masses: MW=" << myinput.get(MWc) << ", MZ="
         << myinput.get(MZc) << endl << endl;
```

```
// compute matrix element for ee->dd with vector coupling in initial
// state and vector coupling in final state
int ini = ELE, fin = DQU, iff = VEC, off = VEC;

cout << "=== Matrix element for ee->dd (i=e, f=d) ===" << endl << endl;

// compute vertex form factors:
FA_SMNNLO FAi(ini, myinput), FAf(fin, myinput);
SW_SMNNLO SWi(ini, myinput), SWf(fin, myinput);
cout << "F_A^i (NNLO+) = " << FAi.result() << endl;
cout << "F_A^f (NNLO+) = " << FAf.result() << endl;
cout << "sineff^i (NNLO+) = " << SWi.result() << endl;
cout << "sineff^f (NNLO+) = " << SWf.result() << endl;
cout << endl;
```

```
double cme,          // center-of-mass energy
       cost = 0.5; // scattering angle
Cplx res1, res2;

cout << "SM matrix element M_VV for cos(theta)=" << cost << ": " << endl;
// compute matrix element for ee->dd using SM form factors:
mat_SMNLO M(ini, fin, iff, off, FAi, FAf, SWi, SWf, cme*cme, cost,
myinput);
cout << "sqrt(s)\t\ttot. result\t\ttoff-resonance contrib." << endl;
for(cme = 10.; cme <= 190.; cme += 20.)
{
    M.setkinvar(cme*cme, cost);
    res1 = M.result();
    res2 = M.resoffZ();
    cout << cme << " \t" << res1 << " \t" << res2 << endl;
}
cout << endl;

return 0;
}
```


Complex-pole masses: MW=80.35, MZ=91.1535

=== Matrix element for ee->dd (i=e, f=d) ===

F_A^i (NNLO+) = (0.034499,0)

F_A^f (NNLO+) = (0.0345443,0)

sineff^i (NNLO+) = (0.231172,0)

sineff^f (NNLO+) = (0.230985,0)

SM matrix element M_VV for cos(theta)=0.5:

sqrt(s)	tot. result	off-resonance contrib.
10	(0.000316739,-5.58082e-06)	(0.000309429,-5.53734e-06)
30	(3.53793e-05,-5.99317e-07)	(2.84458e-05,-5.59139e-07)
50	(1.25851e-05,-1.90789e-07)	(6.4247e-06,-1.59184e-07)
70	(6.07798e-06,-5.97311e-08)	(1.19433e-06,-4.81728e-08)
90	(-7.31188e-07,-3.55673e-06)	(8.7104e-09,-1.80673e-09)
110	(3.14635e-06,-1.62001e-07)	(4.59289e-07,1.10821e-08)
130	(2.12596e-06,-7.90095e-08)	(1.82894e-06,1.92144e-08)
150	(1.5668e-06,-5.34561e-08)	(3.83515e-06,2.49419e-08)
170	(1.20884e-06,-3.97403e-08)	(6.35319e-06,2.97998e-08)
190	(9.60973e-07,-3.33532e-08)	(9.31833e-06,3.12732e-08)

□ Numerical Results:

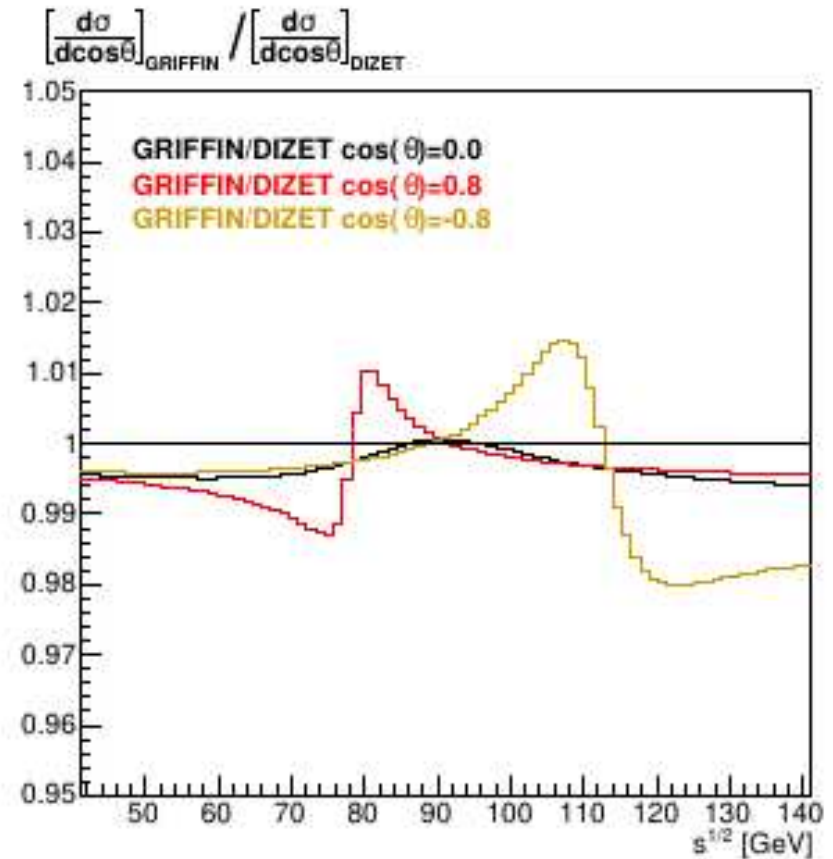
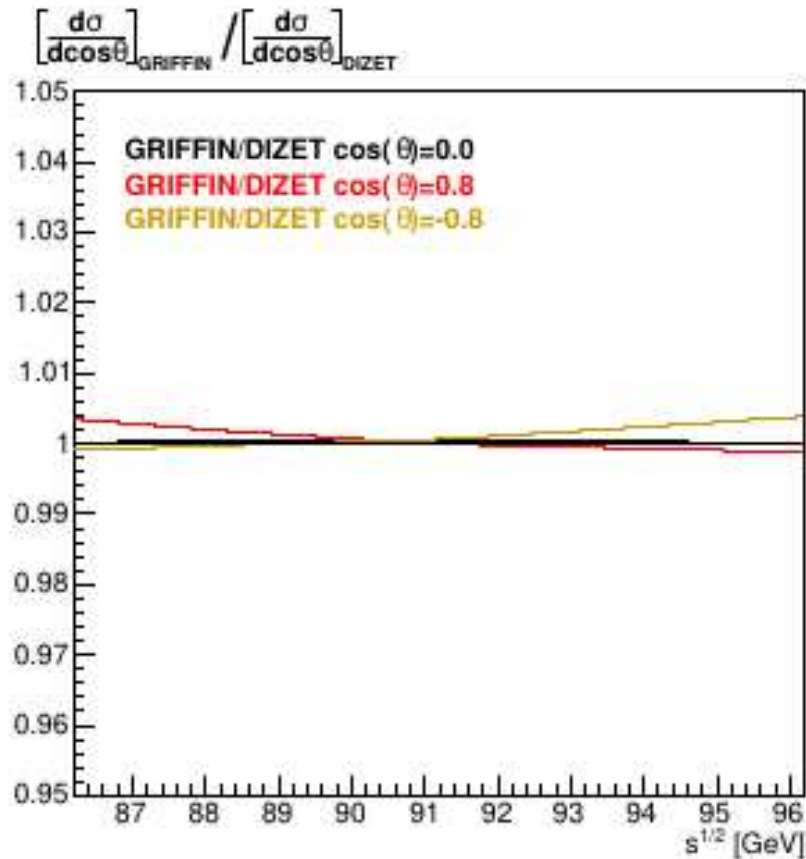
$$|\rho_Z^f| = \frac{2\sqrt{2}F_A^f}{G_\mu M_Z^2}$$

	$ \rho_Z^f $		$\sin^2 \theta_{\text{eff}}^f$		$\Gamma_{Z \rightarrow f\bar{f}}$	
	DIZET 6.45	GRIFFIN	DIZET 6.45	GRIFFIN	DIZET 6.45	GRIFFIN
$\nu\bar{\nu}$	1.00800	1.00814	0.231119	NAN	0.167206	0.167197
$\ell\bar{\ell}$	1.00510	1.00519	0.231500	0.231534	0.083986	0.083975
$u\bar{u}$	1.00578	1.00573	0.231393	0.231420	0.299938	0.299958
$d\bar{d}$	1.00675	1.00651	0.231266	0.231309	0.382877	0.382846
$b\bar{b}$	0.99692	0.99420	0.232737	0.23292	0.376853	0.377432

	DIZET 6.45	GRIFFIN all orders	GRIFFIN $\mathcal{O}(\alpha, \alpha^2, \alpha_t \alpha_s, \alpha_t \alpha_s^2)$
Δr	3.63947×10^{-2}	3.68836×10^{-2}	3.63987×10^{-2}

- Not a **one-one-one match**. (no leading N3LO implemented in dizet v.6.45)
- most numbers are in agreement up to at least **4-digit**. The actual discrepancy is in the realm of missing N3(4)LO.
- fictitious discrepancies stem from the input scheme/definition of the form factors/EWPOs.

Ratios of differential cross-sections for $e^+e^- \rightarrow \mu^+\mu^-$ for different θ :



- $\lesssim \mathcal{O}(10^{-3})$ agreement near Z-pole (\sim NNLO precision)
- %-level agreement away from Z pole (NLO prec., different implementations)

[Note: enhanced corrections when tree-level matrix element is small]

GRIFFIN v1.0: $f\bar{f} \rightarrow f'\bar{f}'$ with NLO EW corrections and h.o. @ Z-pole

Future upgrades:

- Bhabha scattering ($f = f'$)
- Higher-order off-resonance corrections, e.g.
 $\mathcal{O}(\alpha\alpha_s)$, Heller, v.Manteuffel, Schabinger, Spiesberger '20
 $\mathcal{O}(N_f\alpha^2)$ Bonciani et al. '21
- SMEFT $d=6$ operator effects
- W production and decay (a.k.a. charged-current DY)

Try out the code: github.com/lisongc/GRIFFIN/releases

Feedback welcome!

Backup slides

Implementation of higher-order corrections

Corrections entering through $\delta\rho$:

	drho2aas	$\mathcal{O}(\alpha_t\alpha_s)$	[3, 4]
	drho2a2	$\mathcal{O}(\alpha_t^2)$	[5–9]
*	drho3aas2	$\mathcal{O}(\alpha_t\alpha_s^2)$	[10, 11]
*	drho3a2as	$\mathcal{O}(\alpha_t^2\alpha_s)$	[12, 13]
*	drho3a3	$\mathcal{O}(\alpha_t^3)$	[12, 13]
*	drho4aas3	$\mathcal{O}(\alpha_t\alpha_s^3)$	[14–16]

Full corrections to $F_A^f, \sin^2\theta_{\text{eff}}^f$:

*	res2ff	$\mathcal{O}(\alpha_f^2)$	[17–19]
*	res2fb	$\mathcal{O}(\alpha_f\alpha_b)$	[17–20]
*	res2bb	$\mathcal{O}(\alpha_b^2)$	[21–25]
*	res2aas	$\mathcal{O}(\alpha\alpha_s)$	[26–28] (correction to internal gauge-boson self-energies)
*	res2aasnf	$\mathcal{O}(\alpha\alpha_s)$	[29–34] (non-factorizable final-state corrections for $f = q$)
*	res3fff	$\mathcal{O}(\alpha_f^3)$	[35]
*	res3ffa2as	$\mathcal{O}(\alpha_f^2\alpha_s)$	[36]

Z lineshape

■ Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

Kureav, Fadin '85

Berends, Burgers, v. Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v. Neerven '89

Bardin et al. '91; Skrzypek '92

Montagna, Nicrosini, Piccinini '97

Soft photons (resummed) + collinear photons

$$\mathcal{R}_{\text{ini}} = \sum_n \left(\frac{\alpha}{\pi}\right)^n \sum_{m=0}^n h_{nm} \ln^m\left(\frac{s}{m_e^2}\right)$$

Universal ($m=n$) logs known to $n = 6$,

also some sub-leading terms

Ablinger, Blümlein, De Freitas, Schönwald '20

Exclusive description: MC tools

